

PROBES OF PRIMORDIAL BLACK HOLES AS DARK MATTER

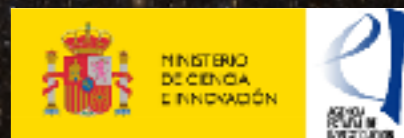
Sergio Palomares-Ruiz

IFIC, CSIC - U. Valencia

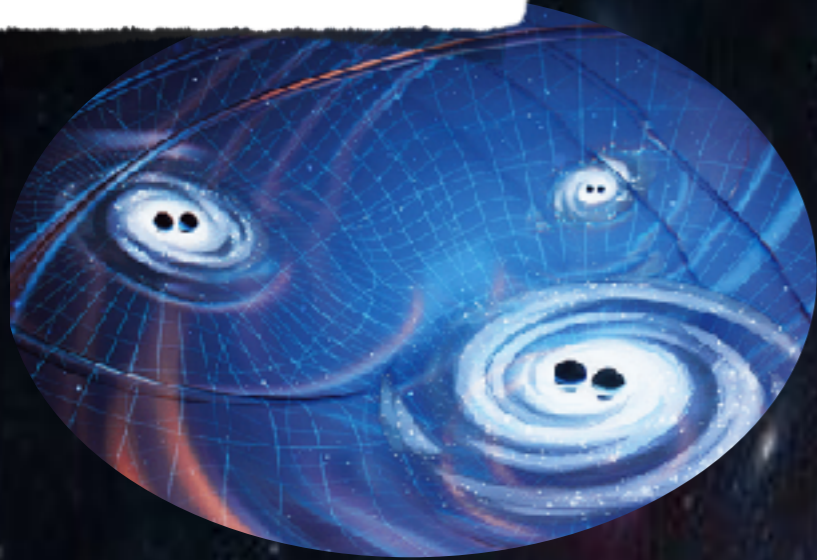


**RICAP-24 Roma International
Conference on Astroparticle Physics**

Frascati,
September 26, 2024



Gravitational waves

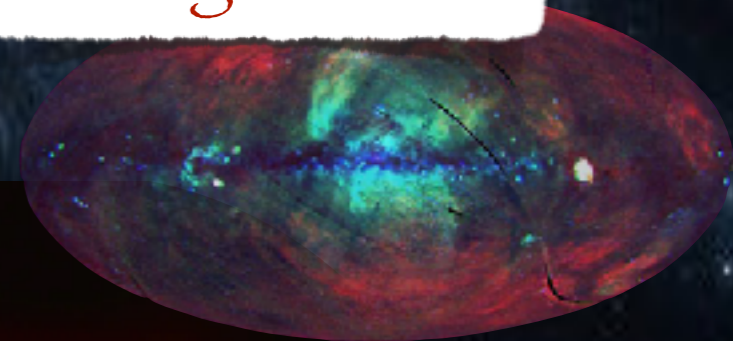


Dark Matter

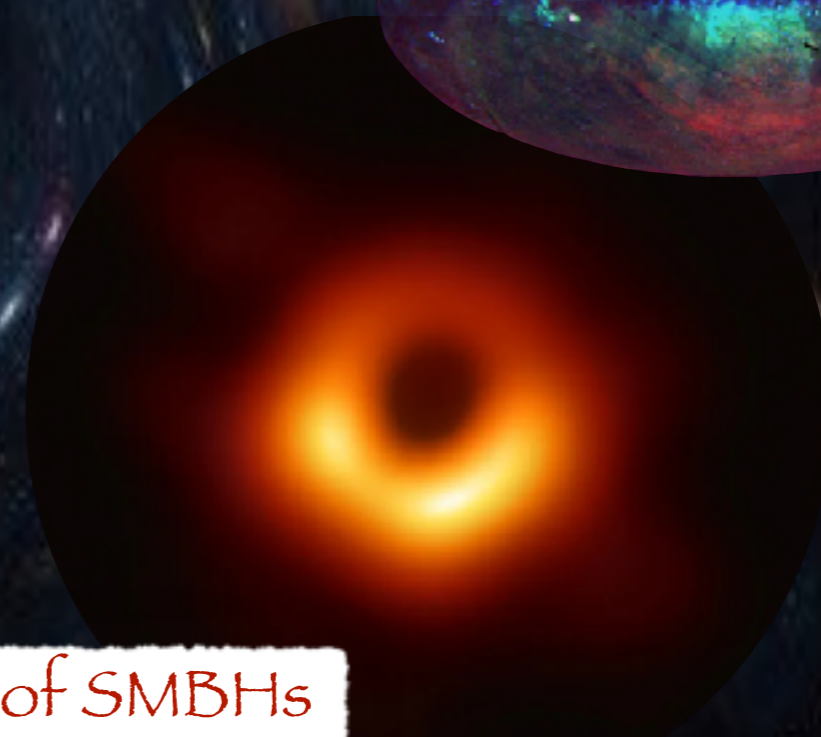
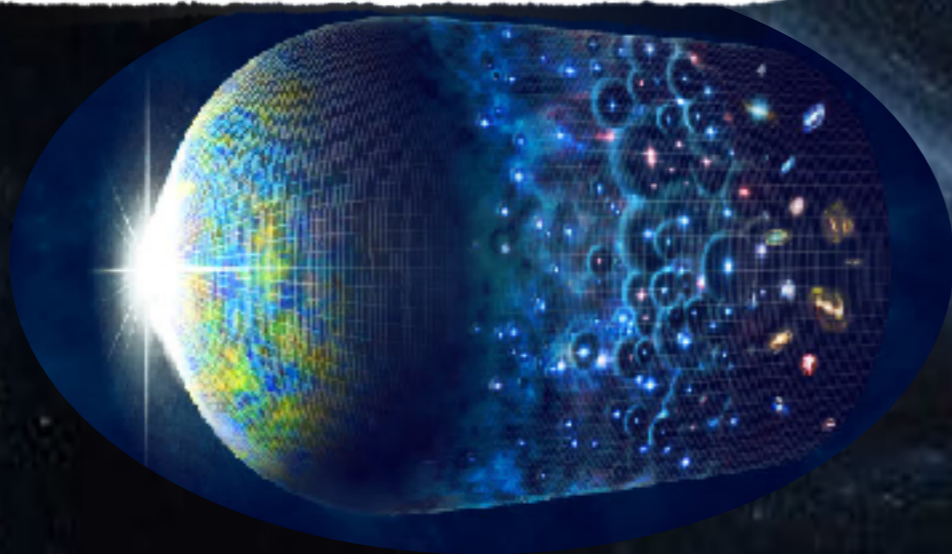


Primordial Black Holes

Cosmic radiation backgrounds



Formation:
Physics of the Early Universe



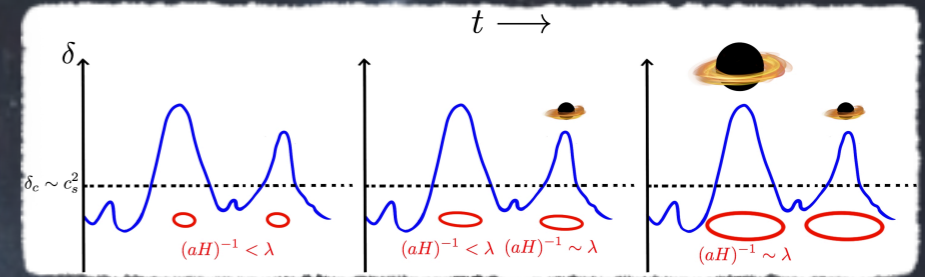
Origin of SMBHs

PRIMORDIAL BLACK HOLES

The early universe is very hot and dense:
ideal environment for black hole formation

Y. B. Zel'dovich and I. D. Novikov, *Sov. Astron.* 10:602, 1967

S. Hawking, *Mon. Not. R. Astron. Soc.* 152:75, 1971



PBHs could form during radiation era from the gravitational collapse of large fluctuations (at horizon entry) with masses of the order of the horizon mass... or via collapse of cosmic string loops or a scalar field, bubble collisions...

$$M_{\text{PBH}} \sim \frac{t}{G} \sim 10^{15} \left(\frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

$$t = 10^{-43} \text{ s} \rightarrow M_{\text{PBH}} \sim 10^{-5} \text{ g}$$

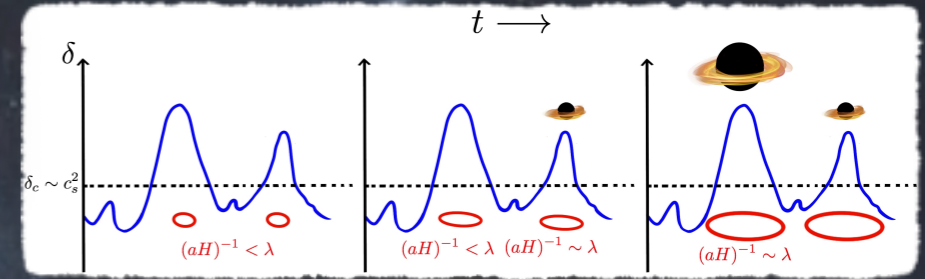
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Black holes radiate thermally, so they eventually evaporate

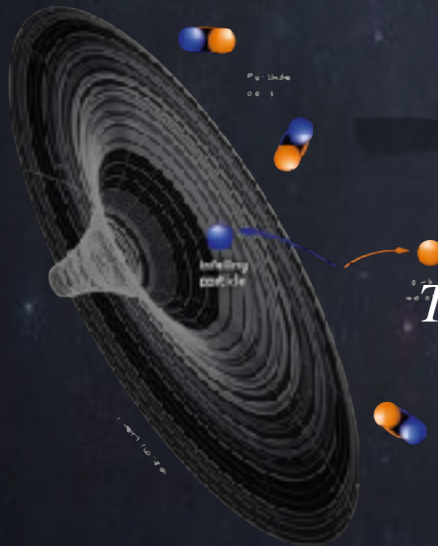
S. W. Hawking, *Commun. Math. Phys.* 43:199, 1975

D. N. Page, *Phys. Rev. D* 13:198, 1976

$$T_{\text{BH}} \sim \frac{1}{8 \pi G M_{\text{BH}}} \sim 10 \left(\frac{10^{15} \text{ g}}{M_{\text{BH}}} \right) \text{ MeV}$$

$$\tau(M_{\text{BH}}) \sim G^2 M_{\text{BH}}^3 \sim 100 \left(\frac{M_{\text{BH}}}{10^{15} \text{ g}} \right)^3 \text{ Gyr}$$

For masses between $10^{-17} M_{\odot}$ and $10^5 M_{\odot}$, they would be present today

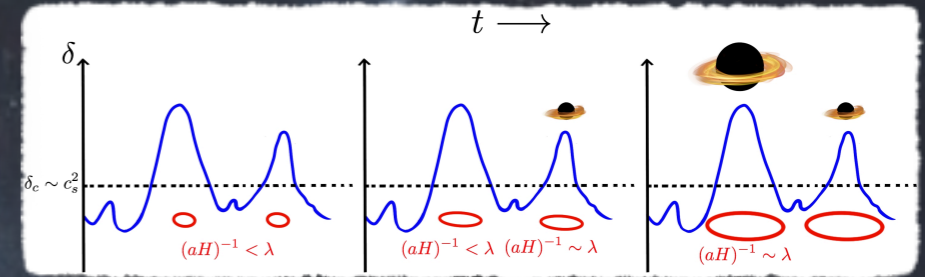


PRIMORDIAL BLACK HOLES

The early universe is very hot and dense:
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Y. B. Zel'dovich and I. D. Novikov, *Sov. Astron.* 10:602, 1967

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For masses between $10^{-17} M_{\odot}$ and $10^5 M_{\odot}$, they would be present today

PBHs would form before BBN, so they would not count as baryonic matter



A DM candidate which is not a new particle (although its formation involves BSM physics)

G. F. Chapline, *Nature* 253:251, 1975

PRIMORDIAL BLACK HOLES

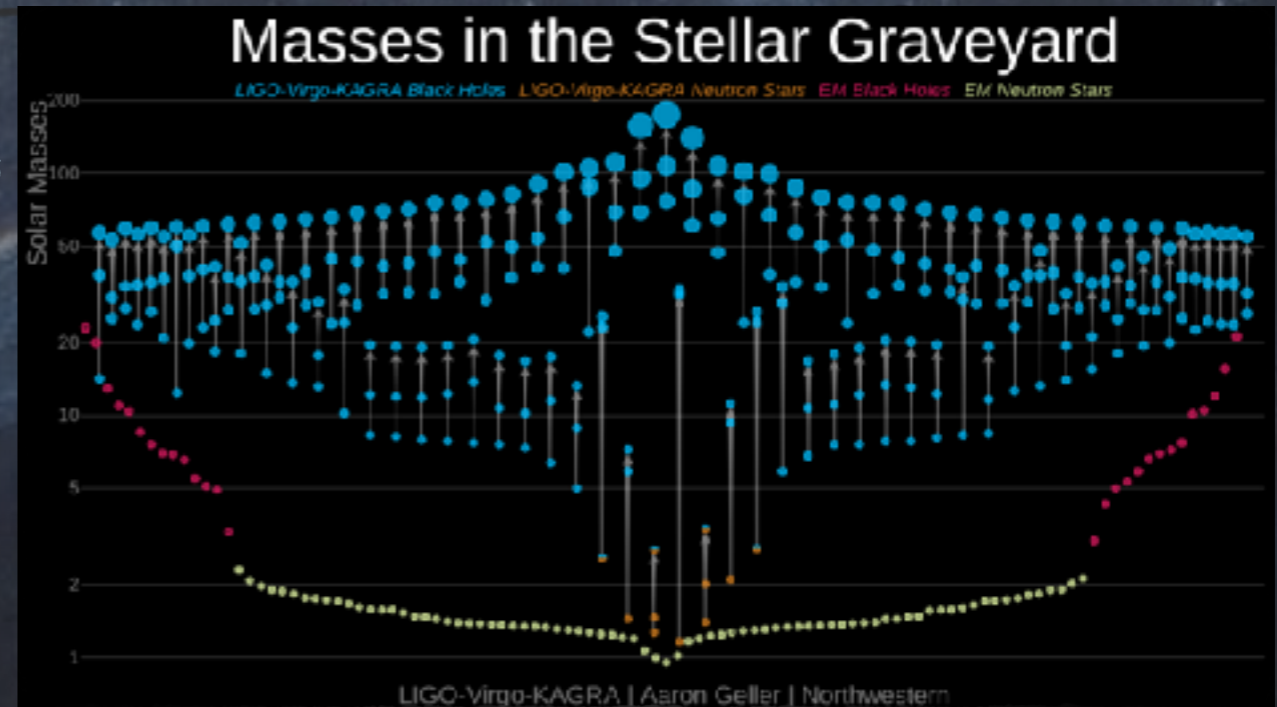
Even if they cannot form all the dark matter... still of great interest

Recent detection of black hole mergers with gravitational waves

B. P. Abbott et al. [LVC], Phys. Rev. Lett. 116:061102, 2016;
Phys. Rev. Lett. 116:241103, 2016; Phys. Rev. Lett. 116:131102, 2016;
Phys. Rev. X6:041015, 2016; Phys. Rev. Lett. 118:221101, 2017;
Astrophys. J. 851:L35, 2017; Phys. Rev. Lett. 119:141101, 2017

Did LIGO detect dark matter?

S. Bird et al., Phys. Rev. Lett. 116:201301, 2016



Insight into early universe physics
(inflation, phase transitions...)

