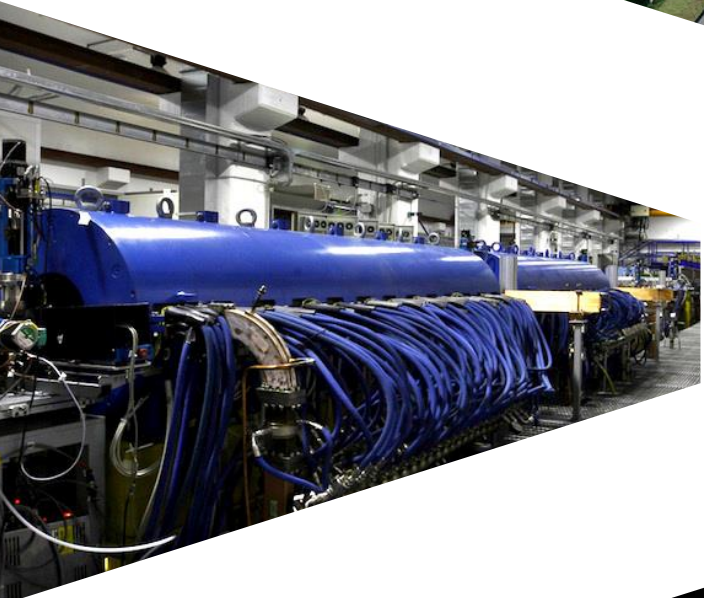
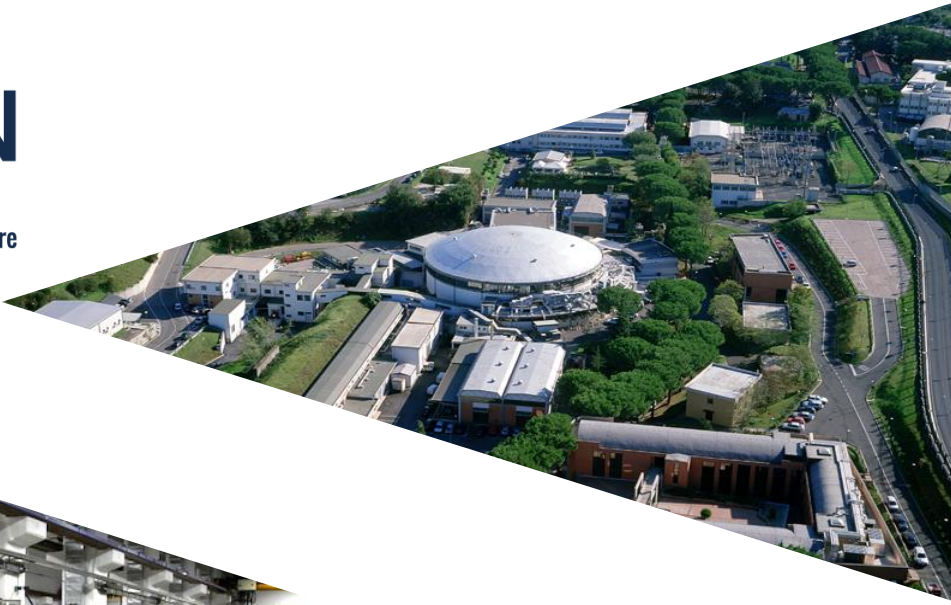




Istituto Nazionale di Fisica Nucleare



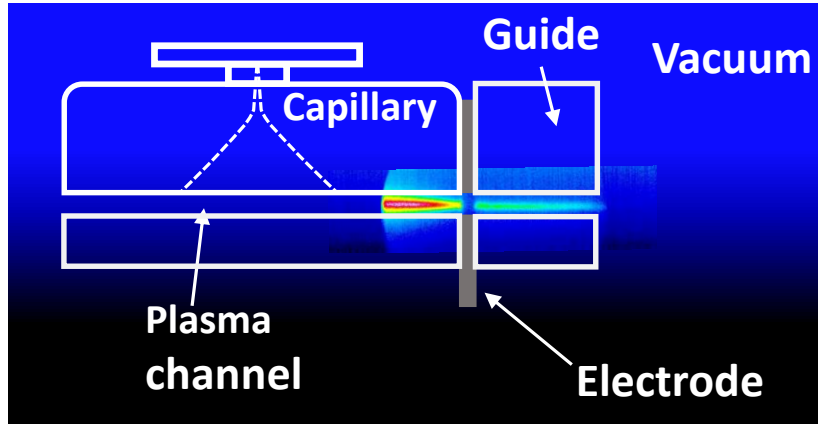
Evaluation of the **transfer matrix** of a **plasma ramp** with squared cosine shape via an approximate solution of **Mathieu differential** **equation**

Stefano Romeo stefano.romeo@Inf.infn.it

Island of Elba 2023-09-18

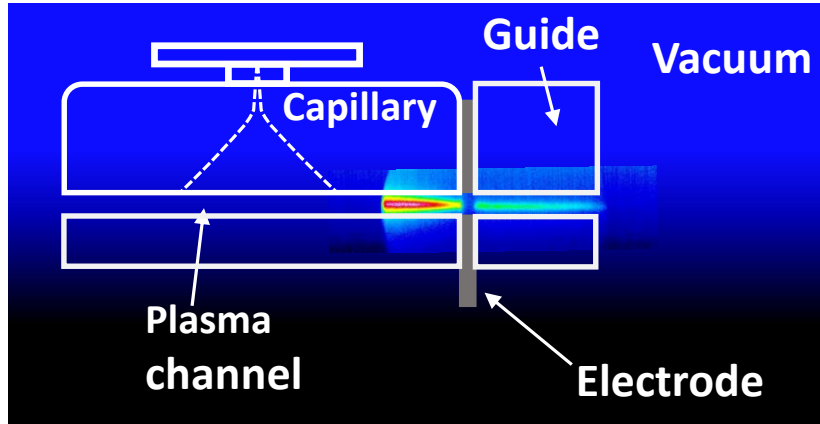


6th European Advanced Accelerator Concepts Workshop



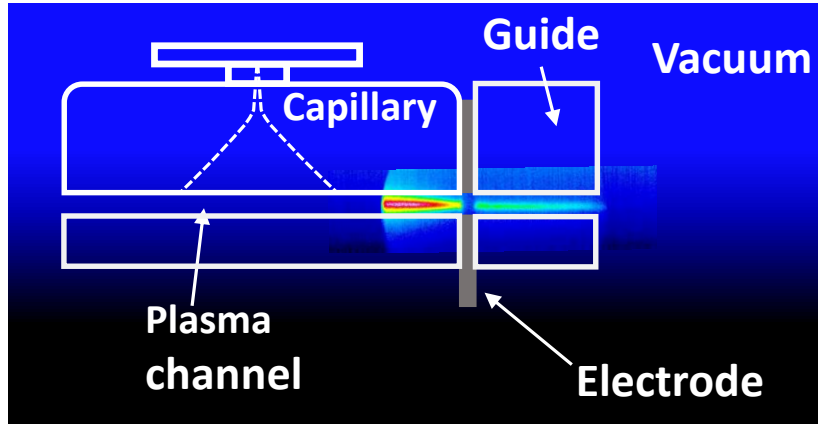
Courtesy of A. Biagioni

- At the edge of discharge capillaries, regions with a gradually decreasing plasma density form, connecting the plasma plateau to the vacuum



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- ❑ At the edge of discharge capillaries, regions with a gradually decreasing plasma density form, connecting the plasma plateau to the vacuum
- ❑ These regions are commonly referred to as plasma ramps



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- Emittance preservation inside plasma is difficult, since in order to balance the extreme focusing field inside plasma the beam needs to be focused up to few micrometers/sub micrometer scale

