

Systematics Uncertainties in Small-scale Experiments

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Disclaimer

- Small-scale experiments —> Single-purpose experiments
 - e.g. Muon $g-2$ experiment, ~ 200 authors
- Technical aspects already covered in previous lectures, I will concentrate on practical applications
 - A very biased selection of experimental examples, mostly from muon and QED precision physics
- Not intended as a review of physics results
 - you could find incomplete references, out-of-date results, etc.

Outline

- General aspects of systematic uncertainties in single-purpose experiments
 - precision measurements
 - rare event searches
- Inclusion of systematics in confidence interval computations
- (Very biased) collection of relevant examples

Generalities

Single-purpose experiments

- Most single-purpose experiments can be classified into two categories:
 - 1) highly accurate measurements of particles properties
 - 2) searches for rare processes (rare decays, interactions of elusive particles)
- Besides accumulating statistics, high sensitivity is achieved through **high precision** (extremely good resolutions, extremely high background rejection) and/or **high accuracy** (no bias in measurements and background estimates)

High Accuracy vs. High precision

PHYSICAL REVIEW LETTERS **131**, 161802 (2023)

Editors' Suggestion

Measurement of the Positive Muon Anomalous Magnetic Moment to 0.20 ppm

Measurement based on the rate of e^+ from μ^+ decays, above a given energy threshold
(**2.5%** e^+ energy resolution)

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ω_a^m (statistical)	...	201
ω_a^m (systematic)	...	25
C_e	451	32
C_p	170	10
C_{pa}	-27	13
C_{dd}	-15	17
C_{ml}	0	3
$f_{\text{calib}} \cdot \langle \omega'_p(\vec{r}) \times M(\vec{r}) \rangle$...	46
B_k	-21	13
B_q	-21	20
$\mu'_p(34.7^\circ)/\mu_e$...	11
m_μ/m_e	...	22
$g_e/2$...	0
Total systematic for \mathcal{R}'_μ	...	70
Total external parameters	...	25
Total for a_μ	622	215

High Accuracy vs. High precision

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Eur. Phys. J. C (2024) 84:216
<https://doi.org/10.1140/epjc/s10052-024-12416-2>

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Regular Article - Experimental Physics

A search for $\mu^+ \rightarrow e^+\gamma$ with the first dataset of the MEG II experiment

MEG II Collaboration

Upper Limit based on the discrimination of 2-body vs. 3-body kinematics with extremely precise measurements
(e.g. **0.2%** e^+ energy resolution)

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Parameter	Impact on limit
$\phi_{e\gamma}$ uncertainty	1.1 %
E_γ uncertainty	0.9 %
$\theta_{e\gamma}$ uncertainty	0.7 %
Normalization uncertainty	0.6 %
$t_{e\gamma}$ uncertainty	0.1 %
E_e uncertainty	0.1 %
RDC uncertainty	< 0.1 %

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Regular Article - Experimental Physics

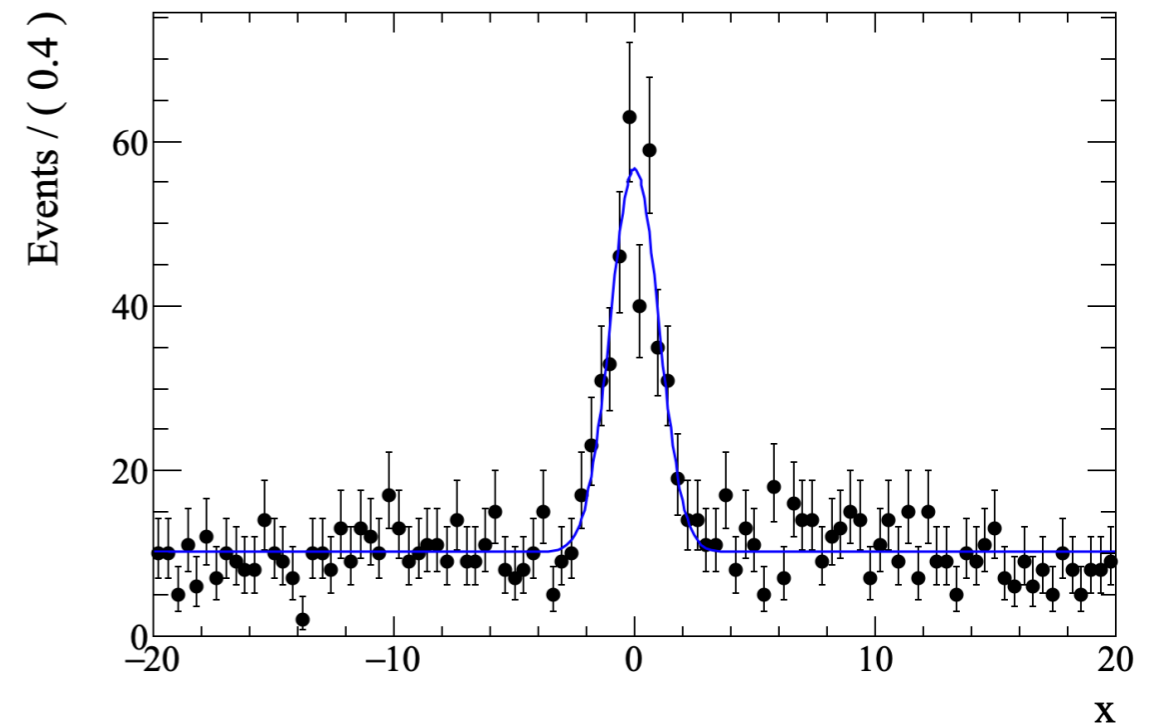
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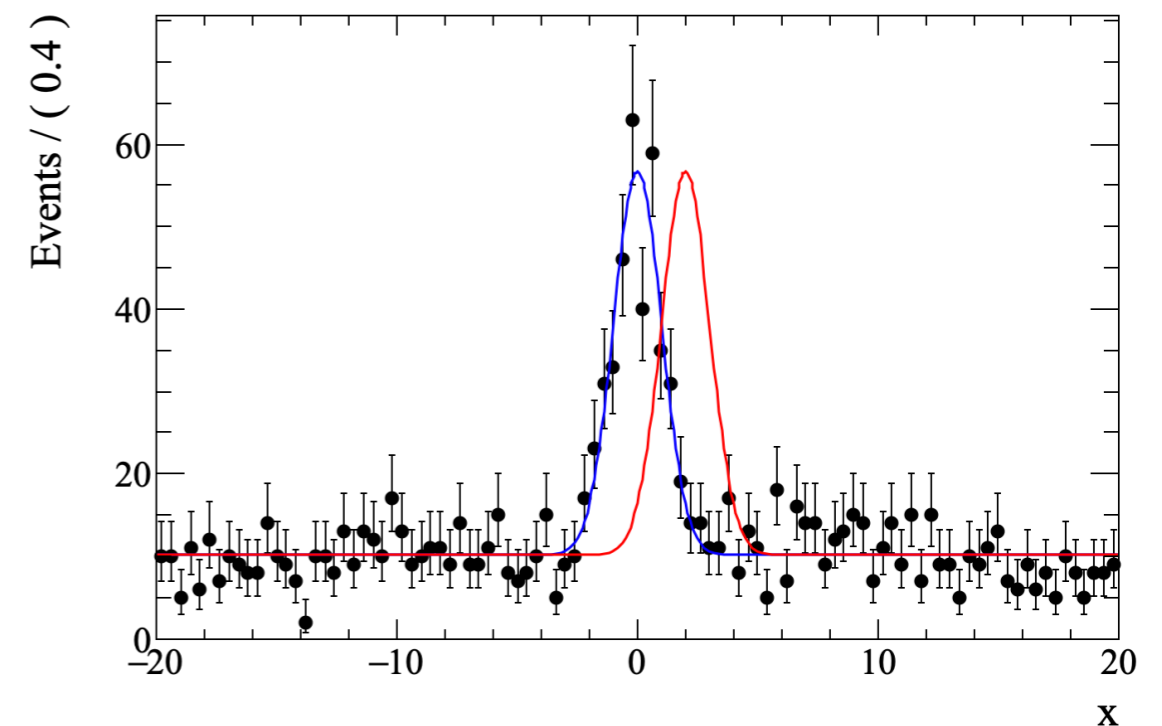
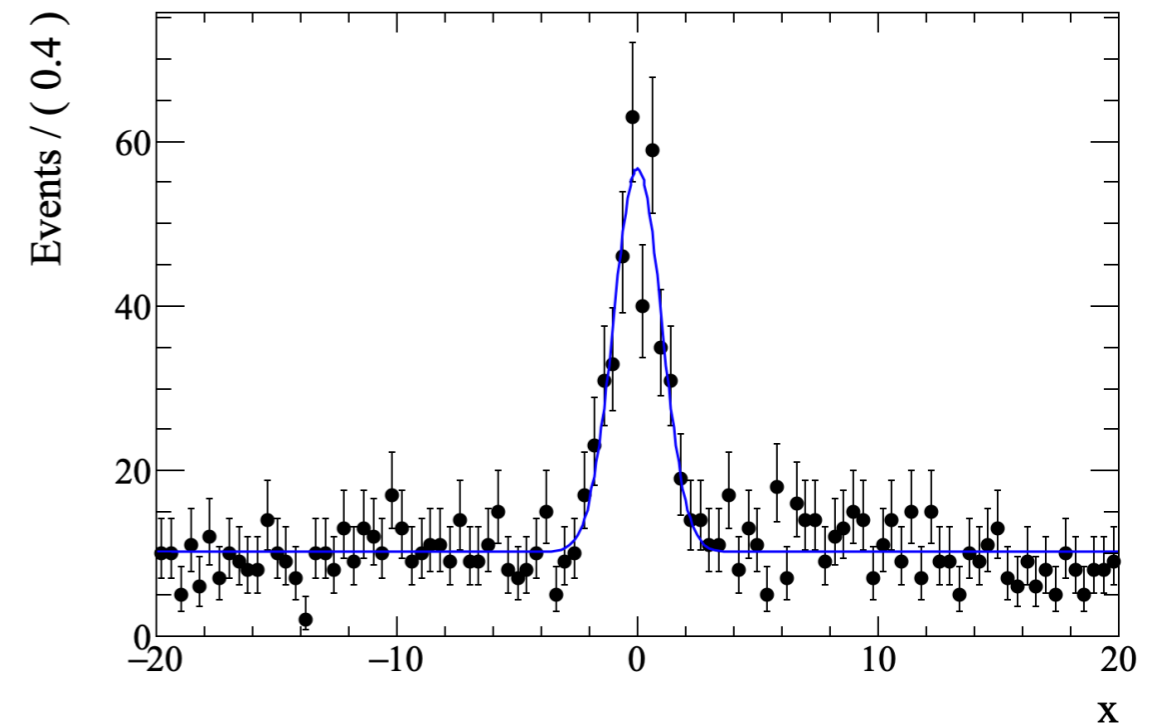
Rare events with high-precision

- Let's suppose to search for a rare process, looking for a peak over the background in a known position of a distribution
 - high sensitivity through high precision



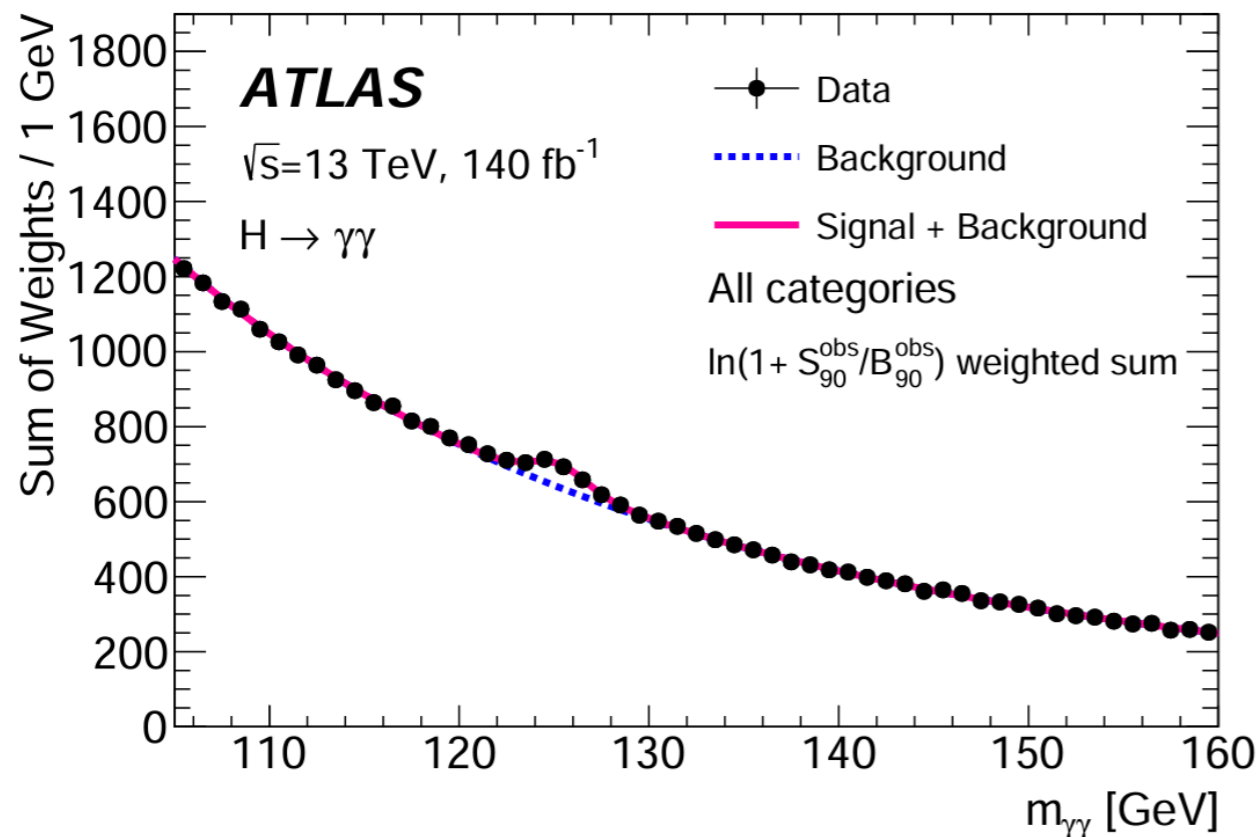
Rare events with high-precision

- Let's suppose to search for a rare process, looking for a peak over the background in a known position of a distribution
 - high sensitivity through high precision
 - the measurement has to be also highly accurate (no relevant bias)

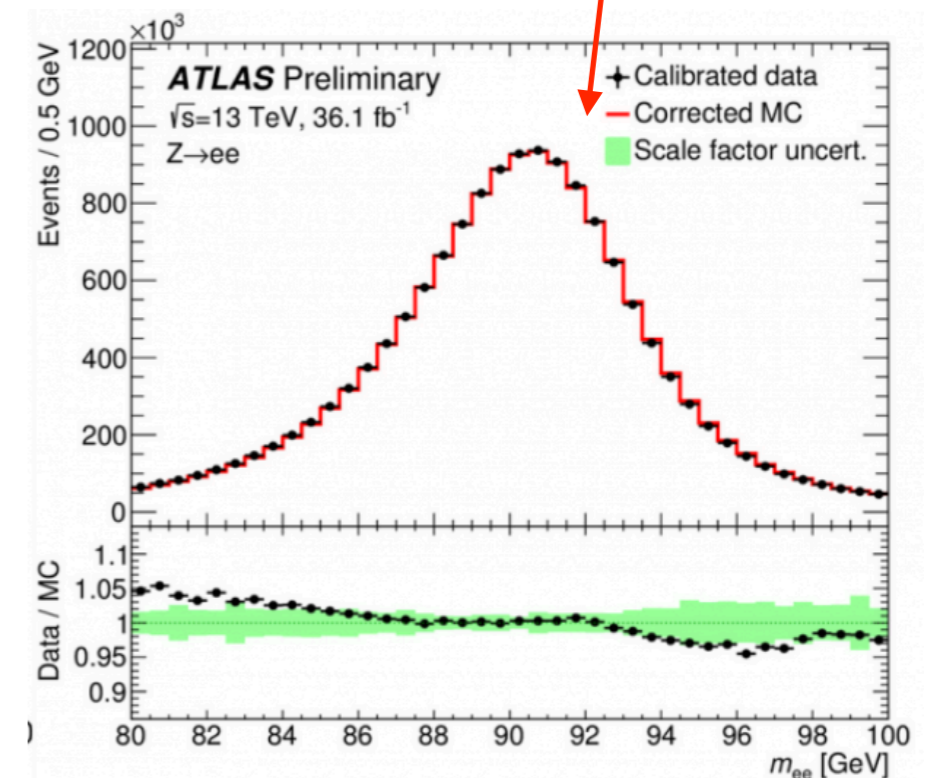


Rare events with high-precision

- In general-purpose experiments, many non-rare physics processes can be used to calibrate the measurements and remove biases



Source	Impact [MeV]
Photon energy scale	83
$Z \rightarrow e^+e^-$ calibration	59
E_T -dependent electron energy scale	44
$e^\pm \rightarrow \gamma$ extrapolation	30
Conversion modelling	24
Signal-background interference	26
Resolution	15
Background model	14
Selection of the diphoton production vertex	5
Signal model	1
Total	90

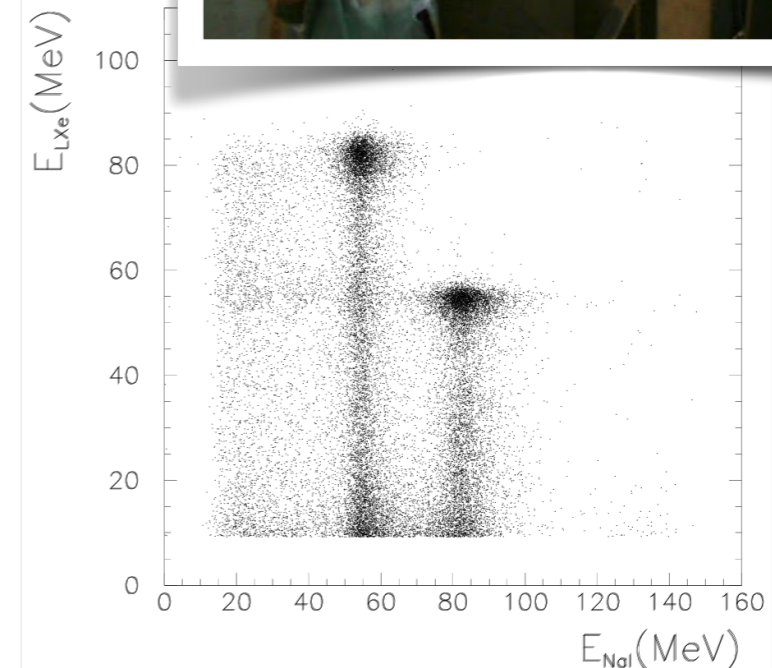
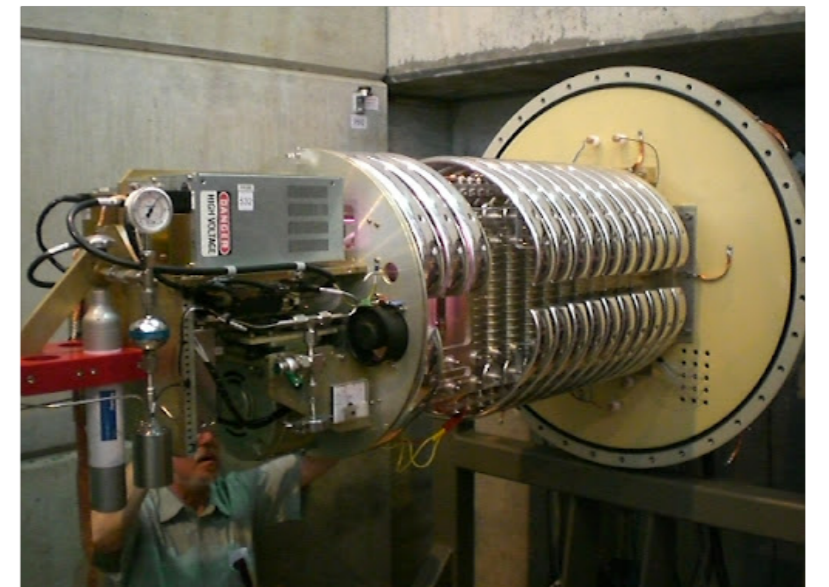


Rare events with high-precision

- Single-purpose experiments typically have extreme resolution (i.e. need for extreme accuracy) and a scarcity of physics processes to be used for calibrations
 - **dedicated tools need to be developed**

Rare events with high-precision - the MEG case

- Single-purpose experiments typically have extreme resolution (i.e. need for extreme accuracy) and a scarcity of physics processes to be used for calibrations
 - **dedicated tools need to be developed**
- In the MEG experiment, where a muon beam is used to search for $\mu \rightarrow e\gamma$, a profusion of calibration tools has been developed, including a **dedicated Cockcroft-Walton accelerator** and a $\pi^-(p, n)\pi^0$ **experiment** with the only purpose of calibrating the photon energy reconstruction.



