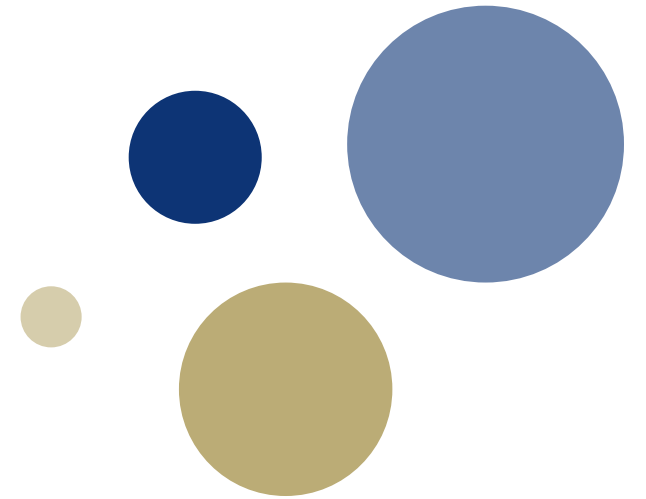




Norwegian University of
Science and Technology



Do the LHAASO Galactic diffuse emission data require a contribution from unresolved sources?

Presented by: Vittoria Vecchiotti

Based on work done in collaboration with: G. Peron, E. Amato, G. Morlino, S. Menchiari, F. L. Villante and G. Pagliaroli

Outline:

1. Galactic gamma-ray diffuse emission model;
2. Source model, compatibility with LHAASO KM2A measurements, unresolved source contribution;
3. Comparisons with LHAASO diffuse emission measurements.

Cataldo et al. Astrophys.J. 904 (2020)

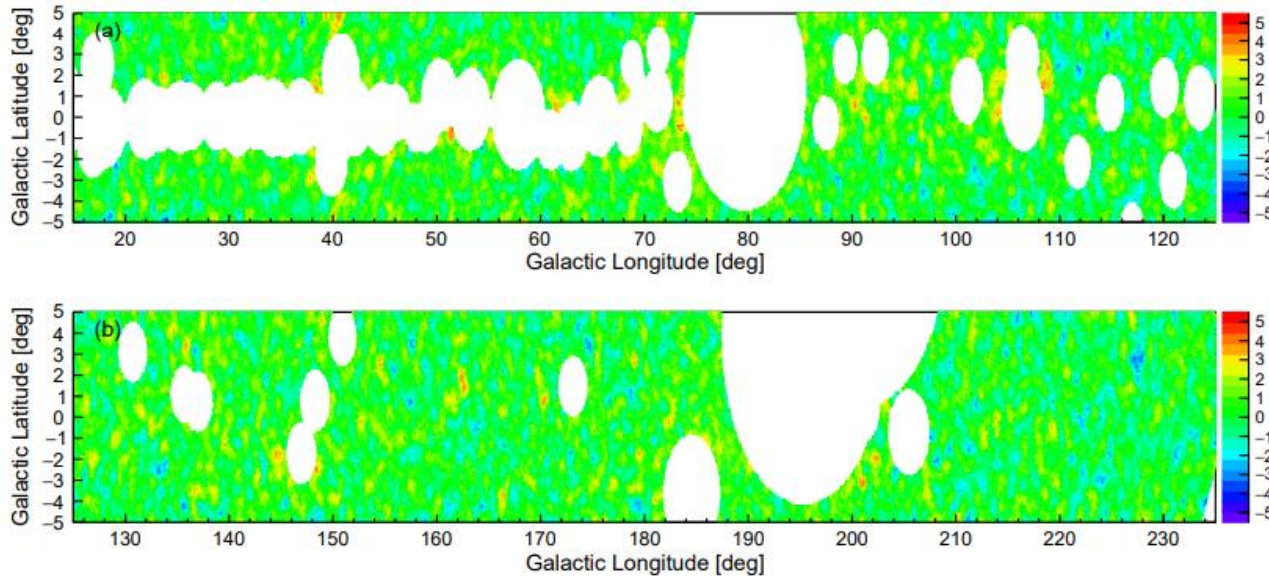
Vecchiotti et al (2024), in preparation

LHAASO diffuse emission measurements:

$$\varphi_{\gamma,\text{diff}}^{\text{LHAASO}}$$



The LHAASO collaboration provided a measurement of the Galactic diffuse γ -ray emission in the energy range 10 TeV to 1 PeV in two sky regions by masking the contribution of known sources. *Z. Cao et al. 2023, Phys. Rev. Lett, 131*



Do the LHAASO Galactic diffuse emission data are contaminated unresolved sources?

$$\varphi_{\gamma,\text{diff}}^{\text{LHAASO}} = \varphi_{\gamma,S}^{\text{UnRes}} + \varphi_{\gamma,\text{diff}}$$

Population study H.E.S.S.

Models:
Assumptions on the CR spatial and energy distributions, cross-section, and ISM.

Diffuse gamma-ray emission:

1. Differential inelastic cross section of pp interaction.

2. Interstellar gas distribution in the Galaxy

$$\varphi_{\gamma}(E_{\gamma}, \hat{n}_{\gamma}) = \int_{E_{\gamma}}^{\infty} dE \frac{d\sigma(E, E_{\gamma})}{dE_{\gamma}} \int_0^{\infty} dl \varphi_{CR}(E, \bar{r}_{\odot} + l\hat{n}_{\gamma}) n_H(\bar{r}_{\odot} + l\hat{n}_{\gamma})$$

3. Cosmic-ray energy and spatial distribution

1. Cross section:

$$f(E_\gamma) = \int_{E_\gamma}^{\infty} dE \frac{d\sigma(E, E_\gamma)}{dE_\gamma} \varphi_{CR}(E)$$

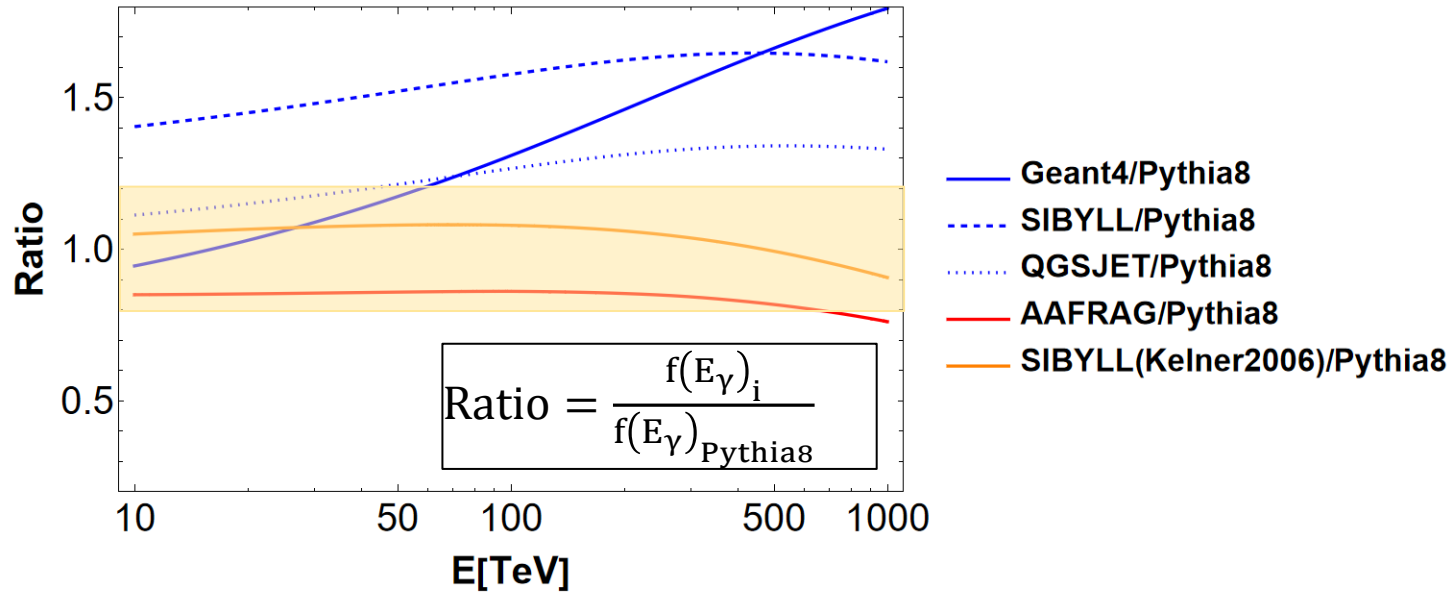
We compare:

- different parameterization based on different MC codes SIBYLL, QGSJET, Pythia8, and Geant4 (*Kafexhiu et al 2014*);
- SIBYLL (*Kelner et al 2006*);
- AAFRAG based on QGSJET-II-04m (*M. Kachelriess et al 2022*)

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$$f(E_\gamma) = \int_{E_\gamma}^{\infty} dE \frac{d\sigma(E, E_\gamma)}{dE_\gamma} \varphi_{CR}(E)$$

