

The physics and astrophysics of merging neutron-star binaries

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Firenze
25.01.17



Plan of the talk

- * The benefits of studying merging binary NSs
- * Anatomy of GW signal: frequencies and EOS
- * GW170817 and radius measurements
- * Dissipative effects: are they important?
- * Ejected mass and nucleosynthesis

The two-body problem in GR

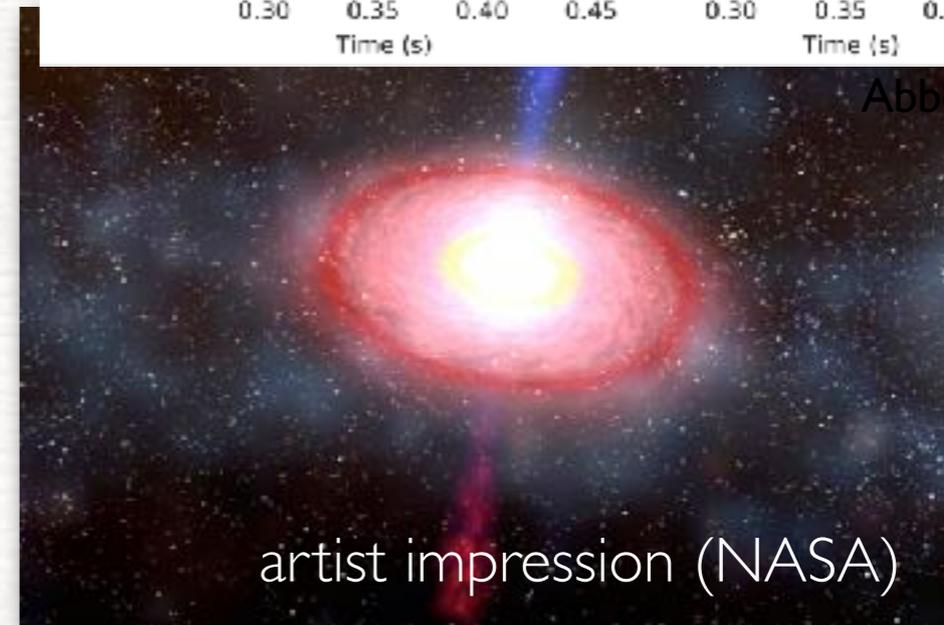
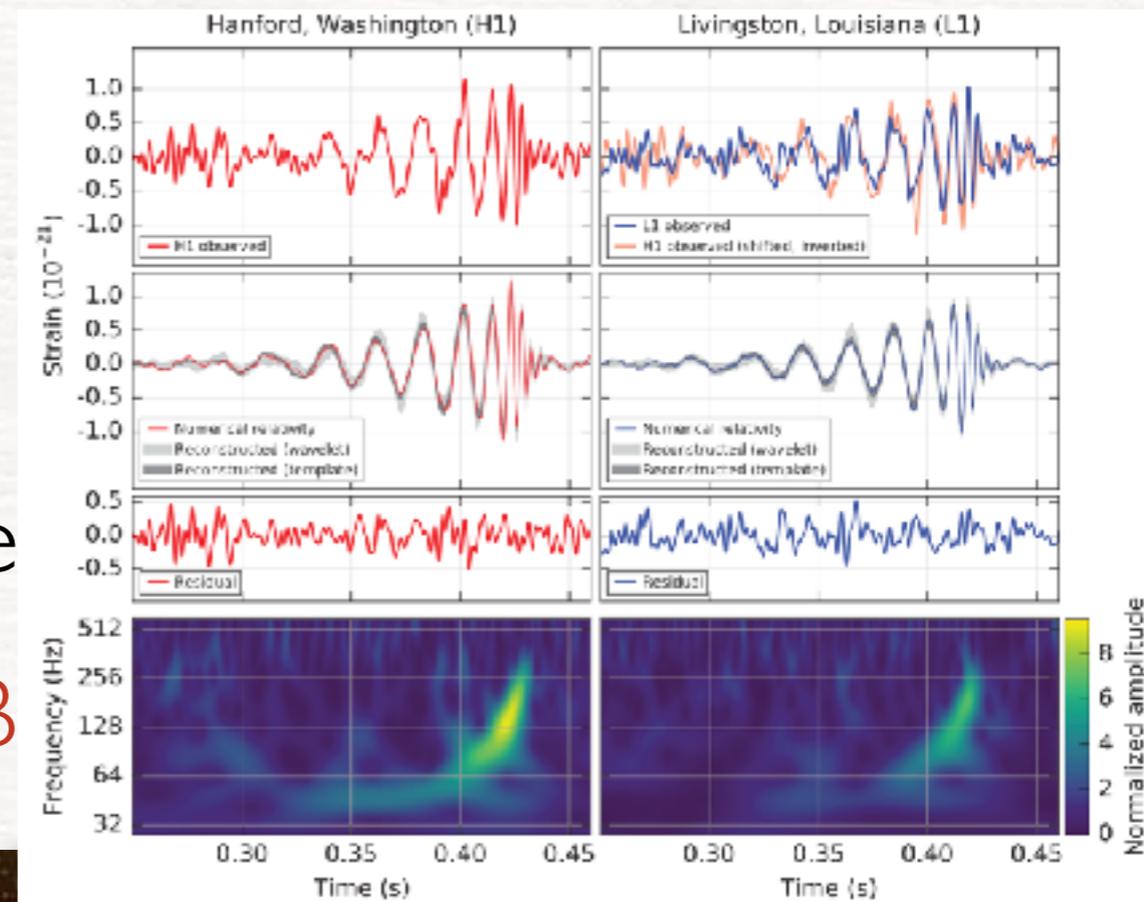
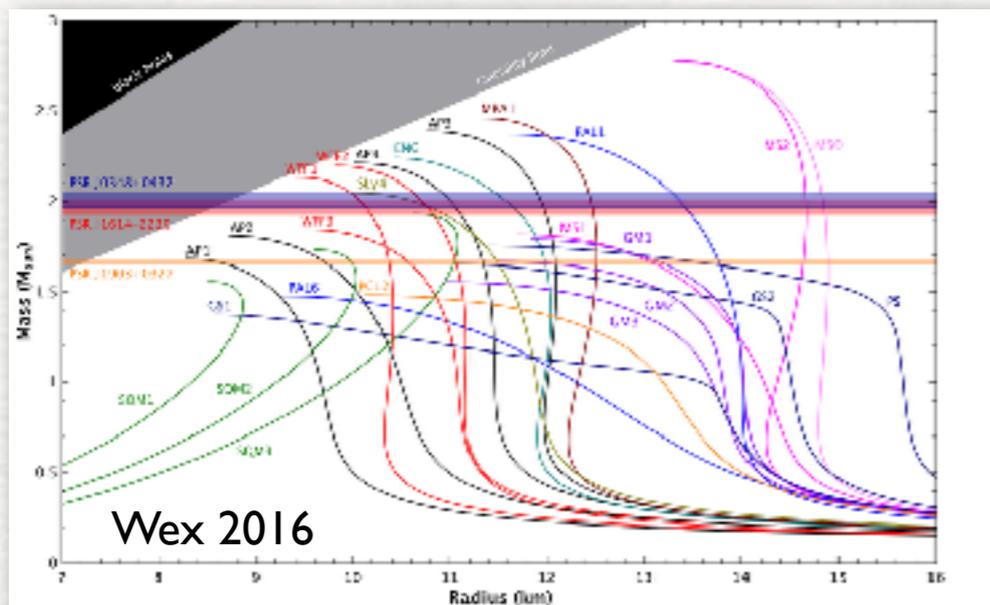
- For BHs we know what to **expect**:

$$\text{BH} + \text{BH} \longrightarrow \text{BH} + \text{GWs}$$

- For NSs the question is more **subtle**: hyper-massive neutron star (HMNS), ie

$$\text{NS} + \text{NS} \longrightarrow \text{HMNS} + \dots ? \longrightarrow \text{BH}$$

- **HMNS** phase can provide clear information on **EOS**



- **BH+torus** system may tell us on the central engine of **GRBs**

The two-body problem in GR

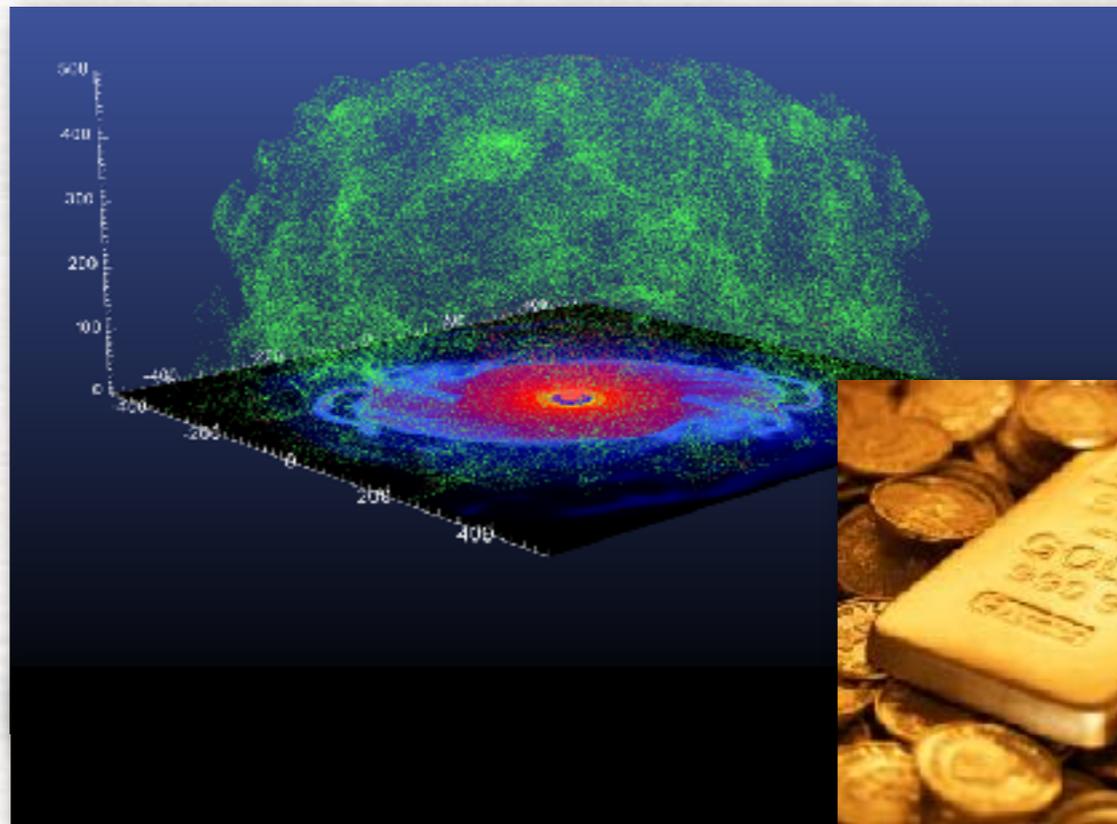
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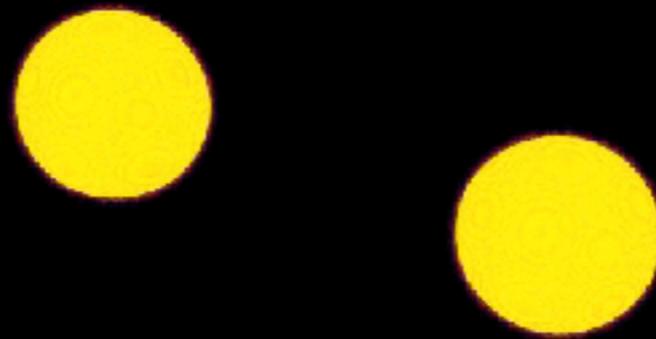


- For NSs the question is more **subtle**: the merger leads to an hyper-massive neutron star (HMNS), ie a metastable equilibrium:



- **ejected matter** undergoes nucleosynthesis of heavy elements



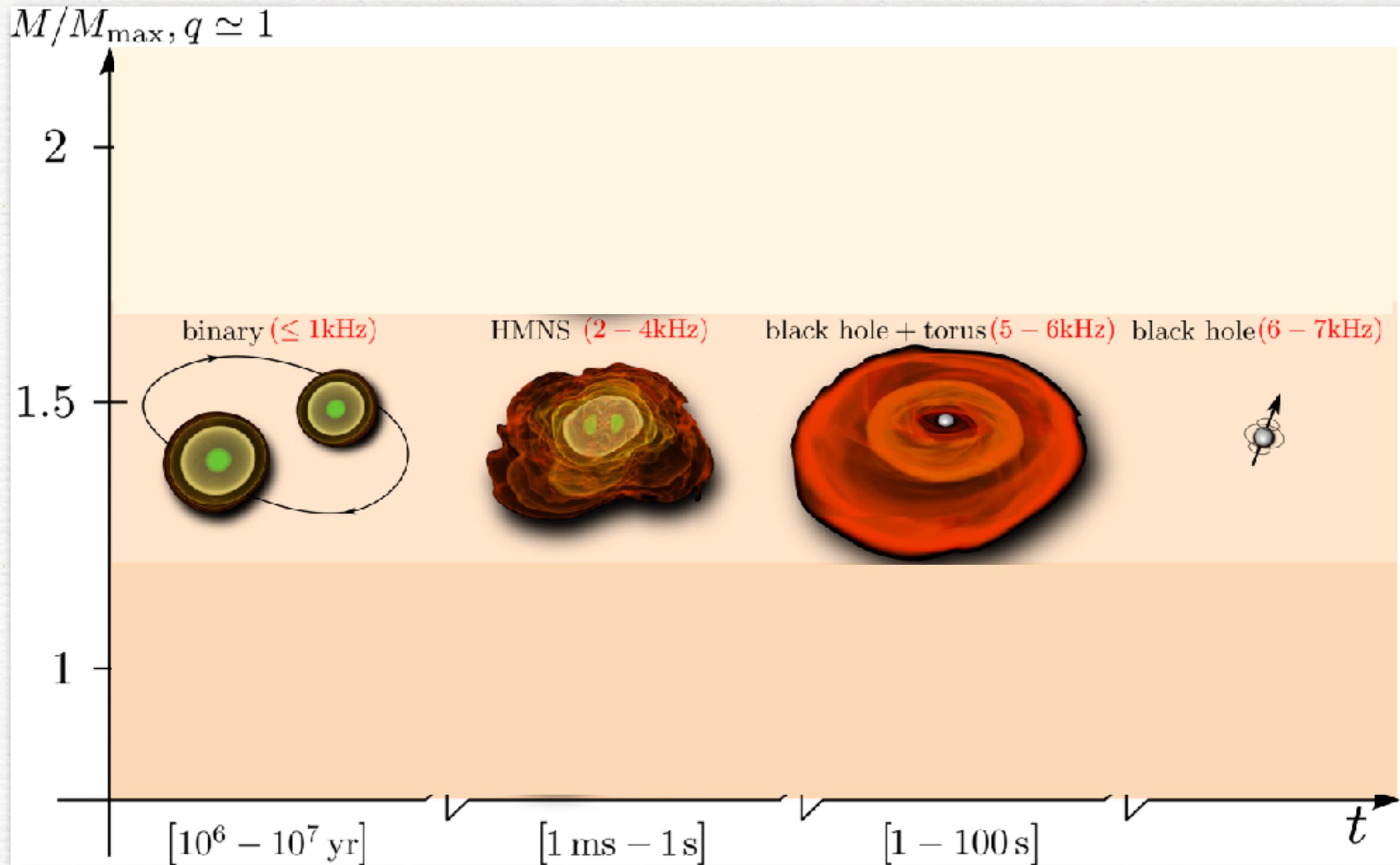


merger → HMNS → BH + torus

Quantitative differences are produced by:

- **total mass** (prompt vs delayed collapse)

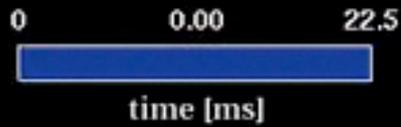
Broadbrush picture



merger → HMNS → BH + torus

Quantitative differences are produced by:

- total mass (prompt vs delayed collapse)
- mass asymmetries (HMNS and torus)



Total mass : $3.37 M_{\odot}$; mass ratio :0.80;



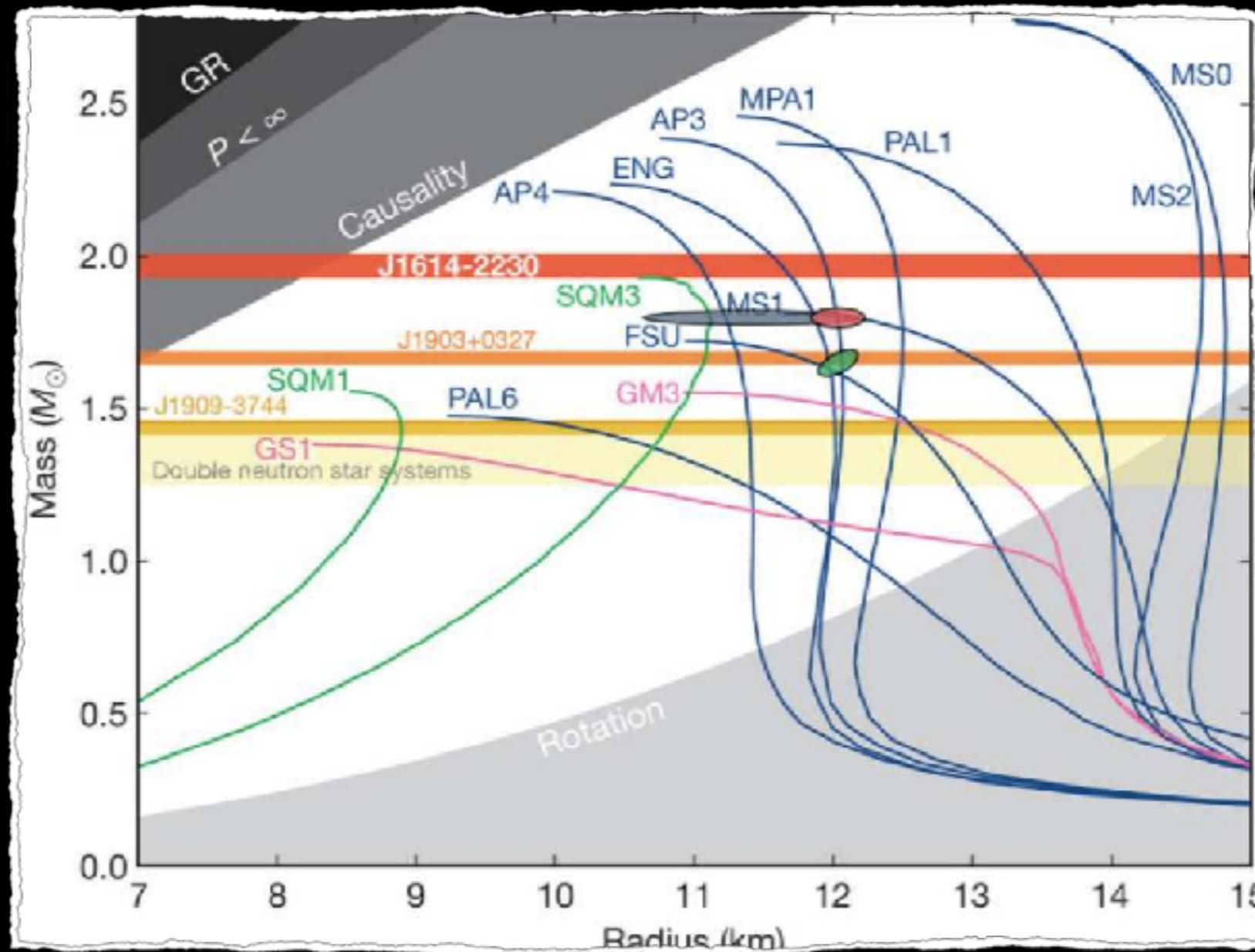
- * the torii are generically **more massive**
- * the torii are generically **more extended**
- * the torii tend to stable **quasi-Keplerian** configurations
- * overall unequal-mass systems have all the ingredients needed to create a GRB

merger → HMNS → BH + torus

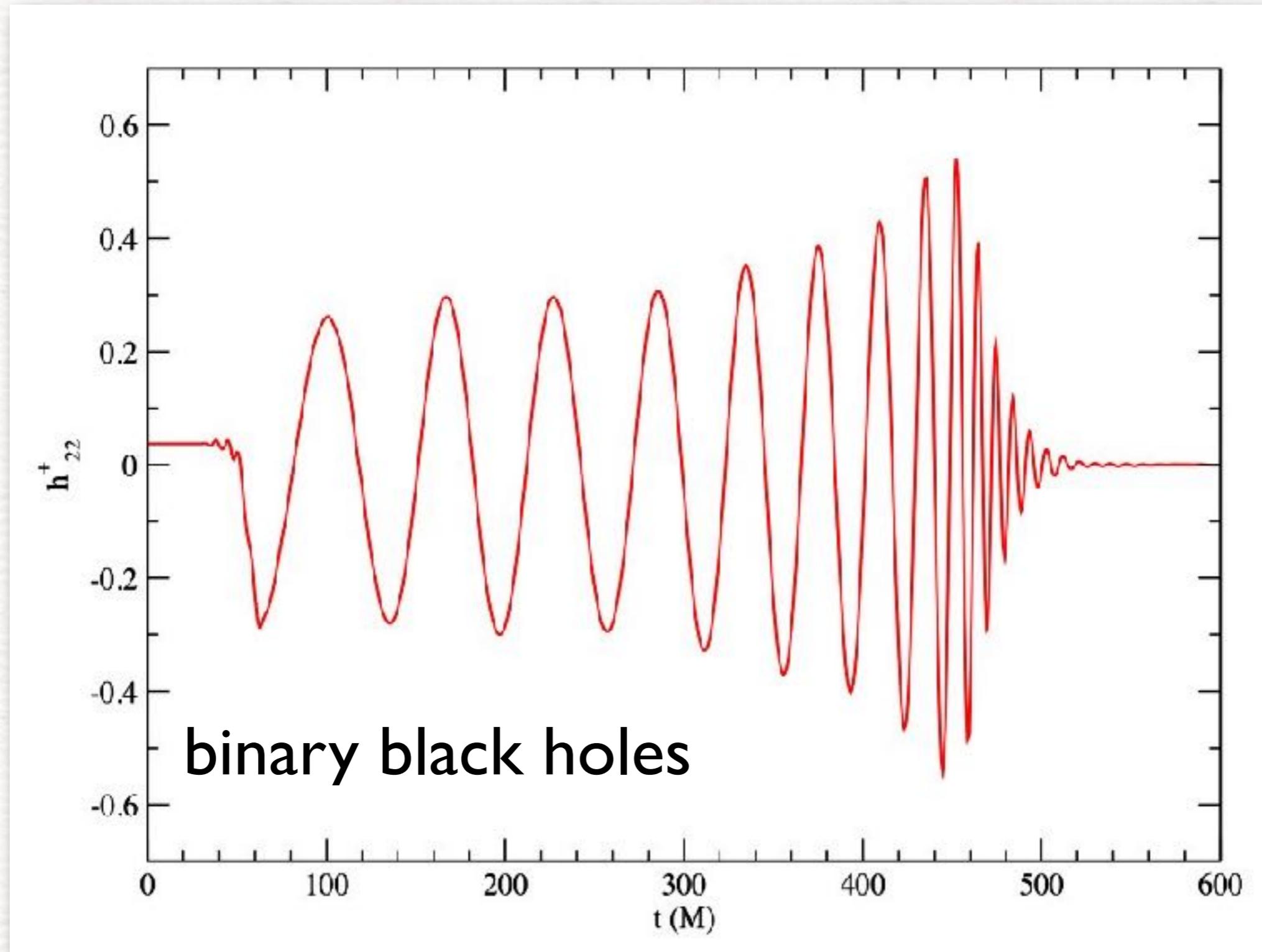
Quantitative differences are produced by:

