

Schemes for Simultaneous Large Transformer Ratio, High Efficiency, Low Energy Spread, High Charge of Accelerated Electron Beams by Tailored Wakefield Plateaus for Long Driver and Witness Bunches

R. T. Ovsianikov, V. I. Maslov, D. O. Shendryk, W. P. Leemans



Universität Hamburg



Plan

00 Problematic

01 Configurations for optimal parameters

- One driver and three witnesses bunch
- One driver and two witnesses bunch
- One driver and one witness bunch
- Two drivers and one witness bunch
- One driver and second bubble's witness witness

02 Instabilities' suppression approaches

03 Conclusions

Simulation info:

- 2.5D particle-in-cell numerical simulations via LCode;
- macro-particles;
- Gaussian electron distribution;

Units of measure

n_0 - for all densities,
 $\sqrt{4\pi n_0 mc^2}$ - for all fields,
 mc^2/e - for wakefield potential,
 $c\omega_p$ - for all distances and emittance,
 mc - for beam and plasma momenta,
 mc^2/ω_p - for beam angular momentum,
 $n_0 mc^2$ - for energy densities,
 $n_0 mc^4/\omega_p^2$ - for integral energies,
 $n_0 mc^4/\omega_p$ - for energy flux densities,
 $n_0 mc^5/\omega_p^2$ - for full energy flux

$$n_0 = 10^{18} \text{cm}^{-3}$$

$$r_b = 1.7 \mu\text{m}$$

Problematic

- *Plasma accelerators are a centimeter-scale source of GeV beams.*
→ Leemans et al., Nature Physics 2, 696 (2006)

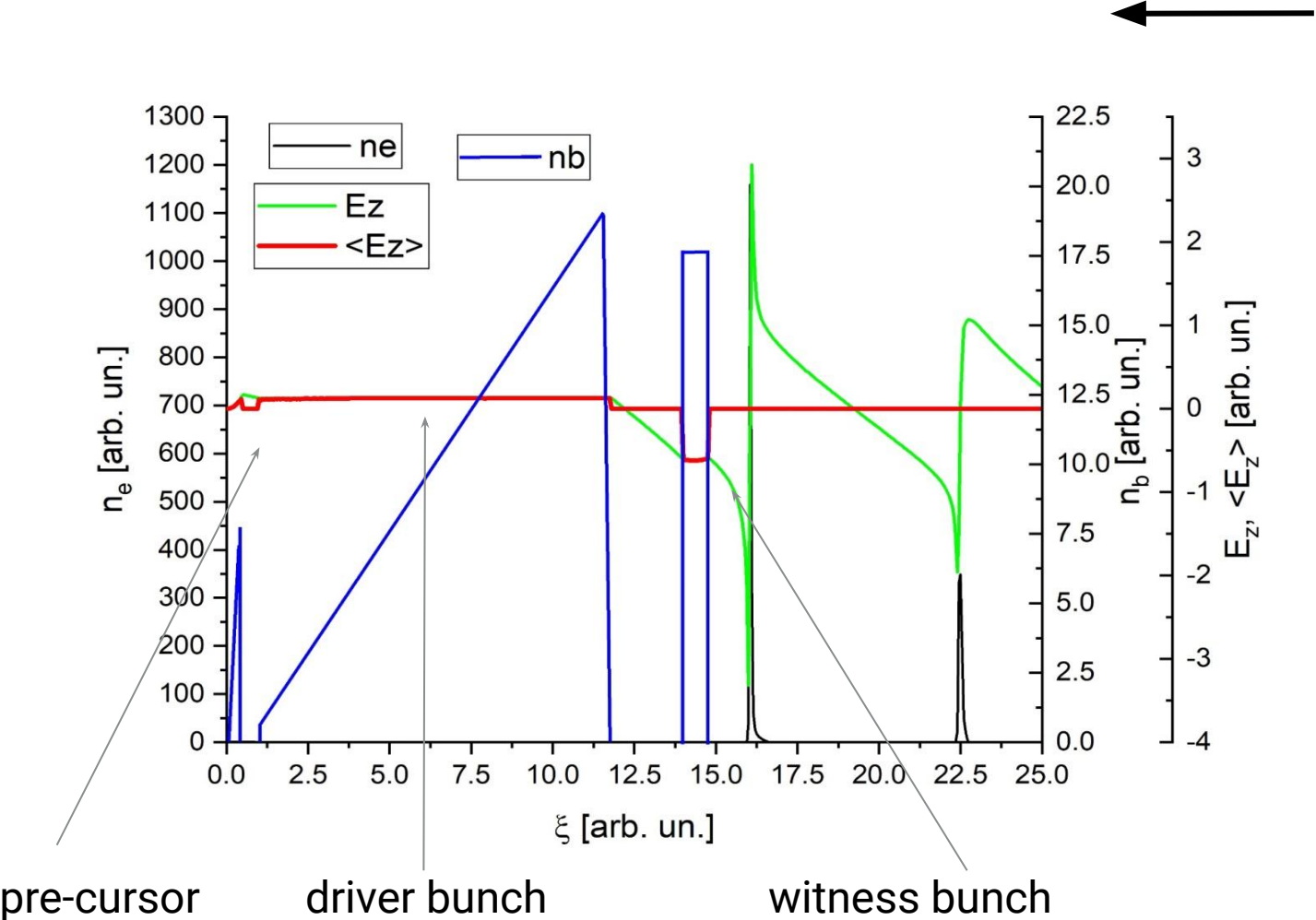
Conventional accelerators

- breakdown of the RF structure (100 MV/m)
- big sizes
- the energy loss by synchrotron radiation

Plasma wakefield accelerators

- gigavolt-per-metre (100 GV/m) electric fields
- compactness
- cheapness
- low repetition rate

Long Driver-bunch, Large Transformer Ratio, Plateau for Driver-bunch and Plateau for Witness-bunch but **Small Charge** of Accelerated Electrons and **Small Accelerator Efficiency**