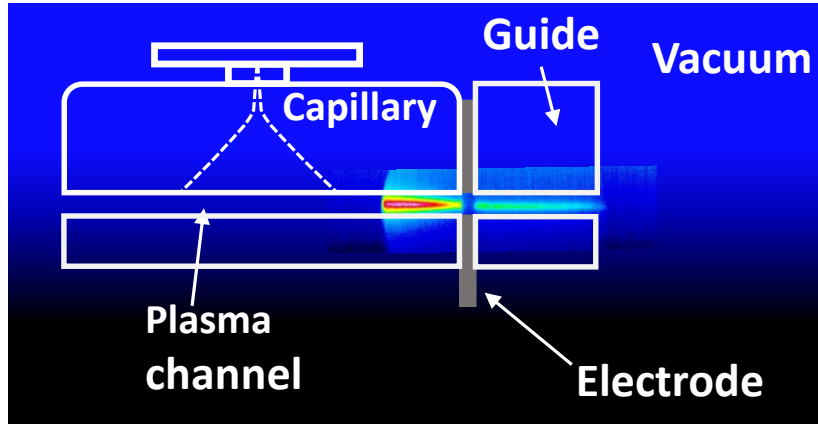
High sensitivity to transverse instability!

$$\sigma_0 = \sqrt[4]{\frac{2}{\gamma}} \sqrt{\frac{\epsilon_n}{k_p}}$$

$$\beta_0 = \frac{\sqrt{2\gamma}}{k_p}$$

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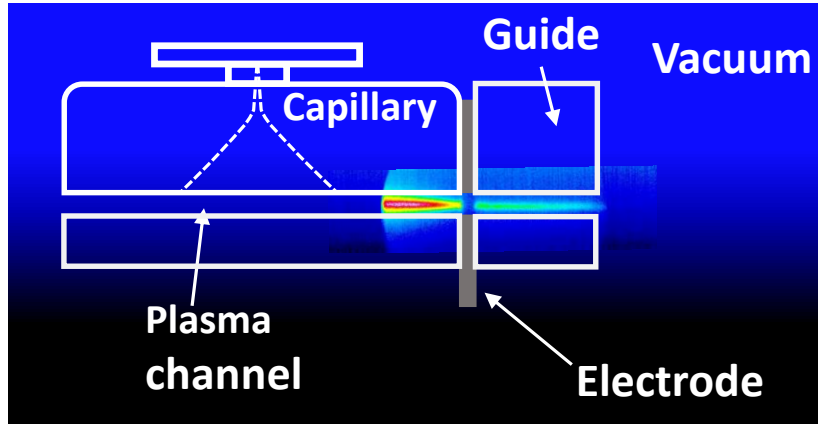
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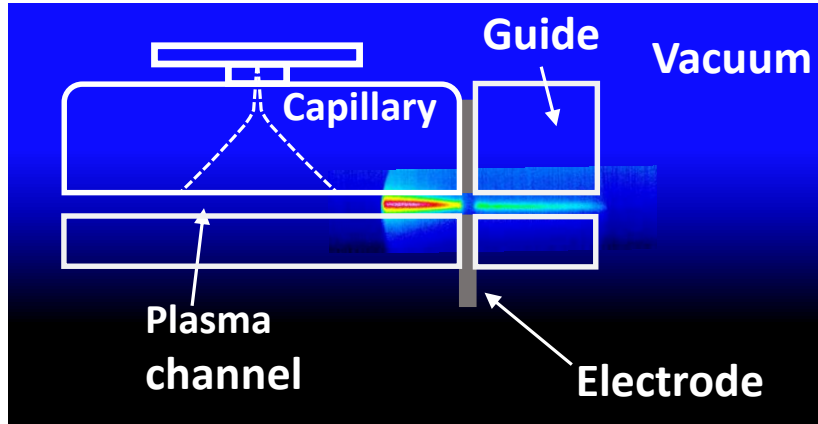
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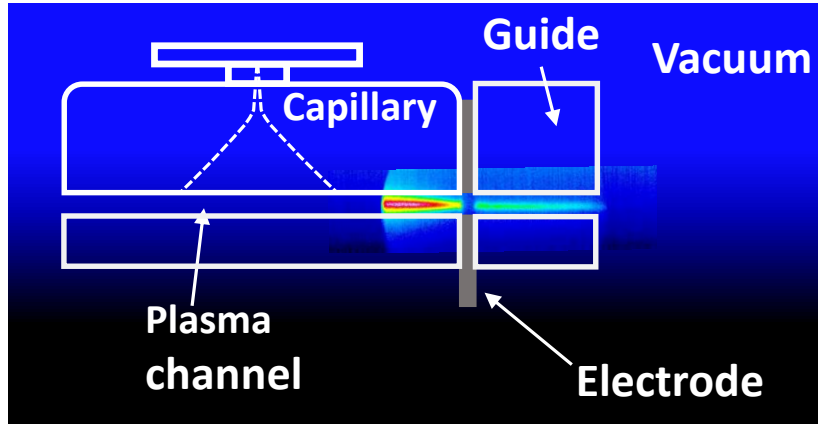
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For any working point perform several simulation scans



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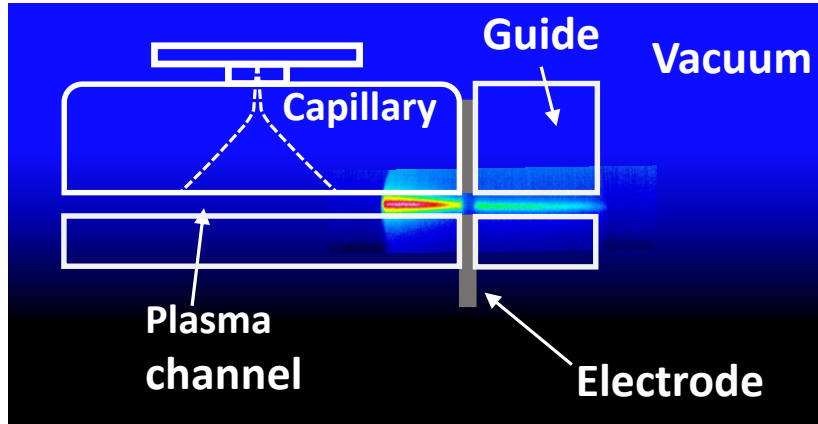
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Analytical solution:

Find a general rule for a suitable shape

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Evaluation of ramp transfer matrix

Analytical solution:

Find a general rule for a suitable shape

Numerical solution:

For any working point perform several simulation scans

Theory: Plasma focusing strength

$$\beta''(z) + 2k_{ext}^2(z)\beta(z) = \frac{2}{\beta(z)} + \frac{[\beta'(z)]^2}{2\beta(z)}$$

Envelope equation

Plasma focusing strength?

$$k_{ext}^2(z) = \frac{k_p^2(z)}{2\gamma}$$

Ion column model
(only the particles
inside the bubble)

- $\beta(z) = \frac{\sqrt{2\gamma}}{k_p}$
- $\alpha(z) = 0$
- $\gamma(z) = \frac{1}{\beta(z)}$
- $\beta''(z) = 0$

Matching conditions
on the plateau

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$$\beta''(z) + 2k_{ext}^2(z)\beta(z) = \frac{2}{\beta(z)} + \frac{[\beta'(z)]^2}{2\beta(z)}$$

Envelope equation

Plasma focusing strength?

$$k_{ext}^2(z) = \frac{k_p^2(z)}{2\gamma}$$

Ion column model
(only the particles inside the bubble)

Neglecting constants, the focusing force inside a plasma ramp depends only from the plasma density, with the energy taken as a parameter

$$k_{ext}^2(z) = \frac{e^2}{2\epsilon_0 m_e c^2} \frac{n_p(z)}{\gamma} = f(z; \gamma)$$

- $\beta(z) = \frac{\sqrt{2\gamma}}{k_p}$
 - $\alpha(z) = 0$
 - $\gamma(z) = \frac{1}{\beta(z)}$
 - $\beta''(z) = 0$
- Matching conditions on the plateau

$$y''(z) + f(z)y(z) = 0$$

Hill's Equation

$$k_{ext}^2(z) = f(z) \propto n_p(z)$$

Negligible deceleration

$f(z)$ is a continuous and derivable function

$$\begin{bmatrix} \beta(z) \\ \alpha(z) \\ \gamma(z) \end{bmatrix} = \begin{bmatrix} C^2 & -2CS & S^2 \\ -CC' & CS' + C'S & -SS' \\ C'^2 & -2C'S' & S'^2 \end{bmatrix} \begin{bmatrix} \beta_0 \\ 0 \\ 1 \\ \frac{1}{\beta_0} \end{bmatrix}$$

