



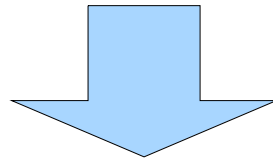
E-989 $g-2$ experiment -

**Study of systematic errors affecting the
electron spectrum
&
Straw Detectors**

The purpose of g-2 experiment is to measure the muon anomalous magnetic momentum to ± 0.14 ppm precision

$$\overline{\mu}_\mu = g_\mu \left(\frac{q}{2m} \right) \overline{s}$$

Dirac theory $\rightarrow g_\mu = 2$ is expected for a structureless, spin-1/2 particle of mass m and charge $q = \pm e$



Anomalous magnetic momentum

$$a_\mu = \left(\frac{g_\mu - 2}{2} \right)$$

Radiative corrections

Leading RC → lowest-order QED process involving the exchange of a virtual photon

$$a_\mu = \alpha / 2\pi = 1.16 \cdot 10^{-3}$$

Complete standard model value of a_μ

QE
D

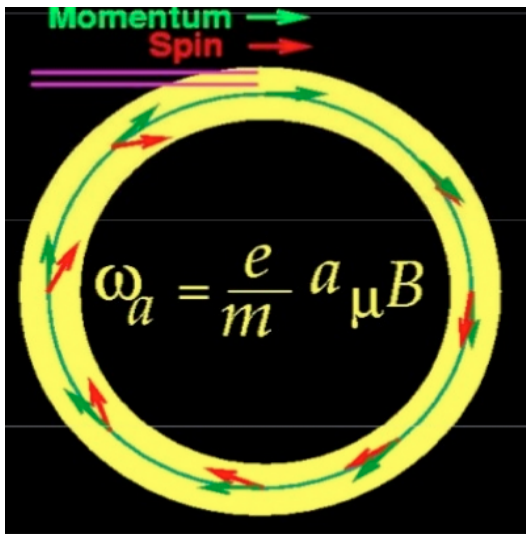
Weak

Hadronic

0.5 ppm precision

E821 Experimental technique

Muons moving in the horizontal plane of a magnetic storage ring, have an anomalous precession spin frequency:



$$\bar{\omega}_a = \omega_s - \omega_c = -a_\mu \frac{q \bar{B}}{m}$$

Spin precession frequency

Cyclotron frequency

$$\omega_s = -\frac{g q \bar{B}}{2m} - (1 - \gamma) \frac{q \bar{B}}{m \gamma}$$

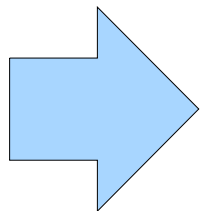
$$\omega_c = -\frac{q \bar{B}}{m \gamma}$$

E821 Experimental technique-I

Because electric quadrupoles are used to provide vertical focusing in the storage ring, their electric field is seen in the muon rest frame as a motional magnetic field that can affect the spin precession frequency:

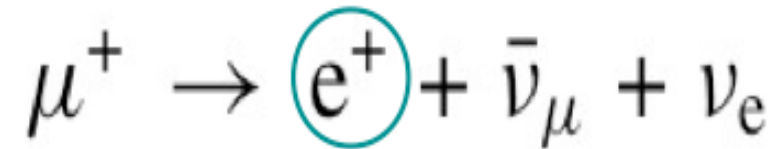
$$\bar{\omega}_a = -\frac{q}{m} \left[a_\mu \bar{B} - \left(a_\mu - \frac{1}{\gamma^e - 1} \right) \frac{\bar{B} \times \bar{E}}{c} \right]$$

→ = 0 if $\gamma = 29.3$

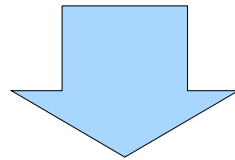


Magic momentum $p_\mu = 3.094 \text{ Gev} / c$

E821 Experimental technique-II



Because of parity violation in the weak decay of the muon, a correlation exists between the muon spin and decay positron direction. Positrons of low energy are emitted preferentially along the muon spin direction and those of high energy are more likely emitted opposite to the spin.



We can measure the muon anomaly from the time distribution of positrons of a certain energy

E821 Experimental technique-III

The integrated decay electron distribution in the lab frame is:

$$N(t) = \frac{N}{\gamma \tau_\mu} e^{\frac{-t}{\gamma \tau_\mu}} \left[1 + A \cos(\omega_a t + \phi) \right]$$

Normalization coefficient

Time dilated muon lifetime

Asymmetry

Phase

For a threshold energy $E_{\text{th}} = 1.8$ GeV, the asymmetry is $A \approx 0.4$ and the average figure-of-merit is maximized, so the statistical uncertainty on ω_a is minimized.

Because I want to study the systematic errors, I chose this configuration.

ROOT simulation-I

$$esp = \frac{N}{\gamma \tau_{\mu}} \cdot e^{\frac{-t}{\gamma \tau_{\mu}}} \cdot [1 + A \cdot \cos(\omega \cdot t + \phi)]$$

Guess:

$$A = 0.4$$

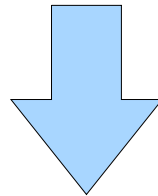
$$\gamma \cdot \tau_{\mu} = 64.4 \mu\text{s}$$