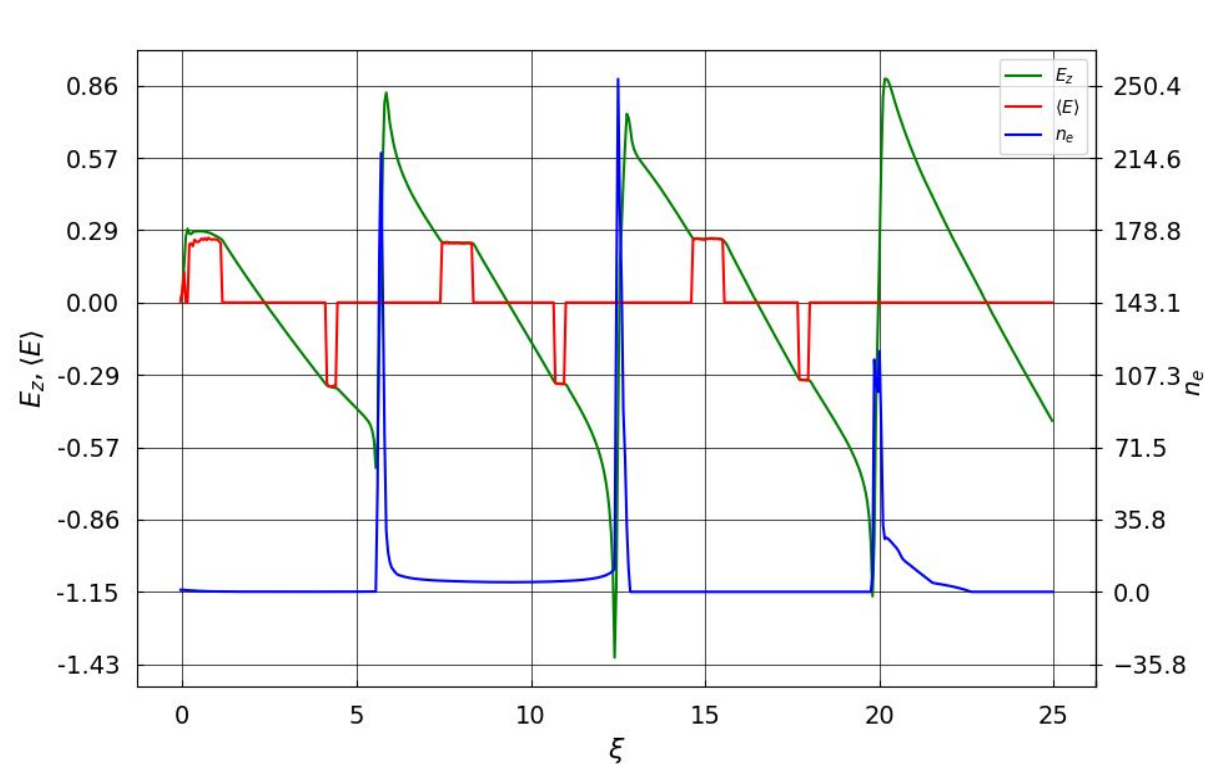
Simultaneous requirements:

- the long profiled driver-bunches;
- a large total charge of accelerated electrons;
- large transformer ratio;
- high efficiency.

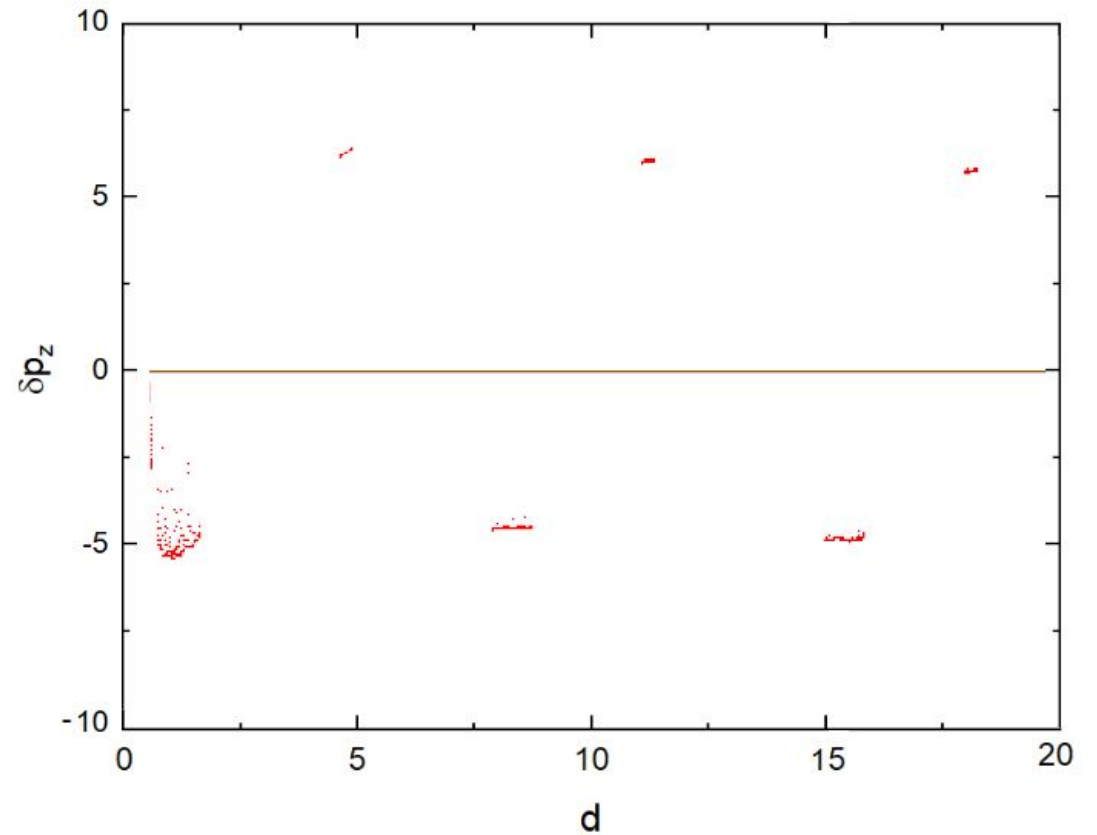
Three pairs driver-witness in three bubbles.

Identical plateaus on the decelerating and accelerating wakefields for the drivers and witnesses.

