

# Feynman Integral

## Synergies Between Particle Physics and Gravitational Waves

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INFN Padova

Fellini Seminar

29th May, 2023



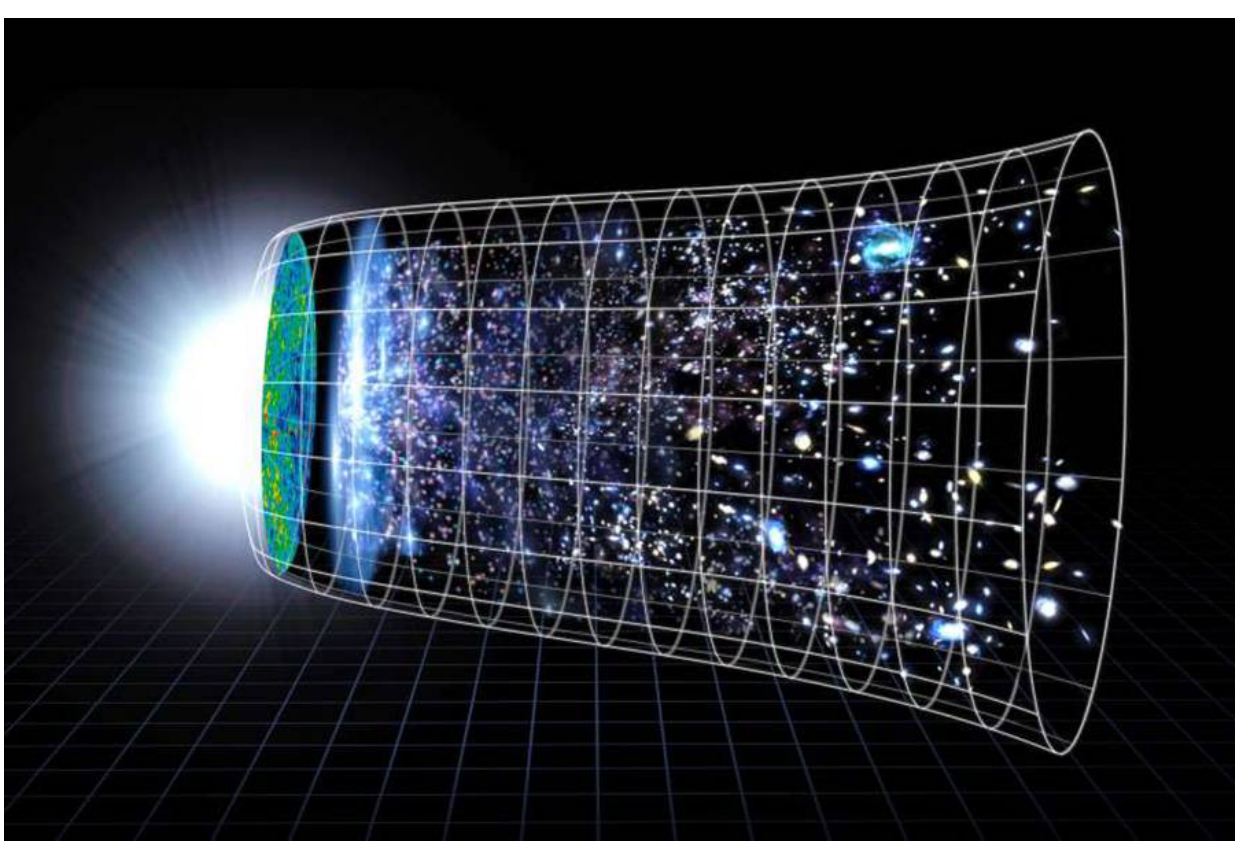
# Scattering Amplitudes



Collider Phenomenology



Gravitational Waves



Cosmology

Precise Observables  $\longleftrightarrow$  Scattering Amplitudes

# Scattering Amplitude: Connecting Theory and Experiment

## Perturbative Expansion of Cross-Section

$$\sigma = \sigma^{(0)} + \alpha_s \sigma^{(1)} + \alpha_s^2 \sigma^{(2)} + \dots$$

**LO**                      **NLO**                      **NNLO**

**Cross-section  
Measured in  
Experiment**

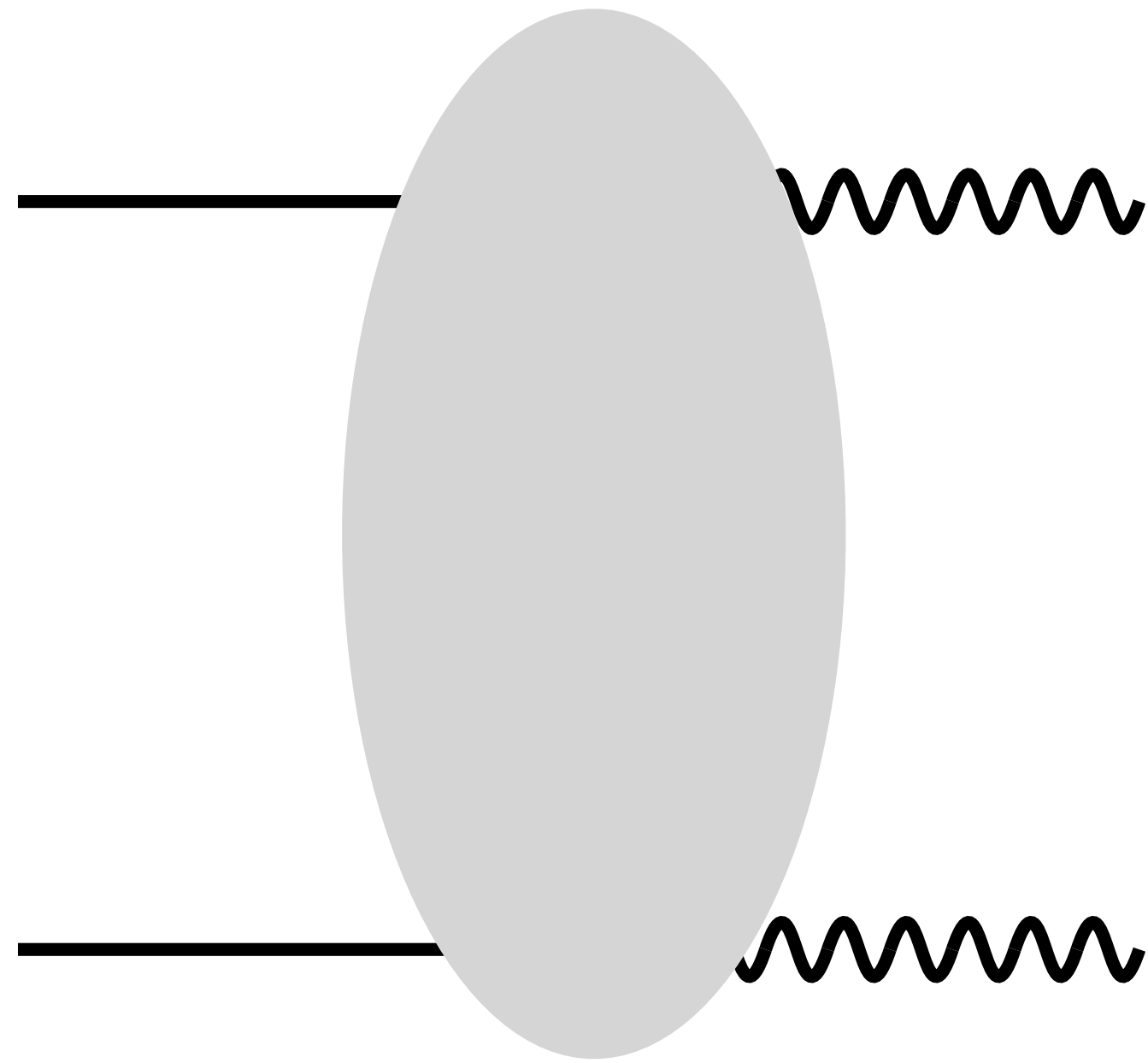
$$\sigma^0 \approx \int |\mathcal{M}_N^{(0)}|^2 d\Phi_N$$

**Theory**

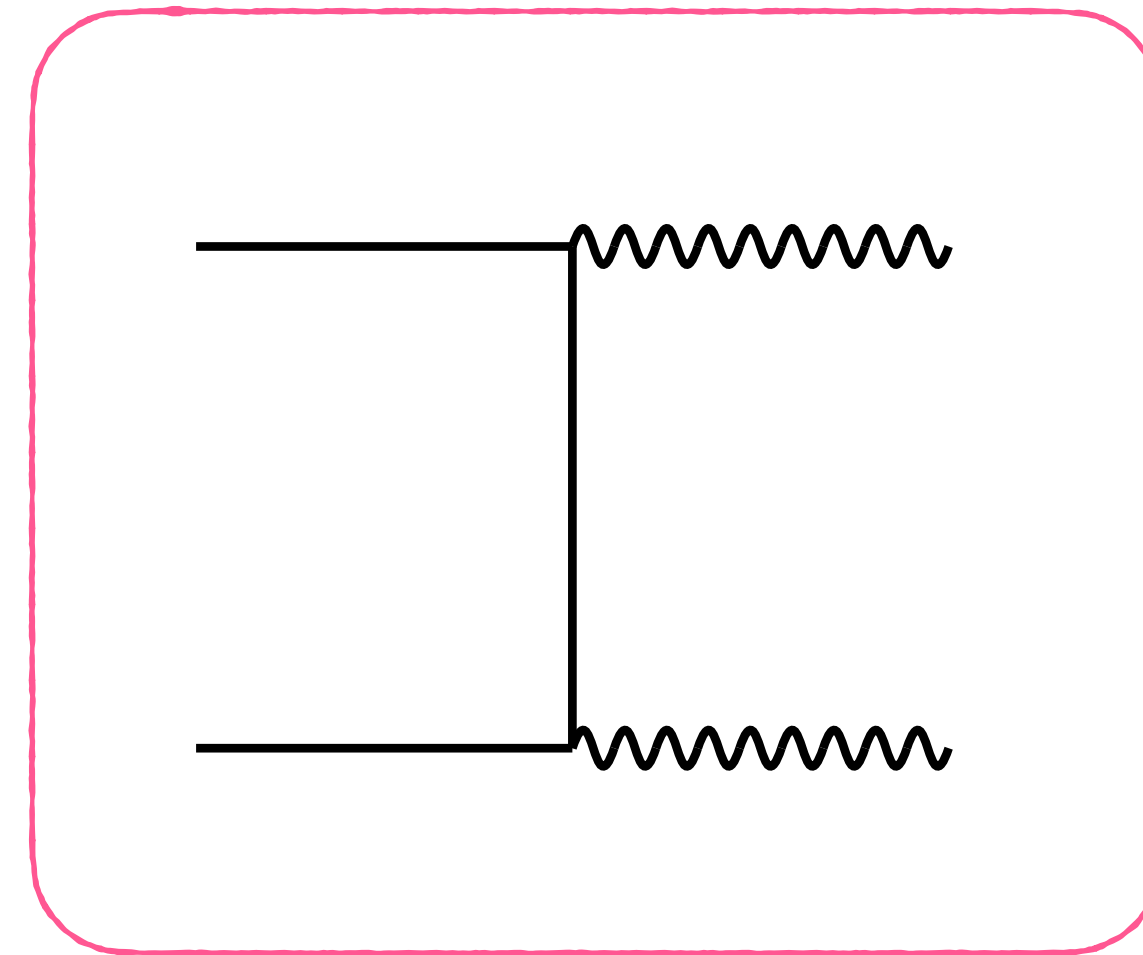
**Scattering Amplitudes**

Sum of Feynman Diagrams

# Scattering Amplitude

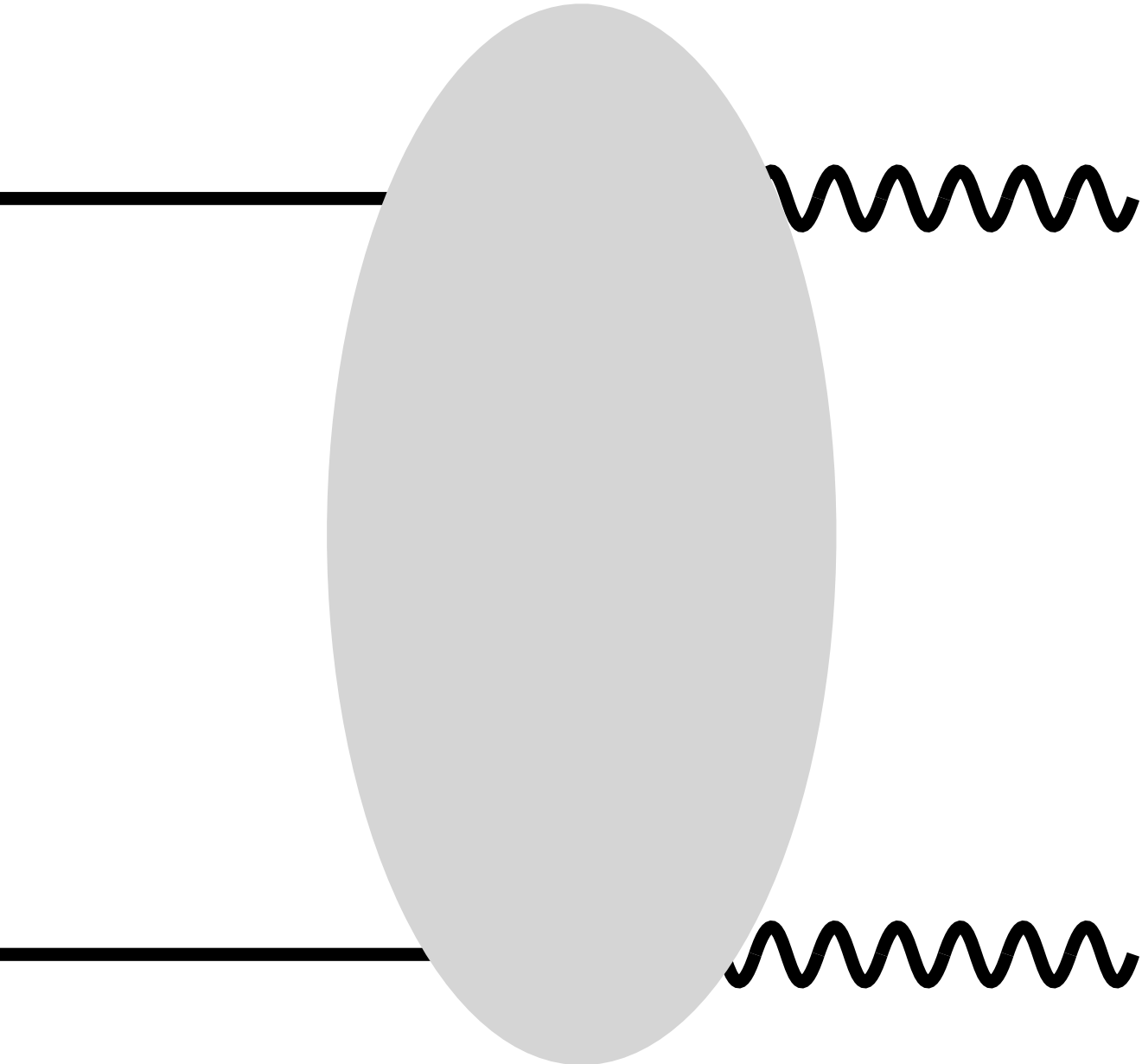


LO

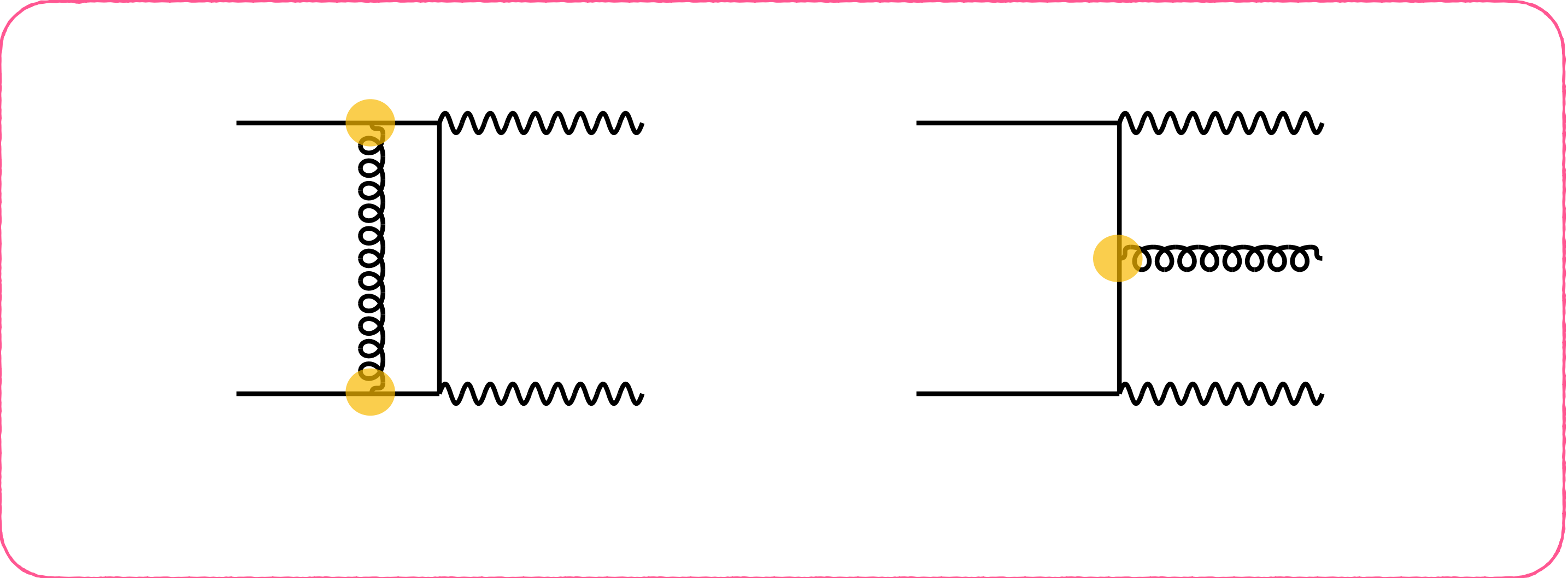


$$\sigma^0 \approx \int |\mathcal{M}_N^{(0)}|^2 d\Phi_N$$

# Scattering Amplitude



NLO

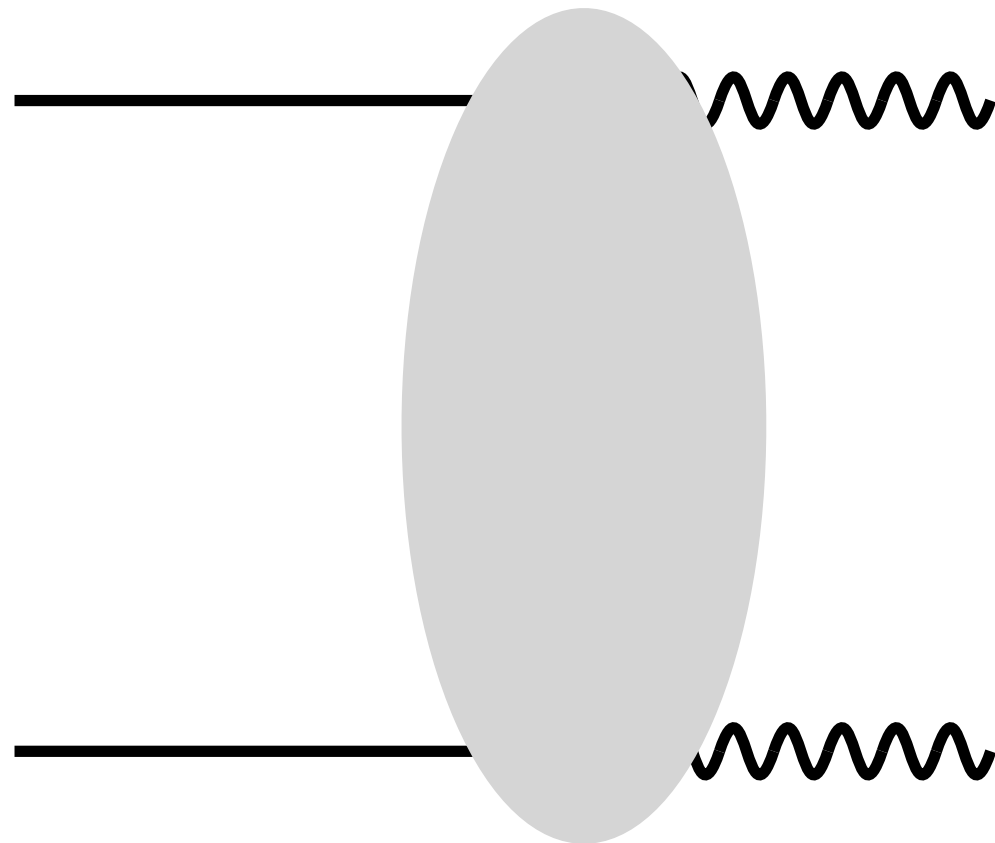


$$\sigma_N^{(1)} \approx \int 2\text{Re} \left( \mathcal{M}_N^{(0)*} \mathcal{M}_N^{(1)} \right) d\Phi_N + \int |\mathcal{M}_{N+1}^{(0)}|^2 d\Phi_{N+1}$$

