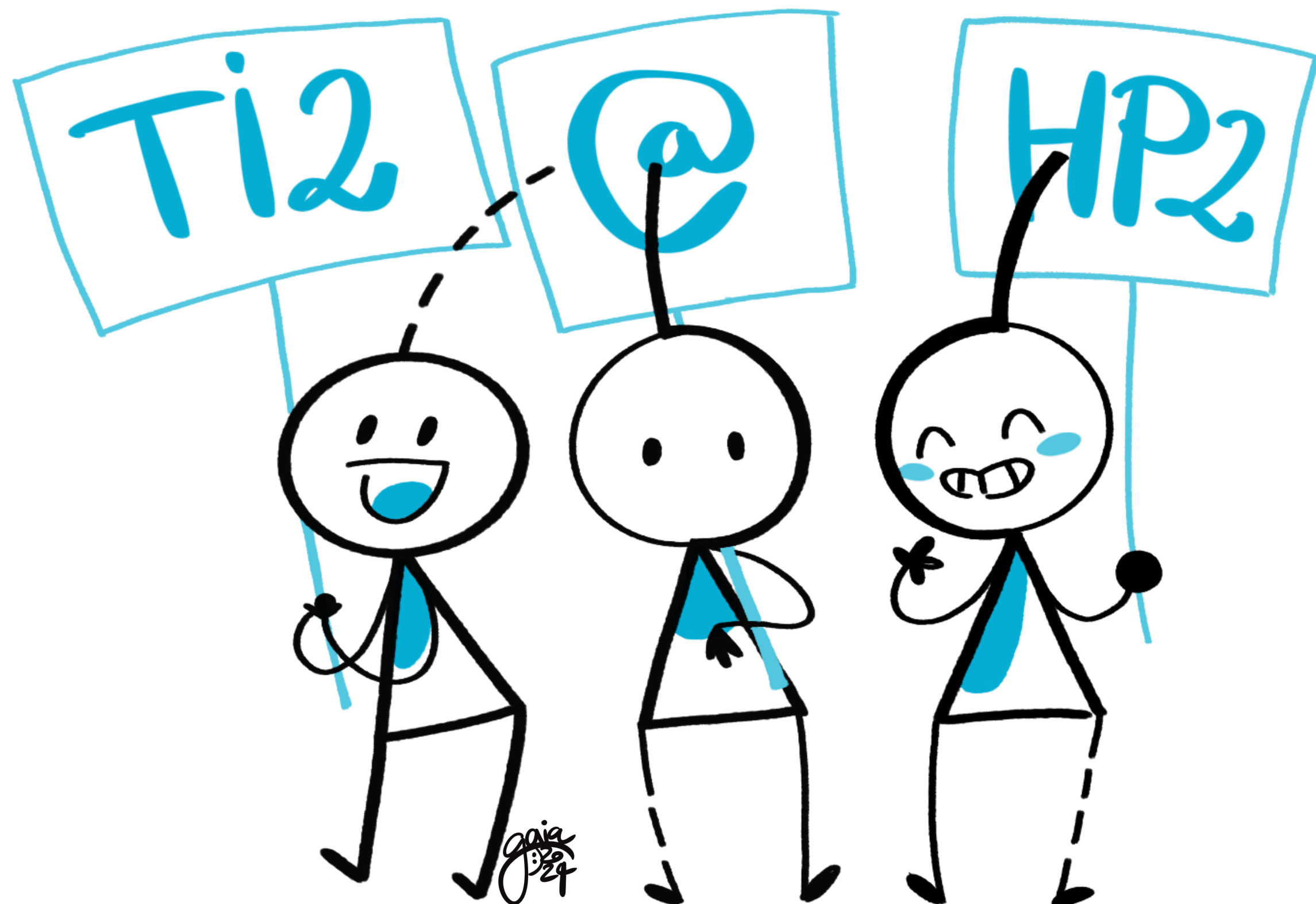


Taming IBPs with Transverse Integration

@ High precision for hard processes



Gaia Fontana, University of Zürich
with Vsevolod Chestnov & Tiziano Peraro

[arXiv:2409.04783](https://arxiv.org/abs/2409.04783)



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Zürich**^{UZH}

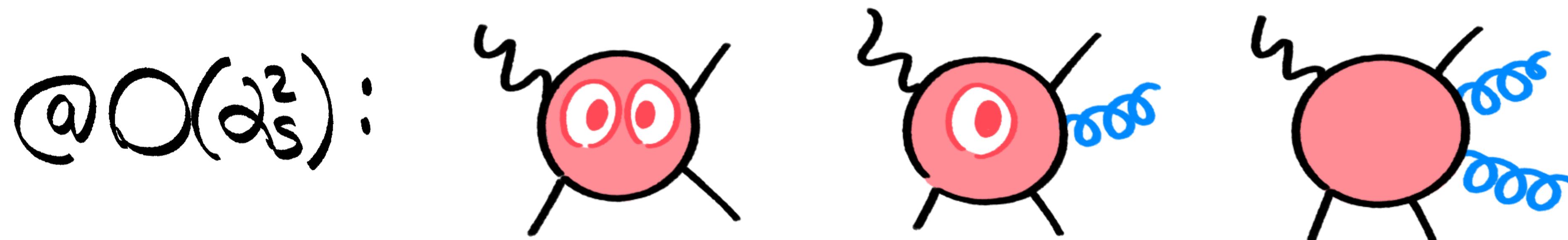
HP2 2024, 12/09/2024

Precision era @ colliders

- Precision physics as
 - test of the Standard model
 - gate to new physics



- High-Lumi upgrade of LHC :
- theory and experiments must have comparable uncertainties
 - needed: %-level accuracy:
 - perturbation theory @ NNLO and often N3LO
 - diagrams with increasing no. of loops, legs & mass scales



Part 0: Background



A dictionary for Feynman integrals

- LEGO® blocks of perturbative QFT beyond tree level
- Key ingredient of phenomenological predictions
- Rich and interesting mathematical structures

Integral Family :

defined by a list of generalised denominators $F \leftrightarrow \{D_{F,1}, \dots, D_{F,n}\}$

- Integral belonging to a family

$$I_{F;\vec{a}}[N] = I_{F;a_1 \dots a_n}[N] = \int \prod_{j=1}^{\ell} d^D k_j \frac{N}{\prod_{j=1}^n D_{F,j}^{a_j}}, \quad N = \text{polynomial in } k_i$$

Numerators are removed via tensor reduction

→ $N = 1$ for IBPs

$$I_{F;\vec{a}} = \int \prod_{j=1}^{\ell} d^D k_j \frac{1}{\prod_{j=1}^n D_{F,j}^{a_j}}$$

- Generalised denominators have the form

$$D_{F,j} = l_j^2 - m_j^2$$

$$D_{F,j} = l_j \cdot v_j - m_j^2$$

l_j linear combination of k_j , v_j linear combination of p_j

We distinguish:

- Proper denominators: $D_{F,j}$ such that $a_j > 0$
- Irreducible scalar products (ISPs): $D_{F,j}$ such that $a_j \leq 0$

- Sectors, $S_{F,\vec{a}}$: integrals with the same set of proper denominators
 - Iteratively, one can define also subsectors/parent sectors
- Corner integral of a sector: integral with $a_j \in \{0,1\}$

Integral decomposition

why?



- Extremely large number of integrals contributing to an amplitude
- Properties/symmetries of an amplitude manifest only after the reduction
- Important for the calculation of the integrals

Integral decomposition

Reduction into a basis of linearly independent
master integrals $\{G_j\} \subset \{I_j\}$

$$I_j = \sum_k c_{jk} G_k$$

master integrals

rational coefficients

$\{G_j\}$ = minimal linearly independent set



Laporta algorithm

Feynman integrals in dimensional regularization obey linear relations,
e.g. **Integration By Parts** identities

+ Lorentz Invariance ids, symmetry relations, ...

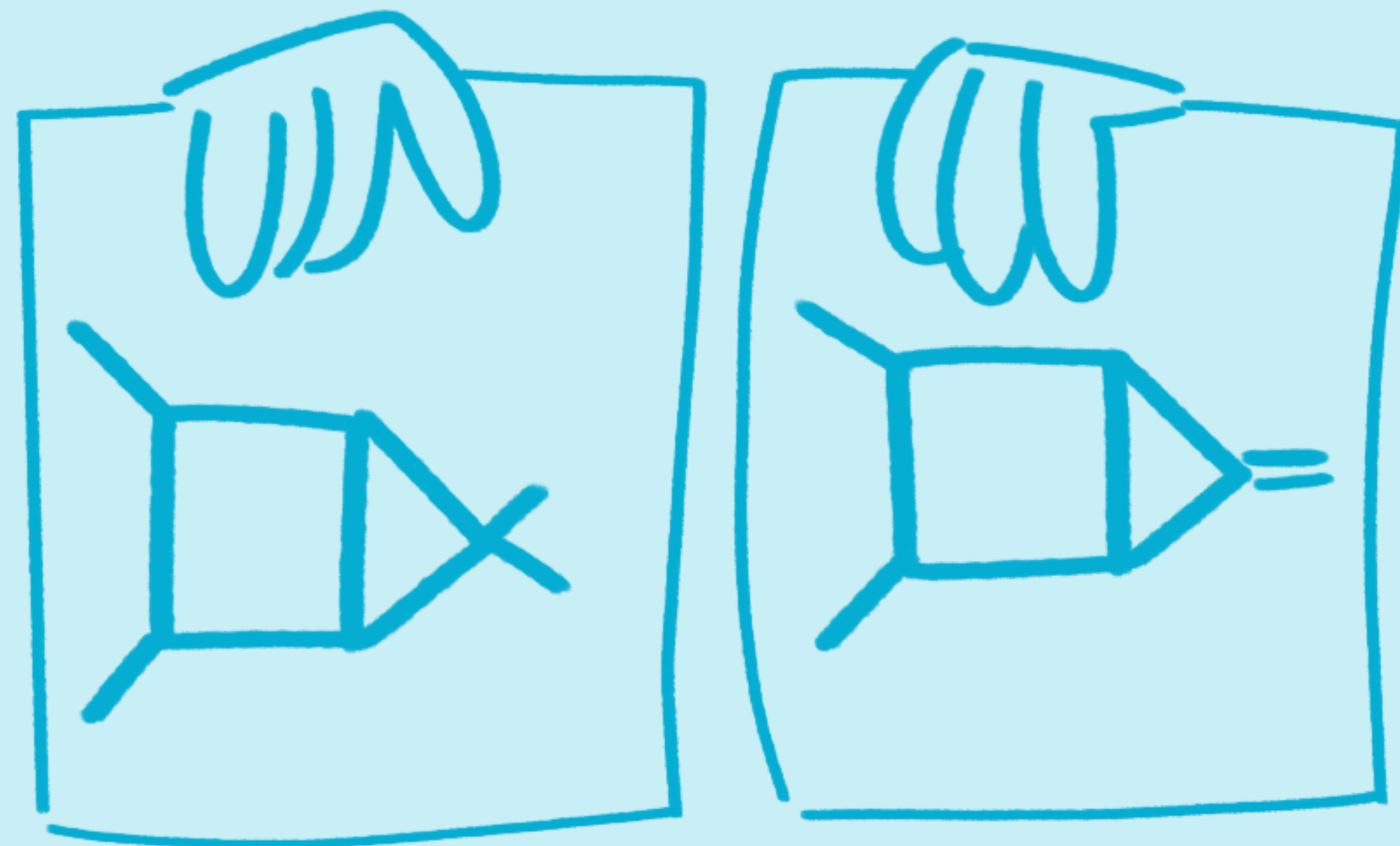
[Chetyrkin, Tkachov (1981), Laporta (2000)]