Large charge of accelerated electrons but **small accelerator efficiency.**



*The dispersion of the witness
z-momentum is 0.02%*

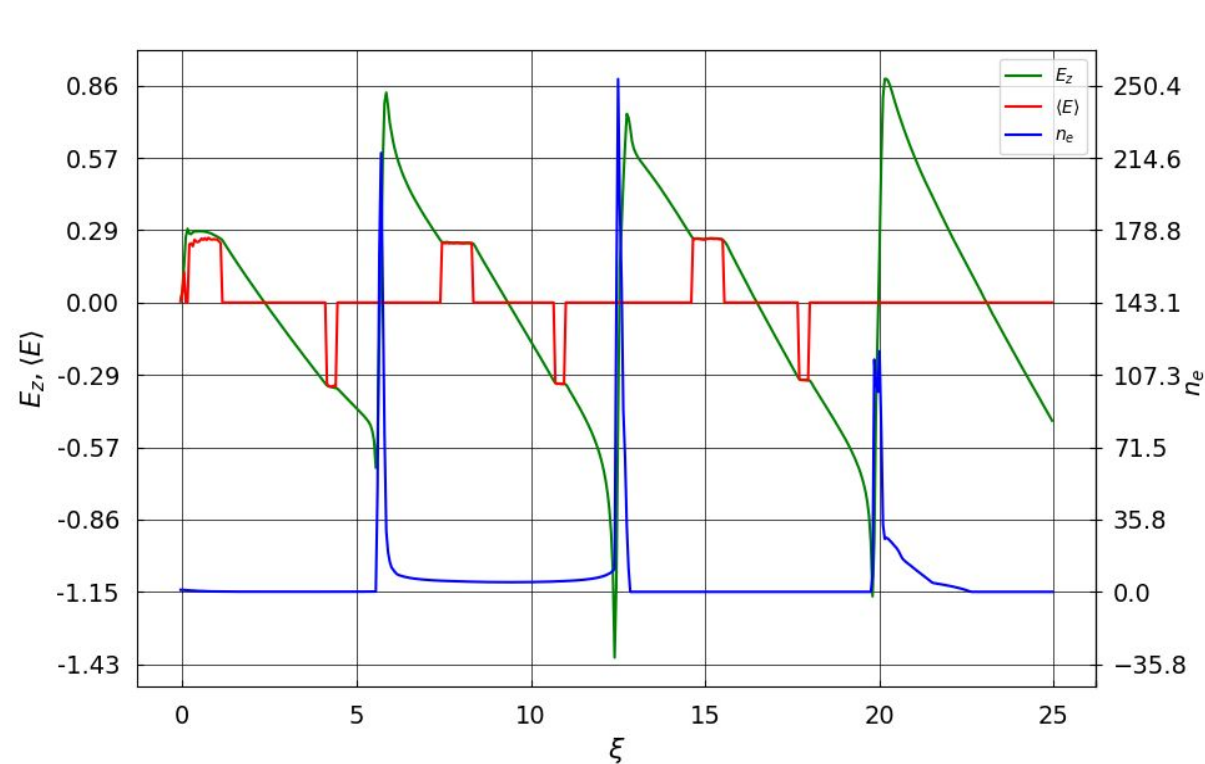


*Phase space portrait of the bunches at the
some instant. $\delta p = p_z - p_{z0}$*

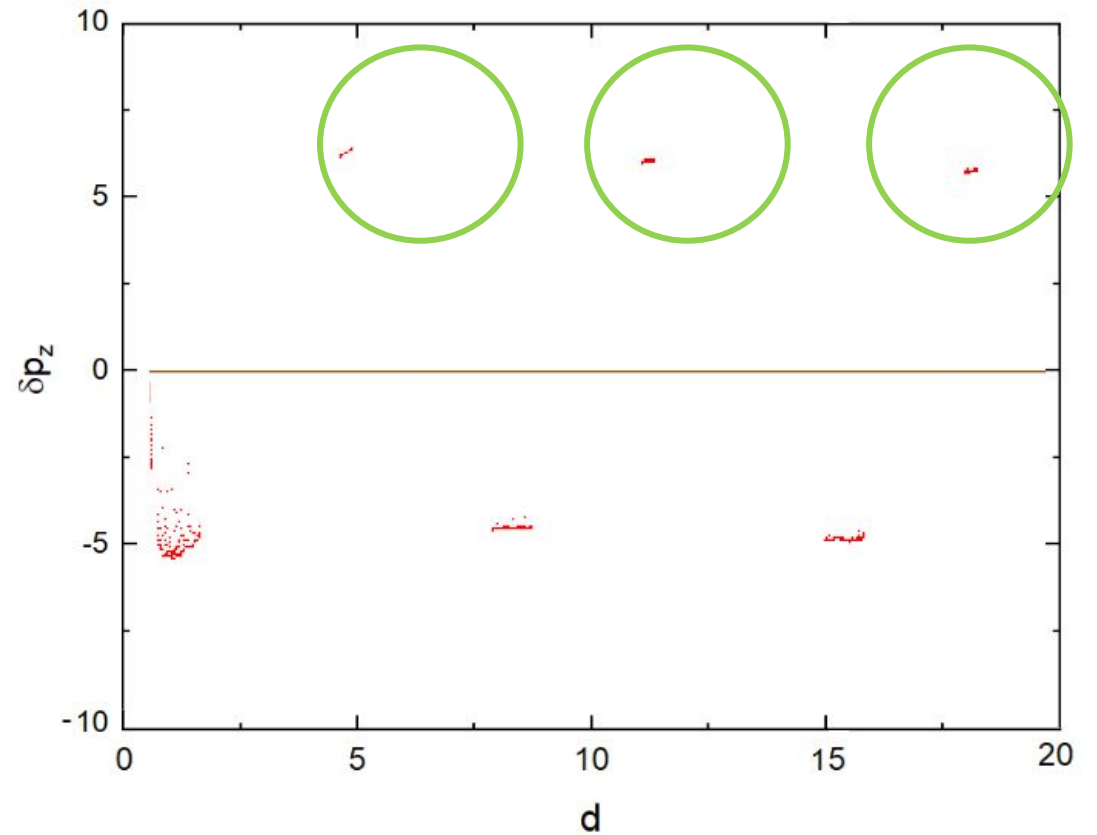
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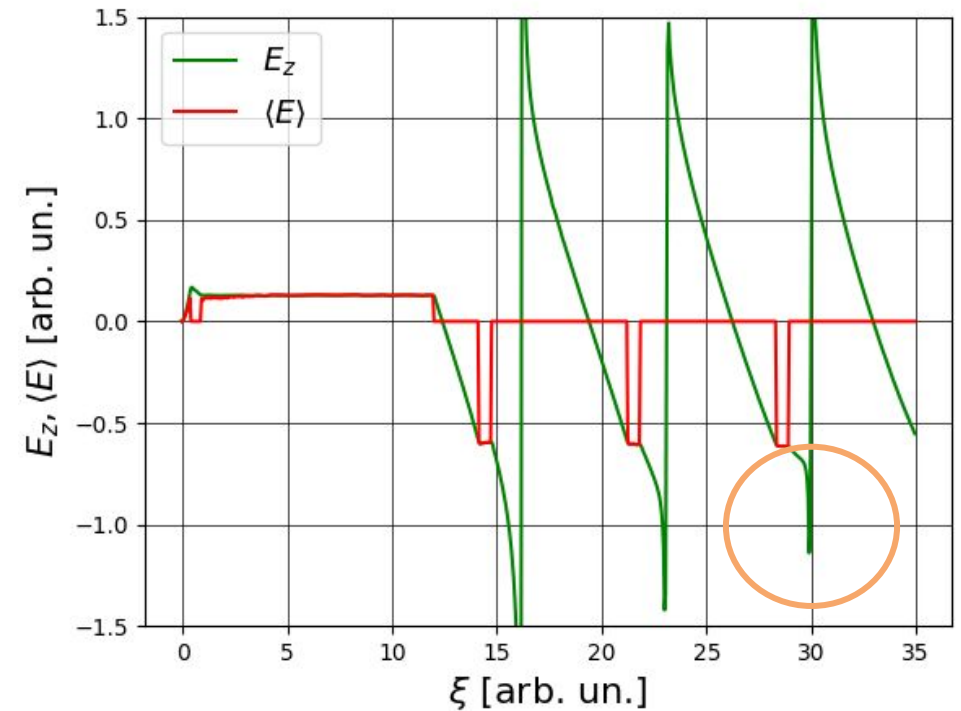
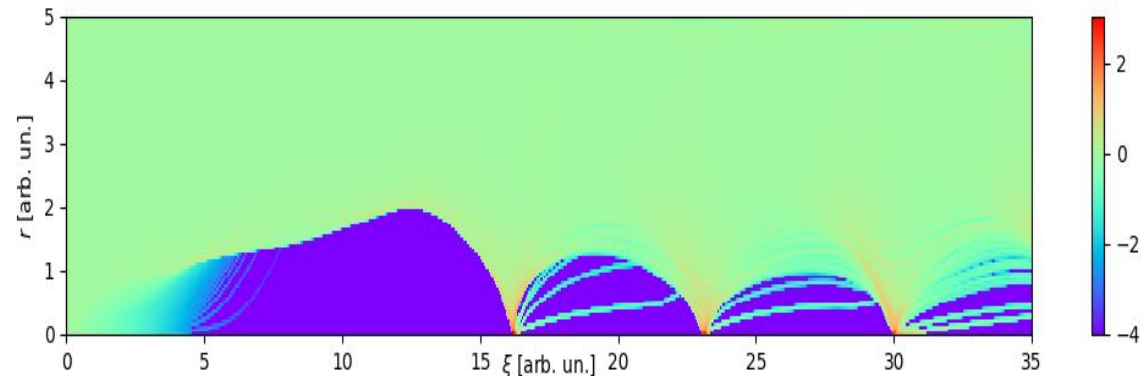
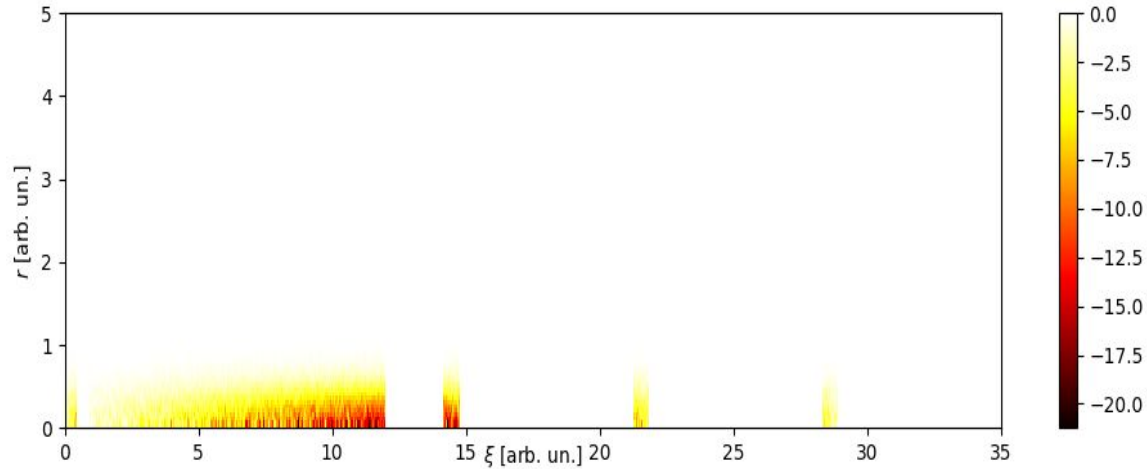
*The dispersion of the witness
z-momentum is 0.02%*