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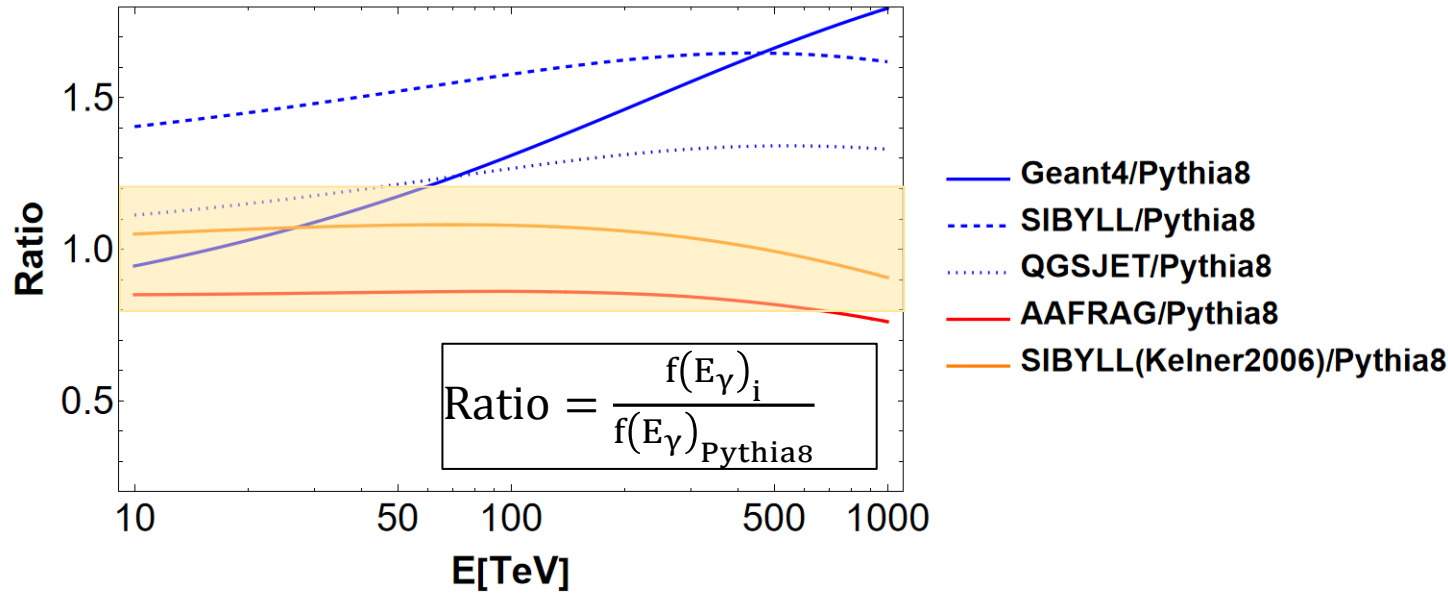
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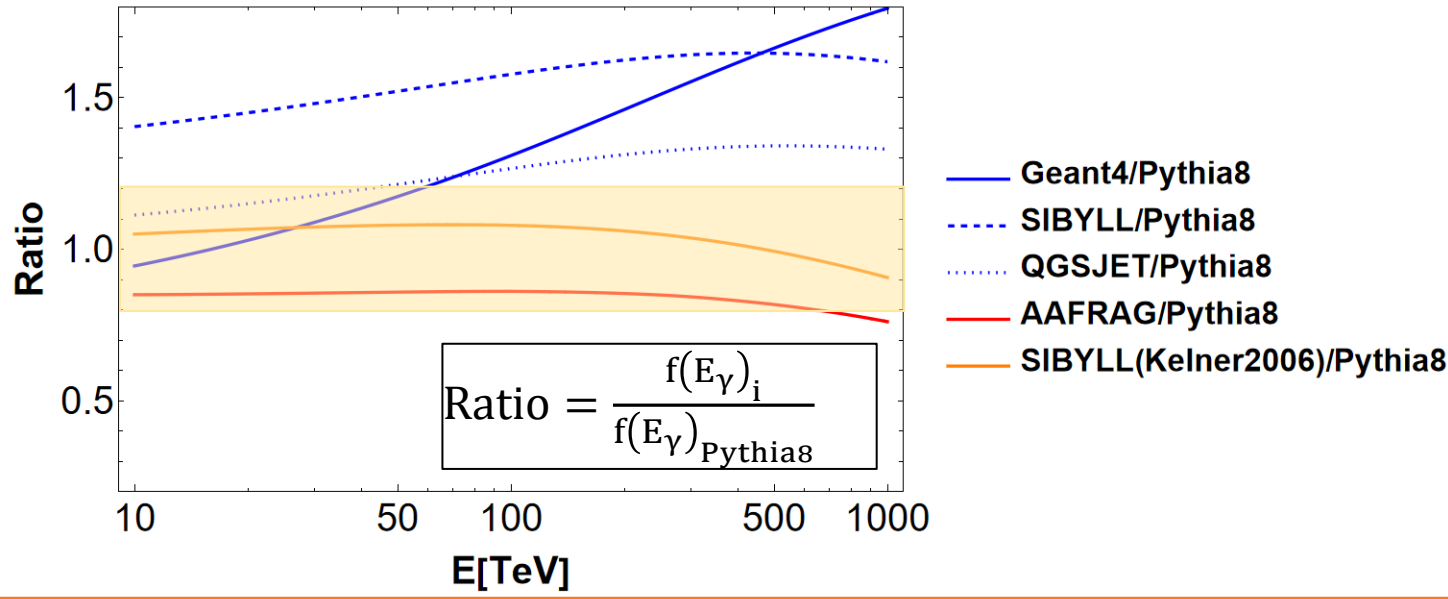
Assumptions cross section:

We take AAFRAG for the fiducial case and SIBYLL (*Kafexhiu et al 2014*) to include uncertainties.

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2. ISM:

$$n_c = \frac{\int_{\Delta\Omega} d\Omega \int_0^\infty dl n_H (\bar{r}_{Sun} + l \hat{n}_\gamma)}{\Delta\Omega}$$

- Dust (masked): $n_{c,inner} = 7.85068 \times 10^{21} \text{ cm}^{-2}$
PLANCK
(Aghanim et al. 2016)
- Galprop (masked): $n_{c,inner} = 9.47769 \times 10^{21} \text{ cm}^{-2}$
 $n_{c,outer} = 8.72069 \times 10^{21} \text{ cm}^{-2}$

Galprop provides ~ 20% more target than the dust in the inner region and ~ 50 % in the outer region

Assumptions ISM:

We take Galprop for the fiducial case and the Dust to include uncertainties.

3. Cosmic ray distribution:

$$\varphi_{CR}(E, \vec{r}) = \varphi_{CR, Sun}(E) g(\vec{r}, R) h(E, \vec{r})$$

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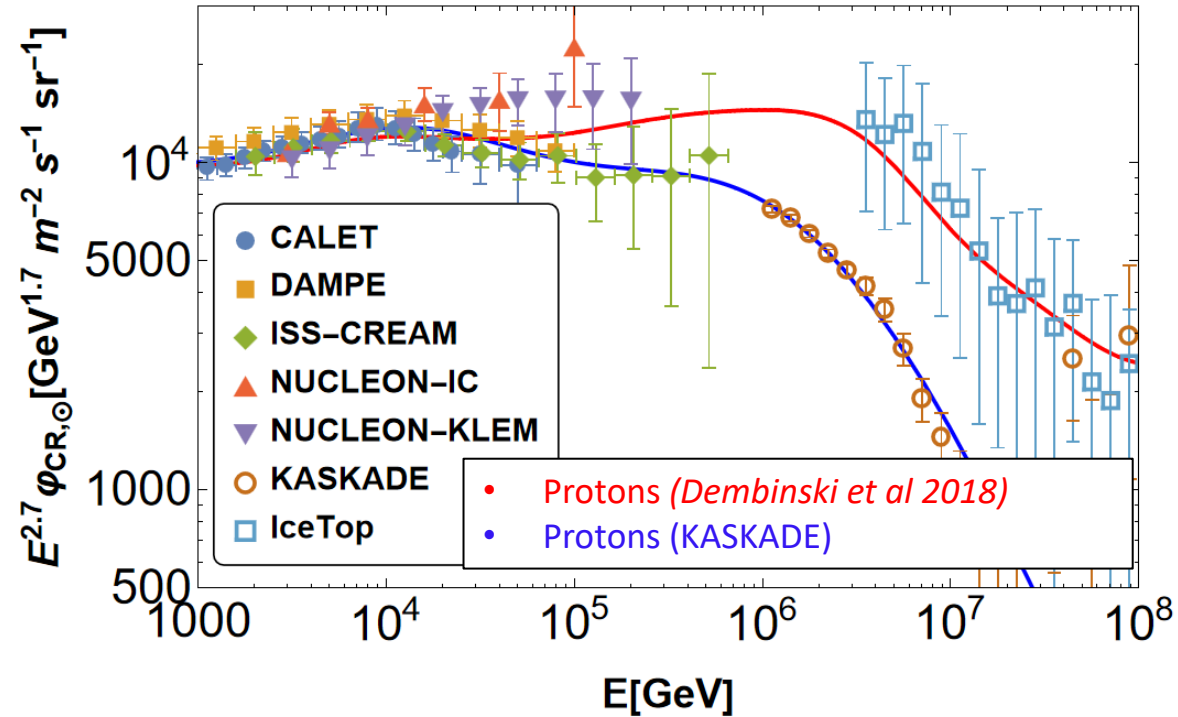
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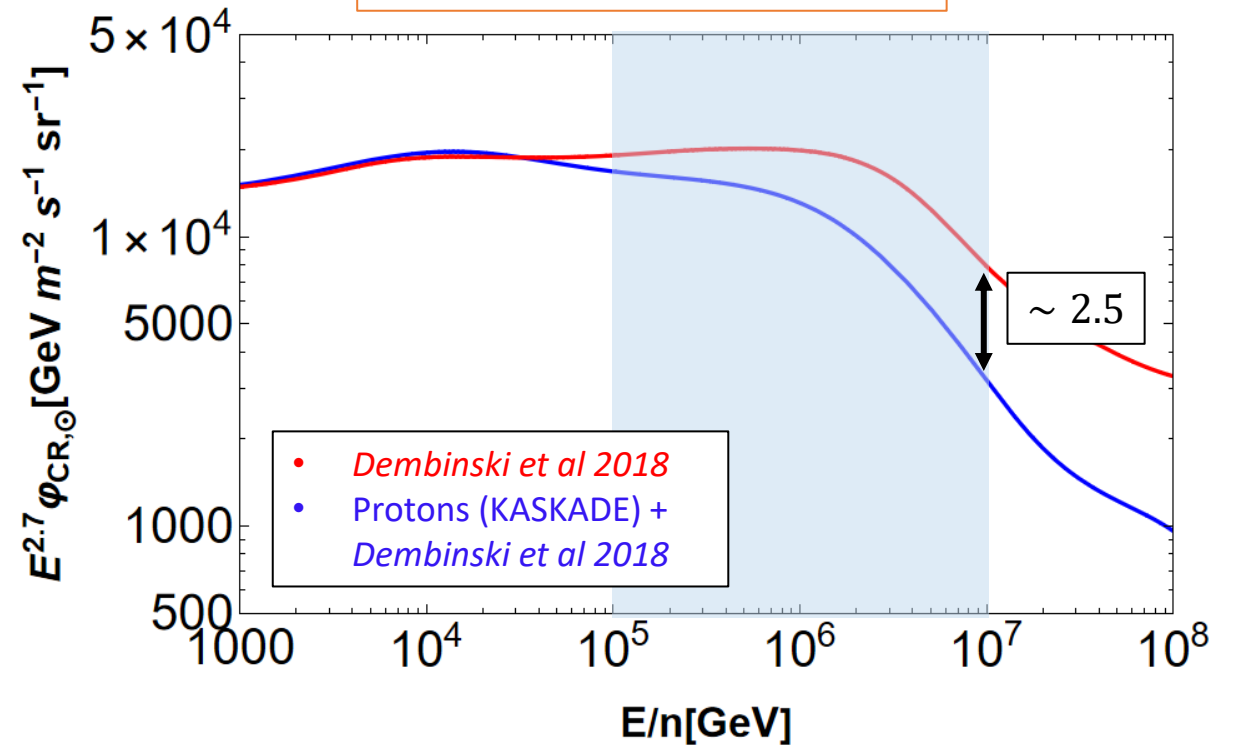
Assumption CR spectrum:

