# Clustering and two-body correlations within extended density functional approaches

#### Workshop MONSTRE

Milano - 11<sup>th</sup> May, 2023



Speaker: S. Burrello

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INFN - Laboratori Nazionali del Sud

Stefano Burrello Clustering and two-body correlations in EDF-based models

### Outline of the presentation

#### Theoretical approaches for nuclear many-body problem

- Ab-initio vs phenomenological models based on energy density functionals (EDF)
- Effective interaction and nuclear matter (NM) Equation of State (EoS)

#### 2 Extended EDF-based models: recent developments and results

- Bridging ab-initio with phenomenological EDF approaches
  - Benchmark on microscopic pseudo-data for low-density neutron matter
  - Power counting analysis based on many-body perturbative expansion
- Beyond mean-field: many-body correlations and clustering phenomena
  - Neutron star (NS) crust modelization for a global and unified EoS
  - Embedding short-range correlations within relativistic approaches

#### Summary and perspectives within MONSTRE

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Mean-field models for nuclear structure and reaction studies Link to ab-initio: low-density expansion and power counting

### Theoretical models for EoS and finite nuclei

- Ab-initio approaches based on many-body expansion
  - Realistic or effective field theory (EFT) interactions
    - $\Rightarrow$  Diagrammatic hierarchy (power counting)







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• Energy Density Functional (EDF) theory

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Mean-field models for nuclear structure and reaction studies Link to ab-initio: low-density expansion and power counting

### Theoretical models for EoS and finite nuclei

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  - Realistic or effective field theory (EFT) interactions
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⇒ Description of ground state and excitations

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Mean-field models for nuclear structure and reaction studies Link to ab-initio: low-density expansion and power counting

#### Nuclear structure: neutron skin and pygmy resonance

- Non-relativistic Skyrme-like EDF
- Structure of neutron-rich nuclei
   [Zheng et al., PRC 94(1), 014313 (2016)]
   [S. Burrello et al., PRC 99(5), 054314 (2019)]
- Neutron skin thickness  $\Delta r_{np} \Leftrightarrow L$





- Time-Dependent-Hartree-Fock **(TDHF)**  $i\hbar\dot{\rho}(t) + \left[\hat{\rho}, \hat{\mathcal{H}}_{eff}[\rho]\right] = 0$
- Isovector dipole (collective) excitations:
  - Pygmy Dipole Resonance (PDR)



Clustering and two-body correlations in EDF-based models

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#### Merging nuclear structure and reaction studies



- Same framework as for nuclear structure  $\Rightarrow$  Merging with reaction studies
- Role of different terms of effective interaction (and EoS) on final outcomes
  - Importance of **momentum** dependent + surface terms (+ symmetry energy)
- Heavy ion collisions are reliable tools to extract information of EoS!

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Theoretical approaches for nuclear many-body problem

Ab-initio vs phenomenological models based on energy density functionals (EDF)
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#### 2 Extended EDF-based models: recent developments and results

- Bridging ab-initio with phenomenological EDF approaches
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### Pure neutron matter (PNM) low-density expansion

#### • Dilute PNM ( $a_s = -18.9 \text{ fm}$ ) $\Rightarrow$ close to unitary limit of interacting Fermi gas

• Lee-Yang expansion in  $(a_s k_F)$  from EFT  $(\nu_i = 2, 4$  for PNM, symmetric NM)

$$\frac{E}{N} = \frac{\hbar^2 k_F^2}{2m} \left[ \frac{3}{5} + (\nu_i - 1) \frac{2}{3\pi} (k_F a_s) + (\nu_i - 1) \frac{4}{35\pi^2} (11 - 2\ln 2) (k_F a_s)^2 + \dots \right]$$



• New class of EDFs inspired by EFT 🗸

Application to drops & nuclei ⇒ surface
 [S. Burrello et al., PRC 103(6), 064317 (2021)]
 Improving neutron effective mass prediction
 Implementation in dynamical models

Finite temperature (T)  $\Rightarrow X$  impact on NS modelization ("pasta" formation) [S. Burrello & M. Grasso, EPJA 58(2), 22 (2022)]

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$$\frac{E}{N} = \frac{n \kappa_F}{2m} \left[ \frac{5}{5} + (\nu_i - 1) \frac{2}{3\pi} (k_F a_s) + (\nu_i - 1) \frac{4}{35\pi^2} (11 - 2 \ln 2) (k_F a_s)^2 + \dots \right]$$