*Phase space portrait of the bunches at the
some instant. $\delta p = p_z - p_{z0}$*

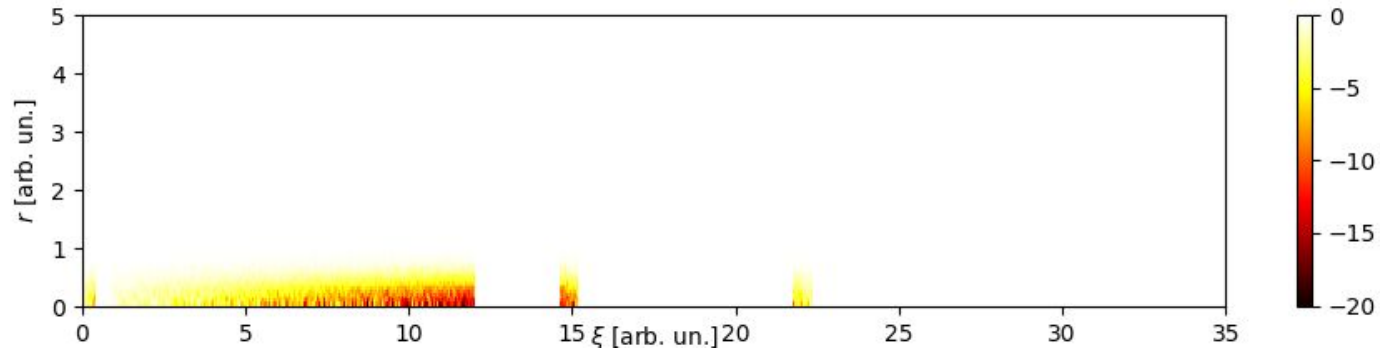
Large Charge of Accelerated Electrons, Large Transformer Ratio, and High Efficiency with Plateaus

for the one Long Driver-bunch and for the three Witness-bunches

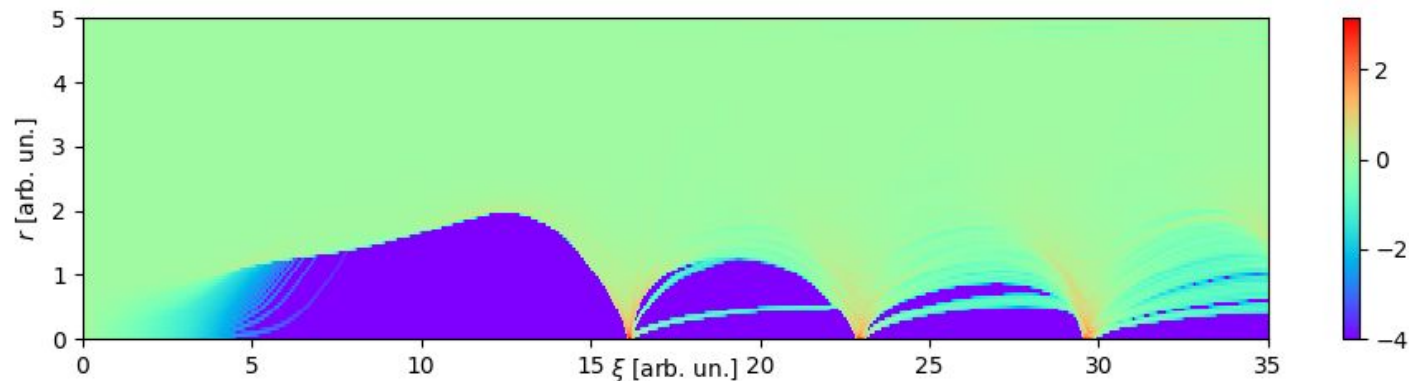


The maximum current of bunch-driver: $I_b = 10.7 \text{ kA}$.
Transformer ratio: 4