We take the data-driven CR spectrum from **Dembinski et al 2018** for the fiducial case and **Protons (KASKADE) + Dembinski et al 2018** to include uncertainties.

Protons



All Particle Spectrum



3. Cosmic ray distribution: $\varphi_{CR}(E, \vec{r}) = \varphi_{CR, Sun}(E) g(\vec{r}, R) h(E, \vec{r})$

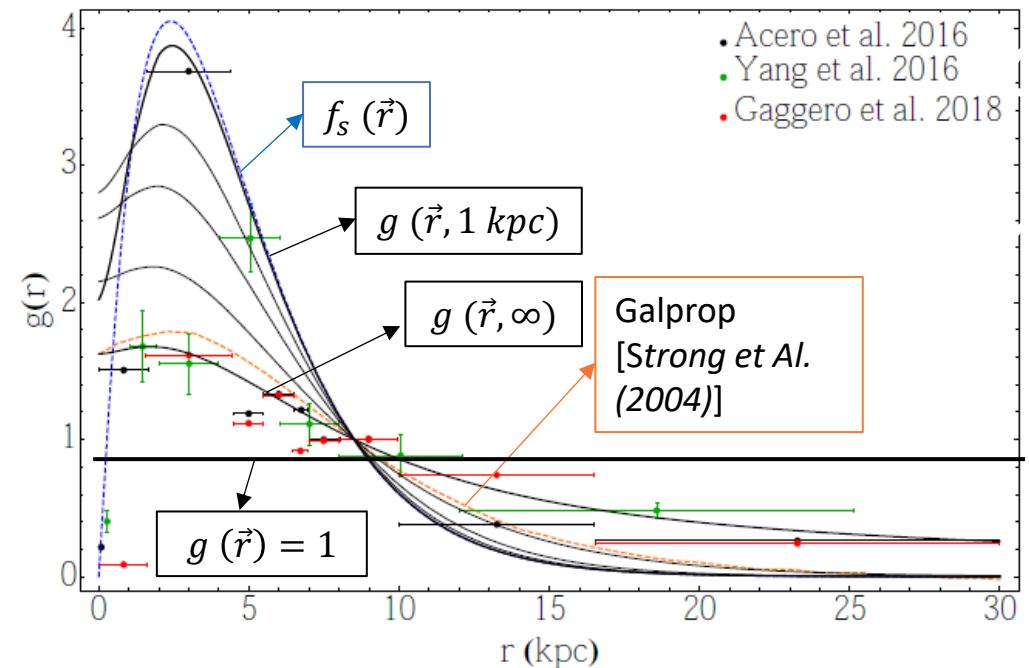


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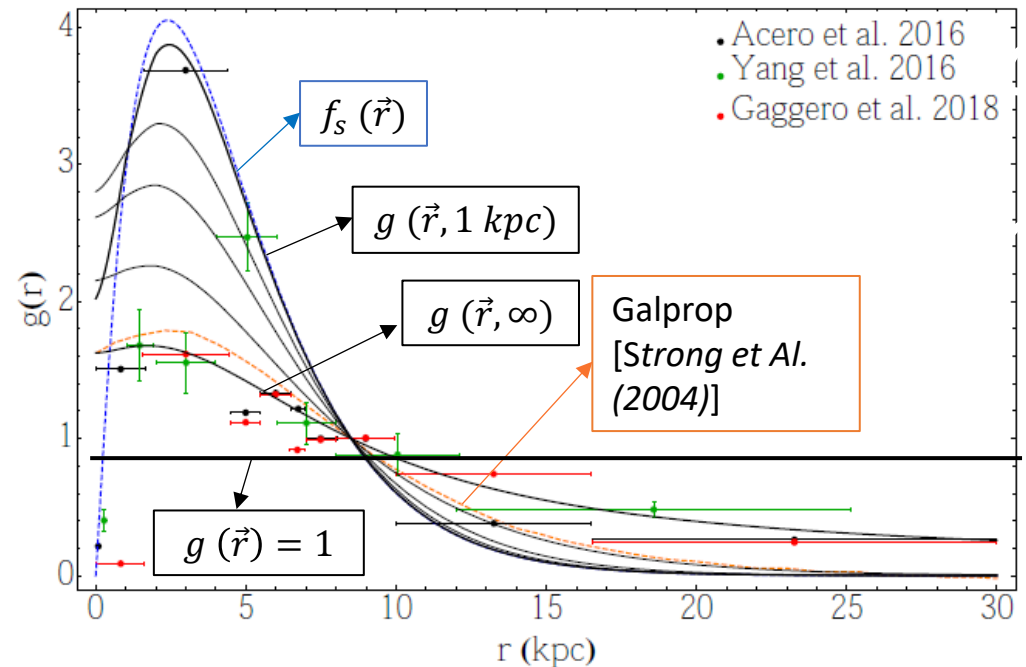
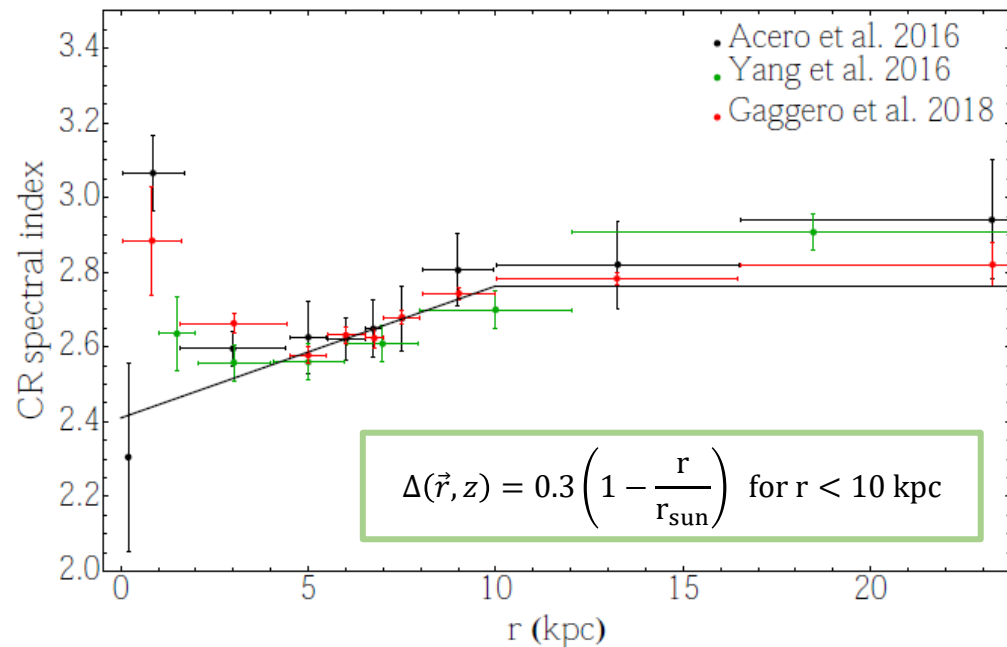
$g(r)$ is determined by the distribution of the CR sources $f_s(\vec{r})$ (proportional to the SNR number density by Green et al. (2015)), and by the propagation of CR in the Galactic magnetic field.