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Summary and perspectives within MONSTRE

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Mean-field models for nuclear structure and reaction studies Link to ab-initio: low-density expansion and power counting

### Beyond MF: towards a power counting in EDF

- Beyond MF (BMF) ⇒ correlations taken into account (double-counting)
  - Hierarchy of interaction (and EoS) contributions  $\Rightarrow$  power counting in EDF
- EoSs at next-to-leading order (NLO) for symmetric NM (SNM) and PNM



- (t<sub>0</sub>, t<sub>3</sub>) Skyrme-like V<sub>LO</sub>
- Renormalizability analysis
  - ✓ perturbative scheme
- Next-to-NLO (EFT-analysis):
  - 🗴 Expansion parameter
  - 🗡 Breakdown scale



[S. Burrello, C.J. Yang, M. Grasso, PLB 811, 13593 (2020)]

✓ BMF study of closed-shell nuclei [C.J. Yang et al., PRC 106 (1), L011305 (2022)]

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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

### Pairing correlations and nuclear superfluidity

• **Pairing** effects on mechanical (spinodal) instability in low-density nuclear matter ⇒ variation on compressibility and isotopic content of the clusterized matter

Stefano Burrello

[S. Burrello, M. Colonna, F. Matera, PRC 89 (2014)]



- Homogenous stellar matter: impact of superfluidity on ν-scattering ⇒ cooling process in proto-NS (PNS) or pre-bounce of supernova explosions
  - [S. Burrello, M. Colonna, F. Matera, PRC 94 (2016)]



### Clustering phenomena and neutron star crust

- Many-body (short-range) correlations (SRCs) below  $\rho_0$ 
  - Formation of **bound** state of nucleons (clustering)
- Phenomenological models with clusters
  - Dilute matter as a mixture of nucleons and nuclei
    - $\Rightarrow$  Nuclear statistical equilibrium (NSE) model
    - [A. R. Raduta, F. Gulminelli, PRC 82, 065801 (2010)]
  - Unified description of NS EoS & crust-core transition
    - Composition and heat capacity of NS inner crust.
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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

#### Clustering phenomena and neutron star crust

- Many-body (short-range) correlations (SRCs) below  $\rho_0$ 
  - Formation of **bound** state of nucleons (clustering)
- Phenomenological models with clusters
  - Dilute matter as a mixture of nucleons and nuclei
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Clustering and two-body correlations in EDF-based models

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Clustering and two-body correlations in EDF-based models

Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

#### In-medium effects and cluster dissolution

- Cluster dissolution approaching saturation from below
   ⇒ Mott effect ruled by Pauli blocking
  - Geometrical excluded-volume mechanism
  - Microscopic in-medium effects
- Generalized relativistic density functional (GRDF)
   ⇒ Meson exchange with density dependent couplings
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Image: Image:

- Mass-shift obtained by solving the in-medium many-body Schrödinger equation
   Parameterization as function of density (n<sub>b</sub>), isospin asymmetry (β). T
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Clustering and two-body correlations in EDF-based models

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### Short-range correlations and EoS at high-density

- Nucleon knock-out in inelastic electron scattering
   [O. Hen et al. (CLAS coll.), Science 346, 614 (2014)]
  - SRCs from tensor components or repulsive core
  - Smearing of Fermi surface (high-k tail at T=0)
  - Two-body (2B) correlations in np  ${}^{3}S_{1}$  channel
  - $\bullet\,$  Pairs amount to  $\approx 20\%\,$  of the nucleon density



- Embedding SRCs in relativistic MF with quasi-deuterons => EoS at high-density
   [S. Burrello & S. Typel, EPJA 58, 120 (2022)]
- X Extension to finite T (+ momentum of clusters with respect to medium)
- 🗡 Inclusion of quasi-deuterons within a kinetic approach [coll. with R. Wang + INFN CT]

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### Covariant formulation of 2B quantal problem

- Single-nucleon momentum distribution  $\Rightarrow$  in-medium 2B wave function (wf)
- Self-consistent calculation with relativistic MF effective interaction
- Covariant formulation of 2B quantal problem
  - (existence of negative-norm "ghost" states)
  - (singular operators unmanageable non-perturbatively)
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- Covariant description of deuteron **bound** and **scattering** states through 2BDEs
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Clustering and two-body correlations in EDF-based models