$$\omega_a = 4.37 \mu\text{s}$$

$$\Phi = 1.5 \text{ rad}$$

N chosen so that the histogram was filled by 10^{12} points

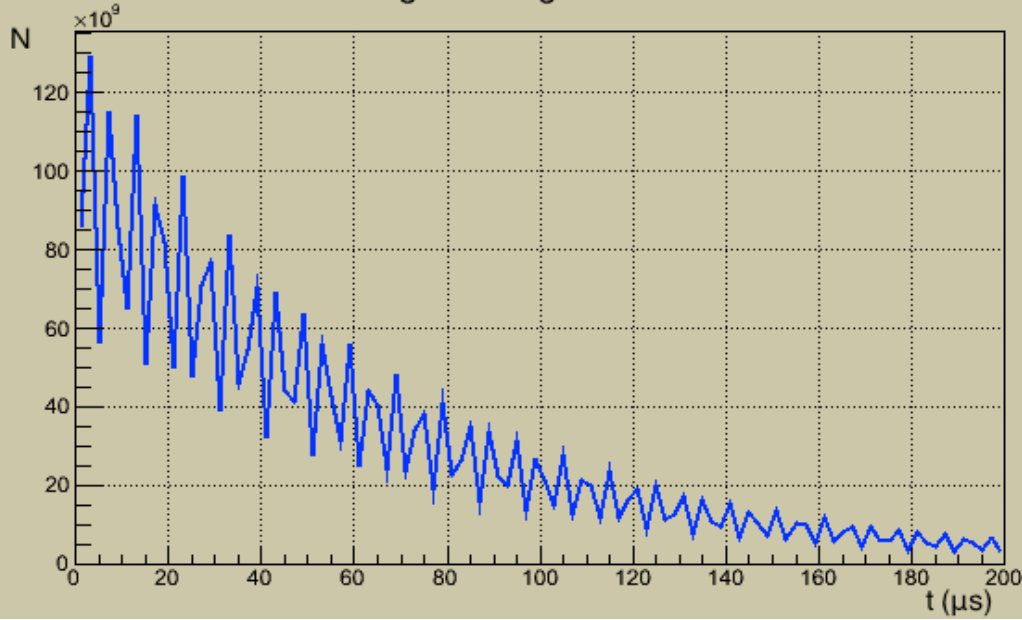
IDEAL CASE



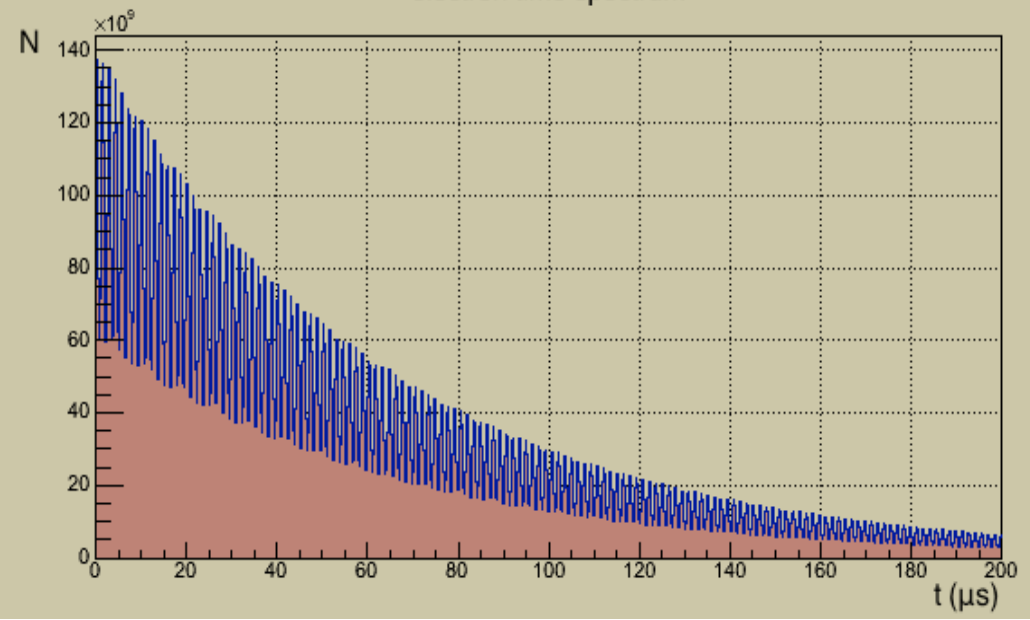
Statistical errors negligible

An example of electron decay time distribution histogram

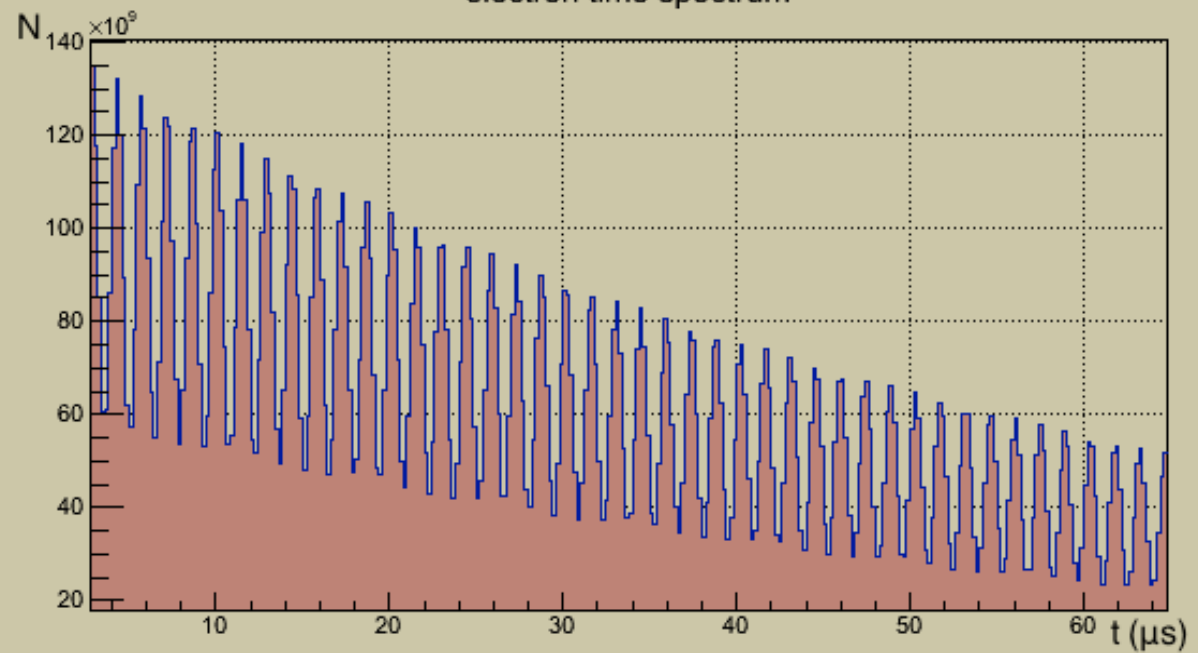
generating function



electron time spectrum



electron time spectrum



ROOT simulation-II

Systematic errors can be represented, at the first order, setting the asymmetry as a linear function of time in the generating function:

$$A = c_1 (1 + c_2 t)$$

Generating function

All parameters, except A, are setted to their values at $E_{th} = 1.8$ GeV

$$esp = \frac{N}{\gamma \tau_\mu} \cdot e^{\frac{-t}{\gamma \tau_\mu}} \cdot [1 + c_1 (1 + c_2 \cdot t) \cdot \cos(\omega \cdot t + \phi)]$$

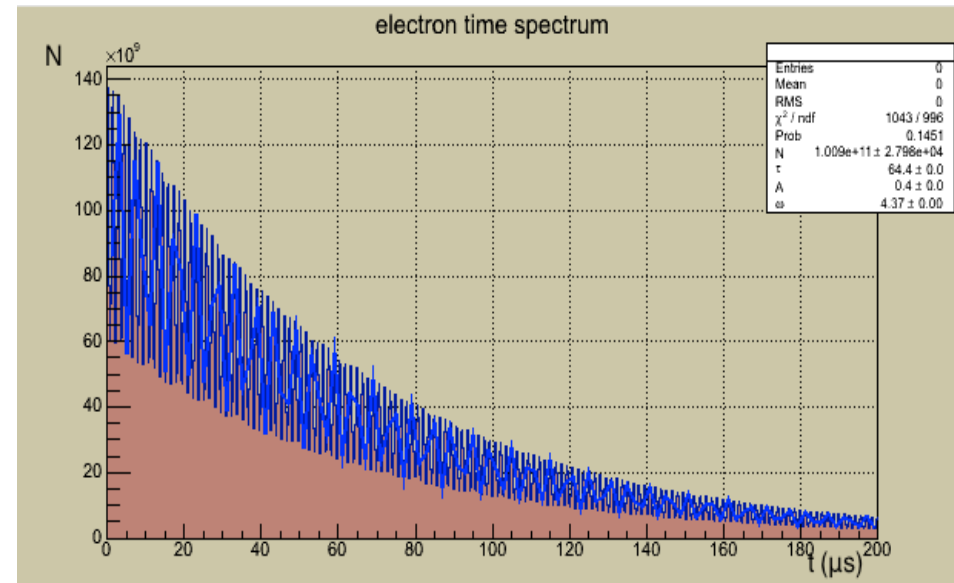
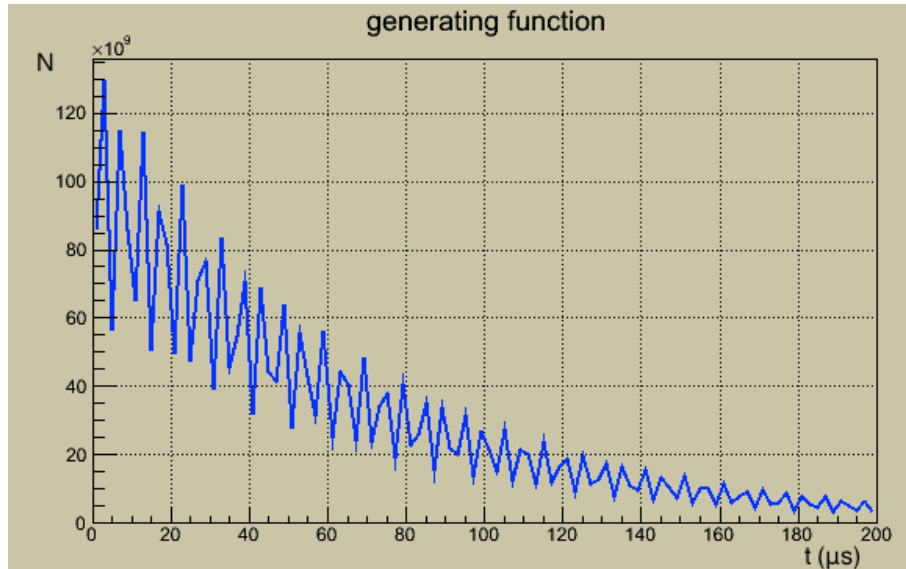
$$esp1 = \frac{N}{\gamma \tau_\mu} \cdot e^{\frac{-t}{\gamma \tau_\mu}} \cdot [1 + A \cdot \cos(\omega \cdot t + \phi)]$$

