

Département de physique nucléaire et corpusculaire









Picosecond Time Stamping in Fully-Monolithic Highly-Granular Pixel Detectors

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10th October 2025 - FATA 2025 - INFN Catania

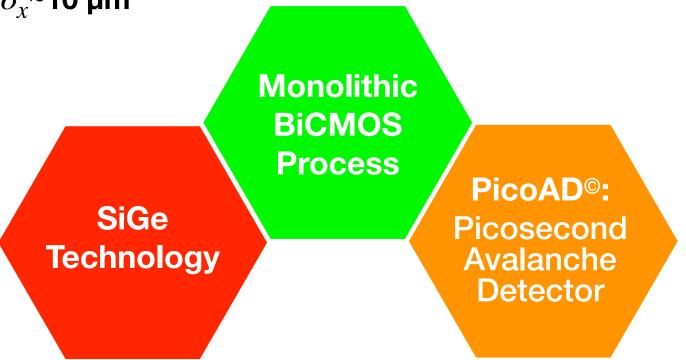


The MONOLITH ERC Project



- We want to measure charged particles with:
 - time resolution $\sigma_{\rm t}$ ~10 ps

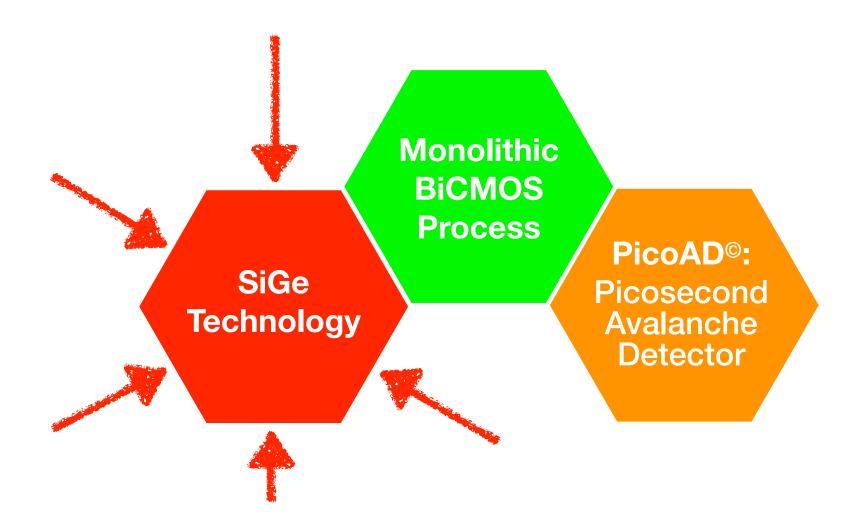
• space resolution σ_{χ} ~10 µm





SiGe Technology





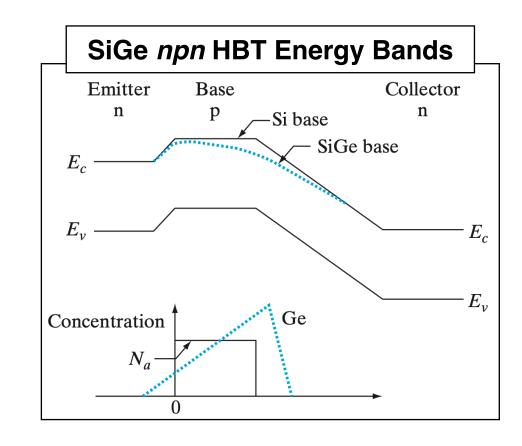


SiGe Technology



- A grading of Germanium creates a small electric field in the base of a Bipolar Junction Transistor (BJT):
 - The charges move via drift
 - The transit time in the base is shorter
 - Thicker base than a BJT
 - Smaller base resistance R_b than a BJT

$$\sigma_{\text{Jitter}} = \frac{\sigma_V}{dV/dt} \propto C_{\text{tot}} \sqrt{\frac{k_1}{\beta} + k_2 \mathbf{R}_{\text{b}}}$$





_eading-edge IHP SG13G2 technology: 130 nm process featuring SiGe HBT



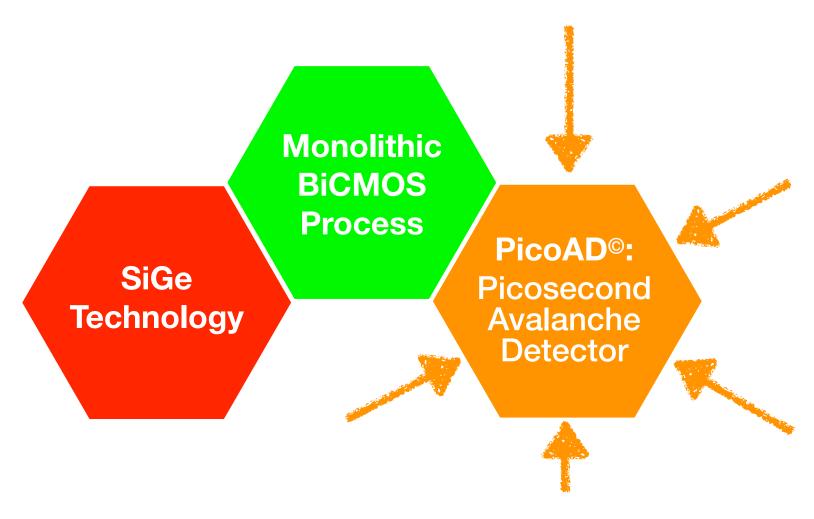
nucléaire et corpusculaire

PicoAD[©] Sensor Concept





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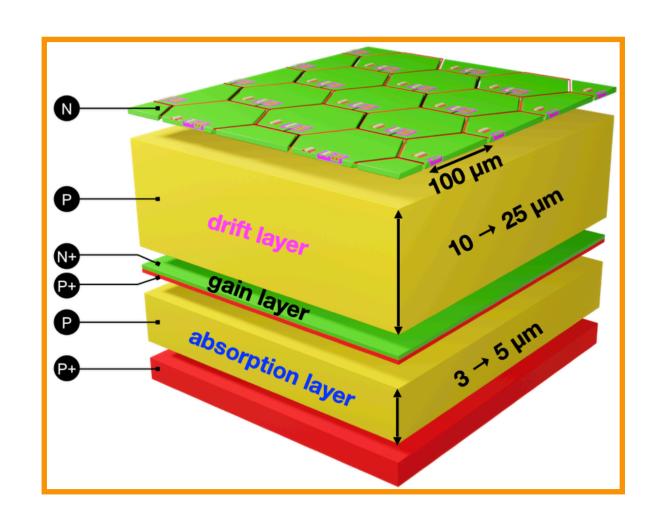




PicoAD© Sensor Concept



- Multi-Junction Picosecond Avalanche Detector[©]
- Continuous and deep gain layer:
 - Decorrelation from implant size
 - High pixel granularity
 - Better spatial resolution
 - Only a small fraction of the charge gets amplified
 - Reduced charge-collection noise $\sigma_{\rm Landau}$

