

Firenze, 7 November 2017

Study of quantum electrodynamics and chromodynamics in atomic bound systems

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Exotic atomic systems

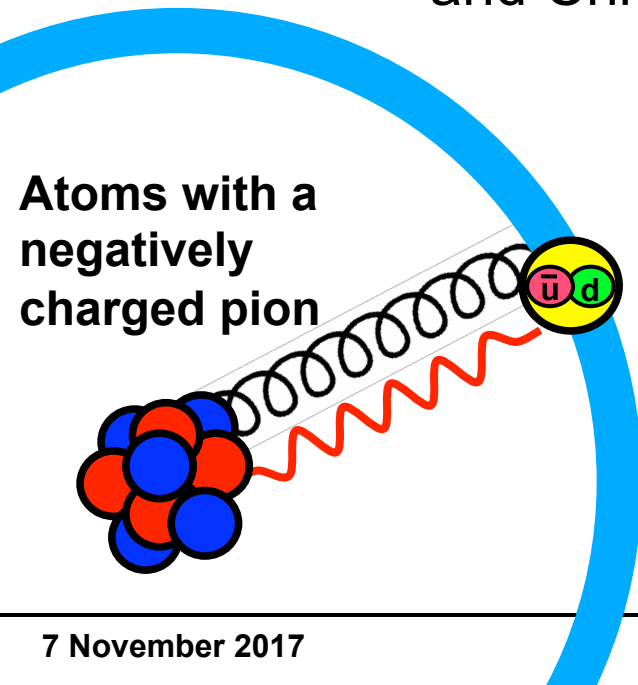
- Simple atomic systems
 - accurate theoretical predictions
 - simple spectra

→ accurate test of fundamental interactions:

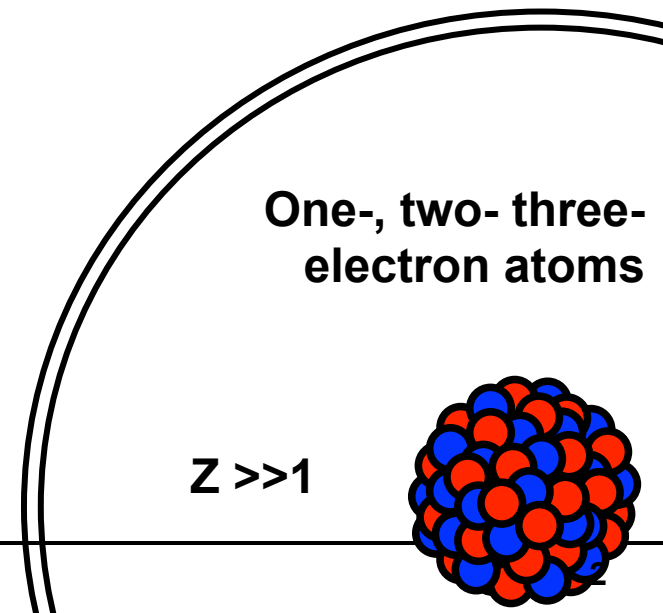
Quantum Electrodynamics (in strong field)
and Chromodynamics (at low energy)

Highly
charged
ions

Pionic
atoms



Atoms with a
negatively
charged pion



One-, two- three-
electron atoms

$Z \gg 1$

Summary

- Introduction
 - Test of QED in strong field
 - Study of strong interaction force at low energy
- Heavy highly charged ions
 - Production
 - Lamb shift measurement
 - Vacuum decay experiment
- Pionic atoms
 - Production
 - Pionic hydrogen and deuterium spectroscopy
 - Pion mass measurement
- Conclusions and outlooks

Summary

- **Introduction**
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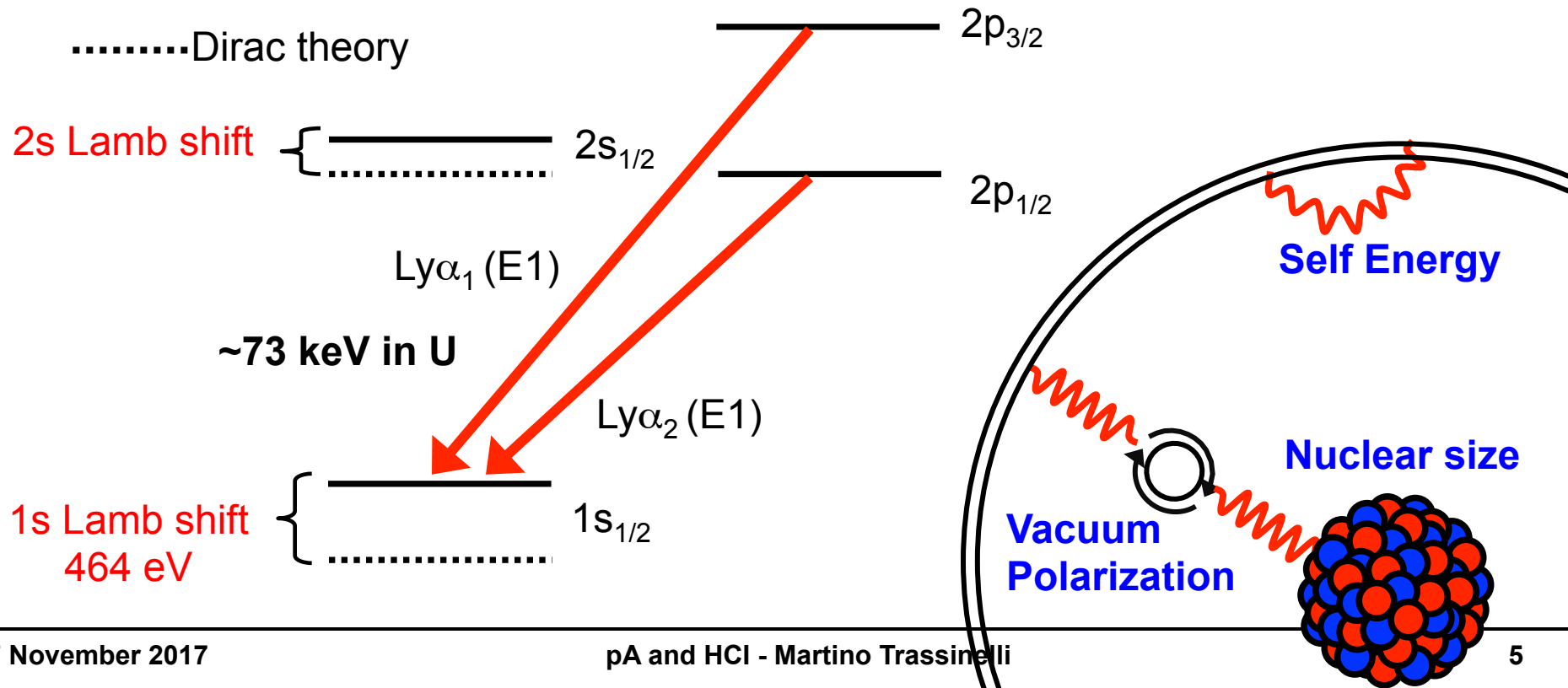
Highly charged ions

Dirac equation (spin-1/2 particles) →

$$E^{nj} = \frac{mc^2}{\sqrt{1 + \frac{(Z\alpha)^2}{\left[n - j - 1/2 + \sqrt{(j + 1/2)^2 - (Z\alpha)^2}\right]^2}}}$$

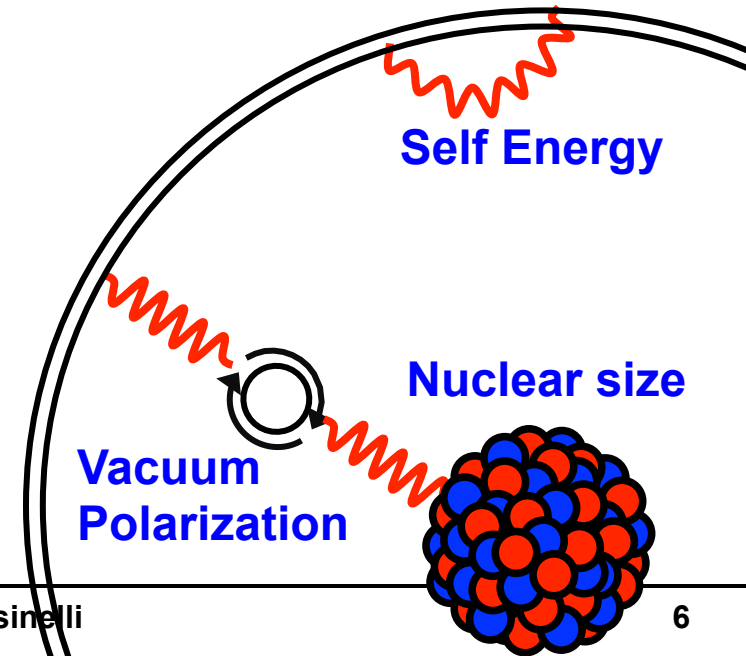
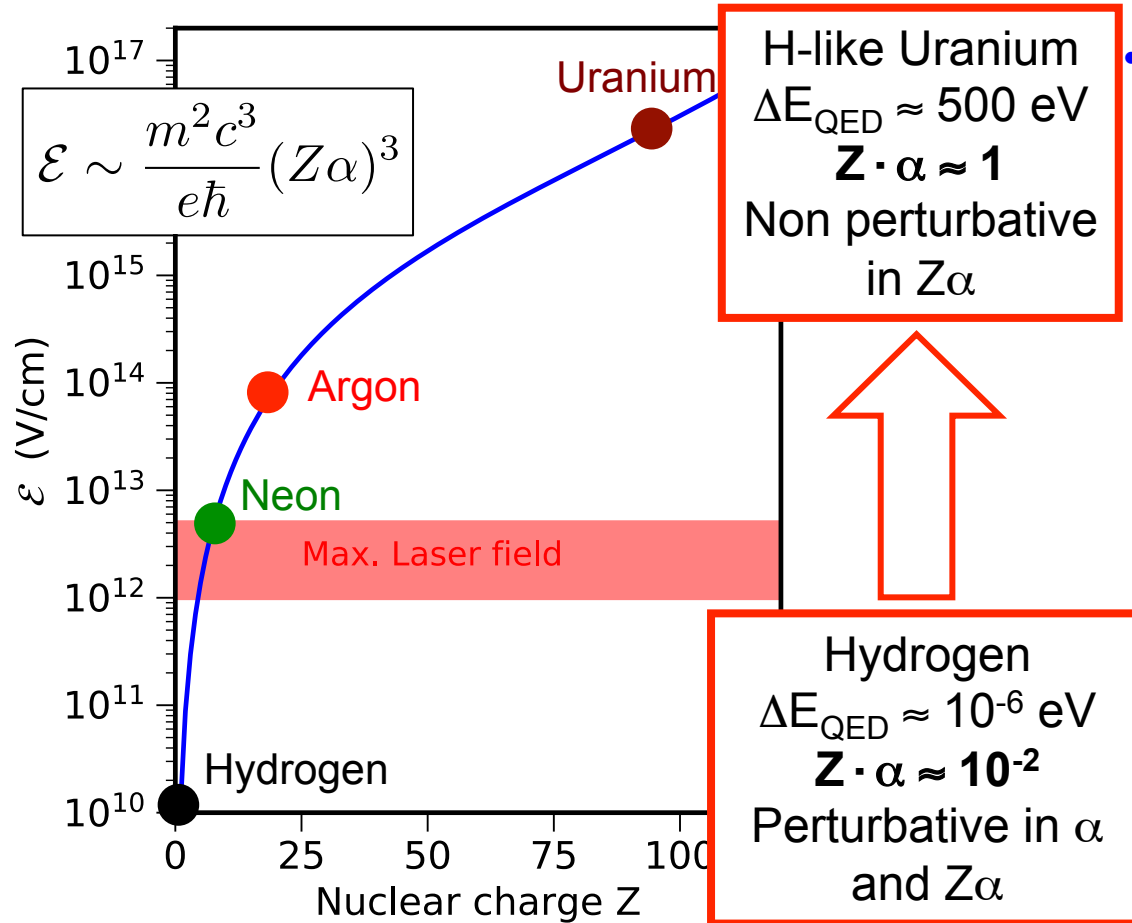
• **Simple atomic systems**

- accurate theoretical predictions
- simple spectra



Few electrons heavy charged ions

- Simple atomic systems
 - accurate theoretical predictions
 - simple spectra
- Strong electric field
 - Quantum Electrodynamic effects enhanced
 - Large perturbation during collisions with other atoms

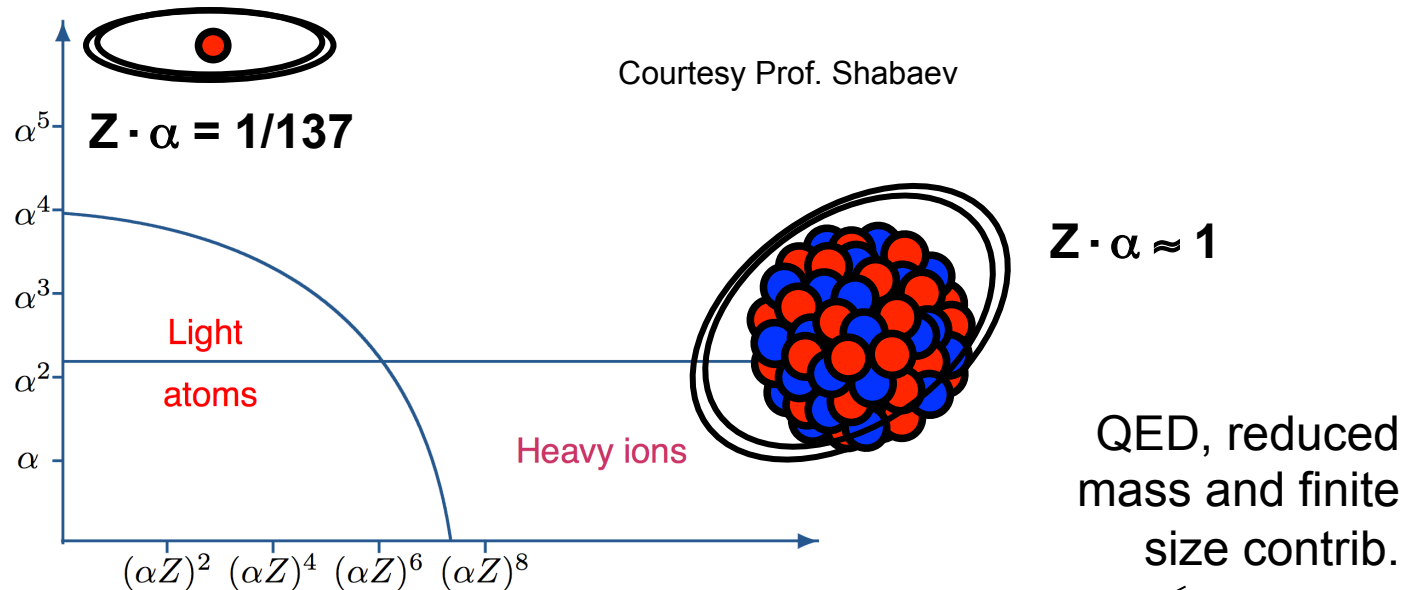


Tests of QED calculations at strong-coupling regime

α and $Z\alpha$ expansion

$$E^{nj} = mc^2 - mc^2 \frac{(Z\alpha)^2}{2n^2} + mc^2 (Z\alpha)^4 [a(n, j) + \alpha b(n, j) + \alpha^2 c(n, j) + \dots] + mc^2 (Z\alpha)^6 [a'(n, j) + \alpha b'(n, j) + \alpha^2 c'(n, j) + \dots] + \dots$$

M.I. Eides *et al.*, Phys. Rep. **342**, 63-261 (2001)
P.J. Mohr *et al.*, Rev. Mod. Phys. **88**, 035009 (2016).



Dirac eq.

α expansion,
all-order in $Z\alpha$

$$E^{nj} = \frac{mc^2}{\sqrt{1 + \frac{(Z\alpha)^2}{\left[n - j - 1/2 + \sqrt{(j + 1/2)^2 - (Z\alpha)^2}\right]^2}}} + \alpha F(Z\alpha) + \alpha^2 G(Z\alpha) + \dots$$

Non-perturbative bound system QED

First order in α

Self- energy

P.J. Mohr, Annals of Physics **88**, 26-51 (1974)



Vacuum polarization

N.L. Manakov *et al.*, Sov. Phys. JETP **68**, 673-679 (1989)

G. Soff *et al.*, Phys. Rev. A **38**, 5066-5075 (1988)



(Free particle QED renormalization in 1940-50)

Dirac eq. (1928)

$$E^{nj} = \frac{mc^2}{\sqrt{1 + \frac{(Z\alpha)^2}{\left[n - j - 1/2 + \sqrt{(j + 1/2)^2 - (Z\alpha)^2}\right]^2}}} + \alpha F(Z\alpha) + \alpha^2 G(Z\alpha) + \dots$$

Non-perturbative bound system QED

First order in α

Self- energy

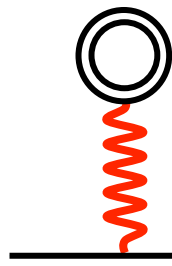
P.J. Mohr, Annals of Physics **88**, 26-51 (1974)



Vacuum polarization

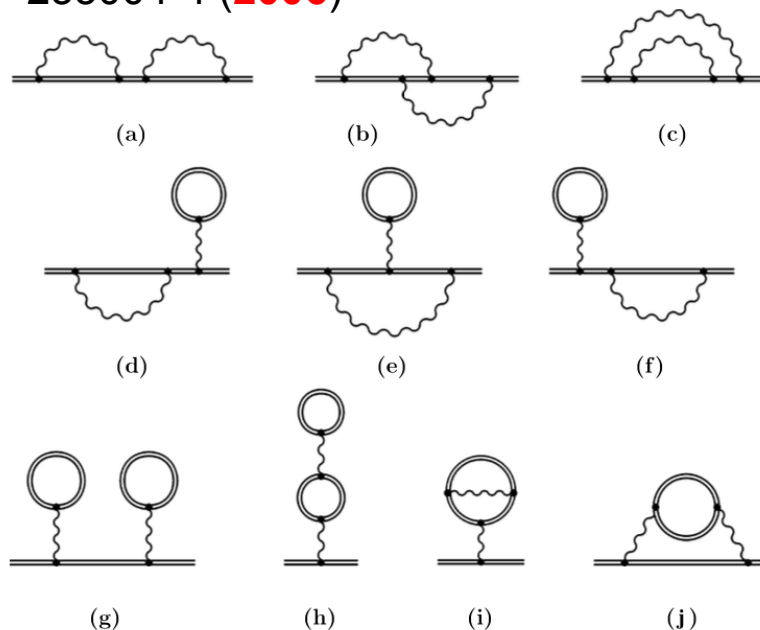
N.L. Manakov *et al.*, Sov. Phys. JETP **68**, 673-679 (1989)

G. Soff *et al.*, Phys. Rev. A **38**, 5066-5075 (1988)



Second order in α

V.A. Yerokhin *et al.*, Phys. Rev. Lett. **97**, 253004-4 (2006)



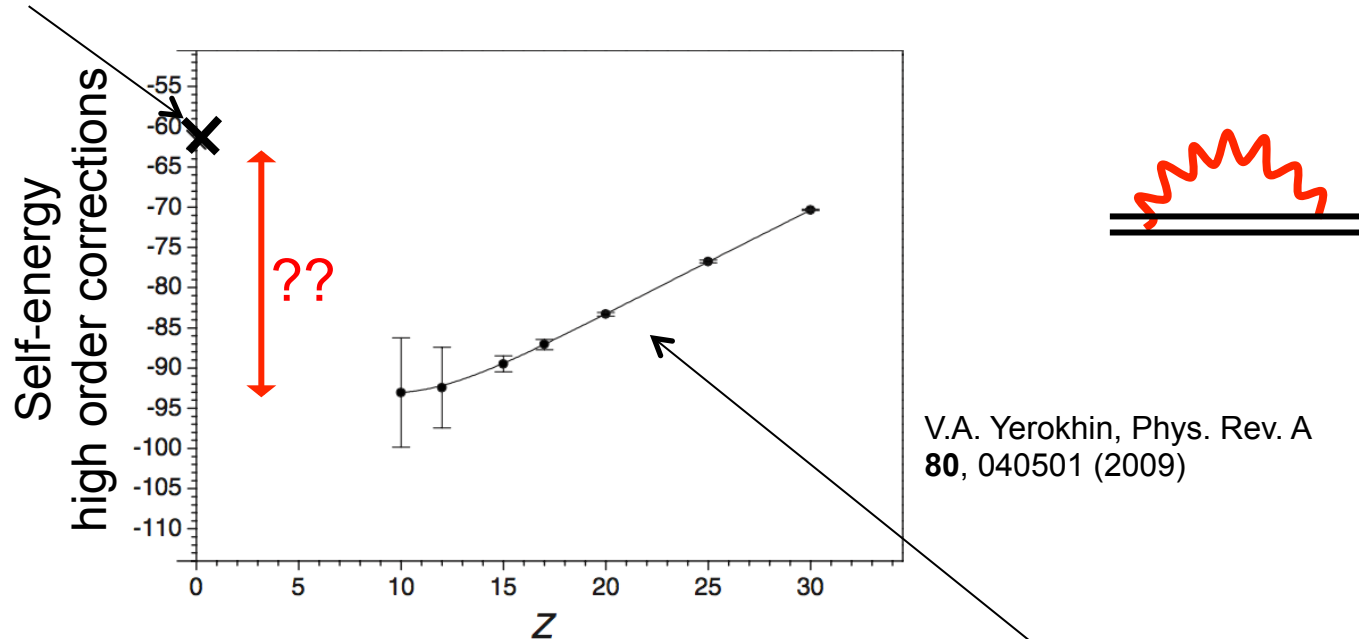
Dirac eq. (1928)

$$E^{nj} = \frac{mc^2}{\sqrt{1 + \frac{(Z\alpha)^2}{\left[n - j - 1/2 + \sqrt{(j + 1/2)^2 - (Z\alpha)^2}\right]^2}}} + \alpha F(Z\alpha) + \alpha^2 G(Z\alpha) + \dots$$

Self-energy predictions

α and $Z\alpha$ expansion

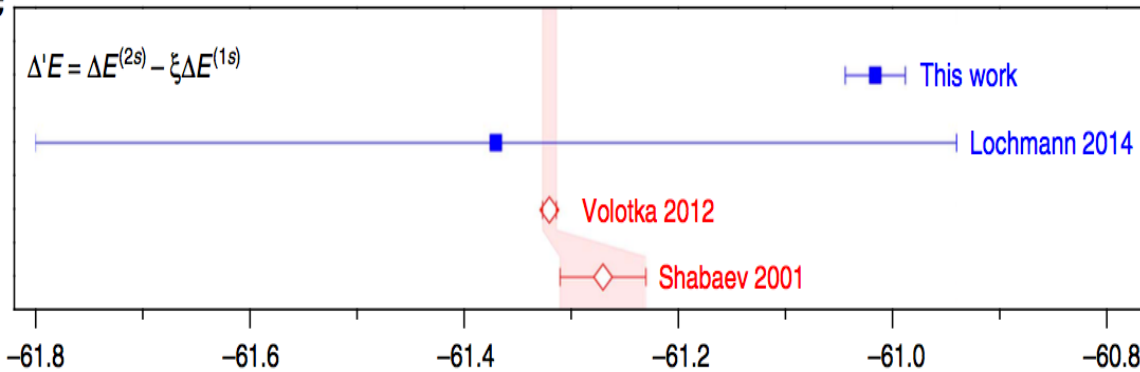
$$E^{nj} = mc^2 - mc^2 \frac{(Z\alpha)^2}{2n^2} + mc^2 (Z\alpha)^4 [a(n, j) + \alpha b(n, j) + \alpha^2 c(n, j) + \dots] + mc^2 (Z\alpha)^6 [a'(n, j) + \alpha b'(n, j) + \alpha^2 c'(n, j) + \dots] + \dots$$



α expansion,
all-order in $Z\alpha$

$$E^{nj} = \frac{mc^2}{\sqrt{1 + \frac{(Z\alpha)^2}{\left[n - j - 1/2 + \sqrt{(j + 1/2)^2 - (Z\alpha)^2}\right]^2}}} + \alpha F(Z\alpha) + \alpha^2 G(Z\alpha) + \dots$$

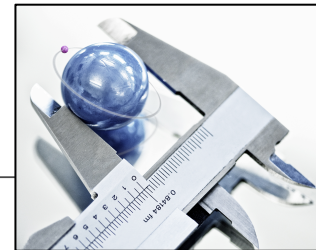
Hyperfine structure of H- and Li-like Bismuth



Laser spectroscopy of stored heavy ions

J. Ullmann *et al.*, Nat. Commun. **8**, 15484 (2017)

Proton radius puzzle

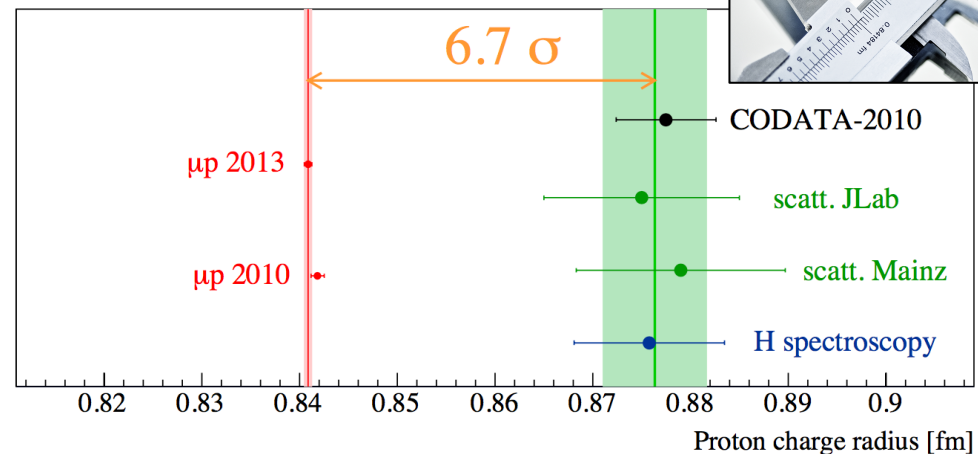


Strong disagreement between **hydrogen** and **muonic hydrogen** spectroscopy

$$m_{\mu} = 200 m_e$$

$$\mathcal{E} \sim \frac{m^2 c^3}{e \hbar} (Z\alpha)^3$$

$$r_n \sim \frac{\hbar n^2}{mcZ\alpha}$$



A. Antognini *et al.*, Science **339**, 417-420 (2013)

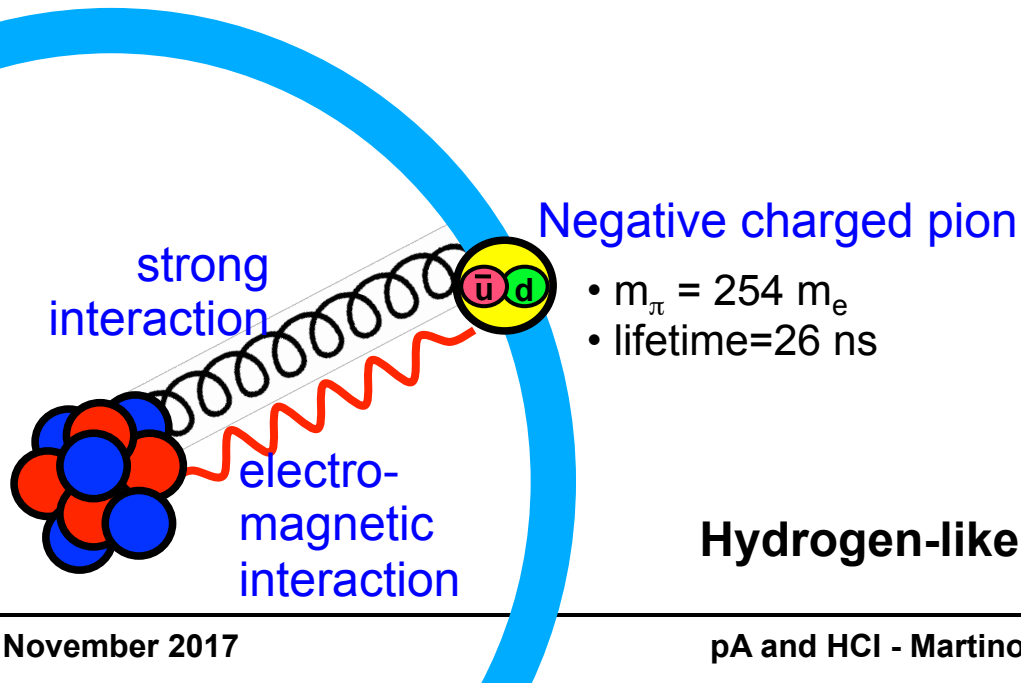
R. Pohl *et al.*, Nature **466**, 213-216 (2010)

Summary

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 - **Study of strong interaction force at low energy**
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Pionic atoms

FAQ: What is a pion?



Hydrogen-like atoms

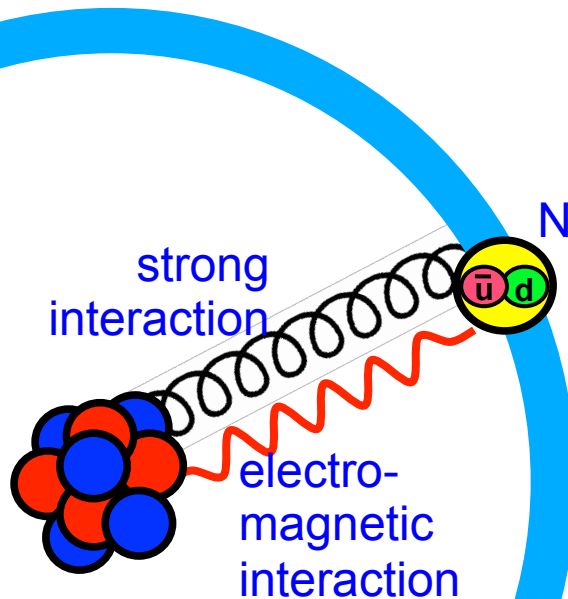
Pionic atoms

- Simple atomic systems
 - accurate theoretical predictions
 - simple spectra
- Strong electric field
 - Quantum Electrodynamic effects enhanced
- Presence of strong interaction
- Small Bohr radius

$$E^{nl} = \frac{mc^2}{\sqrt{1 + \frac{(Z\alpha)^2}{\left[n - l - 1/2 + \sqrt{(l + 1/2)^2 - (Z\alpha)^2}\right]^2}}}$$

$$\mathcal{E} \sim \frac{m^2 c^3}{e\hbar} (Z\alpha)^3$$

$$r_n \sim \frac{\hbar n^2}{mcZ\alpha}$$



Negative charged pion

- $m_\pi = 254 m_e$
- lifetime=26 ns

Hydrogen-like atoms

Pionic atoms

- Electromagnetic interaction $\rightarrow$ bound system

$$E_n = mc^2 \frac{(Z\alpha)^2}{2n^2} + \mathcal{O}[(Z\alpha)^4]$$

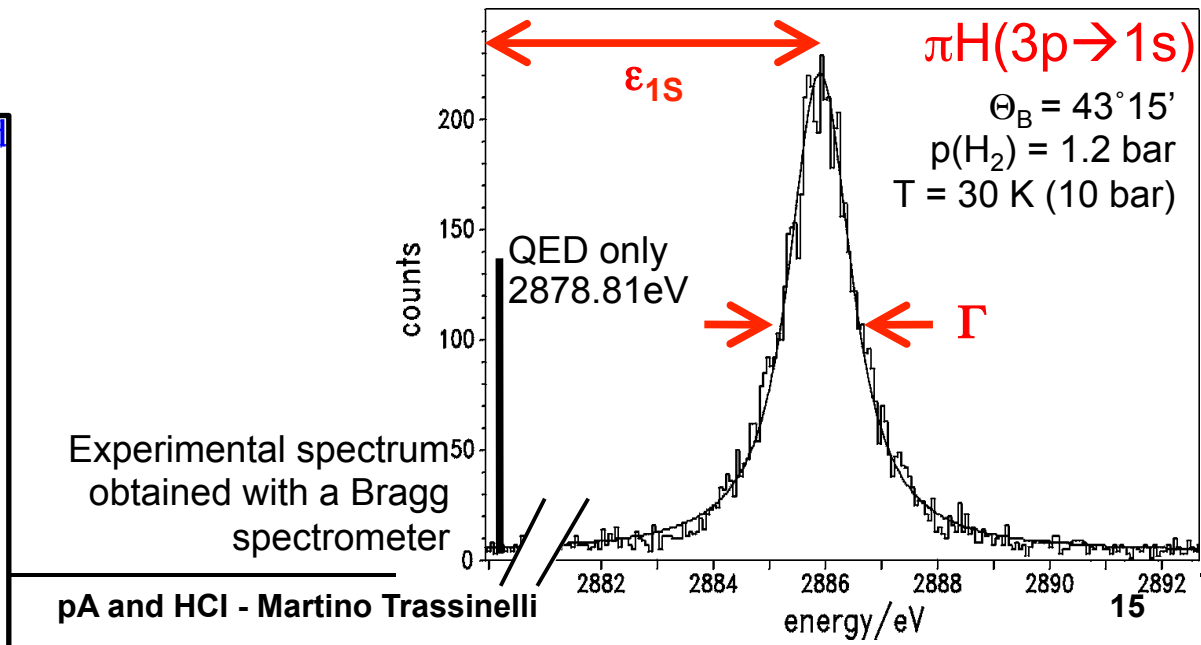
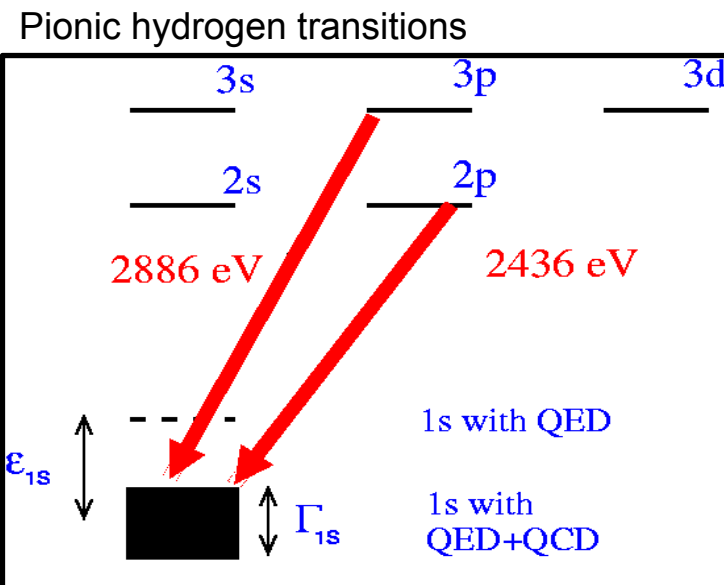
- **Strong interaction!!**

Attraction or repulsion

$\rightarrow$ shift of the energy transition ϵ

Reaction with the nucleus $\rightarrow$ ground state instable

- lifetime of the atom not infinite
- non-zero linewidth Γ of the atomic transitions



Strong interaction Lagrangians

High energy ————— Low energy

Quantum ChromoDynamics Lagrangian

$$\mathcal{L}_{\text{QCD}} = \bar{\psi}_i (i\gamma^\mu (D_\mu)_{ij} - m \delta_{ij}) \psi_j - \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a$$

$$= \bar{\psi}_i (i\gamma^\mu \partial_\mu - m) \psi_i - g \bar{\psi}_i \gamma^\mu T_{ij}^a \psi_j A_\mu^a - \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a,$$

Chiral perturbation theory
Lagrangian (4th order dev.)