The computation is performed over the inverse matrix in order to have simplified matching conditions

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Hill's Equation

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Negligible deceleration

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- $C^2(0)\beta_0 + \frac{s^2(0)}{\beta_0} = \beta_0$
- $C'^2(0)\beta_0 + \frac{s'^2(0)}{\beta_0} = \frac{1}{\beta_0}$
- $-C(0)C'(0)\beta_0 - \frac{s(0)s'(0)}{\beta_0} = 0$
- $\beta''(0) = 0$
- Liouville's Theorem

$$\begin{bmatrix} \beta(z) \\ \alpha(z) \\ \gamma(z) \end{bmatrix} = \begin{bmatrix} C^2 & -2CS & S^2 \\ -CC' & CS' + C'S & -SS' \\ C'^2 & -2C'S' & S'^2 \end{bmatrix} \begin{bmatrix} \beta_0 \\ 0 \\ 1 \\ \frac{1}{\beta_0} \end{bmatrix}$$

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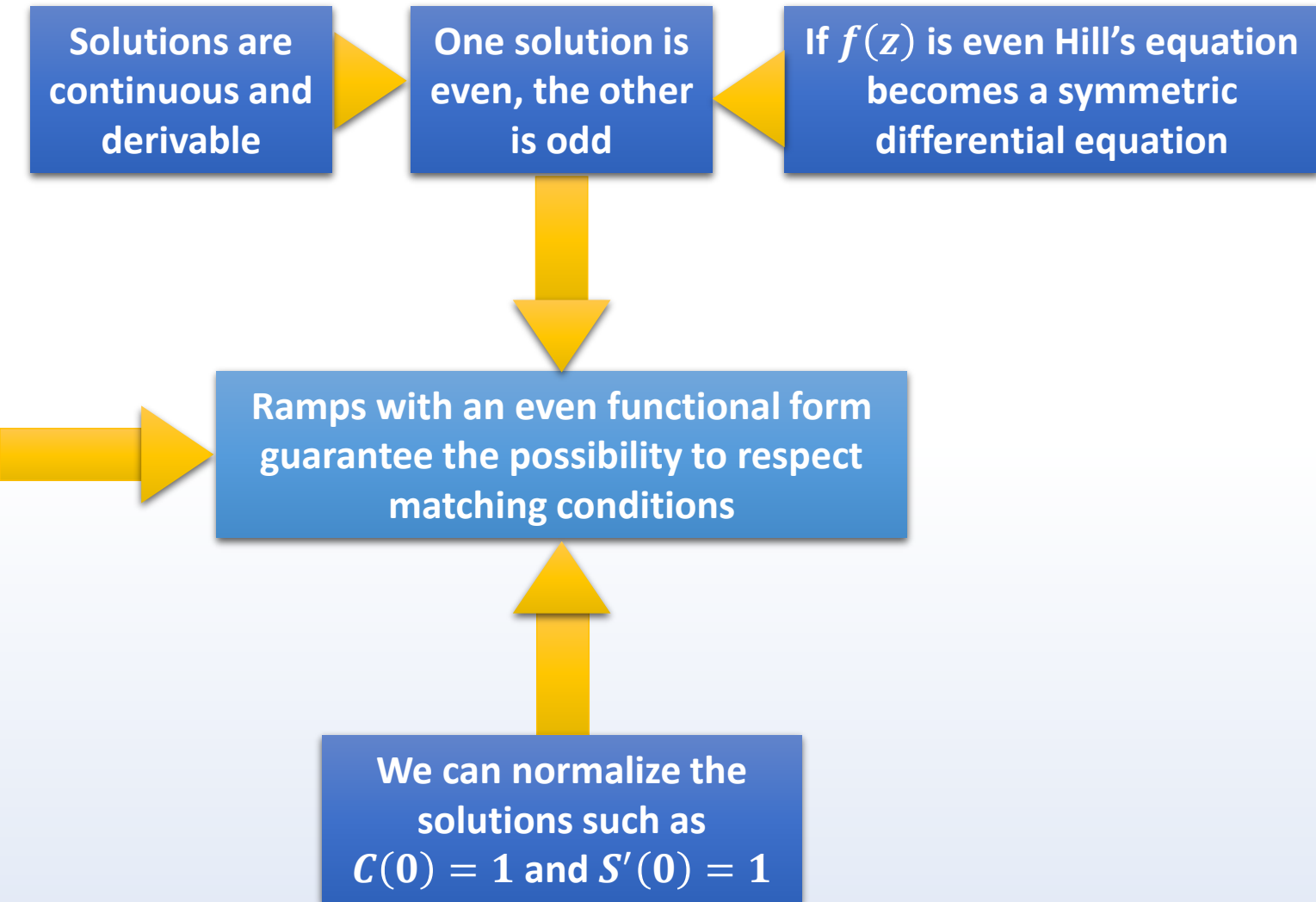
We can normalize the
solutions such as
 $C(0) = 1$ and $S'(0) = 1$

Solutions are
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One solution is
even, the other
is odd

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solutions such as
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Framework: Squared Cosine plasma ramp



THE CHOICE

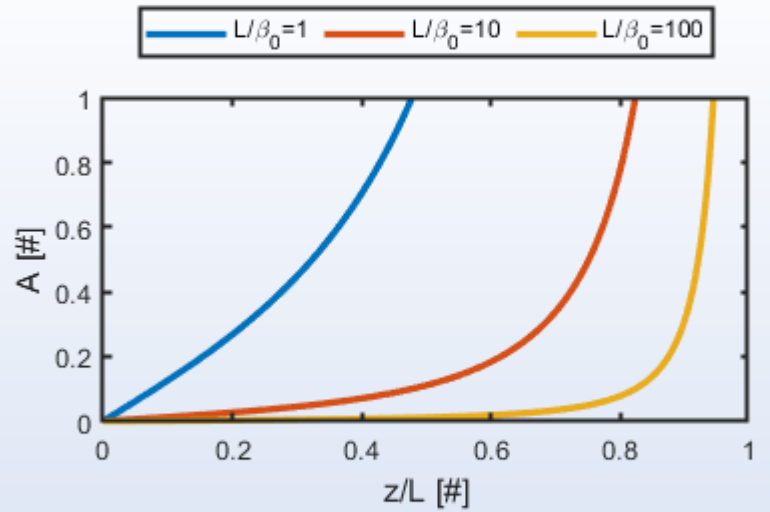
$$f(z) = \frac{1}{\beta_0^2} \cos^2\left(\frac{\pi z}{2L}\right)$$

- Realistic ramp
- No discontinuity at any derivative order
- Partially concave, partially convex
- Never adiabatic

Ramps with an even functional form guarantee the possibility to respect matching conditions

We can normalize the solutions such as $C(0) = 1$ and $S'(0) = 1$

$$\mathcal{A} = \frac{1}{4f^{3/2}(z)} \left| \frac{df(z)}{dz} \right|$$



Hill's Equation

$$y''(z) + \frac{1}{\beta_0^2} \cos^2\left(\frac{\pi z}{2L}\right) y(z) = 0$$

Mathieu Equation

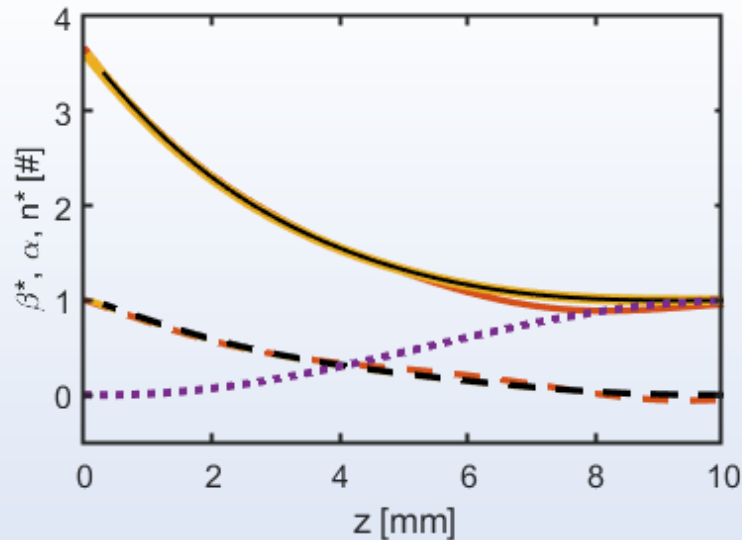
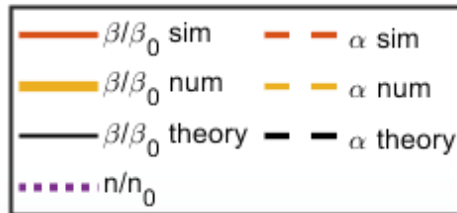
$$y''(\xi) + \frac{2L^2}{\pi^2 \beta_0^2} [1 + \cos(2\xi)] y(\xi) = 0$$