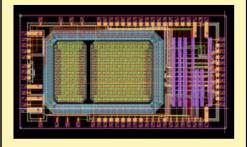




Timing Prototypes without Gain Layer





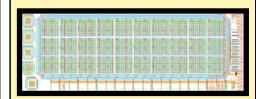


- Square pixel 900 μm
- Discriminated output

200 ps

JINST 13 P04015

2017

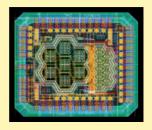


- Square pixels 500 µm
- 100 ps TDC + I/O logic

110 ps

JINST 14 P07013 JINST 14 P02009

2018

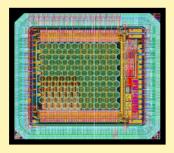


- Hexagonal pixels 130 µm
- Discriminated outputs

45 ps

JINST 14 P11008 JINST 15 P11025

2020

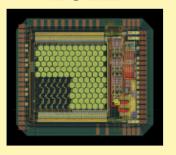


- Hexagonal pixels 65 µm
- 30 ps TDC + I/O logic
- 4 analog outputs

36 ps

JINST 17 P02019

2022



- Hexagonal pixels 65 µm
- Improved frontend
- 4 analog outputs

20 ps

<u>JINST 18 P03047 JINST 19 P01014</u> JINST 19 P04029 JINST 19 P07036

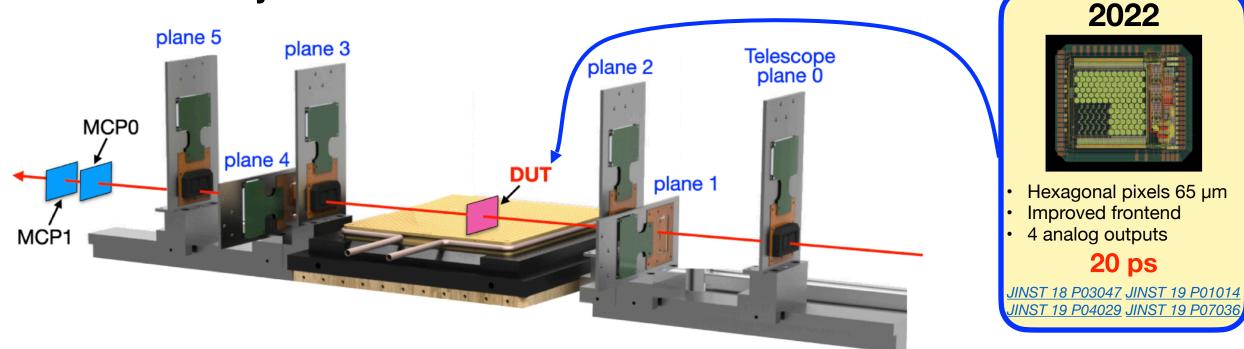
In this presentation



Test Beam: Experimental Setup



- October 2022: SPS Testbeam with 180 GeV/c pions (MIP)
- Measure efficiency and time resolution



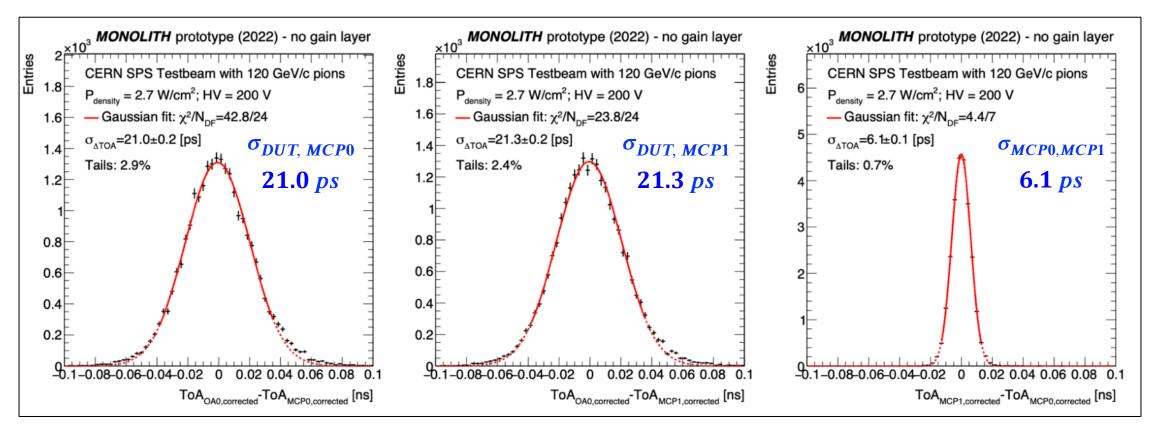
- UNIGE FE-I4 telescope to provide the spatial information ($\sigma_{x,v} \sim 10 \mu m$)
- Two PMT-MCPs ($\sigma_t \sim 5$ ps) to provide the timing reference





Time Resolution Distributions





- Very Gaussian distributions after time walk correction
- Simultaneous fit to extract the time resolution of DUT, MCP0, MCP1:

$$\begin{cases} \sigma_{\text{DUT,MCP0}}^2 = \sigma_{\text{DUT}}^2 + \sigma_{\text{MCP0}}^2 \\ \sigma_{\text{DUT,MCP1}}^2 = \sigma_{\text{DUT}}^2 + \sigma_{\text{MCP1}}^2 \\ \sigma_{\text{MCP0,MCP1}}^2 = \sigma_{\text{MCP0}}^2 + \sigma_{\text{MCP1}}^2 \end{cases}$$

MCP0:
$$\sigma_t = (3.6 \pm 1.5)$$
 ps

MCP1:
$$\sigma_t = (5.0 \pm 1.1)$$
 ps

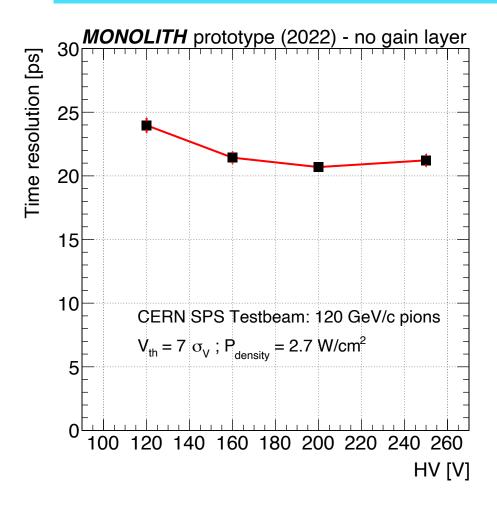
DUT: $\sigma_t = (20.1 \pm 0.3) \text{ ps}$





Time Resolution Results



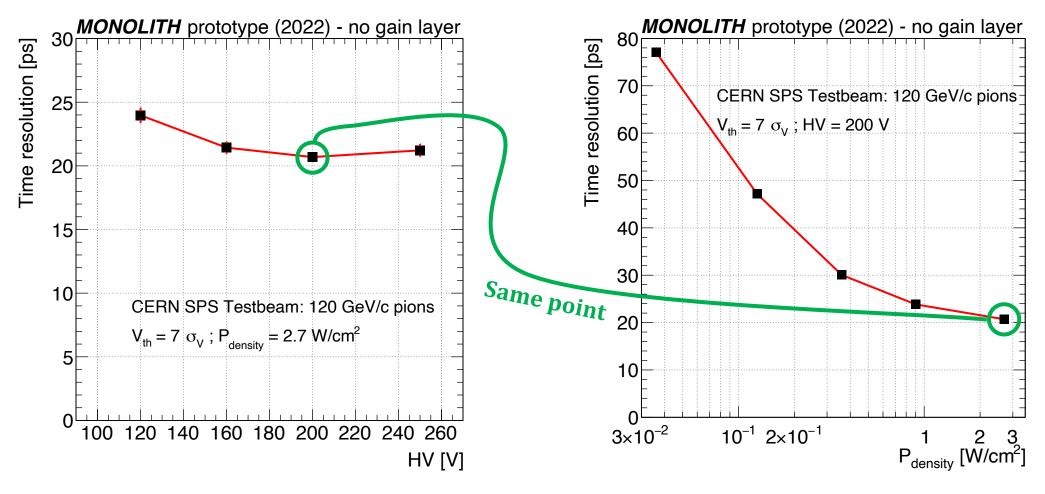


Large plateau of 100 V with ~20 ps



Time Resolution Results





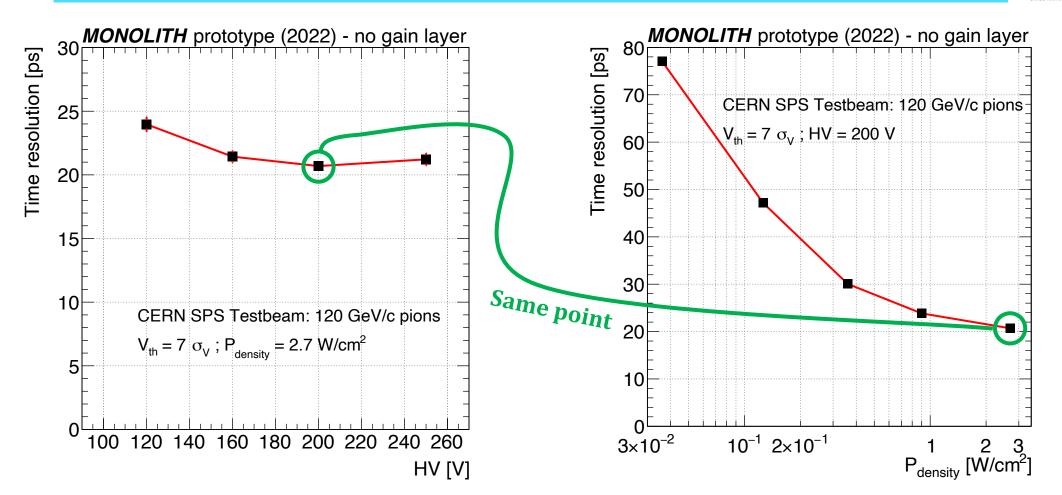
Large plateau of 100 V with ~20 ps

20 ps at 2.7 W/cm² 50 ps at 0.1 W/cm²



Time Resolution Results





What happens if we increase the generated charge?