Fit function

Free parameters in the fit

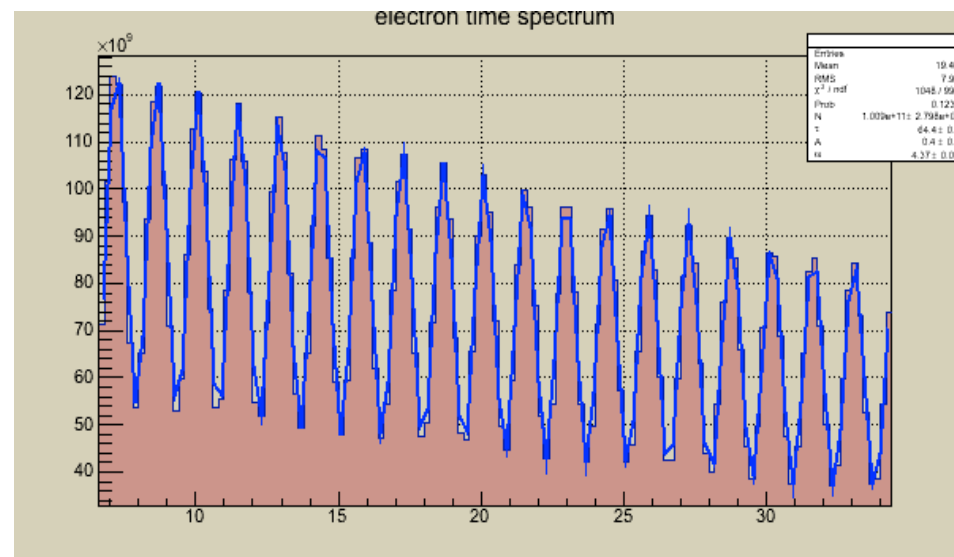
ROOT simulation-III

To verify the goodness of the implemented code and the guesses over the fit function and the generating function, we start setting $A = c_1$

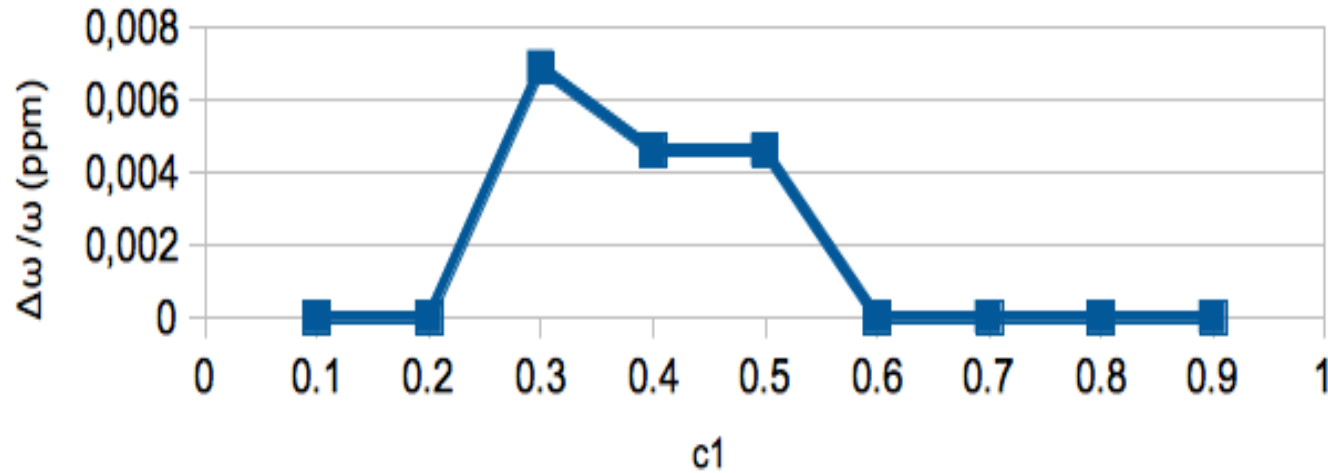


$$c_1 = 0.4$$

$$c_2 = 0$$



ROOT simulation-III



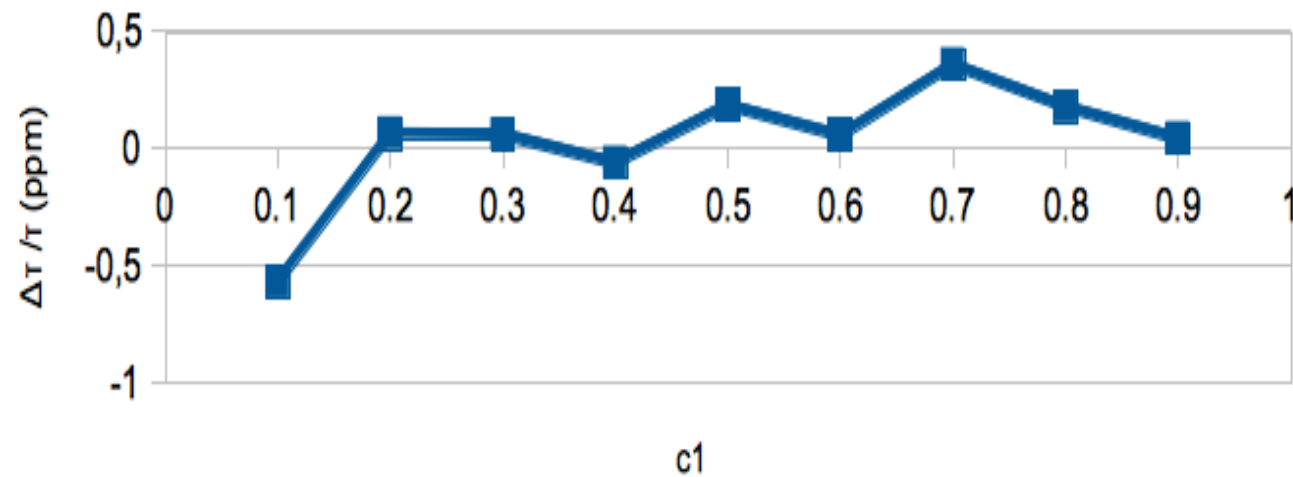
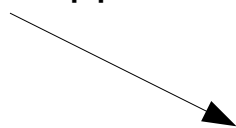
$$0.1 \leq c_1 \leq 0.9$$