WIMPs and PBHs relation: no go

B. Lackí and J. F. Beacom, Astrophys. J. 720:L67, 2010
R. Saito and S. Shiraí, Phys. Lett. B697:95, 2011
D. Zhang, Mon. Not. R. Astron. Soc. 418:1850, 2011

PBHs as DM (or other exotics) generators

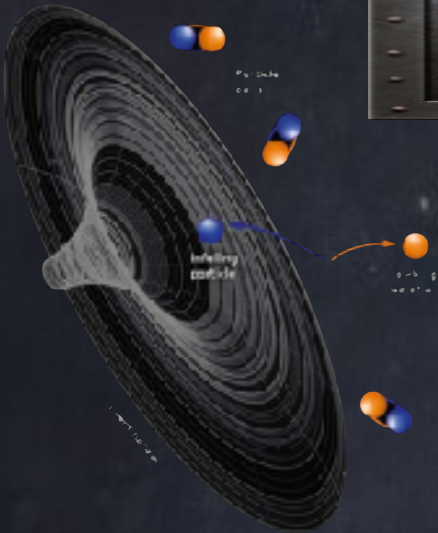
Timing problem: Could PBHs be connected to the origin of SMBHs?

e.g., A. Smith and V. Bromm, Contemp. Phys. 60:111, 2019

PBHs AS DM: ABUNDANCE CONSTRAINTS

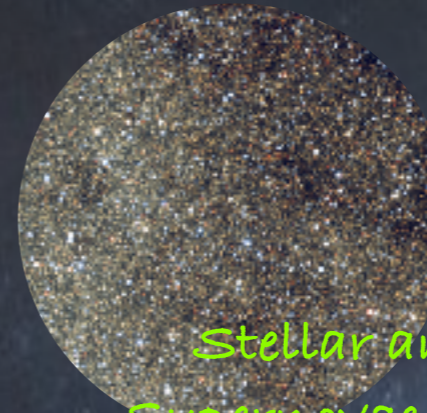
Evaporation

Hawking radiation:
Multi-messenger signals
Heating and ionization of the IGM



Lensing

Stellar and quasar micro-lensing
Supernovae magnification distribution
Strong lensing of FRBs
Femtolensing of GRBs



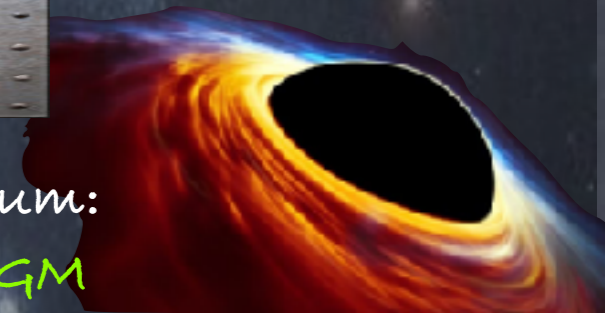
Gravitational Waves

Binary PBHs merger rates
Stochastic GW background
From PBH production



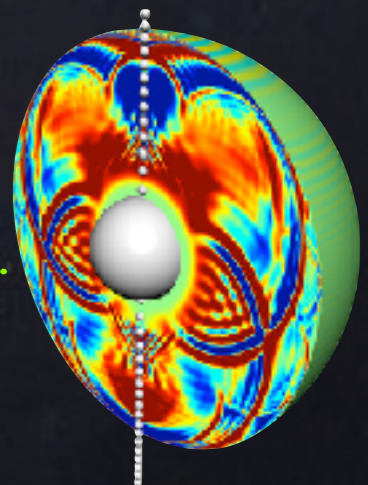
Accretion

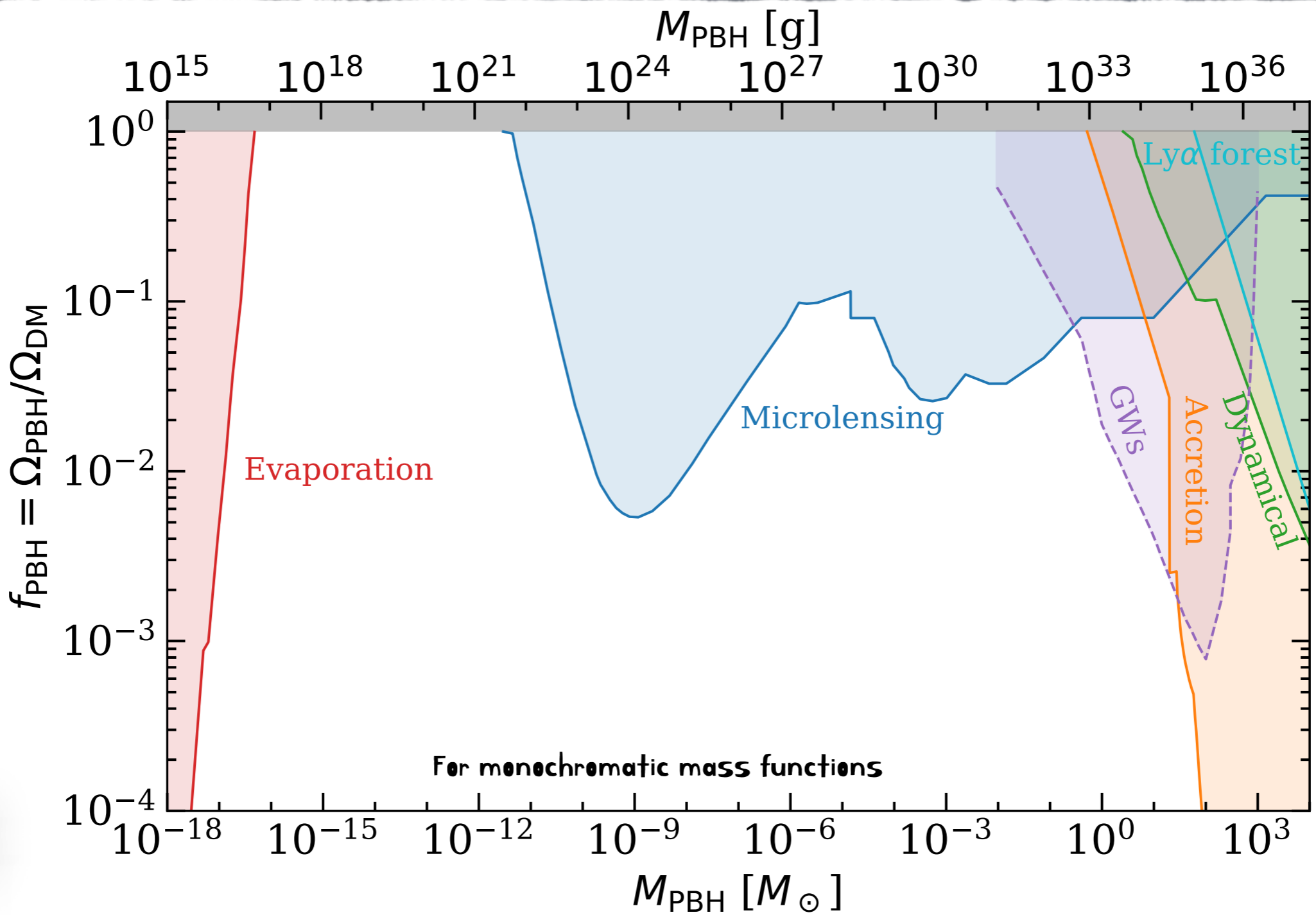
Emission of a broad band spectrum:
Heating and ionization of the IGM
Contribution to the (X-ray, infrared, radio) astrophysical spectra



Dynamical

Disruption of bound systems: galaxies, globular clusters, stellar binaries...
Galactic disk and dynamical friction
Accretion by compact objects which might subsequently be destroyed



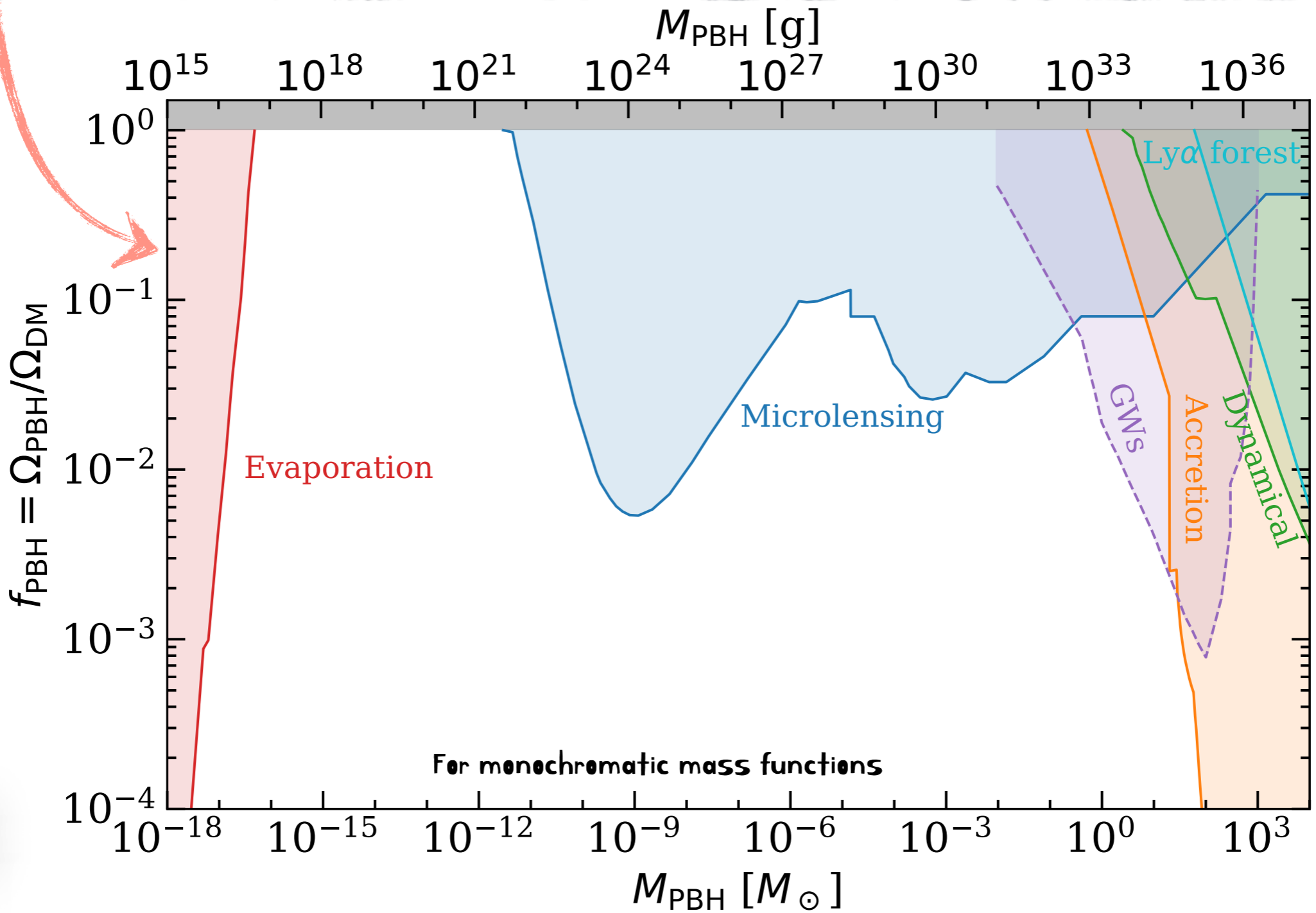


Using
 PBHbounds
 B. J. Kavanagh

P. Villanueva-Domínguez, O. Mena and SPR, *Front. Astron. Space Sci.* 8:87, 2021

Partial evaporation

Hawking radiation: cosmic-ray, γ -ray, ν bkg; ionization and thermal history



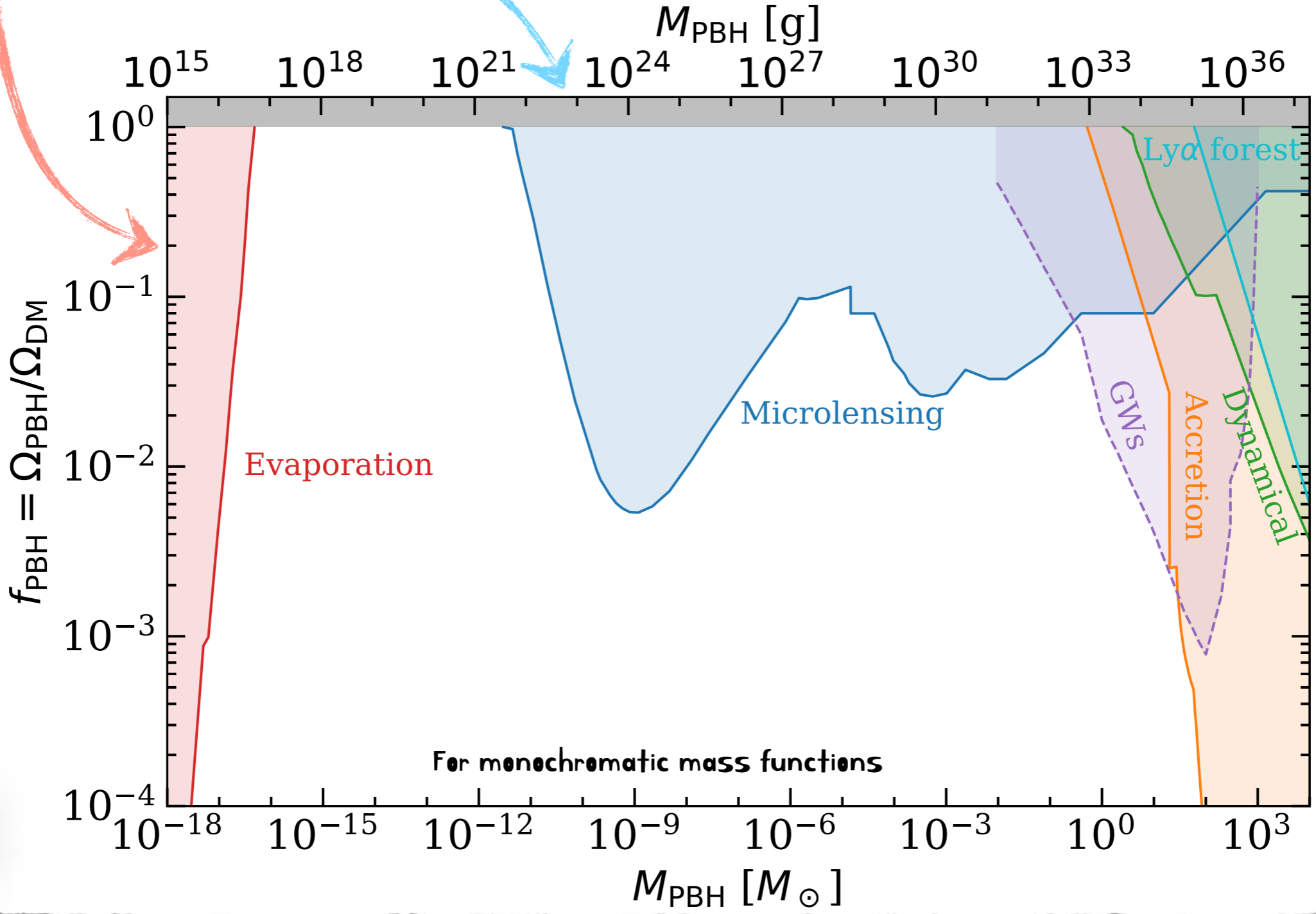
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Partial evaporation