Systematic Uncertainties for High Accuracy

- **High Accuracy for zero tests**

e.g. particles EDMs

$$\Delta \stackrel{?}{=} 0$$

$$\Delta_{meas} = k\Delta + \delta \quad \text{with } k = 1 \pm \sigma_k, \delta = 0 \pm \sigma_\delta$$

$$\Rightarrow \sigma_{\Delta}^{syst} = \frac{\sigma_k}{k}\Delta + \sigma_\delta \sim \sigma_\delta$$

- Additive uncertainties dominate over multiplicative uncertainties

Systematic Uncertainties for High Accuracy

- **High Accuracy for non-zero measurements**

e.g. particles MDMs,
coupling constants

- multiplicative uncertainties are also important
- comparison with SI units is critical —> **a metrology problem**

$$\frac{g_\mu - 2}{2} \sim \frac{m\omega}{qB}$$

precession frequency

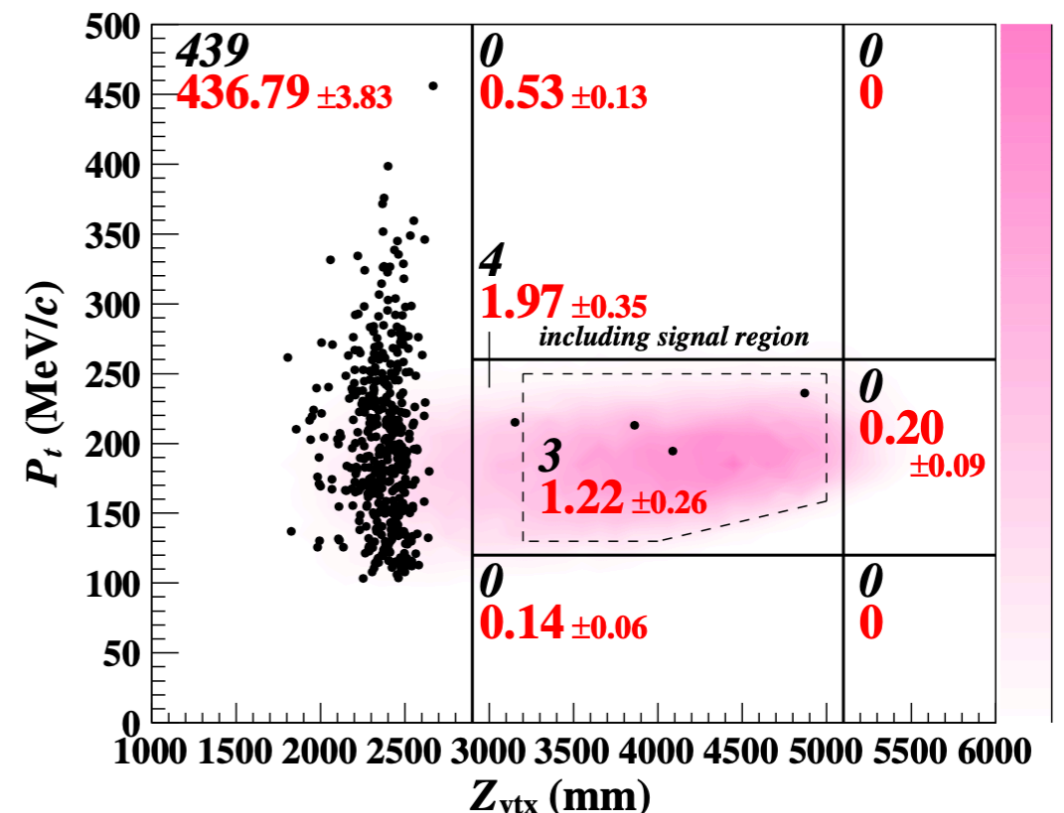
magnetic field

Measured dimensional quantities need to be calibrated against SI standards with << ppm accuracy

Systematic Uncertainties in Rare Event Searches

- **Rare event searches through event patterns**

- rare event searches (rare decays, dark matter, etc.) where background rejection is mostly achieved through **particle identification, vetos, event topology**, etc.
- dominant systematic uncertainties typically from the control of the background rejection efficiency

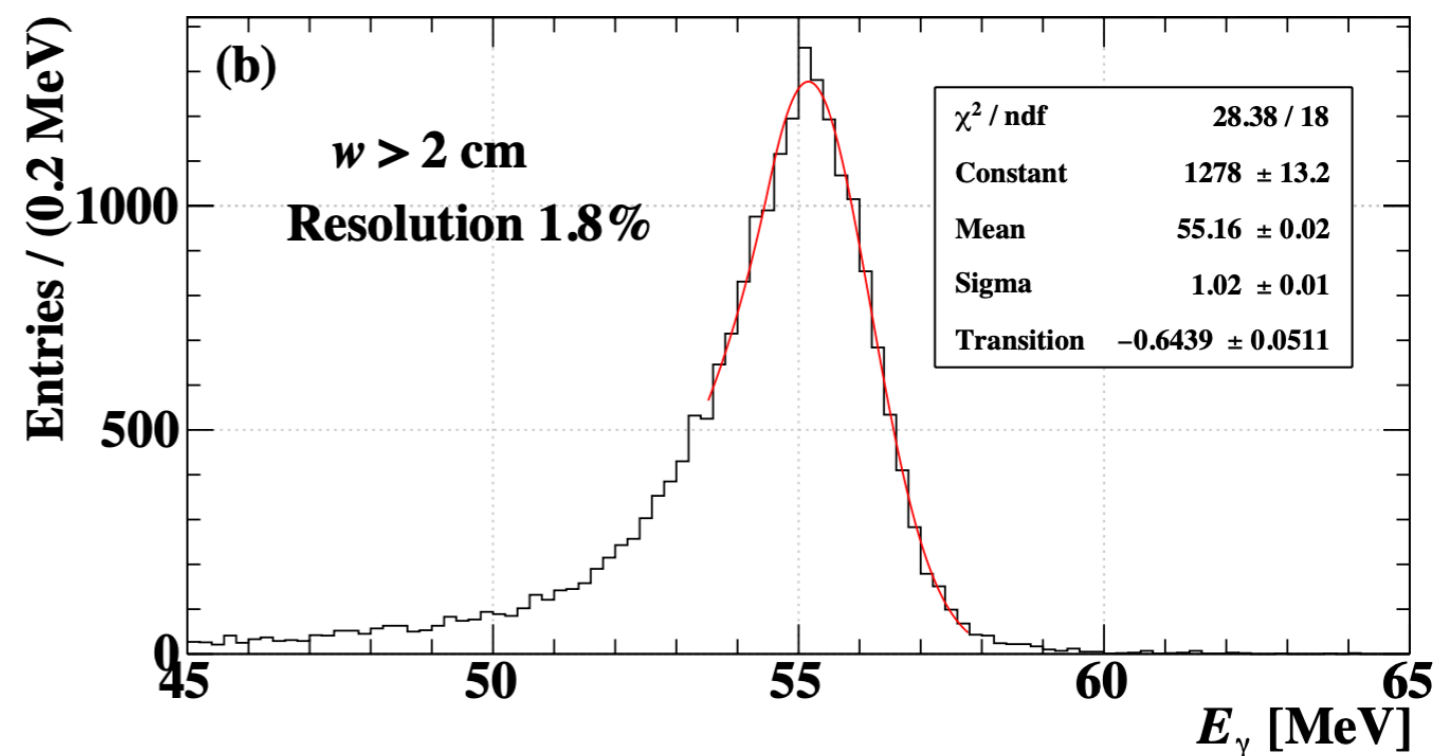


Systematic Uncertainties in Rare Event Searches

- **Rare event searches through precision**

- rare event searches (e.g. rare decays) where the precise measurement of some observable (e.g. kinematics) is required to discriminate signal and background
- high precision requires also high accuracy, which can be only achieved with dedicated tools

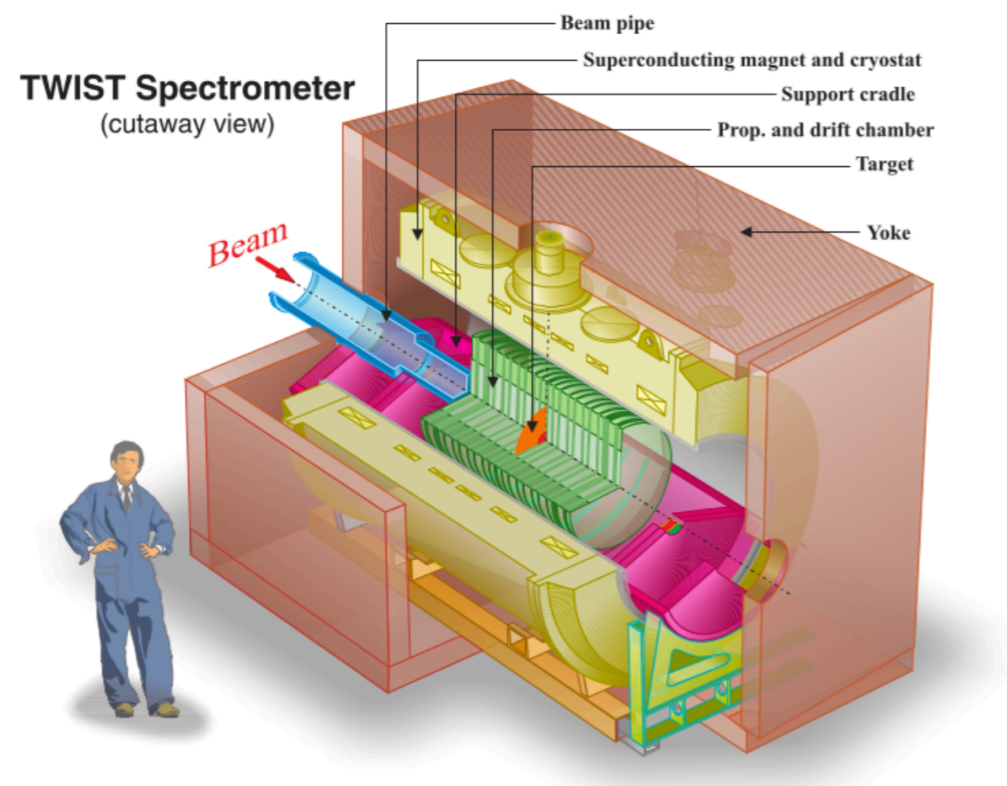
Photon energy calibration in MEG with a pion beam, $\pi^- + p \rightarrow \pi^0(\gamma\gamma) + n$



A ognuno il suo (*Each to their own*)

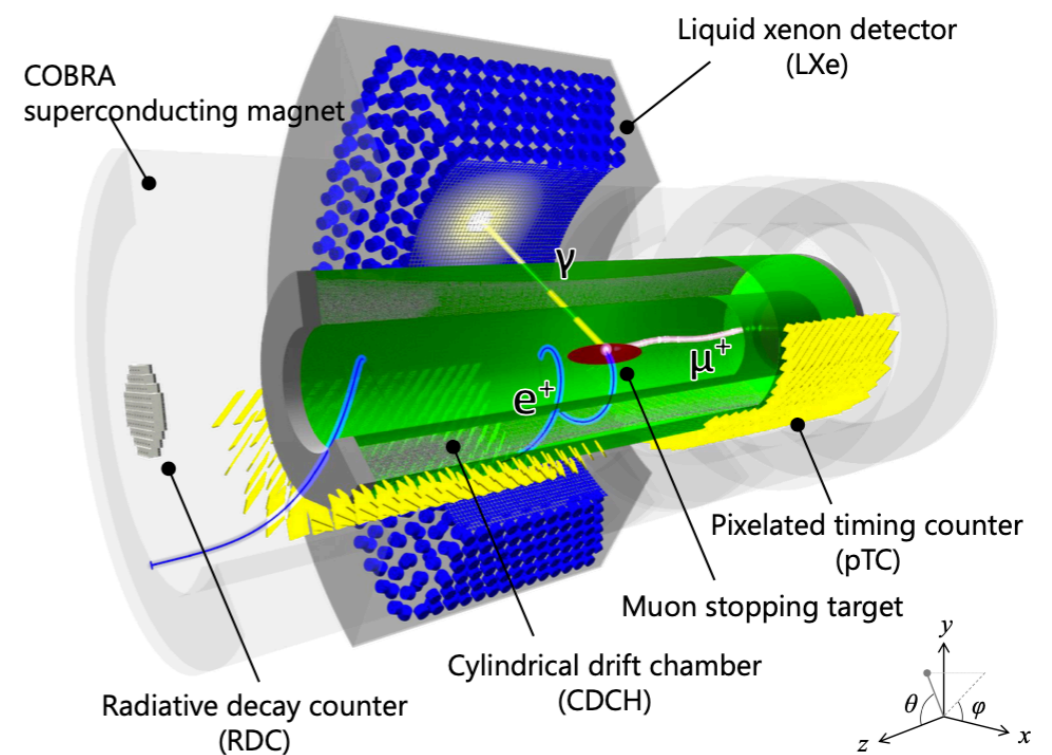
- MEG vs. TWIST

Measurement of parameters of the e^+ energy spectrum in μ^+ decays



Very accurate knowledge of the magnetic field is necessary to measure the spectrum parameters

Measurement of e^+ energy spectrum for the search of $\mu^+ \rightarrow e^+ \gamma$



Magnetic field adjusted from data, exploiting the theoretical knowledge of the Michel spectrum

The role of Monte Carlo simulations

- The extremely high resolutions and accuracies of single-purpose experiments pose strong challenges to Monte Carlo simulations
- Additionally, for rare event searches, MC productions resulting in sufficient statistics of reconstructed background events are computationally unachievable or unreliable
 - e.g. MEG II 2021 data

Potential background events	$\sim 10^{12}$
After the trigger	2×10^7
In the analysis region	66
Within 1σ to the signal	< 1

The role of Monte Carlo simulations

- The extremely high resolutions and accuracies of single-purpose experiments are often not achievable in Monte Carlo simulations
- Additionally, the amount of data available for training in Monte Carlo simulations is typically very limited and, when MC inputs are unavoidable, the related systematic uncertainties are large
 - e.g.

Potential background events	$\sim 10^{12}$
After the trigger	2×10^7
In the analysis region	66
Within 1σ to the signal	< 1

Confidence intervals and systematics

The Feldman-Cousins approach

- In FC, the confidence belt is built using the likelihood ratio test statistics, defined as:

$$\mathcal{R}(\mathbf{p}) = \frac{\mathcal{L}(\mathbf{p})}{\mathcal{L}(\hat{\mathbf{p}})}$$

\mathbf{p} : set of parameters
 $\hat{\mathbf{p}}$: set of parameters maximizing the likelihood
(i.e. fitted value)

- Given a set of values of the parameters \mathbf{p} , the expected distribution of \mathcal{R} is computed
- When the experiment is performed, the value of \mathcal{R} is computed for different hypothetical sets of \mathbf{p} :
 - a set of \mathbf{p} is included in a confidence interval at C.L. = $1 - \alpha$ if more than a fraction α of the experiments is expected to give a larger $\mathcal{R}(\mathbf{p})$ than data

Dealing with physical limits on a parameter

- For rare event searches ($p = N_{\text{sig}}$), the so-called “conditioning” is also included to properly treat the physical constraint $N_{\text{sig}} > 0$:

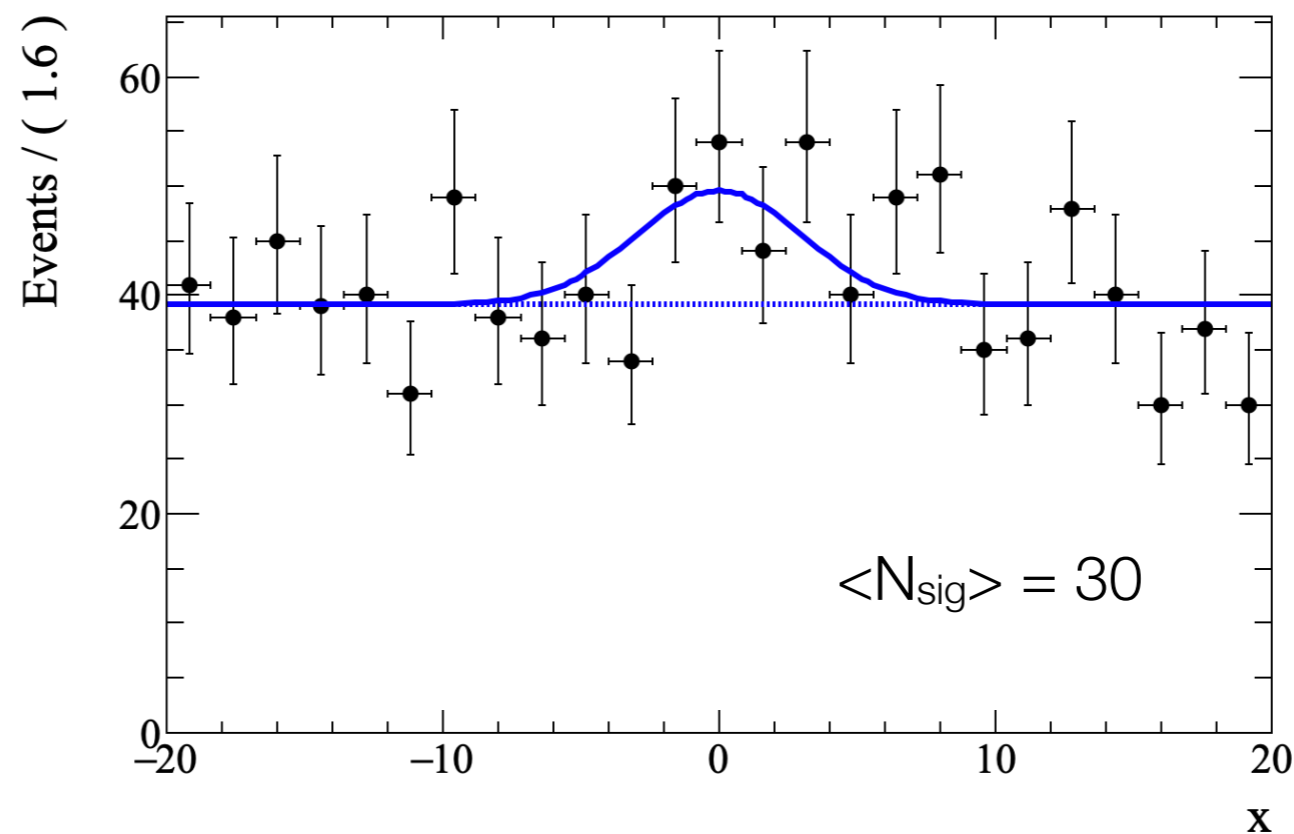
$$\mathcal{R} = \begin{cases} \frac{\mathcal{L}(N_{\text{sig}})}{\mathcal{L}(N_{\text{sig,best}})} & \text{if } N_{\text{sib,best}} \geq 0 \\ \frac{\mathcal{L}(N_{\text{sig}})}{\mathcal{L}(0)} & \text{if } N_{\text{sib,best}} < 0 \end{cases}$$

The Feldman-Cousins approach

- The FC approach requires to evaluate the expected distribution of $\mathcal{R}(\mathbf{p})$
- Typically obtained by generating pseudo-experiments (toy Monte Carlo exp.) according to the expected PDFs