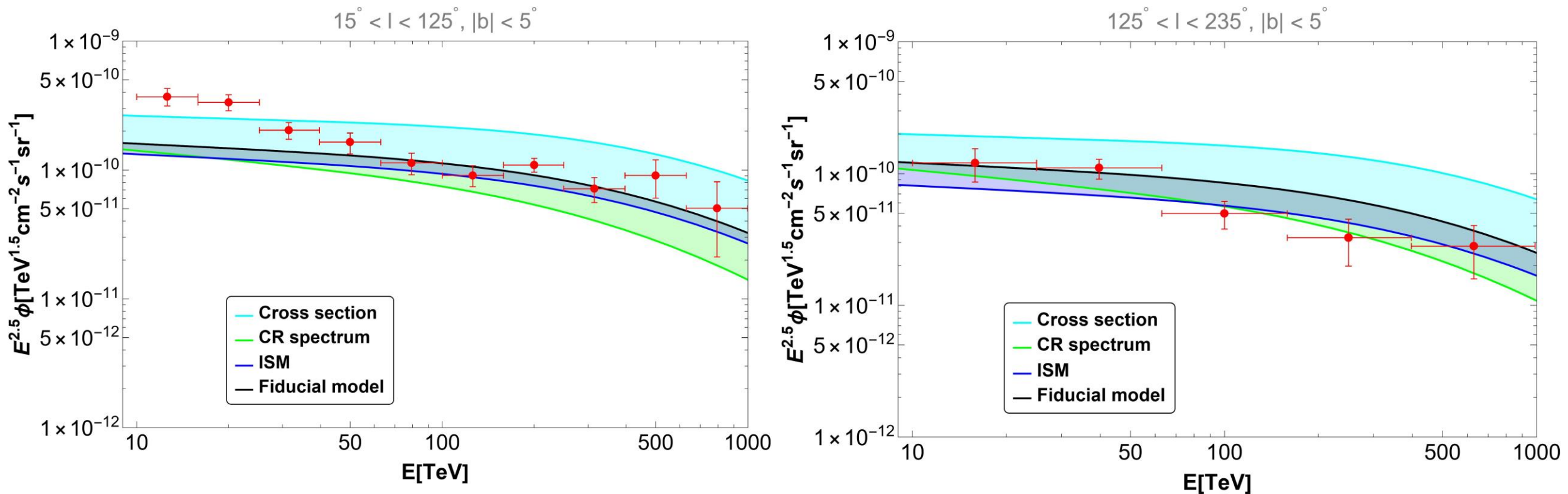
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- ★ Assumption CR spectrum: We take the data-driven CR spectrum from **Dembinski et al 2018** for the fiducial case and **Protons (KASCADE) + Dembinski et al 2018** to include uncertainties.
- ★ $g(r)$ is determined by the distribution of the CR sources $f_s(\vec{r})$ (proportional to the SNR number density by Green et al. (2015), and by the propagation of CR in the Galactic magnetic field.
- ★ **2 cases: with and without spatially dependent CR spectral index** (from the analysis of the FermiLAT data at ~ 20 GeV
Acero et al. (2016), Yang et al. (2016), Gaggero et al. (2018)

$$h(E, \vec{r}) = \left(\frac{E}{20 \text{ GeV}} \right)^{\Delta(\vec{r})}$$



Comparison with LHAASO (standard diffusion):



- LHAASO data can be explained by the “truly” diffuse emission in the outer region and in the inner region above ~ 30 TeV;
The contribution from unresolved sources must be negligible.

Pulsar Wind Nebulae population:

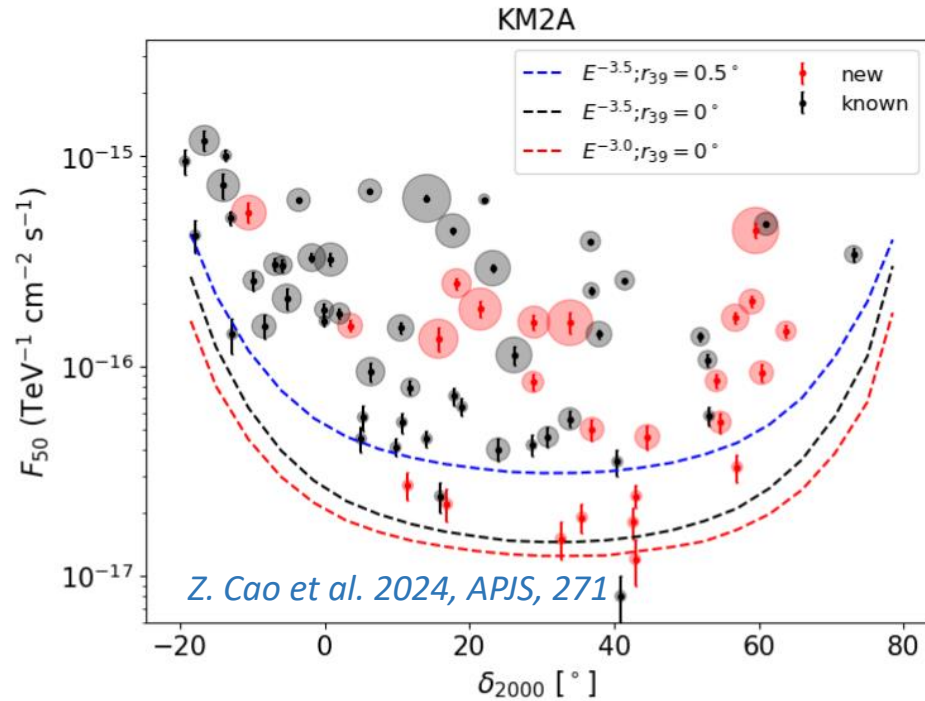
Cataldo et al. *Astrophys.J.* 904 (2020); Pagliaroli et al, *Universe*, 9,881 (2023)

We built a synthetic population of PWNe using the best fit of the maximum luminosity in the energy range 1 – 100 TeV: $L_{max} = 2.2 \times 10^{35}$ erg/s and the spin-down timescale: $\tau_{sd} = 2.9$ kyrs derived from fitting the brightest sources of the HGPS.

Assumptions:

- Latitude, longitude, and radius are extracted from the Lorimer distribution that scales as $\exp(-|z|/H)$ with $H = 0.05$ kpc (it is the value that provides the best chi-square in the fit of HGPS data);
- The age of sources t_{age} is extracted uniformly in the interval $[1, 10^6]$ yr;
- The luminosity is calculated from: $L = L_{max} \left(1 + \frac{t_{age}}{\tau_{sd}}\right)^{-2}$
- Spectrum: power-law with exponential cut-off ($E_{cut} = 100$ TeV), spectral index fixed to: 2.4.

Source contributions:



“The flux sensitivity is defined as the flux normalization required to have 50% probability of detecting a source at 5σ level”

At 50 TeV the differential threshold of point-like sources depends only mildly on the spectral assumption.

