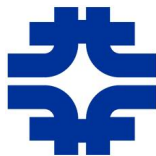


HEP Top Tagger for the search of stops

Giulio Dujany

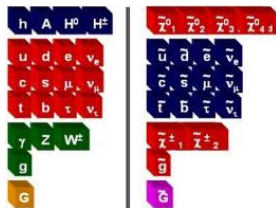
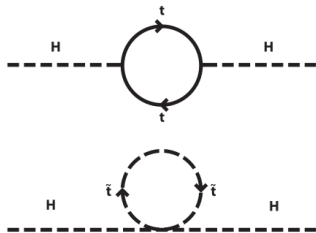
Summer intern 2012
Supervised by: Daryl Hare,
Rick Cavanaugh, Chris Silkworth

26 September 2012



Stop

SUSY can solve the Hierarchy problem



Stop should be the lightest squark

- Top is the heaviest fermion: his loop contribute the most to Higgs' mass correction.
- In MSSM the running masses of squarks are degenerate at high scale.
- Due to large Yukawa coupling stop mass run down at lower energies.

$$pp \rightarrow \tilde{t}\tilde{t}^* \rightarrow (t\chi_1^0)(\bar{t}\chi_1^0) \rightarrow (bjj\chi_1^0)(\bar{b}jj\chi_1^0)$$

Fully Hadronic Decay Channel

Advantages

- Fully reconstructs top kinematics (no \cancel{E}_T from neutrinos)
- \cancel{E}_T due only to neutralinos
- Higher branching ratio (color factor)

Challenges

Needs an efficient identification of boosted tops



We expect a moderate stop mass

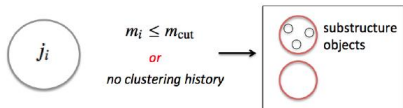


Rate for highly boosted tops in the final state is significantly lower than the overall signal rate

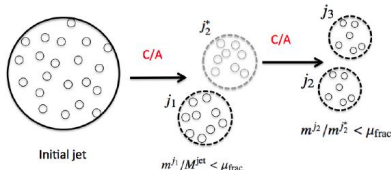


We need a top-tagger for moderately boosted tops

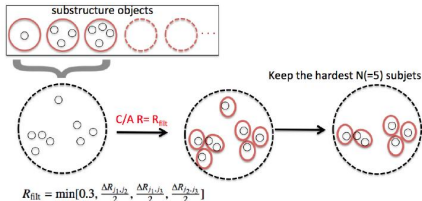
HEP Top Tagger: Top Candidate



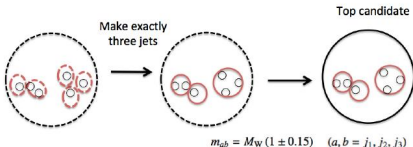
(a) Every object encountered in the de-clustering process is considered a 'substructure object' if it is of sufficiently low mass or has no clustering history.



(b) The mass-drop criterion is applied iteratively, following the highest subject-mass line through the clustering history, resulting in N_i substructure objects.



(c) For every triplet-wise combination of the substructure objects, recluster into subjects and select the N_{subject} leading- p_T subjects, with $3 \leq N_{\text{subject}} \leq N_i$ (here, $N_{\text{subject}} = 5$).



(d) Recluster the constituents of the N_{subject} subjects into exactly three subjects to make the top candidate for this triplet-wise combination of substructure objects.

HEP Top Tagger: Selection requirements

pt cut: combined $p_T > 200$ GeV.

top mass cut: $m_{top} - 25 \text{ GeV} < m_{123} < m_{top} + 25 \text{ GeV}$.

di-mass cut: One of the three:

- $0.2 < \arctan \frac{m_{13}}{m_{12}} < 1.3$ and $R_{min} < \frac{m_{23}}{m_{123}} < R_{max}$

- $R_{min}^2 \left(1 + \left(\frac{m_{13}}{m_{12}} \right)^2 \right) < 1 - \left(\frac{m_{23}}{m_{123}} \right)^2 <$