Large Charge of Accelerated Electrons at Large Transformer Ratio and High Efficiency at Plateaus for One Driver-bunch and Two Witness-bunches

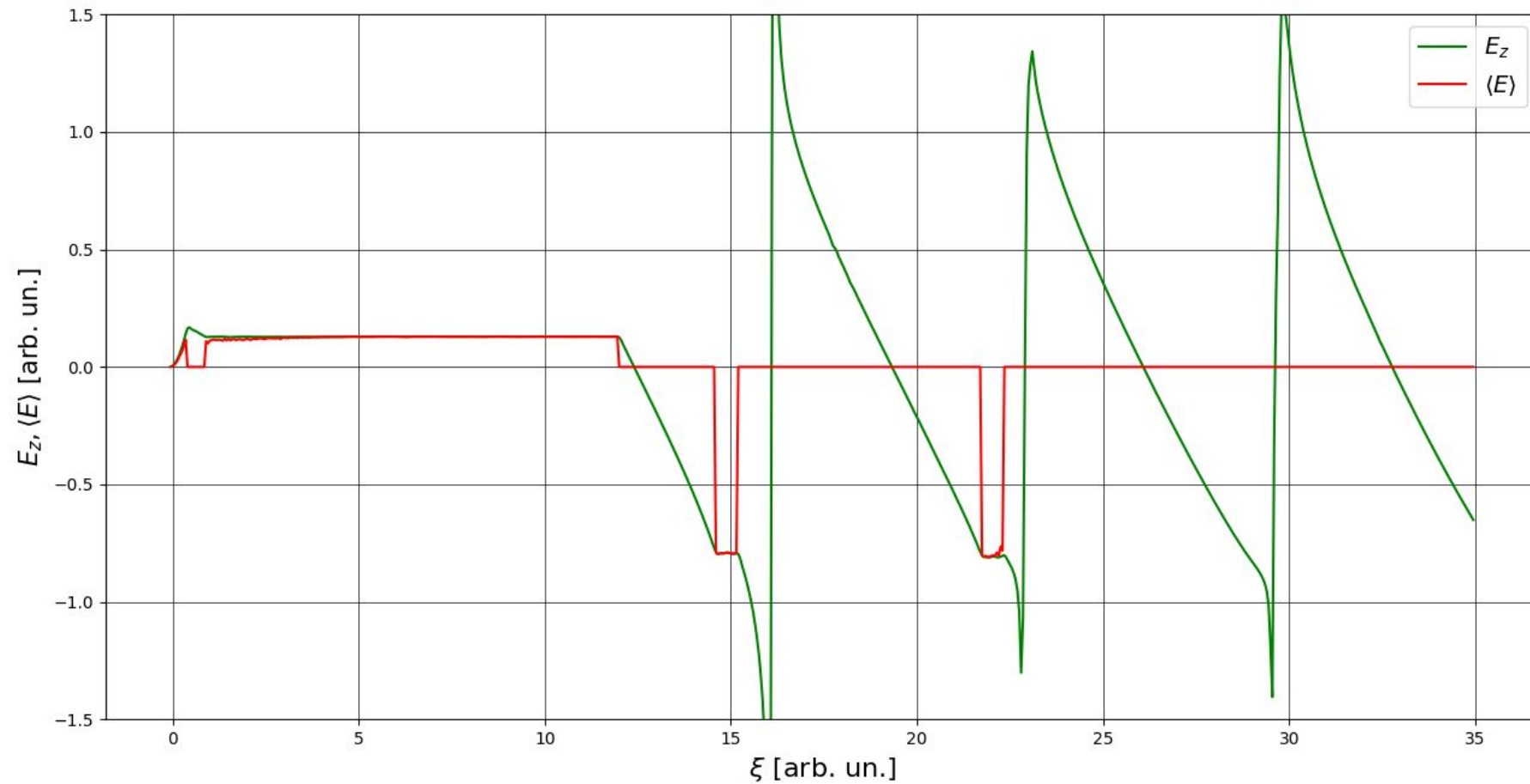


Spatial distribution of the density of two long profiled driver-bunches and of the profiled witness-bunch