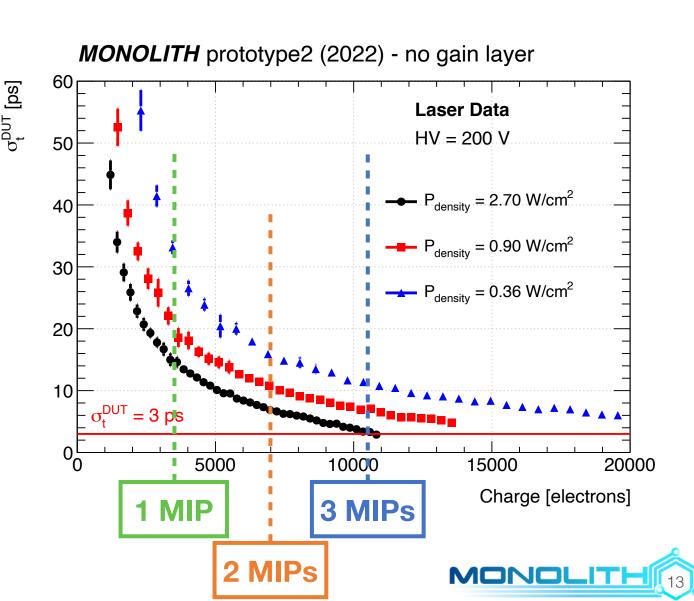
Laser Measurements



- A Minimum Ionizing Particle produces around 3500 electrons (0.5 fC) in the 50 µm-thick active area
- At a power density of 2.7 W/cm²:
 - The time resolution for the equivalent charge of 1 MIP is ~15 ps
- With the laser, the only contribution to the time resolution is the **electronics time** jitter: σ_V 1

 $\sigma_{\text{Jitter}} = \frac{\sigma_V}{\frac{dV}{dt}} \propto \frac{1}{\text{charge}}$

 With charges higher than 3 MIPs (1.5 fC, or 38 keV), the time jitter is less than 3 ps

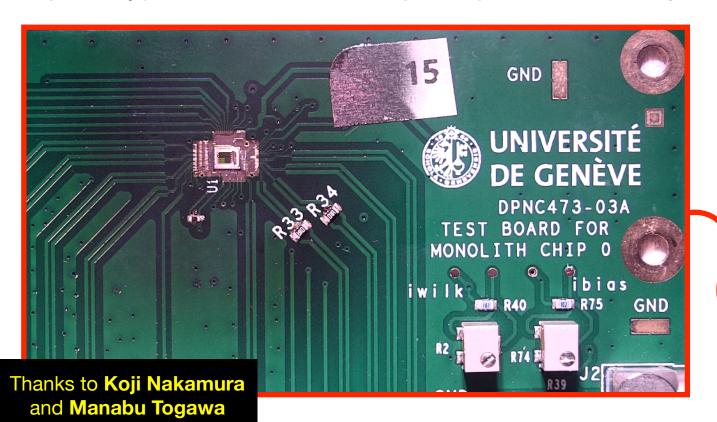




Radiation Hardness



- Assess the performance of efficiency and time resolution after proton irradiation
- 8 prototypes irradiated in Japan up to 1×10¹⁶ n_{eq}/cm²



Board Name	Fluence [1 MeV n _{eq} /cm ²]
M23	2 · 10 ¹³
M22	9 · 10 ¹³
M21	6 · 10 ¹⁴
M19	6 · 10 ¹⁴
M18	3 · 10 ¹⁵
M17	3 · 10 ¹⁵
M16	1 · 10 ¹⁶
M15	1 · 10 ¹⁶

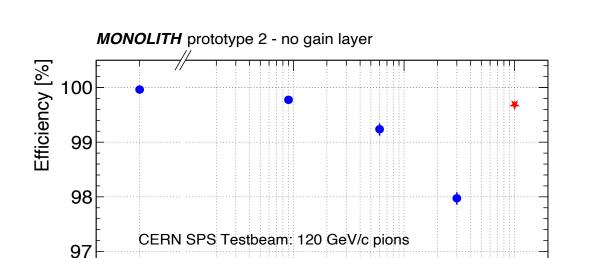




Efficiency and Time Resolution After Irradiation







10¹⁴

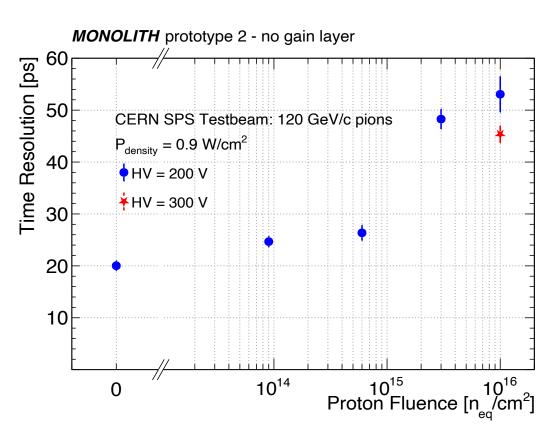
 $P_{density} = 0.9 \text{ W/cm}^2$

♦ HV = 200 V

∔HV = 300 V

0

96



At 10^{16} n_{eq} /cm², with HV = 200 V and $P_{density} = 0.9$ W/cm²:

 $10^{15} 10^{16}$ Proton Fluence $[n_{eq}/cm^2]$

- The **efficiency** is 96.2% -> **99.7%** increasing the High Voltage
- The time resolution is 53 ps -> 45 ps increasing the High Voltage

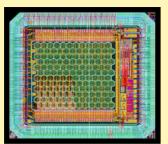




Prototypes With PicoAD©







- Hexagonal pixels 65 µm
- 30 ps TDC + I/O logic
- 4 analog outputs

36 ps

JINST 17 P02019

In this presentation

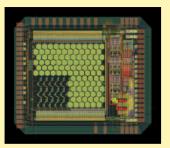
2022

With PicoAD®

17 ps

JINST 17 P10032 JINST 17 P10040

2022



- Hexagonal pixels 65 µm
- Improved frontend
- 4 analog outputs

20 ps

<u>JINST 18 P03047 JINST 19 P01014</u> JINST 19 P04029 JINST 19 P07036, In this presentation

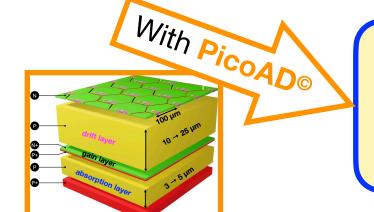
2024

With PicoAD©

12 ps

JINST 20 P04001



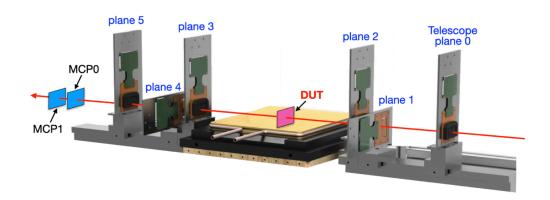




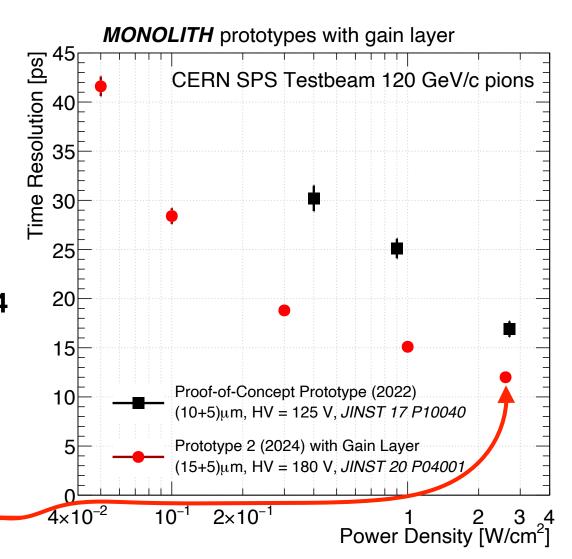


Testbeam Results





- Testbeam measurements with FE-I4 telescope
 - PicoAD[®] proof-of-concept in October 2021
 - Second prototype with gain layer in May 2024
- The efficiency is compatible with 99.9% for many working points
- The time resolution improved by almost a factor of two:
 - The best performance with the second prototype is 12.1 ps with HV of 180 V and Pdensity of 2.6 W/cm²

