Compaction function profiles from stochastic inflation

Rome, December 13, 2023 Eemeli Tomberg, Lancaster University

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Single-field inflation is simple

Action:

$$S = \int \mathrm{d}^4 x \sqrt{-g} \left[\frac{1}{2} R - \frac{1}{2} \partial^\mu \varphi \partial_\mu \varphi - V(\varphi) \right]$$

Background equations of motion: $\ddot{\varphi} + 3H\dot{\varphi} + V'(\varphi) = 0$, $3H^2 = \frac{1}{2}\dot{\varphi}^2 + V(\varphi)$

Slow-roll parameters: $\epsilon_1 \equiv -\partial_N \ln H$, $\epsilon_2 \equiv \partial_N \ln \epsilon_1$













Why this picture is inaccurate

Perturbations in the tail are not Gaussian

 ${\cal R}$ is not the correct statistic for PBH formation

Approximations in two regimes



Patched together at the coarse-graining scale $k = k_{\sigma} \equiv \sigma a H$







Stochastic inflation

$$\begin{split} \phi' &= \pi + \xi_{\phi} \,, \quad \pi' = -\left(3 - \frac{1}{2}\pi^2\right)\pi - \frac{V'(\phi)}{H^2} + \xi_{\pi} \,, \quad H^2 = \frac{V(\phi)}{3 - \frac{1}{2}\pi^2} \\ \delta\phi_k'' &= -(3 - \frac{1}{2}\pi^2)\delta\phi_k' - \left[\frac{k^2}{a^2H^2} + \pi^2(3 - \frac{1}{2}\pi^2) + 2\pi\frac{V'(\phi)}{H^2} + \frac{V''(\phi)}{H^2}\right]\delta\phi_k \end{split}$$

$$\langle \xi_{\phi}(N)\xi_{\phi}(N')\rangle = \frac{1}{6\pi^2} \frac{\mathrm{d}k_{\sigma}^3}{\mathrm{d}N} |\delta\phi_{k_{\sigma}}(N)|^2 \delta(N-N') \langle \xi_{\pi}(N)\xi_{\pi}(N')\rangle = \frac{1}{6\pi^2} \frac{\mathrm{d}k_{\sigma}^3}{\mathrm{d}N} |\delta\phi'_{k_{\sigma}}(N)|^2 \delta(N-N') \langle \xi_{\phi}(N)\xi_{\pi}(N')\rangle = \frac{1}{6\pi^2} \frac{\mathrm{d}k_{\sigma}^3}{\mathrm{d}N} \delta\phi_{k_{\sigma}}(N) \delta\phi'^*_{k_{\sigma}}(N) \delta(N-N')$$

 $\mathcal{R}_{< k} = \Delta N = N - \bar{N}$

Full numerical computations















Simplified stochastic equation: $d\phi = \frac{\epsilon_2}{2}(\phi - \phi_0)dN + \frac{\epsilon_2}{2}\phi_0 e^{\frac{\epsilon_2}{2}N}\sqrt{\mathcal{P}_{\mathcal{R}}(k_\sigma)dN}\,\hat{\xi}_N$ $\langle \hat{\xi}_N \hat{\xi}_{N'} \rangle = \delta_{NN'}$

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$$\phi(N) = \phi_0\left(1 - e^{\frac{\epsilon_2}{2}N}\right) + \frac{\epsilon_2}{2}\phi_0 e^{\frac{\epsilon_2}{2}N}X_{< k_{\sigma}}$$

$$\langle \hat{\xi}_N \hat{\xi}_{N'} \rangle = \delta_{NN'}$$

$$X_{$$

 ΔN distribution



 ΔN distribution





 ΔN distribution



$$X_{$$

$$p(\Delta N_{< k}) = \frac{1}{\sqrt{2\pi\sigma_k}} \exp\left[-\frac{2}{\sigma_k^2 \epsilon_2^2} \left(1 - e^{-\frac{\epsilon_2}{2}\Delta N_{< k}}\right)^2 - \frac{\epsilon_2}{2}\Delta N_{< k}\right]$$
$$\Delta N_{< k} = \mathcal{R}_{< k}$$

Comparison to numerics



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Compaction function: right tool for determining the collapse threshold

$$\mathcal{C} \equiv 2 \frac{M_{\rm MS} - M_{\rm bg}}{R}$$

Collapse:
$$C_{\max} > C_c \approx 0.4$$

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In inflationary variables:

$$\mathcal{C}(r) = \frac{2}{3}(1 - [1 + r\zeta'(r)]^2)$$

Assume spherical symmetry

$$r\zeta'(r) = \sum_{k} \frac{2k^2 \, \mathrm{d}k}{\sqrt{2\pi}} \, \zeta_k \left[\cos(kr) - \frac{\sin(kr)}{kr} \right]$$
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Master formula



Alternative collapse measure: averaged compaction function $R = are^{\zeta}$ $\bar{\mathcal{C}}(r) \equiv \frac{3}{R(r)^3} \int_0^{R(r)} \mathrm{d}\tilde{R}\tilde{R}^2 \mathcal{C} \quad \leftarrow$ $= -\frac{2}{r^3 e^{3\zeta(r)}} \int_0^r \mathrm{d}\tilde{r}\,\tilde{r}^2 e^{3\zeta} [2\tilde{r}\zeta' + 3(\tilde{r}\zeta')^2 + (\tilde{r}\zeta')^3]$











Problems

Collapse simulations have smooth peaks.

Us: Stochastic peaks?

• Physics? Smoothing? Window functions?

Multiple peaks?

- "Outermost peak" gives final collapse?
- Overlapping peaks?

Conclusions

Stochastic inflation introduces non-Gaussian corrections to PBH statistics

Compaction function formalism needed for accurate results

Spiked radial profiles: what to do?



[2205.13540]

Initial PBH fractions

Gaussian approximation, $\mathcal{R}_{< k} > 1$, fixed $k: \beta \approx 5 \times 10^{-16}$

Non-Gaussian statistics, $\mathcal{R}_{< k} > 1$, fixed $k \colon \ eta pprox 2.2 imes 10^{-11}$

 $\bar{\mathcal{C}}_{\max} > 0.4: \quad \beta \approx 1.4 \times 10^{-8}$

 $C_{\rm max} > 0.4: \quad \beta \approx 0.016$