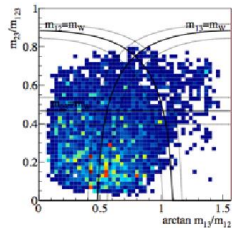
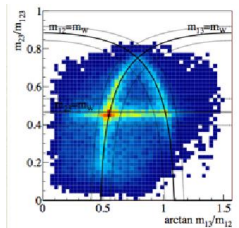
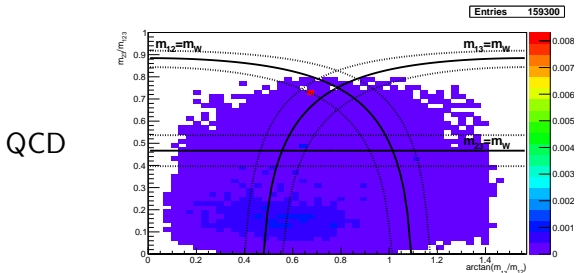
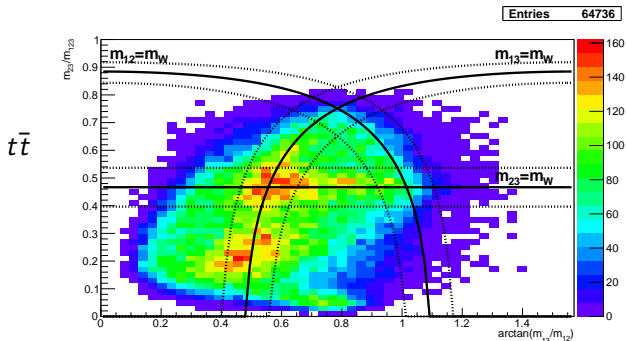
- $R_{max}^2 \left(1 + \left(\frac{m_{13}}{m_{12}} \right)^2 \right)$ and $\frac{m_{23}}{m_{123}} > 0.35$

- $R_{min}^2 \left(1 + \left(\frac{m_{12}}{m_{13}} \right)^2 \right) < 1 - \left(\frac{m_{23}}{m_{123}} \right)^2 <$

- $R_{max}^2 \left(1 + \left(\frac{m_{12}}{m_{13}} \right)^2 \right)$ and $\frac{m_{23}}{m_{123}} > 0.35$

where $R_{min} = 85\% m_W / m_{top}$ and $R_{max} = 115\% m_W / m_{top}$.

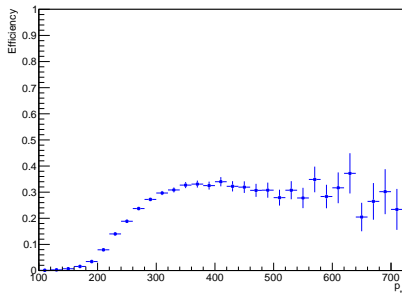
Di-mass plots



Plehn, Spannowsky,
Takeuchi and Zerwas
arXiv 1006.2833

Efficiency and fake rate

Efficiency



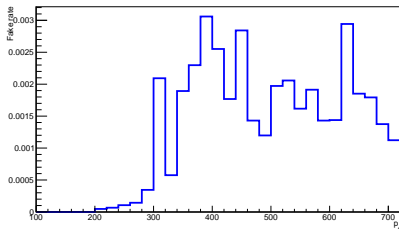
$$\text{Efficiency} = \frac{\# \text{ tagged tops}}{\# \text{ generated tops}}$$

From a $t\bar{t}$ Monte Carlo sample with 450k events.

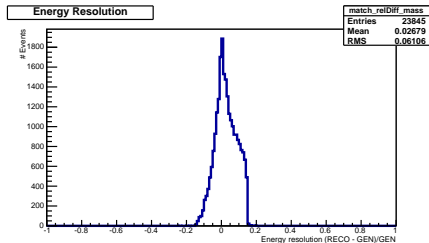
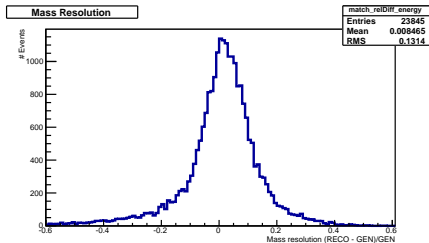
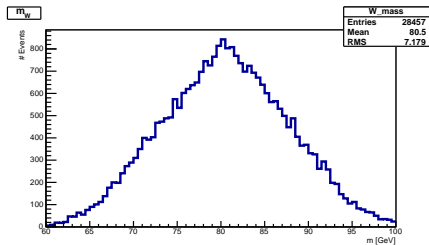
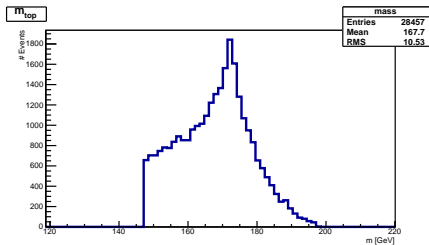
$$\text{Fake rate} = \frac{\# \text{ mistagged jets}}{\# \text{ ac1.5 jets}}$$