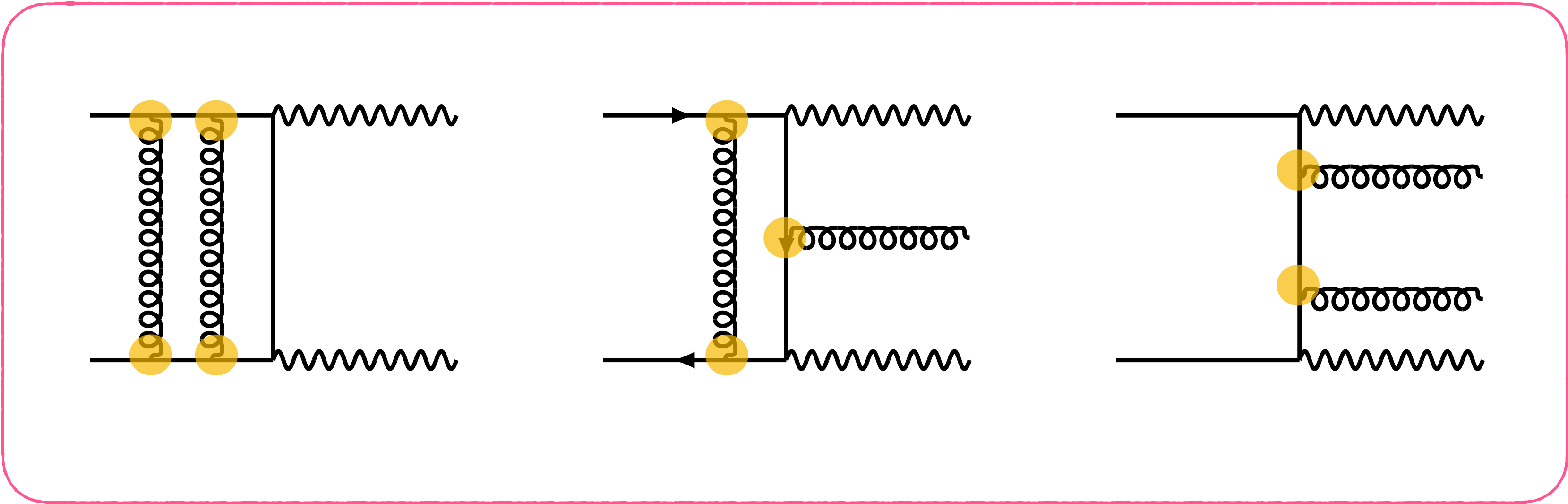
$$\int \left[ \frac{V_2}{\epsilon^2} + \frac{V_1}{\epsilon^1} + V_0 \right] d\phi_2$$

$$\int [R_0] d\phi_3$$

# Scattering Amplitude



NNLO



$$\sigma_N^{(2)} \approx \int 2\text{Re} \left( \mathcal{M}_N^{(0)*} \mathcal{M}_N^{(2)} \right) d\Phi_N + \int 2\text{Re} \left( \mathcal{M}_{N+1}^{(0)*} \mathcal{M}_{N+1}^{(1)} \right) d\Phi_{N+1} + \int |\mathcal{M}_{N+2}^{(0)}|^2 d\Phi_{N+2}$$

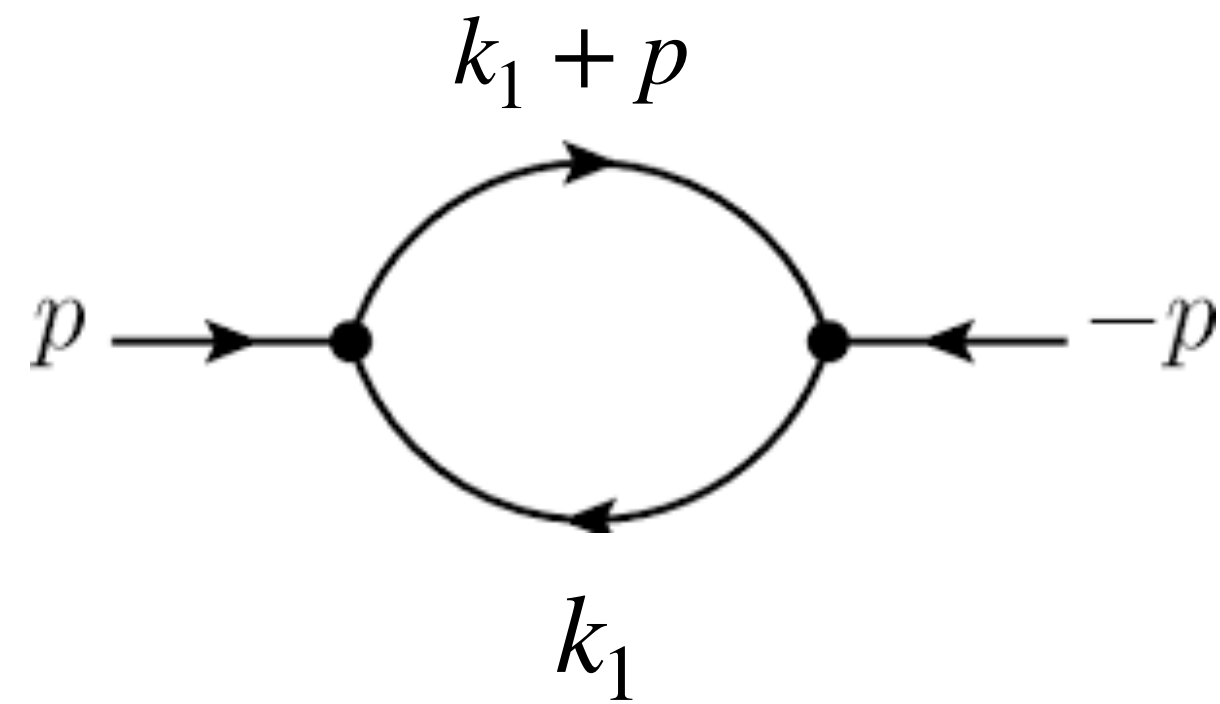
$$\int \left[ \frac{VV_4}{\epsilon^4} + \frac{VV_3}{\epsilon^3} + \frac{VV_2}{\epsilon^2} + \frac{VV_1}{\epsilon^1} + VV_0 \right] d\phi_2$$

$$\int \left[ \frac{RV_2}{\epsilon^2} + \frac{RV_1}{\epsilon^1} + RV_0 \right] d\phi_3$$

$$\int [RR_0] d\phi_4$$

# Loop Integral: An example

One Loop Massless Bubble



$$I(a_1, a_2) = \int \frac{d^d k_1}{(k_1^2)^{a_1} (k_1 + p)^2)^{a_2}}$$

$$D_1 = k_1^2$$

$$D_2 = (k_1 + p)^2$$

# Notion of Loop Integral

Number of Loops

$$I(a_1 \cdots a_N) = \int \cdots \int d^d k_1 \cdots d^d k_l \frac{\mathcal{N}(\{k_i\}, \{p_j\})}{D_1^{a_1} \cdots D_N^{a_N}}$$

Loop Momenta

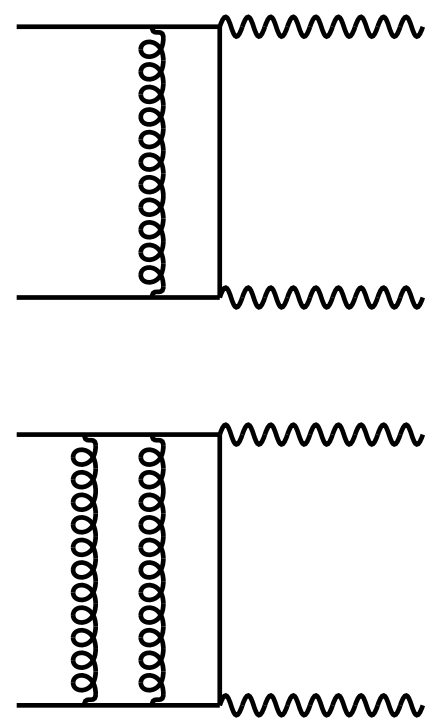
Number of Propagators

$$D_i = q_i^2 - m_i^2$$

$$q_i = \sum_j k_j + \sum_m p_m$$



# Computation of the Loop Amplitude



Generation of the Diagrams via QGRAF



Dirac algebra, Color sum, Trace in the numerators



Reduction to scalar integrals

$$\mathcal{M} = \sum_i a_i I_i \quad i = \mathcal{O}(10^5)$$

# Integration-By-Parts Identity

Chetyrkin, Tkachov

Loop momenta

$$\int \prod_{\alpha=1}^l d^d k_{\alpha} \frac{\partial}{\partial k_{j,\mu}} \left( \frac{v^{\mu}}{D_1^{a_1} \cdots D_N^{a_N}} \right) = 0$$

Loop and external momenta

$$\int_{\alpha=1}^l \prod d^d k_{\alpha} \frac{\partial}{\partial k_{j,\mu}} \left( \frac{v^{\mu}}{D_1^{a_1} \cdots D_N^{a_N}} \right) = \int_{\alpha=1}^l \prod d^d k_{\alpha} \left[ \frac{\partial v^{\mu}}{\partial k_{j,\mu}} \left( \frac{1}{D_1^{a_1} \cdots D_N^{a_N}} \right) - \sum_{j=1}^N \frac{a_j}{D_j} \frac{\partial D_j}{\partial k_{j,\mu}} \left( \frac{v^{\mu}}{D_1^{a_1} \cdots D_N^{a_N}} \right) \right]$$

$$C_1 I(a_1, \cdots a_N - 1) + \cdots + C_r I(a_1 + 1, \cdots a_N) = 0$$

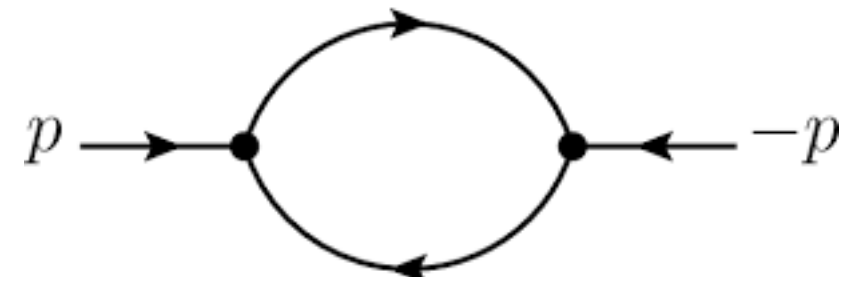
- ✱ Gives relations between different scalar integrals with different exponents
- ✱ **I(l+E)** number of equations
- ✱ Solve the system symbolically : Recursion relations
- ✱ Solve for specific integer value of the exponents : Laporta Algorithm

LiteRed

Fire, Reduze, Kira,...

# Integration-By-Parts Identity (Example)

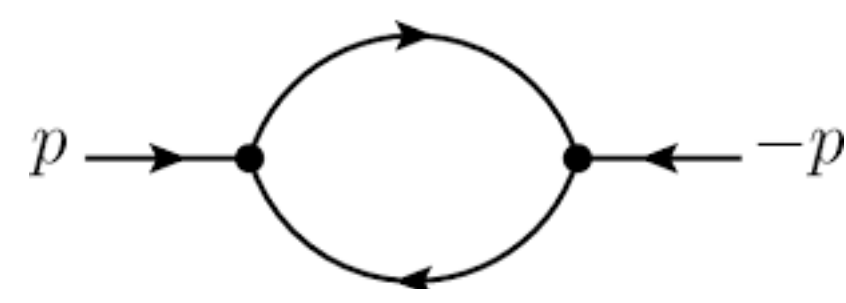
One Loop Massless Bubble



$$I(a_1, a_2) = \int \frac{d^d k_1}{(k_1^2)^{a_1} (k_1 + p)^2)^{a_2}}$$

# Integration-By-Parts Identity (Example)

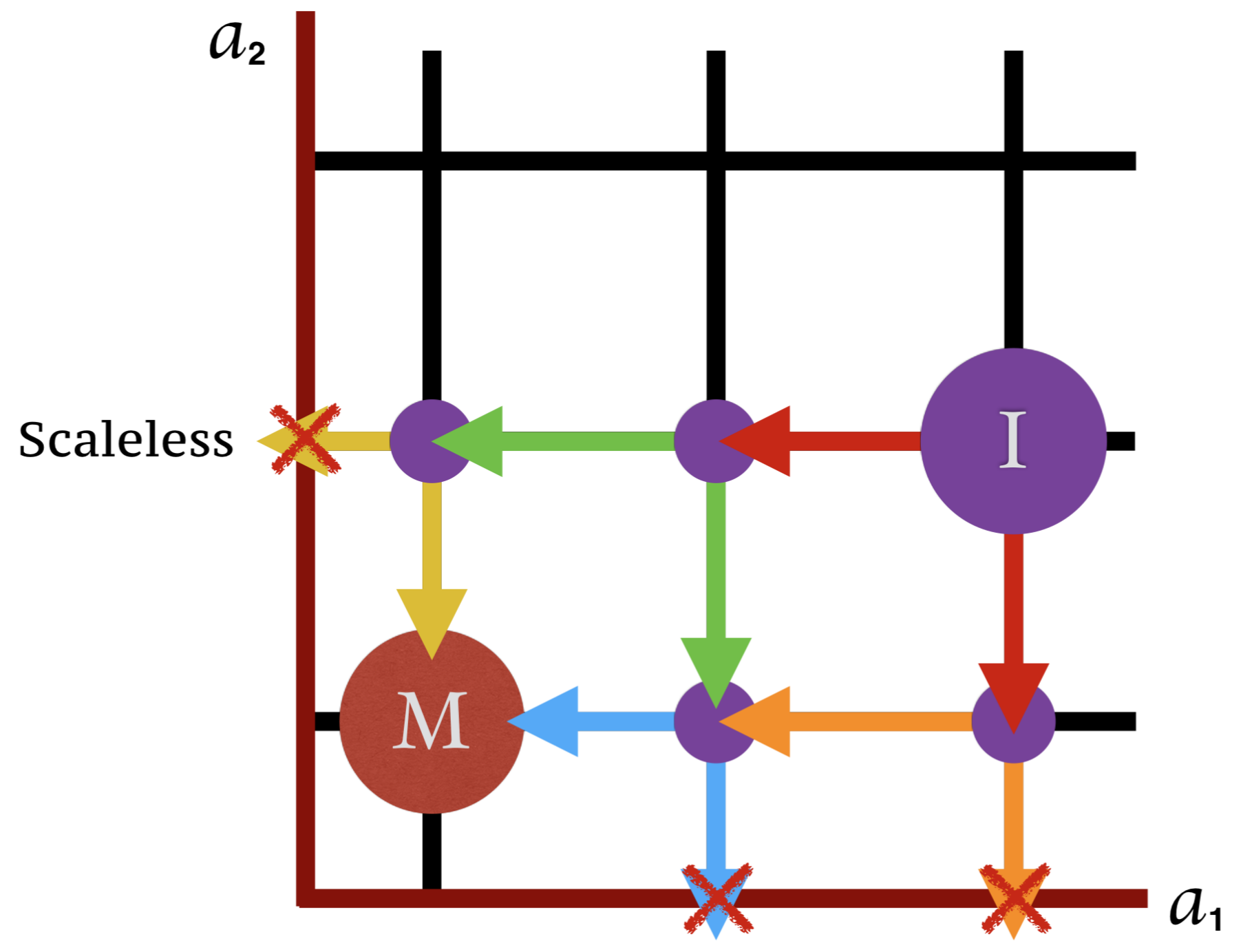
One Loop Massless Bubble



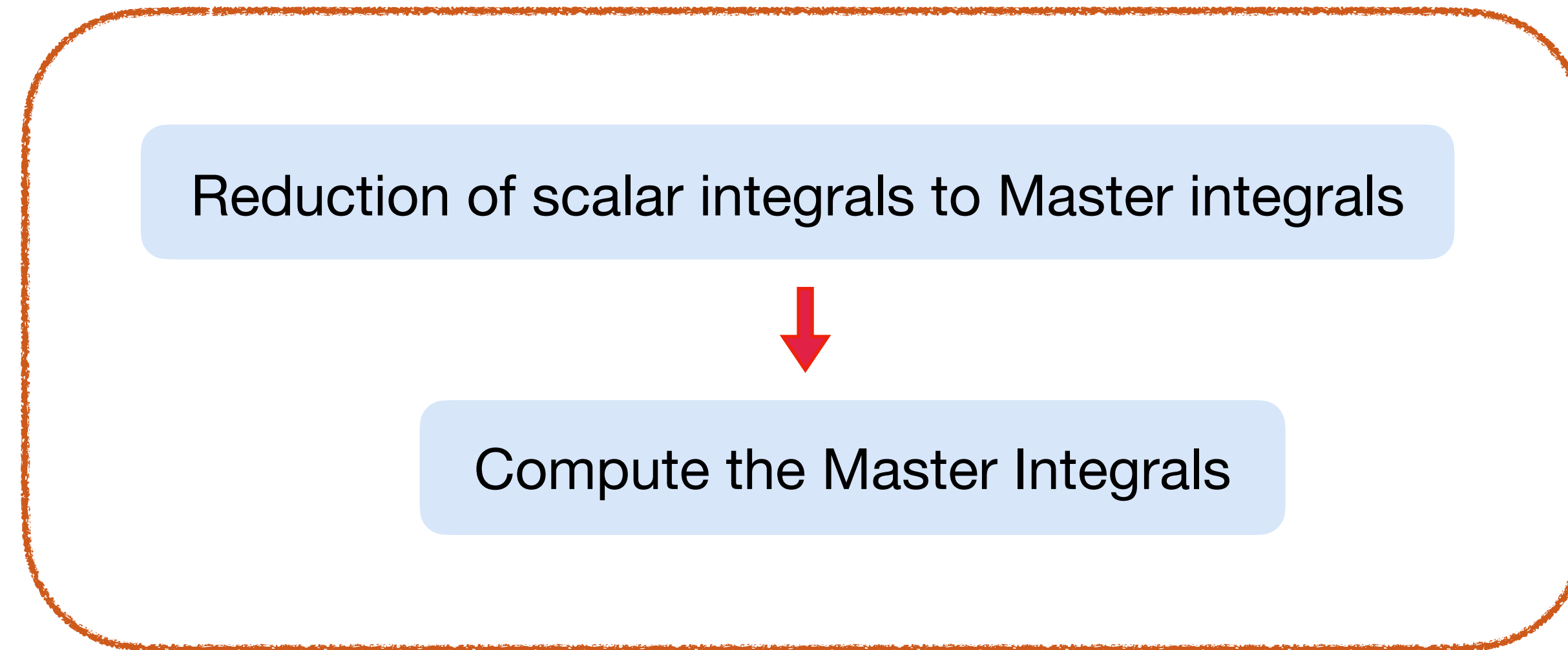
$$I(a_1, a_2) = \int \frac{d^d k_1}{(k_1^2)^{a_1} (k_1 + p)^2)^{a_2}}$$

