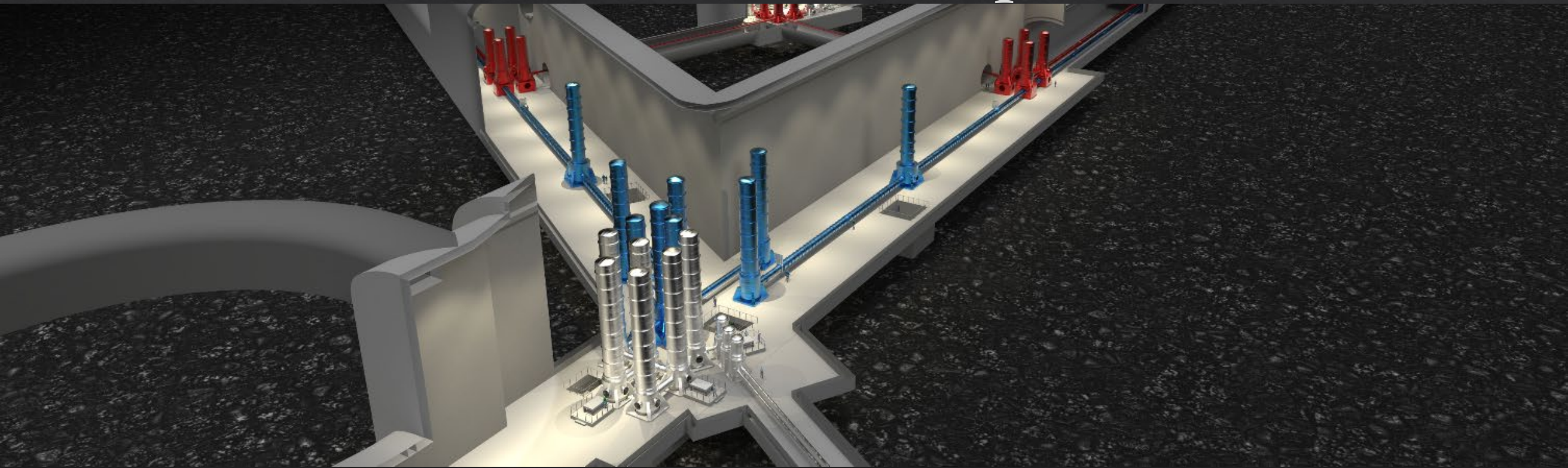


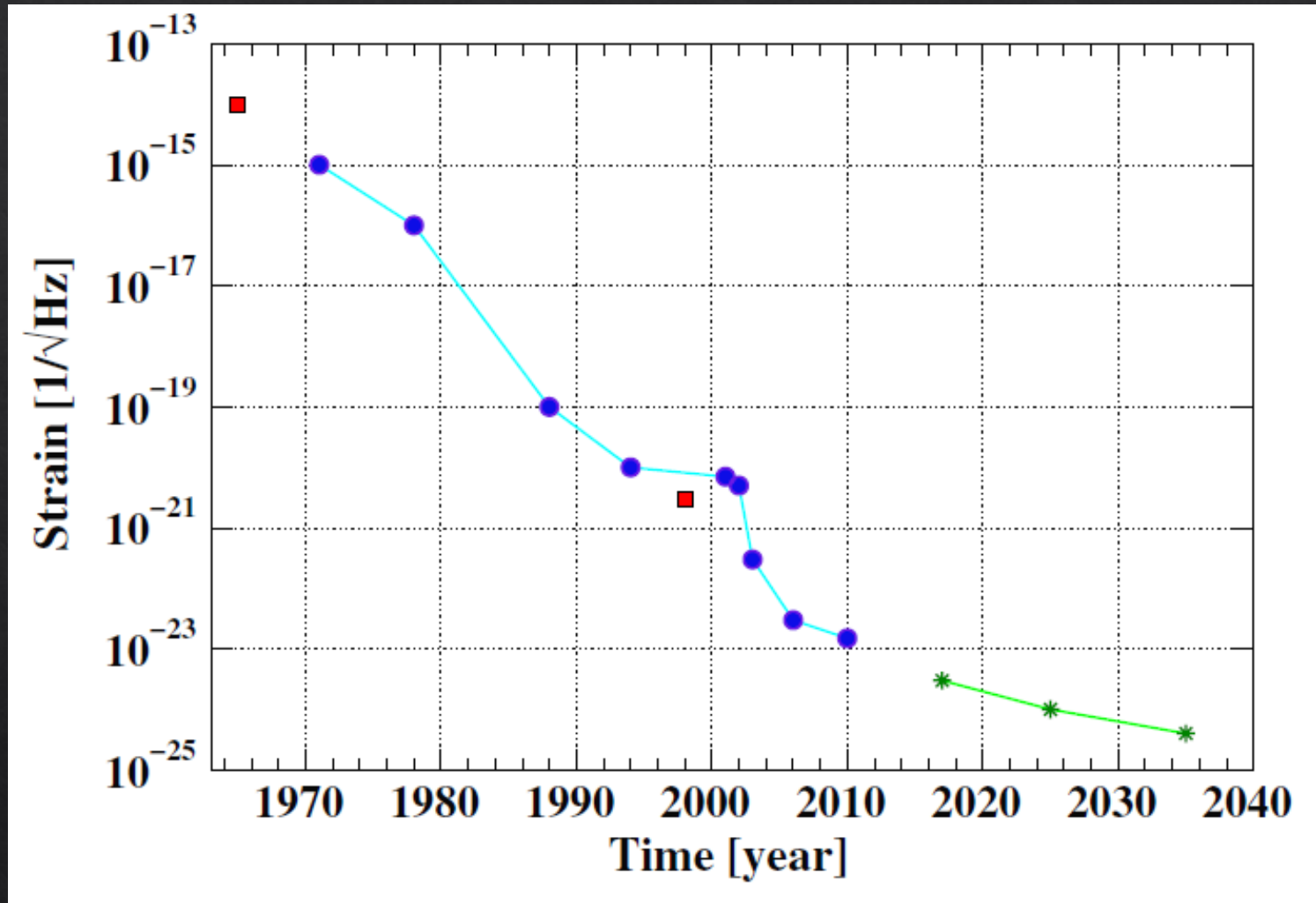
Einstein Telescope



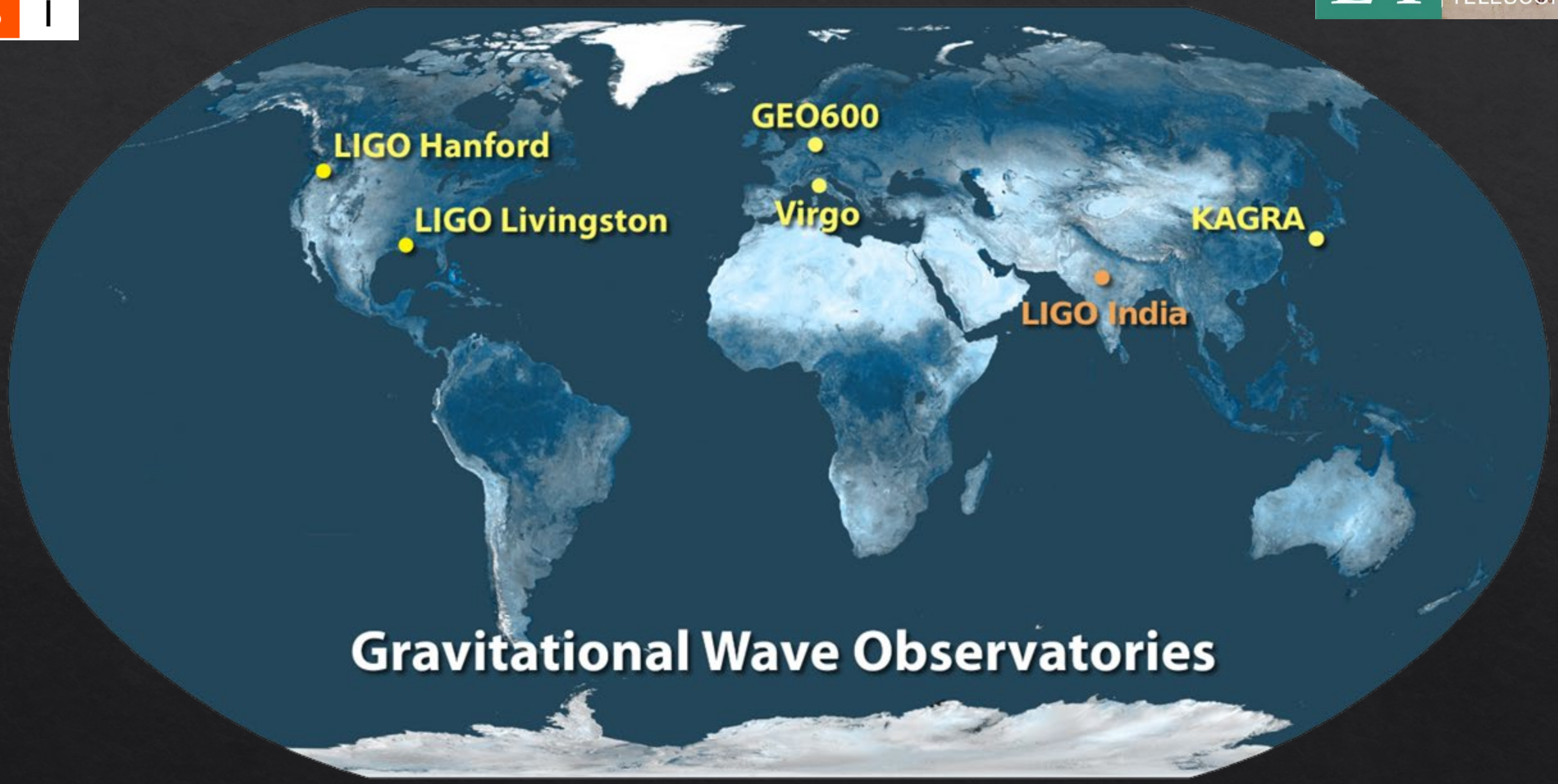
Jan Harms

Gran Sasso Science Institute
INFN - National Laboratory of Gran Sasso

History of GW Detector Sensitivity

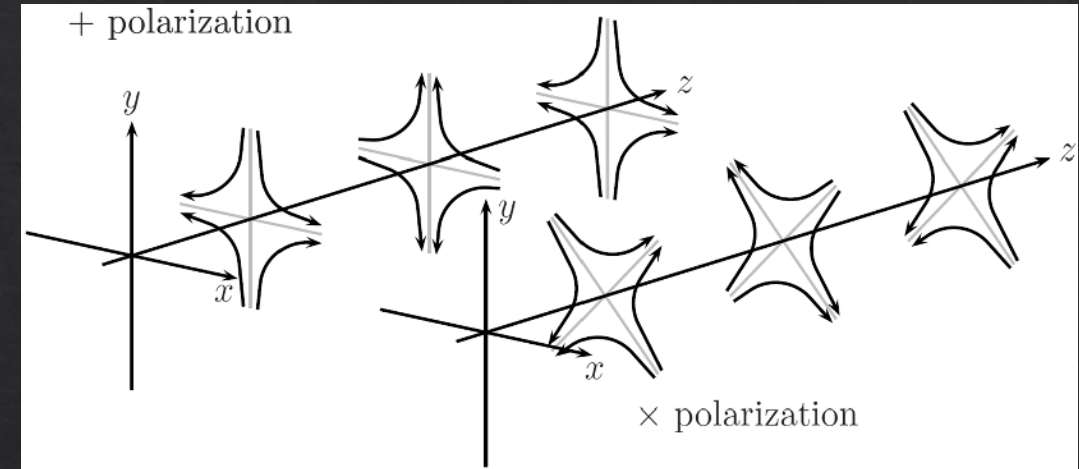
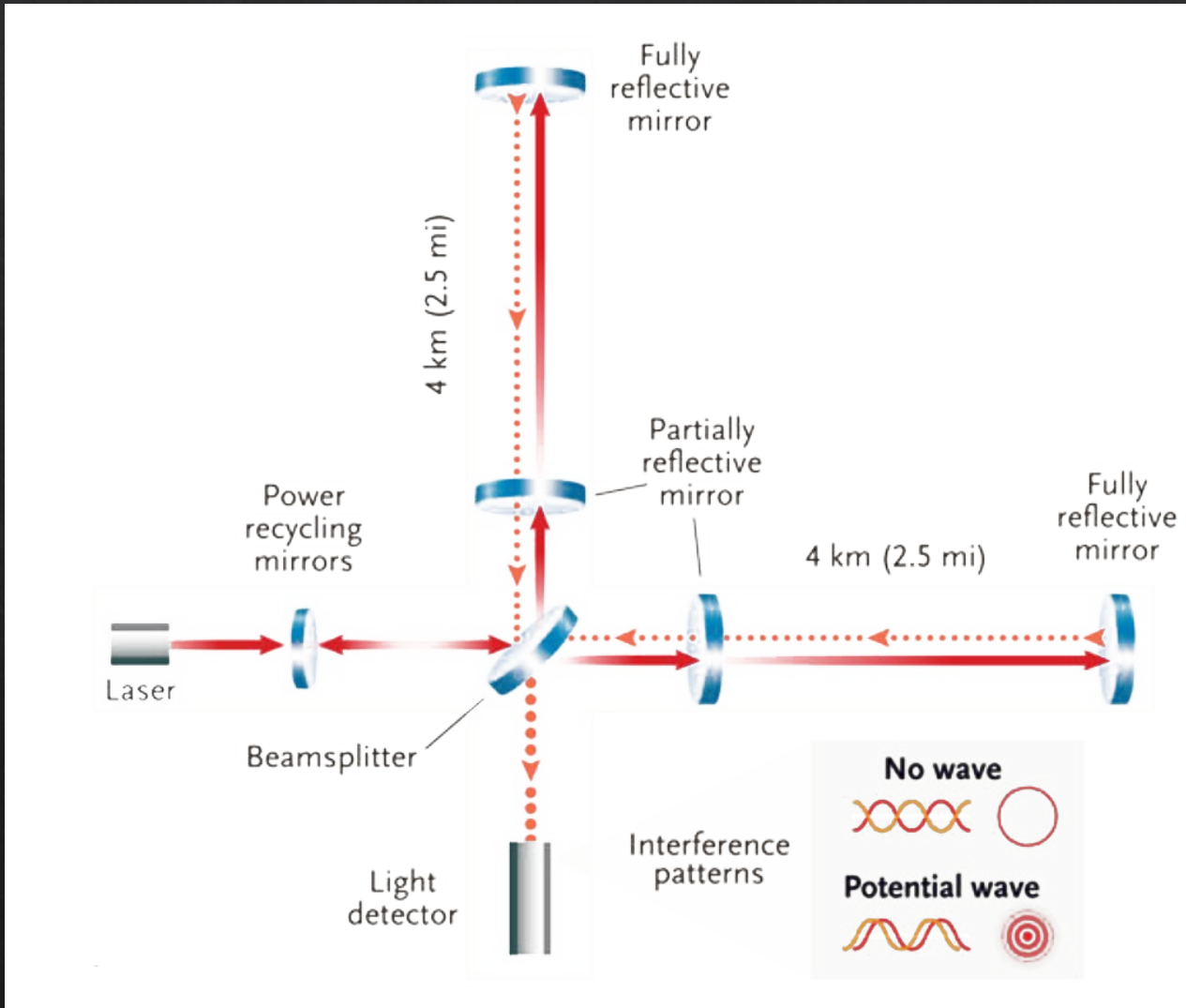


Adhikari, 2014



Gravitational Wave Observatories

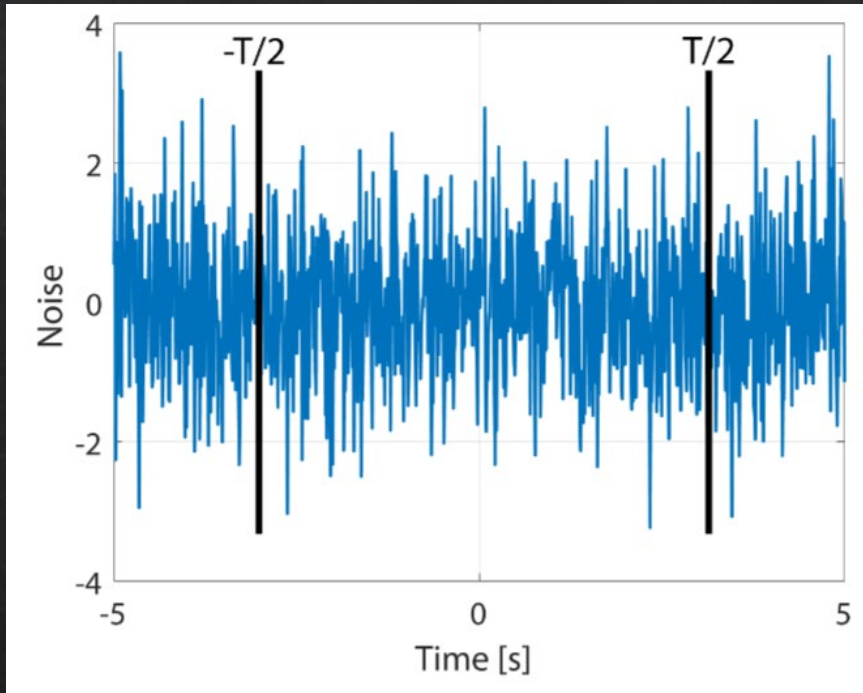
Principle of GW Detection



Gravitational waves change the distance between suspended test masses, which leaves an imprint on the phase of the laser beam.

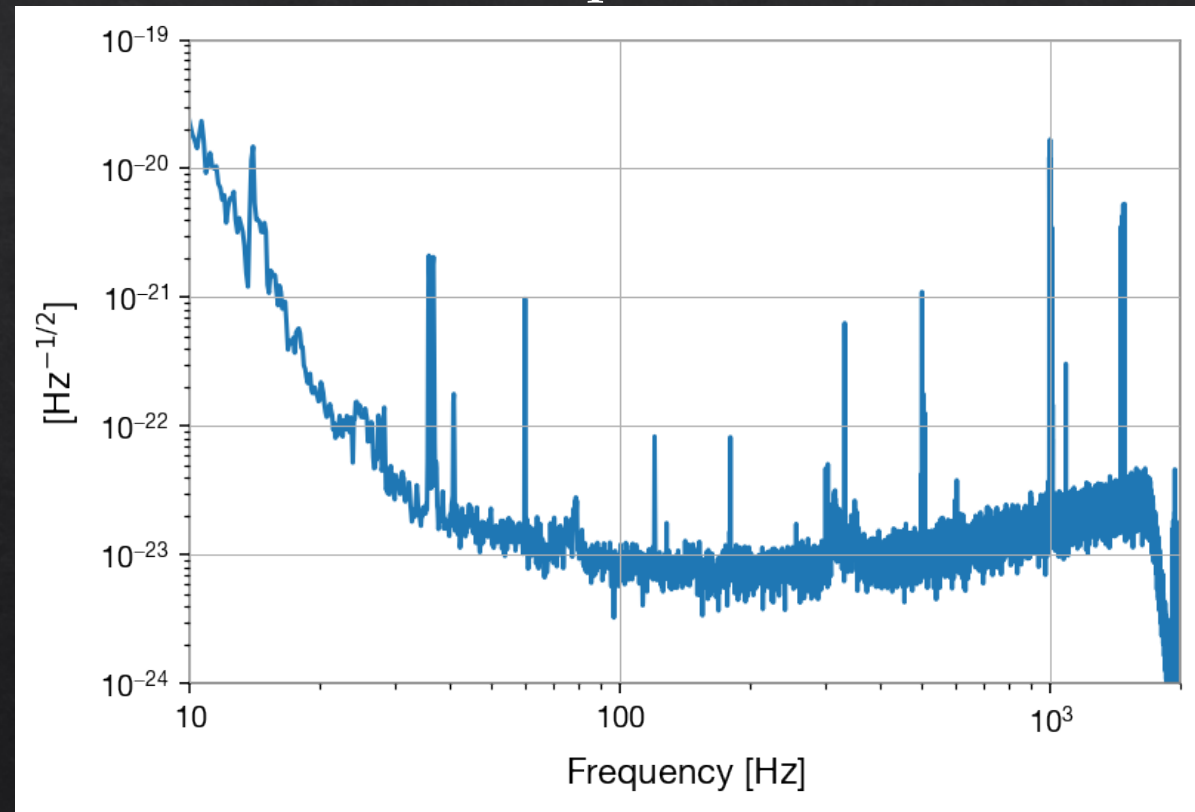
Power spectral density (PSD)

Time series



How to represent stationary noise in frequency domain?

Noise spectrum

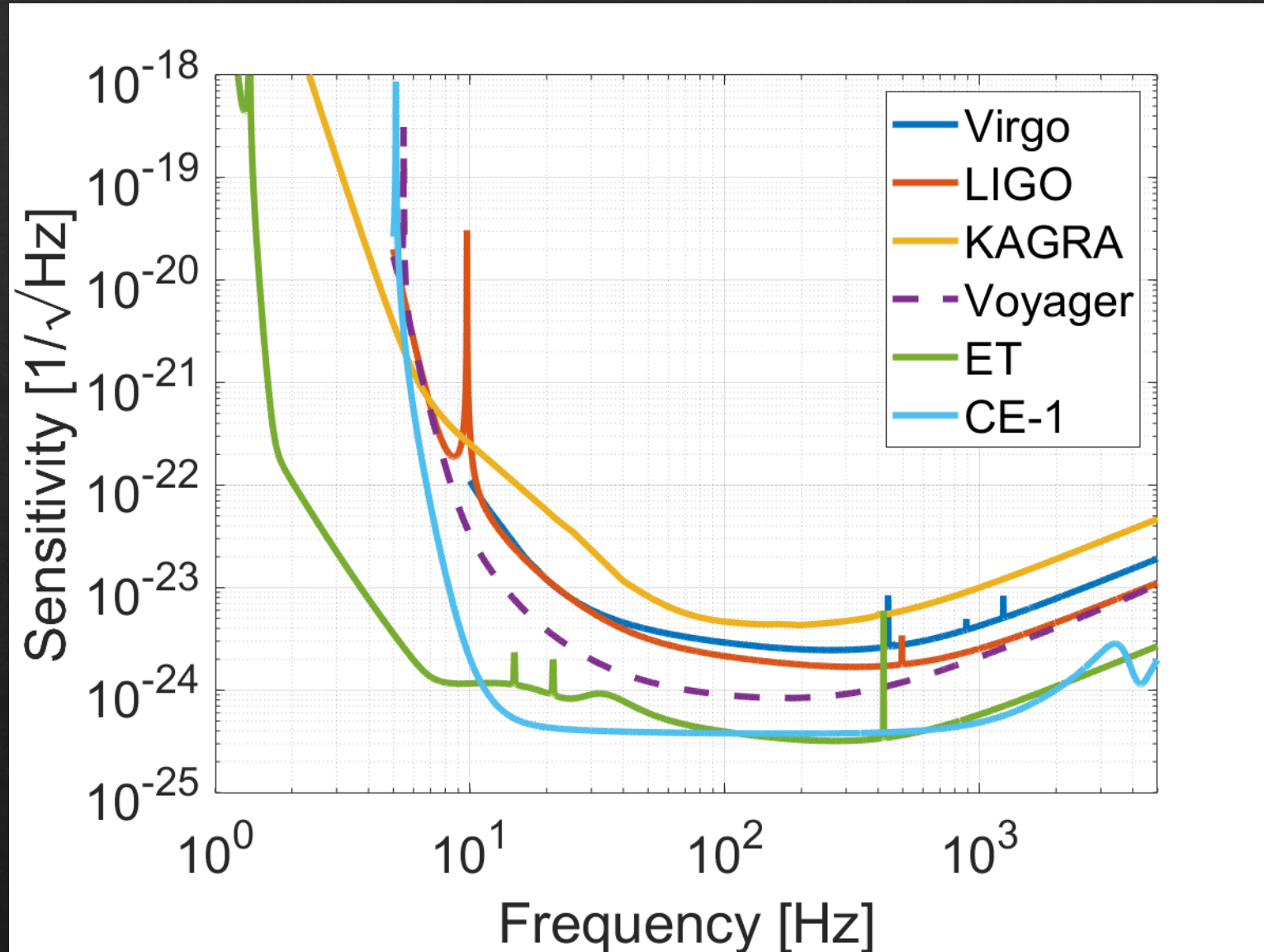


PSD

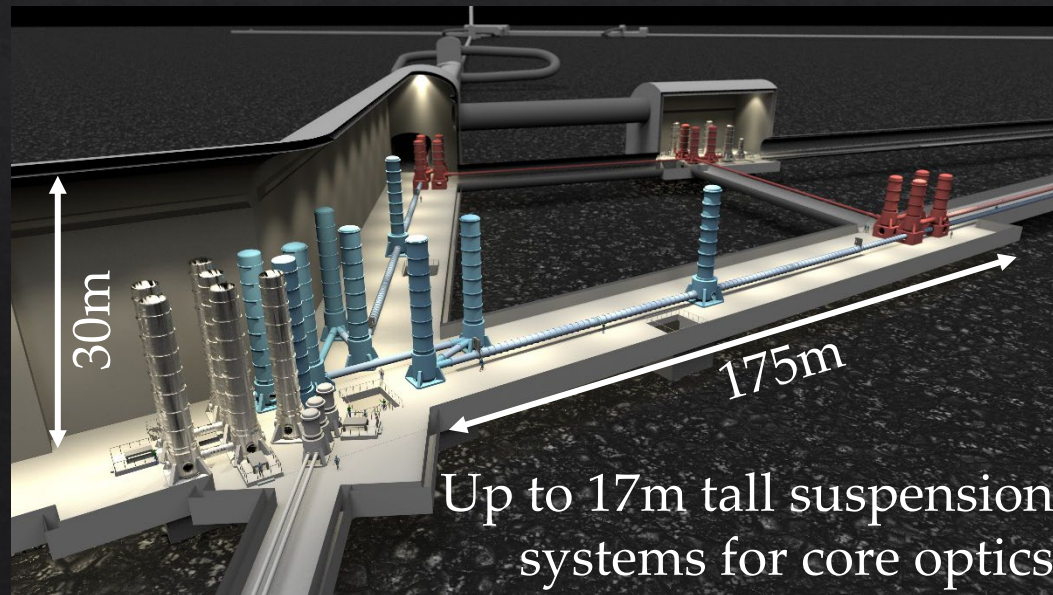
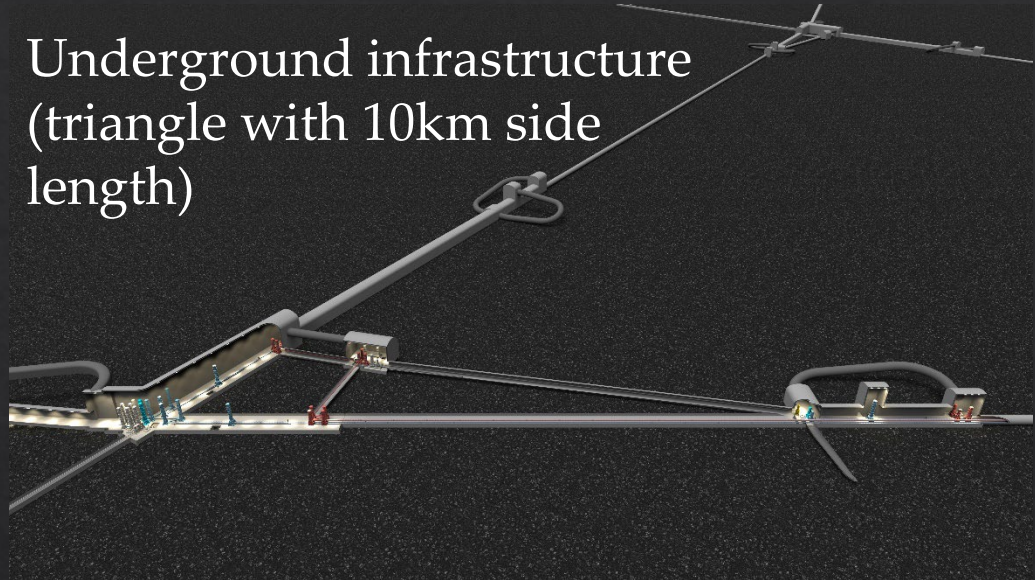
Finite-time Fourier transform