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_\pi^{(p^2)} + \mathcal{L}_\pi^{(e^2)} + \mathcal{L}_\pi^{(e^2 p^2)} + \mathcal{L}_N^{(p)} + \mathcal{L}_N^{(p^2)} + \mathcal{L}_N^{(p^3)} + \mathcal{L}_N^{(e^2)} + \mathcal{L}_N^{(e^2 p)} + \mathcal{L}_\gamma,$$

$$\mathcal{L}_\pi^{(p^2)} + \mathcal{L}_\pi^{(e^2)} + \mathcal{L}_\gamma = \frac{F^2}{4} \langle d^\mu U^\dagger d_\mu U + \chi^\dagger U + U^\dagger \chi \rangle + Z F^4 \langle Q U Q U^\dagger \rangle$$

$$- \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} (\partial_\mu A^\mu)^2,$$

Nucleon propagator in ChPT with pion
and other propagator contributions

J. Gasser *et al.*, Nuclear Physics B 307, 779 (1988)

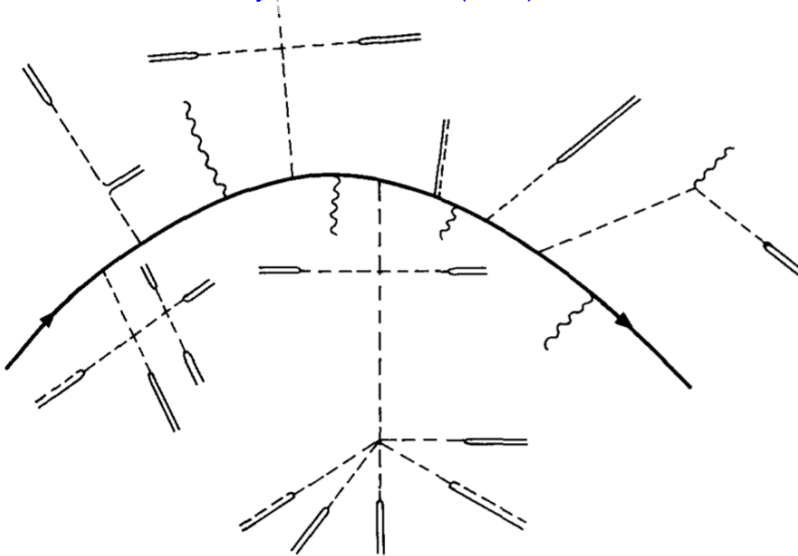


Fig. 1. A typical term in the expansion (3.7) of the nucleon propagator. — nucleon; --- pions; ~~~ vector current; == axial vector current; -o- pseudoscalar density; == scalar density.

$$= F^2 \{ \langle d^\mu U^\dagger d_\mu U \rangle (k_1 \langle Q^2 \rangle + k_2 \langle Q U Q U^\dagger \rangle) + k_4 \langle d^\mu U^\dagger Q U \rangle \langle d_\mu U Q U^\dagger \rangle$$

$$+ k_3 (\langle d^\mu U^\dagger Q U \rangle \langle d_\mu U^\dagger Q U \rangle + \langle d^\mu U Q U^\dagger \rangle \langle d_\mu U Q U^\dagger \rangle) \},$$

$$\bar{\psi} \left\{ i \not{D} - m + \frac{1}{2} g \not{\psi} \gamma_5 \right\} \psi,$$

$$\bar{\psi} \left\{ c_1 \langle \chi_+ \rangle - \frac{c_2}{4m^2} \langle u_\mu u_\nu \rangle D^\mu D^\nu + \text{h.c.} + \frac{c_3}{2} \langle u_\mu u^\mu \rangle \right.$$

$$\left. + \frac{i}{4} c_4 \sigma^{\mu\nu} [u_\mu, u_\nu] + c_5 \hat{\chi}_+ \right\} \psi,$$

V. Baru *et al.*, Nucl. Phys.
A 872, 69-116 (2011)

$$\frac{i}{2m} \bar{\psi} \{ d_5 [\chi_-, u_\mu] D^\mu \} \psi + \text{h.c.},$$

$$F^2 \bar{\psi} \{ f_1 \langle \hat{Q}_+^2 - Q_-^2 \rangle + f_2 \langle Q_+ \rangle \hat{Q}_+ \} \psi,$$

$$+ \frac{F^2}{2} \bar{\psi} \{ g_1 \langle Q_+^2 - Q_-^2 \rangle \gamma^\mu \gamma_5 u_\mu + g_2 \langle Q_+ \rangle^2 \gamma^\mu \gamma_5 u_\mu \} \psi$$

$$+ \frac{i F^2}{2m} \bar{\psi} \{ g_6 \langle Q_+ \rangle \langle Q_- u_\mu \rangle D^\mu + g_7 \langle Q_+ u_\mu \rangle Q_- D^\mu$$

$$+ g_8 \langle Q_- u_\mu \rangle Q_+ D^\mu \} \psi + \text{h.c.},$$

Strong interaction Lagrangians

High energy — — — — — Low energy

Quantum ChromoDynamics Lagrangian

Chiral perturbation theory
Lagrangian (4th order dev.)

$$\mathcal{L}_{\text{QCD}} = \bar{\psi}_i (i\gamma^\mu (D_\mu)_{ij} - m \delta_{ij}) \psi_j - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

$$= \bar{\psi}_i (i\gamma^\mu \partial_\mu - m) \psi_i - g \bar{\psi}_i \gamma^\mu T_{ij}^a \psi_j A_\mu^a - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu},$$

1 coupling constant

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_\pi^{(p^2)} + \mathcal{L}_\pi^{(e^2)} + \mathcal{L}_\pi^{(e^2 p^2)} + \mathcal{L}_N^{(p)} + \mathcal{L}_N^{(p^2)} + \mathcal{L}_N^{(p^3)} + \mathcal{L}_N^{(e^2)} + \mathcal{L}_N^{(e^2 p)} + \mathcal{L}_\gamma,$$

$$\mathcal{L}_\pi^{(p^2)} + \mathcal{L}_\pi^{(e^2)} + \mathcal{L}_\gamma = \frac{F^2}{4} \langle d^\mu U^\dagger d_\mu U + \chi^\dagger U + U^\dagger \chi \rangle + Z F^4 \langle Q U Q U^\dagger \rangle$$

$$- \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} (\partial_\mu A^\mu)^2,$$

$$\mathcal{L}_\pi^{(e^2 p^2)} = F^2 \{ \langle d^\mu U^\dagger d_\mu U \rangle (k_1 \langle Q^2 \rangle + k_2 \langle Q U Q U^\dagger \rangle) + k_4 \langle d^\mu U^\dagger Q U \rangle \langle d_\mu U Q U^\dagger \rangle$$

$$+ k_3 \langle d^\mu U^\dagger Q U \rangle \langle d_\mu U^\dagger Q U \rangle + \langle d^\mu U Q U^\dagger \rangle \langle d_\mu U Q U^\dagger \rangle \},$$

$$\mathcal{L}_N^{(p)} = \bar{\Psi} \left\{ i \not{D} - m + \frac{1}{2} g \not{A} \gamma_5 \right\} \Psi,$$

18 and more

low energy coupling constants (LEC)
(depending on the expansion order)

$$\mathcal{L}_N^{(p^2)} = \bar{\Psi} \left\{ c_1 \langle \chi_+ \rangle - \frac{c_2}{4m^2} \langle u_\mu u_\nu \rangle D^\mu D^\nu + \text{h.c.} + \frac{c_3}{Z} \langle u_\mu u^\mu \rangle \right.$$

$$\left. + \frac{i}{4} c_4 \sigma^{\mu\nu} [u_\mu, u_\nu] + c_5 \chi_+ \right\} \Psi,$$

$$\mathcal{L}_N^{(p^3)} = \frac{i}{2m} \bar{\Psi} \{ d_5 [\chi_-, u_\mu] D^\mu \} \Psi + \text{h.c.},$$

$$\mathcal{L}_N^{(e^2)} = F^2 \bar{\Psi} \{ f_1 \langle \hat{Q}_+^2 - Q_-^2 \rangle + f_2 \langle Q_+ \rangle \hat{Q}_+ \} \Psi,$$

$$\mathcal{L}_N^{(e^2 p)} = \frac{F^2}{2} \bar{\Psi} \{ g_1 \langle Q_+^2 - Q_-^2 \rangle \gamma^\mu \gamma_5 u_\mu + g_2 \langle Q_+ \rangle^2 \gamma^\mu \gamma_5 u_\mu \} \Psi$$

$$+ \frac{i F^2}{2m} \bar{\Psi} \{ g_6 \langle Q_+ \rangle \langle Q - u_\mu \rangle D^\mu + g_7 \langle Q + u_\mu \rangle Q - D^\mu$$

$$+ g_8 \langle Q - u_\mu \rangle Q + D^\mu \} \Psi + \text{h.c.},$$

V. Baru et al., Nucl. Phys.
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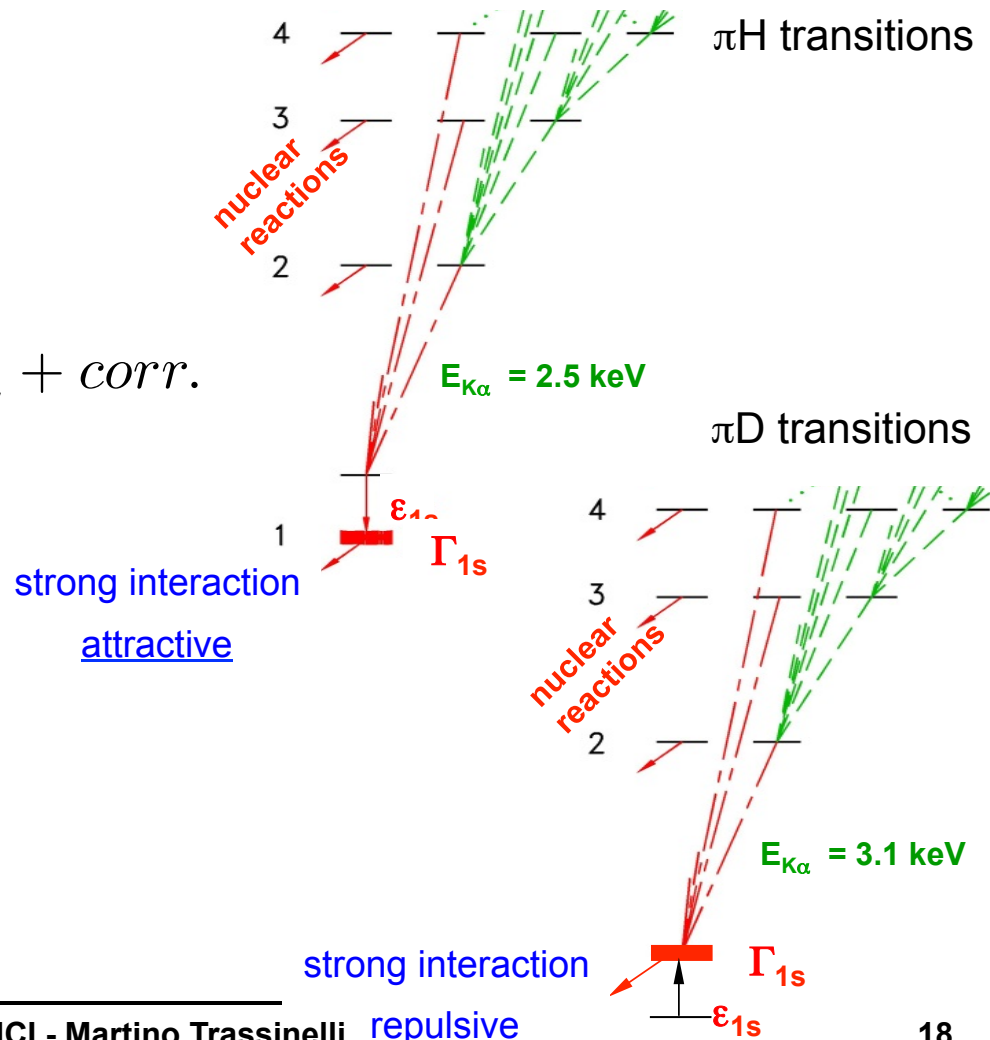
Need of experimental experiment to
measure the LECs and test the
theory!!

Strong interaction in simple bound systems

Shift and width related to the pion-nucleus scattering length and cross sections

$$\pi H \begin{cases} \epsilon_{1s}^{\pi H} \propto a_{\pi^- p \rightarrow \pi^- p} + corr. \\ \Gamma_{1s}^{\pi H} \propto (a_{\pi^- p \rightarrow \pi^0 n})^2 + corr. \end{cases}$$

$$\pi D \begin{cases} \epsilon_{1s}^{\pi D} \propto a_{\pi^- p \rightarrow \pi^- p} + a_{\pi^- n \rightarrow \pi^- n} + corr. \\ \Gamma_{1s}^{\pi D} \propto \Im m(a_{\pi d \rightarrow \pi d}) \end{cases}$$



Strong interaction in simple bound systems

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Basic pion-nucleon
interaction properties

Pion production
cross section from
p-p collisions

Highly charged ions and pionic atoms

- Simple atomic systems

$$E_n = mc^2 \frac{(Z\alpha)^2}{2n^2} + \mathcal{O}[(Z\alpha)^4]$$

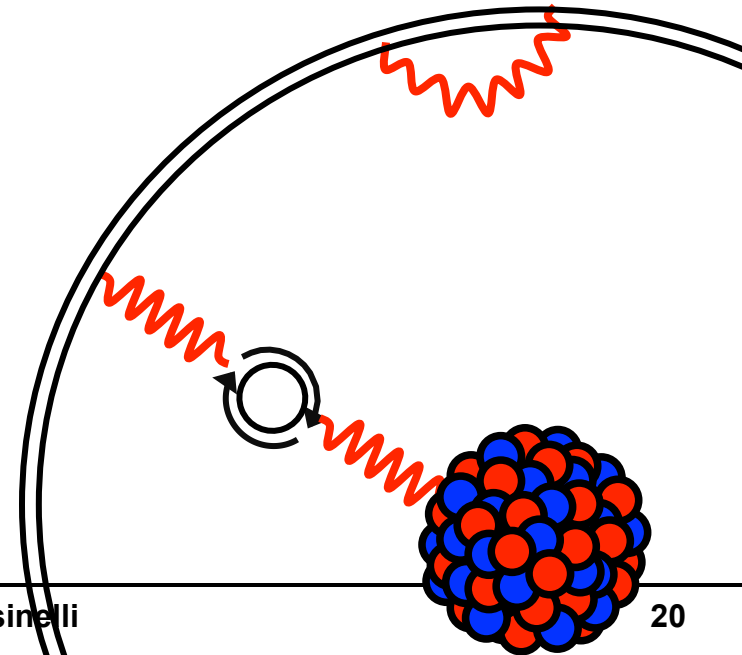
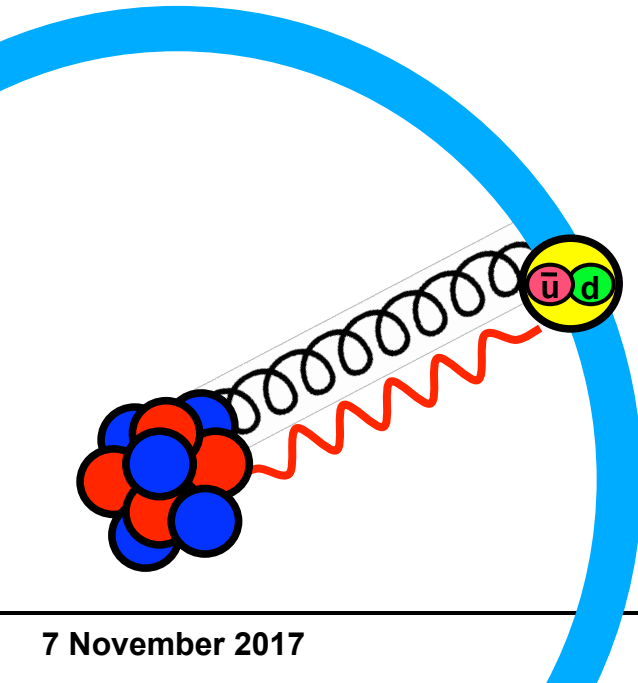
➡ X-ray spectroscopy

- Strong electric field

$$\mathcal{E} \sim \frac{m^2 c^3}{e\hbar} (Z\alpha)^3$$

- Small Bohr radius

$$r_n \sim \frac{\hbar n^2}{mcZ\alpha}$$

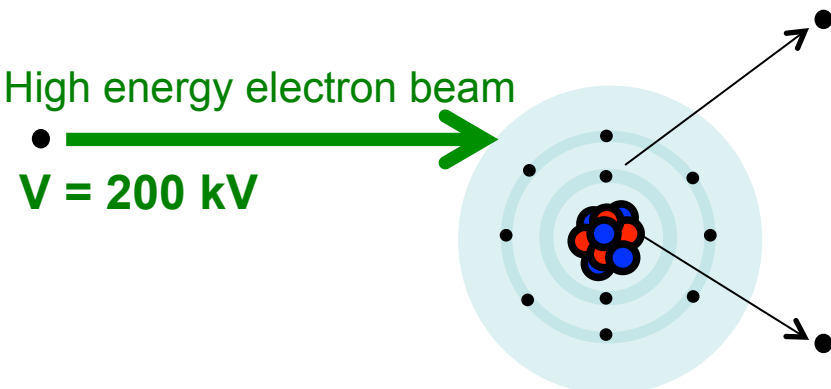


Summary

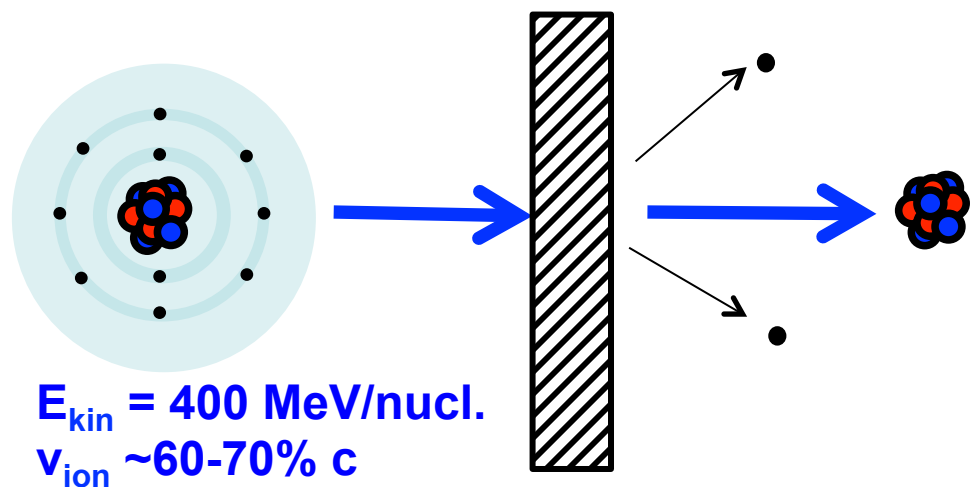
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Highly charged ion production and studies

Electron impact and trapping

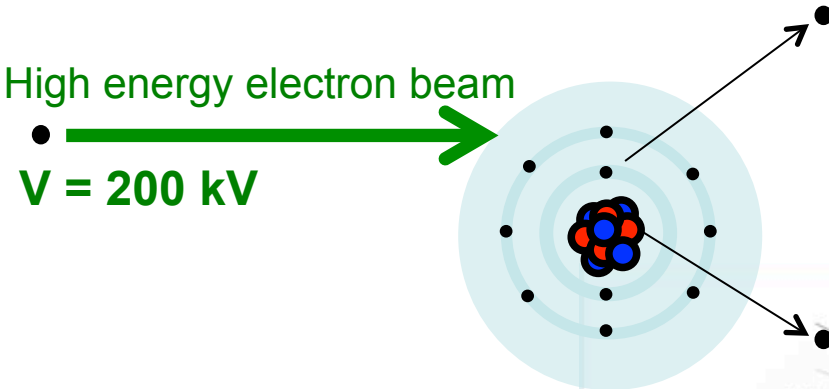


Acceleration and stripping

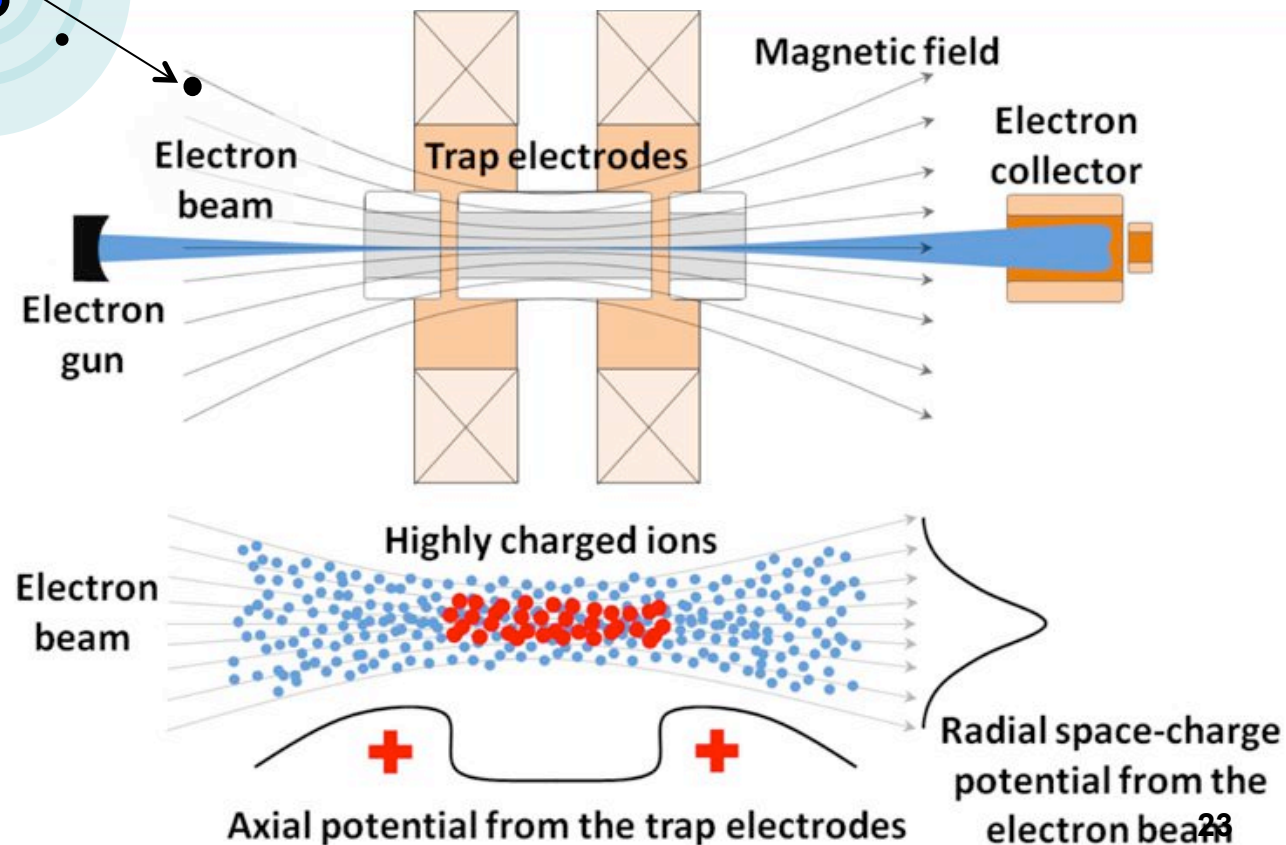


Highly charged ion production and studies

Electron impact and trapping

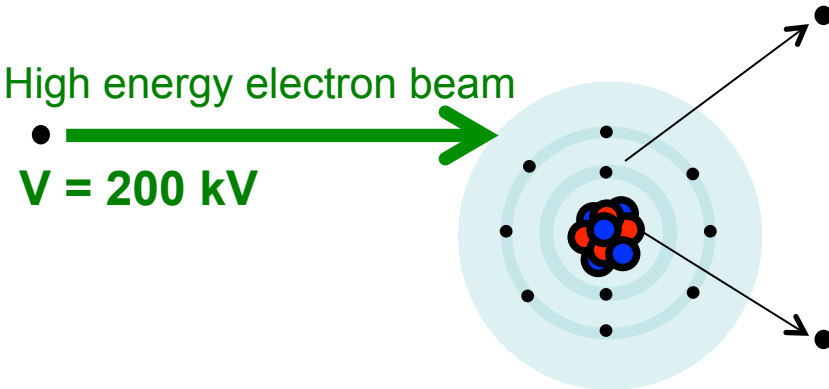


Electron Beam Ion Trap (EBIT) principle

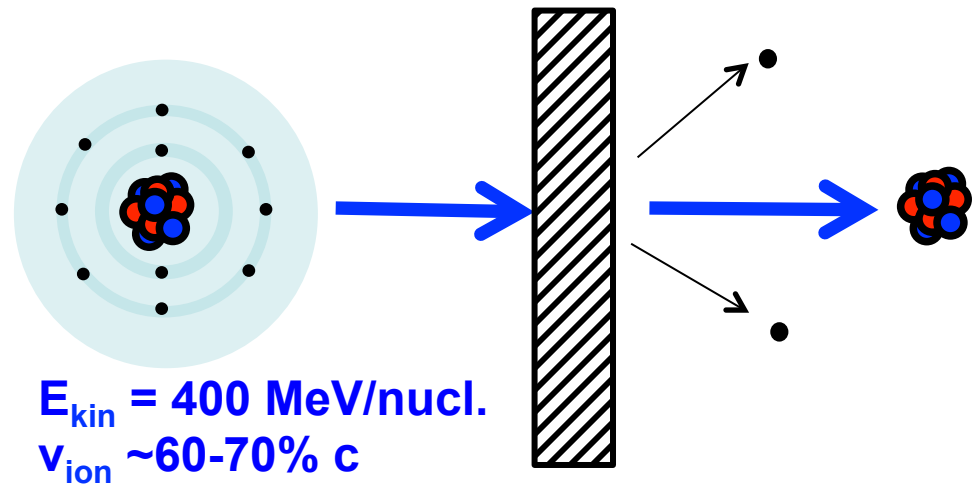


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Electron impact and trapping



Acceleration and stripping



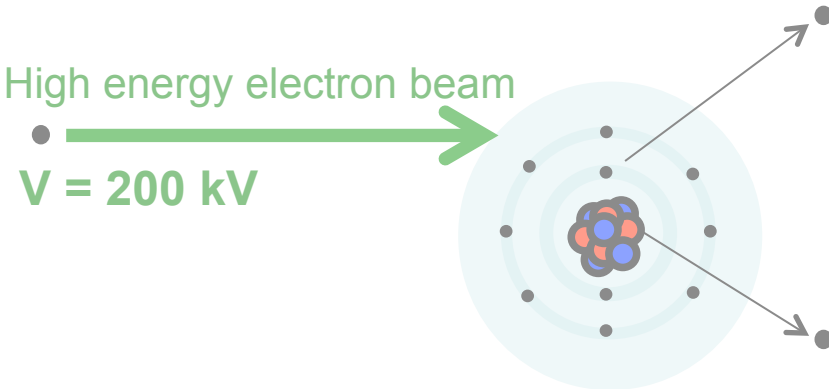
Solid target (graphite, aluminum, copper, ...)
 Thickness $\sim 100 \text{ mg/cm}^2$

- Max Planch Inst., Heidelberg, Germany
- NIST, Gaithersburg, USA
- ...

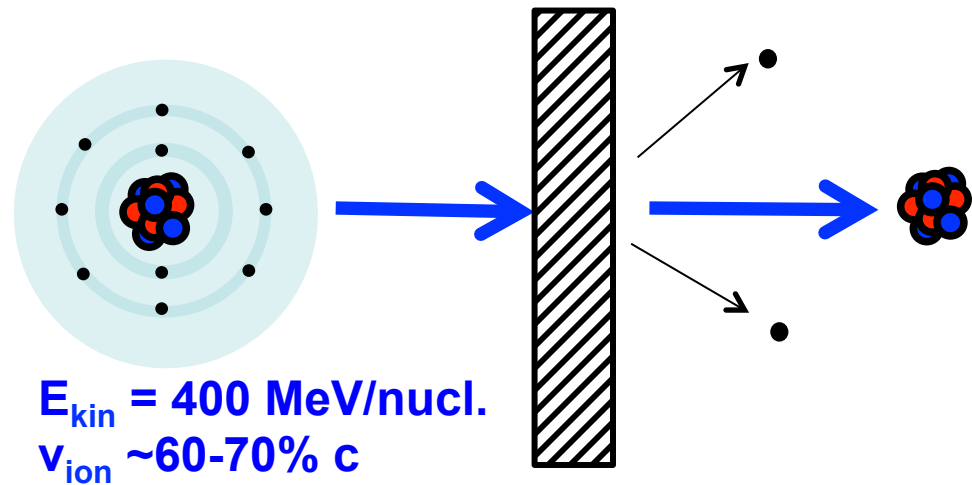
- GSI, Darmstadt, Germany
- IMP, Lanzhou, China
- ...

Highly charged ion production and studies

Electron impact and trapping



Acceleration and stripping

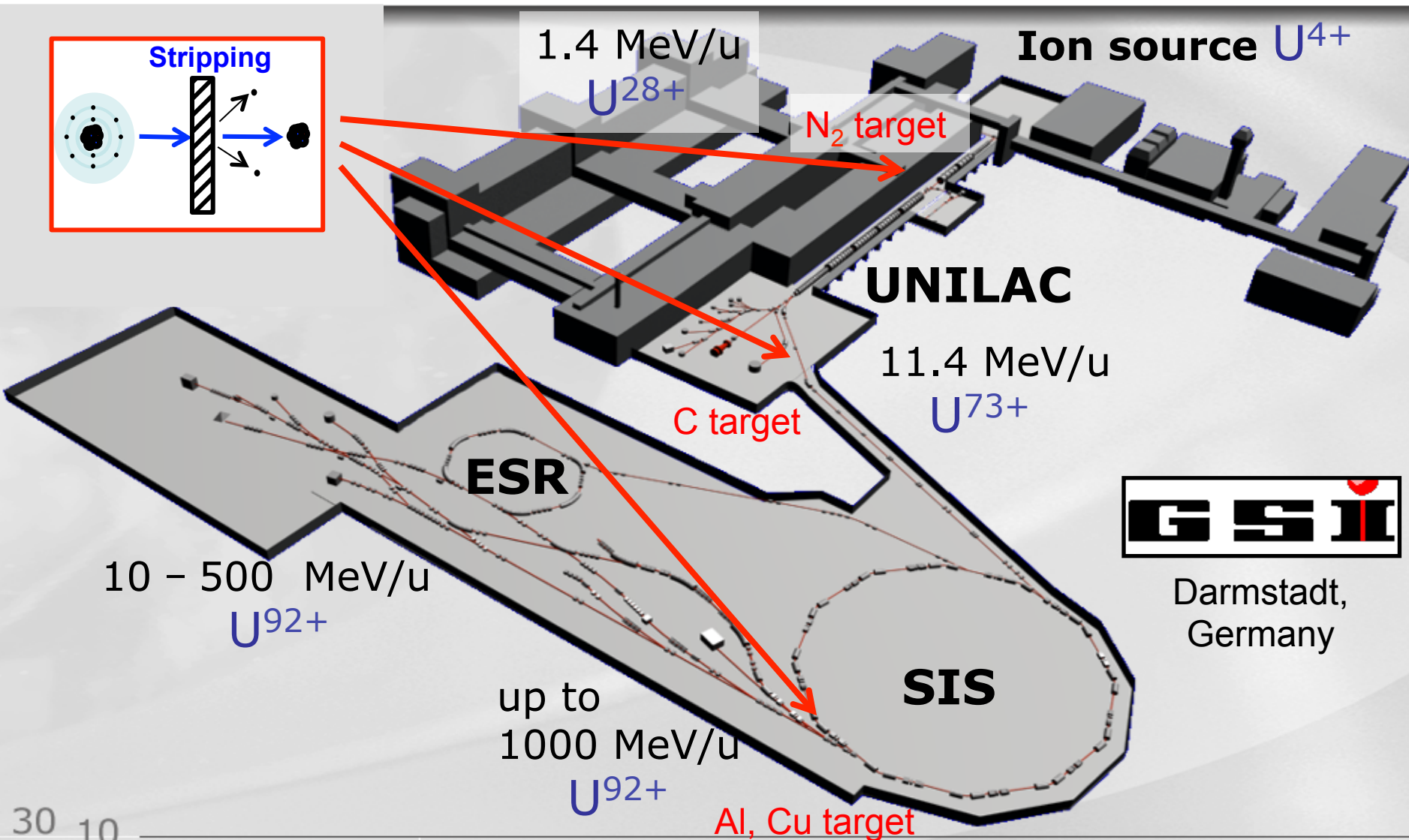


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Highly charged ion production

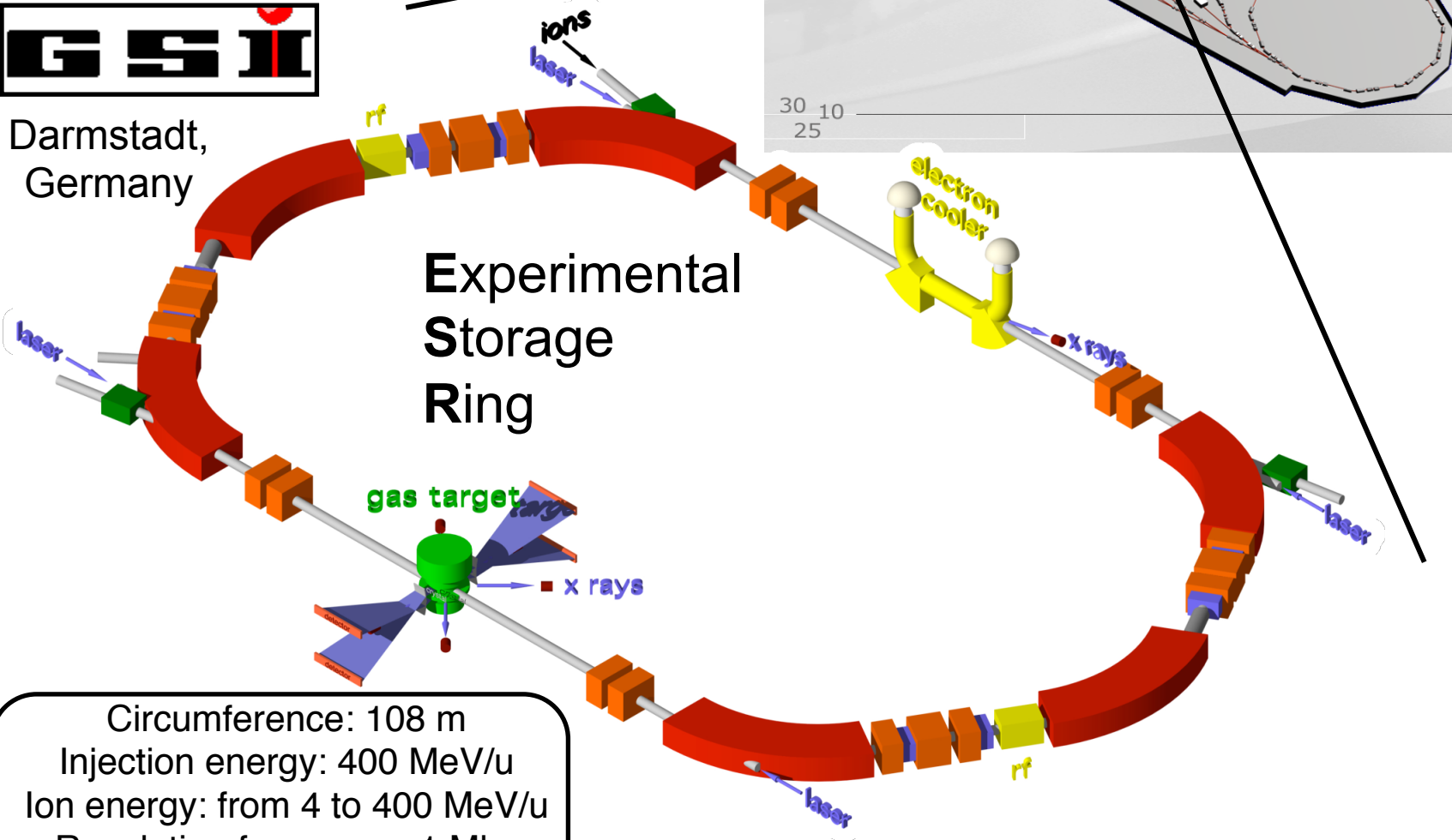


ESR storage ring at GSI



Darmstadt,
Germany

Experimental Storage Ring



Circumference: 108 m

Injection energy: 400 MeV/u

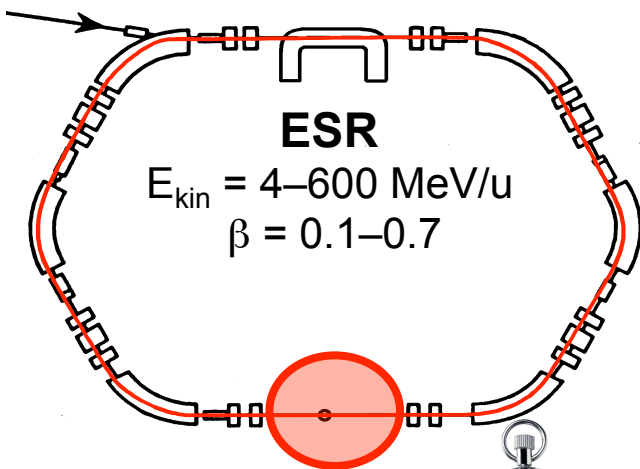
Ion energy: from 4 to 400 MeV/u

Revolution frequency: 1 Mhz

Stored ions: $\sim 10^8$

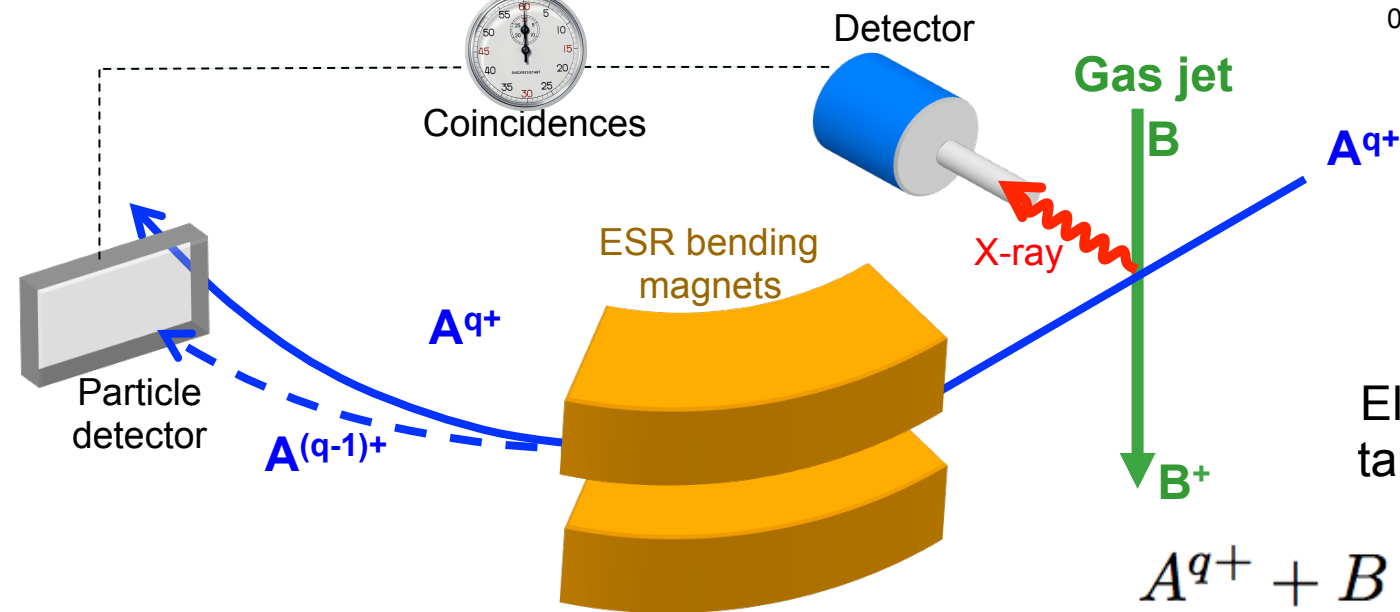
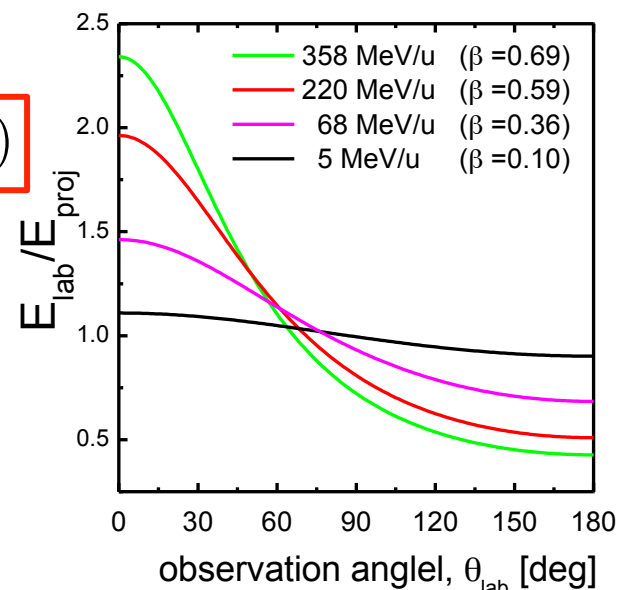
Pressure = 10^{-10-11} mbar