Hill's Equation

$$y''(z) + \frac{1}{\beta_0^2} \cos^2\left(\frac{\pi z}{2L}\right) y(z) = 0$$

Mathieu Equation

$$y''(\xi) + \frac{2L^2}{\pi^2 \beta_0^2} [1 + \cos(2\xi)] y(\xi) = 0$$



Even and odd normalized Mathieu Functions

$$C(z) = \frac{C e_n\left(\frac{2L^2}{\pi^2 \beta_0^2}, -\frac{L^2}{\pi^2 \beta_0^2}, \frac{\pi z}{2L}\right)}{C e_n\left(\frac{2L^2}{\pi^2 \beta_0^2}, -\frac{L^2}{\pi^2 \beta_0^2}, 0\right)}$$

$$S(z) = \frac{2L}{\pi} \frac{S e_n\left(\frac{2L^2}{\pi^2 \beta_0^2}, -\frac{L^2}{\pi^2 \beta_0^2}, \frac{\pi z}{2L}\right)}{S e_n'\left(\frac{2L^2}{\pi^2 \beta_0^2}, -\frac{L^2}{\pi^2 \beta_0^2}, 0\right)}$$

Simplify: Normalized transfer matrix

Normalization

$$R = \frac{L}{\pi\beta_0}$$

Geometric factor

$$c_n = C\left(\frac{\pi}{2}\right) = \frac{C e_n(2R^2, -R^2, \frac{\pi}{2})}{C e_n(2R^2, -R^2, 0)}$$

$$s_n = \frac{S\left(\frac{\pi}{2}\right)}{\beta_0} = 2R \frac{S e_n(2R^2, -R^2, \frac{\pi}{2})}{S e_n'(2R^2, -R^2, 0)}$$

$$c'_n = C'\left(\frac{\pi}{2}\right) \beta_0 = \frac{1}{2R} \frac{C e_n'(2R^2, -R^2, \frac{\pi}{2})}{C e_n(2R^2, -R^2, 0)}$$

$$s'_n = S'\left(\frac{\pi}{2}\right) = \frac{S e_n'(2R^2, -R^2, \frac{\pi}{2})}{S e_n'(2R^2, -R^2, 0)}$$



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Extraction

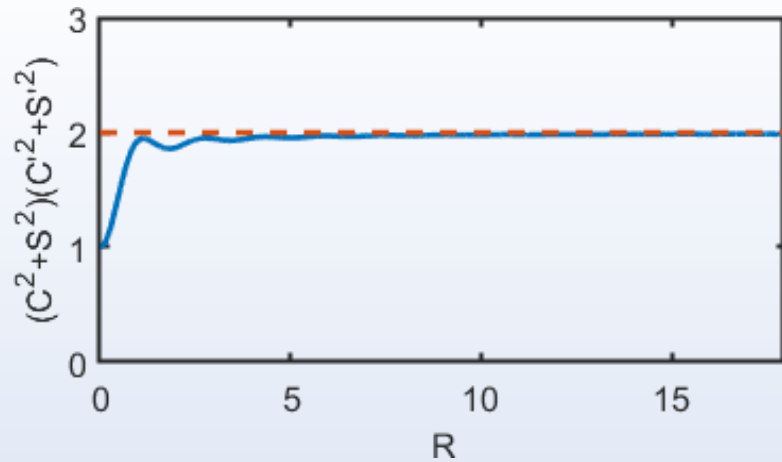
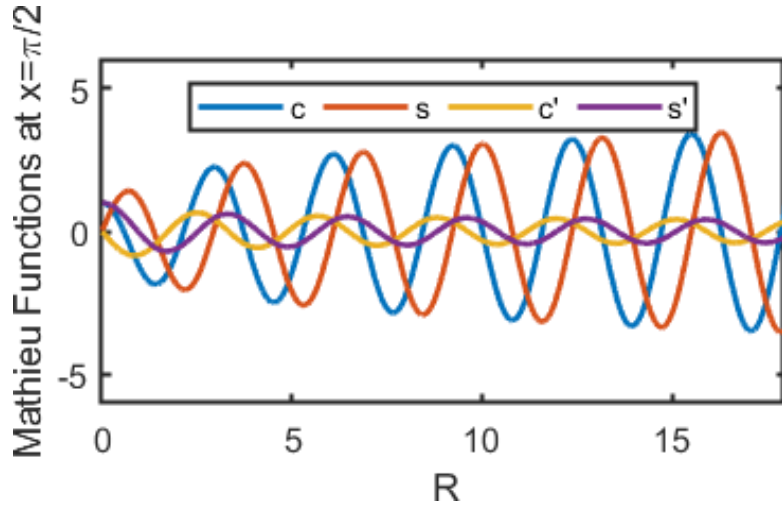
$$\begin{bmatrix} \beta_1/\beta_0 \\ \alpha_1 \\ \gamma_1\beta_0 \end{bmatrix} = \begin{bmatrix} c^2 & -2cs & s^2 \\ -cc' & cs' + c's & -ss' \\ c'^2 & -2c's' & s'^2 \end{bmatrix} \begin{bmatrix} \beta_2/\beta_0 \\ \alpha_2 \\ \gamma_2\beta_0 \end{bmatrix}$$

Normalized transfer matrices

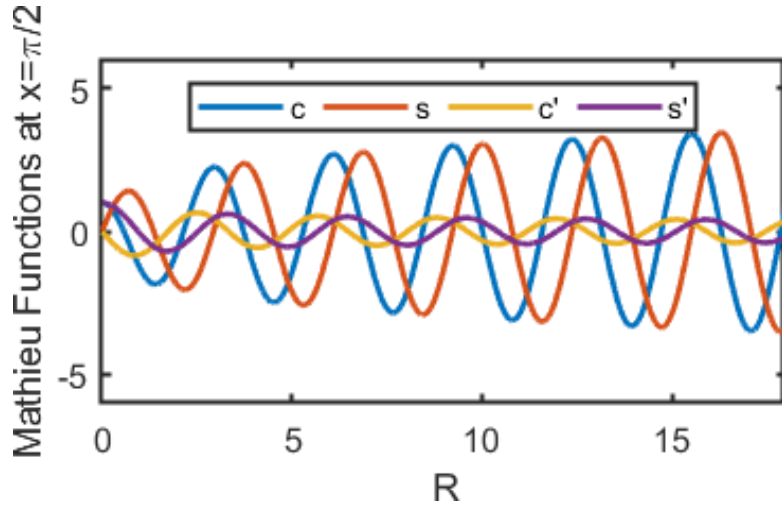
Injection

$$\begin{bmatrix} \beta_2/\beta_0 \\ \alpha_2 \\ \gamma_2\beta_0 \end{bmatrix} = \begin{bmatrix} s'^2 & -2ss' & s^2 \\ -c's' & cs' + c's & -cs \\ c'^2 & -2cc' & c^2 \end{bmatrix} \begin{bmatrix} \beta_1/\beta_0 \\ \alpha_1 \\ \gamma_1\beta_0 \end{bmatrix}$$