Microlensing of stars, SN, QSO

Gravitational lensing



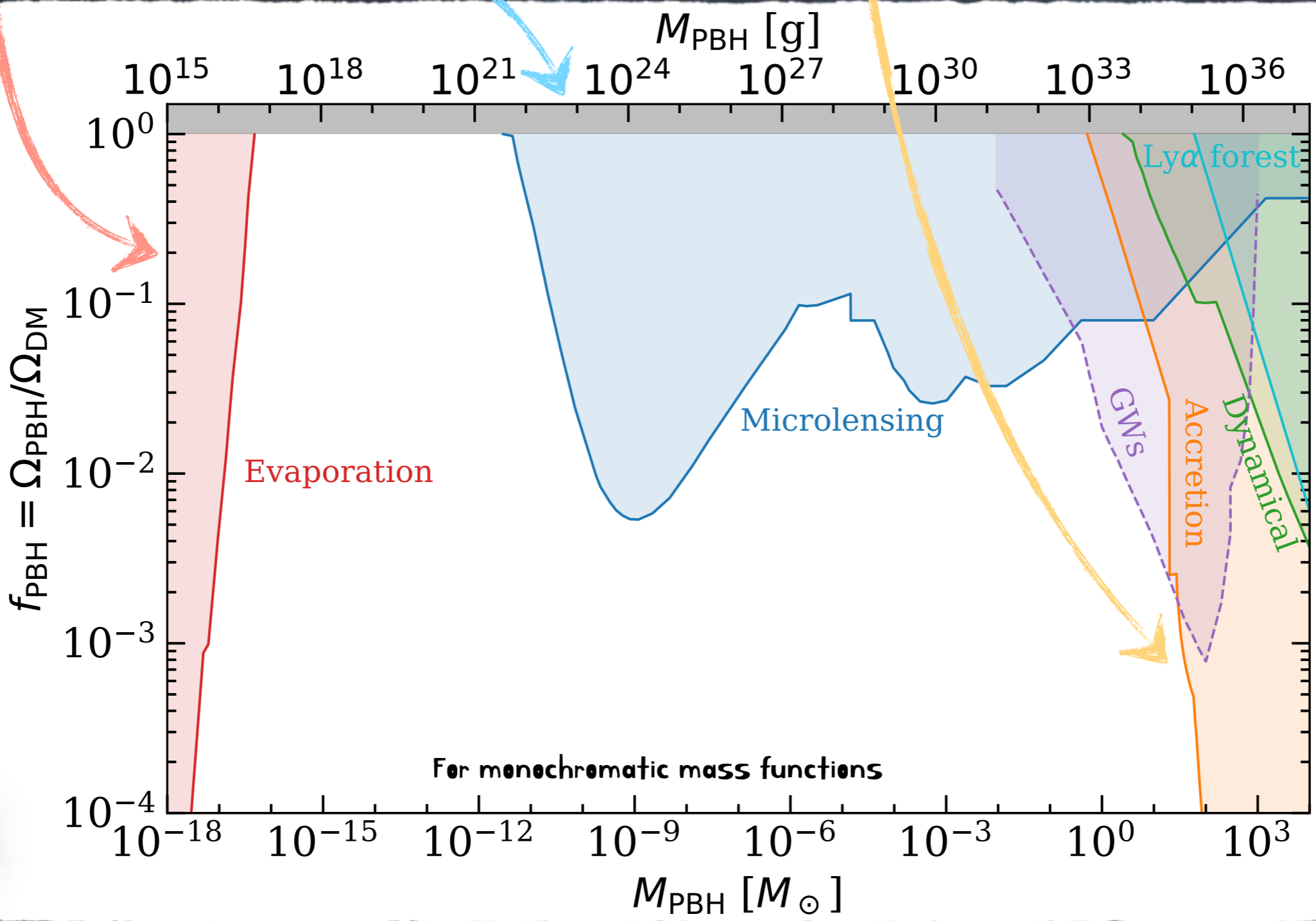
Using PBH bounds
B. J. Kavanagh

P. Villanueva-Domínguez, O. Mena and SPR, Front. Astron. Space Sci. 8:87, 2021

Partial evapo

Ionization and thermal history,
radio and X-ray emission

Accretion



Using
PBHbounds
B. J. Kavanagh

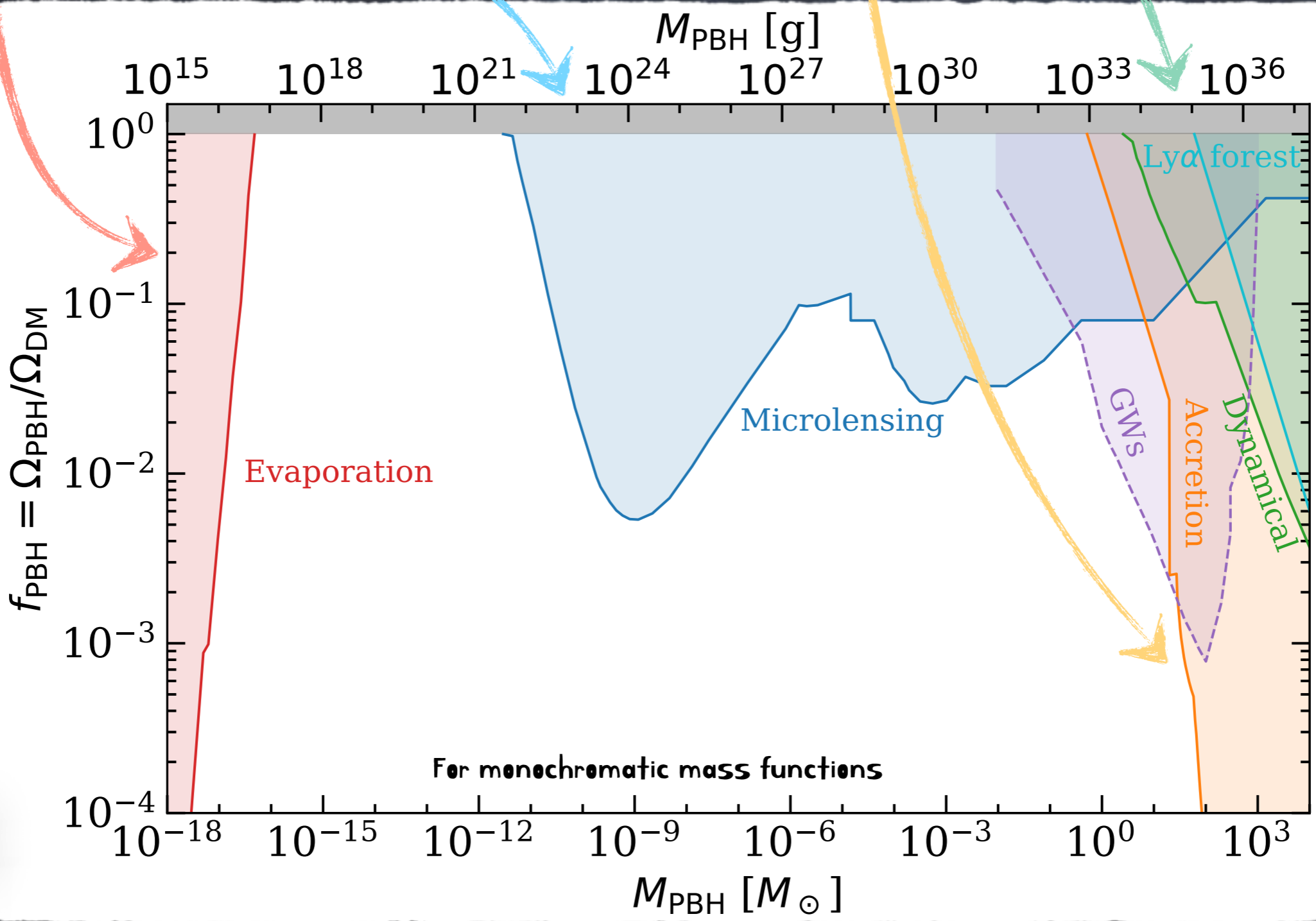
P. Villanueva-Domínguez, O. Mena and SPR, Front. Astron. Space Sci. 8:87, 2021

Partial evaporation

Destruction of astronomical systems by the passage of PBHs

Gravitation

Dynamical effects



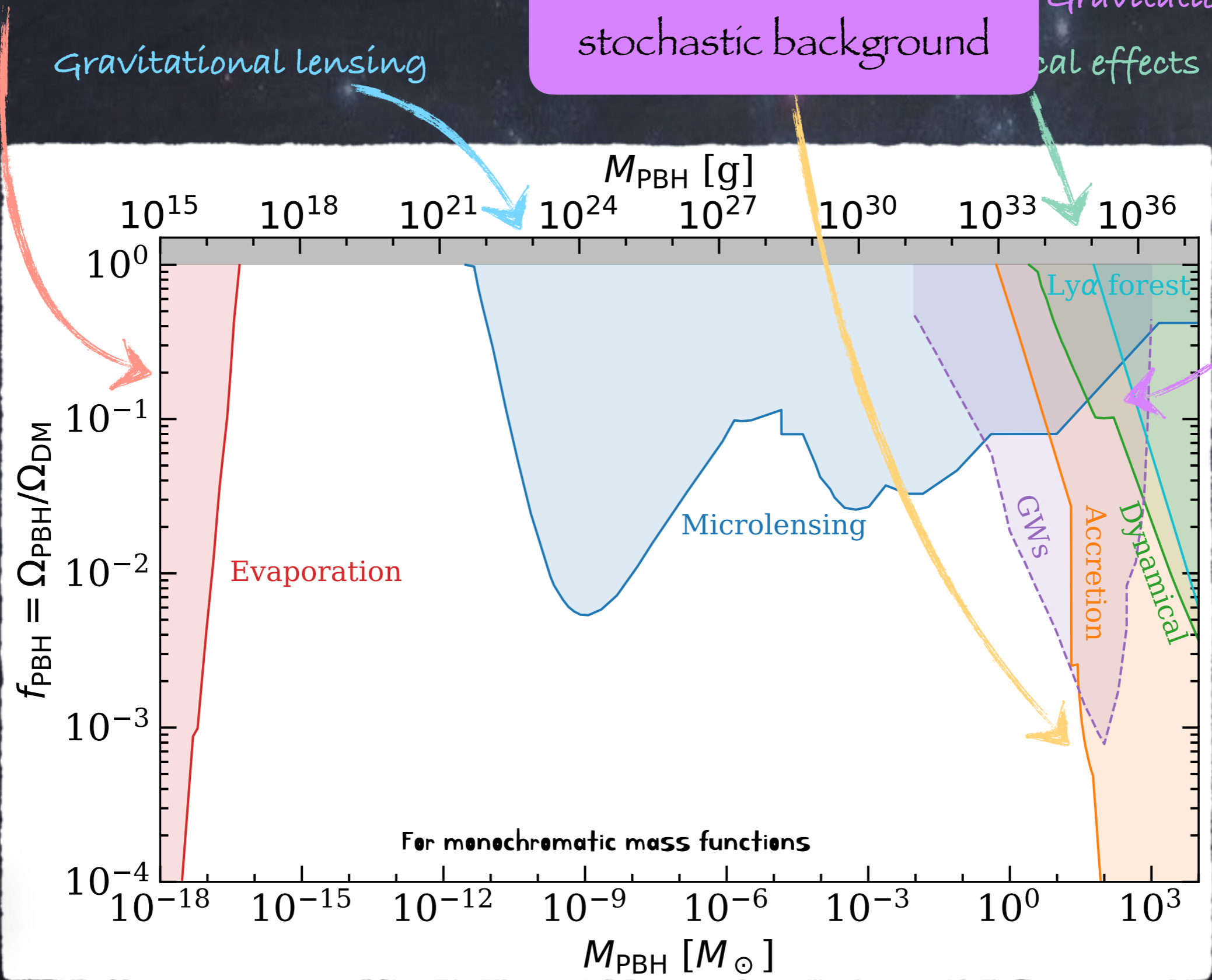
Using PBH bounds
B. J. Kavanagh

P. Villanueva-Domínguez, O. Mena and SPR, Front. Astron. Space Sci. 8:87, 2021

Partial evaporation

LIGO/VIRGO events
stochastic background

Gravitational waves
tidal effects

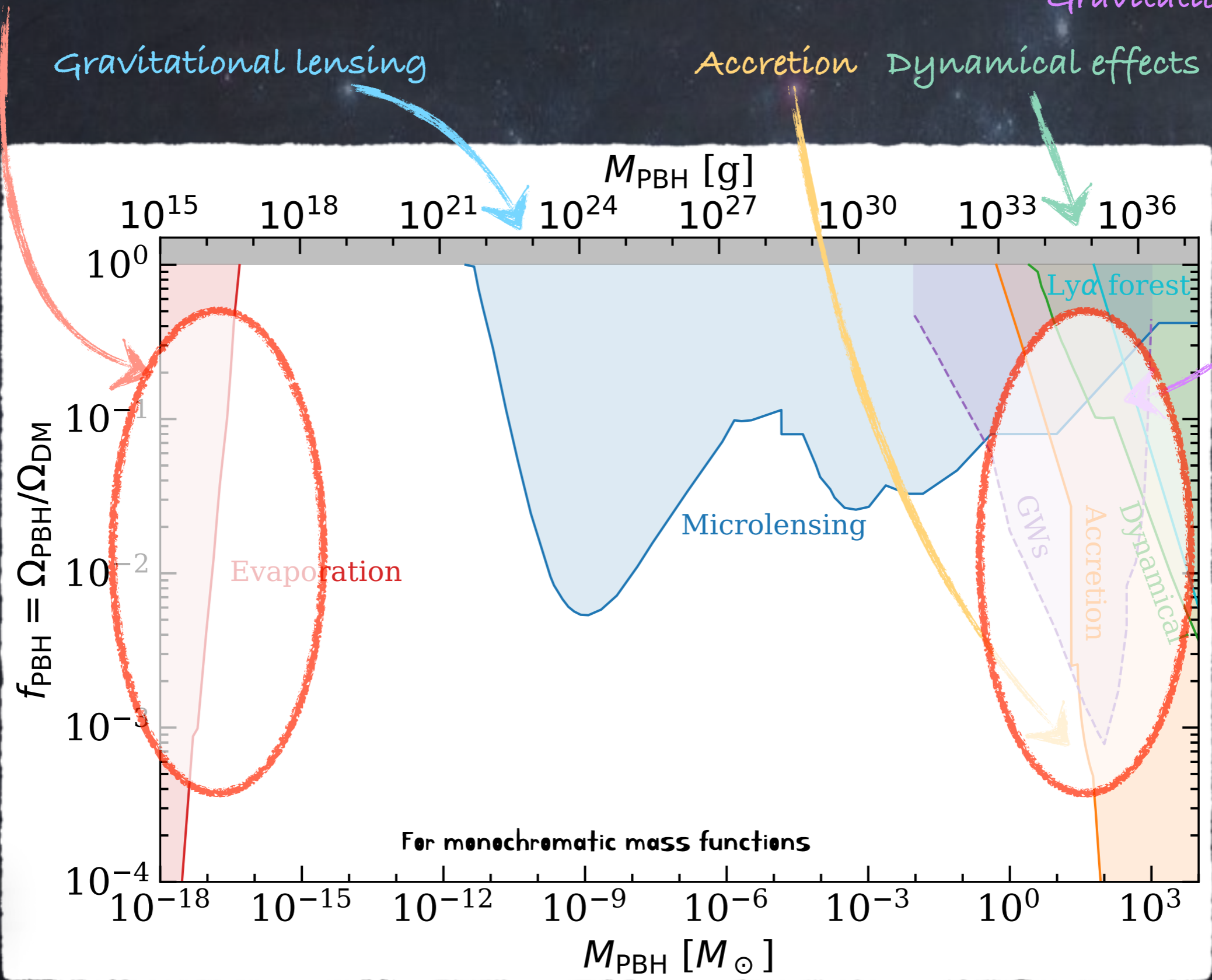


Using
PBH bounds
B. J. Kavanagh

P. Villanueva-Domínguez, O. Mena and SPR, Front. Astron. Space Sci. 8:87, 2021

Partial evaporation

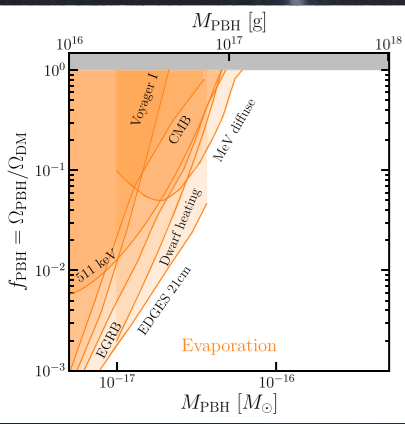
Gravitational waves



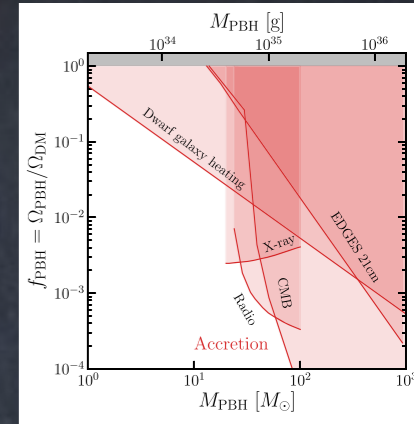
Using PBH bounds
B. J. Kavanagh

P. Villanueva-Domínguez, O. Mena and SPR, Front. Astron. Space Sci. 8:87, 2021

PBHs: EVAPORATION AND ACCRETION



A. M. Green and B. J. Kavanagh, J. Phys. G. 48:043001, 2021



A. M. Green and B. J. Kavanagh, J. Phys. G. 48:043001, 2021

Types of signals

Photons

CMB

Neutrinos

Electrons

21cm

Lyman- α



Messenger production

Energy injection
in the IGM

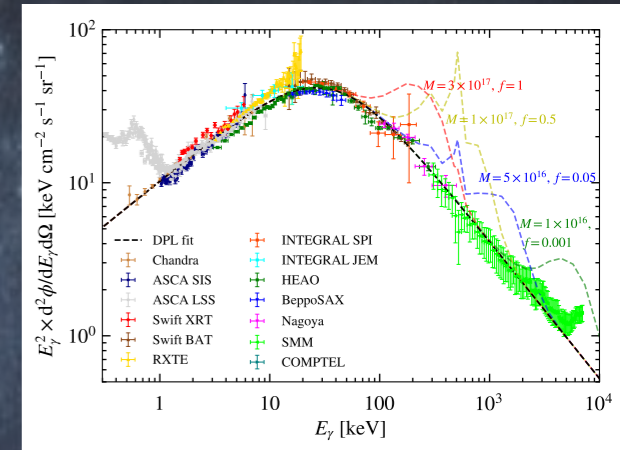
Antimatter

MESSENGER PRODUCTION BY PBHS

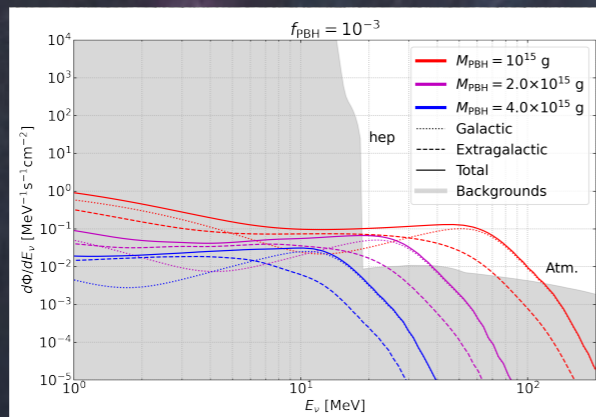
Evaporation

Many references...

