

THz Shaping of Electrons

Advancing Accelerators with Light

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CFEL, DESY

June 12th, 2023

Why explore higher frequencies?

Higher driving frequency allows:

- higher acceleration fields: $|\vec{E}|_{breakdown} \sim \frac{1}{\tau^{1/6}}$ → smaller accelerator

[1] Kilpatrick, W. D., Rev. Sci. Inst. 28, 824 (1957).

[2] Loew, G.A., et al., 13th Int. Symp. on Discharges and Electr. Insulation in Vacuum, Paris, France. 1988.

[3] S. V. Dolgashev, et al. Appl. Phys. Lett. 97, 171501 2010.

[4] M. D. Forno, et al. PRAB. 19, 011301 (2016).

- lower pulse energy for same E-field in the cavity:

$$U_{Pulse} \sim |\vec{E}|^2 V_{cavity} \sim \lambda^3 \sim f^{-3}$$

- reduced pulsed heating: $\Delta T \sim \frac{U_{Pulse}}{A_{Surface}} \sim \lambda \sim f^{-1}$ → higher rep rates!

- higher field gradients:

- stronger compression: $\frac{dF}{dz} \sim \frac{|\vec{E}|}{\lambda} \sim f^2$ → shorter bunches
- quicker acceleration → reduce space charge effects

Wavelengths

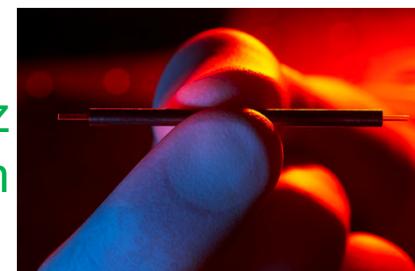
RF Accelerator

100 mm



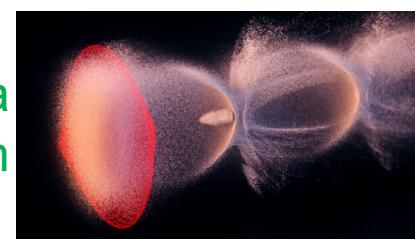
THz

1 mm



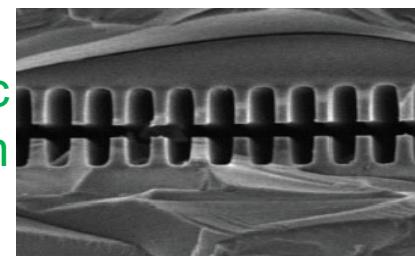
Laser-Plasma

0.01 mm



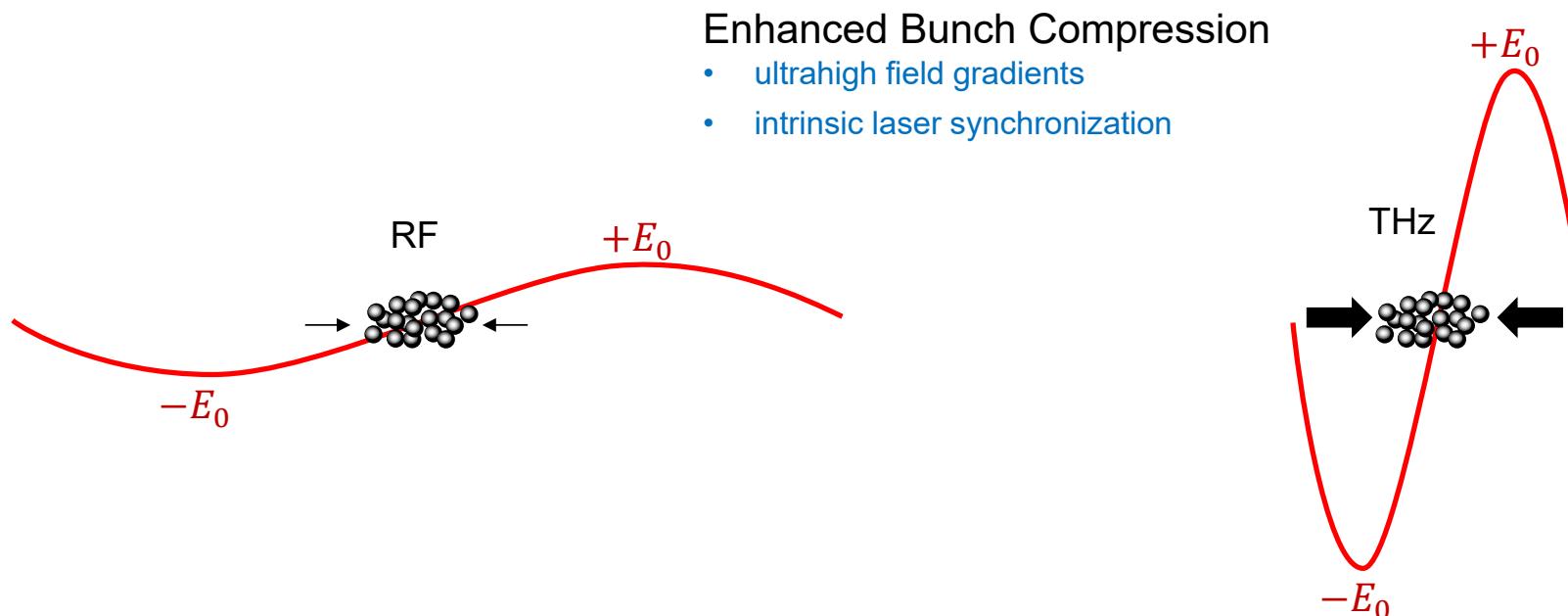
Laser-Dielectric

0.001 mm

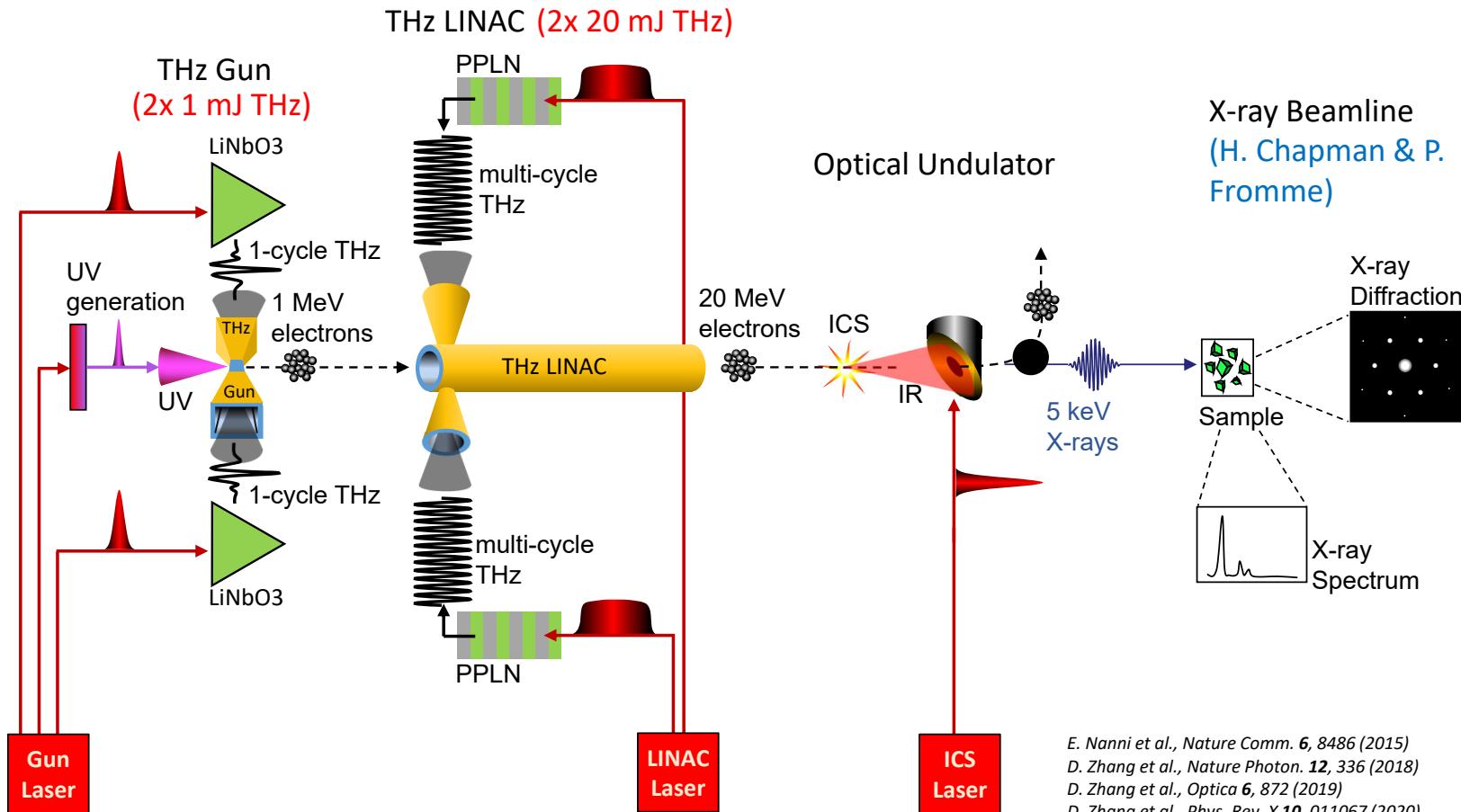


THz Benefits: shorter waves, higher fields, stronger gradients

Parameter	RF Accelerators	THz Accelerators	THz Enhancement
Field Strengths, E_0	10 – 100 MV/m	100 – 1000 MV/m	~10x
Wavelength, λ	3 – 10 GHz	100 – 500 GHz	~100x
Field Gradients, dE_0/dz	~10 GV/m ²	~10,000 GV/m ²	~1000x



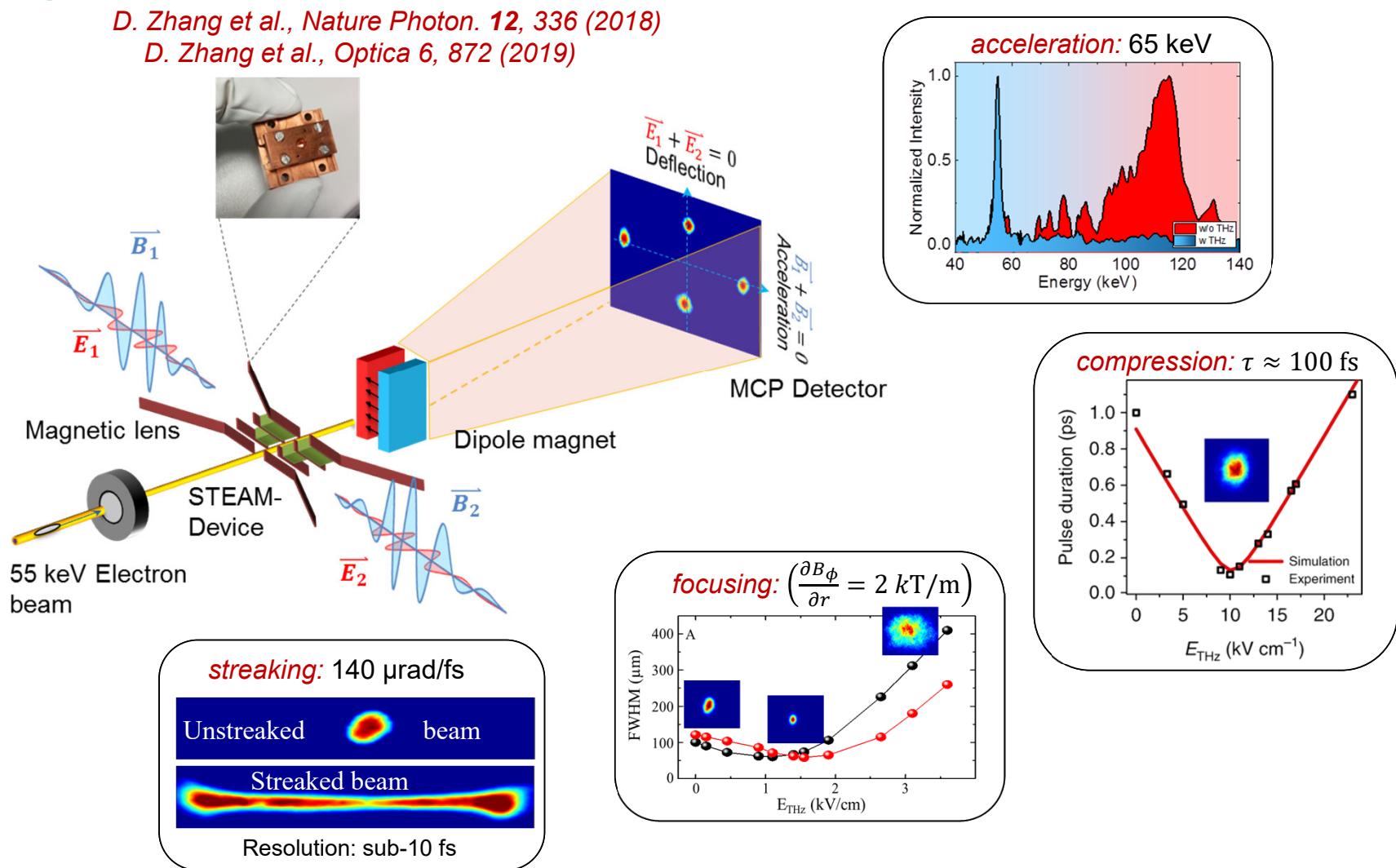
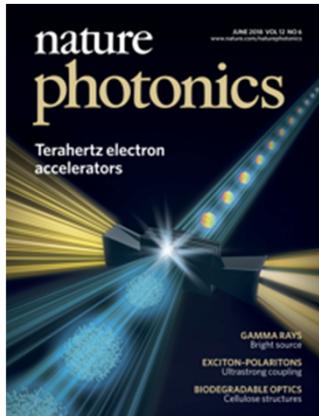
THz-powered electron and X-ray light source: AXIS



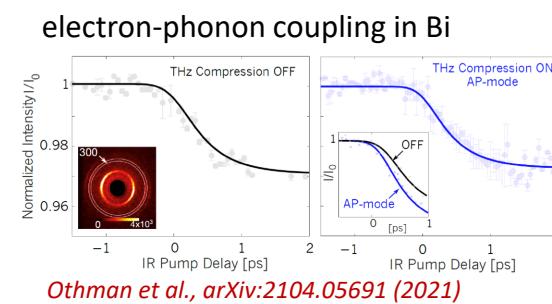
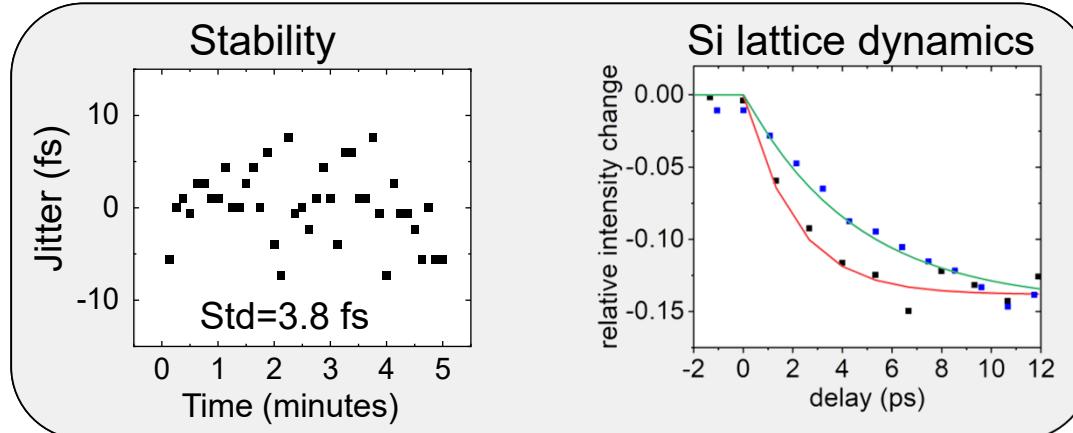
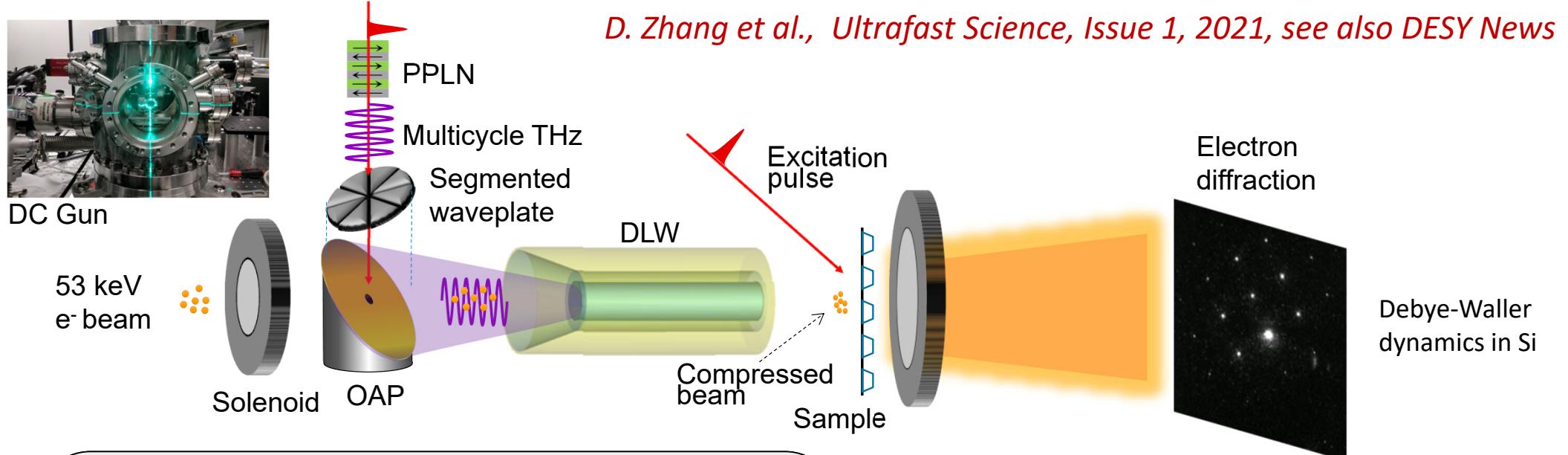
F.X. Kärtner et al., *NIMPRA* **829**, 24 (2016)

STEAM device yields practical beam acceleration & manipulation

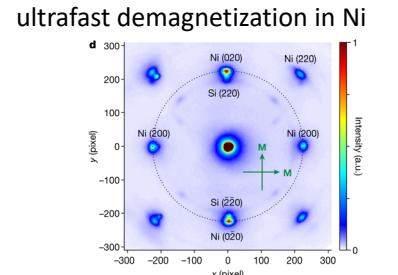
Segmented
THz
Electron
Accelerator &
Manipulator



THz-compression enhances UED temporal resolution



Othman et al., arXiv:2104.05691 (2021)

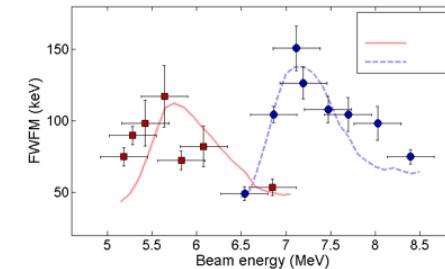
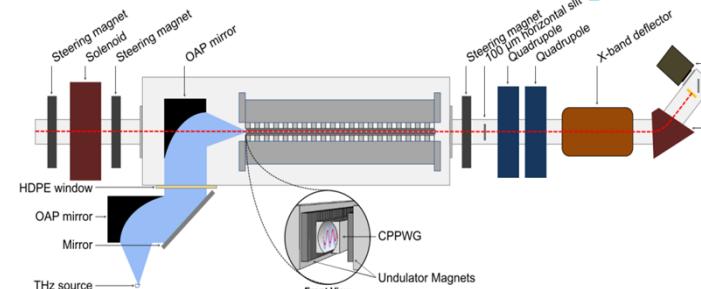


Tauchert et al., Nature 602, 73 (2022)

THz acceleration and beam manipulation takes up speed!