$$\int \left(\prod_{i=1}^{\ell} d^D k_i \right) \frac{\partial}{\partial k_i^\mu} \left(\frac{v_j^\mu}{D_1^{a_1} \dots D_n^{a_n}} \right) = 0, \quad v^\mu = \begin{cases} p_i^\mu = \text{external} \\ k_i^\mu = \text{loop} \end{cases}$$

reduction as solution of a **large**
and **sparse** system of identities

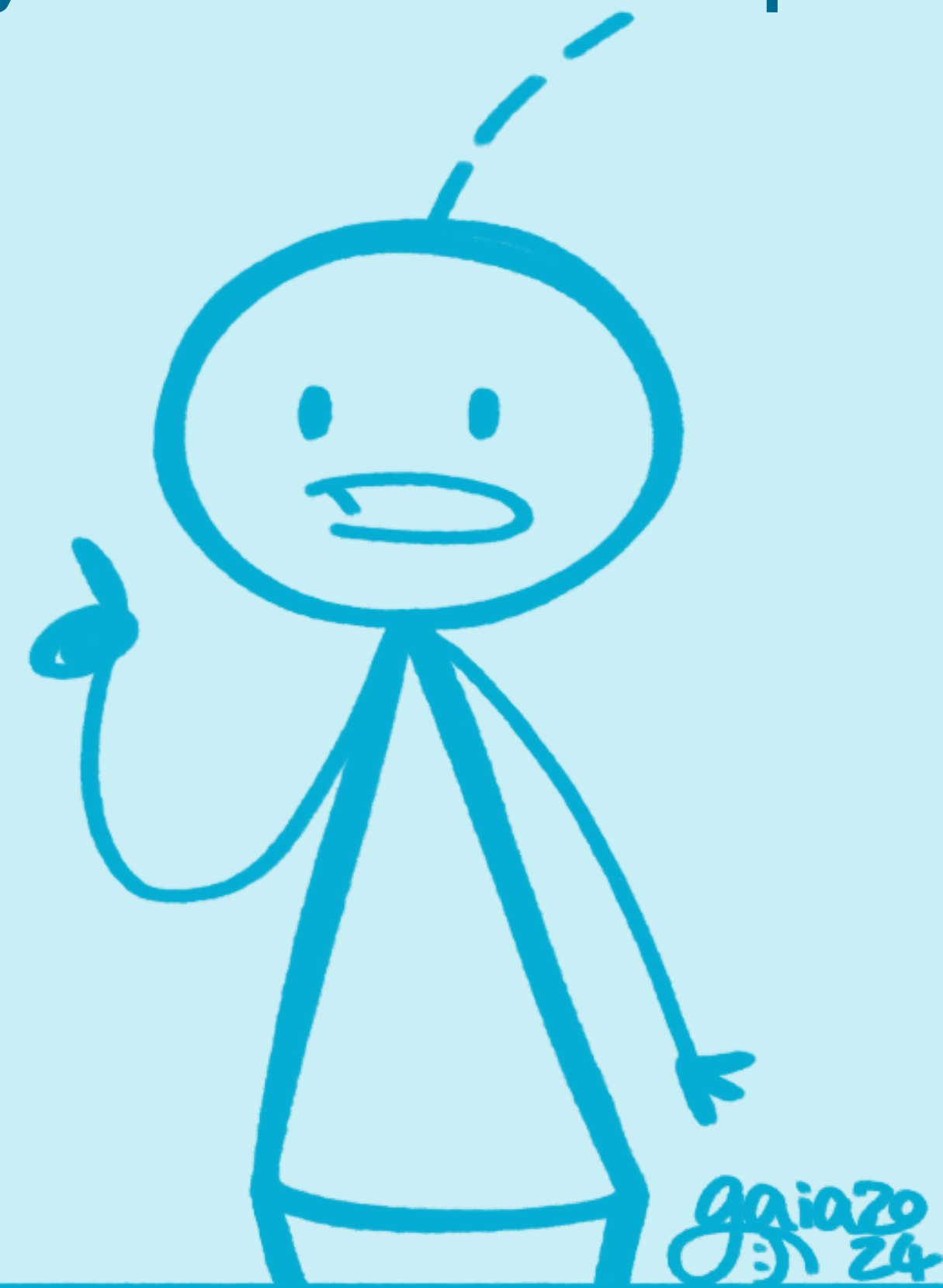
**Computational bottleneck in
state-of-the-art calculations**

Part 1: The main idea



We need you to find the differences between these two pictures

They're the same picture



Transverse integration id.s

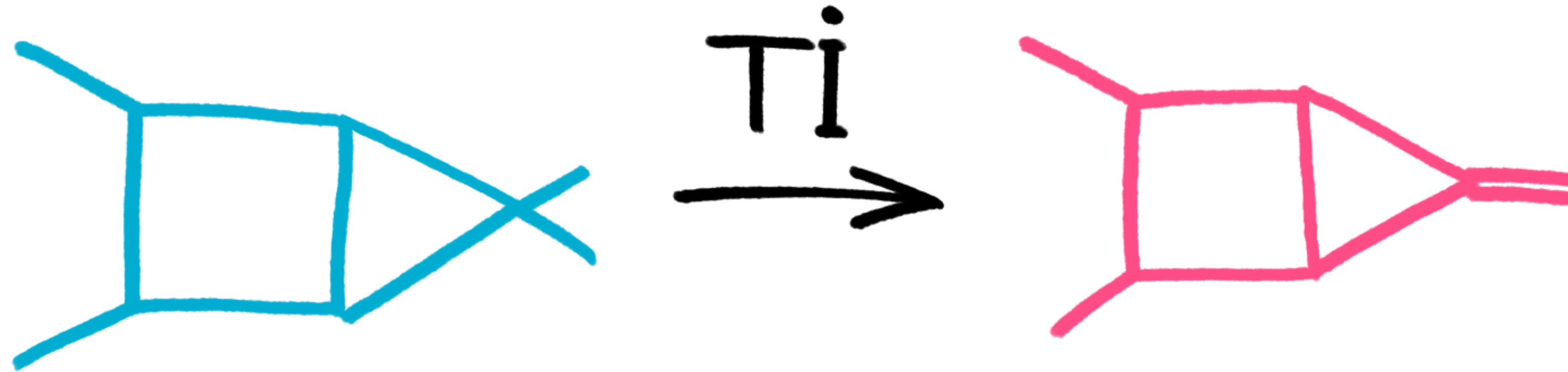
- A way to simplify the identities in the Laporta system
- Formulation in terms of angular integrations in [\[Mastrolia, Peraro, Primo 2017\]](#)

Idea :

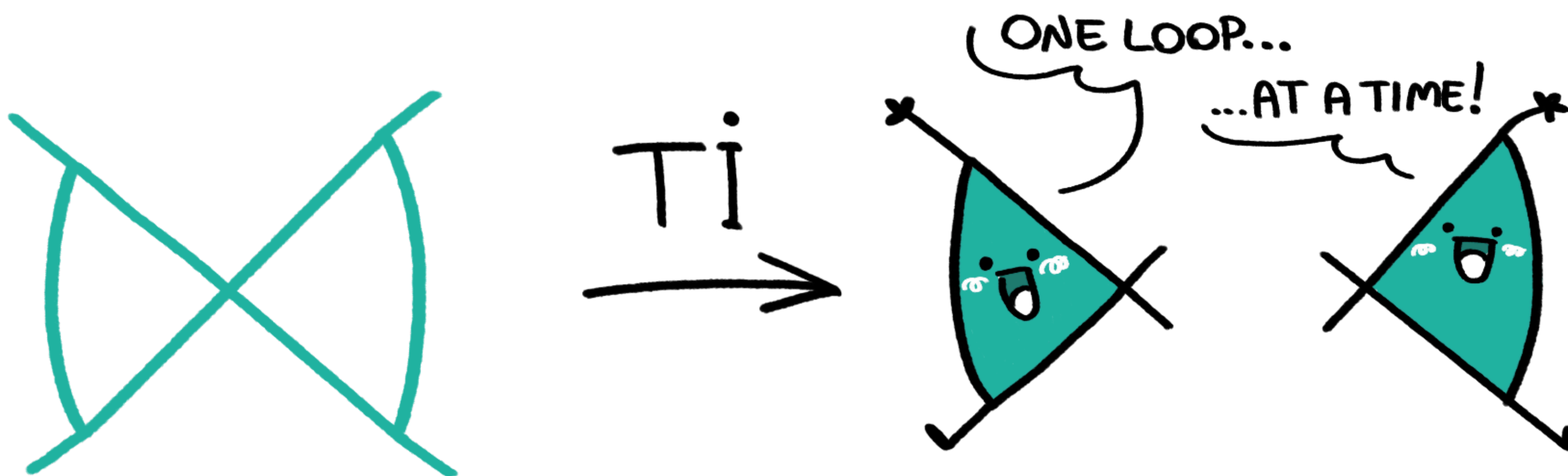
Given a family, map its sectors with fewer external legs (or that are factorizable into fewer loops products) to new families having fewer invariants & fewer irreducible scalar products
⇒ simpler identities

- Already used in tensor/ integrand reduction and numerical unitarity
- Impact on IBP reduction still unexplored