### Outline of the presentation

#### Theoretical approaches for nuclear many-body problem

- Ab-initio vs phenomenological models based on energy density functionals (EDF)
- Effective interaction and nuclear matter (NM) Equation of State (EoS)

#### Extended EDF-based models: recent developments and results

🗢 Bridging ab-initio with phenomenological EDF approaches

#### Beyond mean-field many-body correlations and clustering phenomenaa

#### Summary and perspectives within MONSTRE

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### Final remarks and conclusions

#### Main topic

- Bridging ab-initio with phenomenological EDF approaches
- Beyond mean-field extension: many-body correlations and clustering

#### Main results

- Application to neutron drops and nuclei of ab-initio-benchmarked EDFs
- NLO perturbativity of renormalized scheme compatible with power counting
- Neutron star crust composition and effects of clusters on cooling process
- Embedding SRCs through quasi-deuterons within relativistic approach

#### Further developments and outlooks

- Improving properties of EFT-inspired EDFs and dynamical implementation
- Inclusion of SRCs at finite T and light clusters within a kinetic approach
- Momentum distribution from in-medium wf + comparison with experiments

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#### THANK YOU FOR YOUR ATTENTION!

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## **Back-up slides**

#### Extended EDF-based models: recent developments and results

#### Bridging ab-initio with phenomenological EDF approaches

- Benchmark on microscopic pseudo-data for low-density neutron matter
- Power counting analysis based on many-body perturbative expansion

#### Beyond mean-field: many-body correlations and clustering phenomena

- Neutron star (NS) crust modelization for a global and unified EoSS
- Embedding short-range correlations within relativistic approaches

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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

#### Lee-Yang-based EDFs: YGLO and ELYO

#### • Dilute PNM ( $a_s = -18.9 \text{ fm}$ ) $\Rightarrow$ close to unitary limit of interacting Fermi gas



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$$\mathcal{E}_{Y} = Y_{i}[\rho]\rho^{2} + D_{i}\rho^{8/3} + F_{i}\rho^{(\alpha+2)}, \qquad Y_{i}[\rho] = \frac{B_{i}}{1 - R_{i}\rho^{1/3} + C_{i}\rho^{2/3}}$$
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- Spread of data at higher density  $\Rightarrow$  YGLO (FP) and YGLO (Akmal)
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- Mapping with Skyrme functional  $\mathcal{E}_{Sk} = \mathcal{E}_0 + \mathcal{E}_3 + \mathcal{E}_{eff}$  (except for  $\mathcal{E}_0 \leftrightarrow Y_i$ )
- D<sub>i</sub> term may originate from different sources
  - momentum-dependent term  $\mathcal{E}_{ extsf{eff}}$  y
  - extra density dependent term  ${\cal E}'_{3,Y}$  (with  $lpha'_Y=2/3)$
  - any combination of both
- Splitting parameter  $W \Rightarrow$  weights the contribution of m\* without modifying EoS

$$\mathcal{E}_{\text{eff},Y} = WD_i, \qquad \mathcal{E}'_{3,Y} = (1 - W)D_i$$

## Mapping with Skyrme and splitting parameter

Potential part of YGLO functional (a<sub>i</sub> = -18.9(-20.0) fm, i = S, N)

$$\mathcal{E}_{Y} = Y_{i}[\rho]\rho^{2} + D_{i}\rho^{8/3} + F_{i}\rho^{(\alpha+2)}, \qquad Y_{i}[\rho] = \frac{B_{i}}{1 - R_{i}\rho^{1/3} + C_{i}\rho^{2/3}}$$
$$B_{i} = \frac{2\pi\hbar^{2}}{m}\frac{\nu_{i} - 1}{\nu_{i}}a_{i}, \qquad R_{i} = \frac{6}{35\pi}\left(\frac{6\pi^{2}}{\nu_{i}}\right)^{1/3}\left(11 - 2\ln 2\right)a_{i}, \qquad \alpha = 0.7$$

- Mapping with Skyrme functional  $\mathcal{E}_{Sk} = \mathcal{E}_0 + \mathcal{E}_3 + \mathcal{E}_{eff}$  (except for  $\mathcal{E}_0 \leftrightarrow Y_i$ )
- D<sub>i</sub> term may originate from **different** sources:
  - momentum-dependent term *E*<sub>eff, Y</sub>
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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