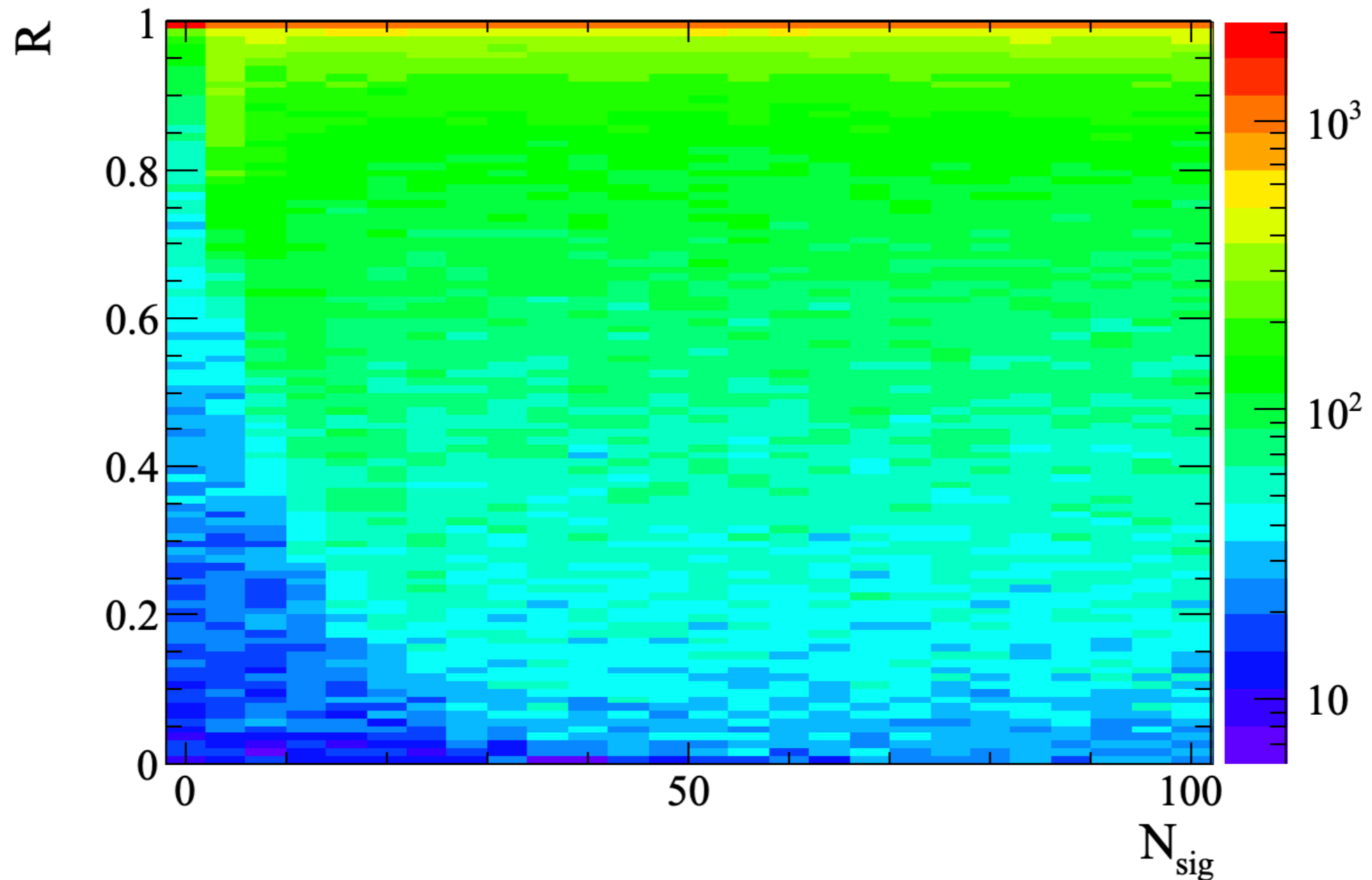
The Feldman-Cousins approach

- Let's consider the case of:
 - a single variable x
 - a gaussian signal ($\mu = 0, \sigma = 3$)
 - a flat background of 1000 events in $x \in [-20,20]$

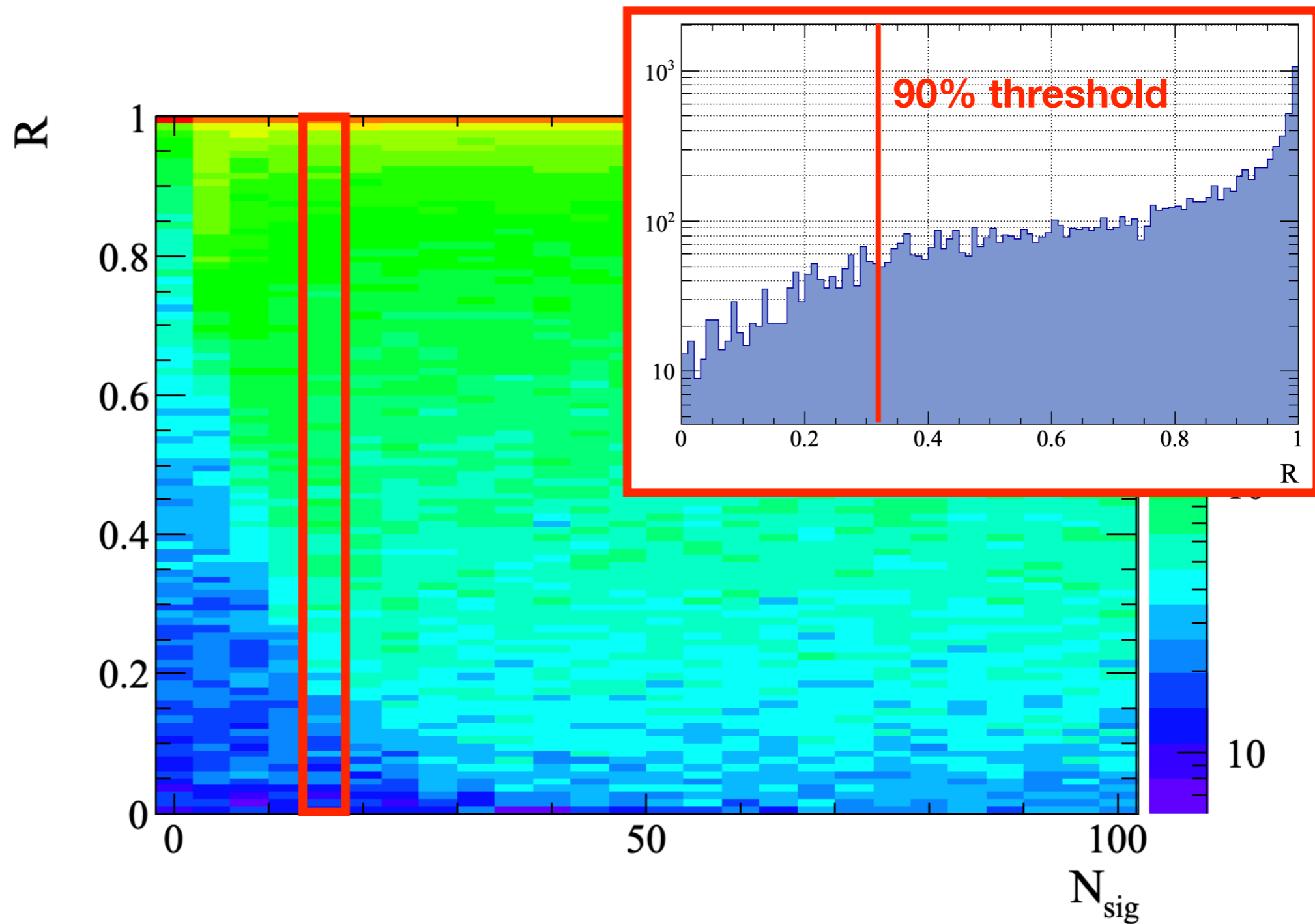


The Feldman-Cousins approach

Distribution of R from toy MCs generated with different $\langle N_{\text{sig}} \rangle$

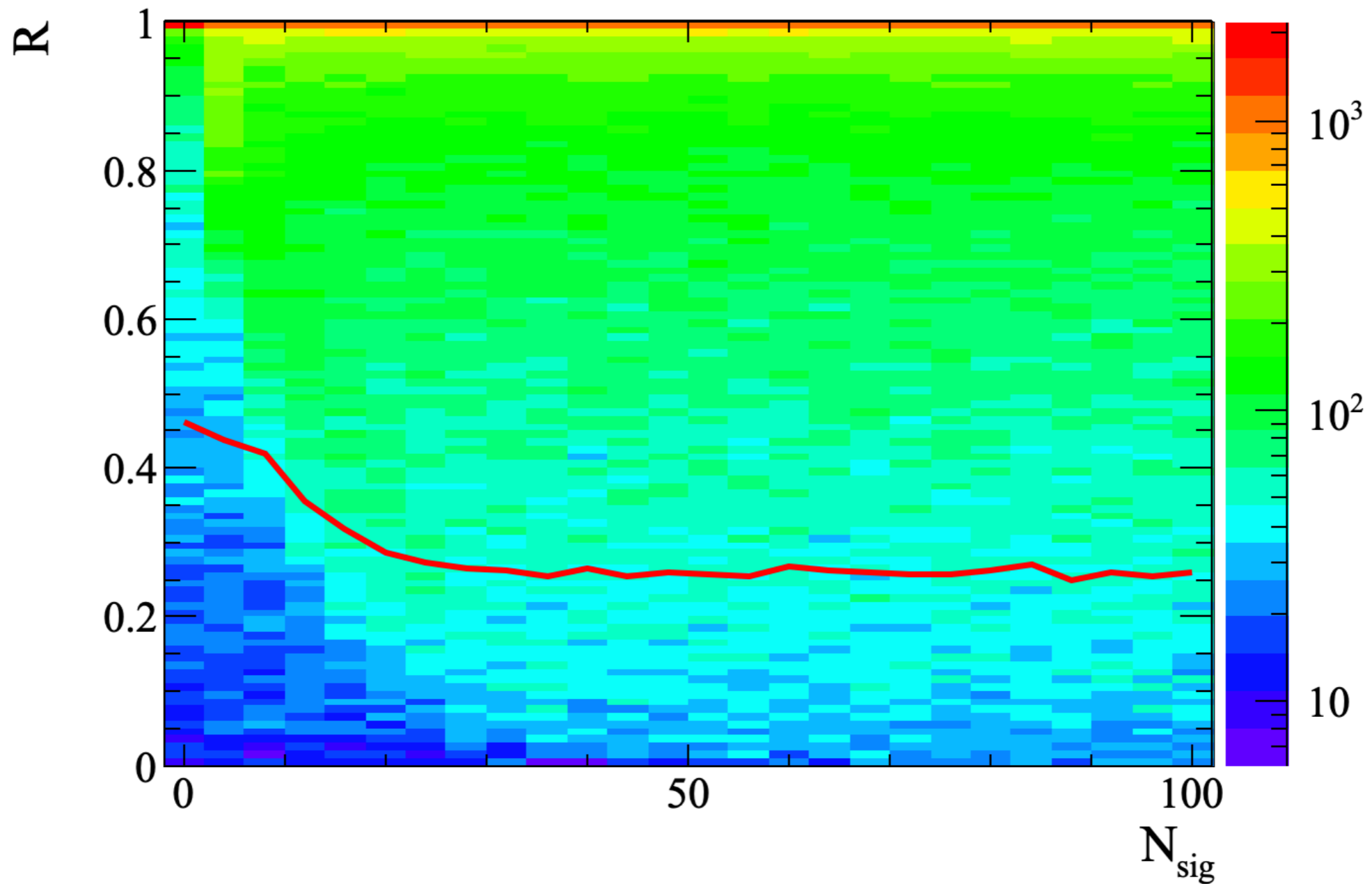


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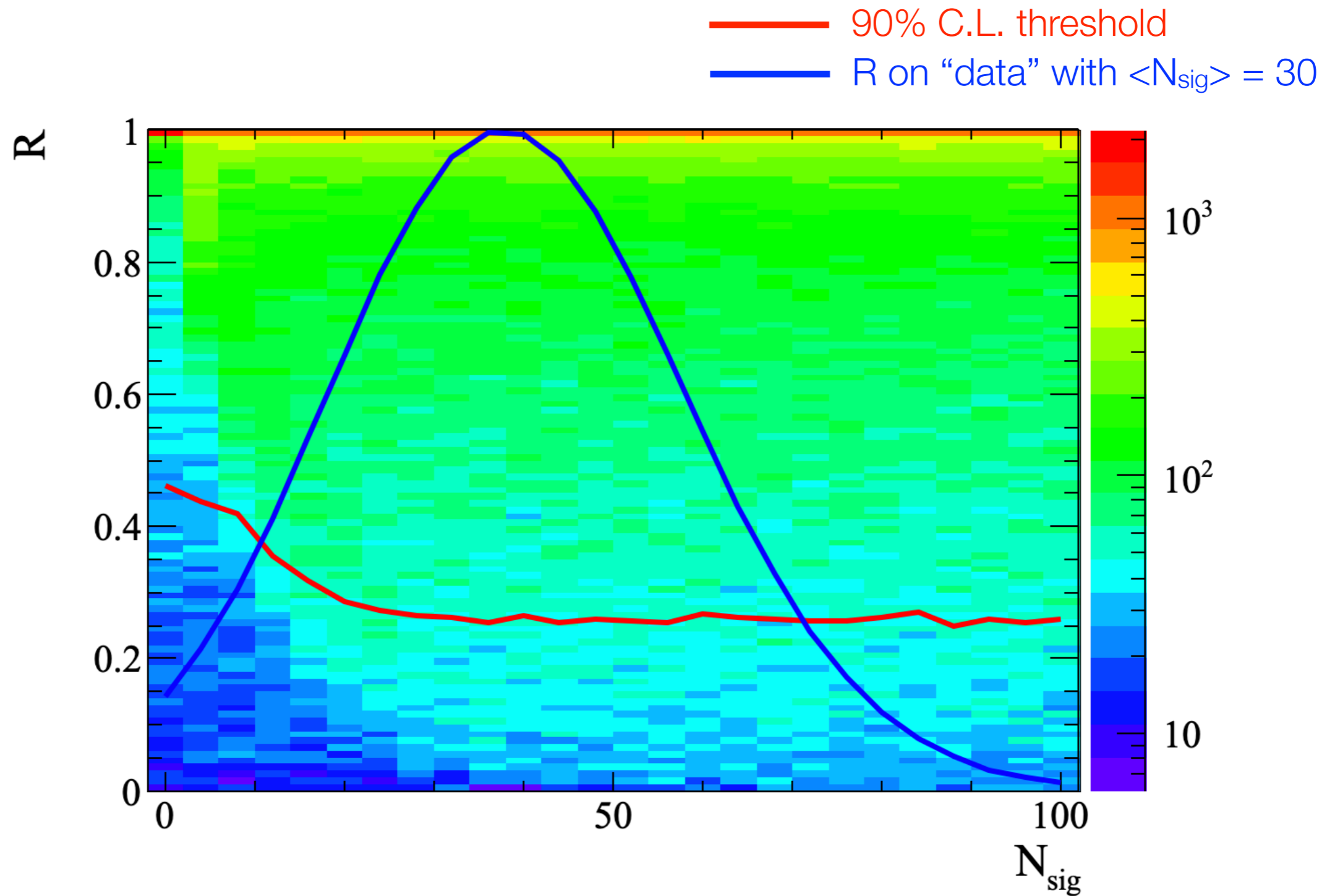


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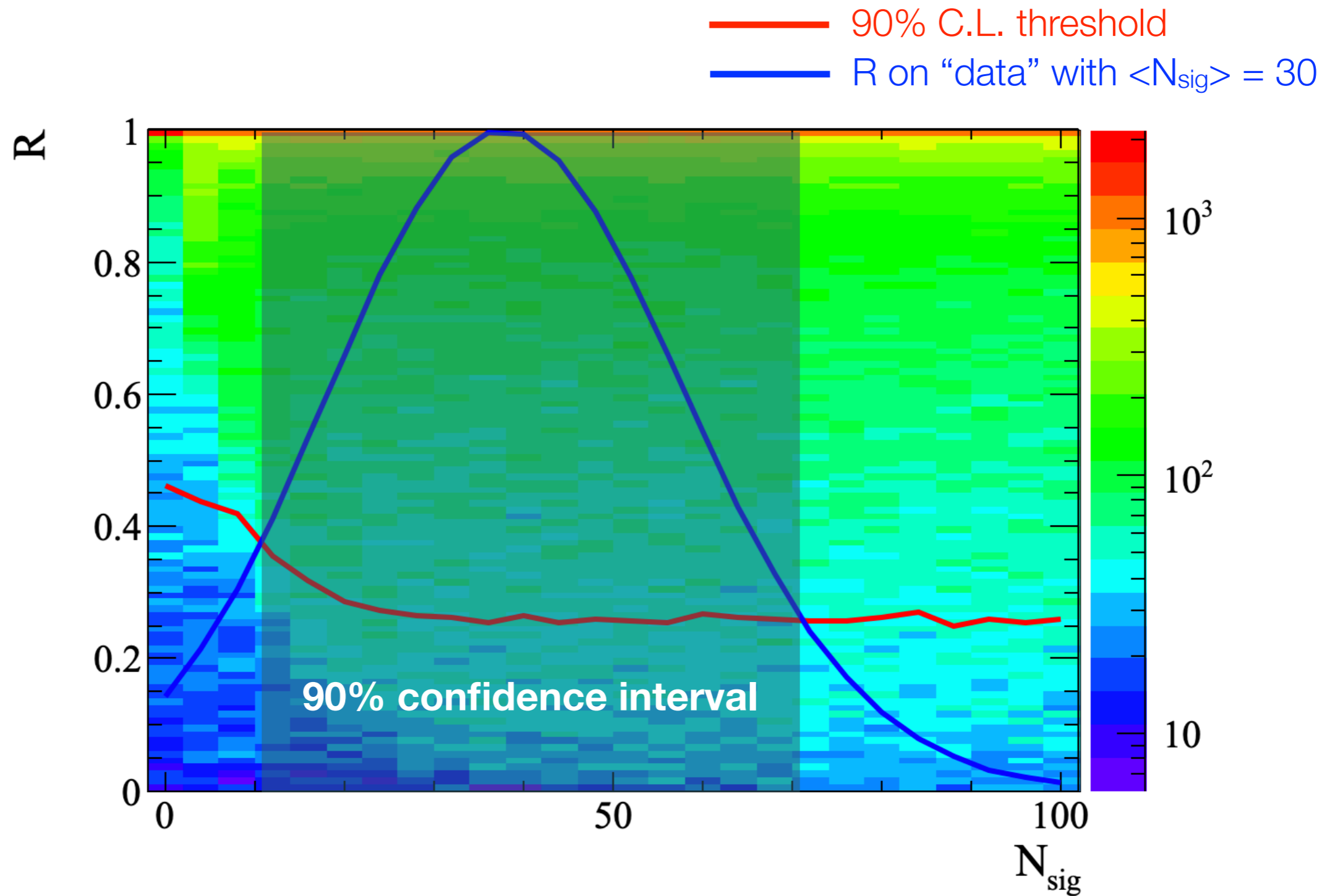
— 90% C.L. threshold



The Feldman-Cousins approach



The Feldman-Cousins approach

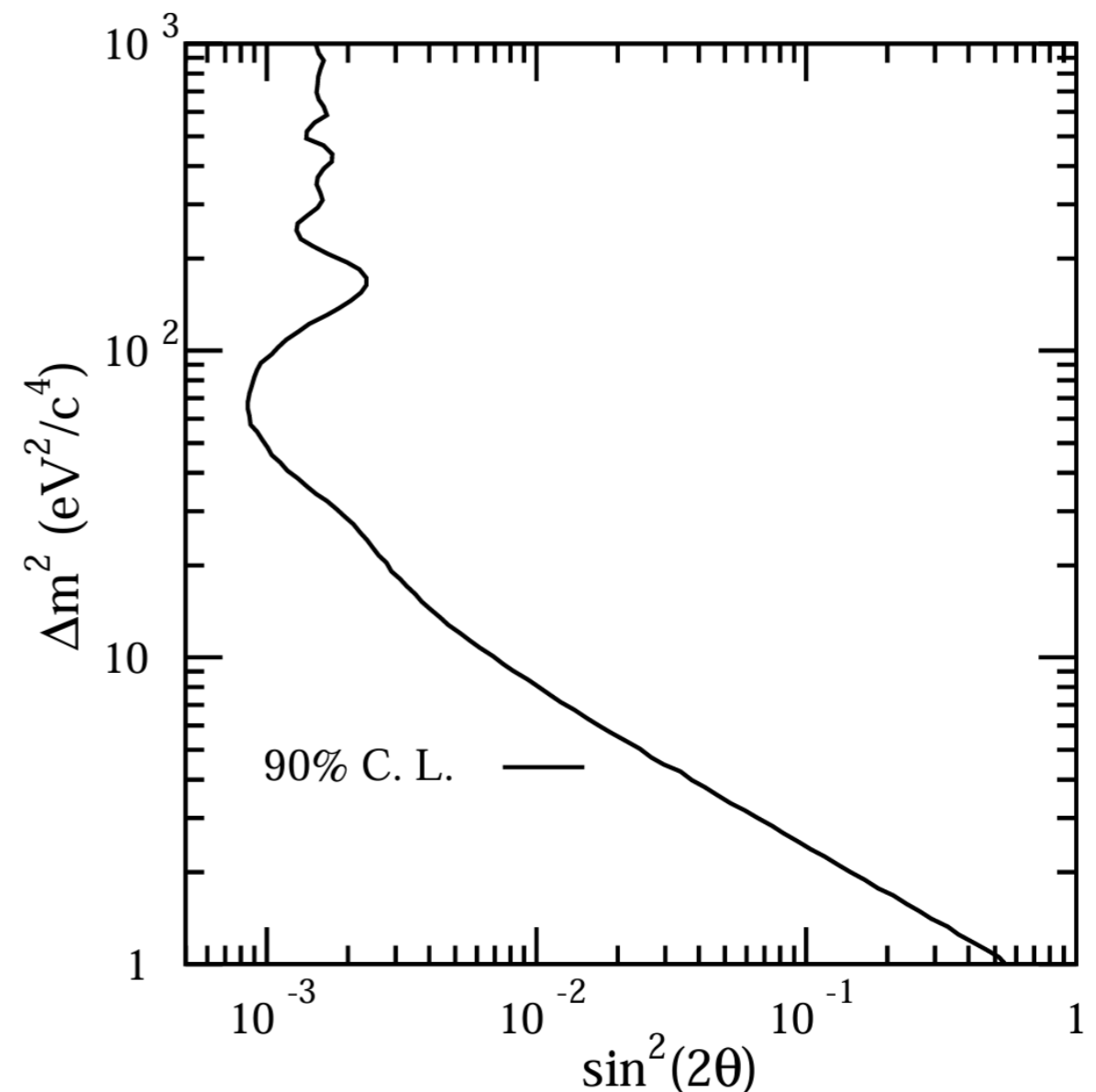


Multi-dimensional case

- The original FC problem (neutrino oscillation parameters)
- Counting experiment:
 - the observable is the number of countings n_{obs} (Poisson PDF)
 - there is an expected number of background events, n_{bkg}
 - for each point in the 2D space, there is an expected number of signal events, n_{sig}
 - for each point, the R distribution is derived from toy MCs with Poisson distribution
 - R from data is compared to the R distribution in toy MCs