- total **mass** (prompt vs delayed collapse)
- mass **asymmetries** (HMNS and torus)
- soft/stiff **EOS** (inspiral and post-merger)
- **magnetic fields** (equil. and EM emission)
- **radiative** losses (equil. and nucleosynthesis)

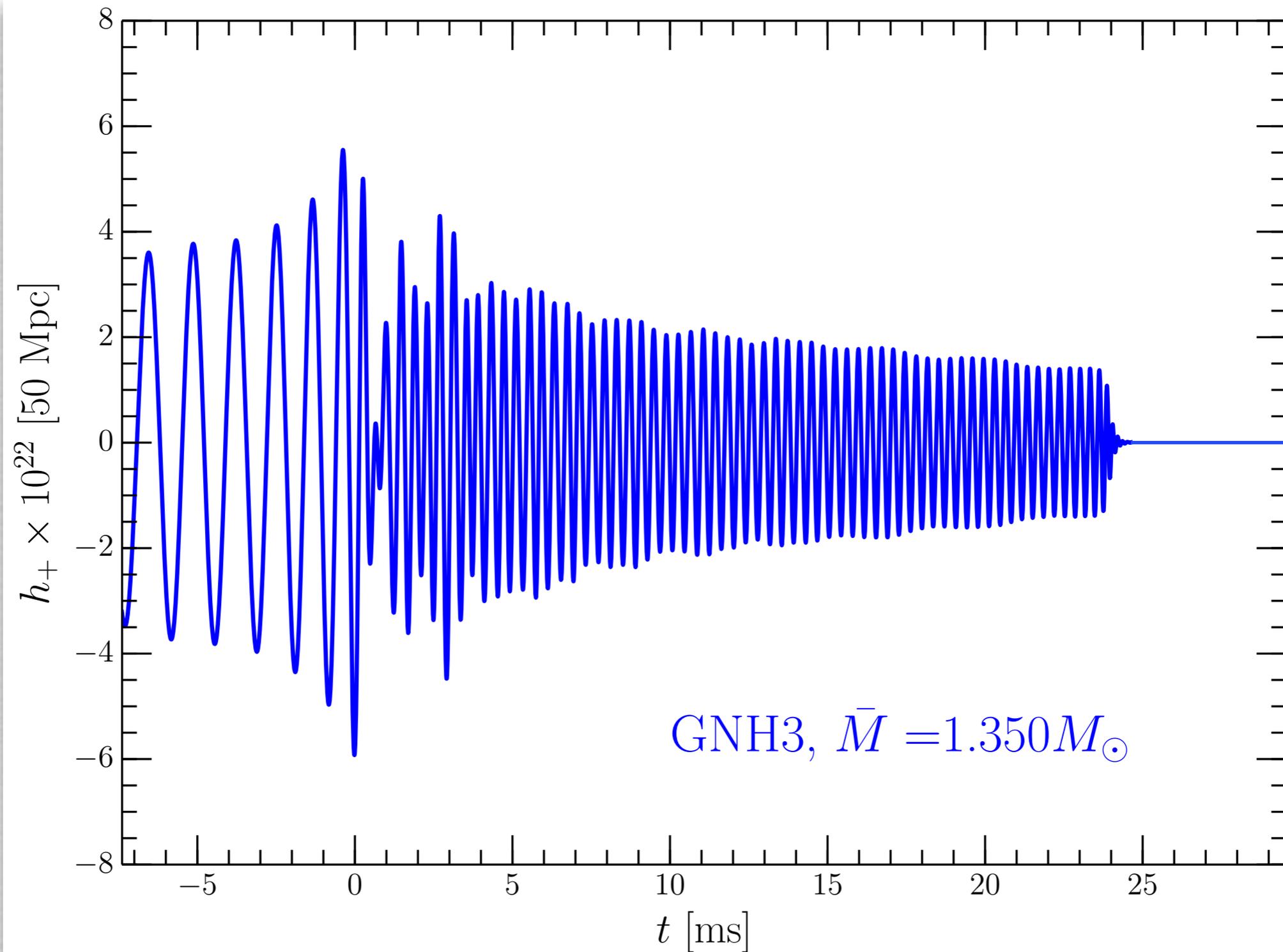
How to constrain the EOS



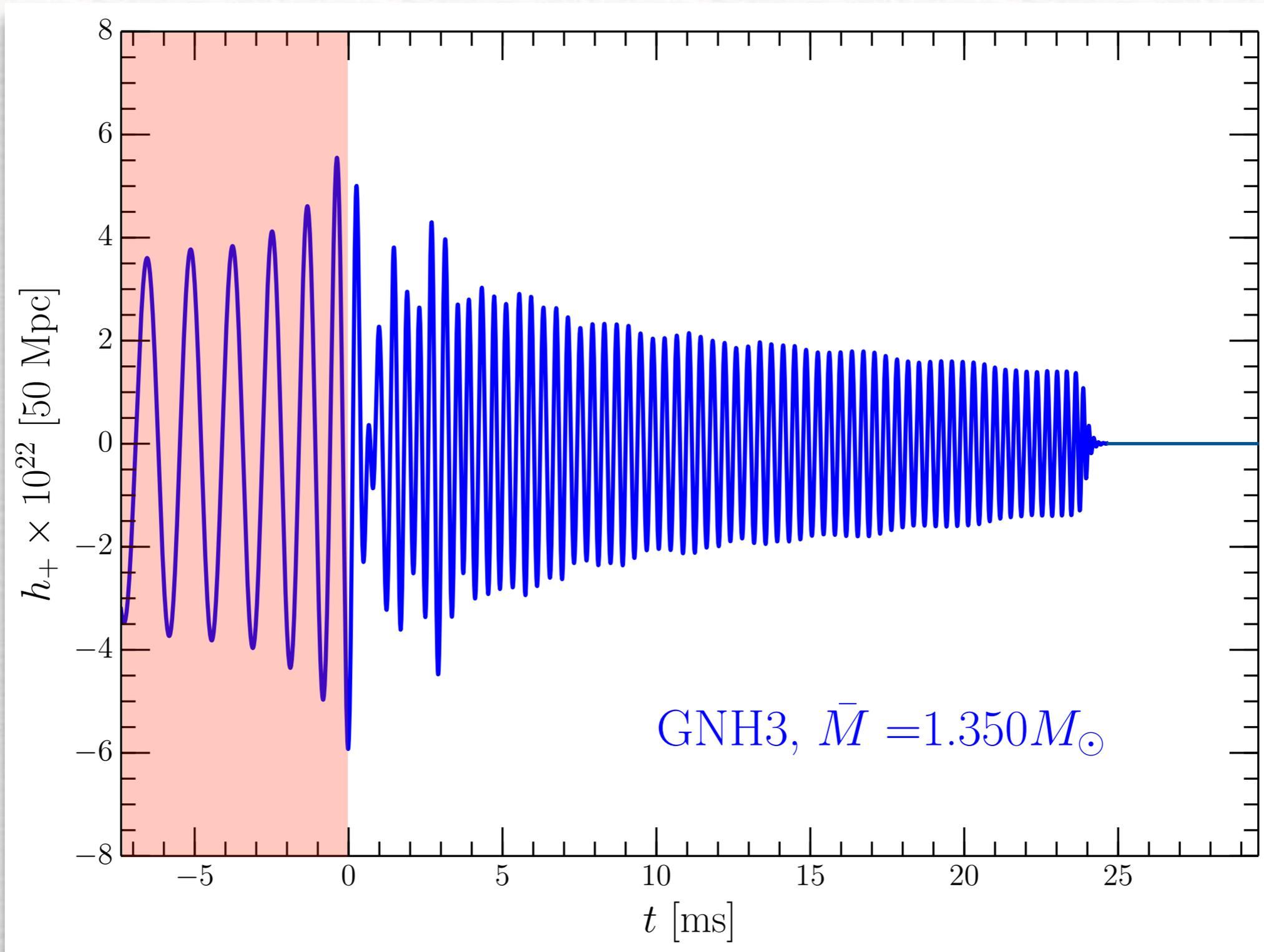
Anatomy of the GW signal



Anatomy of the GW signal

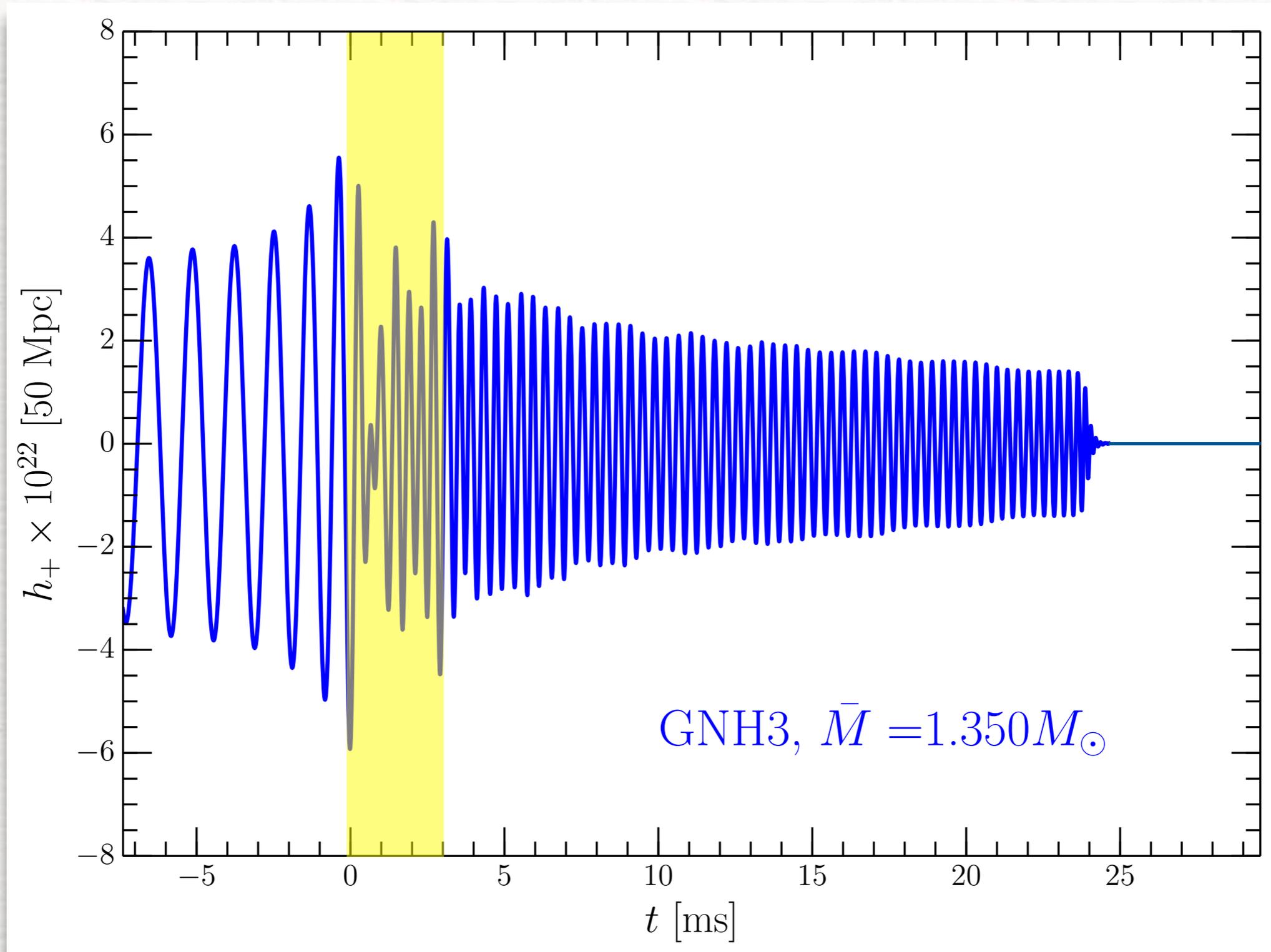


Anatomy of the GW signal



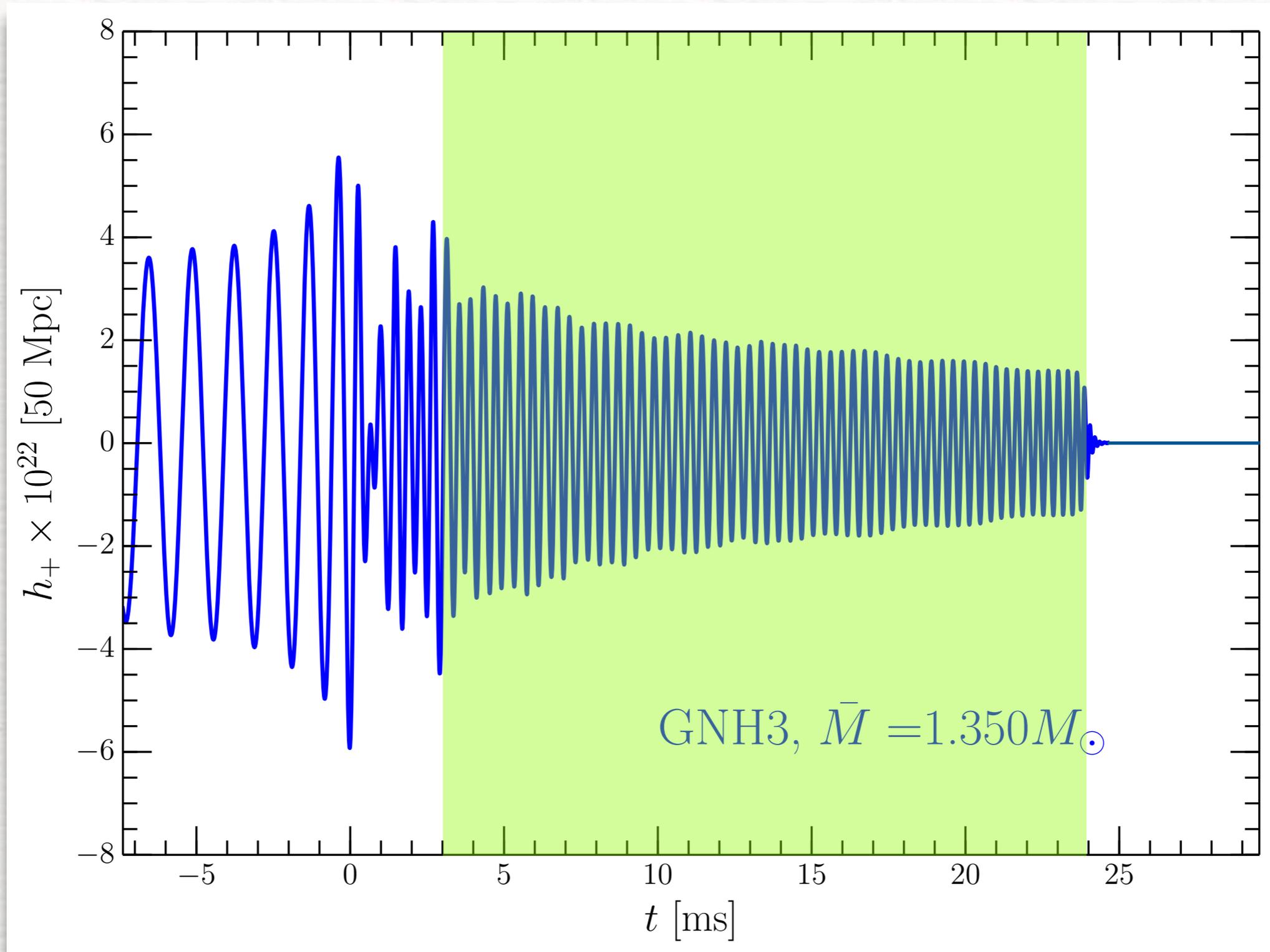
Inspiral: well approximated by PN/EOB; tidal effects important

Anatomy of the GW signal



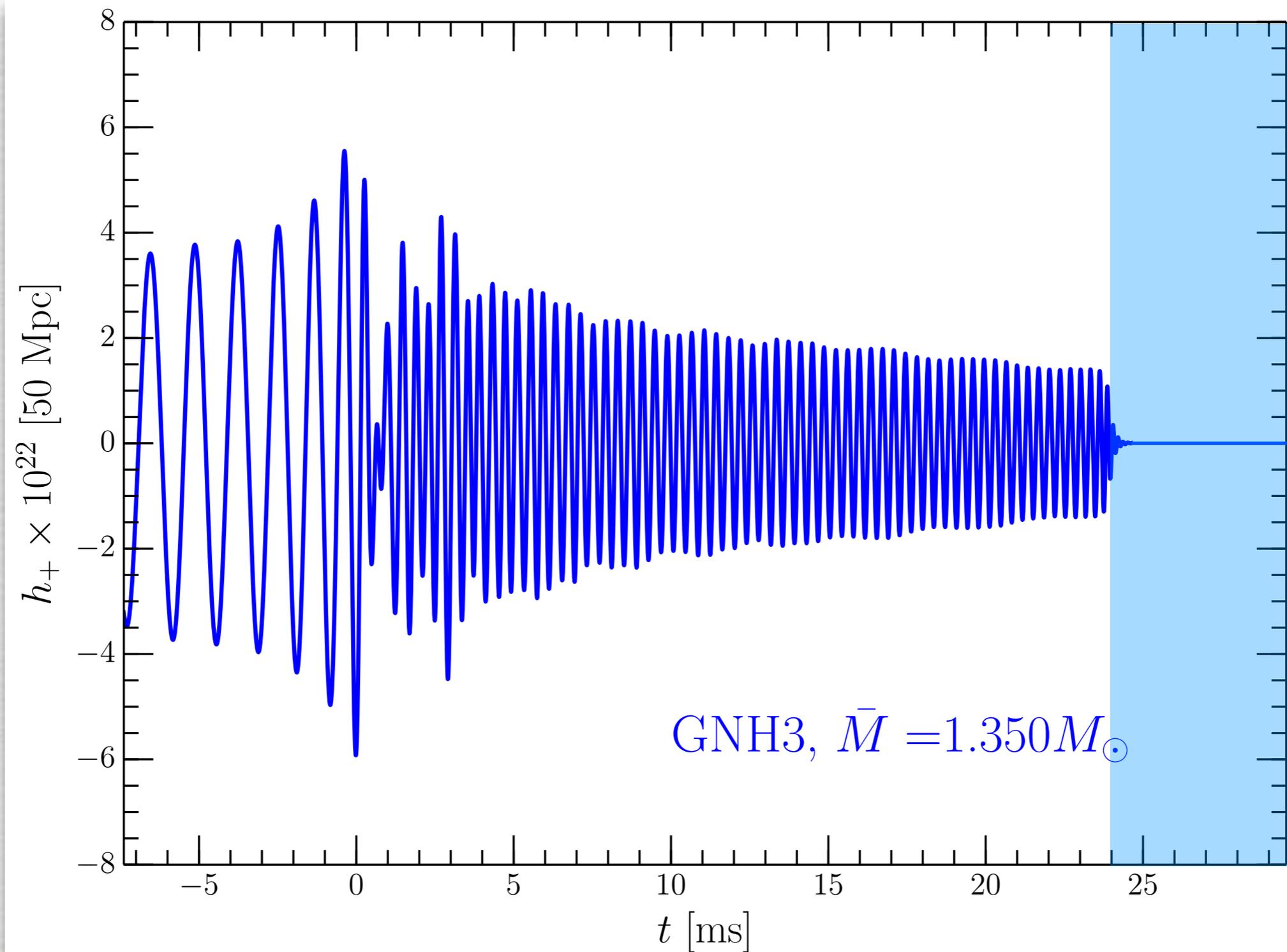
Merger: highly nonlinear but analytic description possible