From a weighted QCD Monte Carlo sample with 750k events.

Fake rate

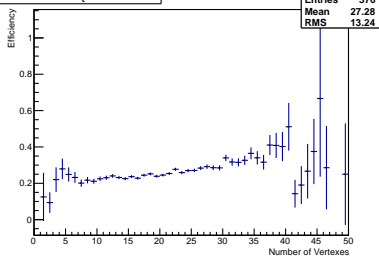


Performance

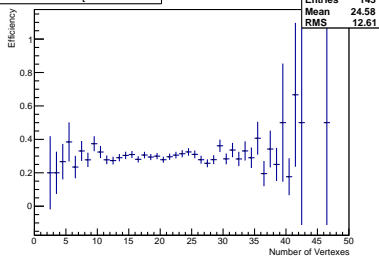


Efficiency pile-up dependence

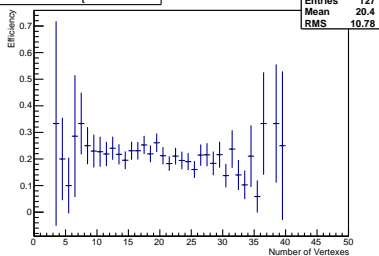
200 GeV < p_t < 300 GeV



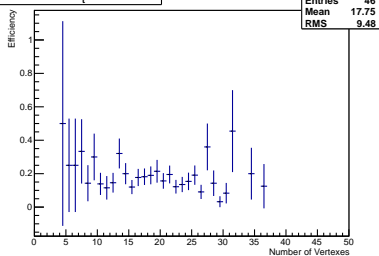
300 GeV < p_t < 400 GeV



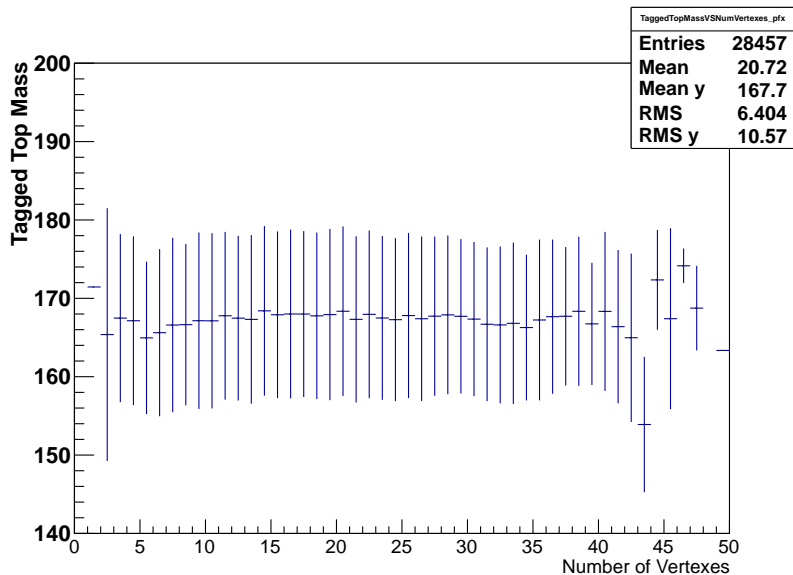
400 GeV < p_t < 500 GeV



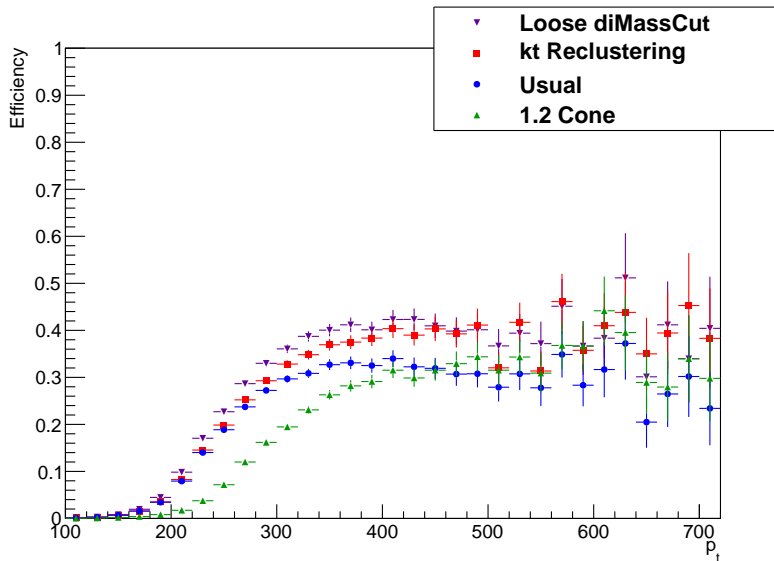
500 GeV < p_t < 600 GeV



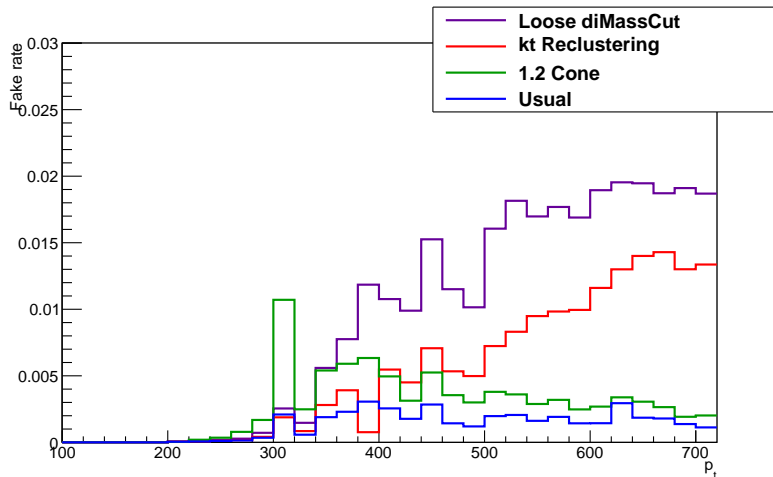
Top mass pile-up dependence



Tuning the HEP Top Tagger: Efficiency

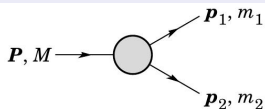


Tuning the HEP Top Tagger: Fake rate



m_T and m_{T2} : definitions

m_T

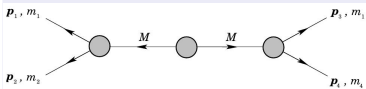


- Used for semi-invisible final state: m_1 visible, m_2 invisible (eg. a neutrino)
- Its distribution can be used to evaluate M . (Used to measure W mass)

$$m_T^2(p_{1T}, m_1, p_{2T}, m_2) = m_1^2 + m_2^2 + 2(E_{1T}E_{2T} - p_{1T}p_{2T})$$

$$\text{where } E_{1T} = \sqrt{m_1^2 + p_{1T}^2}$$

m_{T2}



Used when a pair of particle decays semi-invisibly: particles 1 and 3 are invisible (LSP) while particles 2 and 4 are visible ($t\bar{t}$).