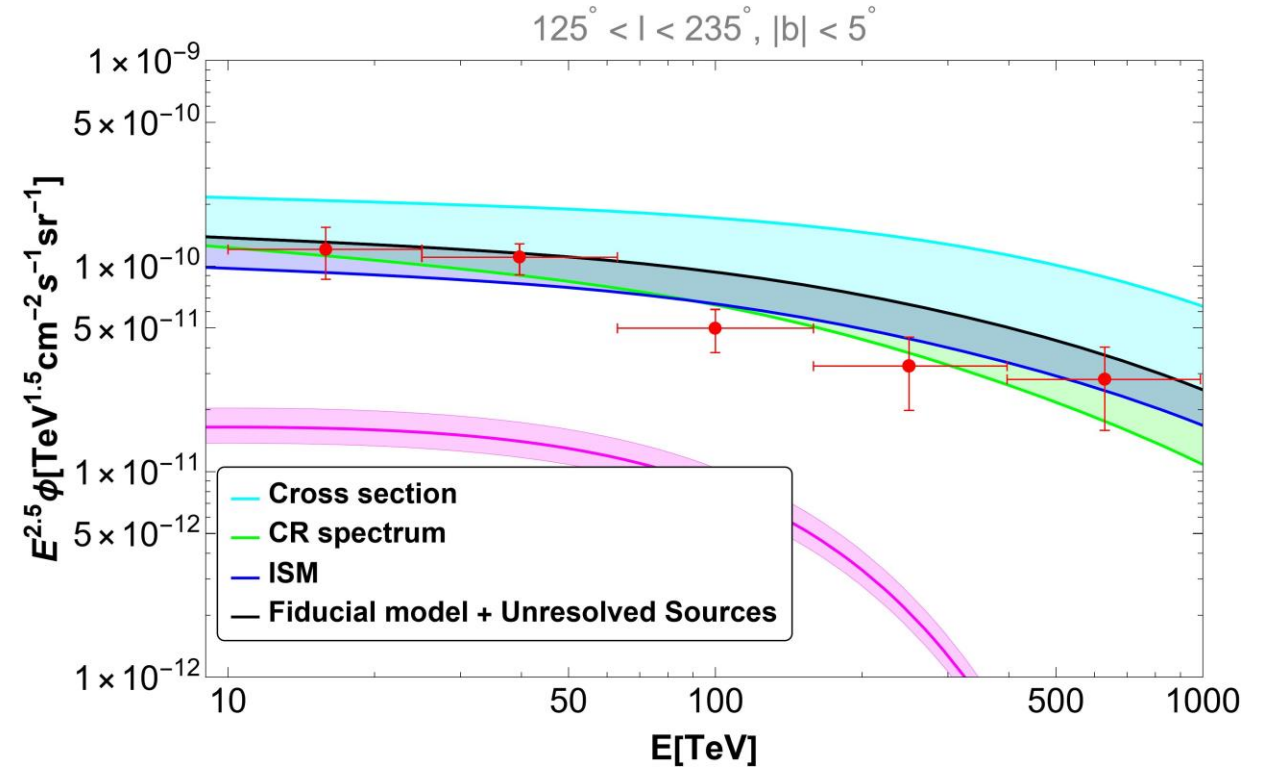
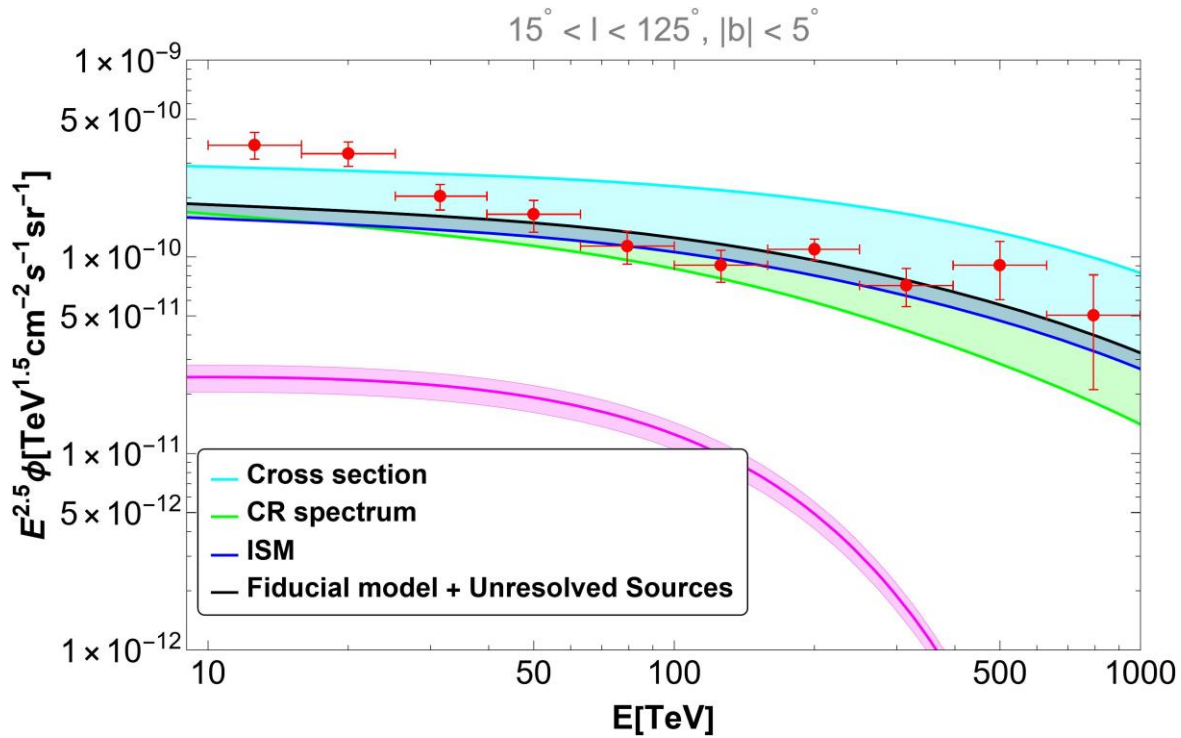
Def: Resolved sources: $\phi_{50} > \phi_{th,50}$, $\phi [TeV^{-1} cm^{-2} s^{-1}]$

		N_R	φ_R	φ_{UNR}	$\varphi_{UNR,H}$
$15^\circ < l < 235^\circ$	MC	84_{-5}^{+7}	$1.69_{-0.43}^{+0.62} \times 10^{-14}$	$2.82_{-0.14}^{+0.15} \times 10^{-15}$	—
	KM2A	76	1.72×10^{-14}	—	—
$15^\circ < l < 125^\circ$	MC	73_{-8}^{+5}	$1.32_{-0.33}^{+0.37} \times 10^{-14}$	$2.56_{-0.16}^{+0.14} \times 10^{-15}$	$2.23_{-0.36}^{+0.34} \times 10^{-16}$
	KM2A	59	1.44×10^{-14}	—	—
$125^\circ < l < 235^\circ$	MC	12_{-2}^{+3}	$2.82_{-1.1}^{+1.8} \times 10^{-15}$	$2.53_{-0.35}^{+0.46} \times 10^{-16}$	$2.08_{-0.34}^{+0.49} \times 10^{-16}$
	KM2A	16	2.74×10^{-15}	—	—

Results:

- The predicted number and flux of resolved sources are compatible with the KM2A quantities within 2σ ;
- The unresolved source flux is suppressed by 91 % and 18 % in the inner and outer regions, respectively.

Comparison with LHAASO (standard diffusion):

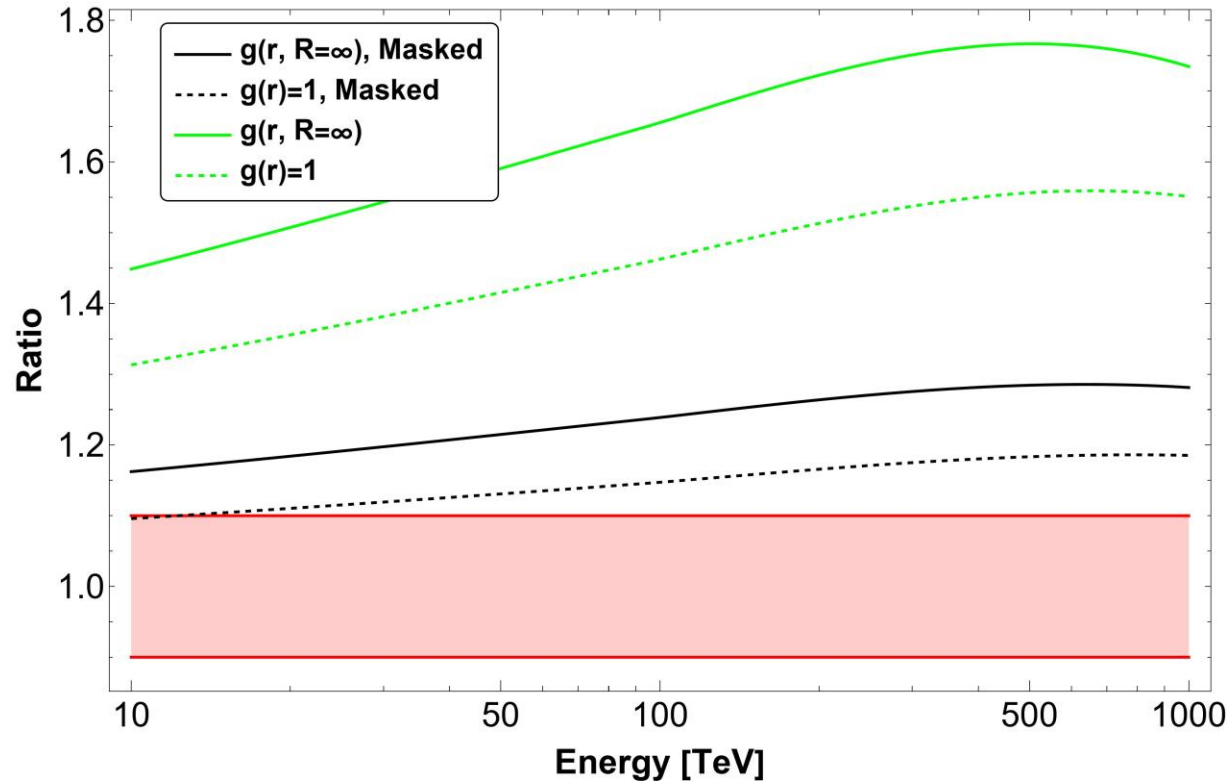


- The masks cancel out most of the unresolved source contributions. Unresolved sources contribute $\sim 15\%$ of the fiducial model at 50 TeV in both regions;

Effect of LHAASO masks on the hardening:

Def: hardening= spatially dependent CR spectral index

Inner Region: $15^\circ < l < 125^\circ$



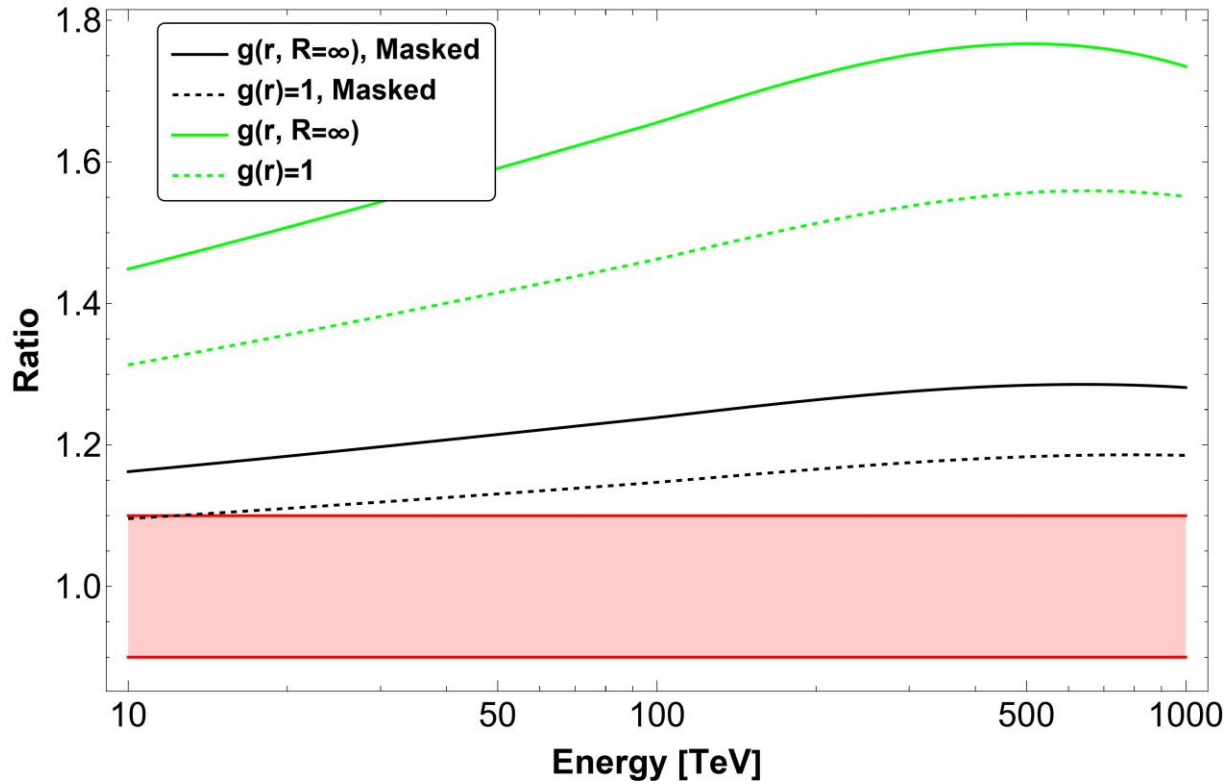
$$\text{Ratio} = \frac{\Phi_{\text{hardening}}}{\Phi_{\text{standard}}}$$

The ratio is *independent* of the cross-section, ISM and CR spectrum but it *depends* on the CR spatial distribution

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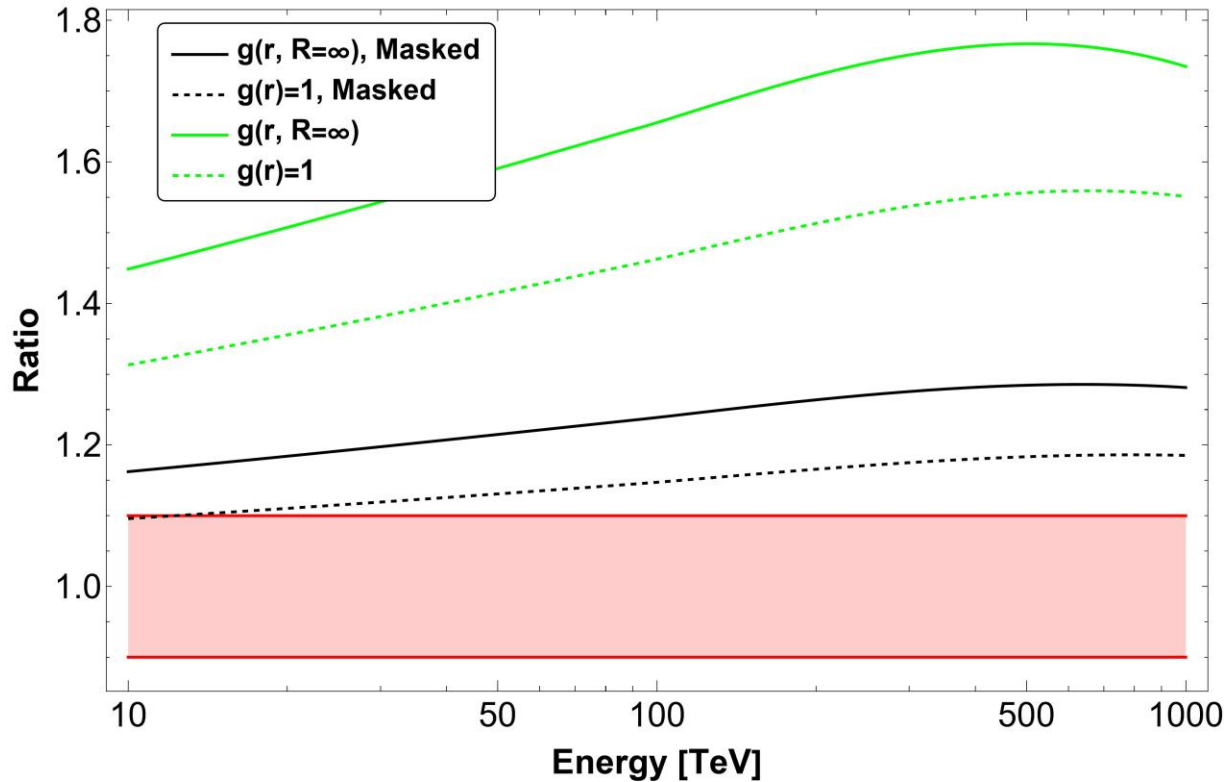
The ratio is *independent* of the cross-section, ISM and CR spectrum but it *depends* on the CR spatial distribution

- $g(r, R = \infty)$: The $\Phi_{\text{hardening}}$ produces 76 % more signal than Φ_{standard} at 500 TeV ;
- $g(r) = 1$: The $\Phi_{\text{hardening}}$ produces 55% more signal than Φ_{standard} at 500 TeV

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After masking:

- $g(r, R = \infty)$: $\Phi_{\text{hardening}}$ produces 28 % more signal than Φ_{standard} at 500 TeV
- $g(r) = 1$: $\Phi_{\text{hardening}}$ produces 18% more signal Φ_{standard} at 500 TeV

Conclusions:

1. The total flux and the number of sources derived in *Cataldo et al 2019* based on the HGPS are compatible with the observation of KM2A within 2σ ;
2. The LHAASO masks cancel most of the effect due to unresolved sources in the inner region (suppressed by 91 %). In the outer region, unresolved sources already produce a negligible contribution to the diffuse emission that is further suppressed by the LHAASO mask of about 18 %;
3. The cross-section, CR spectrum, and ISM uncertainties are non-negligible. However, the LHAASO data seems compatible with the “truly” diffuse emission within uncertainties except for the 2 low energy points in the inner region which could be explained by introducing other classes of unresolved sources.
4. The LHAASO masks significantly reduce the effect of a spatial-dependent CR spectral index. As a consequence, it is challenging to test this hypothesis using LHAASO data.