$$S(y; f) = \lim_{T \rightarrow \infty} \frac{2}{T} |\tilde{y}_T(f)|^2$$

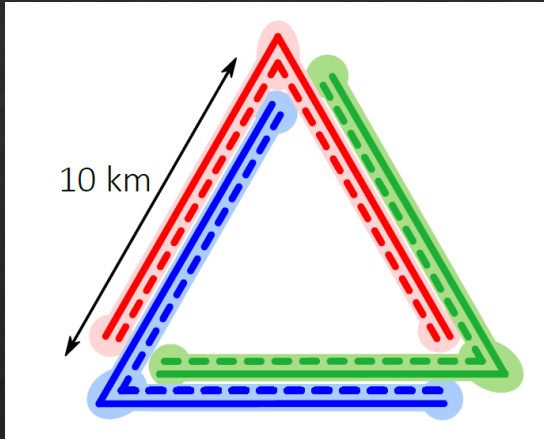
Sensitivity Models



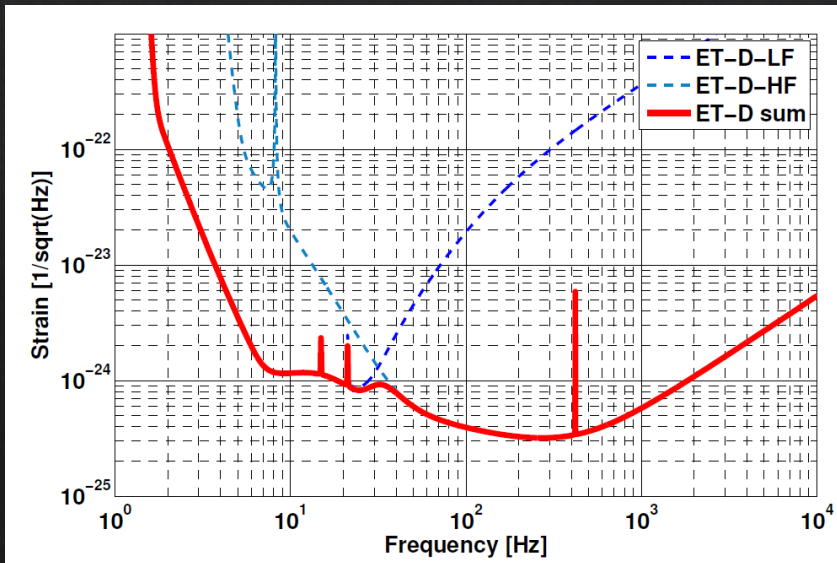
Einstein Telescope



Einstein Telescope as Xylophone

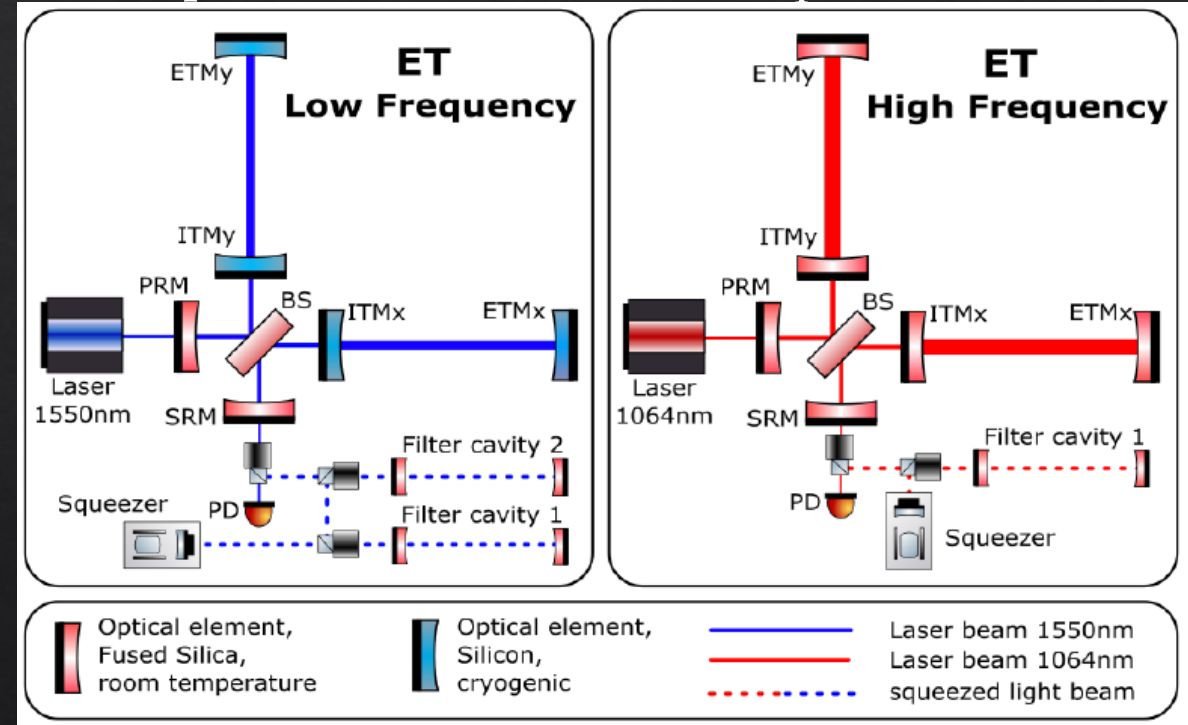


Each vertex is the center of a pair of interferometers, i.e., 6 interferometers in total.

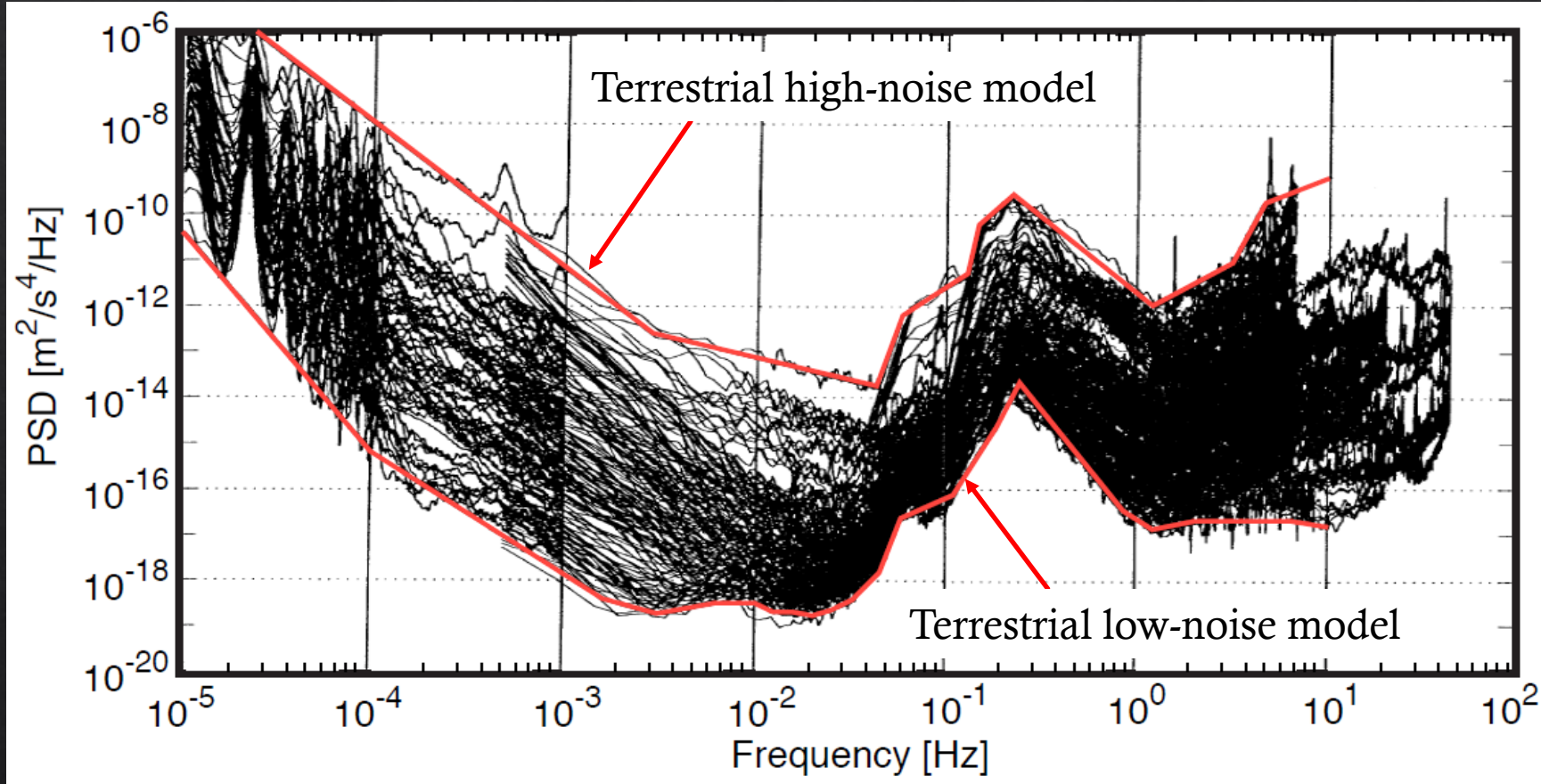


Power: 18kW
Temperature: 10-20K

Power: 3MW
Temperature: 290K



Seismic Background On Earth



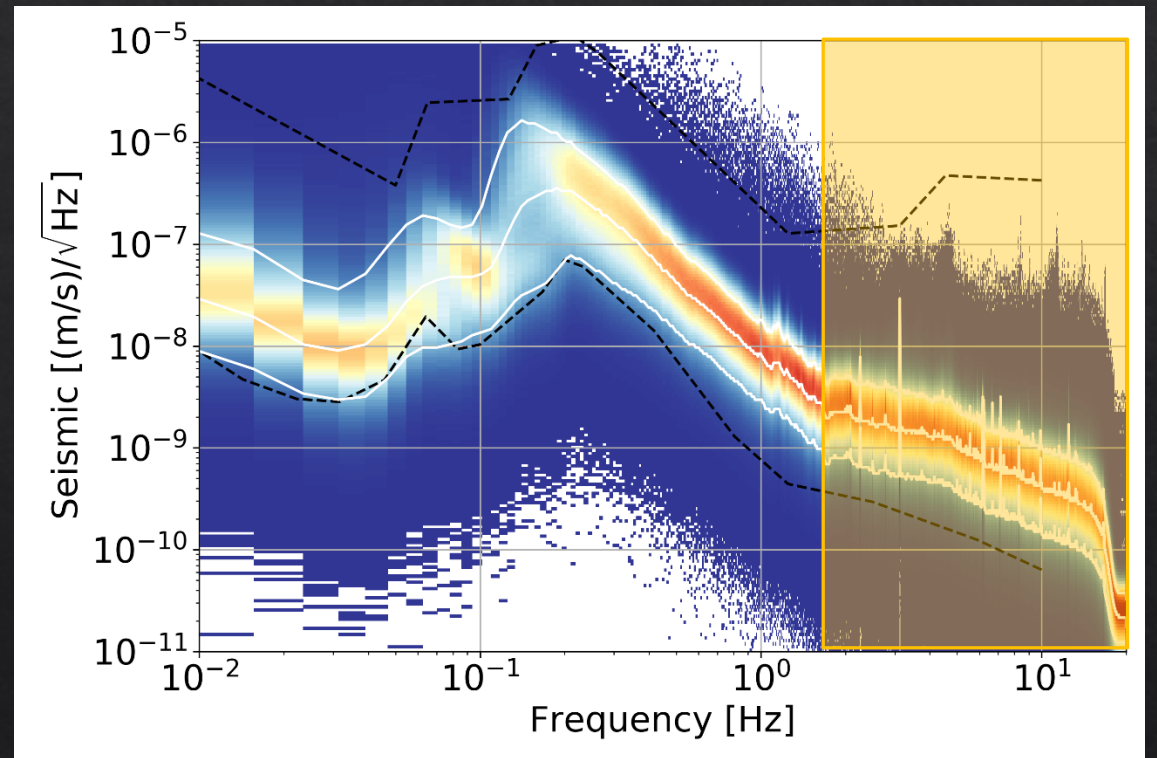
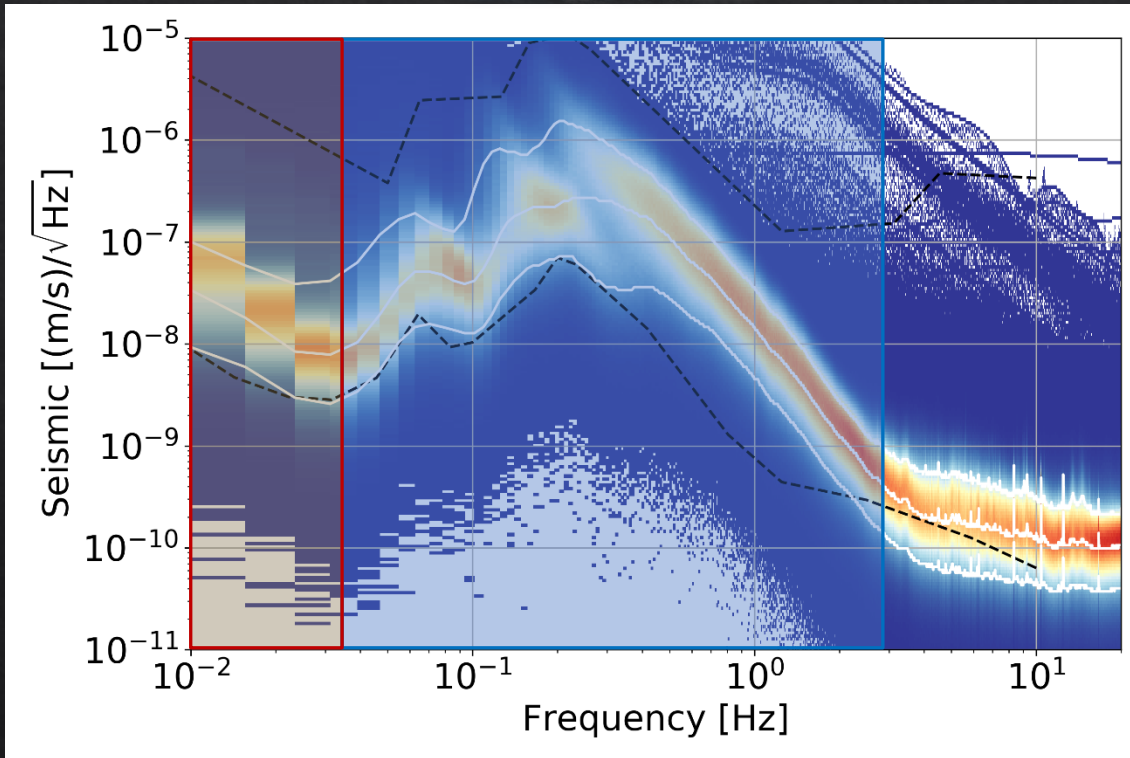
Seismic Noise at ET Candidate Sites



Sardinia



Euregio Meuse-Rhine



Tunnels and vacuum pipes



KAGRA

Vacuum pumps



Virgo

Ventilation



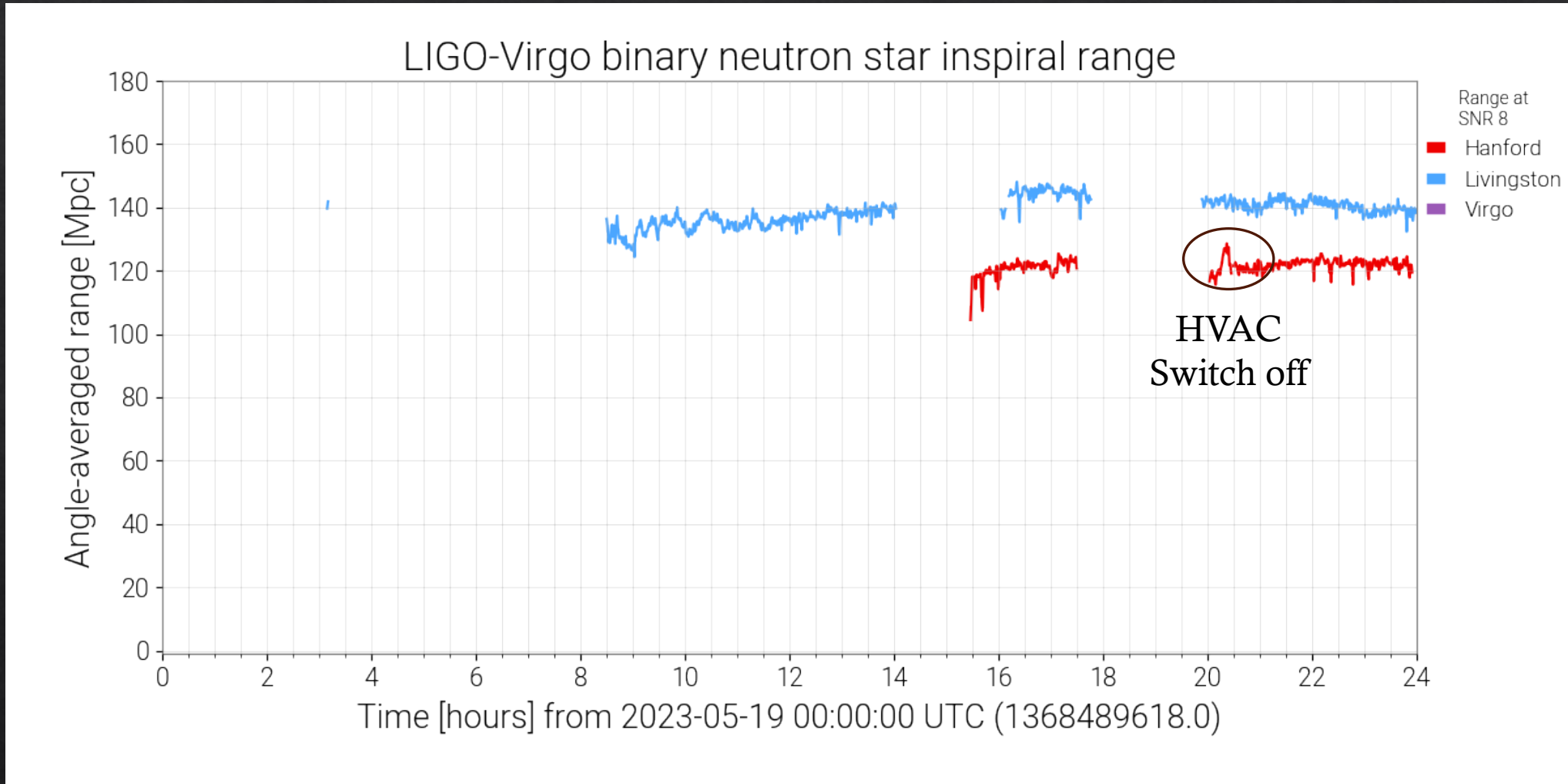
Sanford Underground Research Facility

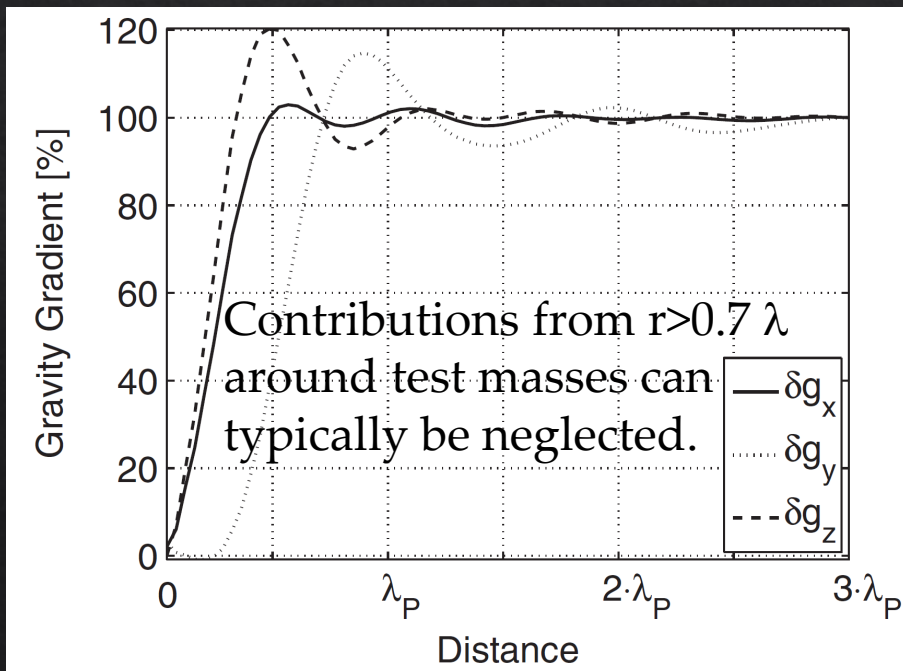
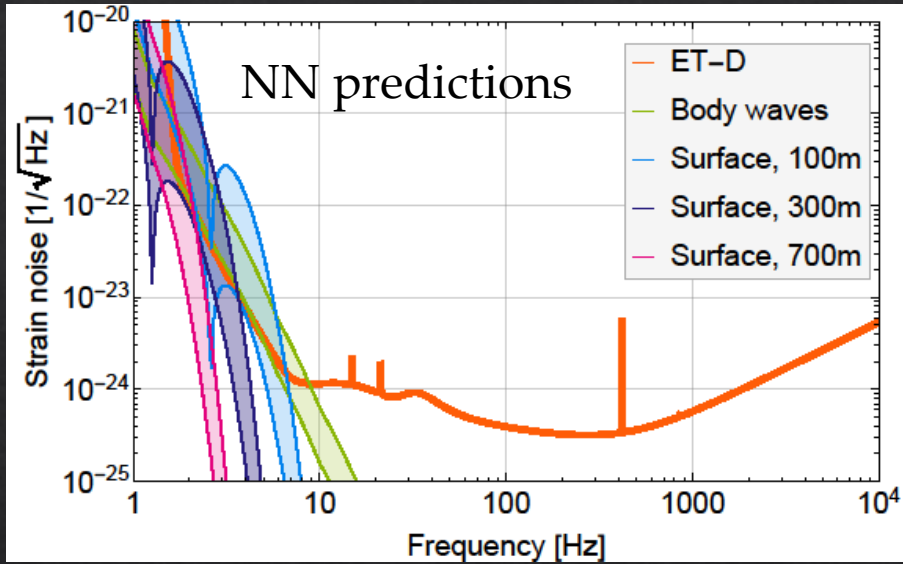
Cryocooler



KAGRA

HVAC: LIGO Hanford





1) Density perturbation by a seismic field

$$\delta\rho(\vec{r}, t) = -\nabla \cdot \left(\rho(\vec{r}) \vec{\xi}(\vec{r}, t) \right)$$

2) Associated gravity perturbation

$$\delta\phi(\vec{r}_0, t) = G \int dV \frac{\nabla \cdot \left(\rho(\vec{r}) \vec{\xi}(\vec{r}, t) \right)}{|\vec{r} - \vec{r}_0|}$$

3) Solve for a specific seismic field

$$\delta\vec{a}(\vec{r}_0, t) = \frac{4\pi G\rho}{3} \left(2\vec{\xi}^P(\vec{r}_0, t) - \vec{\xi}^S(\vec{r}_0, t) \right)$$

Newtonian-noise Cancellation

Wiener filter

Correlations
between all inputs y

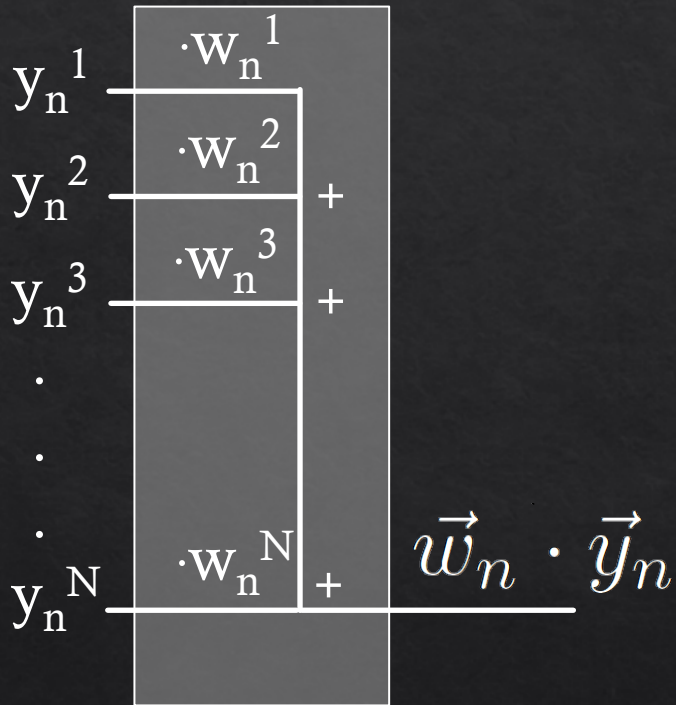
Linear discrete filter

$$\vec{w}_n = \langle x_n \vec{y}_n^\dagger \rangle \cdot \langle \vec{y}_n \vec{y}_n^\dagger \rangle^{-1}$$

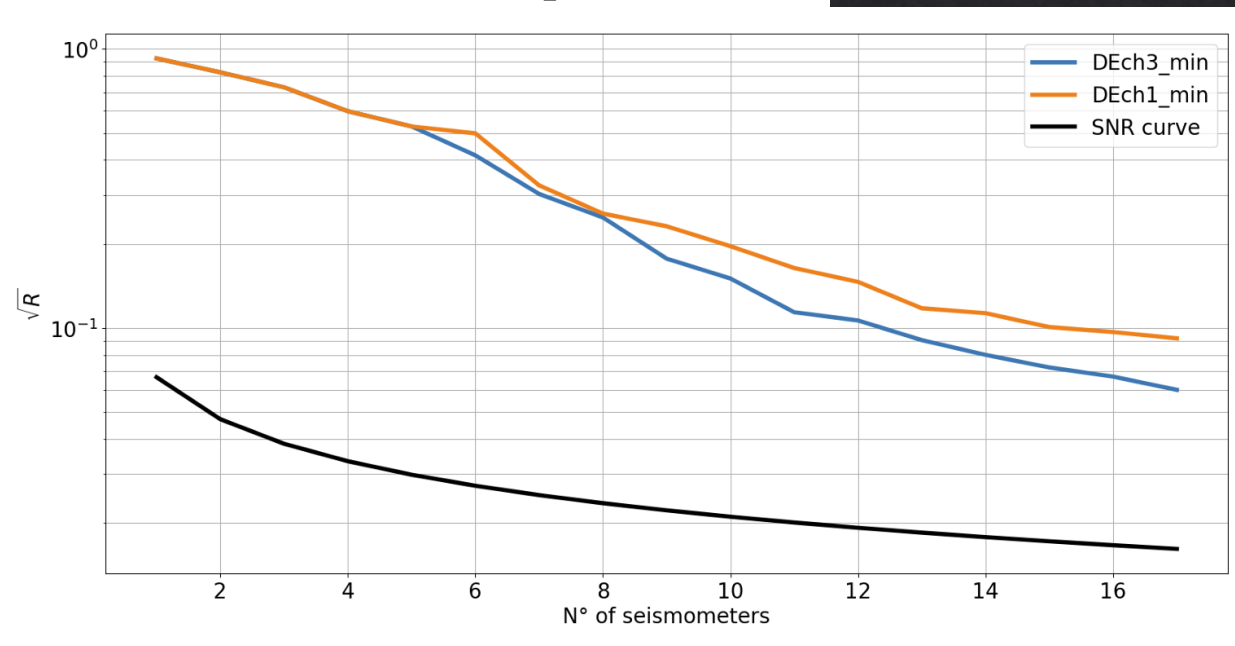
Correlations between
inputs y and target x