Experiments at the gas jet-target

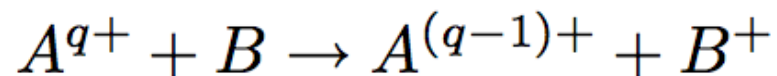


Relativistic Doppler shift

$$E_{\text{ion}} = E_{\text{lab}} \gamma (1 - \beta \cos \theta)$$



Electron transfer from the target atom to the fast ion

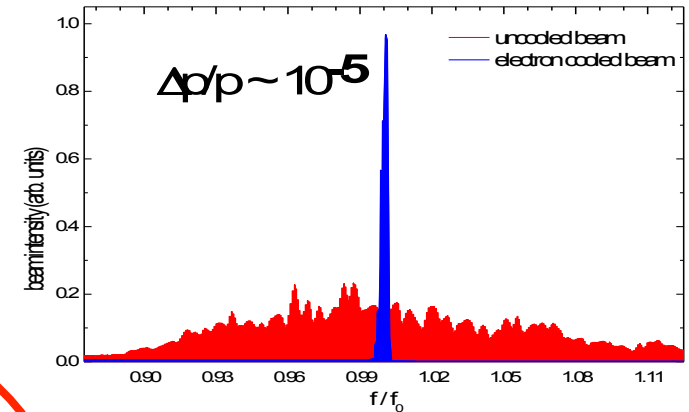
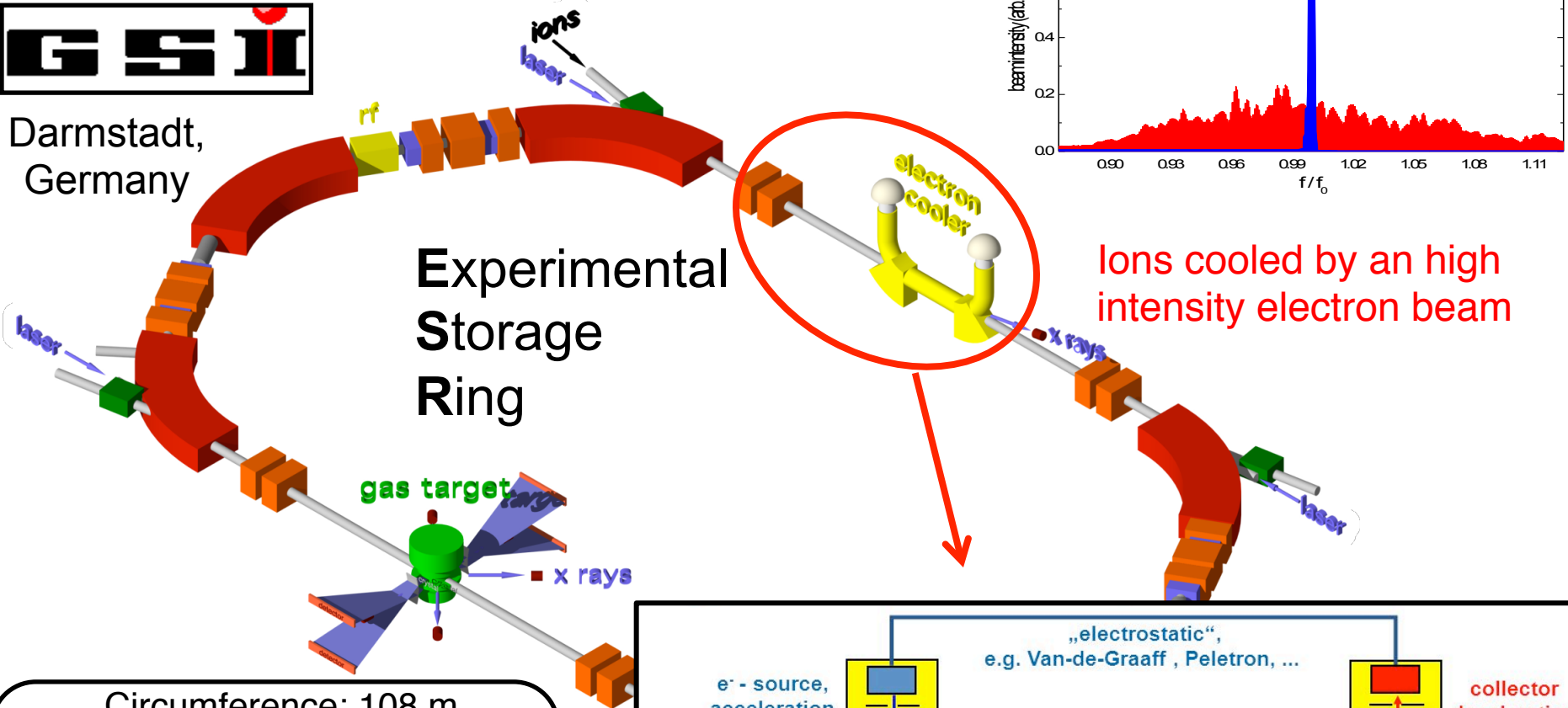


ESR storage ring at GSI



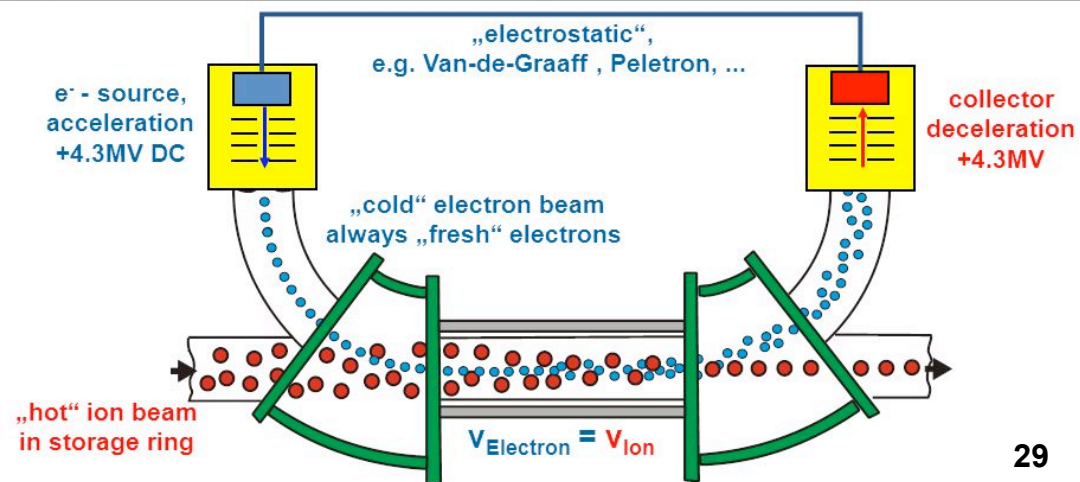
Darmstadt,
Germany

Experimental Storage Ring

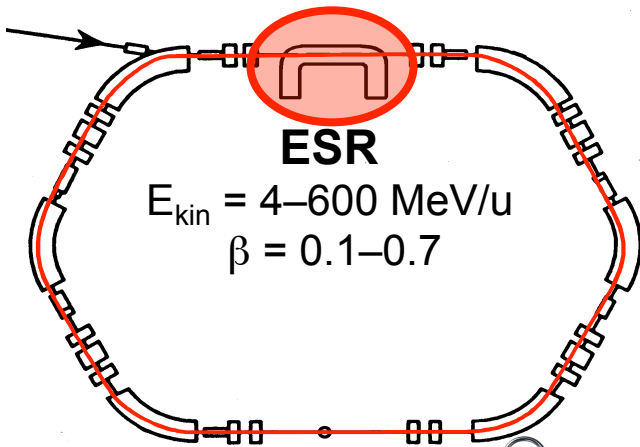


Ions cooled by an high
intensity electron beam

Circumference: 108 m
Injection energy: 400 MeV/u
Ion energy: from 4 to 400 MeV/u
Revolution frequency: 1 Mhz
Stored ions: $\sim 10^8$
Pressure = 10^{-10-11} mbar

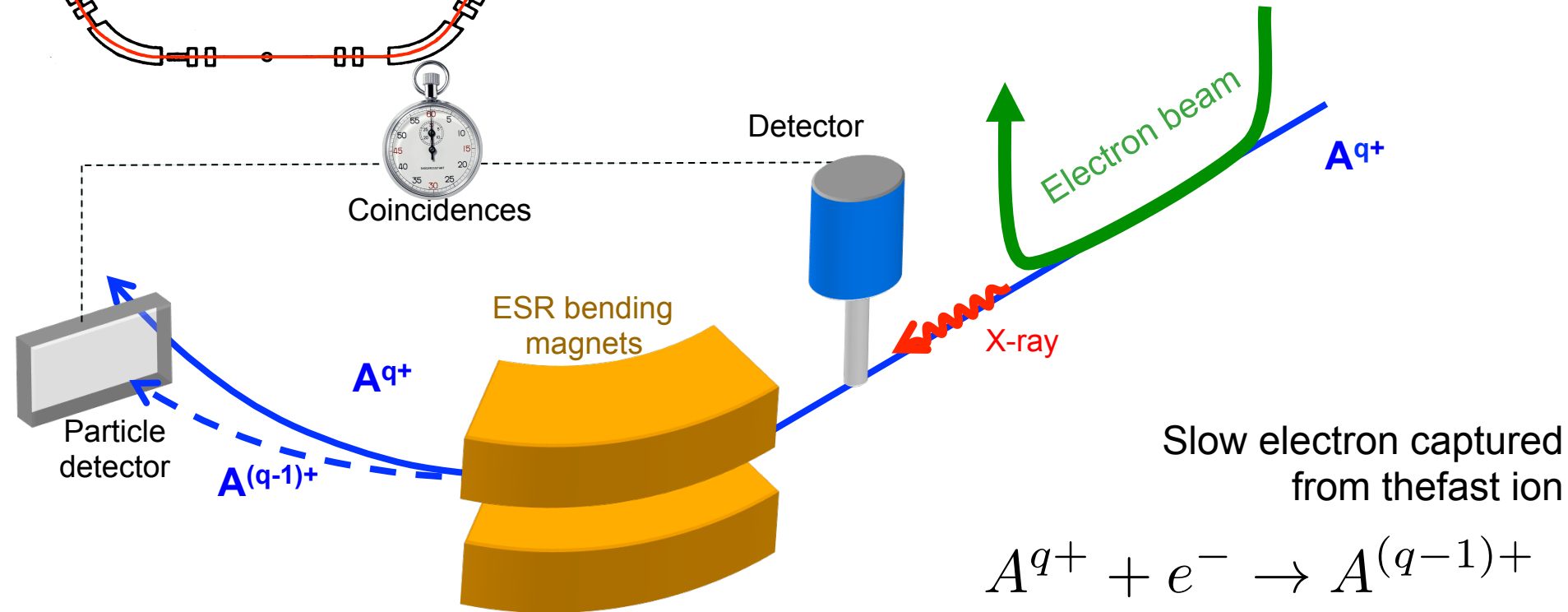
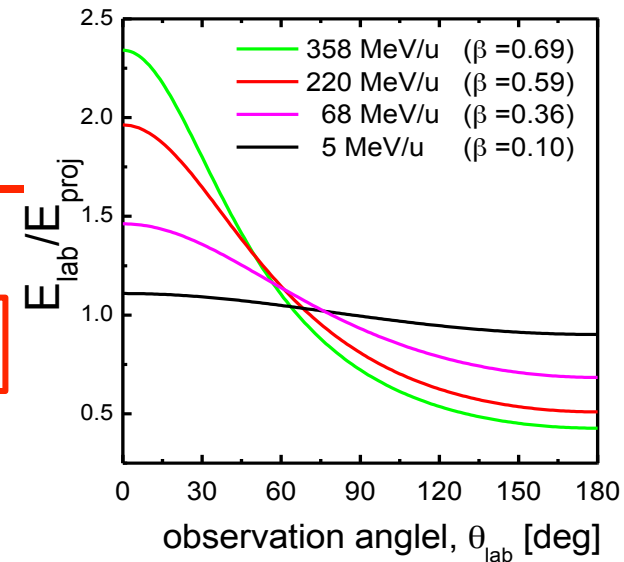


Experiments at e-cooler



Relativistic Doppler shift

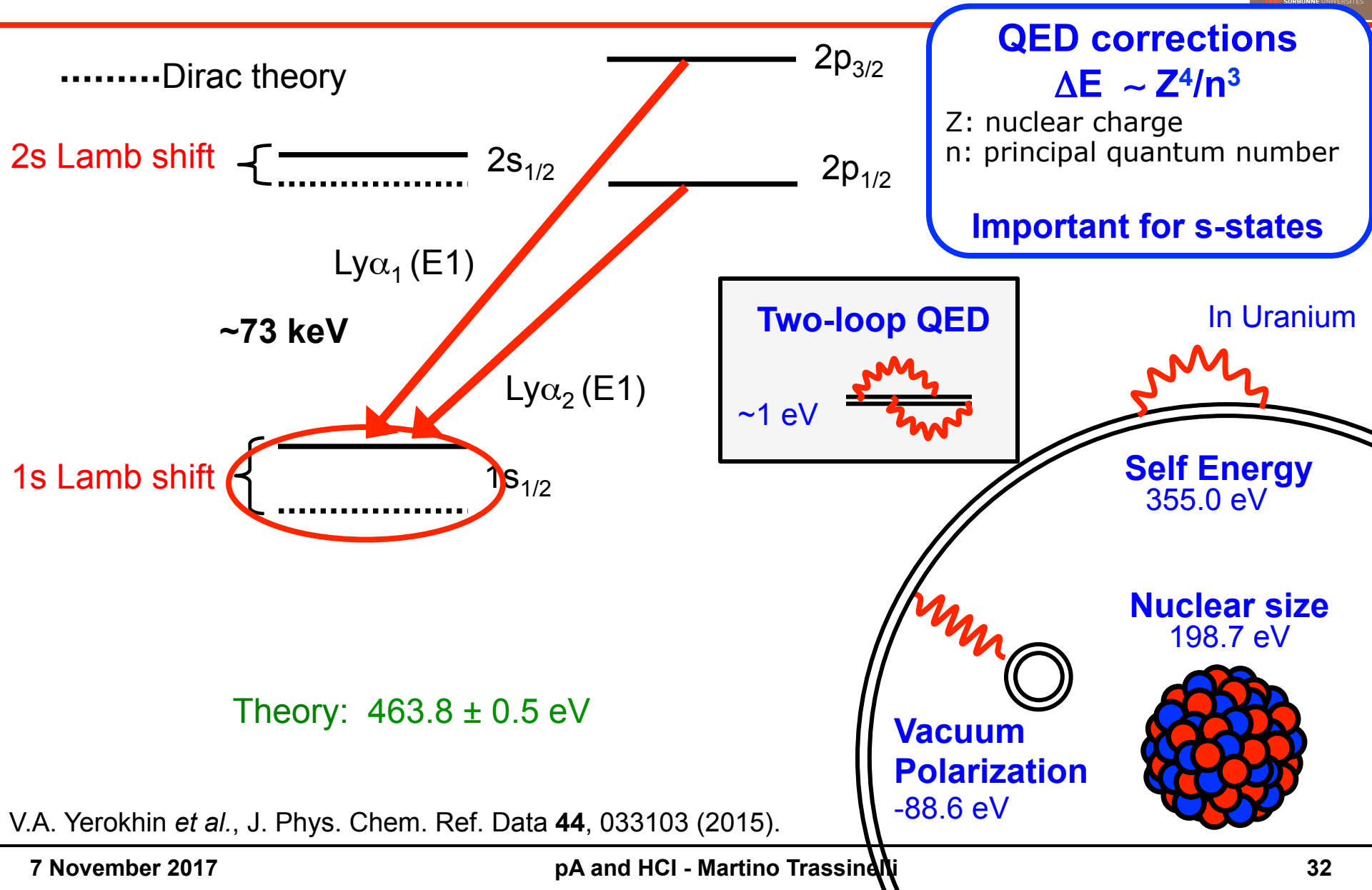
$$E_{\text{ion}} = E_{\text{lab}} \gamma (1 - \beta \cos \theta)$$



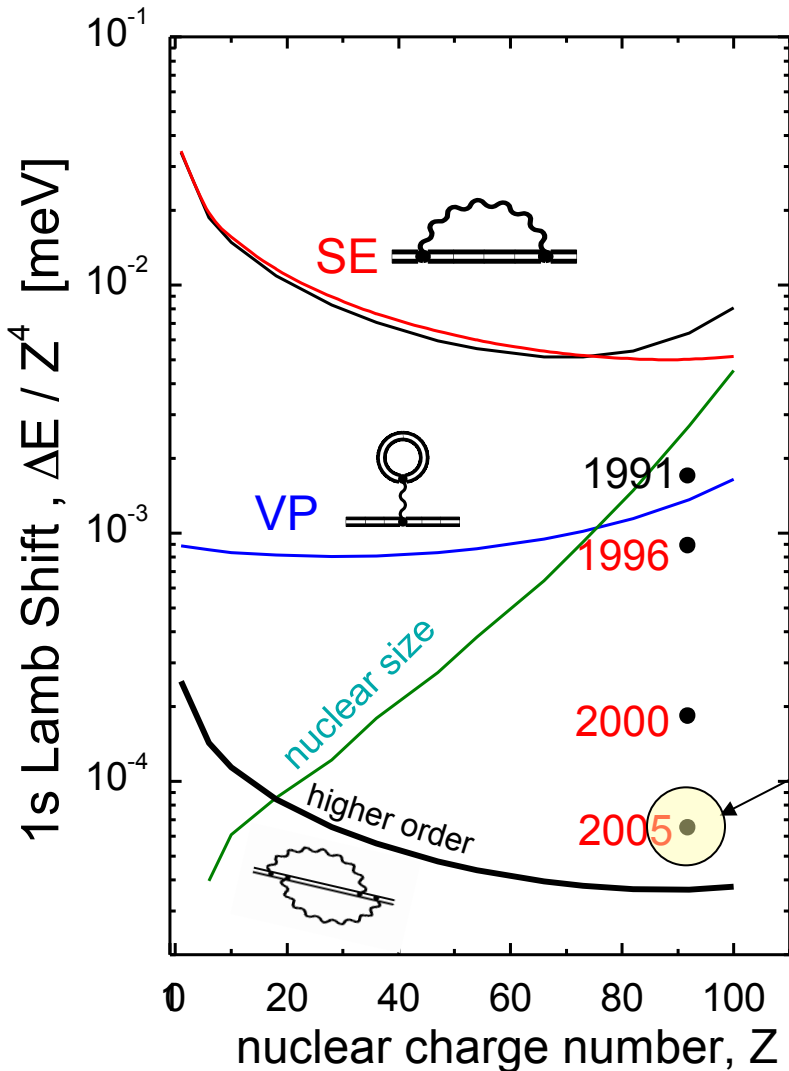
Summary

- Introduction
 - Test of QED in strong field
 - Study of strong interaction force at low energy
- Heavy highly charged ions
 - Production
 - **Lamb shift measurement**
 - Vacuum decay experiment
- Pionic atoms
 - Production
 - Pionic hydrogen and deuterium spectroscopy
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- Conclusions and outlooks

Lamb shift of H-like Uranium



Lamb shift of H-like Uranium



QED corrections

$$\Delta E \sim Z^4/n^3$$

Z : nuclear charge

n : principal quantum number

Important for s-states

Two-loop QED

~ 1 eV

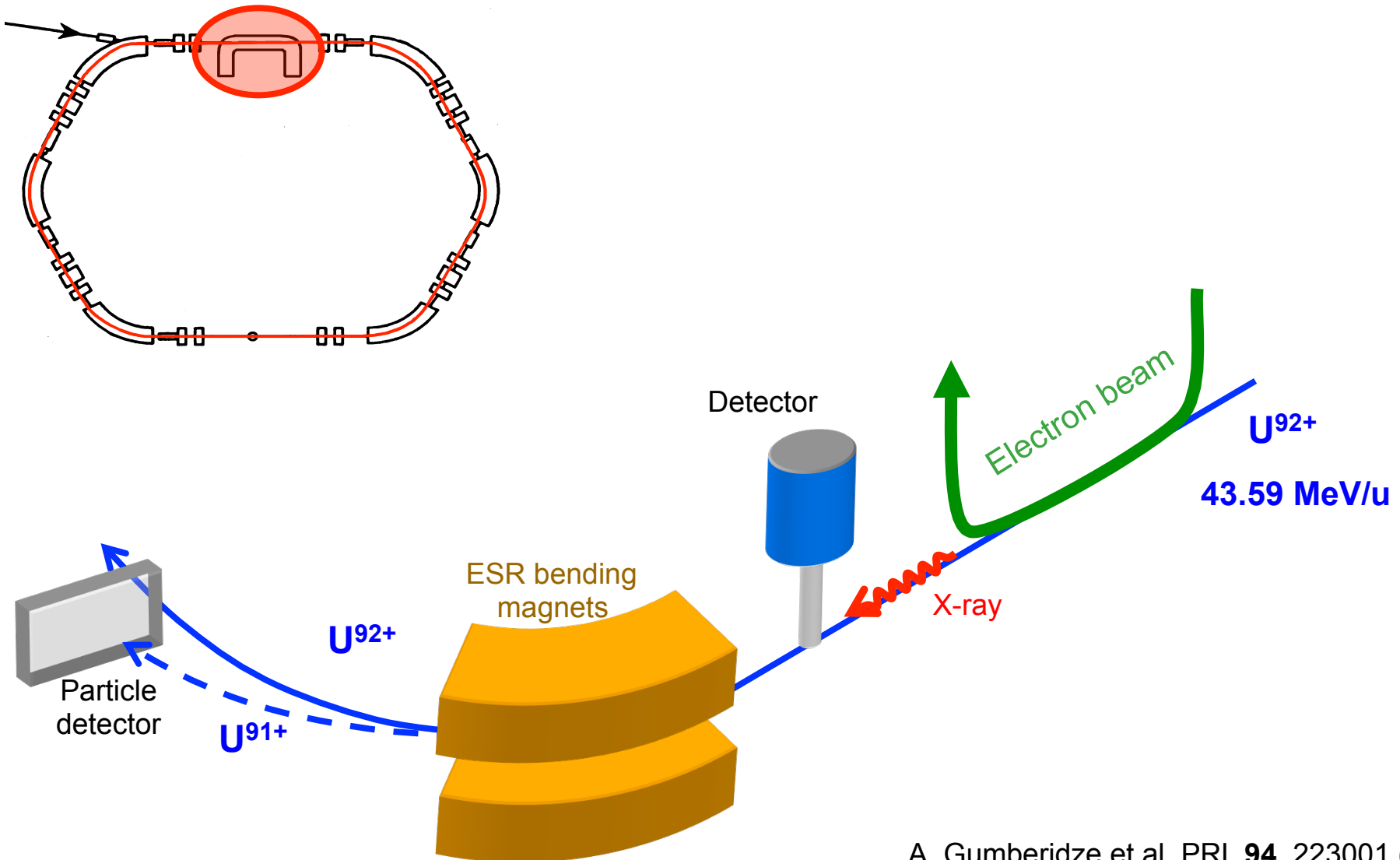
In Uranium

Self Energy
355.0 eV

Nuclear size
198.7 eV

Vacuum
Polarization
-88.6 eV

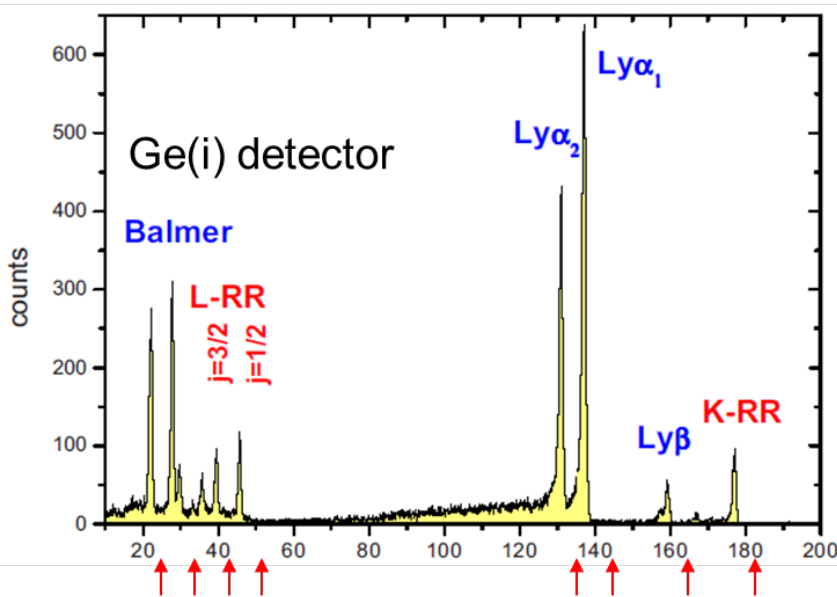
Lamb shift of H-like Uranium



A. Gumberidze et al., PRL **94**, 223001 (2005)

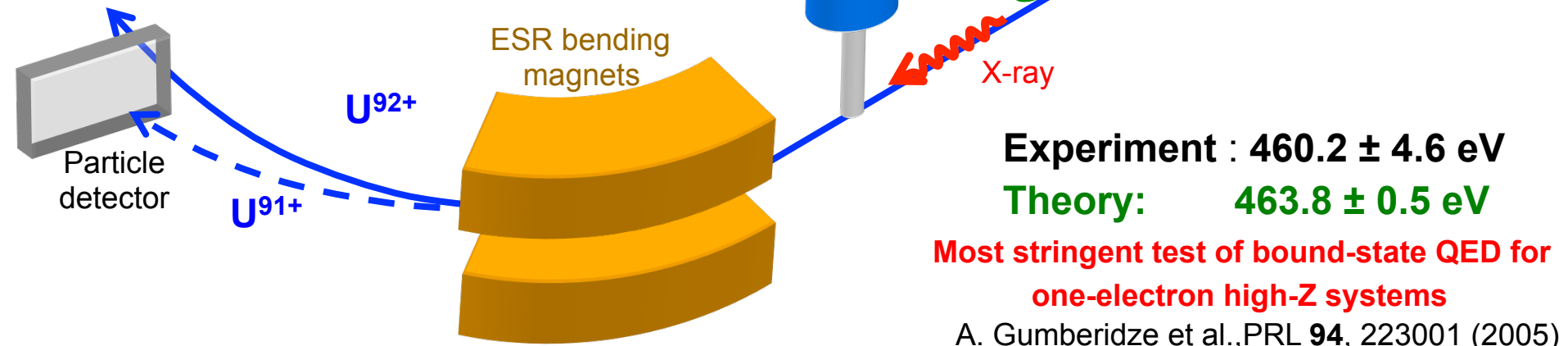
Lamb shift of H-like Uranium

Many lines providing intrinsic energy calibration and redundant data.

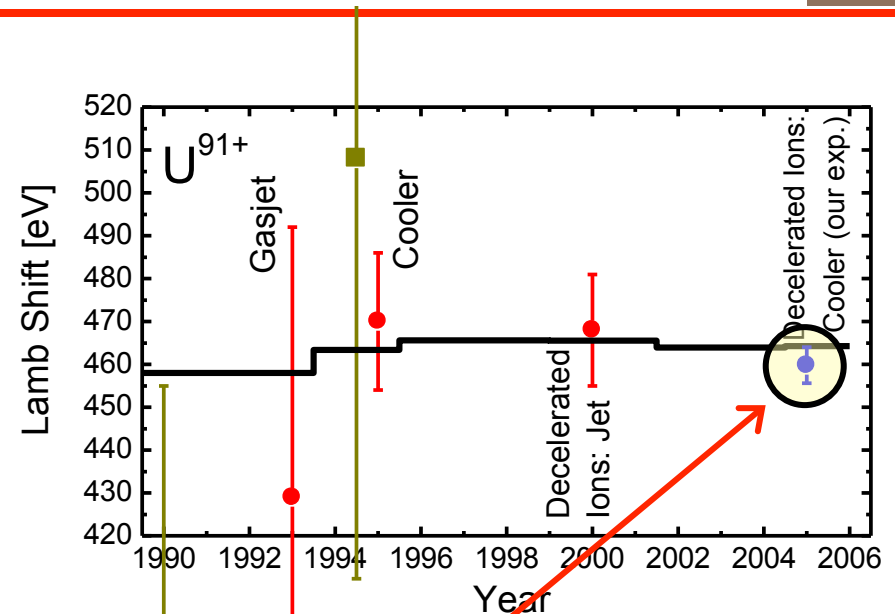
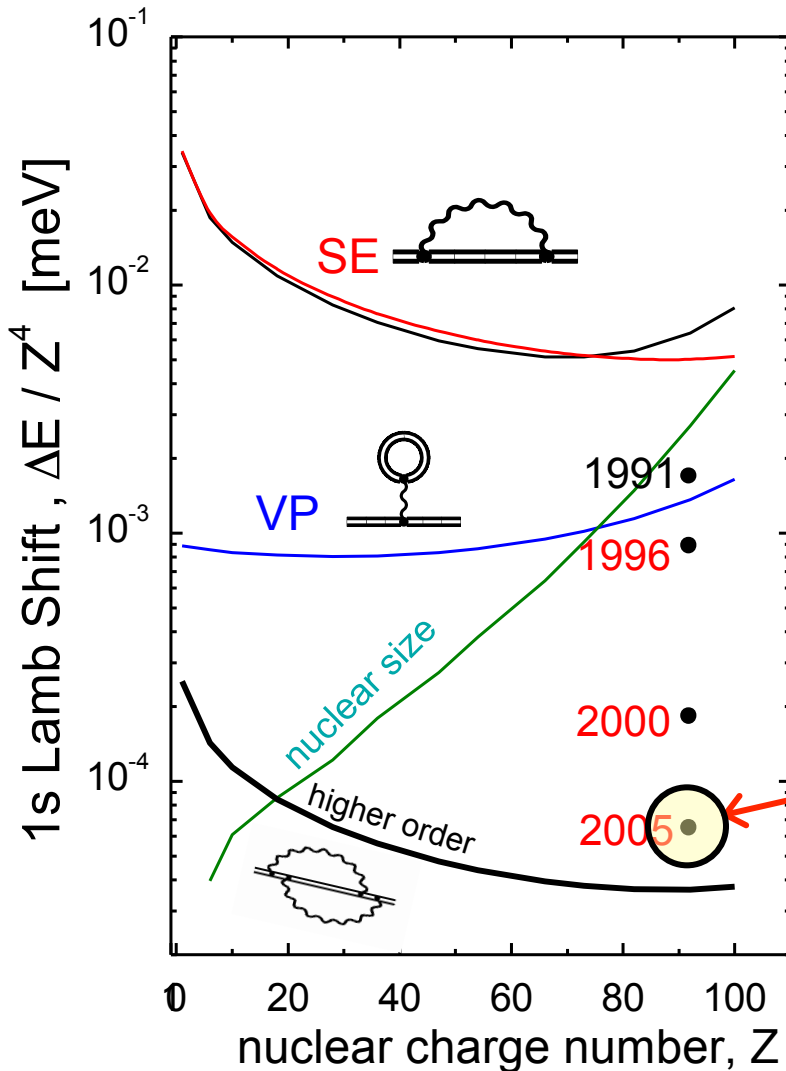