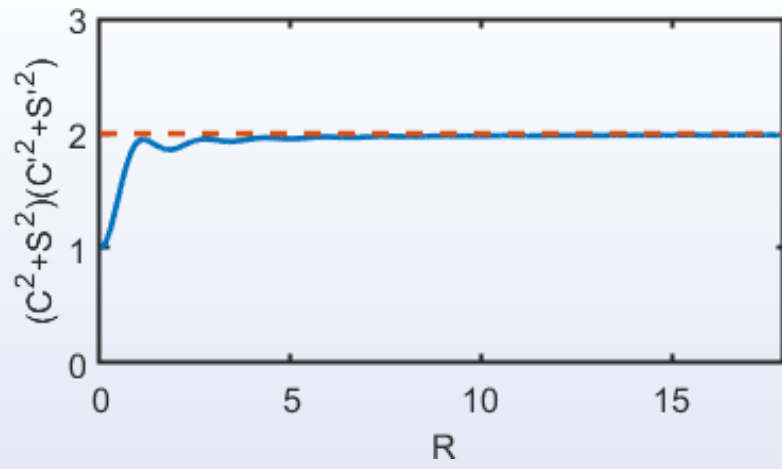
Approximation: Behavior of Mathieu functions





Approximation: Behavior of Mathieu functions



- Sinusoidal behavior
- Phase delay between Mathieu functions and their first derivative
- The product $(c_n^2 + s_n^2)(c'_n{}^2 + s'_n{}^2)$ saturates for $R \gg 1$
- Transfer matrix is unimodular (Liouville's theorem)



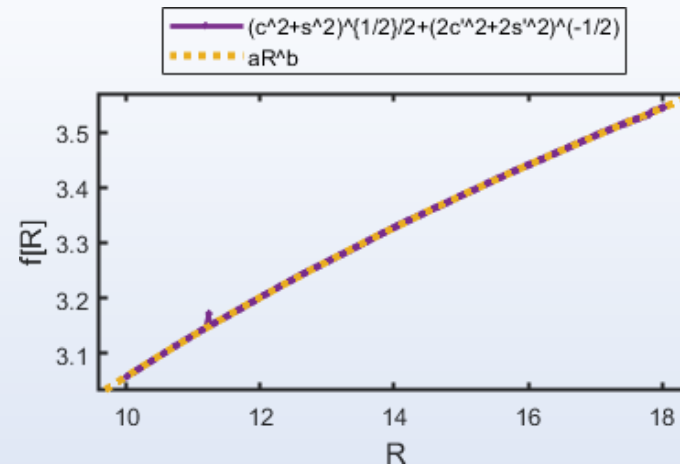
$$\begin{aligned}
 c_n &\approx a R^b \cos(2R + \omega) \\
 s_n &\approx a R^b \sin(2R + \omega) \\
 c'_n &\approx -\frac{1}{a \cos^2 \omega - \sin^2 \omega} R^{-b} \sin(2R - \omega) \\
 s'_n &\approx \frac{1}{a \cos^2 \omega - \sin^2 \omega} R^{-b} \cos(2R - \omega)
 \end{aligned}$$


$$(c_n^2 + s_n^2)(c_n'^2 + s_n'^2) = \left(\frac{1}{\cos^2 \omega - \sin^2 \omega} \right)^2 \approx 2 \quad \omega = \frac{\pi}{8}$$

Alignment: Evaluation of the constants

$$(c_n^2 + s_n^2)(c_n'^2 + s_n'^2) = \left(\frac{1}{\cos^2 \omega - \sin^2 \omega} \right)^2 \approx 2 \quad \omega = \frac{\pi}{8}$$

$$\frac{1}{2} \sqrt{c_n^2 + s_n^2} + \frac{1}{2} \sqrt{\frac{2}{c_n'^2 + s_n'^2}} = aR^b$$



$$a = \sqrt{3}$$
$$b = \frac{1}{4}$$

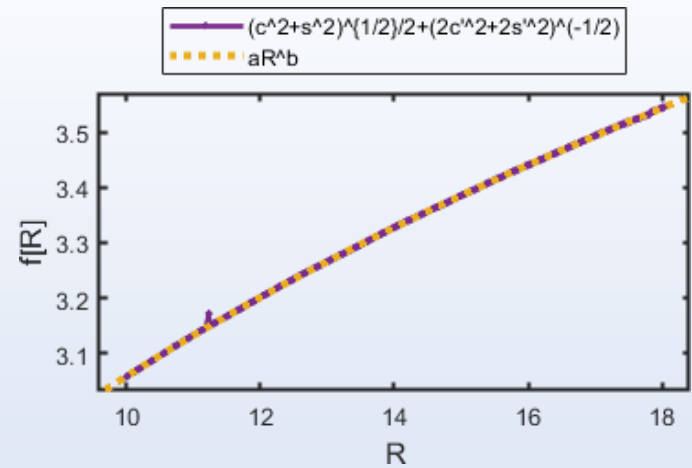
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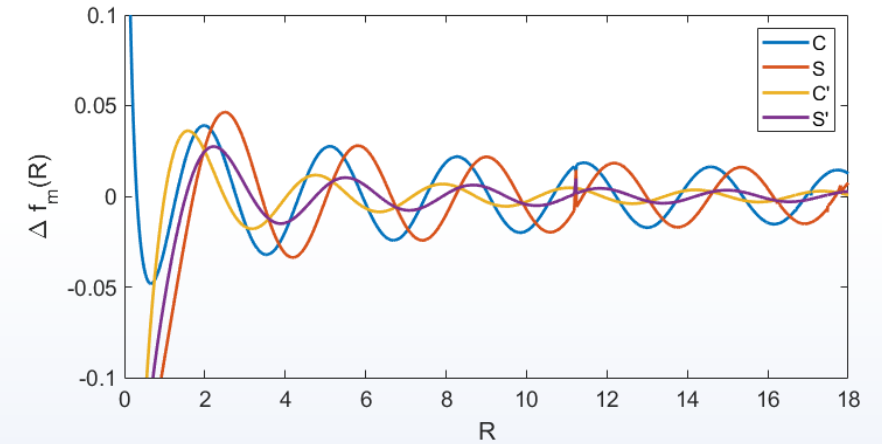
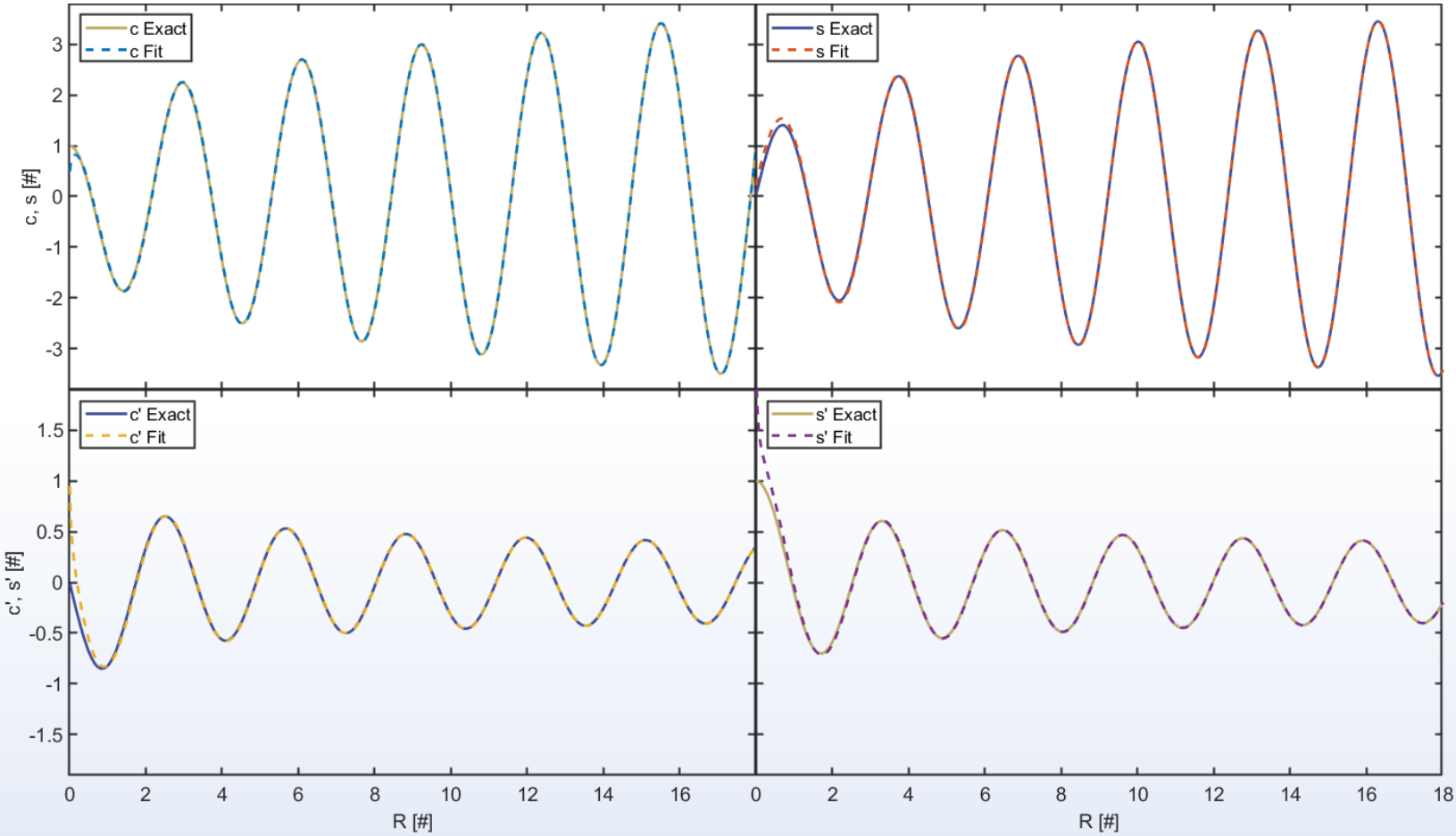
$$c_n \approx \sqrt{3} R^{\frac{1}{4}} \cos\left(2R + \frac{\pi}{8}\right)$$
$$s_n \approx \sqrt{3} R^{\frac{1}{4}} \sin\left(2R + \frac{\pi}{8}\right)$$
$$c_n' \approx -\sqrt{\frac{2}{3}} R^{-\frac{1}{4}} \sin\left(2R - \frac{\pi}{8}\right)$$
$$s_n' \approx \sqrt{\frac{2}{3}} R^{-\frac{1}{4}} \cos\left(2R - \frac{\pi}{8}\right)$$