$$c_2 = 0$$

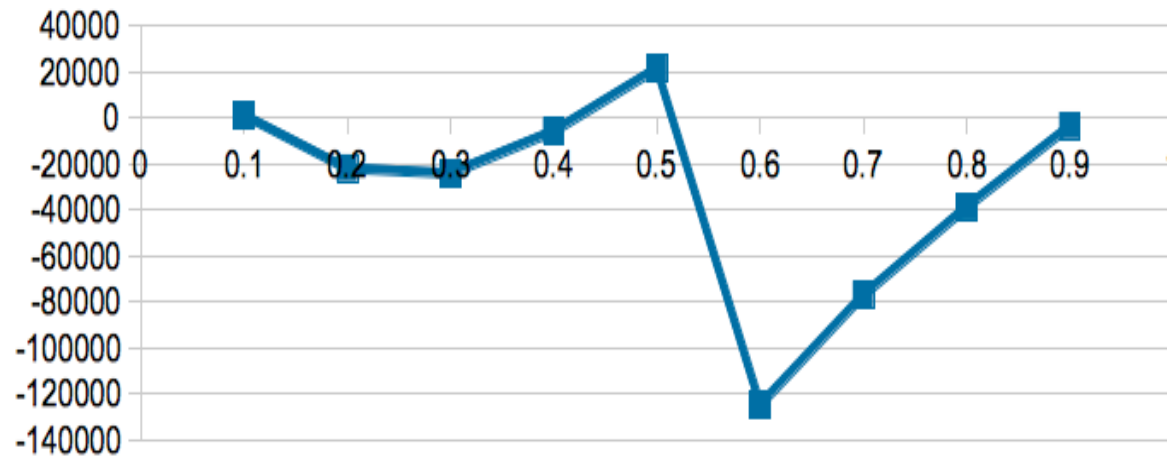
Maximum variation 10^{-3} ppm



Maximum variation 1 ppm

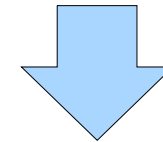


ROOT simulation-IV

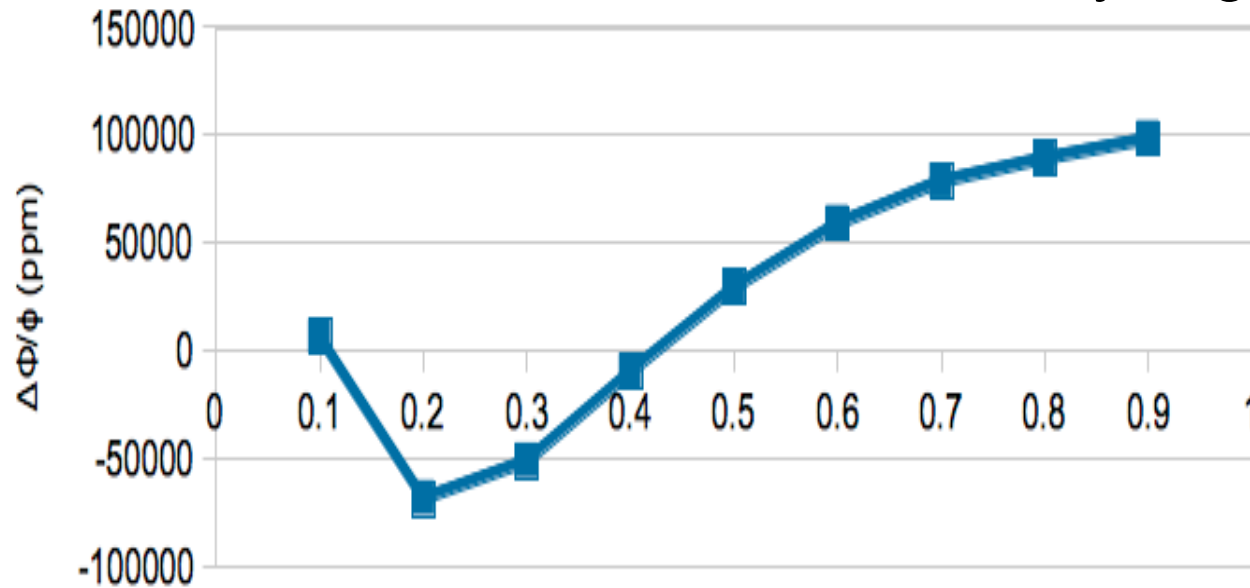


c1

Very large values for $\Delta A/A$ and $\Delta\Phi/\Phi$



There is a problem!



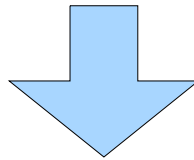
c1

ROOT simulation-V

Is the code or the fit function choice the problem?

To understand this, I have removed parameters in the fit function one at a time and found out it was the phase

Using a fit function and a generating function without the phase also the asymmetry variations are negligible

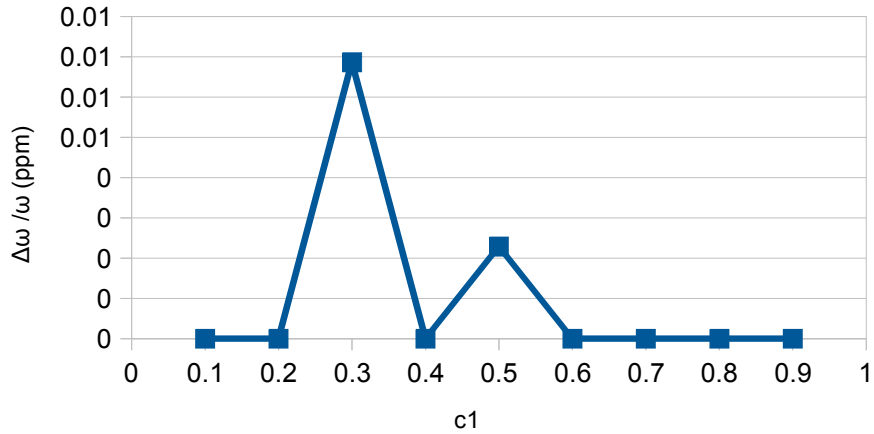


The problem is the phase!

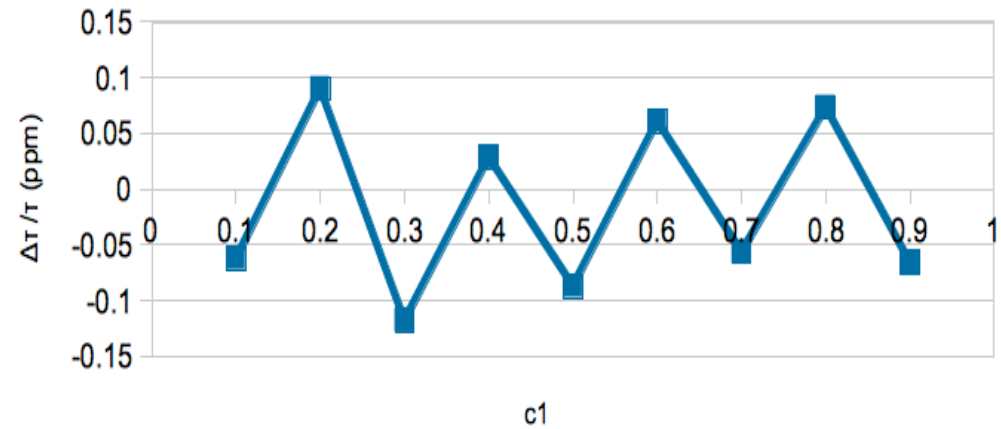
ROOT simulation-VI

$\Delta\omega / \omega$ vs c_1

for $A = c_1 (1 + c_2)$ and $c_2 = 0$

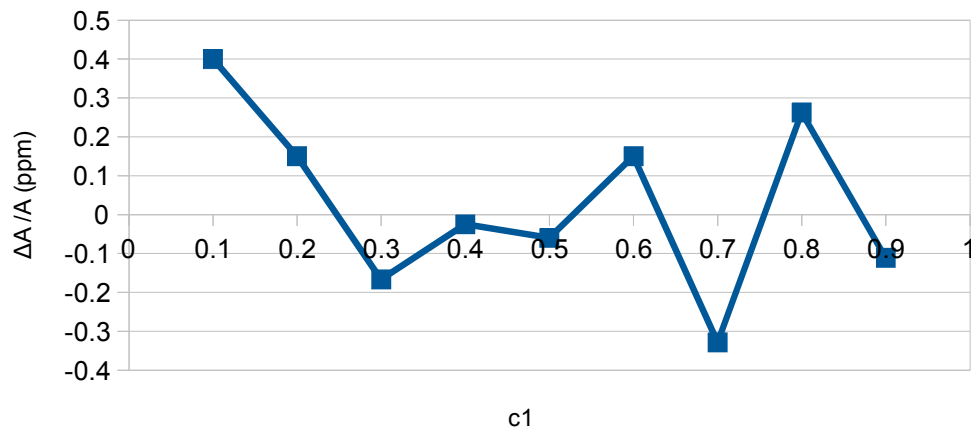


Getting rid of the phase



$\Delta A / A$ vs c_1

for $A = c_1 (1 + c_2)$ and $c_2 = 0$



Conclusions-I

- We found that we could not perform the fit with Φ as a free parameter
- When we add a linear offset to the generated A , we get expected offsets in the measured value of A and we found that this is so
- The extracted value of ω seems to be insensitive to a linear offset in A

Straw Chambers

Straw chambers are basically proportional chambers constructed with a single anode wire centered in an aluminized mylar tube forming the grounded cathode. The cathode is filled with gas at 1÷4 atmospheres.



This mylar tube is several mils thick wrapped with two strips glued together in a barber pole strip fashion

Those we used are 1m long and 5.1 mm wide and filled with Ar-CO₂ .

Straw Chambers-I

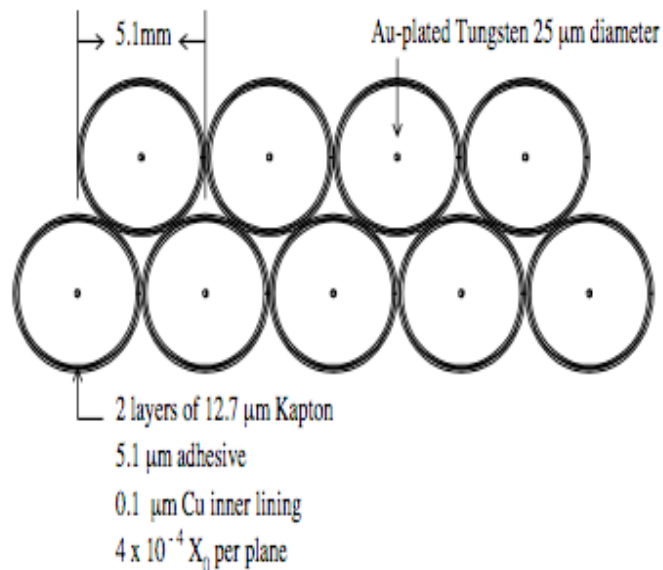
The advantages of a straw chambers when compared to the similar multiwire chambers are:

- The straw detector is inexpensive, robust and relatively simple to construct
- The damage and possible down time caused by wire breakage is minimal since the broken wire is isolated in the tube cell and will only need to be disconnected
- The effects of signal cross talk are minimized as the straw cathode provides a complete ground shield between nearby wires
- The problems of electrostatic alignment distortions are minimal when the anode is kept reasonably centered in the straw.

Straw Chambers-II

There are three main components to performing the g-2 measurement: measurement of the precession frequency, measurement of the magnetic field and measurement of the spacial distribution of muons within the field.