### Density-dependent scattering length and p-wave

- ELYO: Density-dependent scattering length
  - Tuned by **low-density** condition  $|a_s(k_F)k_F| = 1$
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$$t_2(1-x_2) = W_2 \frac{4\pi\hbar^2}{m} a_p^3(\rho)$$

[J. Bonnard, M. Grasso, D. Lacroix, PRC 101, 064319 (2020)]
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### Energy of neutron drops and effective mass

#### ● Application of YGLO and ELYO on finite systems ⇒ neutron drops

- [J. Bonnard, M. Grasso, D. Lacroix, PRC 98, 034319 (2018); PRC 103, 039901(E) (2021)]
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- Adjustment on energy values of drops available from ab-initio calculations
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Stefano Burrello



Clustering and two-body correlations in EDF-based models

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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

#### Exploring neutron dripline: isotopic chain

#### • Hartree-Fock calculations with YGLO: reproduction of ground state properties

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Correlation between tail of density profiles and symmetry energy slope at low-density

Image: A marked black

Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

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- Crucial role of low-density constraints for pasta-phases formation in NS crust [H. Dinh Thi, T. Carreau, A. F. Fantina, F. Gulminelli, A&A 654, A114 (2021)]
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## **Back-up slides**

#### Extended EDF-based models: recent developments and results

#### Bridging ab-initio with phenomenological EDF approaches

- Benchmark on microscopic pseudo-data for low-density neutron matter
- Power counting analysis based on many-body perturbative expansion

#### Beyond mean-field: many-body correlations and clustering phenomena

- Neutron star (NS) crust modelization for a global and unified EoS
- Embedding short-range correlations within relativistic approaches

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- Standard Skyrme interaction, without gradient, spin-orbit and tensor parts
   Minimal to get obtain the in EoS of symmetric nuclear matter (SNM)
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   [C. J. Yang, M. Grasso, X. Roca-Maza, G. Colò and K. Moghrabi, Phys. Rev. C94 (3), 034311 (2016)
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$$\begin{split} E^{(2)} &= -\frac{1}{4} \frac{\Omega}{(2\pi)^9} \int d\mathbf{k}_1 \int d\mathbf{k}_2 \int d\mathbf{q} \frac{|<\mathbf{k}_1 \mathbf{k}_2 |V| \mathbf{k}_1' \mathbf{k}_2' > |^2}{\epsilon_1' + \epsilon_2' - \epsilon_1 - \epsilon_2} \\ \mathbf{k}_1' &= \mathbf{q} + \mathbf{k}_1, \qquad \mathbf{k}_2' = \mathbf{q} + \mathbf{k}_2, \qquad \epsilon_i^{(\prime)} = \frac{\hbar^2 k_i^{(\prime)2}}{2m_i^*} \end{split}$$



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### BMF renormalizability with Skyrme interaction

• EoSs at LO for SNM and pure neutron matter (PNM)

$$\frac{E_{SNM}^{(LO)}}{A} = \frac{3}{10} \frac{\hbar k_F^2}{m} + \frac{\hbar k_F^3}{4\pi^2} t_0 + \frac{\hbar k_F^{3+3\alpha}}{4\pi^2} T_3, \quad T_3 = \left(\frac{2}{3\pi^2}\right)^{\alpha} \frac{t_3}{6}$$
$$\frac{E_{PNM}^{(LO)}}{N} = \frac{3}{10} \frac{\hbar k_{F,n}^2}{m} + \frac{\hbar k_{F,n}^3}{12\pi^2} t_0 (1-x_0) + \frac{\hbar k_{F,n}^{3+3\alpha}}{12\pi^2} \frac{T_3}{2^{\alpha}} (1-x_3)$$

- Renormalizability for SNM by constraining interaction parameters [C. J. Yang, M. Grasso, K. Moghrabi and U. Van Kolck, PRC 95, 054325 (2017)]
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### BMF renormalizability with Skyrme interaction