- decelerated ions
- close to 0-deg observation geometry
- standard germanium detector

$$E_{\text{ion}} = E_{\text{lab}} \gamma (1 - \beta \cos \theta)$$



Lamb shift of H-like Uranium



Latest experiment

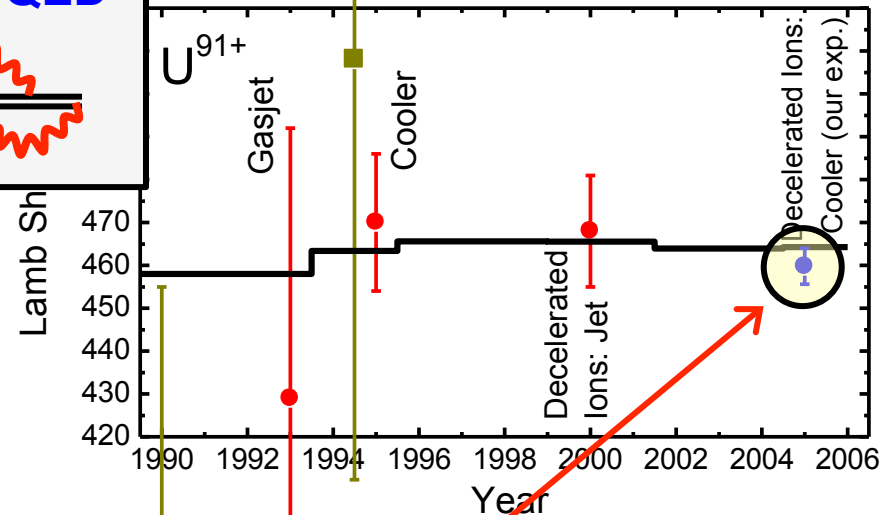
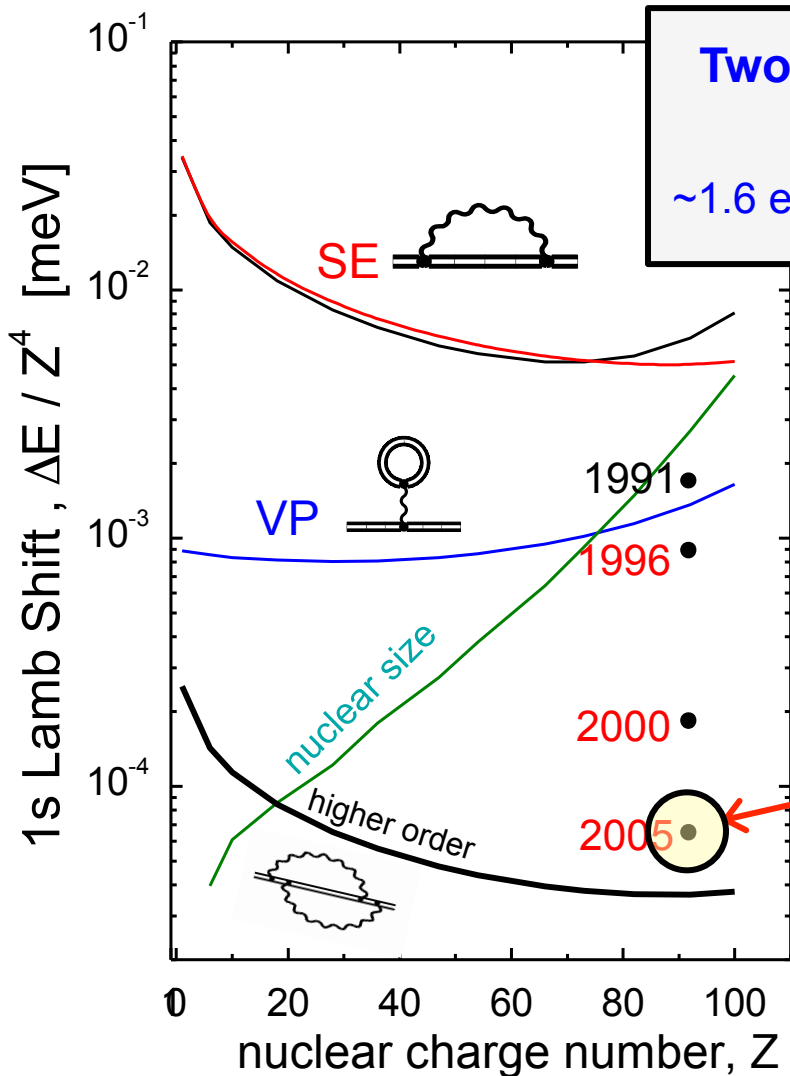
Experiment : 460.2 ± 4.6 eV

Theory: 463.8 ± 0.5 eV

Most stringent test of bound-state QED for one-electron high- Z systems

A. Gumberidze et al., PRL **94**, 223001 (2005)

Lamb shift of H-like Uranium



Latest experiment

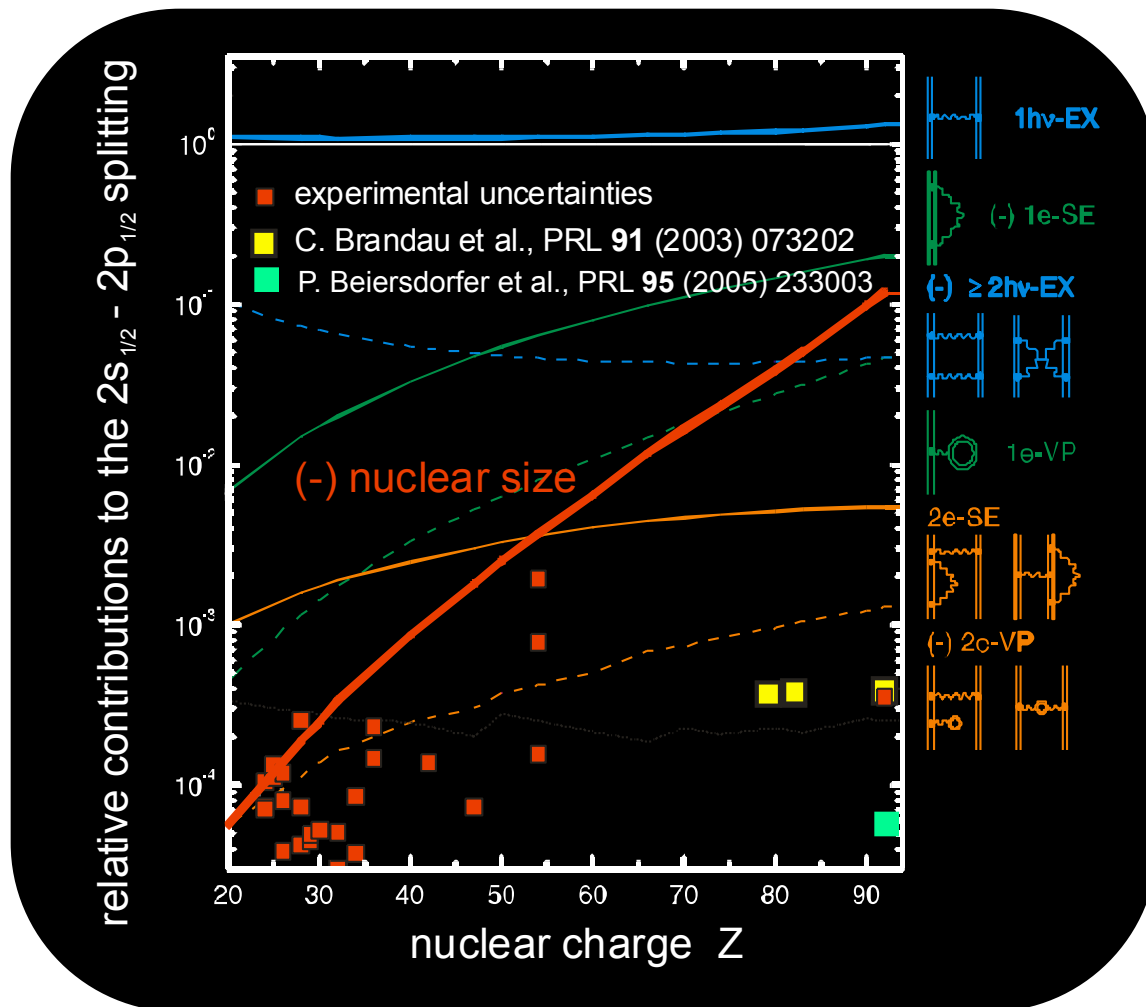
Experiment : 460.2 ± 4.6 eV

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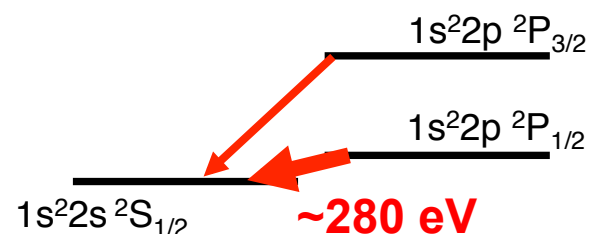
Most stringent test of bound-state QED for one-electron high-Z systems

A. Gumberidze et al., PRL **94**, 223001 (2005)

Most precise spectroscopy tests in Li-like systems



Li-like Uranium (3 el.)



XUV spectroscopy

Detailed knowledge about the nucleus is vital for strong field QED studies (nuclear size but also nuclear structure $\Rightarrow$ nuclear polarization)

Summary

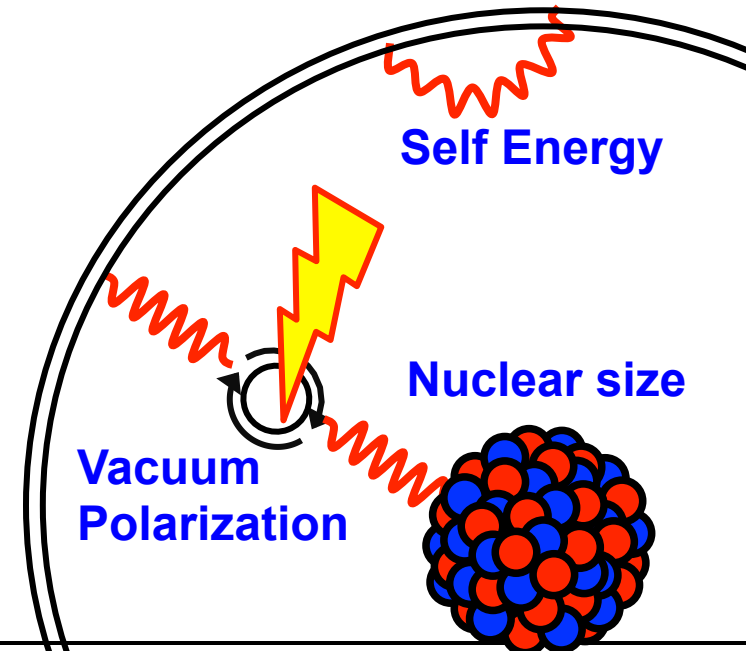
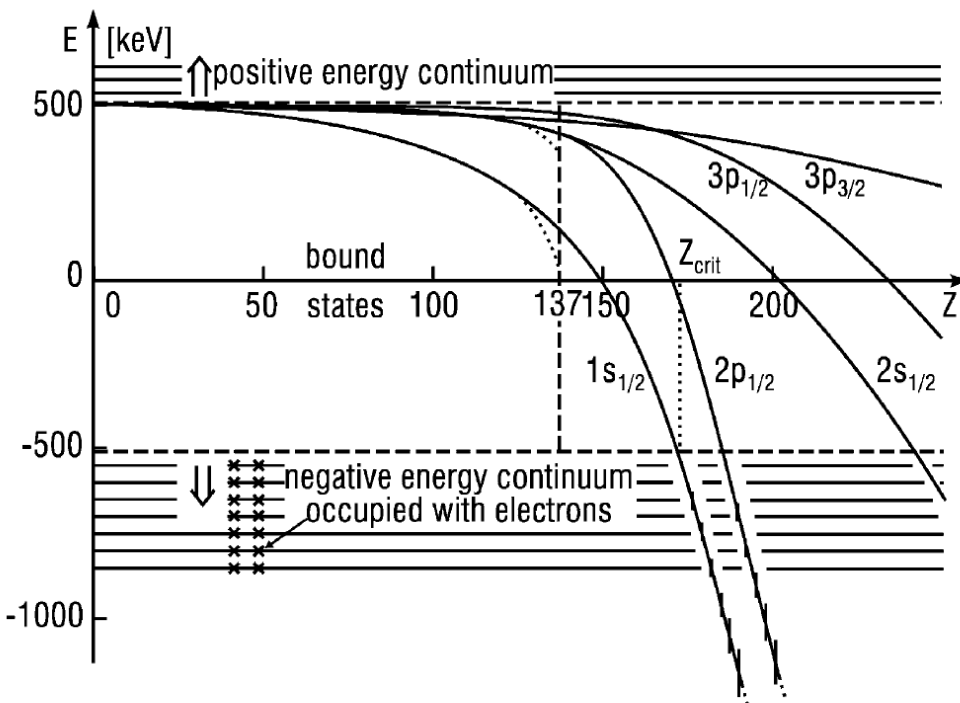
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Supercritical fields

Dirac equation: $E_{1s} = mc^2 \sqrt{1 - (Z\alpha)^2}$

$Z > 1/\alpha = 137 \rightarrow$ electron/positron pair creation

Finite nucleus size $\rightarrow Z_{\text{crit}} = 173$

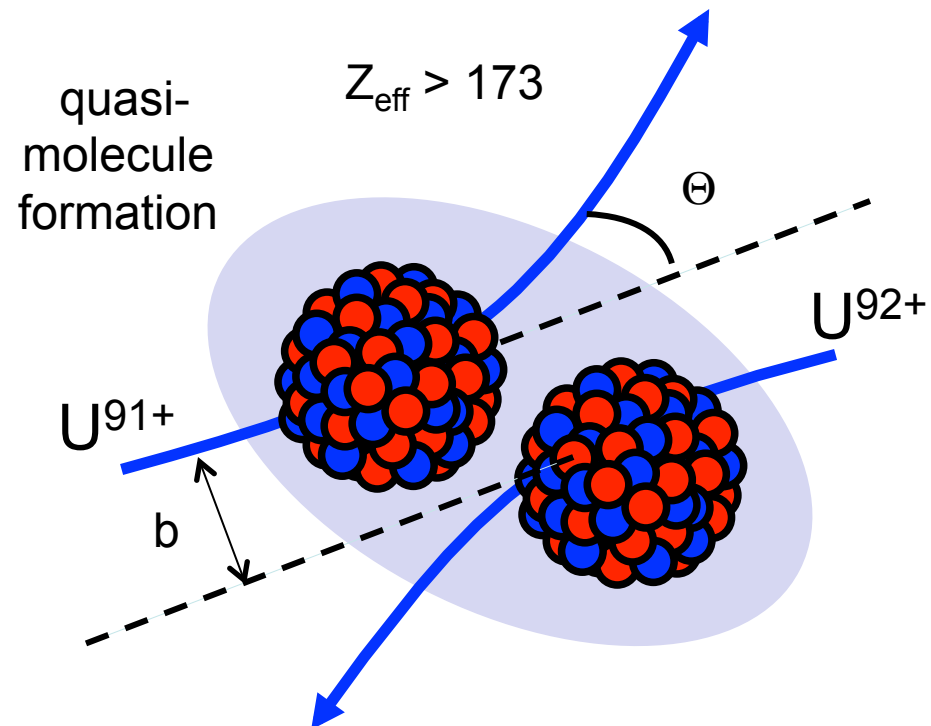
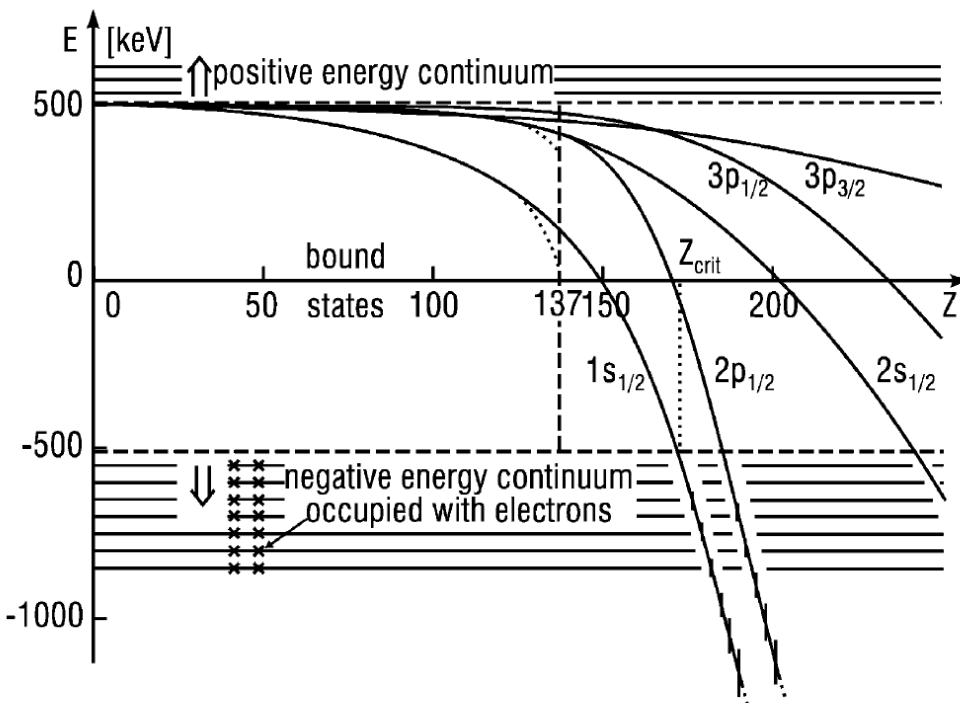


Supercritical fields

Dirac equation: $E_{1s} = mc^2 \sqrt{1 - (Z\alpha)^2}$

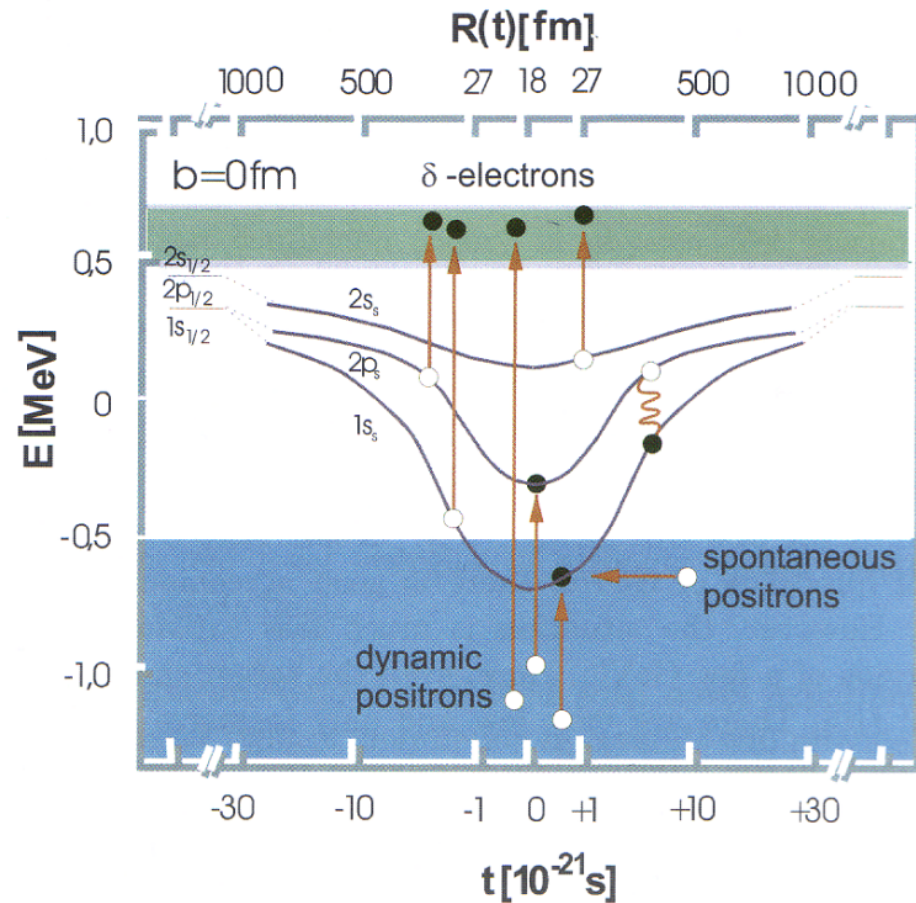
$Z > 1/\alpha = 137 \rightarrow$ electron/positron pair creation

Finite nucleus size $\rightarrow Z_{\text{crit}} = 173$



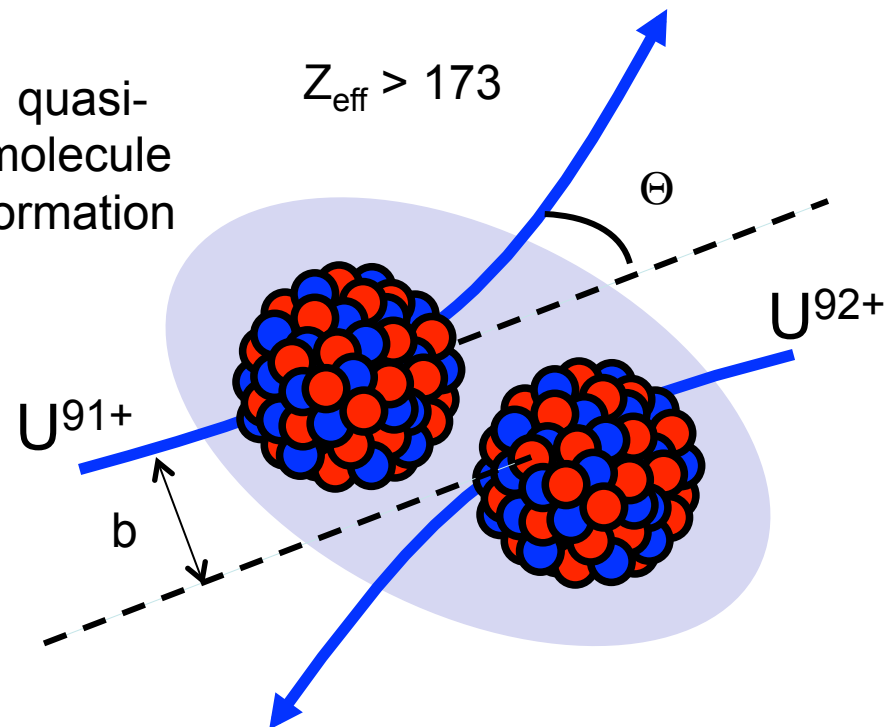
Supercritical Fields: formation of super heavy quasi-molecules

Slow ion – ion collision
~ 5 MeV/u

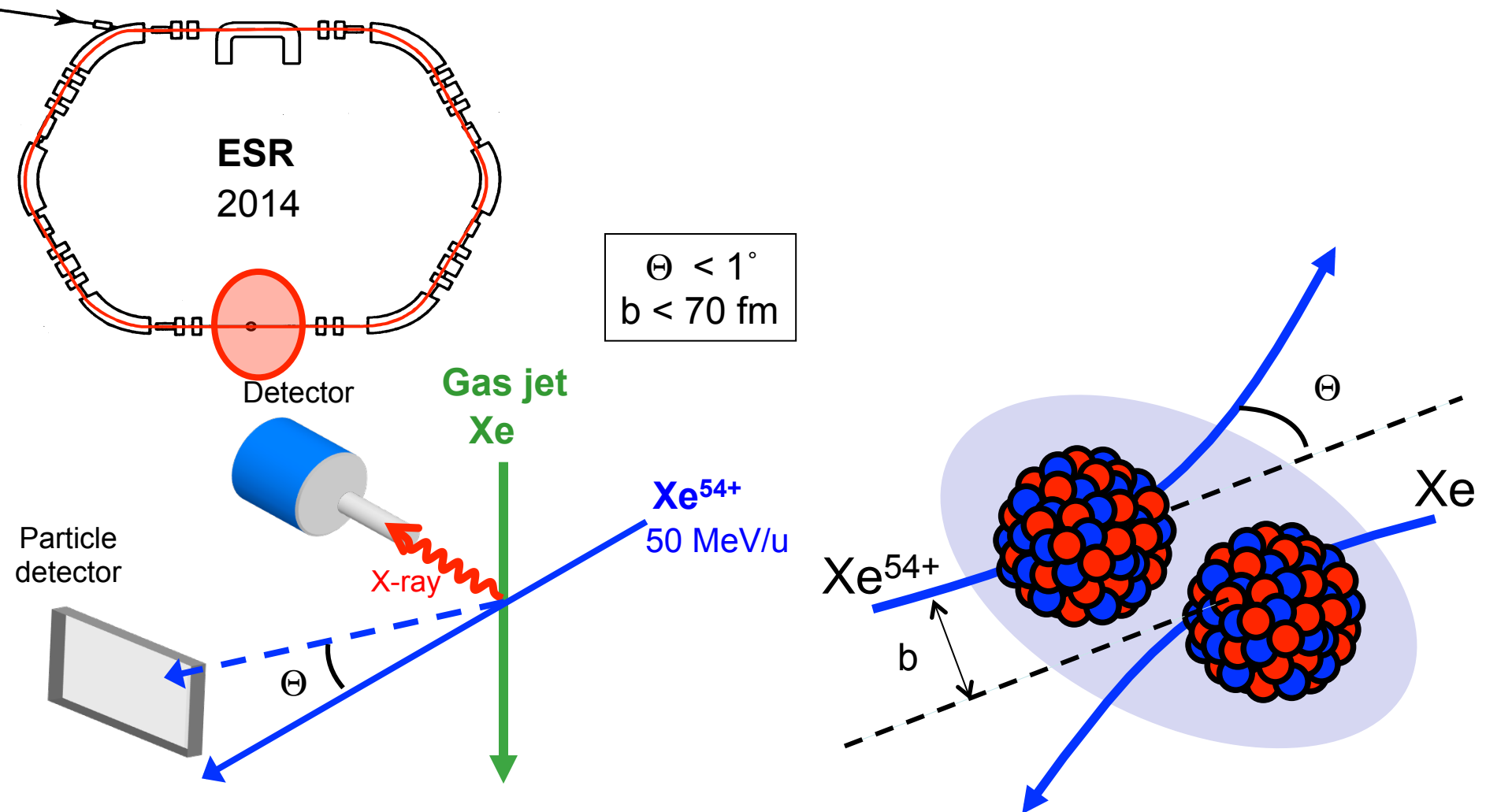


quasi-molecule formation

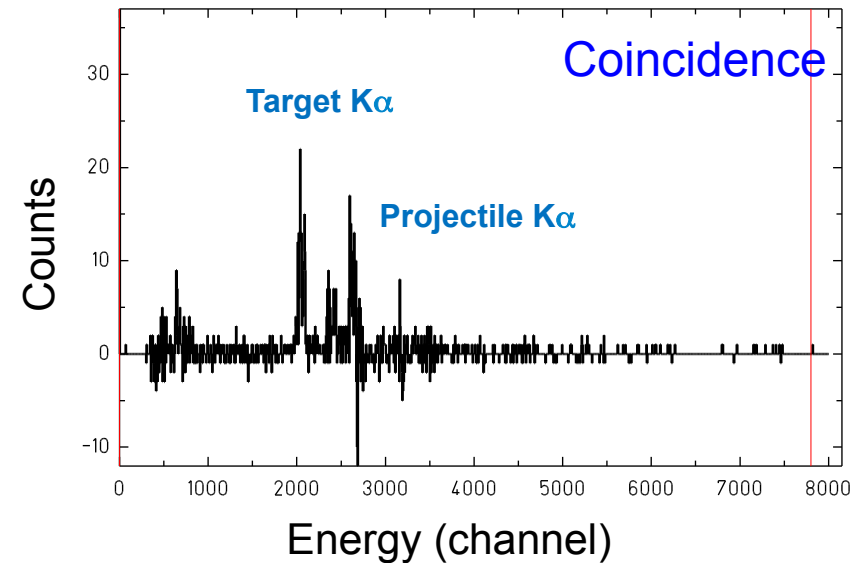
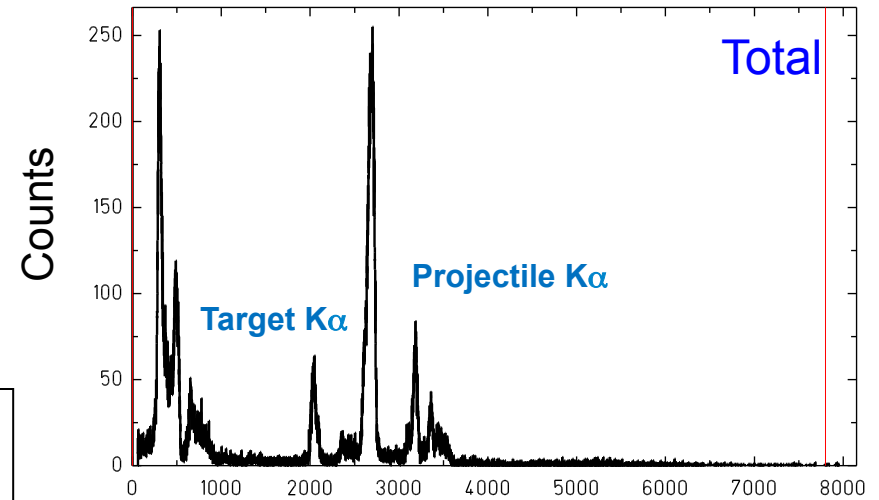
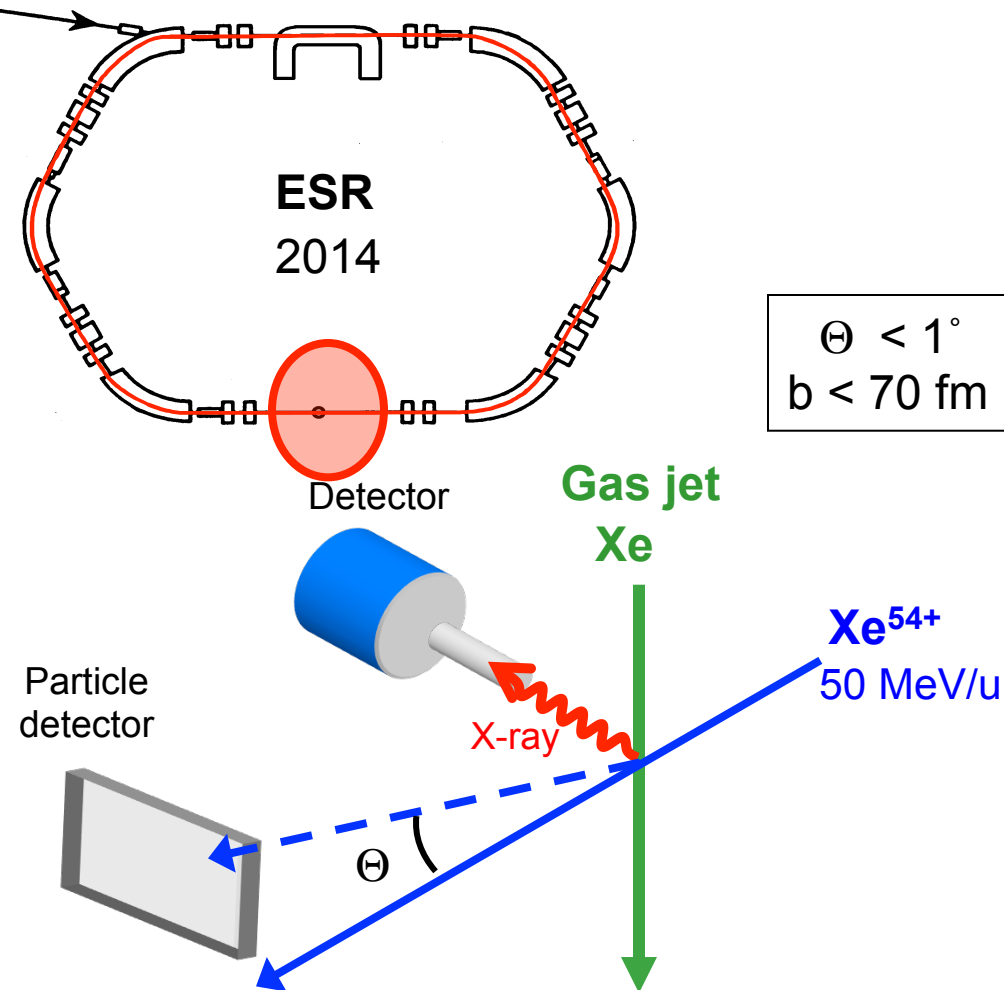
$Z_{eff} > 173$



First steps to supercritical fields



First steps to supercritical fields



Summary

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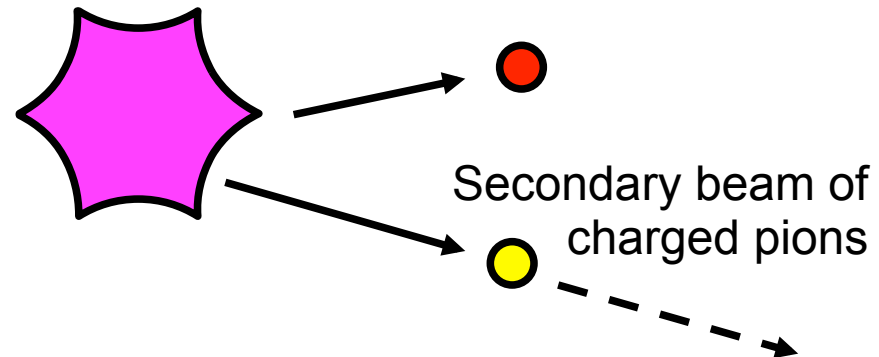
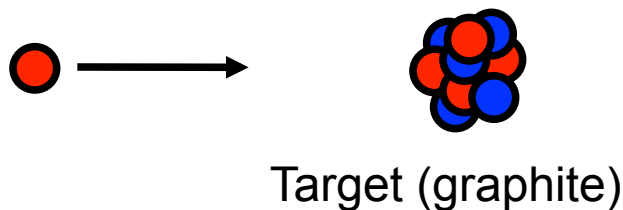
Pion beam production

Production at the Paul Scherrer Institut (Villigen, Switzerland)

- Proton beam : $E_{\text{kin}} = 590 \text{ MeV/c}$, $I = 1.9 \text{ mA}$
- Graphite target
- 10^8 pions/sec, $E_{\text{kin}} = 110 \text{ MeV/c}$



Accelerate proton



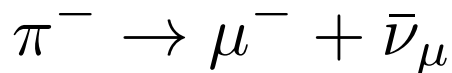
Pionic and muonic atoms production

- Cyclotron trap to stop the pions:
 - strong magnetic field ($B_{\max}=3.5$ Tesla)
 - plastic degraders (energy loss)
- Gaseous target:
 - T: from 14 K to room temperature
 - effective pressure: from ~ 0 to 40 bars