Anatomy of the GW signal



post-merger: quasi-periodic emission of bar-deformed HMNS

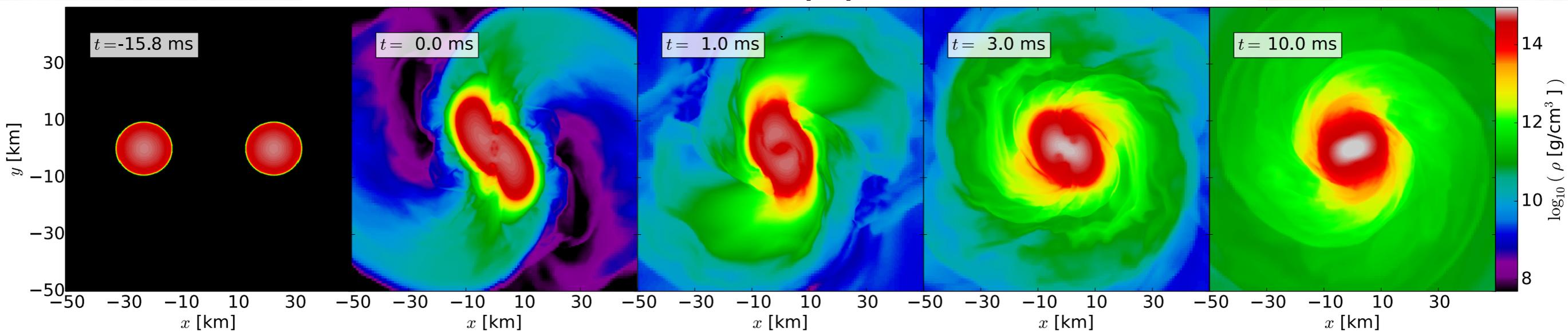
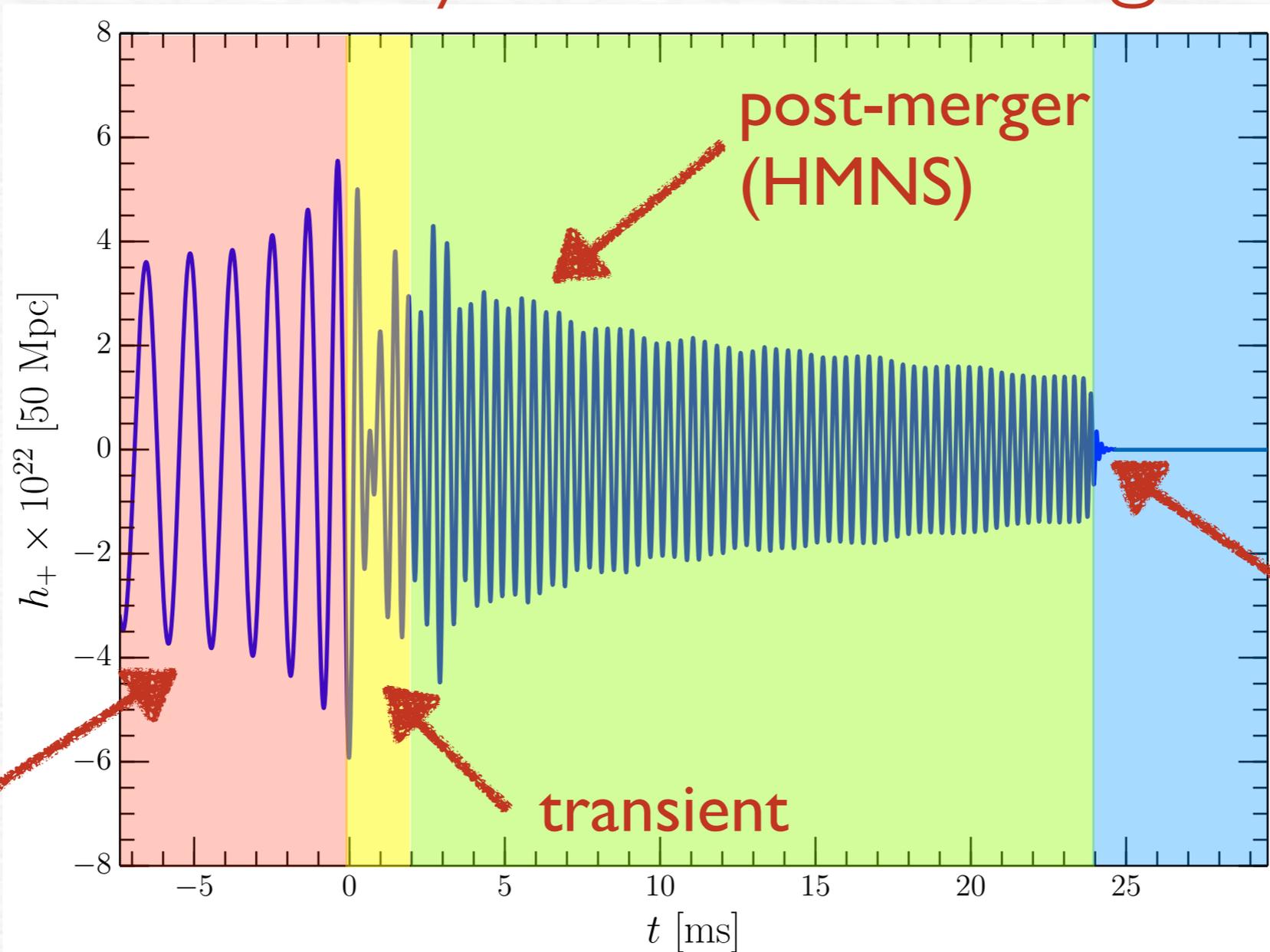
Anatomy of the GW signal



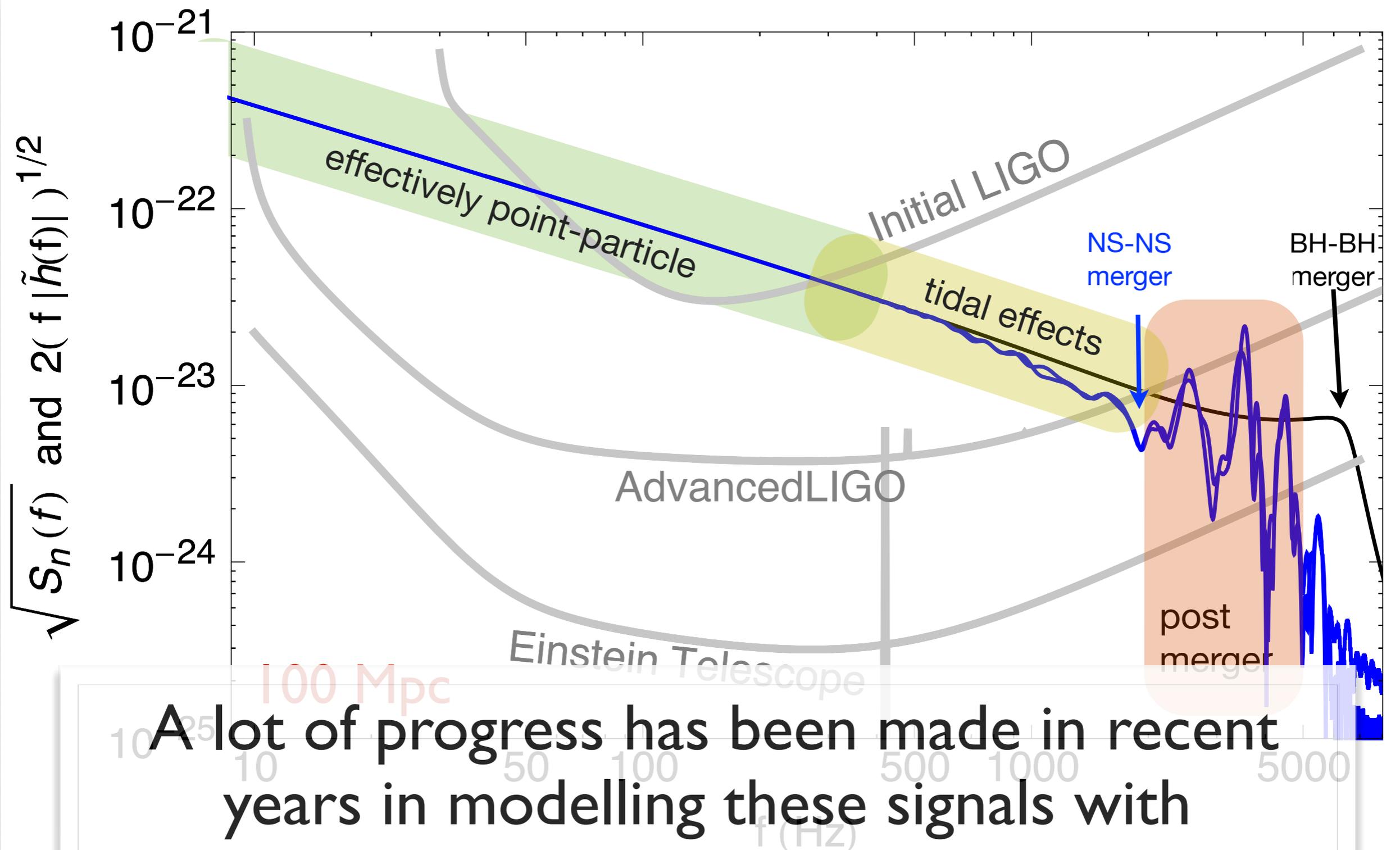
Collapse-ringdown: signal essentially shuts off.

Anatomy of the GW signal

Chirp signal



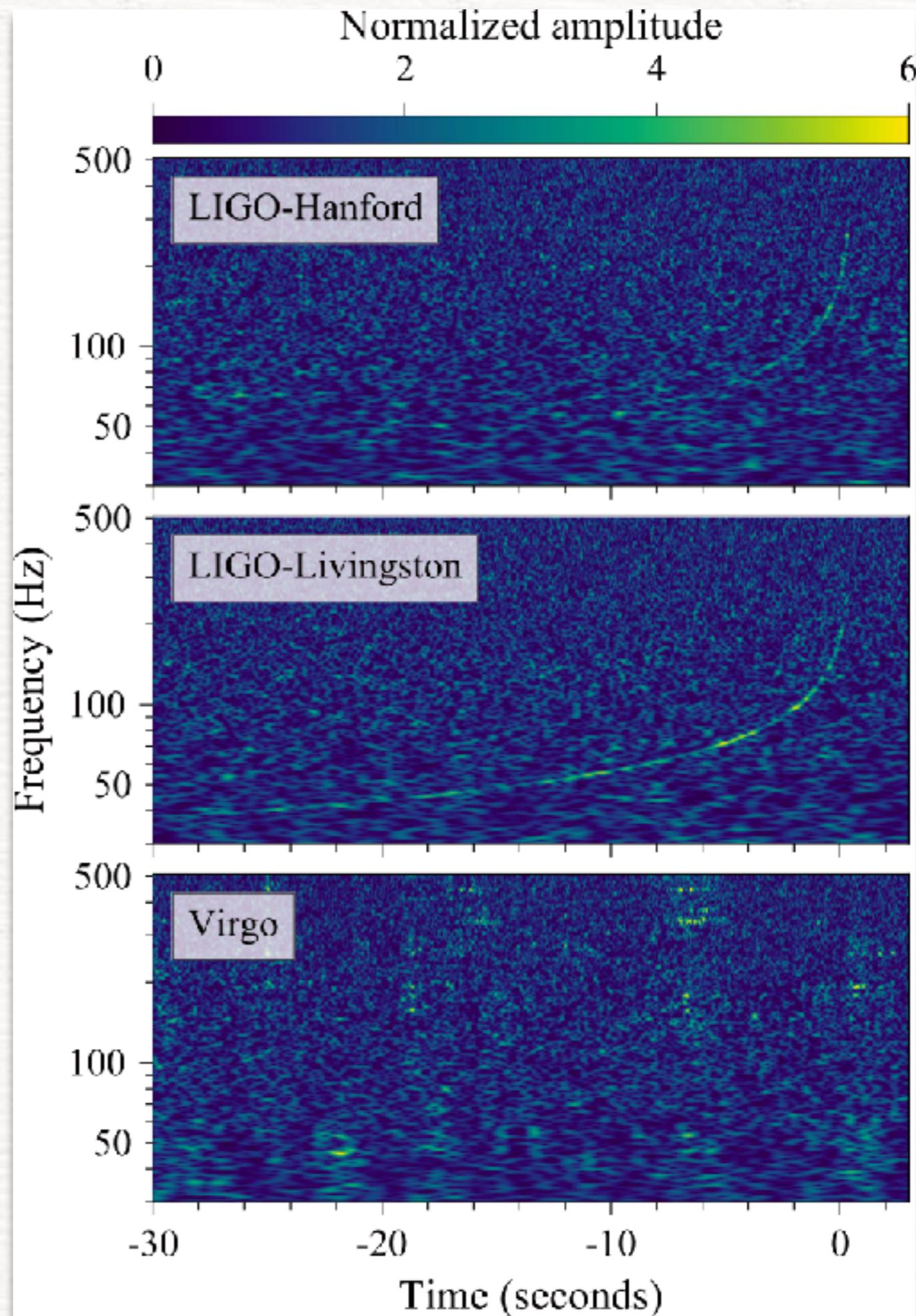
In frequency space



A lot of progress has been made in recent years in modelling these signals with numerical simulations of different binaries

GW170817

- On 16 October 2017 the LSC/Virgo collaboration announced detection of the gravitational signal from merging binary neutron-star system.



GW170817

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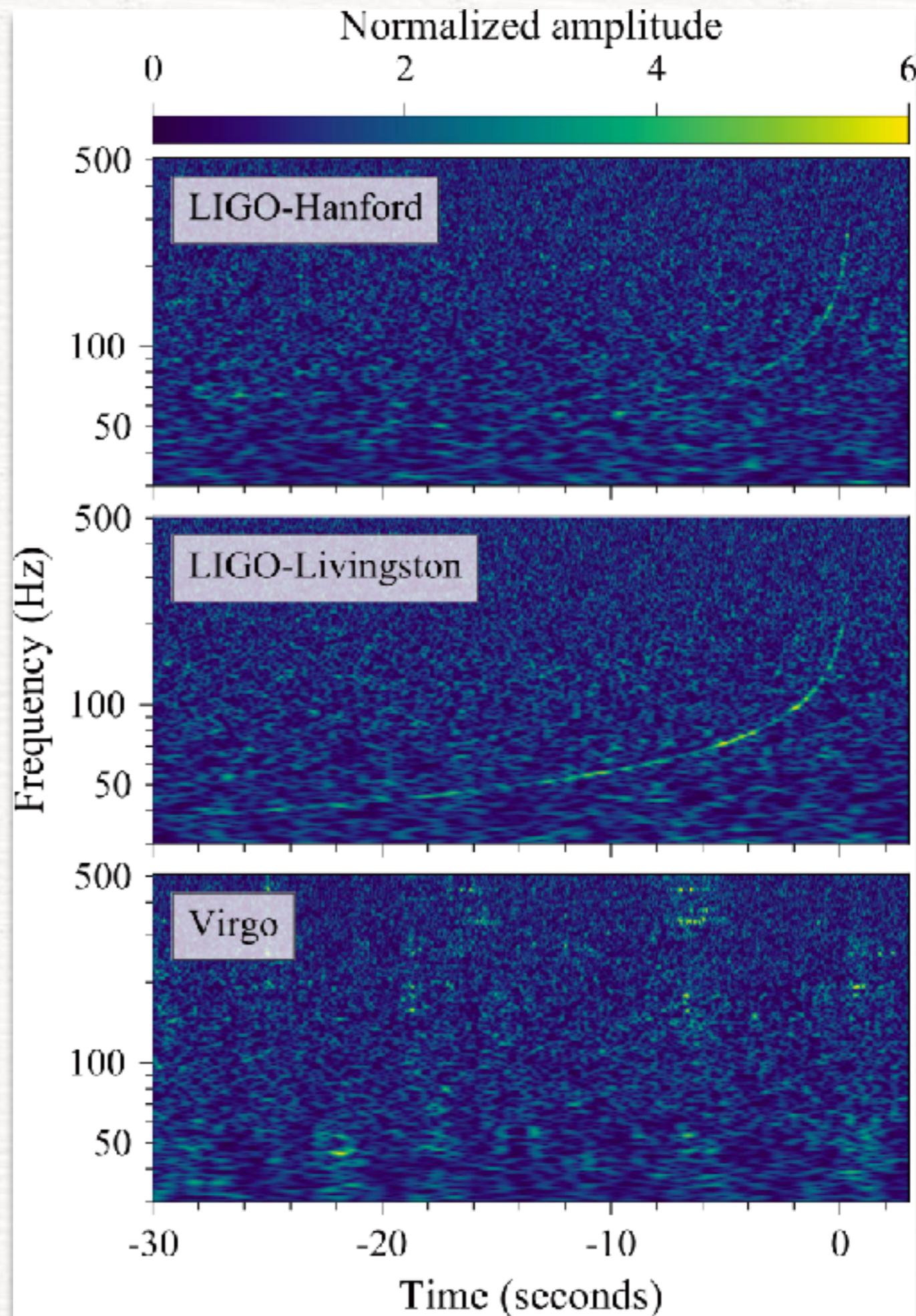
- Total mass:

$$M_1 + M_2 = 2.74_{-0.01}^{+0.04} M_{\odot}$$

- Individual masses:

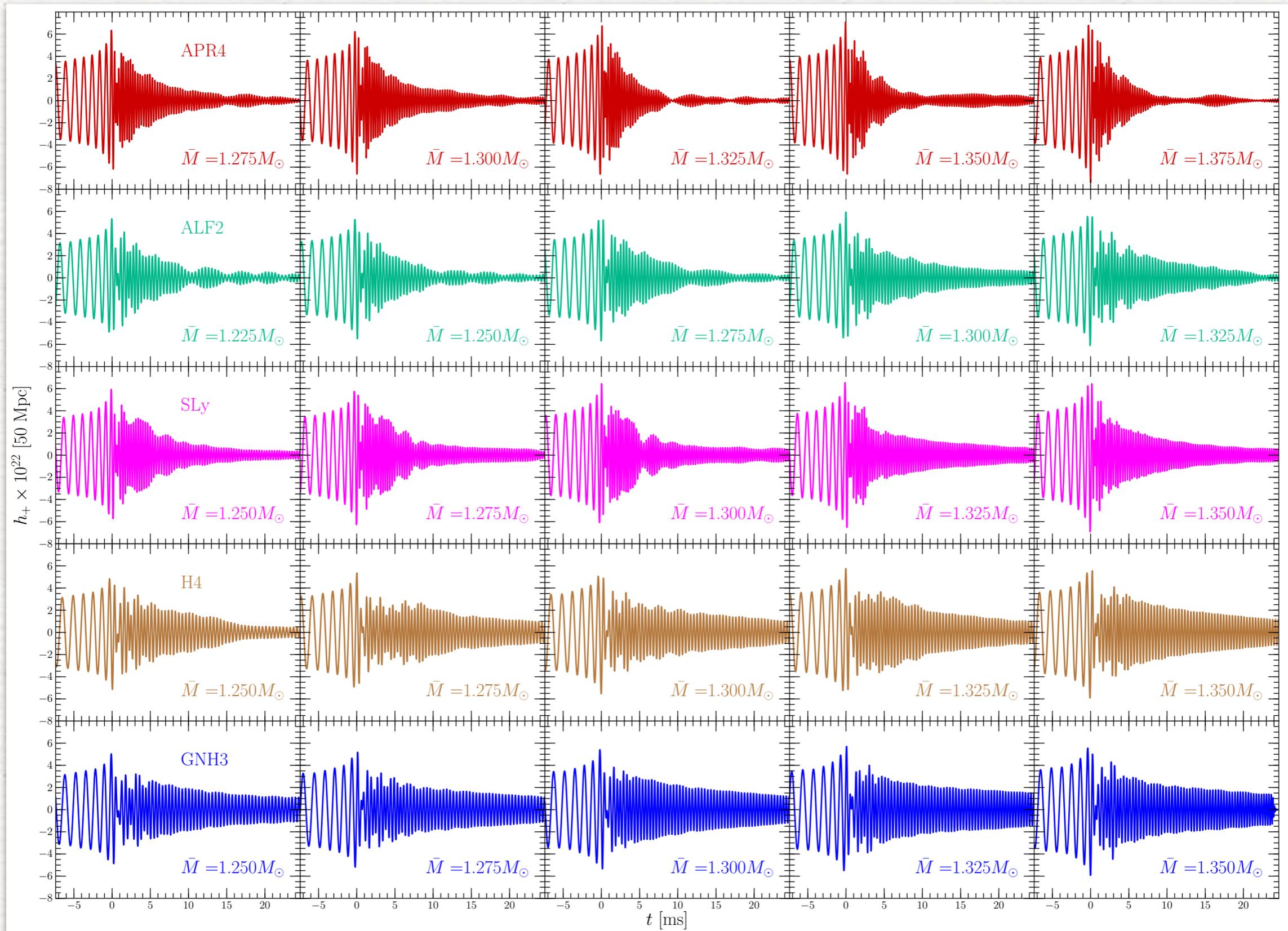
$$M_1 = 1.36 - 1.60 M_{\odot}$$

$$M_2 = 1.17 - 1.36 M_{\odot}$$



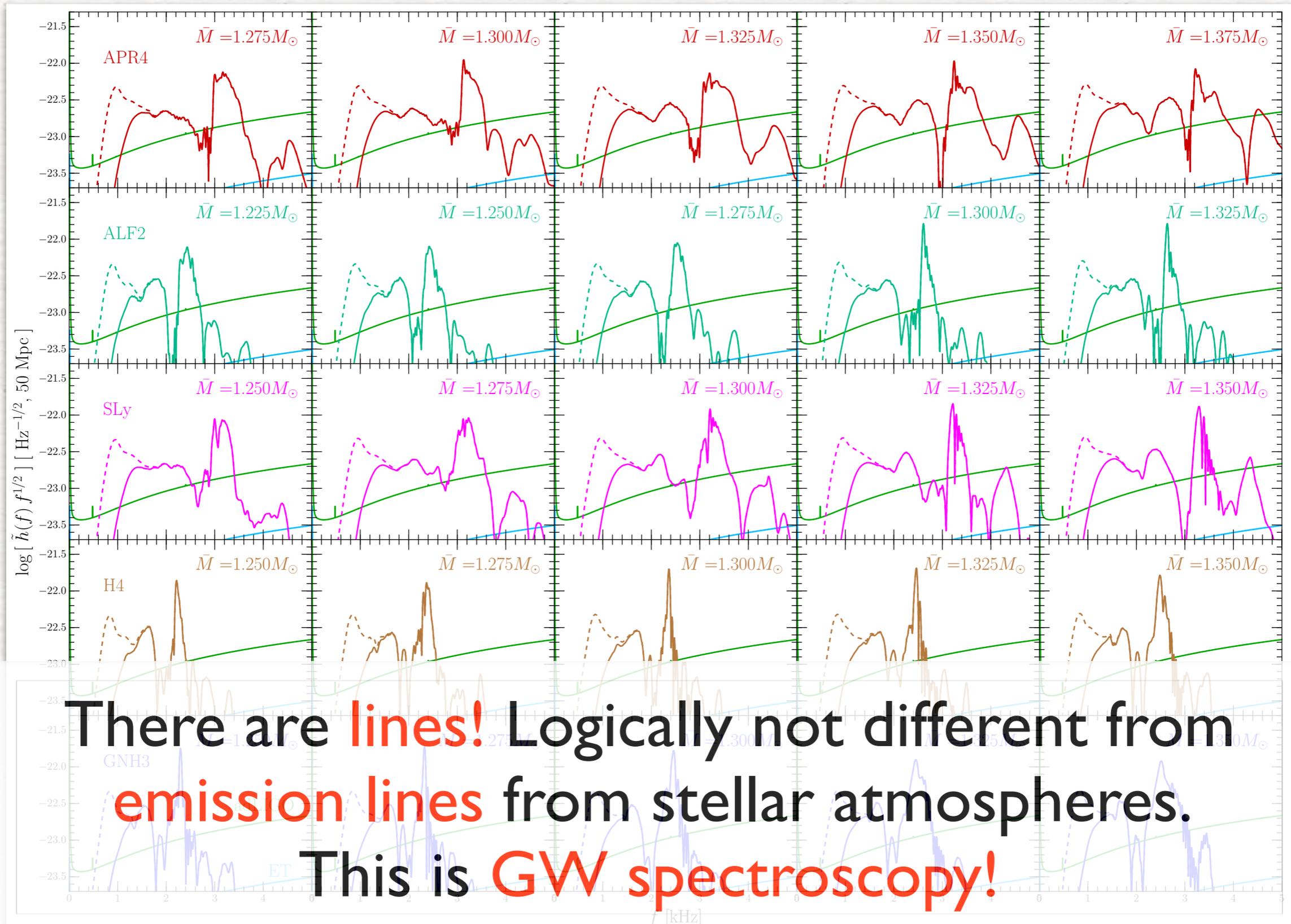
What we can do nowadays

Takami, LR, Baiotti (2014, 2015), LR+ (2016)



Extracting information from the EOS

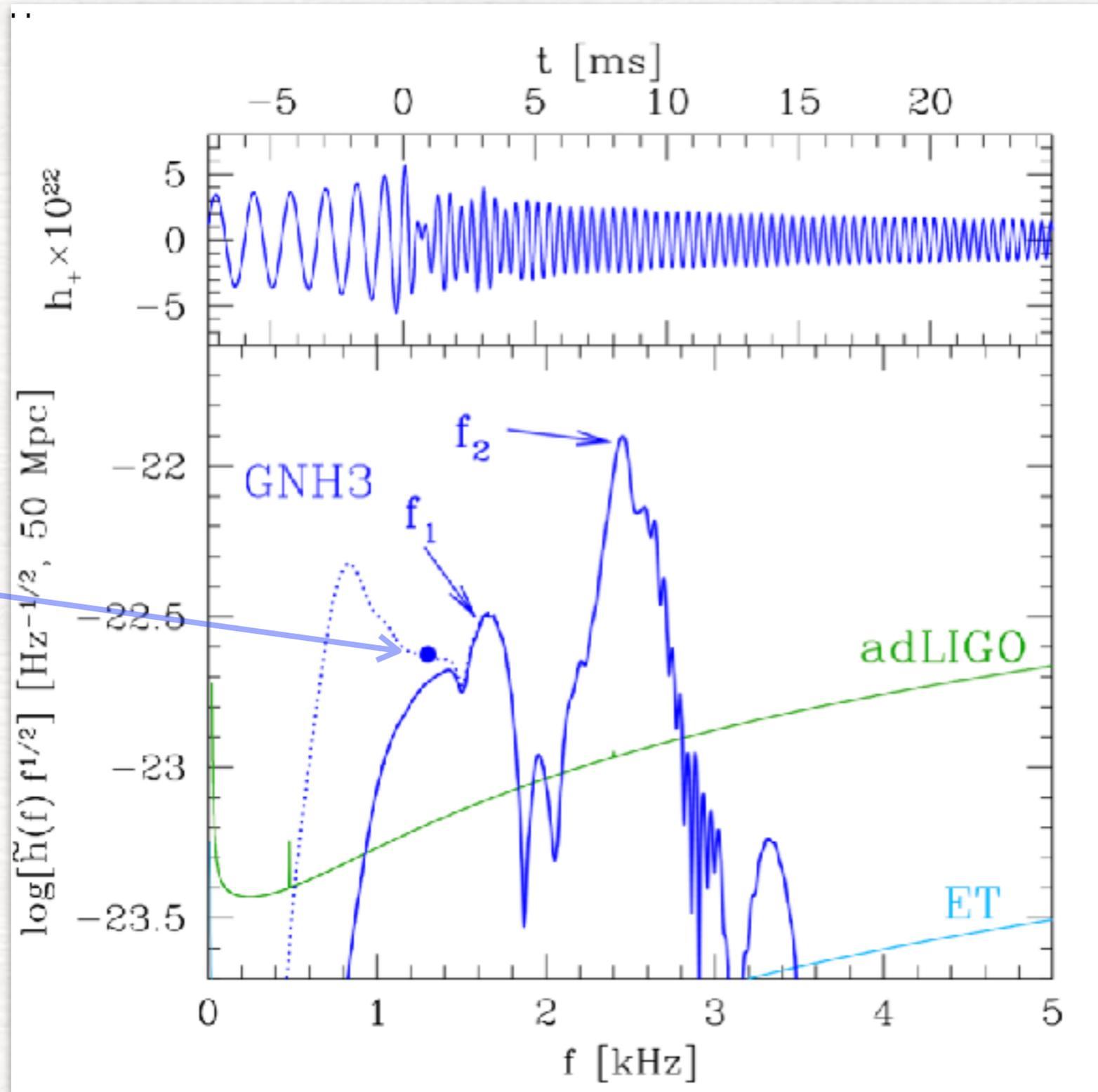
Takami, LR, Baiotti (2014, 2015), LR+ (2016)



A new approach to constrain the EOS

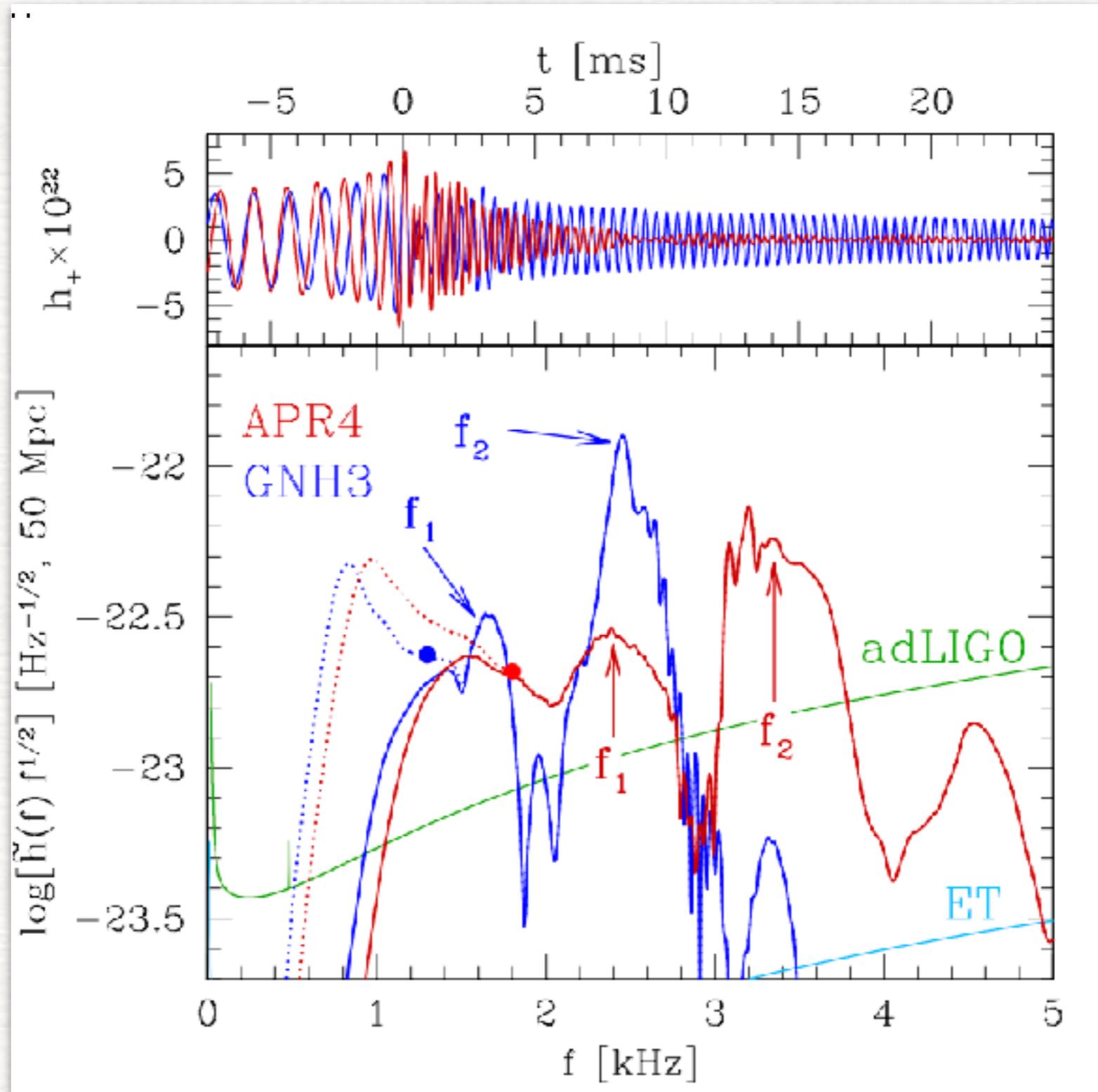
Oechslin+2007, Baiotti+2008, Bauswein+ 2011, 2012, Stergioulas+ 2011, Hotokezaka+ 2013, Takami 2014, 2015, Bernuzzi 2014, 2015, Bauswein+ 2015, Clark+ 2016, LR+2016, de Pietri+ 2016, Feo+ 2017, Bose+ 2017 ...

merger
frequency



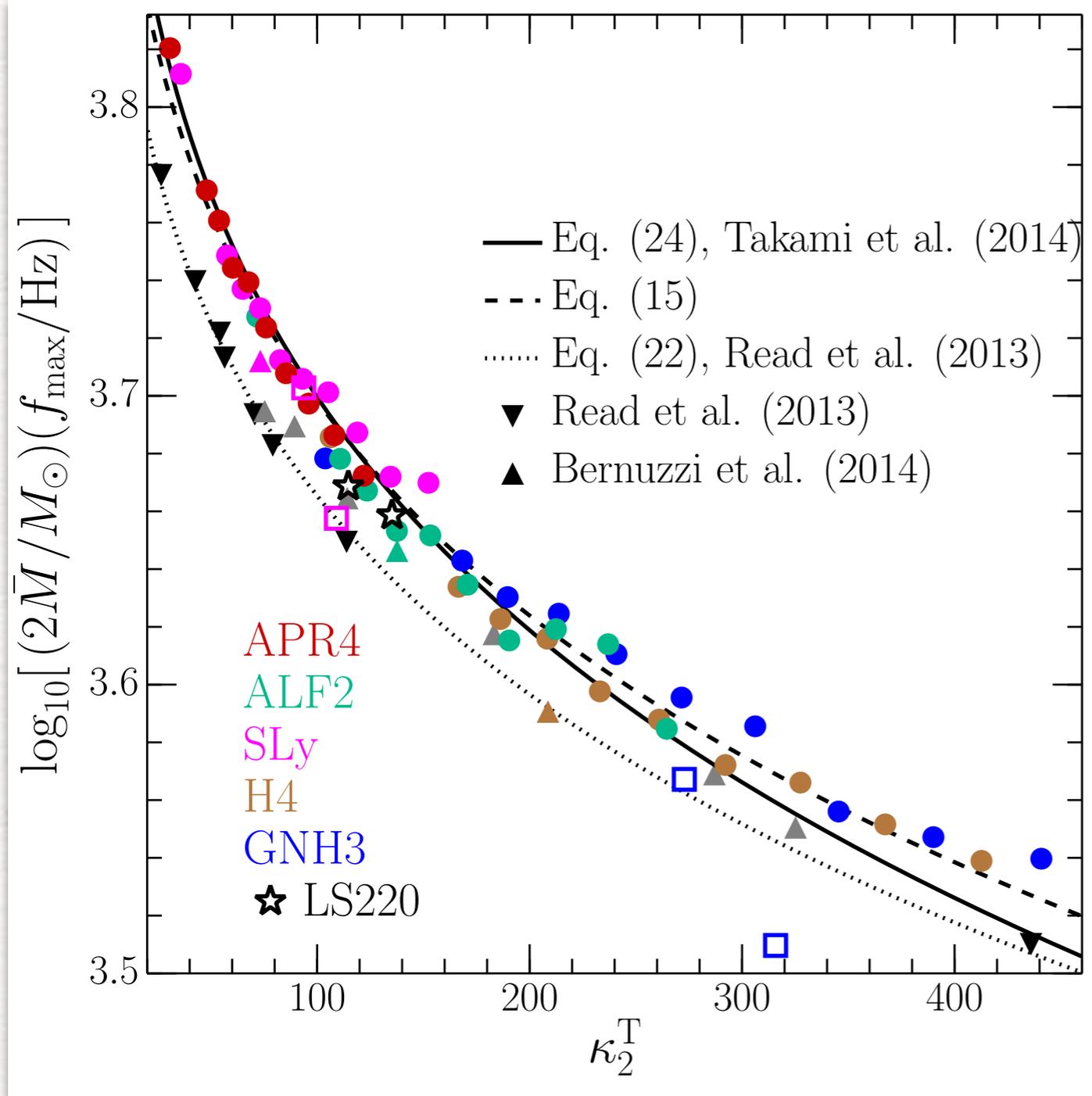
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Quasi-universal behaviour

Quasi-universal behaviour: **inspiral**



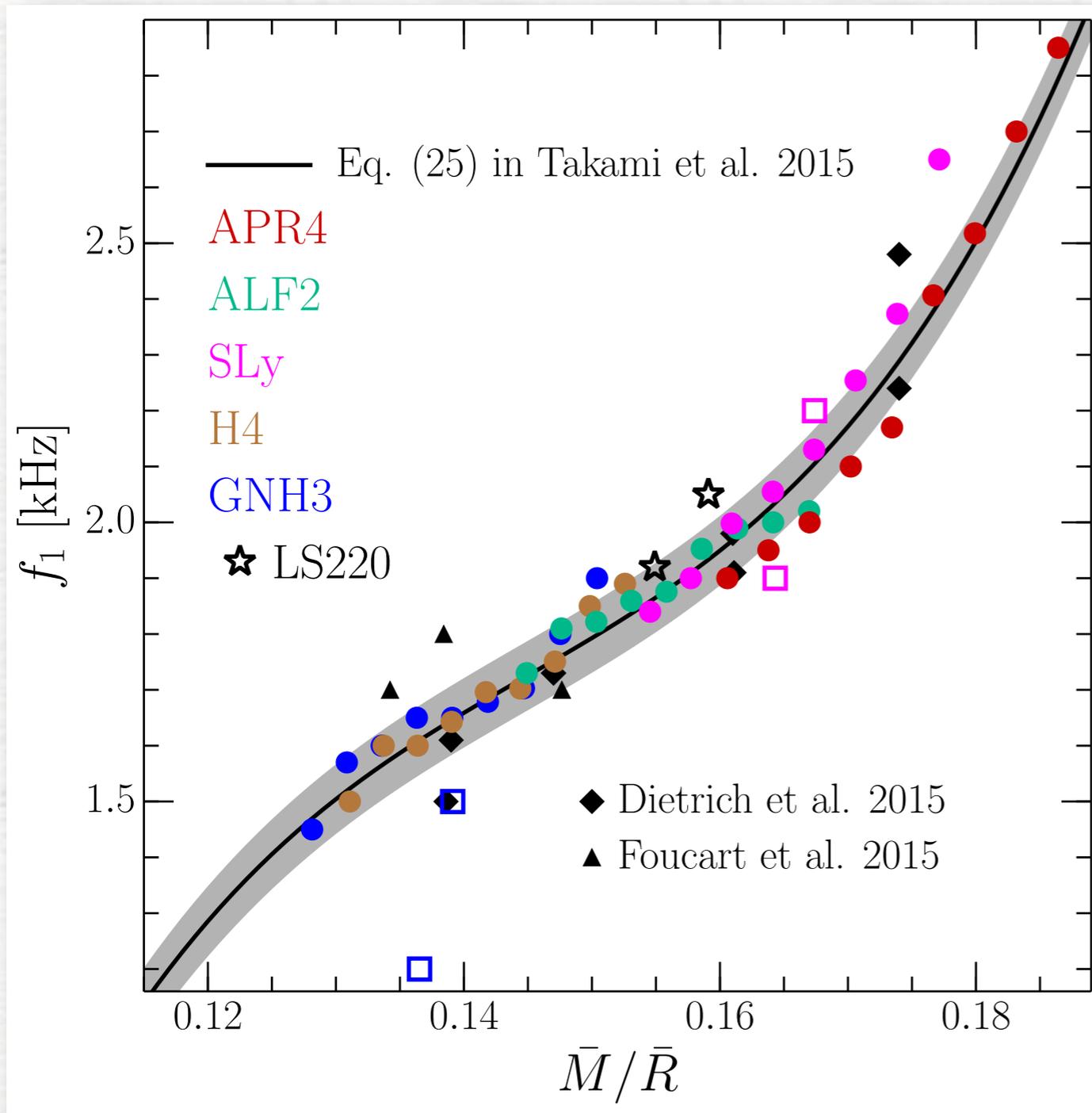
“surprising” result: **quasi-universal** behaviour of GW frequency at amplitude peak (Read+2013)

Many other simulations have confirmed this (Bernuzzi+ 2014, Takami+ 2015, LR+2016).