Residual noise power
after subtraction

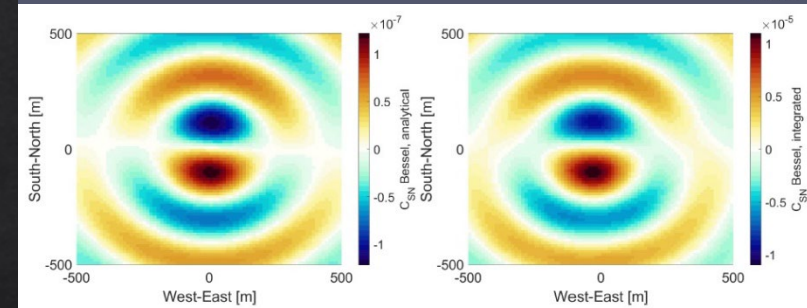
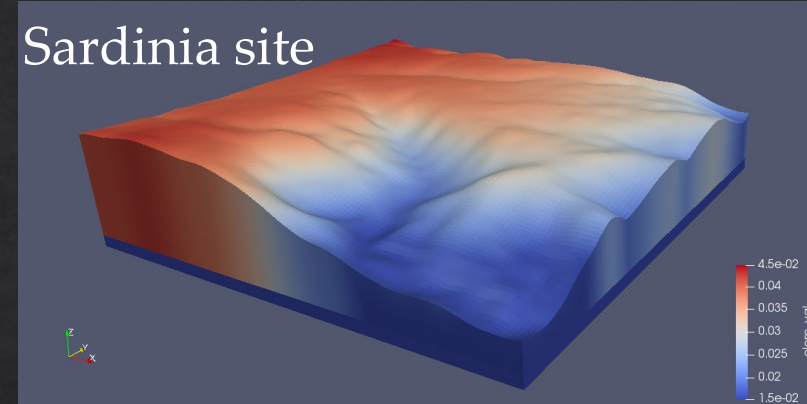
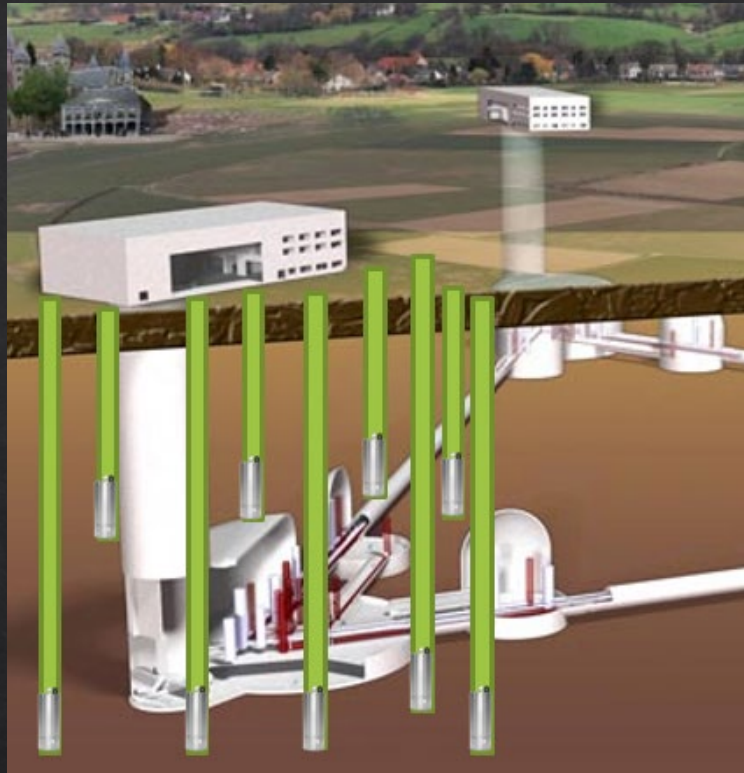
$$\langle |\vec{w}_n \cdot \vec{y}_n - x_n|^2 \rangle$$



Assuming optimal sensor placement:
10 borehole seismometers per test mass

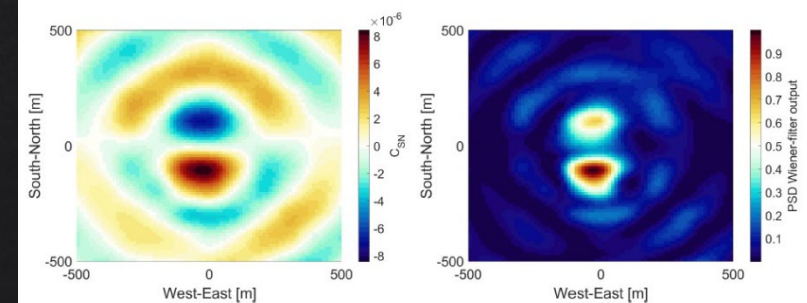


Simulations + data for optimization



(a) Flat-surface, isotropic.

(b) Topographic kernel, isotropic.

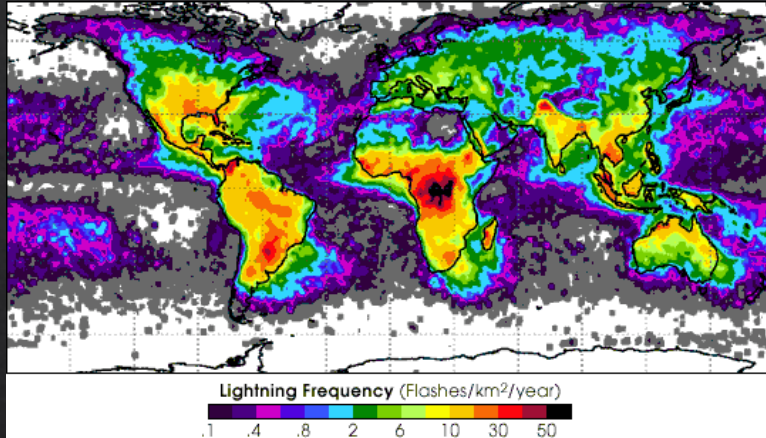


(c) SPEC-FEM3D simulation.

(d) PSD Wiener filter.

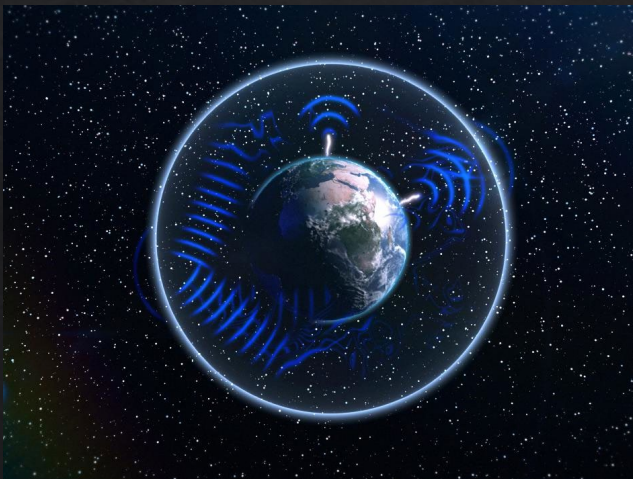
Open questions:

- 1) Where to place the sensors and how many?
- 2) How to design the noise-cancellation filter?
- 3) What types of seismic sensors should be used?

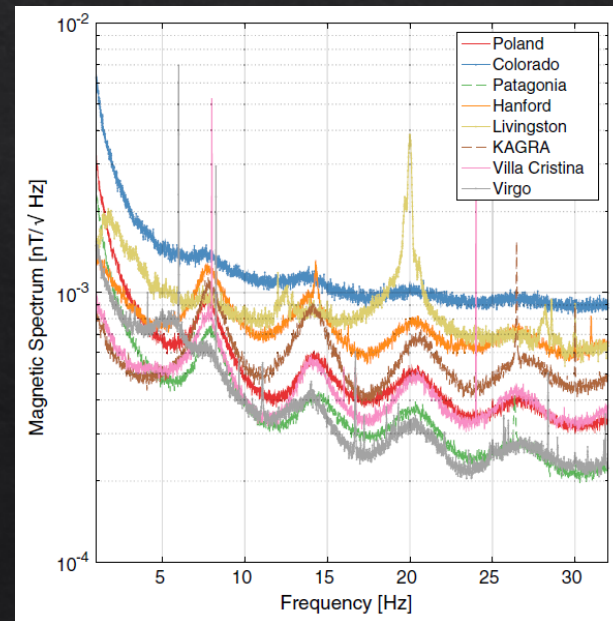


Magnetic disturbances can appear coherently in a global detector network.

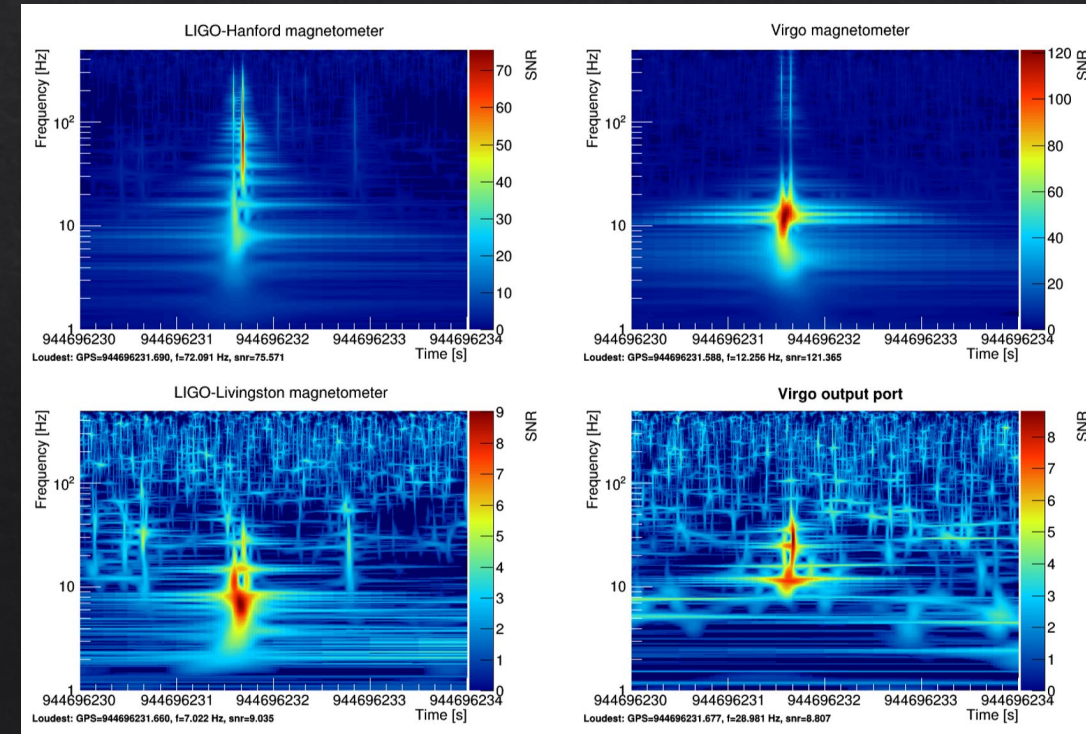
Globally coherent disturbances and noise



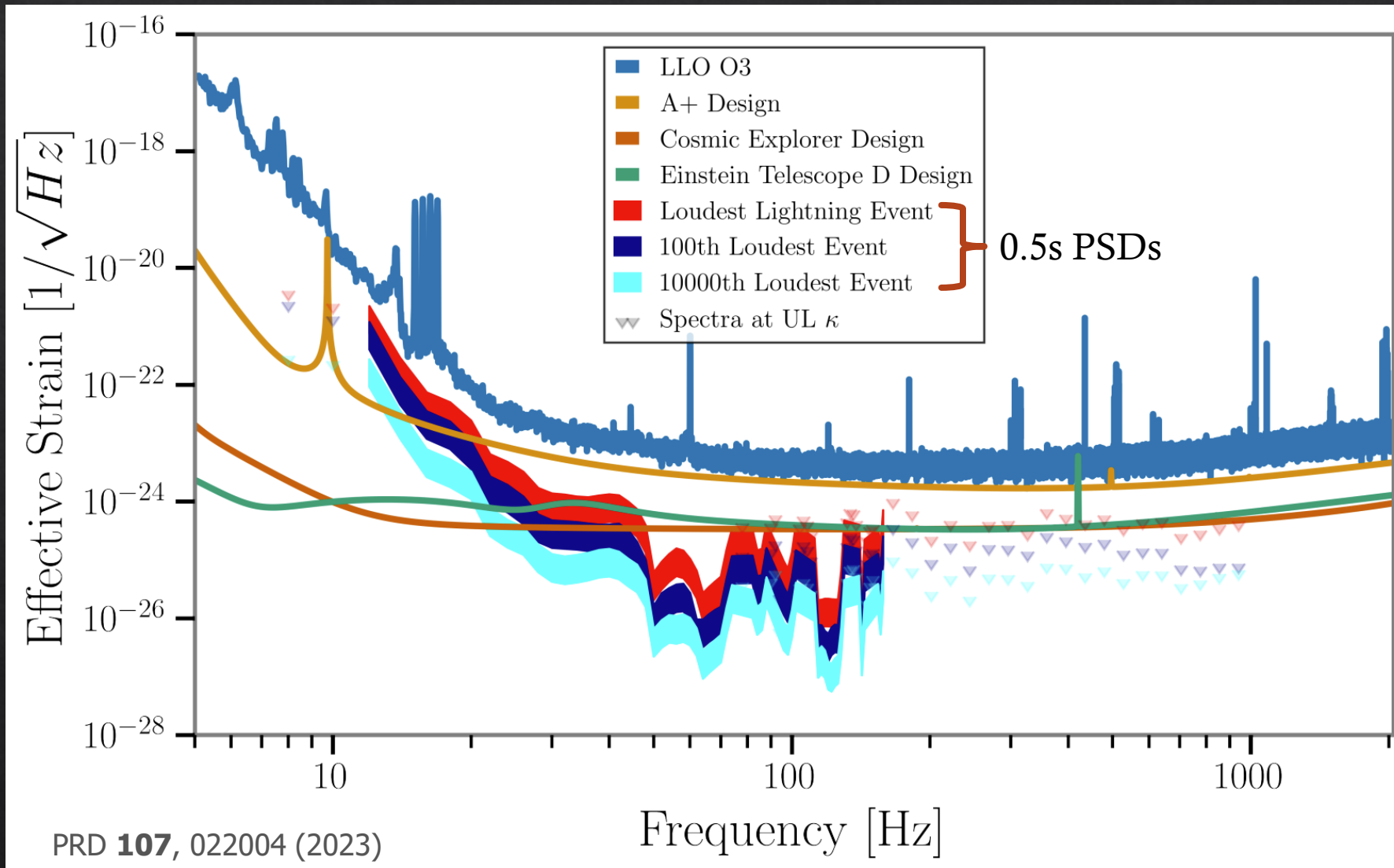
Schumann resonances



Lightning transient can be coincident



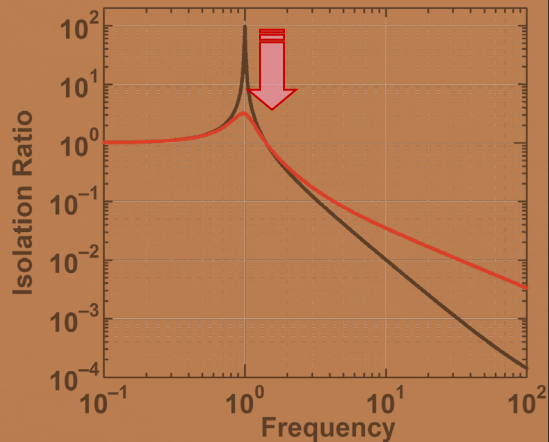
Magnetic Coupling



Virgo Superattenuator

Damping

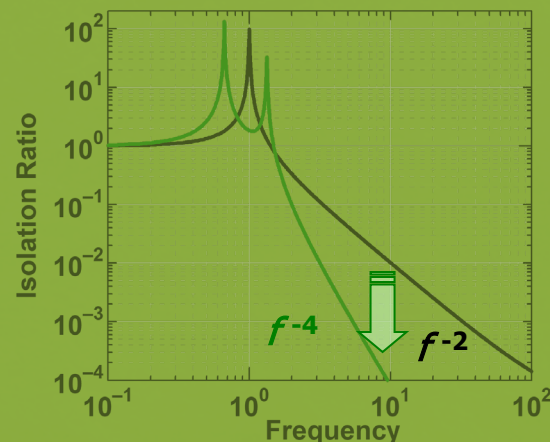
Lower peak height



Less isolation

Cascaded

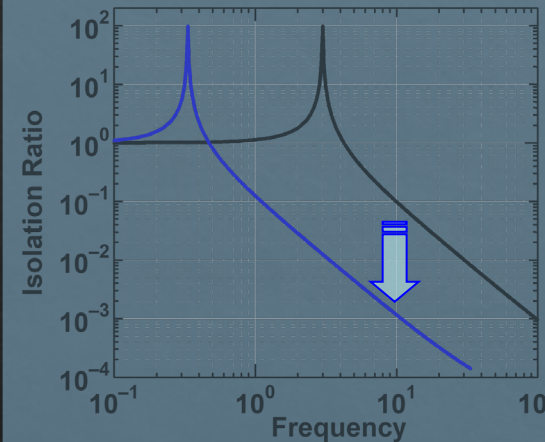
Steeper isolation curve



More peaks

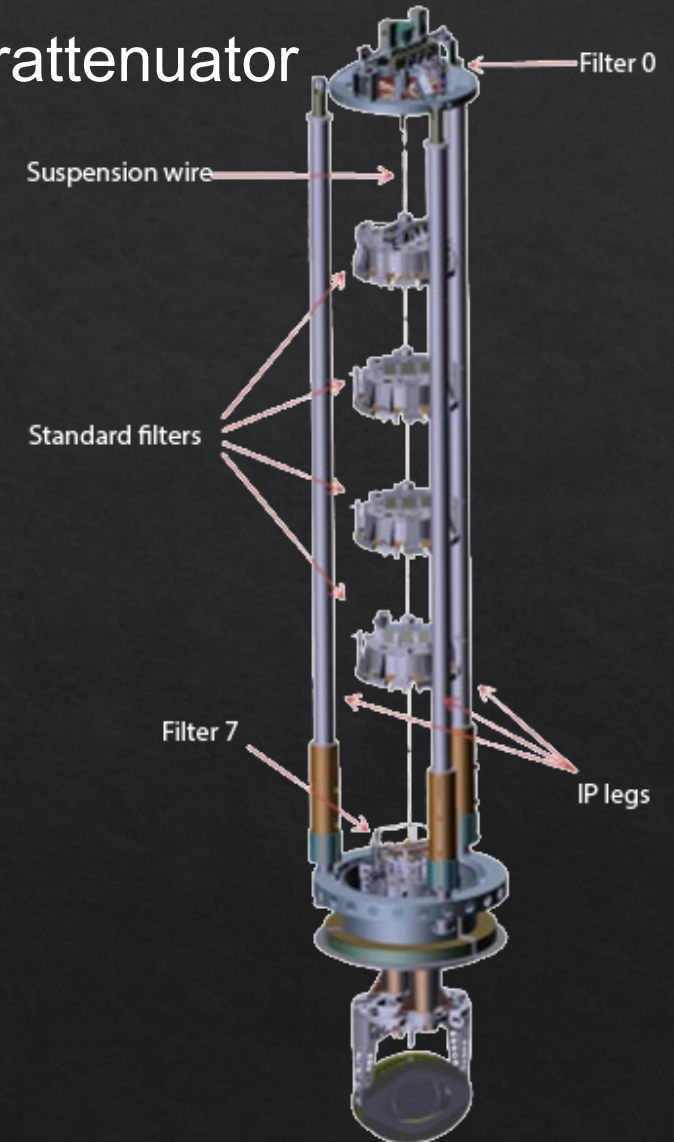
Larger structure

Lower resonance frequency

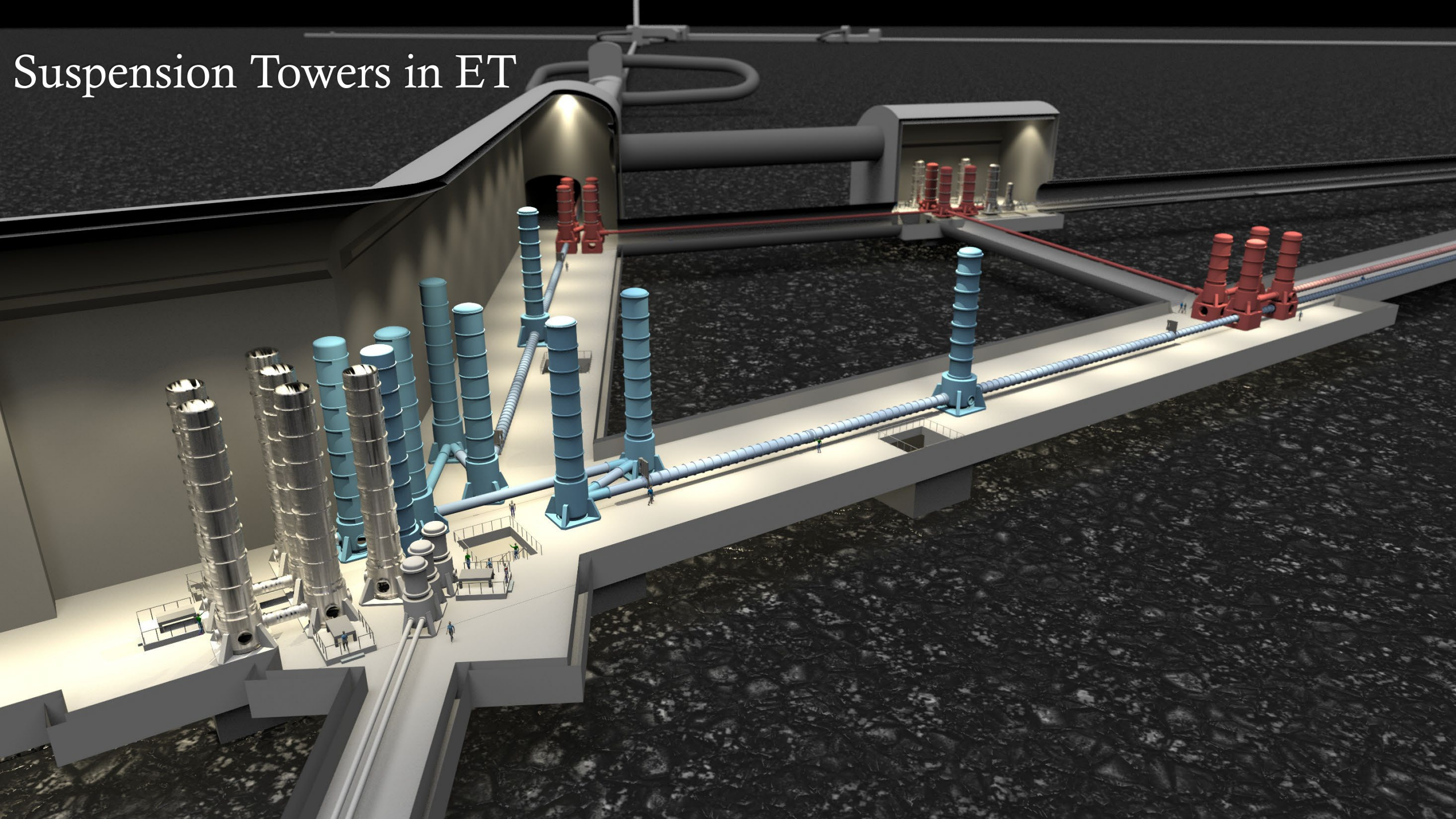


Difficult to realize

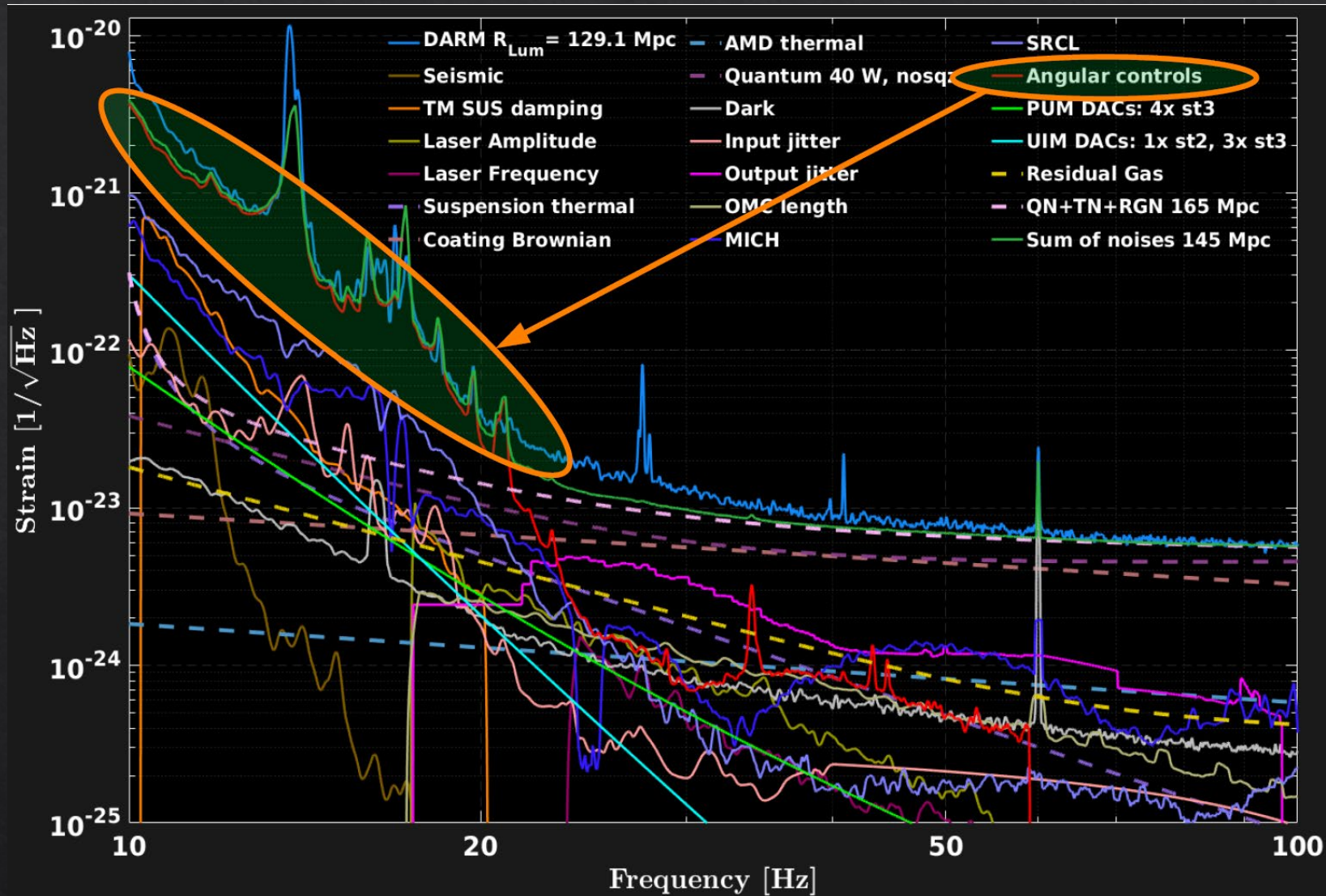
In practice: use combination of these methods