$$\frac{1}{2} \sqrt{c_n^2 + s_n^2} + \frac{1}{2} \sqrt{\frac{2}{c_n'^2 + s_n'^2}} = aR^b$$



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Check: Numerical vs. Analytical Approximation



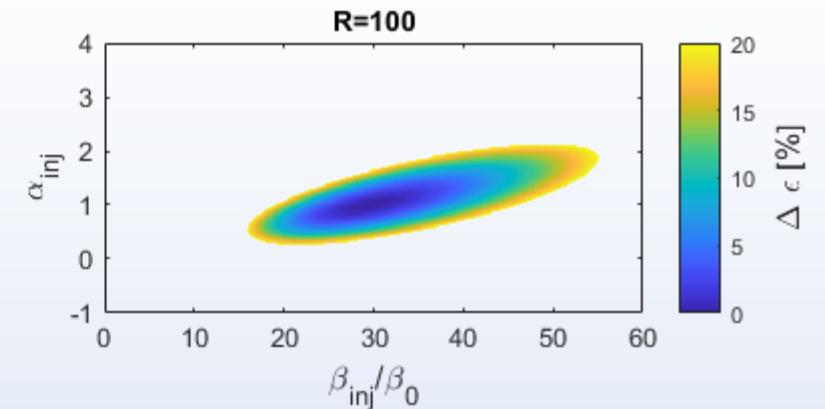
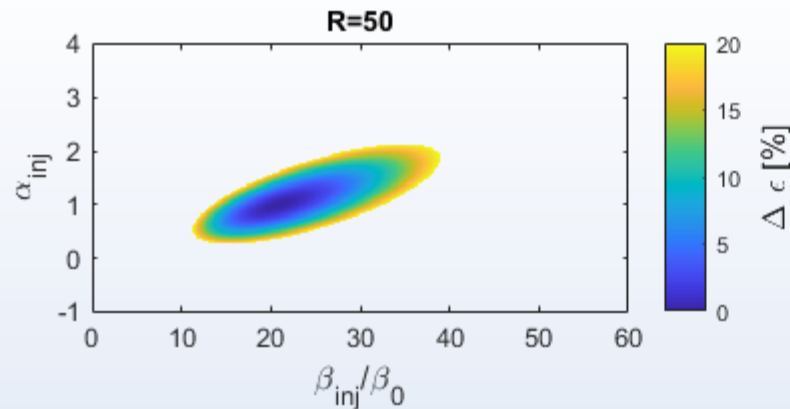
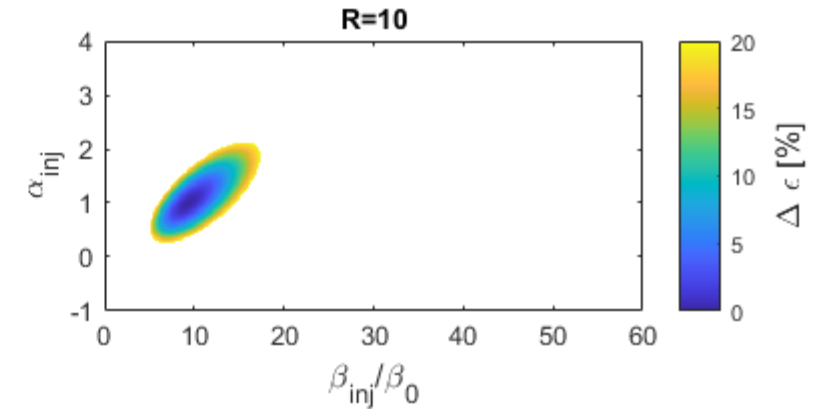
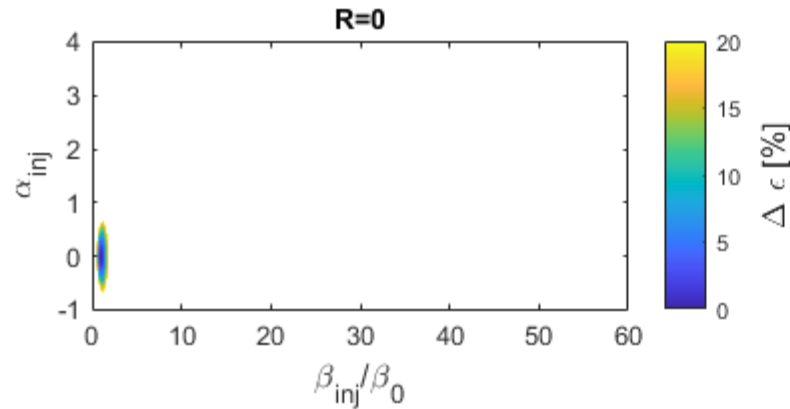
Transfer matrix

$$\begin{bmatrix} \beta^* \\ \alpha \\ \gamma^* \end{bmatrix} = \mathbf{M} \begin{bmatrix} \beta_{inj}/\beta_0 \\ \alpha_{inj} \\ \gamma_{inj}\beta_0 \end{bmatrix}$$