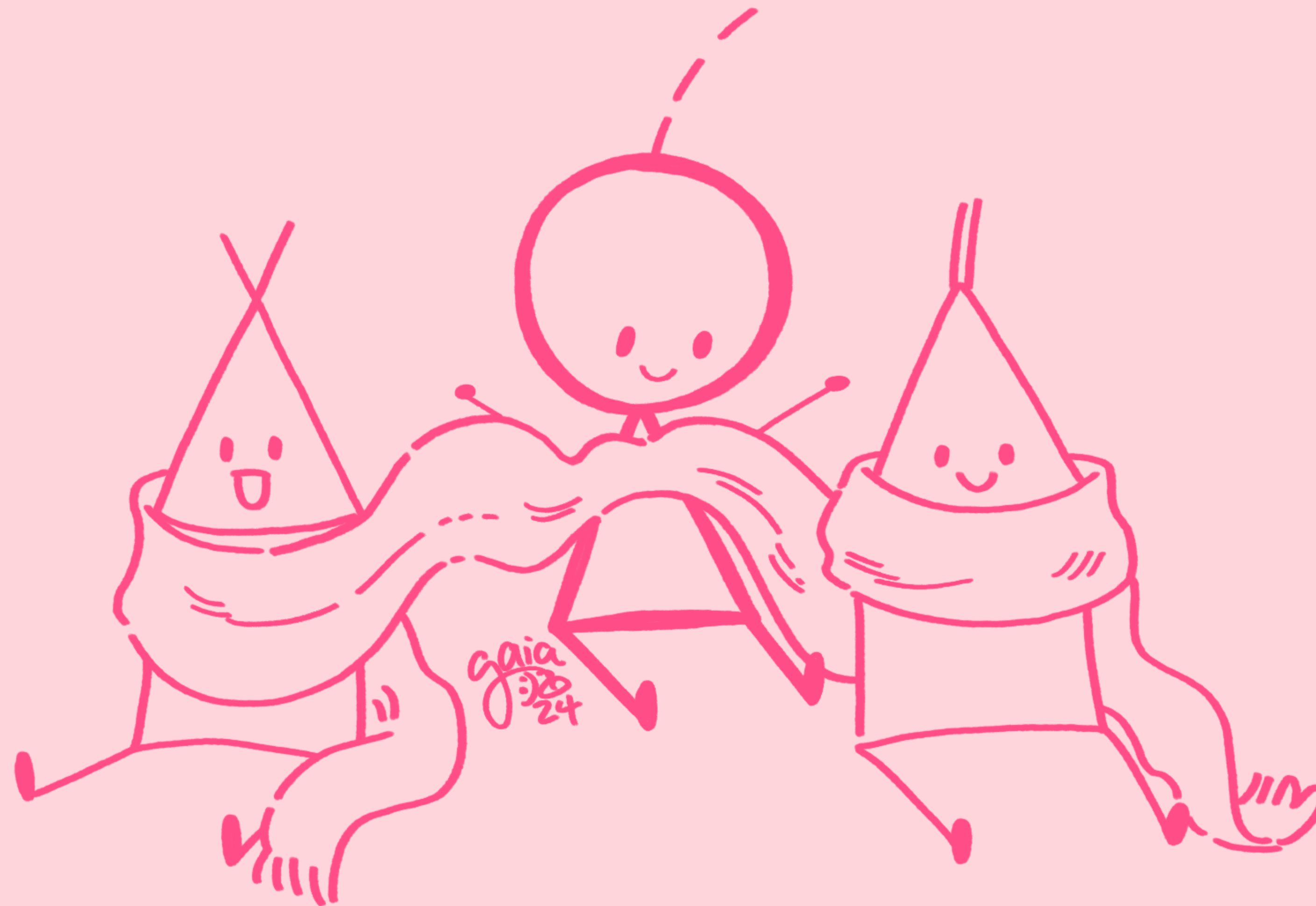
Application/1 → tested on cutting edge examples



Application/2 → only tested in simple cases (for now!)



Part 2: Practical example

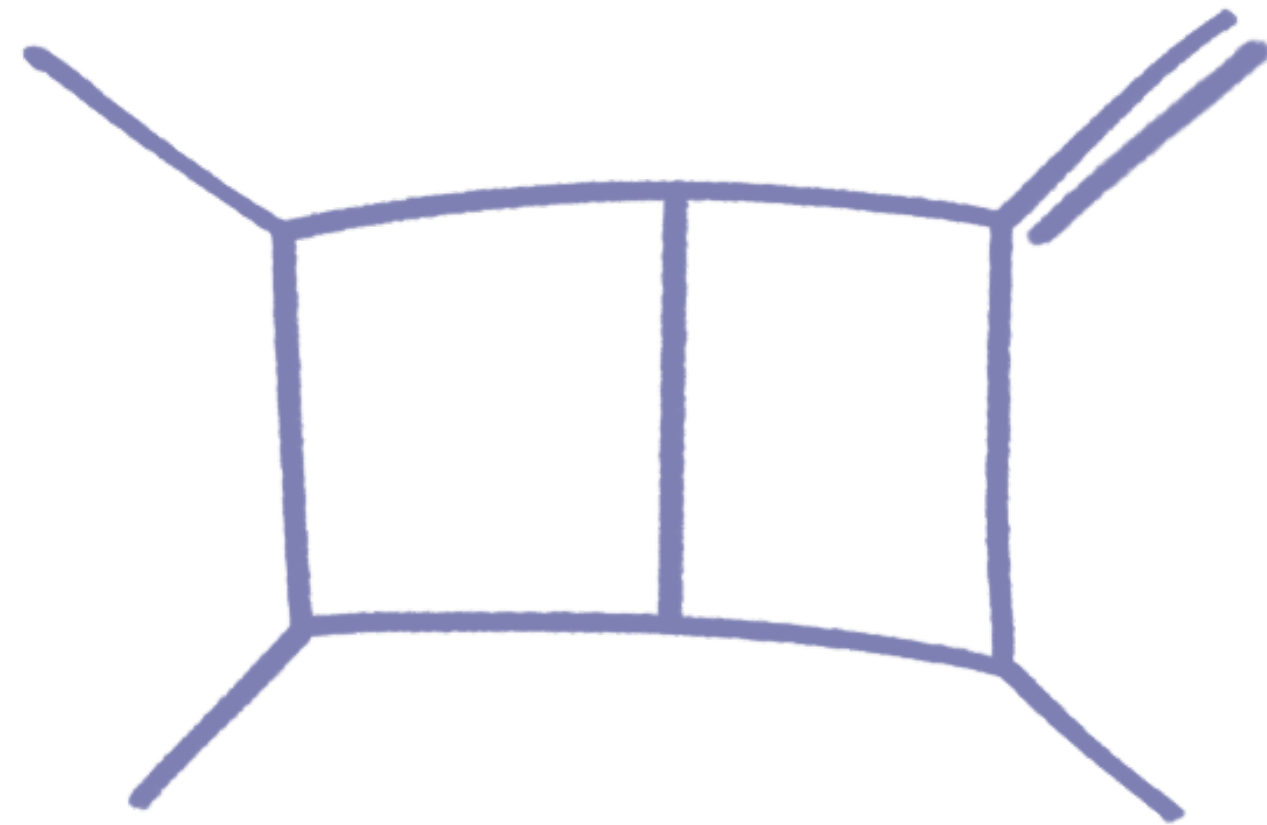


Practical example

Double box family with one external mass integral

$$s = (p_1 + p_2)^2, \quad t = (p_1 + p_3)^2, \quad m^2 = p_4^2, p_1^2 = p_2^2 = p_3^2 = 0$$

Top sector $S_{\text{db};111111100}$



9 generalised den.s
7 proper denominators
2 ISPs
3 invariants

Double box (db)

$$D_{\text{db},1} = k_1^2$$

$$D_{\text{db},2} = (k_1 + p_1)^2$$

$$D_{\text{db},3} = (k_1 + p_1 + p_2)^2$$

$$D_{\text{db},4} = (k_1 + k_2)^2$$

$$D_{\text{db},5} = k_2^2$$

$$D_{\text{db},6} = (k_2 - p_1 - p_2 - p_3)^2$$

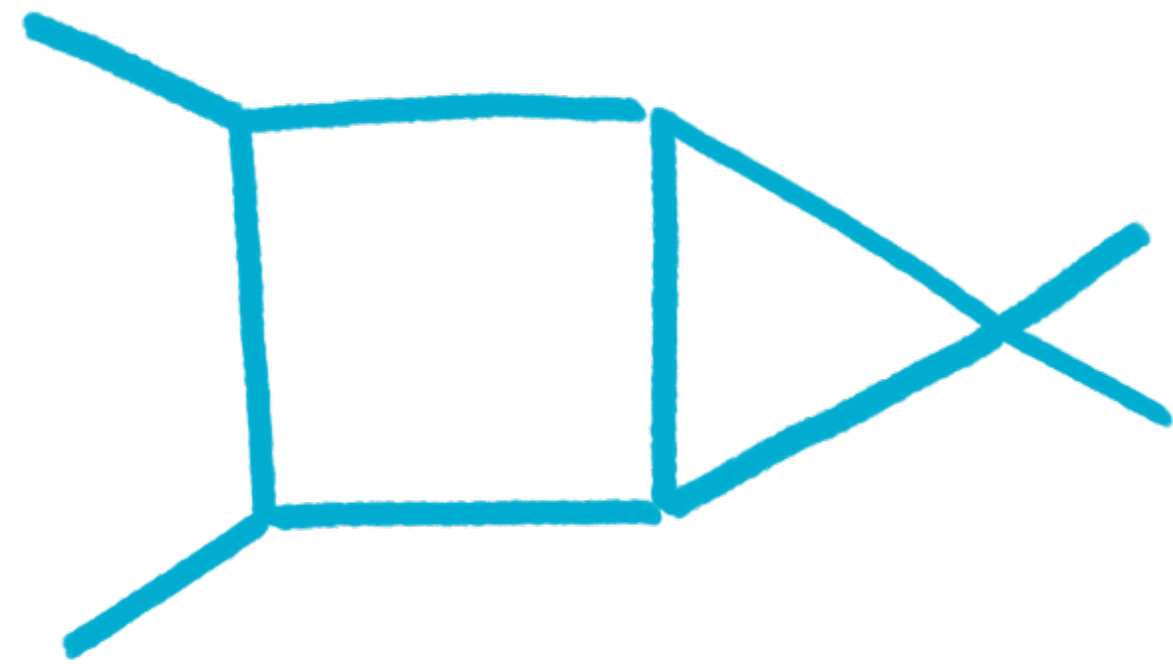
$$D_{\text{db},7} = (k_2 - p_1 - p_2)^2$$