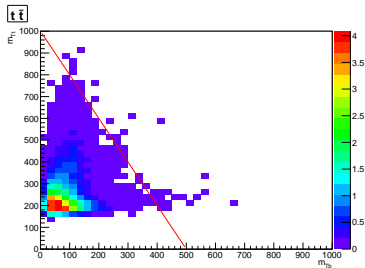
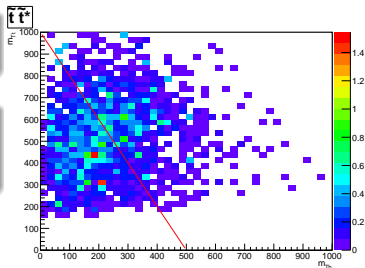
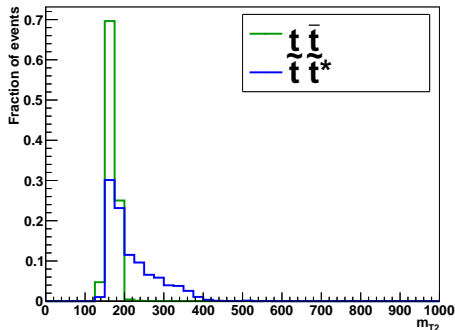
$$m_{T2}^2(p_{2T}, m_2, p_{4T}, m_4, \not{p}_T, m_1) = \min_{\not{A}_T^{(1)} + \not{A}_T^{(2)} = \not{p}_T} [\max \{ m_T^2(p_{2T}, m_2, \not{p}_T, m_1), m_T^2(p_{4T}, m_4, \not{p}_T, m_1) \}]$$

m_T and m_{T_2} : signal selection

- $m_{T_t} = m_T(p_t, m_t, \cancel{E}_T, 0)$
- $m_{T_b} = m_T(p_b, m_b, \cancel{E}_T, 0)$

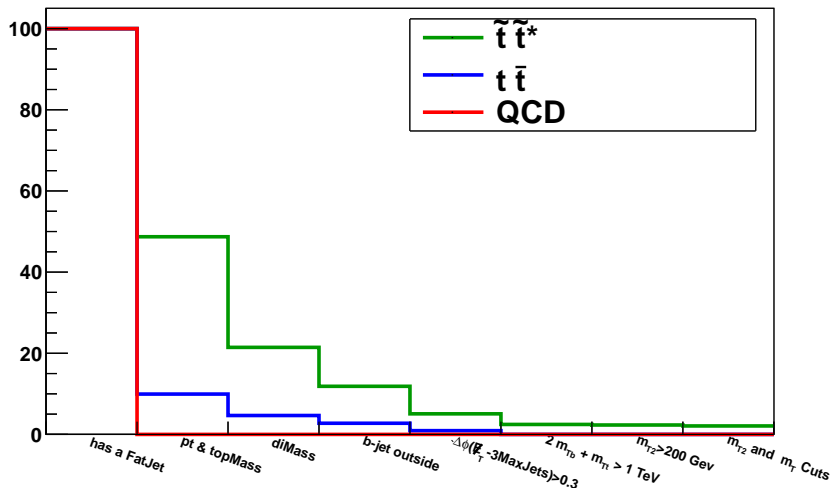
Cuts

- $m_{T_2} > 200$ GeV
- $2m_{T_b} + m_{T_t} > 1$ TeV



Analysis: a first glimpse

CutFlow

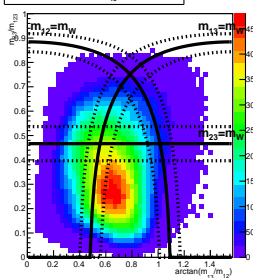


Conclusions

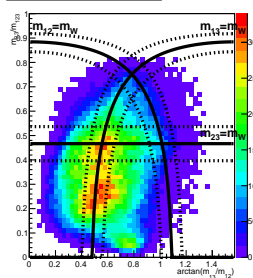
- The Hep Top Tagger is now implemented in the CMS software.
- The tagger's Efficiency and Fake rate were evaluated from Monte Carlo:
 - **efficiency:** between 20% and 40%
 - **fake rate:** on the order of 0.5%
- Pile-up has little effect on the tagger's performance.
- Looked into several possible improvements.
- We have made modest progression towards a new baseline search for stops.

BACKUP

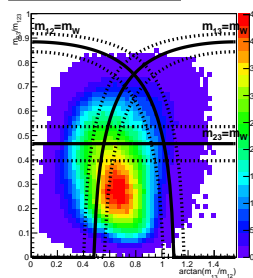
m_{22}/m_{122} over $\arctan(m_{13}/m_{12})$ No Cuts Entries 1551954