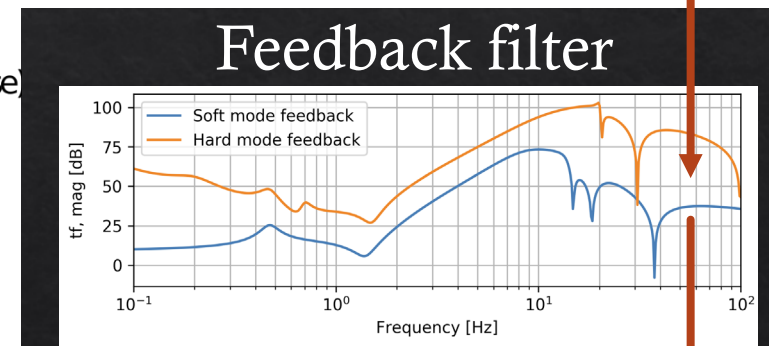
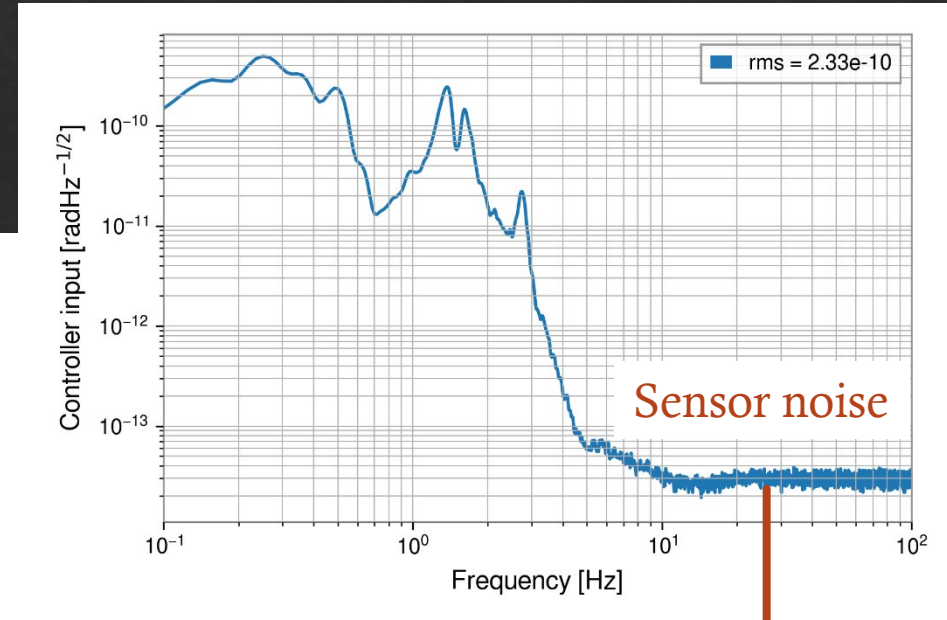
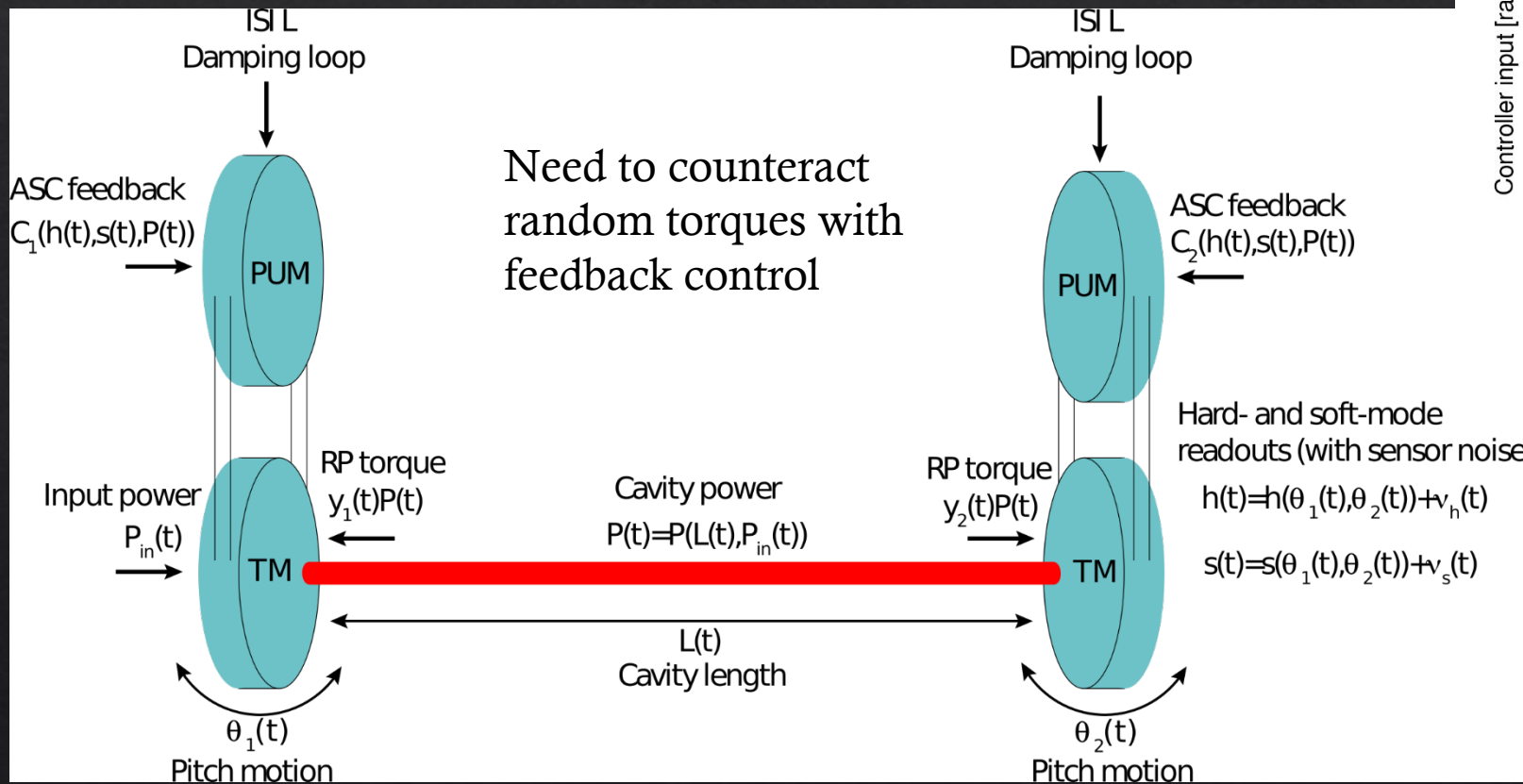
Suspension Towers in ET



LIGO noise budget



Torques continuously act on test masses:
Radiation pressure, noise from suspensions

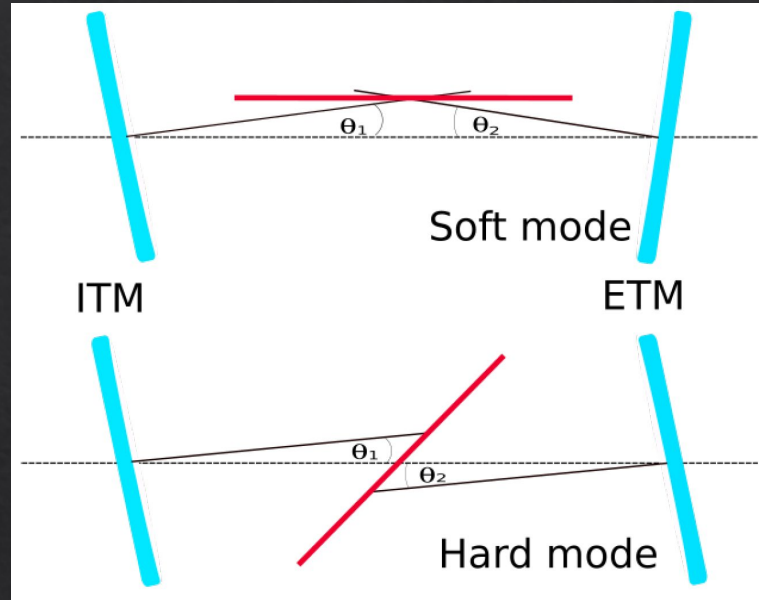
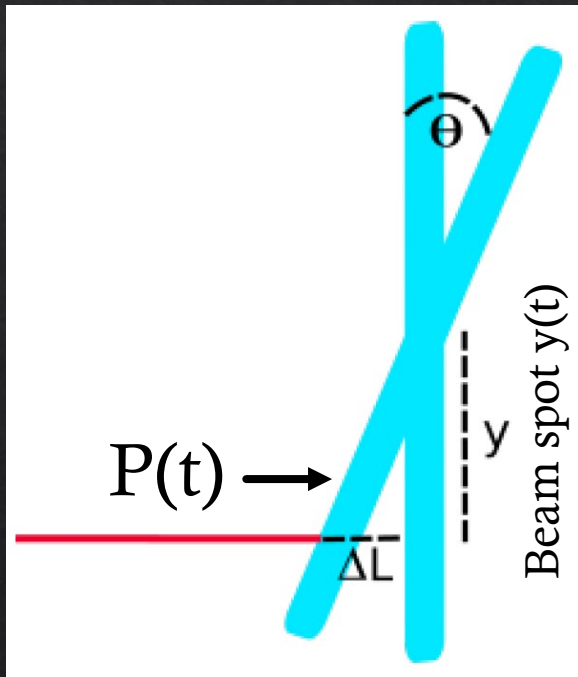


Back to suspensions and test masses

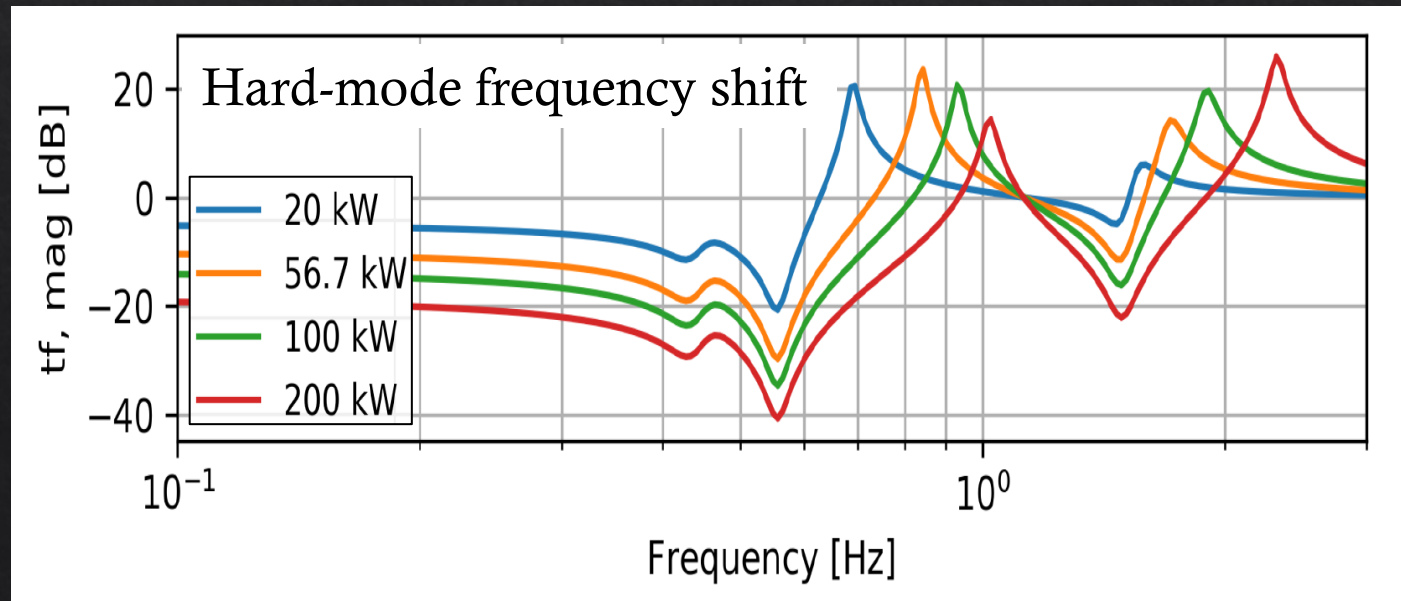
Angular Control: ET Strategy

Radiation-pressure torque

$$\tau_{RP}(t) = \frac{2P(t)}{c}y(t)$$



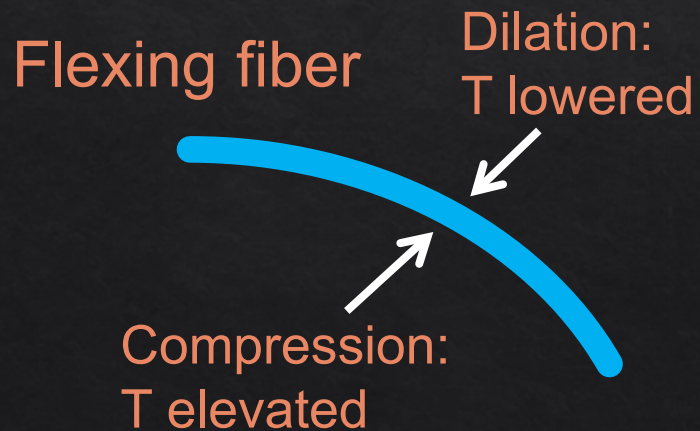
Optomechanical system



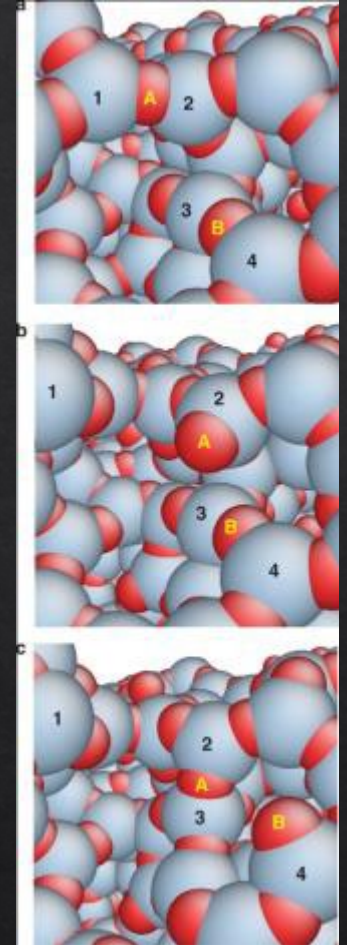
ET-LF: low power, high mass

Every dissipation mechanism leads to a coupling of the system to a thermal reservoir and thermal noise.

Thermo-elastic noise:
Irreversible heat flux across temperature gradients



Brownian noise:
Many possible causes (for example, change in atomic bonds)

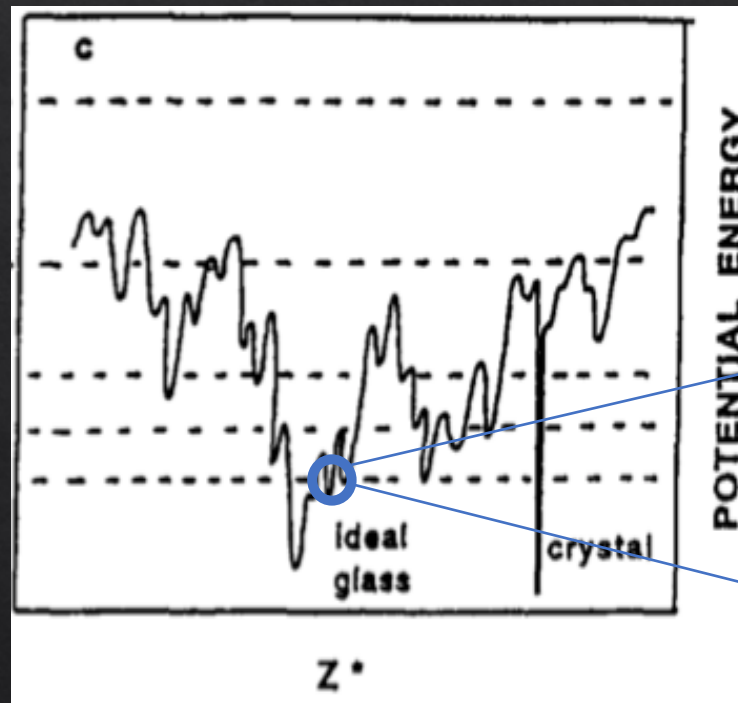


Fluctuation-dissipation theorem:
Thermal-noise spectrum proportional to mechanically dissipated power

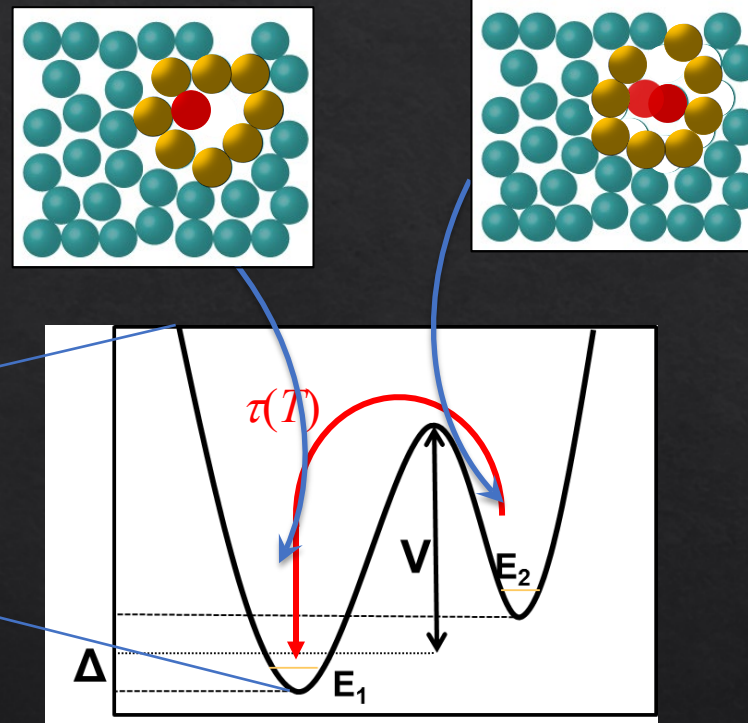
$$S_x(\Omega) = \frac{8\pi kT}{\Omega^2} \frac{W_{\text{diss}}}{F_p^2}$$

$$W_{\text{diss}} \propto \frac{1}{Q}$$

“nearby” minima lead to tunneling or thermally-activated motion of groups of atoms



Angell, 1997

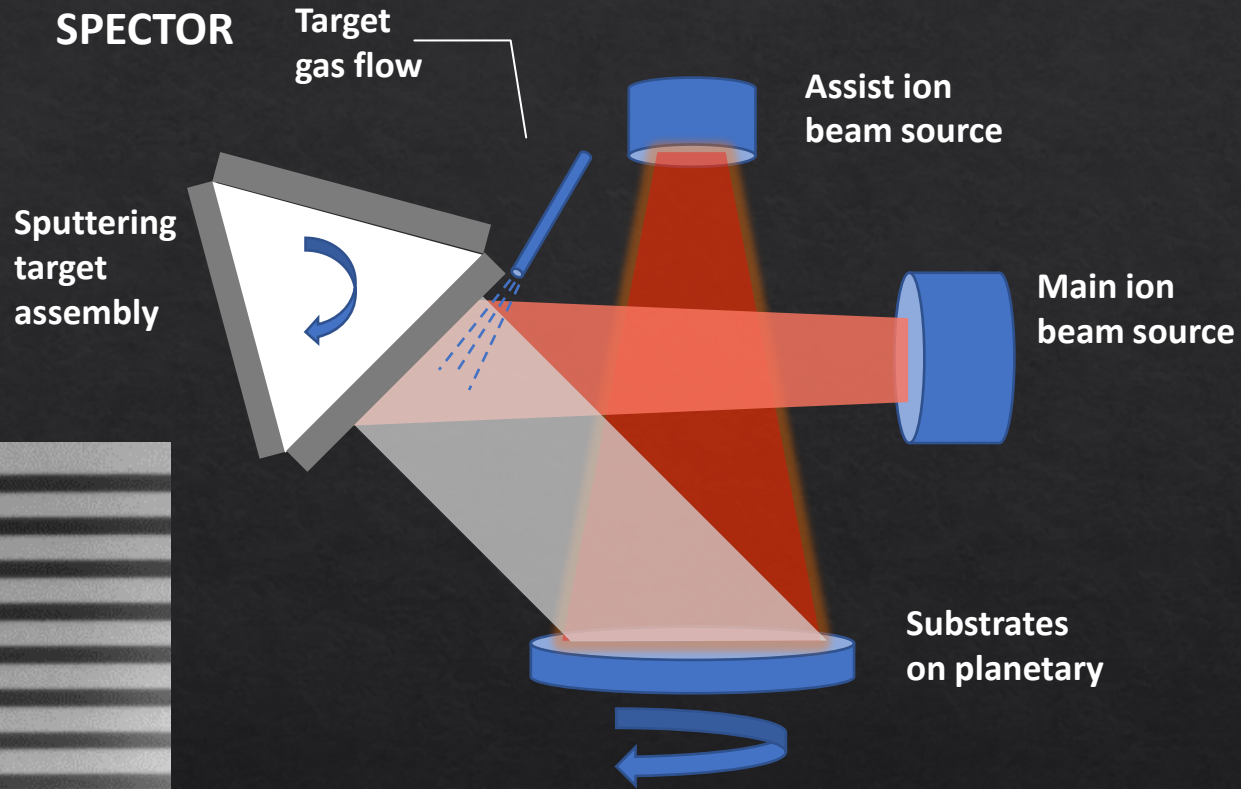
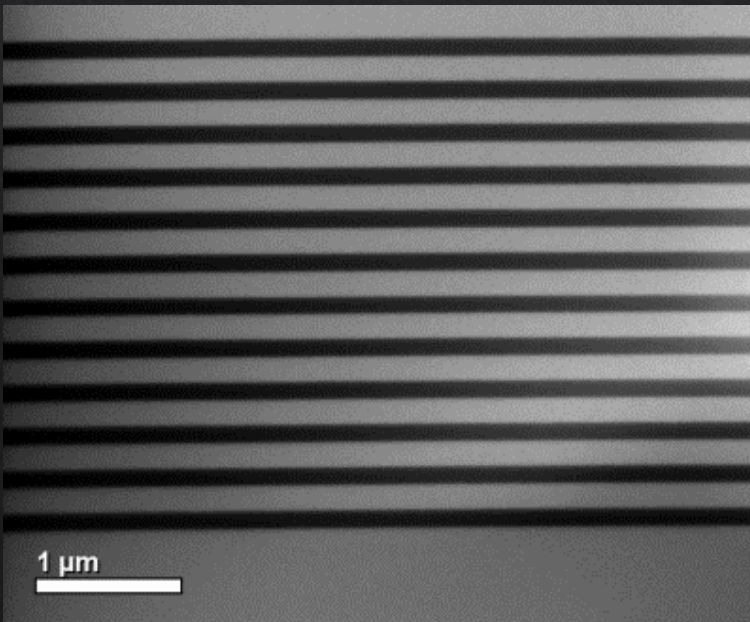


Two-level systems from neighboring energy minima in structural landscape:

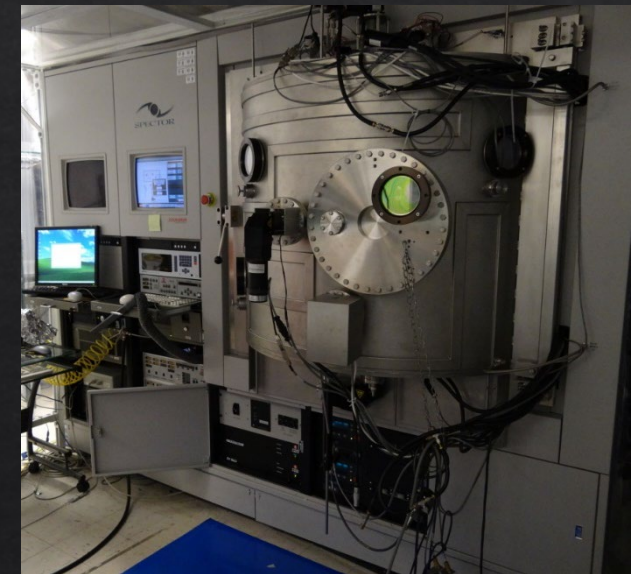
- At low T , atomic structure **tunnels** between the $\gtrsim \mu\text{eV}$ energy splitting $E_{1,2} \pm \Delta$
- At higher T , atomic motion is **thermally activated**, requiring $k_B T \sim$ barrier height V

Coating Deposition

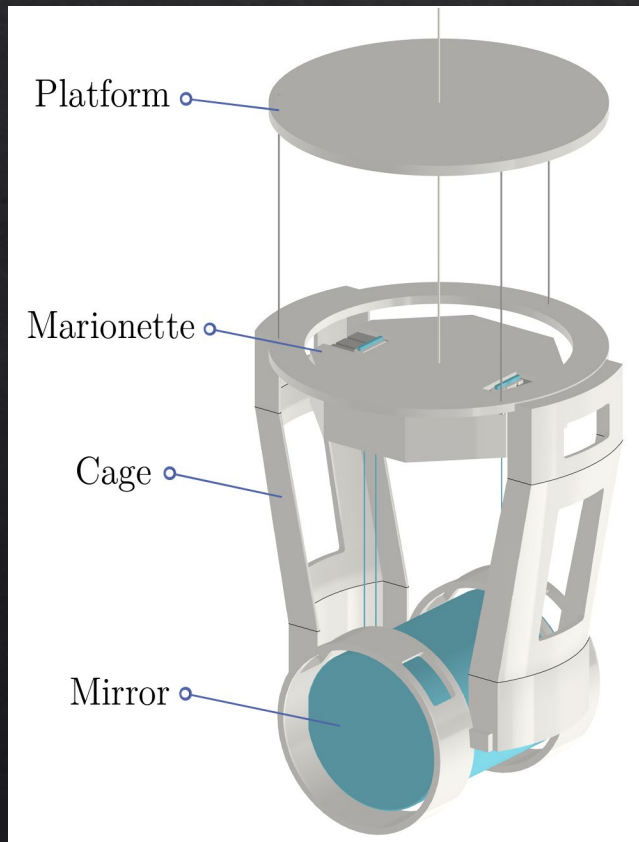
Dielectric stacks grown by DIBS



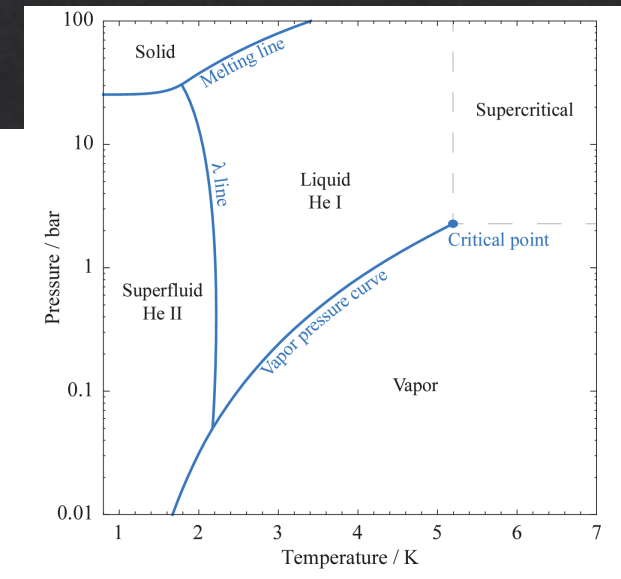
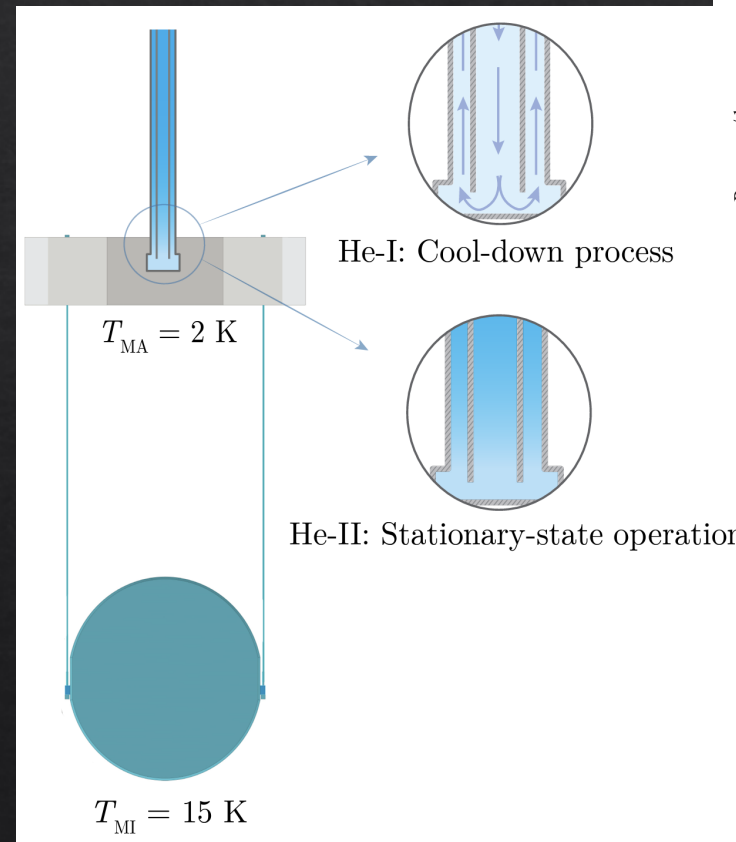
Dual ion-beam sputtering system