$$D_{\text{db},8} = k_2 \cdot p_1$$

$$D_{\text{db},9} = k_1 \cdot (-p_1 - p_2 - p_3)$$

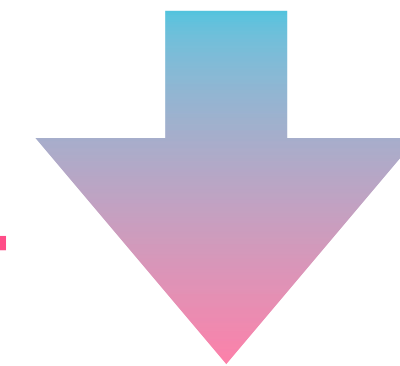
Sector with fewer ext legs

Sector $S_{db;11110100}$

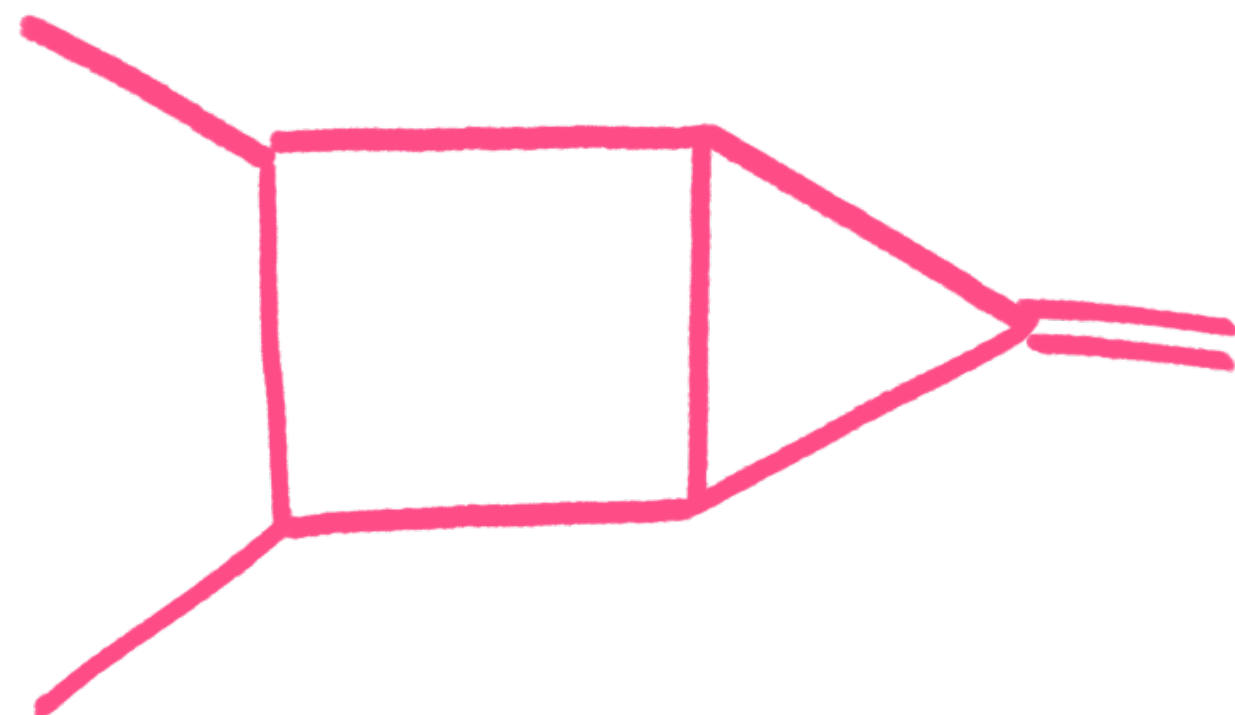


$$a_{6,8,9} \leq 0.$$

9 generalised den.s
6 proper denominators
3 ISPs
3 invariants



BUT if we consider the boxtriangle as a **NEW family = TI family ...**



7 generalised den.s
6 proper denominators
1 ISPs
1 invariant

Box triangle (bt)

$$D_{bt,1} = k_1^2$$

$$D_{bt,2} = (k_1 + p_1)^2$$

$$D_{bt,3} = (k_1 + p_1 + p_2)^2$$

$$D_{bt,4} = (k_1 + k_2)^2$$

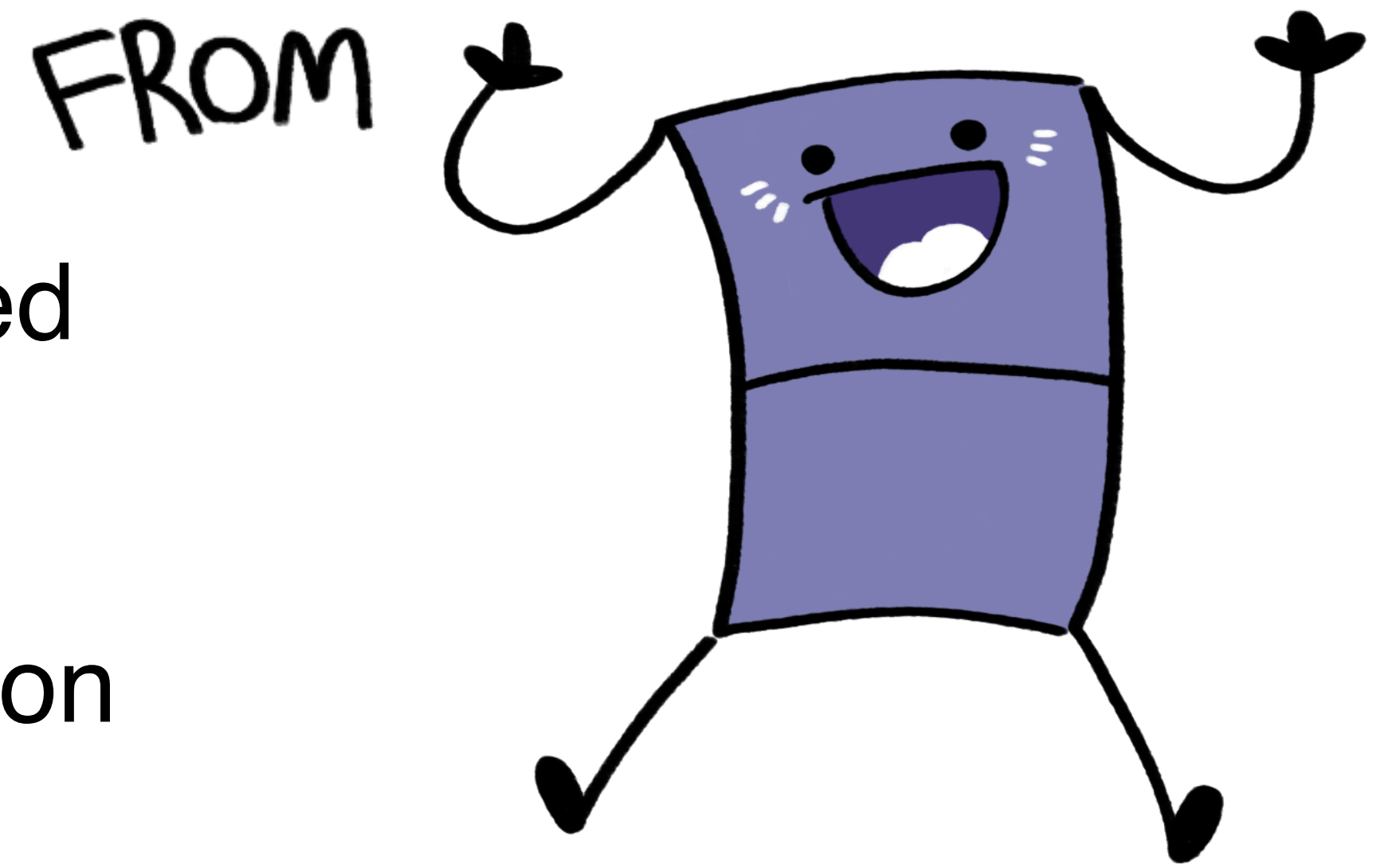
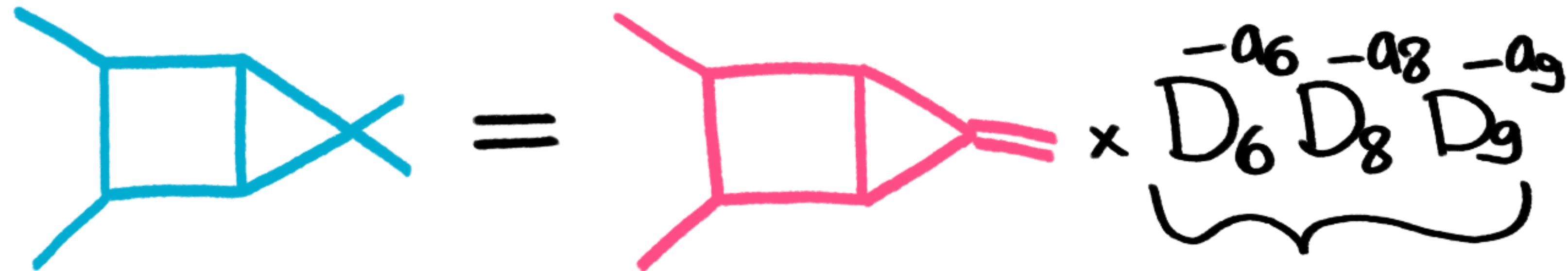
$$D_{bt,5} = k_2^2$$

$$D_{bt,6} = (k_2 - p_1 - p_2)^2$$

$$D_{bt,7} = k_2 \cdot p_2$$

We have the map...

$$I_{\text{db}; a_1 a_2 a_3 a_4 a_5 a_6 a_7 a_8 a_9} = I_{\text{bt}; a_1 a_2 a_3 a_4 a_5 a_7 0} [D_{\text{db},6}^{-a_6} D_{\text{db},8}^{-a_8} D_{\text{db},9}^{-a_9}]$$

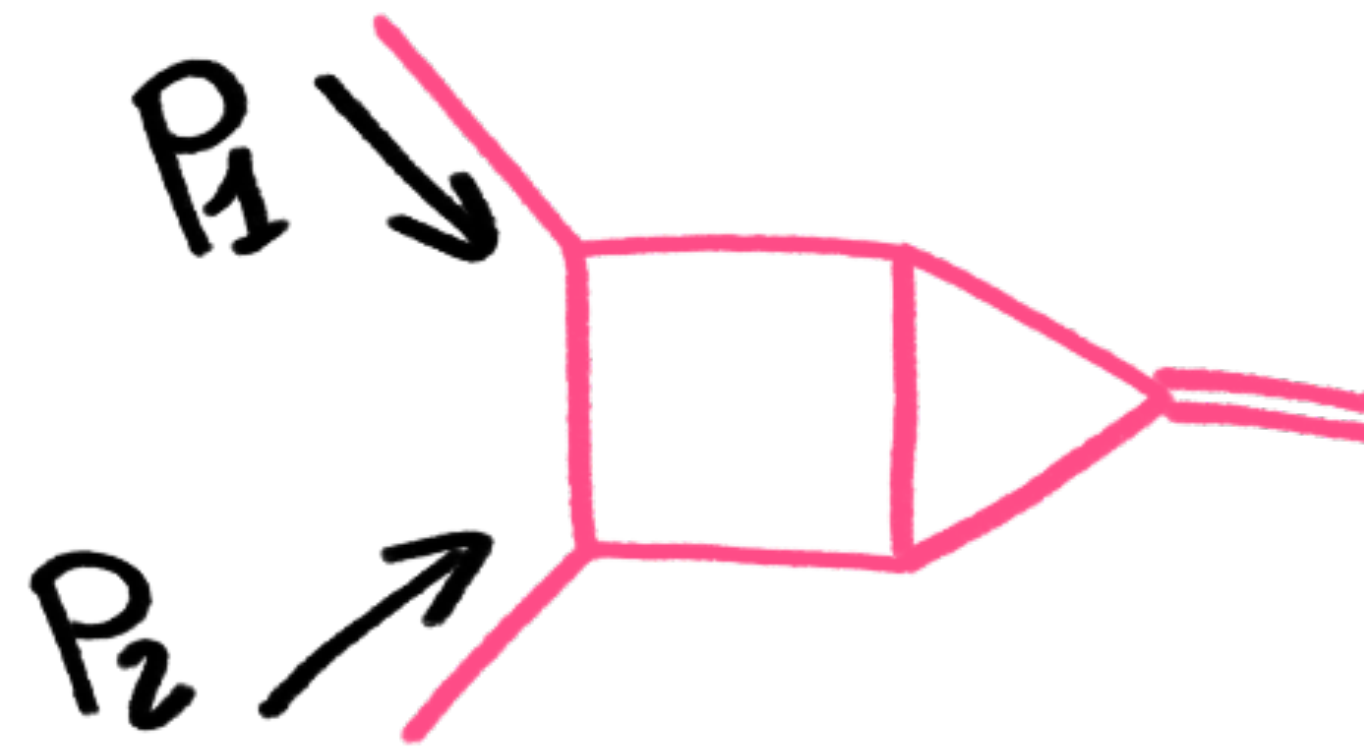


- Numerator needs to be mapped to generalised denominators of new family **bt**
- Mapping can be done via transverse integration

