# Polarization effect in $Cd_{1-x}Zn_xTe$ and $Cd_{1-x}Mn_xTe$ detectors under various biasing conditions O. Amzallag<sup>1</sup>, L. Chernyak<sup>2</sup>, A. Ruzin<sup>1</sup>

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Wide bandgap alloys of II-VI group materials are widely used for gamma-photon detection. The more established CdZnTe alloy is recently being challenged by CdMnTe, which shown improved uniformity and wider tunable bandgap. These compounds mostly owe their high resistivity to deep level compensation process (Fermi level "pinning"). Therefore, such compound "semi-insulators" have high densities of traps. In spite of that fact, such detectors exhibit reasonable charge collection. However, when they are exposed to high fluxes, considerable polarization is often observed. The later occurs due to high volume trapping, leading to modification of internal electric field. In this work we present the polarization differences in CdZnTe and CdMnTe devices grown by the same method. The study is performed by reconstruction of electric filed, using improved TCT method.

CMT



A pulse laser illuminates the detector through the front contact and creates charge carriers near it. The charge carriers drift in the presence of a field. Their movement creates induced current in contacts, which is the output of the system. We applied negative voltage to measure



electron movement.

### Transient Current Technique system



## Populating the traps

#### The reconstructed electric field from the measured current pulses

According to Shockley-Ramo theorem, the relationship between the instantaneous induced current and the electric field is:

 $I_{e,h}(t) = q\vec{v}(\vec{r}(t)) \cdot \vec{E}_{w}(\vec{r}(t)) = N_{e,h}(t) \cdot e^{\vec{\tau}_{eff,e,h}} \cdot \mu_{e,h} \cdot \vec{E}(\vec{r}(t)) \cdot \vec{E}_{w}(\vec{r}(t))$ 

where,  $\vec{v}$ - charge carrier velocity,  $\vec{r}$ - location of the charge carrier,  $\mu$ - mobility,  $\vec{E}_w$ - the weighting field given by the in the detector.

life time,  $\tau_{eff,e}$ , was extracted from lower voltage measurements. Estimated electron's mobility:  $\mu_e = 1000 \frac{v \cdot s}{cm^2}$ .





### Recombination in the front electrode (near the generation area)

![](_page_0_Figure_18.jpeg)

significant for a shallow radiation absorption low energy gamma rays, alpha, low energy x rays.

- U-shape of the field distribution is expected, as the detectors produced uniformly and underwent annealing treatment.
- The asymmetry in the electric field around the center of the detector (0.25 [cm]) can be a result of limit recombination velocity in the contact or no uniform absorption and mobility.
- In CZT, the stabilization under voltage is faster than in CMT.
- The duration of depolarization is similar in both detectors (16/17)hours).
- In CZT, the electric field near the front contact becomes lower and higher near the back contact under voltage.
- In CMT, the main change in the electric field is in the detector's depth.

The effect of populating the traps under voltage is significant for a deep absorption -

high energy gamma rays and high energy x rays.

The CCE in CZT is 25% greater than in CMT. The effect of polarization on the characteristics of CZT and CMT detectors, must be taken into account. Especially when it comes to a shallow radiation absorption.

Future research: Polarization effect in  $Cd_{1-x}Zn_{x}Te$  and  $Cd_{1-x}Mn_{x}Te$  detectors under radiation.