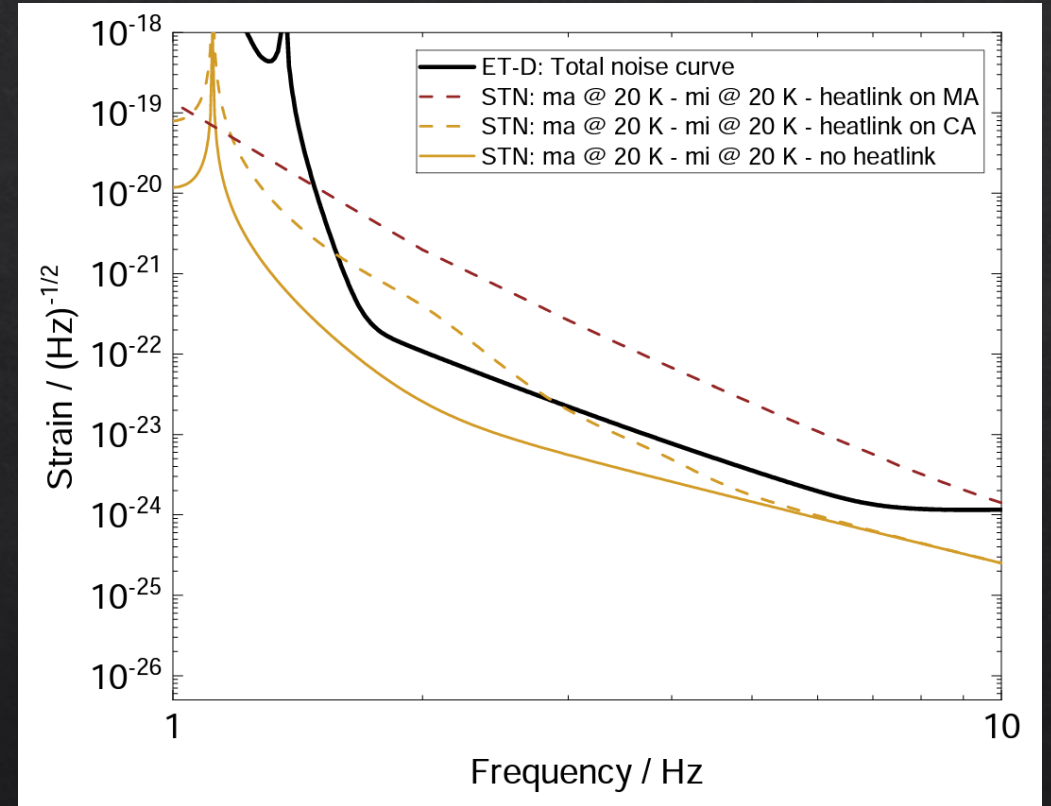
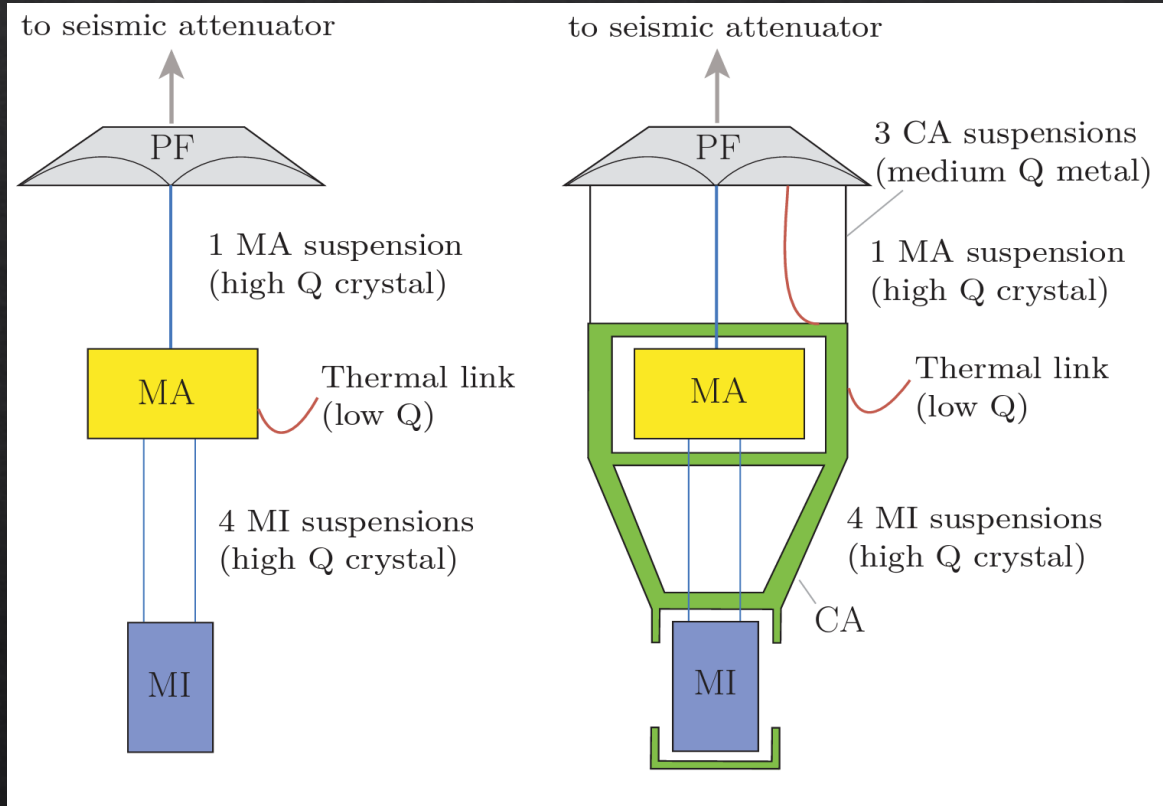
Cooling the ET-LF test masses is one of the biggest technological challenges of ET.



Conductive tube for initial He-I cool down, and stationary heat-transport with He-II is under investigation.



Suspension Thermal Noise



Quantum Noise



Heisenberg uncertainty principle

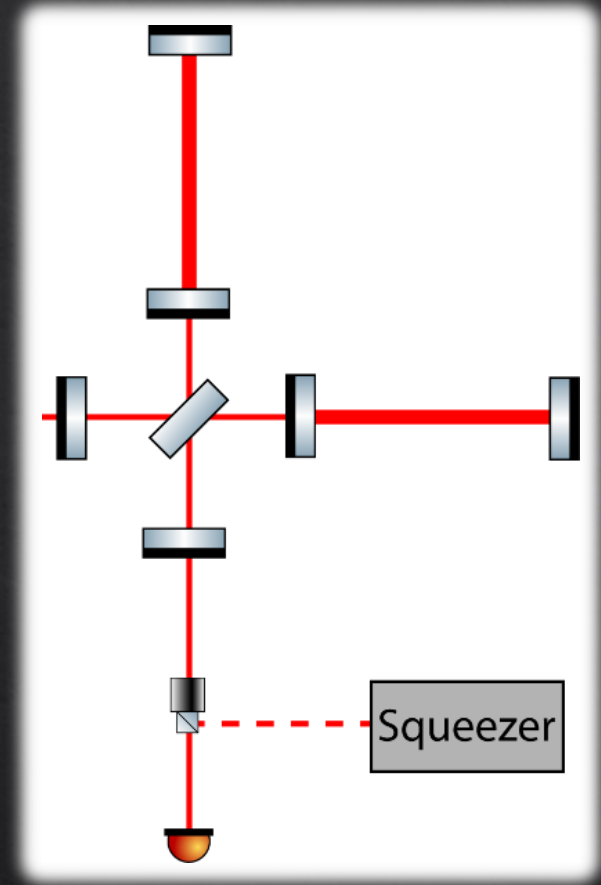
$$\Delta p \Delta x \geq \frac{\hbar}{2}$$

What are the position and momentum variables in the case of light?

Multiple answers, but for GW detectors, the conjugate variables are the quadratures of the EM field:

$$E(t) = E_1(t)\cos(\omega_0 t) + E_2(t)\sin(\omega_0 t)$$

Fundamental measurement in ET: Counting photons

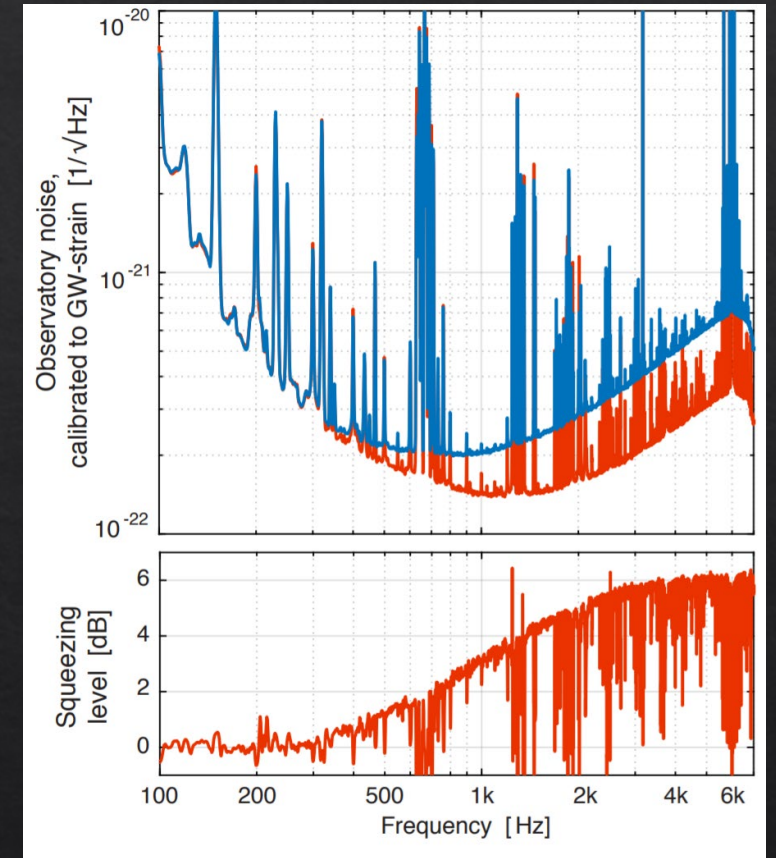
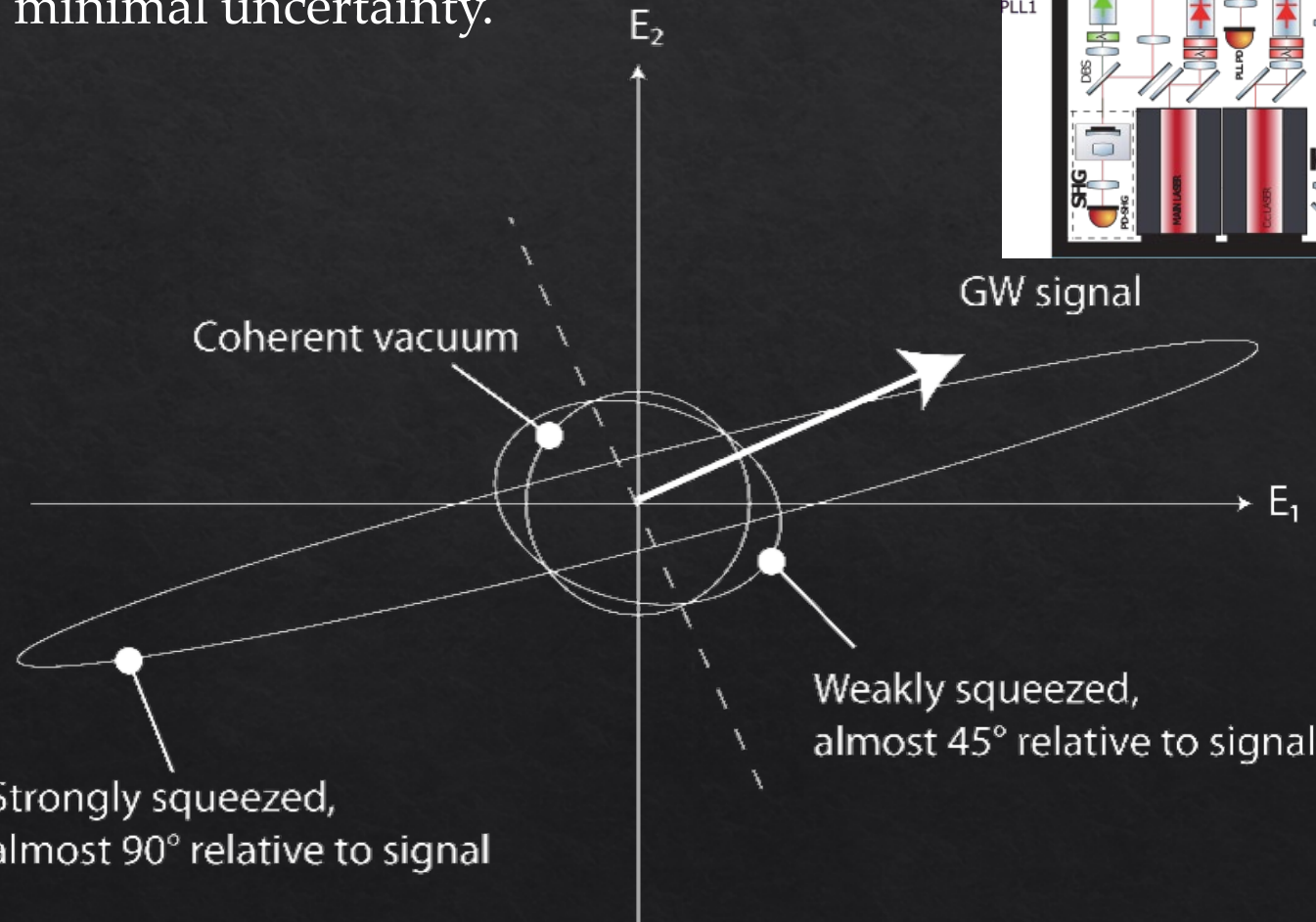
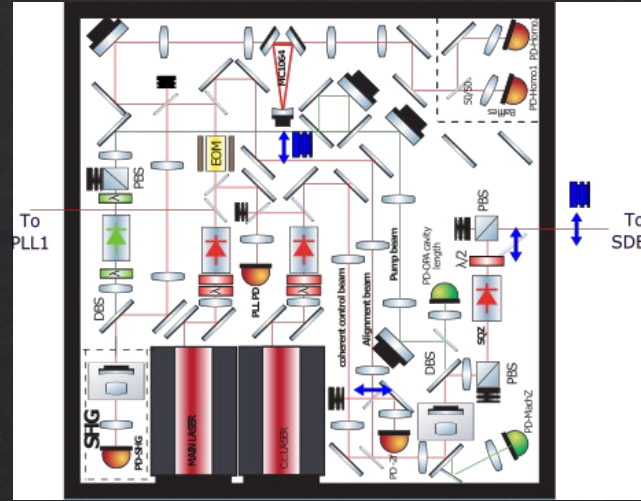


Caves: manipulate quantum state at the dark port.

Squeezing

Squeezer developed at AEI Hannover

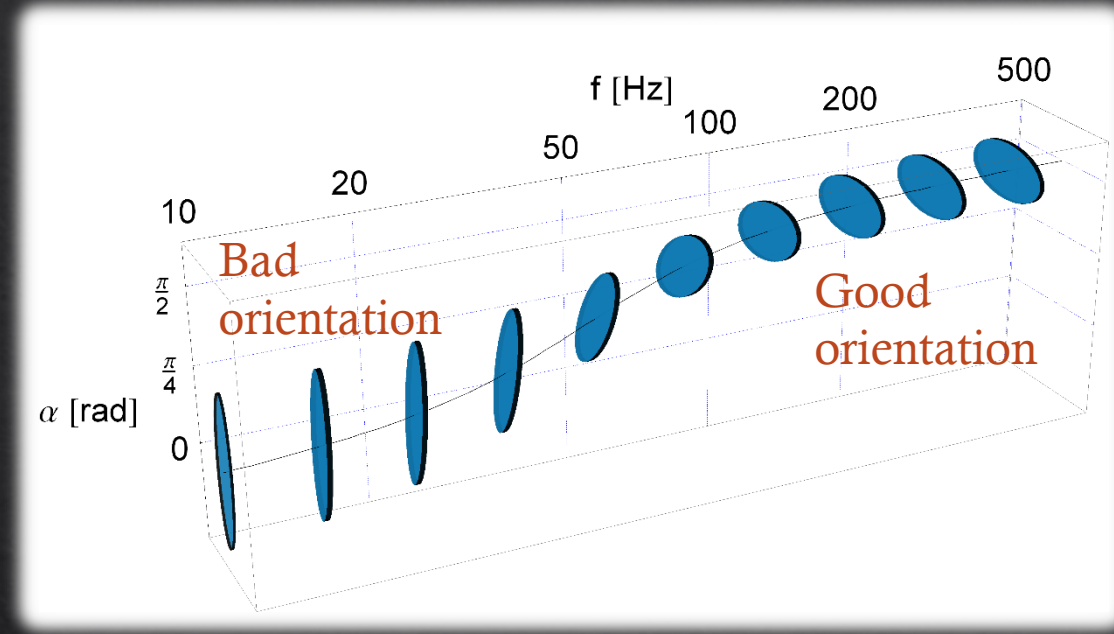
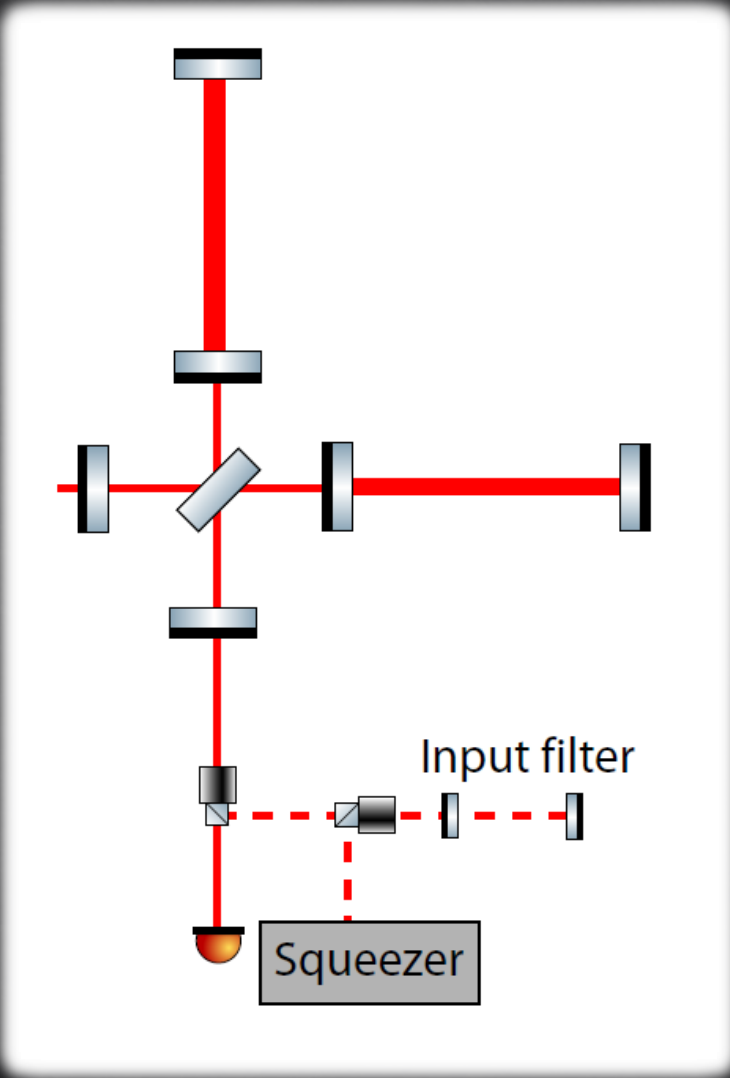
Coherent states and squeezed states both have minimal uncertainty.



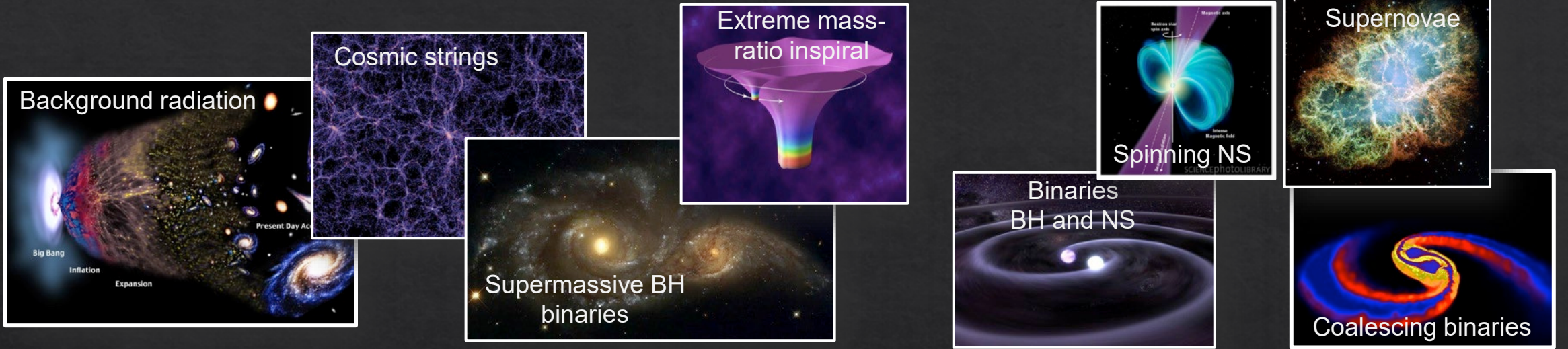
Lough et al 2020

Quantum Filter

Squeezing ellipse gets rotated by the laser interferometer.



A filter is needed to compensate the rotation of the squeezing ellipse.

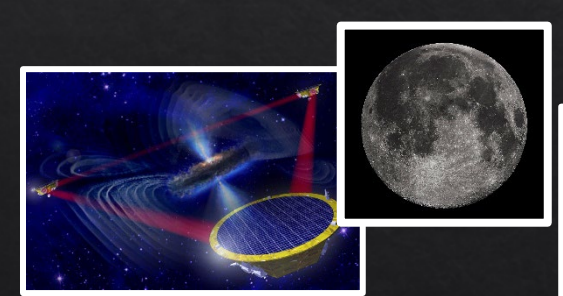
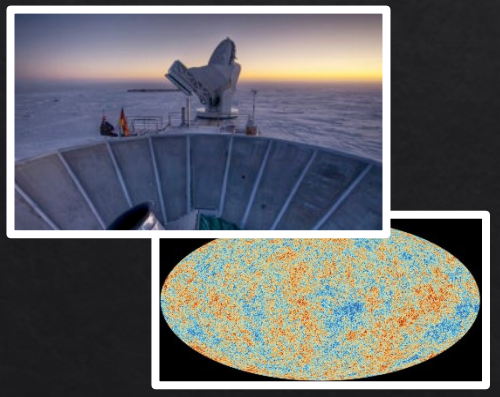