How to do transverse integration

Decomposition of a vector in parallel and transverse component

$$v^\mu = v_{\parallel}^\mu + v_{\perp}^\mu, \quad v_{\parallel}^\mu = c_1 p_1^\mu + c_2 p_2^\mu$$

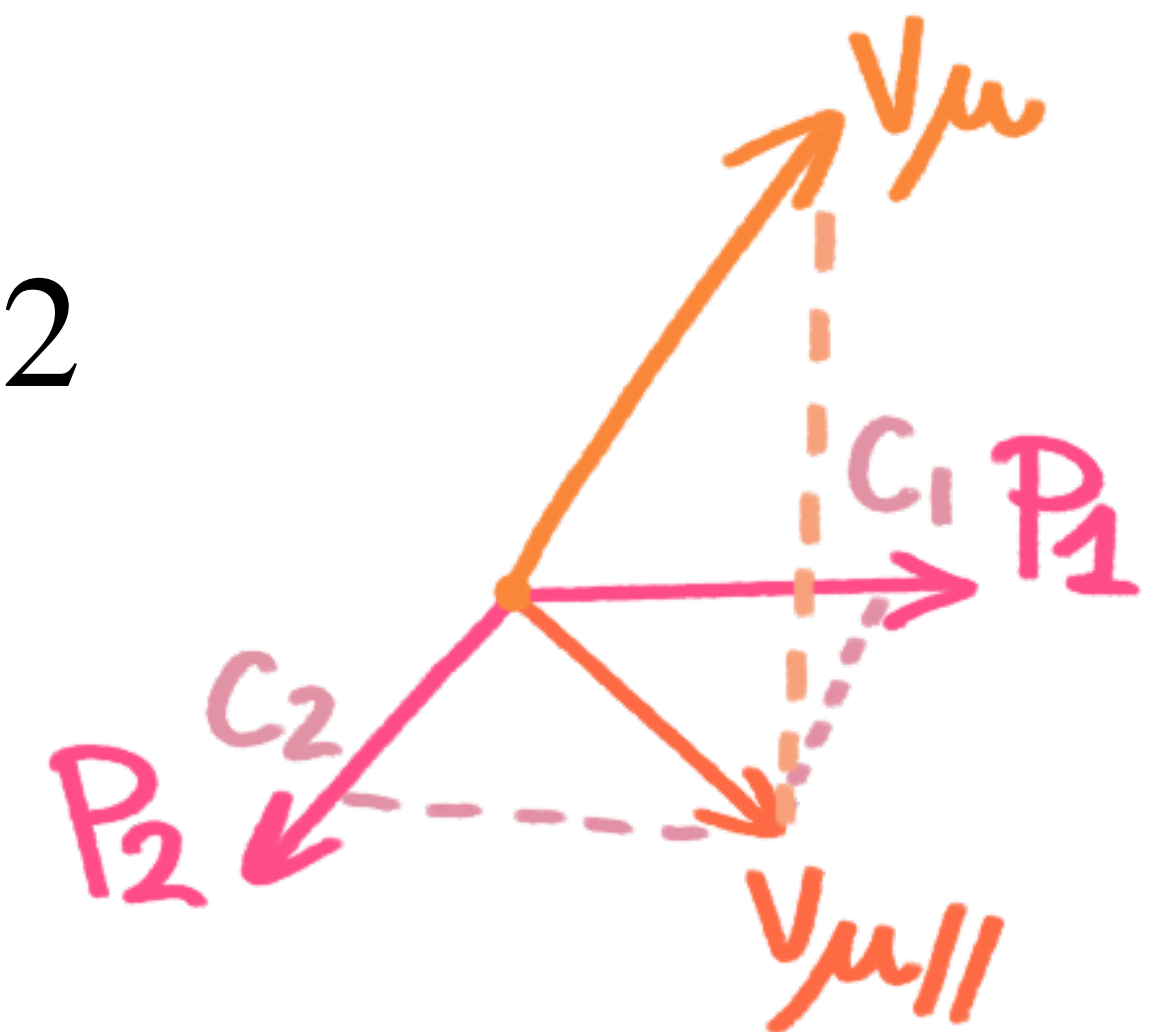


Parallel space spanned by the external legs of the new **bt** family

With

$$v_{\perp} \cdot p_i = 0, \quad v \cdot p_i = v_{\parallel} \cdot p_i \quad \text{for } i = 1, 2$$

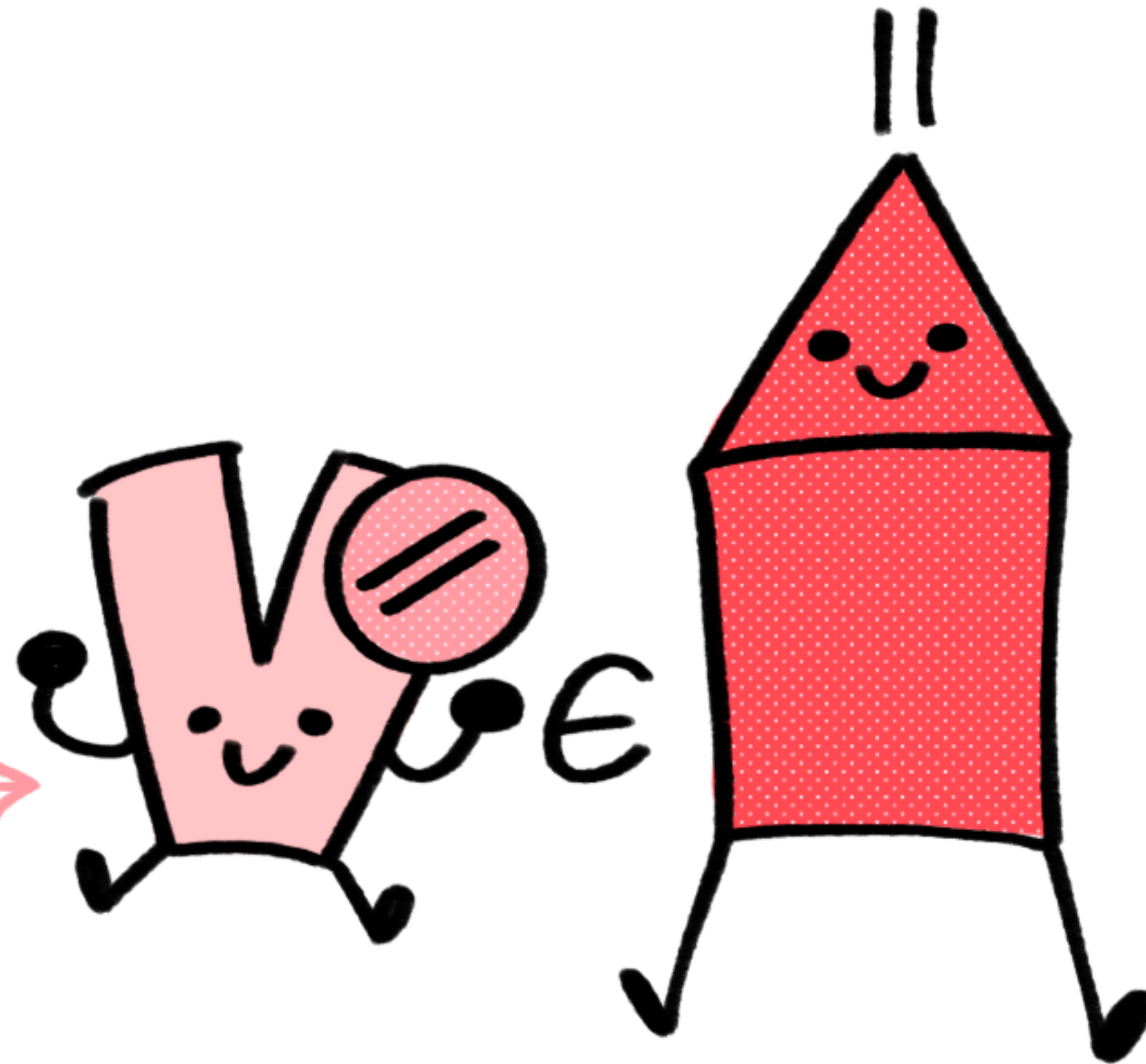
$$p_{1,\perp}^\mu = p_{2,\perp}^\mu = 0$$



Coefficients of the parallel space decomposition found as

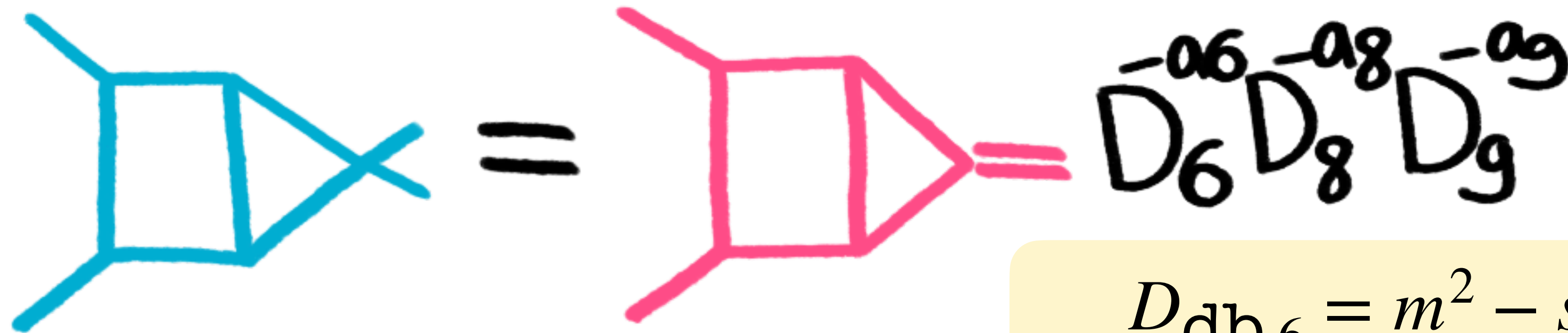
$$\begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \frac{2}{s} \begin{pmatrix} p_2 \cdot v \\ p_1 \cdot v \end{pmatrix} \quad v_{\parallel}^{\mu} = c_1 p_1^{\mu} + c_2 p_2^{\mu}$$

$$v^{\mu} = v_{\parallel}^{\mu} + v_{\perp}^{\mu}$$



In practice

First, rewrite the extra scalar products of **db** as functions of the ones of **bt**



$$D_{\text{db},6} = m^2 - s + D_{\text{bt},6} - 2(k_2 \cdot p_3)$$

$$D_{\text{db},8} = s/2 + D_{\text{bt},5}/2 - D_{\text{bt},6}/2 - D_{\text{bt},7}$$

$$D_{\text{db},9} = s/2 + D_{\text{bt},1}/2 - D_{\text{bt},3}/2 - (k_1 \cdot p_3)$$