Quasi-universal behaviour in the **inspiral** implies that once **f_{\max}** is measured, so is tidal deformability, hence $I, Q, M/R$

$$\Lambda = \frac{\lambda}{\bar{M}^5} = \frac{16}{3} \kappa_2^T \quad \text{tidal deformability or Love number}$$

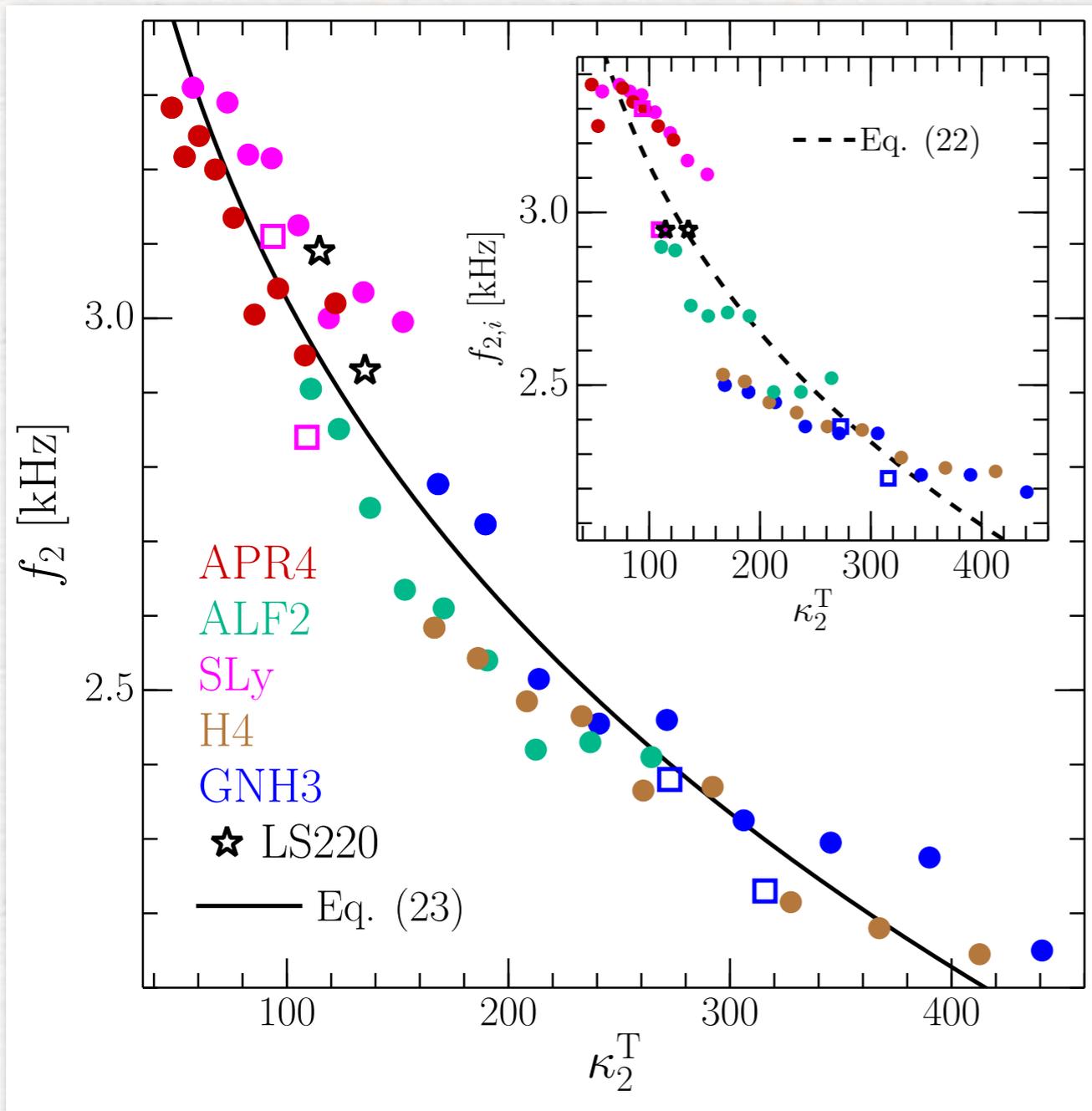
Quasi-universal behaviour: post-merger



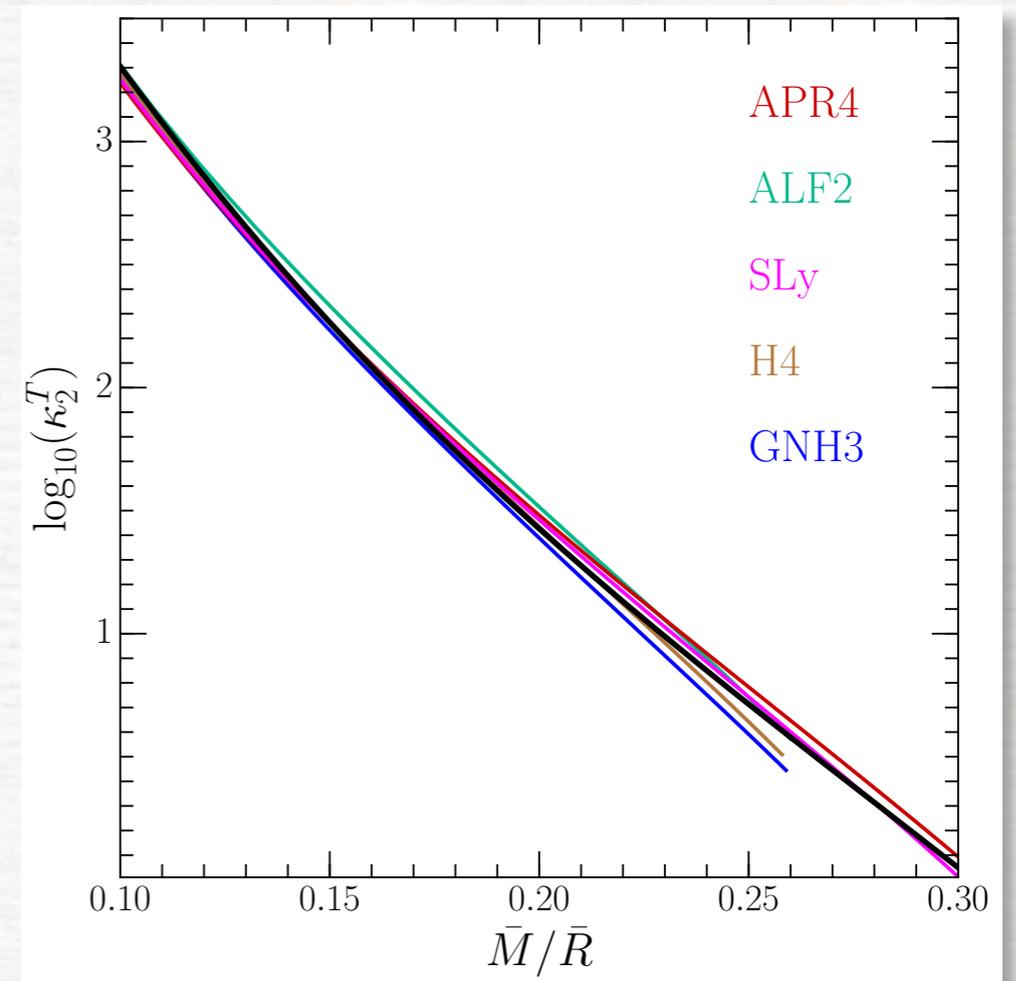
We have found **quasi-universal behaviour**: i.e., the properties of the spectra are only weakly dependent on the EOS.

This has profound implications for the analytical modelling of the GW emission: “what we do for one EOS can be extended to all EOSs.”

Quasi-universal behaviour: post-merger



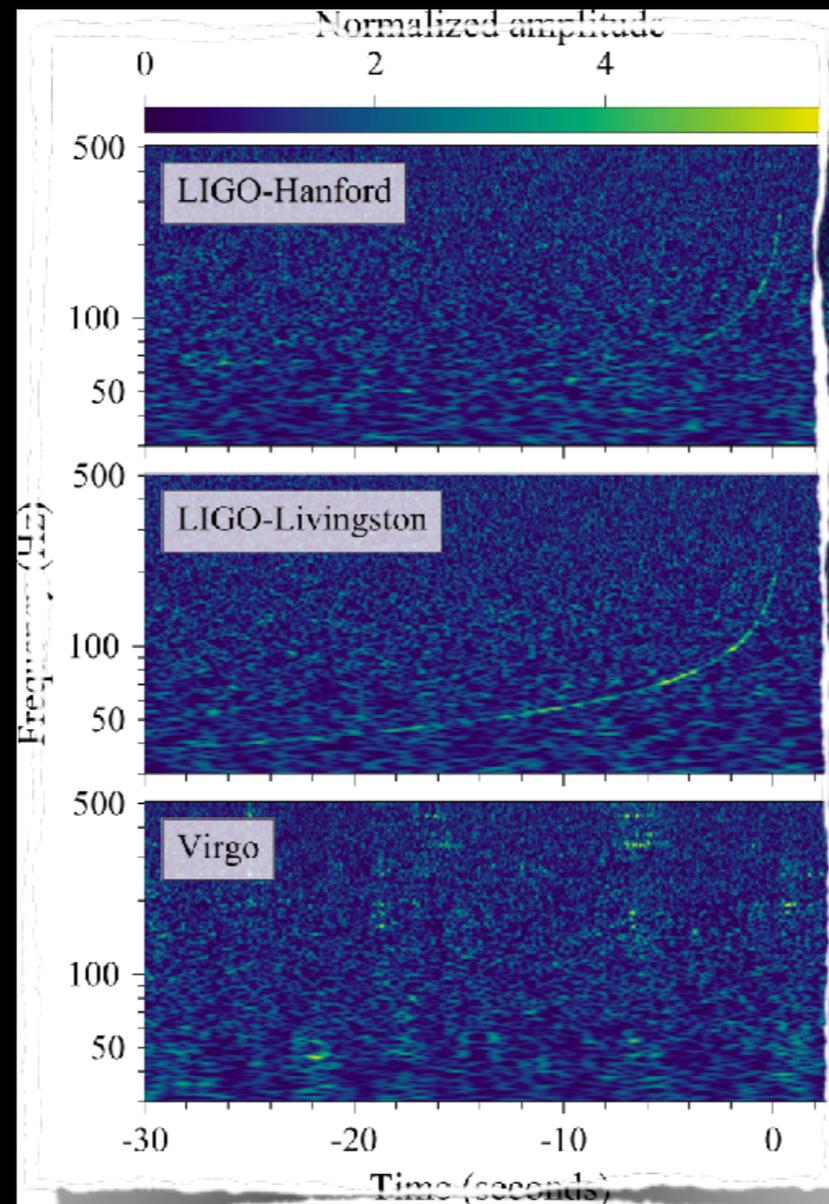
- Correlations with Love number found also for high frequency peak f_2 .
- This and other correlations are **weaker** but equally useful.



- Important correlation also between **compactness** and **deformability**

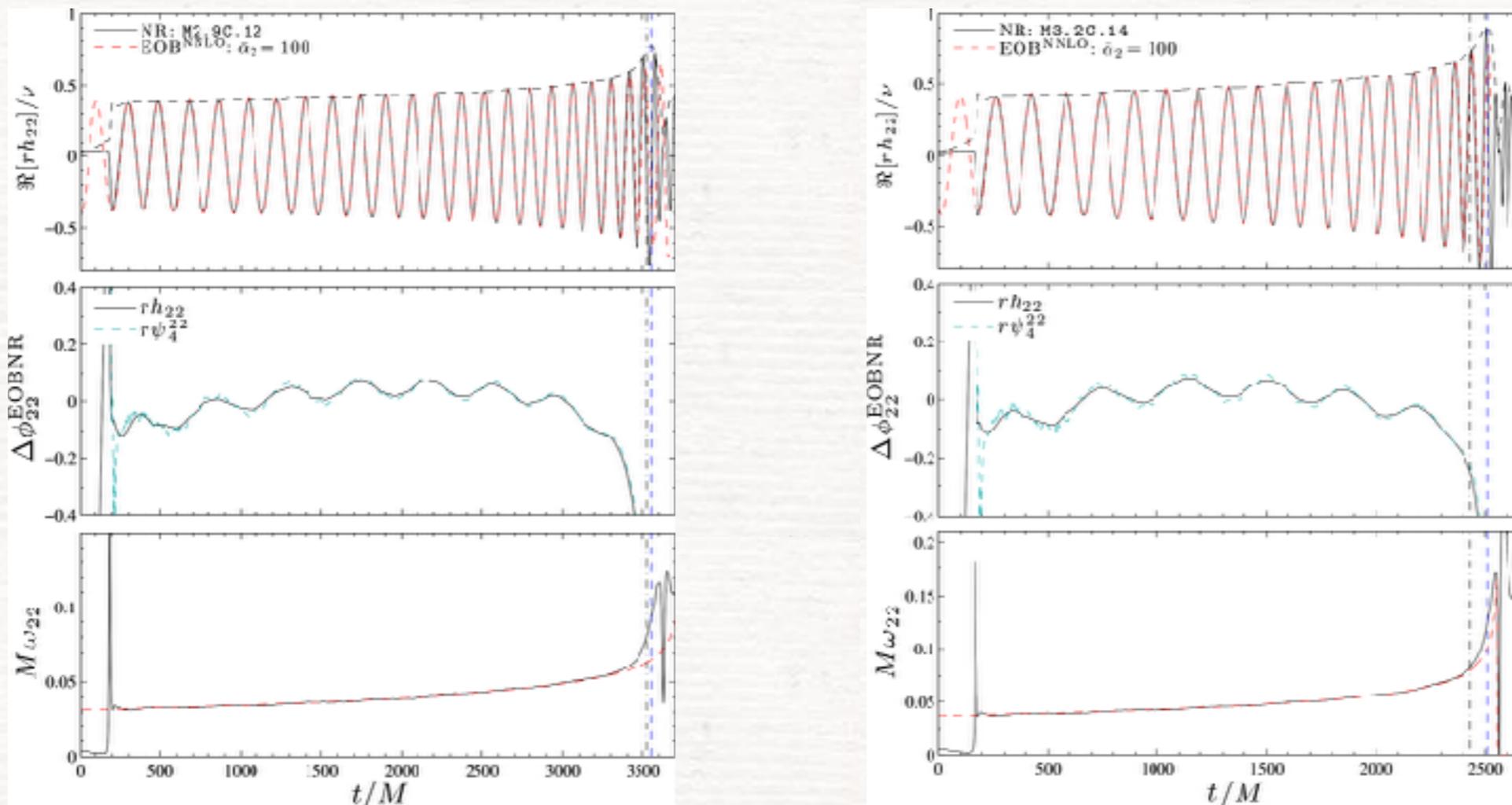
Radius estimate from binary population

Bose, Chakravarti, LR, Sathyaprakash, Takami (2018)



Analytical modelling of waveform

- **Analytical** modelling of signal is essential for statistical analysis
- **Inspiral** well reproduced with PN or EOB (Baiotti+11, Bernuzzi+15a, 15b, Lackey+16, Hotokezaka+15, Hinderer+16)



- **Inspiral** also benefits from very accurate codes, **WhiskyTHC**: highest demonstrated convergence order (Radice+13, 13b)

Spectroscopy is useful but hard

- Measuring post-merger not easy: H4, $M = 2 \times 1.35 M_{\odot}$

$$|2\tilde{h}(f) f^{1/2}| \simeq 10^{-22} / \sqrt{\text{Hz}} \text{ at } f = f_2 \simeq 2470 \text{ Hz}$$

$$\text{SNR} \simeq |2\tilde{h}(f) f^{1/2}| [\delta f / (f S_h(f))]^{1/2} \simeq 1.8$$

- However, multiple detections will provide excess power and SNRs can add up statistically providing information.
- Idea not new and implemented in the past (Del Pozzo+ 2013, Agathos+ 2015, Clark+ 2015, 2016); there's a new twist here.

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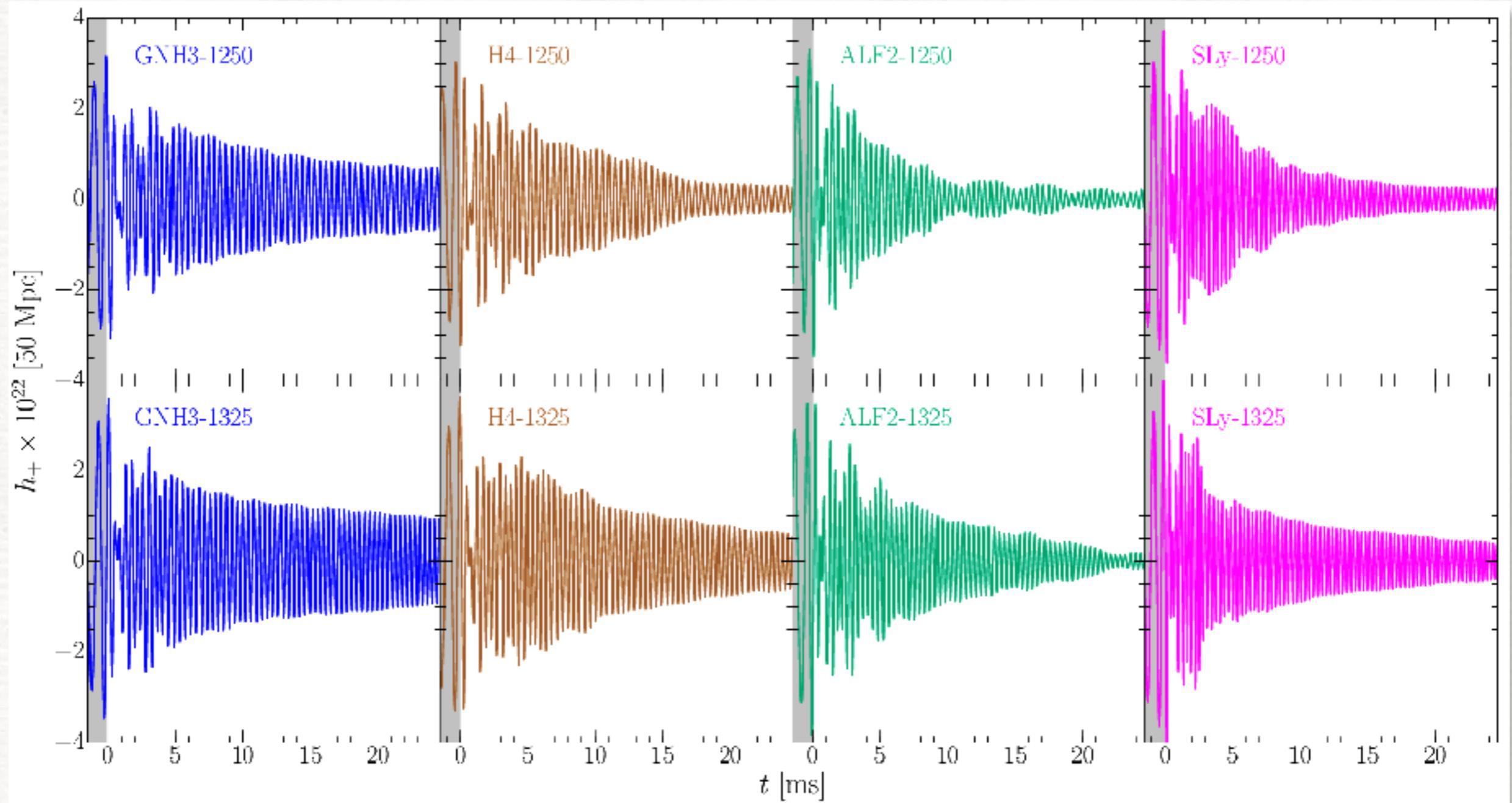
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- However, multiple detections will provide excess power and SNRs can add up statistically providing information.
- Idea not new and implemented in the past (Del Pozzo+ 2013, Agathos+ 2015, Clark+ 2015, 2016); there's a new twist here.
- **Analytical modelling** of post-merger GW via main frequencies $(f_{\text{max}}, f_1, f_2)$ and **universal relations** allow us to relate **inspiral** and **postmerger** to progenitor stars.