$$\mathcal{R}(\Delta m^2, \sin^2(2\theta)) = \frac{\text{Poisson}(n_{\text{obs}}; n_{\text{sig}} + n_{\text{bkg}})}{\text{Poisson}(n_{\text{obs}}; n_{\text{sig,best}} + n_{\text{bkg}})}$$

with $n_{\text{sig,best}} > 0$

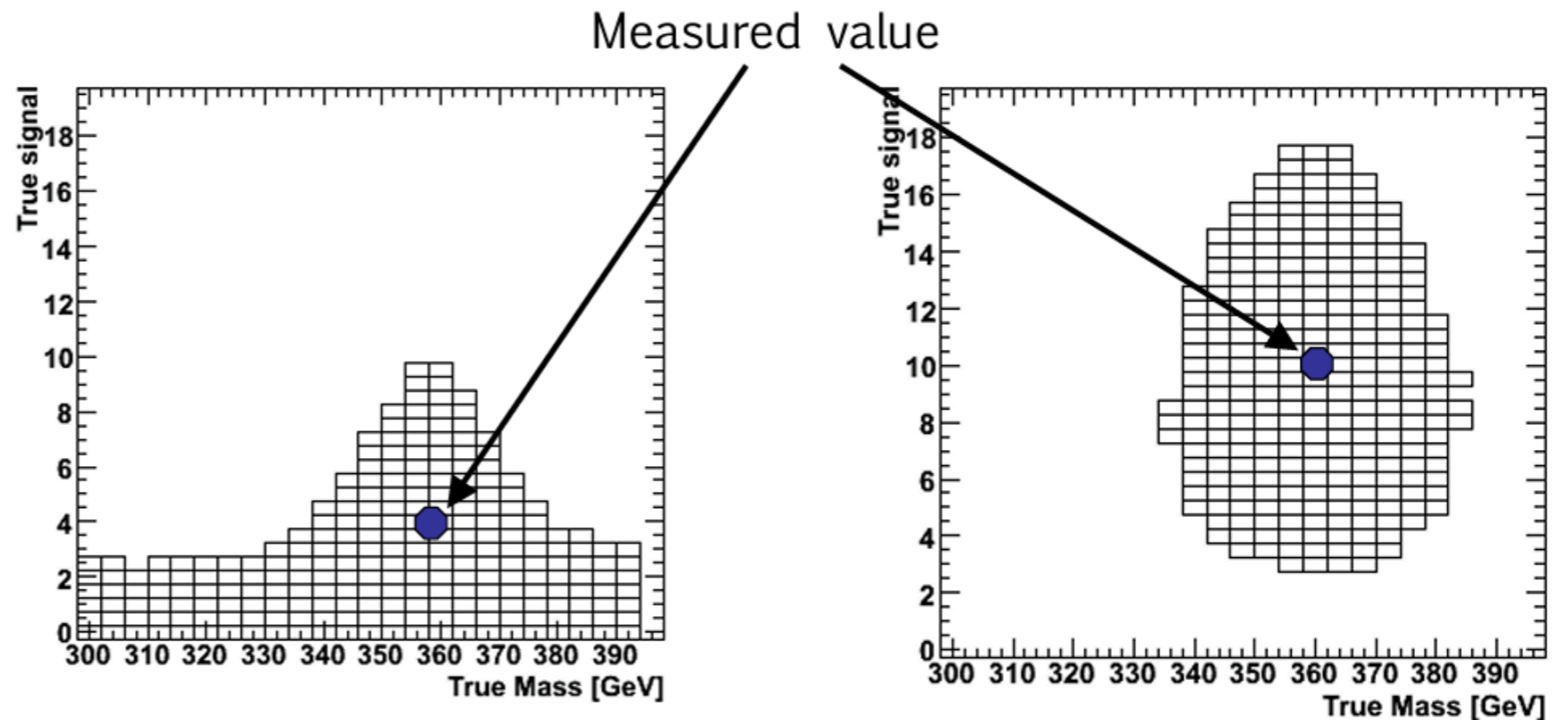


Multi-dimensional case

$$\mathcal{R} = \frac{\mathcal{L}(N_{\text{sig}})}{\mathcal{L}(\hat{N}_{\text{sig}})}$$



$$\mathcal{R} = \frac{\mathcal{L}(N_{\text{sig}}, M)}{\mathcal{L}(\hat{N}_{\text{sig}}, \hat{M})}$$



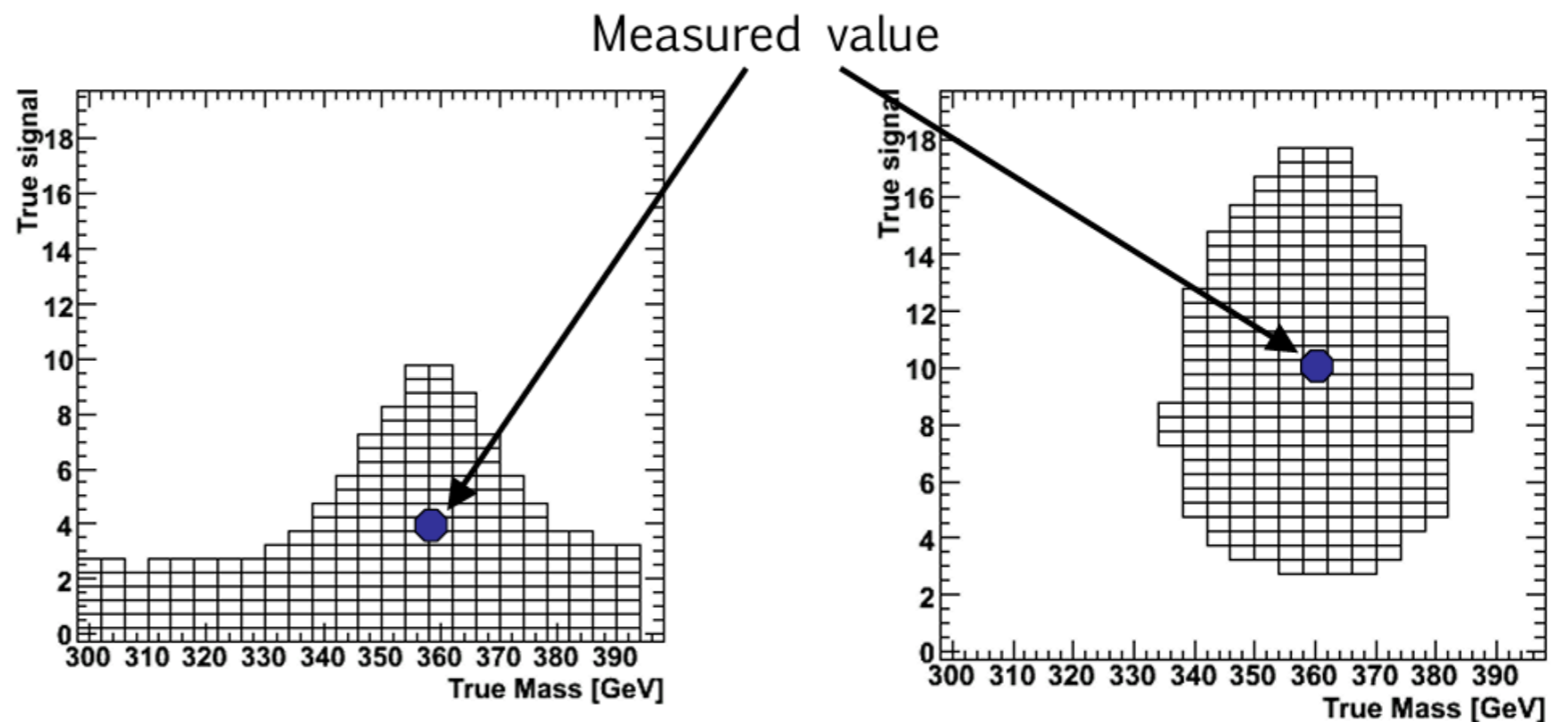
- Coverage guaranteed for each pair of true (N_{sig}, M)
 - no need of look-elsewhere corrections

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The Feldman-Cousins approach

- The FC approach requires to evaluate the expected distribution of $\mathcal{R}(\mathbf{p})$
- Typically obtained by generating pseudo-experiments (toy Monte Carlo exp.) according to the expected PDFs
- CAVEAT: can be computationally heavy for multi-dimensional parameter space, complex likelihoods and very small p-values:
 - for a 5σ test, need to generate $\sim 10^9$ toy MC experiments

Inclusion of systematics

- For the inclusion of systematics, the most popular approaches are:
 - semi-bayesian approach (Highland-Cousins): the likelihood is integrated over the nuisance parameters before applying the desired statistical approach

e.g. likelihood for poisson-distributed yields, integrated over a gaussian uncertainty on the expected background b

$$q(n)_{s+b} = \frac{1}{\sqrt{2\pi\sigma_b}} \int_0^\infty p(n)_{s+b'} e^{-(b-b')^2/2\sigma_b^2} db'$$

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e.g. likelihood for poisson-distributed yields, integrated over gaussian uncertainties on the expected background b and signal efficiency ϵ

$$q(n)_{s+b} = \frac{1}{2\pi\sigma_b\sigma_\epsilon} \int_0^\infty \int_0^\infty p(n)_{b'+\epsilon's} \times e^{-(b-b')^2/2\sigma_b^2} e^{-(1-\epsilon')^2/2\sigma_\epsilon^2} db' d\epsilon'$$

Inclusion of systematics

- For the inclusion of systematics, the most popular approaches are:
 - profile likelihood ratio: the likelihood is maximized with respect to the nuisance parameters when building the likelihood ratio:

$$\mathcal{L}(\mathbf{p}, \mathbf{q}) = P(\text{data} | \mathbf{p}, \mathbf{q})P(\mathbf{q})$$

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External constraint: PDF (e.g. gaussian)
representing the uncertainty on the
nuisance parameters

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Likelihood maximized over \mathbf{q} for fixed \mathbf{p}

$$\mathcal{R}(\mathbf{p}) = \frac{\mathcal{L}(\mathbf{p}, \hat{\mathbf{q}})}{\mathcal{L}(\hat{\mathbf{p}}, \hat{\mathbf{q}})}$$

Likelihood maximized over \mathbf{p} and \mathbf{q}

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$$\mathcal{R} = \frac{\text{Poisson}(n_{\text{obs}}; n_{\text{sig}} + n_{\text{bkg}})P(n_{\text{bkg}})}{\text{Poisson}(n_{\text{obs}}; n_{\text{sig,best}} + n_{\text{bkg,best}})P(n_{\text{bkg,best}})}$$

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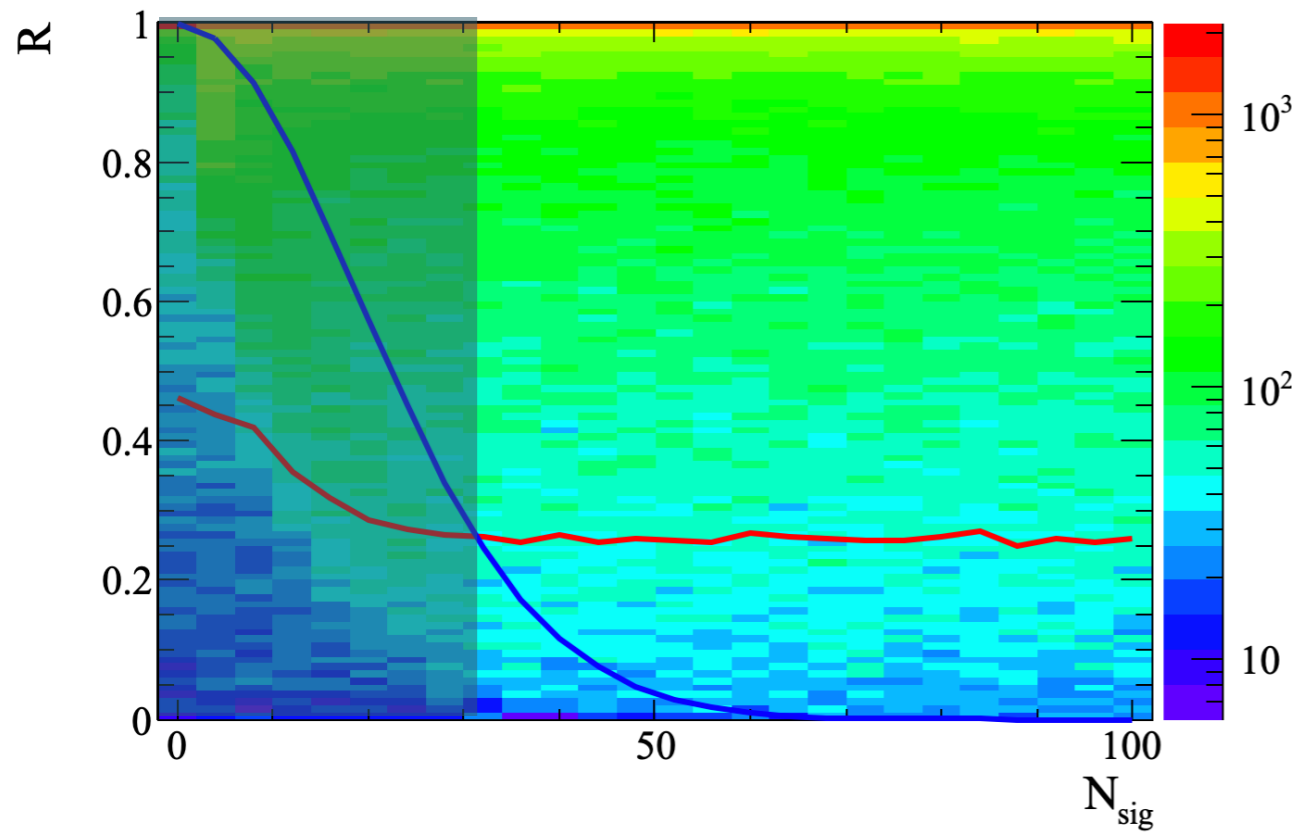
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Set of nuisance parameters which maximizes the likelihood for a given N_{sig}