IBP Identity

$$I(a_1, a_2) = \frac{a_1 + a_2 - d - 1}{p^2(a_2 - 1)} I(a_1, a_2 - 1) + \frac{1}{p^2} I(a_1 - 1, a_2)$$



# Loop Amplitude



Number of Master Integrals

$$\mathcal{M} = \sum_i c_i J_i \quad i = \mathcal{O}(10^2)$$

# Integral Decomposition and Intersection Theory

Frellesvig, Gasparotto, **MKM**, Mastrolia, Mattiazzi, Mizera (2019)

Frellesvig, Gasparotto, Laporta, **MKM**, Mastrolia, Mattiazzi, Mizera (2019)

Frellesvig, Gasparotto, **MKM**, Mastrolia, Mattiazzi, Mizera (2020)

Chestnov, Frellesvig, Gasparotto, **MKM**, Mastrolia (2022)

Chestnov, Gasparotto, **MKM**, Mastrolia, Matsubara-Heo, Munch, Takayama (2022)

# Decomposition of Feynman Integral

$$I = \sum_{i=1}^{\nu} c_i J_i$$

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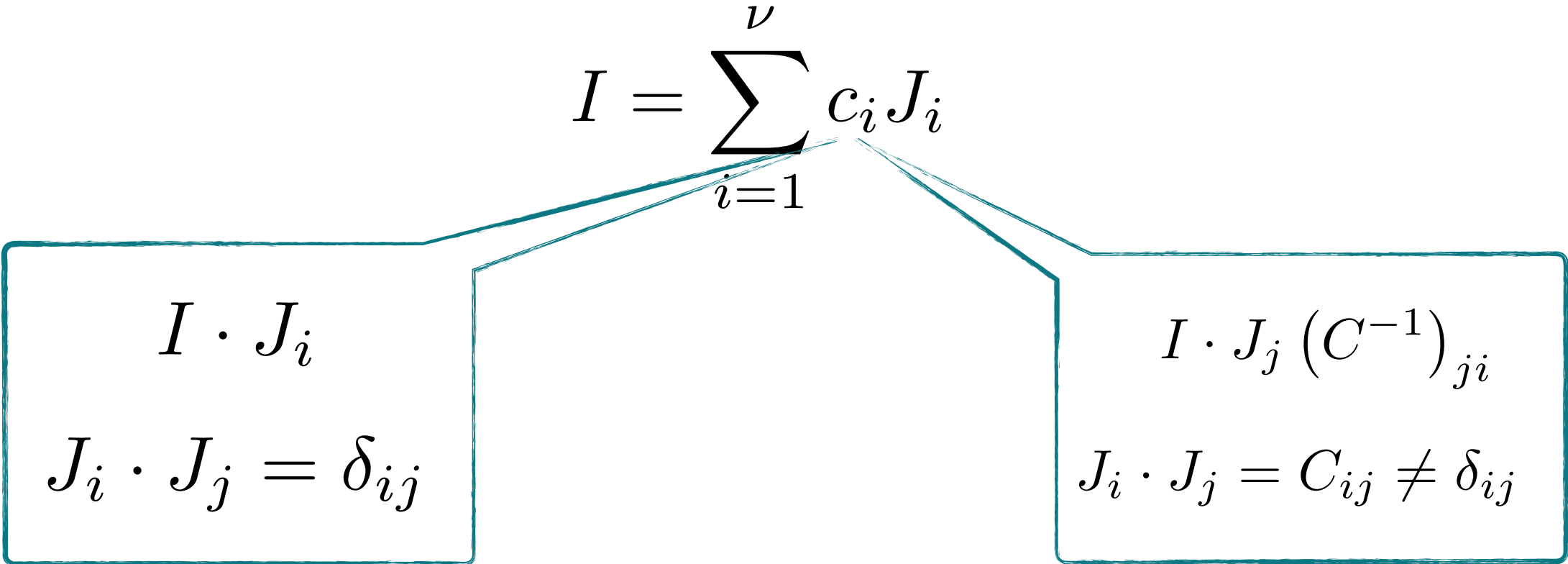
$$I \cdot J_i$$
$$J_i \cdot J_j = \delta_{ij}$$

$$I \cdot J_j (C^{-1})_{ji}$$
$$J_i \cdot J_j = C_{ij} \neq \delta_{ij}$$

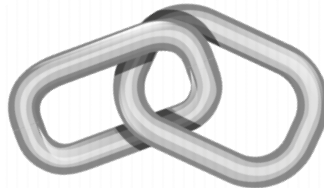


# Intersection Theory and Feynman Integral

Mastrolia, Mizera (2018)



**Intersection Theory**



**Feynman Integral**

# Intersection Theory and Feynman Integral

$$I = \sum_{i=1}^{\nu} c_i J_i$$

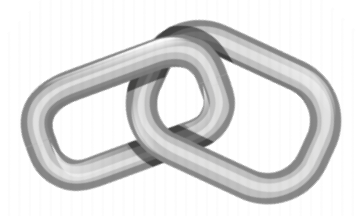
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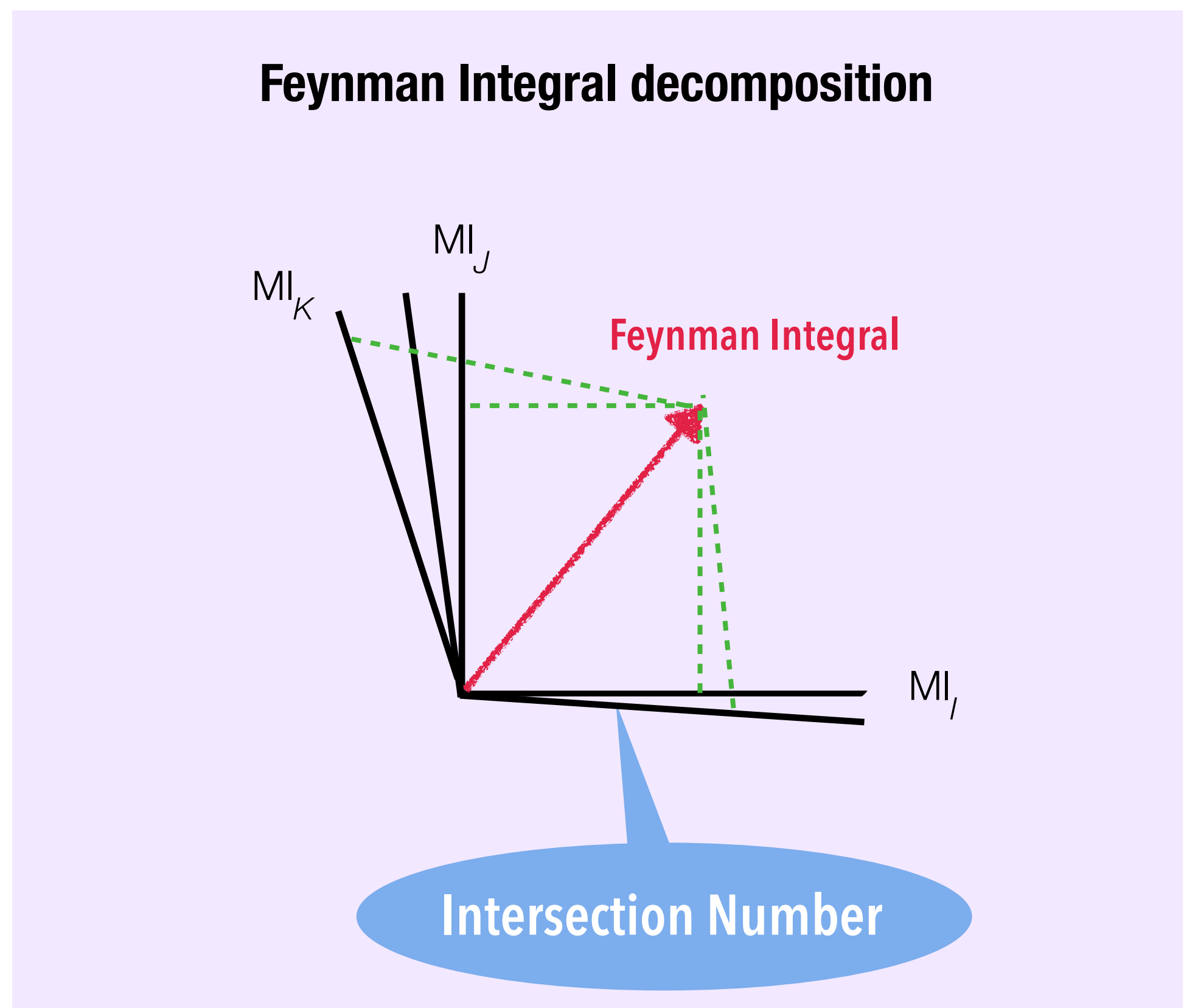
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Intersection Theory



Feynman Integral



Mastrolia, Mizera (2018)  
 Frellesvig, Gasparotto, MKM, Mastrolia, Mattiazzi, Mizera (2019)  
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 Chestnov, Gasparotto, MKM, Mastrolia, Matsubara-Heo, Munch, Takayama (2022)

**What is the Vector Space ?**

**How to define the scalar product ?**

# Intersection Theory

Aomoto, Gelfand, Kita, Cho, Matsumoto,  
Mimachi, Mizera, Yoshida

Single valued differential form

$$I = \int_{\mathcal{C}} z^b (1-z)^{c-b} (1-tz)^{-a} \frac{dz}{z}$$

Multi-valued Function

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Single valued differential form

$$I = \int_{\mathcal{C}} z^b (1-z)^{c-b} (1-tz)^{-a} \frac{dz}{z}$$

Multi-valued Function

$$\langle \varphi | \mathcal{C} \otimes u \rangle$$

Pairing

$$I = \int_{\mathcal{C}} u(\mathbf{z}) \overbrace{\varphi(\mathbf{z})}^{\text{Twisted Co-cycle}}$$

Twisted Cycle

$u(\mathbf{z})$  is a multi-valued function  
 $u(\mathbf{z})$  vanishes on the boundaries of  $\mathcal{C}$ ,  $u(\partial\mathcal{C}) = 0$

# Basics of Intersection Theory

$$0 = \int_{\mathcal{C}} d(u\xi) = \int_{\mathcal{C}} (du \wedge \xi + u d\xi) = \int_{\mathcal{C}} u \left( \frac{du}{u} \wedge + d \right) \xi \equiv \int_{\mathcal{C}} u \nabla_{\omega} \xi.$$

$$\omega \equiv d \log u$$

$$\nabla_{\omega} \equiv d + \omega \wedge$$

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**Equivalence Class**

$${}_{\omega}\langle \varphi | : \varphi \sim \varphi + \nabla_{\omega} \xi$$

$$\int_{\mathcal{C}} u \varphi = \int_{\mathcal{C}} u (\varphi + \nabla_{\omega} \xi)$$

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**Vector Space of n-forms**

$$H_\omega^n \equiv \{n\text{-forms } \varphi_n \mid \nabla_\omega \varphi_n = 0\} / \{\nabla_\omega \varphi_{n-1}\}$$

**Twisted Cohomology Group**

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Vector Space of n-forms

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Twisted Cohomology Group

Dual space

$$H_{-\omega}^n, \quad \nabla_{-\omega} = d - \omega \wedge$$



# Dimension of the Vector Space: Number of MIs

$$\chi(X) = \sum_{k=0}^{2n} (-1)^k \dim H_{\omega}^k. \quad H_{\omega}^{k \neq n} \text{ vanish.}$$

Aomoto (1975)

$$\begin{aligned} \nu &= (-1)^n \chi(X) \\ &= (-1)^n (n+1 - \chi(\mathcal{P}_{\omega})) \\ &= \{\text{number of solutions of } \omega=0\} \end{aligned}$$

# Decomposition of differential forms

Number of Linearly independent forms (twisted co-cycle) is  $\nu$

$$\text{Basis} \quad \langle e_i | \quad i = 1, 2, \dots, \nu$$

$$\text{Dual Basis} \quad |h_j\rangle \quad j = 1, 2, \dots, \nu$$

$$\text{Monomial Basis:} \quad \langle e_i | = \langle \phi_i | \equiv z^{i-1} dz$$

$$\text{d-Log Basis:} \quad \langle e_i | = \langle \varphi_i | \equiv \frac{dz}{z - z_i}$$

**Metric Matrix:**

$$C_{ij} = \langle e_i | h_j \rangle$$

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Metric Matrix:

$$\mathbf{C}_{ij} = \langle e_i | h_j \rangle$$

Master Decomposition Formula:

$$\langle \varphi | = \sum_{i,j=1}^{\nu} \langle \varphi | h_j \rangle (\mathbf{C}^{-1})_{ji} \langle e_i |$$

$$\mathbf{M} = \begin{pmatrix} \langle \varphi | \psi \rangle & \langle \varphi | h_1 \rangle & \langle \varphi | h_2 \rangle & \dots & \langle \varphi | h_\nu \rangle \\ \langle e_1 | \psi \rangle & \langle e_1 | h_1 \rangle & \langle e_1 | h_2 \rangle & \dots & \langle e_1 | h_\nu \rangle \\ \langle e_2 | \psi \rangle & \langle e_2 | h_1 \rangle & \langle e_2 | h_2 \rangle & \dots & \langle e_2 | h_\nu \rangle \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \langle e_\nu | \psi \rangle & \langle e_\nu | h_1 \rangle & \langle e_\nu | h_2 \rangle & \dots & \langle e_\nu | h_\nu \rangle \end{pmatrix} \equiv \begin{pmatrix} \langle \varphi | \psi \rangle & \mathbf{A}^\top \\ \mathbf{B} & \mathbf{C} \end{pmatrix}$$

$$\det \mathbf{M} = \det \mathbf{C} \left( \langle \varphi | \psi \rangle - \mathbf{A}^\top \mathbf{C}^{-1} \mathbf{B} \right) = 0$$

$$\langle \varphi | \psi \rangle = \mathbf{A}^\top \mathbf{C}^{-1} \mathbf{B}$$

$$= \sum_{i,j=1}^{\nu} \langle \varphi | h_j \rangle (\mathbf{C}^{-1})_{ji} \langle e_i | \psi \rangle$$

# Factorization of Identity

$$(i)(-i) = \mathbb{I}$$

Complex Number

$$\sum_n |n\rangle\langle n| = \mathbb{I}$$

Quantum Mechanics

$$\sum_{i,j=1}^{\nu} |e_j\rangle (C^{-1})_{ij} \langle e_i| = \mathbb{I}_c$$

Feynman Integral ?

$$\sum_{i,j=1}^{\nu} |c_j\rangle (H^{-1})_{ij} \langle c_i| = \mathbb{I}_h$$

# Decomposition of Uni-variate Integral

Mastrolia, Mizera (2018)

Frellesvig, Gasparotto, MKM, Mastrolia, Mattiazzi, Mizera (2019)

## Integrals

$$I = \int_{\mathcal{C}} u \varphi = \langle \varphi | \mathcal{C} \rangle$$

## Number of MIs

$$\omega \equiv d \log u$$

$\nu = \{\text{the number of solutions of } \omega = 0\}$

$$I = \sum_{i=1}^{\nu} c_i J_i \quad J_i = \langle e_i | \mathcal{C} \rangle$$

Matsumoto (1998)

## Computation of Intersection Number

Uni-Variate

## Choice of Bases / MIs

$$\langle e_i | = \langle \phi_i | \equiv z^{i-1} dz \quad \langle e_i | = \langle \varphi_i | \equiv \frac{dz}{z - z_i}$$

## Master Decomposition Formula

$$\langle \varphi | = \sum_{i,j=1}^{\nu} \langle \varphi | h_j \rangle (\mathbf{C}^{-1})_{ji} \langle e_i |$$

Metric Matrix

$$\mathbf{C}_{ij} = \langle e_i | h_j \rangle$$

# Decomposition of Multi-Variate Integral

Frellesvig, Gasparotto, MKM, Mastrolia, Mattiazzi, Mizera (2019)

Frellesvig, Gasparotto, Laporta, MKM, Mastrolia, Mattiazzi, Mizera (2020)

## Integrals

$$I = \int_{\mathcal{C}} u \varphi = \langle \varphi | \mathcal{C} \rangle$$

## Number of MIs

$$\omega \equiv d \log u(\mathbf{z}) = \sum_{i=1}^n \hat{\omega}_i dz_i$$

$\nu$  = Number of solutions of the system of equations

$$\hat{\omega}_i \equiv \partial_{z_i} \log u(\mathbf{z}) = 0, \quad i = 1, \dots, n$$

$$I = \sum_{i=1}^{\nu} c_i J_i \quad J_i = \langle e_i | \mathcal{C} \rangle$$

## Choice of Bases

$$e_i(\mathbf{z}) \quad h_i(\mathbf{z})$$

$$\mathbf{C}_{ij} = \langle e_i | h_j \rangle$$

$$\langle \varphi | = \sum_{i,j=1}^{\nu} \langle \varphi | h_j \rangle (\mathbf{C}^{-1})_{ji} \langle e_i |$$