THz-powered iFEL

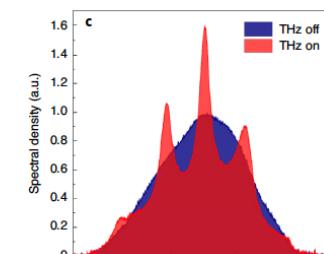
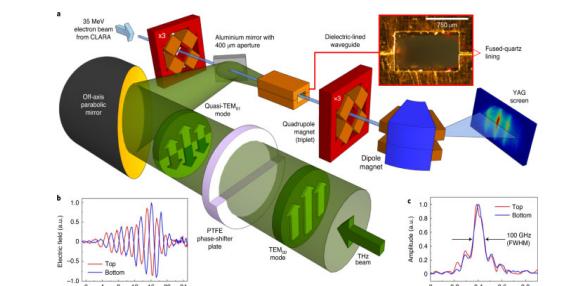
E. Curry et al., PRL 120, 094801 (2018)



$$\begin{aligned} E_{in} &= 5 \text{ MeV} \\ \Delta E &= \pm 75 \text{ keV} \\ E_{THz} &= 1 \mu\text{J} \end{aligned}$$

Acceleration

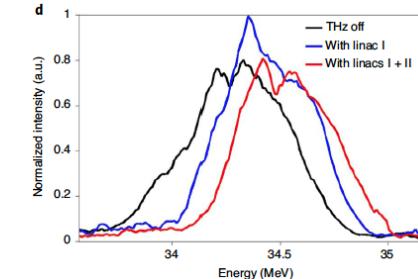
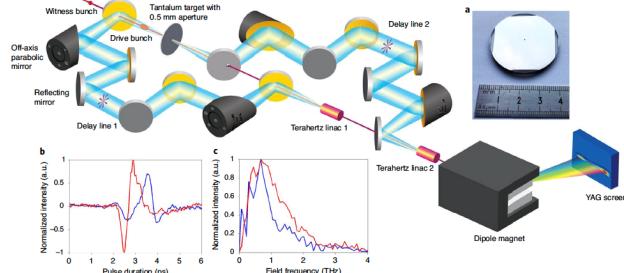
M. Hibberd et al., Nat. Photon. (2020)



$$\begin{aligned} E_{in} &= 35 \text{ MeV} \\ \Delta E &= \pm 10 \text{ keV} \\ E_{THz} &= 2 \mu\text{J} \end{aligned}$$

Acceleration & Staging

H. Xu et al., Nat. Photon. (2021)



$$\begin{aligned} E_{in} &= 34.3 \text{ MeV} \\ \Delta E &= 150 \text{ keV} \\ E_{THz} &= 132 \mu\text{J} \end{aligned}$$

Others:

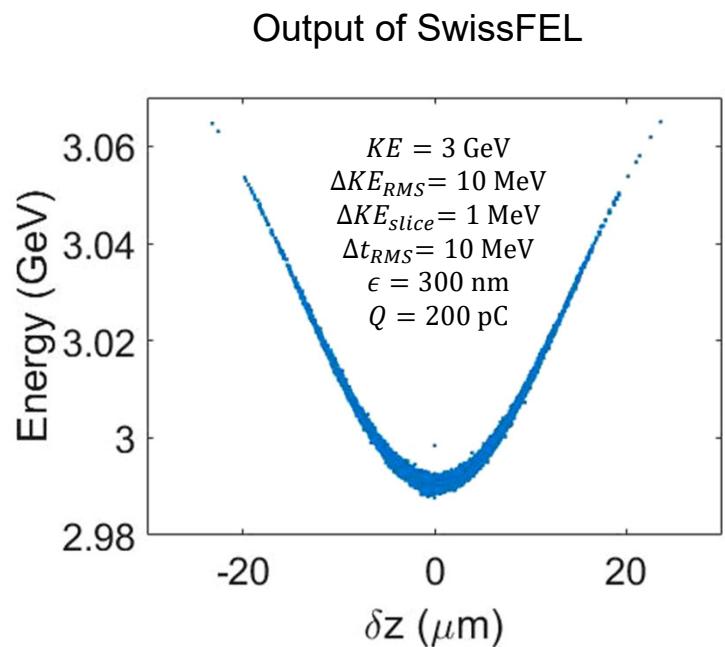
DESY.

Nicholas H. Matlis - WP13 EuPRAXIA electron and photon diagnostics - 2023.06.12

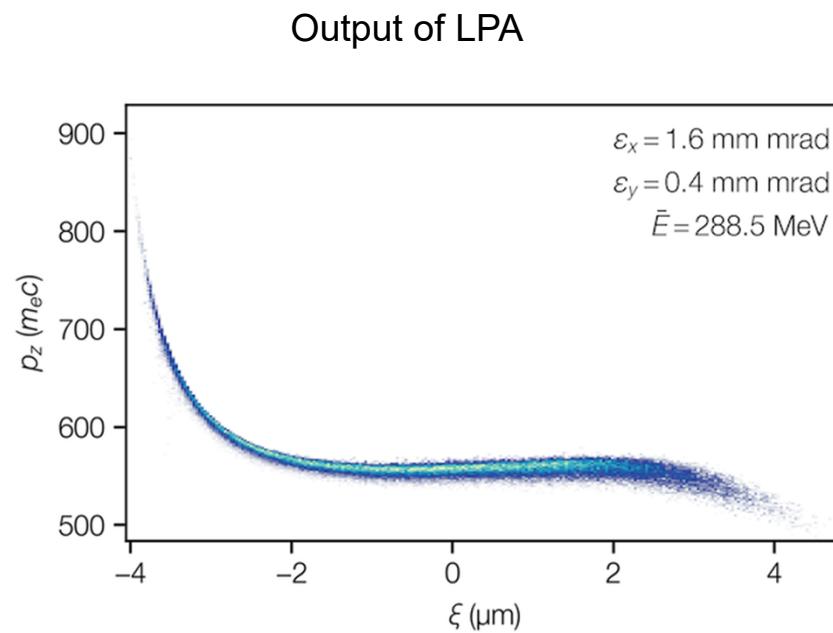
- *Kealhofer et al., Science 359, 459 (2016)*
- *Walsh et al., Nature Comm. 8, 421 (2017)*
- *Zhao et al., PRX 8, 021061 (2018)*
- *Li et al., Phys. Rev. Accel. Beams 22, 012803 (2019)*

Applications of THz for LPAs and Conventional Accelerators

- THz Streaking
- THz Compression
- THz Focusing
- THz Beam Shaping



Discussion, R. Ischebeck et al., SwissFEL, PSI

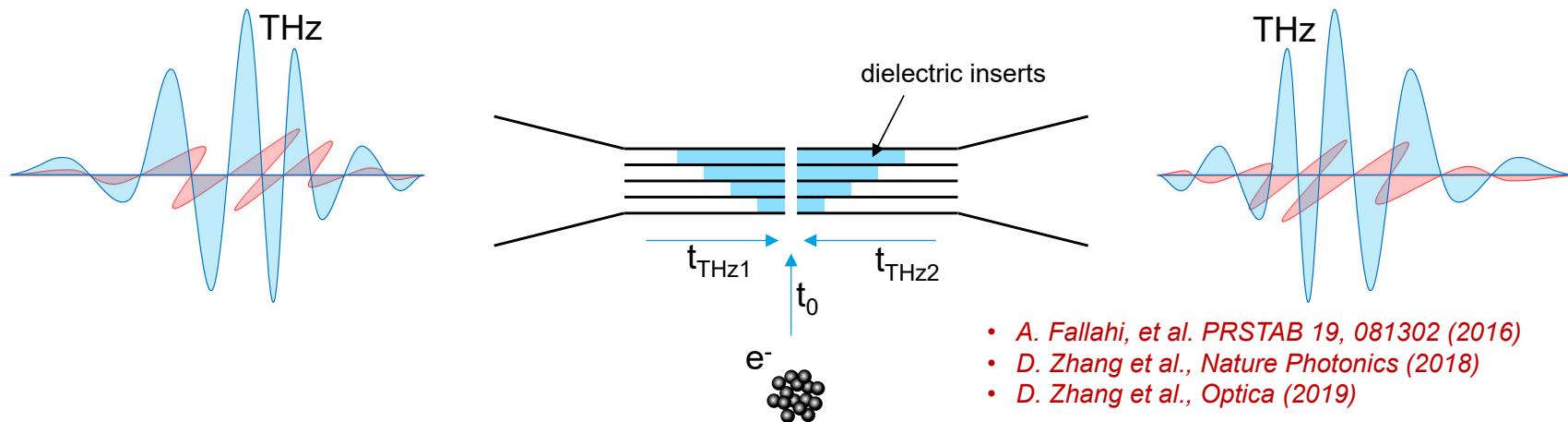


M. Kirchen et al., Phys. Rev. Lett. **126**, 174801 (2021)

2-geometries for THz-electron interaction

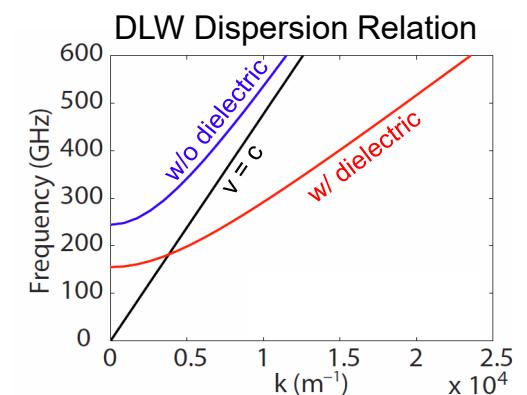
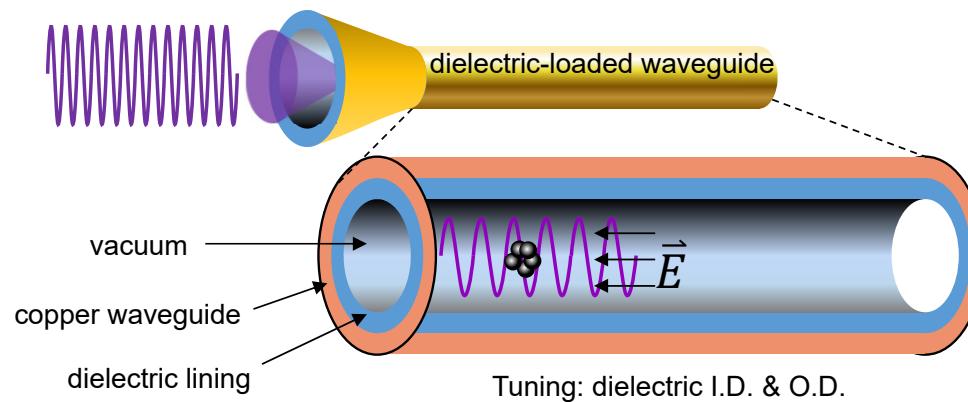
Transverse Geometry

- Single (few)-cycle THz
- Tuning: dielectric L, h



Co-propagating Geometry

- Narrowband (multicycle) THz
- Tuning: dielectric $R_{I.D.}, R_{O.D.}$



L.J. Wong et al., Opt. Exp. 21, 9792 (2013)

Benefits and Challenges for Applying THz to High-Energy Beams

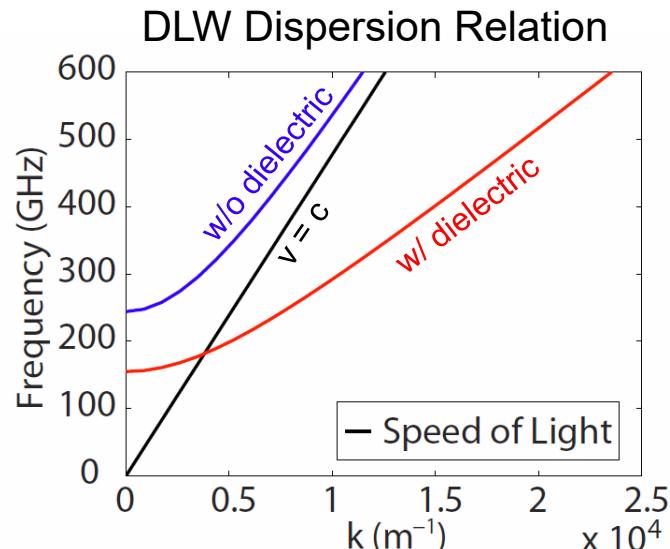
Challenges

- High-energy requires lots of “umpf”

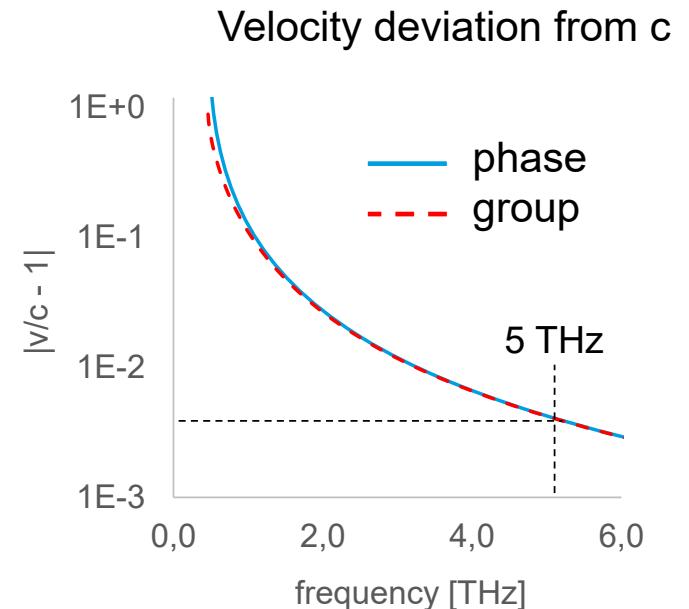
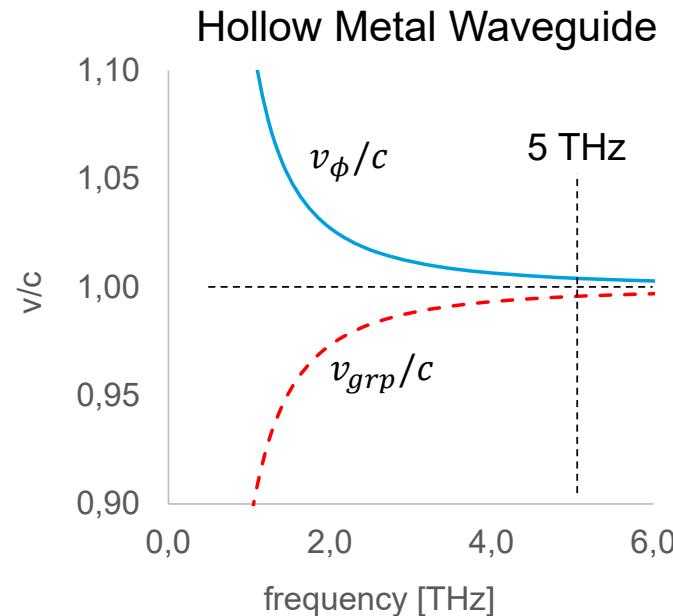
Benefits

- Beams are moving really close to c
- Beam velocity doesn't change with energy
- Long interaction lengths possible
- THz pulse durations can be short
- Waveguide dispersion is minimized

Near luminal operation brings key advantages

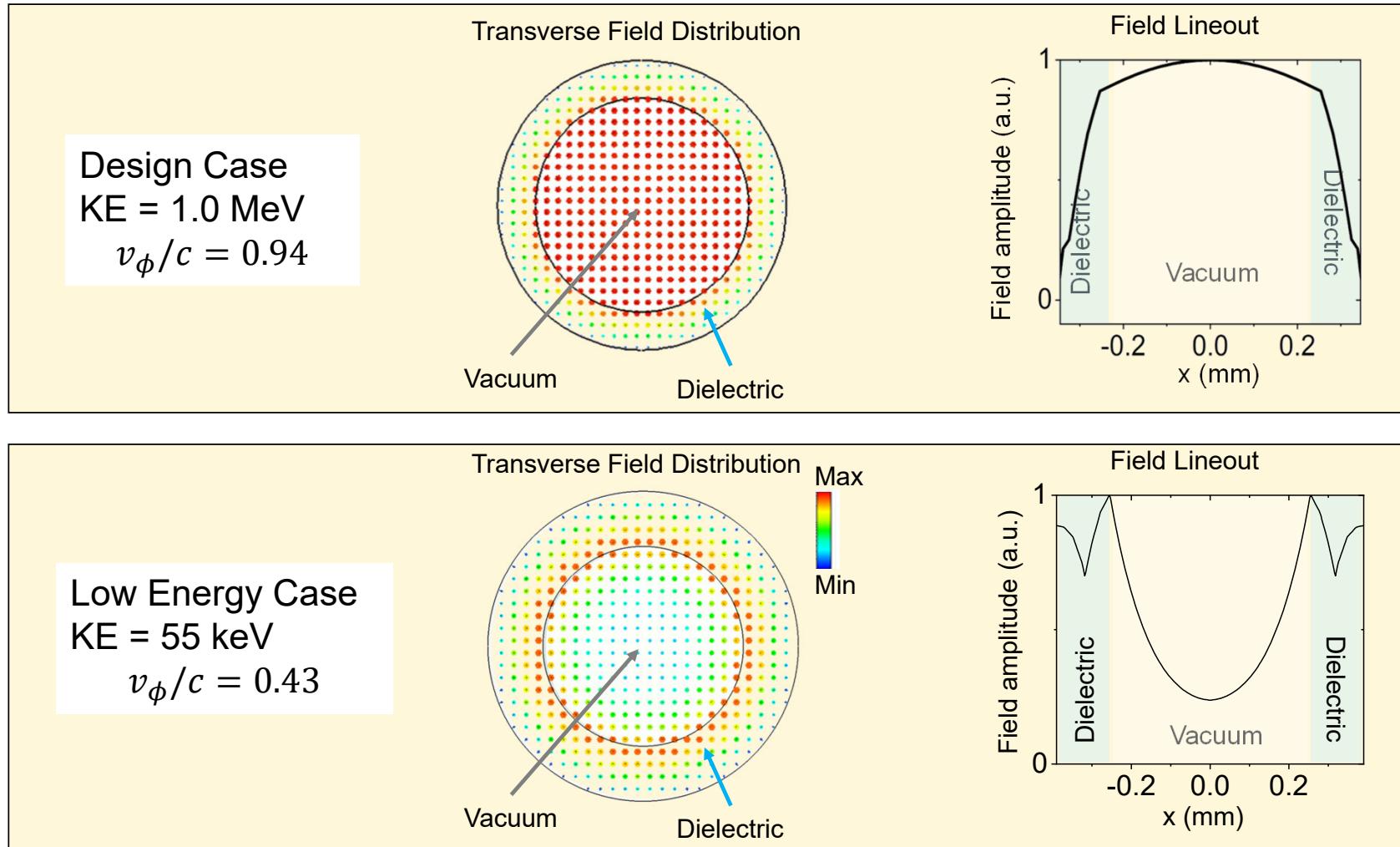


L.J. Wong et al., Opt. Exp. 21, 9792 (2013)