Microwave background

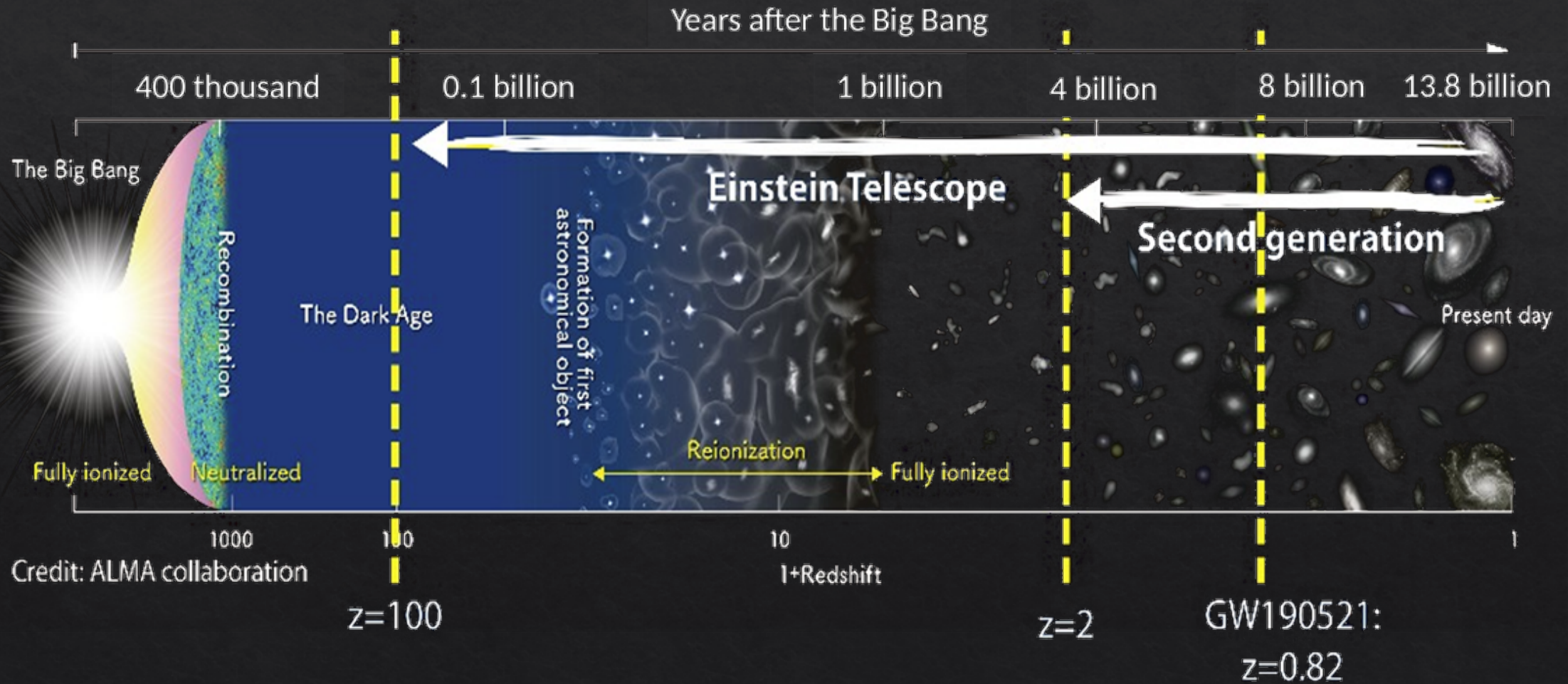
Pulsar timing

Space detectors

Terrestrial



Binary BH Detection Horizon



Credit: ALMA collaboration

Detection Numbers

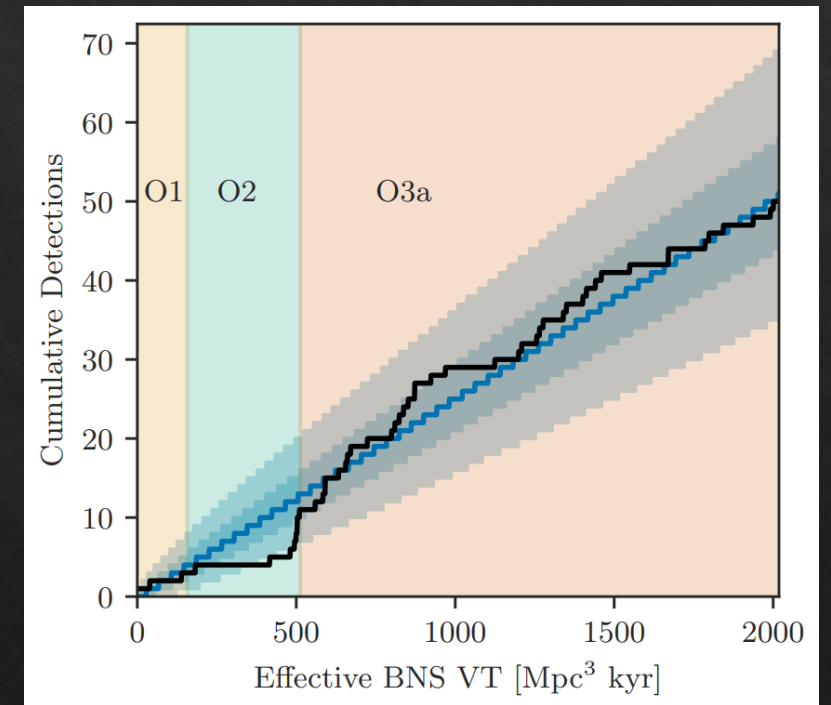
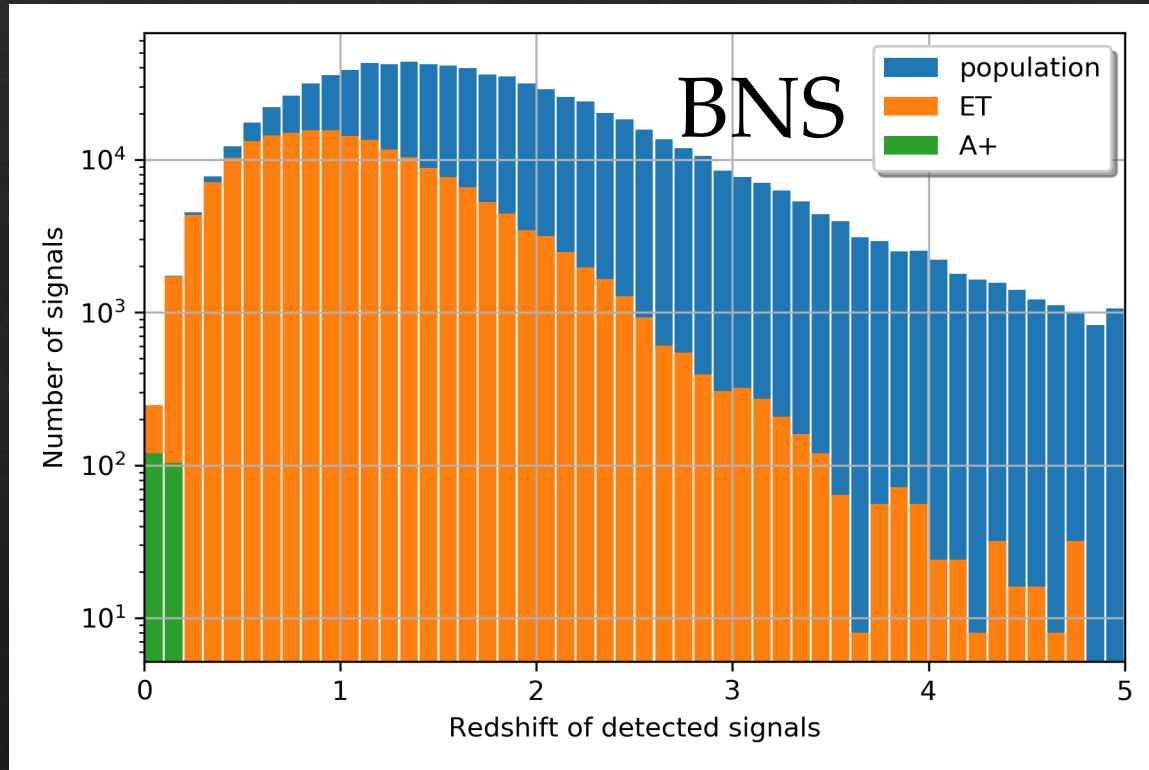
ET is expected see about 100000
BNS and 100000 BBH per year

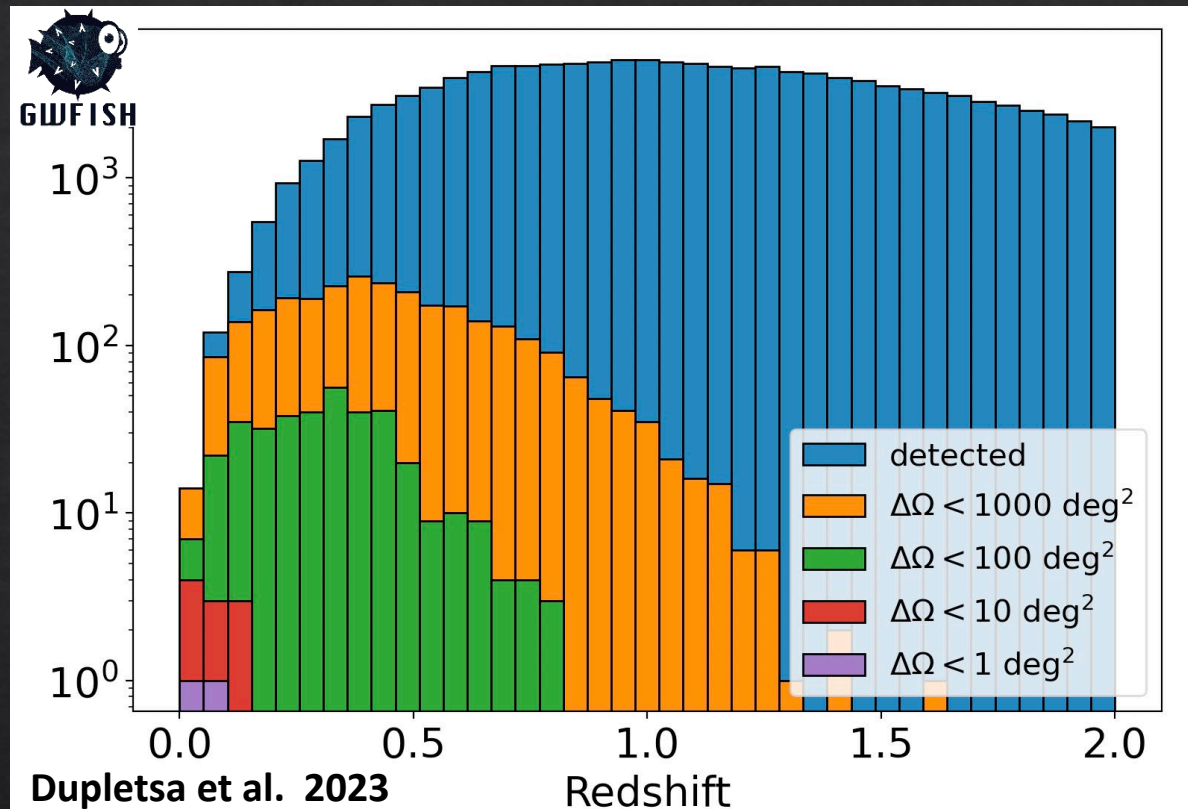
O1: 12.9.2015 – 19.1.2016

O2: 30.11.2016 – 25.8.2017

O3a: 1.4.2019 – 30.9.2019

O3b: 1.11.2019 – 26.3.2020

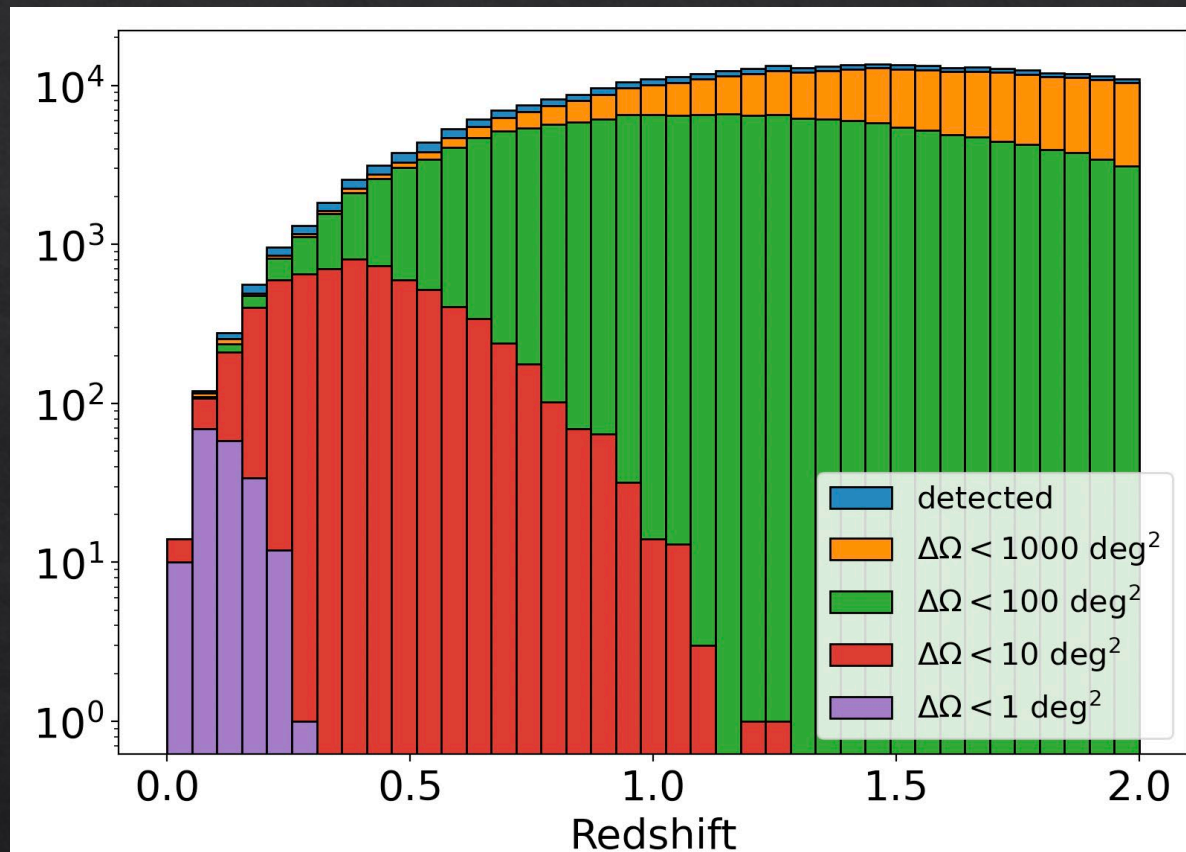




ET low frequency sensitivity makes it possible to localize BNS:

Modulation of signal amplitude and phase when observing it for many hours.

- O(100) detections per year with sky-localization (90% c.r.) $< 100 \text{ sq. deg}$
- Early warning alerts!



Cosmic Explorer
A proposed detector in the USA

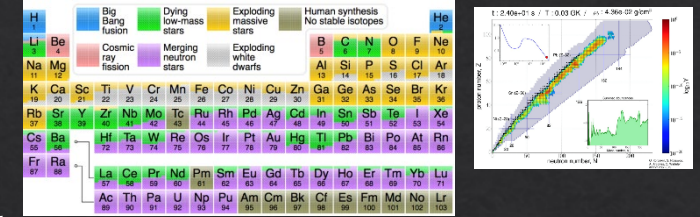


- $O(1000)$ detections per year with sky-localization (90% c.r.) $< 10 \text{ sq. deg}$

Radioactively powered transients



Nucleosynthesis

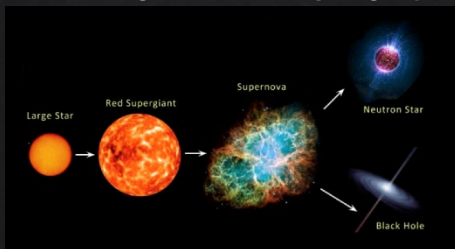


Relativistic astrophysics



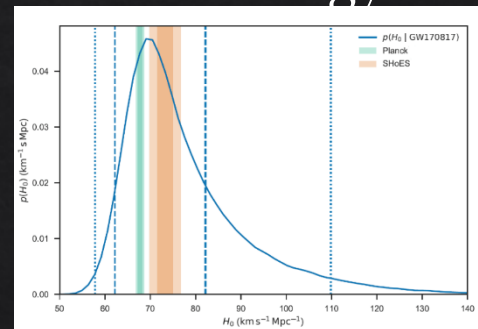
Credit: M Branchesi

Compact-object formation and evolution

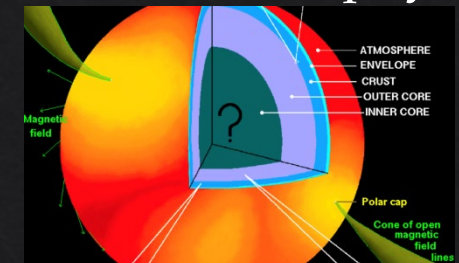


Cogne; J Harms

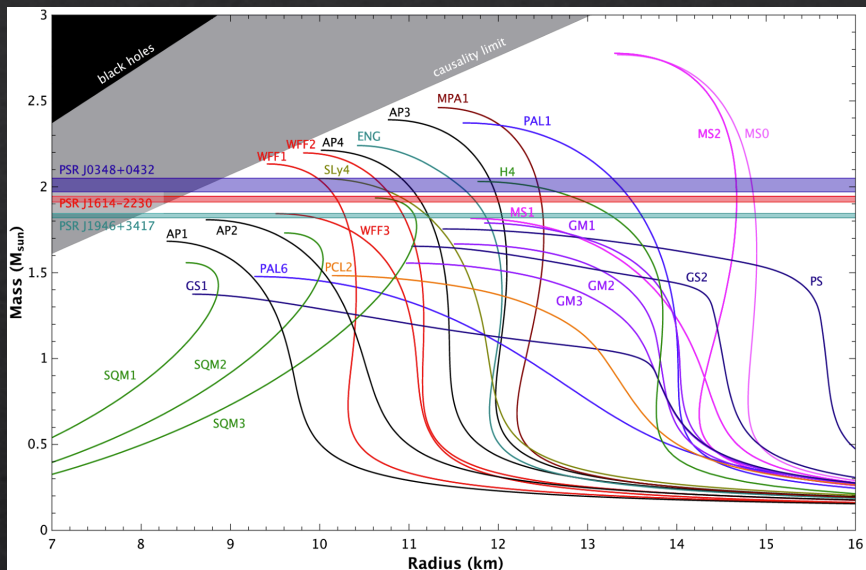
Cosmology



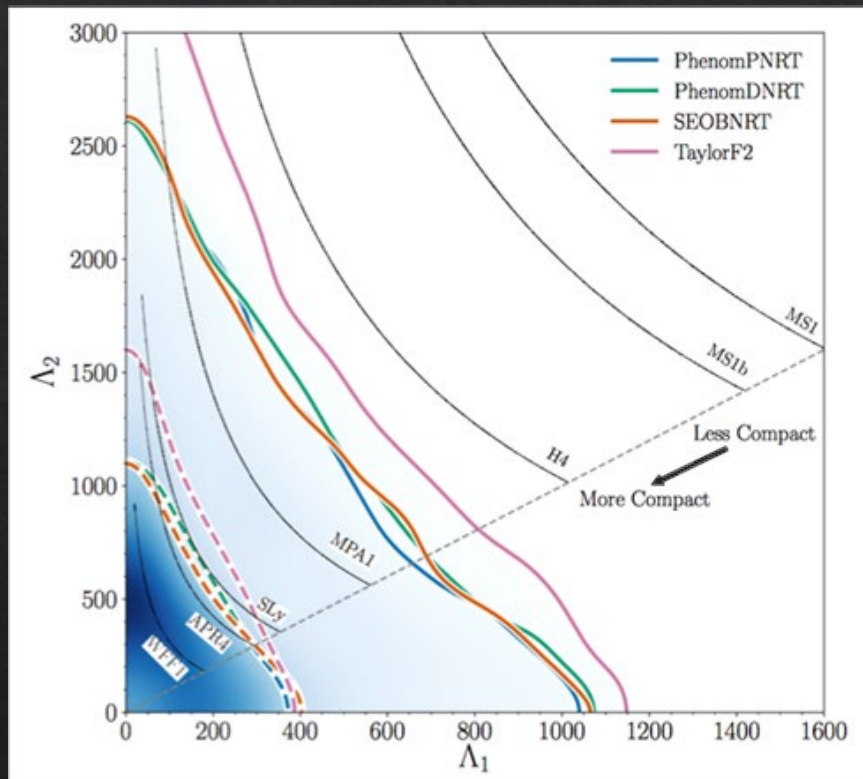
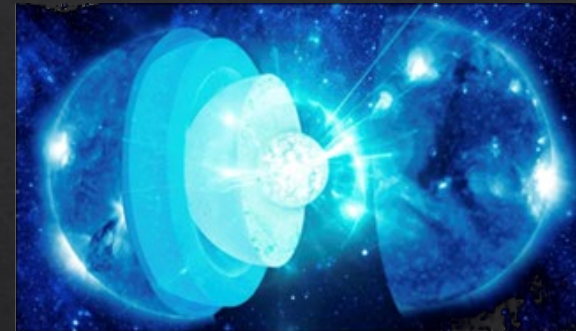
Nuclear matter physics



Intrinsic mass – NS radius relation

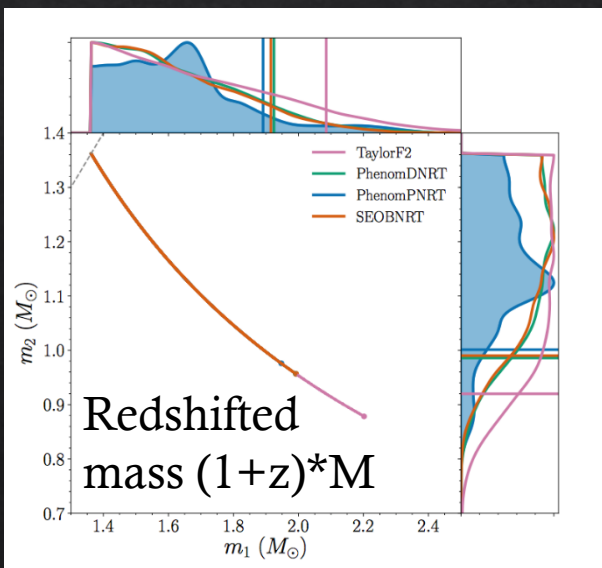
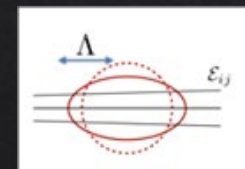


Observing deformations of neutron stars just before merger



TIDAL DEFORMABILITY

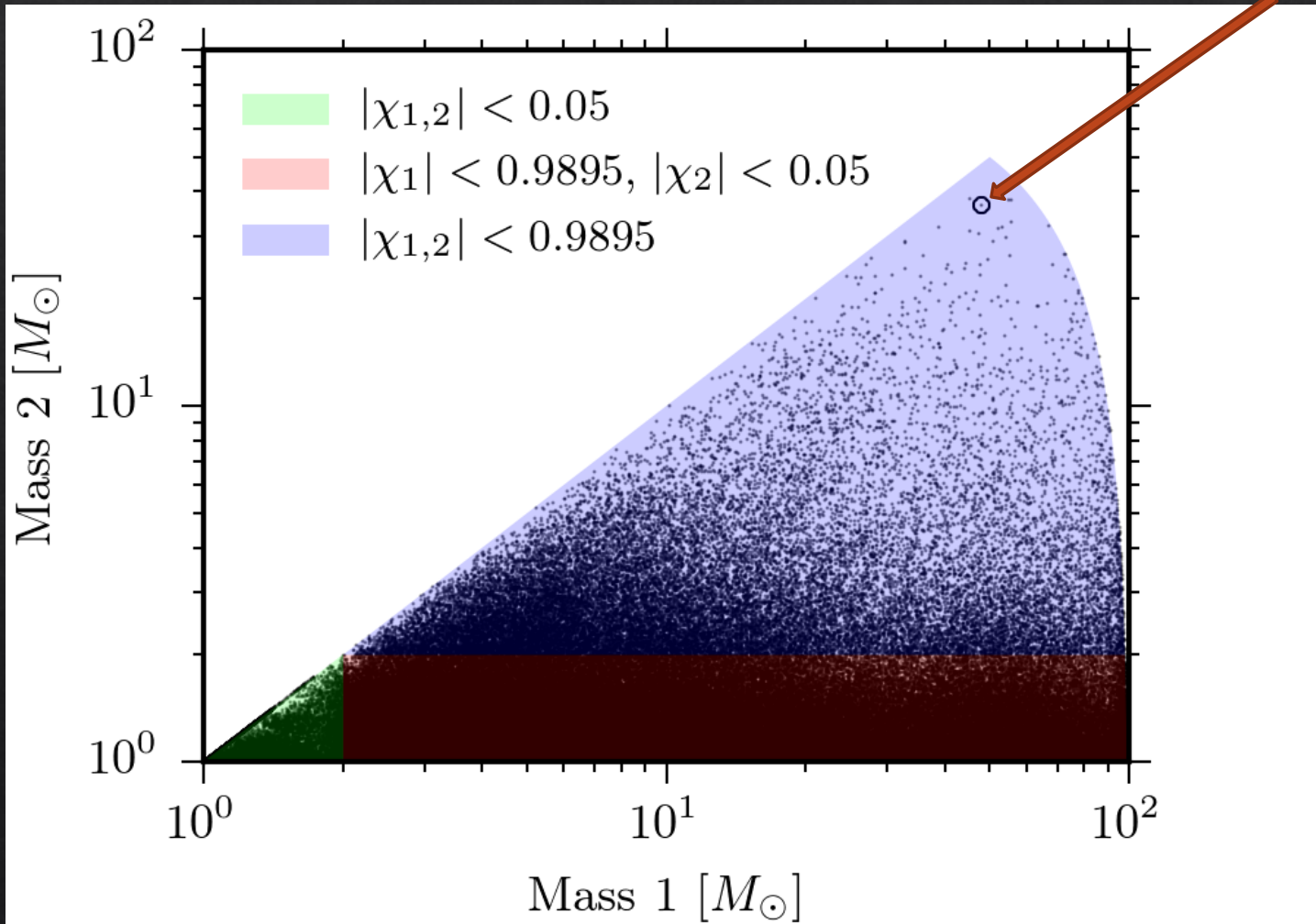
$$\Lambda = (2/3)k_2[(c^2/G)(R/m)]^5$$



How to Detect a CBC: Matched Filtering

Example: template bank (in mass space) when GW150914 was discovered

Best match with GW150914



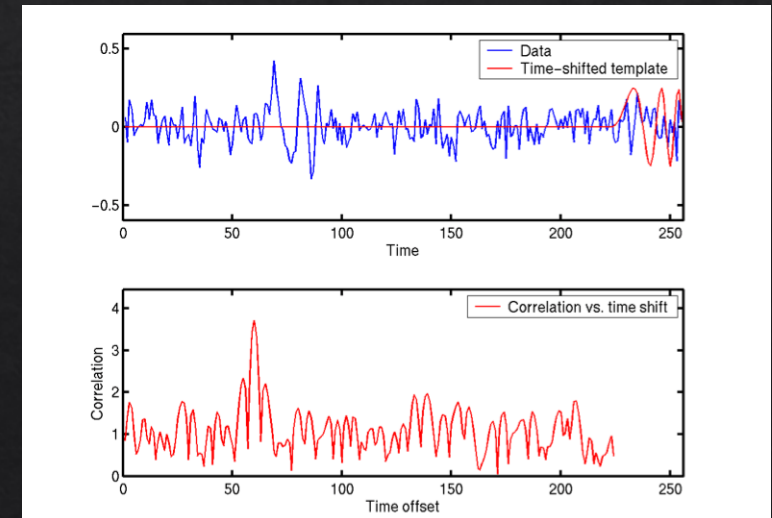
Matched filtering

FFT of data

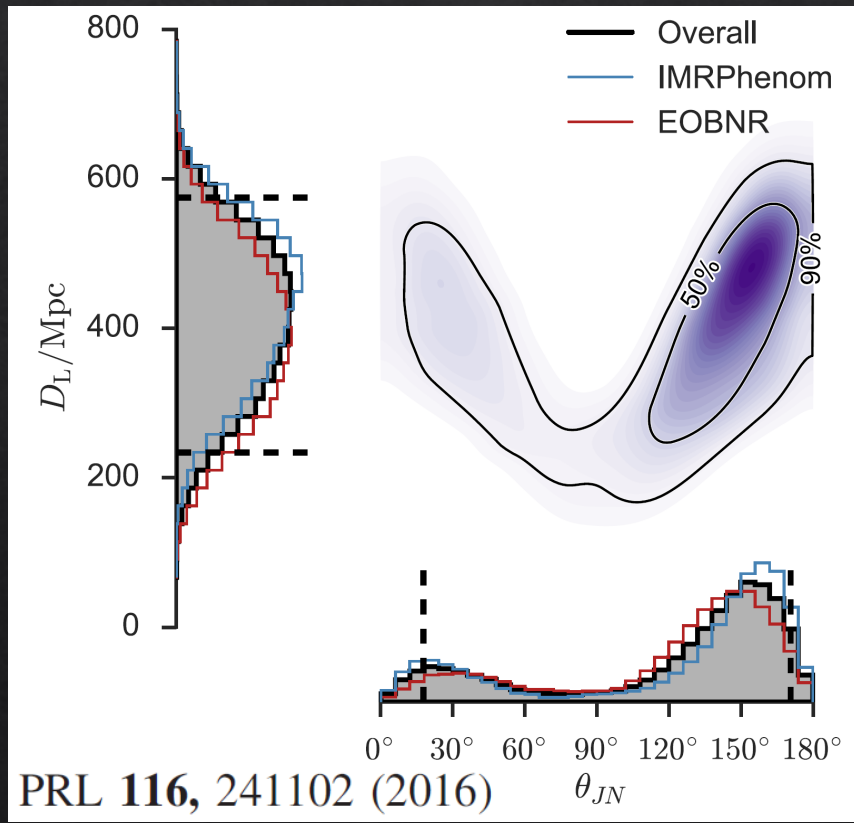
Template; can be generated in frequency domain using stationary phase approximation

$$C(t) = 4 \int_0^{\infty} \frac{\tilde{s}(f) \tilde{h}^*(f)}{S_n(f)} e^{2\pi i f t} df$$

Noise power spectral density



GW150914



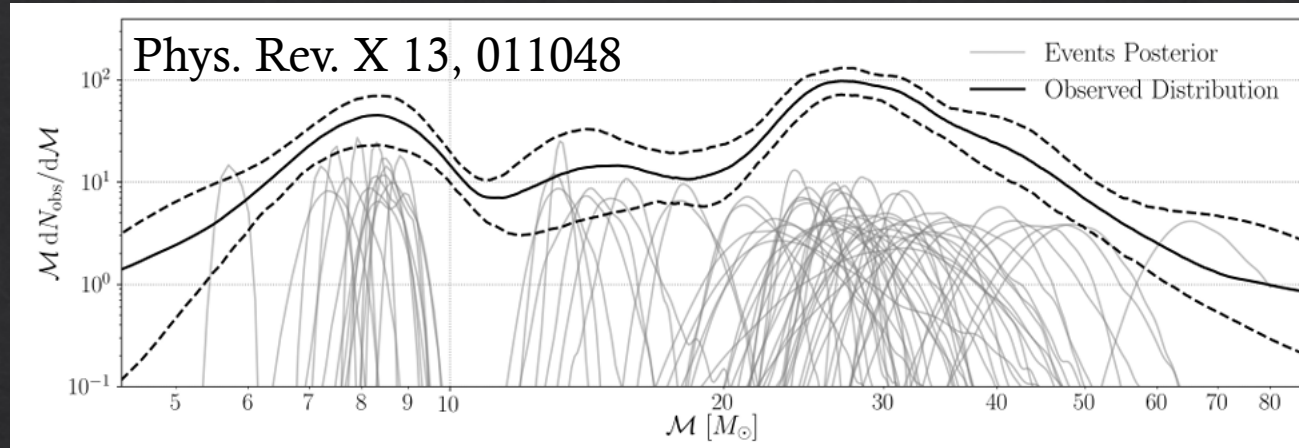
$$\underbrace{p(\theta|d)}_{\text{Posterior}} = \frac{\underbrace{p(d|\theta)}_{\text{Likelihood}} \underbrace{p(\theta)}_{\text{Prior}}}{\underbrace{p(d)}_{\text{Evidence}}}$$