Backup slides

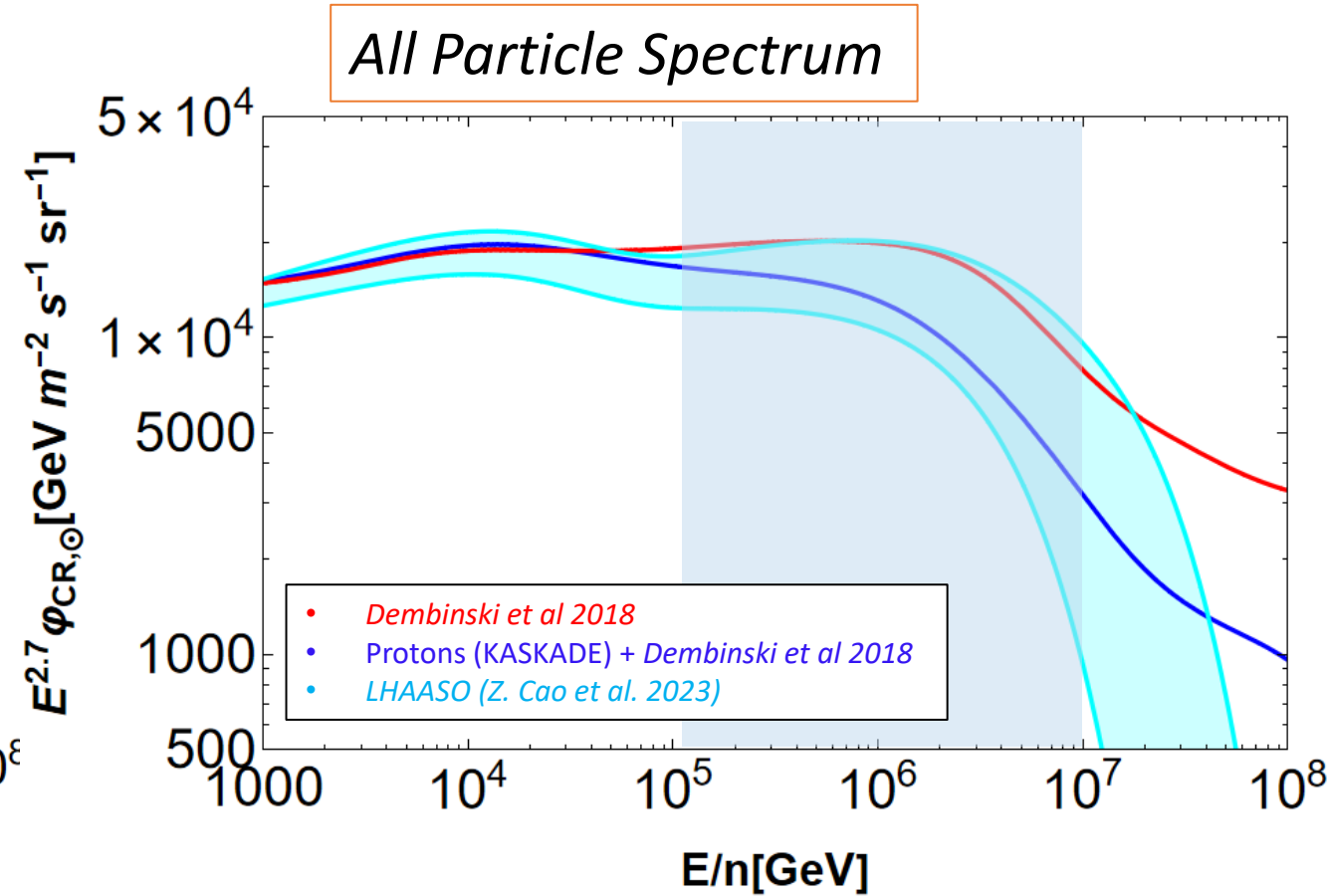
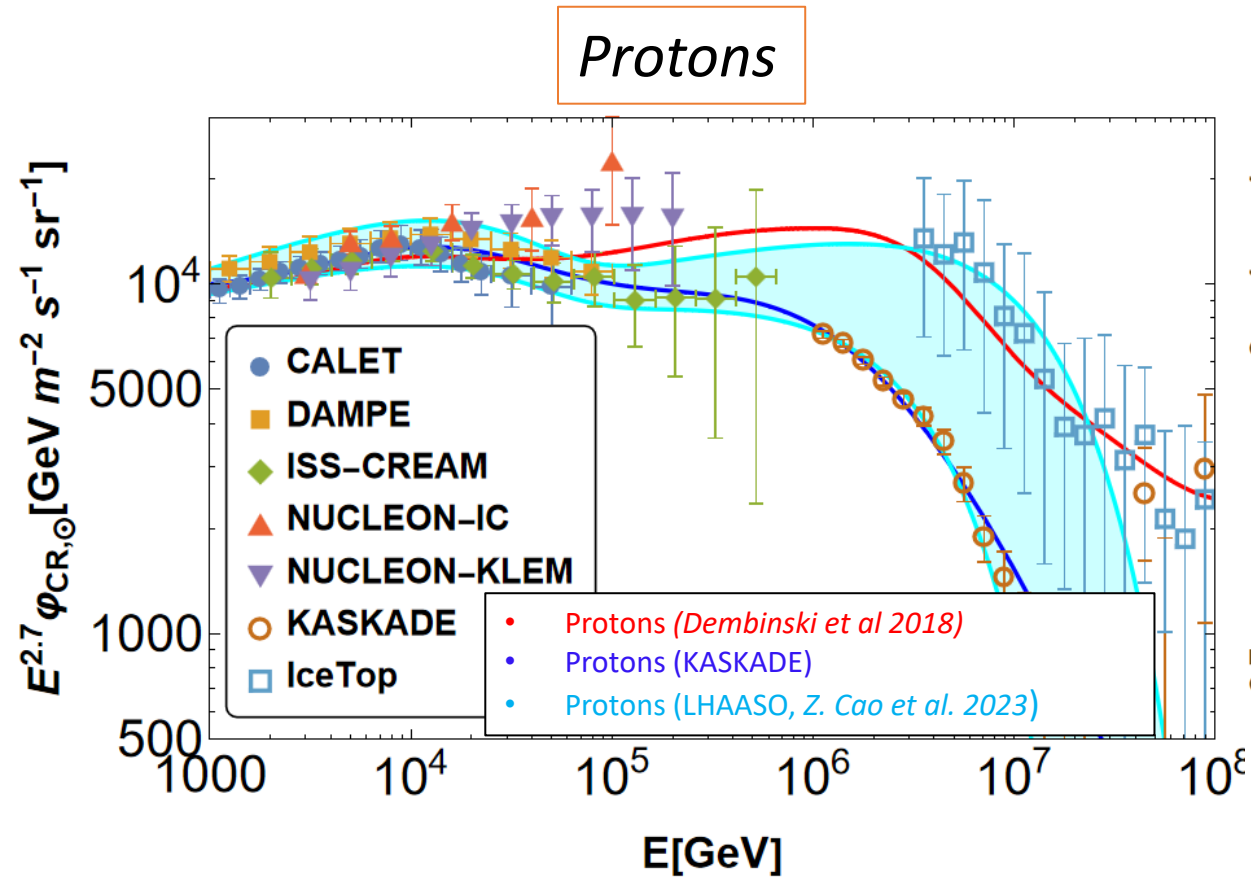
3. Cosmic ray distribution:

$$\varphi_{CR}(E, \vec{r}) = \varphi_{CR,Sun}(E) g(\vec{r}, R) h(E, \vec{r})$$

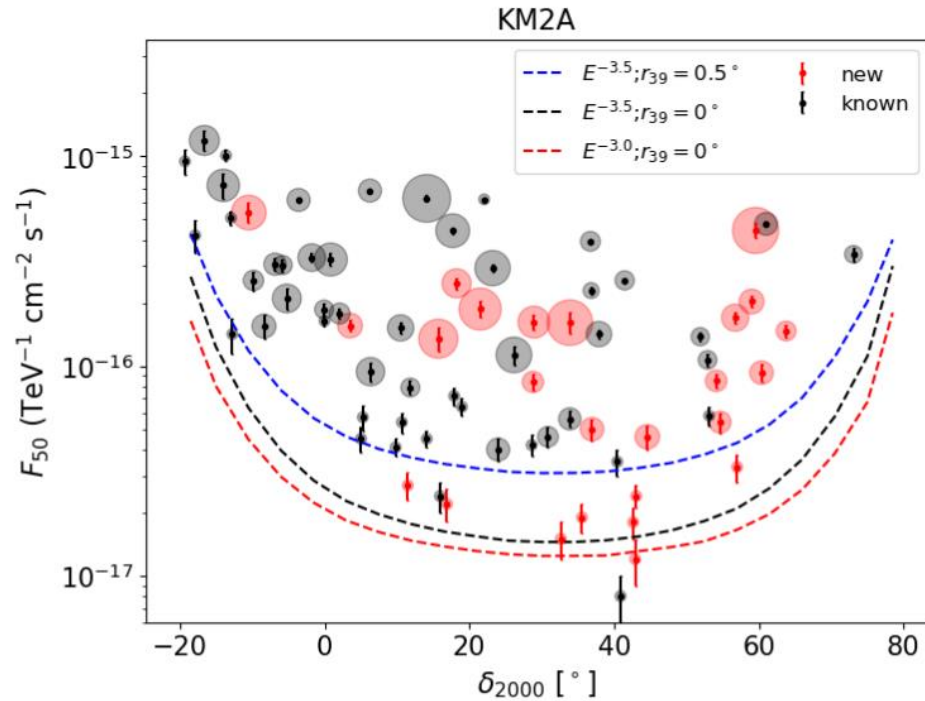


Assumption CR spectrum:

We take the data-driven CR spectrum from **Dembinski et al 2018** for the fiducial case and **Protons (KASKADE) + Dembinski et al 2018** to include uncertainties.



Source contributions (Size 40 pc):



Best fit for size 40 pc:
 $L_{max} = 2 \times 10^{35}$ erg/s
 $\tau_{sd} = 4.6$ kyrs

Def: Resolved sources: $\phi_{50} > \phi_{th,50} \sqrt{(\sigma_{psf}^2 + \sigma_s^2) / \sigma_{psf}^2}$ where $\sigma_{psf} = 0.2^\circ$ and σ_s is the angular size of the source.

		N_R	φ_R	φ_{UNR}	$\varphi_{UNR,H}$
$15^\circ < l < 235^\circ$	MC	69_{-5}^{+5}	$2.28_{-0.56}^{+0.74} \times 10^{-14}$	$7.17_{-0.45}^{+0.37} \times 10^{-15}$	—
	KM2A	76	1.72×10^{-14}	—	—
$15^\circ < l < 125^\circ$	MC	62_{-6}^{+4}	$1.76_{-0.45}^{+0.47} \times 10^{-14}$	$6.10_{-0.41}^{+0.37} \times 10^{-15}$	$8.19_{-1.64}^{+1.74} \times 10^{-16}$
	KM2A	59	1.44×10^{-14}	—	—
$125^\circ < l < 235^\circ$	MC	7_{-1}^{+2}	$4.21_{-1.8}^{+3.4} \times 10^{-15}$	$1.05_{-0.27}^{+0.25} \times 10^{-15}$	$8.32_{-1.99}^{+2.49} \times 10^{-16}$
	KM2A	16	2.75×10^{-15}	—	—

Results:

- The predicted number and flux of resolved sources are compatible with the KM2A quantities within 2σ (except in the outer region);
- The unresolved source flux is suppressed by 86 % and 20 % in the inner and outer regions, respectively.