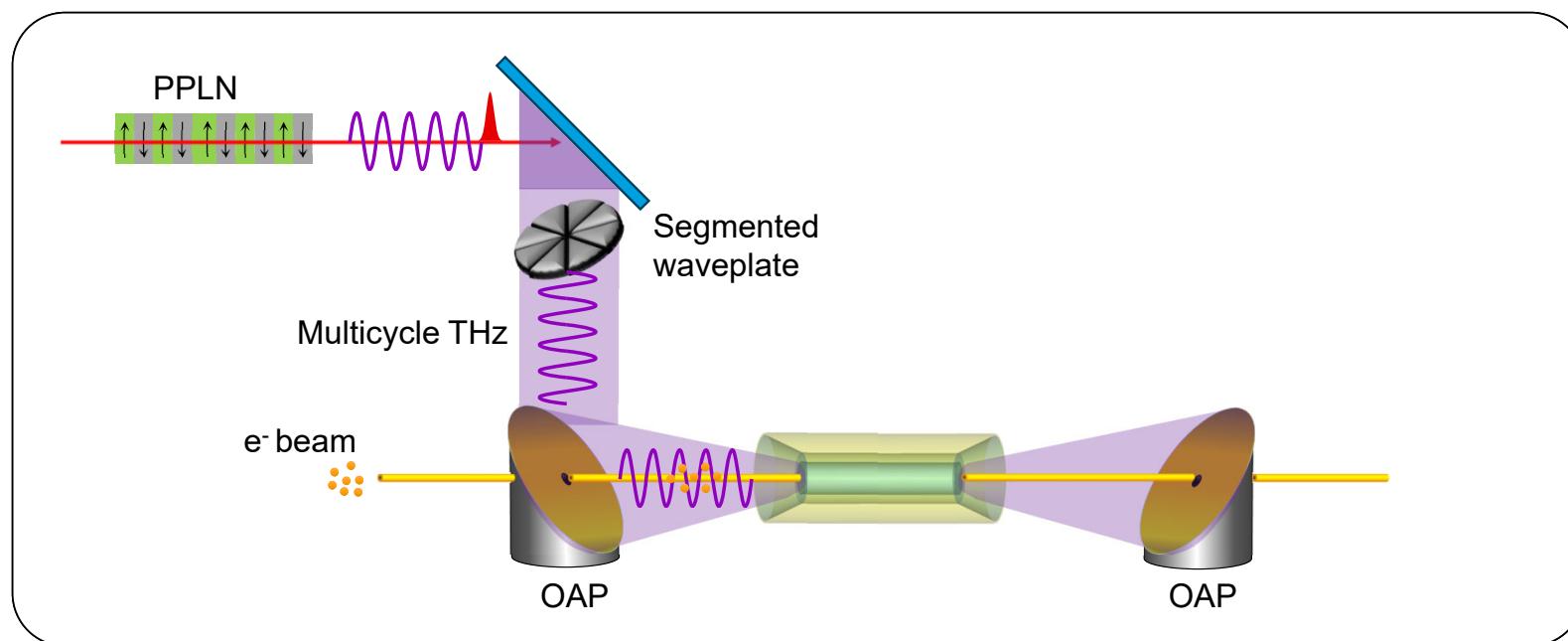
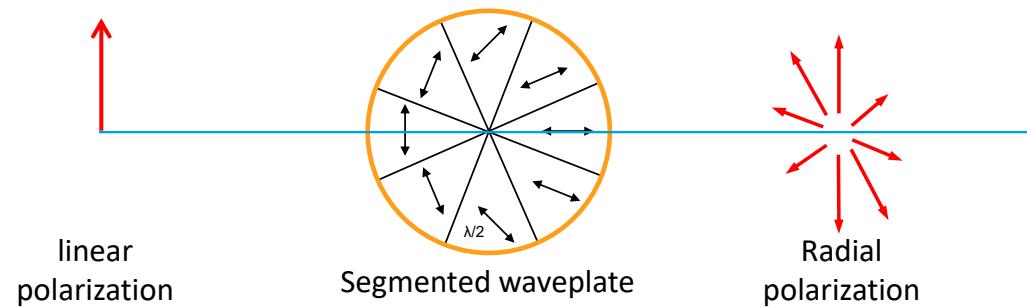


- THz dispersion is very low → can use short pulses
- THz pulse walk-off is very low → can use short pulses

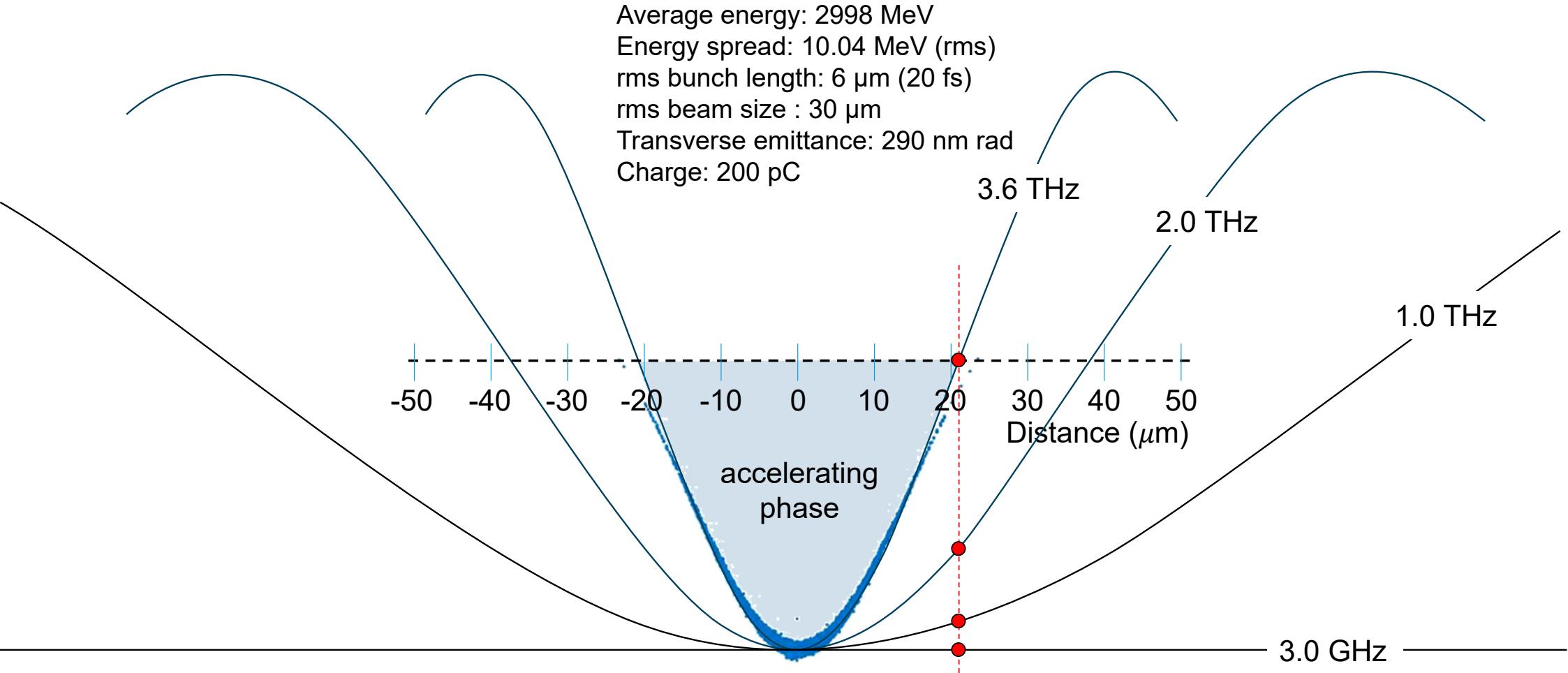
Benefits of high electron energy: uniform field profile



Acceleration mode TM_{01} in DLW needs radially polarized input



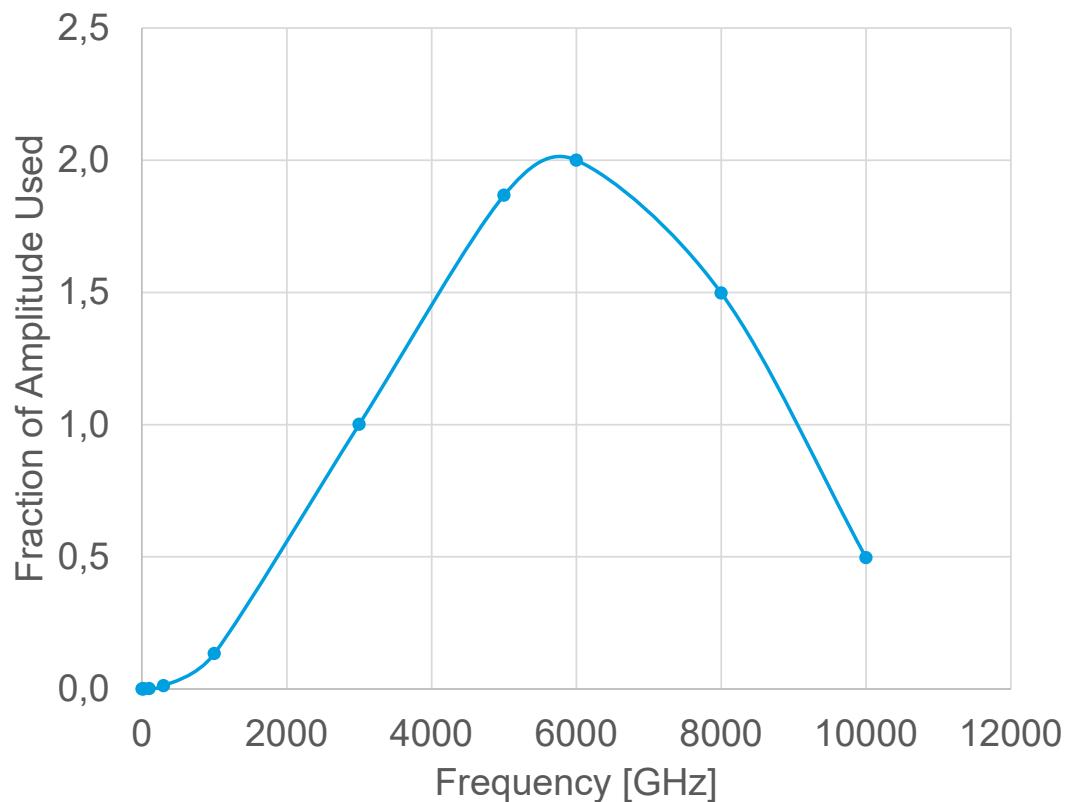
Tuning the frequency to the length-scale of the feature



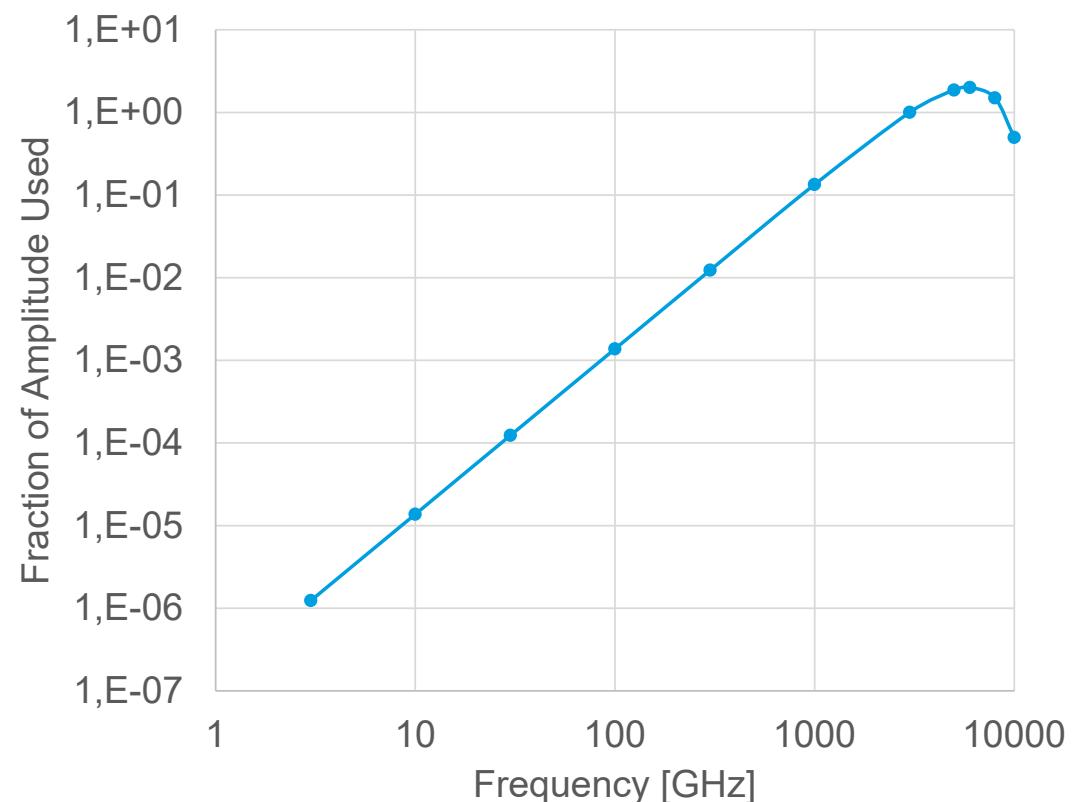
Quantifying the effectiveness vs frequency

RF is ~1 million times less effective than THz

Corrective Effectiveness for 50um Bunch

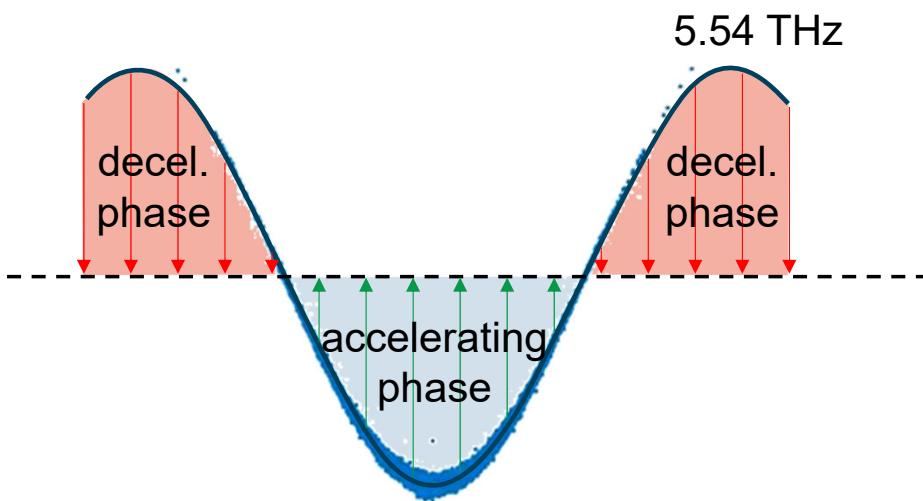


Corrective Effectiveness for 50um Bunch



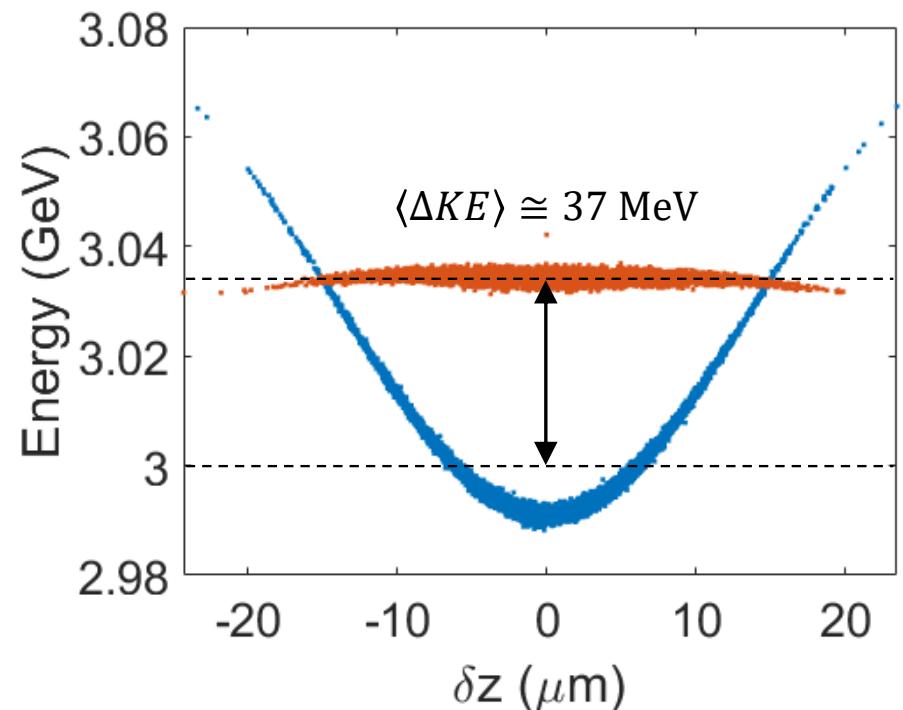
Optimized frequency minimizes average energy gain

Matched frequency uses full peak to peak amplitude of THz



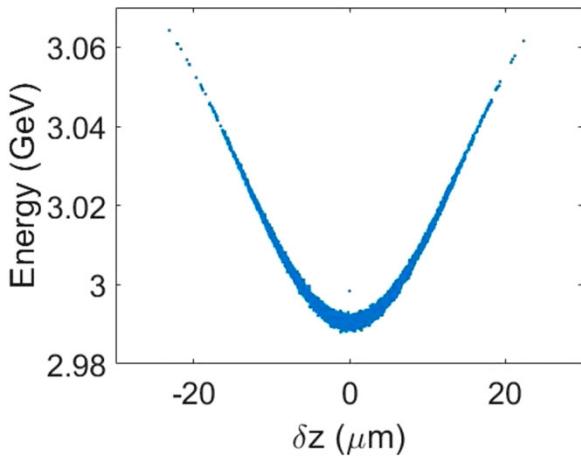
Energy increase of beam:

- $Q = 200 \text{ pC}$
- $\langle \Delta KE \rangle \cong 37 \text{ MeV}$
- $\delta E = 7.4 \text{ mJ}$