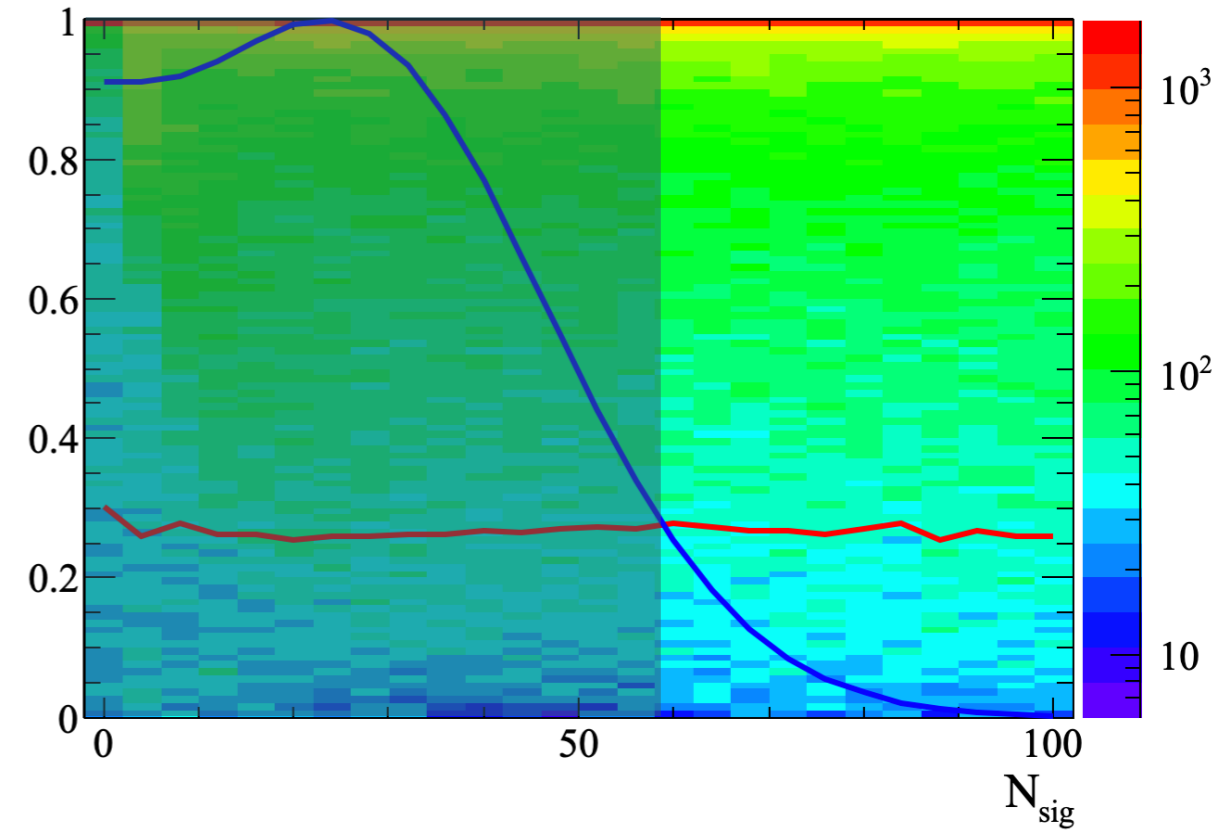
Inclusion of systematics

— 90% C.L. threshold
— R on “data” with $\langle N_{\text{sig}} \rangle = 0$

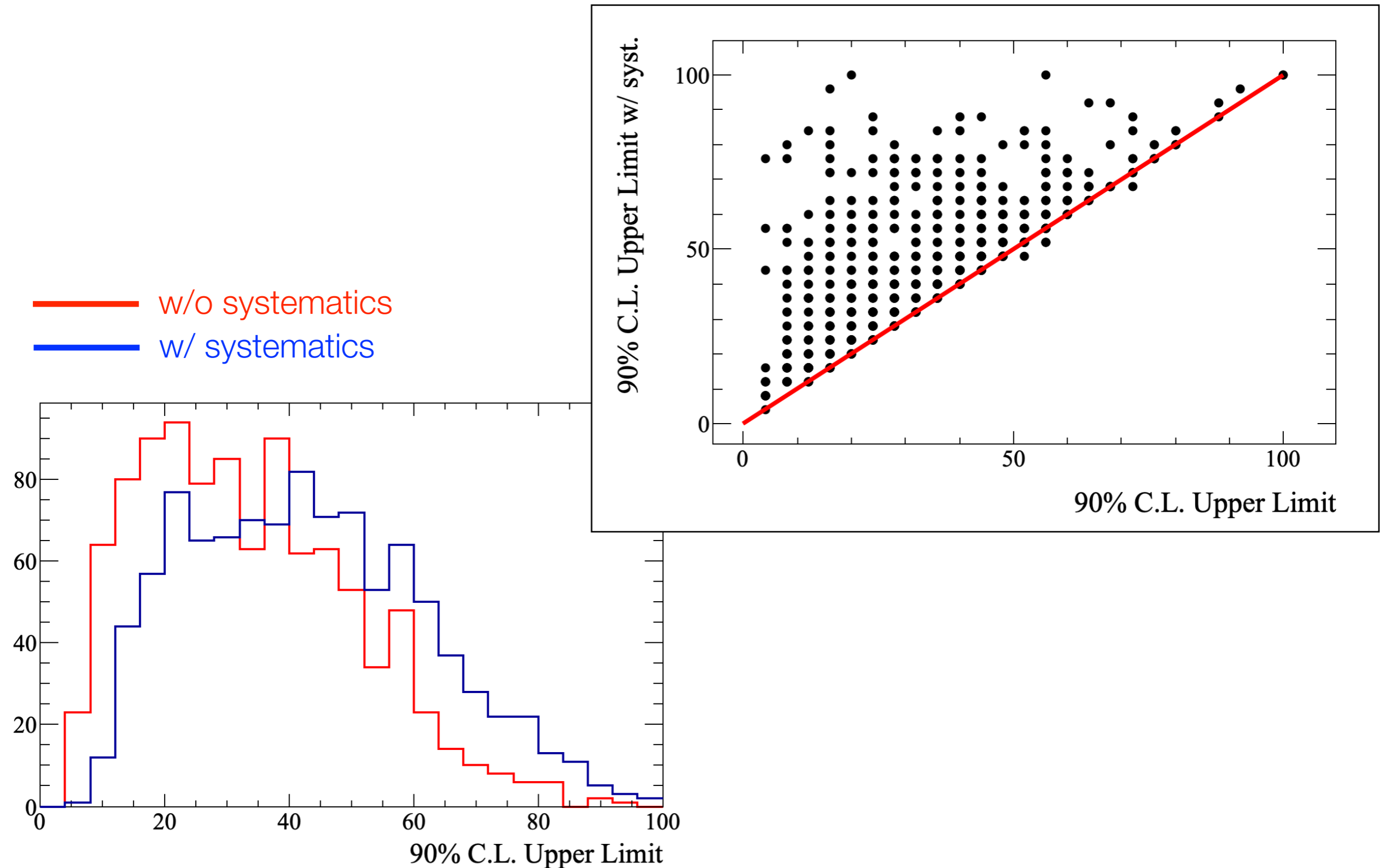
w/o systematics



w/ ± 3 uncertainty on μ



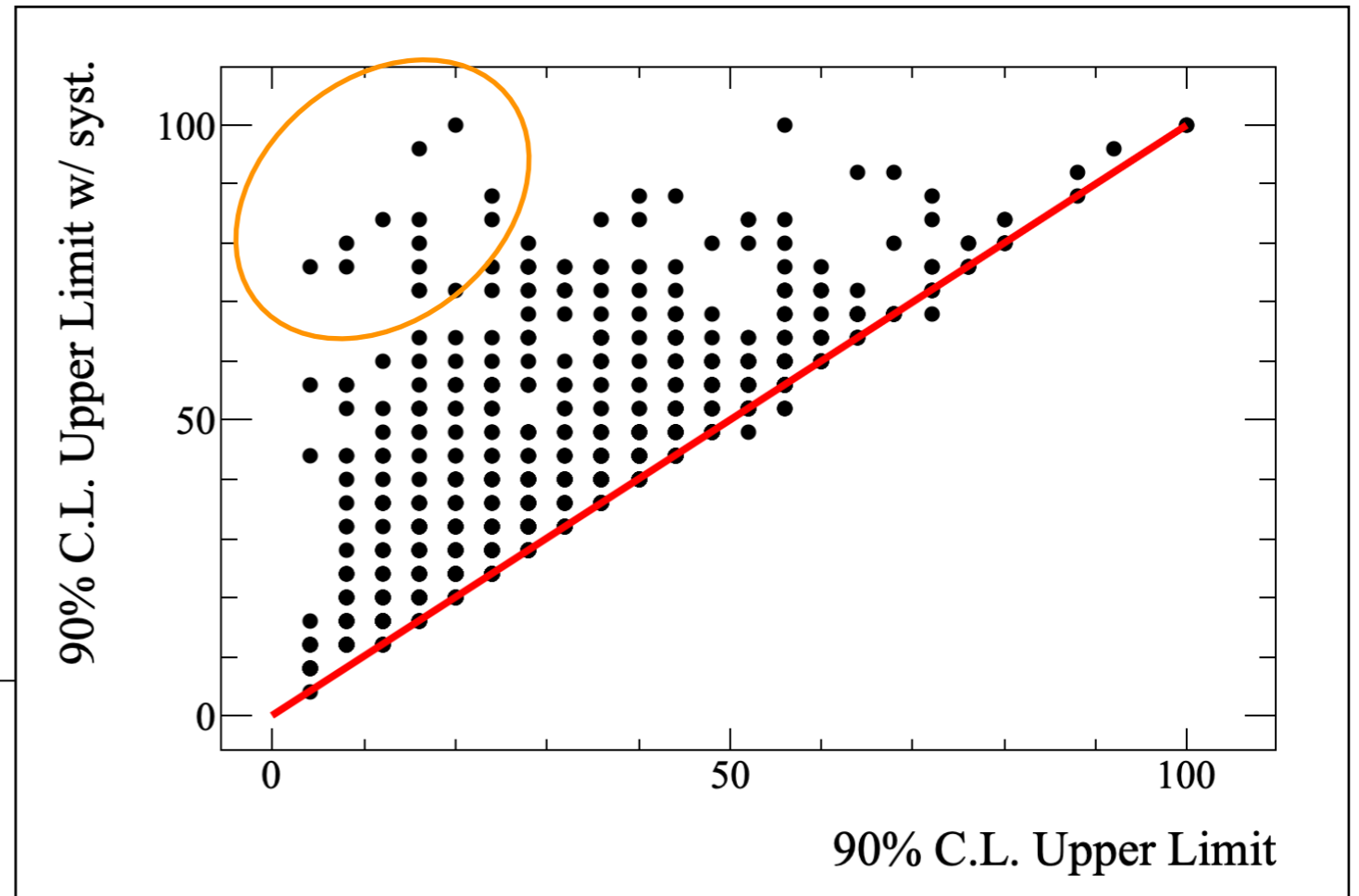
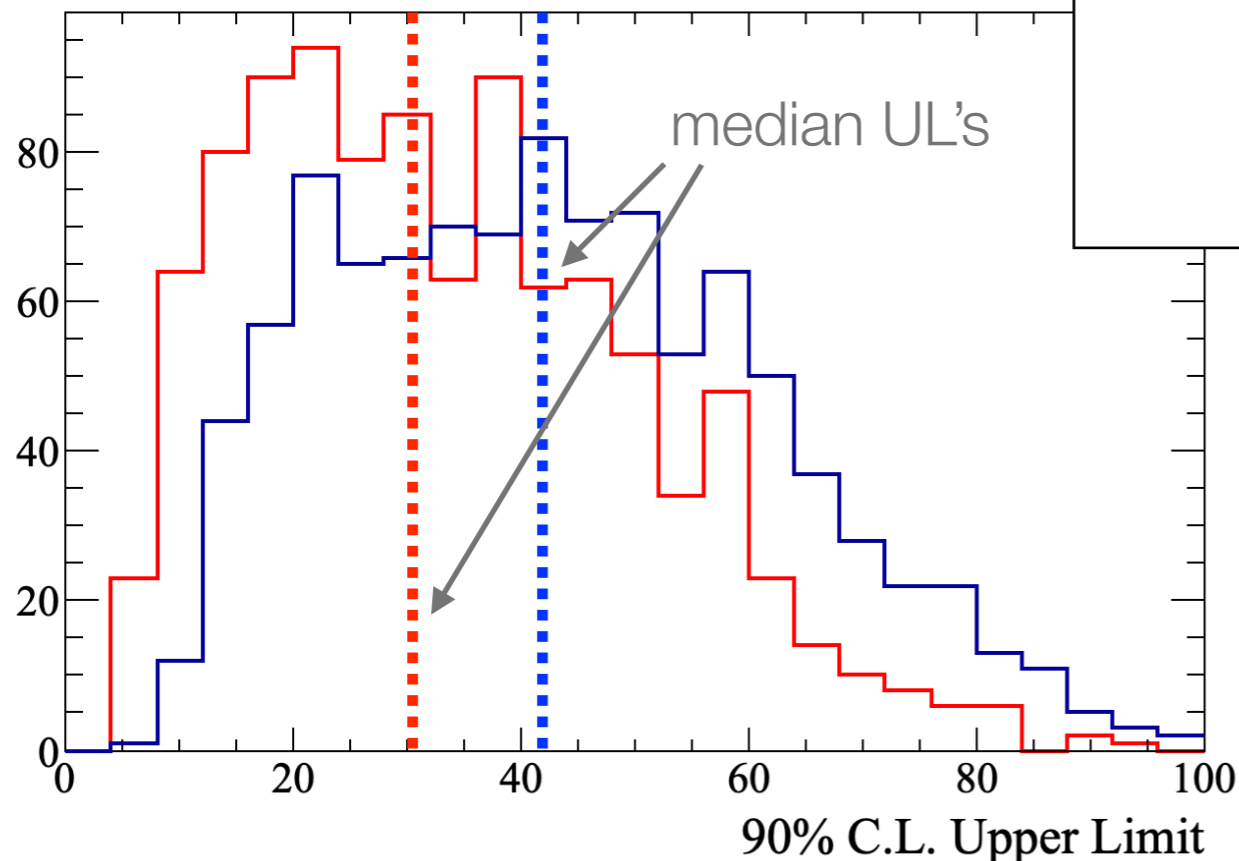
Inclusion of systematics



Inclusion of systematics

With an uncertainty in the mean equal to the gaussian width, the impact in some experiments can be huge (up to a factor 10!!!)

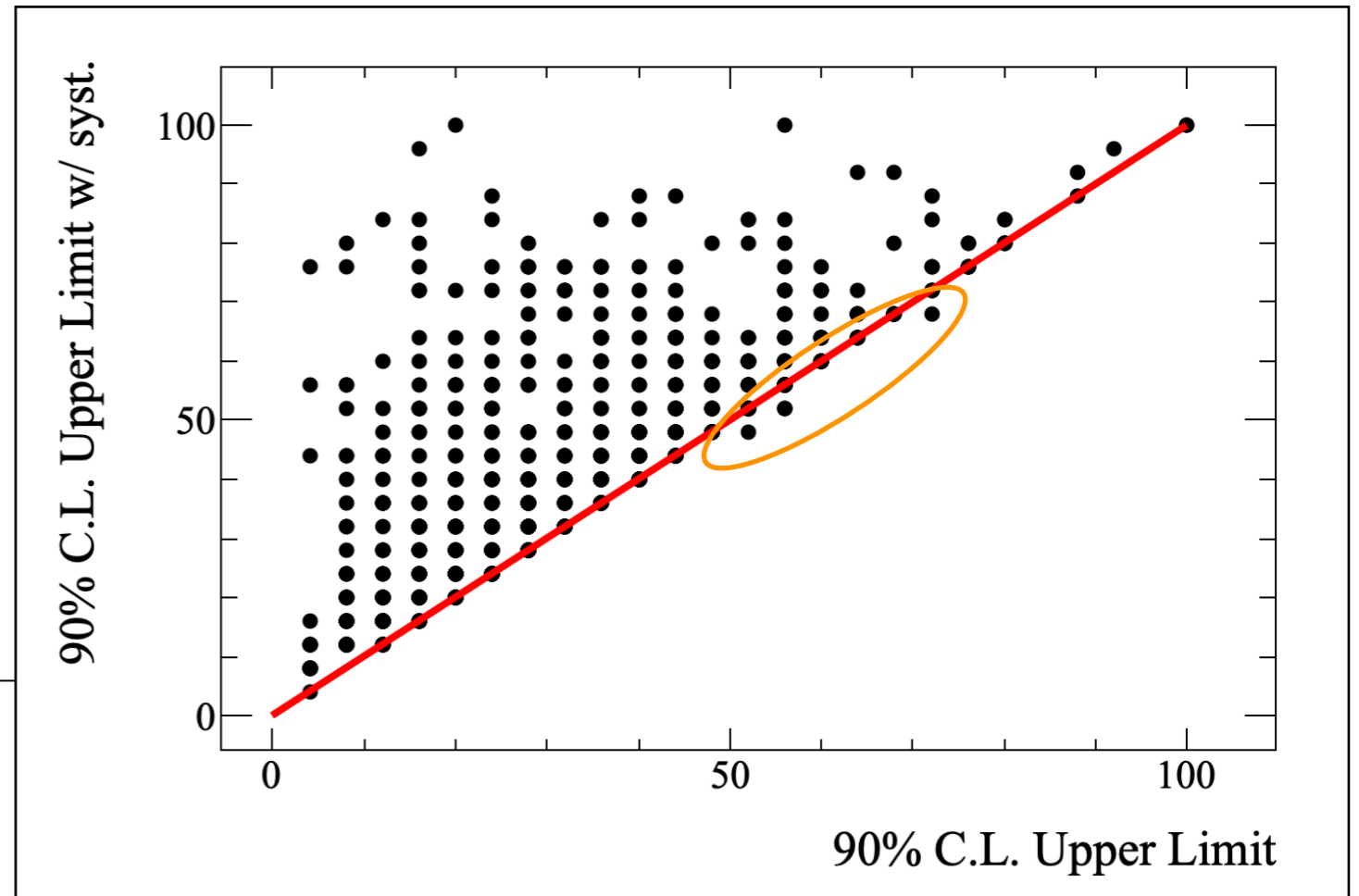
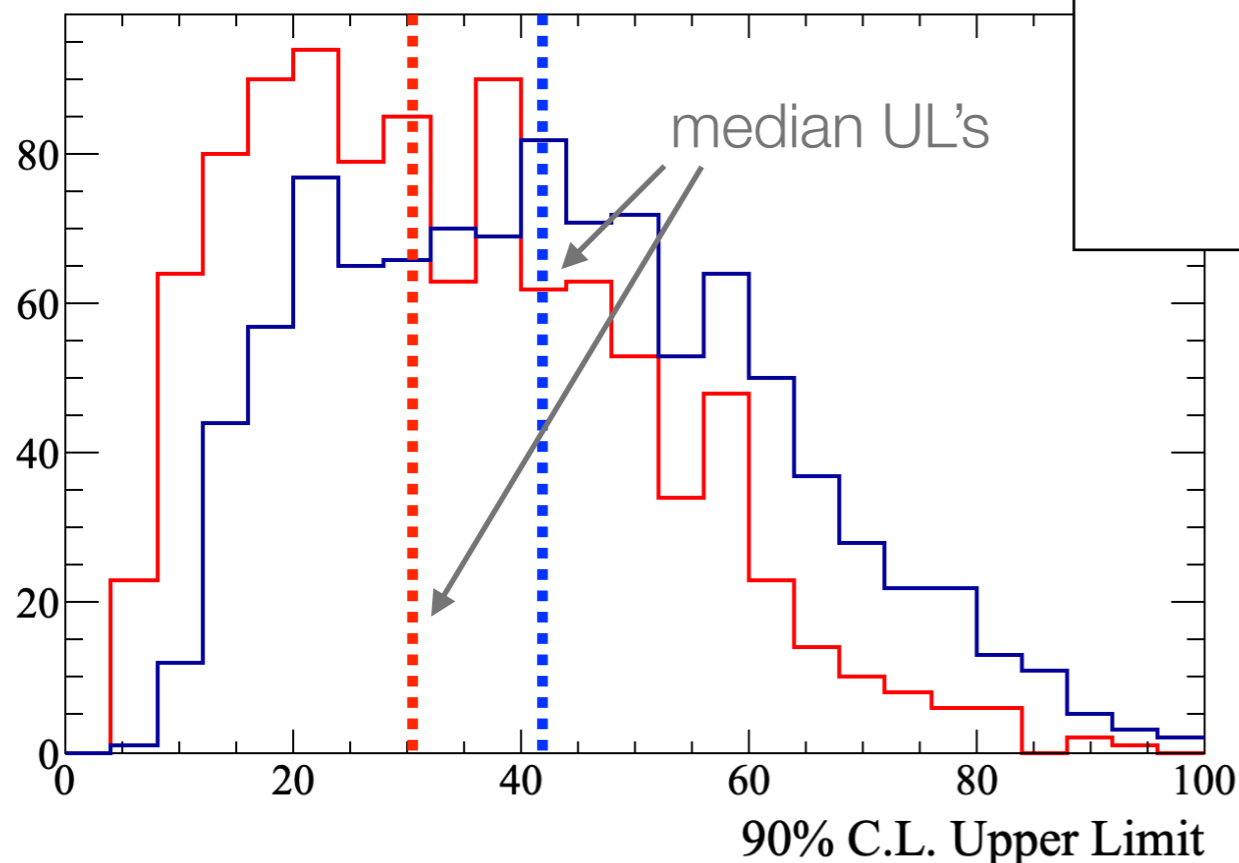
— w/o systematics
— w/ systematics



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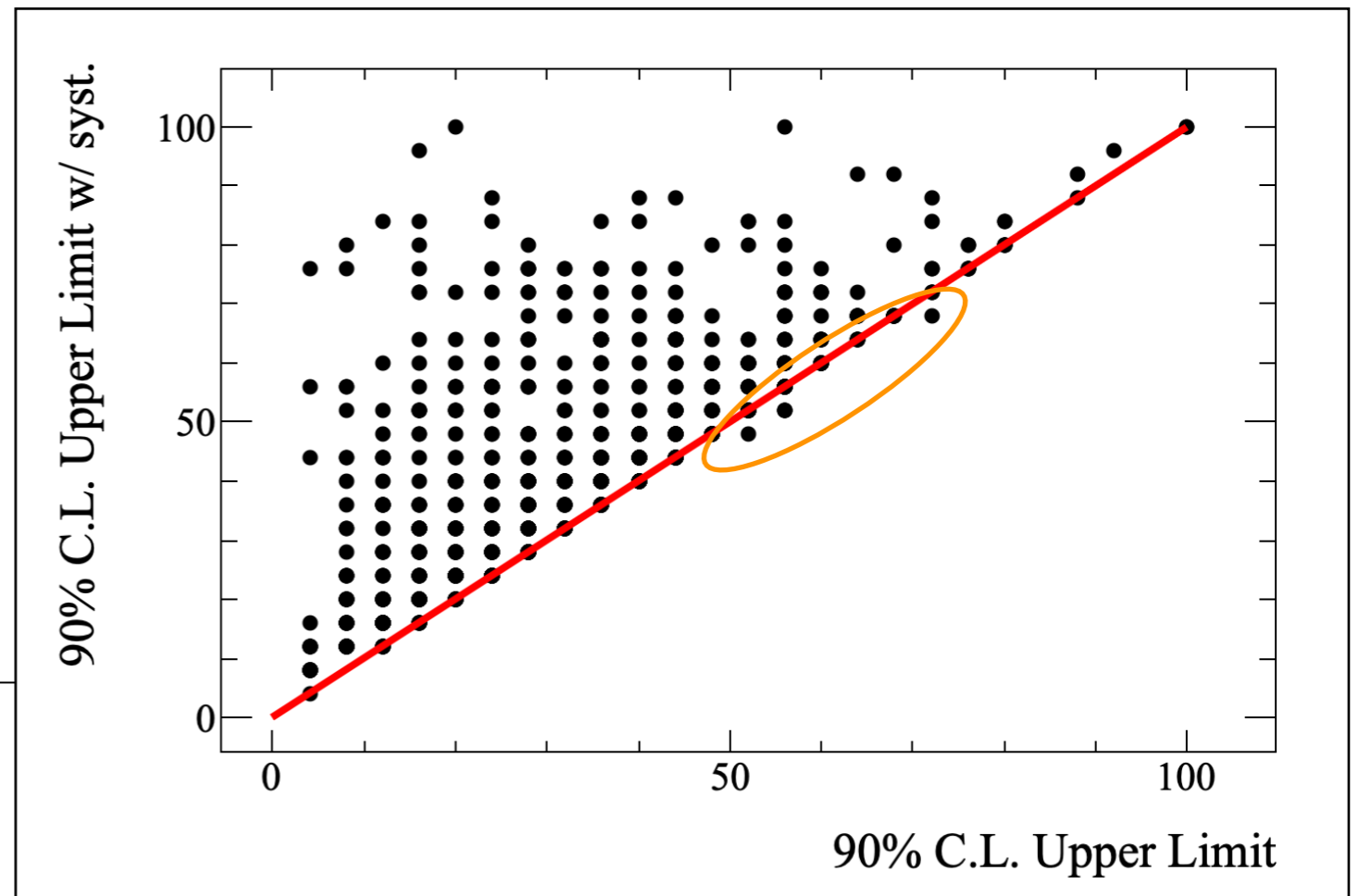
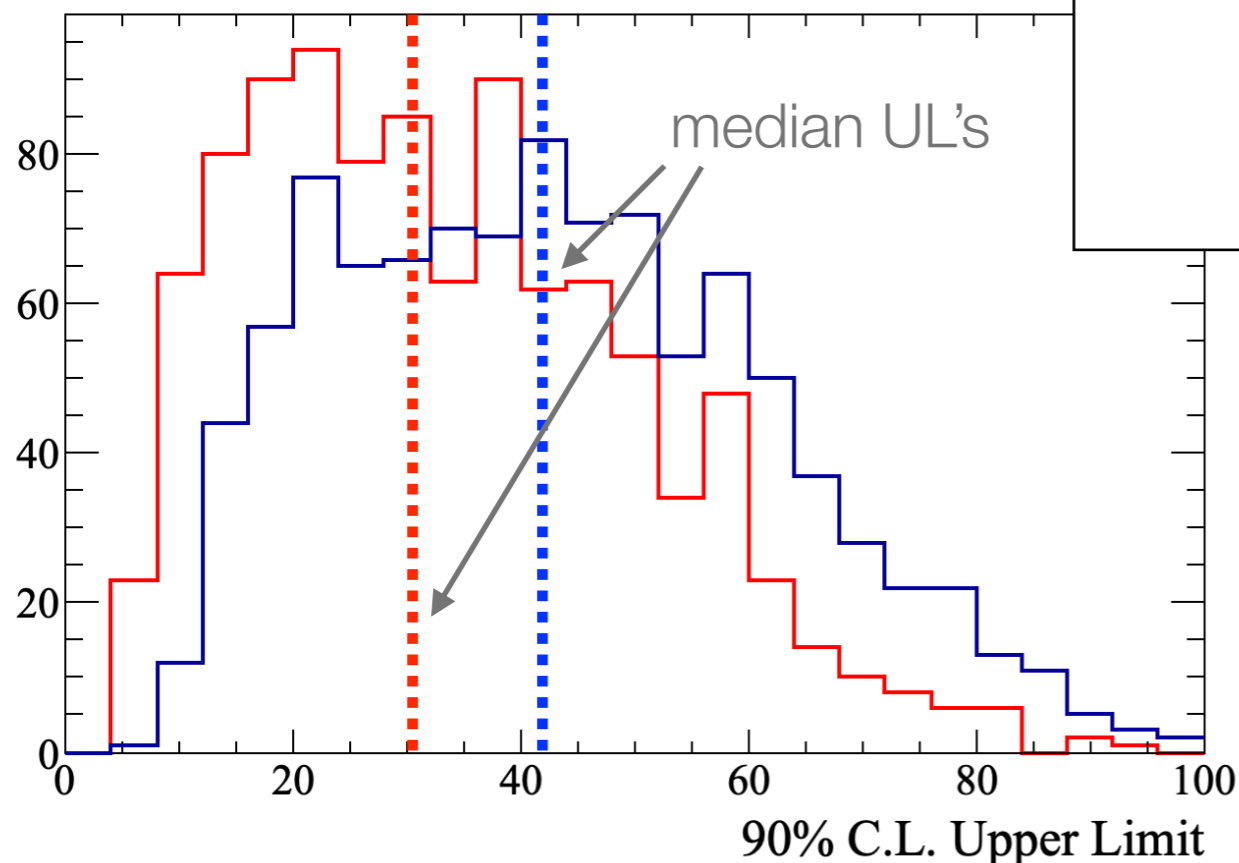


N.B. due to stat. fluctuations, in some experiments the UL w/ systematics can be lower than w/o
Is it a problem?