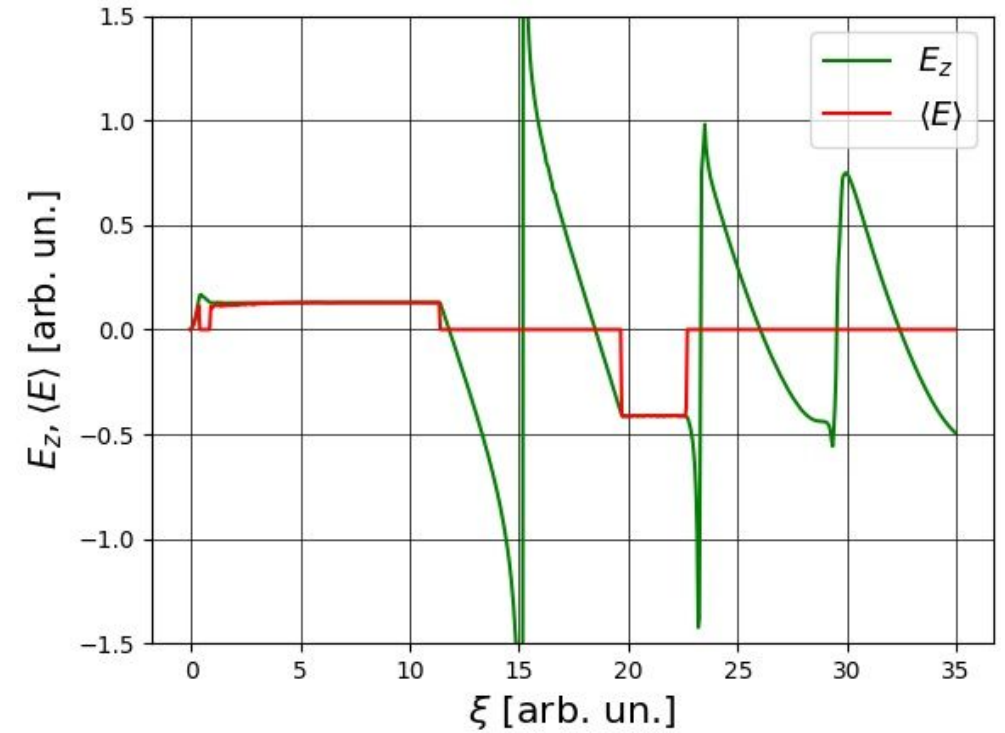
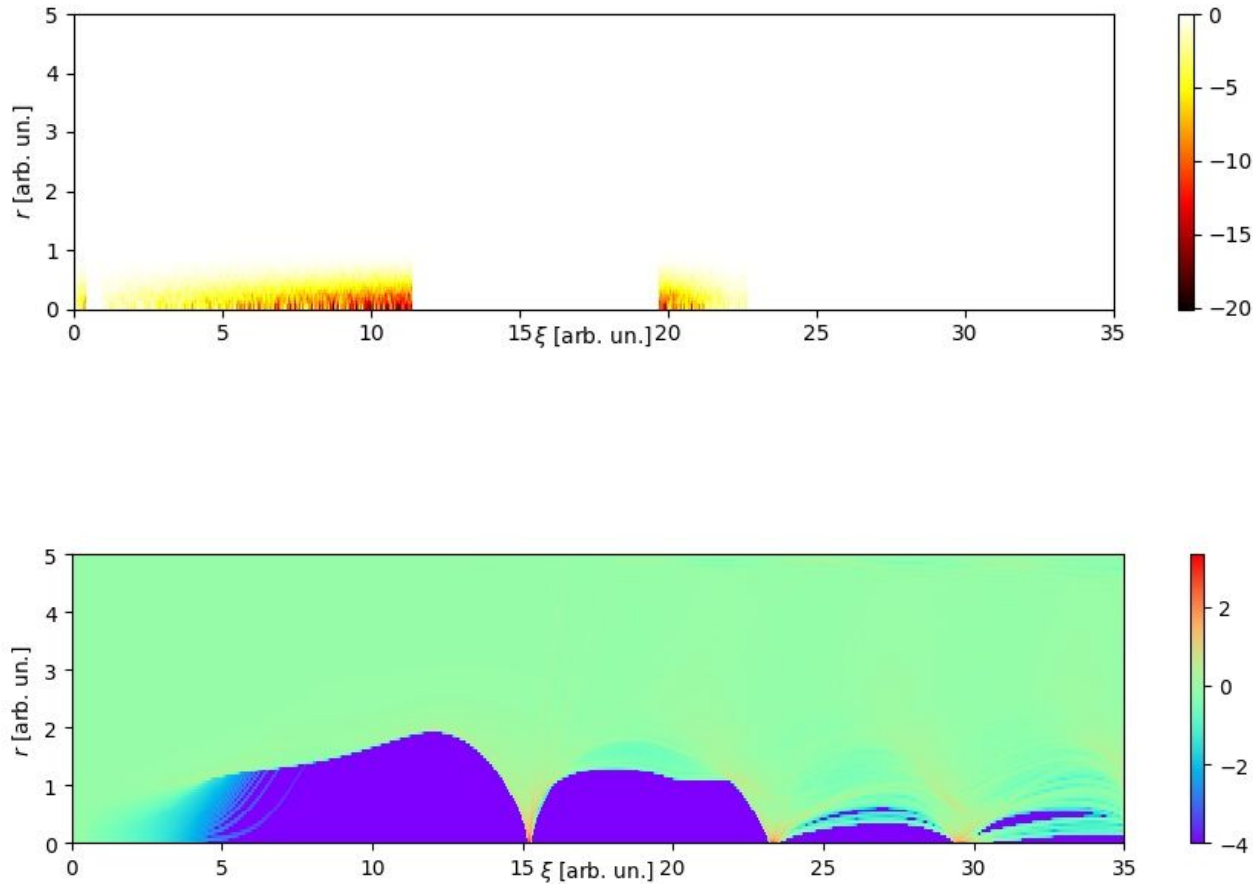
Spatial distribution of the plasma electron density, excited by two long profiled driver-bunches and reconstructed by the profiled witness-bunch in blowout regime

Large Charge of Accelerated Electrons at Large Transformer Ratio and Low Efficiency at Plateaus for One Driver-bunch and Two Witness-bunches



The maximum current of bunch- driver: $I_b = 10.7 \text{ kA}$.
Transformer ratio: 6.1

Large Charge of Accelerated Electrons at Large Transformer Ratio and High Efficiency at Plateaus for One Driver-bunch and One Witness-bunch

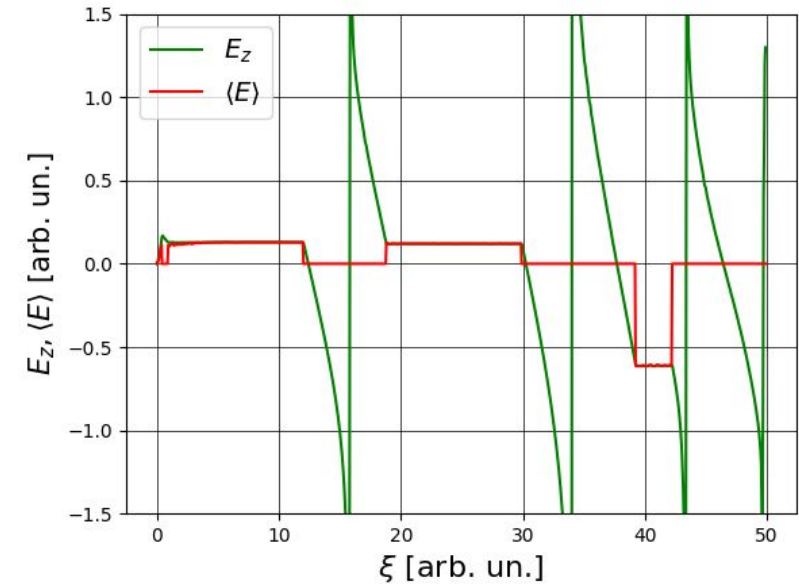
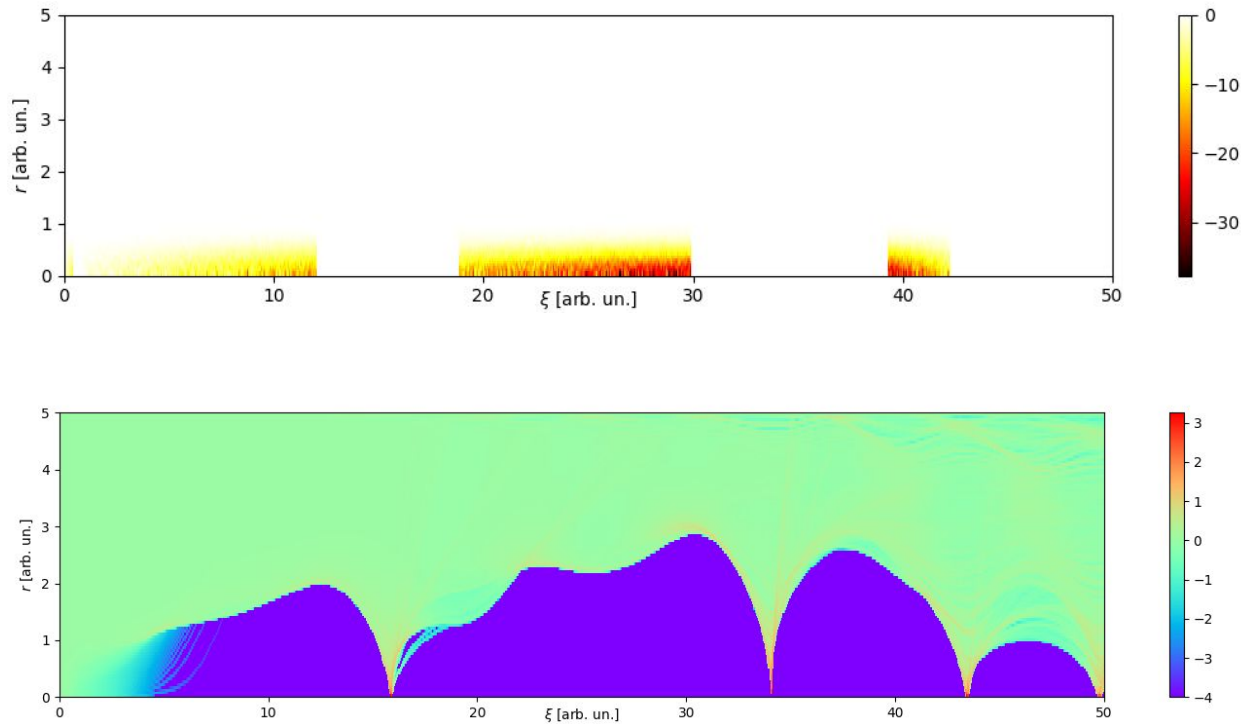