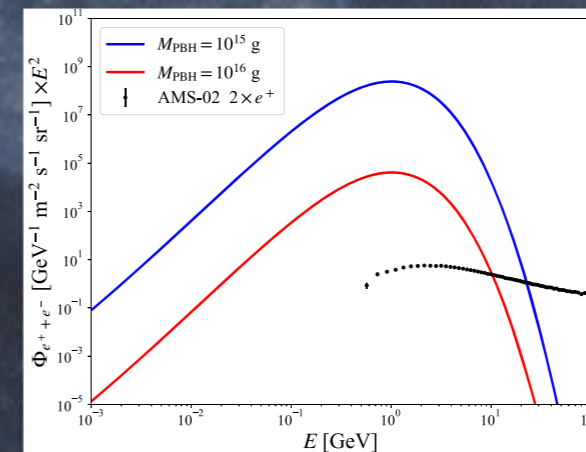
- Extragalactic and galactic γ /X-ray backgrounds
- Local electron-positron flux
- Electron-induced near-UV, synchrotron and IC flux
- 511 keV electron-positron annihilation line
- Antiproton flux
- Extragalactic and galactic neutrino flux



X.-H. Tan and J.-Q. Xia, arXiv:2404.17119



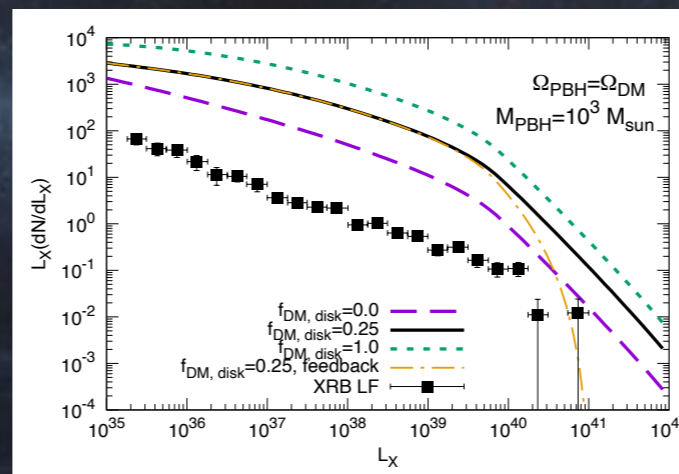
N. Bernal, V. Albornoz-Muñoz, SPR and P. Villanueva-Domínguez, JCAP 10:068, 2022



B.-Y. Su et al., Eur. Phys. J. C84:606, 2024

Accretion

Galactic IR/X/radio backgrounds



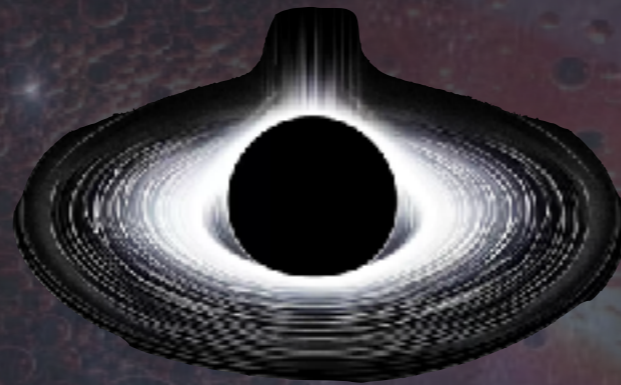
Y. Inoue and A. Kusenko, JCAP 10:034, 2017

See also:

- Kashlinsky, Astrophys. J. 823:L25, 2016
- D. Gaggero et al., Phys. Rev. Lett. 118:241101, 2017
- G. Hasinger, JCAP 07:022, 2020

ENERGY INJECTION FROM PBHS EVAPORATION

Evaporation rate



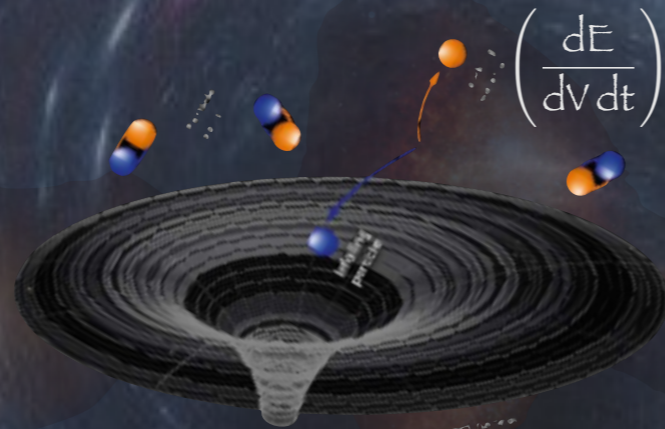
$$\frac{dM_{\text{PBH}}}{dt} = - \frac{f(M_{\text{PBH}})}{M_{\text{PBH}}^2}$$

D. Page, Phys. Rev. D13:198, 1976; Phys. Rev. D14:3260, 1976; Phys. Rev. D16:2402, 1977

Emission rate equal to that of a thermal body with temperature:

$$T_{\text{BH}} \sim \frac{1}{8\pi G M_{\text{BH}}} \sim 10 \left(\frac{10^{15} \text{ g}}{M_{\text{BH}}} \right) \text{ MeV} \quad \tau(M_{\text{BH}}) \sim G^2 M_{\text{BH}}^3 \sim 100 \left(\frac{M_{\text{BH}}}{10^{15} \text{ g}} \right)^3 \text{ Gyr}$$

Energy injection



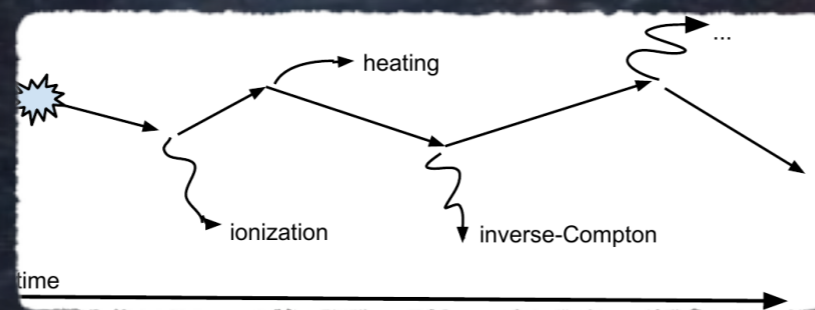
$$\left(\frac{dE}{dV dt} \right)_{\text{inj}} = n_{\text{PBH}} \int E \left(\frac{dN}{dt dE} \right)_{\text{evap}} dE = \frac{f_{\text{PBH}} \rho_{\text{DM}}}{M_{\text{PBH}}} \int E \frac{g}{2\pi} \frac{\Gamma(E, M_{\text{PBH}})}{e^{E/T_{\text{PBH}} \pm 1}} dE$$

Emission of a thermal quasi-black-body spectrum

J. H. McGibbon and B. R. Webber, Phys. Rev. D41:3052, 1990
J. H. McGibbon, Phys. Rev. D44:376, 1991

Energy deposition

Excitations, ionization and heating



From A. C. Vincent

Not on-the-spot

$$\left(\frac{dE}{dV dt} \right)_{\text{dep,c}} = f_c(z) \left(\frac{dE}{dV dt} \right)_{\text{inj}}$$

T. R. Slatyer, Phys. Rev. D93:023521, 2016

ENERGY INJECTION FROM ACCRETION BY PBHs

Accretion rate



usually Bondi-Hoyle-Lyttleton model

$$\dot{M}_{\text{PBH}} = 4\pi \lambda \rho_{\infty} \frac{(GM_{\text{PBH}})^2}{(c_s^2 + v_{\text{rel}}^2)^{3/2}}$$

F. Hoyle and R. A. Lyttleton, Mon. Not. R. Astron. Soc. 101:227, 1941
 H. Bondi and F. Hoyle, Mon. Not. R. Astron. Soc. 104:273, 1944
 H. Bondi, Mon. Not. R. Astron. Soc. 112:195, 1952

...but also see Park-Ricotti accretion

K. Park and M. Ricotti, Astrophys. J. 739:2, 2011;
 Astrophys. J. 747:9, 2012; Astrophys. J. 767:163, 2013

Energy injection



Luminosity: $L_{\text{acc}} = \epsilon \dot{M}_{\text{PBH}}$

$$\left(\frac{dE}{dV dt} \right)_{\text{inj}} = L_{\text{acc}} n_{\text{PBH}} = L_{\text{acc}} \frac{f_{\text{PBH}} \rho_{\text{DM}}}{M_{\text{PBH}}}$$

if cooling is inefficient, a thick disk could form: ADAF

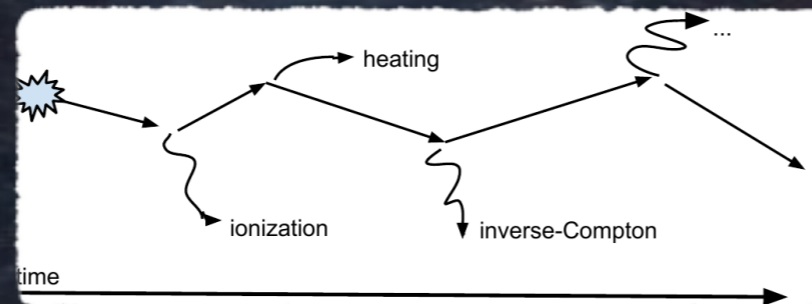
$$\epsilon = \epsilon_0 \left(\frac{10 \dot{M}_{\text{PBH}}}{L_{\text{Edd}}} \right)^a$$

See, e.g., F.-G. Xie and F. Yuan, Mon. Not. R. Astron. Soc. 427:1580, 2012

Excitations, ionization and heating

Not on-the-spot

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From A. C. Vincent

$$\left(\frac{dE}{dV dt} \right)_{\text{dep,c}} = f_c(z) \left(\frac{dE}{dV dt} \right)_{\text{inj}}$$

T. R. Slatyer, Phys. Rev. D93:023521, 2016

Energy spectrum: synchrotron, IC and bremsstrahlung

R. Mahadevan, Astrophys. J. 477:585, 1997 Probes of PBHs as DM

EFFECTS OF ENERGY INJECTION FROM PBHS

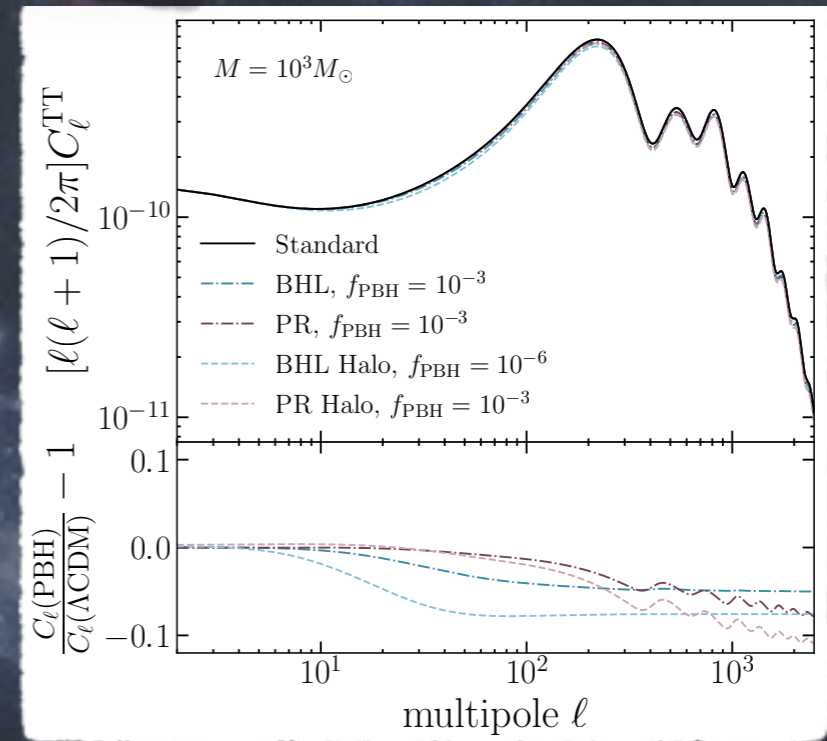
Damping of CMB anisotropies

Evaporation

- V. Poulin, J. Lesgourgues and P. D. Serpico, JCAP 03:043, 2017
- S. Clark et al., Phys. Rev. D95:083006, 2017
- P. Stöcker, M. Kramer, J. Lesgourgues and V. Poulin, JCAP 03:018, 2018
- H. Poulter et al., arXiv:1907.06485
- S. K. Acharya and R. Katri, JCAP 06:018, 2020

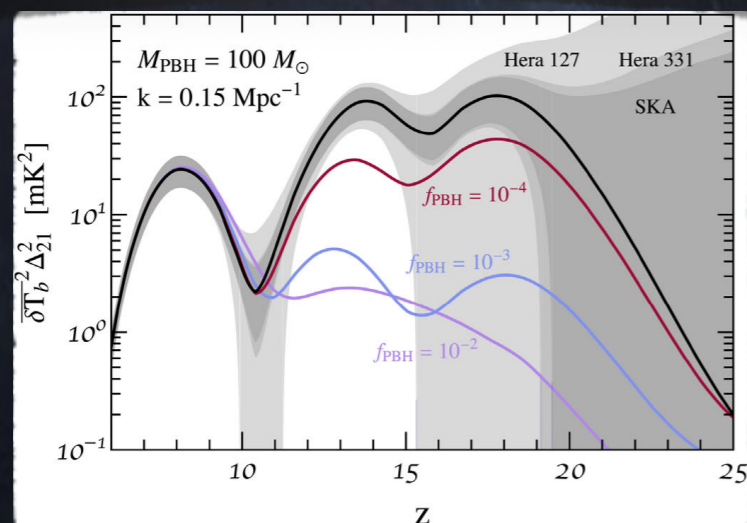
Accretion

- B. J. Carr, Mon. Not. R. Astron. Soc 194:639, 1981
- M. Ricotti, J. P. Ostriker and K. J. Mack, Astrophys. J. 680:829, 2008
- B. Horowitz, arXiv:1612.07264
- Y. Ali-Haïmoud and M. Kamionkowski, Phys. Rev. D95:043534, 2017
- V. Poulin et al., Phys. Rev. D96:083524, 2017
- P. Serpico et al., Phys. Rev. Res. 2:023204, 2020
- L. Píga et al., JCAP 12:016, 2022
- G. Facchinetti, M. Lucca and S. Clesse, Phys. Rev. D107:043537, 2022
- D. Agius et al., JCAP 07:003, 2024



D. Agius et al., JCAP 07:003, 2024

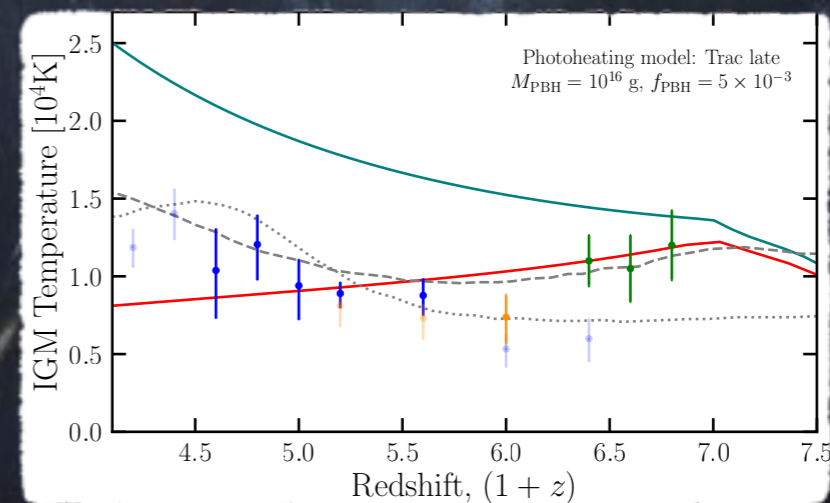
Heating of the 21cm temperature



O. Mena, SPR, P. Villanueva-Domínguez and S. J. Witte, Phys. Rev. D100:043540, 2019

Sergio Palomares-Ruiz

Broadening of Lyman- α absorption features



A. K. Saha, A. Singh, P. Parashari and R. Laha, arXiv:2409.10617

Example of evaporation constraint: Neutrino flux

B. J. Carr, K. Kohri, Y. Sendouda and J. Yokoyama, *Phys. Rev. D*81:104019, 2010

B. Dasgupta, R. Laha and A. Ray, *Phys. Rev. Lett.* 125:101101, 2020

S. Wang et al., *Phys. Rev. D*103:043010, 2021

V. De Romerí, P. Martínez-Miravé and M. Tórtola, *JCAP* 10:051, 2021

R. Calabrese et al., *Phys. Lett. B*829:137050, 2022

N. Bernal, V. Albornoz-Muñoz, *SPR* and P. Villanueva-Domínguez, *JCAP* 10:068, 2022

NEUTRINOS FROM PBHS EVAPORATION

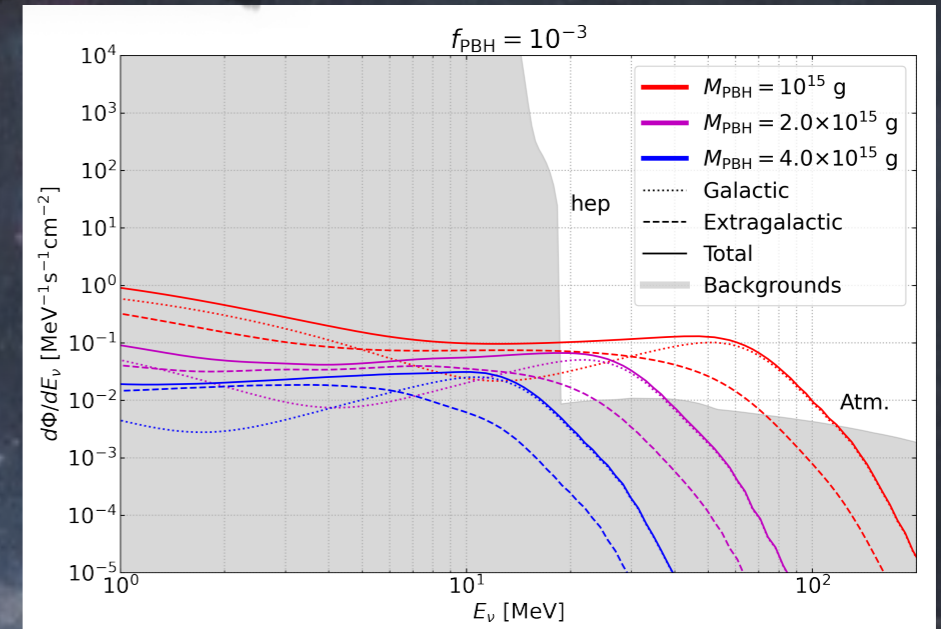
with BlackHawk A. Arbey and J. Auffinger, Eur. Phys. J. C79:693, 2019