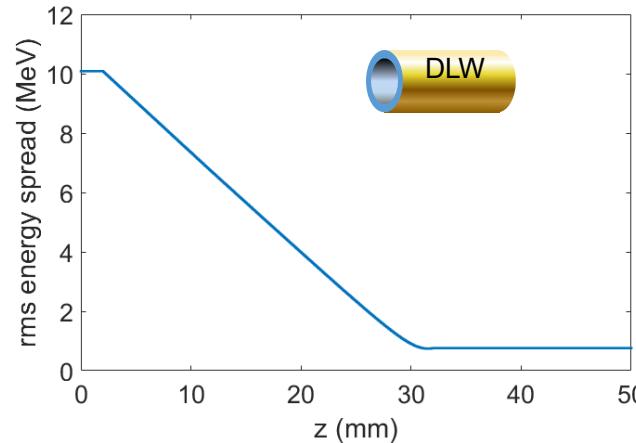


Near luminal v_ϕ and v_{grp} allow long interaction lengths

$$E_{THz} = 1.5 \text{ GV/m}$$



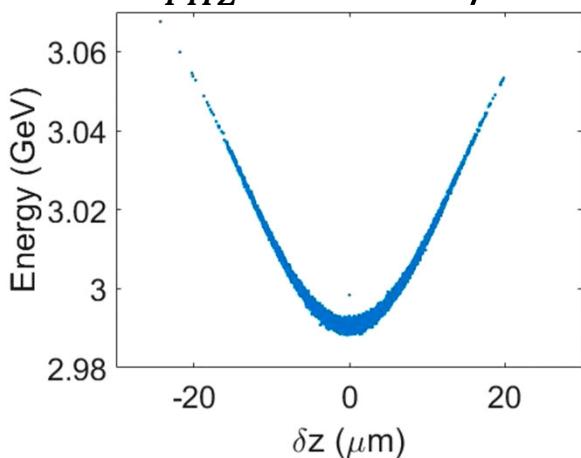
$$L = 3 \text{ cm}$$



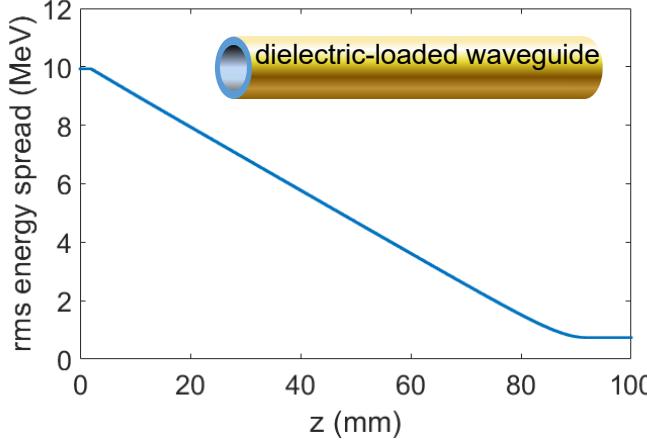
Parameters

- $d_{waveguide} = 250 \mu\text{m}$
- $v_\phi \cong c$
- $\phi_{inj} = 135^\circ$
- $\langle \Delta KE \rangle_{RMS} = 0.8 \text{ MeV}$

$$E_{THz} = 0.5 \text{ GV/m}$$



$$L = 9 \text{ cm}$$



Parameters

- $d_{waveguide} = 250 \mu\text{m}$
- $v_\phi \cong c$
- $\phi_{inj} = 135^\circ$
- $\langle \Delta KE \rangle_{RMS} = 0.8 \text{ MeV}$

$$E_{pulse} = 1 \text{ mJ}$$

$$u = \epsilon_0 |E_{field}|^2$$

$$V = \frac{1}{4} \pi d^2 N \Lambda_{THz}$$

$$E_{pulse} = uV$$

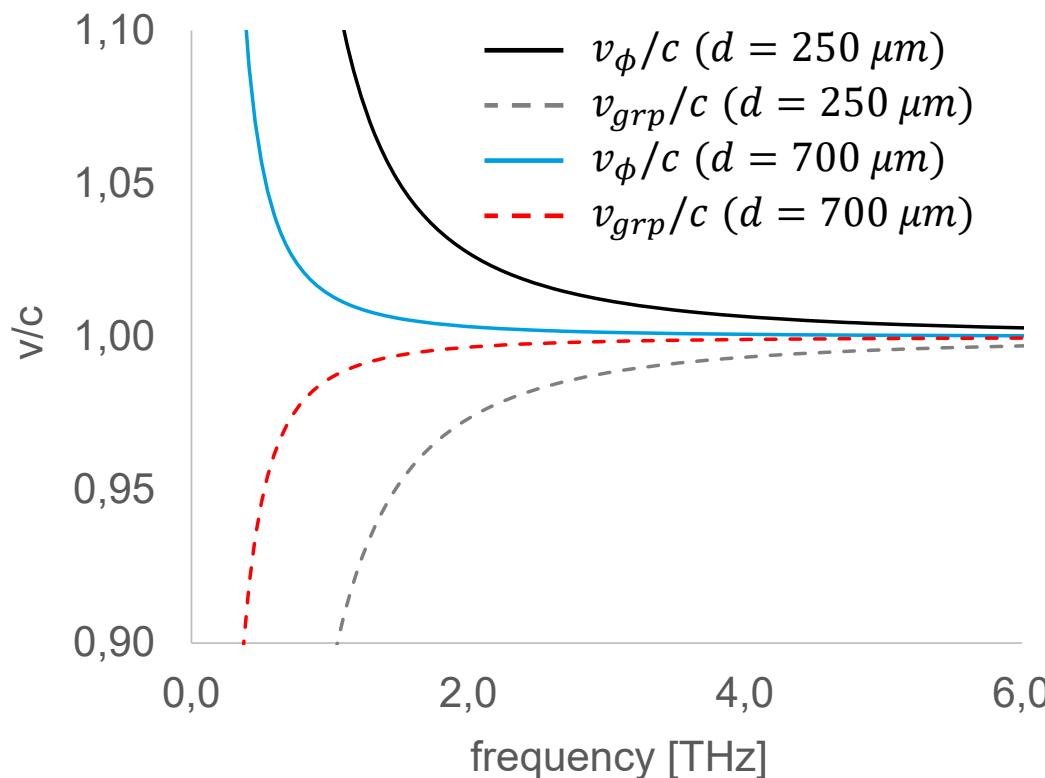
$$E_{pulse} = 100 \mu\text{J}$$

Phase and group velocities approach zero at higher frequency

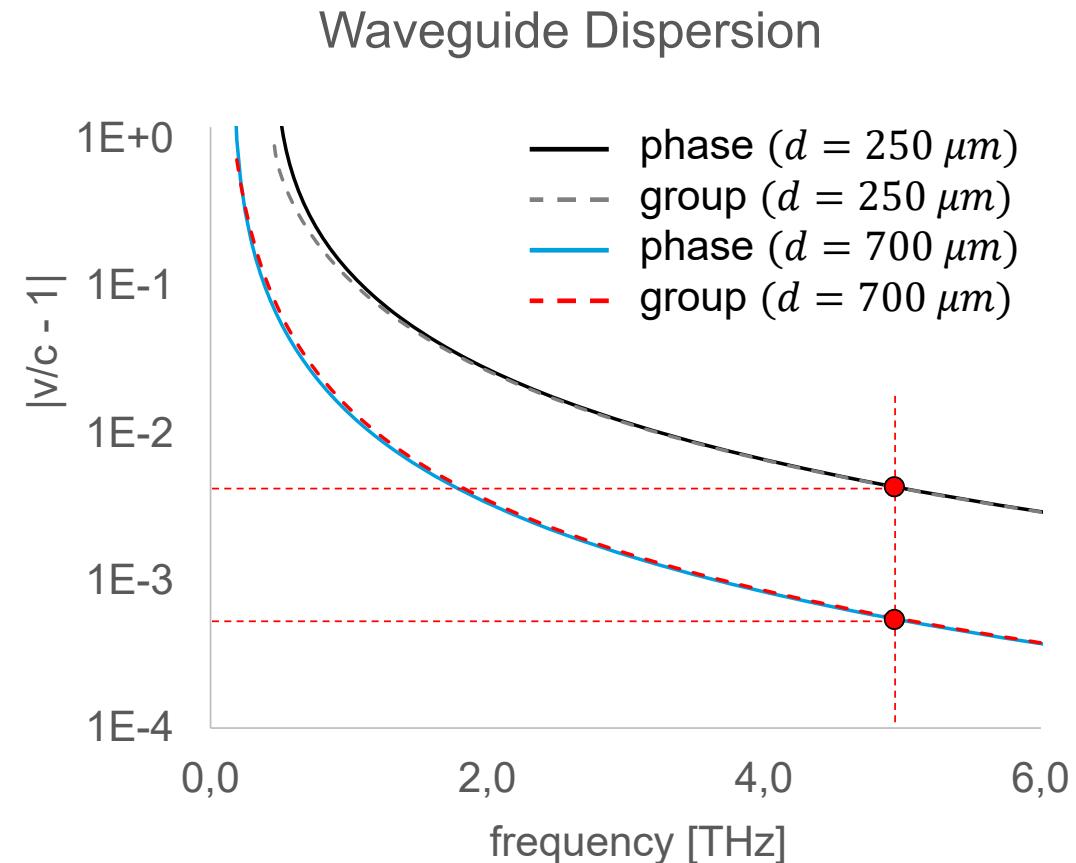
$$\frac{v_\phi}{c} = \frac{f}{\sqrt{f^2 - f_c^2}}$$



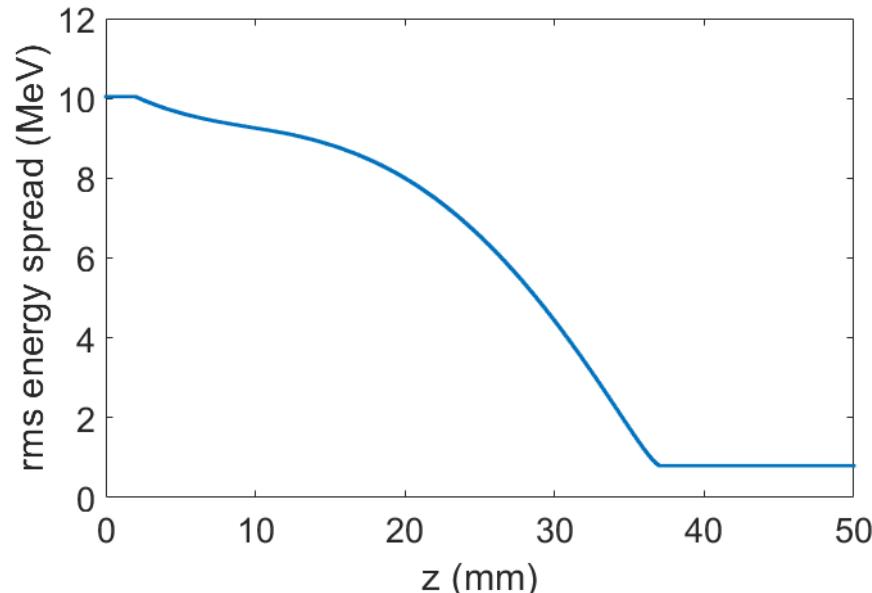
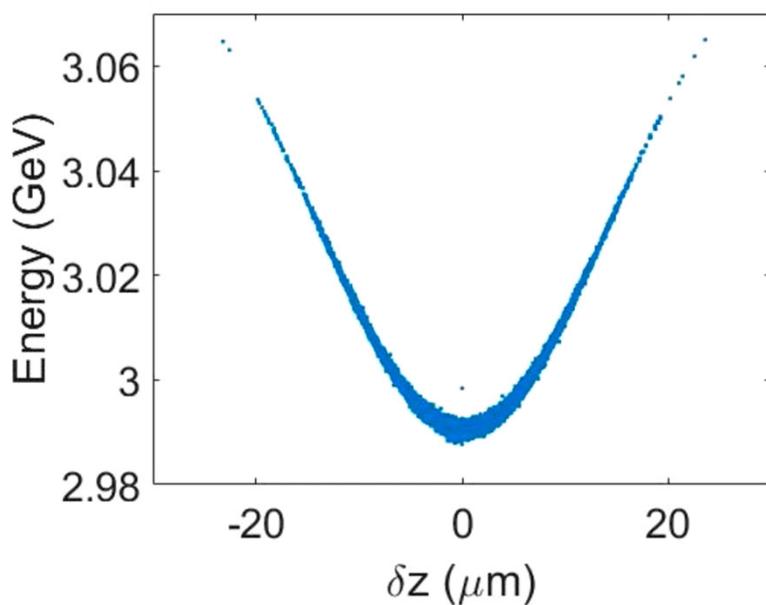
Waveguide Dispersion



Waveguide Dispersion



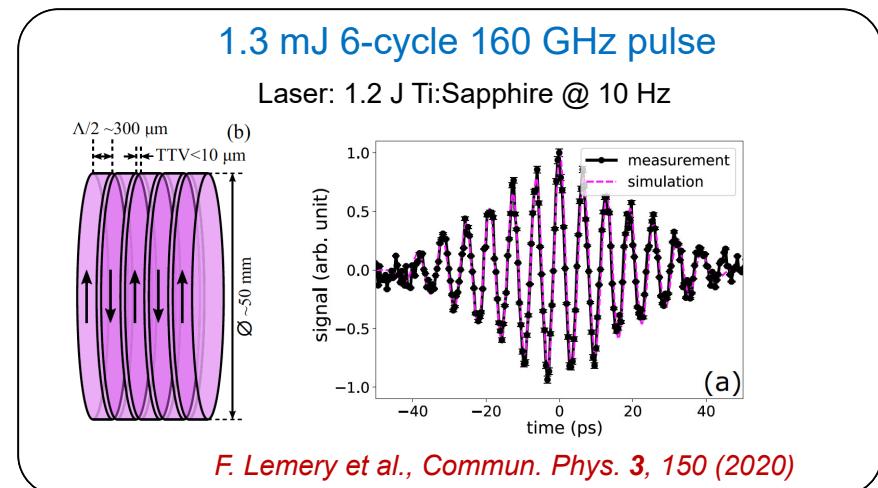
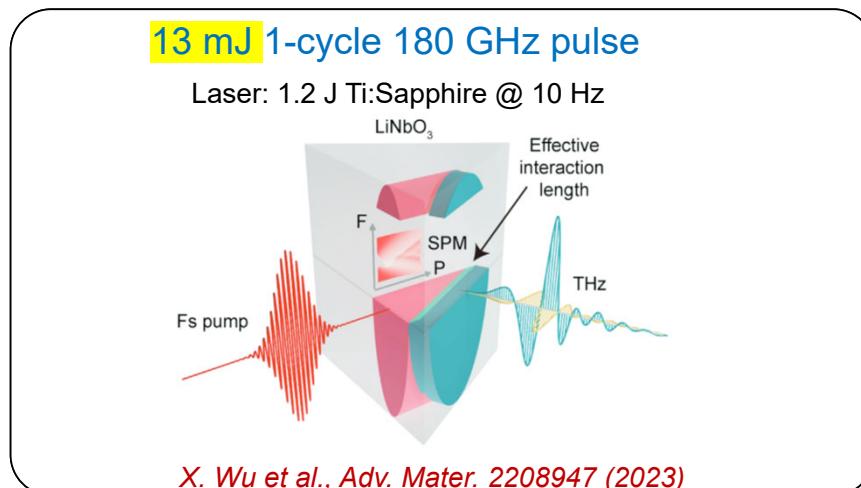
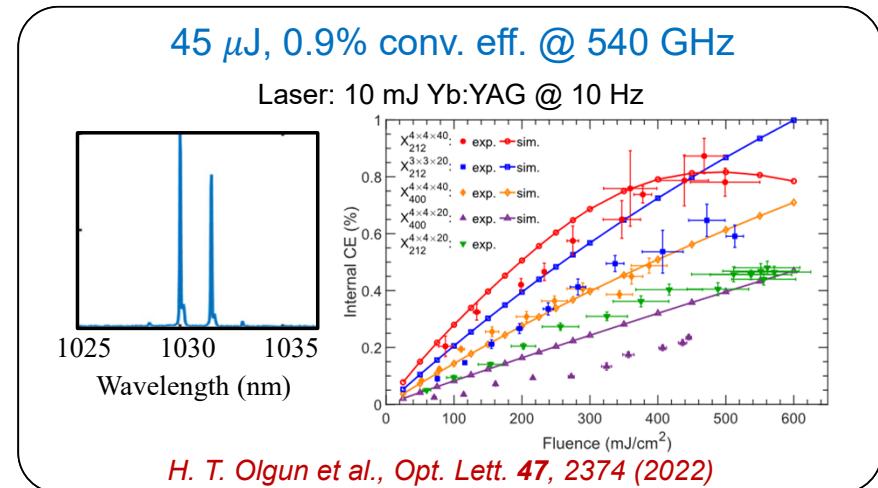
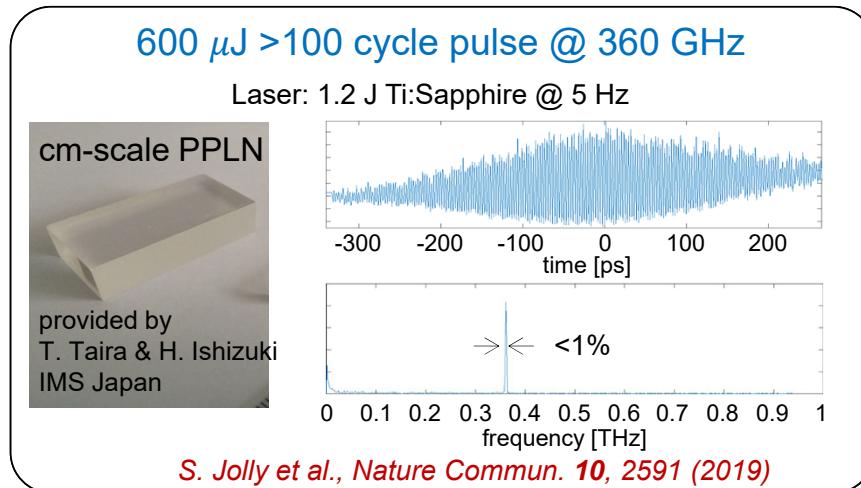
Larger waveguides may work without dielectric



dielectric-loaded waveguide

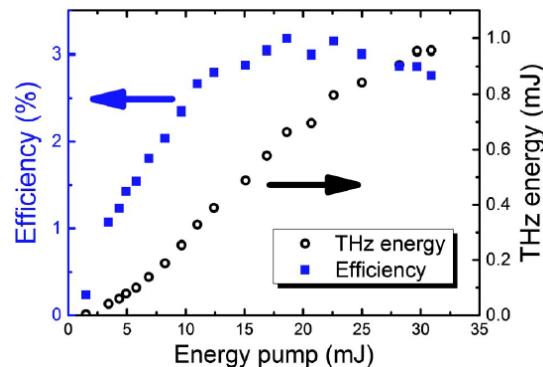
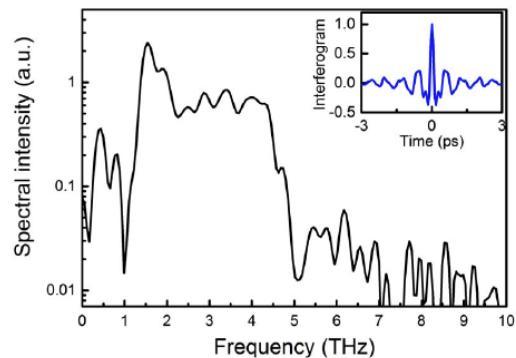
- Parameters
- $d_{\text{waveguide}} = 700 \mu\text{m}$
 - $E_{\text{field}} = 1.5 \text{ GV/m}$
 - $E_{\text{pulse}} \approx 10 \text{ mJ}$
 - $v_\phi \cong 1.0006 c$
 - $\phi_{\text{inj}} = 85^\circ$
 - $L = 37 \text{ mm}$
 - $\langle \Delta KE \rangle_{\text{RMS}} = 0.8 \text{ MeV}$