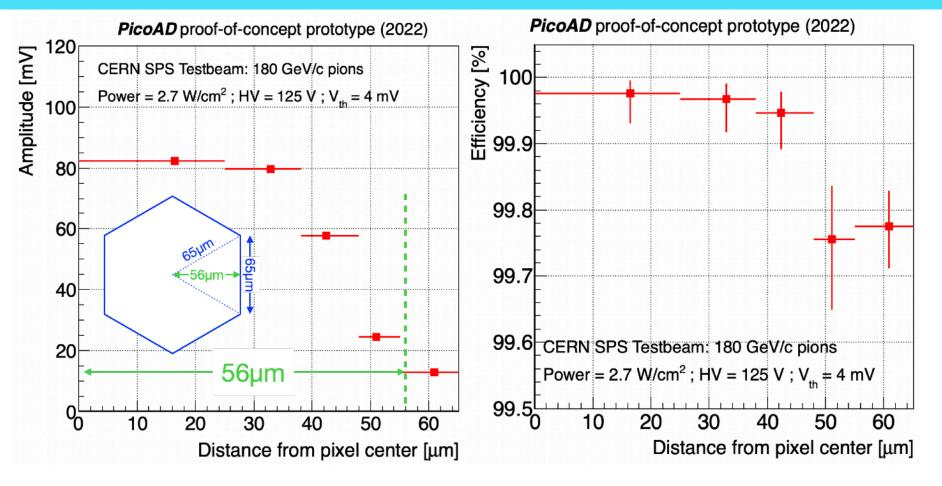




Performance Within the Pixel







- We saw a slight degradation of the performance towards the pixel's edge with the PicoAD[©] proof-of-concept:
 - Lower amplitudes and lower efficiencies

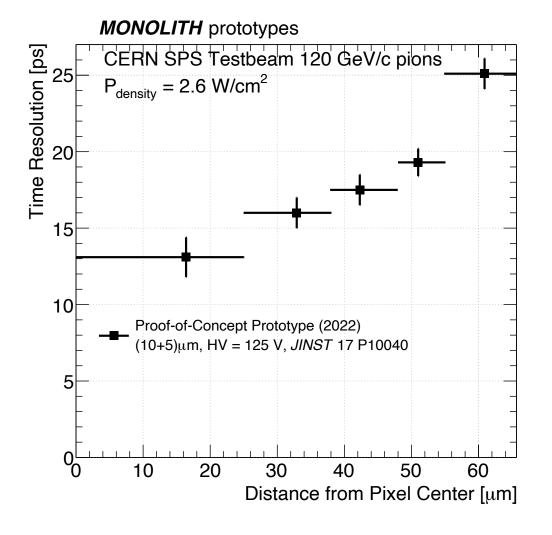




Time Resolution Within the Pixel



- The degradation of the amplitudes leads to a degradation of the time resolution:
 - The **time resolution** of the **PicoAD**[©] proof-of-concept varies between **13 and 25 ps** (almost a factor 2 worse), with an average of 17 ps

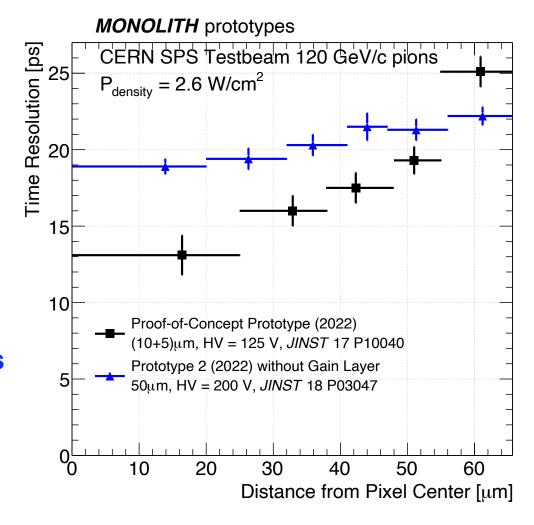




Time Resolution Within the Pixel



- The degradation of the amplitudes leads to a degradation of the time resolution:
 - The **time resolution** of the **PicoAD®** proof-ofconcept varies between **13 and 25 ps** (almost a factor 2 worse), with an average of 17 ps
- Some **improvements** on the ASIC and the substrate make the situation more stable:
 - The time resolution of the second prototype
 without a gain layer varies between 19 and 24 ps



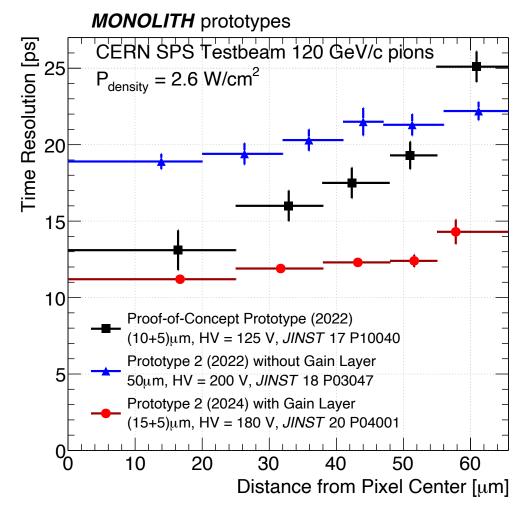




Time Resolution Within the Pixel



- The degradation of the amplitudes leads to a degradation of the time resolution:
 - The **time resolution** of the **PicoAD®** proof-ofconcept varies between **13 and 25 ps** (almost a factor 2 worse), with an average of 17 ps
- Some **improvements** on the ASIC and the substrate make the situation more stable:
 - The time resolution of the second prototype
 without a gain layer varies between 19 and 24 ps
 - The time resolution of the second prototype with a gain layer varies between 11 and 14 ps





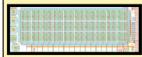
Summary





- Square pixel 900 µm Discriminated output
 - 200 ps

2017



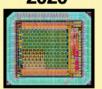
- Square pixels 500 µm 100 ps TDC + I/O logic
 - 110 ps

2018



- Hexagonal pixels 130 µm
- Discriminated outputs
 - 45 ps

2020



- Hexagonal pixels 65 µm
- 30 ps TDC + I/O logic 4 analog outputs
- 36 ps



- Hexagonal pixels 65 µm Improved frontend
- 4 analog outputs
 - 20 ps

- 20 ps time resolution with a MIP
- Laser measurements:
 - 3 ps with the equivalent charge of 3 MIPs
- After irradiation at 10¹⁶ n_{eq}/cm²:
 - 99.7% efficiency
 - 45 ps time resolution

2022

With PicoAD®

17 ps

2024 With PicoAD® 12 ps

12 ps time resolution with a MIP

17 ps time resolution with a MIP

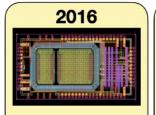
MONOLIT



Other Projects at UNIGE

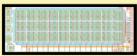






- Square pixel 900 µm Discriminated output
- 200 ps

2017



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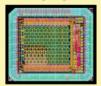
2018



- Hexagonal pixels 130 µm
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- Hexagonal pixels 65 μm
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36 ps

2022



- Hexagonal pixels 65 μm
- Improved frontend
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20 ps

Medical Application: 100µPET



2024



2022

With PicoAD[©]
17 ps

See Roberto Cardella's talk 2024

With PicoAD[©]
12 ps



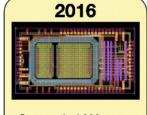


Future Perspectives





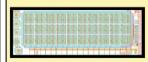
High-radiation level timing layers for particle physics



Square pixel 900 µm
Discriminated output

200 ps

2017



Square pixels 500 μm 100 ps TDC + I/O logic

110 ps

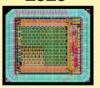
2018



Hexagonal pixels 130 µm
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45 ps

2020



- Hexagonal pixels 65 μm
 30 ps TDC + I/O logic
- 30 ps TDC + I/O logic4 analog outputs