Galactic contribution

$$\frac{d\Phi_G(E_\nu)}{dE_\nu} = \frac{d^2N_\nu(E_\nu, t_0)}{dE_\nu dt} \frac{f_{\text{PBH}}}{M} \int \frac{d\Omega}{4\pi} \int_0^{\ell_{\text{max}}} dl \rho_{\text{DM}}(r(l, \phi))$$

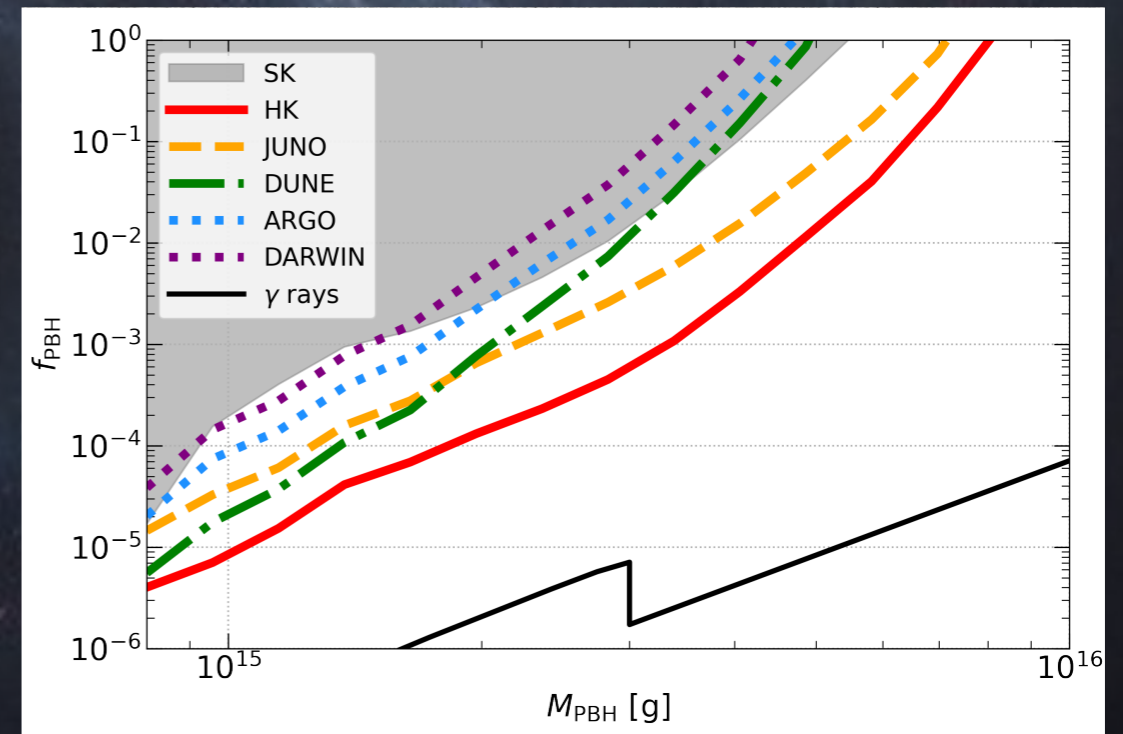
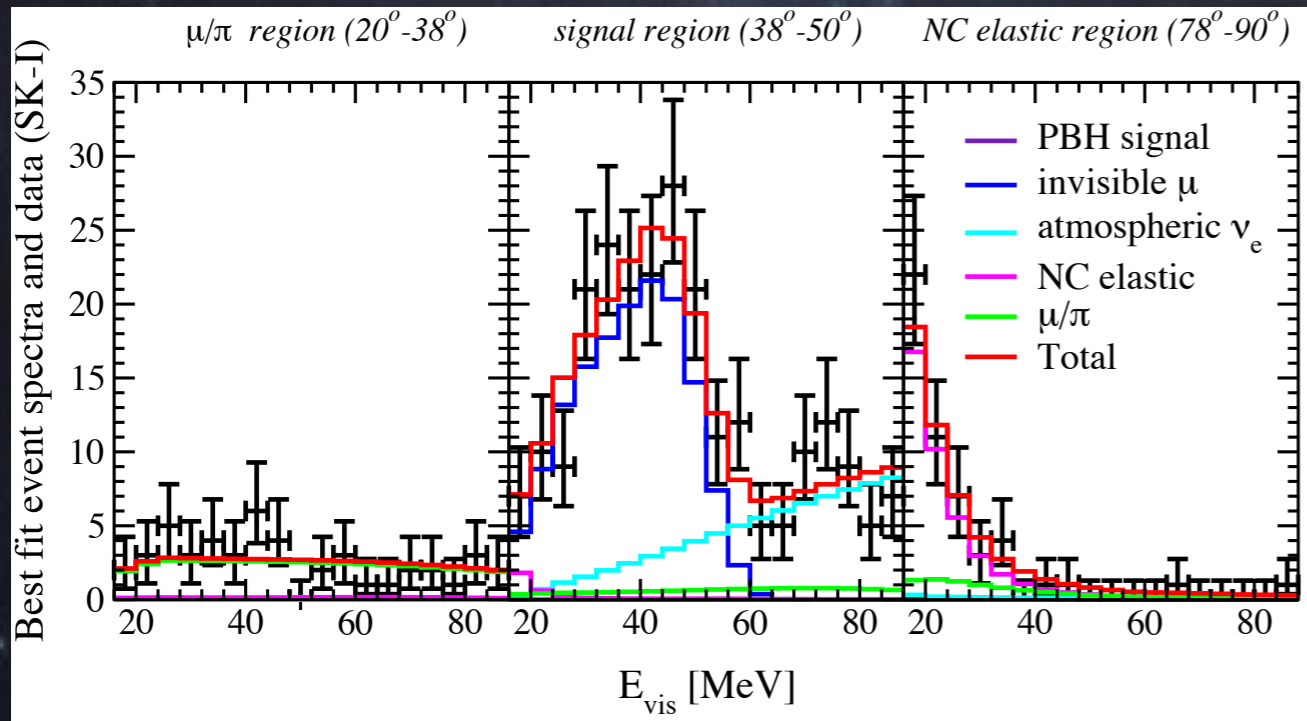
Extragalactic contribution

$$\frac{d\Phi_{\text{EG}}(E_\nu)}{dE_\nu} = n_{\text{PBH}}(t_0) \int_{z_{\text{max}}}^{z_{\text{min}}} dz \left| \frac{dt}{dz} \right| (1+z) \frac{d^2N_\nu(E_\nu(1+z), t)}{dE_\nu dt}$$



Analysis of SK DSNB data

...and prospects for HK, JUNO, DUNE, ARGO, DARWIN



N. Bernal, V. Albornoz-Muñoz, SPR and P. Villanueva-Domínguez, JCAP 10:068, 2022

Example of accretion constraint: 21cm signal

In the context of EDGES (accretion):

A. Hektor et al., Phys. Rev. D98:023503, 2018
Y. Yang, Phys. Rev. D104:063528, 2021

Forecasts (accretion):

O. Mena, SPR, P. Villanueva-Domínguez and S. J. Witte, Phys. Rev. D100:043540, 2019

In the context of EDGES (evaporation):

S. Clark et al., Phys. Rev. D98:043006, 2018
Y. Yang, Phys. Rev. D102:083538, 2020
A. Halder and M. Pandey, MNRAS 508:3446, 2021
A. Halder and S. Banerjee, Phys. Rev. D103:0530044, 2021
S. Mittal et al., JCAP 03:030, 2022
U. Mukhopadhyay, D. Majumdar and A. Halder, JCAP 10:099, 2022
A. K. Saha and R. Laha, Phys. Rev. D105:103026, 2022

Forecasts (evaporation):

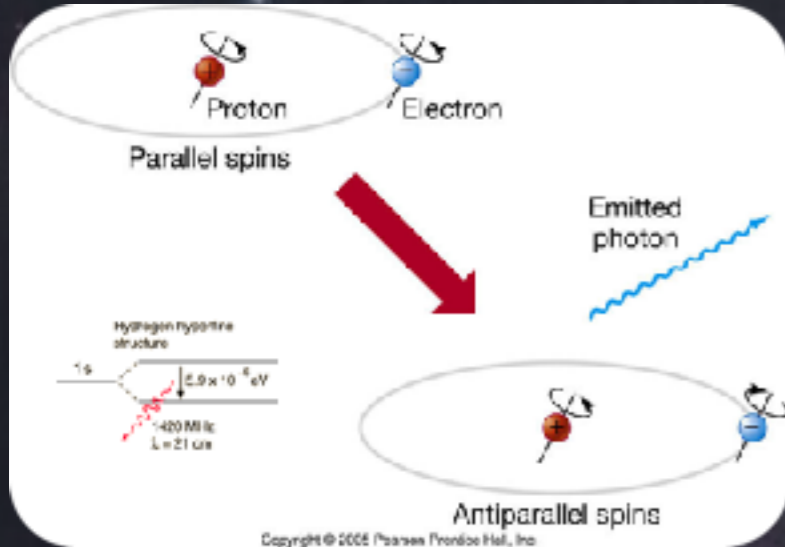
P. K. Natwariya, A. C. Nayak and T. Srivastava, MNRAS 510, 4236, 2021
J. Cang, Y. Gao and Y.-Z. Ma, JCAP 03:012, 2022
Y. Yang, Phys. Rev. D106:123508, 2022

For evaporation

constraints, see also:

THE 21CM LINE

Predicted by H. van de Hulst in 1944 and first observed by H. I. Ewen and E. M. Purcell in 1951

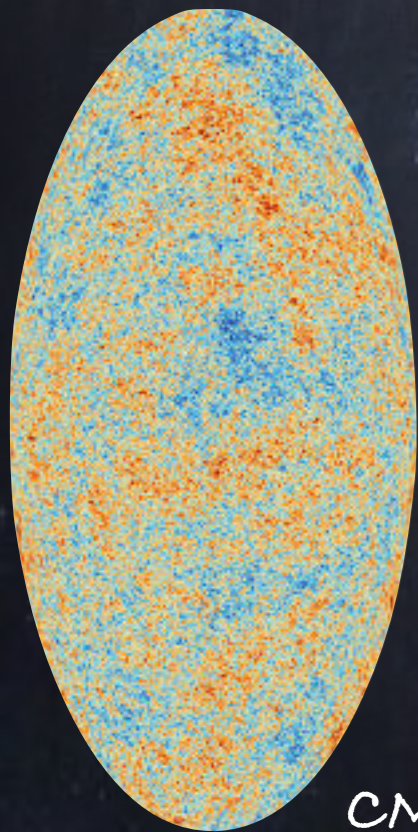


Hyperfine transition: $\nu = 1420 \text{ MHz}$

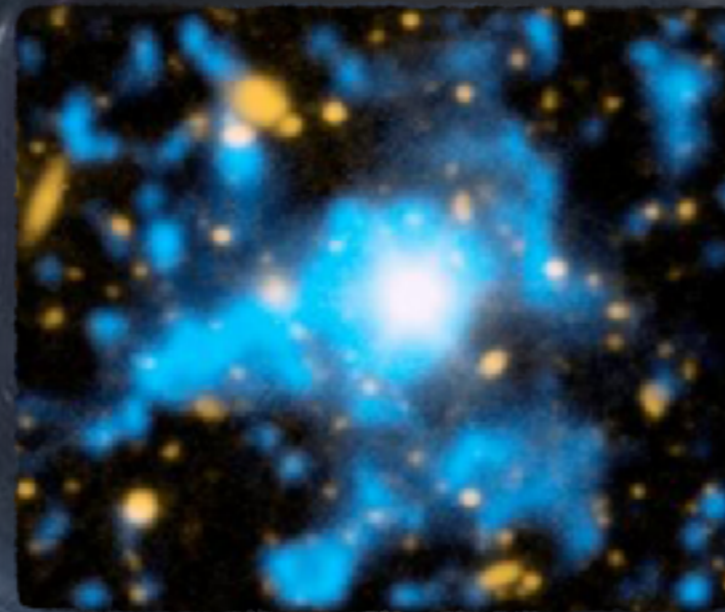
21cm photon from HI clouds during cosmic dawn: $\nu \sim 100 \text{ MHz}$

$$\nu = \frac{\nu_0}{(1+z)}$$

neutral hydrogen gas (intergalactic medium: IGM)



CMB photons as backlight



emission/absorption



observer

$z \sim 1000$

Population of ground and excited states controlled by:

absorption and stimulated emission of background radiation
collisions of neutral hydrogen
excitation/de-excitation by Lyman- α photons

$z=0$

THE 21CM SIGNAL

Differential brightness temperature

$$\delta T_b(\nu) \approx 27 x_{HI} (1 + \delta) \left(1 - \frac{T_{CMB}}{T_S} \right) \left(\frac{1+z}{10} \right)^{1/2}$$