Analytical modelling of postmerger waveform

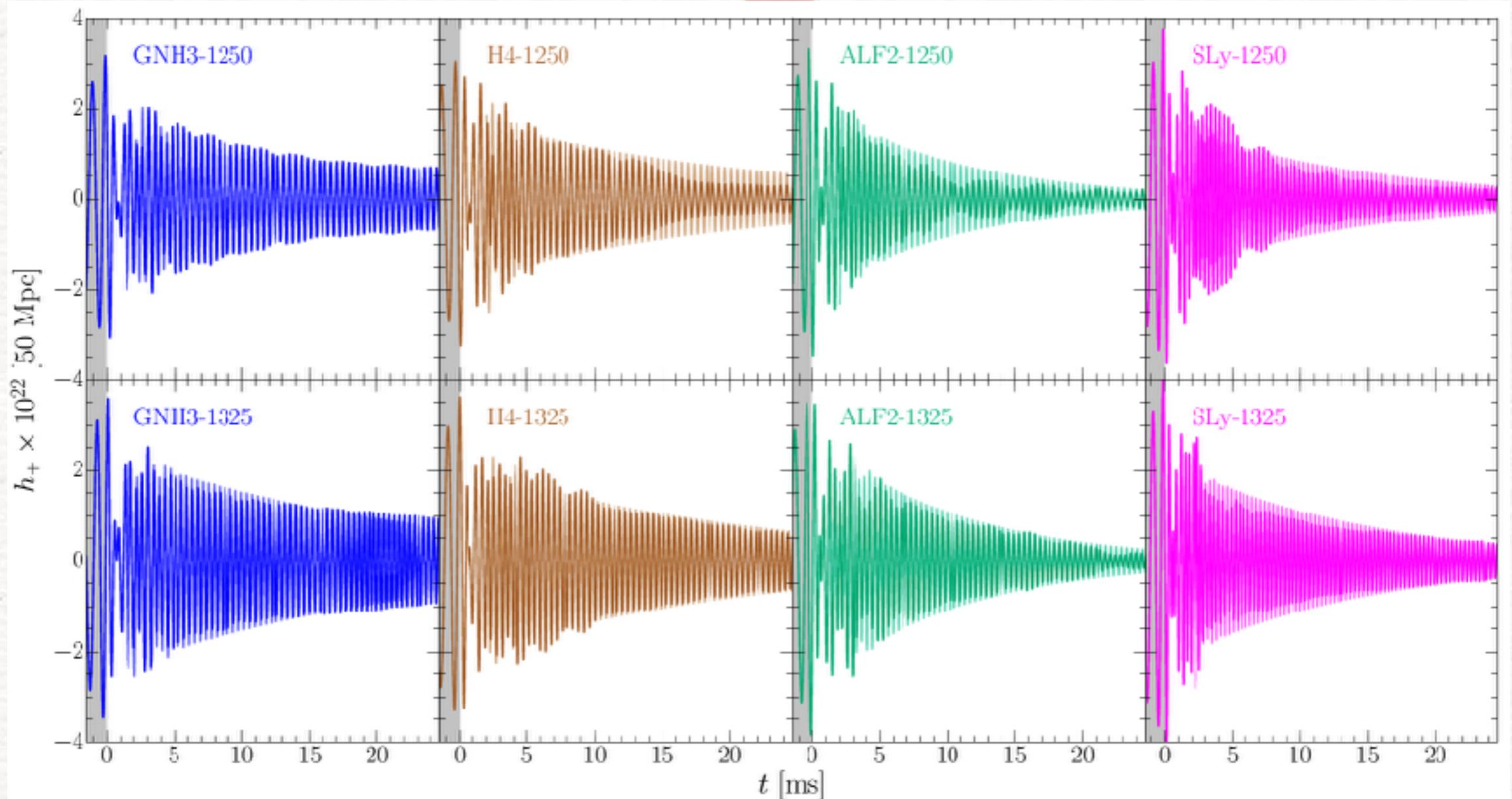
- **Postmerger** appears hopeless but isn't (Clark+14, 16; Bose+17)



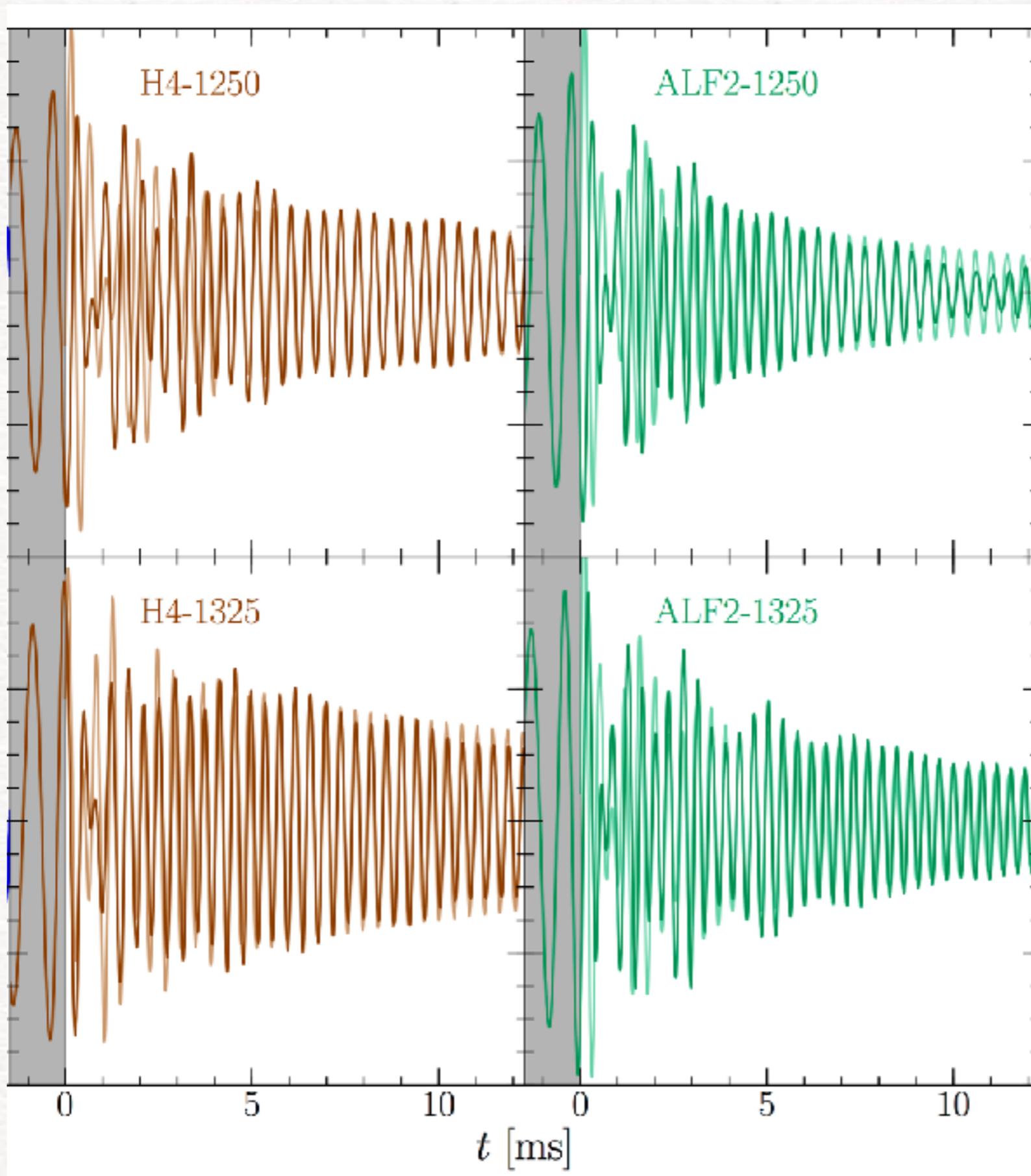
Analytical modelling of postmerger waveform

- Knowledge of spectral properties provides **analytic ansatz**

$$h(t) = \alpha \exp(-t/\tau_1) [\sin(2\pi f_1 t) + \sin(2\pi(f_1 - f_{1\epsilon})t) + \sin(2\pi(f_1 + f_{1\epsilon})t)] + \exp(-t/\tau_2) \sin(2\pi f_2 t + 2\pi\gamma_2 t^2 + \pi\beta_2).$$

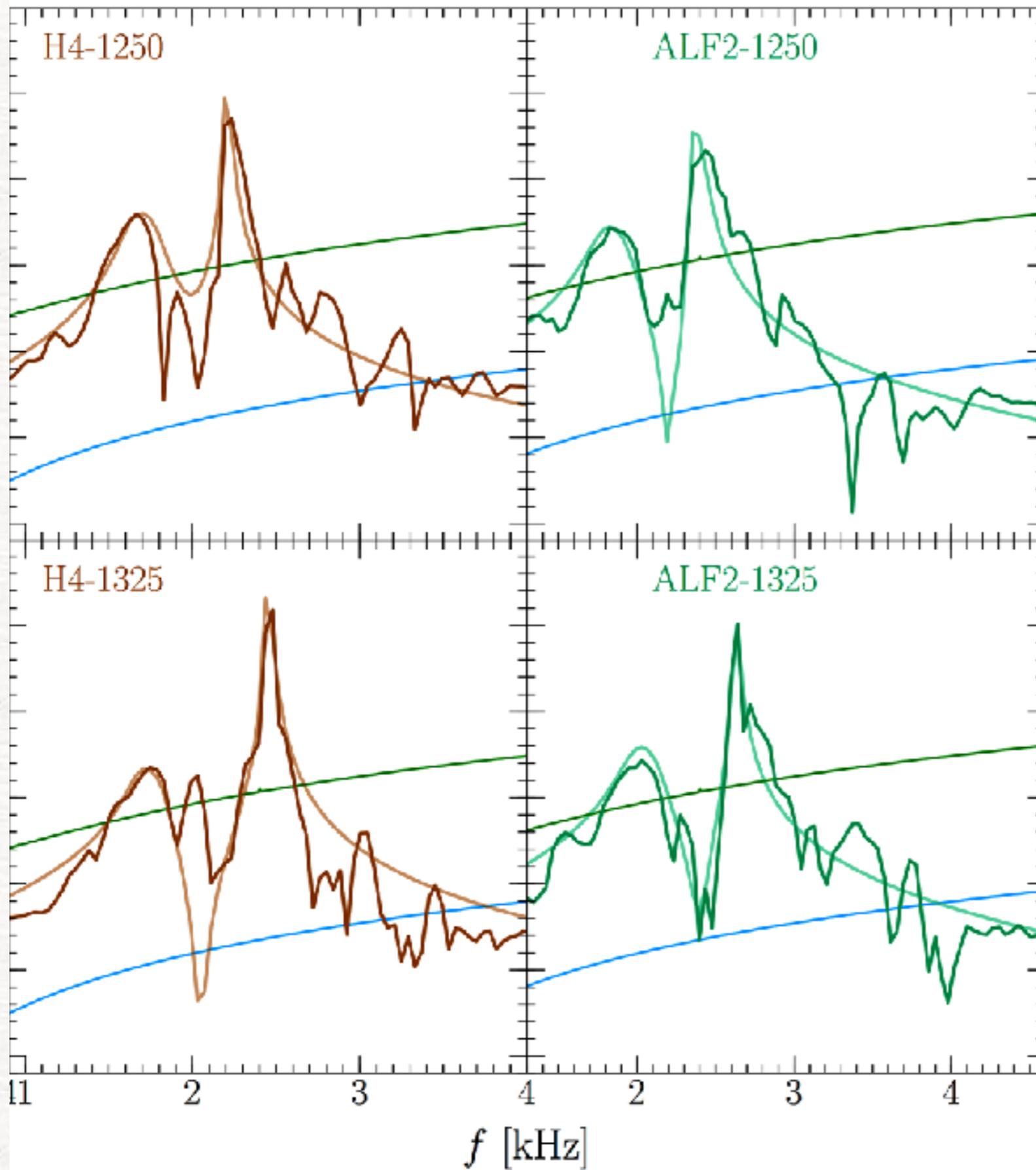


Analytical modelling of postmerger waveform



- Overall pretty decent fit in **phase**
- Fit in **amplitude** is less good but also less important

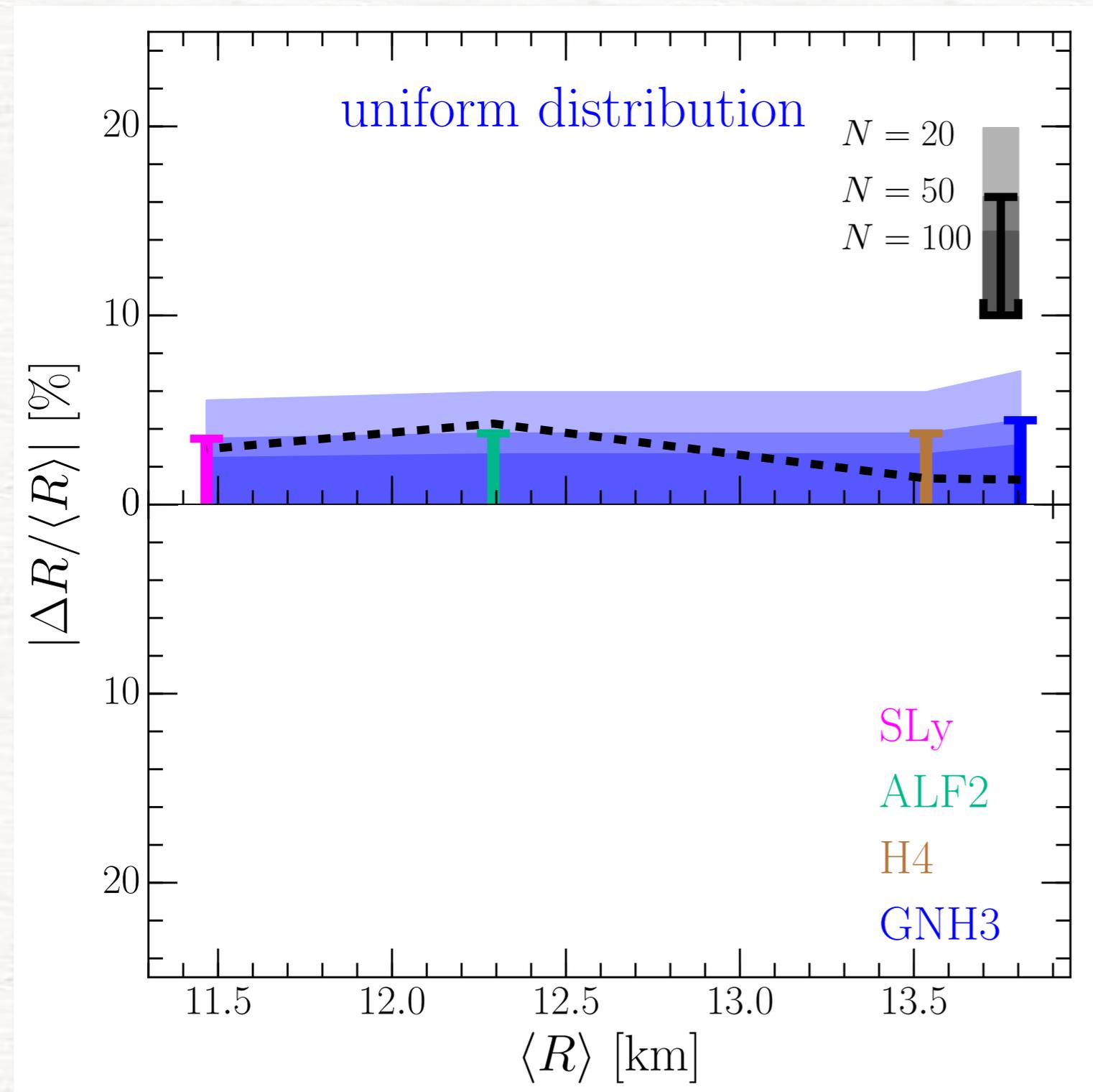
Analytical modelling of postmerger waveform



- Good match is clear also in **frequency space**

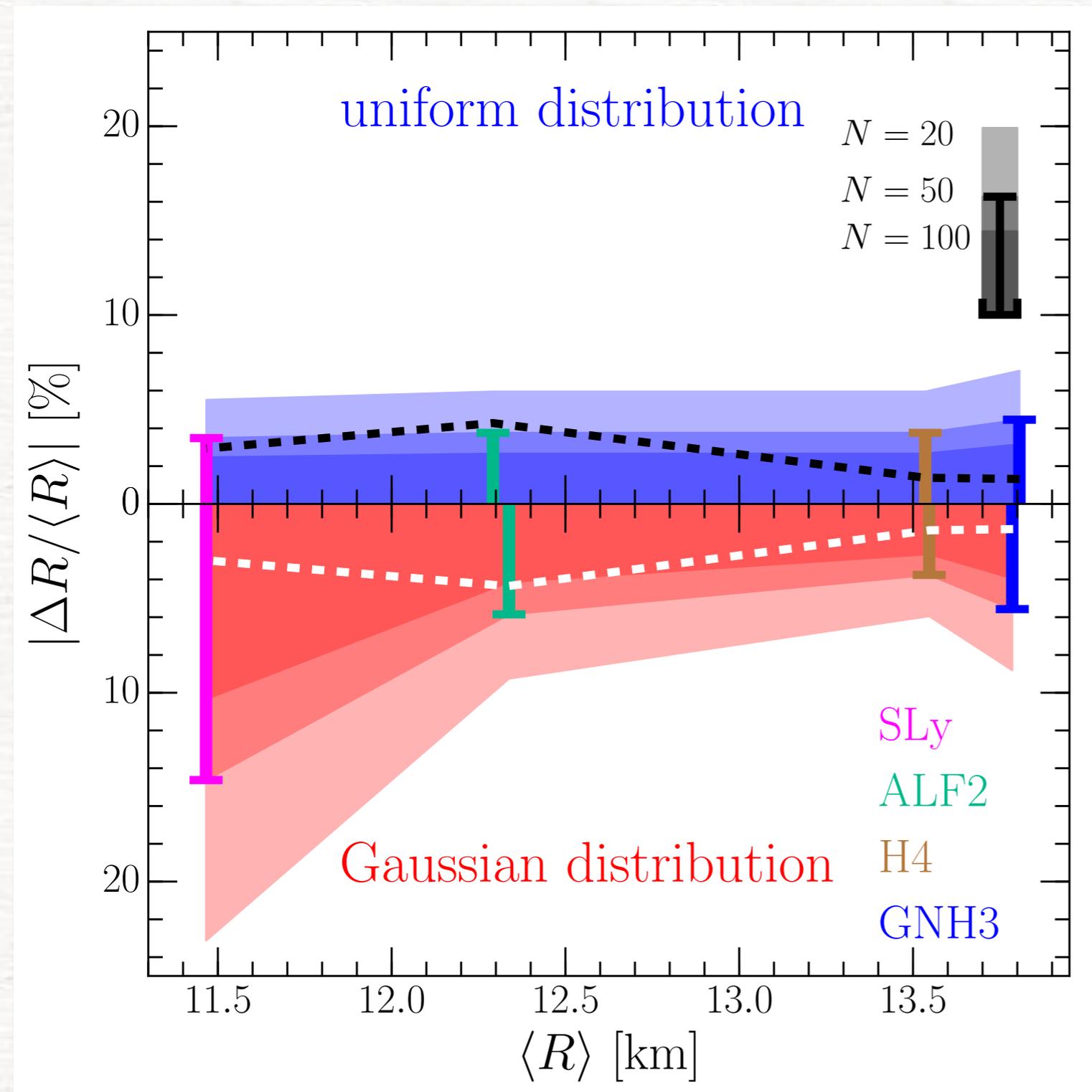
In summary: despite the complex signal, an **analytic** description of the **full GW signal** is now possible.

Constraining the radius: MonteCarlo vs Fisher



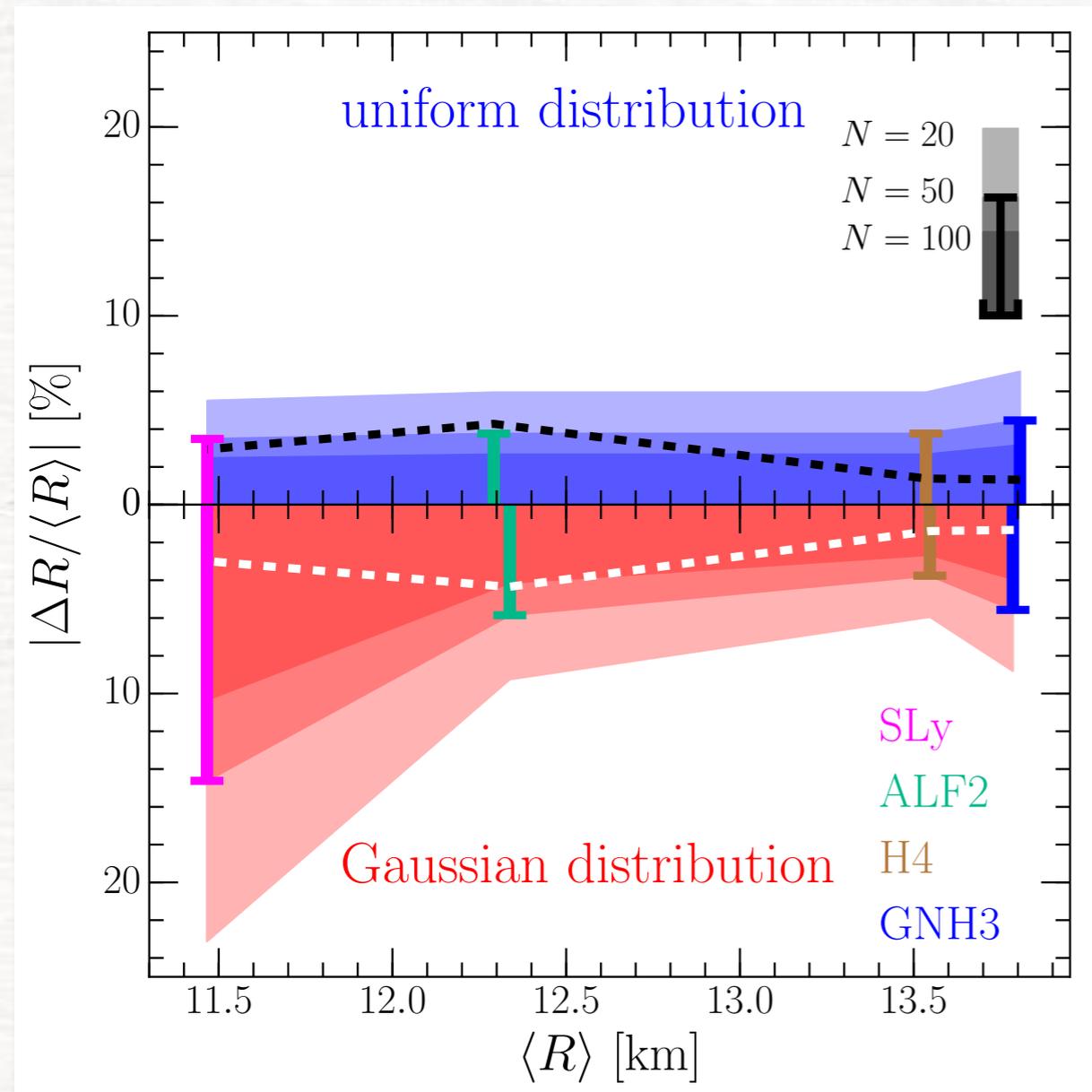
- **uniform** distribution in mass $[1.21, 1.38] M_{\odot}$ between 100 and 300 Mpc; isotropic distribution in space.
- dashed lines for results of Fisher-matrix analysis with $N=50$
- errors scale like \sqrt{N}

Constraining the radius: MonteCarlo vs Fisher



- **Gaussian** distribution in mass $[1.21, 1.38] M_{\odot}$ centred at $1.35 M_{\odot}$ with variance 0.05 Binaries are between 100 and 300 Mpc; isotropic distribution in space.
- dashed lines for results of Fisher-matrix analysis with $N=50$
- errors scale like \sqrt{N}