• EoSs at next-to-LO (NLO) for SNM and neutron matter (PNM)

$$\frac{E_{SNM}^{(\text{NLO})}}{A} = \frac{3}{10} \frac{\hbar k_F^2}{m} + \frac{\hbar k_F^3}{4\pi^2} t_0^{\wedge} + \frac{\hbar k_F^{3+3\alpha}}{4\pi^2} T_3^{\wedge} + \frac{E_{SNM,f}^{(2)}(k_F)}{A} + \frac{E_{SNM,d}^{(2)}(k_F,\Lambda)}{A}, \quad T_3 = \left(\frac{2}{3\pi^2}\right)^{\alpha} \frac{t_3}{6} \\ \frac{E_{PNM}^{(\text{NLO})}}{N} = \frac{3}{10} \frac{\hbar k_{F,n}^2}{m} + \frac{\hbar k_{F,n}^3}{12\pi^2} t_0^{\wedge}(1-x_0^{\wedge}) + \frac{\hbar k_{F,n}^{3+3\alpha}}{12\pi^2} \frac{T_3^{\wedge}}{2\alpha}(1-x_3^{\wedge}) + \frac{E_{PNM,f}^{(2)}(k_{F,n})}{N} + \frac{E_{PNM,d}^{(2)}(k_{F,n},\Lambda)}{N}$$

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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

# Perturbativity and power counting in EDF theories

Convergence of energy contributions











- $2^{nd}$  order finite part progressively suppressed  $\Rightarrow$  perturbative problem
- Expansion parameter and uncertainties analysis  $\Rightarrow$  Next-to-NLO (NNLO)
- First application to finite nuclei [C.J. Yang, W.G. Jiang, S. Burrello, M. Grasso, arXiv:2110.0195]

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# **Back-up slides**

#### Extended EDF-based models: recent developments and results

Bridging ab-initio with phenomenological EDF approaches

#### Beyond mean-field: many-body correlations and clustering phenomena

- Neutron star (NS) crust modelization for a global and unified EoS
- Embedding short-range correlations within relativistic approaches

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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

# Quasi-deuterons as surrogate for SRCs in GRDF

- Effective resonances (quasi-clusters) for treatment of SRCs at supra-saturation • Embedded in GRDF model through in-medium modifications of  $\Delta m_{\perp}^{(\rm high)}$
- Two-body correlations in np  ${}^3S_1$  channel  $\Rightarrow$  quasi-deuteron
- $T = 0 \Rightarrow$  **boson condensate** of deuterons under chemical potentials **equilibrium**

 $\begin{array}{l} \left( a_{1}a_{2}+a_{2}a_{3}\right) \left( b_{1}a_{2}=0\right) &= a_{1}a_{2}\left( b_{1}a_{2}=0\right) \\ \left( \left( b_{1}a_{2}+b_{2}a_{3}\right) \left( b_{1}a_{3}\right) + \left( b_{1}a_{3}\right) \right) &= \left( b_{1}a_{1}a_{2}\right) \\ \left( \left( b_{1}a_{2}a_{3}+b_{3}a_{3}\right) + \left( b_{1}a_{3}\right) + \left( b_{1}a_{3$ 

- $m_{\text{nuc}}^* \ge 0 \Rightarrow 0 \le X_d \le \min\left\{X_d^{(\max)}, 1 |\beta|\right\}, X_d^{(\max)} = \frac{m_{\text{nuc}}}{\gamma_d \int_{-\infty}^{\infty} m_d \gamma_d}$
- Crucial role of scaling factor  $\chi_d \equiv \chi$  for bound nucleon-meson coupling strenght

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$$\mu_d = \mu_n + \mu_p \Rightarrow \qquad m_d^* + \Delta m_d^{(high)} + V_d' = \sqrt{k_a^2 + (m_b^*)^2 + V_p' + \sqrt{k_p^2 + (m_p^*)^2 + V_p'}}$$

 $m_{i}^{*} = m_{i} - S_{i} \qquad S_{i} = \chi_{i} A_{i} C_{\sigma} n_{\sigma} \qquad V_{i} = \chi_{i} A_{i} (C_{\omega} n_{\omega} + C_{\mu} n_{\rho})$   $V_{i}^{I} = V_{i} + W_{i} + W_{i}^{(e)} \qquad W_{i} = \frac{1}{2} \left( C_{\omega}^{i} n_{\omega}^{2} + C_{\mu}^{i} n_{\rho}^{2} - C_{\sigma}^{i} n_{\sigma}^{2} \right)$   $V_{i}^{(e)} = n_{d} \frac{\partial \Delta m_{d}^{(\text{high})}}{\partial n_{i}} \qquad C_{j} = \frac{\Gamma_{j}^{2}(n_{b})}{m_{e}^{2}} \qquad C_{j}^{i} = \frac{dC_{j}}{dn_{b}}, \qquad j = \sigma, \omega, \rho$ 