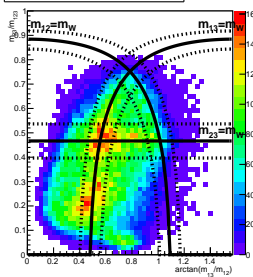
Matched Fat Jets No Cuts Entries 115716



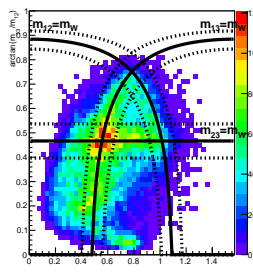
Unmatched Fat Jets No Cuts Entries 1436236



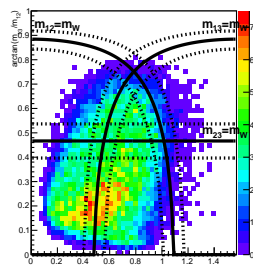
m_{22}/m_{122} over $\arctan(m_{13}/m_{12})$ With Cuts Entries 64736

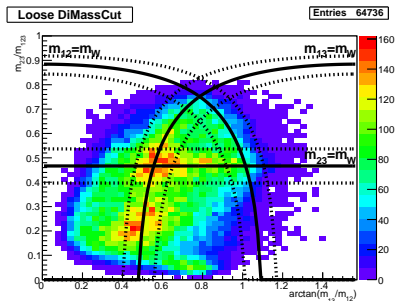
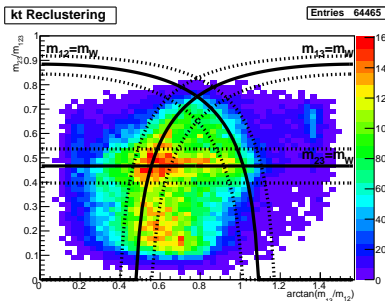
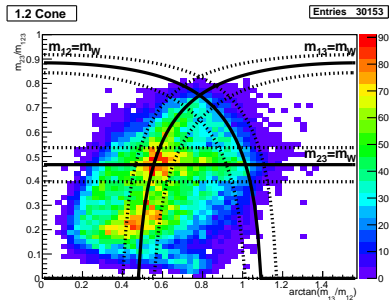
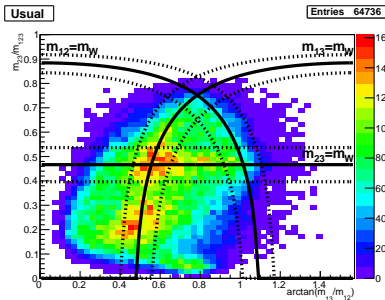


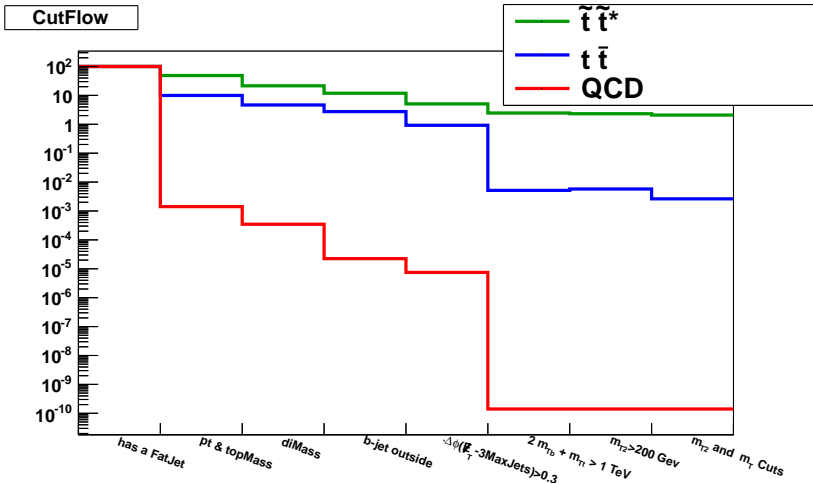
Matched Fat Jets With Cuts Entries 36677



Unmatched Fat Jets With Cuts Entries 28059







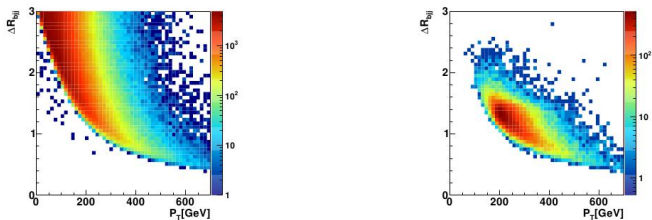
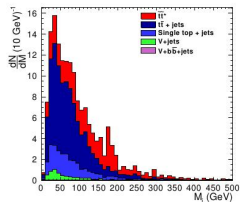


Figure 2: Left: partonic ΔR_{bjj} vs p_T distribution for a Standard Model $t\bar{t}$ sample. Right: the same correlation, but only for tagged top quarks and based on the reconstructed kinematic properties.