*Metric Matrix*

*Master Decomposition Formula*

Computation of Intersection Number

Multi-Variate

# Computation of Intersection Number

Fibration Method	Matsumoto (1998)
	Goto (2015)
	Mizera (2019) Frellesvig, Gasparotto, Laporta, MKM, Mastrolia, Mattiazzi, Mizera (2019)
	Wienzierl (2020) Frellesvig, Gasparotto, MKM, Mastrolia, Mattiazzi, Mizera (2020)
	Caron-Huot, Pokraka (2021)
Secondary Equation	Matsubara-Heo (2019)
	Chestnov, Gasparotto, MKM, Mastrolia, Matsubara-Heo, Munch, Takayama (2022)
Multivariate Differential Equation	Matsumoto (1998)
	Chestnov, Frellesvig, Gasparotto, MKM, Mastrolia (2022)

# Univariate Intersection Number

Matsumoto, Mizera

$$\langle \varphi_L | \varphi_R \rangle = \frac{1}{2\pi i} \int_X \varphi_L \wedge \varphi_R$$



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Matsumoto, Mizera

$$\langle \varphi_L | \varphi_R \rangle = \frac{1}{2\pi i} \int_X \varphi_L \wedge \varphi_R$$

$$\langle \varphi_L | \varphi_R \rangle_\omega = \sum_{p \in \mathcal{P}} \operatorname{Res}_{z=p} (\psi_p \varphi_R)$$

$$\nabla_{\omega_p} \psi_p = \varphi_{L,p}$$

**First Order Differential Equation**

# Univariate Intersection Number

Matsumoto, Mizera

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**First Order Differential Equation**

$$\nabla_{\omega_p} \psi_p = \varphi_{L,p}$$

**Laurent Expansion around the poles of  $\omega$**

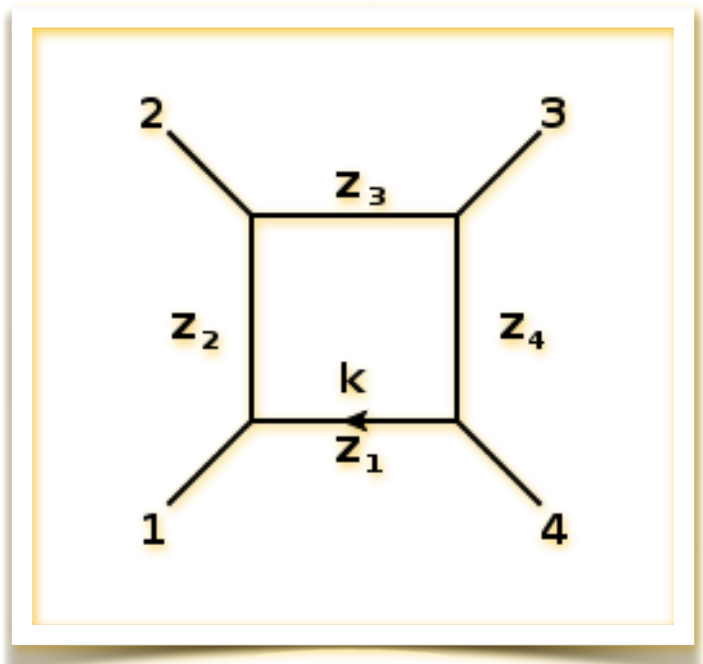
$$\tau \equiv z - p$$

**Known :**  $\varphi_{L,p}$

$$\text{Ansatz : } \psi_p = \sum_{j=\min}^{\max} \psi_p^{(j)} \tau^j + \mathcal{O}(\tau^{\max+1})$$

The coefficients are obtained by solving the differential equation

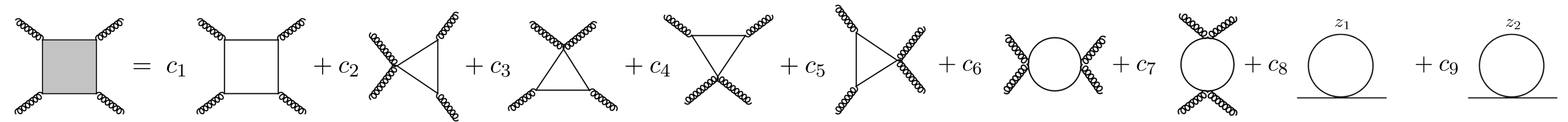
# Examples of decomposition

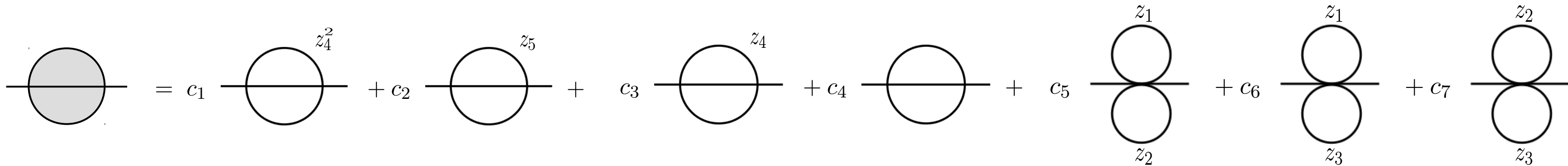


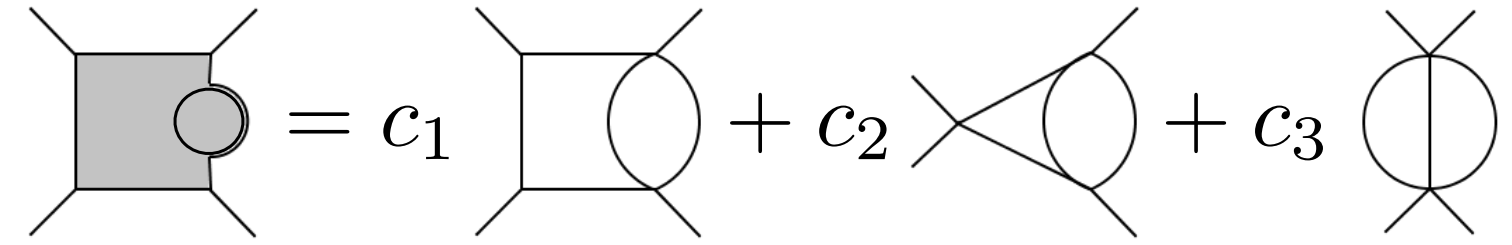
$$\text{shaded square} = c_1 \text{square} + c_2 \text{circle} + c_3 \text{circle}$$

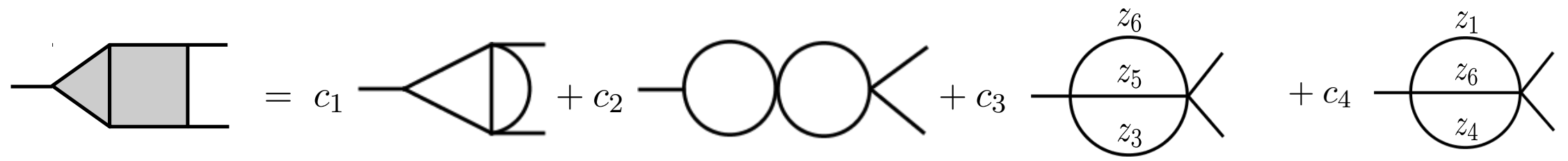
$$(c_1, c_2, c_3) = \left( \langle \text{shaded square} | \text{square} \rangle \langle \text{shaded square} | \text{circle} \rangle \langle \text{shaded square} | \text{circle} \rangle \right) \begin{pmatrix} \langle \text{square} | \text{square} \rangle & \langle \text{square} | \text{circle} \rangle & \langle \text{square} | \text{circle} \rangle \\ \langle \text{circle} | \text{square} \rangle & \langle \text{circle} | \text{circle} \rangle & \langle \text{circle} | \text{circle} \rangle \\ \langle \text{circle} | \text{square} \rangle & \langle \text{circle} | \text{circle} \rangle & \langle \text{circle} | \text{circle} \rangle \end{pmatrix}^{-1}$$

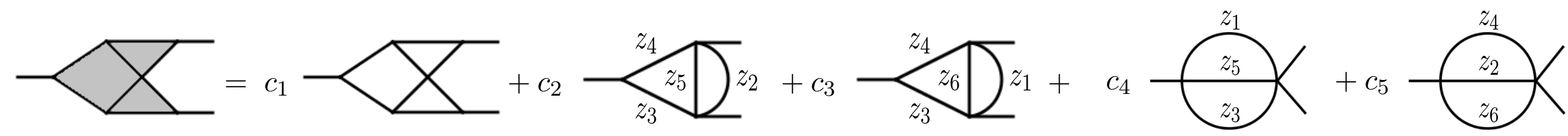
# Further Applications

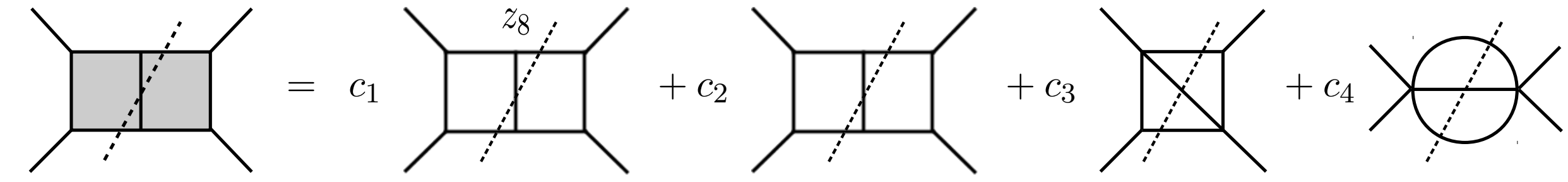












# Gravitational Wave Observables

**MKM, Mastrolia, Patil, Steinhoff (2022)**

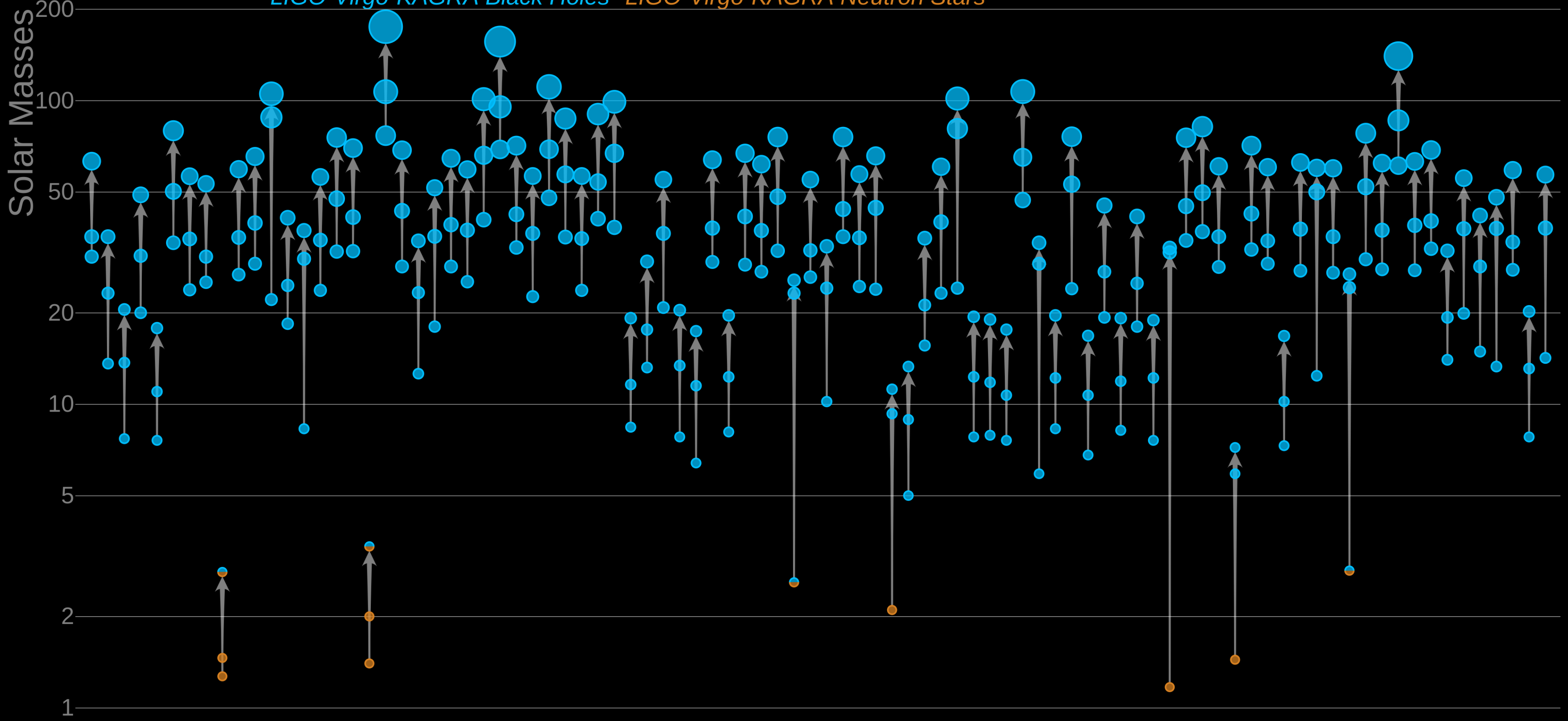
**MKM, Mastrolia, Patil, Steinhoff (2022)**

**MKM, Mastrolia, O Silva, Patil, Steinhoff (2023)**

# GW observations

## Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars

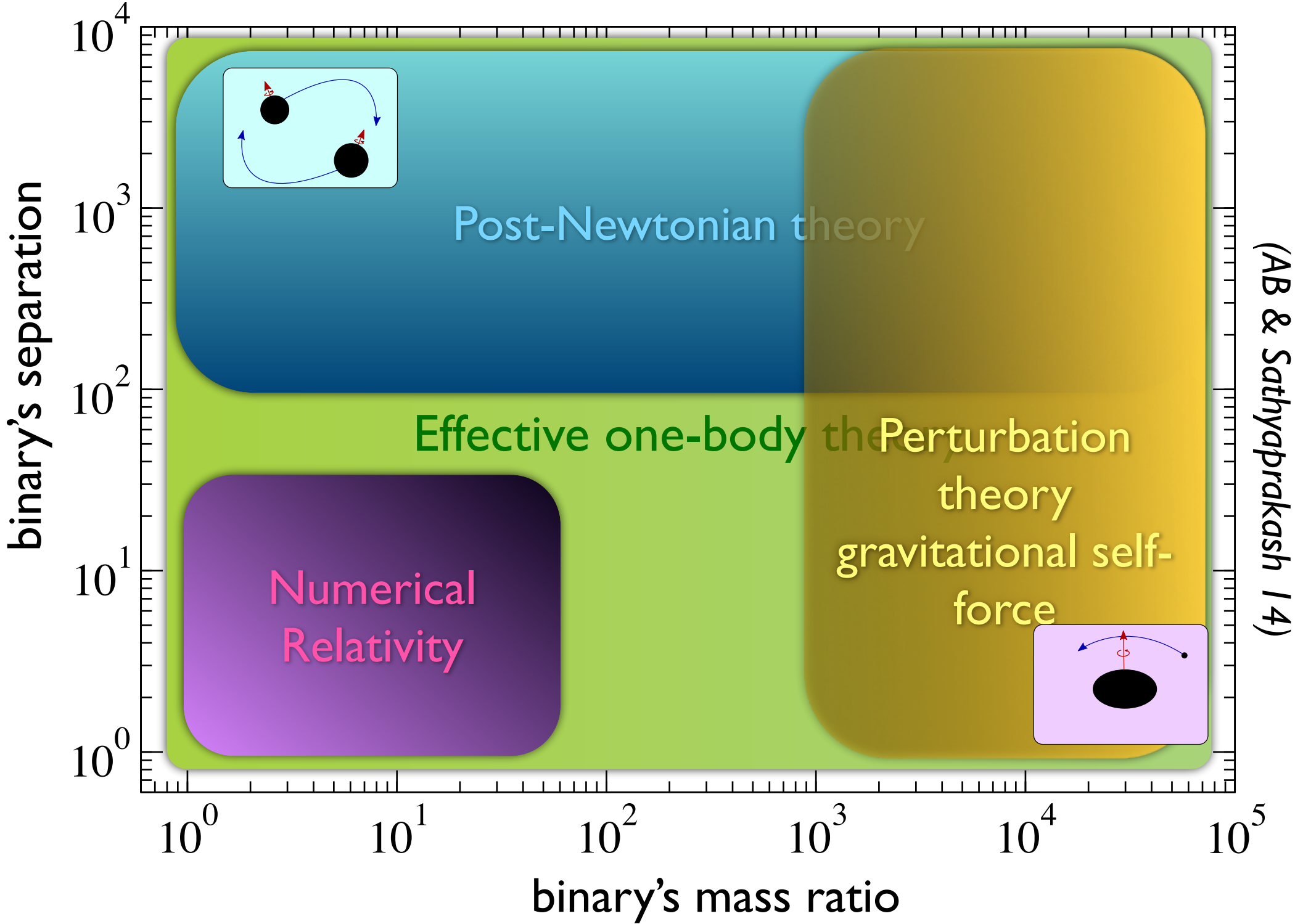
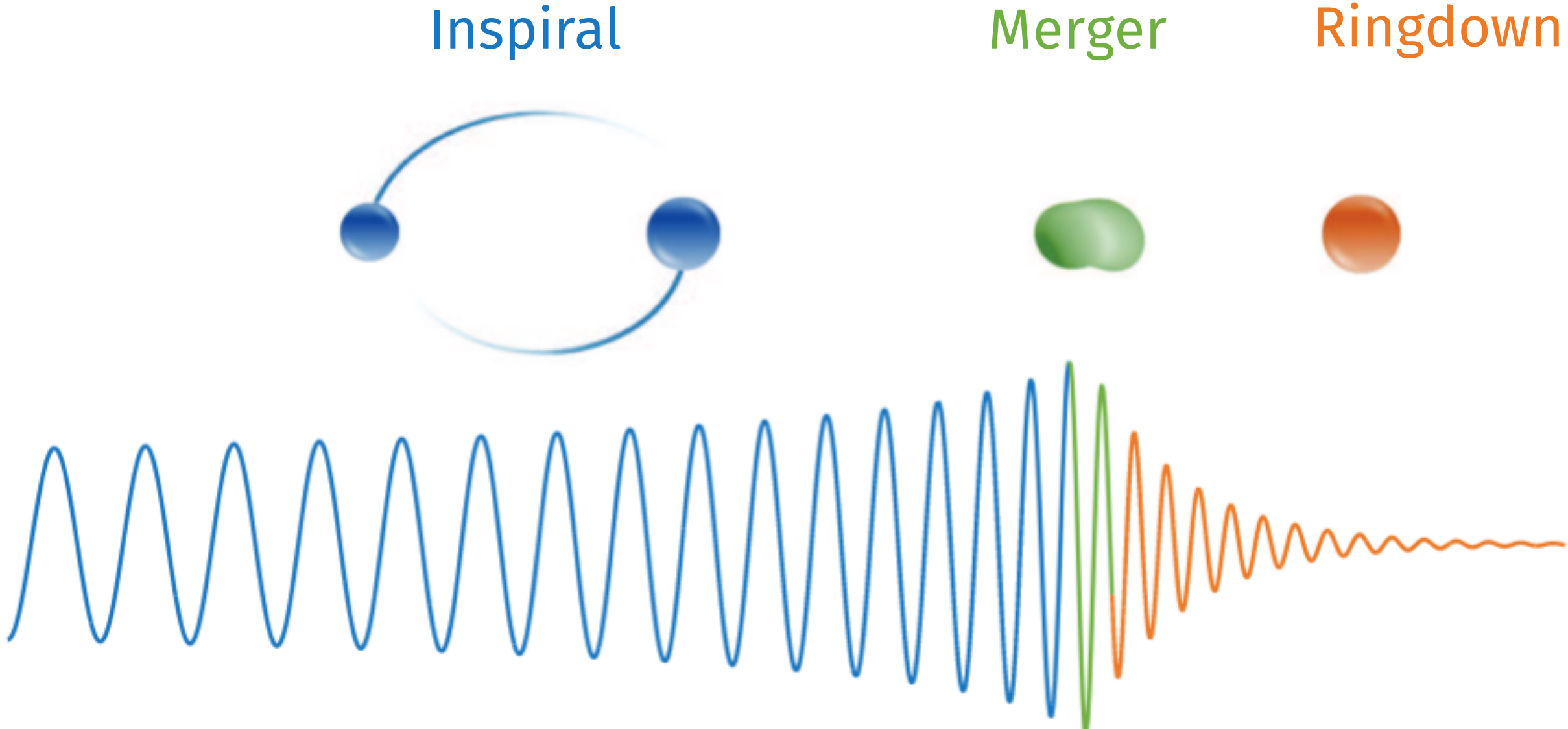