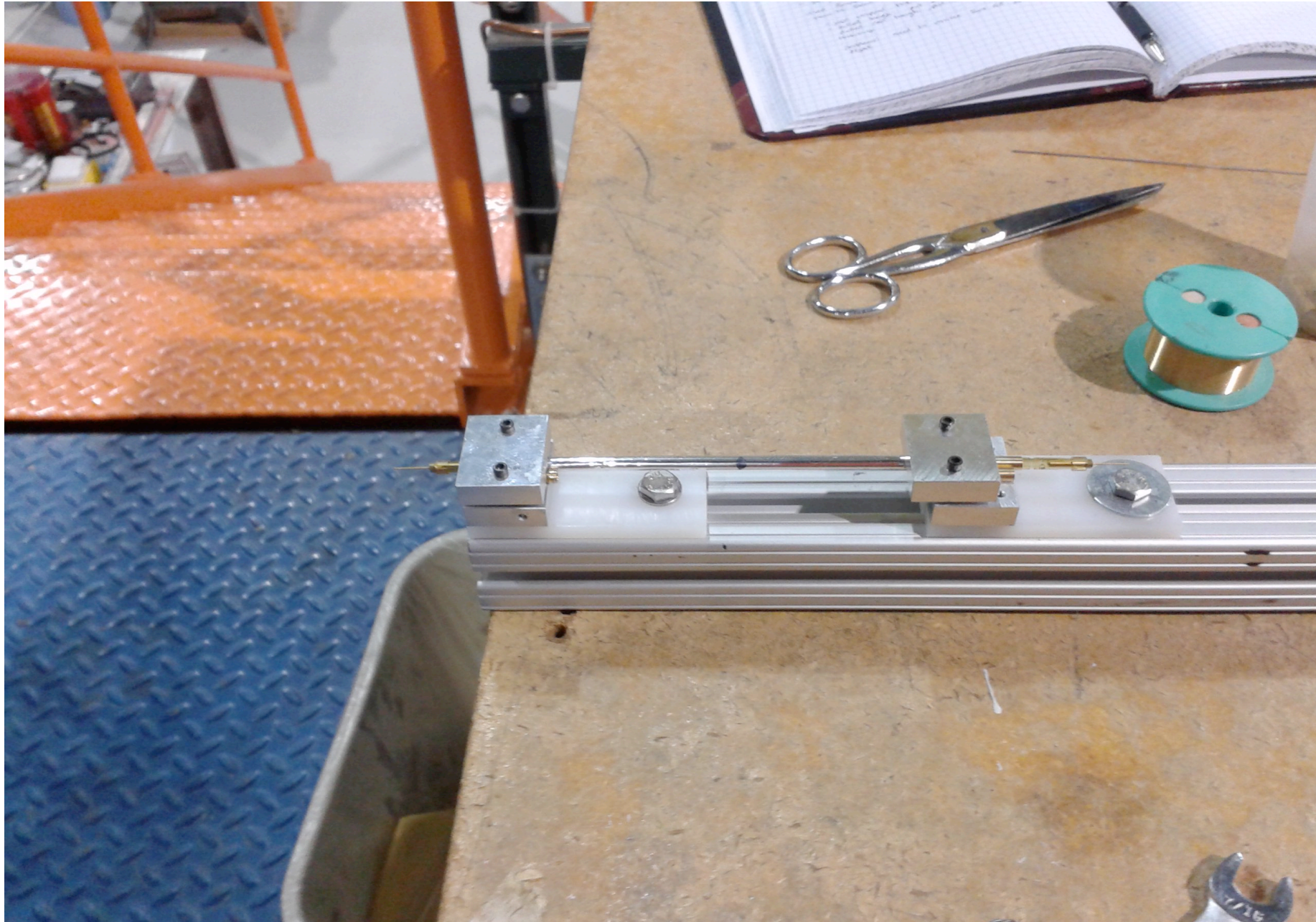
The muon spacial distribution can be mapped by measuring the position trajectories and extrapolating back to the point where the trajectory is tangent to the muon orbit. This is do using straw detectors.



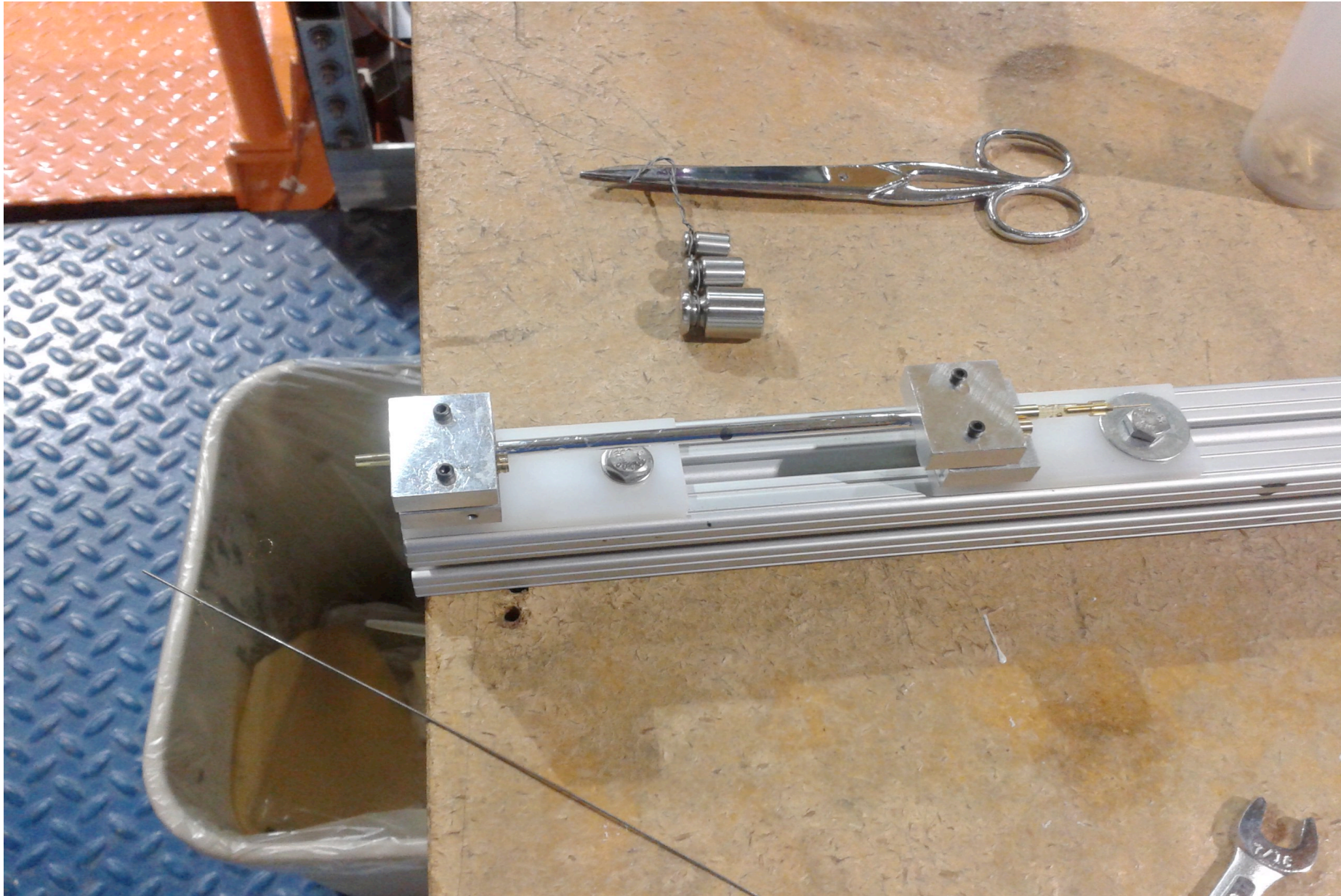
The prototype system consist of 20 straws arranged ad in figure, split equally between vertical and horizontal configuration.

The decay point of the muon is required to be knew within 3 mm in both the radial and vertical position, this led to a requirement of better than 100 μm position resolution per straw.