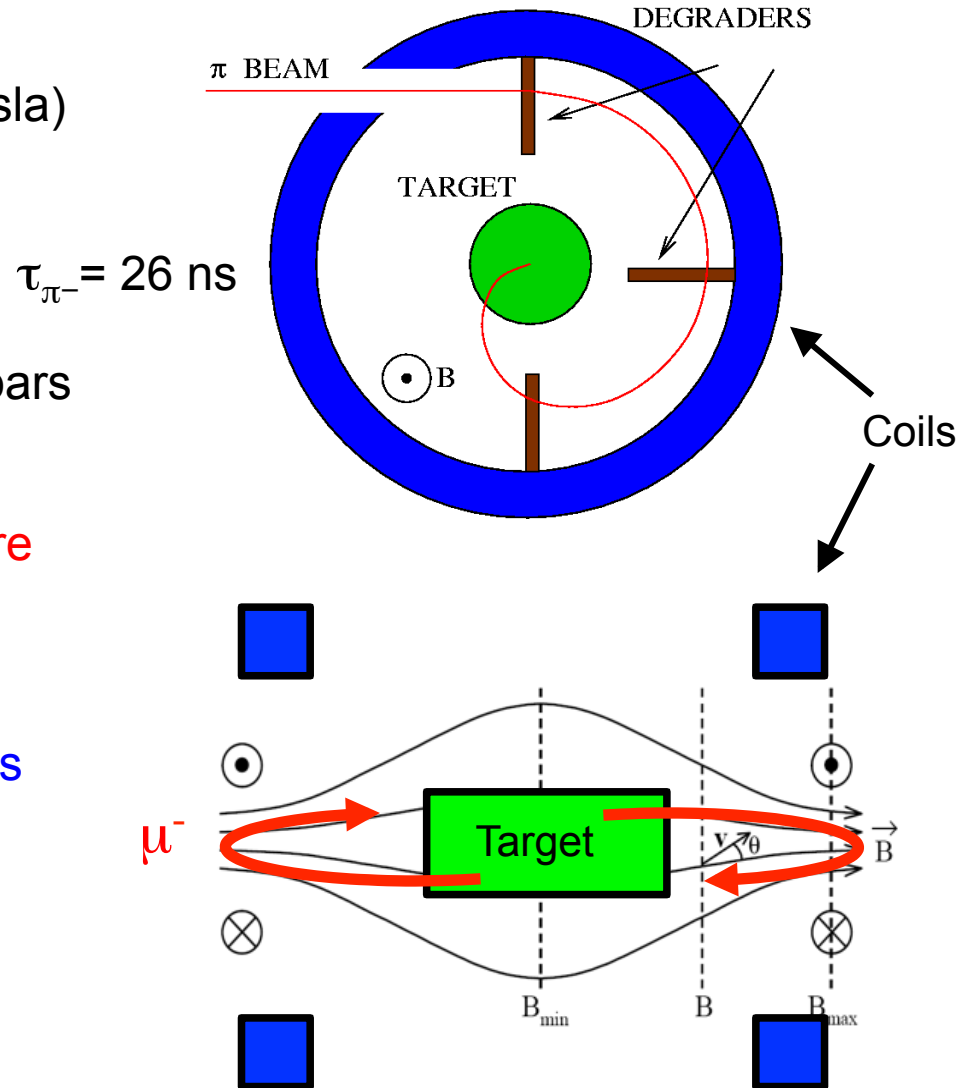


0.5% (per bar) of the incoming pions are stopped inside the target

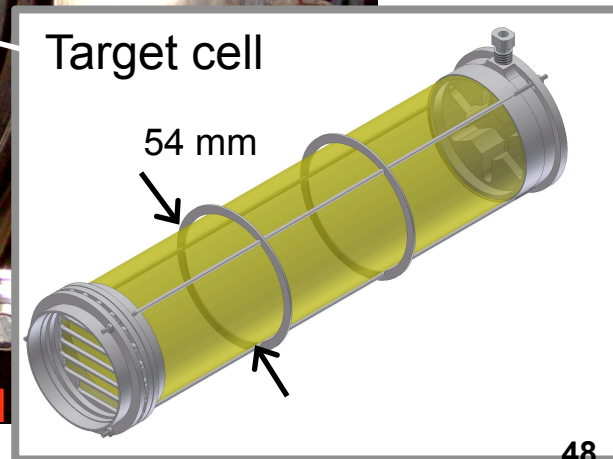
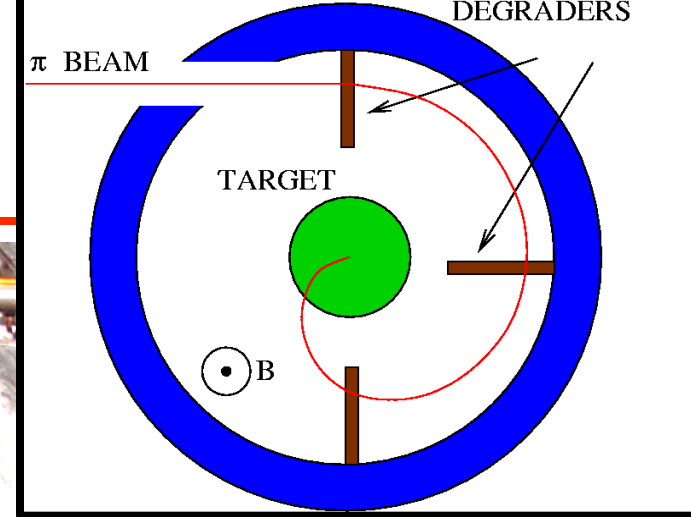
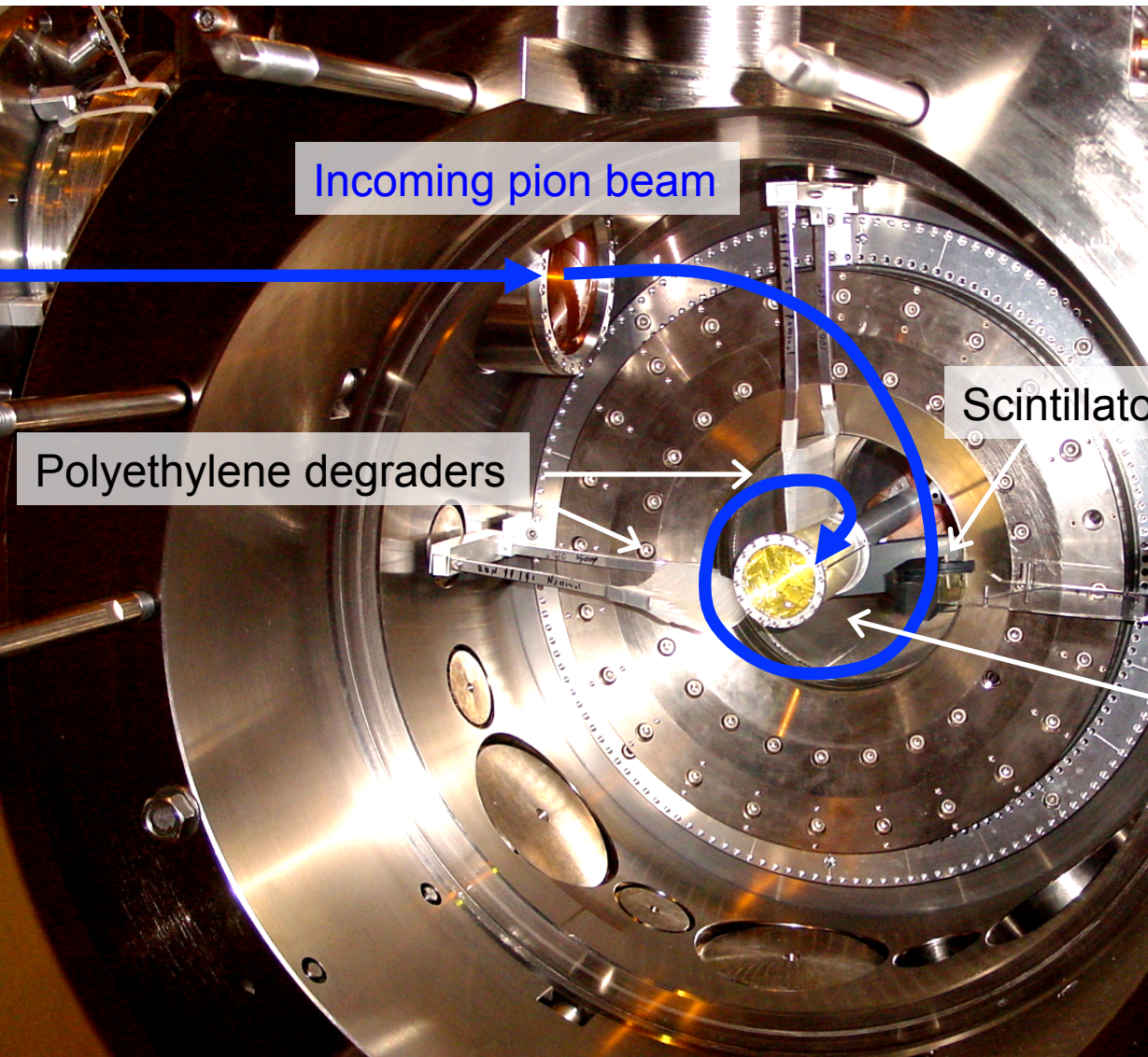
- Production and trapping of the muons



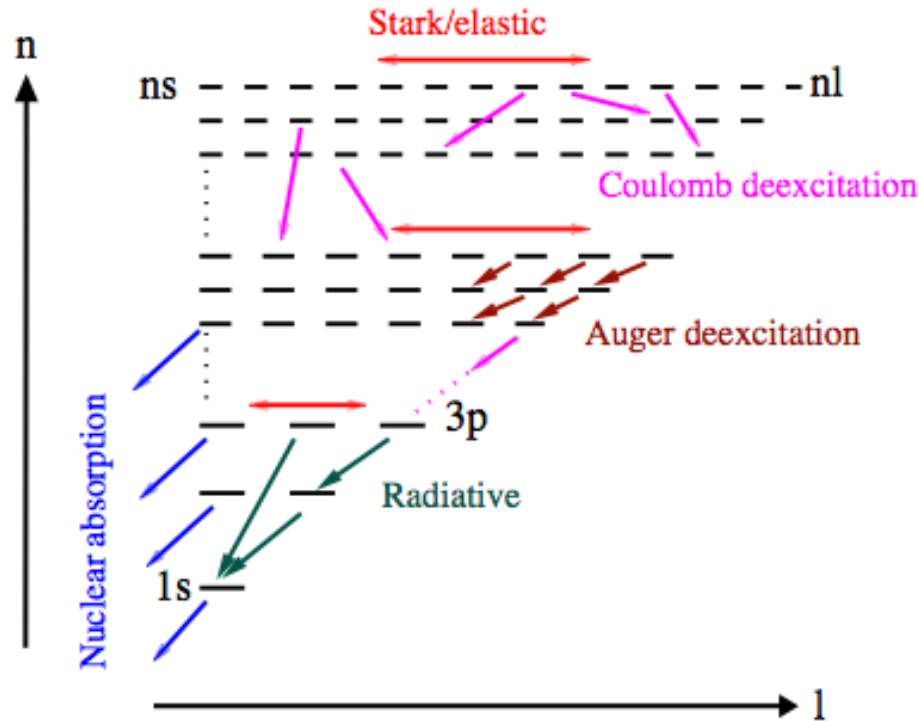
Formation of muonic and pionic atoms



Cyclotron trap



Exotic atoms formation



$$\Gamma_{\text{Auger}} \propto \frac{1}{\sqrt{2\Delta E + 38\text{eV}}}$$

Dominates for large n

- Deceleration and stop particle in target
 $E_{\text{kin}} \sim 10 \text{ eV}$

- Capture at radii of outmost electrons

$$n_{\pi} = n_{el} \sqrt{\frac{m_{\pi}}{m_{el}}} \approx 16 \quad \text{in the case of } \pi\text{H}$$

Highly excited state!

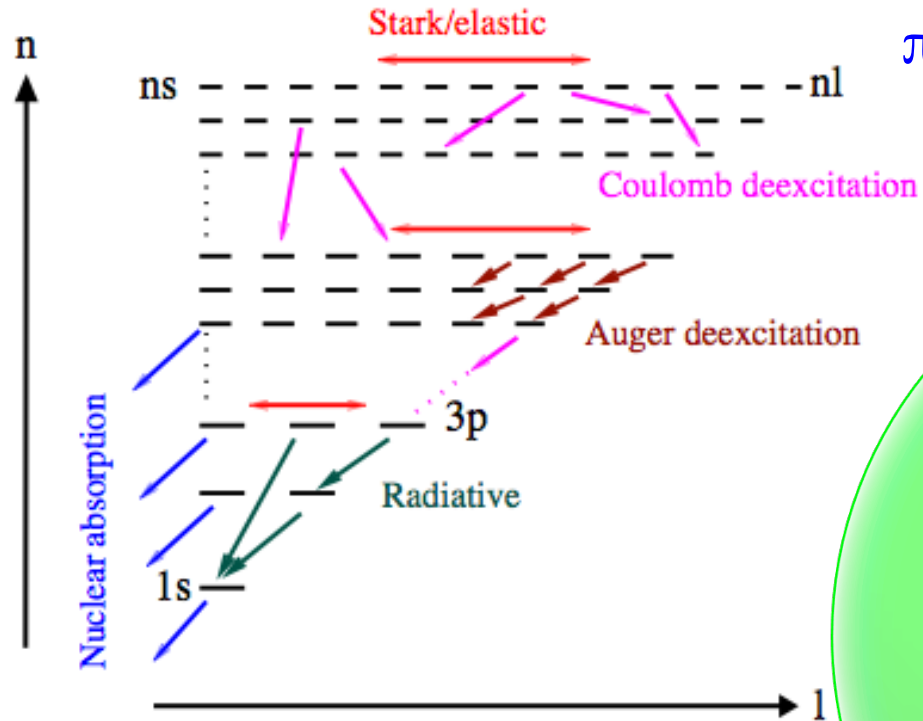
- De-excitation by competing Auger and X-ray emission

$$\Gamma_{\text{rad}} \propto \Delta E^3$$

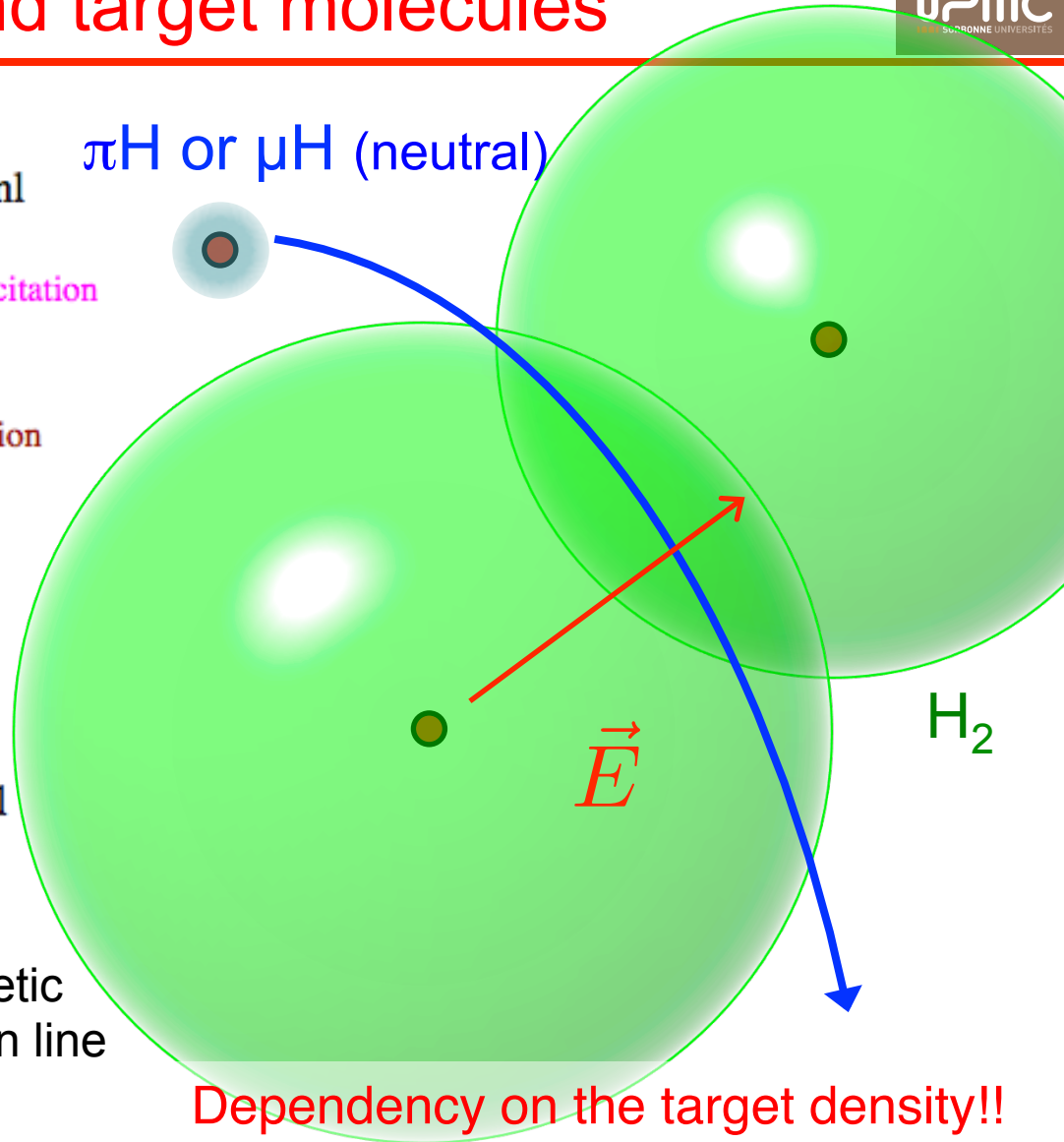
Dominates for small n

Low Z + dilute targets => no electrons remaining!!

Transitions induced by collisions between exotic hydrogen and target molecules



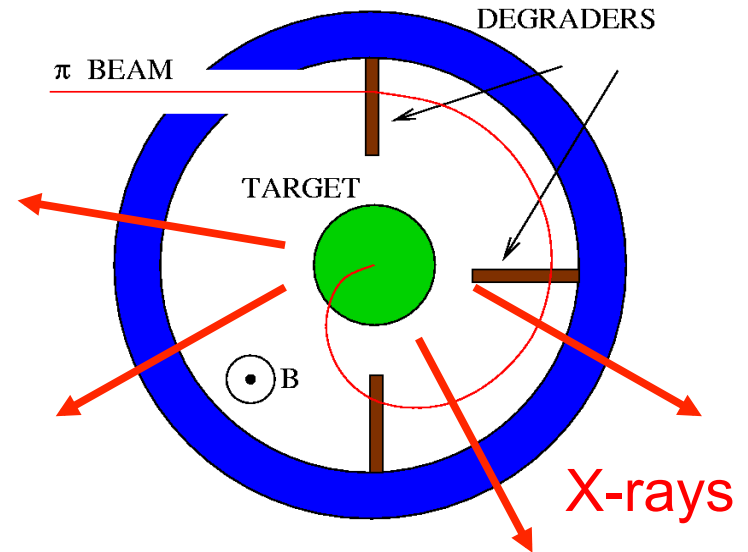
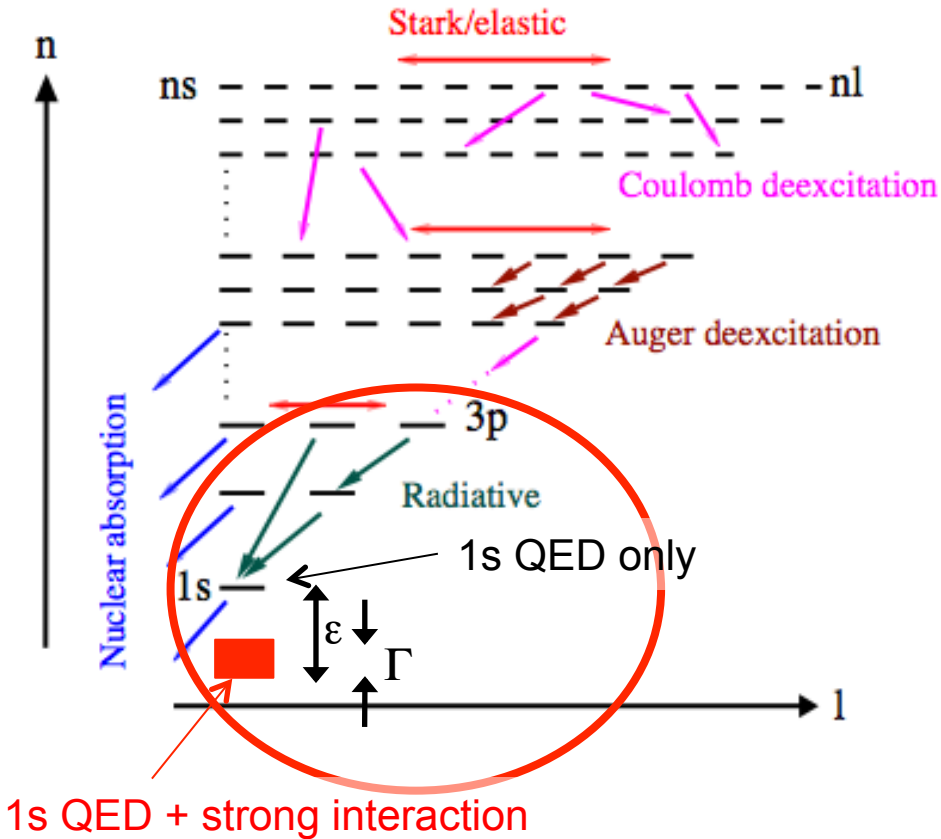
πH or μH (neutral)



Dependency on the target density!!

- Stark mixing
- Coulomb de-excitation $\rightarrow$ gain of kinetic energy $\rightarrow$ Doppler shift of the transition line
- Molecular formation

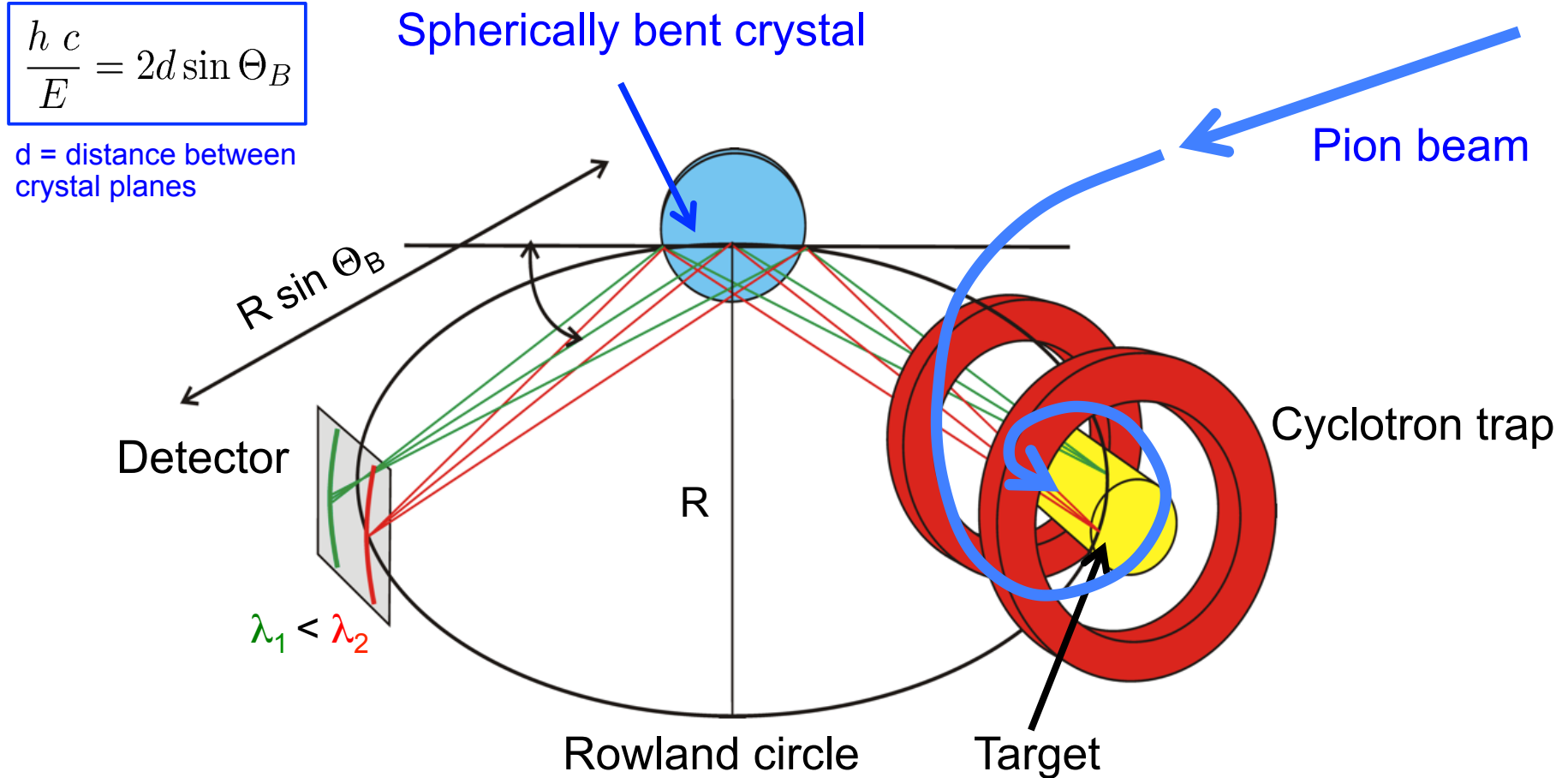
X-ray emission and detection



High power resolution for 2-4 keV
X-ray required



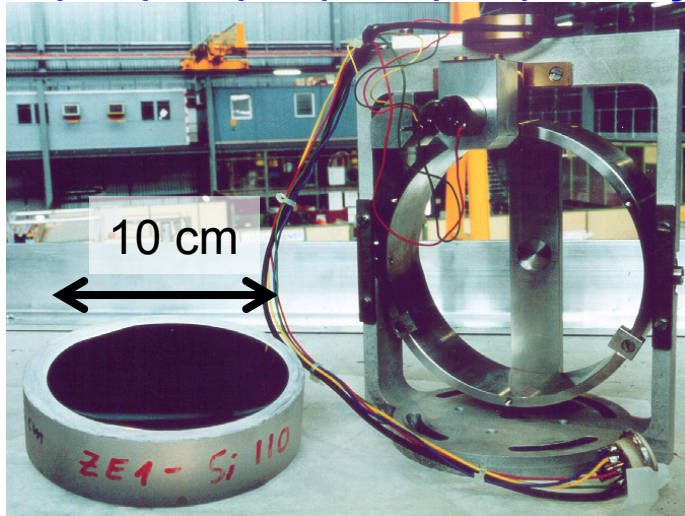
Bragg crystal spectroscopy



- [1] H.H. Johann, Zeitschrift für Physik **69**, 185 (1931)
- [2] J. Eggs *et al.*, Zeitschrift für angewandte Physik **20**, 118 (1965)
- [3] D. Gotta, Progress in Particle and Nuclear Physics **52**, 133 (2004)
- [4] D. Gotta *et al.*, Spectrochim. Acta, Part B **120**, 9 (2016)

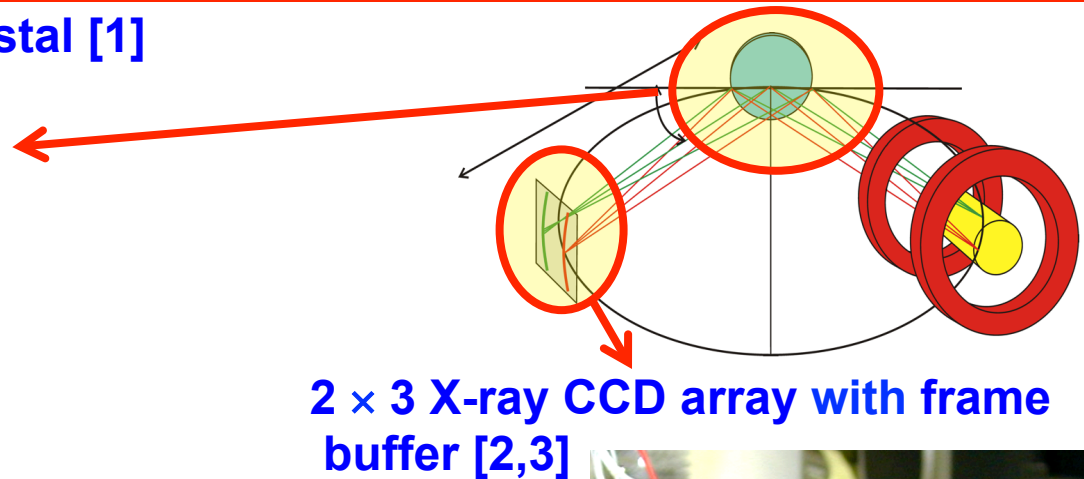
Diffraction crystal and position sensitive detector

Si(110), Si(111), Qz(100), ... crystal [1]



Radius of curvature: ~ 3 m
Diameter: 10 cm
Thickness: 290 μm
Support: polished quartz lens

Produced by Zeiss (Oberkochen, Germany)

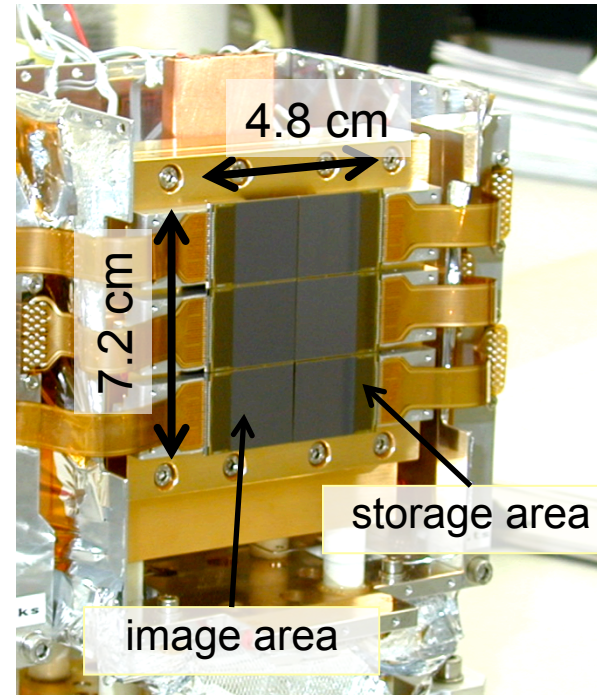


2 × 3 X-ray CCD array with frame buffer [2,3]

Total transmission:
 $\sim 5 \cdot 10^{-8}$

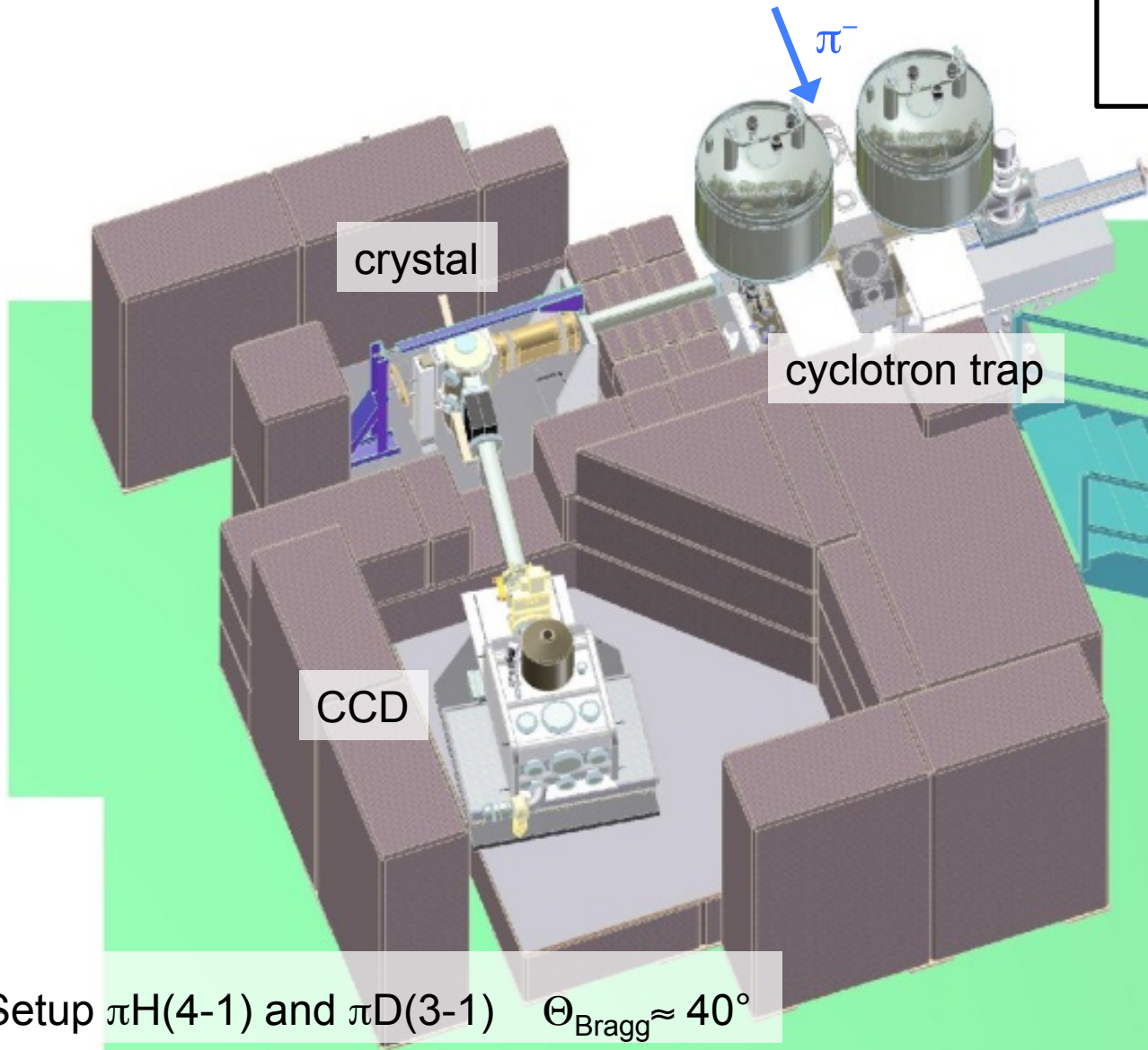
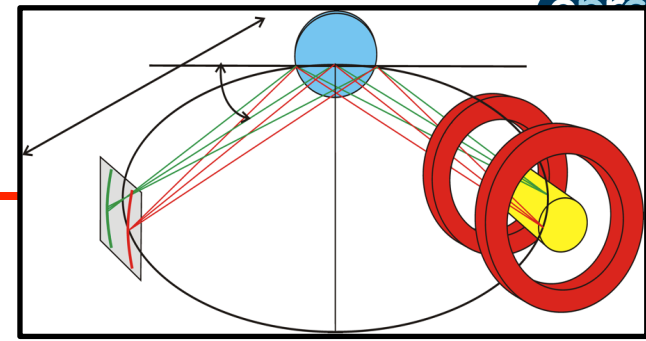
pixel size 40 $\mu\text{m} \times 40 \mu\text{m}$
600 × 600 pixels per chip
frame transfer ≈ 10 ms
data processing ~ 24 s
operates at -100°C

$\Delta E \approx 150$ eV @ 4 keV
Efficiency $\approx 90\%$



- [1] D.S. Covita *et al.*, Rev. Sci. Instrum. **79**, 033102-3 (2008)
- [2] N. Nelms *et al.*, Nucl. Instrum. Methods A **484**, 419 (2002)
- [3] P. Indelicato *et al.*, Rev. Sci. Instrum. **77**, 043107 (2006)

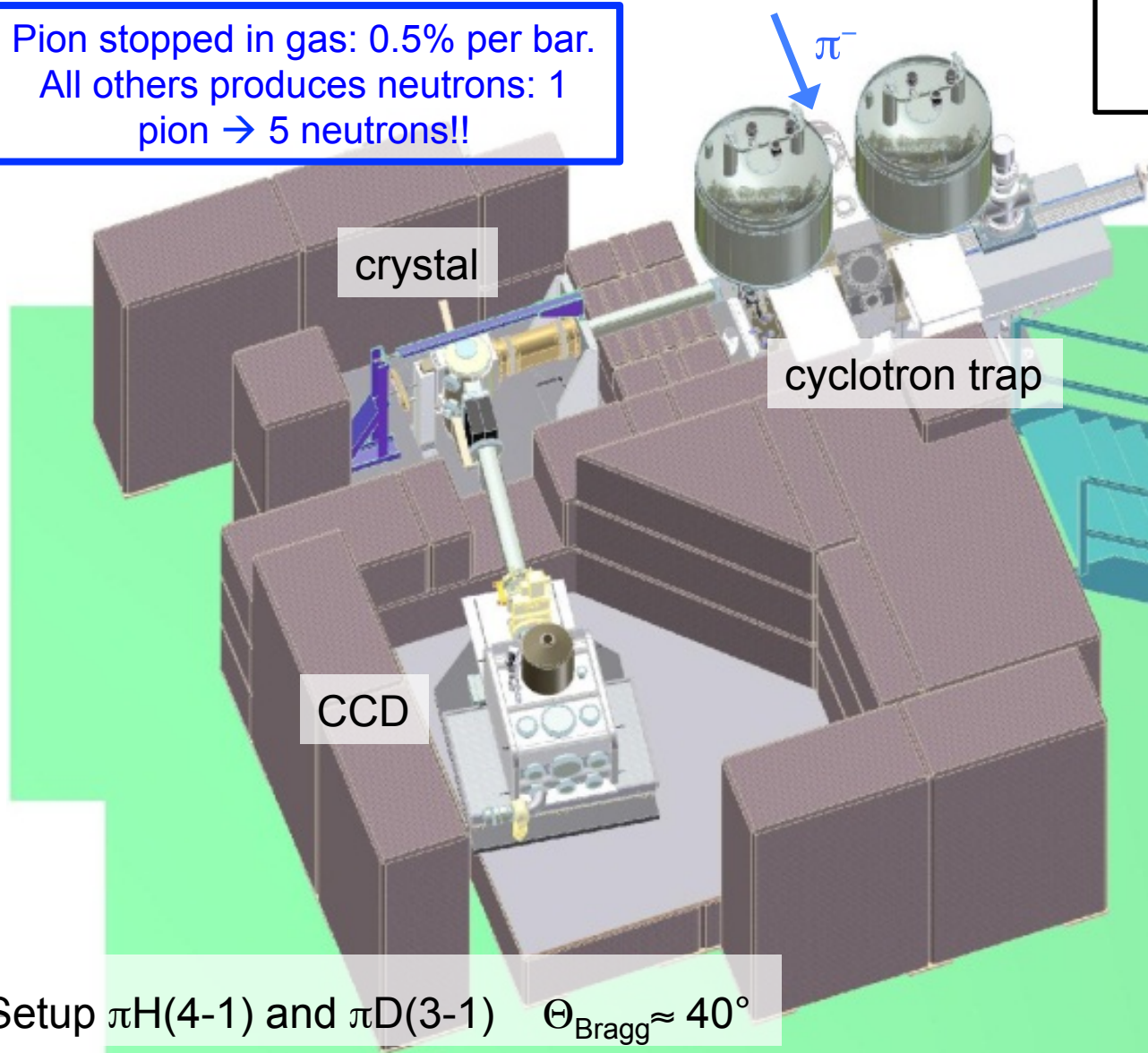
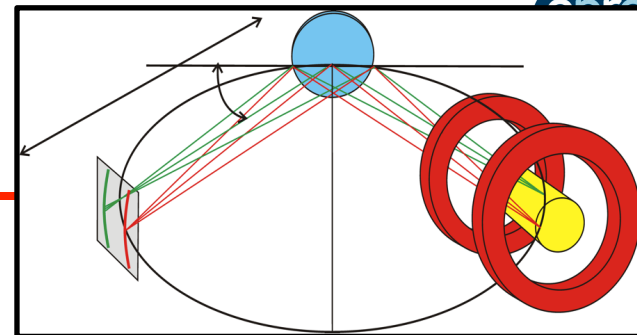
Typical set-up at PSI



Setup $\pi\text{H}(4-1)$ and $\pi\text{D}(3-1)$ $\Theta_{\text{Bragg}} \approx 40^\circ$

Typical set-up at PSI

Pion stopped in gas: 0.5% per bar.
All others produces neutrons: 1
pion $\rightarrow$ 5 neutrons!!