### Tasks

- 👤 Supplement conventional Analysis
- 👤 Increase Theoretical Precision
- 👤 Perform Gravity phenomenology

# Solving two-body problem in GR

Antelis, moreno (2016)



Post-Newtonian (PN)

Numerical Relativity

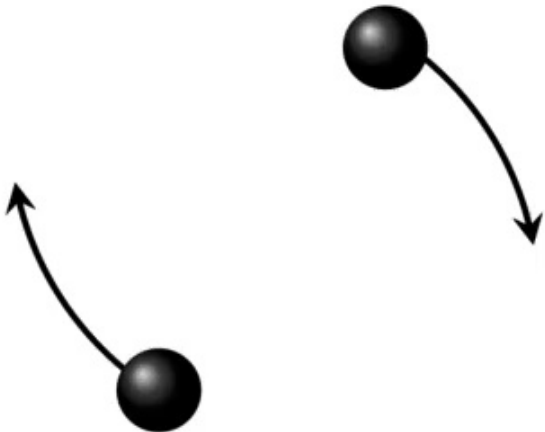
Perturbation Theory

Post-Minkowskian (PM)

# Analytical Approximation Methods

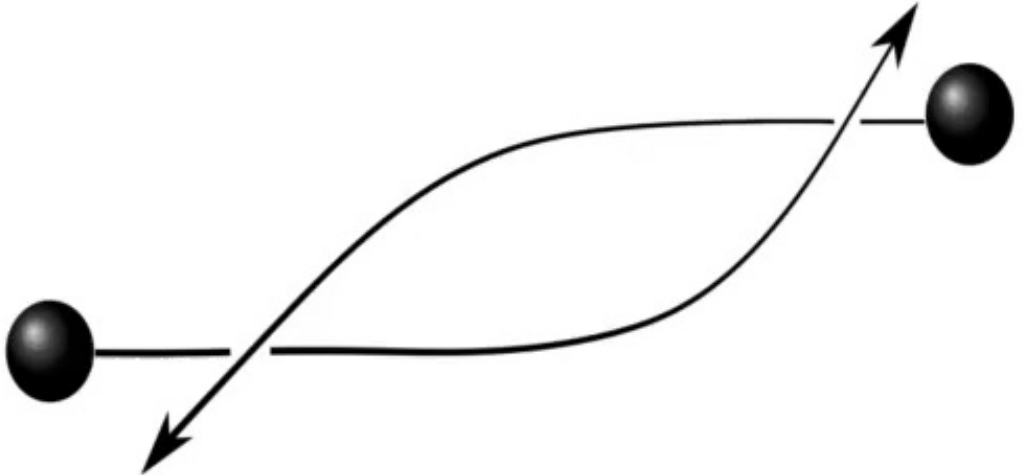
## Post-Newtonian (PN)

$$\frac{v^2}{c^2} \sim \frac{GM}{rc^2} \ll 1$$



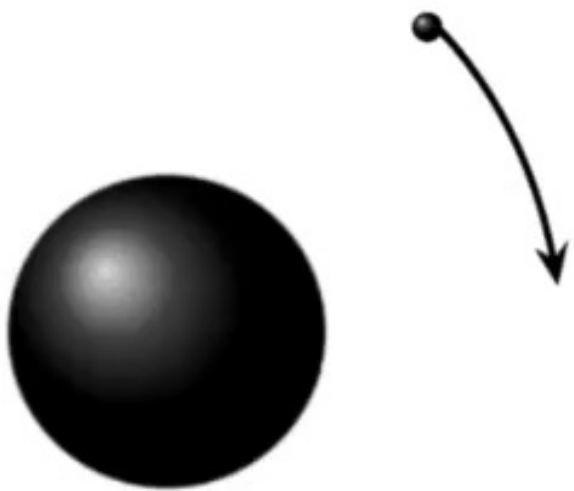
## Post-Minkowskian (PM)

$$\frac{GM}{rc^2} \ll 1$$

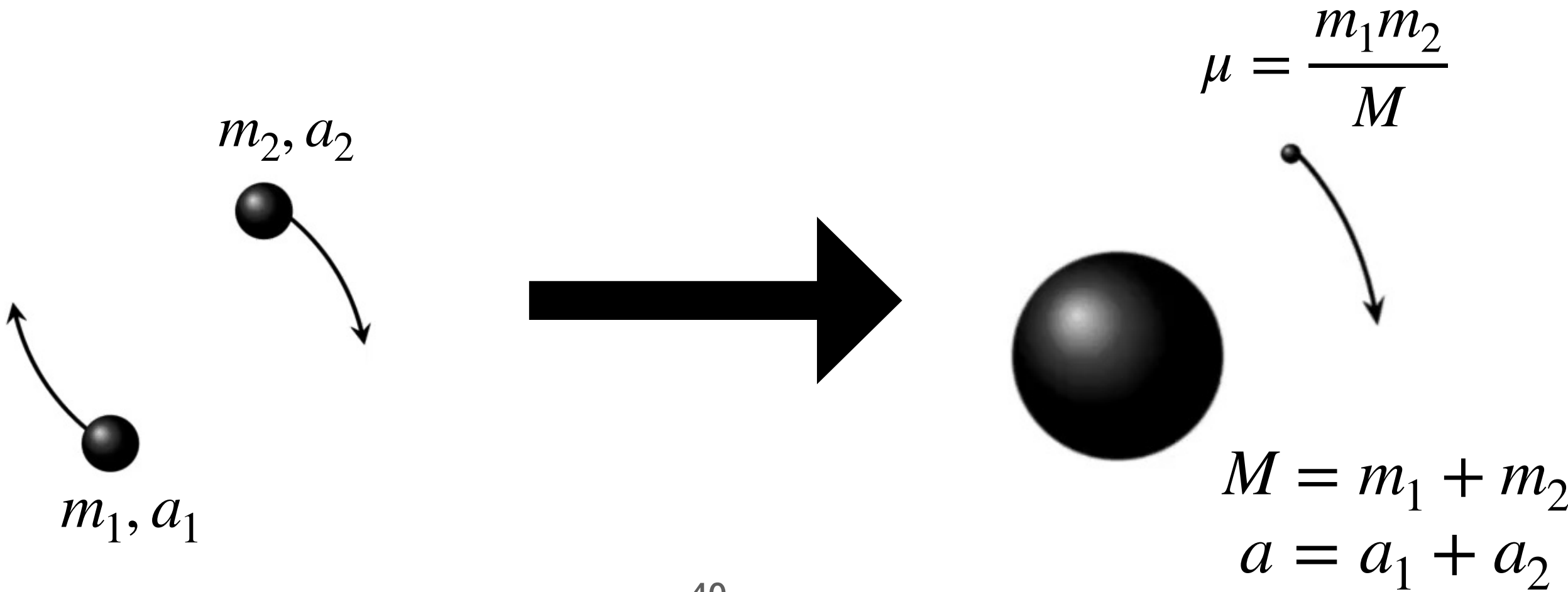


## Self-Force (SF)

$$\frac{m_1}{m_2} \ll 1$$

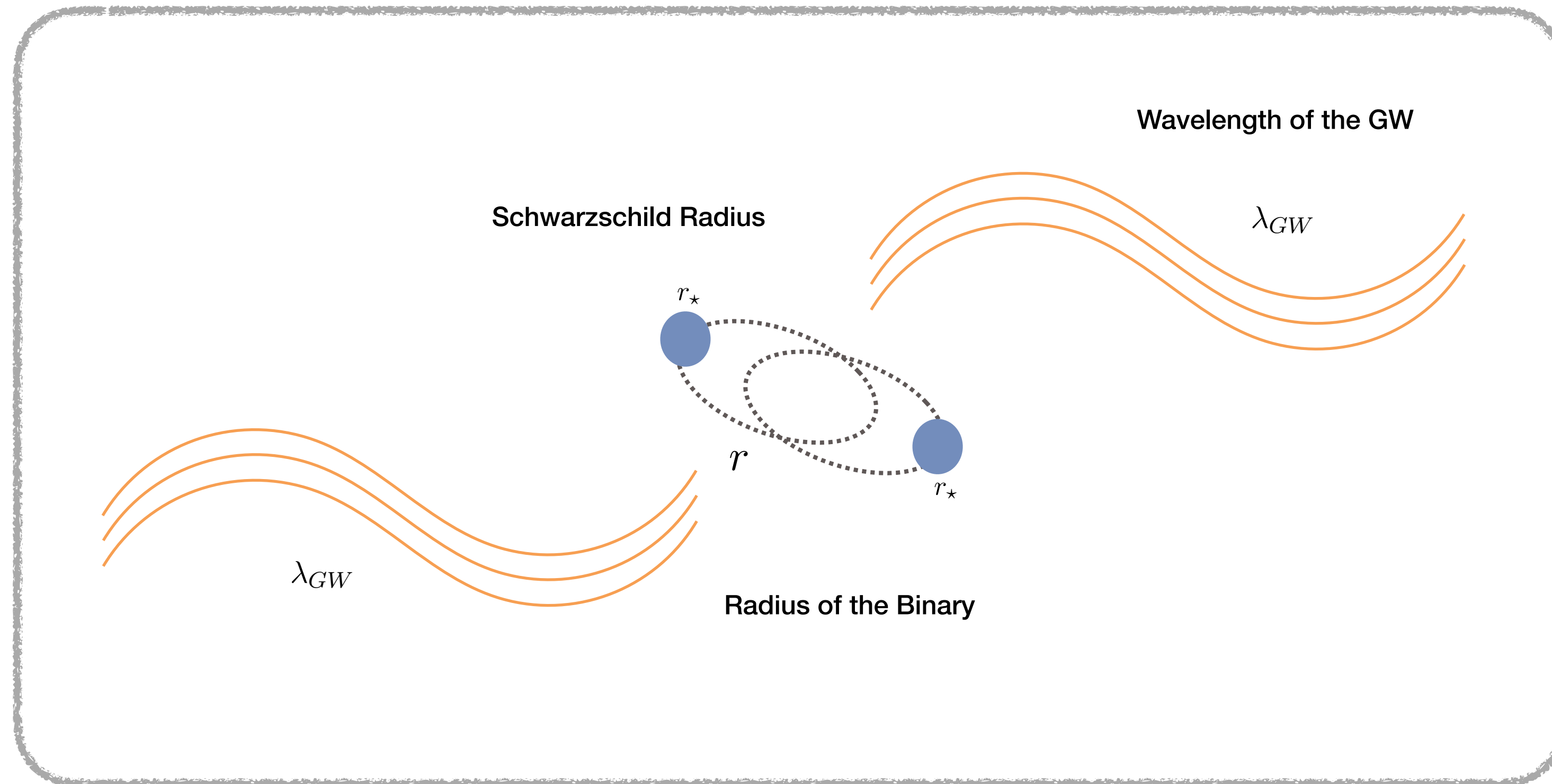


## Effective One-Body (EOB)





# Post-Newtonian Expansion EFT set up



## Equations of Motion

$$\begin{aligned} \dot{r} &= \frac{d\mathcal{H}}{dp_r} & \dot{p}_r &= -\frac{d\mathcal{H}}{dr} + \mathcal{F}_r \\ \dot{\phi} &= \frac{d\mathcal{H}}{dp_\phi} & \dot{p}_\phi &= -\frac{d\mathcal{H}}{d\phi} + \mathcal{F}_\phi \end{aligned}$$

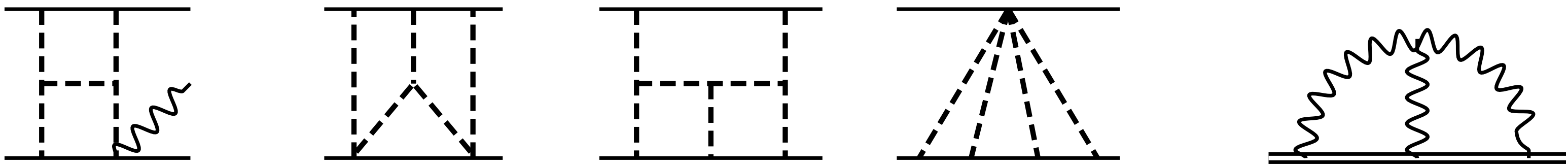
Need:

Hamiltonian  $\mathcal{H}$

Radiation Reaction  $\mathcal{F}$

# Advantage of QFT techniques

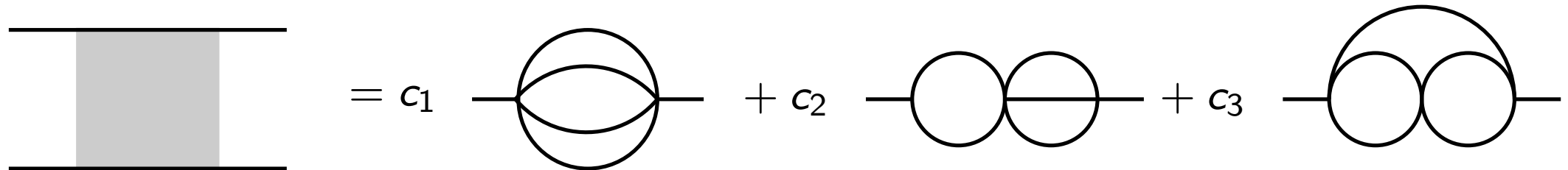
**Use of Feynman diagrams**



**Dimensional regularization**

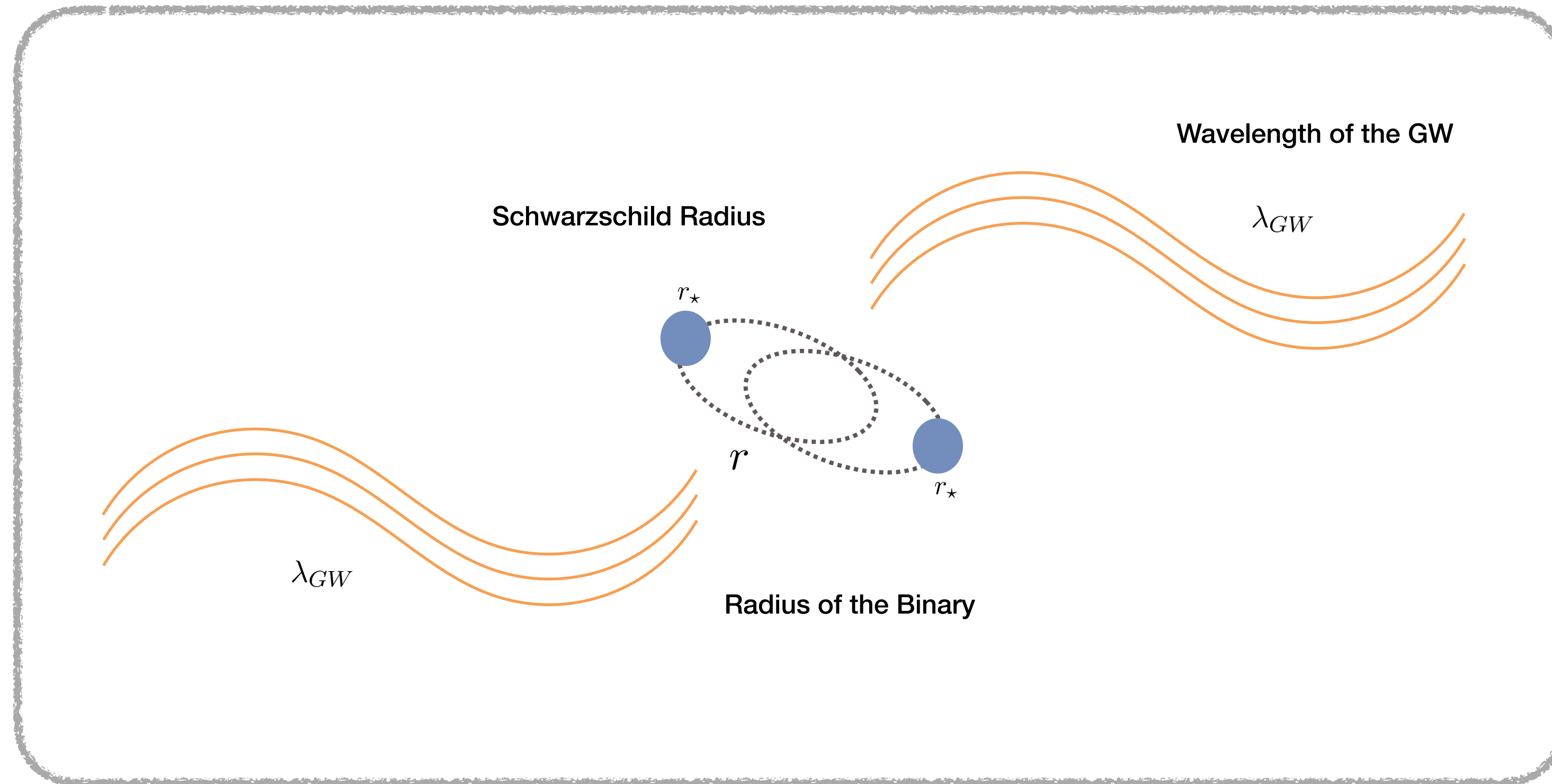
Better to handle spurious divergences

**Multi-loop Techniques**



- IBP relations
- Differential Equations

# Post-Newtonian Expansion EFT set up



Hierarchy of scales

$$r_* \ll r \ll \lambda_{GW}$$

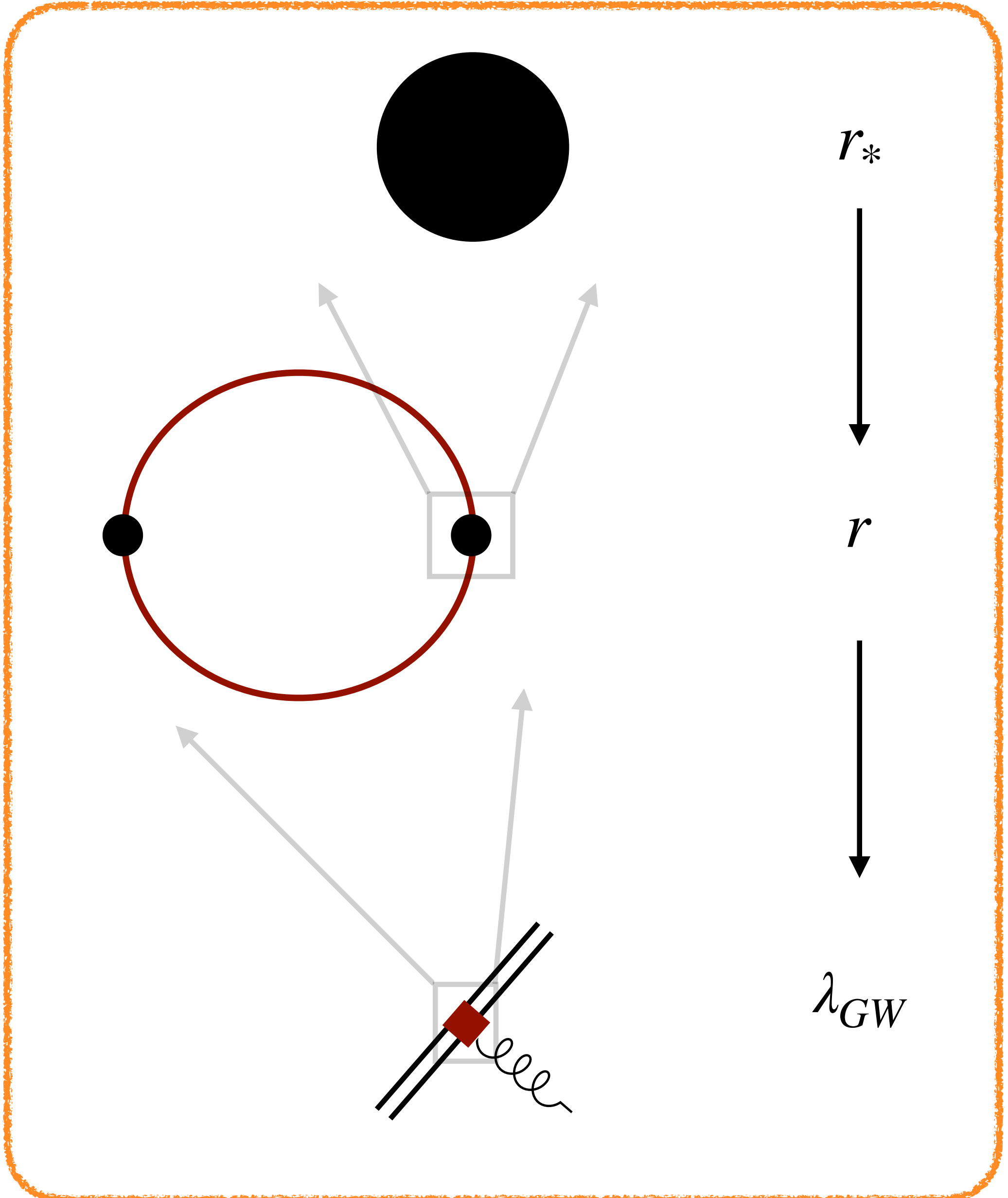
# Post-Newtonian Expansion EFT set up

Hierarchy of scales  
 $r_* \ll r \ll \lambda_{GW}$

## Tower of EFTs

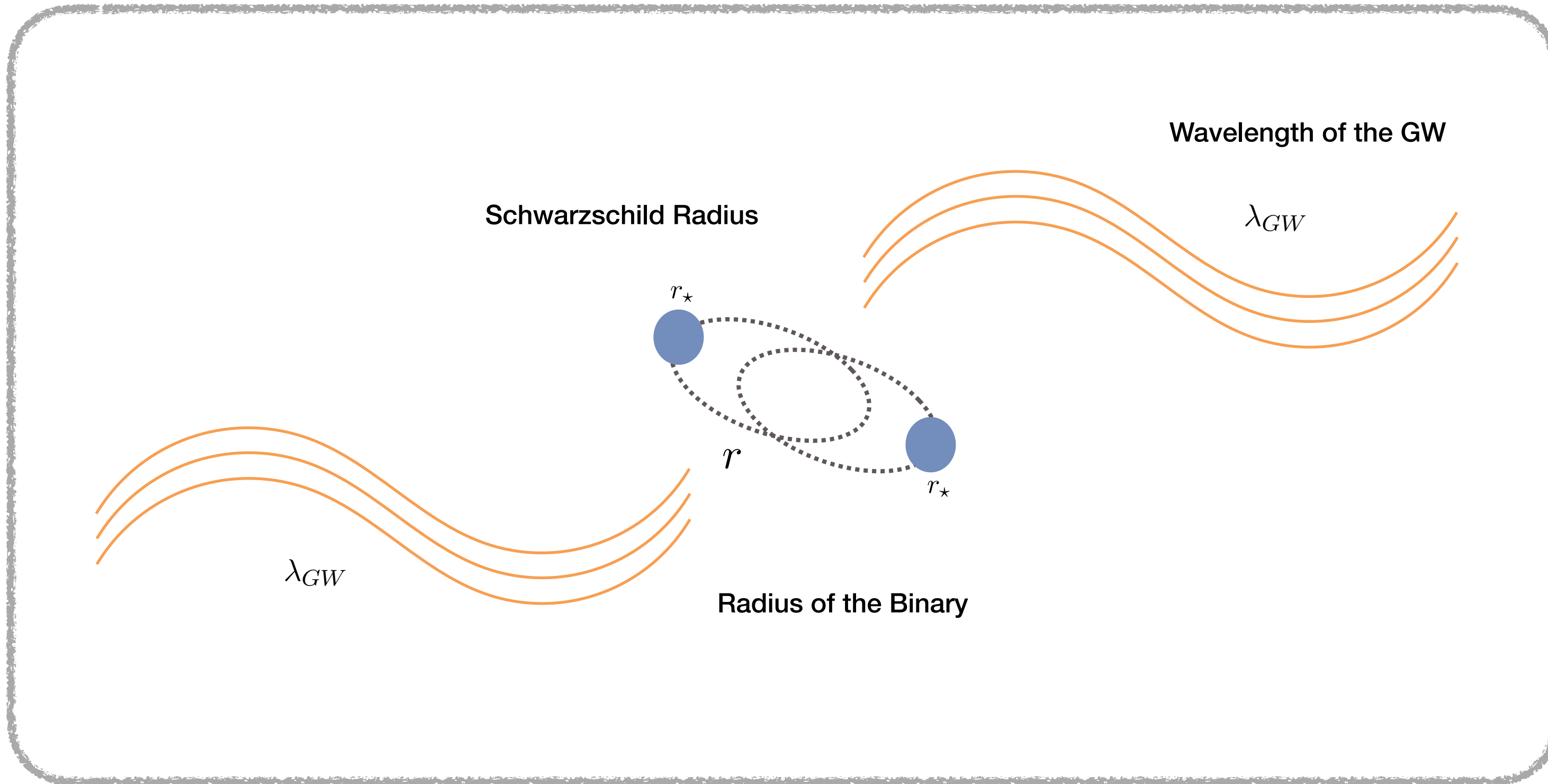
Goldberger, Rothstein

- 1. One-Particle EFT for Compact Object
- 2. EFT of Composite Particle for Binary
- 3. Effective Theory of Dynamical Multipoles



# Post-Newtonian Expansion EFT set up

Goldberger, Rothstein



Hierarchy of scales

$$r_* \ll r \ll \lambda_{GW}$$

Tower of EFTs

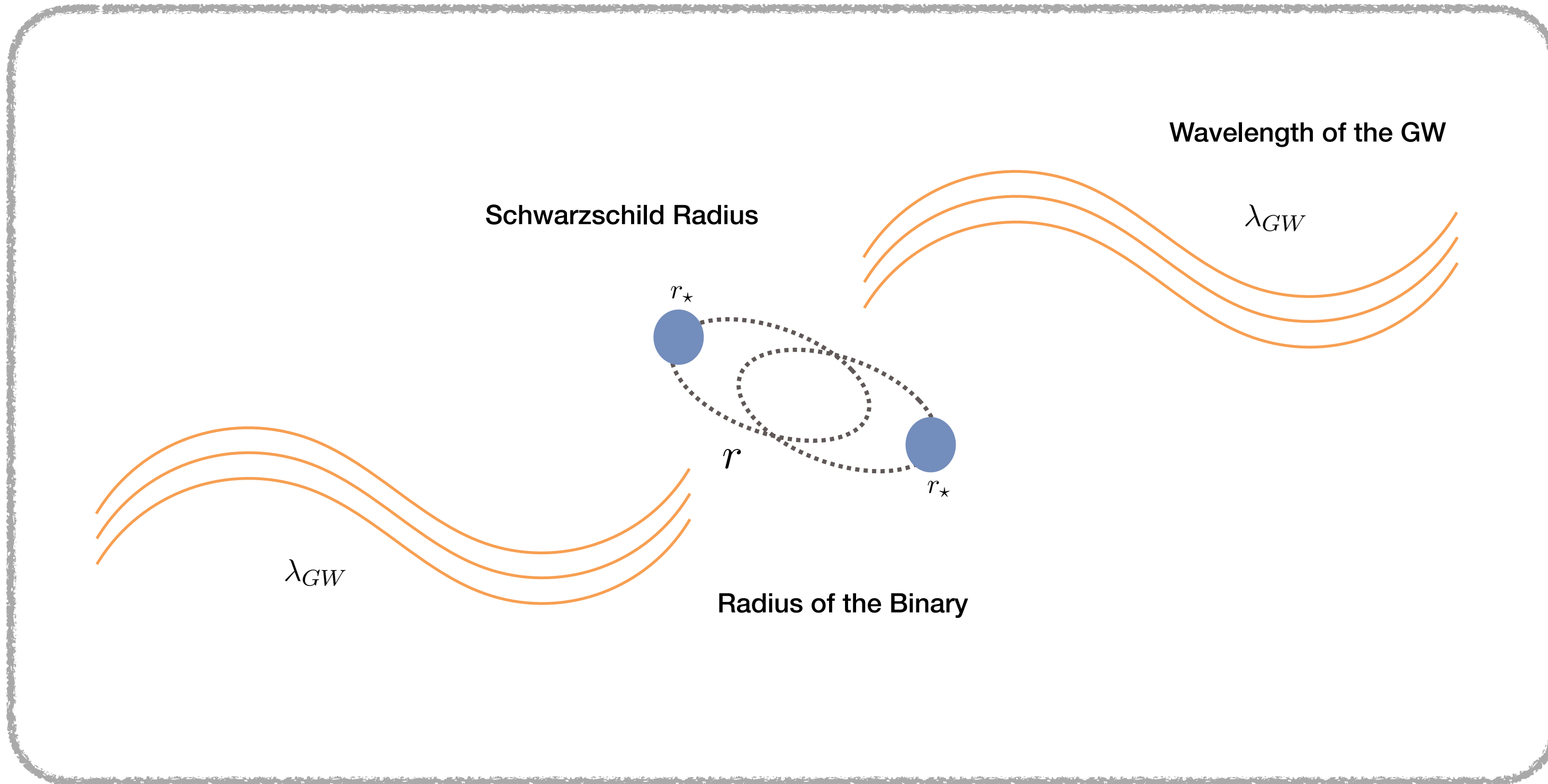
1. One-Particle EFT for Compact Object

$$S[g_{\mu\nu}] = -\frac{1}{16\pi G} \int d^4x \sqrt{g} R$$

$$S_{pp}[g_{\mu\nu}] = -m \int d\sigma \sqrt{u^2}$$

# Post-Newtonian Expansion EFT set up

Goldberger, Rothstein



Hierarchy of scales

$$r_* \ll r \ll \lambda_{GW}$$

Tower of EFTs

2. EFT of Composite Particle for Binary

$$S[g_{\mu\nu}] = -\frac{1}{16\pi G} \int d^4x \sqrt{g} R$$

$$g_{\mu\nu} = \eta_{\mu\nu} + H_{\mu\nu} + h_{\mu\nu}$$

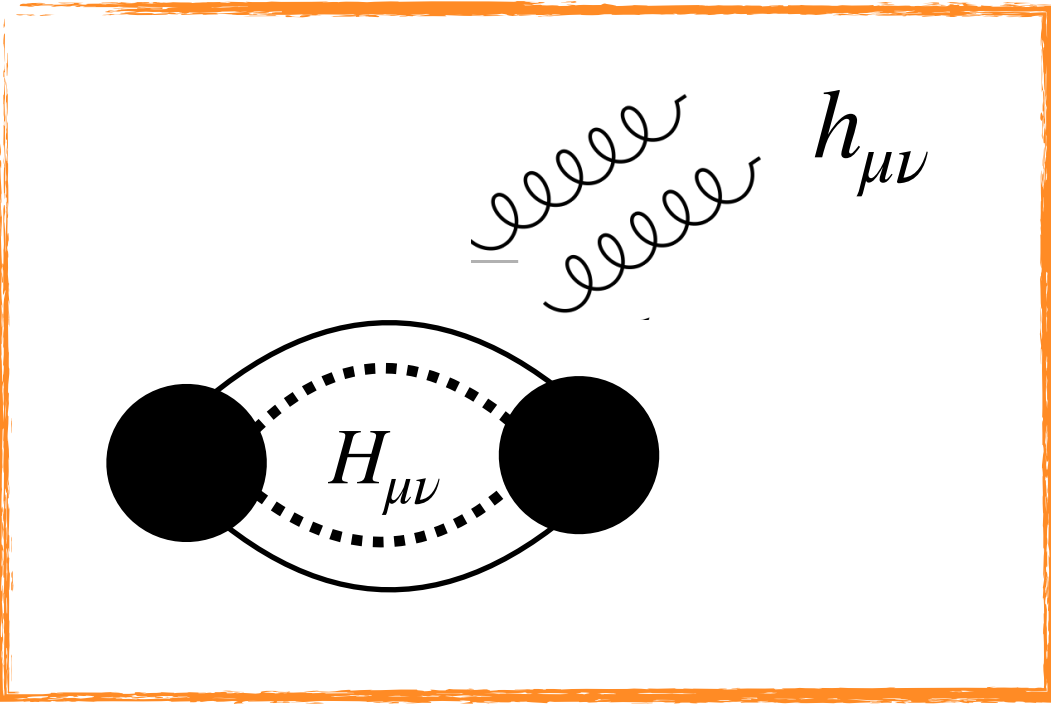
$$S_{pp}[g_{\mu\nu}, x_K] = \sum_{K=1}^2 -m_K \int d\sigma \sqrt{u_K^2}$$

Method of Regions

potential gravitons  $H_{\mu\nu}$  with scaling  $(k_0, \mathbf{k}) \sim (v/r, 1/r)$

radiation gravitons  $\bar{h}_{\mu\nu}$  with scaling  $(k_0, \mathbf{k}) \sim (v/r, v/r)$

# EFT at the orbital scale: Conservative Dynamics



$$e^{i S_{eff}[x_K]} = \int \mathcal{D}\bar{h}_{\mu\nu} \int \mathcal{D}H_{\mu\nu} \exp \left\{ iS[\eta + \bar{h} + H] + i \sum_{K=1}^2 S_{pp}[x_K(t), \eta + \bar{h} + H] \right\}$$

Effective Action for Dynamical Multipoles

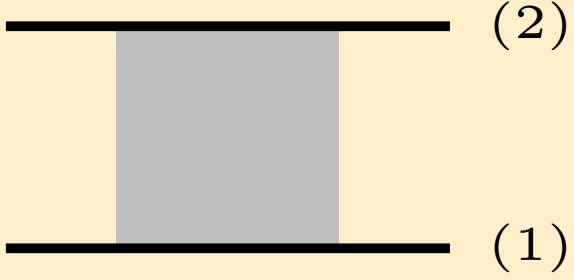
$$e^{i S_{eff}[x_K]} = \int \mathcal{D}\bar{h}_{\mu\nu} \exp \left\{ iS[\eta + \bar{h}] + \text{diagrams} + \dots \right\}$$