Fraction of neutral H

Reionization suppresses the signal

Baryon overdensity

$$\frac{n_1}{n_0} = 3 e^{-T_{21}/T_S}$$

Spin temperature:
occupation of the two states

$$\delta T_b \approx 0 \quad \text{if} \quad T_S \sim T_{CMB}$$

no signal

$$\delta T_b > 0 \quad \text{if} \quad T_S > T_{CMB}$$

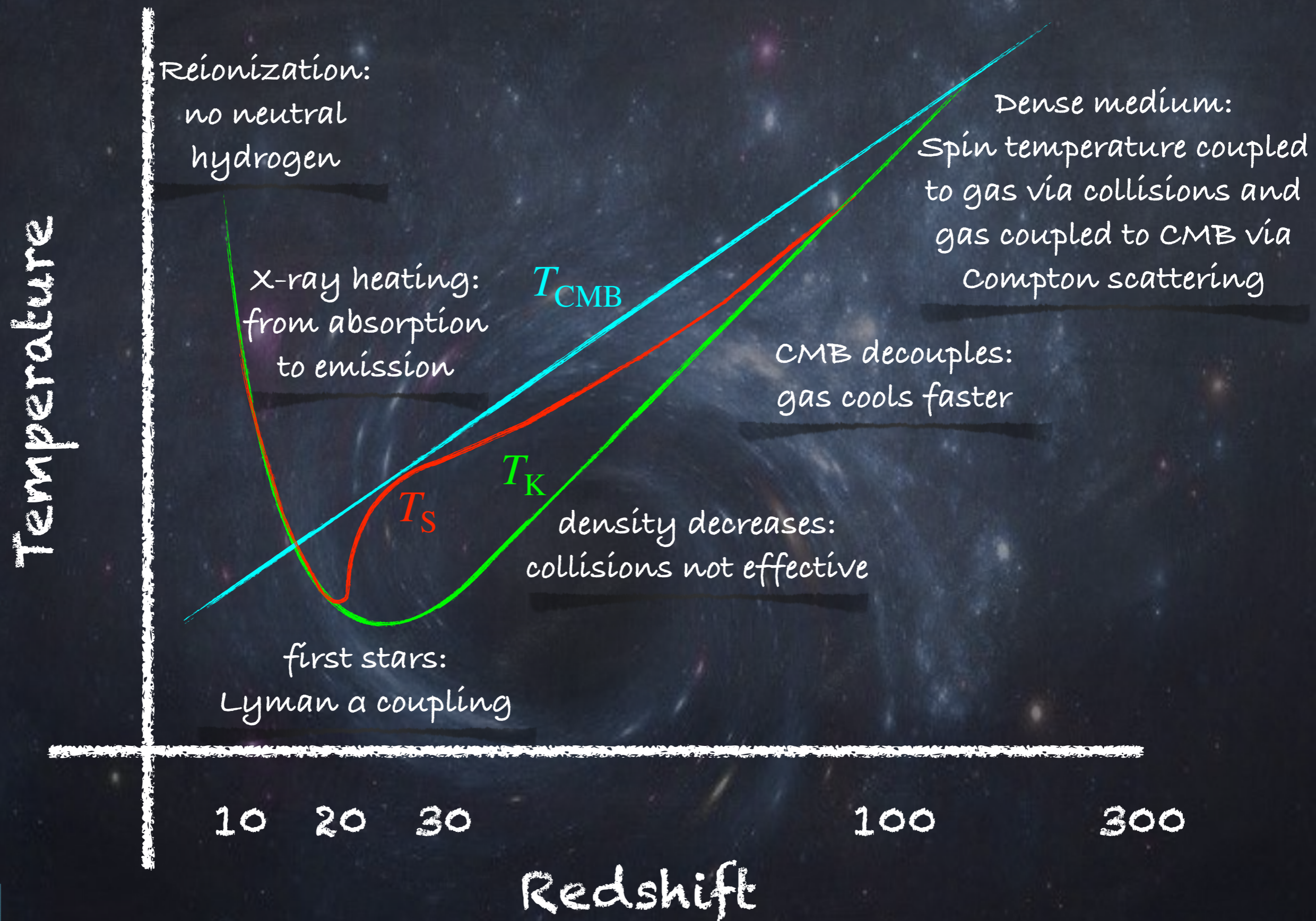
signal in emission, can saturate

$$\delta T_b < 0 \quad \text{if} \quad T_S < T_{CMB}$$

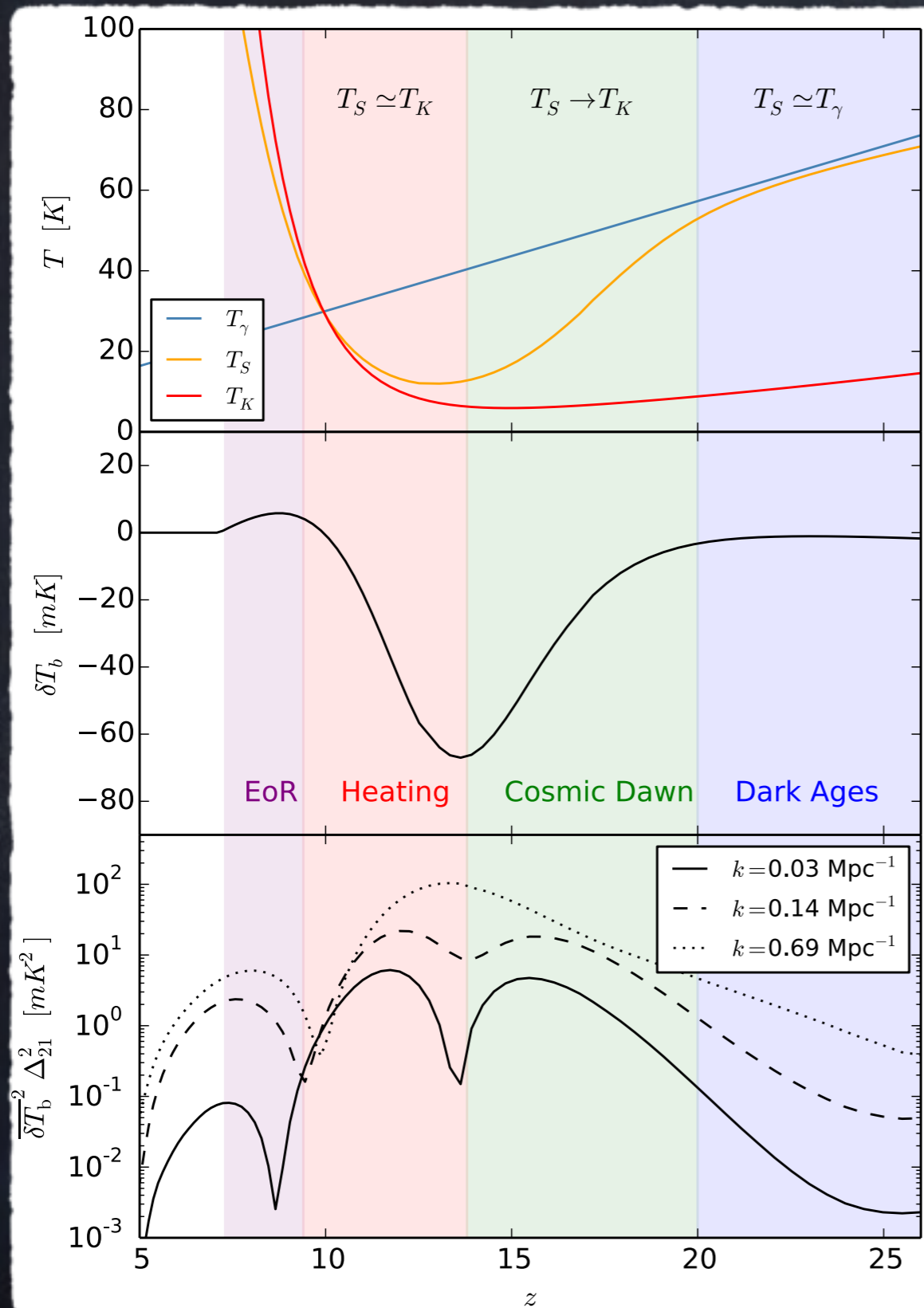
signal in absorption,
limited by gas temperature

Astrophysical processes decouple T_S from T_{CMB}

THE 21CM SIGNAL: TIME EVOLUTION



THE 21CM SIGNAL: TIME EVOLUTION



Dense medium:
Spin temperature coupled
to gas via collisions and
gas coupled to CMB via
Compton scattering

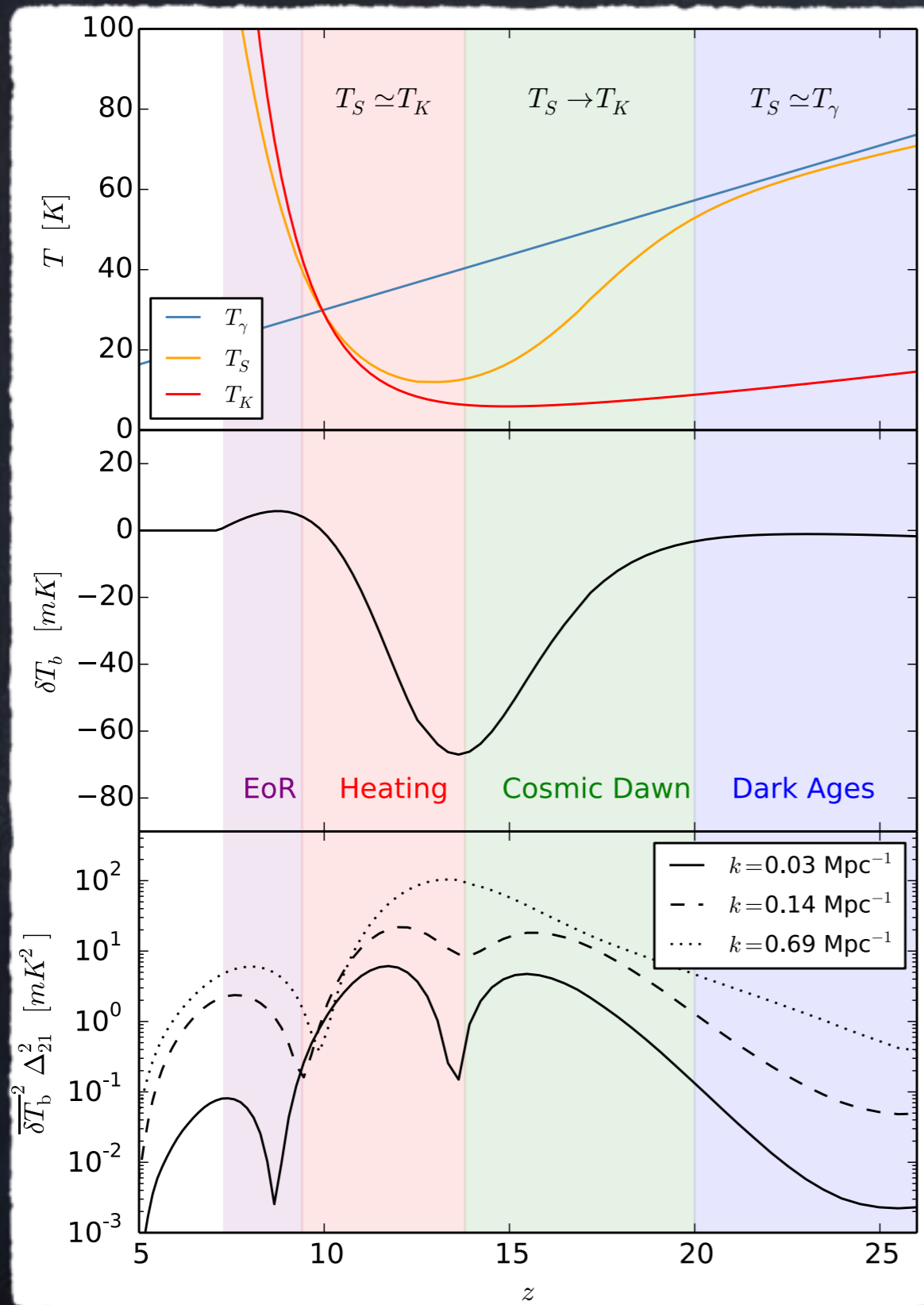
CMB decouples:
gas cools faster

Density decreases:
collisions not effective

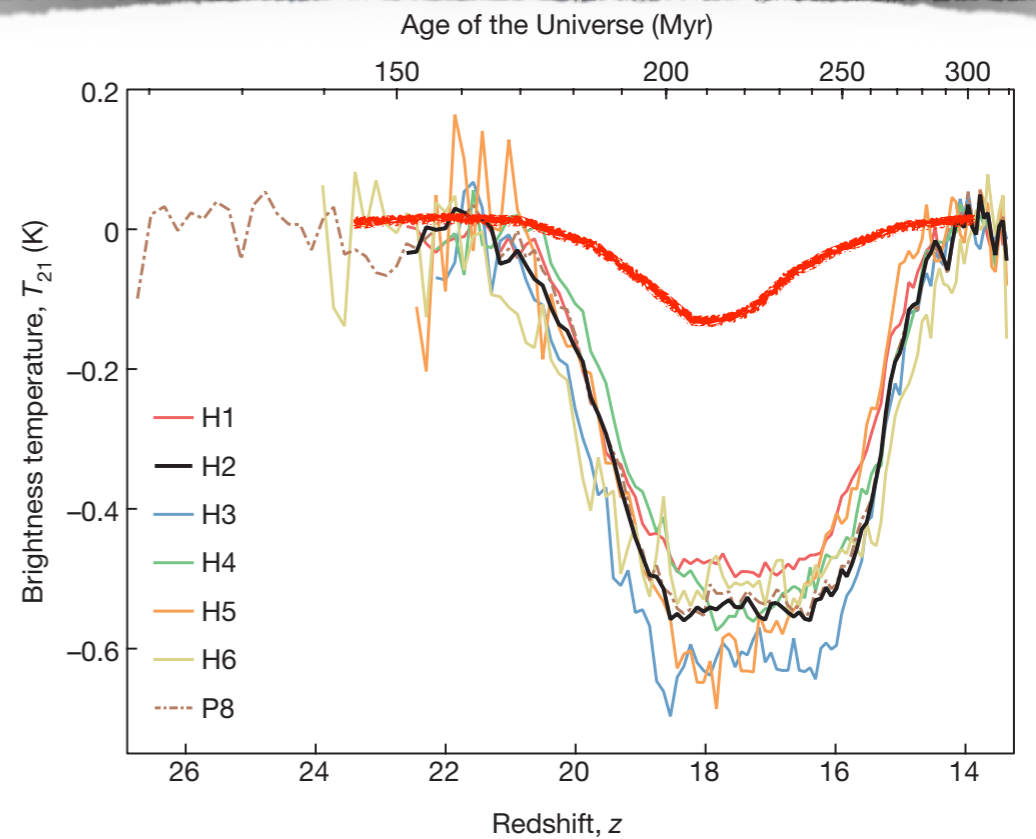
100

300

THE 21CM SIGNAL: TIME EVOLUTION



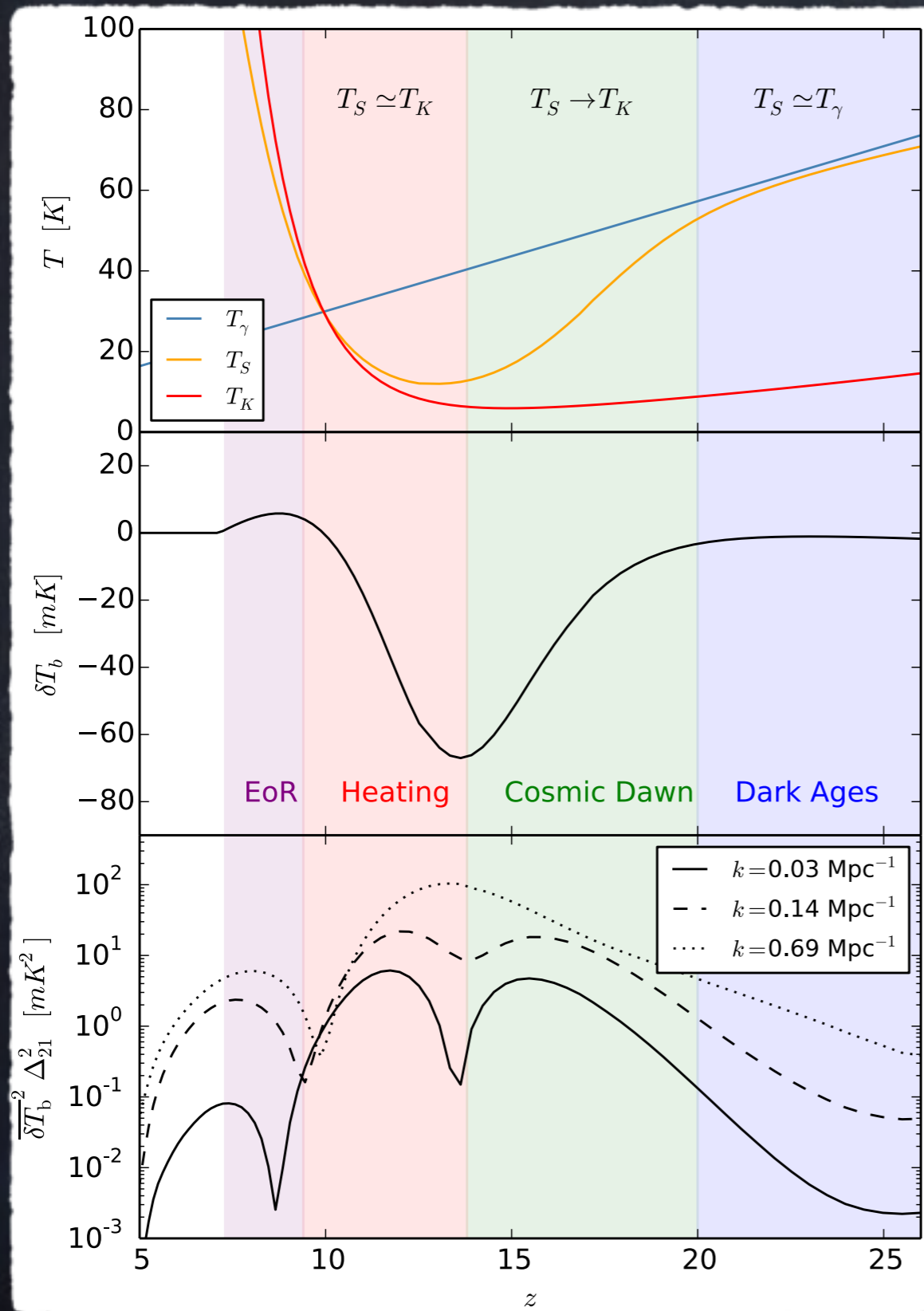
EDGES:
new physics? astrophysics?
systematics?