36 ps

2022



- Hexagonal pixels 65 µmImproved frontend
- 4 analog outputs

20 ps

Moderate radiation level, extremely thin timing layers for particle physics

Medical Application: 100μPET

2024

2022 With PicoAD® 17 ps With PicoAD[©]
12 ps

2024

Photonics



People Involved











Giuseppe lacobucci

- Project P.I.
- System Design



Thanushan Kugathasan

- Lead Chip Design
- **Analog Electronics**



Talk on 100µPET

yesterday

Roberto Cardella

- Analog and Digital Electronics
- Sensor Design



Leonardo Cecconi

- Chip Design
- Firmware



Luca lodice

- Chip Design
- Firmware



Carlo Alberto Fenoglio

- Chip Design
- Firmware



Andrea Pizarro Medina

- **Laboratory Tests**
- Data Analysis



- Magdalena Munker
- Jordi Sabater Iglesias



Lorenzo Paolozzi

- Sensor Design
- **Analog Electronics**



Stefano Zambito Laboratory Tests

Data Analysis



Mateus Vicente

- System Integration
- **Laboratory Tests**



Viros Sriskaran

- Analog Electronics
- **Digital Electronics**



Chiara Magliocca

- Laboratory Tests
- Data Analysis



Théo Moretti

- Laboratory Tests
- Data Analysis



Jihad Saidi

- **Laboratory Tests**
- Data Analysis



- Ivan Semendyaev
- Rafaella Kotitsa



Didier Ferrere

- System integration
- Laboratory test



Yannick Favre

- Board design
- RO system



Sergio Gonzalez-Sevilla

- System integration
- Laboratory test



Stéphane Débieux

- Board design
- RO system

Main research partners:



Roberto Cardarelli INFN Rome2 & UNIGE



Marzio Nessi **CERN & UNIGE**



Holger Rücker IHP Mikroelektronik



Matteo Elviretti IHP Mikroelektronik

Thanks for Your Attention

Questions?





The MONOLITH ERC Project



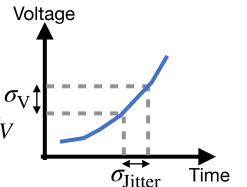
• A precise measure of time (σ_t ~30 ps) is achieved by reducing the two main contributions to the time resolution:

• The electronics time jitter:

$$\sigma_{\text{Jitter}} = \frac{\sigma_V}{dV/dt}$$

• Related to the signal rise σ time dV/dt and the noise σ_V

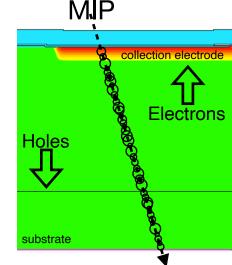
 We want fast and lownoise amplification



• The charge-collection noise:

$$\sigma_{\rm Landau} = w(d) \, \sigma_{\tau, \rm SC} \propto \frac{d}{\sqrt{ln(d)}}$$

- Related to the detector thickness d
- We want a small detector thickness



SiGe Technology Monolithic BiCMOS Process

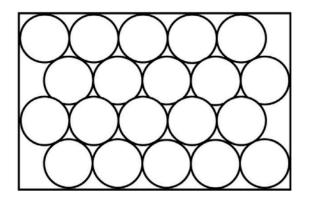
PicoAD©:
Picosecond
Avalanche
Detector

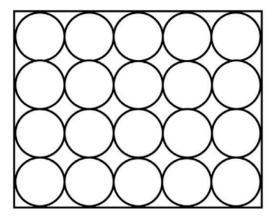


Hexagonal Pixels



- Three possible regular shapes to use:
 - equilateral triangles
 - squares
 - regular hexagons
- Hexagons have the highest angles (120°)
 - electric fields in the corners are better under control
- Moreover, the same amount of pixels can fits in less space than squares





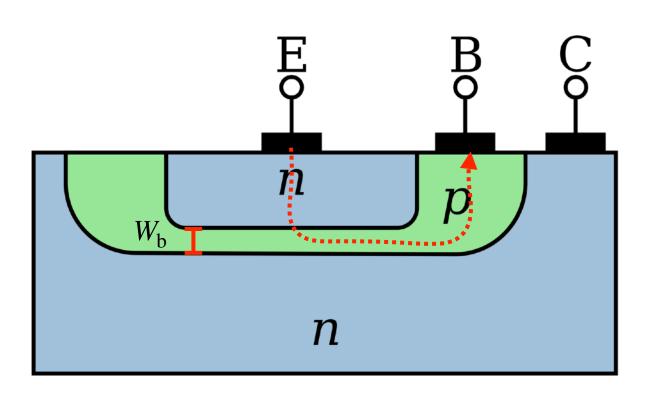


Thick Base = Small Base Resistance



- When a grading of Ge is added to the base of a BJT, it is possible to make the base thicker
- A thicker base thickness $W_{\rm b}$ means a smaller base resistance $R_{\rm b}$ than a BJT
- The electrons coming from the emitter to the base see the base thickness $W_{\rm b}$ as the transversal dimension:

$$R_{\rm b} \propto \frac{\rho_{\rm b} l_{\rm b}}{W_{\rm b}}$$





Charge-Collection Noise in Detail

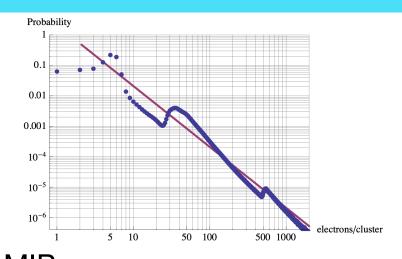


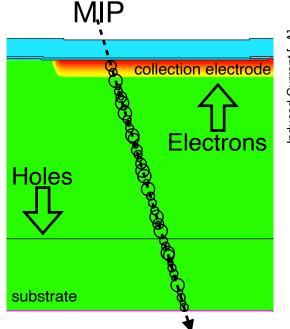


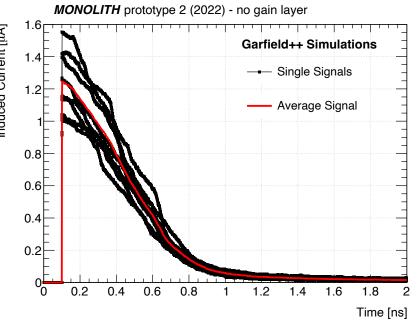
- The charge generated by a Minimum lonizing Particle follows a pattern of clusters with different charges
- The probability of electrons for each cluster is known
- Electrons and holes drift thanks to the electric field and induce a current:

$$i_{ind} = -q \vec{v}_{drift} \cdot \vec{E}_{weight}$$

- A cluster stops inducing current when it's collected, creating noise on the induced current signal
- The contribution of this noise to the time resolution is called chargecollection noise (or Landau noise)



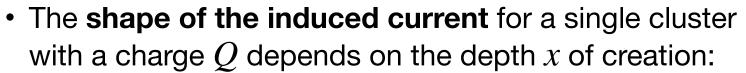






Charge-Collection Noise in Detail





$$i_{SC}(x,t) = \frac{Q v_e}{d} \Theta\left(\frac{x}{v_e} - t\right) + \frac{Q v_h}{d} \Theta\left(\frac{d - x}{v_h} - t\right)$$

And so it does the center-of-gravity time:

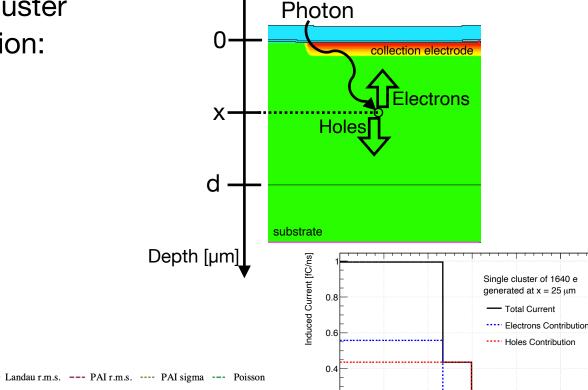
$$\tau = \frac{\int_0^\infty t \, i_{SC}(x,t) \, dt}{\int_0^\infty i_{SC}(x,t) \, dt} = \frac{1}{2d} \left[\frac{x^2}{v_e} + \frac{(d-x)^2}{v_h} \right]$$

• With its variance:

$$\sigma_{\tau,SC}^2 = \frac{4}{180} \frac{d^2}{v_e} - \frac{7}{180} \frac{d^2}{v_e v_h} + \frac{4}{180} \frac{d^2}{v_h}$$