Multi-mJ laser-based THz sources now exist



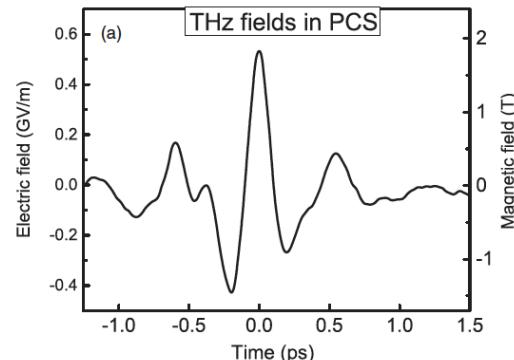
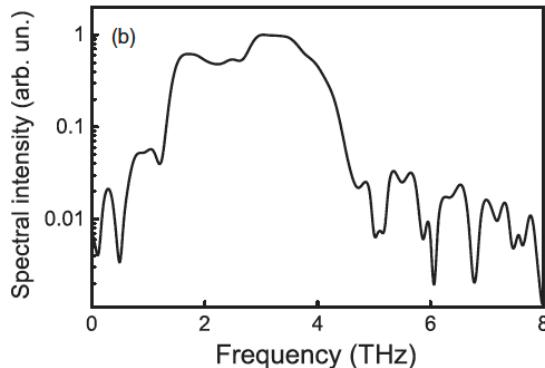
Organic crystals promising for 5 THz, GV/m fields

wavelength difference for 5 THz generation: 15 nm



- Laser: Cr:forsterite, 1250 nm, 30 mJ, 10 Hz
- Bandwidth: 20 nm
- Crystal: DSTMS

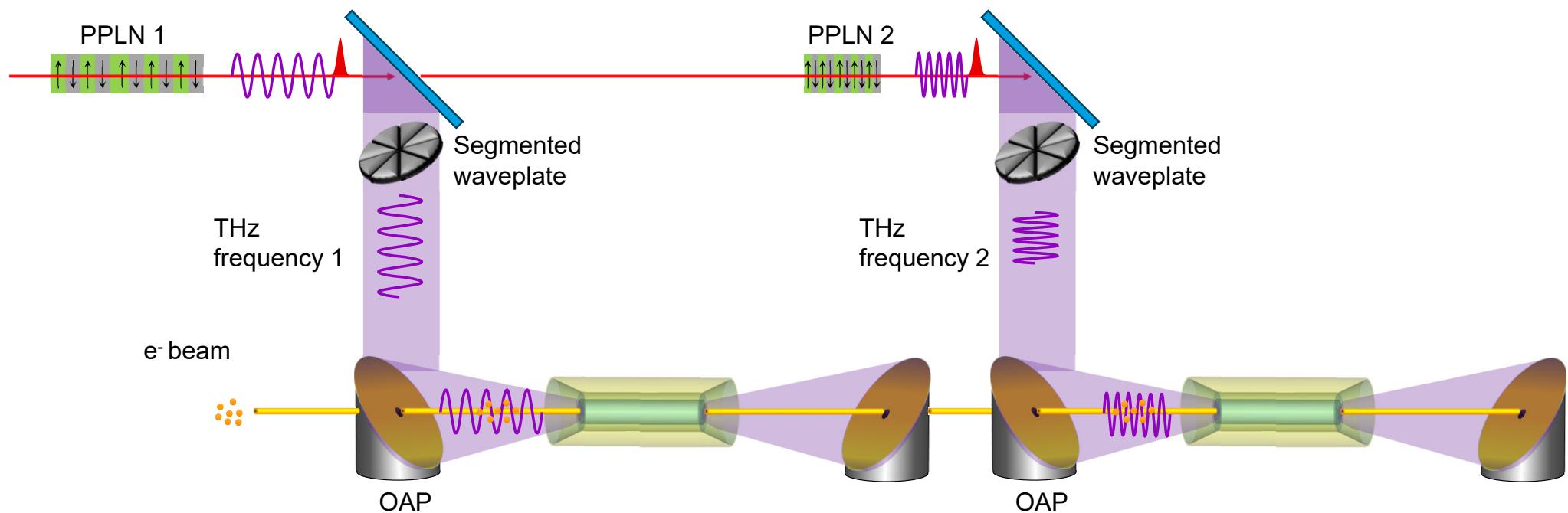
C. Vicario et al., Opt. Lett. **39**, 6632 (2014)



- Laser: Ti:Sapphire, 20 mJ, 100 Hz
- OPA output: 3 mJ, 1500 nm
- Bandwidth: >100 nm
- Crystal: DSTMS

C. Vicario et al., PRL **112**, 213901 (2014)

Staged THz interaction with laser recycling (Fourier Concept)



- Cost of insertion is low
- Main cost (laser) can be reused and flexibly repurposed

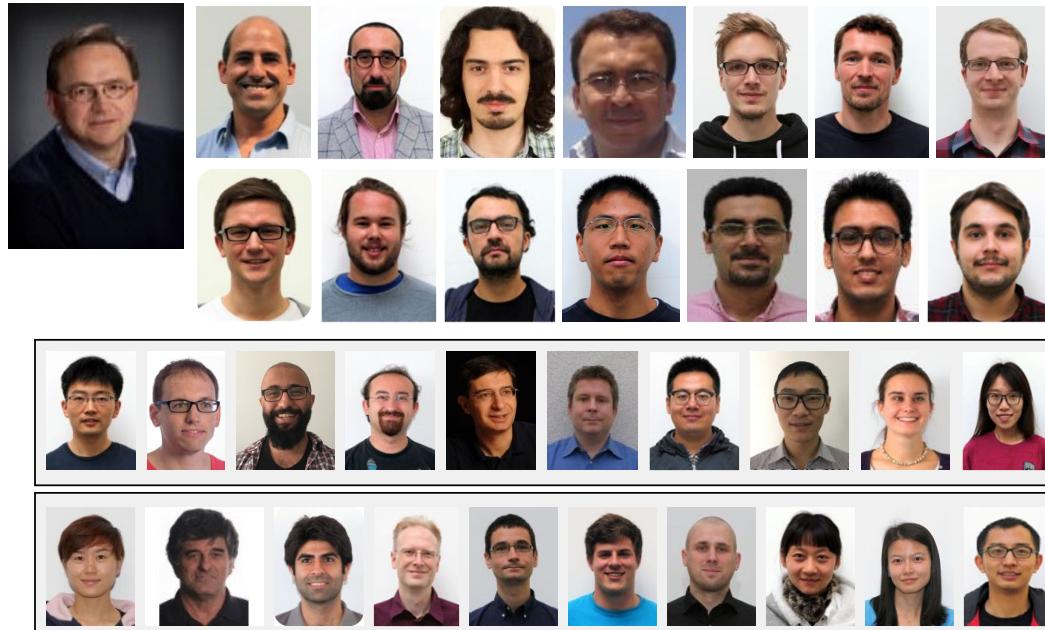
Next steps...

- Explore parameter space for Chirp Correction
 - Ask for sample problems
 - Full-scale simulations
 - Beam loading effects
 - Efficiency optimization
 - Dispersion management
- Apply for funding to do feasibility study
- Look for candidates to do proof-of principle experiments
 - LPAs are good options because of laser

AXSIS Team and Collaboration

Ultrafast Optics and X-rays Group

Alumni

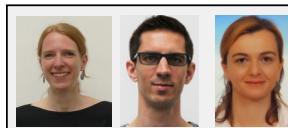


erc SYNERGY Grant

Coherent Diffractive Imaging Group

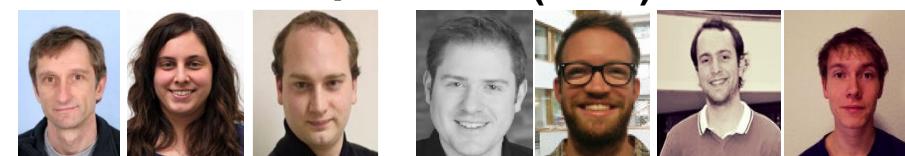


Bio-Physical-Chemistry Group

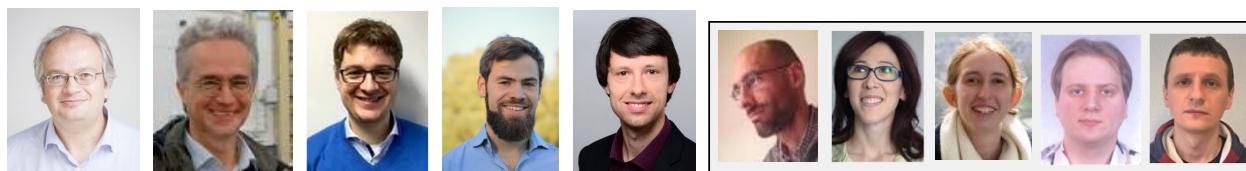


Detector Group

UHH (LUX)



Accelerator Division: DESY



Conclusions

- THz waves present unique options for electron acceleration and manipulation
- Energy requirements are reasonable or near reasonable for addressing key problems in plasma and conventional accelerators
- Worthwhile to explore the options
- Vision of the future: Synergy between THz, plasma and conventional accelerators



Attosecond X-ray Science: Imaging and Spectroscopy

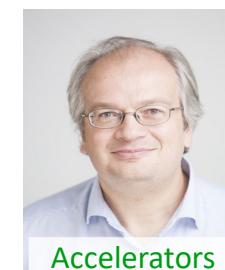
THz-driven Accelerator and Light Source



 **Ultrafast Optics & X-Rays**
Center for Free-Electron Laser Science
Group Leader:
Dr. Franz X. Kärtner



X-Rays
Henry Chapman
DESY & UHH



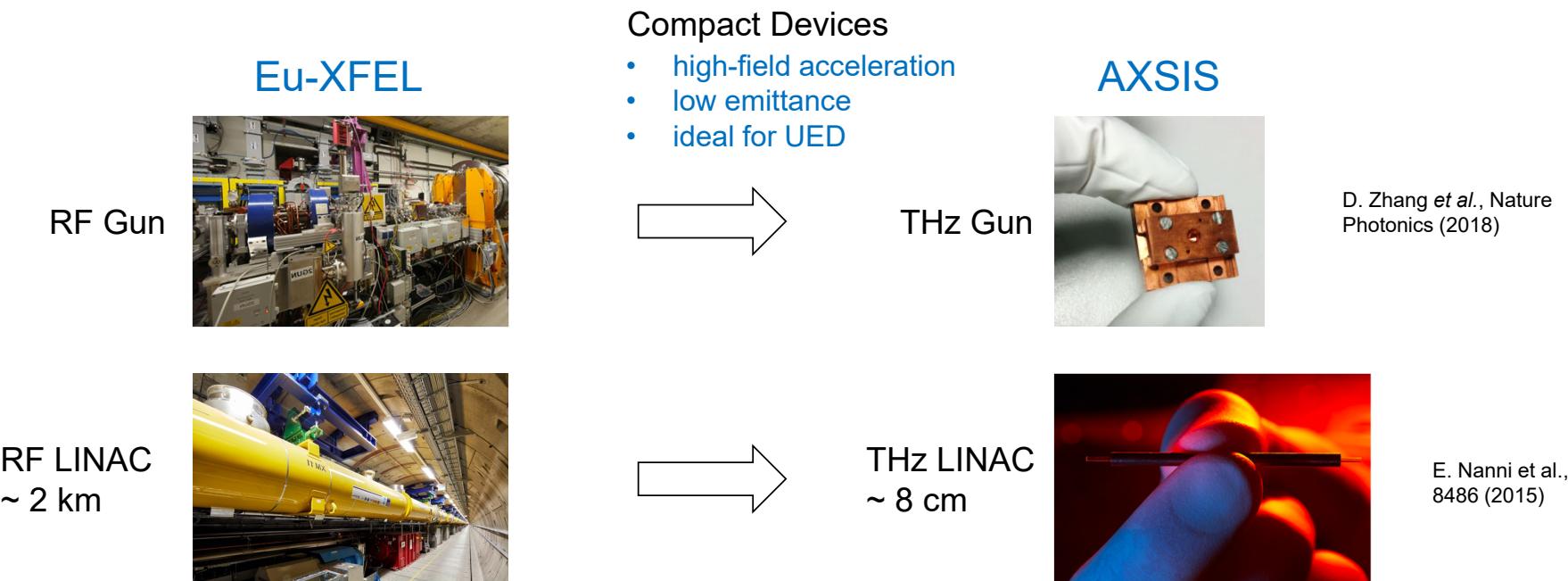
Accelerators
Ralph Assmann
DESY



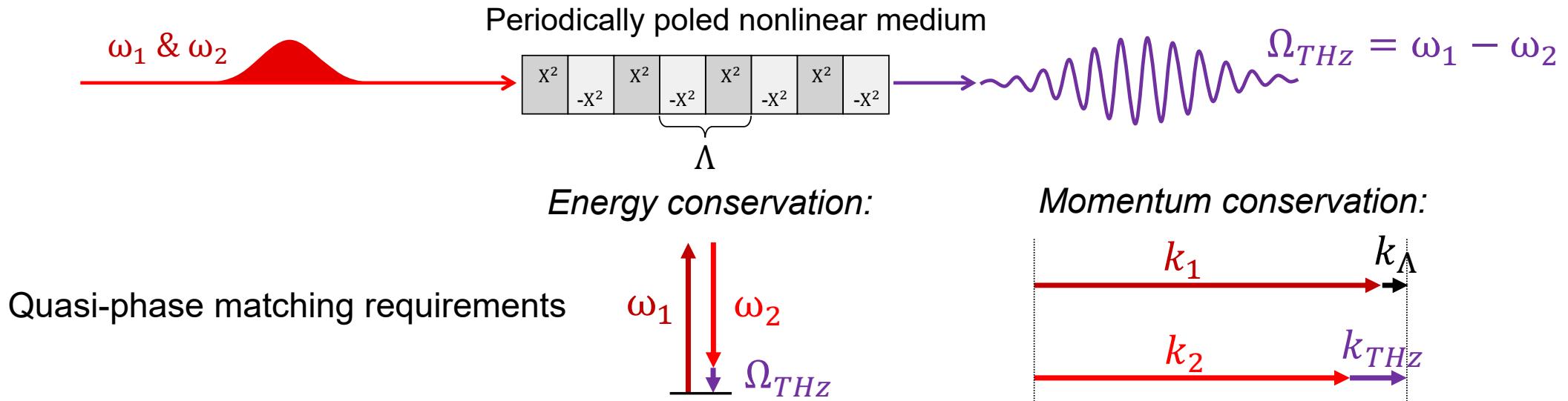
Bio-Phys-Chem
Petra Fromme
DESY & ASU

THz Accelerator Technology: Scaling from cm to mm and below

Parameter	RF Accelerators	THz Accelerators	THz Enhancement
Field Strengths, E_0	10 – 100 MV/m	100 – 1000 MV/m	~10x
Wavelength, λ	3 – 10 GHz	100 – 500 GHz	~100x
Field Gradients, dE_0/dz	~10 GV/m ²	~10,000 GV/m ²	~1000x



Quasi Phase-Matching Enables Efficient Multicycle THz Generation

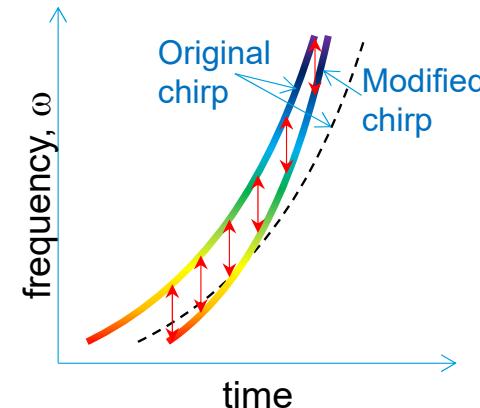
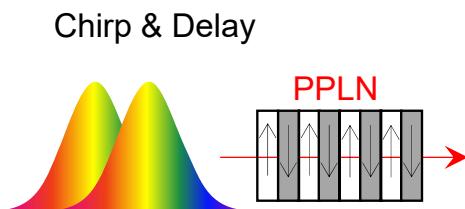


Recent results from our group in THz generation with QPM

Record Result	Energy	Material	Conv. Eff.	Article
Energy	1 μJ	LiNbO_3	0.12%	Carbajo <i>et al.</i> , Opt. Lett. 40, 5762 (2015)
Energy	40 μJ	LiNbO_3	0.13%	Ahr <i>et al.</i> , Opt. Lett. 42, 2118 (2017)
Energy	600 μJ	LiNbO_3	0.24%	Jolly <i>et al.</i> , Nat. Commun. 10, 2591 (2019)
Energy	1,300 μJ	LiNbO_3	0.14%	Lemery, Commun Phys 3, 150 (2020)
Material	0.7 μJ	KTP	0.16%	W. Tian <i>et al.</i> , Opt. Lett. 46, 741 (2021)
Efficiency	45 μJ	LiNbO_3	0.89%	H. T. Olgun <i>et al.</i> , Opt. Lett. 47, 2374 (2022)

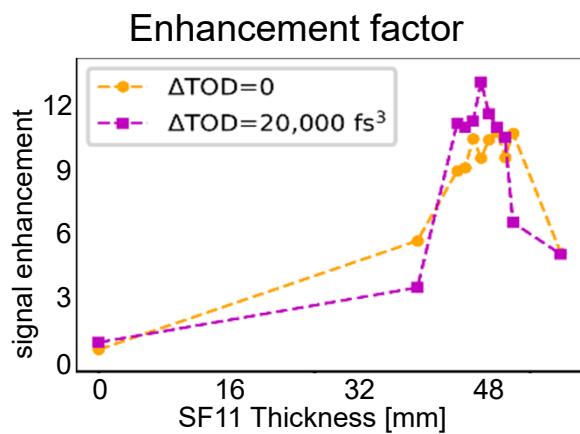
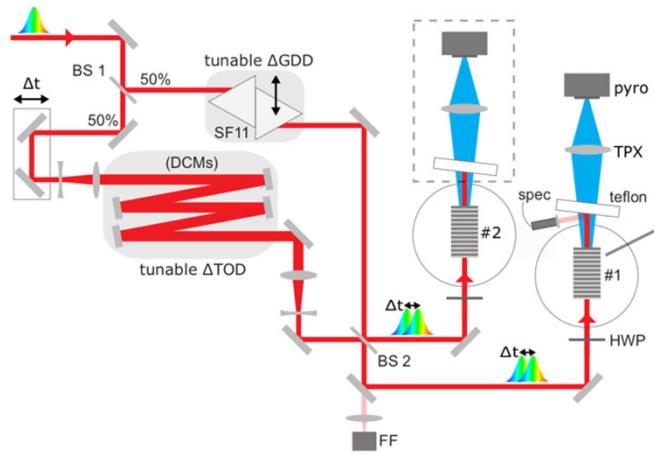
Record Energy Multicycle THz Generation using Chirp & Delay

