

Optimised density tailoring for dephasing mitigation in laser wakefield accelerators

in simulations...

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Dephasing in Laser Wakefield Accelerators



The usable phase region for an electron LWFA is accelerating and focusing.

For the linear regime

 $a_0 \ll 1 \qquad \omega_p \ll \omega_L$

The group velocity of the laser is

$$v_g = c\sqrt{1 - \omega_p^2/\omega_L^2}$$

This sets the phase velocity of the wake and so the dephasing length is

$$L_{\phi} = \frac{\pi c}{2\omega_p} \frac{1}{\beta_g^{-1} - 1}$$
$$L_{\phi} \approx \pi c \frac{\omega_L^2}{\omega_p^3}$$

Dephasing in Laser Wakefield Accelerators



• Maximum electron energy reached for electrons injected behind peak field position

FBPIC simulation: $a_0 = 0.95$, guiding channel with $n_e = 5 \times 10^{17}$ cm⁻³

FBPIC simulations: R. Lehe et al., CPC, 2016

Dephasing can be removed in a plasma density gradient





Non-Linear regime

$$a_0 \ge 1 \qquad \omega_p \ll \omega_L$$

$$v_g < c\sqrt{1 - \omega_p^2/\omega_L^2}$$

$$L_{\phi} \neq \frac{\pi c}{2\omega_p} \frac{1}{\beta_g^{-1} - 1}$$

Dephasing can be removed in a plasma density gradient



Increasing plasma density reduces wavelength of plasma wave and rephases electrons

Dephasing can be prevented in a plasma density gradient

For $n_e \ll n_c$, the requirement for preventing dephasing [1-4]. $\frac{d}{dz} \left(\frac{\phi}{\omega_r(z)}\right) = \frac{1}{c} - \frac{1}{v_c} \qquad \phi = \text{electron phase in LWFA}$

For $a_0 \ll 1$, Pukhov [3] calculated the ideal density profile as, $n_e(z) = \frac{n_{e0}}{(1 - z/L_{\rm inh})^{2/3}}$ $L_{\rm inh} = \frac{c}{\omega} \left(\frac{n_c}{n_{e0}}\right)^{3/2} \frac{2\phi}{3}$



For highly non-linear LWFA, the above theory is not valid and the pulse evolution makes analytical treatment very difficult. Some experiments and simulations have shown approximate tailoring can still help [5-8].

- [1] T. Katsouleas, *Physical Review A* **33**, 2056–2064 (1986).
- [2] P. Sprangle et al. Phys. Rev. E 63, 056405 (2001).
- [3] A. Pukhov *Physical Review E* **77**, 025401 (2008).
- [4] W. Rittershofer et al. Phys. Plasmas 17 (6) 063104 (2010)
- [5] Z. Zhang et al New J. Phys. 17 103011 (2015)
- [6] E. Guillaume, et al. Phys. Rev. Lett. 115, 155002 (2015)
- [7] A. Döpp et al. Phys. Plasmas 23 (5): 056702 (2016)
- [8] C. Aniculaesei *et al.* Sci Rep 9, 11249 (2019).

Determining the ideal density profile "on-the-fly"



Calculating the next density step



Case A: Quasi linear LWFA with external injection



- $a_0 = 0.95$, guiding channel
- Initial on-axis plasma density $n_e = 5 \times 10^{17} \text{ cm}^{-3}$
- Initial laser energy = 1 J

Case A: Quasi linear LWFA with external injection



- $a_0 = 0.95$, guiding channel
- Initial on-axis plasma density n_e = 5 x 10¹⁷ cm⁻³
- Initial laser energy = 1 J



Ideal density matches theory until a_0 increases

Large increase in energy gain – especially for acceleration in 2nd period

Case B: Non-linear LWFA with external injection



Case C: Highly Non-linear LWFA with ionisation injection



4

3

2

0

 $-n_e = 1.2 \times 10^{18} \text{ cm}^{-3}$

10

8

6 z [mm] 12

 $-\phi = 3/4 \text{ OTF}$

 Produces 4 GeV, compared to 1.5 GeV from best constant density case

Summary and Outlook

- Simulations with on-the-fly density tuning can prevent dephasing even in highly non-linear LWFA
- Optimal tuning results in much higher electron energies
- This could allow LWFA to be purely depletion limited implying high efficiency
- Further work required to look at beam loaded cases
- Could potentially find tuning for other beam quantities
 - E.g. emittance preservation
 - Optimise beam loading
- This approach can help optimise experimental design practical realisation of density profiles is an important challenge