Inclusion of systematics

With an uncertainty in the mean equal to the gaussian width, the impact in some experiments can be huge (up to a factor 10!!!)

— w/o systematics
— w/ systematics

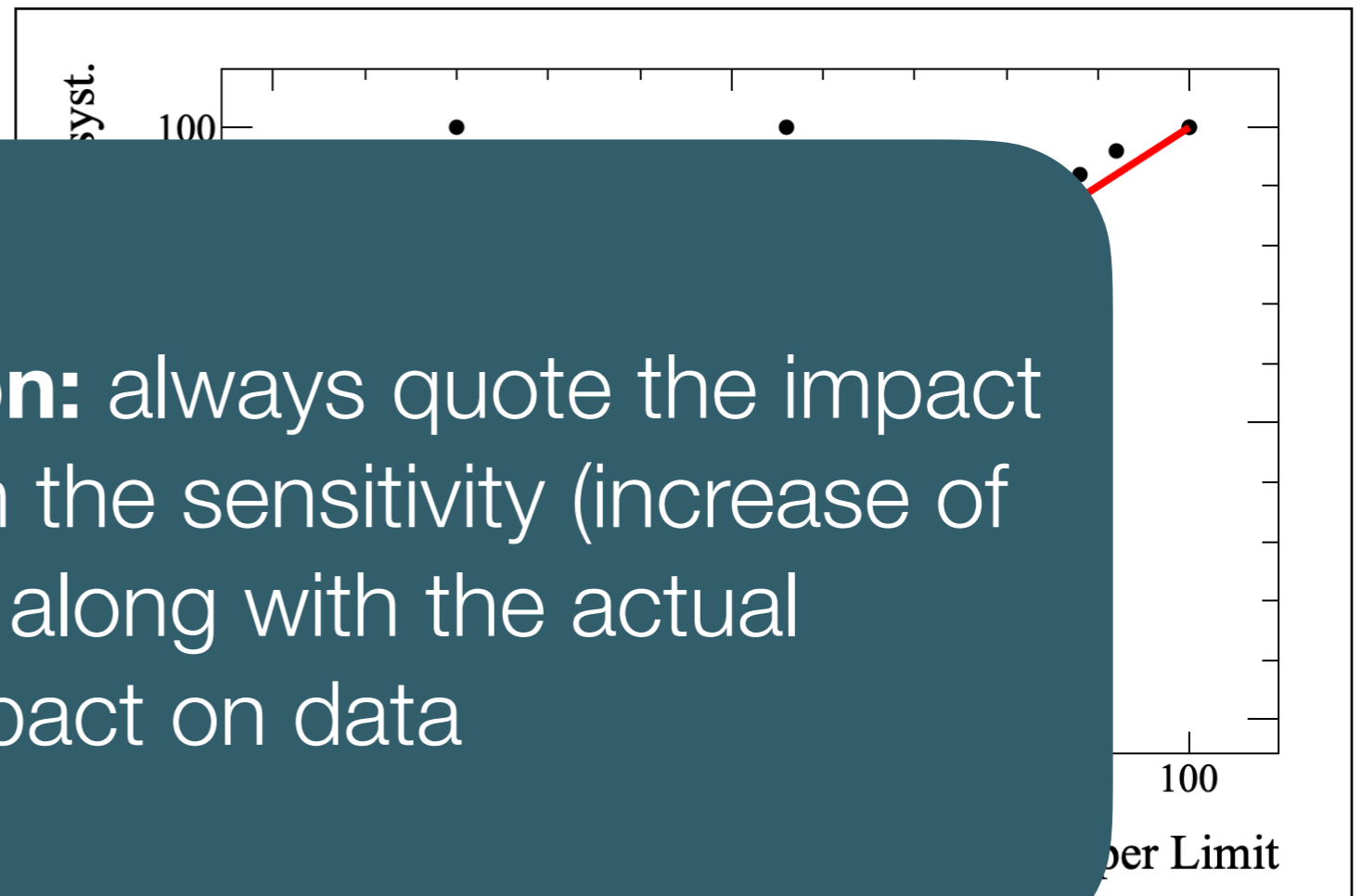


N.B. due to stat. fluctuations, in some experiments the UL w/ systematics can be lower than w/o systematics

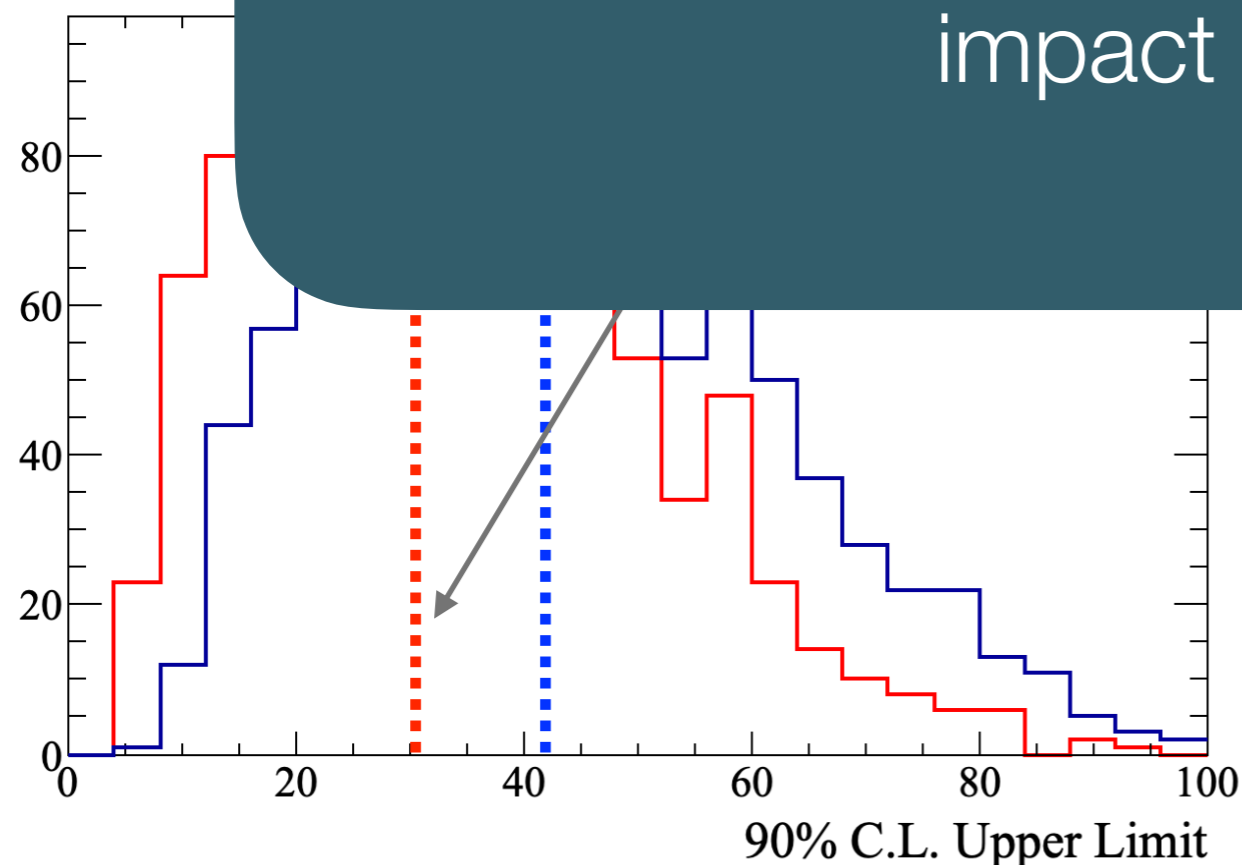
Is it a problem? *No, systematics are treated here in a frequentistic way, what matters is the coverage*

Inclusion of systematics

With an uncertainty in the mean equal to the measurement width, the impact of systematics can be huge



Recommendation: always quote the impact of systematics on the sensitivity (increase of median UL) along with the actual impact on data



N.B. due to stat. fluctuations, in some experiments the UL w/ systematics can be lower than w/o

Is it a problem? *No, systematics are treated here in a frequentistic way, what matters is the coverage*

Inclusion of systematics

- For the inclusion of systematics, the most popular approaches are:
 - simplified approach: in the toy MCs used to build the likelihood ratio distributions, nuisance parameters are randomly fluctuated in the generation or fit of the toy MC experiments (conceptually similar to the semi-bayesian approach)

The generation of the toy MC experiments

- What values of the nuisance parameters should be used when generating the toy MC samples?
- Several options:
 - **a priori estimate:** fixed values decided a priori
 - can have significant under- or over-coverage

The generation of the toy MC experiments

- What values of the nuisance parameters should be used when generating the toy MC samples?
- Several options:
 - **conservative:** generate with the values giving the worst upper limit
 - can have very large over-coverage

The generation of the toy MC experiments

- What values of the nuisance parameters should be used when generating the toy MC samples?
- Several options:
 - **Highland-Cousins:** extract a random value of the nuisance parameters for each toy MC experiment, according to an a-priori distribution
 - can have some over-coverage when the nuisance parameter have a true single value (not varying from experiment to experiment)

The generation of the toy MC experiments

- What values of the nuisance parameters should be used when generating the toy MC samples?
- Several options:
 - **a-posteriori Highland-Cousins:** extract a random value of the nuisance parameters for each toy MC experiment, according to an a-posteriori distribution derived from data
 - can still have some over-coverage, but less than the a-priori method

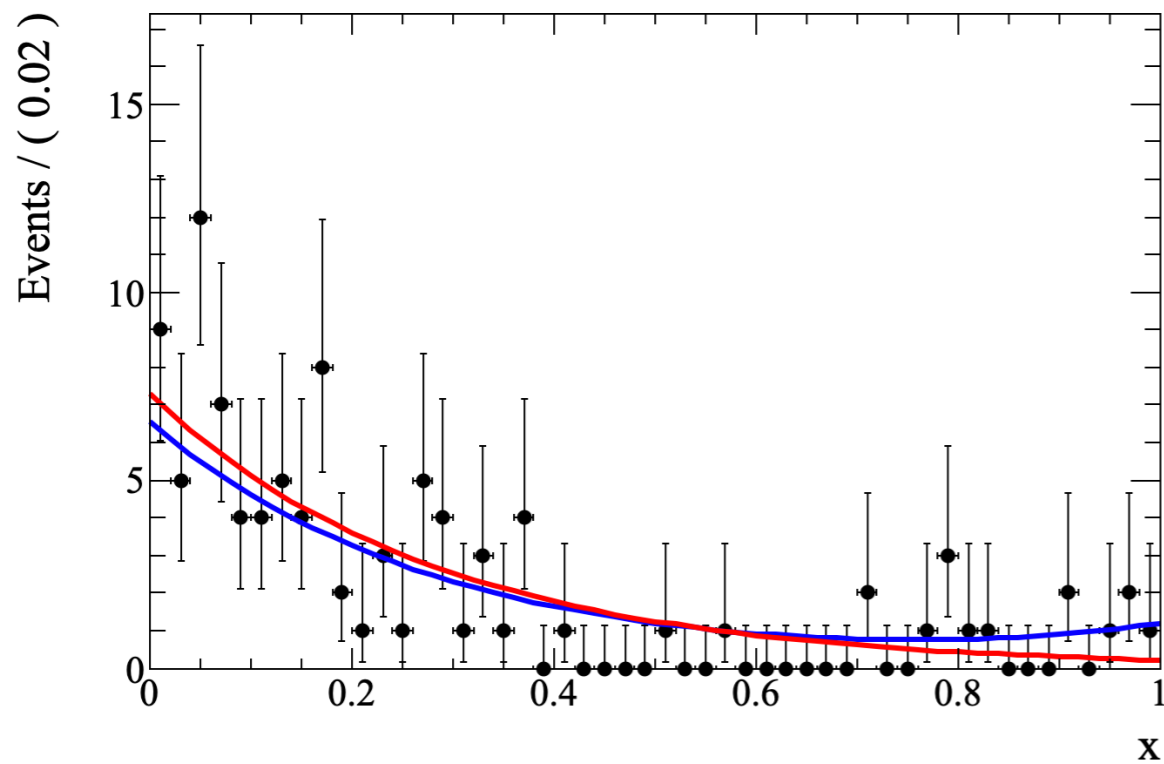
PDFs from small Monte Carlo samples

- We have seen that the generation of MC samples for extremely rare events can be problematic
- Nonetheless, sometimes the use of MC to extract PDFs is unavoidable
 - if the PDF shape is not known a priori, and the MC sample is too small to infer a reliable parameterization, there is a strong risk of overestimating or underestimating the systematic uncertainties, due to the inclusion of unnecessary shape uncertainties

PDFs from small Monte Carlo samples

- True distribution
- Exponential estimated from the MC

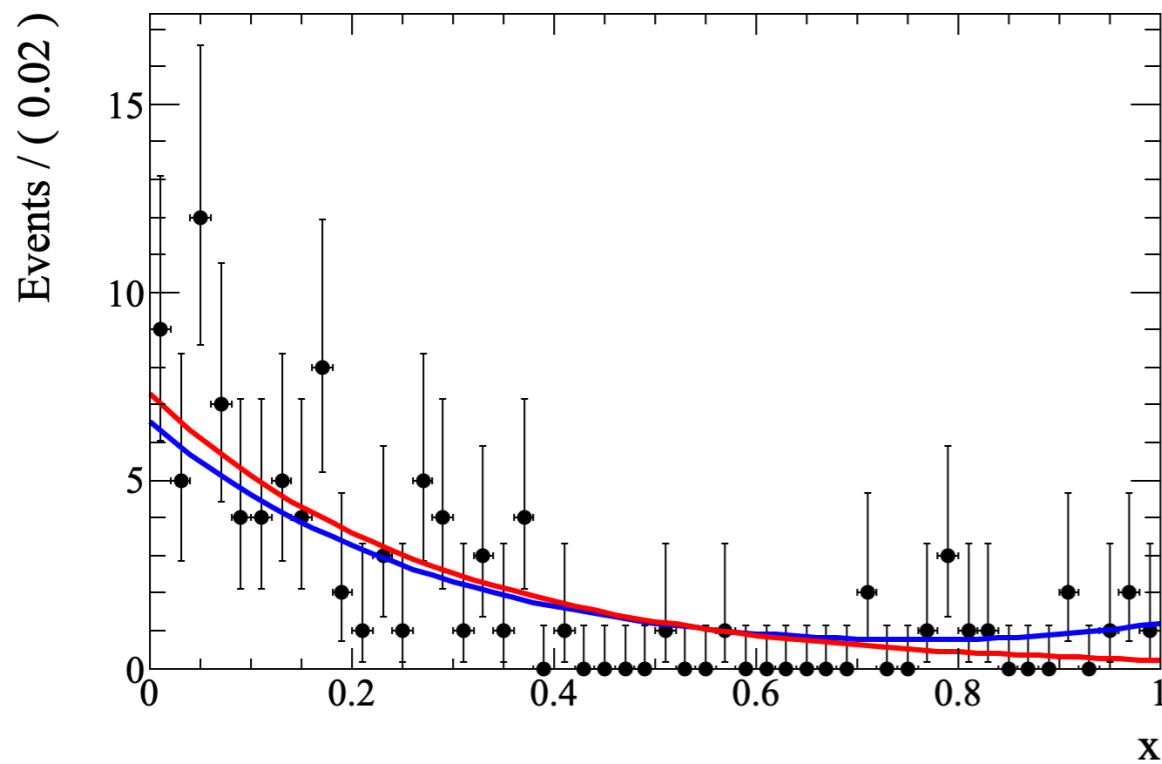
MC



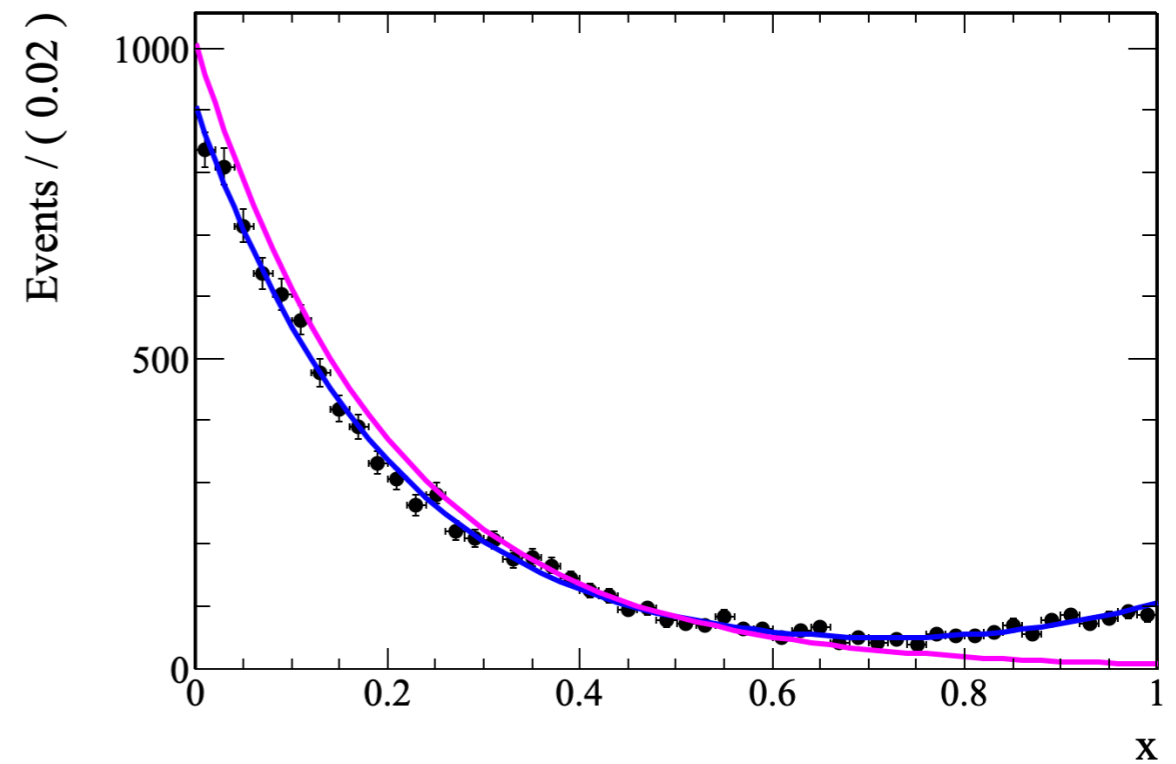
PDFs from small Monte Carlo samples

- True distribution
- Exponential estimated from the MC
- Exponential estimated on data with constraint from the MC

MC



Data



PDFs from small Monte Carlo samples

- We have seen that the generation of MC samples for extremely rare events can be problematic
- Nonetheless, sometimes the use of MC to extract PDFs is unavoidable
 - if the PDF shape is not known a priori, and the MC sample is too small to infer a reliable parameterization, there is a strong risk of overestimating or underestimating the systematic uncertainties, due to the inclusion of unnecessary shape uncertainties
 - using MC histograms to represent the PDFs, with a proper treatment of uncertainties, could be the solution