S. Jolly et al., *Nature Commun.* **10**, 2591 (2019)
S. Jolly et al., *Opt. Express* **27**, 34769 (2019)

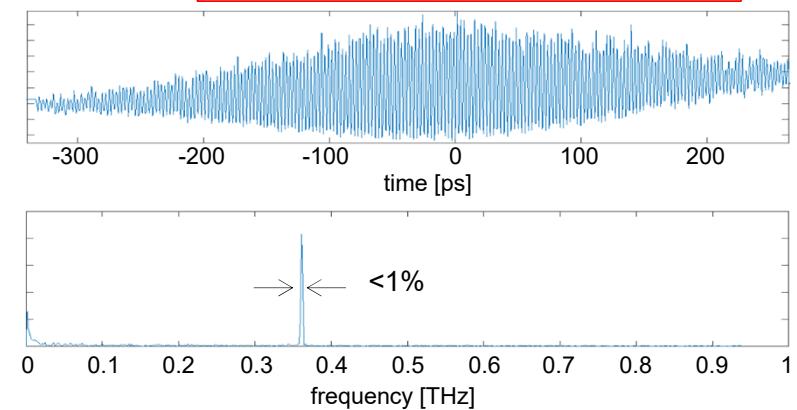


Solution:
Modify chirp of 1 pulse to produce constant $\Delta\omega = \Omega$

Asymmetric dispersion setup



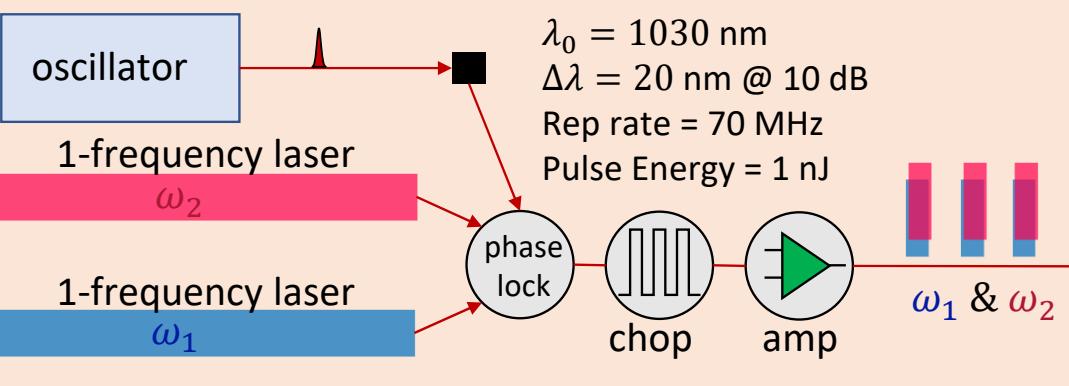
Record Narrowband THz Gen.
• 0.24% conversion eff.
• 1J Ti:S → **600 μJ THz**



Record multicycle THz efficiency using 2-line laser

Custom 2-Line Laser

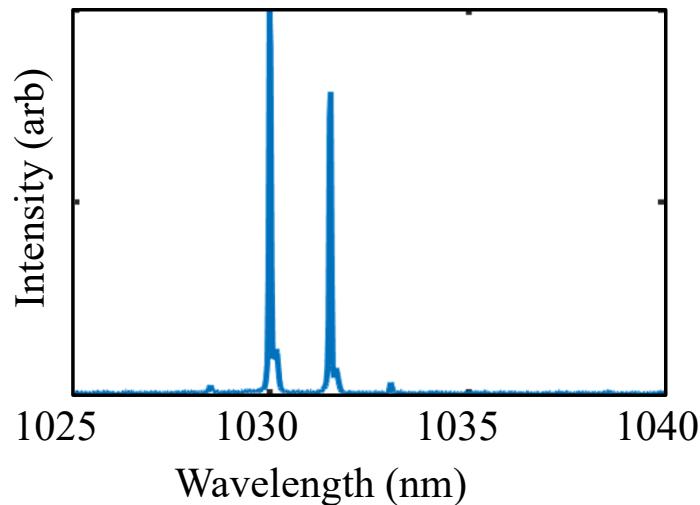
H. T. Olgun et al., Opt. Lett. 47, 2374 (2022)



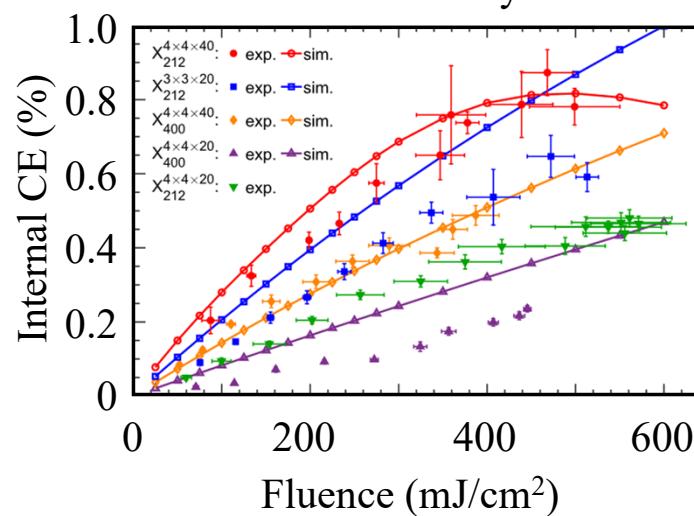
Record Narrowband THz Eff.

- 0.9% conversion eff.
- 10 mJ Yb:YAG \rightarrow 45 μJ THz

Laser Spectrum (before)

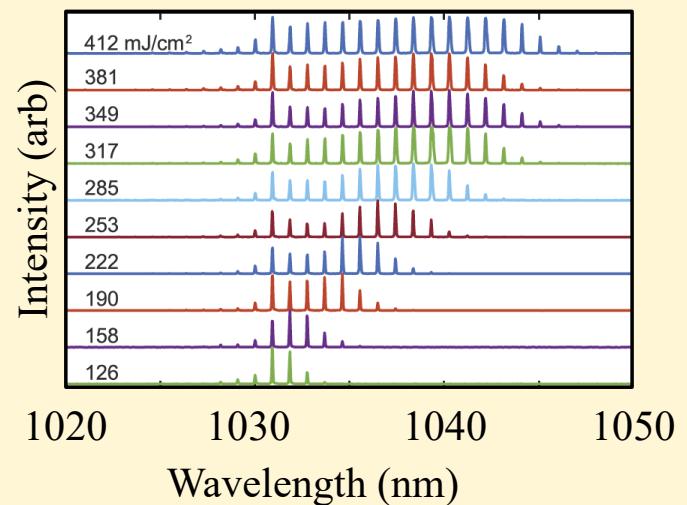


Conversion Efficiency

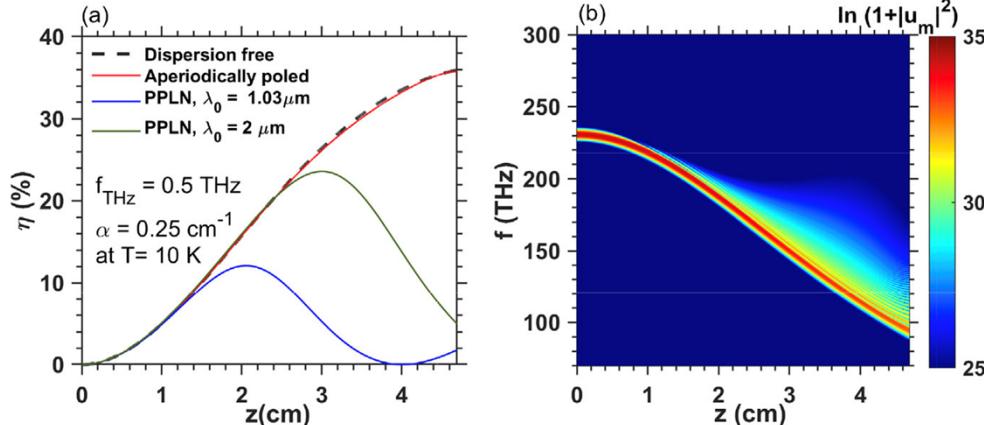


Strong Cascading

Laser Spectrum (after)

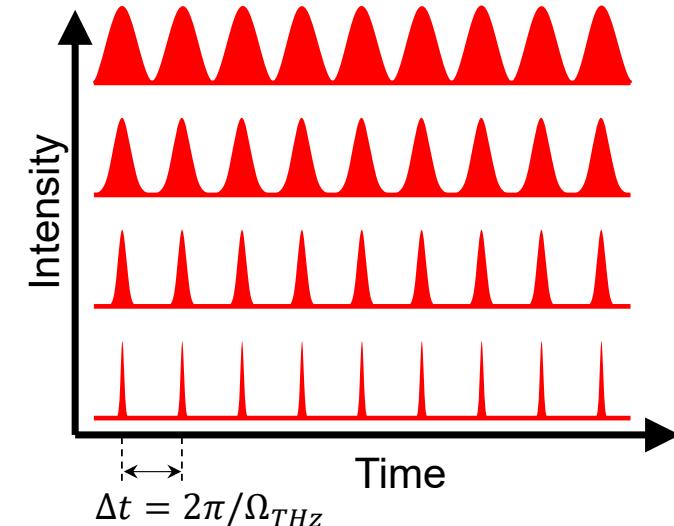
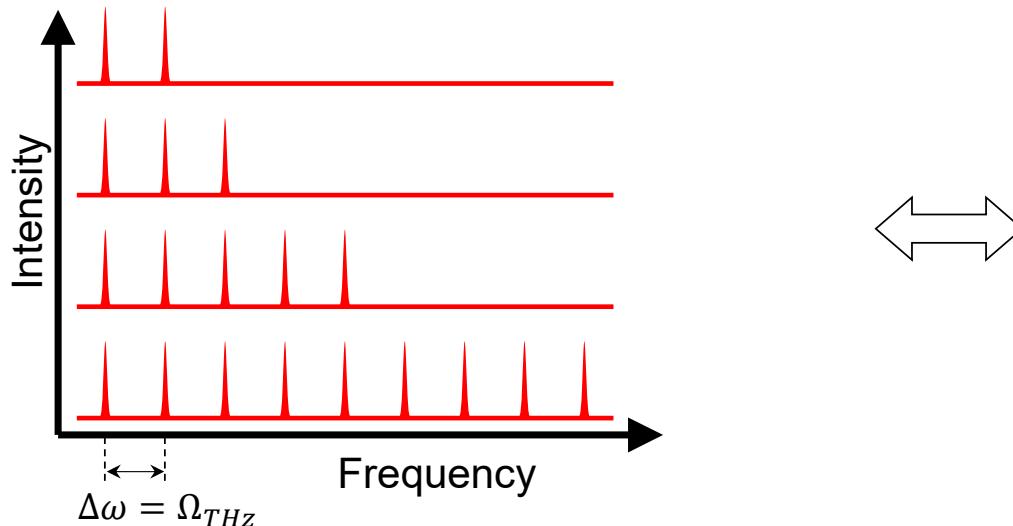


THz Generation becomes more efficient with more spectral lines

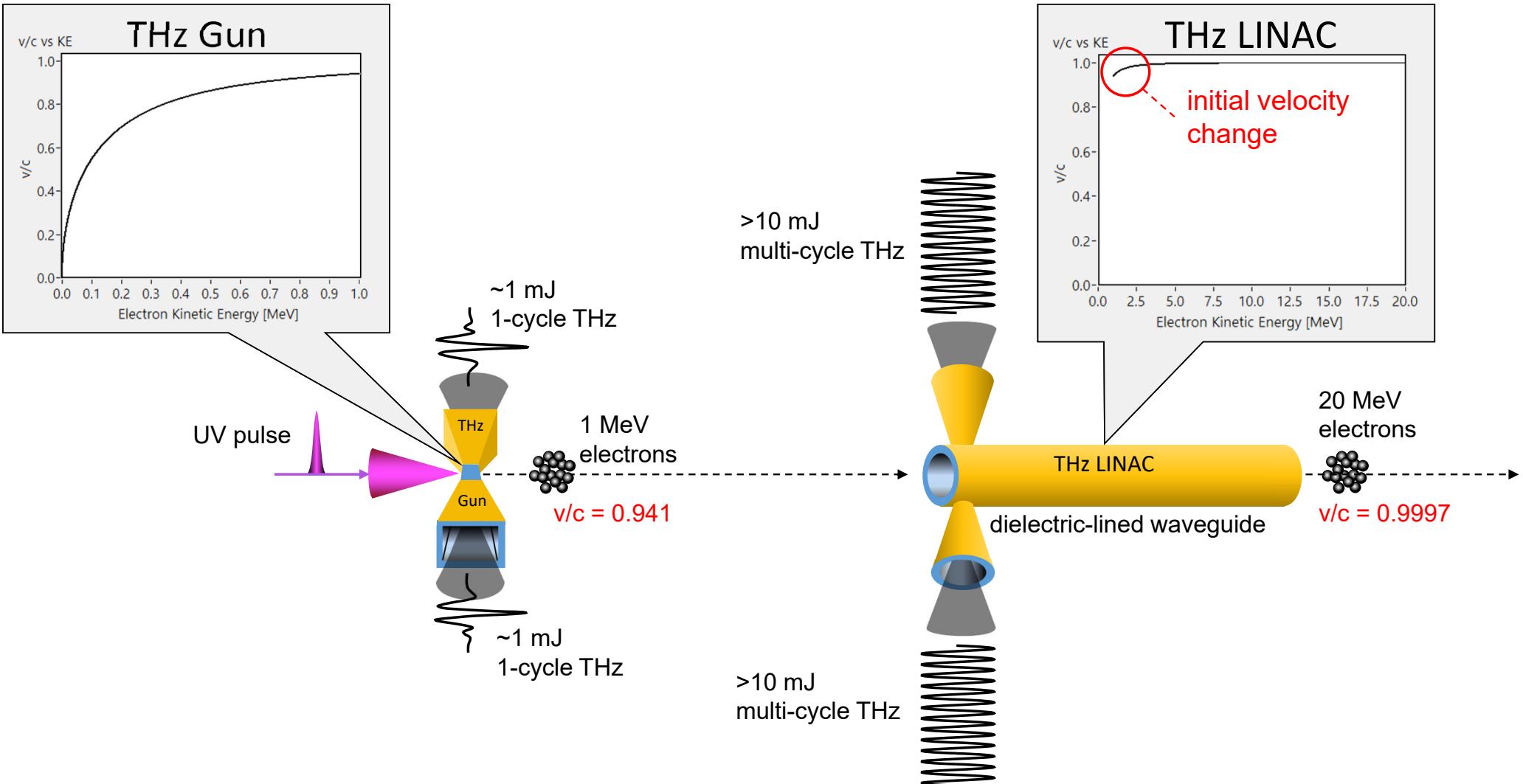


Theory predicts multiple percent efficiency is possible with optimized laser input

- K. Ravi & F. X. Kärtner, *Laser & Photon. Rev.* **14**, 2000109 (2020)
K. Ravi et al., *Opt. Lett.* **24**, 25582 (2016)
A. G. Stepanov, *JETP Lett.*, **85**, 227(2007)
M. Cronin-Golomb, *Opt. Lett.* **29**, 2046 (2004)

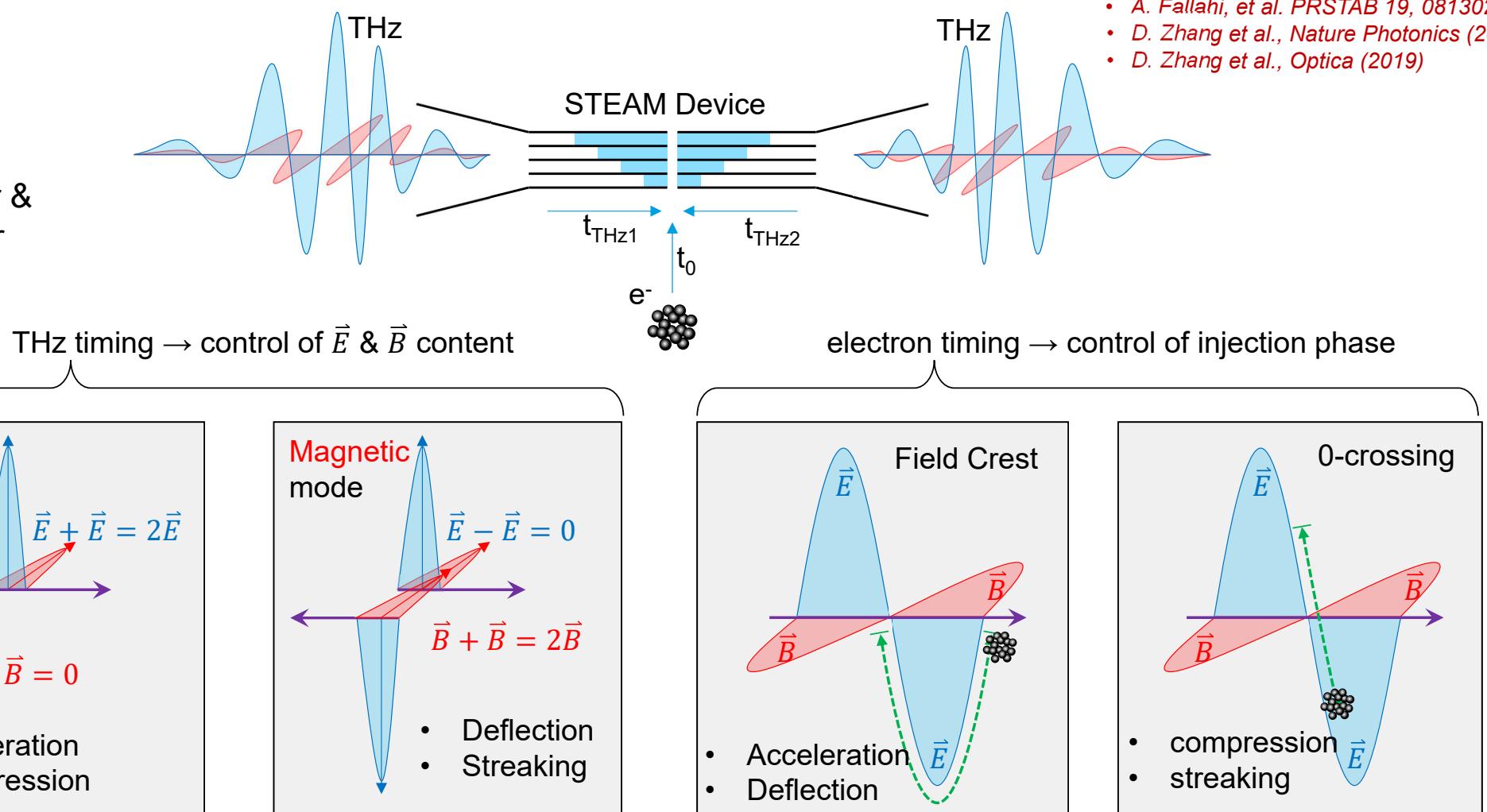


THz Accelerator is composed of THz Gun + THz LINAC



THz and electron timing controls **STEAM** function

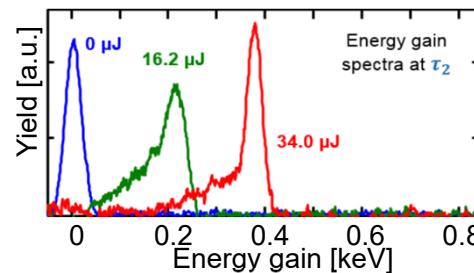
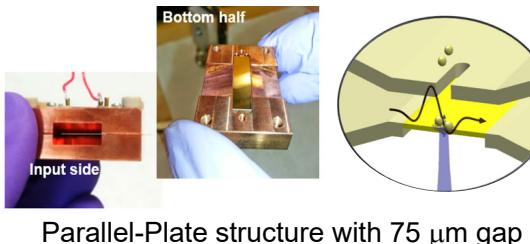
**Segmented
THz
Electron
Accelerator &
Manipulator**