The maximum current of the bunch-driver: $I_b = 10.1$ kA.
Transformer ratio: 3.2

Modified and Developed Schemes

Large Charge of Accelerated Electrons at Large Transformer Ratio and High Efficiency at Plateaus

for Two Long Driver-bunches and for the Witness-bunch in the Accelerating Wakefield



The maximum current of the first bunch-driver is $I_b = 10.7 \text{ kA}$ and for second $I_b = 19.7 \text{ kA}$. The transformer ratio is 4.5

Instabilities in plasma wakefield accelerators

- Betatron oscillations
- Moving ions
- Hosing instability
- Low radius bunches
- Heavy atoms
- Nonlinear or close to zero focusing force

S.Diederichs, C.Benedetti, E.Esarey, M.Thevenet, J.Osterhoff, C. B. Schroeder. Phys. Plasmas. 29, 043101 (2022)

Hosing Instability in the Blow-Out Regime for Plasma-Wakefield Acceleration

C. Huang,¹ W. Lu,¹ M. Zhou,¹ C. E. Clayton,¹ C. Joshi,¹ W. B. Mori,¹ P. Muggli,² S. Deng,² E. Oz,² T. Katsouleas,² M. J. Hogan,³ I. Blumenfeld,³ F. J. Decker,³ R. Ischebeck,³ R. H. Iverson,³ N. A. Kirby,³ and D. Walz³

¹University of California, Los Angeles, California 90095, USA

²University of Southern California, Los Angeles, California 90089, USA

³Stanford Linear Accelerator Center, Menlo Park, California 94025, USA

The electron hosing instability in the blow-out regime of plasma-wakefield acceleration is investigated using a linear perturbation theory about the electron blow-out trajectory in Lu *et al.* [in Phys. Rev. Lett. **96**, 165002 (2006)]. The growth of the instability is found to be affected by the beam parameters unlike in the standard theory Whittum *et al.* [Phys. Rev. Lett. **67**, 991 (1991)] which is strictly valid for preformed channels. Particle-in-cell simulations agree with this new theory, which predicts less hosing growth than found by the hosing theory of Whittum *et al.*

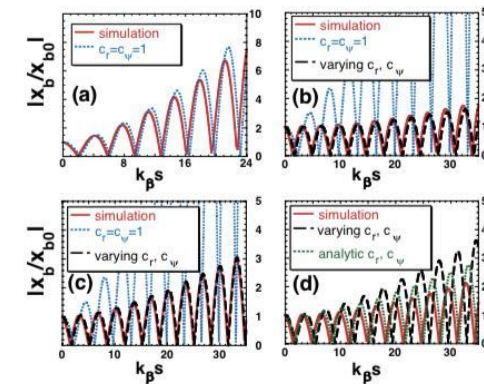
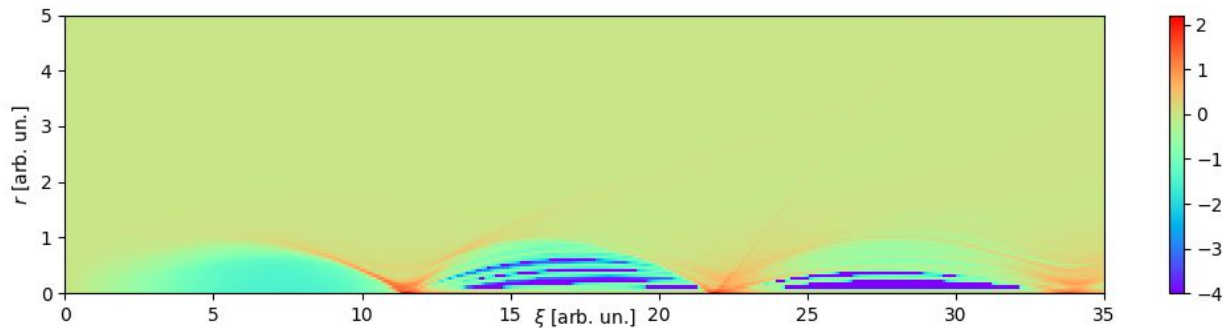
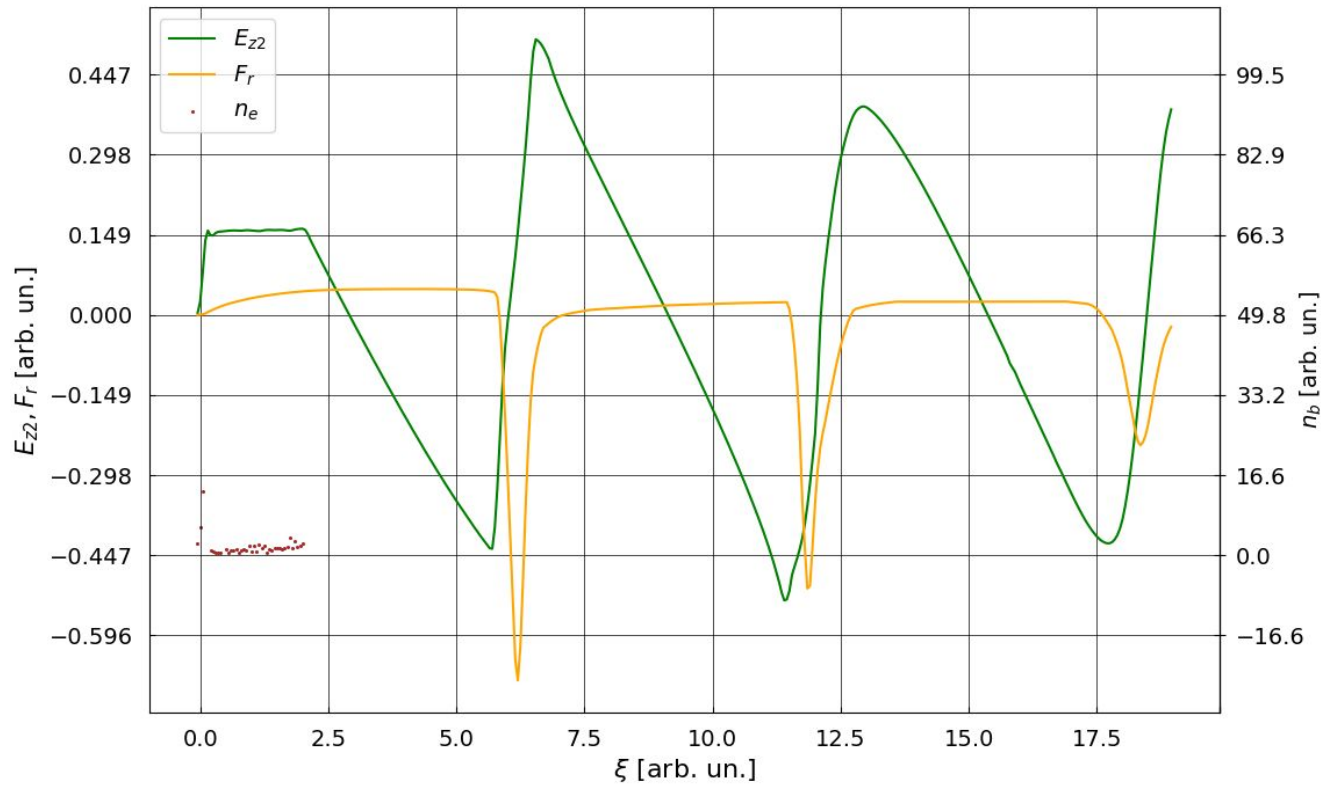