huge background
on the CCD



heavy concrete

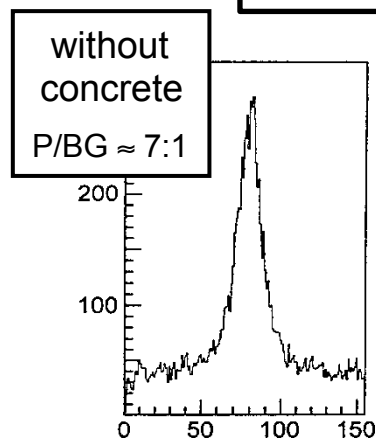
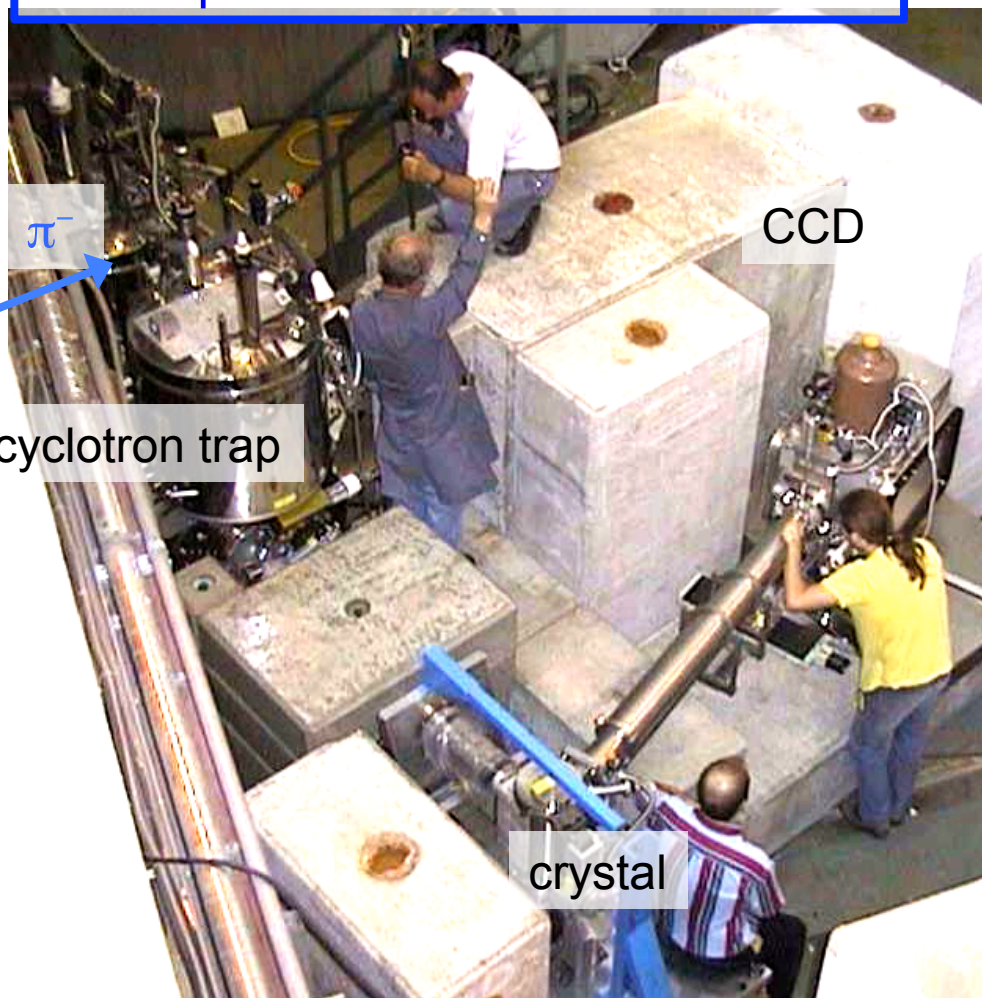
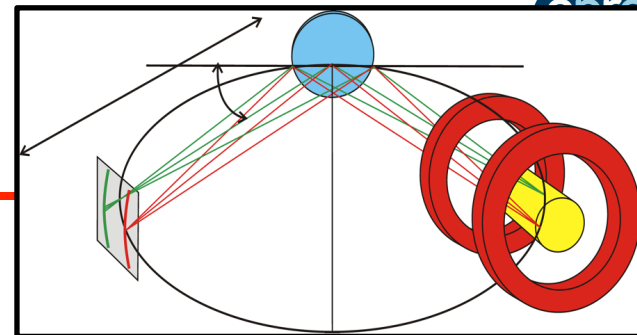


peak-to-background
reduced by a factor
of 10

Setup $\pi\text{H}(4-1)$ and $\pi\text{D}(3-1)$ $\Theta_{\text{Bragg}} \approx 40^\circ$

Typical set-up at PSI

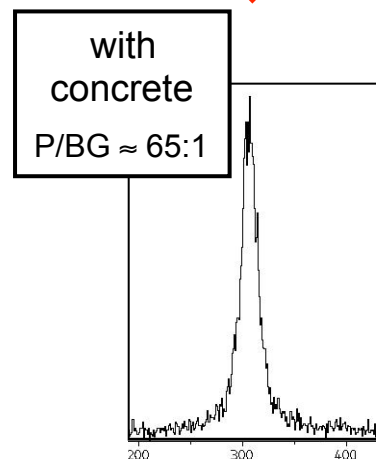
Pion stops in gas: few % all others:
1 pion makes 5 neutrons!!



huge background
on the CCD



heavy concrete

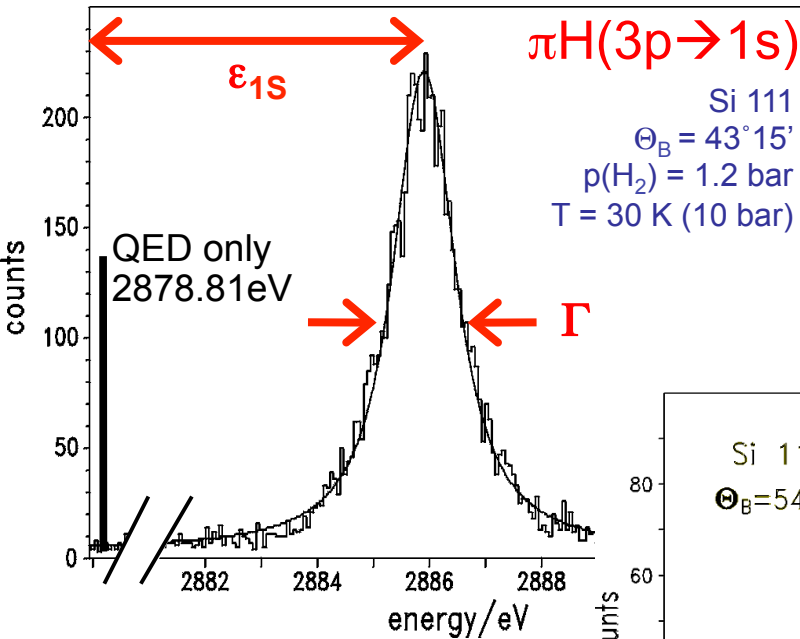


peak-to-background
reduced by a factor
of 10

Summary

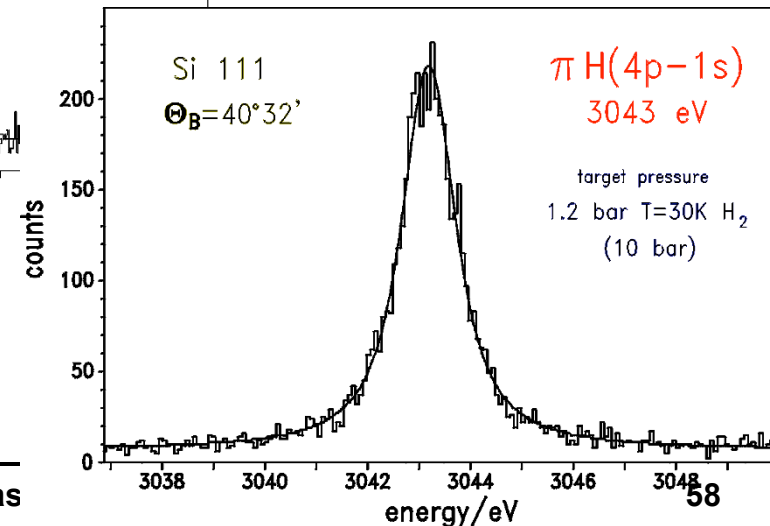
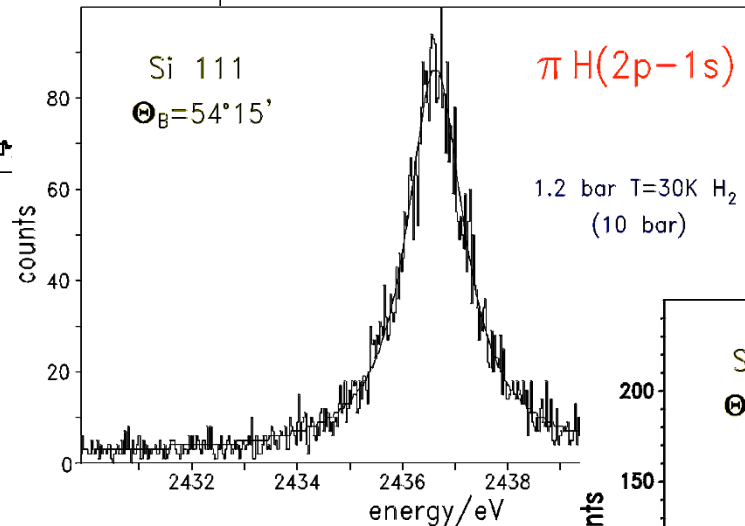
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Pionic hydrogen results



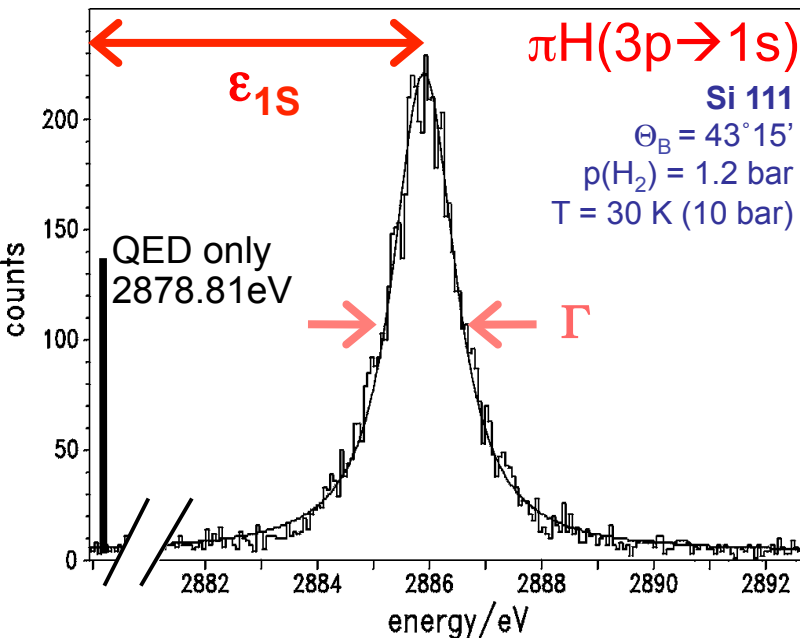
Independent measurement:

- Different transitions
- Different energy calibration
- Different crystals
- **Different target densities**



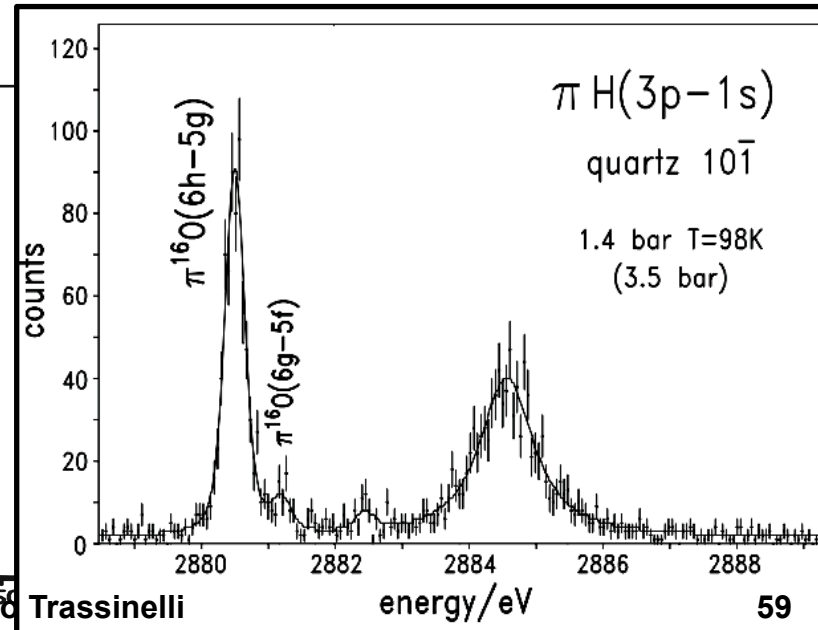
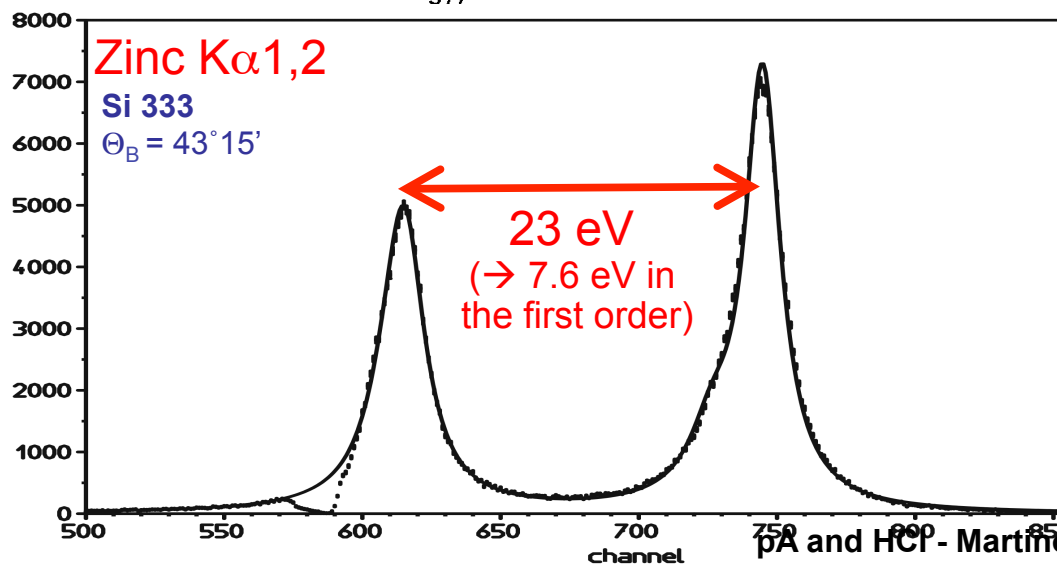
Several weeks of data acquisition for each spectrum

Pionic hydrogen shift measurement

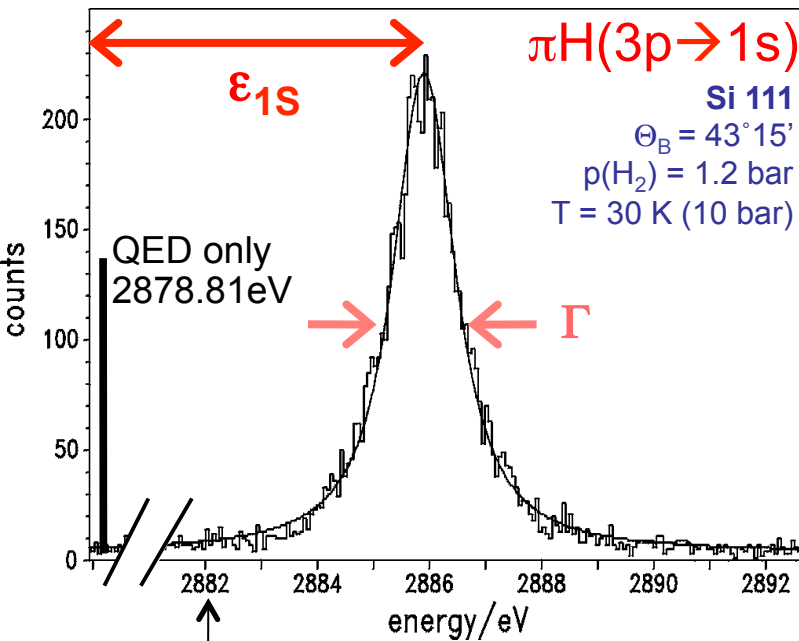


Calibrations

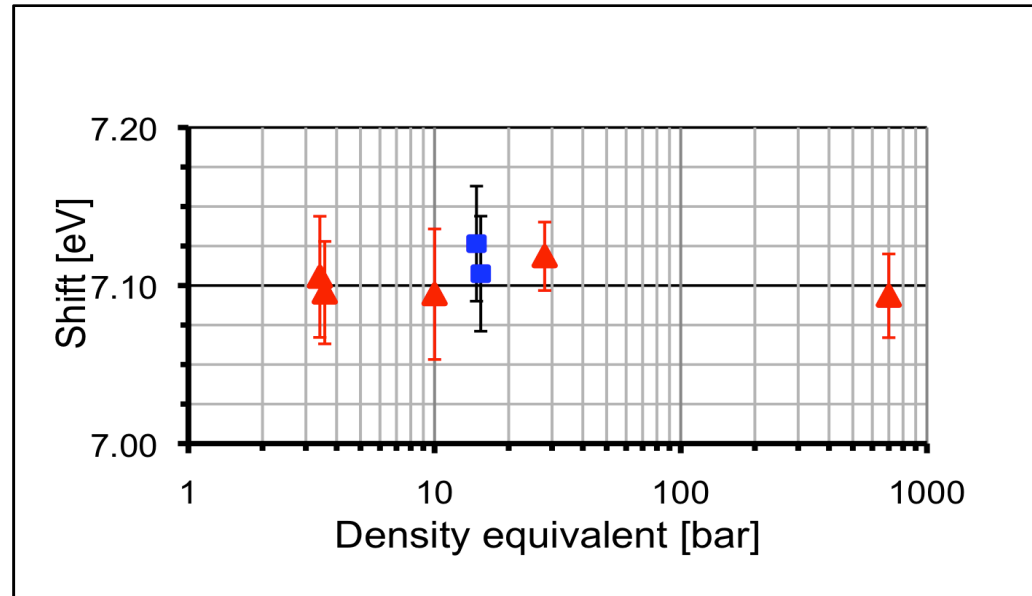
- $K\alpha$ fluorescence from selected target (different reflection order)
- from other exotic atom transitions (where strong interaction effects are negligible)
 - smaller systematic errors from crystal spectrometer aberrations and possible chemical shifts and pion mass accuracy



Pionic hydrogen shift measurement



Satellite position



ETHZ-PSI previous experiment – Ar $K\alpha$
H.-Ch.Schröder et al. Eur. Phys.J.C 1 **473** (2001)

$$\varepsilon_{1s} = + 7.0869 \pm 0.0071 \pm 0.0064 \text{ eV } (\pm 0.13\%) [1]$$

0.001 eV from QED predictions [4] (attractive force)
→ from the pion mass

Results

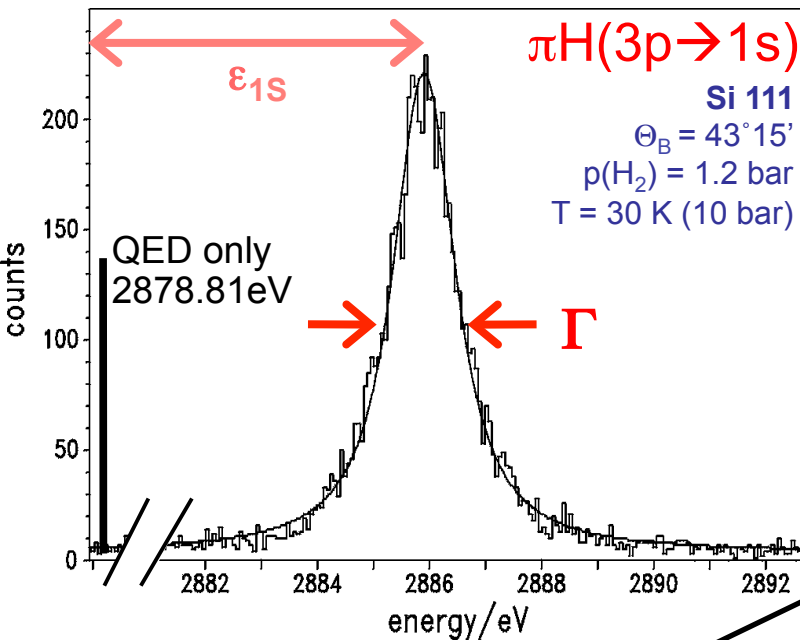
- No satellite from molecular X-ray detected (theoretical predictions [2] not confirmed)
- New value for the strong interaction shift in πH [3] (more 3 times smaller uncertainty)

[1] M. Hennebach et al., Eur. Phys. J. A **50**, 1-10 (2014)

[2] S.Kilic et al., Phys. Rev. A **70** 042506 (2004)

[3] D. Gotta, et al., Lect. Notes Phys. **745** 165 (2008)

[4] J. Schlessler et al., Phys. Rev. C **84**, 015211 (2011)



Much more complex!!

$$\Gamma_{exp} = \Gamma_{crystal} \otimes \Gamma_{Doppler} \otimes \Gamma_{1s}$$

from the response function
of the instrument

from the fast movement
of the ions

from the strong
interaction (goal!)

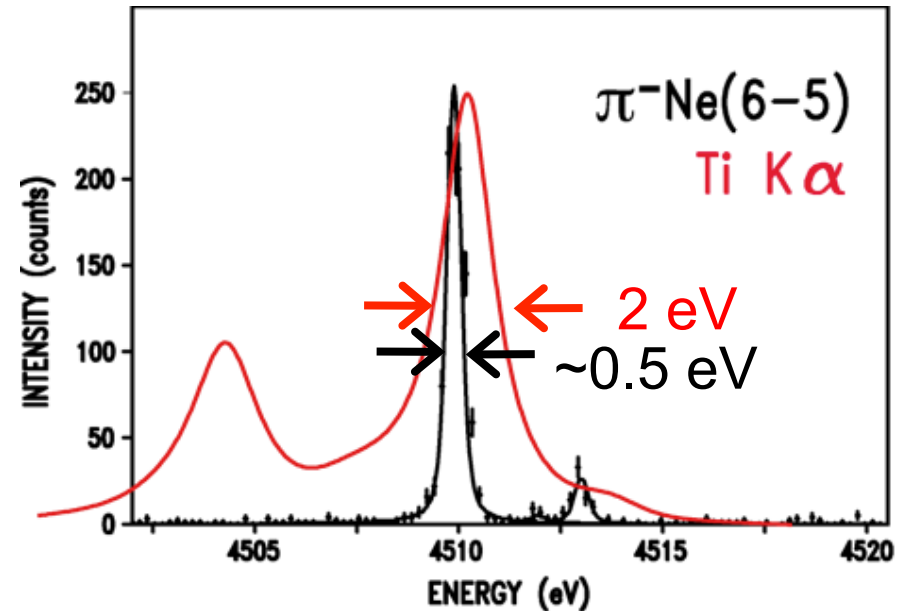
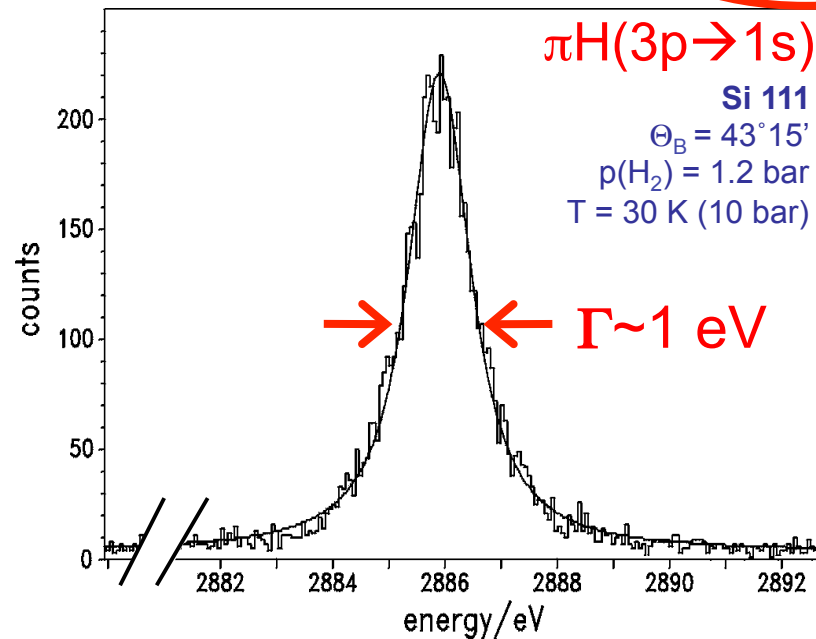
Characterization of the spectrometer [1,2]

- new measurements highly charged ion X-ray transitions
- simulation of the instrument

Accurate investigation to the atomic cascade

- new theoretical approaches
- new experiment with muonic hydrogen

$$\Gamma_{exp} = \Gamma_{crystal} \otimes \Gamma_{Doppler} \otimes \Gamma_{1s}$$



How to characterize the response function of the spectrometer?

Where to find a X-ray line with a $\Gamma < 0.5 \text{ eV}$?

- Fluorescence X-ray: **too broad!!**
- Pionic atoms (no Coulomb explosion): too low **count rate!!**
- ... **highly charged ion radiation!!**

Spectrometer characterization from He-like S, Cl, Ar M1 transitions

X-ray emission from a confined hot plasma (ECR ion source)

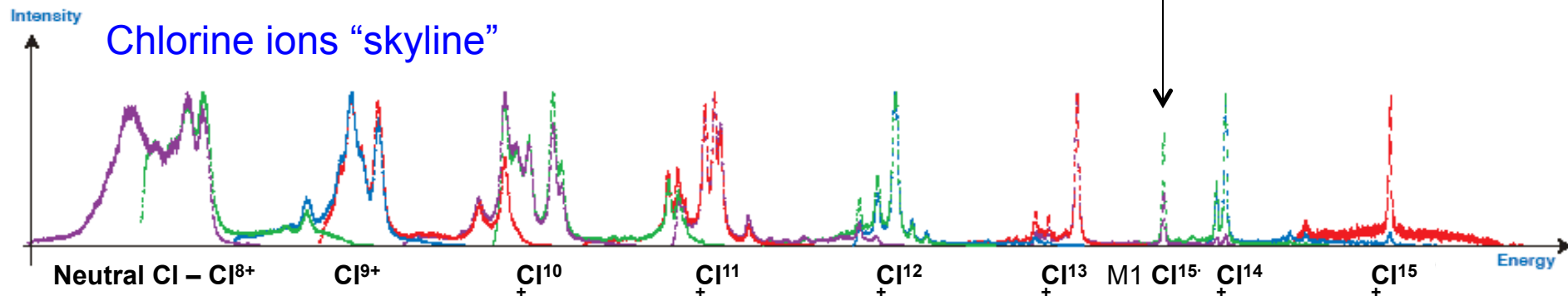
Ideal for spectrometer response
function studies

Measurement with:

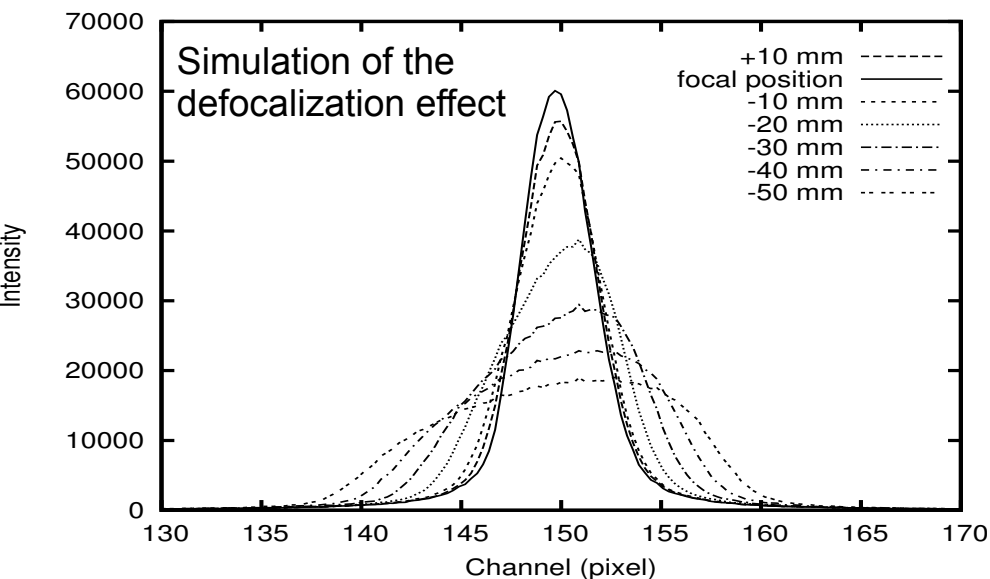
- Different crystal
- Different crystal aperture
- Different focal position
- Different X-ray energies (Ar, Cl, S)

- Transition width $\ll 1\text{meV}$
- Doppler broadening $< 40\text{ meV}$

$1s2s\ ^3S_1 \rightarrow 1s2s\ ^1S_0$
(M1)



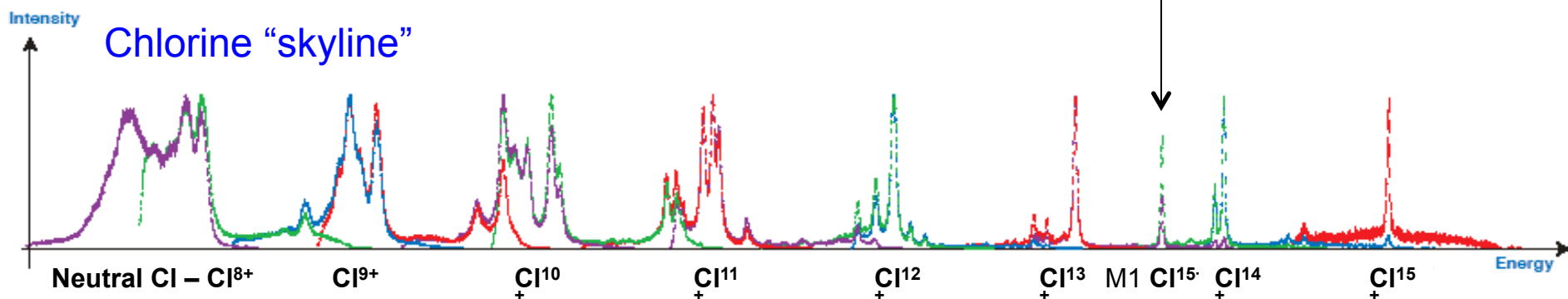
Spectrometer characterization from He-like S, Cl, Ar M1 transitions



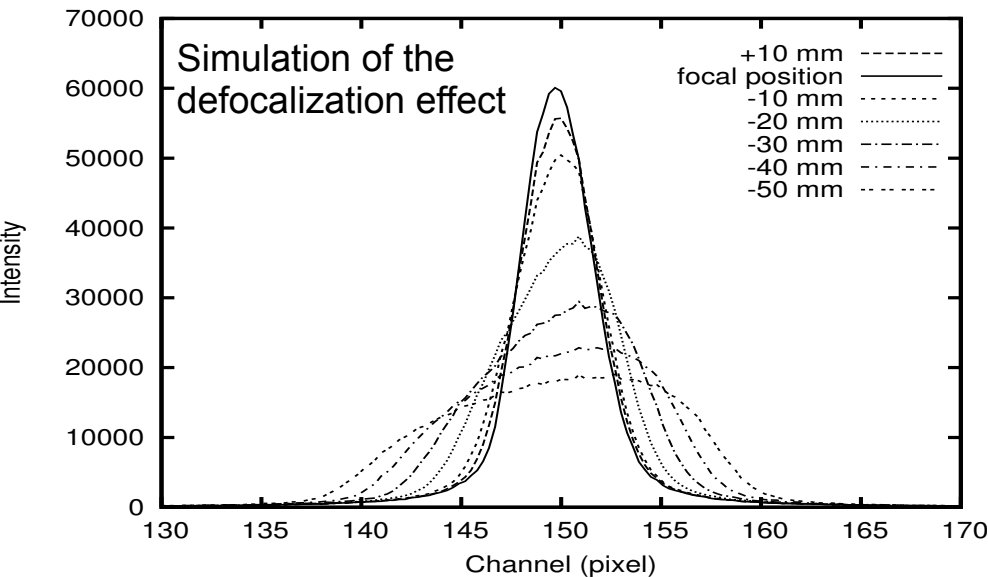
Ideal for spectrometer response function studies

Measurement with:

- Different crystal
- Different crystal aperture
- Different focal position
- Different X-ray energies (Ar, Cl, S)



Spectrometer characterization from He-like S, Cl, Ar M1 transitions



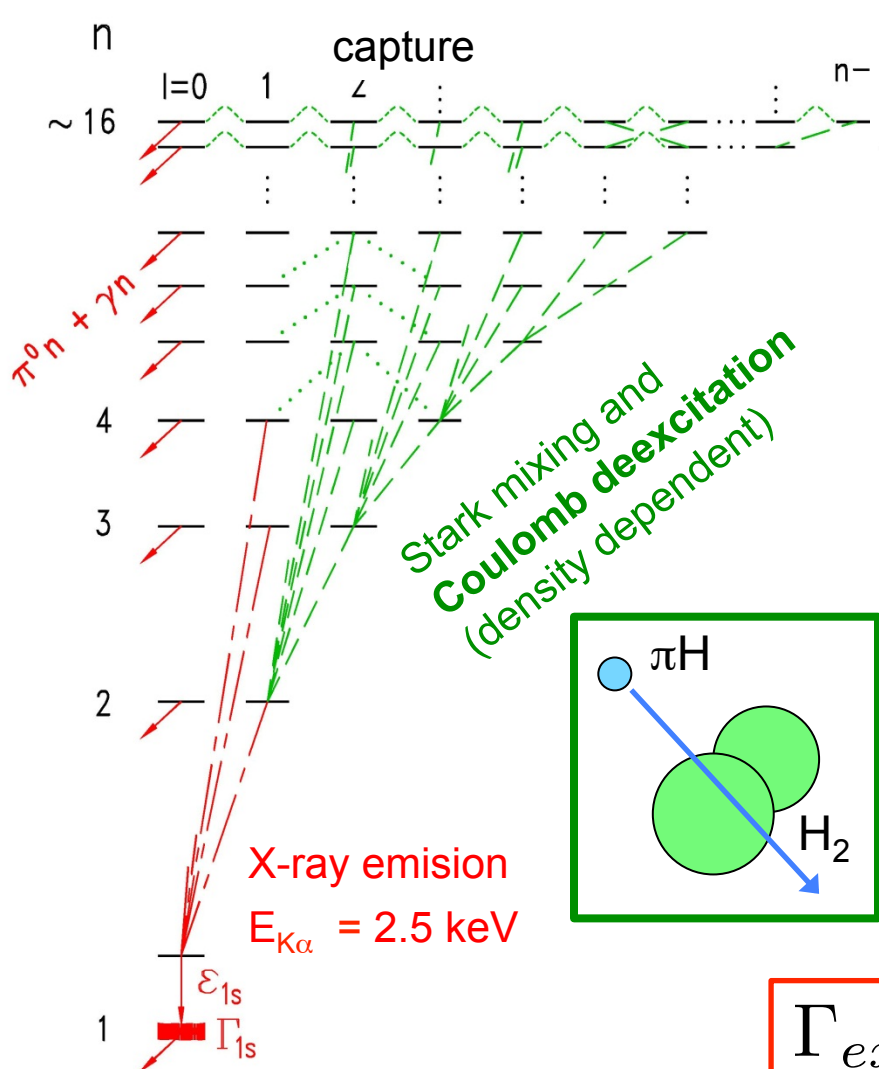
Ideal for spectrometer response
function studies

Measurement with:

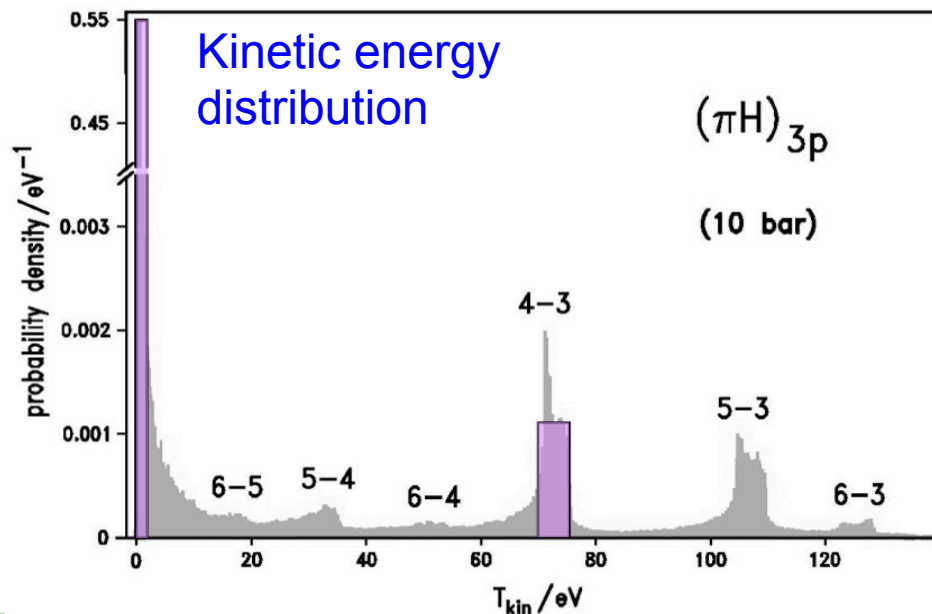
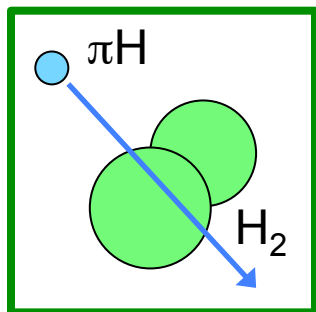
- Different crystal
- Different crystal aperture
- Different focal position
- Different X-ray energies (Ar, Cl, S)

$$\Gamma_{exp} = \Gamma_{crystal} \otimes \Gamma_{Doppler} \otimes \Gamma_{1s}$$

Transitions induced by collisions in πH

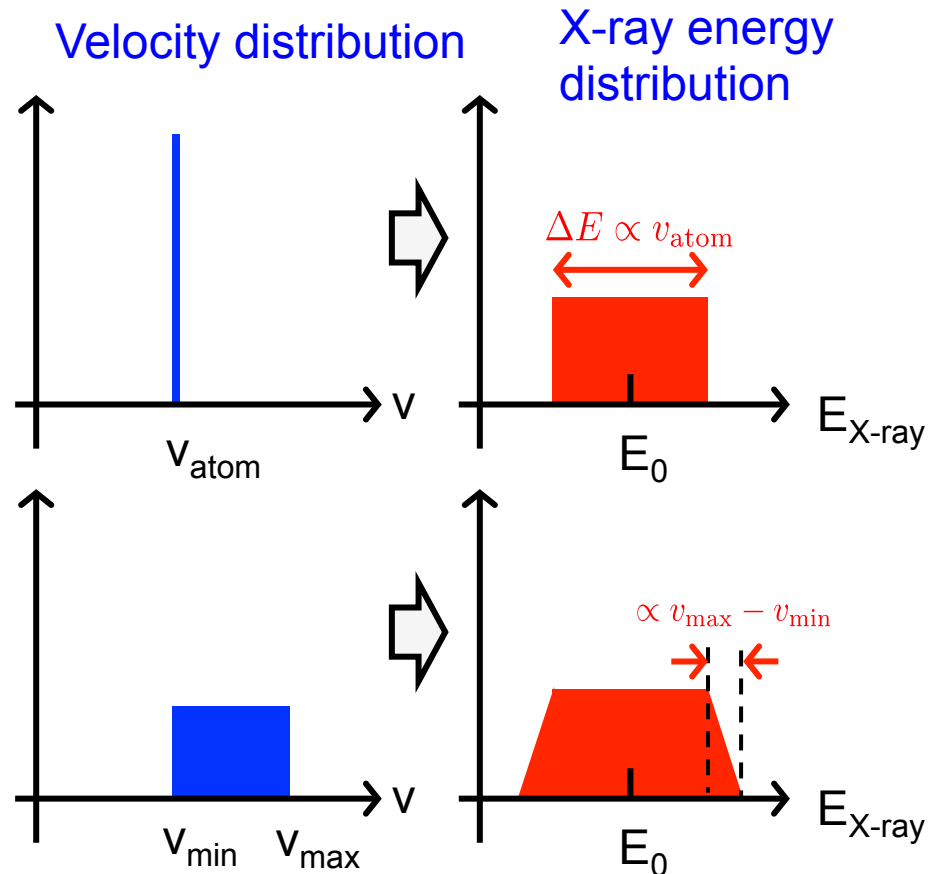
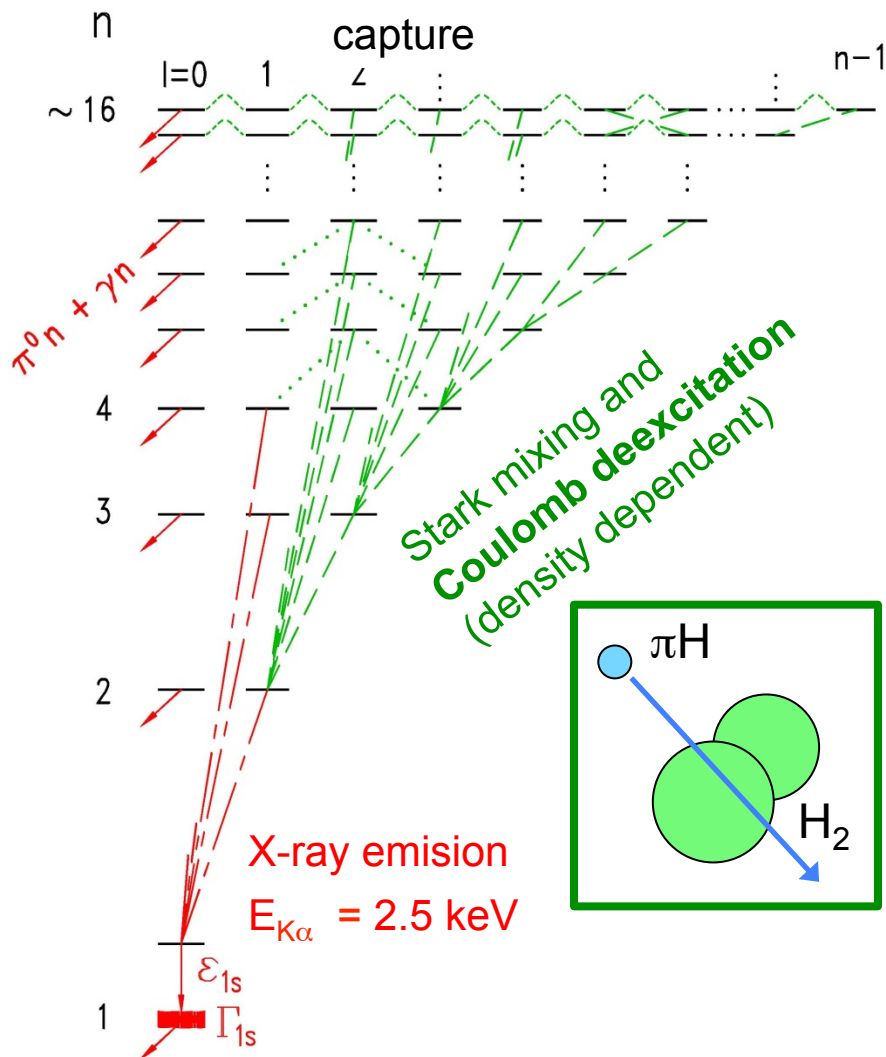


Stark mixing and
Coulomb deexcitation
(density dependent)

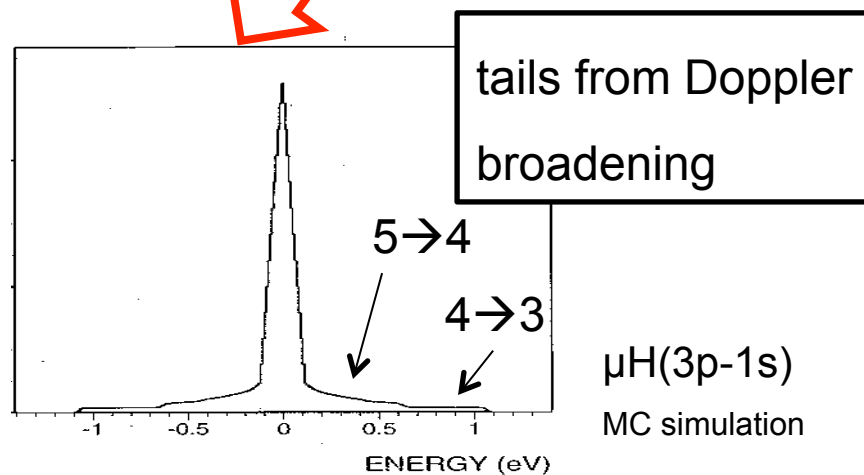
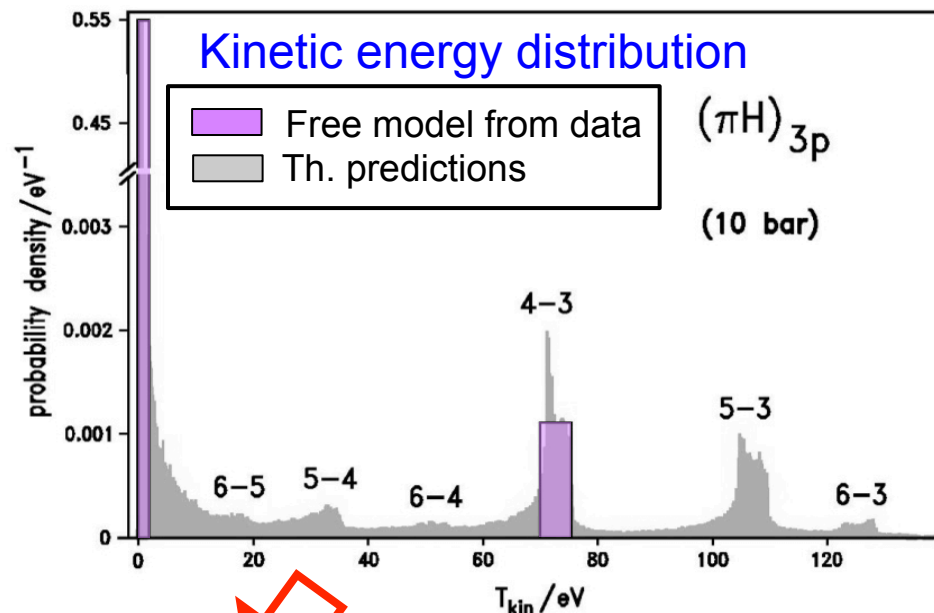
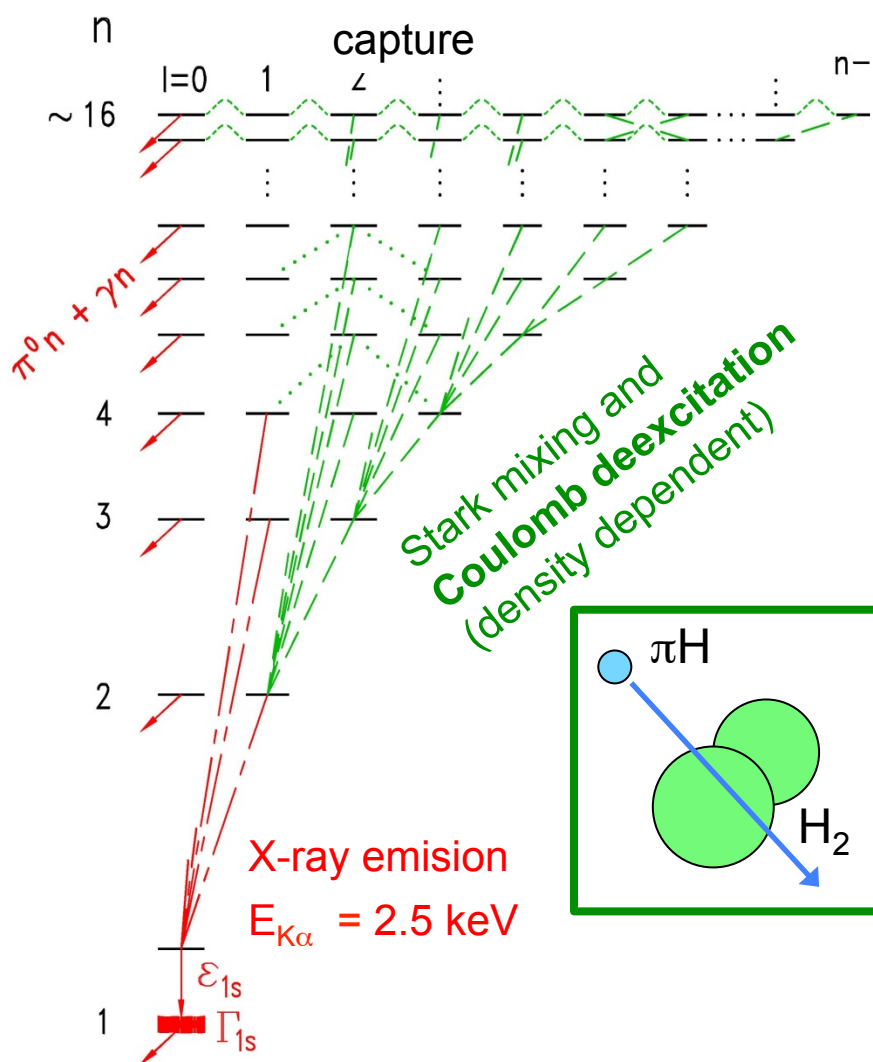


$$\Gamma_{\text{exp}} = \Gamma_{\text{crystal}} \otimes \Gamma_{\text{Doppler}} \otimes \Gamma_{1s}$$

Transitions induced by collisions in πH



Transitions induced by collisions in πH

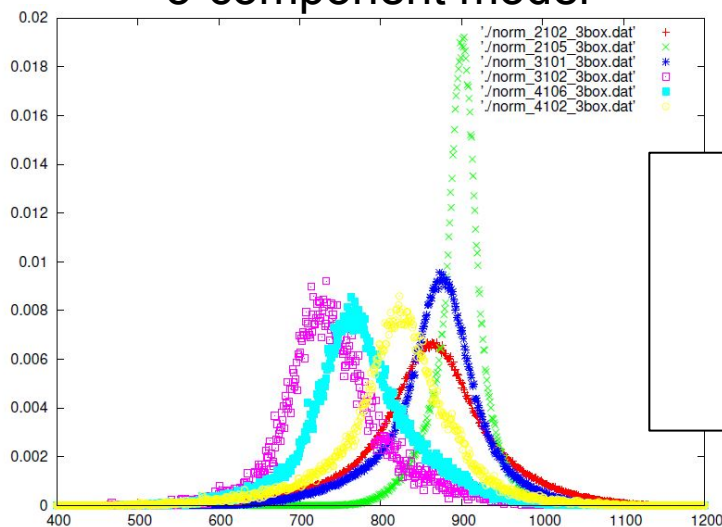


Work in progress

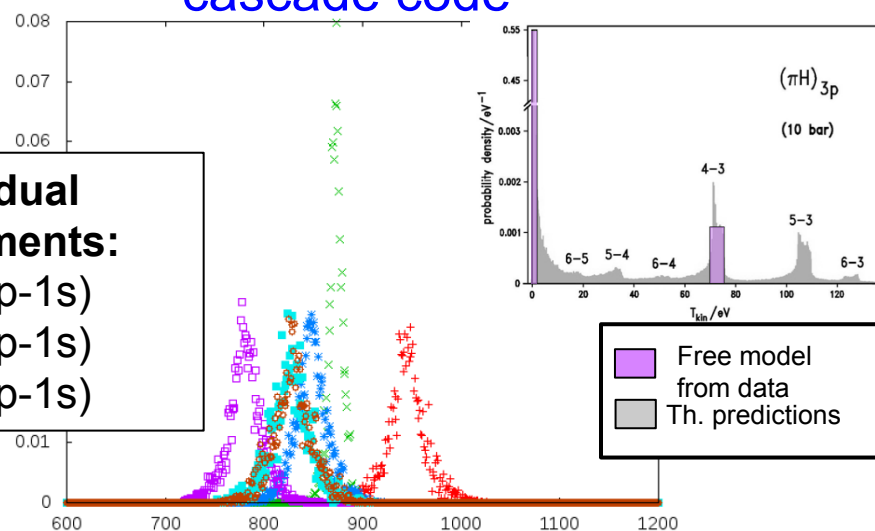
$$\Gamma_{exp} = \Gamma_{crystal} \otimes \Gamma_{Doppler} \otimes \Gamma_{1s}$$

Phenomenological approach
3-component model

Kinetic energy distribution from
cascade code



6 individual measurements:
2 x $\pi\text{H}(2p-1s)$
2 x $\pi\text{H}(3p-1s)$
2 x $\pi\text{H}(4p-1s)$



$$\Gamma_{1s} = 0.85^{+0.04}_{-0.05} \text{ eV} (\pm 4-5\%) \text{ preliminary}$$

Already 3-4 times better accuracy than Sigg et al. Phys. Rev. Lett. **75** 3245 (1995)
 $\Delta\Gamma \rightarrow \Delta\Gamma/2$ with a good cascade understanding

Strong interaction in simple bound systems

Shift and width related to the pion-nucleus scattering length and cross sections

$$\pi^H \left\{ \begin{array}{l} \epsilon_{1s}^{\pi H} \propto a_{\pi^- p \rightarrow \pi^- p} + corr. \\ \Gamma_{1s}^{\pi H} \propto (a_{\pi^- p \rightarrow \pi^0 n})^2 + corr. \end{array} \right. \quad \checkmark \quad \text{Ok...}$$

$$\pi^D \left\{ \begin{array}{l} \epsilon_{1s}^{\pi D} \propto a_{\pi^- p \rightarrow \pi^- p} + a_{\pi^- n \rightarrow \pi^- n} + corr. \\ \Gamma_{1s}^{\pi D} \propto \Im m(a_{\pi d \rightarrow \pi d}) \end{array} \right.$$

Basic pion-nucleon interaction properties

Strong interaction in simple bound systems

Shift and width related to the pion-nucleus scattering length and cross sections

$$\pi^H \left\{ \begin{array}{l} \epsilon_{1s}^{\pi^H} \propto a_{\pi^- p \rightarrow \pi^- p} + corr. \quad \checkmark \\ \Gamma_{1s}^{\pi^H} \propto (a_{\pi^- p \rightarrow \pi^0 n})^2 + corr. \quad \text{Ok} \dots \end{array} \right.$$

$$\pi^D \left\{ \begin{array}{l} \epsilon_{1s}^{\pi^D} \propto a_{\pi^- p \rightarrow \pi^- p} + a_{\pi^- n \rightarrow \pi^- n} + corr. \quad \checkmark \\ \Gamma_{1s}^{\pi^D} \propto \Im m(a_{\pi d \rightarrow \pi d}) \end{array} \right.$$

Basic pion-nucleon interaction properties

Strong interaction in simple bound systems

Shift and width related to the pion-nucleus scattering length and cross sections

$$\begin{aligned}
 \pi H \left\{ \begin{aligned} \epsilon_{1s}^{\pi H} &\propto a_{\pi^- p \rightarrow \pi^- p} + corr. && \propto a^+ + a^- + f_{LEC}(c_1, f_1, f_2) \\ \Gamma_{1s}^{\pi H} &\propto (a_{\pi^- p \rightarrow \pi^0 n})^2 + corr. && \propto (a^-)^2 + g(c_1, f_2) \end{aligned} \right. \\
 \pi D \left\{ \begin{aligned} \epsilon_{1s}^{\pi D} &\propto a_{\pi^- p \rightarrow \pi^- p} + a_{\pi^- n \rightarrow \pi^- n} + corr. && \propto a^+ + h.o.(a^+, a^-) + h_{LEC}(c_1, f_1, f_2) \\ \Gamma_{1s}^{\pi D} &\propto \Im m(a_{\pi d \rightarrow \pi d}) \end{aligned} \right.
 \end{aligned}$$

Low energy constants of the effective Chiral Perturbation Lagrangian

Isospin scattering lengths $a^\pm$

πN isospin symmetry
(2 quark flavour only)

$$1 \otimes 1/2 \Rightarrow 1/2 \oplus 3/2$$



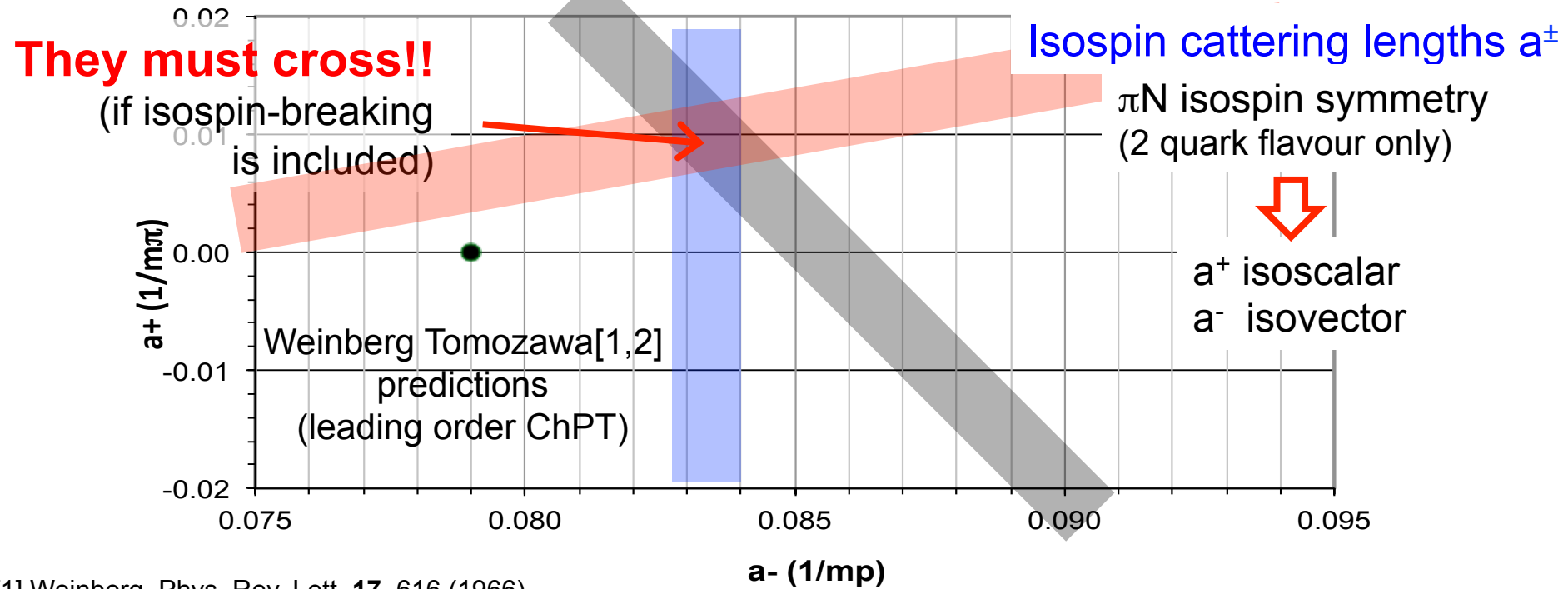
a^+ isoscalar
 a^- isovector

πN scattering lengths results

$$\pi H \left\{ \begin{array}{l} \epsilon_{1s}^{\pi H} \propto a^+ + a^- + f_{LEC}(c_1, f_1, f_2) \\ \Gamma_{1s}^{\pi H} \propto (a^-)^2 + g(c_1, f_2) \end{array} \right.$$

Qualitatively comparison
with the Chiral Pert. Theory
predictions

$$\pi D \quad \epsilon_{1s}^{\pi D} \propto a^+ + h.o.(a^+, a^-) + h_{LEC}(c_1, f_1, f_2)$$



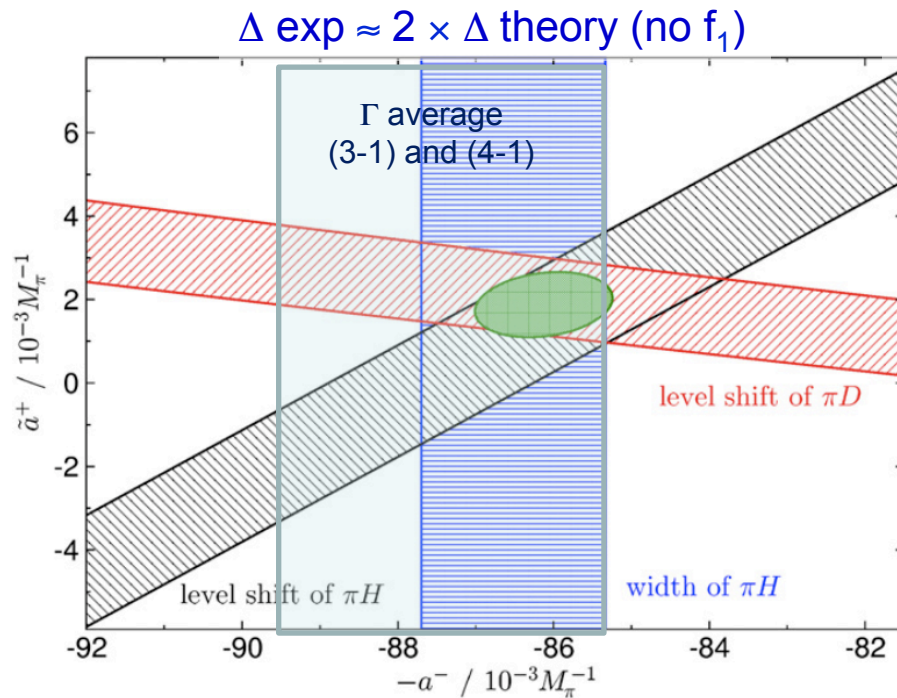
[1] Weinberg, Phys. Rev. Lett. **17**, 616 (1966)
[2] Y. Tomozawa, Nuovo Cim. A **46** 707 (1966)

πN scattering lengths results

$$\pi H \left\{ \begin{array}{l} \epsilon_{1s}^{\pi H} \propto a^+ + a^- + f_{LEC}(c_1, f_1, f_2) \\ \Gamma_{1s}^{\pi H} \propto (a^-)^2 + g(c_1, f_2) \end{array} \right.$$

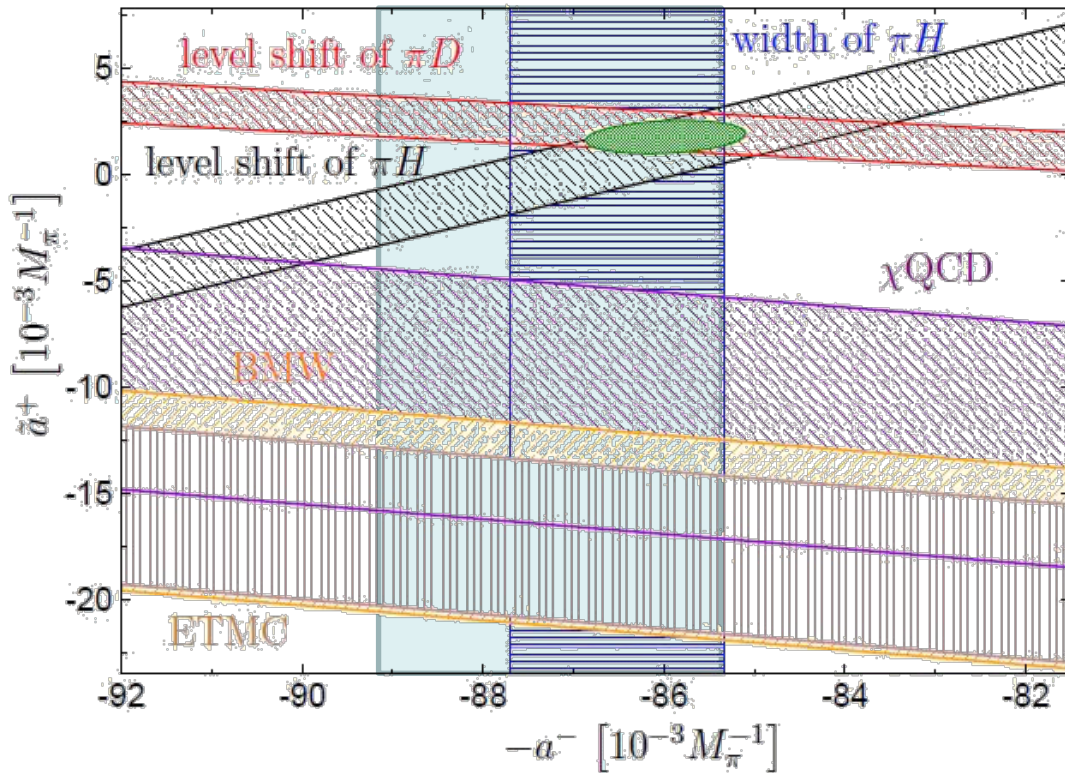
Quantitatively comparison
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$$\pi D \quad \epsilon_{1s}^{\pi D} \propto a^+ + h.o.(a^+, a^-) + h_{LEC}(c_1, f_1, f_2)$$



- consistency ✓
- $a^+ > 0$!!
- Constraints on f_1
- **πD decisive**

Beyond chiral perturbation theory, lattice calculations and dark matter detection



← pionic atoms analysis



← Lattice predictions of the **sigma term**

M. Hoferichter, J. Ruiz de Elvira, B. Kubis, and U.-G. Meißner, Phys. Lett. B **760**, 74-78 (2016)

$$a^+ \rightarrow \text{sigma term } \sigma_{\pi N} = \frac{m_u + m_d}{4M} \langle p | \bar{u}u + \bar{d}d | p \rangle \leftrightarrow \langle p | \bar{s}s | p \rangle \text{ contents} \rightarrow 59.1 \pm 3.5 \text{ MeV}$$

Crucial for the interpretation of dark-matter direct-detection experiments!!

J. Ellis, K.A. Olive, and C. Savage, Phys. Rev. D **77**, 065026 (2008)

$a^- \rightarrow$ pion-nucleus coupling constant via GMO sum rule

Summary

- Introduction
 - Test of QED in strong field
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- Heavy highly charged ions
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 - Vacuum decay experiment
- Pionic atoms
 - Production
 - Pionic hydrogen and deuterium spectroscopy
 - **Pion mass measurement**
- Conclusions and outlooks

Pion mass measurement from πN and μO X-ray spectroscopy

Pion mass (unknown) or muon mass (reference)

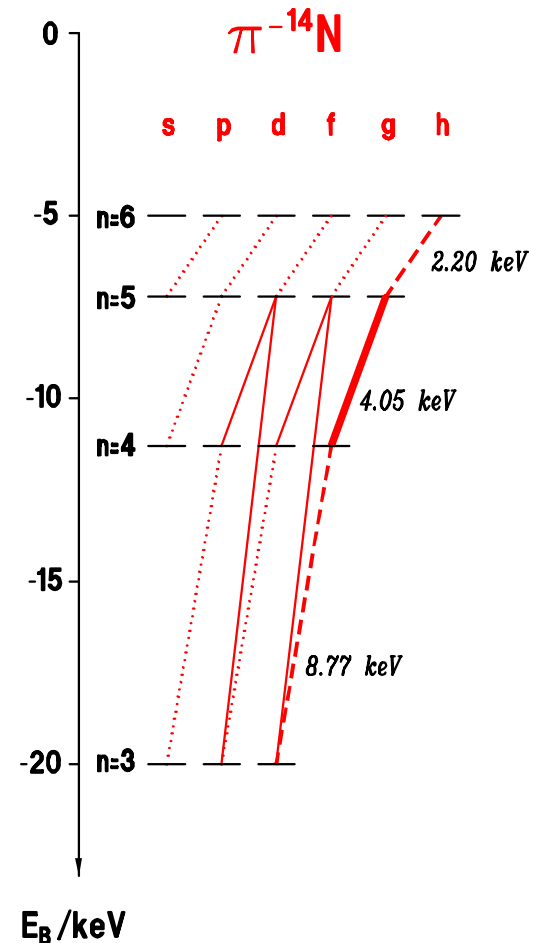
$$E_{\text{X-ray}} = mc^2 \frac{(Z\alpha)^2}{2} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right) + \mathcal{O}(Z^4\alpha^4)$$

QED calculation only

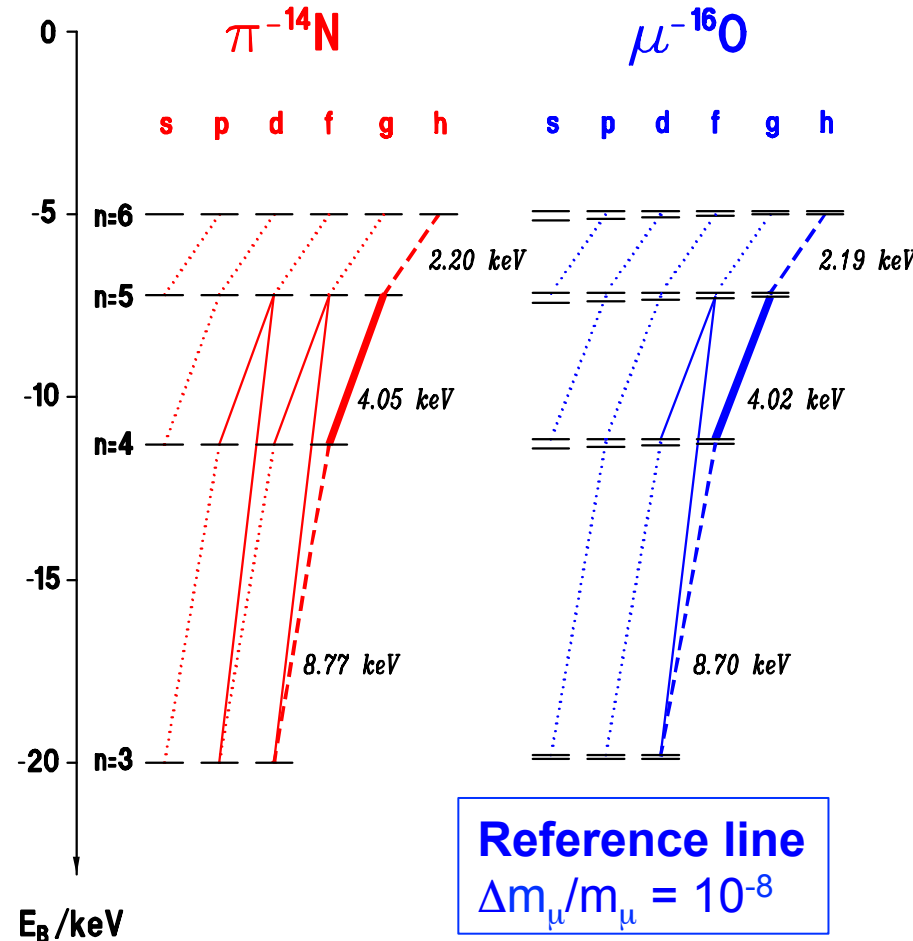
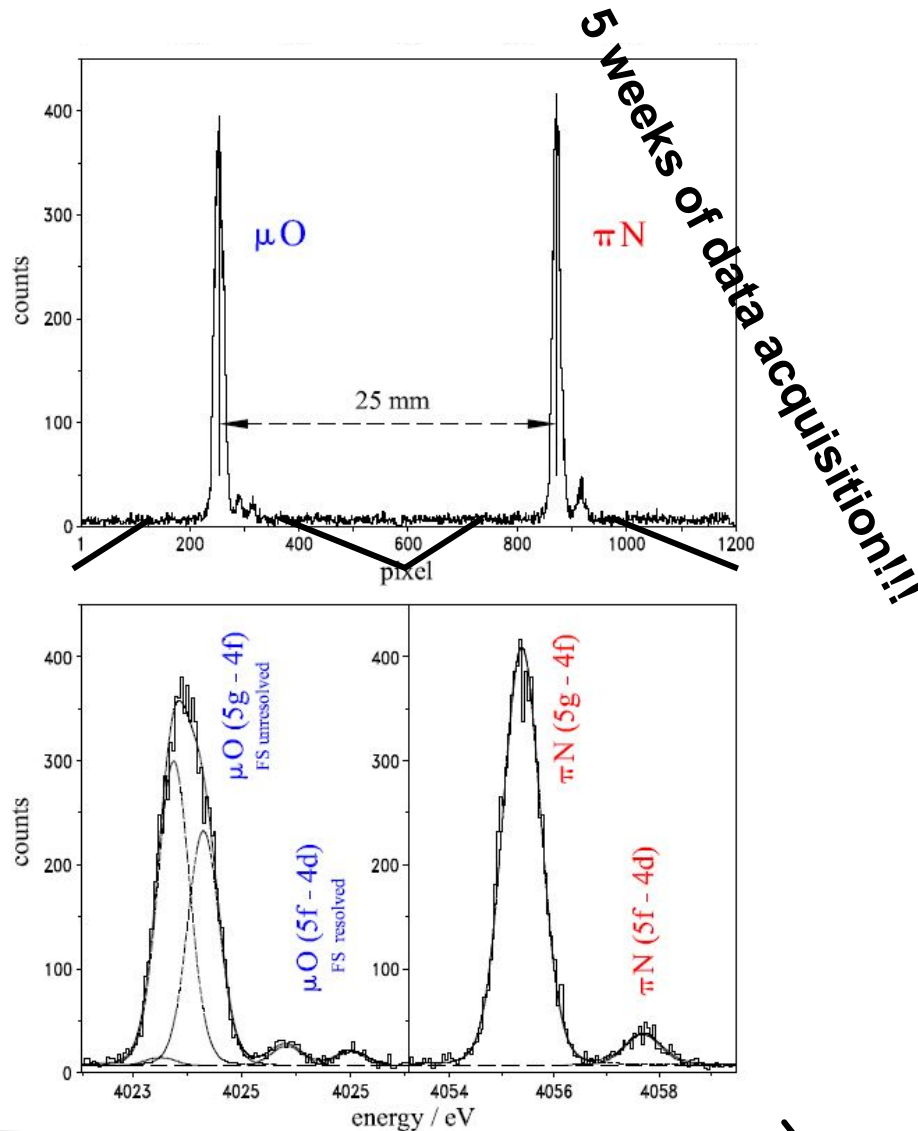
Circular transition enhanced

$$|n, \ell = n\rangle \rightarrow |n-1, \ell = n-1\rangle$$

Strong interaction effects minimized



5g-4f transitions in πN and μO



Muon = $\frac{1}{2}$ -spin particle
→ fine structure in μO

From the line position to the pion mass

From the spatial diff. to the angular position diff.