Gaussian likelihood approximation:
Assume that detector noise is Gaussian

$$p(d|\theta) \propto \exp \left[-\frac{1}{2} \sum_k \langle h_k(\theta) - d_k | h_k(\theta) - d_k \rangle \right]$$

$$\langle a|b \rangle = 4 \int_0^\infty df \frac{\Re(\tilde{a}(f)\tilde{b}^*(f))}{S(n;f)}$$

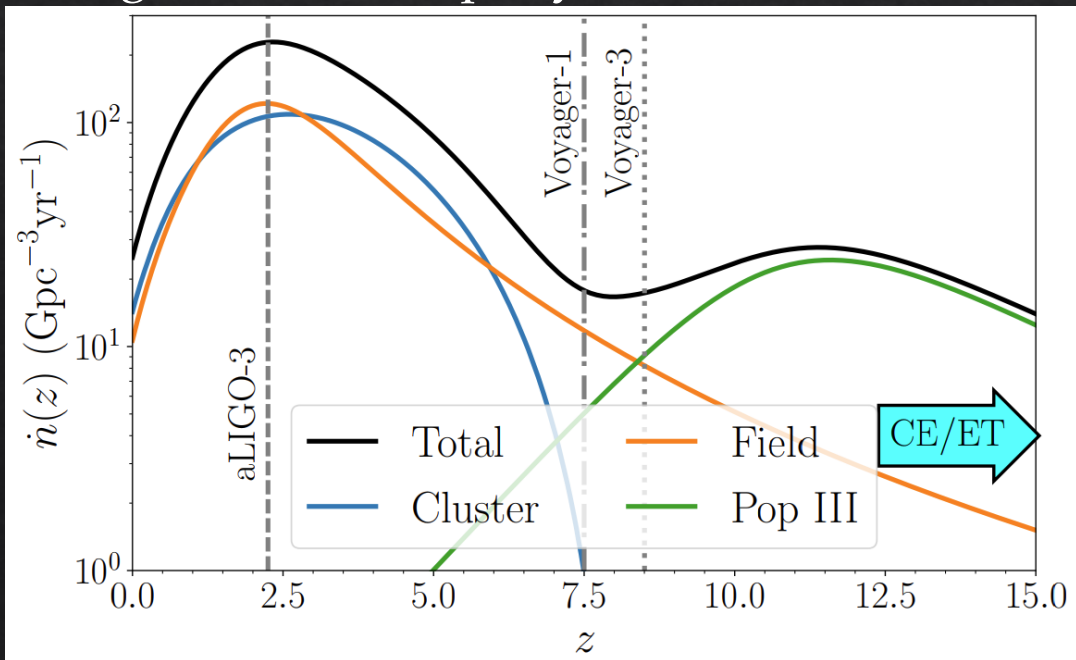
Black-hole Populations



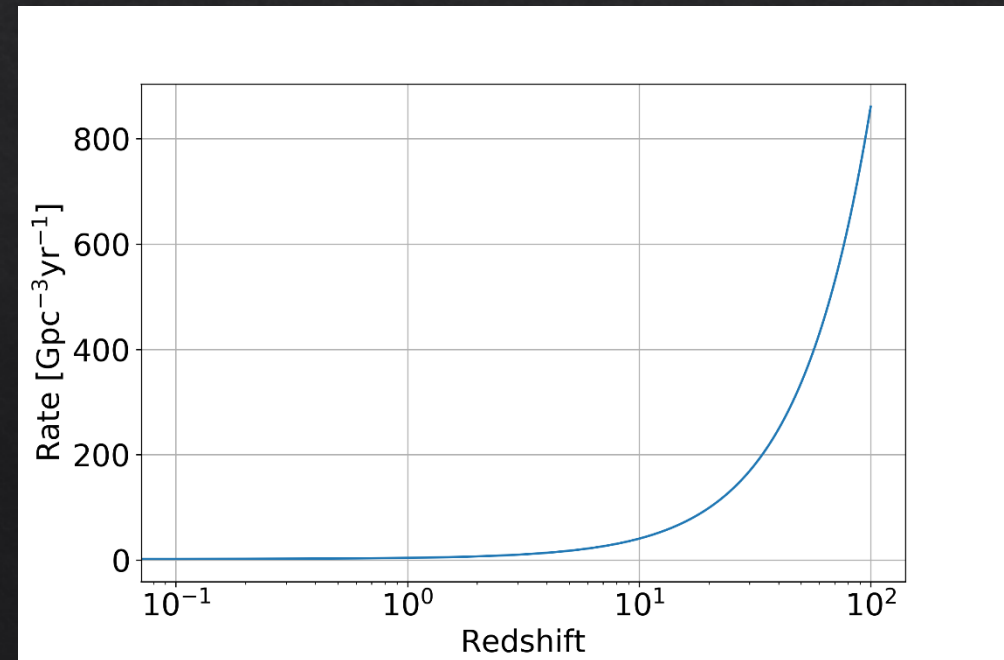
LVK, 2023

Precise population studies from 10^5 BBH mergers observed per year with ET/CE.

Is there a population of primordial black-holes?



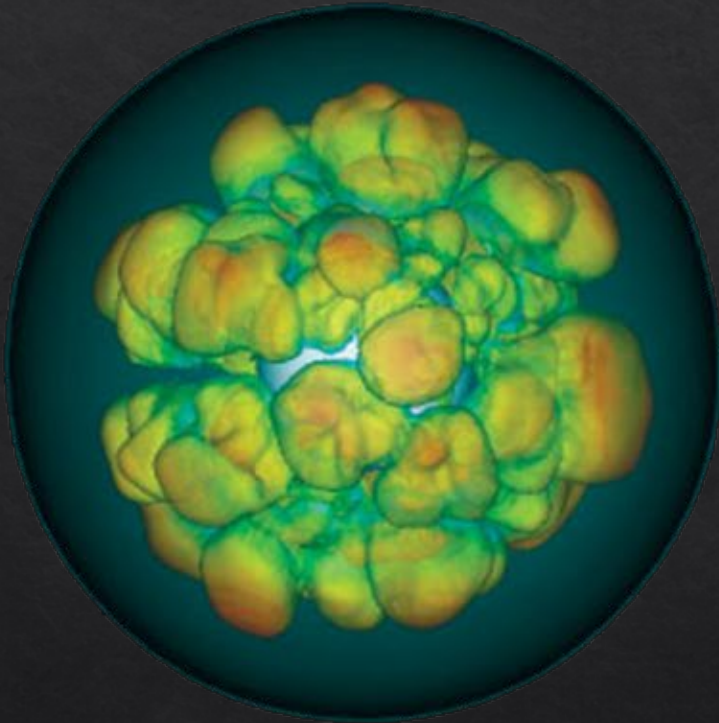
Ng et al, 2021



De Luca et al, 2020

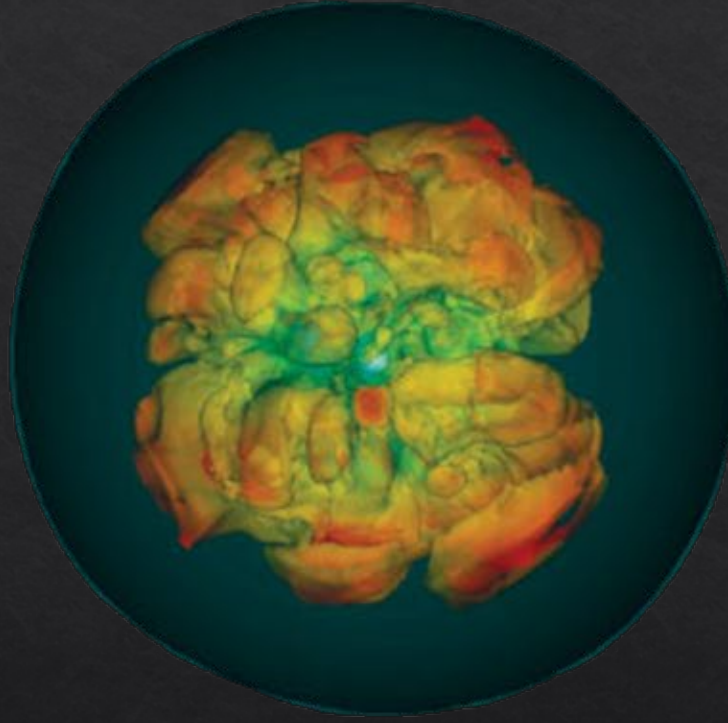
Core-Collapse Supernovae

90 ms



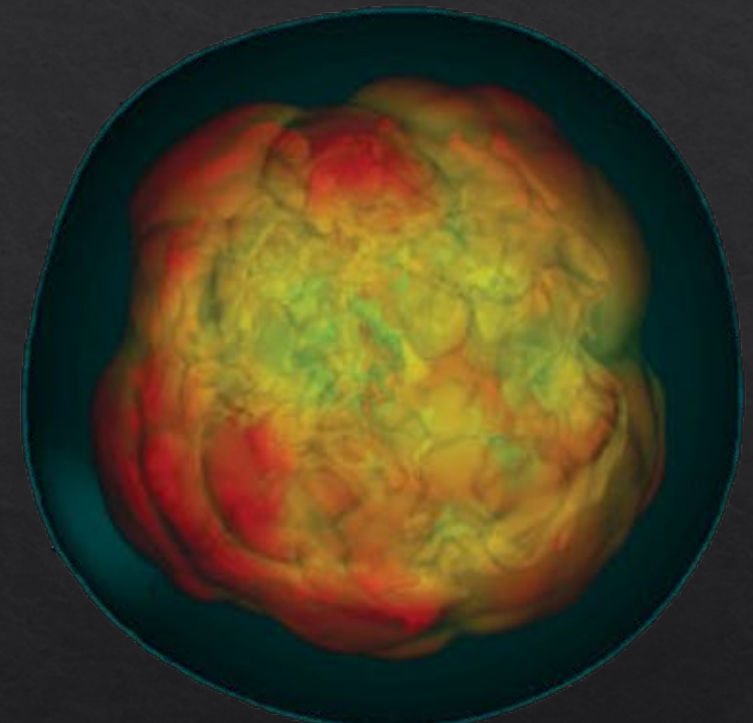
100 km

170 ms



500 km

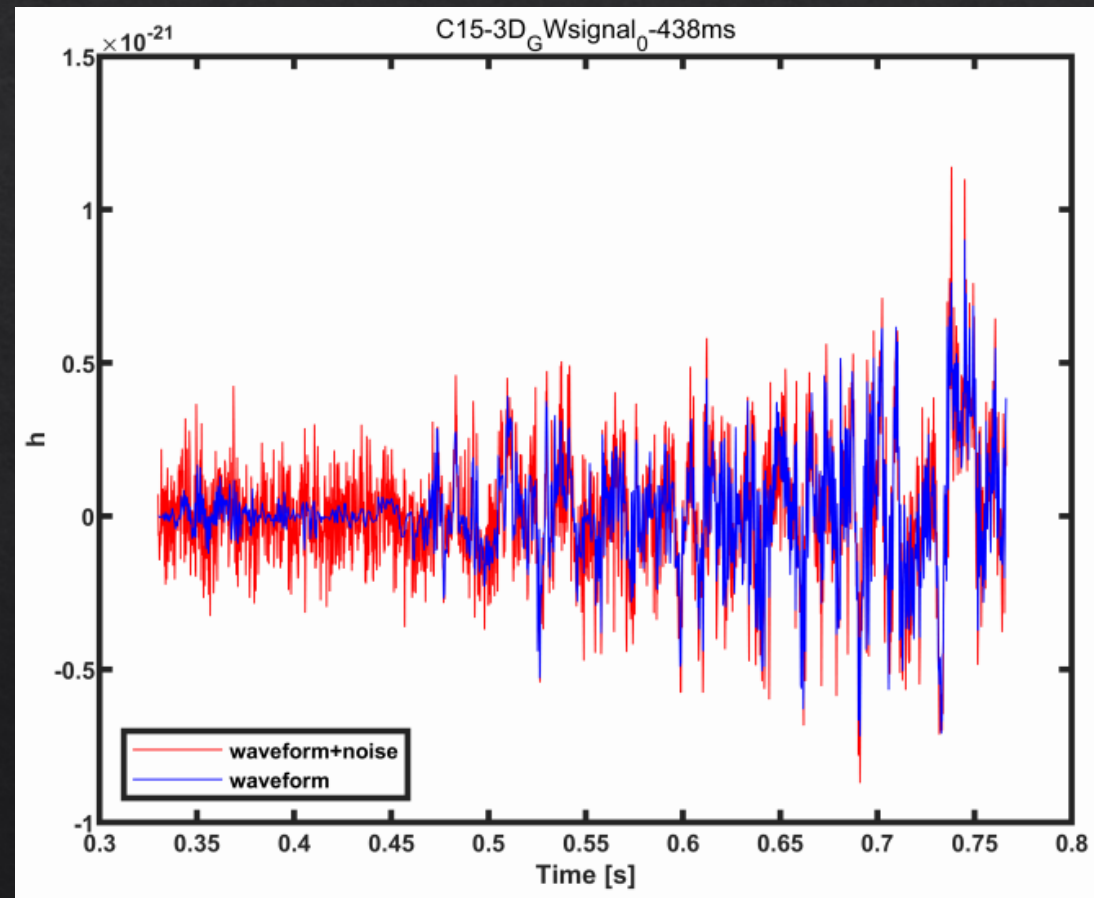
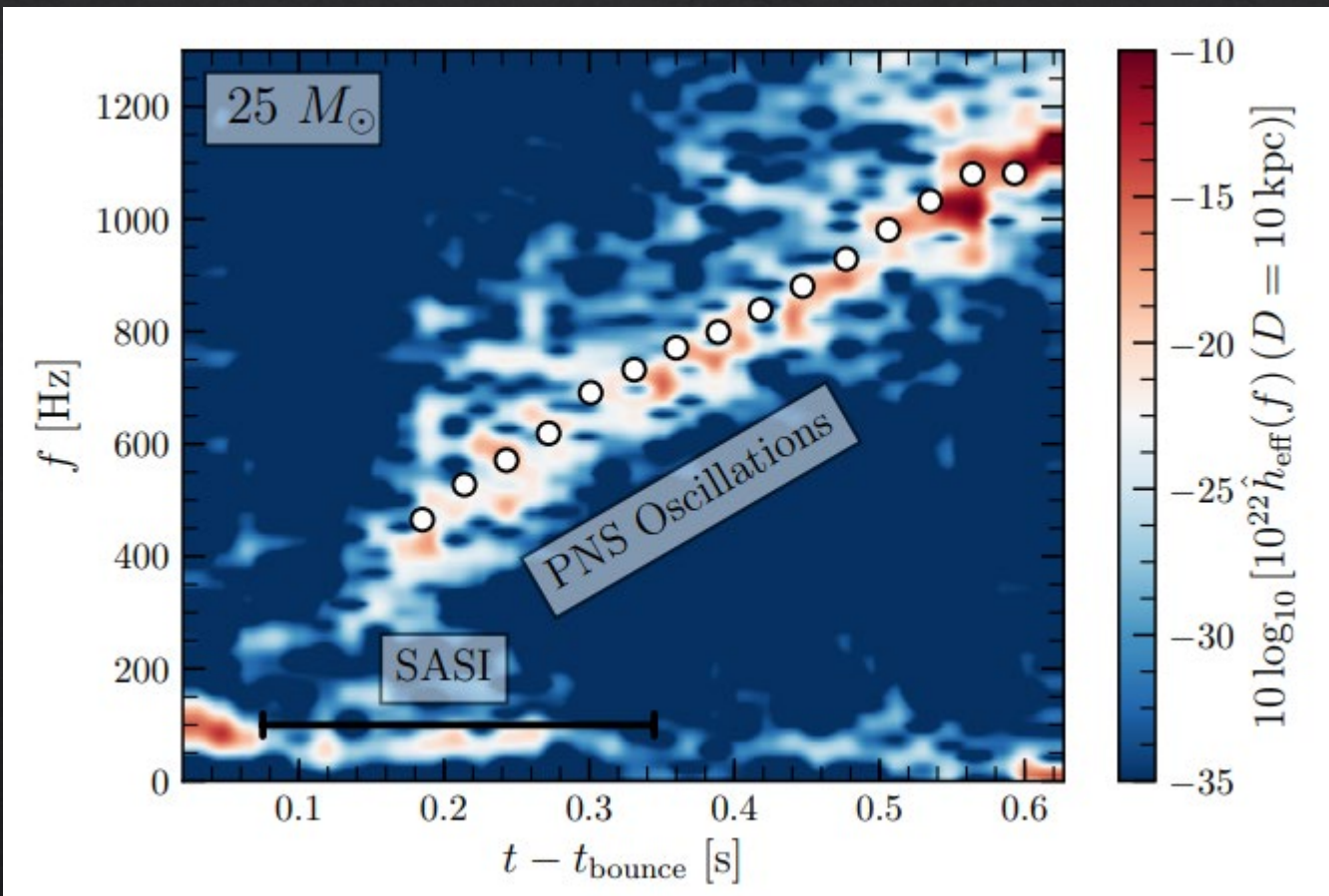
350 ms



3,000 km

GW signals from core-collapse SNe

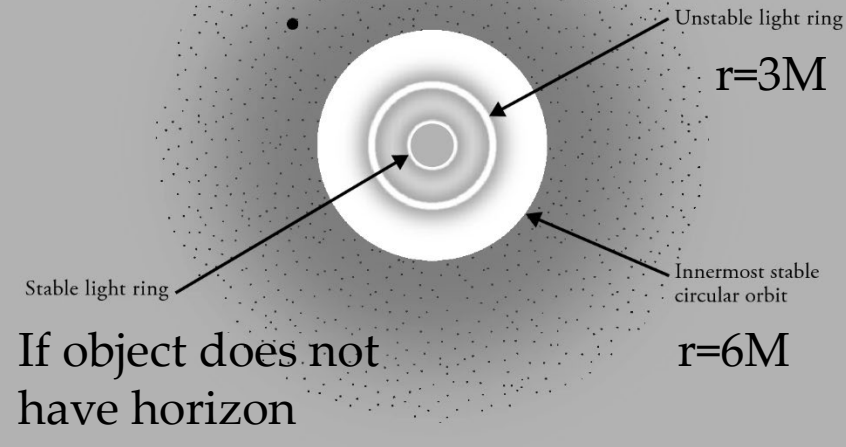
Modeling of SN signals is an extremely complex problem and not yet fully understood.



Black holes: the ultimate engine of discovery
[cit Cardoso, 2020]

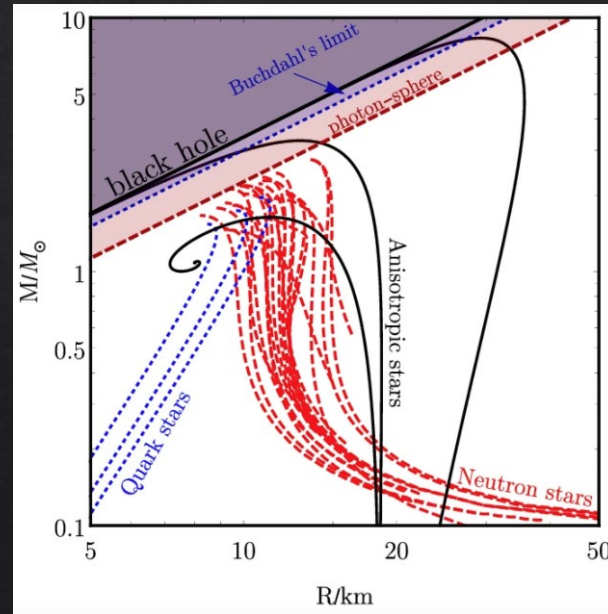
Probing the structure
of BH spacetime

BH horizon at $r=2M$



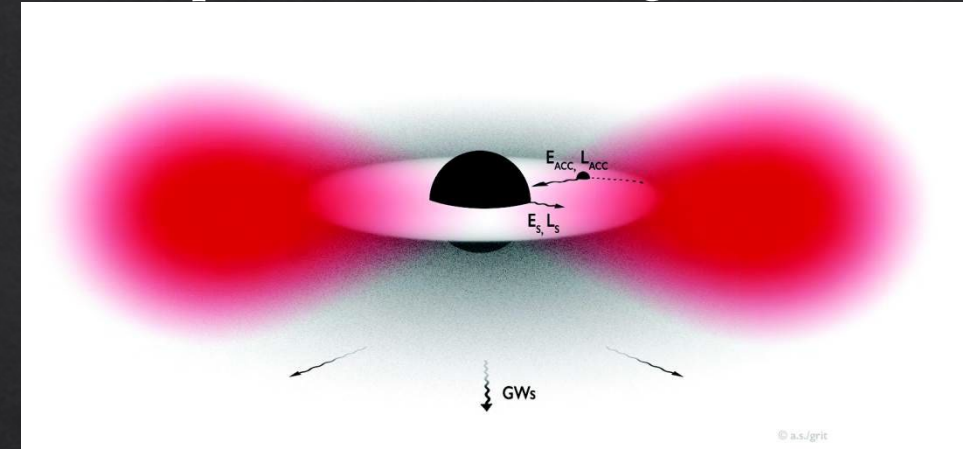
Cardoso, Pani, 2019

Exotic compact objects:
Infer mass-radius relation

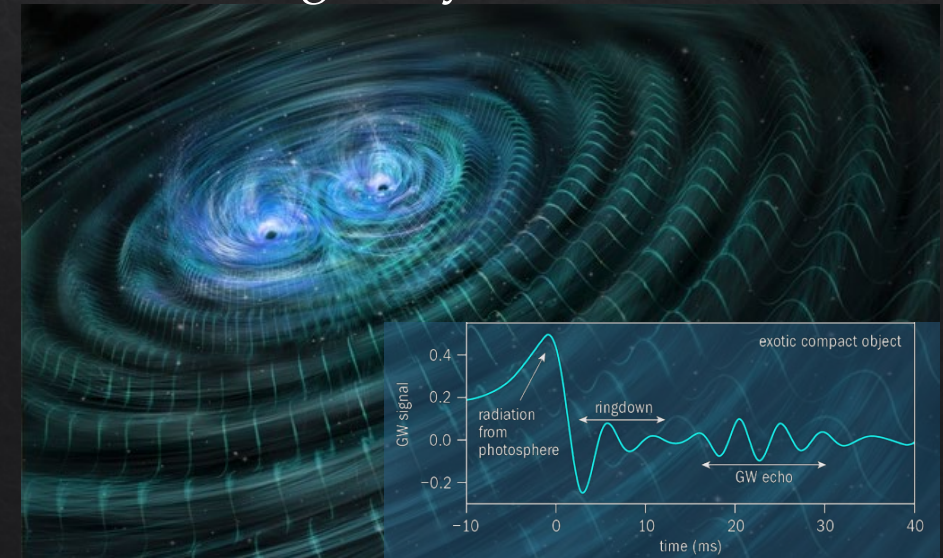


Cardoso, Pani, 2019

BH superradiance with light bosons



Quantum gravity: area discretization



Agullo et al, 2021

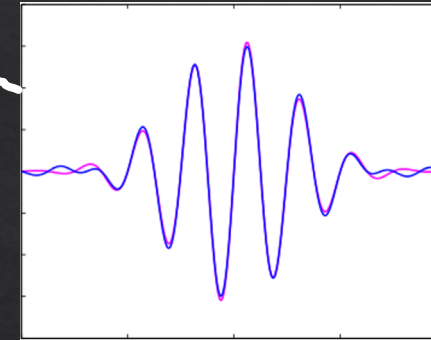
Coherent WaveBurst

wavelet transform

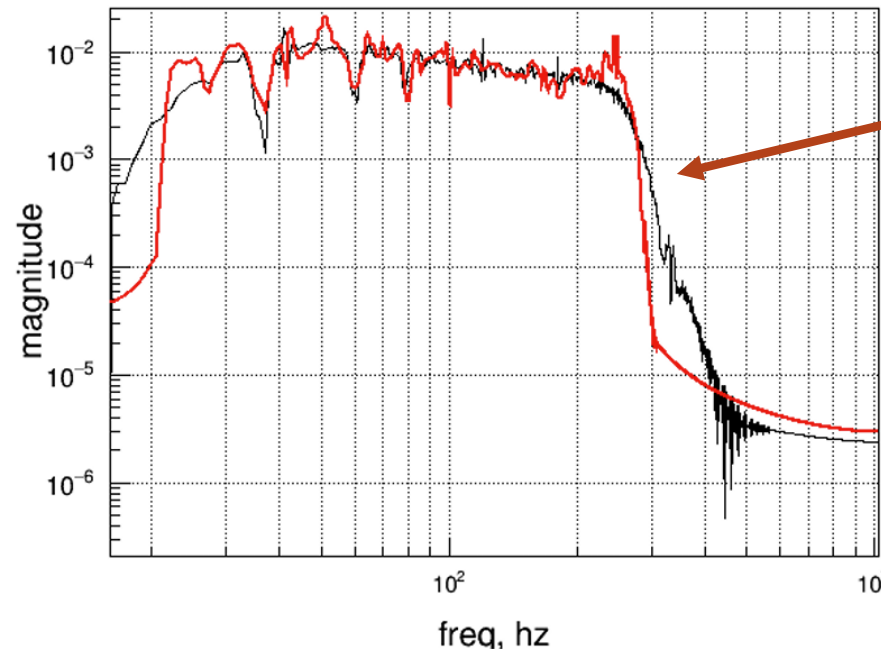
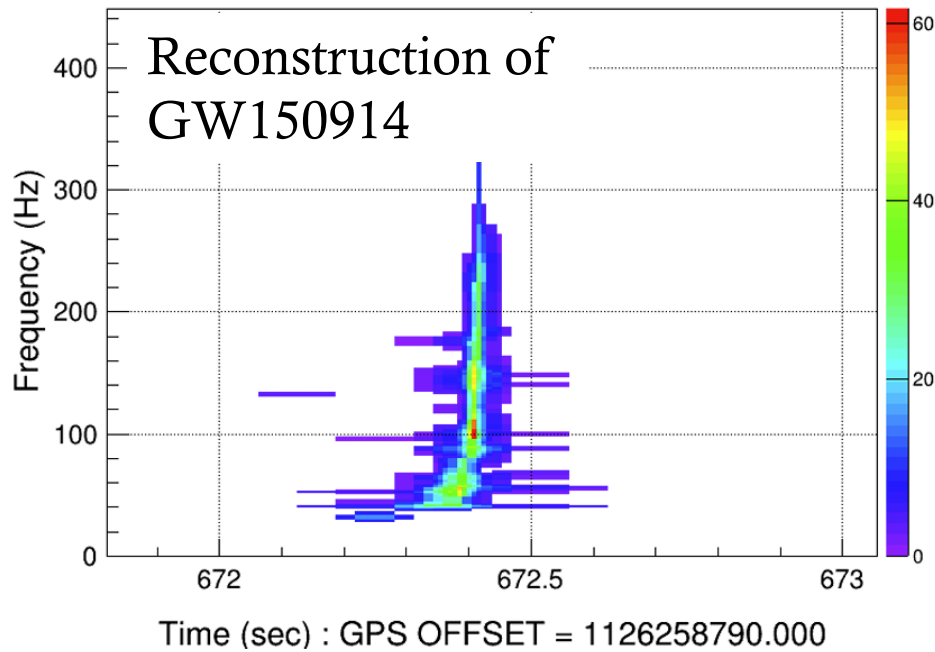
data

wavelet

$$X(\tau, \phi, Q) = \int_{-\infty}^{+\infty} x(t) w(t - \tau, \phi, Q) e^{-i2\pi\phi t} dt$$