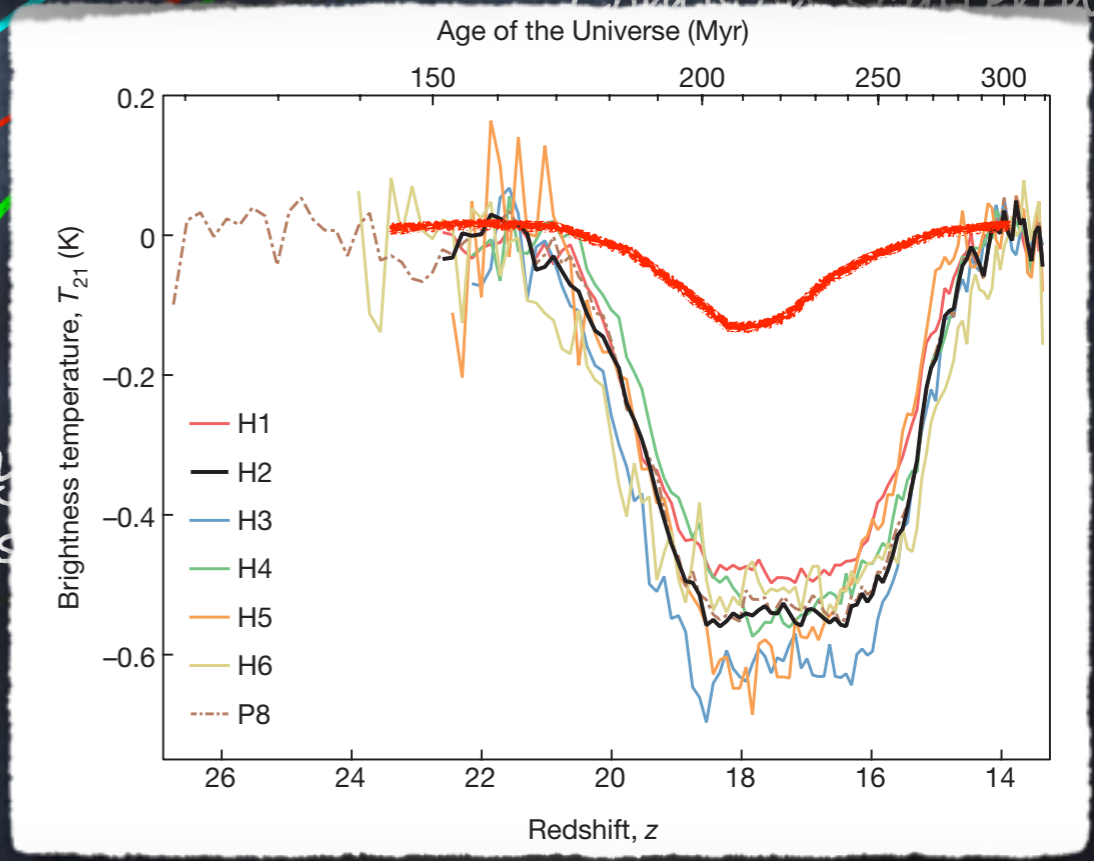
J. D. Bowman et al., Nature 555:67, 2018

P. Villanueva-Domínguez, PhD Thesis, 2021

THE 21CM SIGNAL: TIME EVOLUTION



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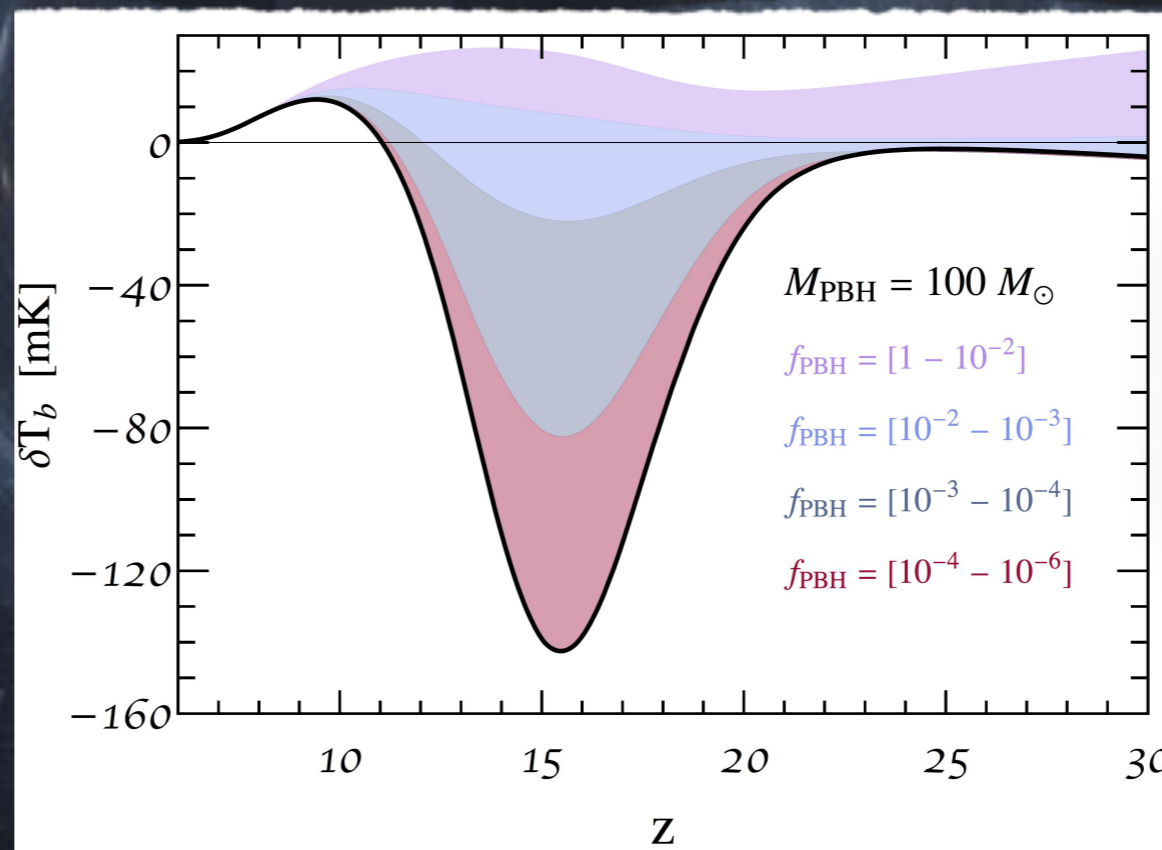
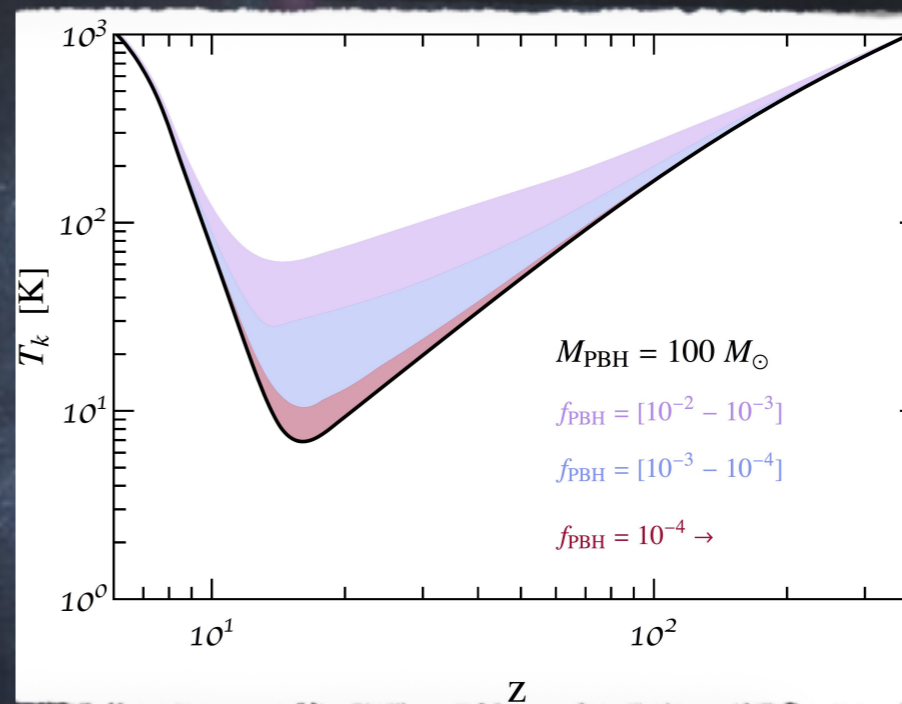
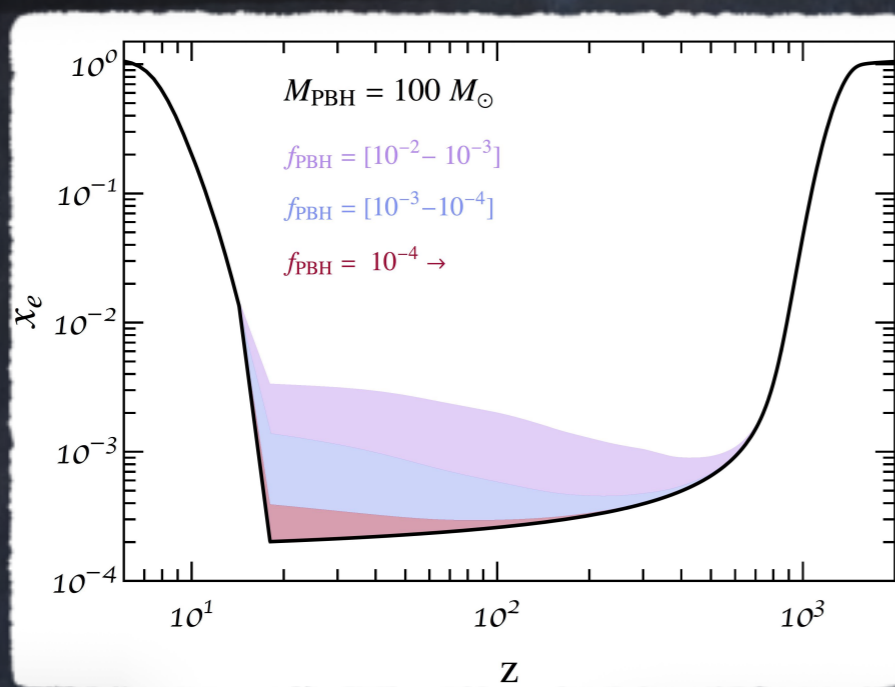


Fully incompatible (amplitude, location AND shape) with standard scenarios

S. Witte, P. Villanueva-Domínguez, S. Gariazzo, O. Mena and SPR, Phys. Rev. D97:103533, 2018

PBHs: BRIGHTNESS TEMPERATURE

Accretion: Injected energy goes into ionizing and heating the IGM



From absorption to emission

We use CosmoRec

J. Chluba and R. M. Thomas,
 Mon. Not. R. Astron. Soc. 412:478, 2011

Impact on the
 brightness temperature:
 suppression of signal

We use 21cmFAST

A. Mesinger, S. Furlanetto and R. Cen,
 Mon. Not. R. Astron. Soc. 411:955, 2011

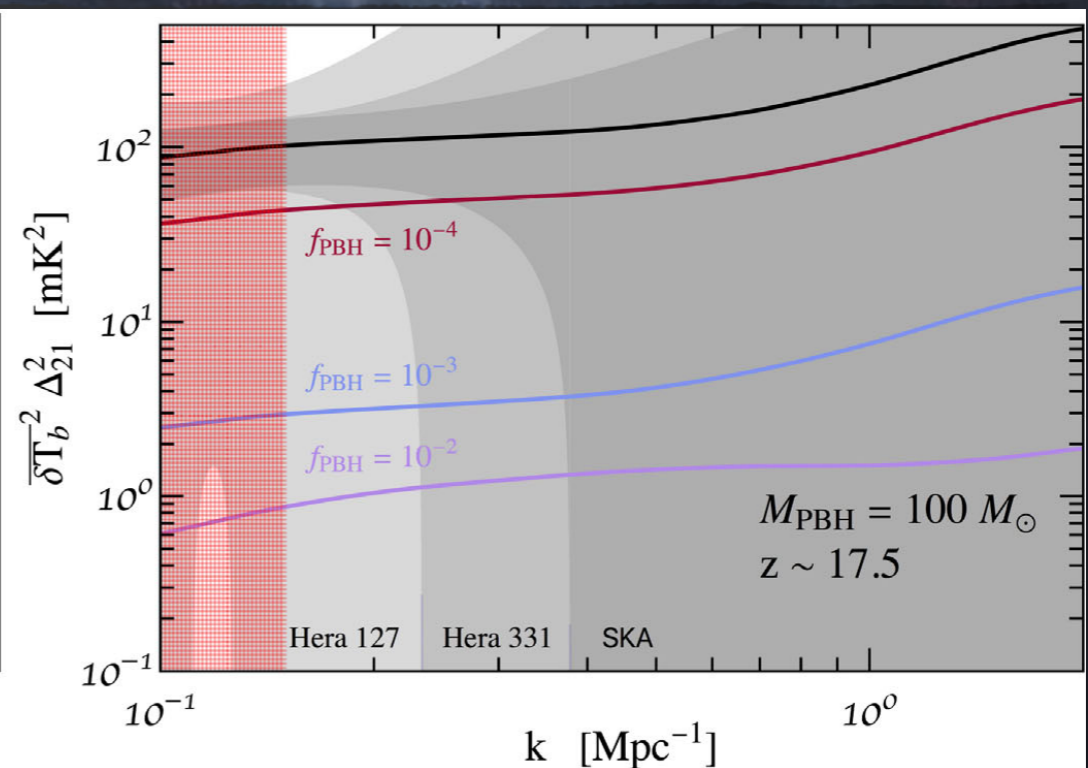
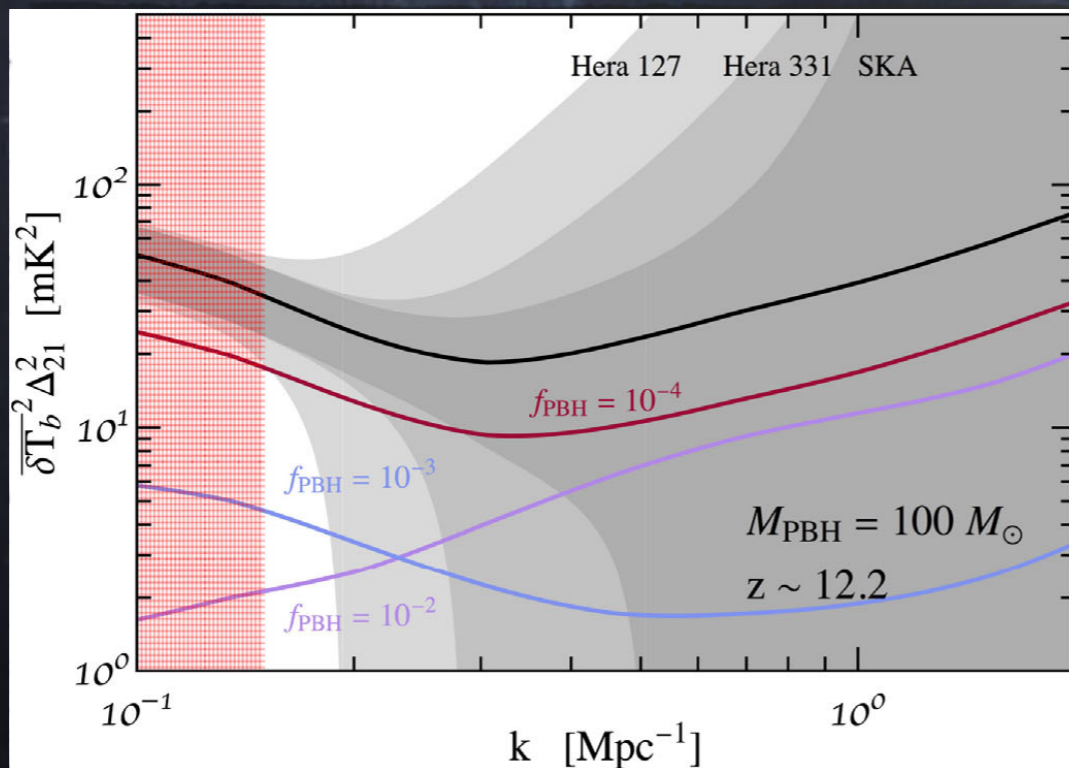
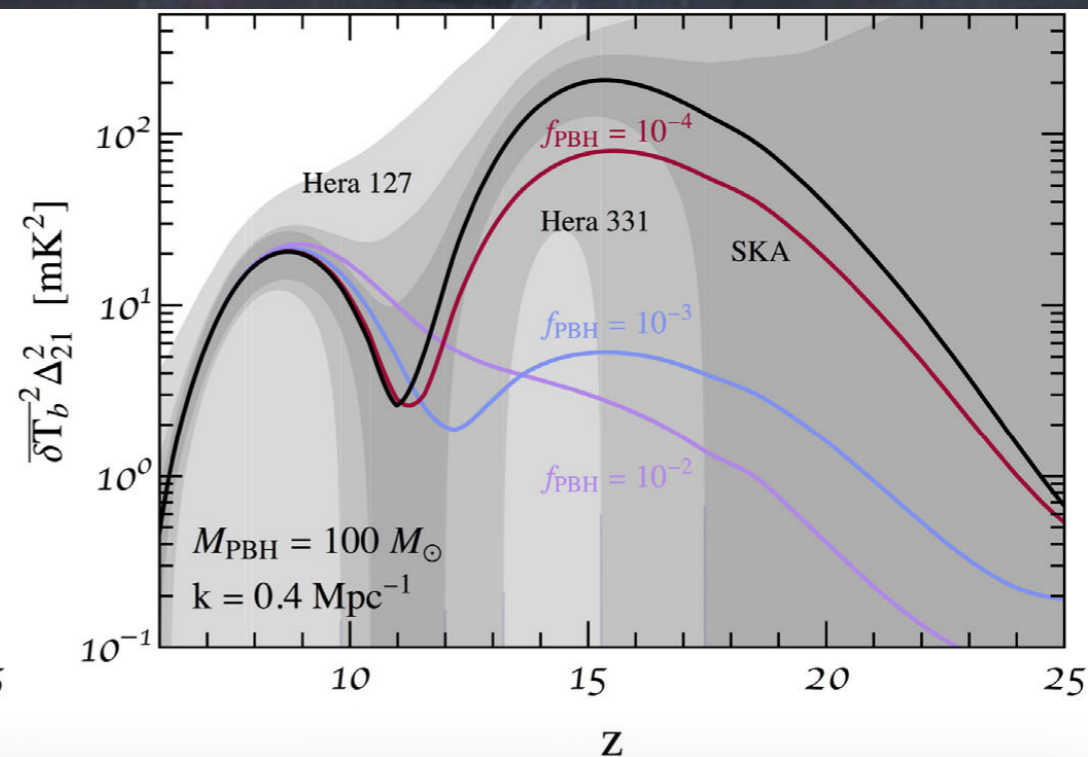
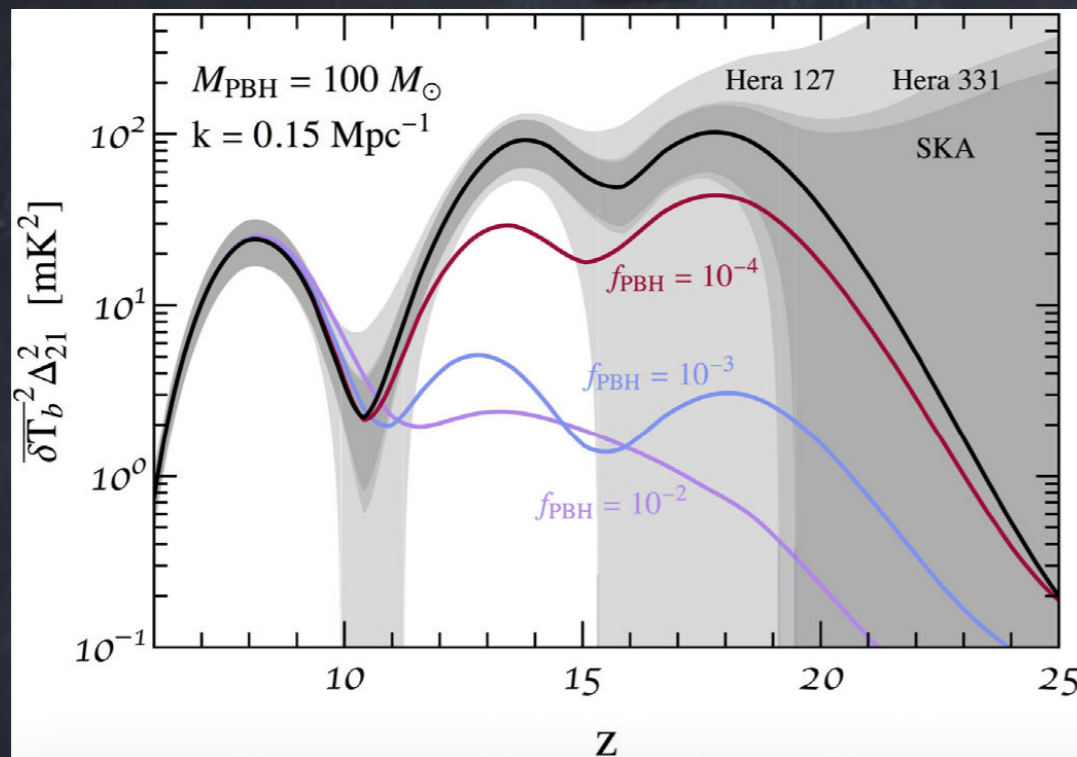
PBHs: 21CM POWER SPECTRUM