- A. Fallahi, et al. PRSTAB 19, 081302 (2016)
- D. Zhang et al., Nature Photonics (2018)
- D. Zhang et al., Optica (2019)

Preliminary work demonstrates elements required for a THz gun

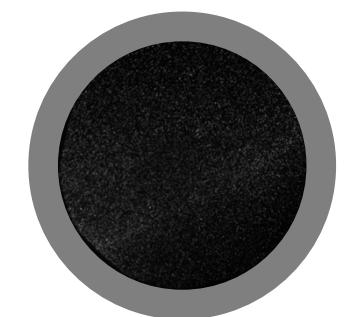
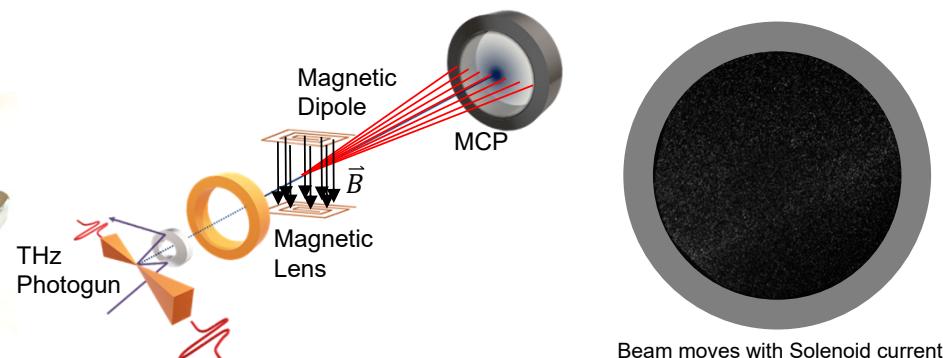
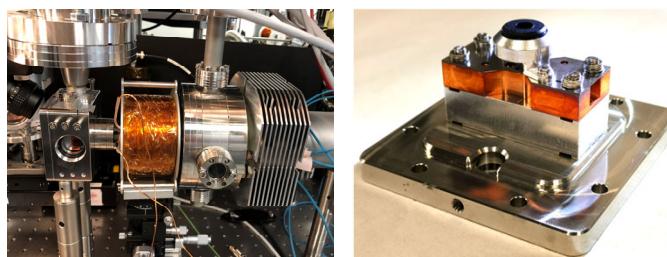
- THz Gun I: $0 \rightarrow 0.8$ keV acceleration



- 1st demonstration of THz gun
- Sub-keV electrons demonstrated
- One-side pumped
- 1D THz field concentration

W. Huang, et al., Optica 3, 1209 (2016)
A. Fallahi, et al., PRSTAB 19, 081302 (2016)

- THz Gun III: $0 \rightarrow 3$ keV acceleration

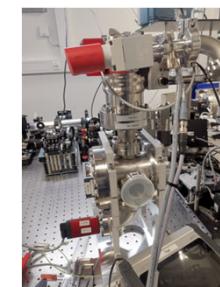
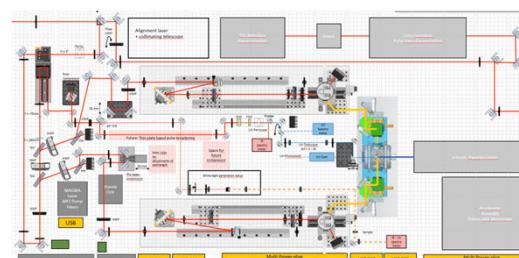
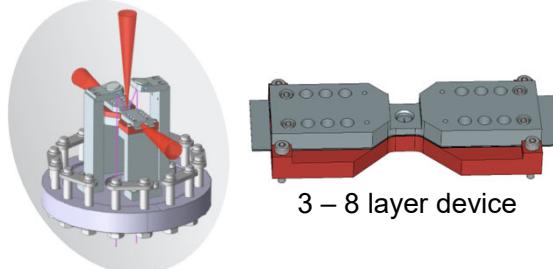


Beam moves with Solenoid current

- $E_{kin} = 2.43$ keV
- Highest energy for THz-driven photogun
- Control over injection demonstrated

N. H. Matlis et al. in preparation

- THz Gun IV: $0 \rightarrow 200$ keV acceleration (in development)

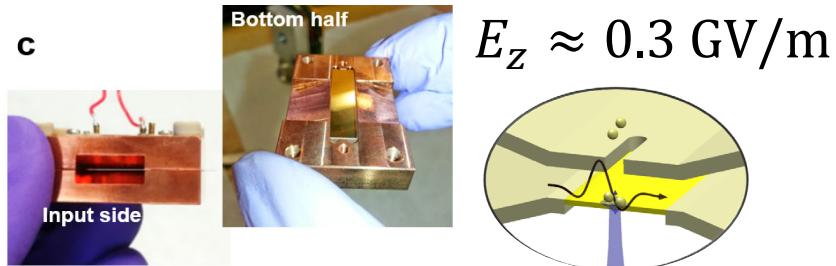


Electron Spec's

- $E_k \sim 190$ keV
- $\Delta E_k/E_k \sim 3.2\%$
- $\Delta t_{RMS} \sim 37$ fs
- $\sigma_x \sim 34$ μm
- $\epsilon_N \sim 0.15$ mm · mrad

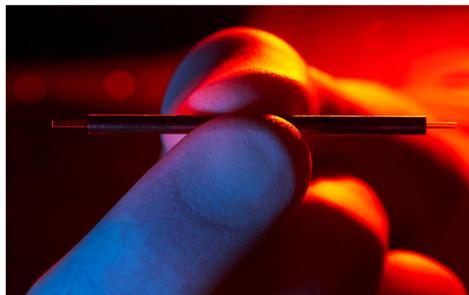
THz acceleration – Nonrelativistic electrons

THz Gun: $0 \rightarrow 0.8$ keV acceleration



W. Huang, et al., Optica 3, 1209 (2016)
A. Fallahi, et al., PRSTAB 19, 081302 (2016)

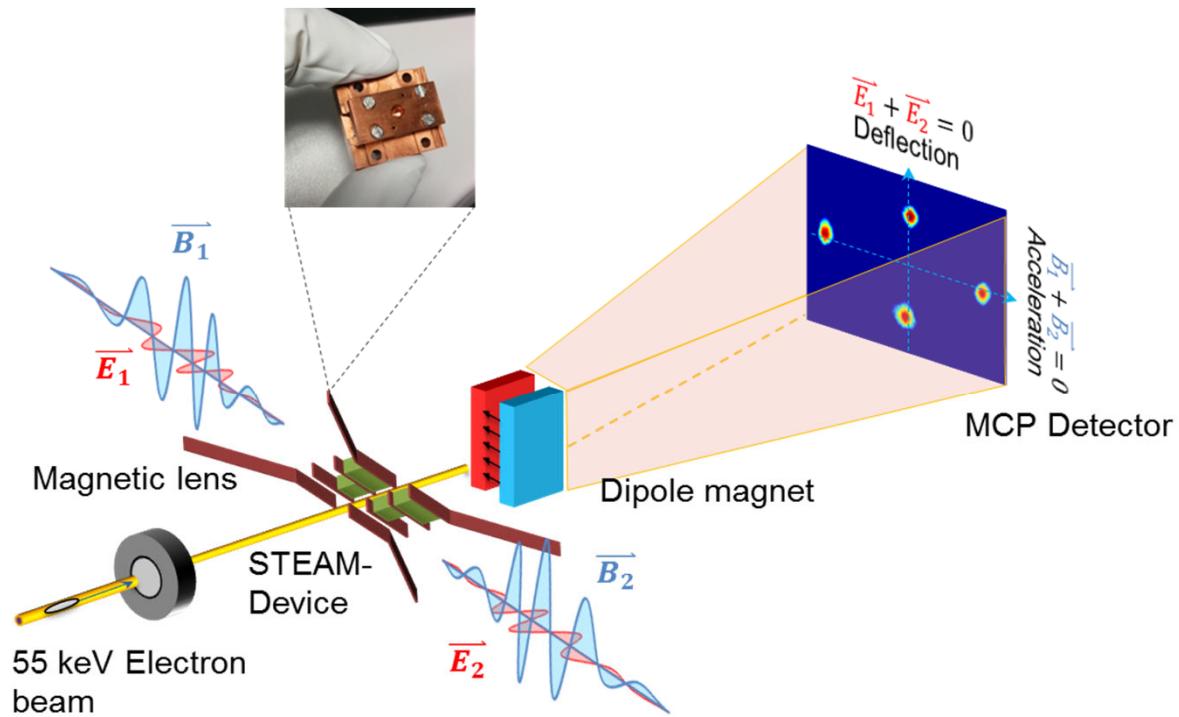
THz LINAC: ± 7 keV energy modulation



mm-scale THz waveguide

E. Nanni et al., Nature Comm. 6, 8486 (2015)

THz Manipulator (STEAM): acceleration, compression, focusing and streaking

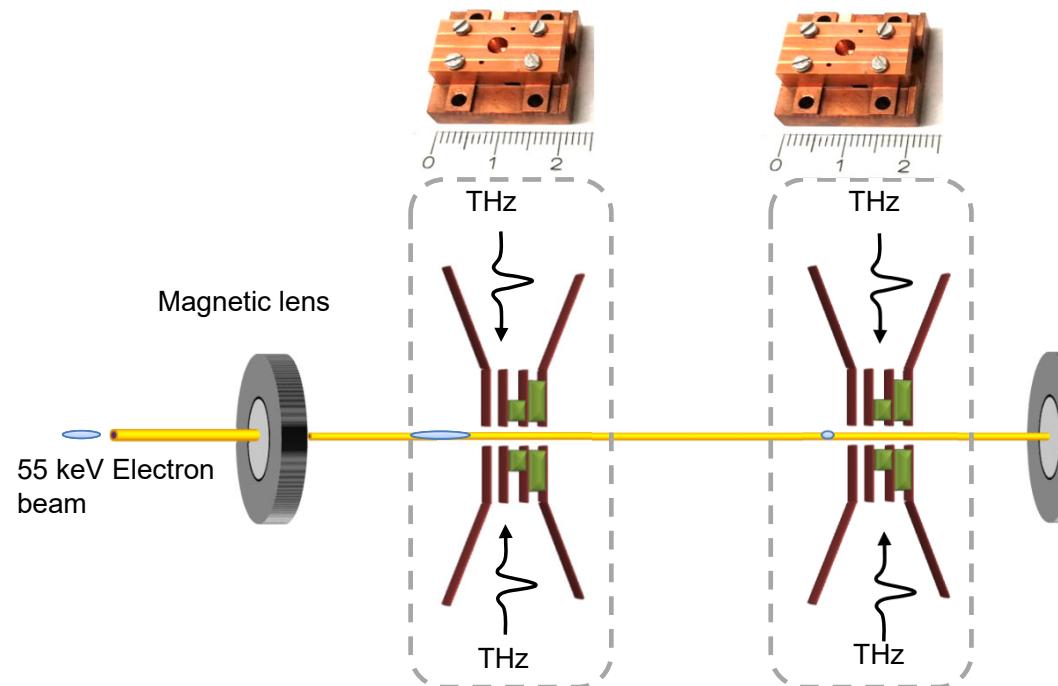


D. Zhang et al., Nature Photonics (2018)
D. Zhang et al., Optica Vol. 6, pp. 872-877 (2019)

STEAM devices combined & configured for diverse applications

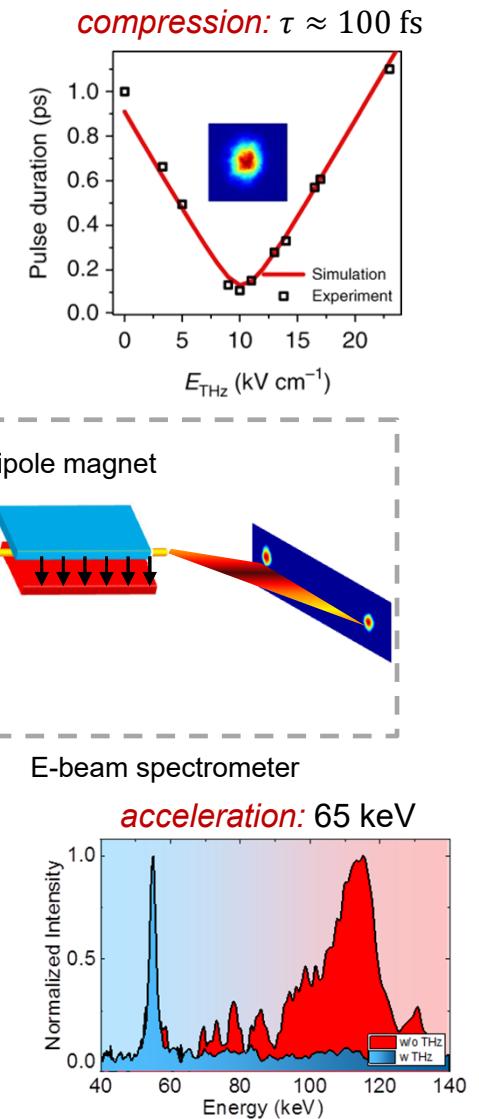
D. Zhang et al., *Nature Photon.* **12**, 336 (2018)

D. Zhang et al., *Optica* **6**, 872 (2019)

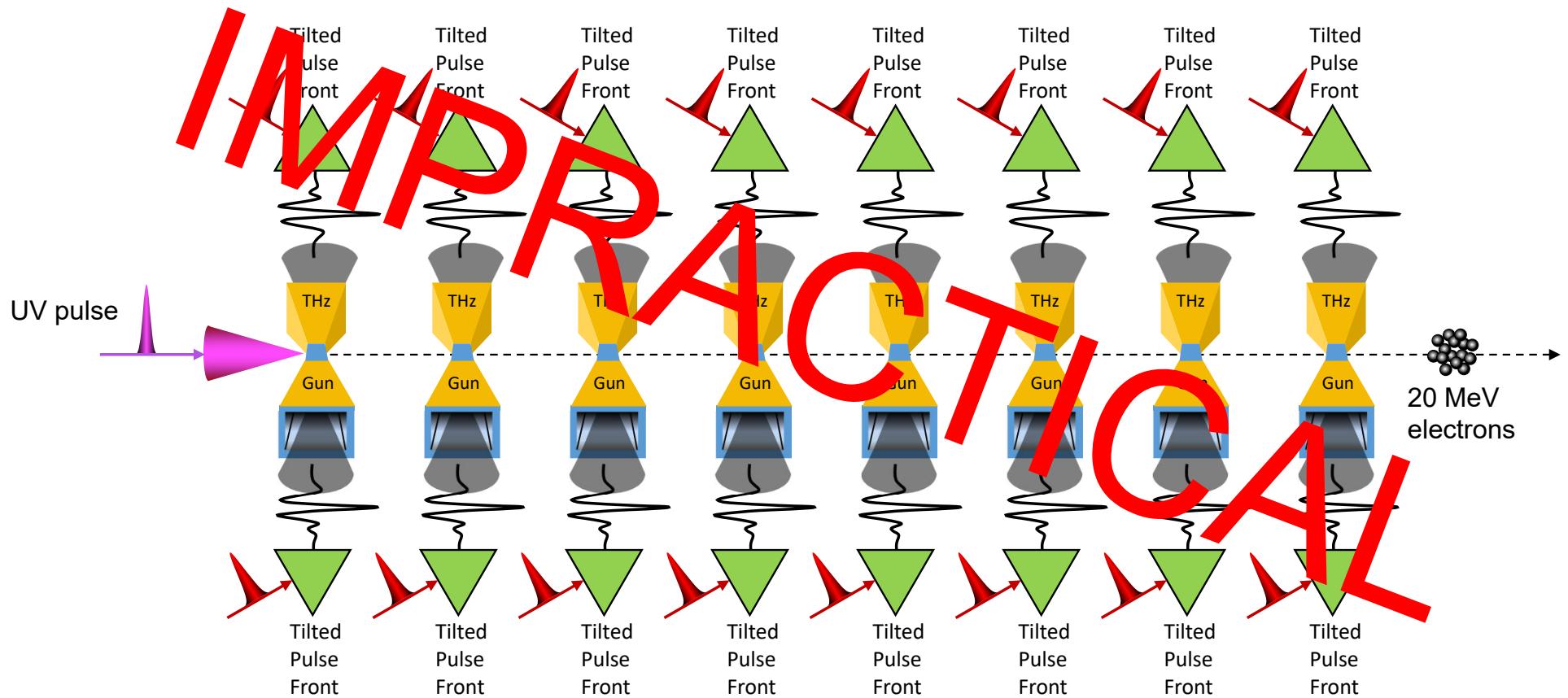


Pulse duration meas: **Buncher** **Streaker**

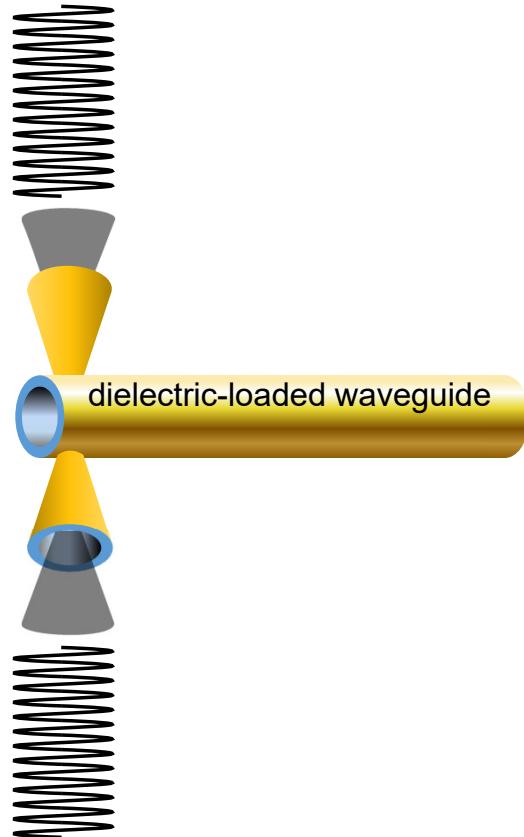
Electron Acceleration: **Buncher** **Accelerator**



Why not use a series of transversely-pumped boosters?