We are left with integrals of the family **bt** of the form

$$I_{\text{bt};\vec{a}}[(k_1 \cdot p_3)^{\beta_1} (k_2 \cdot p_3)^{\beta_2}] \times \left(\frac{k_1 \cdot p_3}{k_2 \cdot p_3} \right)^{\beta_1} \left(\frac{k_2 \cdot p_3}{k_1 \cdot p_3} \right)^{\beta_2}$$

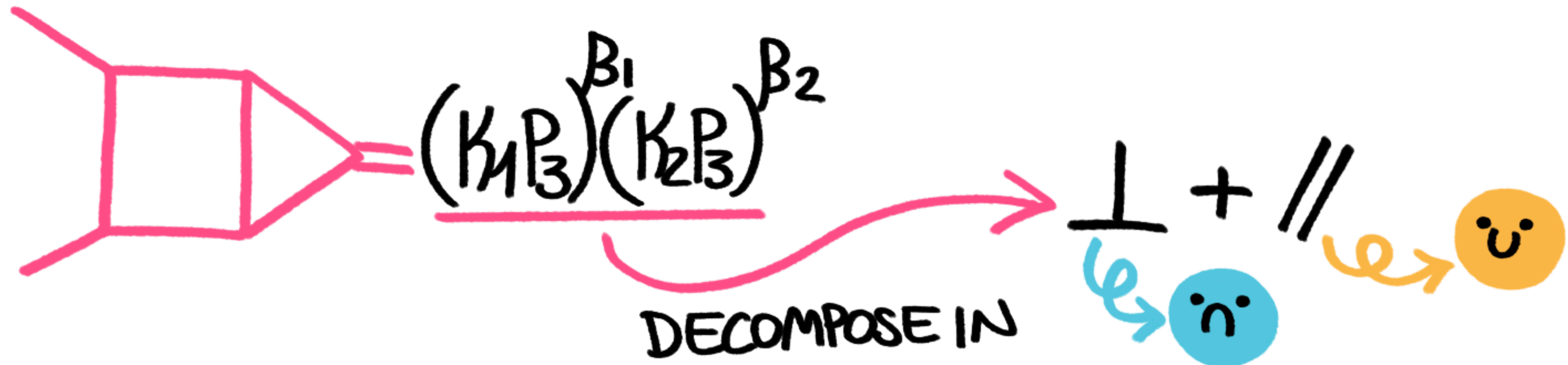
Rewrite the scalar products as

$$(k_1 \cdot p_3) = (k_1 \cdot p_{3,\parallel}) + (k_{1,\perp} \cdot p_3)$$

$$(k_2 \cdot p_3) = (k_2 \cdot p_{3,\parallel}) + (k_{2,\perp} \cdot p_3)$$

First RHS term becomes

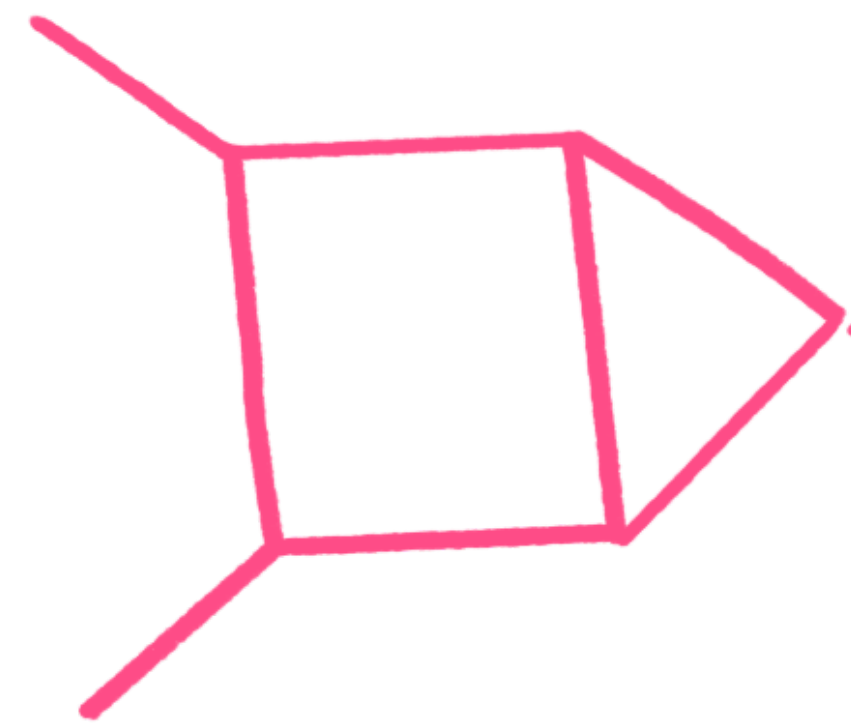
$$(k_i \cdot p_{3,\parallel}) = \frac{2}{s} \left((k_i \cdot p_1)(p_2 \cdot p_3) + (k_i \cdot p_2)(p_1 \cdot p_3) \right)$$



Only scalar products remaining are $(k_{1,\perp} \cdot p_3)$ & $(k_{2,\perp} \cdot p_3)$

$$I_{\text{bt};\vec{a}}[(k_{1,\perp} \cdot p_3)^{\beta_1} (k_{2,\perp} \cdot p_3)^{\beta_2}] = p_{3\mu_1} \cdots p_{3\mu_{\beta_1}} p_{3\nu_1} \cdots p_{3\nu_{\beta_2}} I_{\text{bt};\vec{a}}[k_{1,\perp}^{\mu_1} \cdots k_{1,\perp}^{\mu_{\beta_1}} k_{2,\perp}^{\nu_1} \cdots k_{2,\perp}^{\nu_{\beta_2}}],$$

Tensor integrals can be decomposed in products of tensors and form factors



$$= [k_i^\mu k_j^\nu \dots] = \sum_j C_j T_j^{\mu\nu\dots}$$

Tensor basis only depends on $g_{\perp}^{\mu\nu}$ and not on ext legs!!

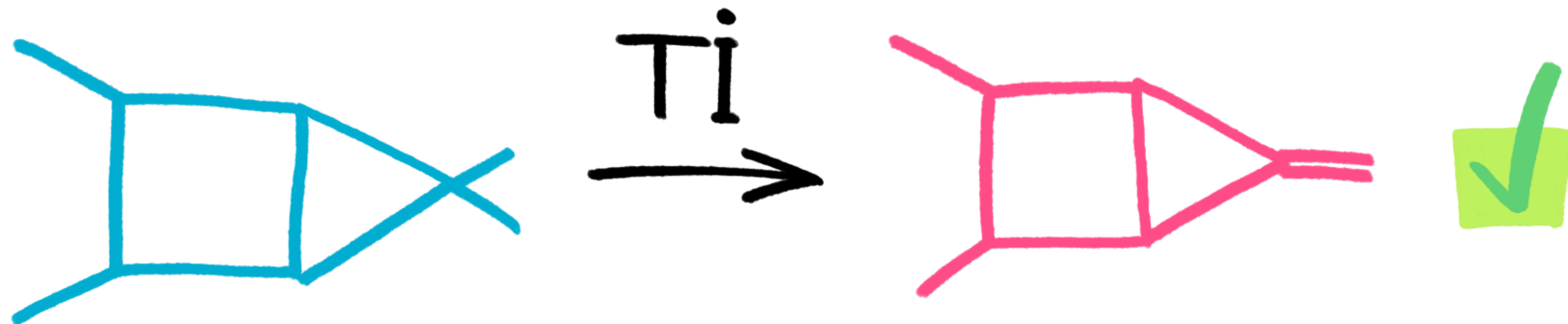
$$I_{\text{bt};\vec{a}}[k_{1,\perp}^{\mu_1} \cdots k_{1,\perp}^{\mu_{\beta_1}} k_{2,\perp}^{\nu_1} \cdots k_{2,\perp}^{\nu_{\beta_2}}] = \sum_j C_j T_{j,\perp}^{\mu_1 \cdots \mu_{\beta_1} \nu_1 \cdots \nu_{\beta_2}}$$

After this step we have only scalar products of $(k_{i,\perp} \cdot k_{j,\perp})$

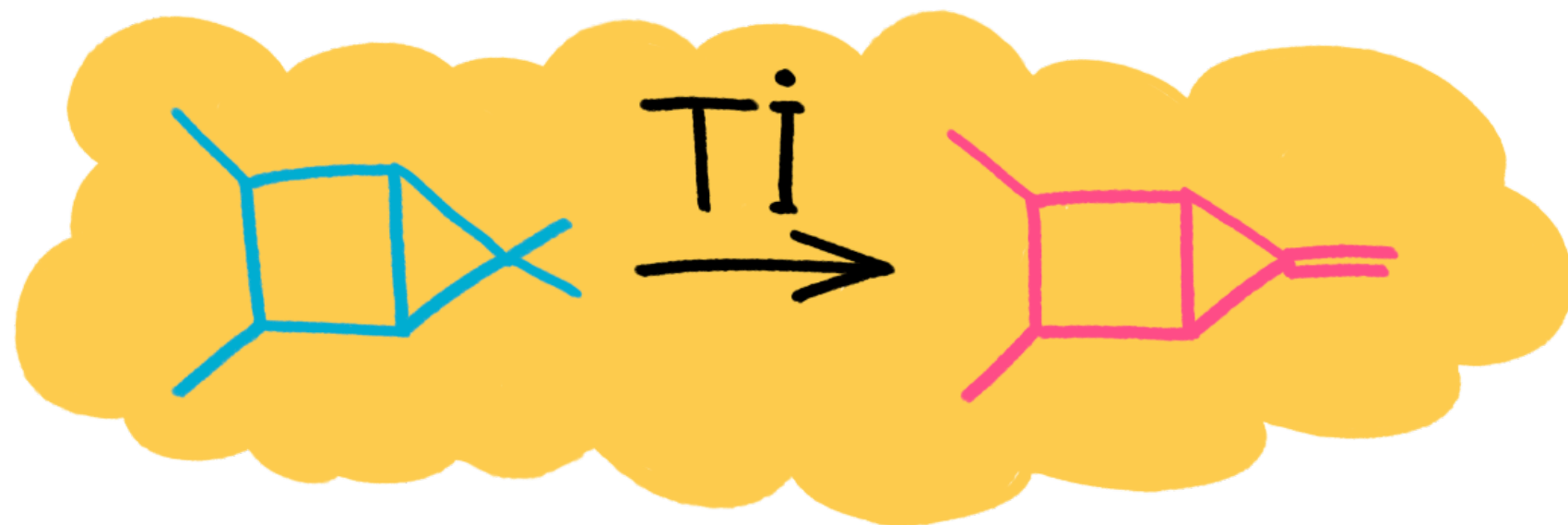
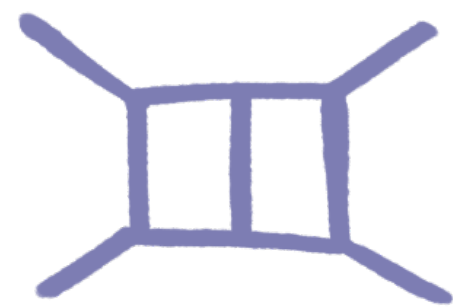
That we can rewrite using $(k_{i,\perp} \cdot k_{j,\perp}) = (k_i \cdot k_j) - (k_{i,\parallel} \cdot k_{j,\parallel})$

with $(k_{i,\parallel} \cdot k_{j,\parallel}) = \frac{2}{s} \left((k_i \cdot p_1)(k_j \cdot p_2) + (k_i \cdot p_2)(k_j \cdot p_1) \right)$