We use 21cmSense

J. C. Pober et al., *Astrophys. J.* 145:65, 2013

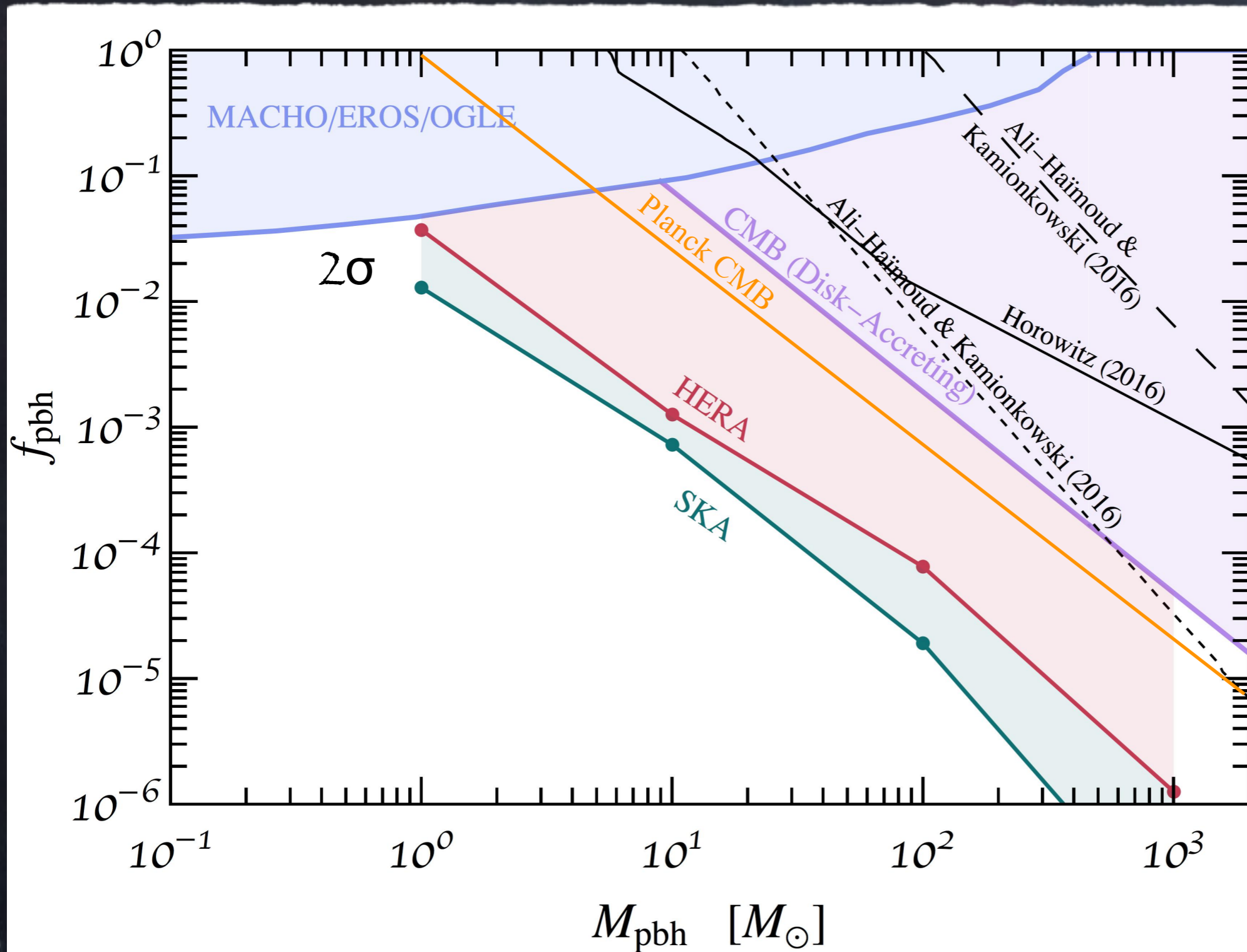
J. C. Pober et al., *Astrophys. J.* 782:66, 2014

Four-parameter astrophysical model



O. Mena, SPR, P. Villanueva-Domingo and S. J. Witte, *Phys. Rev. D* 100:043540, 2019

PBHs ABUNDANCE: SENSITIVITY



Y. Ali-Haimoud and
M. Kamionkowski,
Phys. Rev. D95:043534, 2017

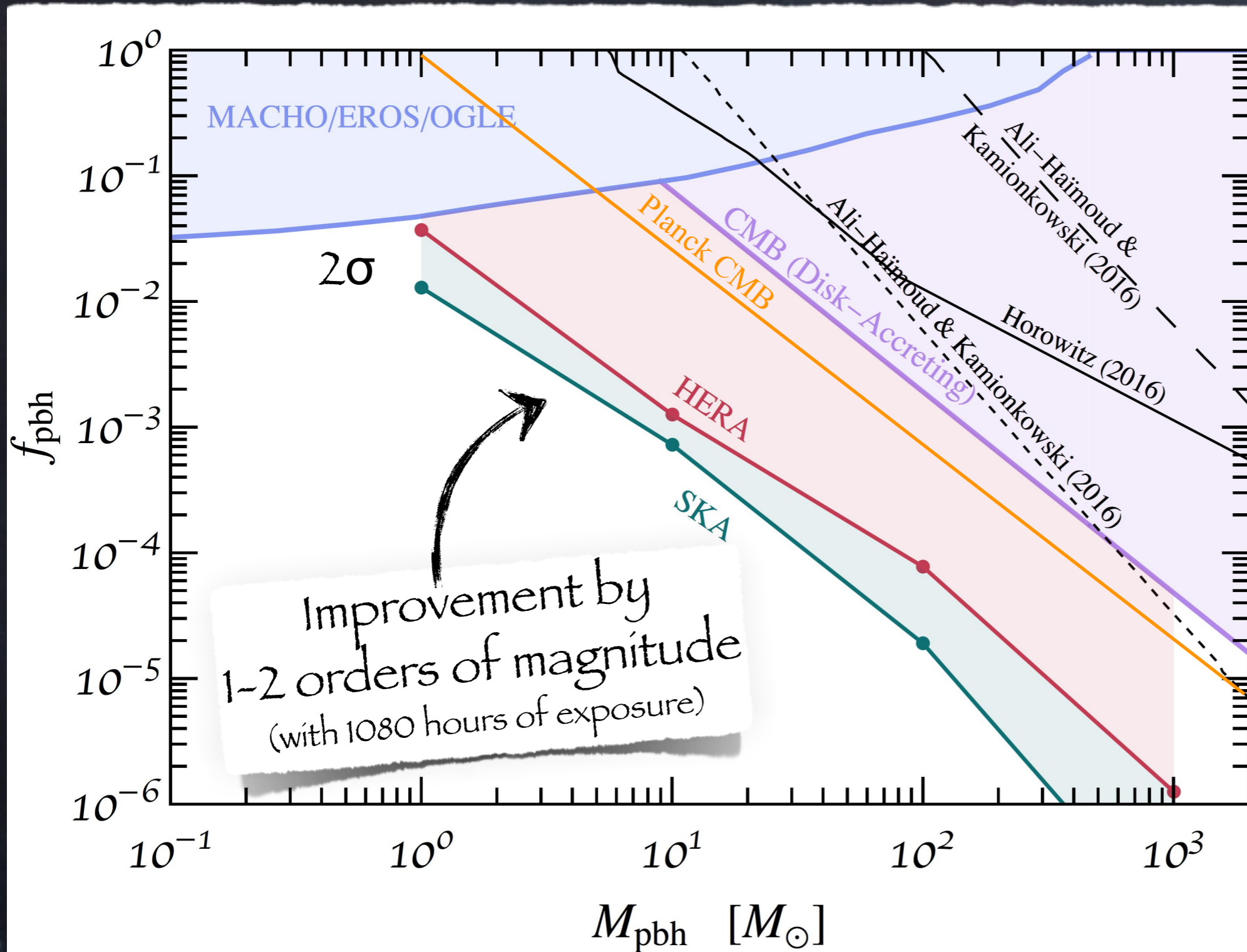
B. Horowitz,
arXiv: 1612.07264

V. Poulin et al.,
Phys. Rev. D96:083524, 2017

P. Serpico et al.,
Phys. Rev. Res. 2:023204, 2020

O. Mena, SPR, P. Villanueva-Domínguez and S. J. Witte, Phys. Rev. D100:043540, 2019

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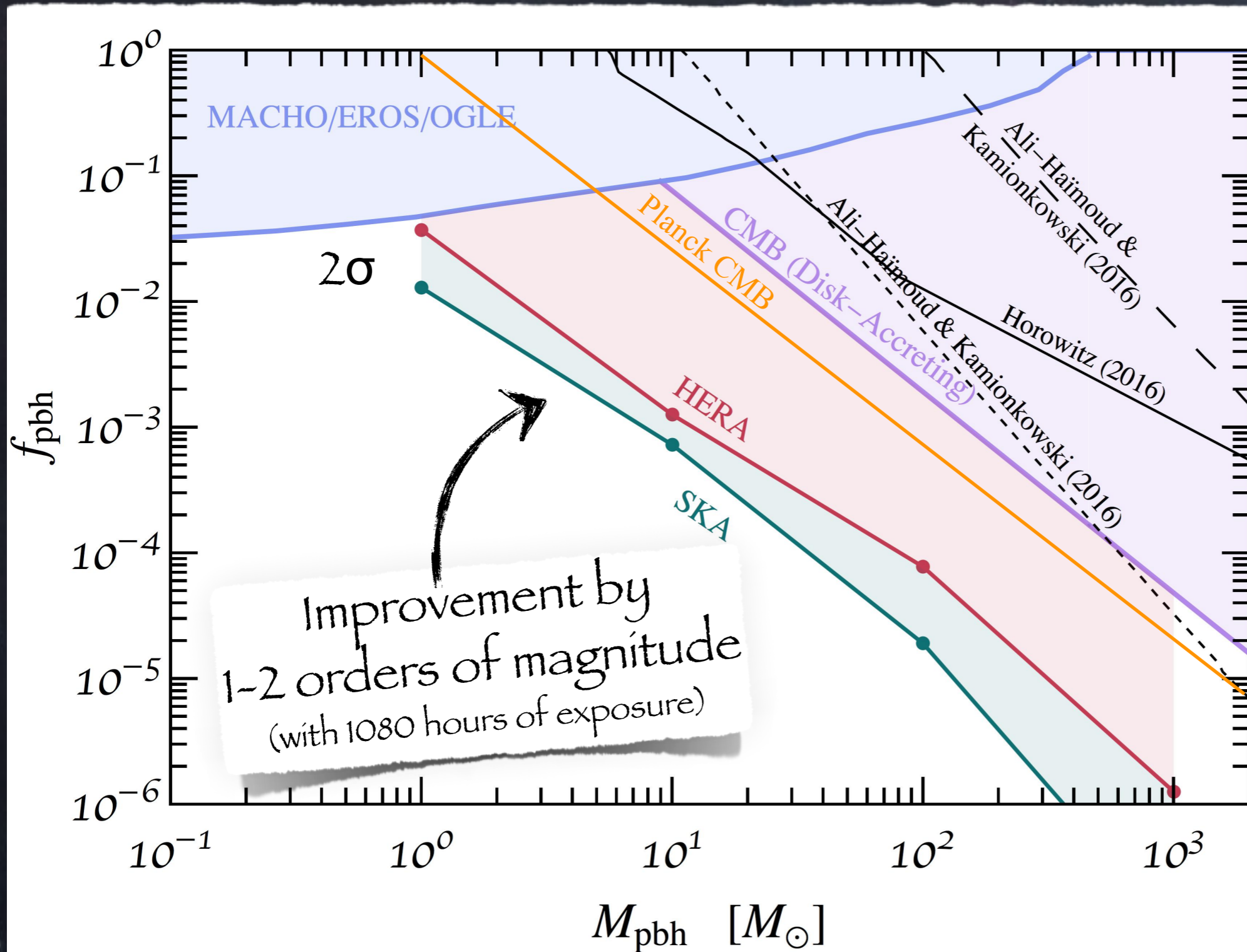
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P. Serpico et al.,
Phys. Rev. Res. 2:023204, 2020



Accretion model
is critical

O. Mena, SPR, P. Villanueva-Domínguez and S. J. Witte, Phys. Rev. D100:043540, 2019

SOME FINAL WORDS OF CAUTION

All constraints have caveats!

They all depend on the mass function!
Extended mass functions are, in general, more constrained

Uncertainties on estimates and observations can be very significant

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Just to name a few...

Evaporation

Memory burden : slowing down of evaporation

Lensing

Size of sources, wave optics, halo properties...

Gravitational Waves

Clustering and PBH binaries survival

Accretion

Accretion modeling

Dynamical

Assumed initial stellar distribution
Are stellar binaries genuine or spurious?