• For a **MIP** with many clusters, and when the amplifier is slow, the charge-collection noise contribution to the time resolution is:

$$\sigma_{\mathrm{Landau}} = \sigma_{\tau,\mathrm{MIP}} = w(d) \, \sigma_{\tau,\mathrm{SC}} \propto \frac{d}{\sqrt{ln(d)}}$$





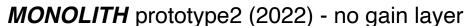
0.2

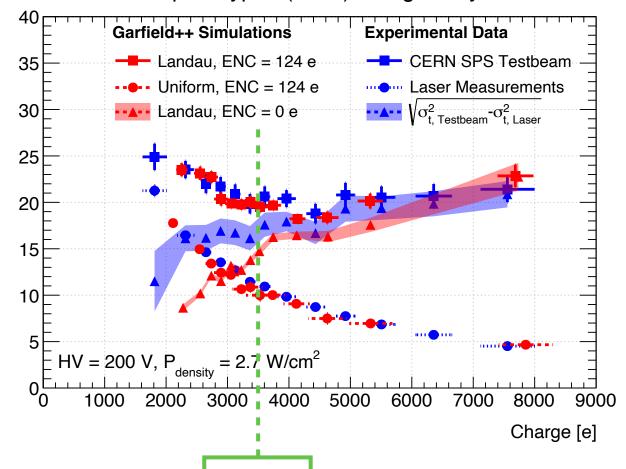


Charge-Collection Noise



- Garfield++ simulation using:
 - The measured Equivalent Noise Charge
 - The electric field (TCAD)
 - The electronics δ-response (Cadence)
- The MIP simulation matches the testbeam data ($\sigma_{
 m Jitter} \oplus \sigma_{
 m Landau}$)
- The charge deposition created by the laser is uniform (no $\sigma_{
 m Landau}$)
- The simulation matches the laser data
- The difference in quadrature between the testbeam and laser should be the $\sigma_{\rm Landau}$ contribution
- The simulated trend is similar to the measured one
 - There are non-negligible correlations between $\sigma_{
 m Jitter}$ and $\sigma_{
 m Landau}$



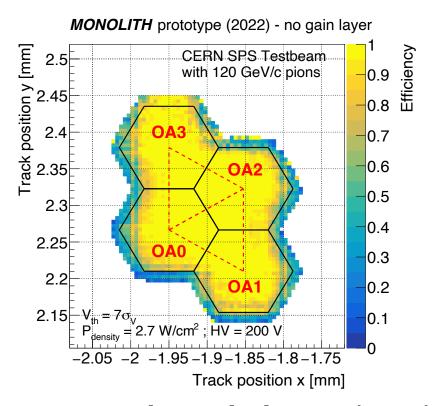






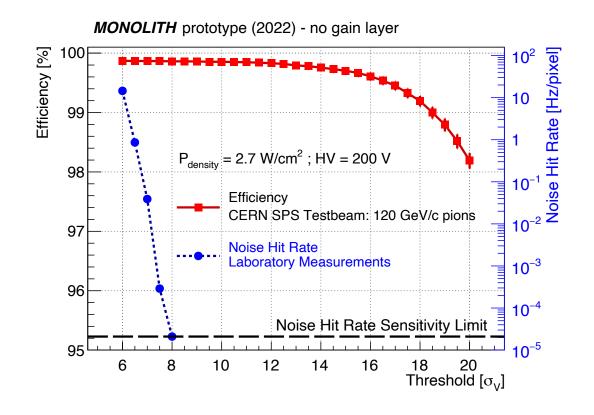
Efficiency Results







- Selection of two triangles:
 - representative of the whole pixel
 - unbiased by the telescope resolution



- Large plateau of 99.8% efficiency
 - $\sigma_V \approx 1.4 \text{ mV} \approx 100 e^-$
- Negligible noise hit rate already at $7\sigma_V$

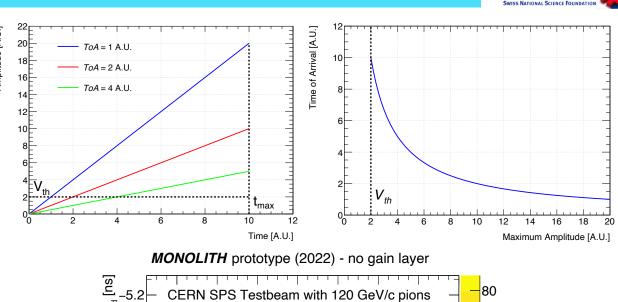


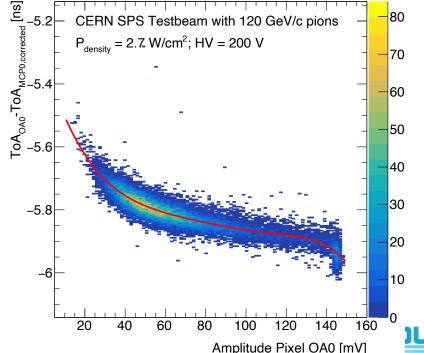


Time-Walk Correction



- The time walk is the spread in the time of arrival created by signals having different amplitudes
- In our prototypes, the time-walk correction is implemented in two possible ways:
 - Either taking the time of arrival as the time to reach a fixed threshold, and subtracting the average time of arrival for each amplitude
 - Either taking the time of arrival as the time to reach a constant fraction of the signal, which automatically corrects the time walk







Département de physique

nucléaire et corpusculaire

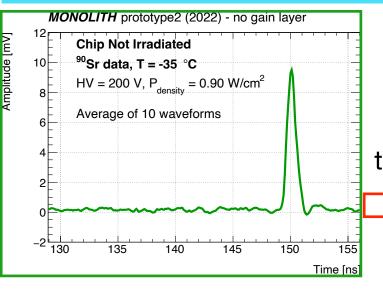
Main Characteristics After Irradiation



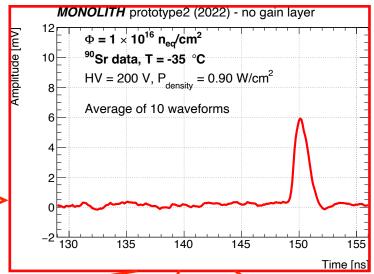


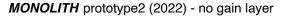


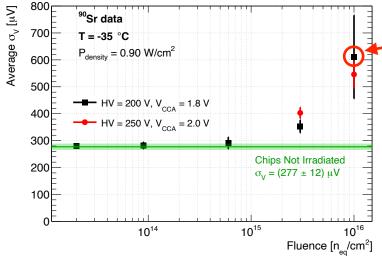


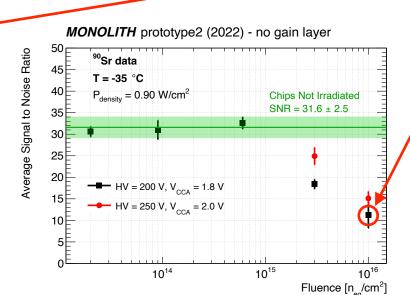


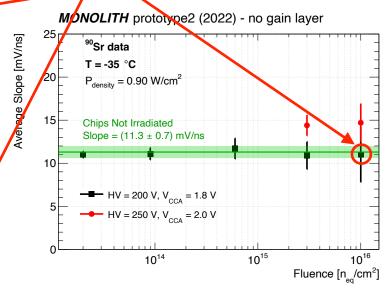
After irradiation to 1×10¹⁶ n_{eq}/cm²









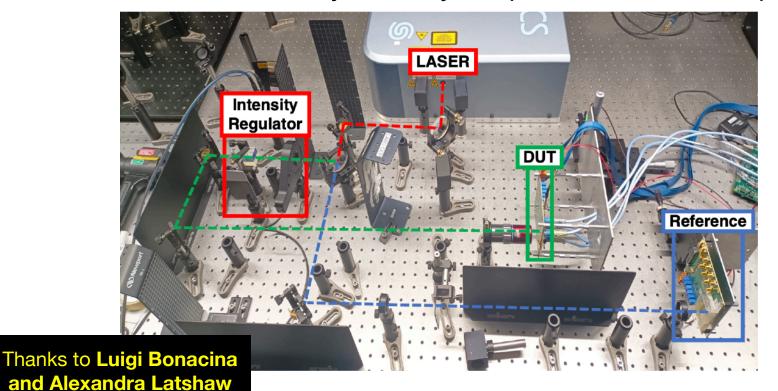




Femtosecond Laser Measurements



- Pulsed infrared femtosecond laser:
 - Uniform charge distribution: **no charge-collection noise** ($\sigma_{Landau} = 0$)
 - To study the pure contribution of the **electronics** σ_{Jitter}
- Time coincidence between two of our samples:
 - "Reference": large laser pulse for a precise reference
 - "DUT": variable intensity to study the performance vs. amplitude