•  $m_{\text{nuc}}^* \ge 0 \Rightarrow 0 \le X_d \le \min\left\{X_d^{(\max)}, 1 - |\beta|\right\}, X_d^{(\max)} = \frac{m_{\text{nuc}}}{\gamma_d C_c n_b} \frac{m_{\text{nuc}}}{n_b \to \infty}$ 

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• With scalar  $(S_i)$ , vector  $(V_i)$  and rearrangement  $(W_i, W_i^{(r)})$  potentials (i = nuc, d)

$$\mu_{d} = \mu_{n} + \mu_{p} \Rightarrow \qquad m_{d}^{*} + \Delta m_{d}^{(\text{high})} + V_{d}' = \sqrt{k_{n}^{2} + (m_{n}^{*})^{2} + V_{n}' + \sqrt{k_{p}^{2} + (m_{p}^{*})^{2} + V_{p}'}$$

$$m_i^* = m_i - S_i \qquad S_i = \chi_i A_i C_\sigma n_\sigma \qquad V_i = \chi_i A_i (C_\omega n_\omega + C_\rho n_\rho)$$
$$V_i' = V_i + W_i + W_i^{(r)} \qquad W_i = \frac{1}{2} \left( C_\omega' n_\omega^2 + C_\rho' n_\rho^2 - C_\sigma' n_\sigma^2 \right)$$
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$$\mu_{d} = \mu_{n} + \mu_{p} \Rightarrow \qquad m_{d}^{*} + \Delta m_{d}^{(\text{high})} + V_{d}' = \sqrt{k_{n}^{2} + (m_{n}^{*})^{2}} + V_{n}' + \sqrt{k_{p}^{2} + (m_{p}^{*})^{2}} + V_{p}'$$

$$\begin{split} m_i^* &= m_i - S_i \qquad S_i = \chi_i A_i C_\sigma n_\sigma \qquad V_i = \chi_i A_i \left( C_\omega n_\omega + C_\rho n_\rho \right) \\ V_i' &= V_i + W_i + W_i^{(r)} \qquad W_i = \frac{1}{2} \left( C_\omega' n_\omega^2 + C_\rho' n_\rho^2 - C_\sigma' n_\sigma^2 \right) \\ W_i^{(r)} &= n_d \frac{\partial \Delta m_d^{(\text{high})}}{\partial n_i} \qquad C_j = \frac{\Gamma_j^2(n_b)}{m_i^2} \qquad C_j' = \frac{dC_j}{dn_b}, \qquad j = \sigma, \omega, \rho \end{split}$$

•  $m_{\text{nuc}}^* \ge 0 \Rightarrow 0 \le X_d \le \min\left\{X_d^{(\max)}, 1 - |\beta|\right\}, \ X_d^{(\max)} = \frac{m_{\text{nuc}}}{\chi_d C_\sigma n_b} \xrightarrow[n_b \to \infty]{} 0$ 

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# Quasi-deuterons mass-shift at high-density

- Scaling factor for deuteron-meson coupling strenght
  - $\chi = 1 \Rightarrow$  same strength as for free nucleons
  - $\chi < 1 \Rightarrow$  in-medium effects and description of chemical equilibrium constant
    - [L. Qin et al., PRL 108, 172701 (2012); R. Bougault et al., J. Phys. G 47, 025103 (2020)]
- $1/\sqrt{2} < \chi_s = (0.85 \pm 0.05)$  universal scaling factor [H. Pais et al., PRC 97, 045805 (2018)] • No crossing  $\Rightarrow \Delta m_d(n_b, X_d)$  invertible function for any density  $n_b$
- Δm<sup>(high)</sup><sub>d</sub> ≪ Δm<sup>GRDF</sup><sub>d</sub> ⇒ Large change beyond Mott density for extended GRDF [S. Typel, EPJ Special Topics 229, 3433-3444 (2020)]]

• Interpolation of low-(Pauli blocking) and high-(condensate model) density limit



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#### **Piecewise interpolation and saturation constraints**

• Piecewise parameterization:  $\Delta m_d(n_b, X_d) = \min \left\{ \Delta m_d^{(\text{low})}(n_b), \Delta m_d^{(\text{high})}(n_b, X_d) \right\}$ 