Conservative Dynamics

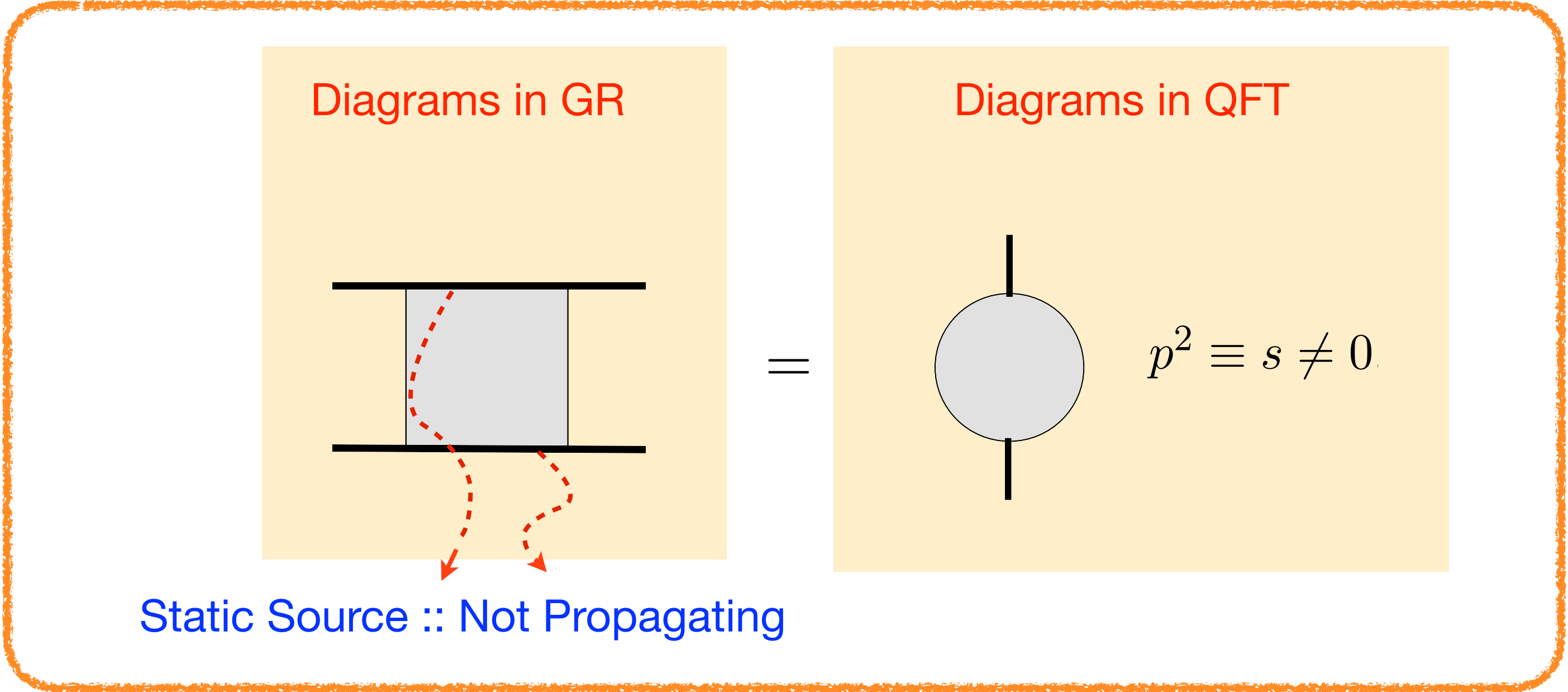
$$\int \mathcal{D}H \exp \left\{ iS[\eta + H, h=0] + iS_{pp}[x_K, \eta + H, h=0] \right\} = e^{i S_{eff}[h=0, x_K]} = e^{i \int dt \mathcal{L}_{eff}}$$

# Potential for the 2-body system

Goldberger, Rothstein, Porto, Levi, ...  
 Foffa, Sturani, Sturm, Mastrolia (2016)

$$\mathcal{V}_{\text{eff}} = \mathbf{i} \lim_{d \rightarrow 3} \int_{\mathbf{p}} e^{\mathbf{i}\mathbf{p} \cdot (\mathbf{x}_{(1)} - \mathbf{x}_{(2)})}$$


## Key Observation





# Status of PN Results

	PN order		1.5	2.5	3.5	4.5	5.5	6.5
	0	1	2	3	4	5	6	
no spin	N	1PN	2PN	3PN	4PN	5PN	6PN	
spin-orbit		LO SO	NLO SO	N2LO SO	N3LO SO	N4LO SO		
spin <sup>2</sup>			LO S2	NLO S2	N2LO S2	N3LO S2		
spin <sup>3</sup>				LO S3	NLO S3	N2LO S3		
spin <sup>4</sup>					LO S4	NLO S4		
spin <sup>5</sup>						LO S5	NLO S5	
spin <sup>6</sup>							LO S6	

need up to

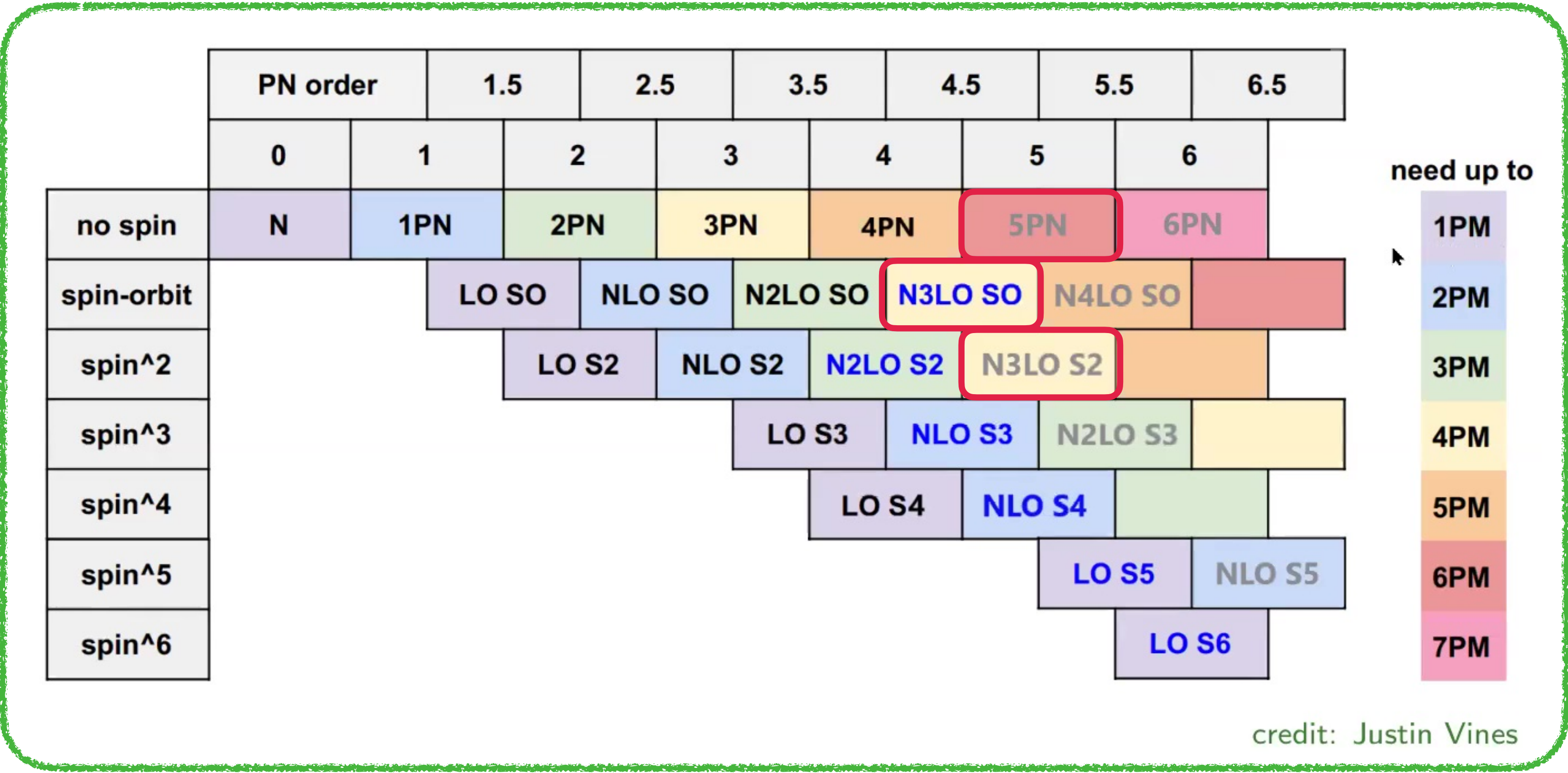
- 1PM
- 2PM
- 3PM
- 4PM
- 5PM
- 6PM
- 7PM

credit: Justin Vines

- Levi, McLeod, Steinhoff, Teng, Von Hippel,...
- Kim, Levi, Yin (2021)
- Kim, Levi, Yin (2022)
- MKM, Mastrolia, Patil, Steinhoff (2022)
- Levi, Yin (2022)
- MKM, Mastrolia, Patil, Steinhoff (2022)

- 1PN [Einstein, Infeld, Hoffman '38].
- 2PN [Ohta *et al.*, '73].
- 3PN [Jaranowski, Schaefer, '97; Damour, Jaranowski, Schaefer, '97; Blanchet, Faye, '00; Damour, Jaranowski, Schaefer, '01]
- 4PN [Damour, Jaranowski, Schäfer, Bernard, Blanchet, Bohe, Faye, Marsat, Marchand, Foffa, Sturani, Mastrolia, Sturm, Porto, Rothstein...]
- 5PN [Foffa, Mastrolia, Sturani, Sturm, Bodabilla, '19; Blümlein, Maier, Marquard, '19; Bini, Damour, Geralico, '19; Blümlein, Maier, Marquard, '19; Almeida, Foffa, Sturani, '22;]

# Status of PN Results



Levi, McLeod, Steinhoff, Teng, Von Hippel, ..

Kim, Levi, Yin (2021)

Kim, Levi, Yin (2022)

MKM, Mastrolia, Patil, Steinhoff (2022)

Levi, Yin (2022)

MKM, Mastrolia, Patil, Steinhoff (2022)

Brunello, MKM, Mastrolia, Patil (W.I.P)

1PN [Einstein, Infeld, Hoffman '38].

2PN [Ohta et al., '73].

3PN [Jaranowski, Schaefer, '97; Damour, Jaranowski, Schaefer, '97; Blanchet, Faye, '00; Damour, Jaranowski, Schaefer, '01]

4PN [Damour, Jaranowski, Schäfer, Bernard, Blanchet, Bohe, Faye, Marsat, Marchand, Foffa, Sturani, Mastrolia, Sturm, Porto, Rothstein...]

5PN [Foffa, Mastrolia, Sturani, Sturm, Bodabilla, '19; Blümlein, Maier, Marquard, '19; Bini, Damour, Geralico, '19; Blümlein, Maier, Marquard, '19; Almeida, Foffa, Sturani, '22;]

# EFT of Spinning Objects

Levi, Steinhoff (2015)

$$S_{\text{EH}} = -\frac{c^4}{16\pi G_N} \int d^4x \sqrt{g} R[g_{\mu\nu}] + \frac{c^4}{32\pi G_N} \int d^4x \sqrt{g} g_{\mu\nu} \Gamma^\mu \Gamma^\nu$$

$$\Omega_{(a)}^{\mu\nu} = \Lambda_{(a)A}^\mu \frac{d\Lambda_{(a)}^{A\nu}}{d\tau}$$

$$S_{\text{pp}} = \sum_{a=1,2} \int d\tau \left( -m_{(a)} c \sqrt{u_{(a)}^2} - \frac{1}{2} S_{(a)\mu\nu} \Omega_{(a)}^{\mu\nu} - \frac{S_{(a)\mu\nu} u_{(a)}^\nu}{u_{(a)}^2} \frac{du_{(a)}^\mu}{d\tau} + \mathcal{L}_{(a)}^{(R)} + \mathcal{L}_{(a)}^{(R^2)} + \dots \right)$$

$$S_{(a)\mu\nu} = -2 \frac{\partial L_{\text{pp}}}{\partial \Omega_{(a)}^{\mu\nu}}$$

$$\mathcal{L}_{(a)}^{(R)} = -\frac{1}{2m_{(a)}c} \left( C_{\text{ES}^2}^{(0)} \right)_{(a)} \frac{E_{\mu\nu}}{u_{(a)}} \left[ S_{(a)}^\mu S_{(a)}^\nu \right]_{\text{STF}} + \dots$$

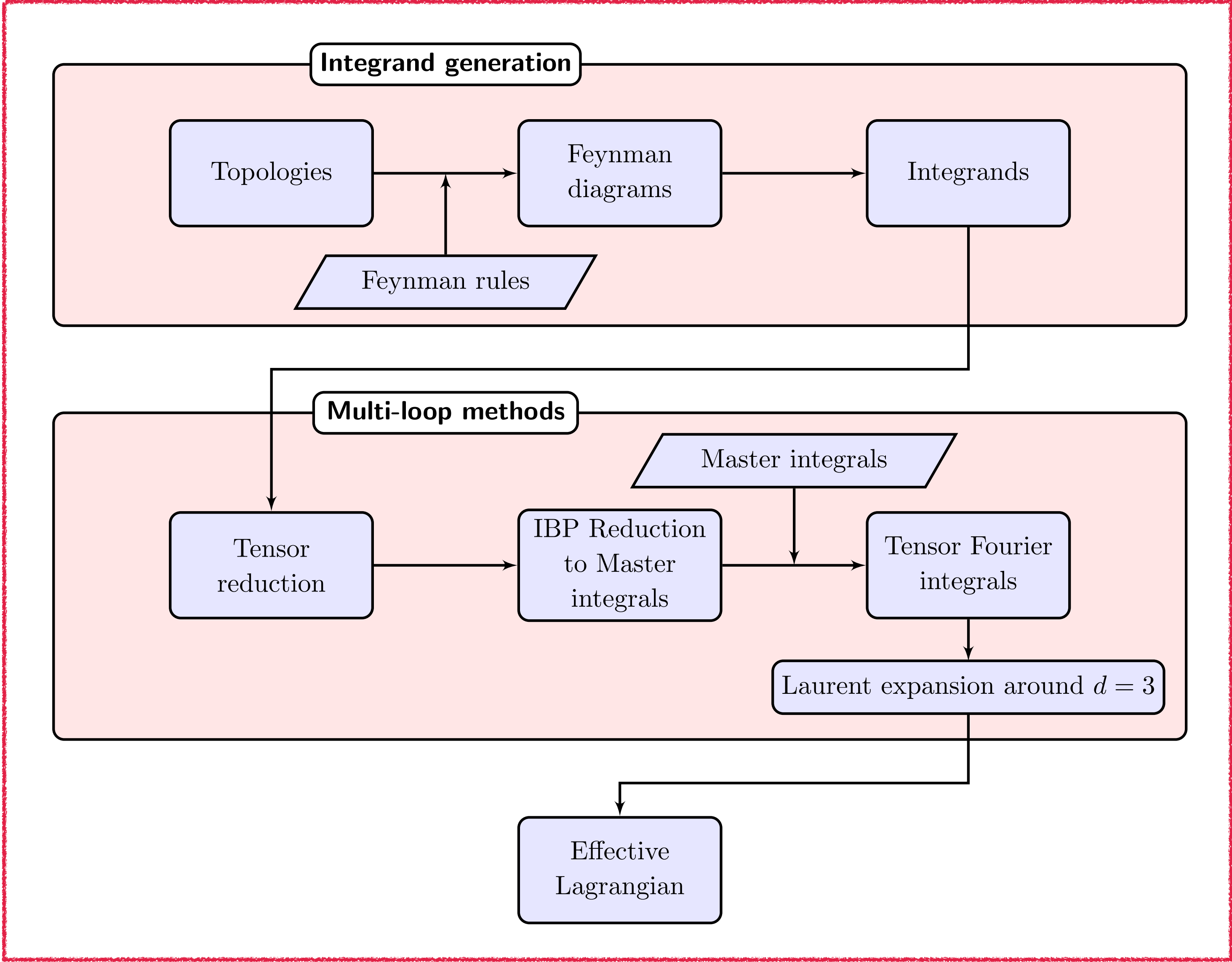
$$\mathcal{L}_{(a)}^{(R^2, S^0)} = \frac{1}{2} \left( C_{\text{E}^2}^{(2)} \right)_{(a)} \frac{G_N^2 m_{(a)}}{c^5} \frac{E_{\mu\nu} E^{\mu\nu}}{u_{(a)}^3} S_{(a)}^2 + \dots$$

$$\mathcal{L}_{(a)}^{(R^2, S^2)} = \frac{1}{2} \left( C_{\text{E}^2 \text{S}^2}^{(0)} \right)_{(a)} \frac{G_N^2 m_{(a)}}{c^5} \frac{E_{\mu\alpha} E_\nu^\alpha}{u_{(a)}^3} \left[ S_{(a)}^\mu S_{(a)}^\nu \right]_{\text{STF}} + \dots$$

# Computational Algorithm : Towards Automation

MKM, Mastrolia, Patil, Steinhoff (2022)

MKM, Mastrolia, Patil, Steinhoff (2022)



☑ Automated in-house codes

📌 Aim to publish the code in future

# Diagrams for Spinning Binaries

MKM, Mastrolia, Patil, Steinhoff (2022)

MKM, Mastrolia, Patil, Steinhoff (2022)

$S^0$

Order	Diagrams	Loops	Diagrams
0PN	1	0	1
1PN	4	1	1
		0	3
2PN	21	2	5
		1	10
		0	6
3PN	130	3	8
		2	75
		1	38
		0	9

(a) Non-spinning sector

$S^1$

Order	Diagrams	Loops	Diagrams
LO	2	0	2
NLO	13	1	8
		0	5
N <sup>2</sup> LO	100	2	56
		1	36
		0	8
N <sup>3</sup> LO	894	3	288
		2	495
		1	100
		0	11

(b) Spin-orbit sector

$S^2$

Order	Diagrams	Loops	Diagrams
LO	1	0	1
NLO	7	1	3
		0	4
N <sup>2</sup> LO	58	2	27
		1	24
		0	7
N <sup>3</sup> LO	553	3	125
		2	342
		1	76
		0	10

(a) Spin1-Spin2 and Spin1<sup>2</sup> (Spin2<sup>2</sup>) sector

Order	Diagrams	Loops	Diagrams
LO	1	0	1
NLO	4	1	1
		0	3
N <sup>2</sup> LO	25	2	7
		1	12
		0	6
N <sup>3</sup> LO	168	3	15
		2	101
		1	43
		0	9


(b) ES<sup>2</sup> sector

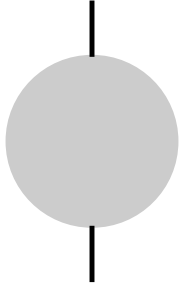
Order	Loops	Diagrams
LO	1	1

(c) E<sup>2</sup> sector

Order	Loops	Diagrams
LO	1	1

(d) E<sup>2</sup>S<sup>2</sup> sector

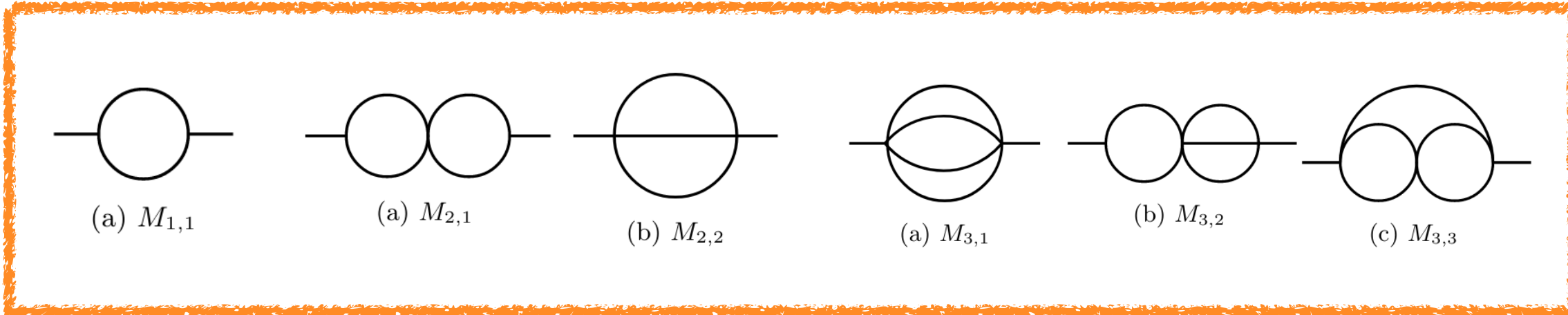
$$\mathcal{L}(x_a, \dot{x}_a, \ddot{x}_a, \dots, S_a, \dot{S}_a, \ddot{S}_a, \dots) = -i \lim_{d \rightarrow 3} \int_{\mathbf{p}} e^{i\mathbf{p} \cdot (\mathbf{x}_{(1)} - \mathbf{x}_{(2)})}$$


$$= -i \lim_{d \rightarrow 3} \int_{\mathbf{p}} e^{i\mathbf{p} \cdot (\mathbf{x}_{(1)} - \mathbf{x}_{(2)})}$$