Background

- $t\bar{t}$
 - ▶ Where a t decays in $b + \text{a hadronic } \tau + \text{a hard neutrino}$.
 - ▶ Plus \cancel{E}_T due to poor jet reconstruction .
 - ▶ Where a t decays leptonically producing a hard neutrino and the lepton veto fail.
- $Z + \text{jets}$ where the Z decays invisibly.
- QCD.



$(m_{\bar{t}}, m_{\chi_1^0}) = (340 \text{ GeV}, 0 \text{ GeV})$ at $\sqrt{s} = 7$
TeV and $\mathcal{L} = 5 \text{ fb}^{-1}$

Signal selection

- Lepton veto.
- $\cancel{E}_T > 175 \text{ GeV}$.
- A top-tagged fat-jet.
- Outside the fat jet a b-tagged jet.

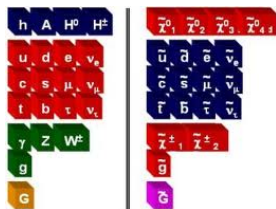
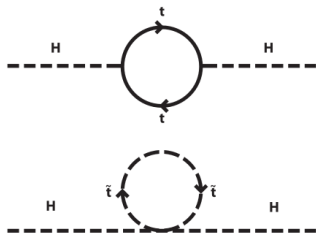
HEP Top Tagger

- Start with a Cambridge/Aachen Fat-Jet with $\Delta R = 1.5$ and undo step by step the clustering algorithm to look for subjets.
- When the mother jet becomes two daughters keep both if the larger mass is less than 80% of the mother mass. Otherwise keep the daughter with larger mass.
Continue as long subjets' mass > 30 GeV. (To cut soft QCD)
- Filter subjets with $\Delta R_{filter} = \min(0.3, \Delta R_{ij}/2)$.
- Keep the five filtered subjets with highest p_T .
- Select the set of three subjets whose combined mass is closest to m_{top} .
- Sort the three selected subjets by p_T (j_1, j_2, j_3) and calculate combined invariant masses m_{12}, m_{13}, m_{23} .
- Check selection requirements.

Hierarchy problem

The SM Higgs is a fundamental scalar \implies his mass is quadratically divergent (with momentum cut off)

- Fine tuning (not “natural”)
- Physics BSM that cut off quadratically divergent loop contribution to Higgs mass
 - ▶ SUSY
 - ▶ Little Higgs
 - ▶ ...



Stop

Top quark is the fermion with the highest Yukawa coupling



His loop is the one that contribute the most to the Higgs' mass correction.

- In MSSM the running masses of squarks are degenerate at high scale
- Due to large Yukawa coupling stop mass run down at lower energies



Stop should be the lightest squark

Stop's decay channels

$$pp \rightarrow \tilde{t}\tilde{t}^*$$

Supposing R-parity, stops at LHC should be pair produced.

$$\tilde{t} \rightarrow t\chi_1^0$$

We suppose stop to be sufficiently heavy to decay into a top and a neutralino but not enough to kinematically allows his decay to a top and a gluino.

Other analyses, to avoid QCD background, ask the top to decay semileptonically:

$$pp \rightarrow \tilde{t}\tilde{t}^* \rightarrow (t\chi_1^0)(\bar{t}\chi_1^0) \rightarrow (bl^+\nu\chi_1^0)(\bar{b}jj\chi_1^0) + (\bar{b}l^-\bar{\nu}\chi_1^0)(bjj\chi_1^0)$$

We, complementary, look for a fully hadronic decay:

$$pp \rightarrow \tilde{t}\tilde{t}^* \rightarrow (t\chi_1^0)(\bar{t}\chi_1^0) \rightarrow (bjj\chi_1^0)(\bar{b}jj\chi_1^0)$$