Benefits and Challenges of Dielectric-Loaded Waveguide LINACs



Benefits

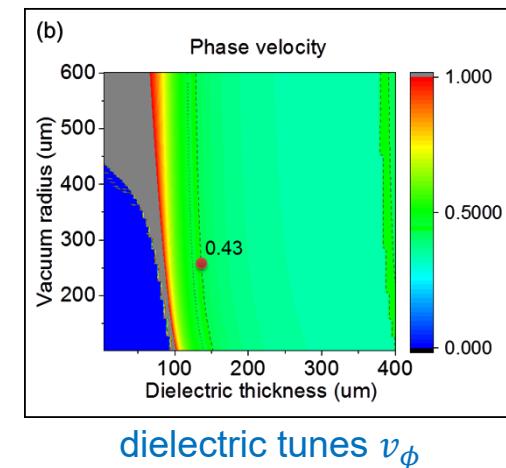
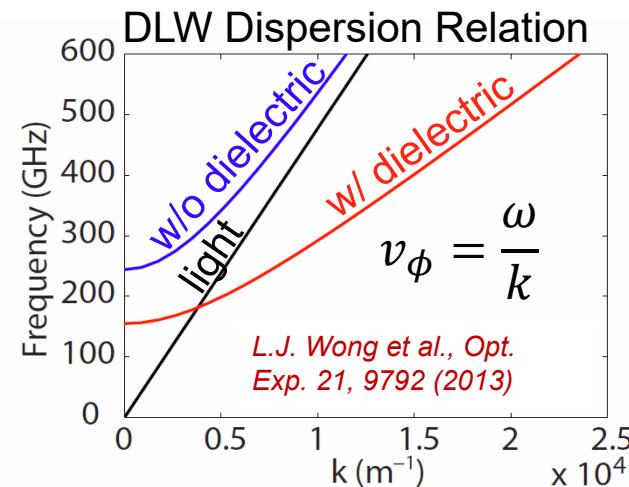
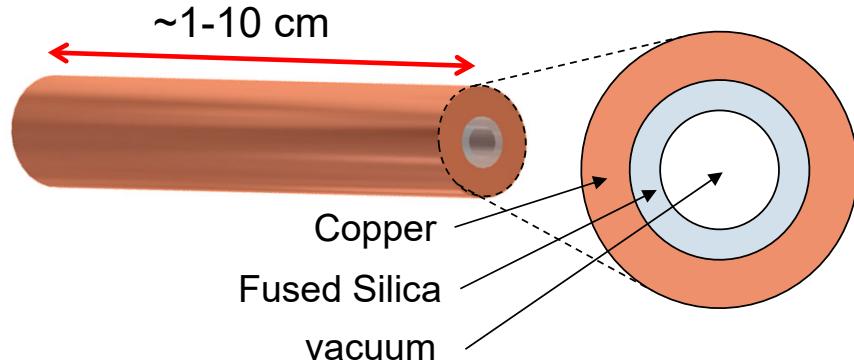
- Simple fabrication
- Tunable by varying the dimensions
- Travelling-wave nature allow efficient driving with pulses
- Dielectric sustains higher fields

Critical Issues

- **Tuning:** adjusting the speed of the wave
- **Dispersion** causes the THz pulse evolves
- **Dephasing:** electrons to move relative to wave crests
- **Walk-off:** electrons move relative to pulse envelope)
- **Inefficient** for low electron energies
- **Coupling** requires mode conversion to TM_{01}

Phase velocity tuned using dielectric layer thickness

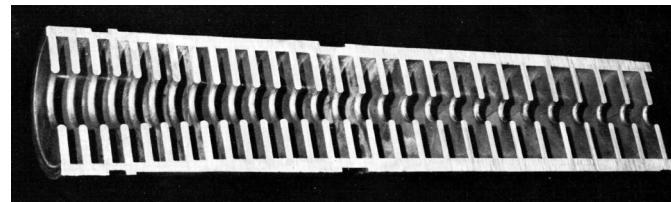
“Dielectrically-Loaded Waveguide”



“Cavity Waveguide”



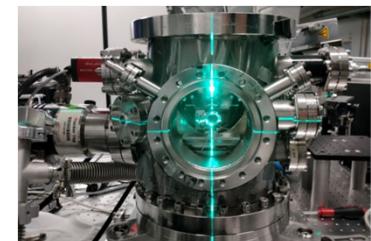
“Disk-Loaded Waveguide”



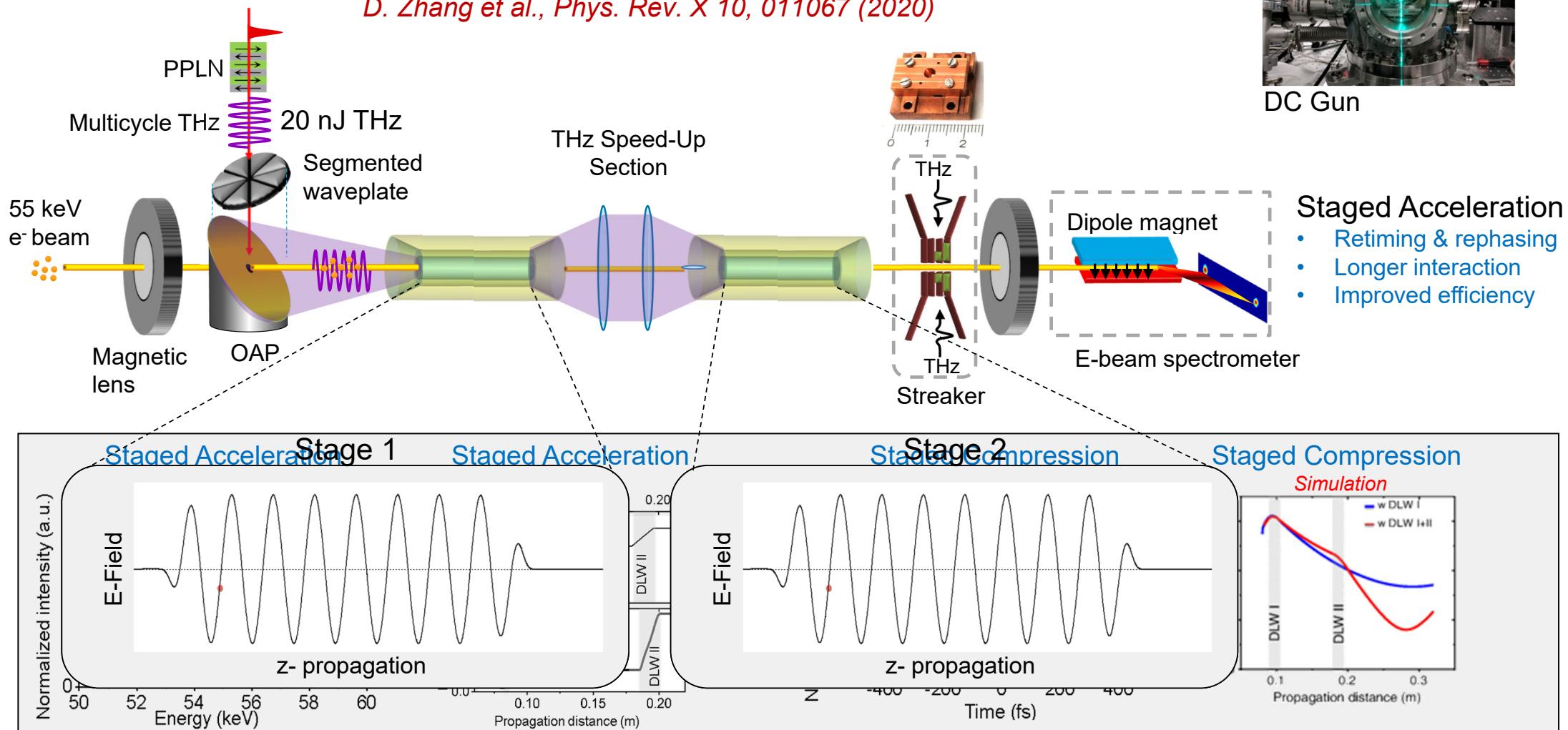
Restrictions used to tune phase velocity

Staged THz acceleration with energy recycling

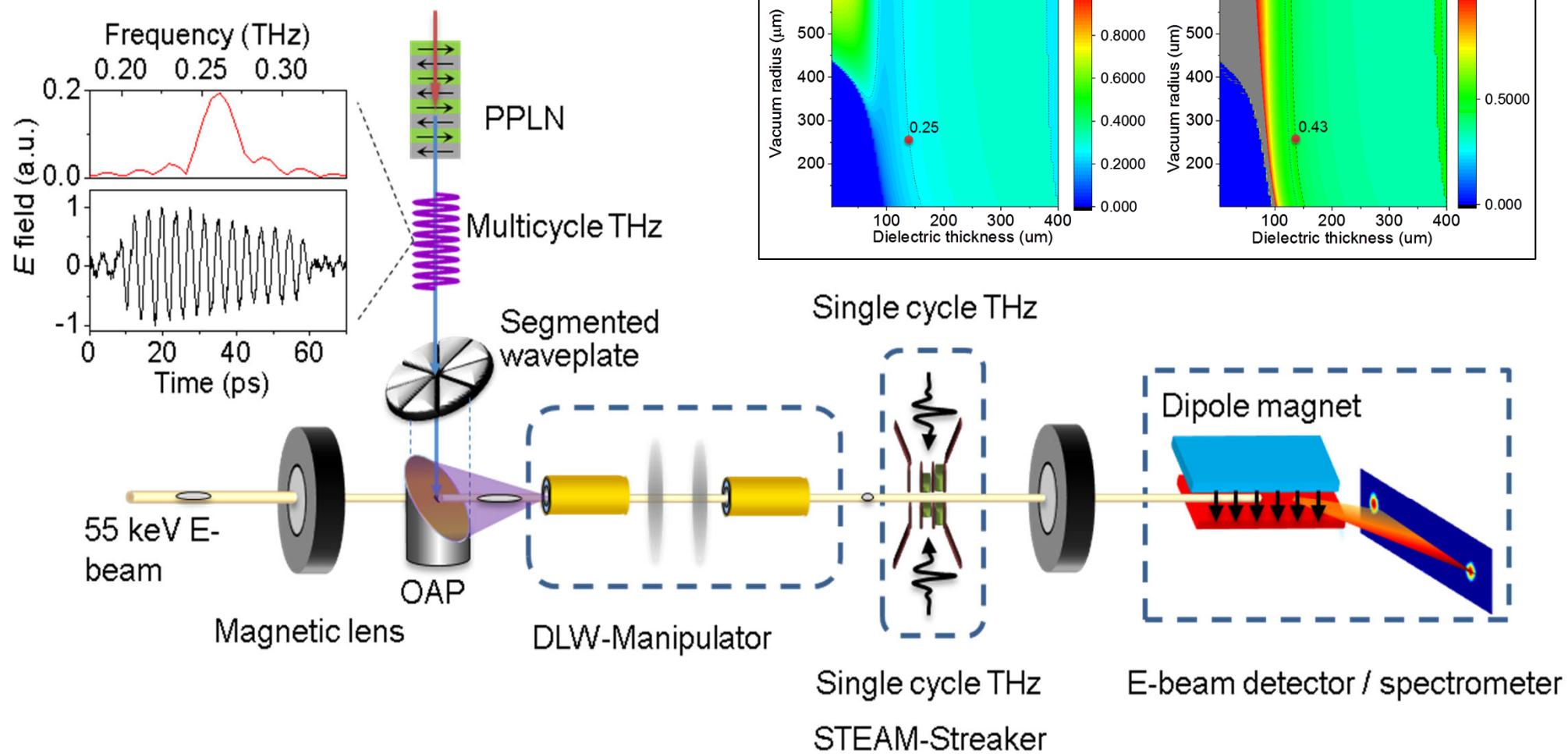
D. Zhang et al., Phys. Rev. X 10, 011067 (2020)



DC Gun

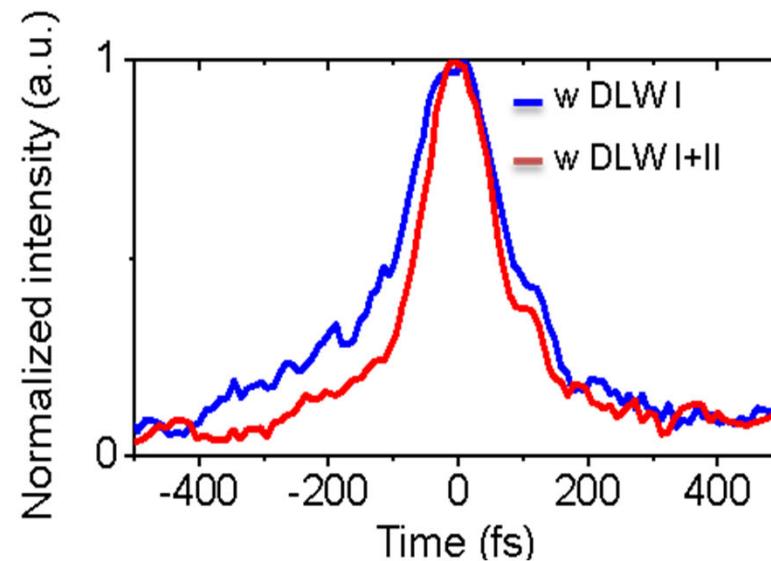
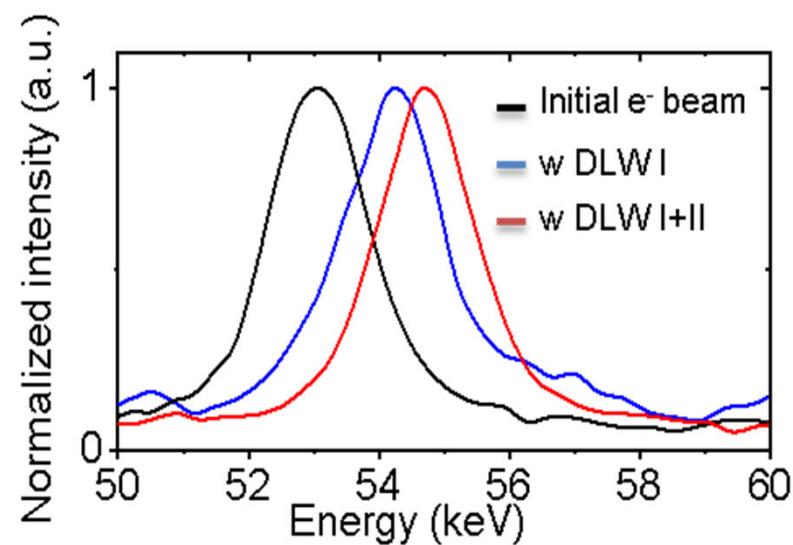
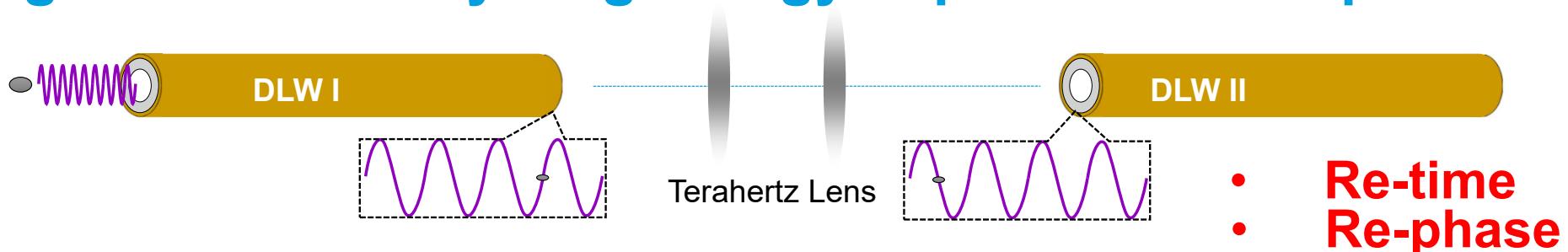


Solution for low injection v_e : staged DLWs & energy recycling



Staging DLWs and recycling energy improves LINAC performance

(a)



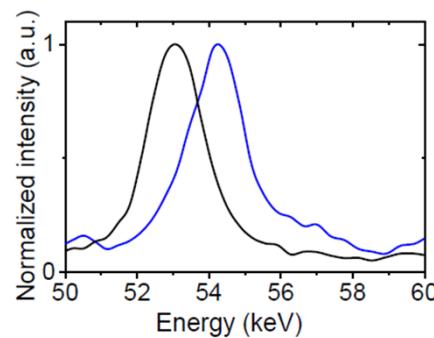
$$v/c (53 \text{ keV}) = 0.423$$

$$v/c (54.2 \text{ keV}) = 0.427$$

D. Zhang et al., Phys. Rev. X 10, 011067 (2020)

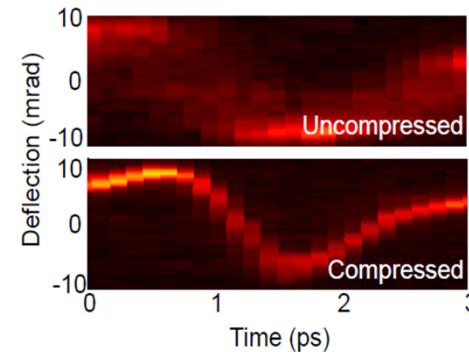
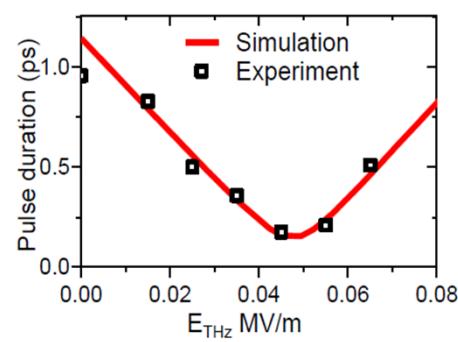
Electron acceleration, compression and focusing with DLW

Electron acceleration

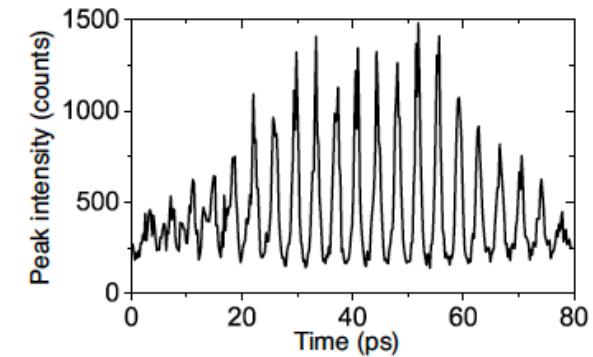
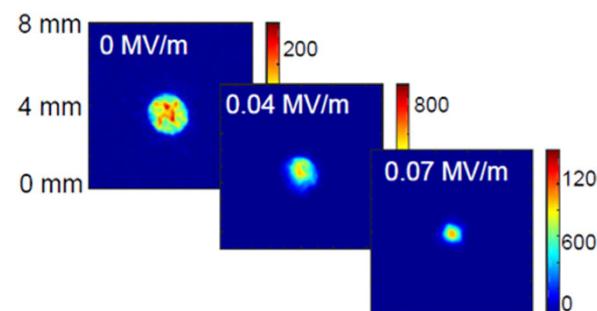


- THz energy ~ 20 nJ
- electron energy gain ~ 1 keV
- minimal dephasing

Electron compression



Electron focusing



TM_{01} - mode of DLW Group and Phase velocity