 $\sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{$ 

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 Constraints on NM at saturation (n<sub>0</sub>) (E/A, m<sub>nuc</sub><sup>\*</sup>, pressure, symmetry energy)
 Experimental results of SRCs in nuclei ⇒ X<sub>d,0</sub> = 0.2 (pairs ≈ 20% of density)

	χ	$\Gamma_{\sigma,0}$	$\Gamma_{\omega,0}$	$\Gamma_{ ho,0}$	$\Delta m_{d,0}$ [MeV]	$\frac{d \Delta m_d}{dn_b} \Big _{n_0} $ [MeV fm <sup>3</sup> ]
	1	10.580042	13.217226	3.556424	104.92	813.98
	$1/\sqrt{2}$	10.919963	13.719324	3.400187	58.23	570.80
DD2	—	10.686681	13.342362	3.626940		

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### **Deuteron mass-shift parametrization**

• Unified mass-shift parameterization ( $\gamma = 1$ ) [S. Burrello, S. Typel, EPJA 58, 120 (2022)]  $\Delta m_d(x) = \frac{ax}{1+bx} + cx^{\eta+1} [1 - \tanh(x)] + fx^{\gamma} \tanh(gx), \qquad x = \frac{n_b}{n_0}$ 

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DD2 - d3

541.726060

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-140.309501

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214.368137

0.75

2.715545

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243.472387

# SNM: impact on EoS and matter incompressibility

- Attraction in presence of quasi-deuterons  $\iff$  attraction/repulsion for  $\Gamma_i$ -refit
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## Effect on symmetry energy and its slope



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## Coupling between IS and IV modes

- Symmetric nuclear matter: IS and IV modes are decoupled
- Neutron-rich systems: n and p oscillate with different amplitudes ⇒ coupling



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# Influence of the effective interaction

- SAMi-J interactions
  - [X. Roca-Maza et al., PRC87, (2013)]
  - $\Rightarrow$  isolate influence of IV channel

$$E_{\rm sym}(\rho) = C(\rho)l^2$$



- Sensitivity of E<sub>IV-GDR</sub> to E<sub>sym</sub> at crossing
- Role of symmetry energy **slope**:
  - IV PDR
- Agreement with Vlasov results

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- Role of symmetry energy **slope**:
  - IV PDR ⇔ neutron skin thickness
- Agreement with Vlasov results

[Zheng, H. et al., PRC 94, (2016)]

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#### Comparison between Vlasov and TDHF model



• Good reproduction of IV GDR and IS GDR

Two contributions in low-energy region: [see M. Urban, PRC85, (2012)]

PDR mode (outer surface).

toroidal mode (inner surface against bulk).

Displacement of PDR peaks ⇒ numerical treatment of surface.

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#### Link between nuclear response and density profiles



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#### Sn isotope chain: N/Z evolution of PDR

#### Dipole response evolution with the neutron/proton content ⇒ Sn isotopes chain

- Question: Why IV PDR fraction of EWSR does not grow from N=70 to N=82?
  [S. Ebata, T. Nakatsukasa, T. Inakura, Phys. Rev. C 90, 024303 (2014)]
- Explanation: it reflects the decrease in the IS fraction and IS dipole strength [S. Burrello et al., Phys. Rev. C 99, 054314 (2019)]

• Need to normalize the mixing effect to the IS PDR strength  $\Rightarrow$   $R_f = rac{f_{PDR}^{NV}}{f_{PDR}^{FS}}$ 

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Neutron star crust and unified equation of state Short-range correlations within relativistic approaches

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- Dipole response evolution with the neutron/proton content ⇒ Sn isotopes chain
- Question: Why IV PDR fraction of EWSR does not grow from N=70 to N=82?
  [S. Ebata, T. Nakatsukasa, T. Inakura, Phys. Rev. C 90, 024303 (2014)]
- Explanation: it reflects the decrease in the IS fraction and IS dipole strength



Stefano Burrello Clustering and two-body correlations in EDF-based models

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#### Comparison between TDHF and RPA

#### • TDHF and RPA equivalent in zero-amplitude limit, despite technical procedures

- Question: which numerical parameters ensure the best agreement?
- Dependence on box size (i.e. discretization of continuum single-particle states)
  [S. Burrello et al., Phys. Rev. C 99, 054314 (2019)]
- Very good agreement when the size is large enough (also for transition densities)

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