Dimensional Regularization  $d = 3 + \epsilon$

IBP Decomposition

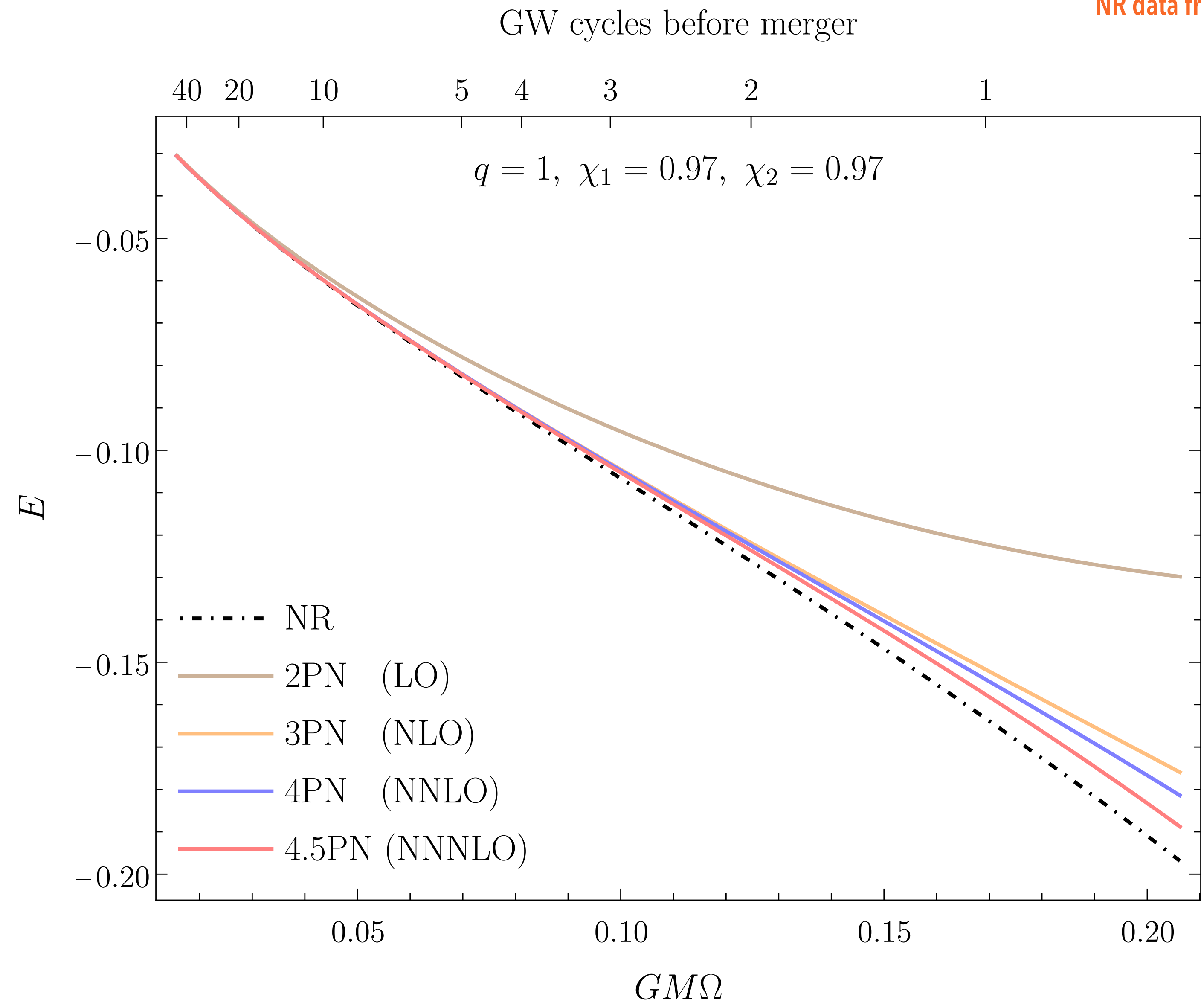
Three Loop MIs



# Binding Energy for Spin-Orbit Coupling

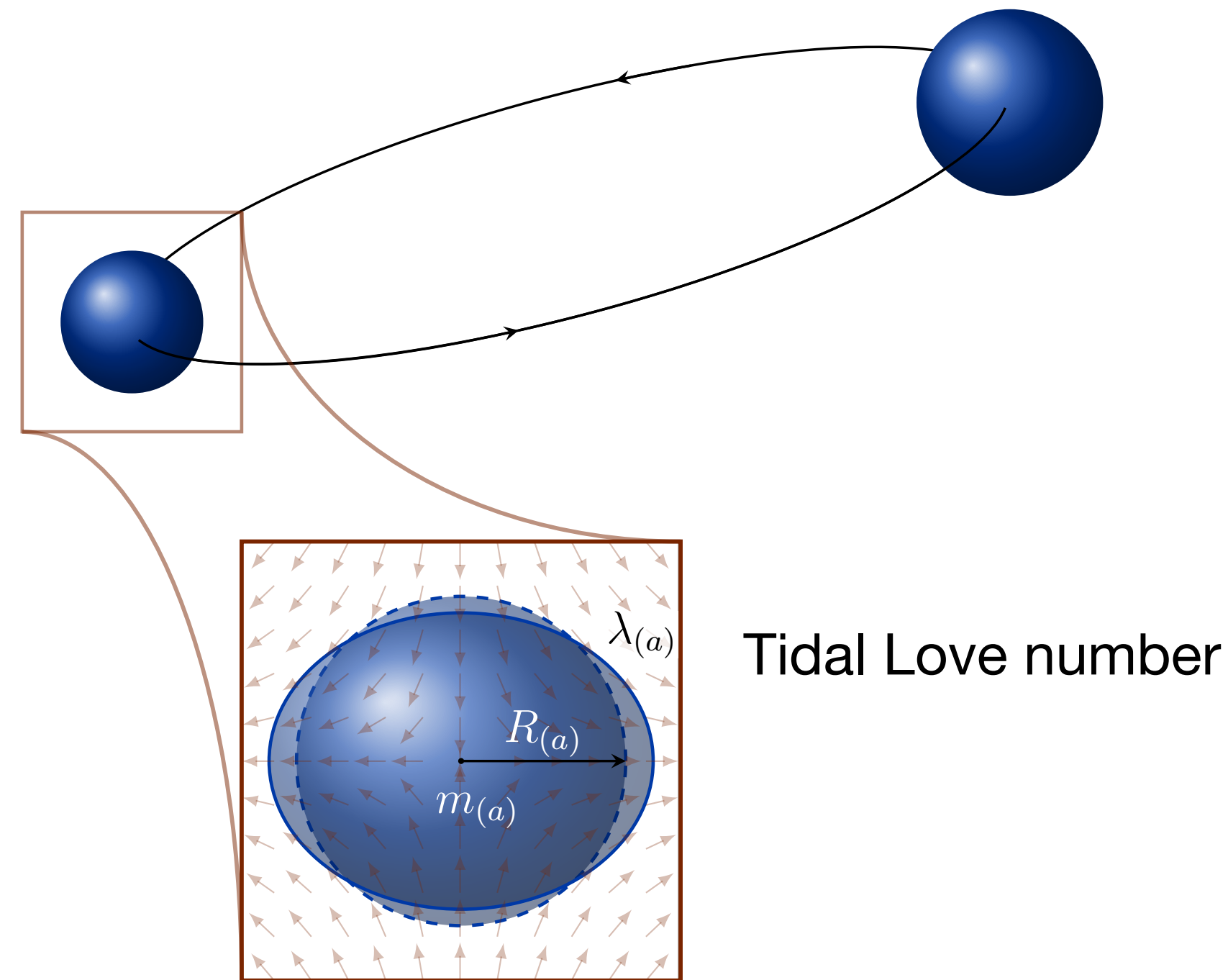
MKM, Mastrolia, Patil, Steinhoff (2022)

NR data from Ossokine, Dietrich, Foley, Katebi, Lovelace (2018)



# Tidal Effects

MKM, Mastrolia, O Silva, Patil, Steinhoff (2023)



- NS features a number of oscillation modes
- The dominant mode is known as *f-mode*, which is the lowest frequency surface gravity waves
- The frequency depend only on the mean density of the star and not on the Equation of State of the NS
- The *f*-modes dynamical tides are important as it significantly affect the inference of the equations of state of NS

Pratten, Schmidt, Williams (2022)

# Dynamical Electric Tides at 2 PN

MKM, Mastrolia, O Silva, Patil, Steinhoff (2023)

$$\mathcal{L}_{\text{DT}} = \frac{z}{4\lambda\omega_f^2} \left[ \frac{c^2}{z^2} \frac{dQ_{\mu\nu}}{d\tau} \frac{dQ^{\mu\nu}}{d\tau} - \omega_f^2 Q_{\mu\nu} Q^{\mu\nu} \right] - \frac{z}{2} E_{\mu\nu} Q^{\mu\nu}$$

Adiabatic limit:  $\omega_f \rightarrow \infty$

$$Q_{\mu\nu} = -\lambda E_{\mu\nu}$$

Tidal deformability

## Binding Energy

$$E_{\text{AT}}(x, \tilde{\lambda}_{(a)}) = -x^6 (9\tilde{\lambda}_{(+)}) + x^7 \left[ \left( \frac{33}{4}\nu - \frac{121}{8} \right) \tilde{\lambda}_{(+)} - \left( \frac{55}{8} \right) \delta\tilde{\lambda}_{(-)} \right] \\ + x^8 \left[ \left( -\frac{91}{16}\nu^2 + \frac{2717}{42}\nu - \frac{20865}{224} \right) \tilde{\lambda}_{(+)} + \left( \frac{715}{48}\nu - \frac{11583}{224} \right) \delta\tilde{\lambda}_{(-)} \right]$$

## Scattering Angle

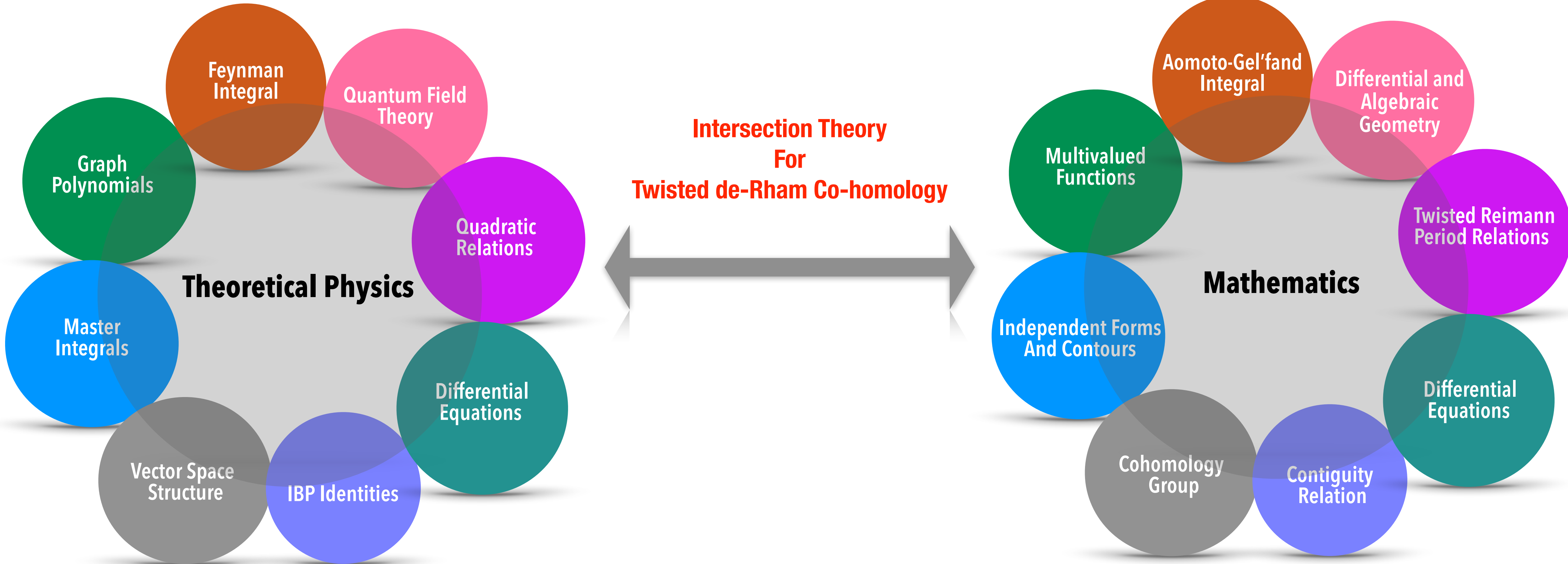
$$\frac{\chi_{\text{AT}}}{\Gamma} = \frac{1}{Mb^4} [\lambda_{(+)} \delta\lambda_{(-)}] \cdot \left\{ \pi \left( \frac{G_N M}{v^2 b} \right)^2 \begin{bmatrix} 1 \\ 0 \end{bmatrix} \left\{ \frac{45}{16} + \frac{135}{32} \left( \frac{v^2}{c^2} \right) + \frac{1575}{256} \left( \frac{v^4}{c^4} \right) \right\} \right. \\ \left. + \left( \frac{G_N M}{v^2 b} \right)^3 \left\{ 48 \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \begin{bmatrix} 732/5 \\ 12 \end{bmatrix} \left( \frac{v^2}{c^2} \right) + \frac{3}{35} \begin{bmatrix} 3073 \\ 593 \end{bmatrix} \left( \frac{v^4}{c^4} \right) \right\} \right. \\ \left. + \pi \left( \frac{G_N M}{v^2 b} \right)^4 \left\{ \frac{315}{8} \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \frac{315}{64} \begin{bmatrix} 51 - 2\nu \\ 5 \end{bmatrix} \left( \frac{v^2}{c^2} \right) + \frac{15}{128} \begin{bmatrix} 5331 - 274\nu \\ 1383 \end{bmatrix} \left( \frac{v^4}{c^4} \right) \right\} \right\}$$



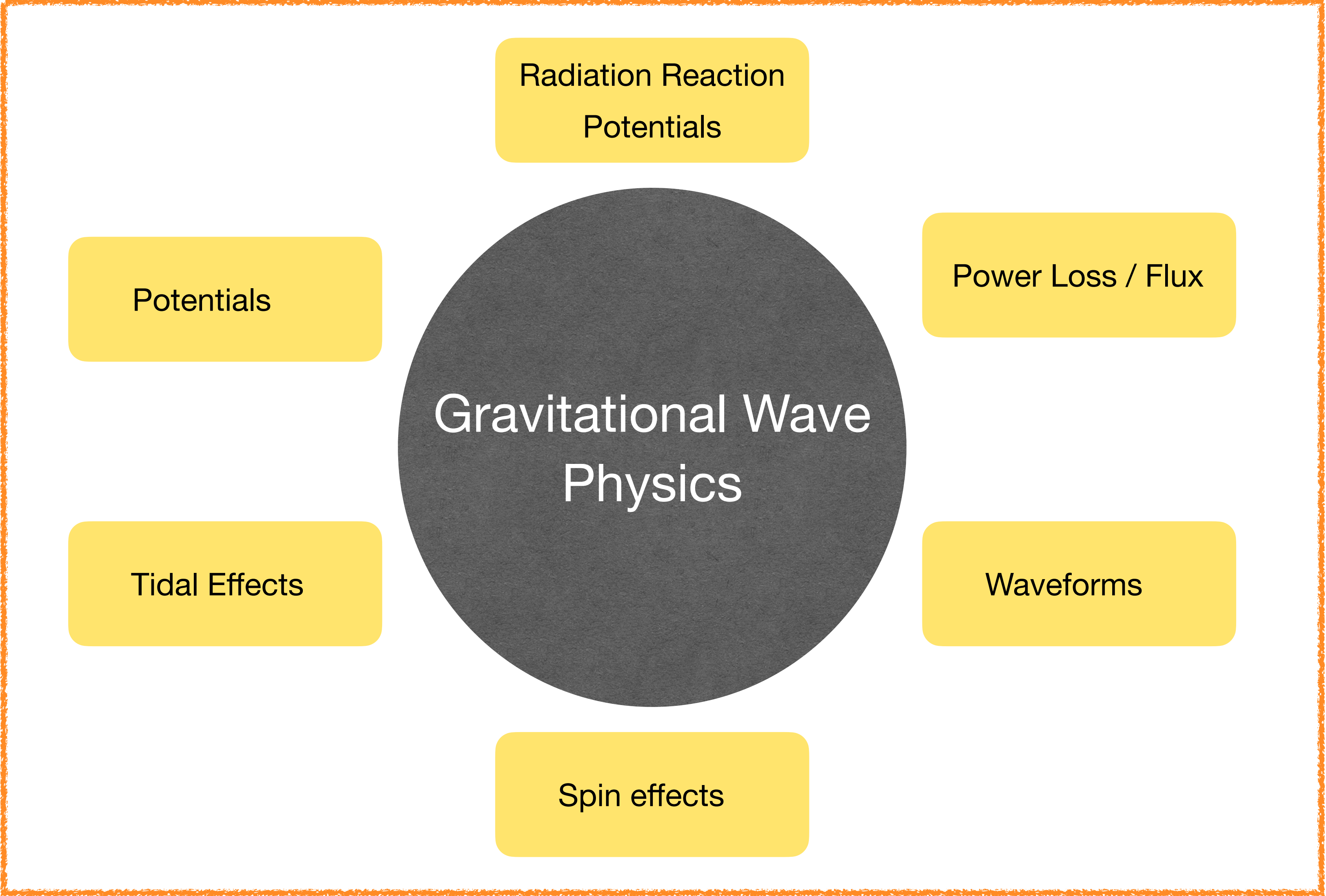
# Conclusion

- ✓ Novel Algebraic Property Unveiled
  - ✓ The algebra of Feynman Integrals is controlled by intersection numbers
  - ✓ Intersection Numbers : Scalar Product/Projection between Feynman Integrals
  - ✓ Useful for both Physics and Mathematics
  
- ✓ Applications to GW phenomenology
  - ✓ progress in understanding spin effects / tidal effects for the compact binaries
  - ✓ A number of observables e.g binding energy, scattering angle has been computed to high precision

# Outlook



# Outlook



Thank You