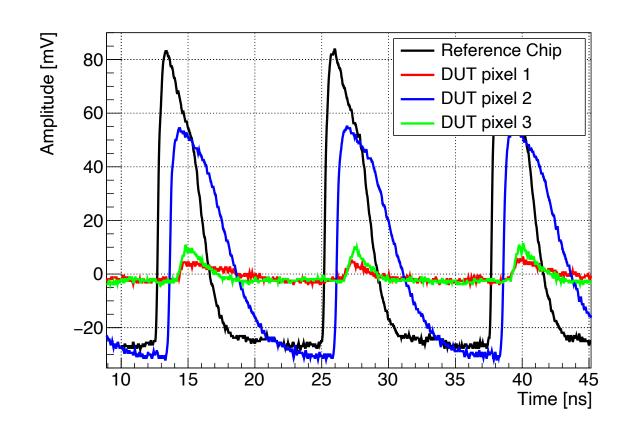


Femtosecond Laser Characteristics





- Pulsed infrared femtosecond laser:
 - Wavelength of 1060 nm
 - Intrinsic jitter of 100 fs
 - Hence the name femtosecond laser
 - Repetition frequency of 80 MHz
 - Focused on the DUT on a Gaussian spot with a diameter of ~30 µm
- The Reference and DUT ASICs can tolerate the high repetition rate of the laser





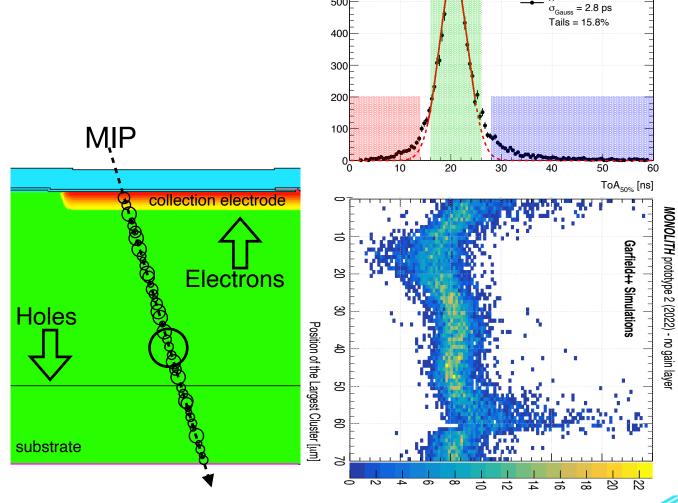
Non-Gaussianity of the ToA Distributions







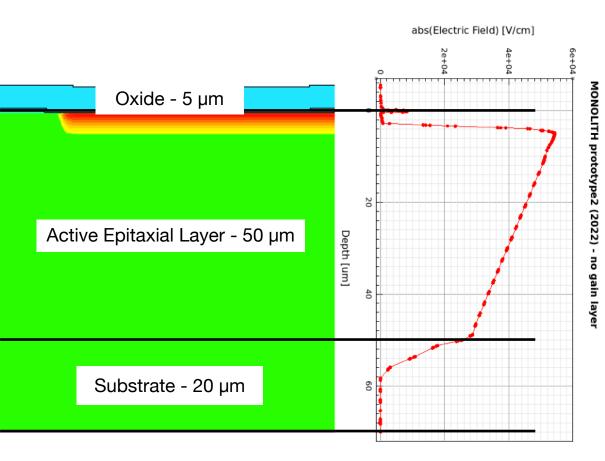
- Sometimes, clusters containing more than 1000 electrons are generated in the detector
- The corresponding voltage signal has a specific time of arrival according to the position of the big cluster:
 - Big clusters between 10 to 20 μm produce faster signals
 - Big clusters in the peripheral regions produce slower signals
- Garfield++ simulation were helpful in understanding this issue

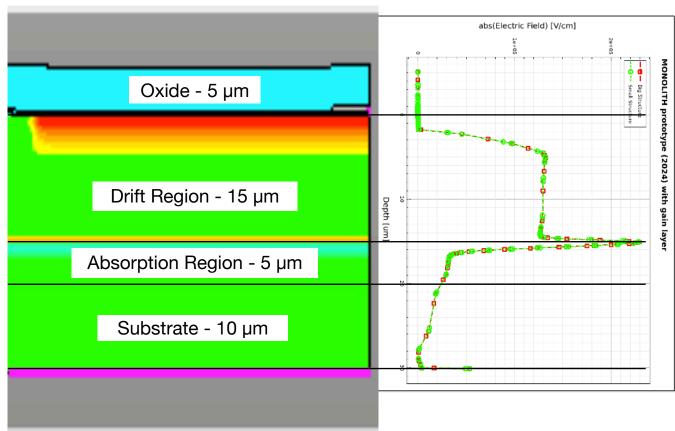




Electric Field Profiles







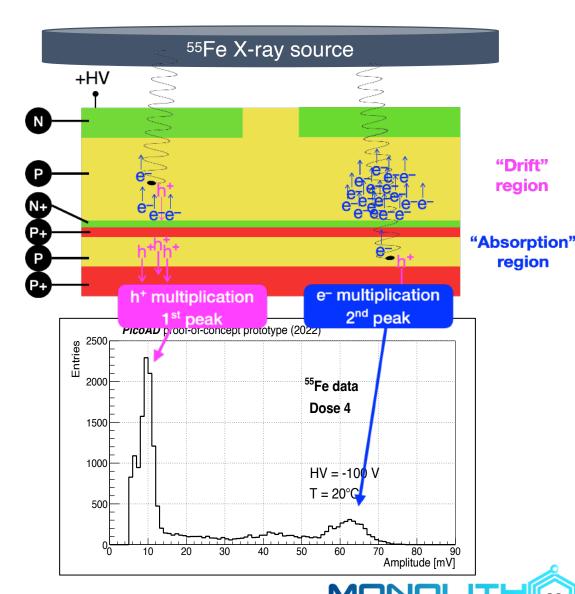


Gain Measurements with a 55Fe Source





- X-rays from ⁵⁵Fe radioactive source:
 - ~5.9 keV photons with point-like charge deposition -> 1640 electrons
- Characteristic double-peak spectrum
 - Photon absorbed in the **drift region**:
 - Holes drift through the gain layer and multiply
 - First peak in the spectrum
 - Photon absorbed in the absorption region:
 - Electrons drift through the gain layer and multiply
 - Second peak in the spectrum
- We can calculate the gain for electrons from the spectra

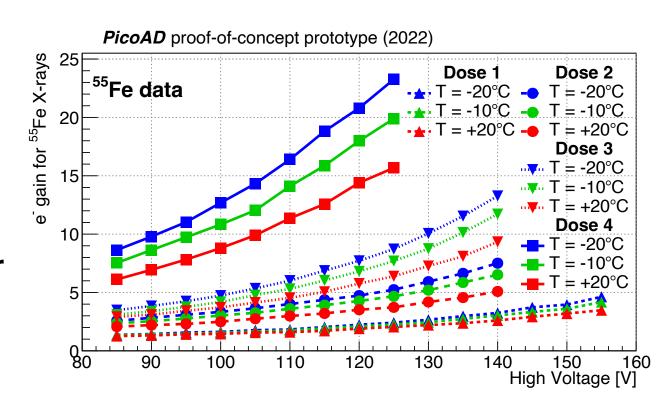




PicoAD[©] Proof-of-Concept Gain Results



- The gain layer of the PicoAD[®] proof-ofconcept prototype works as expected:
 - Higher gain for higher bias voltage
 - Higher gain for higher gain layer doses
 - Higher gain for lower temperatures
- But, the maximum electron gain is 23 for
 55Fe X-rays at 125 V and -20 °C
 - It is smaller than the simulated one with TCAD