Straw Prototype Construction



Straw Prototype Construction



ROOT simulation-VII

Step 2: $c_2 \neq 0$

c_1	c_2	τ	A	ω	X^2/dof
0.3	2.000E-009	64.3999944	0.29999999	4.37	0.95
	2.00E-008	64.3999952	0.30000036	4.36999998	0.96746386
	2.00E-007	64.3999951	0.30000268	4.36999999	1.09702811
	2.00E-006	64.3999629	0.30002255	4.37000001	8.63818273
	2.00E-005	64.3999304	0.30022562	4.37000002	739.460843

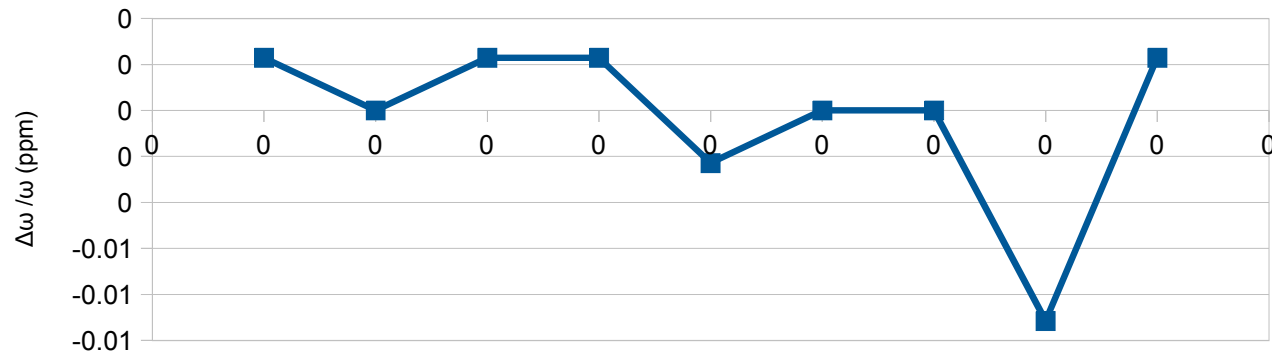
ROOT simulation-VIII

$$c_1 = 0.3$$

$$0.1 \cdot 10^{-9} \leq c_2 \leq 0.9 \cdot 10^{-9}$$

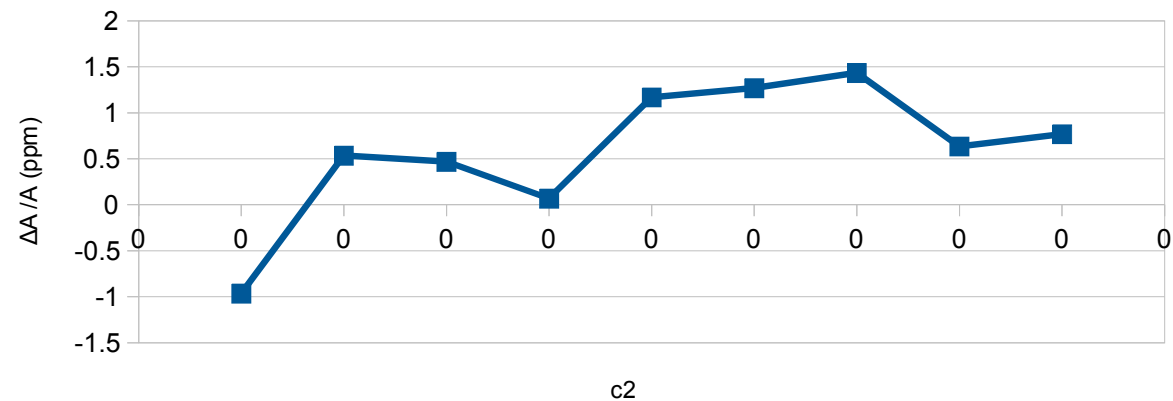
$\Delta\omega / \omega$ vs c_2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



$\Delta A / A$ vs c_2

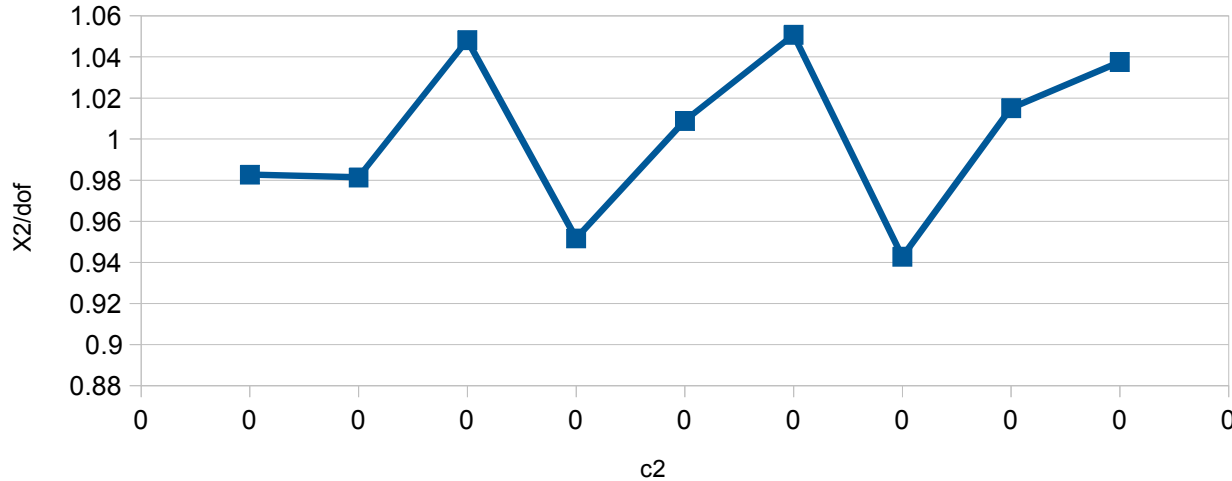
for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



ROOT simulation-IX

X2/dof vs c2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$

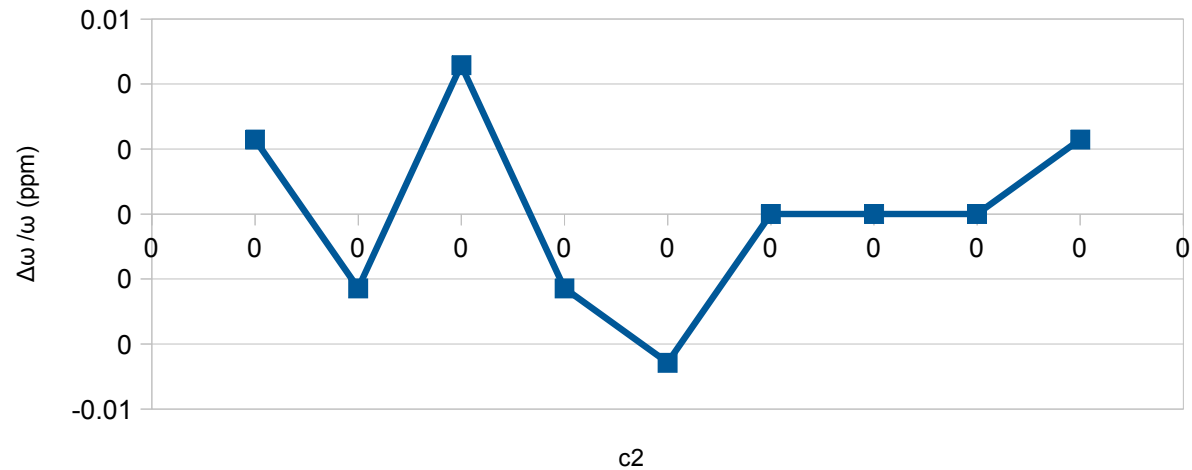


$$c_1 = 0.3$$

$$0.1 \cdot 10^{-8} \leq c_2 \leq 0.9 \cdot 10^{-8}$$

$\Delta\omega / \omega$ vs c2

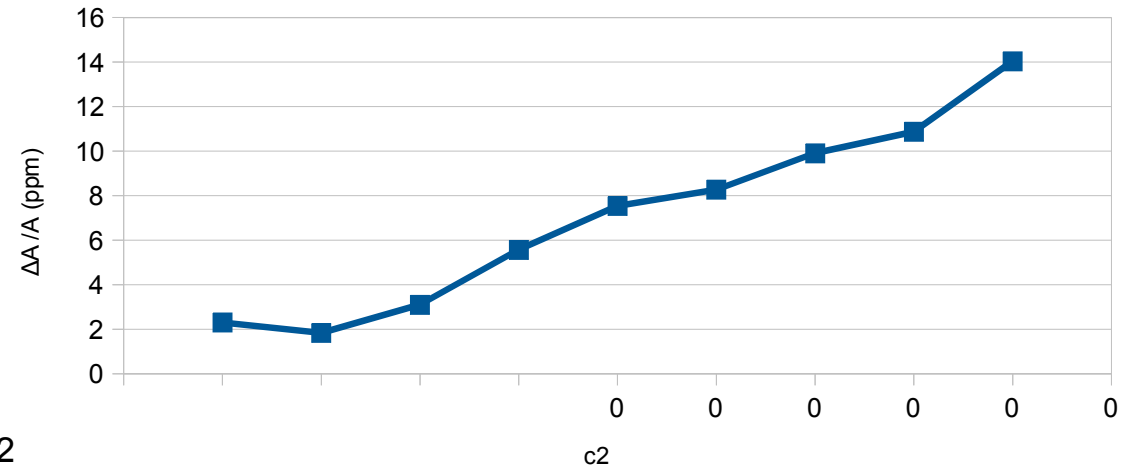
for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