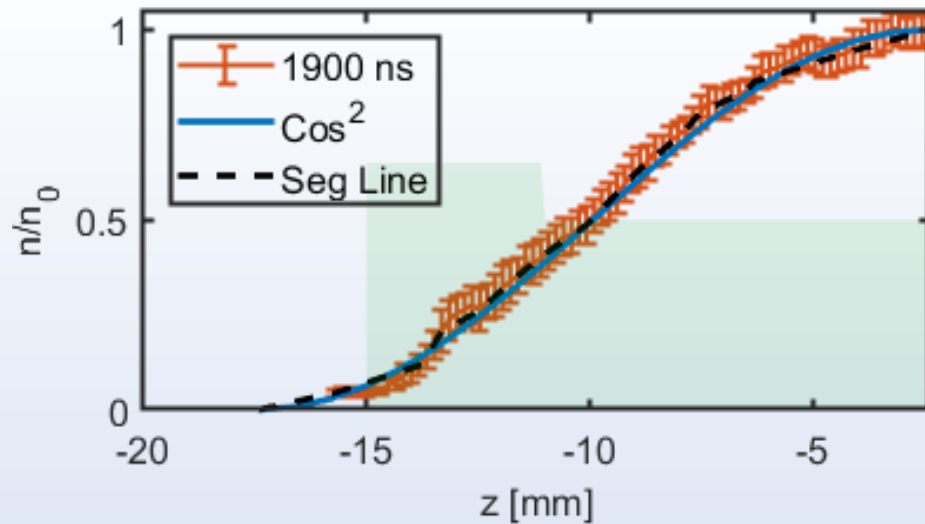
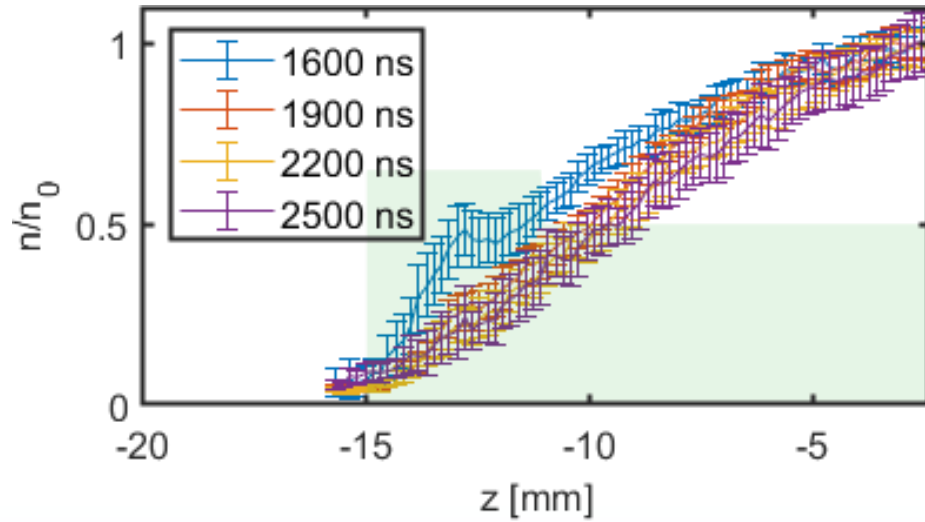
Mehrling-Floettmann equation for emittance growth

$$\varepsilon_{n,fin} = \frac{\varepsilon_{n,inj}}{2} \left(\frac{1 + \alpha^2}{\beta^*} + \beta^* \right)$$

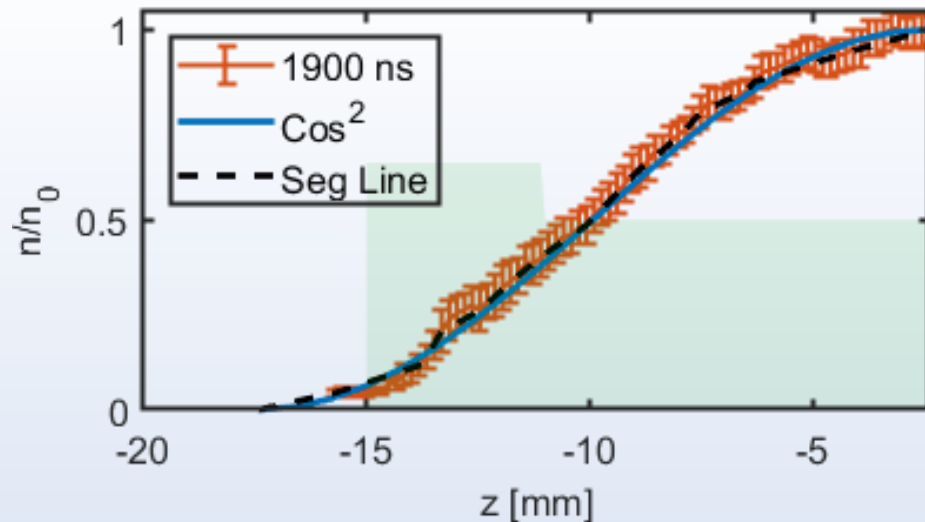
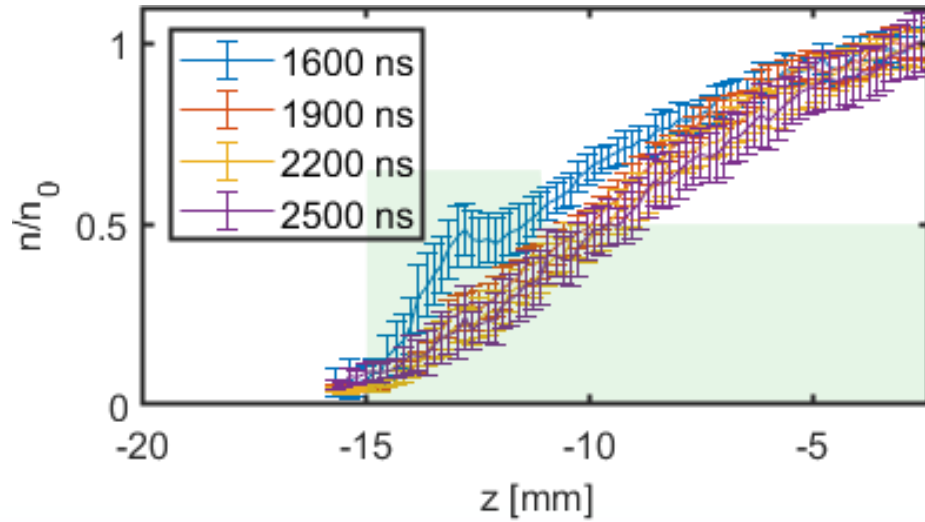
Mehrling, T., et al. "Transverse emittance growth in staged laser-wakefield acceleration." *Physical Review Special Topics-Accelerators and Beams* 15.11 (2012): 111303.