The Beeston-Barlow approach

- Binned fit with different populations (e.g. signal, backgrounds)
- The expected MC content of each bin i for each population j (A_{ji}) is treated as a nuisance parameter, constrained from the actual MC

$$f_i = \sum_{j=0}^m p_j A_{ji}$$

f_i : expected yield in bin i of data

p_j : data/MC scale factor for population j (to be estimated)

$$\ln \mathcal{L} = \sum_{i=1}^n d_i \ln f_i - f_i + \sum_{i=1}^n \sum_{j=1}^m a_{ji} \ln A_{ji} - A_{ji}$$

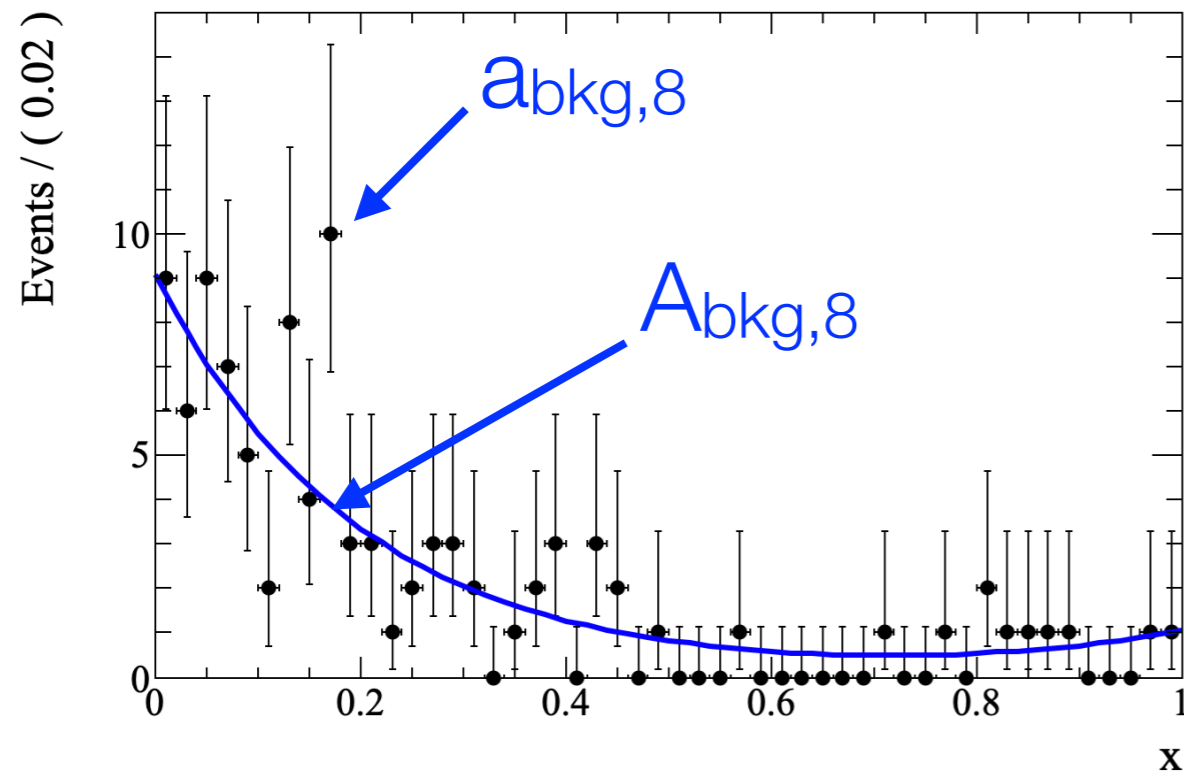
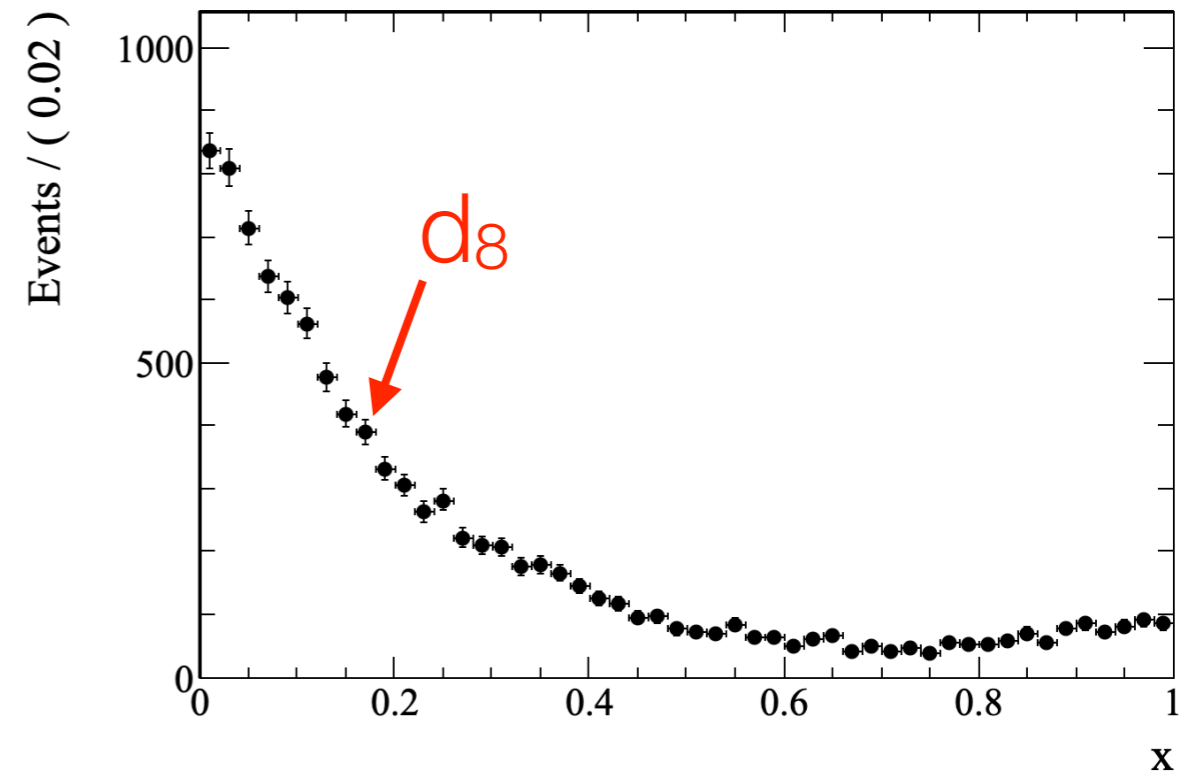
Likelihood for data (Poisson)

d_i : observed yield in bin i of data

Likelihood for MC (Poisson)

a_{ij} : observed yield in bin i of MC for population j
 A_{ji} : expected yield in bin i of MC for population j
 (nuisance parameters)

The Beeston-Barlow approach

MC

Data


$$f_i = \sum_{j=0}^m p_j A_{ji}$$

$$\ln \mathcal{L} = \sum_{i=1}^n d_i \ln f_i - f_i + \sum_{i=1}^n \sum_{j=1}^m a_{ji} \ln A_{ji} - A_{ji}$$

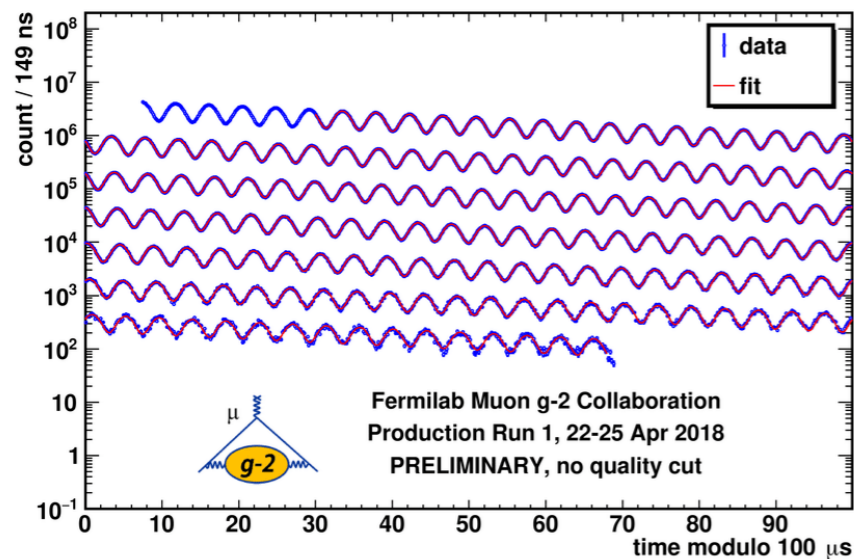
**Can be computationally
unaffordable, lite versions exist**

Practical Examples

Muon g-2 experiment at FNAL

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: measurement of the muon g-2
- Approach: measurement of the spin precession frequency of muons in orbit in a magnetic field



$$\frac{g_\mu - 2}{2} \sim \frac{m\omega}{qB}$$

- Experimental observable: rate of positrons with $E > E_{\text{thr}}$ emitted in forward direction w.r.t. the muon momentum

SYSTEMATICS

- Main criticality: accuracy of the magnetic field

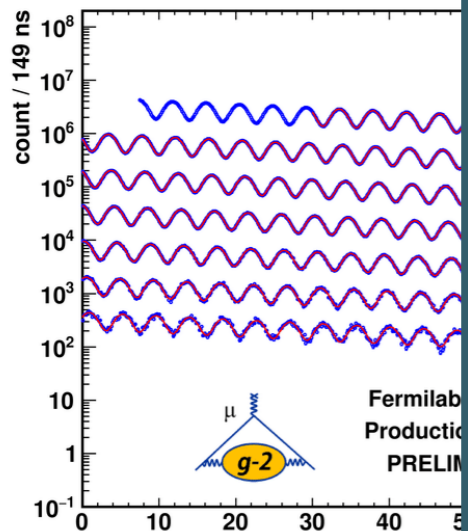
$$a_\mu = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} \frac{\mu'_p(T_r)}{\mu_e(H)} \frac{\mu_e(H)}{\mu_e} \frac{m_\mu g_e}{m_e 2}$$

- Solution: instead of taking absolute field and frequency measurements, a **standard metrologic process** (*precession frequency of shielded protons in a spherical sample, ω_p*) is measured in the same field, and the ratio is used
- Residual uncertainties from beam dynamics, temperature stability, external inputs

Muon g-2 experiment at FNAL

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: measure the anomalous magnetic moment of the muon, a_μ .
- Approach: measure the precession frequency of the muon spin in a magnetic field.



- Experiment uses muons and positrons with $E > E_{\text{thr}}$ emitted in forward direction w.r.t. the muon momentum

Quantity	Correction (ppb)	Uncertainty (ppb)
ω_a^m (statistical)	...	201
ω_a^m (systematic)	...	25
C_e	451	32
C_p	170	10
C_{pa}	-27	13
C_{dd}	-15	17
C_{ml}	0	3
$f_{\text{calib}} \cdot \langle \omega'_p(\vec{r}) \times M(\vec{r}) \rangle$...	46
B_k	-21	13
B_q	-21	20
$\mu'_p(34.7^\circ)/\mu_e$...	11
m_μ/m_e	...	22
$g_e/2$...	0
Total systematic for \mathcal{R}'_μ	...	70
Total external parameters	...	25
Total for a_μ	622	215

accuracy of the

$$\frac{T_r}{H} \frac{\mu_e(H)}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

of taking
and frequency
a **standard**
process

frequency of
in a spherical
measured in the
the ratio is used

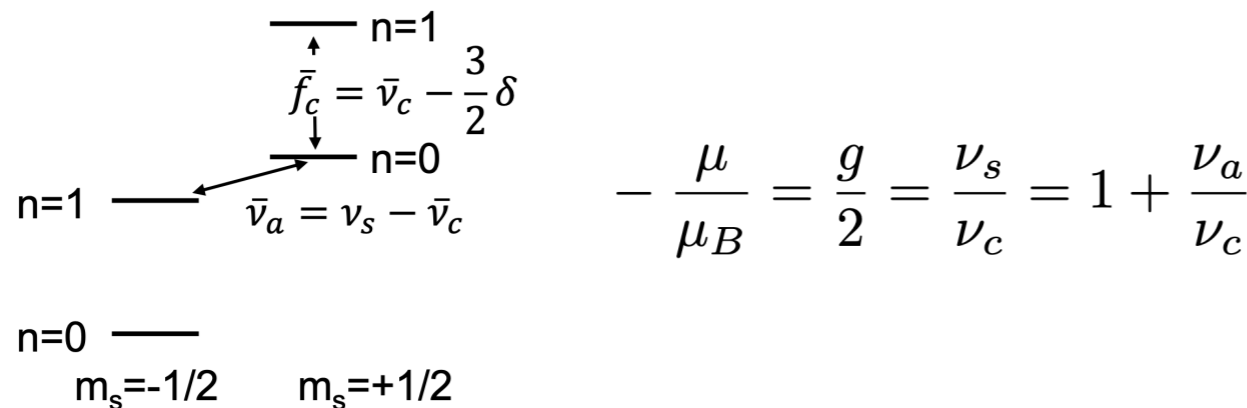
ainties from

beam dynamics, temperature stability, external inputs

Electron g-2

OBJECTIVE AND EXPERIMENTAL APPROACH

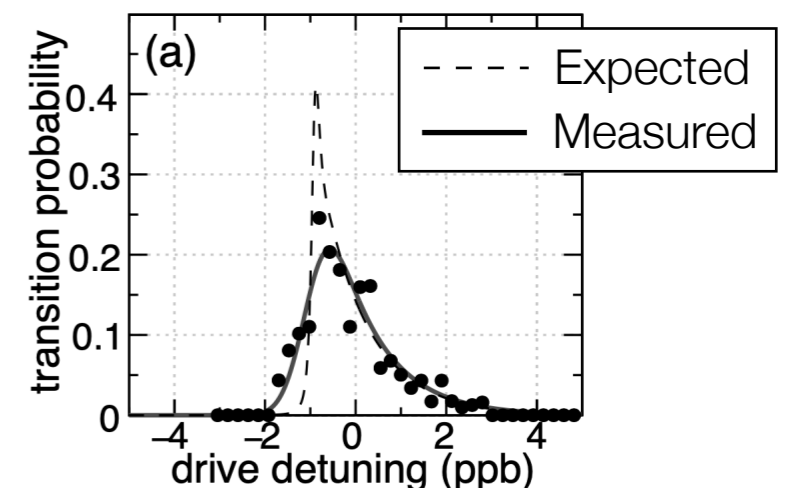
- Objective: measurement of the electron g-2
- Approach: measurement of energy transitions for a single electron in a magnetic field inside a Penning trap



- Experimental observable: quantum jumps $h\nu$, excited with electrodes, induce currents in the electrodes themselves, with a resonance if the frequency of the excitation matches ν

SYSTEMATICS

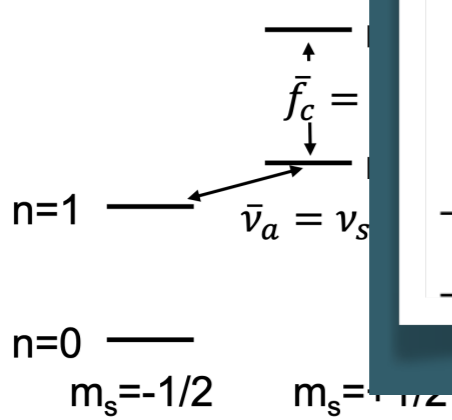
- The necessary quantities ν_a , ν_c are measured in situ with the same approach \rightarrow no metrology issue
- Dominant uncertainty is expected from the correction between frequencies for free and trapped electrons
- Indeed, B field fluctuations are observed and induce additional systematics



Electron g-2

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: measure $g-2$
- Approach: measure transitions between magnetic field states

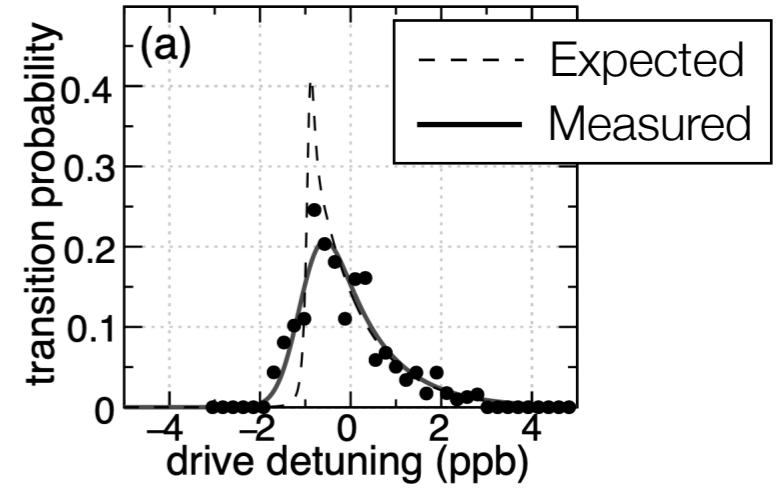


- Experimental observable: quantum jumps $h\nu$, excited with electrodes, induce currents in the electrodes themselves, with a resonance if the frequency of the excitation matches ν

SYSTEMATICS

quantities ν_a, ν_c in situ with the $n \rightarrow n_0$ transition. The uncertainty is dominated by the correction to the frequencies for free electrons. Fluctuations are induced additional

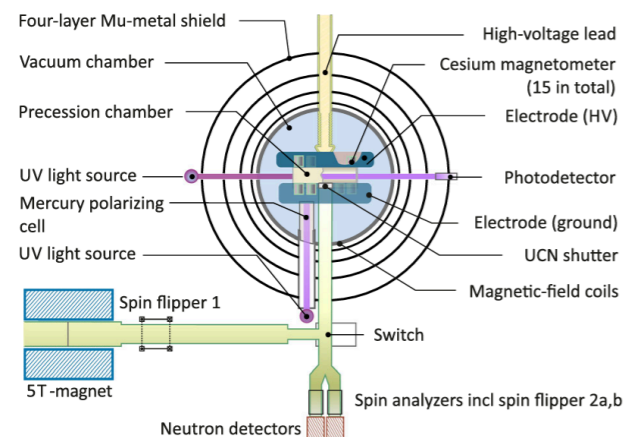
Source	Uncertainty $\times 10^{13}$
statistical	0.29
cyclotron broadening	0.94
cavity correction	0.90
nuclear paramagnetism	0.12
anomaly power shift	0.10
magnetic field drift	0.09
total	1.3



nEDM

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: search for a neutron EDM
- Approach: measurement of spin precession frequency in $E + B$ field



$$f_n = \frac{1}{\pi\hbar} |\mu_n \vec{B}_0 + d_n \vec{E}|$$

- Experimental technique: Ramsey spectroscopy of polarized ultra-cold neutrons within a shielded chamber:
 - $\pi/2$ spin flip by oscillating magnetic field
 - find the frequency f_n that maximizes the asymmetry btw. spin up and down

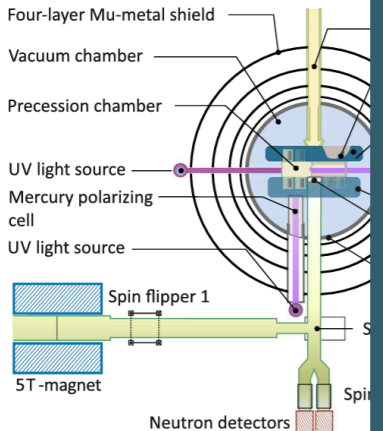
SYSTEMATICS

- Main criticality: this is not a genuine measurement of zero, because $\mu_n B_0$ has to be subtracted from a non-zero measurement:
 - absolute value of B_0 need to be known
 - comagnetometry against a metrologic standard (^{199}Hg)
- Several residual systematics:
 - systematics in the ^{199}Hg measurement
 - magnetic non-uniformities
 - asymmetries in the distribution of neutrons w.r.t. the magnetic field in the chamber
 -

nEDM

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective:
- Approach: precession



- Experimental spectroscopy neutrons with
 - $\pi/2$ spin flipper field
 - find the frequency f_n that maximizes the asymmetry btw. spin up and down

SYSTEMATICS

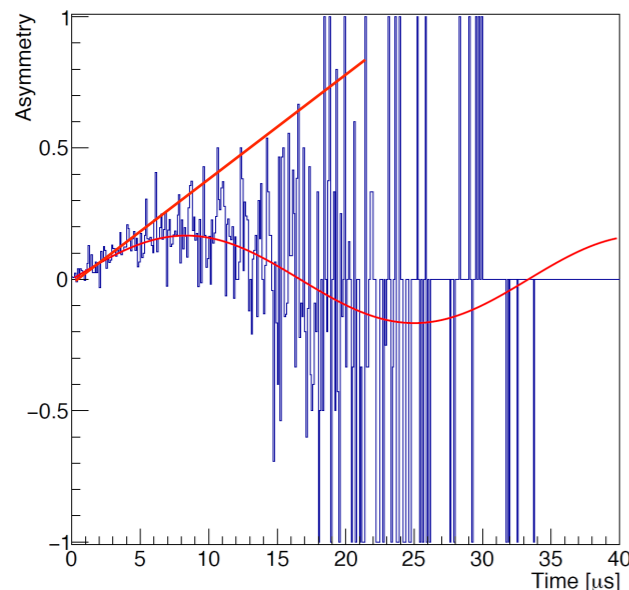
- Main criticality: this is not a measurement of zero, it is to be non-zero
- of B_0 need to
- symmetry against a standard (^{199}Hg)
- systematics: of the ^{199}Hg
- uniformities
- in the distribution r.t. the magnetic field in the chamber
-