ROOT simulation-X

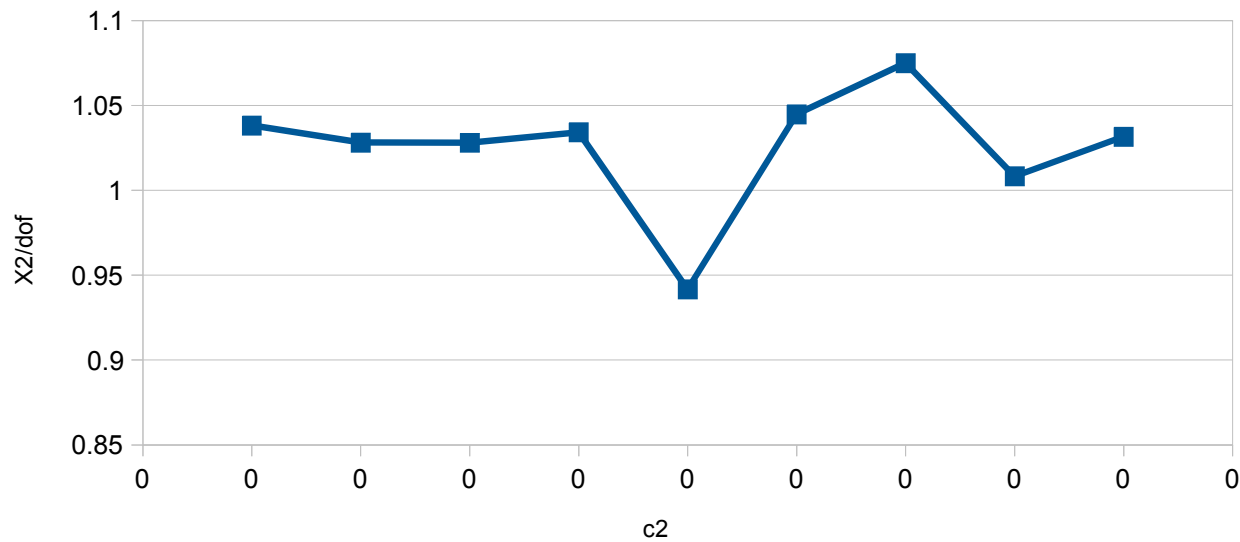
$\Delta A / A$ vs c_2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



X^2/dof vs c_2

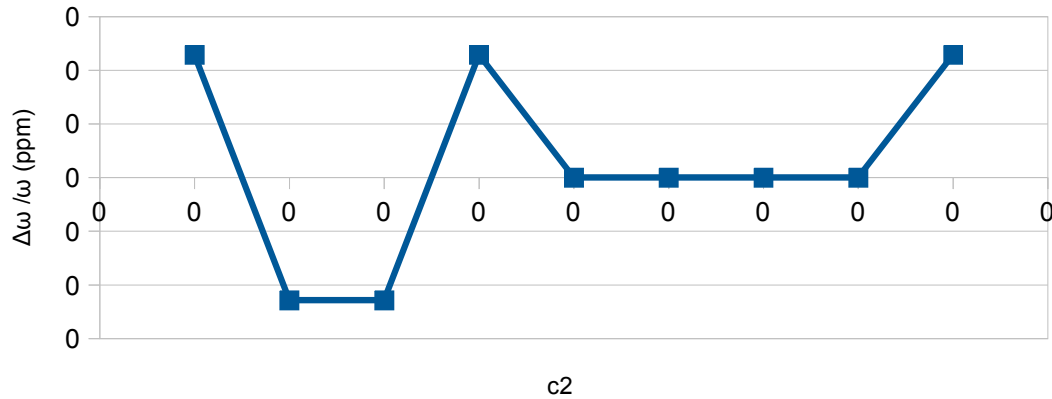
for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



ROOT simulation-XI

$\Delta\omega / \omega$ vs c_2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$

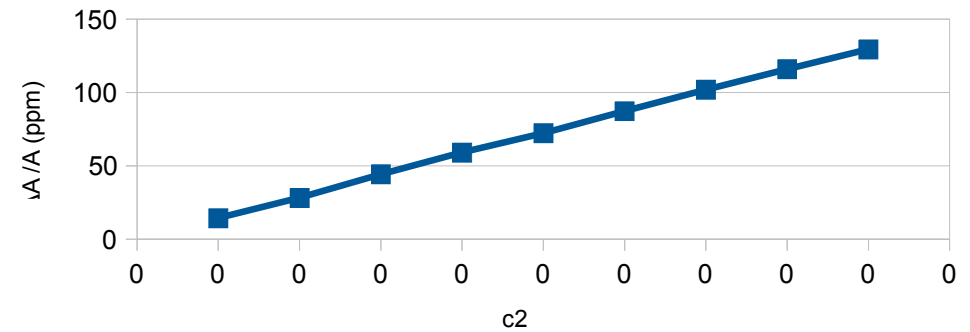


$$c_1 = 0.3$$

$$0.1 \cdot 10^{-7} \leq c_2 \leq 0.9 \cdot 10^{-7}$$

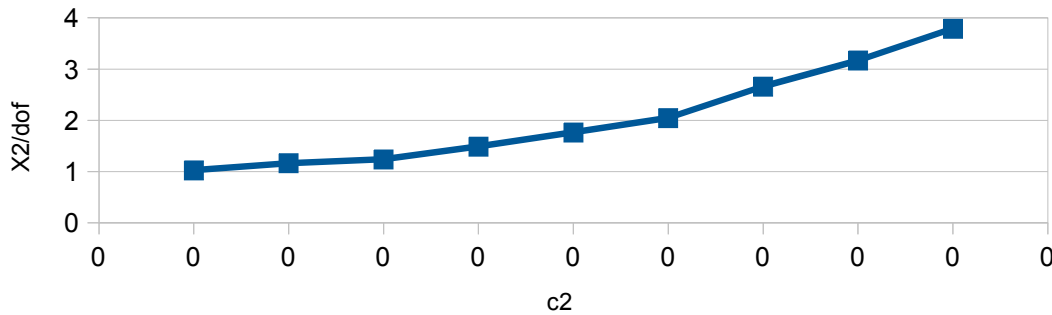
$\Delta A / A$ vs c_2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



X_2/dof vs c_2

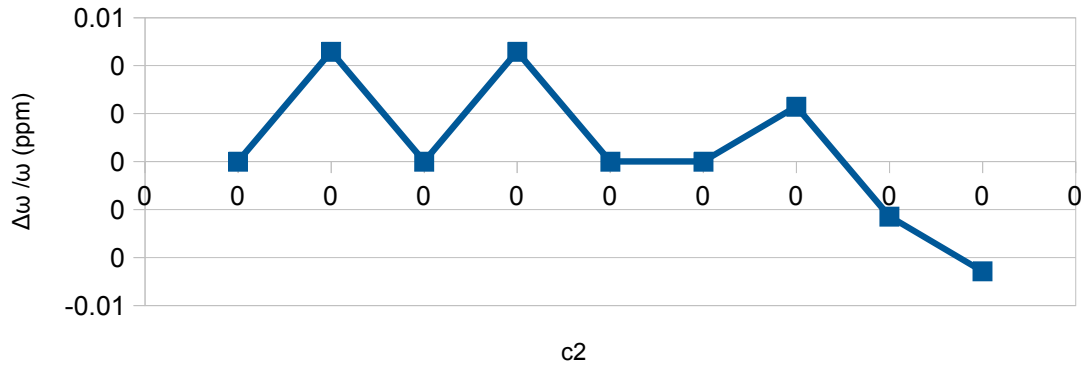
for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



ROOT simulation-XII

$\Delta\omega / \omega$ vs c_2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$