Perspective: Experimental ramps

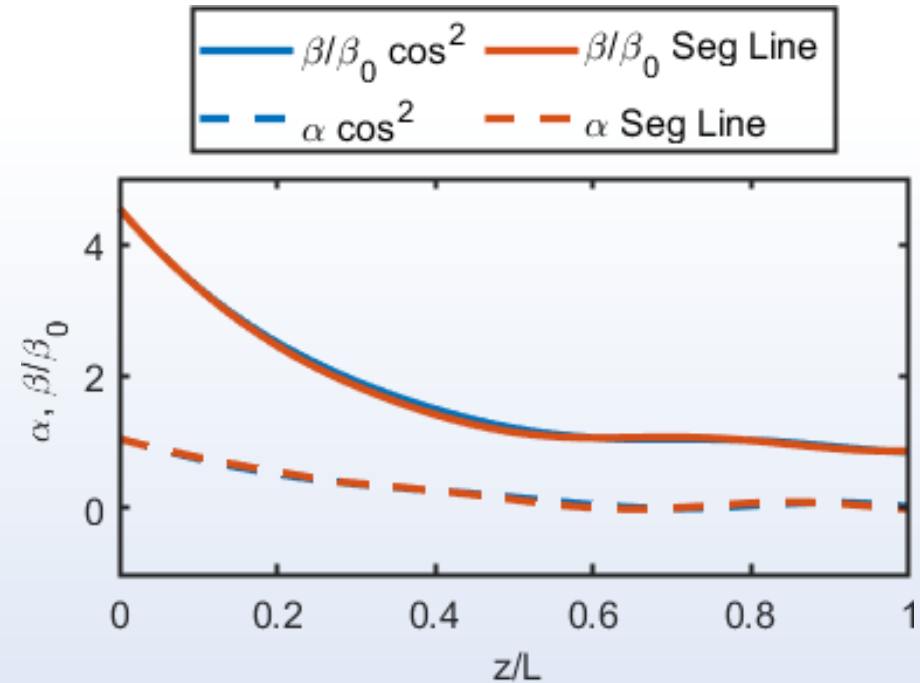


Perspective: Experimental ramps

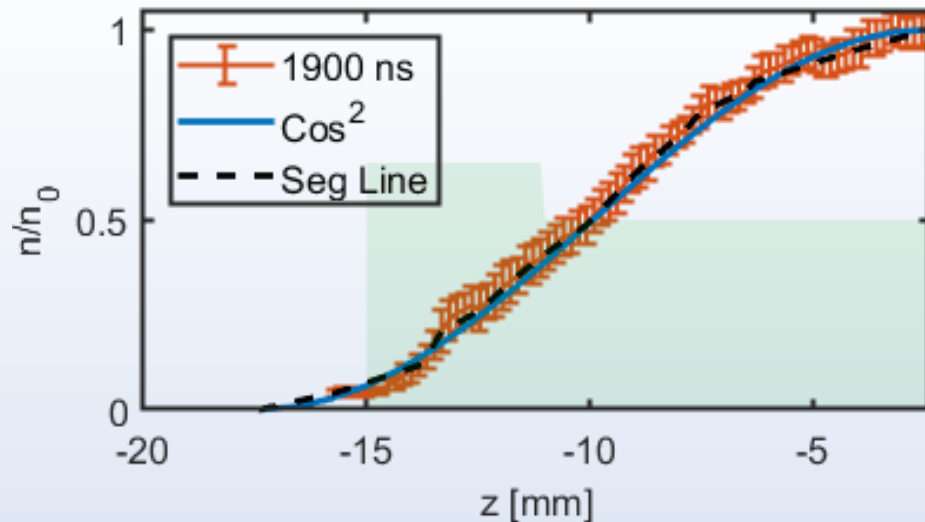
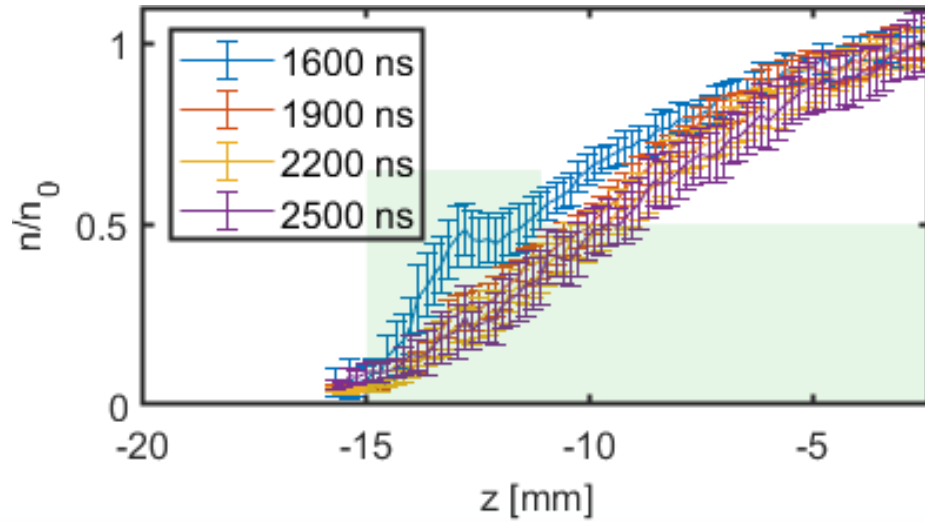


Plasma ramps with squared cosine shape were actually realized at the plasma lab of SPARC_LAB facility by means of tapered capillaries

The measured shapes were replayed in numerical simulations by means of segmented lines and tested in comparison with analytical squared cosine shape

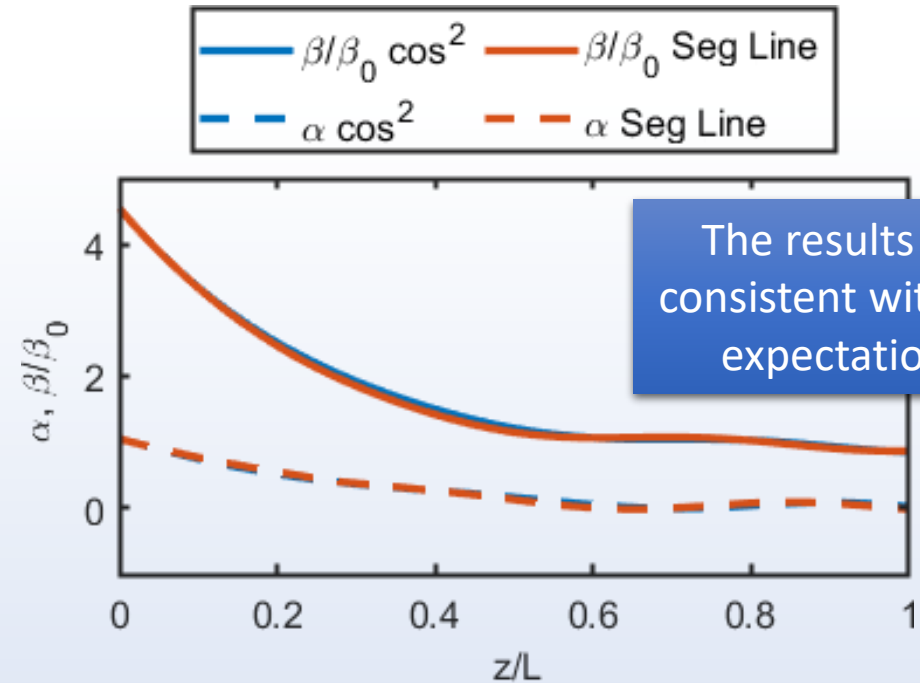


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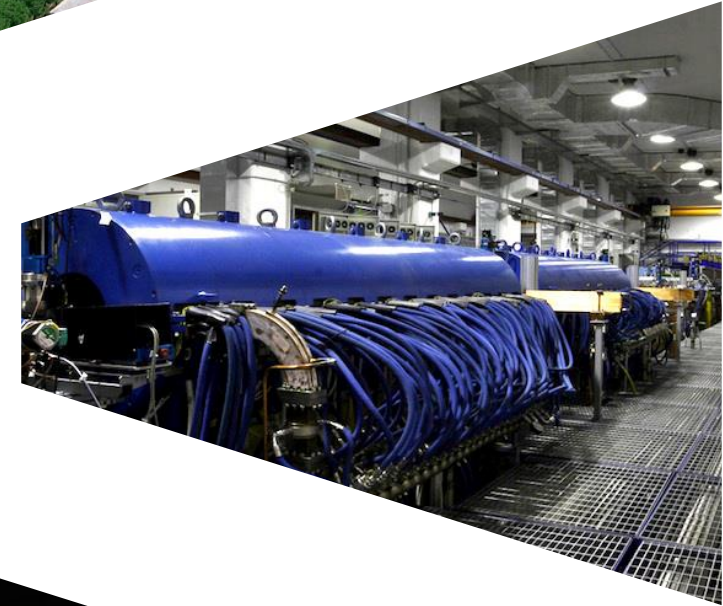
- ❑ We opted for a purely analytical approach trying to evaluate the transfer matrix of a plasma ramp
- ❑ We concluded that a ramp with a symmetric form could simplify the problem
- ❑ We opted for a ramp with a squared cosine shape
- ❑ We evaluated that this ramp is never adiabatic and its application lead to Mathieu Equation
- ❑ We compared the solution of the Mathieu Equation with actual envelope evolution
- ❑ We deduced that the ramp transfer matrix can be normalized to a geometric parameter
- ❑ We found an analytical approximation that allows to write the transfer matrix in a simple way

Tons of applications on their way!

- ❑ We opted for a purely analytical approach trying to evaluate the transfer matrix of a plasma ramp
- ❑ We concluded that a ramp with a symmetric form could simplify the problem
- ❑ We opted for a ramp with a squared cosine shape
- ❑ We evaluated that this ramp is never adiabatic and its application lead to Mathieu Equation
- ❑ We compared the solution of the Mathieu Equation with actual envelope evolution
- ❑ We deduced that the ramp transfer matrix can be normalized to a geometric parameter
- ❑ We found an analytical approximation that allows to write the transfer matrix in a simple way

Tons of applications on their way!

Thank you
for
the attention!



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