FIG. 3 (color). Hosing growth in four regimes. x_{b0} is the initial displacement of the beam centroid. We assume $\Delta_s = 0.1r_0$, $\Delta_L = 1c/\omega_p$ for the analytic curve in case (d). The slightly slower hosing growth in simulation (d) is caused by nonlinearity in beam-channel centroid coupling.

Weakly Nonlinear Mode with Narrow Bunches

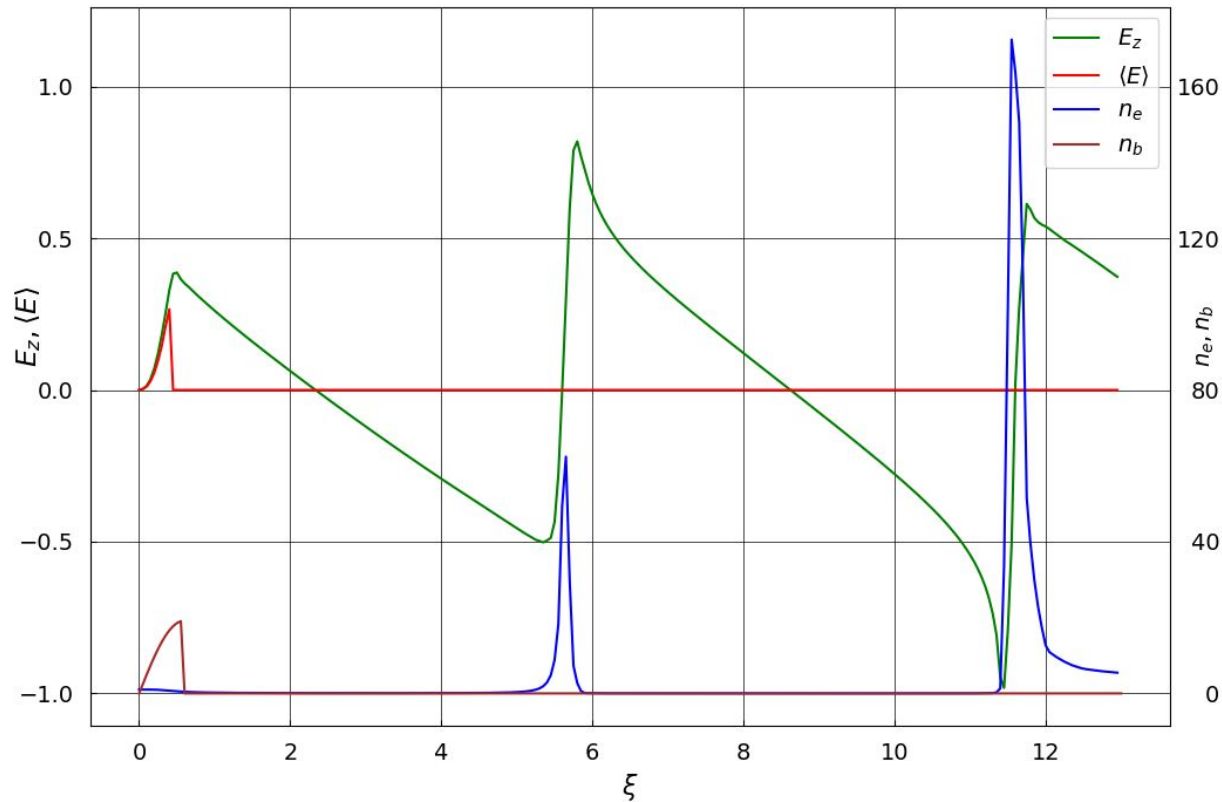


Analytical investigations

T.J.Mehrling, R.A.Fonseca, A.Martinez de la Ossa, J.Vieira. Phys. Rev. Lett. 118(2017)174801;

A.Knetsch, B.Sheeran, L.Boulton et al. Phys. Rev. Accel. Beams. 24, 101302 (2021);

Regularization and Advantages of Wakefield in a Weakly Nonlinear Regime with Narrow Bunches



$$1d \quad \frac{dE_z}{dz} \approx 4\pi en_b \rightarrow E_z \sim \xi - \xi_w(t),$$

$$3d \quad 4\pi r^2 E_r = \frac{4\pi}{3} en_i r^3$$

$$E_r = E_z = \frac{en_i}{3} r \sim r \sim z - z_{bub}.$$

$$E_r = E_z = -\frac{en_b}{3} r \sim -r \sim z_b - z.$$

from V. Maslov, R. Ovsianikov, Numerical simulation of plateau formation by an electron bunch on the distribution of an accelerating wakefield in a plasma

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- R. Assmann, E. Gschwendtner, K. Cassou CERN Yellow Reports: Monographs, v. 1, p. 91, 2022.
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Diana.Shendryk@ruhr-uni-bochum.de