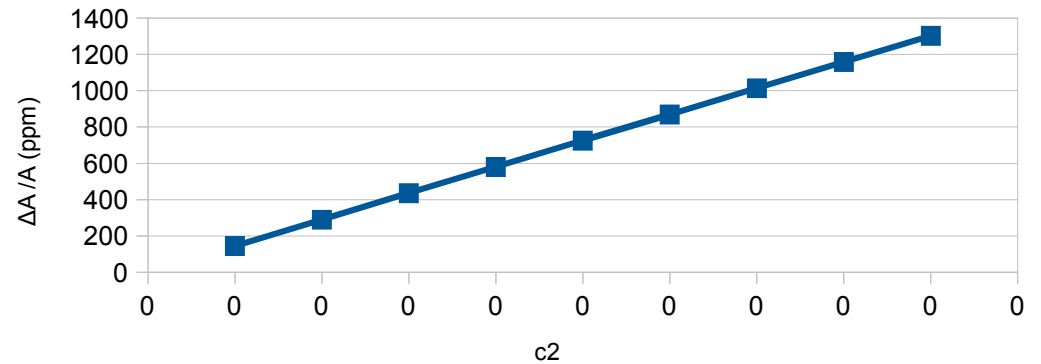


$$c_1 = 0.3$$

$$0.1 \cdot 10^{-6} \leq c_2 \leq 0.9 \cdot 10^{-6}$$

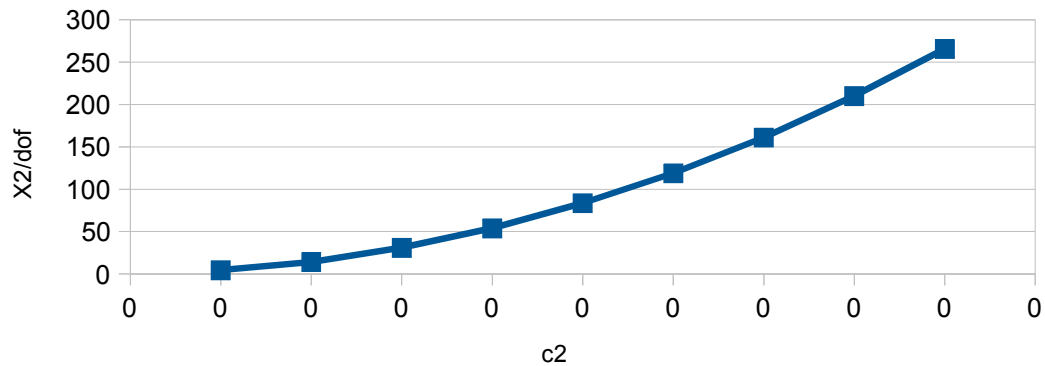
$\Delta A / A$ vs c_2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$

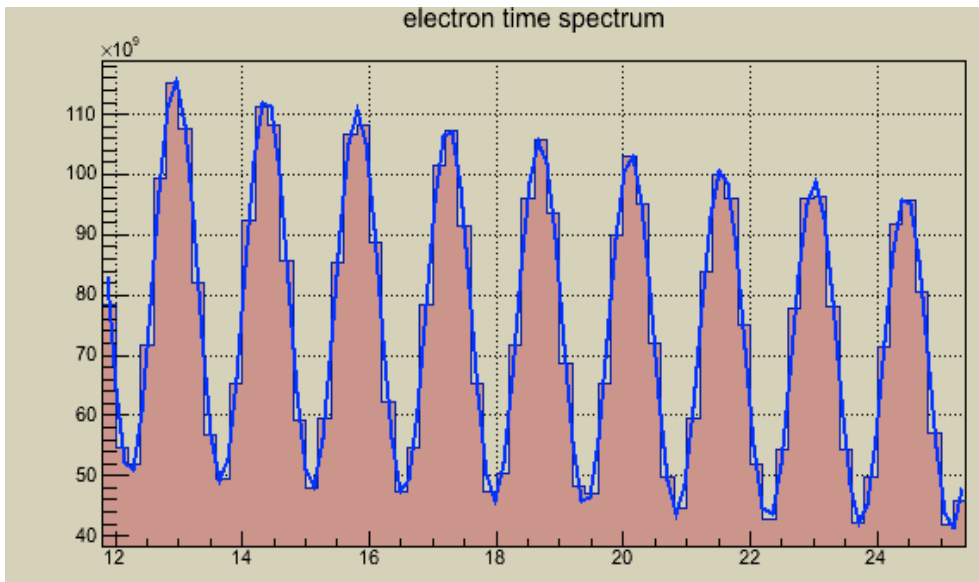
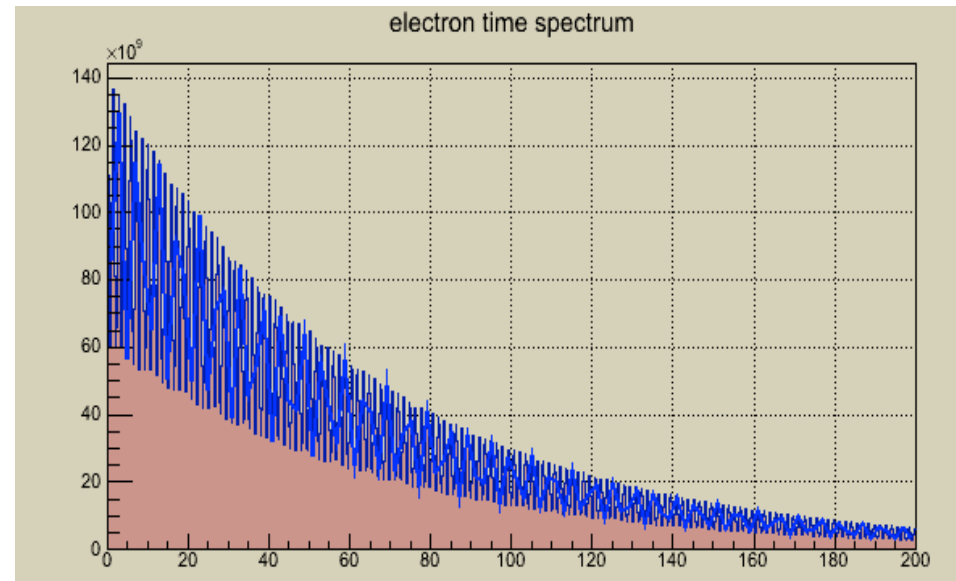
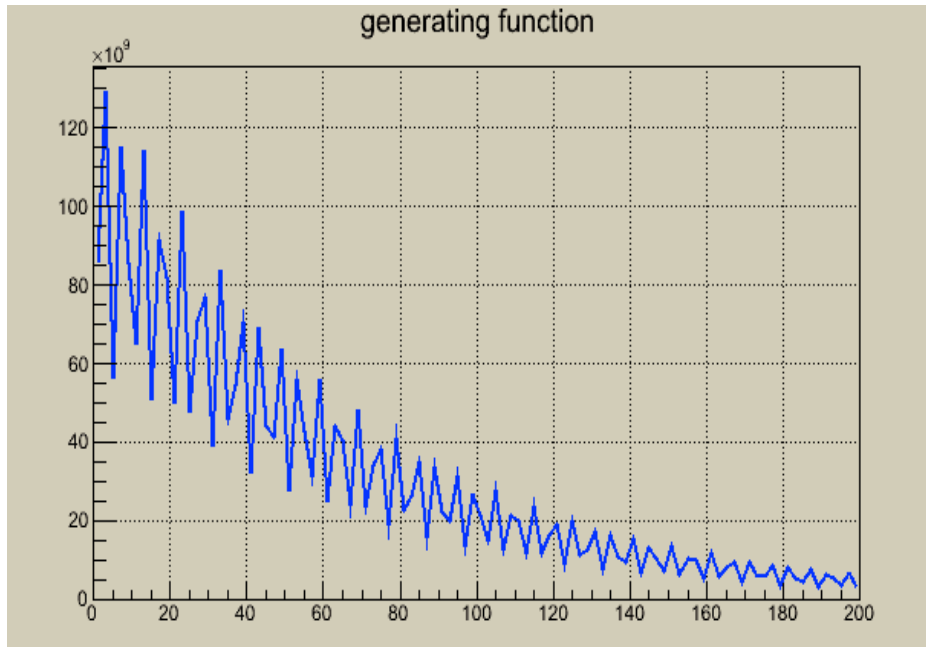


X^2/dof vs c_2

for $A = c_1 (1 + c_2 t)$ and $c_1 = 0.3$



ROOT simulation-XIII



Conclusions-I

- We found that we could not perform the fit with Φ as a free parameter
- When we add a linear offset to the generated A, we get expected offsets in the measured value of A and we found that this is so
- The extracted value of omega seems to be insensitive to a linear offset in A
- The χ^2 of the fit gets very bad very quickly ($c2 > 5 \times 10^{-7}$)
We are not sure how to translate this into a systematic error on omega.
Might be more of a hard specification.