Effect	Shift	Error
Error on $\langle z \rangle$...	7
Higher-order gradients \hat{G}	69	10
Transverse field correction $\langle B_T^2 \rangle$	0	5
Hg EDM [8]	-0.1	0.1
Local dipole fields	...	4
$v \times E$ UCN net motion	...	2
Quadratic $v \times E$...	0.1
Uncompensated G drift	...	7.5
Mercury light shift	...	0.4
Inc. scattering ^{199}Hg	...	7
TOTAL	69	18

muEDM

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: search for a muon EDM
- Approach: detection of non-zero spin precession in a magnetic field, with MDM precession canceled by a suitable combination of E and B fields



$$\vec{d} = \frac{\eta e}{2mc} \vec{s}$$

$$\omega = \frac{\eta q}{2m} \vec{\beta} \times \vec{B}$$

$$A(t) \propto \frac{2P_0 E_f \alpha |d_\mu|}{a \hbar \gamma^2} t$$

- Experimental observable: time-dependent asymmetry of positrons emitted along and opposite to the B field

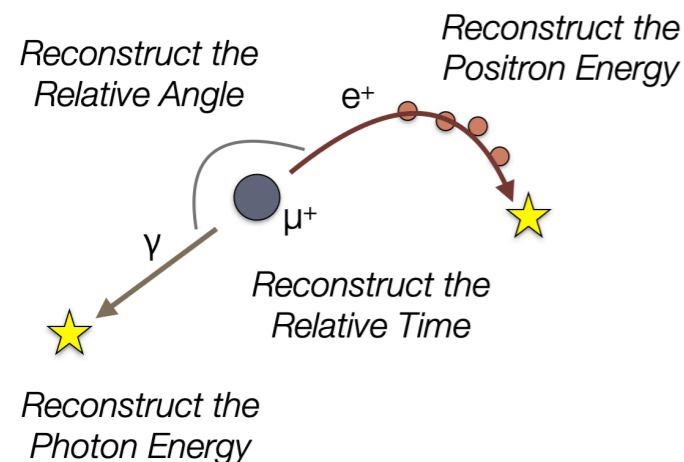
SYSTEMATICS

- Main criticality:
 - it is not necessary to know extremely well the main components of the field, but fake EDM can arise from fringe fields
- EDM is CP-violating, standard electrodynamics is CP-conserving:
 - systematics can be canceled by inverting B and injection direction
 - indeed, it moves the systematics from electrodynamics to the symmetry between the two injection modes
- Detector asymmetries to be kept under control

MEG & MEG II

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: search for $\mu \rightarrow e\gamma$
- Approach: search for e^+ and γ in μ^+ decays at rest with 2-body kinematics (monochromatic e^+ and γ at 180°)



$$E_e = E_\gamma = 52.8 \text{ MeV}$$

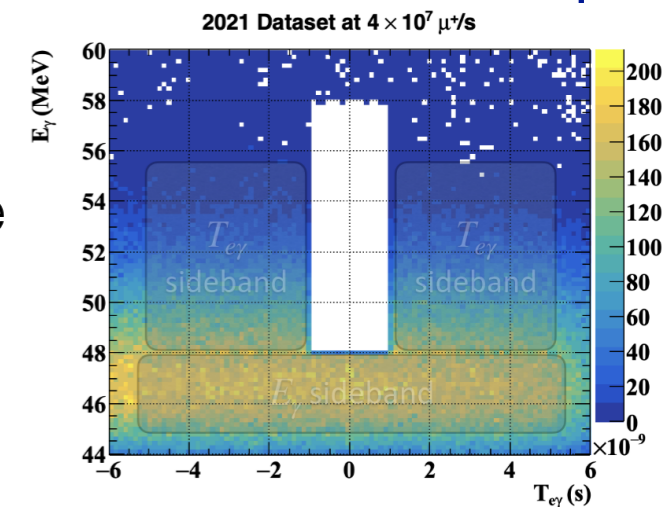
$$\Theta_{e\gamma} = 180^\circ$$

$$T_{e\gamma} = 0$$

- Experimental technique: photon reconstruction in a LXe calorimeter, positron reconstruction in a magnetic spectrometer

SYSTEMATICS

- Robust control of the background from sidebands

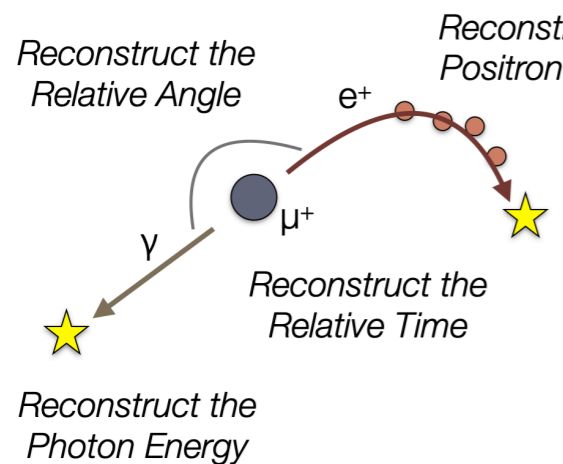


- Main criticalities:
 - calibrating the photon energy scale requires a dedicated $\pi^-(p, n)\pi^0$ experiment
 - **no physics process to calibrate the relative angle \rightarrow rely on detector alignments**
- Dominant systematics from target alignment:
 - tolerable in MEG, required a dedicated target monitoring system with photo cameras in MEG II due to better resolutions

MEG & MEG II

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: search for $\mu \rightarrow e\gamma$
- Approach: search for e^+ and γ in μ^+ decays at rest w (monochromatic)

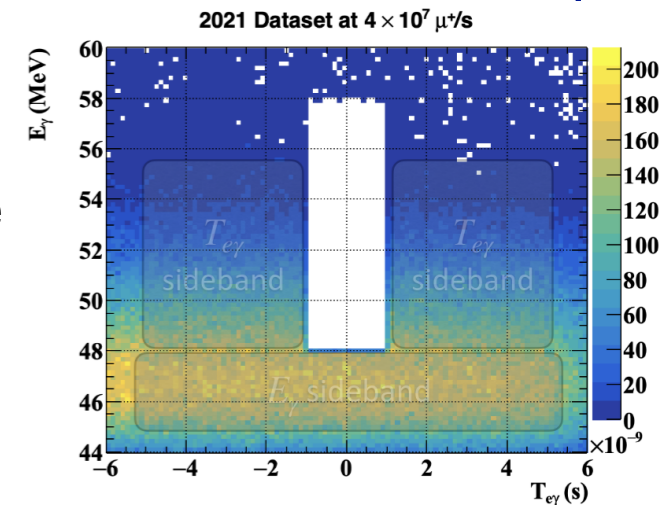


- Experimental technique: photon reconstruction in a LXe calorimeter, positron reconstruction in a magnetic spectrometer

SYSTEMATICS

- Robust control of the background

Parameter	Impact on limit
$\phi_{e\gamma}$ uncertainty	1.1 %
E_γ uncertainty	0.9 %
$\theta_{e\gamma}$ uncertainty	0.7 %
Normalization uncertainty	0.6 %
$t_{e\gamma}$ uncertainty	0.1 %
E_e uncertainty	0.1 %
RDC uncertainty	< 0.1 %



ities:

g the photon energy scale
a dedicated $\pi^-(p, n) \pi^0$
ent

ics process to calibrate
ive angle \rightarrow rely on
alignments

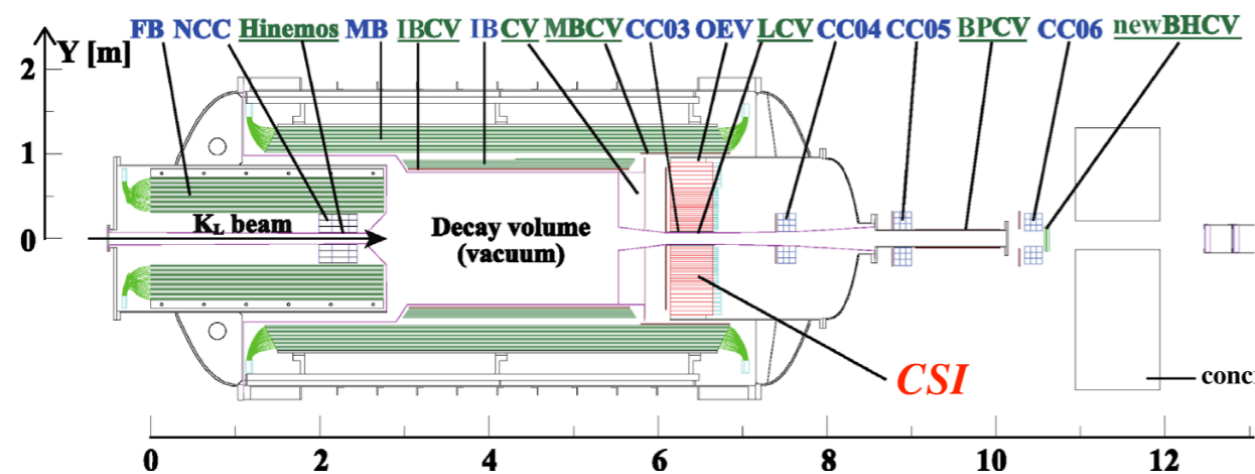
Dominant systematics from target alignment:

- tolerable in MEG, required a dedicated target monitoring system with photo cameras in MEG II due to better resolutions

KOTO

OBJECTIVE AND EXPERIMENTAL APPROACH

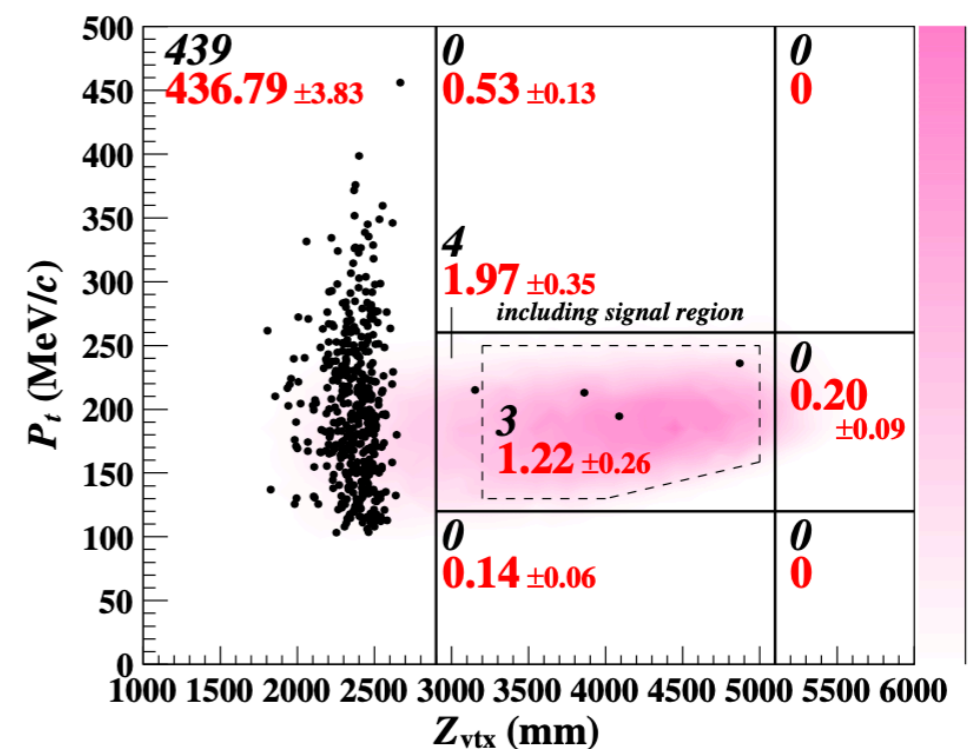
- Objective: search for $K_L^0 \rightarrow \pi^0 \nu \bar{\nu}$
- Approach: search for kaon decays with only 2 photons and nothing else



- Experimental technique: kaons decaying in a volume surrounded by hermetic neutral and charged particle detectors, used to veto background decays

SYSTEMATICS

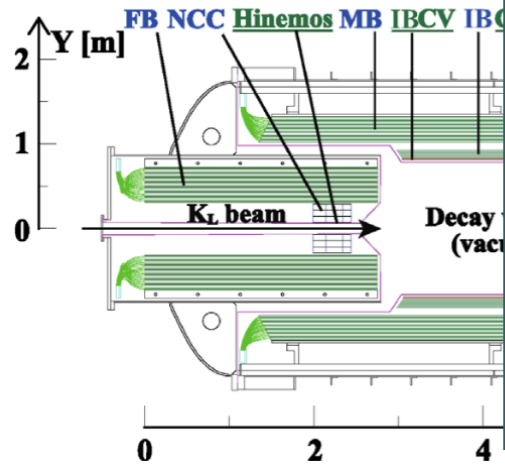
- Main criticality: hermeticity and PID to be precisely controlled to get rid of the Standard Model and beam-halo backgrounds up to 10^7 rejection factor:
 - dominant systematics from the expected background rates



KOTO

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: search for $K_L^0 \rightarrow \pi^0 \nu \bar{\nu}$
- Approach: search for kaon decays with only 2 photons

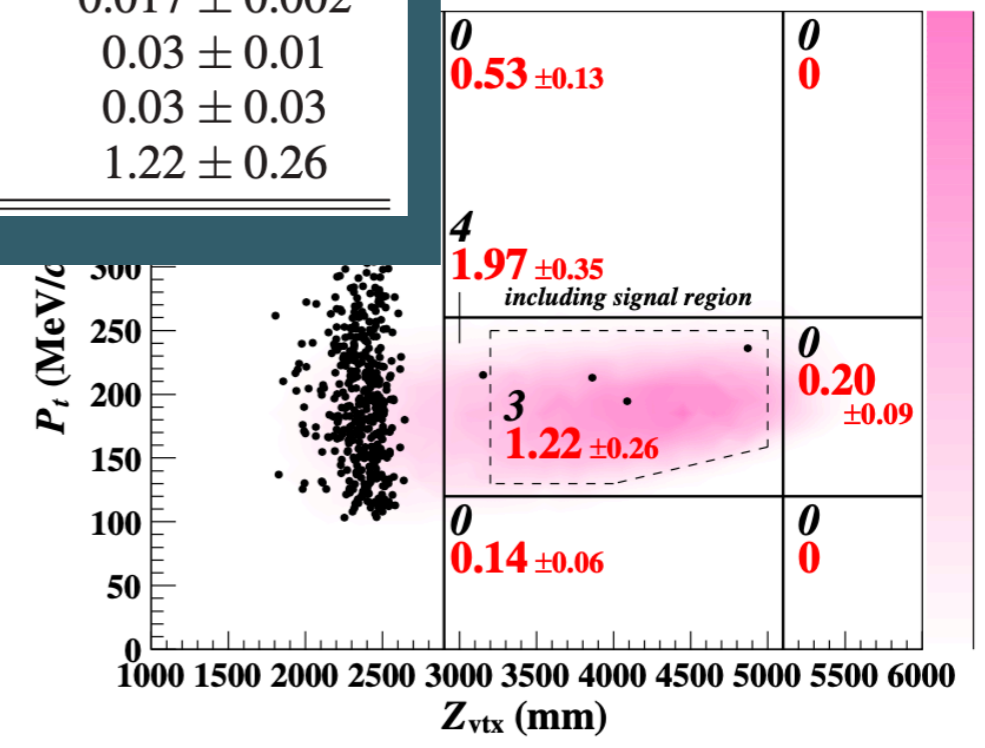


Source		Number of events
K_L	$K_L \rightarrow 3\pi^0$	0.01 ± 0.01
	$K_L \rightarrow 2\gamma$ (beam halo)	0.26 ± 0.07^a
	Other K_L decays	0.005 ± 0.005
K^\pm		0.87 ± 0.25^a
Neutron	Hadron cluster	0.017 ± 0.002
	CV η	0.03 ± 0.01
	Upstream π^0	0.03 ± 0.03
Total		1.22 ± 0.26

SYSTEMATICS

- Main criticality: hermeticity and PID to be precisely controlled to control of the Standard Model
- also backgrounds up factor:
- systematics from the background rates

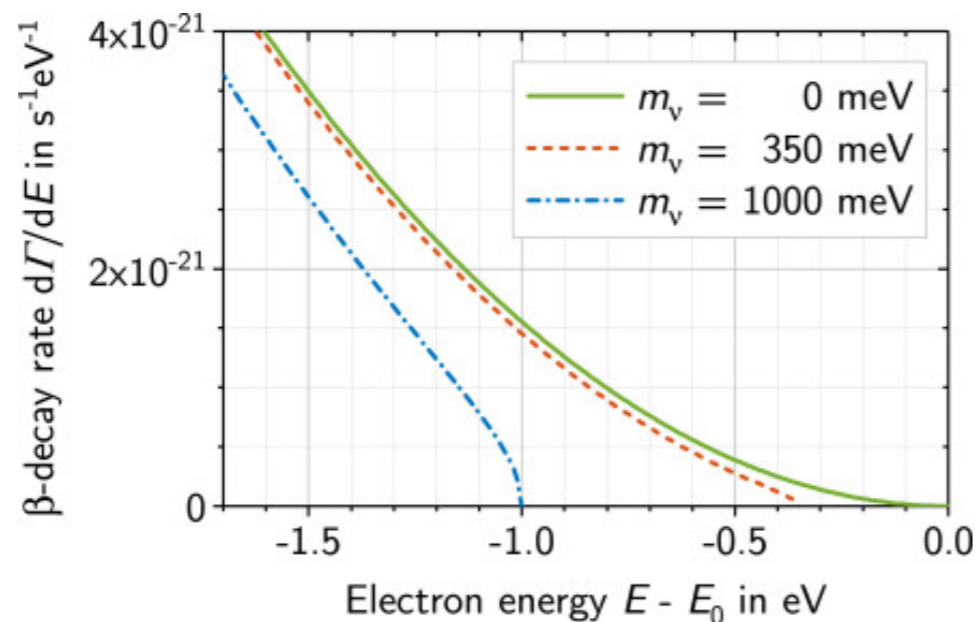
- Experimental technique: kaons decaying in a volume surrounded by hermetic neutral and charged particle detectors, used to veto background decays



KATRIN

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: measurement of the electron anti-neutrino mass
- Approach: measurement of the beta decay end-point



- Experimental technique: electromagnetic filter to count events above a certain energy threshold

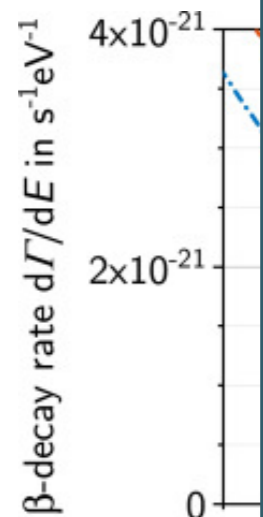
SYSTEMATICS

- Dominant systematics:
 - electric and magnetic field (accuracy & stability)
 - non-Poisson background
- Origin of non-Poisson background:
 - nuclear decays from contaminants produce keV electrons
 - they ionize the residual gas, producing secondaries
 - many secondaries from a very small number of primaries (i.e. correlated background) \rightarrow non-Poisson fluctuations

KATRIN

OBJECTIVE AND EXPERIMENTAL APPROACH

- Objective: measure the anti-neutrino mass
- Approach: endpoint of β -decay



Effect	68.2% CL uncertainty on m_ν^2 (eV ²)
Statistical	0.29
Non-Poissonian background	0.11
Source-potential variations	0.09
Scan-step-duration-dependent background	0.07
qU -dependent background	0.06
Magnetic fields	0.04
Molecular final-state distribution	0.02
Column density \times inelastic scattering cross section	0.01
Activity fluctuations	0.01
Energy-loss function	<0.01
Detector efficiency	<0.01
Theoretical corrections	<0.01
High-voltage stability and reproducibility	<0.01
Total uncertainty	0.34

- Experimental technique: electromagnetic filter to count events above a certain energy threshold

Systematics:
 magnetic field stability)
 background
 Poisson
 decays from
 events produce
 ons
 the residual
 cing
 s
 boundaries from a
 very small number of
 primaries (i.e. correlated
 background) \rightarrow non-
 Poisson fluctuations

Conclusions

- Compared to general-purpose experiments, systematics in single-purpose experiments pose some special challenges:
 - metrology issues
 - need of dedicated (hardware) tools for the control of systematics
- Control of systematic uncertainties is a critical aspect in the design of single-purpose experiments, often requiring special expertise from outside the HEP field