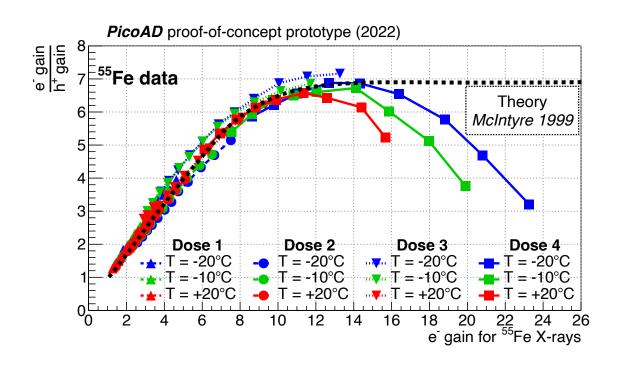




PicoAD© Proof-of-Concept Gain Results





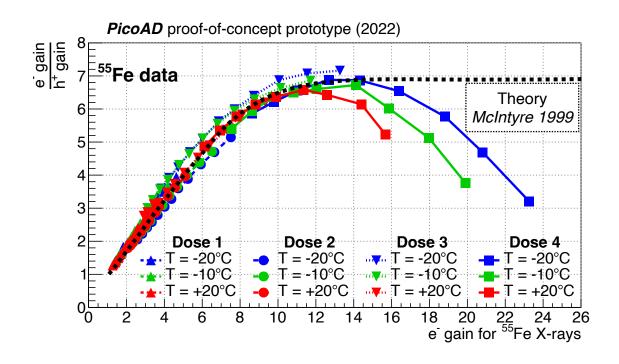


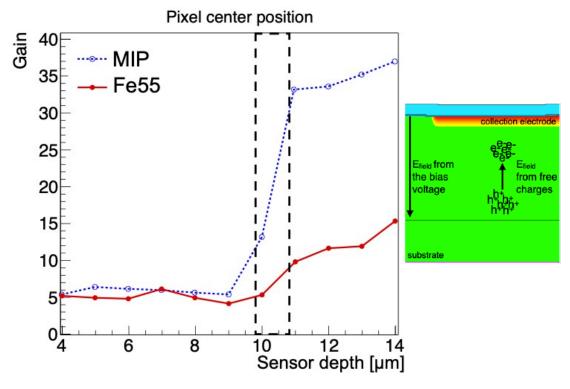
- The gain for holes and electrons can be studied simultaneously with PicoAD®
- The gain ratio as a function of the electron gain deviates from the theoretical trend



PicoAD[©] Proof-of-Concept Gain Results







- The gain for holes and electrons can be studied simultaneously with PicoAD[©]
- The **gain ratio** as a function of the electron gain deviates from the theoretical trend
- Refined TCAD simulations show that:
 - The gain is suppressed by the large cloud of charges generated by the ⁵⁵Fe source
 - This suppression is not present for a MIP

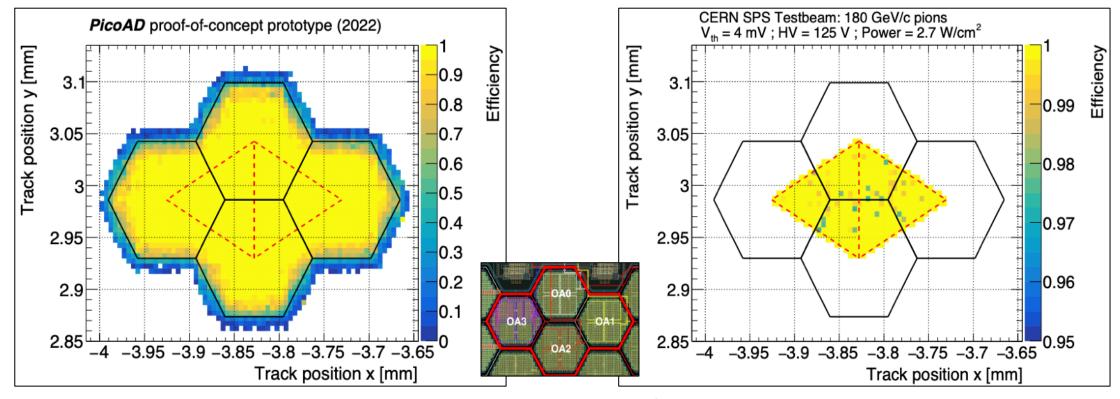




PicoAD[©] Proof-of-Concept Efficiency Maps







 The apparent degradation at the edges is due to the finite resolution of the telescope (~10 μm)

- Selection of two triangles:
 - representative of the whole pixel
 - unbiased from the telescope resolution





PicoAD[©] Proof-of-Concept ΔToA Distributions

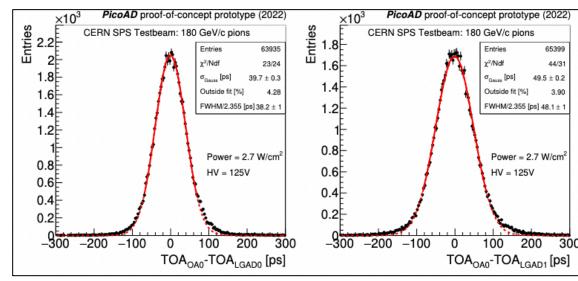


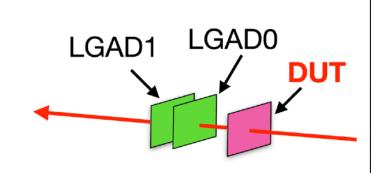


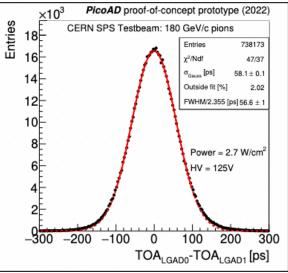




- The distributions are after a timewalk correction
- The distributions are Gaussian
 - Except ~2-4 % of the entries that are in non-Gaussian tails
- The three σ_{Gauss} from the fits give the three timing resolutions of:
 - The DUT
 - The two LGADs







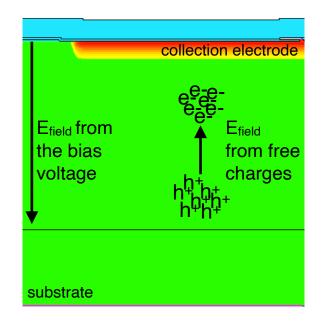


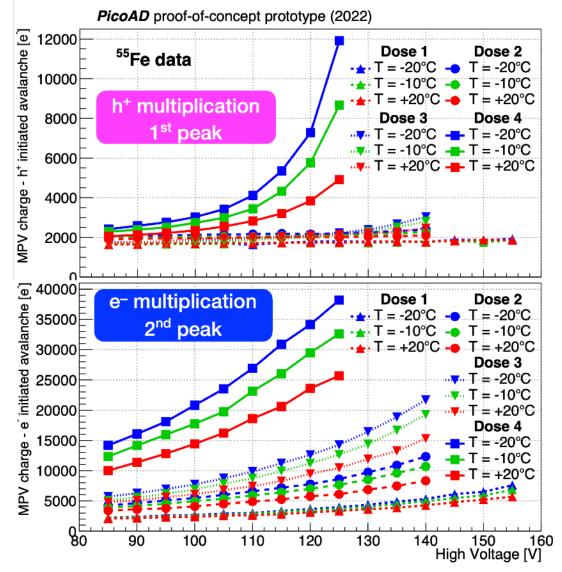
PicoAD© Gain Suppression





- Only the holes have the correct exponential trend that the gain should have as a function of the bias voltage:
 - The high number of charges creates an electric field opposite to the one created by the bias voltage
 - The electron gain is suppressed by the high number of charges





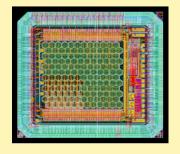


Improvements from 2020 to 2022



- Similar matrix configuration, but:
 - From standard 50 Ω cm on 1 Ω cm substrate to 50 μ m-thick 350 Ω cm epi-layer on 1 Ω cm substrate led to:
 - A smaller pixel capacitance
 - A larger depletion from 23 μm to 50 μm
 - A much larger voltage plateau
 - A saturated drift velocity everywhere
 - Preamp and driver voltage decoupling led to:
 - An optimal amplifier operation
 - The removal of the cross-talk
 - An optimised front-end electronics layout, with a semi-differential output, and high-frequency cables led to:
 - A better rise time (from 600 ps to 300 ps)

2020

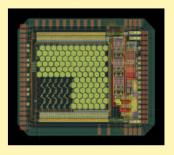


- Hexagonal pixels 65 µm
- 30 ps TDC + I/O logic
- 4 analog outputs

36 ps

JINST 17 P02019

2022



- Hexagonal pixels 65 µm
- Improved frontend
- 4 analog outputs

20 ps

<u>JINST 18 P03047 JINST 19 P01014</u> <u>JINST 19 P04029 JINST 19 P07036</u>