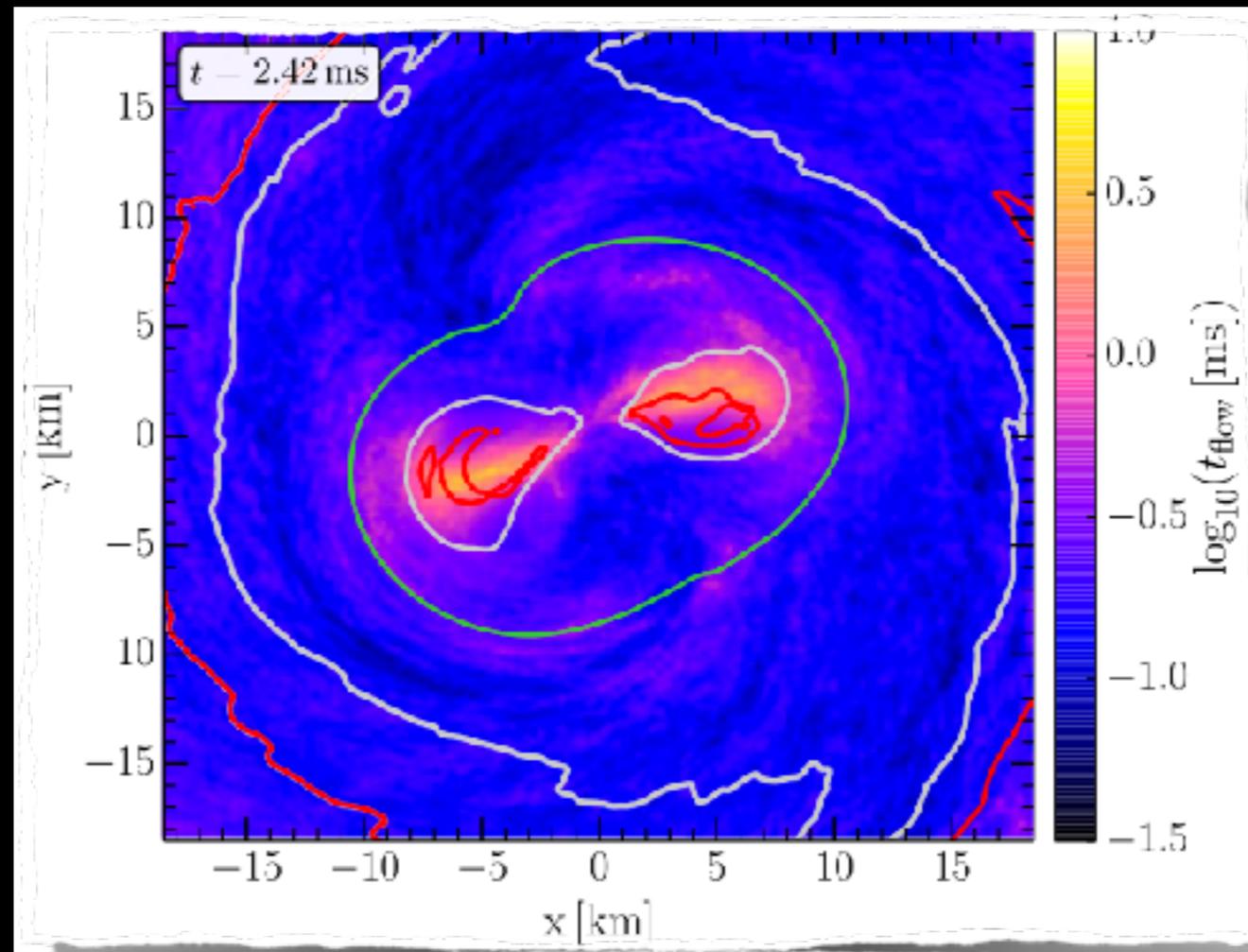
All in all



- stiff EOSs: $|\Delta R / \langle R \rangle| < 10\%$ for $N \sim 20$
- soft EOSs: $|\Delta R / \langle R \rangle| \sim 10\%$ for $N \sim 50$
- discriminating stiff/soft EOSs will be possible even with moderate N
- discriminating two-stiff /two-soft EOSs will be harder
- very soft EOSs remain a challenge
- golden binary: **SNR ~ 6 at 30 Mpc**
 $|\Delta R / \langle R \rangle| \lesssim 2\%$ at 90% confidence

Viscous dissipation in the post-merger

Alford, Bovard, Hanauske, LR, Schwenzer (2018)



Viscous contributions

- Viscous dissipation is normally neglected in numerical modelling on assumption microscopic viscosity too small.
- Possible channels are:
 1. nuclear-matter **shear** viscosity
 2. nuclear-matter **bulk** viscosity
 3. neutrino **shear** viscosity (Guilet+ 2016)
 4. “**MRI-induced**” **viscosity** (Radice2017, Shibata+2017a, b)
- Channels 3. and 4. act on timescales typical of MRI, which depends on B-field and very **uncertain** still.
- Impact on GWs depends on value for viscous angular momentum transport; everyone’s bet?... $\tau \gtrsim 10 - 100$ ms

Viscous contributions: I. shear viscosity

- Low-temperature, **electron-dominated** regime, i.e.

$$T \lesssim 10 \text{ MeV}$$

$$\tau_{\eta}^{(e)} \approx 1.6 \times 10^8 \text{ s} \left(\frac{z_{\text{typ}}}{1 \text{ km}} \right)^2 \left(\frac{T}{1 \text{ MeV}} \right)^{\frac{5}{3}} \left(\frac{n_0}{n_B} \right)^{\frac{5}{9}} \left(\frac{0.1}{x_p} \right)^{\frac{14}{9}},$$

- High-temperature, **neutrino-dominated** regime, i.e.

$$T \gtrsim 10 \text{ MeV}$$

$$\tau_{\eta}^{(\nu)} \approx 54 \text{ s} \left(\frac{0.1}{x_p} \right) \left(\frac{m_n^*}{0.8 m_n} \right)^2 \left(\frac{\mu_e}{2 \mu_\nu} \right)^4 \left(\frac{z_{\text{typ}}}{1 \text{ km}} \right)^2 \left(\frac{T}{10 \text{ MeV}} \right)^2,$$

Hence, shear viscosity not relevant unless neutrinos dominate and flow is turbulent with $z_{\text{typ}} \sim 10 - 100 \text{ m}$; not likely

Viscous contributions: I. shear viscosity

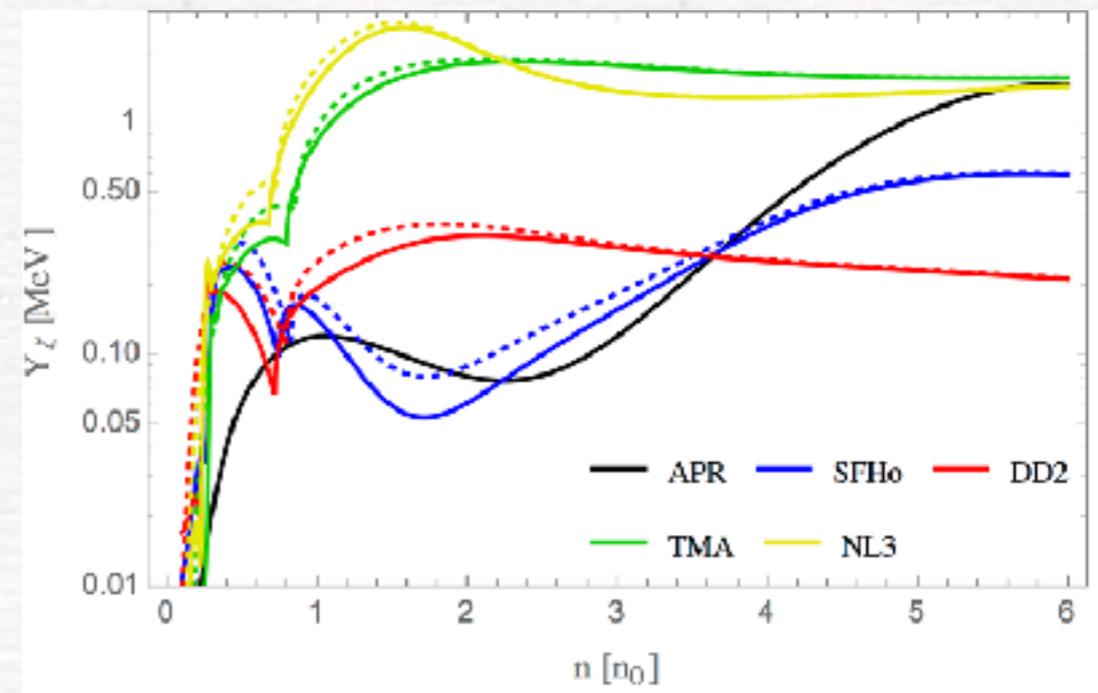
- Impact of bulk viscosity depends sensitively on process responsible for flavor re-equilibration.
- If **direct-Urca** dominates, bulk viscosity will be very **small**: never possible for softer EOSs, hard for stiff EOS at small T.
- If **modified-Urca** dominates, then bulk viscosity

$$\mathcal{E}_{\text{comp}} \approx K \bar{n} (\Delta n / \bar{n})^2 / 18$$

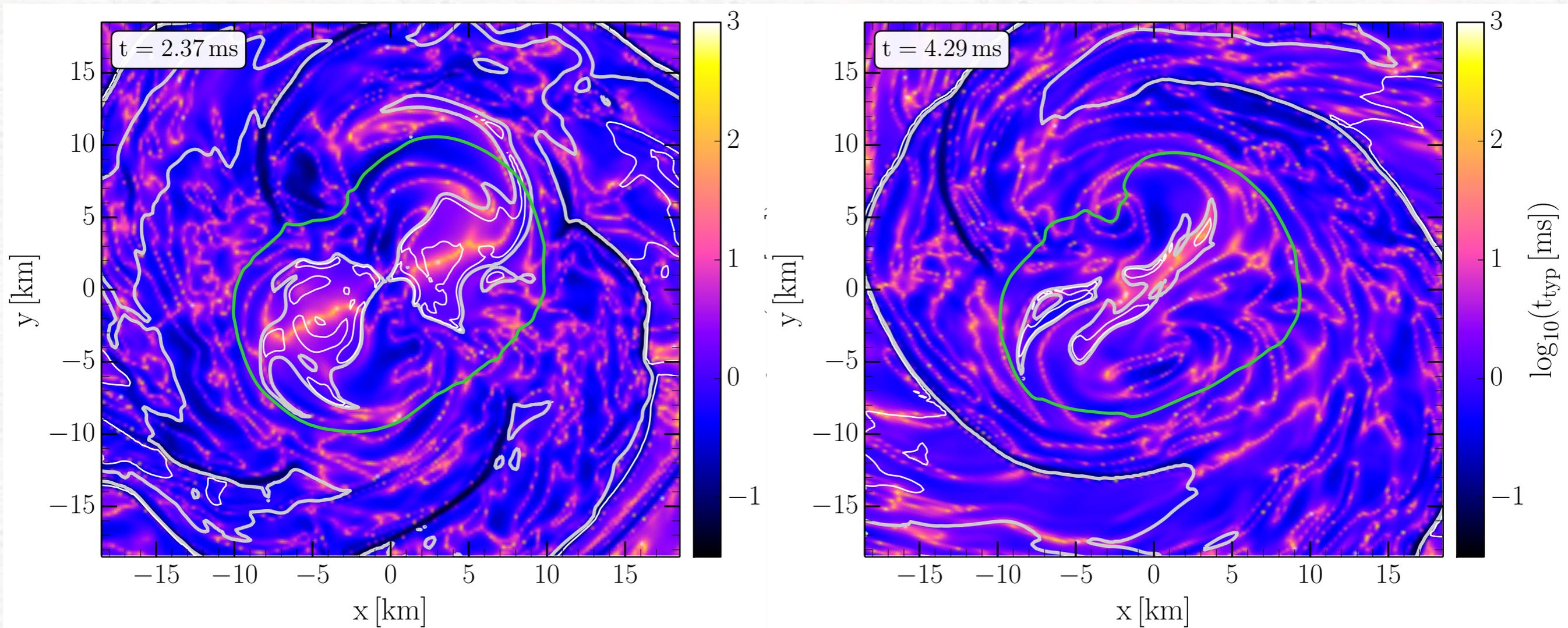
$$\begin{aligned} \tau_{\zeta} &\equiv \mathcal{E}_{\text{comp}} / (d\mathcal{E}/dt)_{\text{bulk}} \approx K \bar{n} t_{\text{exp}}^2 / (36\pi^2 \bar{\zeta}) \\ &\approx 7 \text{ ms} \left(\frac{t_{\text{exp}}}{1 \text{ ms}} \right) \left(\frac{K}{250 \text{ MeV}} \right) \left(\frac{0.1 \text{ MeV}}{Y_{\zeta}} \right) \end{aligned}$$

$t_{\text{exp}} \sim$ bulk-dissipation timescale of internal energy

K : nuclear compressibility at n_0



Viscous contributions



$$t_{\text{exp}}^{\text{inst}} \sim \frac{\rho}{D_t \rho} = \frac{1}{\nabla \cdot \vec{v}}$$

right after merger

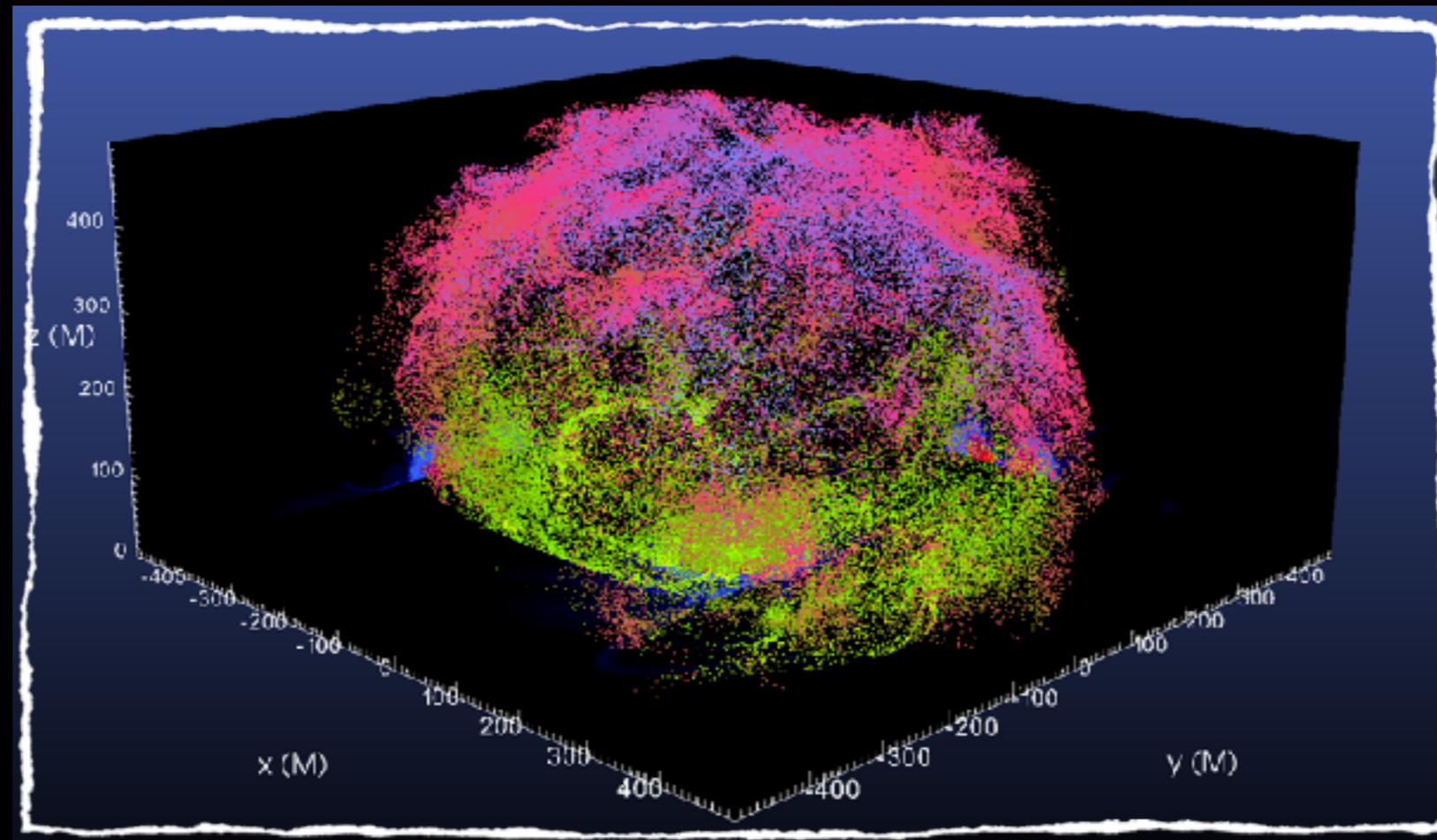
$$t_{\text{exp}} \lesssim \tau_{\text{dyn}} = \frac{R}{c_s}$$

instantaneous bulk-dissipation timescale can be measured in simulations.

bulk-dissipation timescale comparable with dynamical timescale in large portions of the object: cannot be ignored

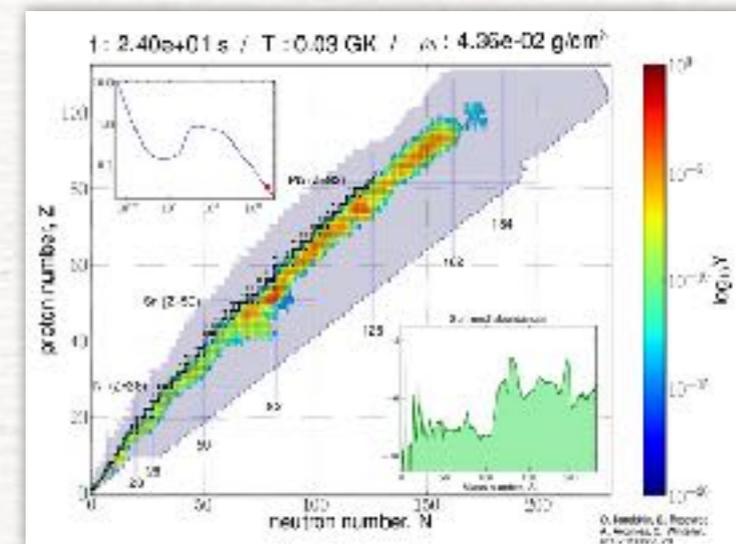
Ejected matter and nucleosynthesis

LR, Most, Weih (2018)



Nucleosynthesis

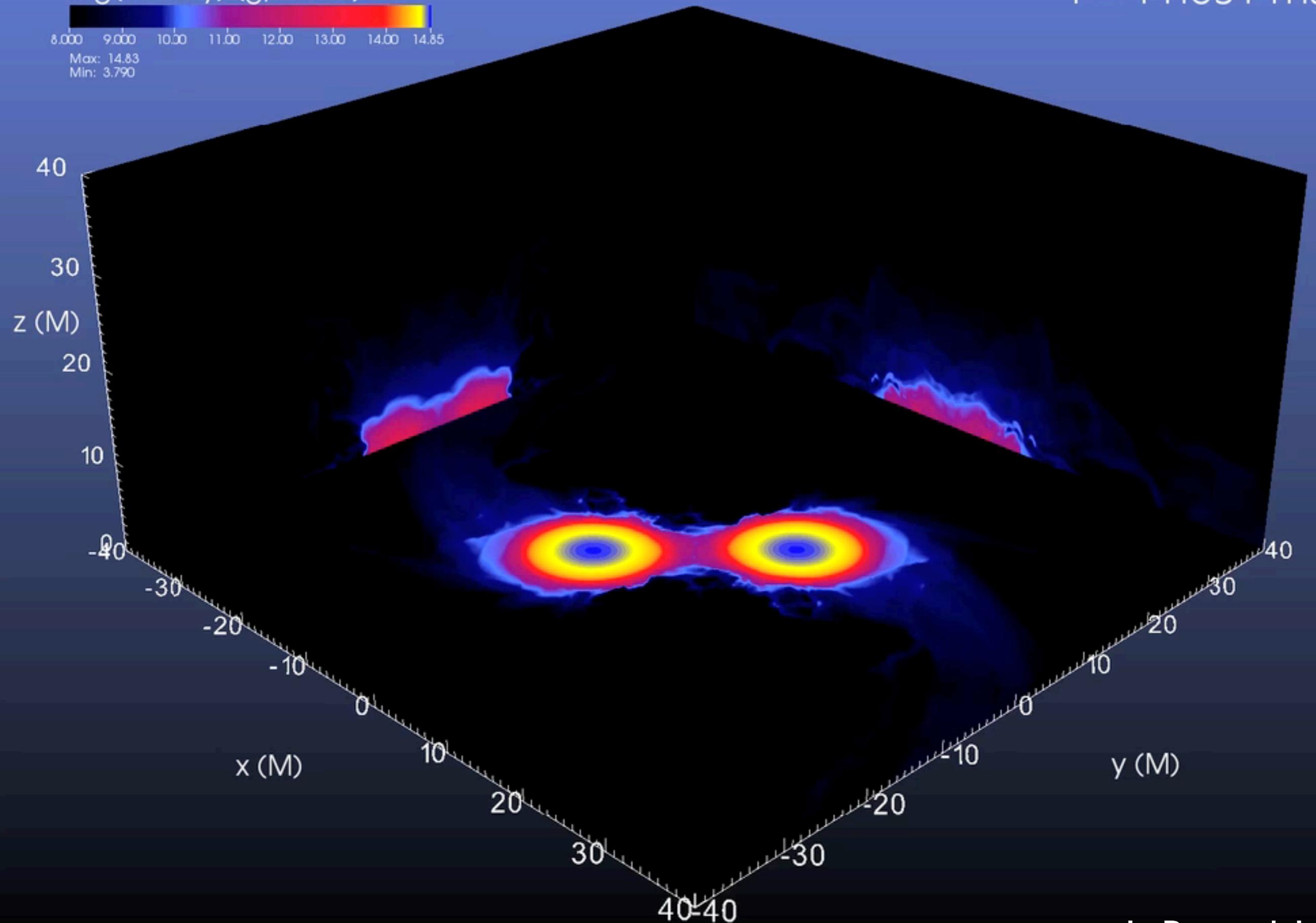
- Already in the 50's, nuclear physicists had tracked the production of elements in stars via nuclear fusion.
- **Heavy elements** ($A \gtrsim 56$) cannot be produced in stellar interiors but can be synthesised during a **supernova**.
- Modern numerical simulations of supernovae have shown that the temperature and energies are not large enough to produce the “**very heavy**” elements ($A \gtrsim 120$).
- To produce such elements one needs very high temperatures and “**neutron-rich**” material.
- **Neutron-star mergers** seem perfect candidates for this process!