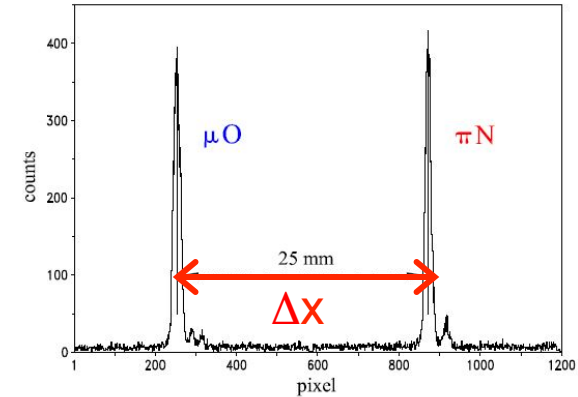
$$\Delta\Theta = -2 \arctan \left(\frac{\Delta x}{2D} \right)$$

From the angular position diff. to the transition energy

$$E_{\pi N} = E_{\mu O} \frac{1}{\cos \Delta\Theta - \frac{\sqrt{1 - [hc/(2dE_{\mu O})]^2}}{hc/(2dE_{\mu O})} \sin \Delta\Theta}$$

where

$$E_{\mu O} = f_{\text{Dirac}}^{\text{QED}}(m_{\mu}) = \tilde{m}_{\mu} c^2 \frac{(Z\alpha)^2}{2} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right) + \mathcal{O}_{\text{Dirac}}^{\text{QED}}(Z^4 \alpha^4)$$



$$\frac{hc}{E} = 2d \sin \Theta_B$$

Bragg law

From the transition energy to the **pion mass**

$$m_{\pi} = f_{\text{Klein-Gordon}}^{\text{QED}}(E_{\pi N})$$



QED calculation (for πN) [1], ...

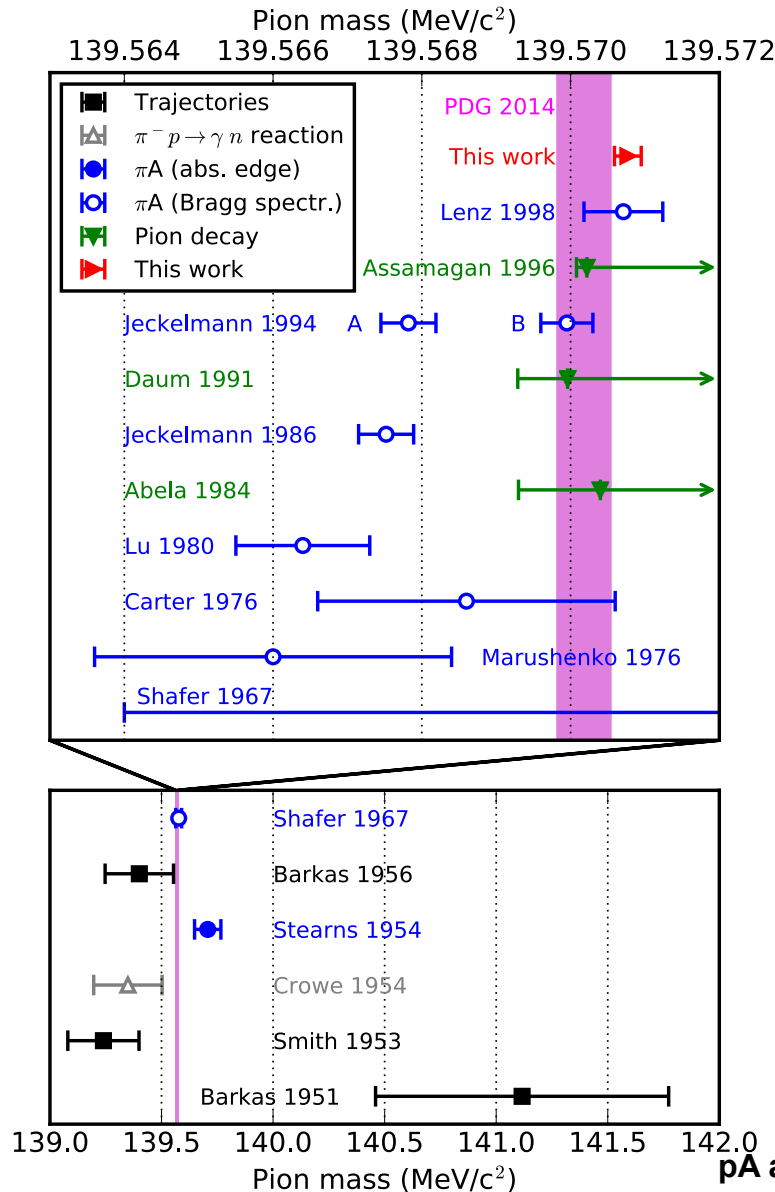
[1] M.T. and P. Indelicato, Phys. Rev. A **76**, 012510(2007)

List of systematic effects

type of uncertainty	μO / arcsec	πN / arcsec	total / arcsec	uncertainty / ppb
index of refraction shift	13.22	12.94	-0.28	± 20
silicon lattice constant				± 2
bending correction	14.01	13.71	0.30	± 20
penetration depth correction	-0.07	-0.07	0	± 4
focal length				± 670
CCD alignment				± 340
pixel distance				± 120
alignment of detector normal				+ 0 - 30
detector height offset				+ 0 - 35
shape of target window				± 100
shape of reflection				± 225
individual curvature correction				± 150
temperature correction				± 30
response function and Doppler broadening				+ 290 - 350
line pattern modelling				+ 190 - 290
fit interval				± 15
QED energy				± 350
conversion constant hc				± 2
4f strong interaction $45 \mu\text{eV}$		0.003	-0.003	± 10
5g strong interaction $0.2 \mu\text{eV}$		0.000	0.000	± 0
K electron screening				± 0
total systematic error				+ 950 - 1000
statistical error				± 820

Total uncertainty: 1.3×10^{-6}
 Statistical: 0.8×10^{-6}
 Systematics: $\sim 1 \times 10^{-6}$

The new measurement of the pion mass



Particle data group [1]:
 $139.57018 \pm 0.00035 \text{ MeV}/c^2$
 $\rightarrow 2.5 \times 10^{-6}$ accuracy

Our work [2]:
 $139.57077 \pm 0.00018 \text{ MeV}/c^2$
 $\rightarrow 1.3 \times 10^{-6}$ accuracy

- No effect of eventual remaining K-shell electrons ($< 10^{-9}$)
- High accuracy calibration line (0.25×10^{-6} from theory calc.)

Improved accuracy for X-ray standards from pionic atoms [3]

- [1] Particle Data Group. Chinese Phys. C **38**, 090001 (2014)
 [2] M.T. *et al.*, Phys. Lett. B **759**, 583-588 (2016)
 [3] D.F. Anagnostopoulos *et al*, Phys. Rev. Lett. **91**, 240801 (2003)

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 - Vacuum decay experiment
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Conclusions and perspectives for pionic atoms

- End of a long measurement campaign
- New input for investigation on strong interaction at low energy
- Deeper understanding (but not yet complete) of the atomic cascade and cross sections in exotic atomic systems
- New measurement of the pion mass

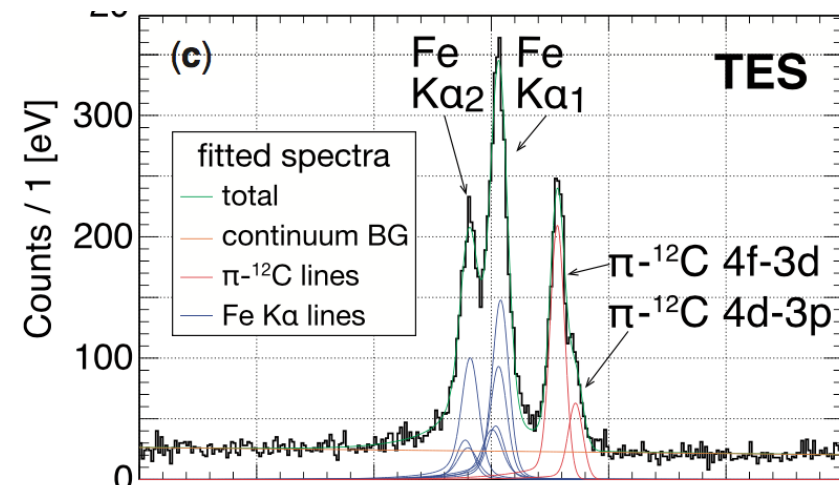
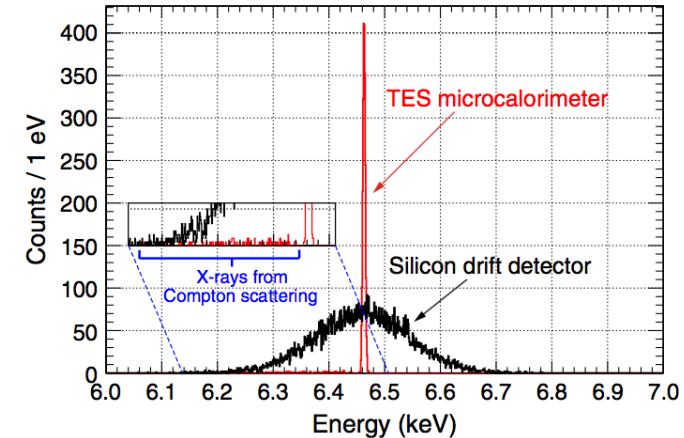
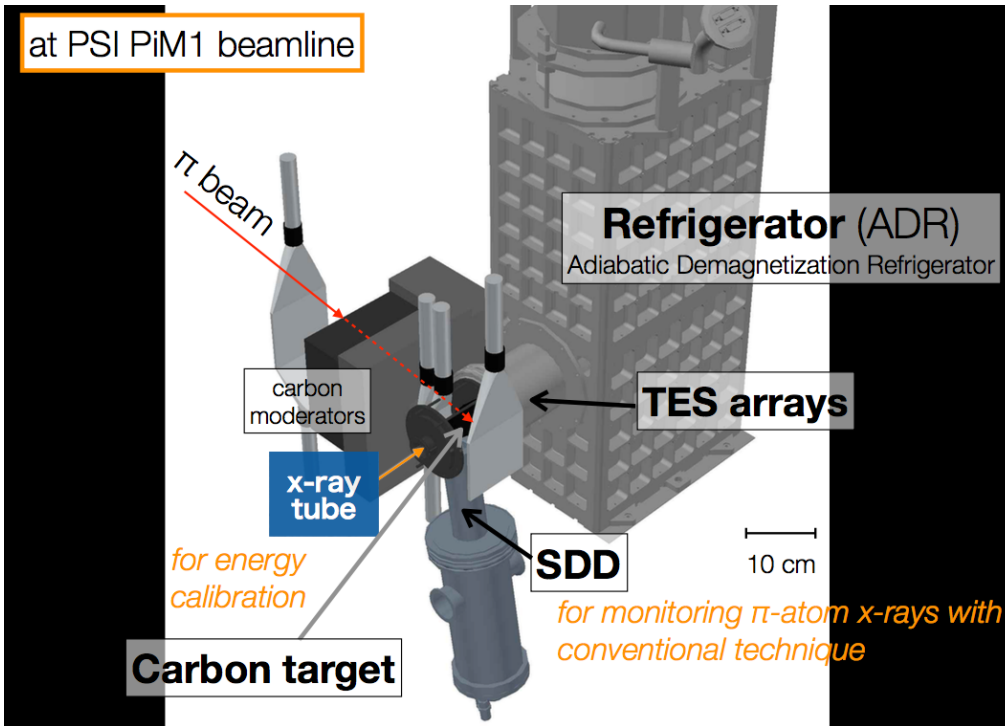
Next...

- Atomic cascade and cross sections calculations has to be improved
- New high-resolution experiments with πT and πHe ? $\rightarrow$ Additional constraints on a^+ a-
- ...
- Resolution improvement?
 - Crystal spectrometer: $0.5 \rightarrow 0.2$ eV?
 - Microcalorimeter [1]? Lower resolution but higher efficiency
- New measurement of the pion mass: laser spectroscopy of πHe^+
Proposed experiment at PSI from Hori and co. [2]

[1] S. Okada *et al.*, Prog. Th.Exp. Phys. **2016**, (2016)

[2] M. Hori *et al.*, Phys. Rev. A **89**, 042515 (2014)

Microcalorimeter measurement of π A and KA



S. Okada, D.A. Bennett, C. Curceanu *et al.*, Progress of Theoretical and Experimental Physics **2016**, (2016)
 S. Okada, D.A. Bennett, W.B. Doriese *et al.*, J. Low Temp. Phys. **176**, 1015-1021 (2014)

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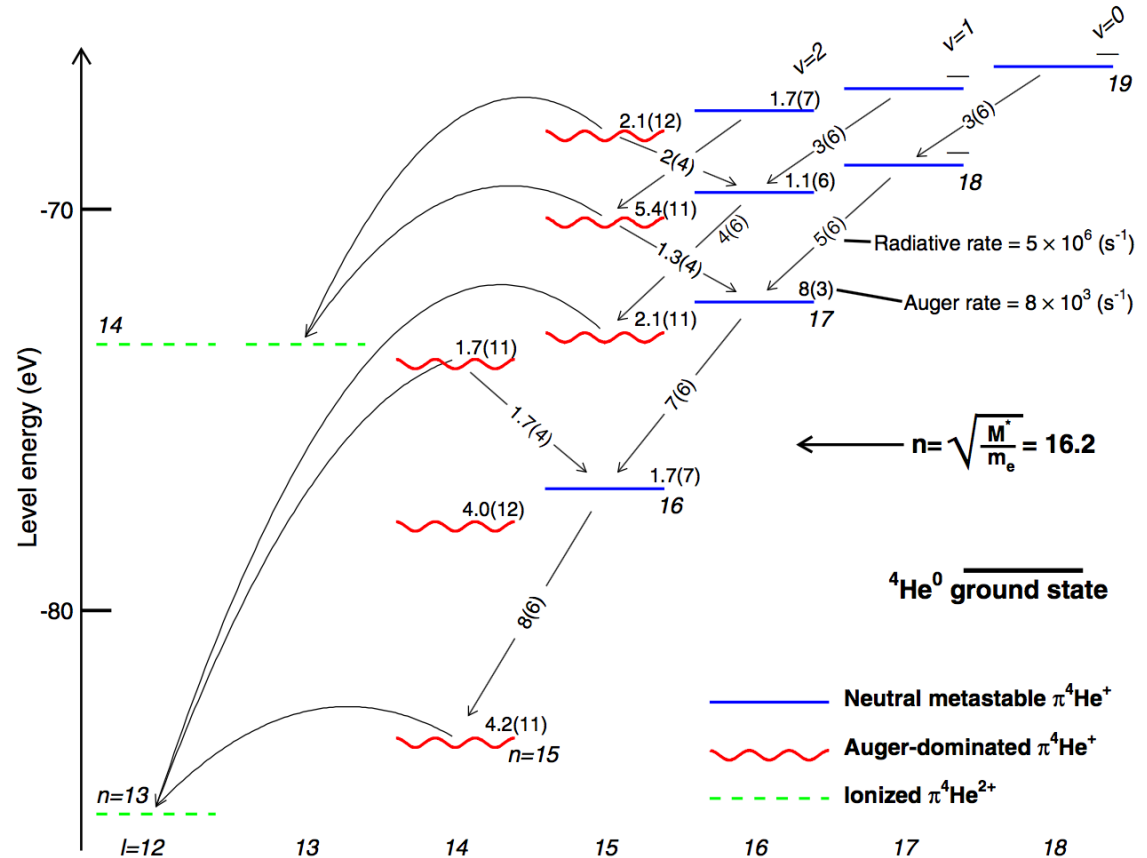
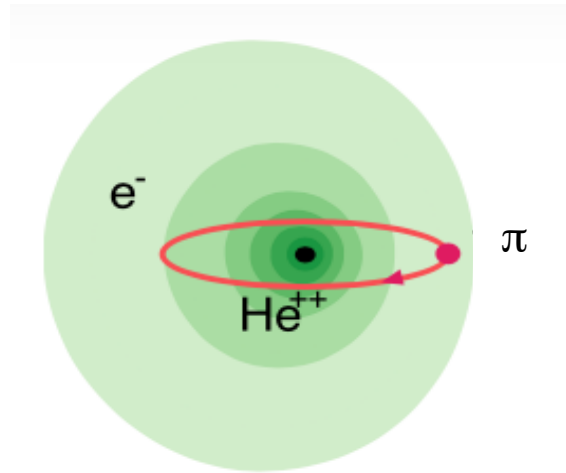
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[2] M. Hori *et al.*, Phys. Rev. A **89**, 042515 (2014)

Pion mass measurement by laser spectroscopy of πHe^+



Goal accuracy: better than 1ppm

M. Hori *et al.*, Hyperfine Interact. **233**, 83-87 (2015).

M. Hori *et al.*, Phys. Rev. A **89**, 042515 (2014).

Conclusions and perspectives for heavy highly charged ion

- New QED tests in the strong Coulomb field (non perturbative regime)
- New phenomenon expected (spontaneous emission of positrons from the vacuum decay)
- Two-loop QED not yet tested in hydrogenlike systems

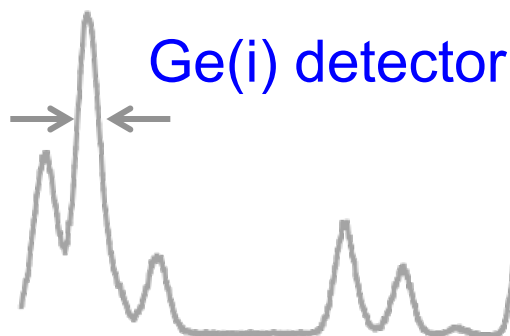
Next...

- Many new experiments already in preparation for the future
- New experimental methods in development
- New FAIR facility in construction and commissioning

Towards an accuracy of 1 eV of the 1s level Lamb shift

Resolution:
400 eV at 60 keV

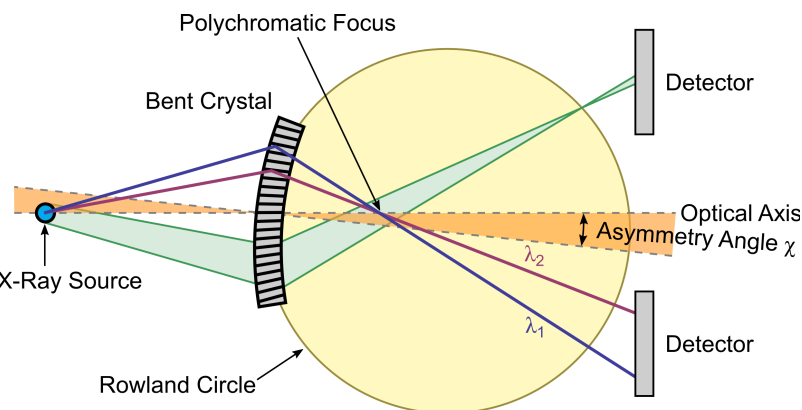
$$\varepsilon = 10^{-4}$$



U^{91+} Lamb shift $\delta E = 4.6$ eV

Crystal spectroscopy

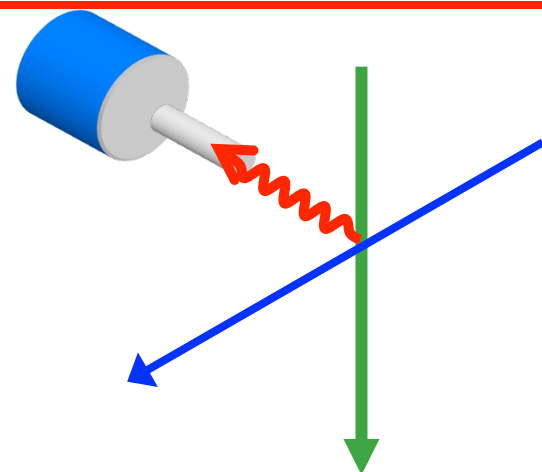
Res: **~ 50 eV** at 60 keV **$\varepsilon = 10^{-8}$**



FOCAL spectrometer
T. Gassner, M.T. et al submitted to PRL

Microcalorimeters

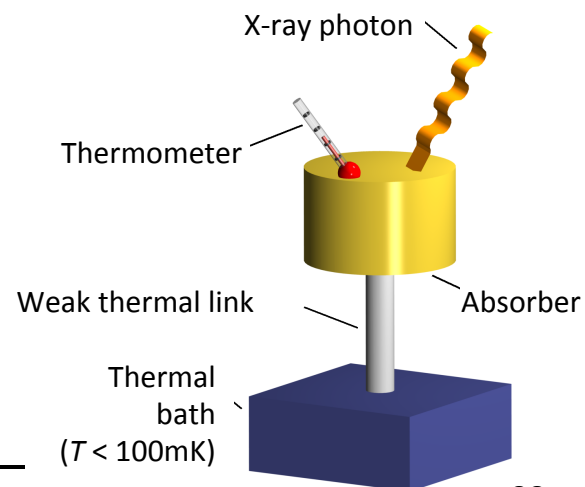
Res: **~ 50 eV** at 60 keV **$\varepsilon = 10^{-4-6}$**



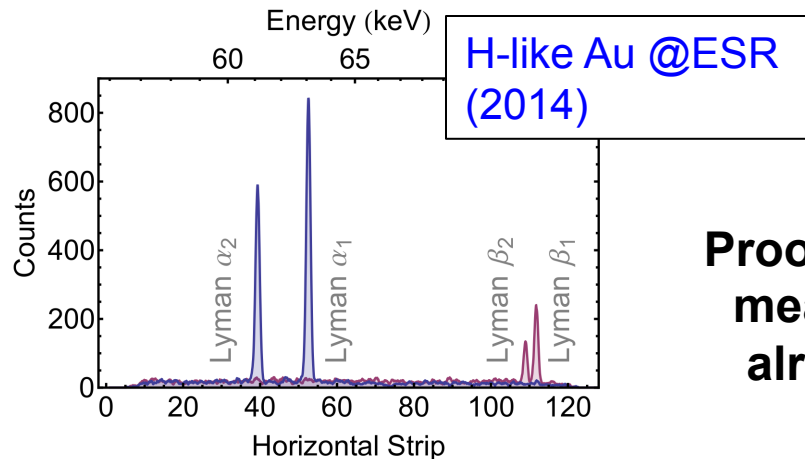
**Resistive and
magnetic
microcalorimeters**

S. Kraft-Bermuth *et al.*, J. Phys. B
50, 055603 (2017)

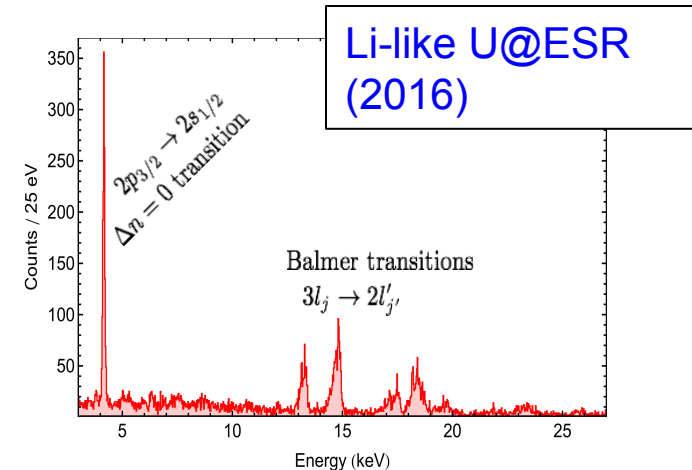
C. Pies *et al.*, J. Low Temp. Phys.
167, 269-279 (2012)



Towards an accuracy of 1 eV of the 1s level Lamb shift

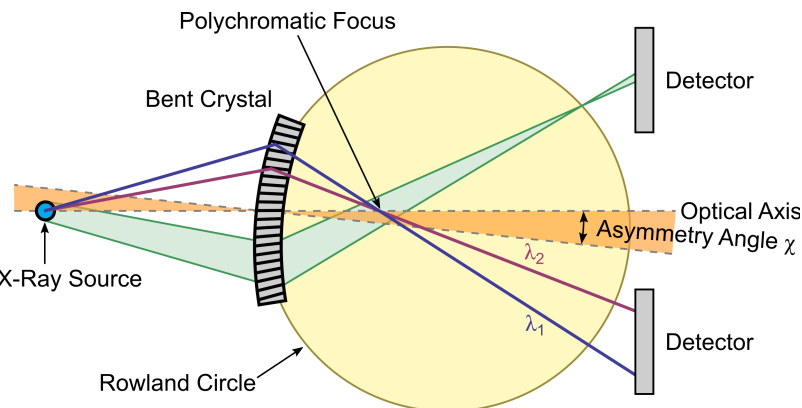


Proof of principle measurements already done!



Crystal spectroscopy

Res: ~50 eV at 60 keV $\epsilon = 10^{-8}$



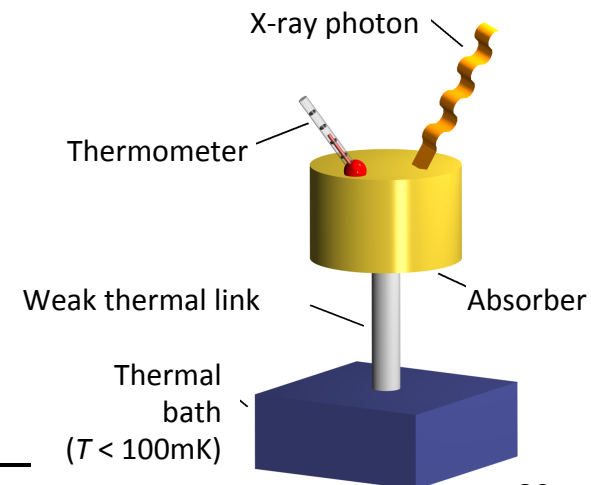
FOCAL spectrometer
T. Gassner, M.T. et al submitted to PRL

Microcalorimeters

Res: ~50 eV at 60 keV $\epsilon = 10^{-4-6}$

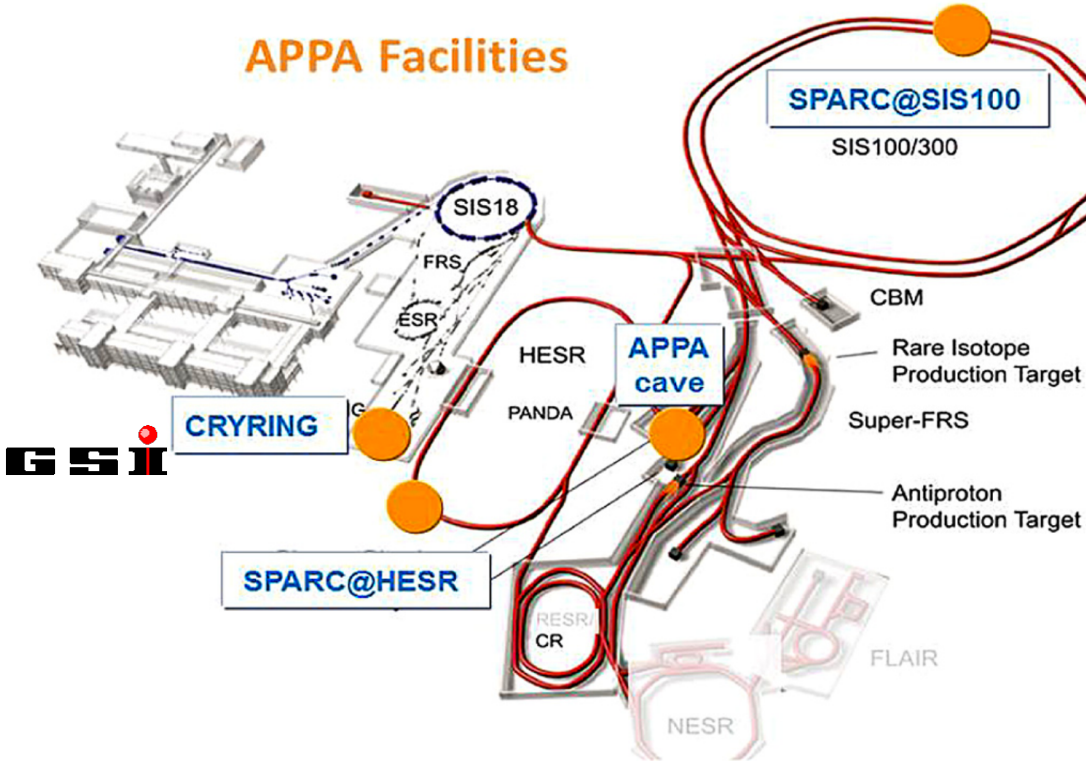
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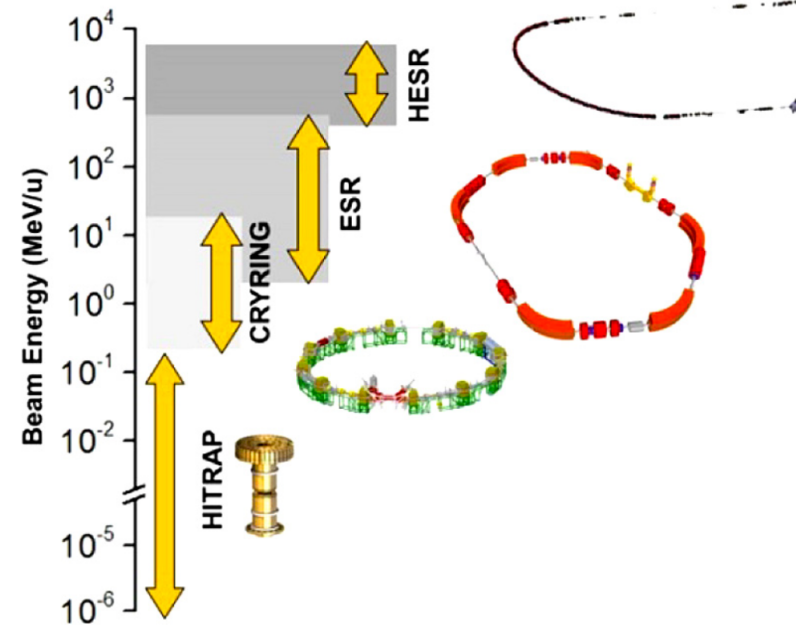


FAIR facility

APPA Facilities



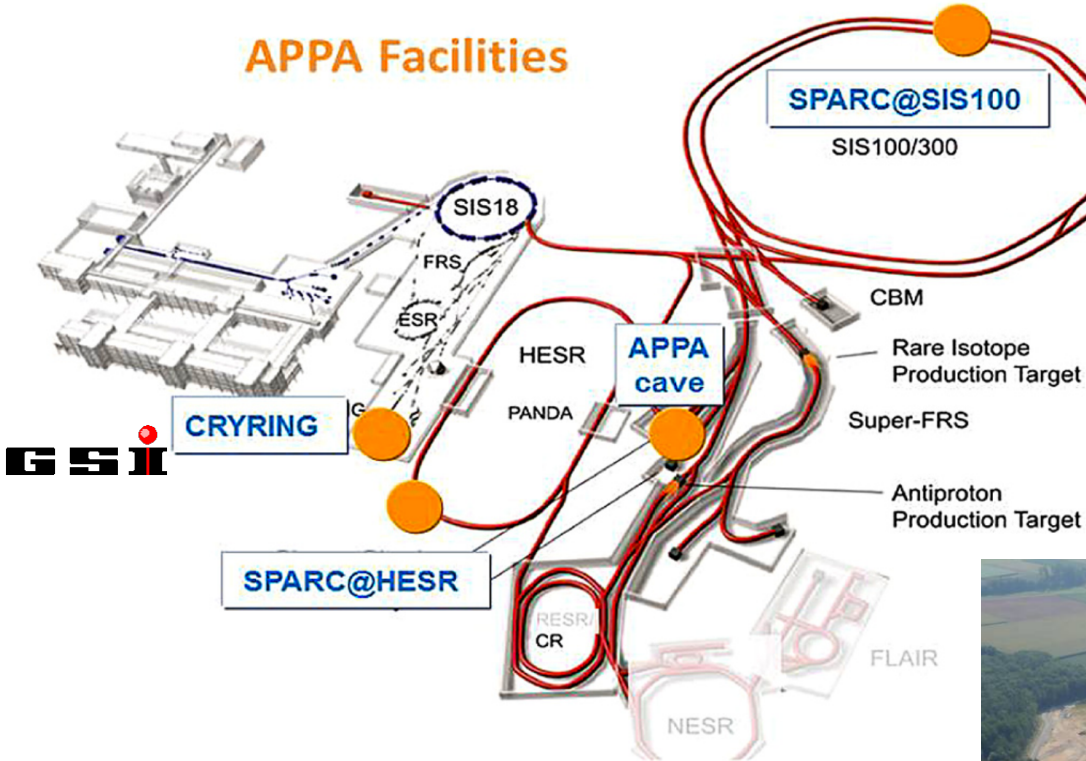
FAIR



Availability of highly charged ions
up to bare U with energies between
10 GeV/u to rest

FAIR facility

APPA Facilities



FAIR

Phase-0: 2018-2019

Phase-1: 2024

Availability of highly charged ions
up to bare U with energies between
10 GeV/u to rest





D. Banas, H.F. Beyer, K.H. Blumenhagen, F. Bosch, C. Brandau, A. Bräuning-Demian, H. Bräuning, W. Chen, Chr. Dimopoulou, E. Förster, T. Gaßner, S. Geyer, A. Gumberidze, S. Hagmann, S. Heß, R. Heß, P.-M. Hillenbrand, P. Indelicato, P. Jagodzinski, T. Kampfer, Chr. Kozhuharov, M. Lestinsky, D. Liesen, Yu.A. Litvinov, R. Loetzsch, B. Manil, R. Märtin, F. Nolden, N. Petridis, M.S. Sanjari, K.S. Schulze, M. Schwemlein, A. Simionovici, U. Spillmann, M. Steck, Th. Stöhlker, C.I. Szabo, S. Trotsenko, I. Uschmann, G. Weber, O. Wehrhan, N. Winckler, D. Winters, N. Winters, E. Ziegler

F. D. Amaro, D. F. Anagnostopoulos, A. Blechmann, S. Biri, P. Bühler, D. S. Covita, J. M. F. dos Santos, H. Fuhrmann, H. Gorke, D. Gotta, A. Gruber, M. Hennebach, A. Hirtl, T. Ishiwatari, Th. Jensen, E.-O. Le Bigot, V. E. Markushin, J. Marton, M. Nekipelov, V. Pomerantsev, V. Popov, S. Schlessler, A. Schmelzbach, Ph. Schmid, L. M. Simons, Th. Strauch, M. Theisen, J. F. C. A. Veloso, J. Zmeskal

S. Cervera, V. Gafton, T. Kirchner, E.V. Gafton, D. Hannani, S. Hidki, E. Lacaze, E. Lamour, A. Lévy, S. Macé, M. Maranglolo, J. Mérot, F. Mezdari, J.-Y. Pacquet, C. Prigent, J.-M. Ramillon, J.-P. Rozet, A. Salehzadeh, S. Steydli, D. Vernhet