Successfully mapped **db** in **bt**



Flowchart



Ext. Legs  < Ext. Legs 

$$\text{Dens}(\text{blue box}) = f(\text{Dens}(\text{pink box}))$$

$$\text{pink box} = (k_i \cdot p_j), p_j \notin \text{Ext. Legs}(\text{pink box})$$

$$k_i \cdot p_j = k_i \cdot p_{j\parallel} + k_{i\perp} \cdot p_j$$

$$(k_{i\perp} \cdot p_j)^2 = k_{i\perp}^\mu k_{i\perp}^\nu \dots p_{j\mu} p_{j\nu} \dots$$

$$k_{i\perp}^\mu k_{i\perp}^\nu \dots = \sum_j C_j T_j^{\mu\nu \dots}$$

$$k_{i\perp} \cdot k_{j\perp}$$

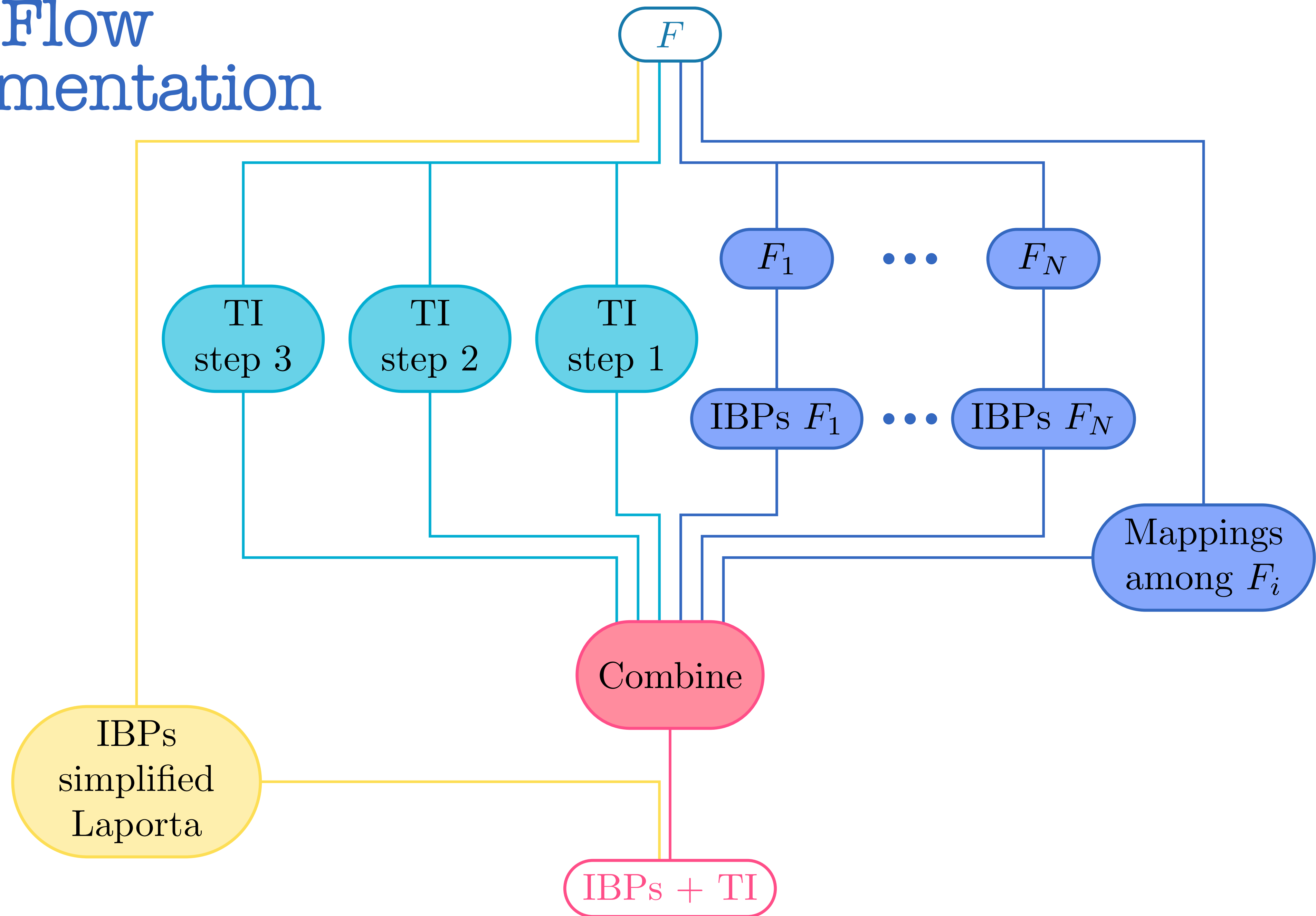
$$k_{i\perp} k_{j\perp} = k_i k_j - k_{i\parallel} k_{j\parallel}$$