log(density) (g/cm³)



$t = 11.801$ ms

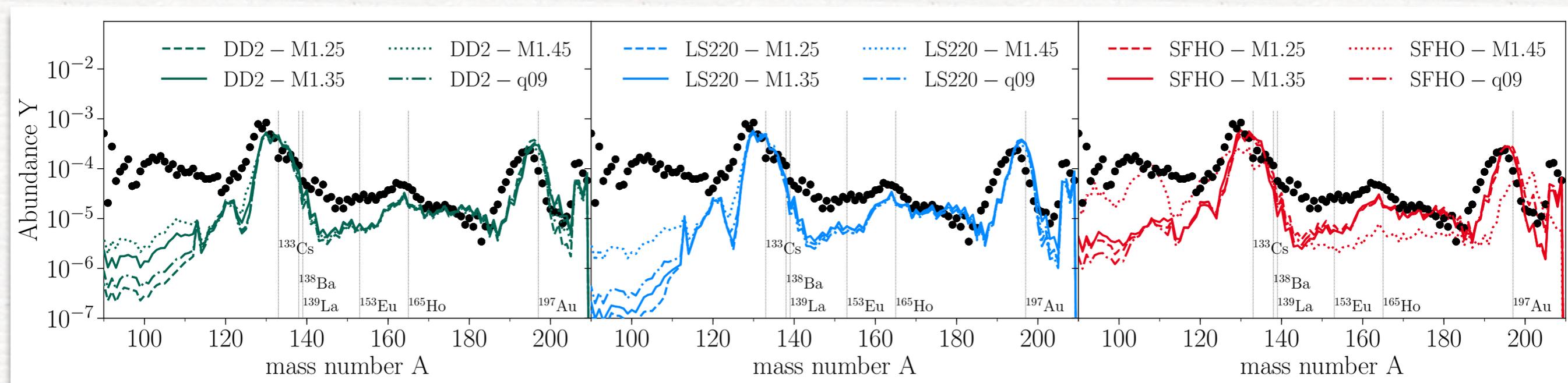


L. Bovard, LR

Relative abundances

Bovard+ 17

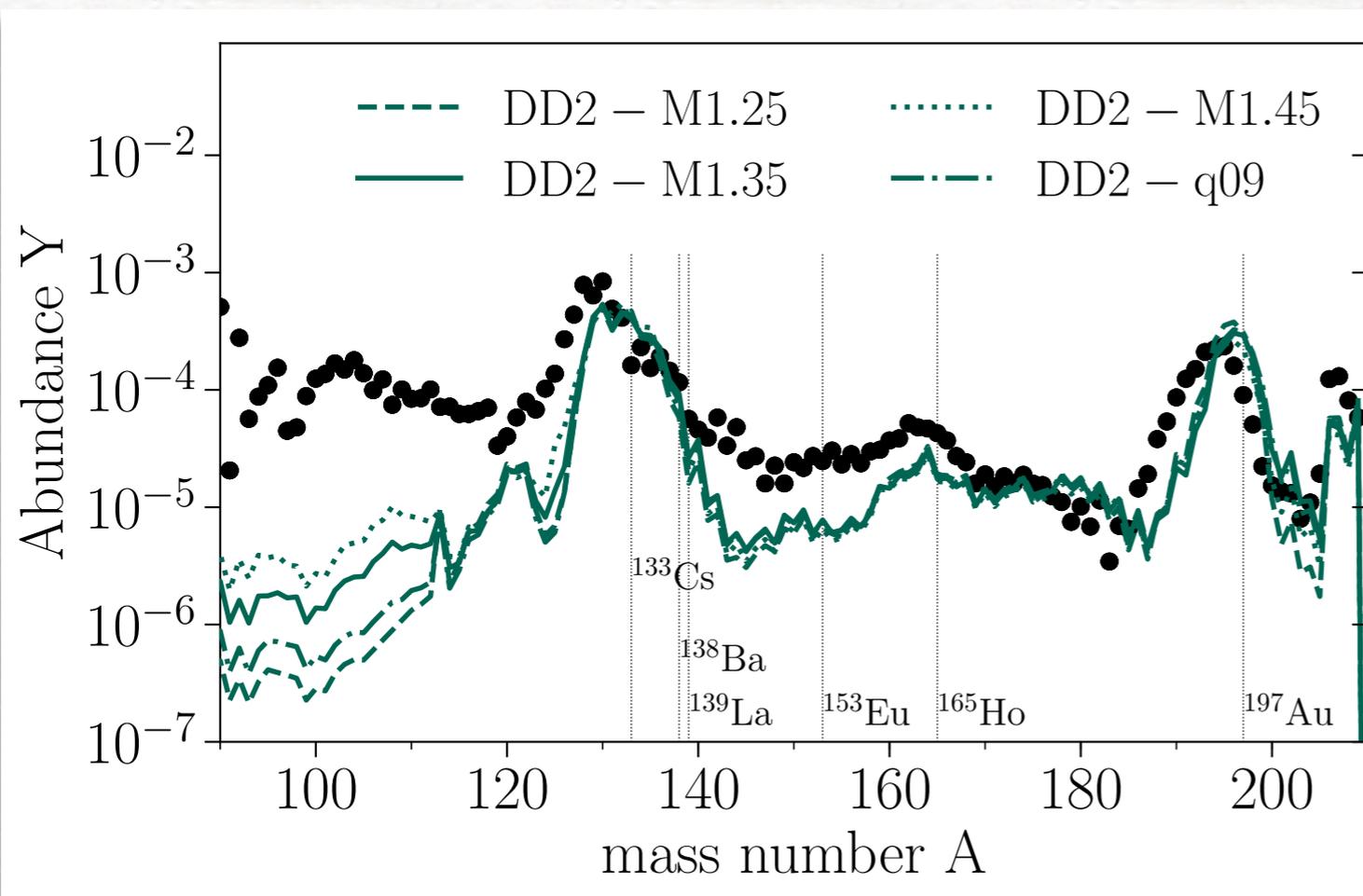
- Abundance pattern for $A > 120$ in good agreement with solar.
- Even **tiny amounts** of ejected matter ($0.01 M_{\odot}$) sufficient to explain observed abundances.
- Extremely **robust** behaviour across different **EOSs, masses, nuclear reactions and merger type**



Relative abundances

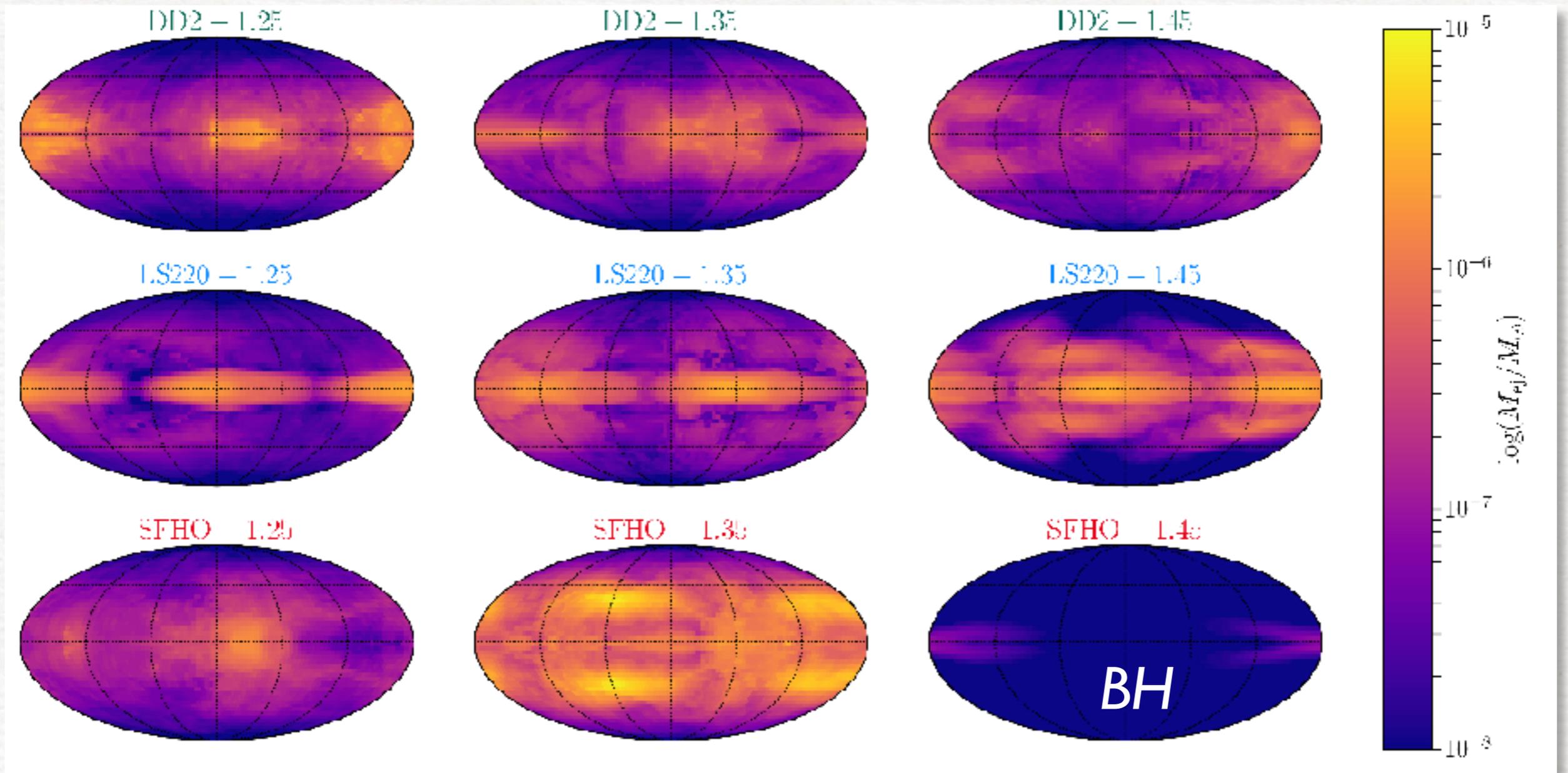
Bovard+ 17

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- Extremely **robust** behaviour across different **EOSs, masses, nuclear reactions and merger type**



- GW170817 produced total of **16,000** times the mass of the Earth in heavy elements (**10** Earth masses in **gold/platinum**)
- We are not only **stellar dust** but also **neutron-star dust!**

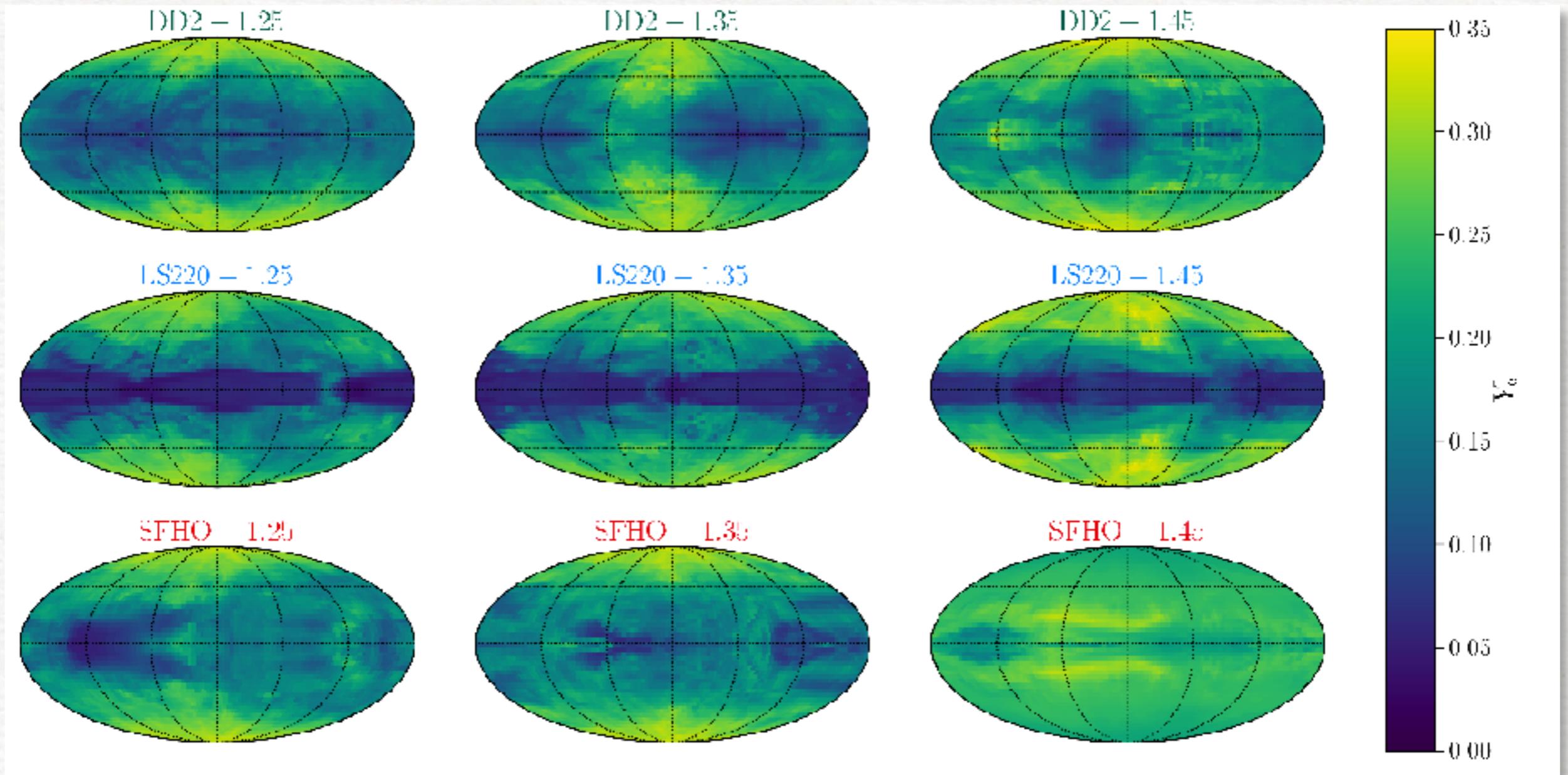
Spatial distributions: M_{ej} Bovard+ 17



Spatial distribution of M_{ej} impacts detectability of EM counterpart:

- ★ most of M_{ej} lost at low latitudes;
- ★ depending on EOS/mass, contamination also in polar regions

Spatial distributions: Y_e Bovard+ 17

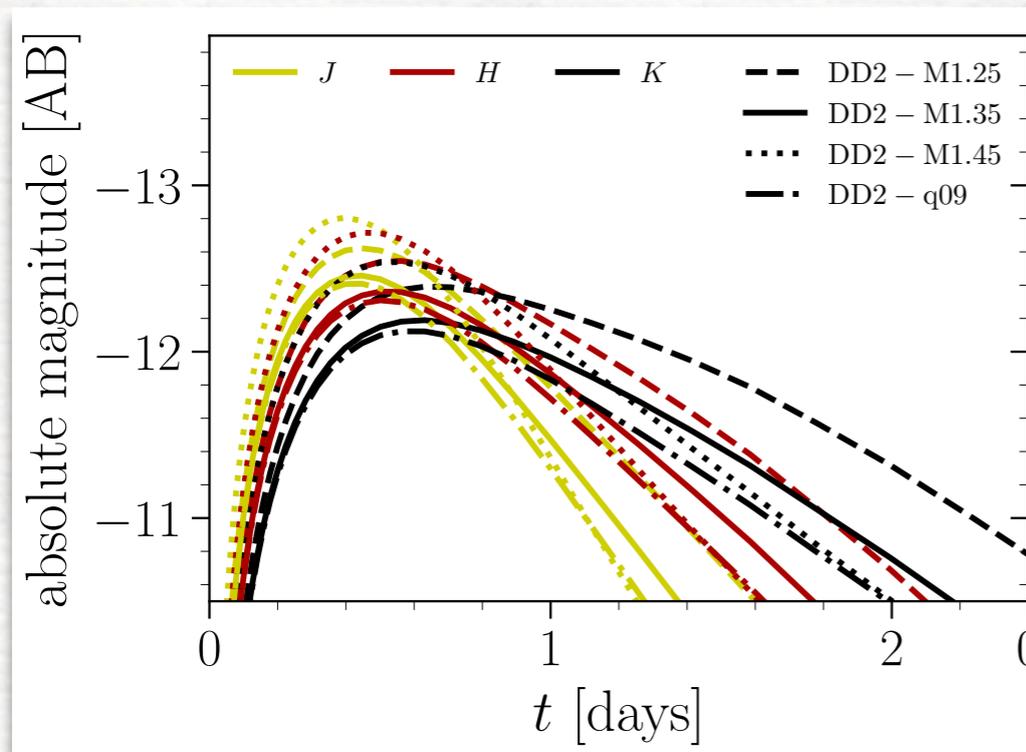


Spatial distribution of Y_e impacts detectability of EM counterpart:

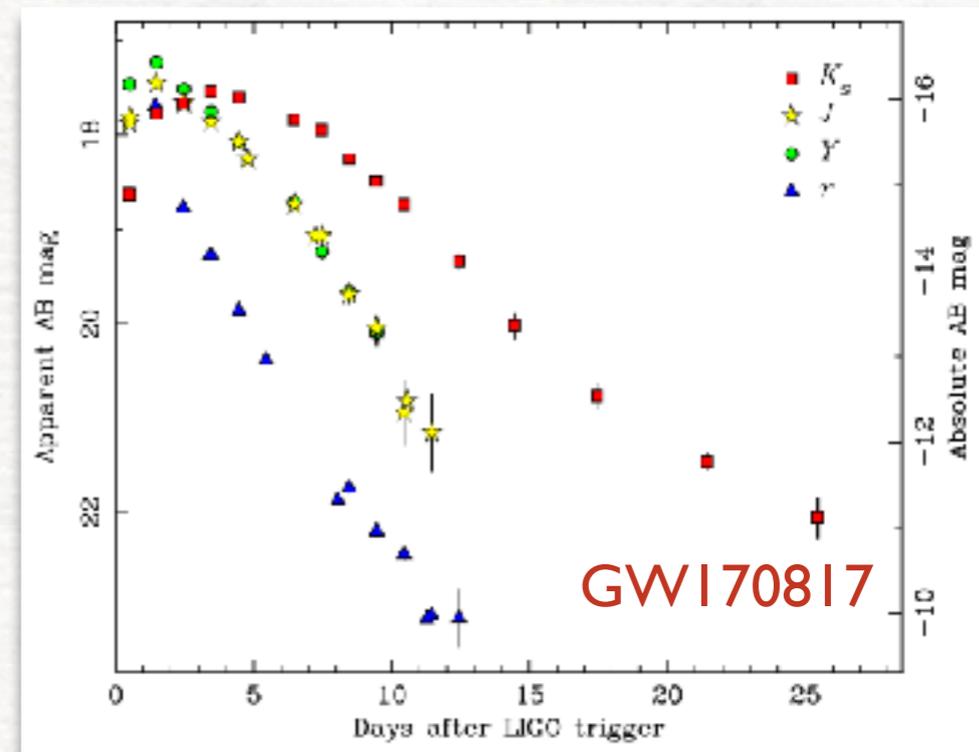
- ★ **high** Y_e in **polar** regions: **blue** (optical) macronova
- ★ **low** Y_e in **equatorial** regions: **red** (FIR) macronova

Kilonova emission

- Ejected matter undergoes **nucleosynthesis** as expands and cools.
- When critical densities and temperatures are reached, matter undergoes radioactive decay emitting light (optical/infrared): **kilonova/macronova** (Li & Paczynski '98).



simulations



observations

- Astronomical observations of GW170817 show **kilonova emission**: evidence connection **GRBs** and **binary neutron stars!**

Conclusions

- * Spectra of post-merger shows clear peaks, some of which are **”quasi-universal”**. that is, independent of the EOS
- * Used together with tens of observations and analytic modelling of post-merger, universal relations set tight constraints on EOS.
- * Magnetic fields unlikely to be detected during the inspiral but **important** after the merger: instabilities and EM counterparts.
- * **Mergers** lead to tiny but important ejected matter and macronova emission. “high-A” nucleosynthesis very robust.
- * A single event **GW170817** has provided wealth of information and new limits on the **maximum mass**

Merging binaries of NSs are Einstein’s **richest laboratory:**

GWs, nuclear physics, astrophysics.

Much more in coming years!...