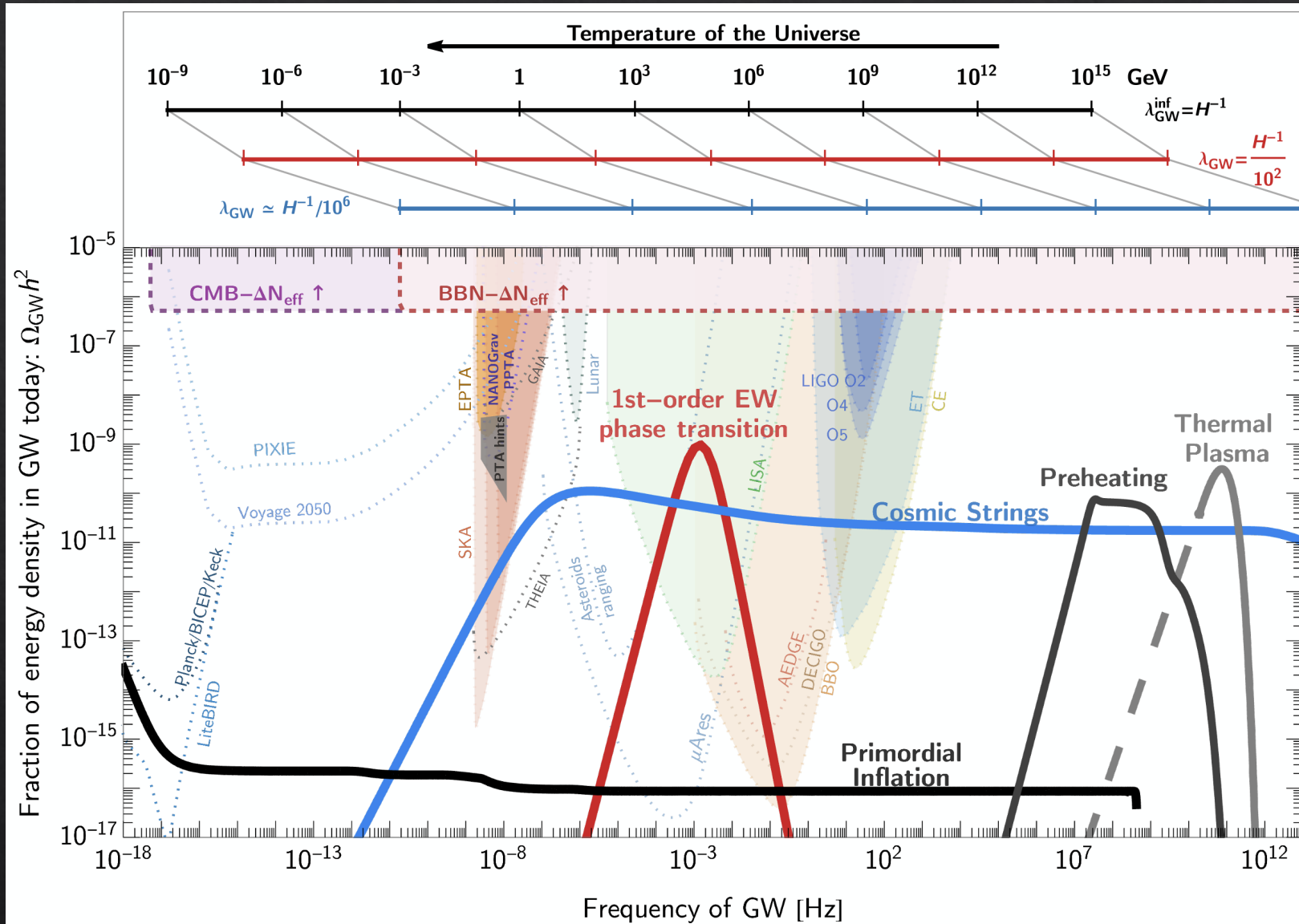


Likelihood 641 - dt(ms) [7.8125:250] - df(hz) [2:64] - npix 131



Comparison in frequency domain

Primordial GW Backgrounds

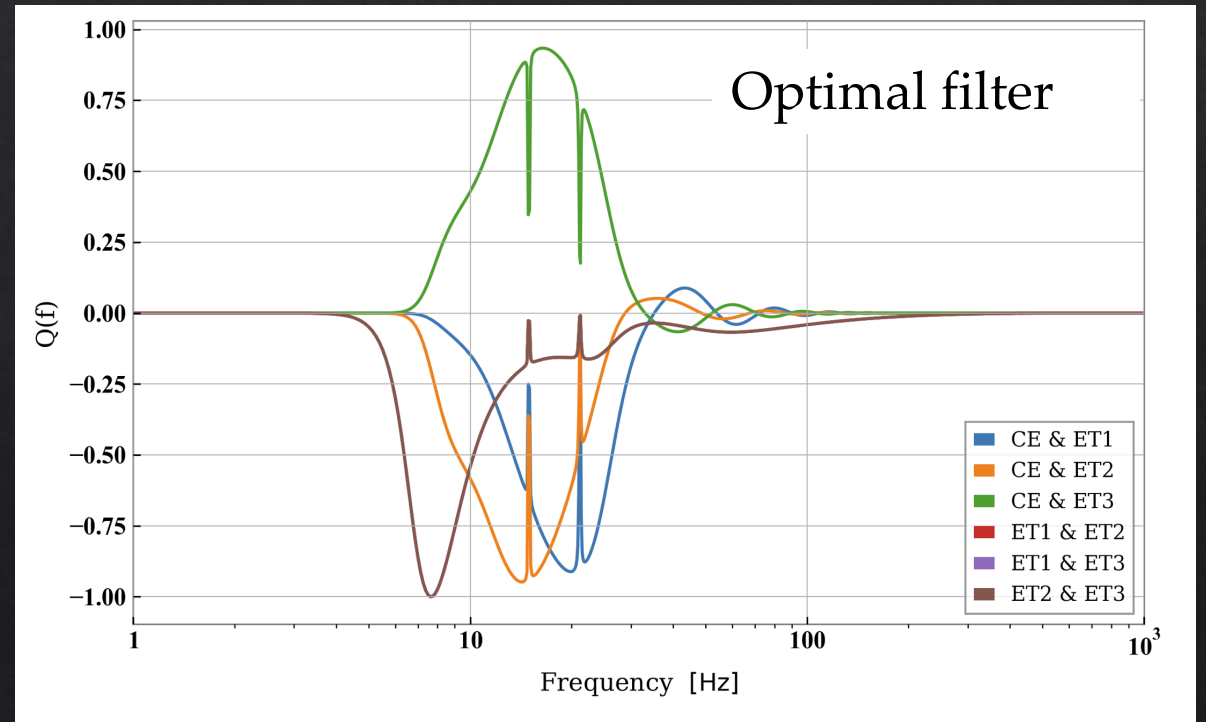
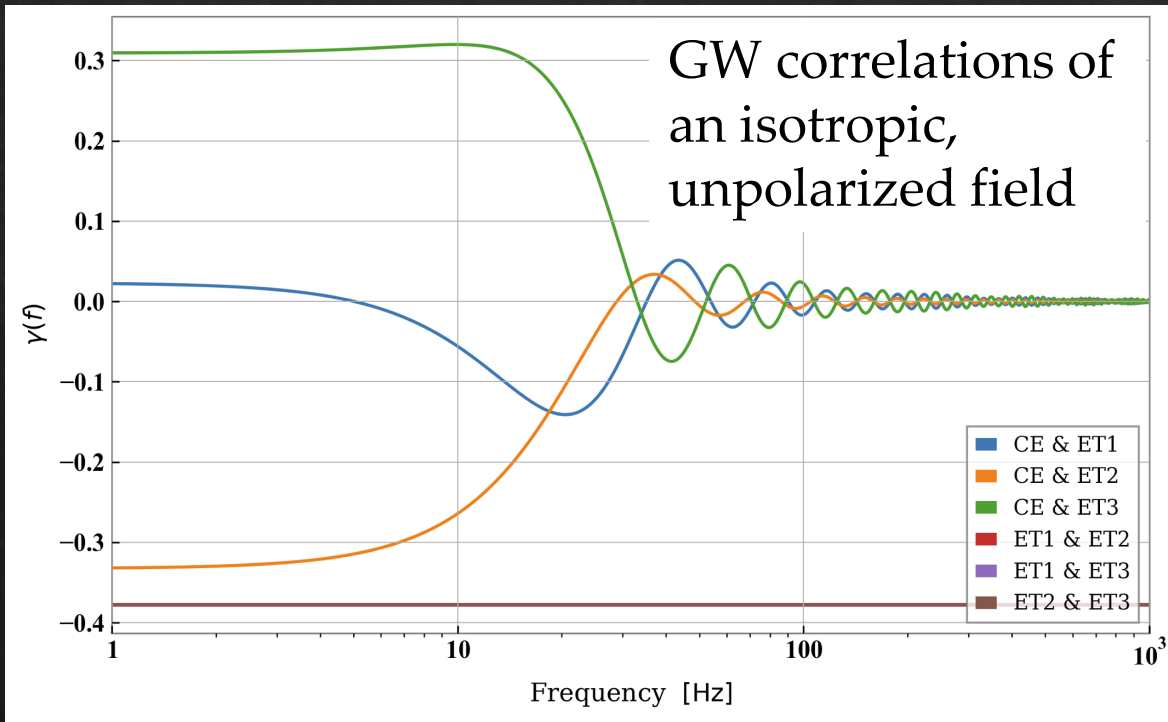


Simakachorn, 2022

Observation of Stochastic Backgrounds

Correlate data between GW detectors!

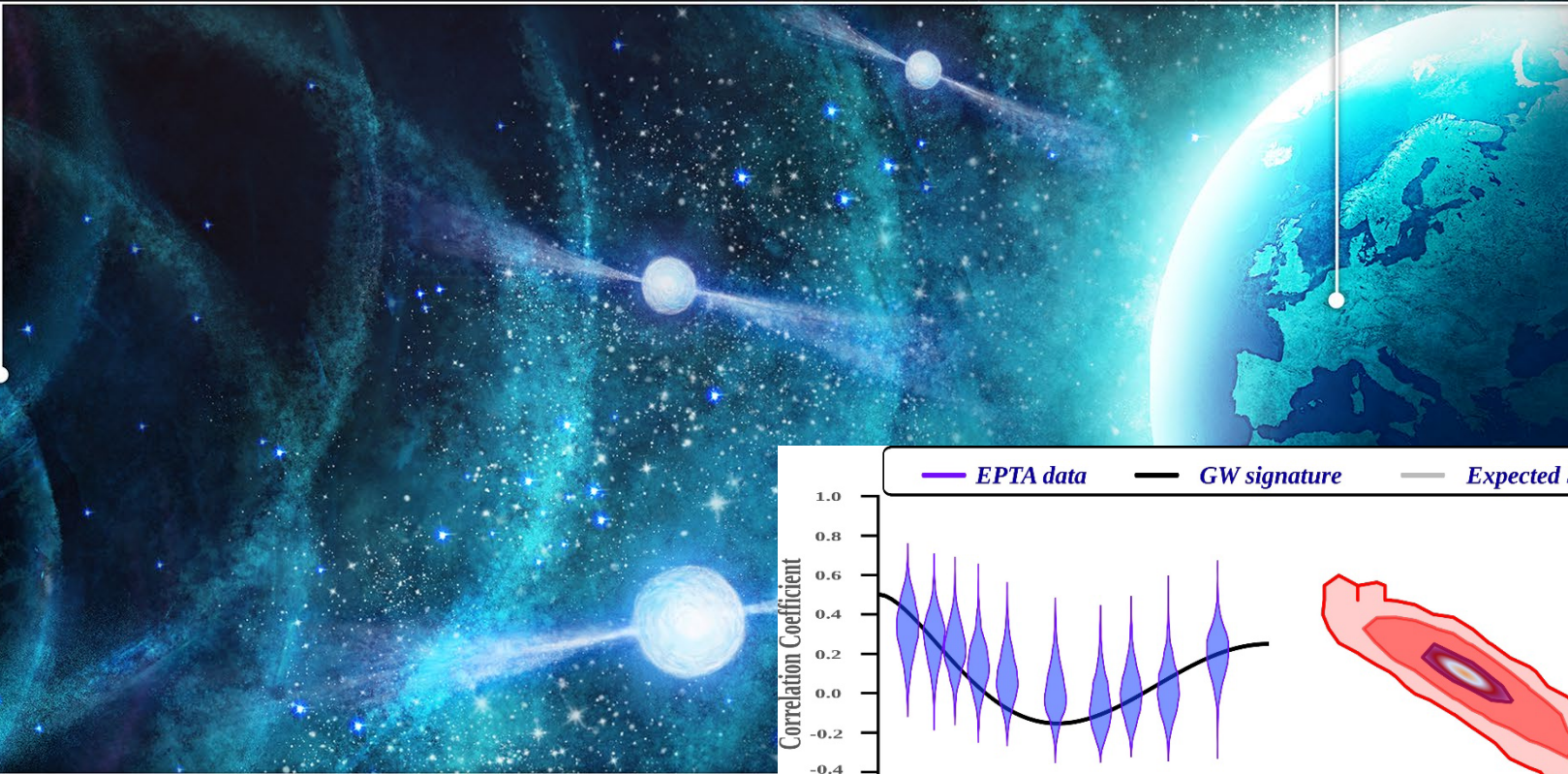
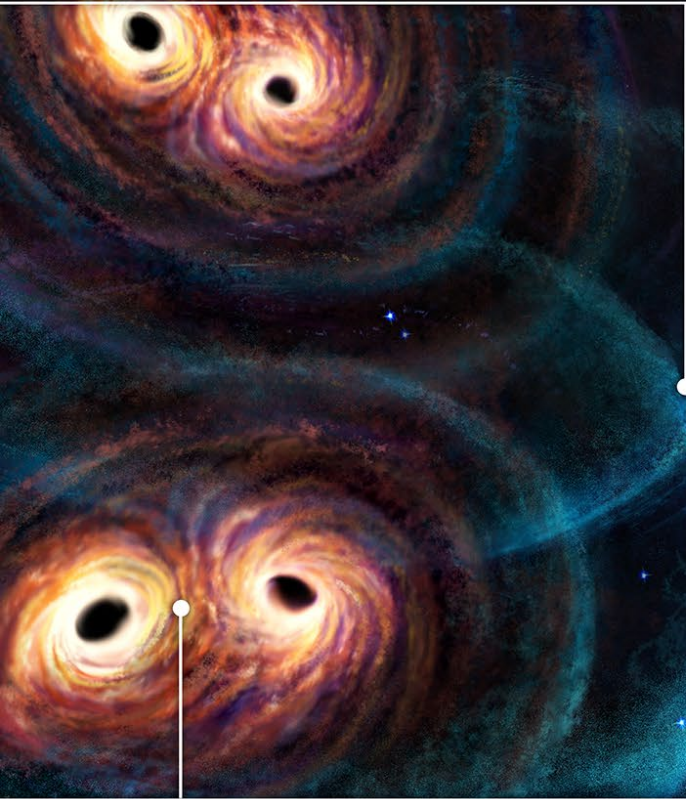
$$C_{ij}(f) = S_{\text{GW}}(f)\gamma_{ij}(f) \quad \langle C_{ij} \rangle = \int_0^{\infty} df \langle C_{ij}(f) \rangle \tilde{Q}_{ij}(f)$$



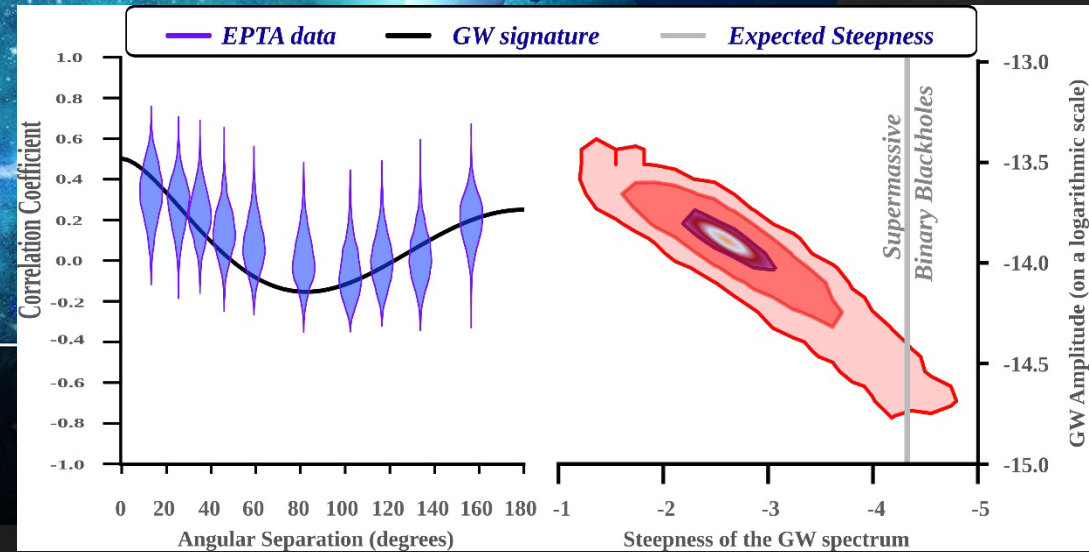
The News of the Day: PTAs

Gravitational waves stretch and squeeze space-time

The largest telescopes on Earth are used to precisely monitor the rotating ticks of these pulsars over decades to reveal the faint echoes of distant black holes

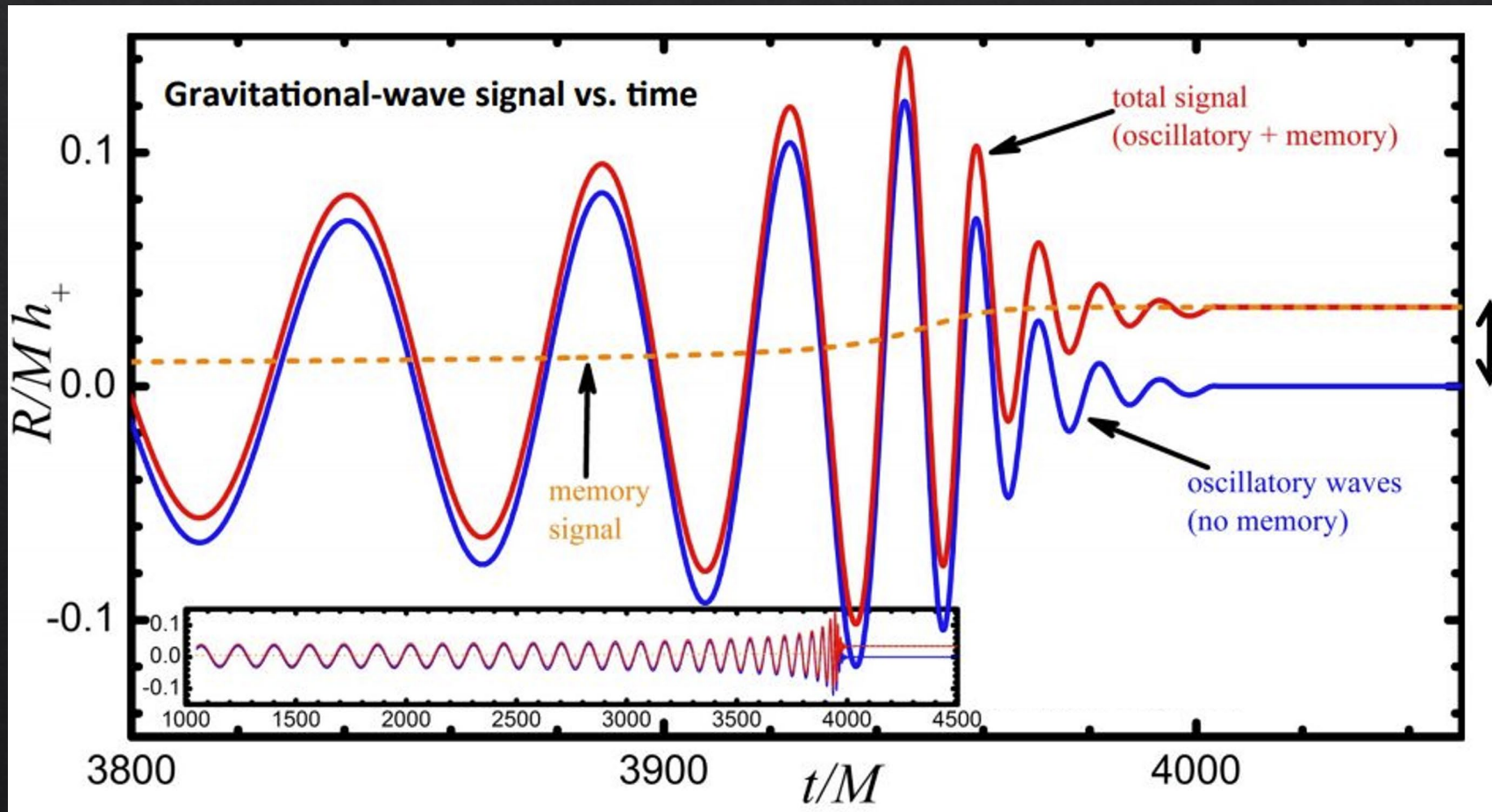


Supermassive black hole binaries in the distant Universe generate gravitational waves



Time-integrated GW energy emitted into solid angle

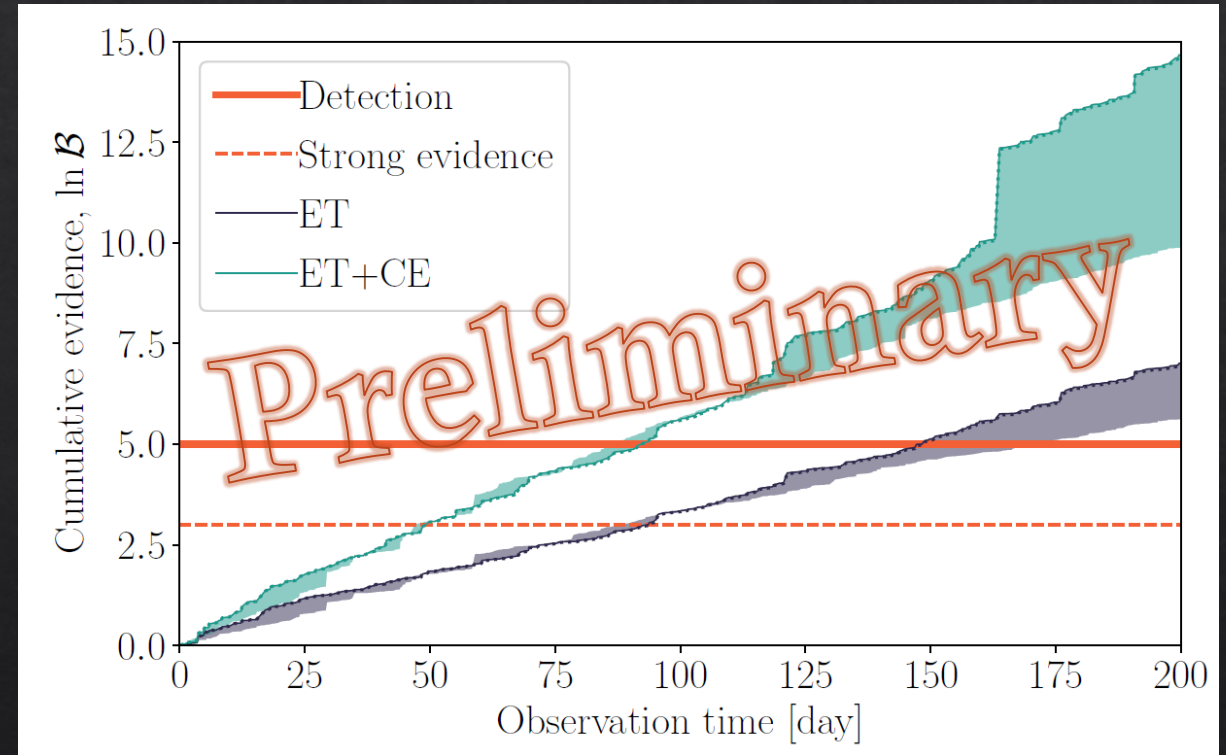
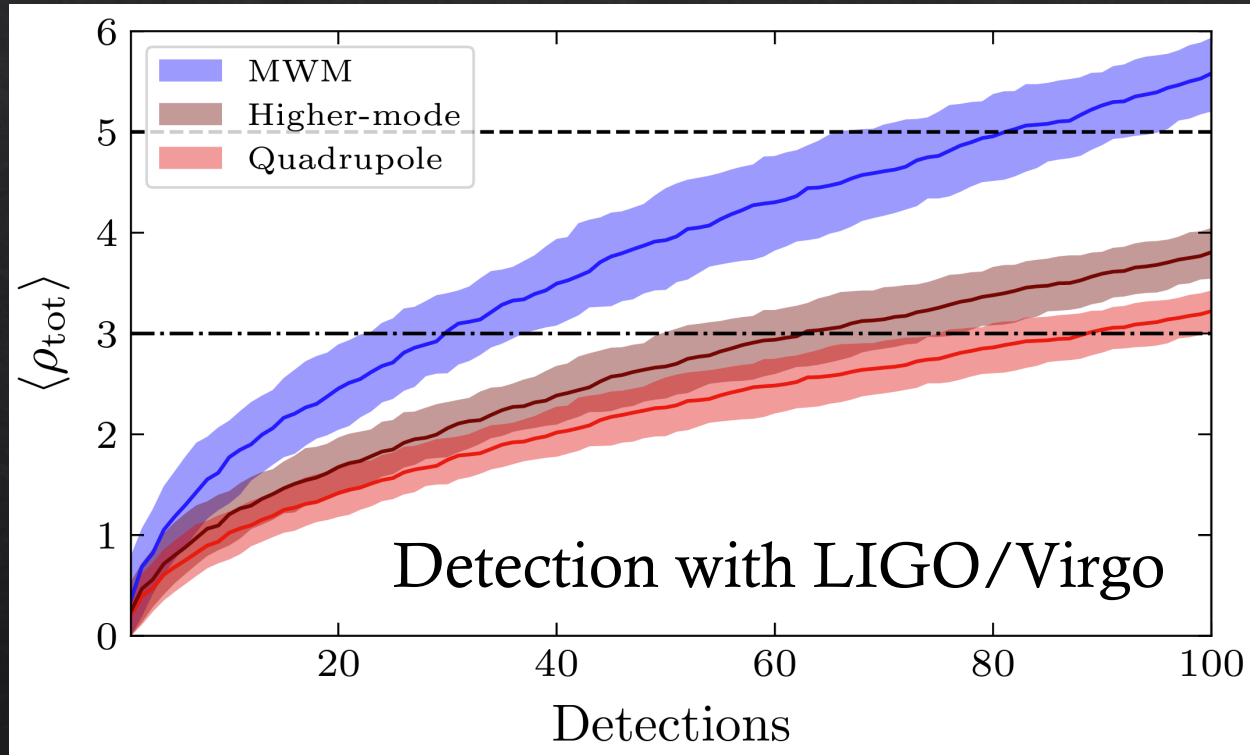
Christodoulou memory effect:
$$\Delta h_{jk}^{\text{TT}} = \frac{4}{r} \int \frac{dE}{d\Omega'} \left[\frac{\xi^{j'} \xi^{k'}}{1 - \cos\theta'} \right]^{\text{TT}} d\Omega'$$



GWs leave a lasting deformation of spacetime behind them.

Deep connection to spacetime symmetries and quantum gravity.

Evidence of detection and for model selection is accumulated with increasing number of GW signals



Boersma et al, 2020