Summary assumptions:

Assumptions for the diffuse emission fiducial model:

- CRs: *Dembinki et al 2018*;
- Gas: Galprop;
- Cross section: AAFRAG;
- CR spatial distribution of CR: $g(\vec{r}, \infty)$

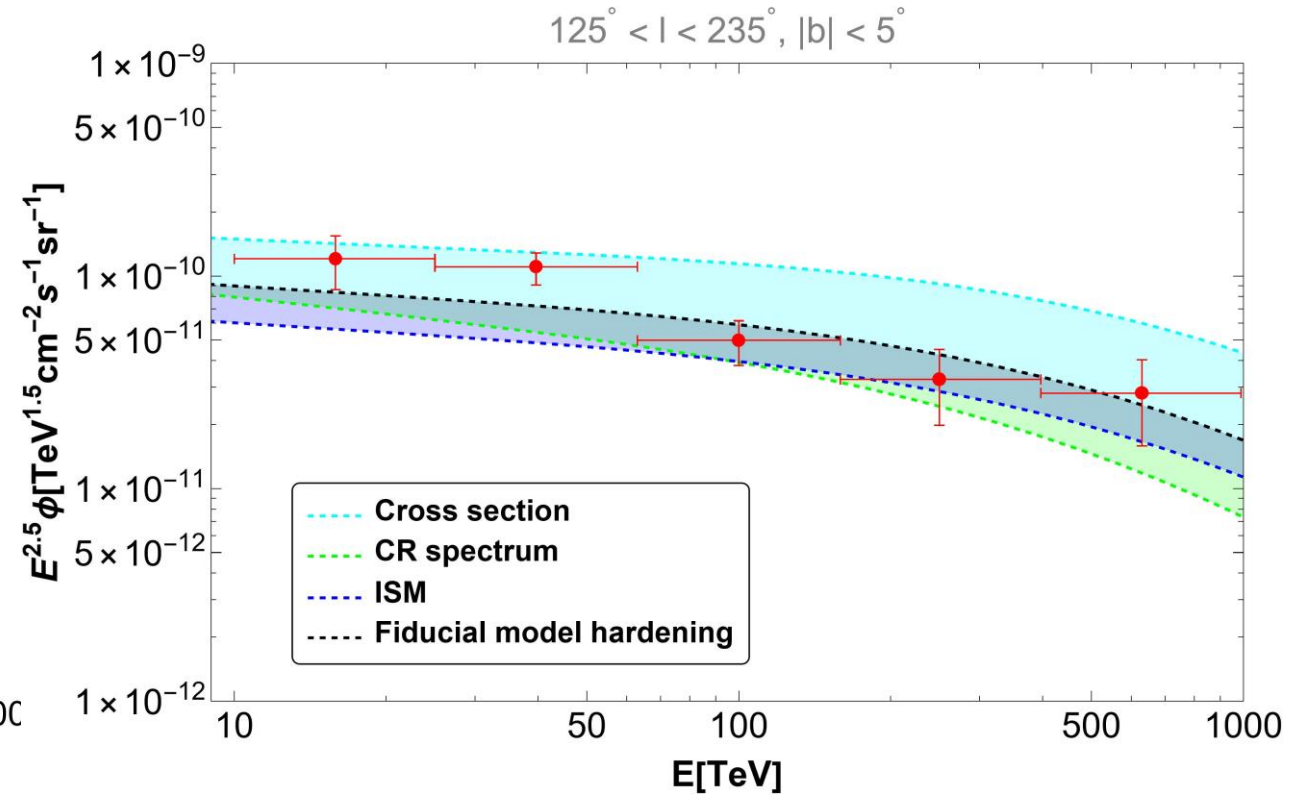
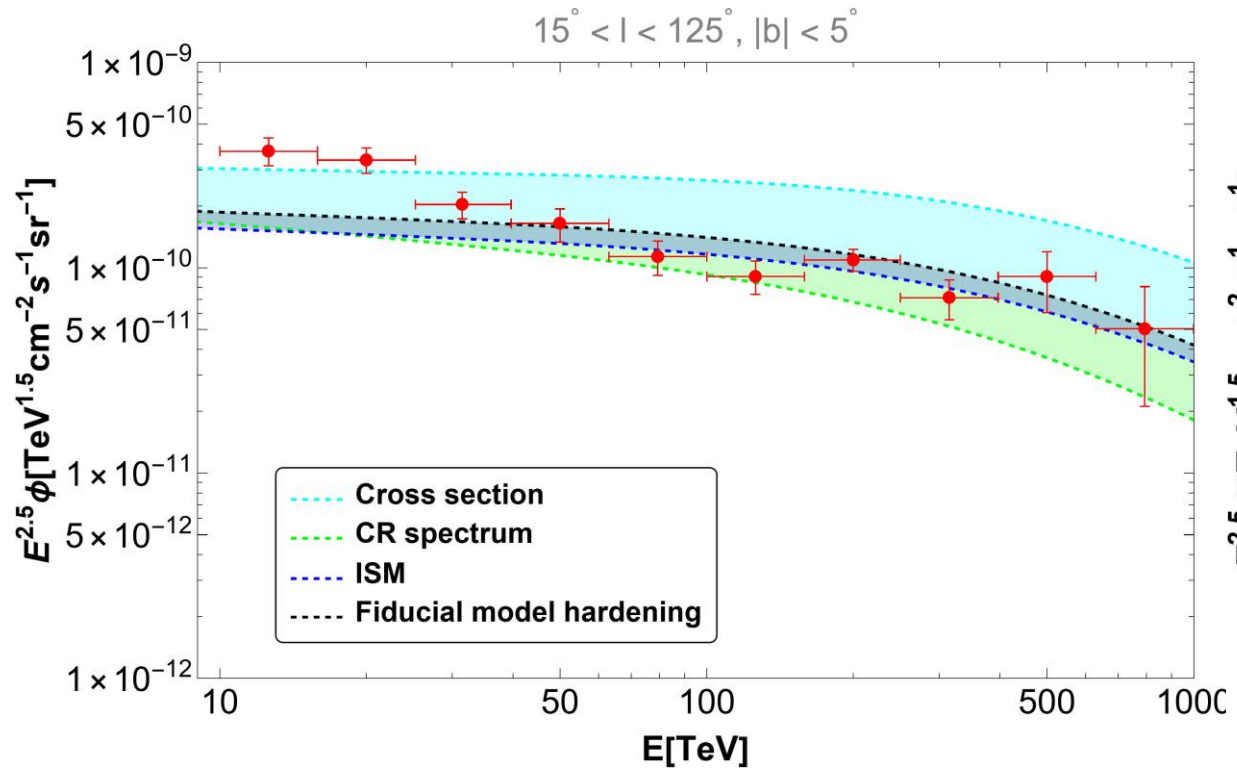
Variation with respect to the fiducial model:

- CRs: fit protons (KASKADE) + All elements (*Dembinki et al 2018*);
- Gas: Dust;
- Cross section: Sybill;
- CR spatial distribution of CR: $g(\vec{r}, \infty)$

Assumptions unresolved sources:

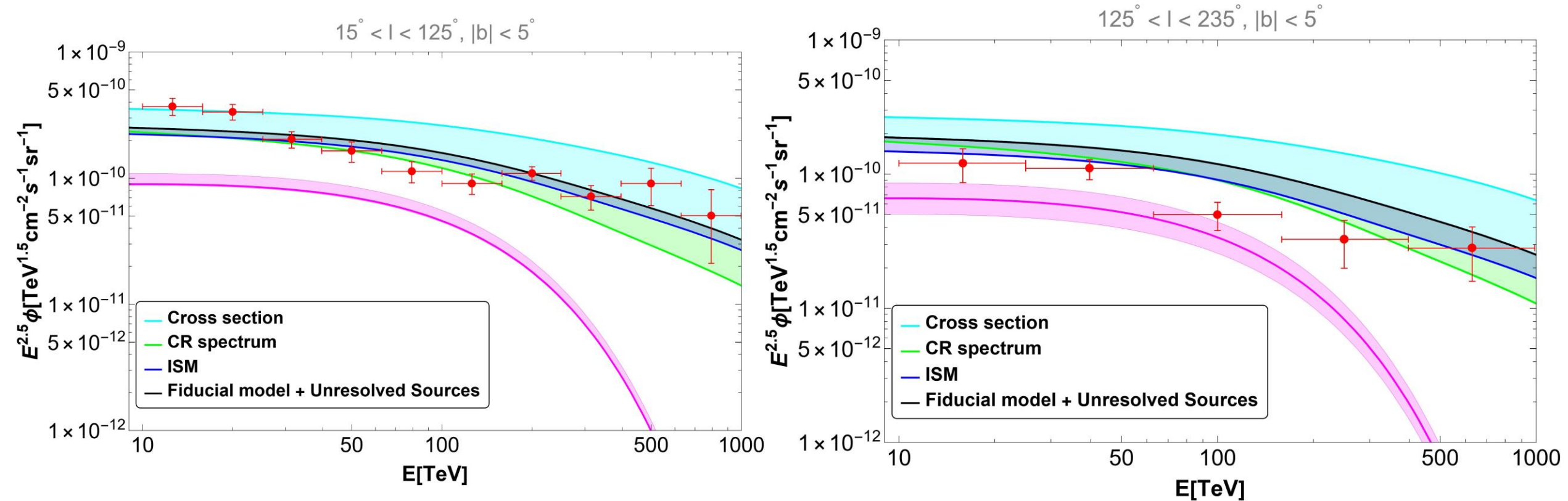
- Spectrum: power-law with exponential cut-off ($E_{\text{cut}} = 100$ TeV), spectral index fixed to: 2.4.
- Thickness of the disk $H=0.05$ kpc.

Comparison with LHAASO (hardening effect):



- LHAASO cannot be used to distinguish the hardening hypothesis from the standard one;

Comparison with LHAASO (standard diffusion and size 40 pc):



- Unresolved sources contribute $\sim 54\%$ of the fiducial model at 50 TeV in both regions;
- The 40 pc size case corresponds to an upper limit for the unresolved source contribution.