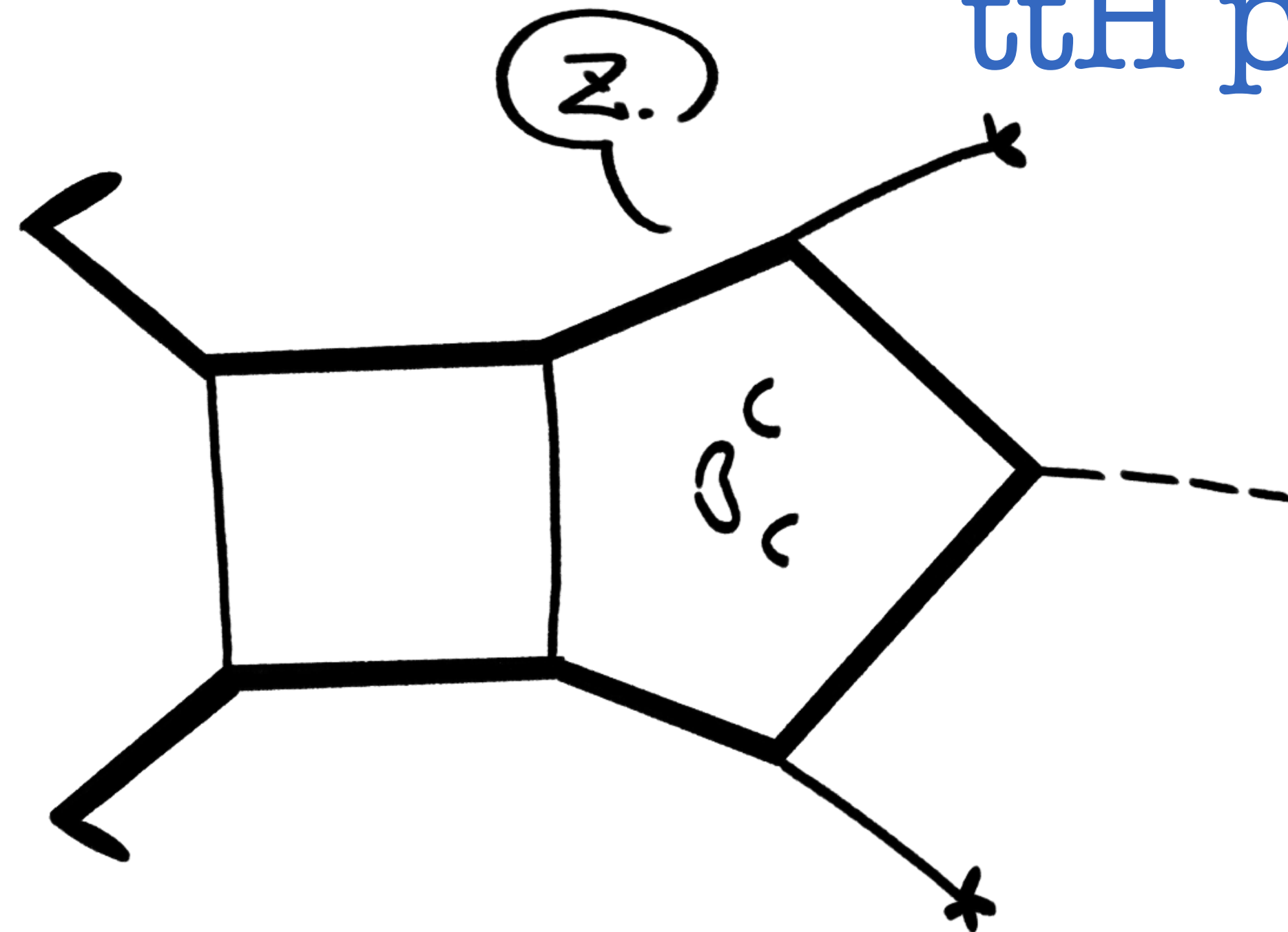
Part 3: Implementation & benchmarks



FiniteFlow implementation

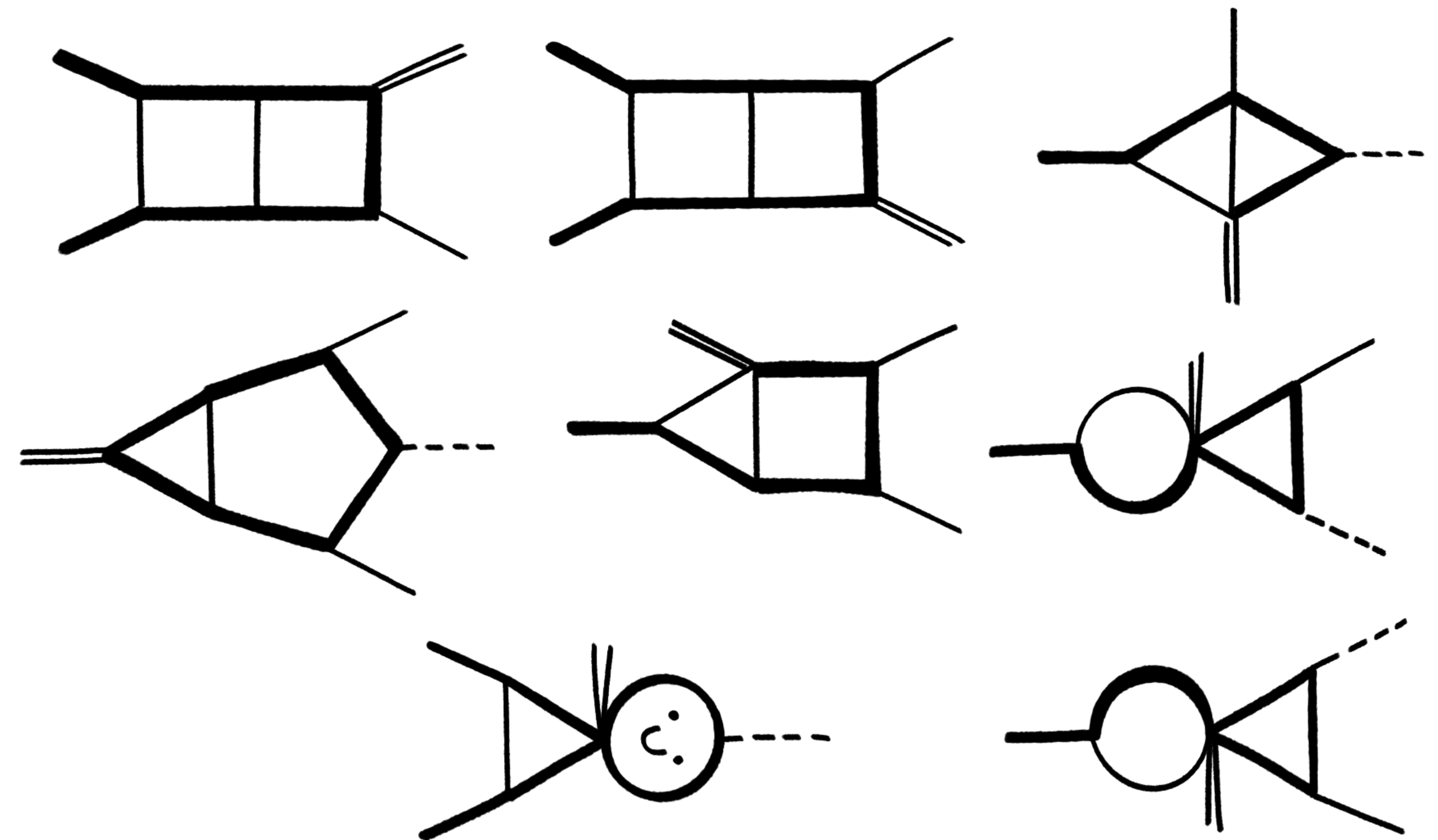


ttH pentabox



All integrals up to rank 5

TI families with 4 external legs



Speedup vs traditional Laporta: 2.3x

Breakdown:

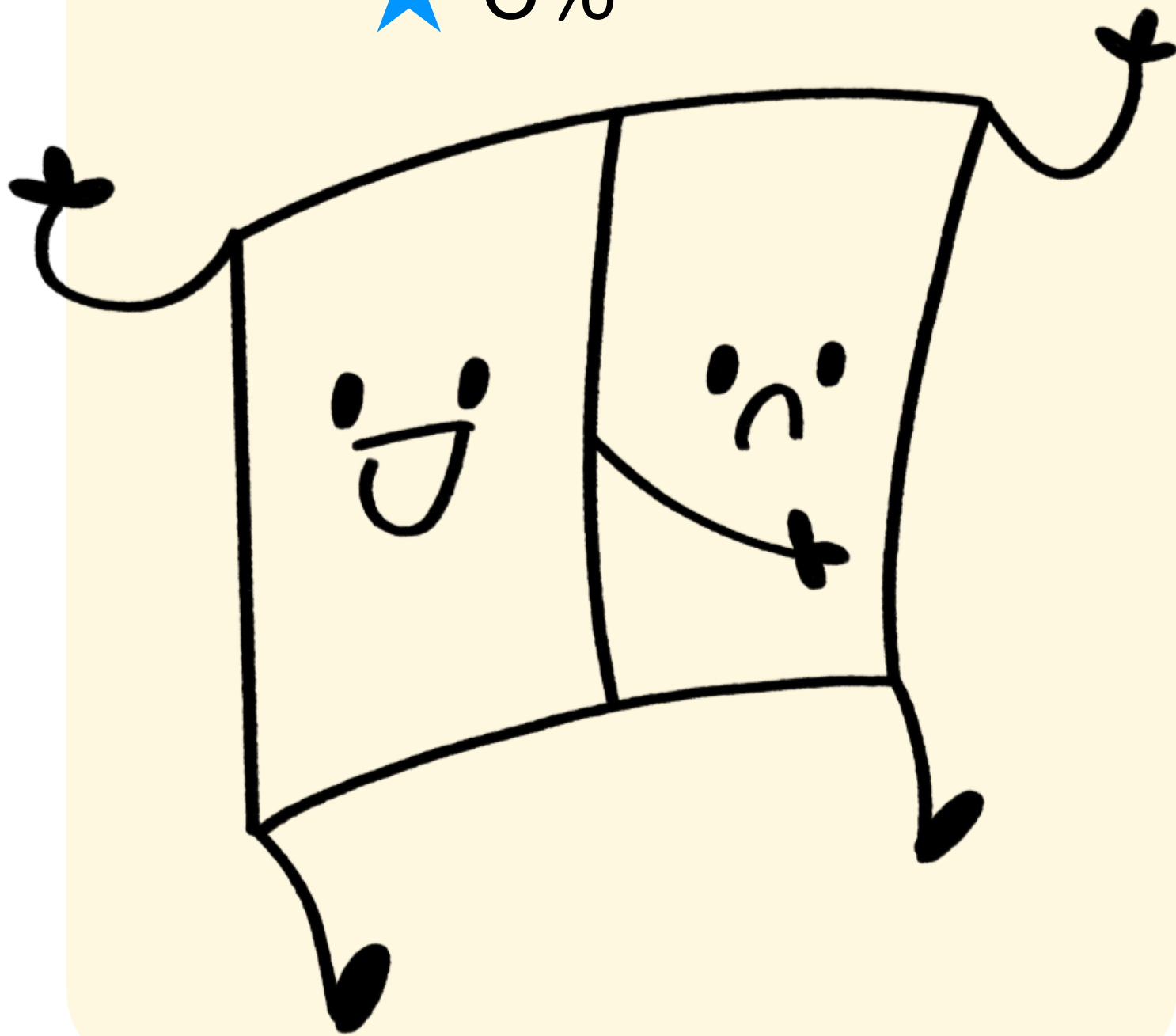
- ★ solving the simplified Laporta IBP system: 48%
- ★ evaluating the coefficients of the TI identities: 22%
- ★ evaluating the solution of the IBP system for the TI families: 23%

More examples ...

Double pentagon

- Speedup: 3.3x

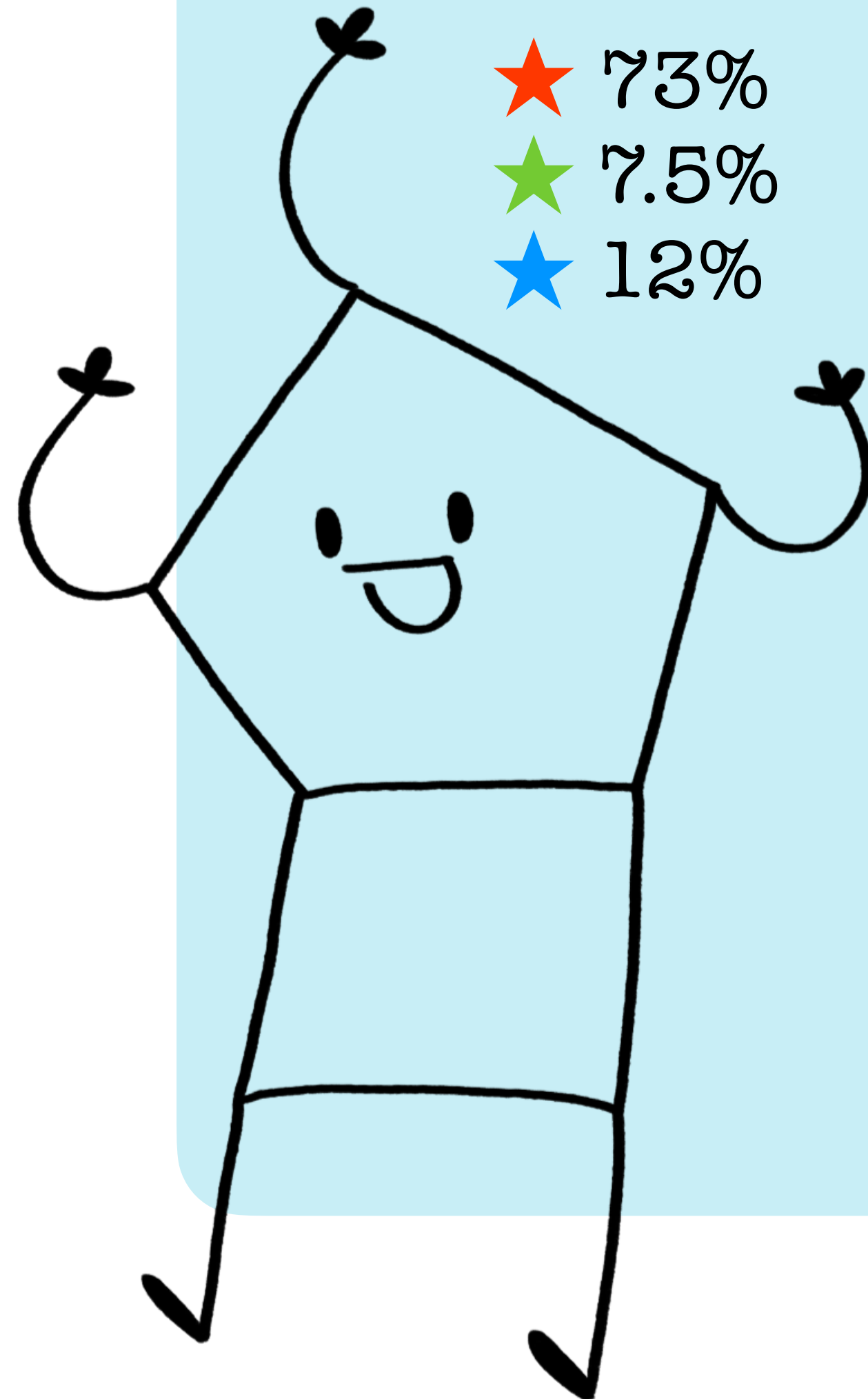
- ★ 70%
- ★ 21%
- ★ 6%



Massless pentabox

- Speedup: 4x

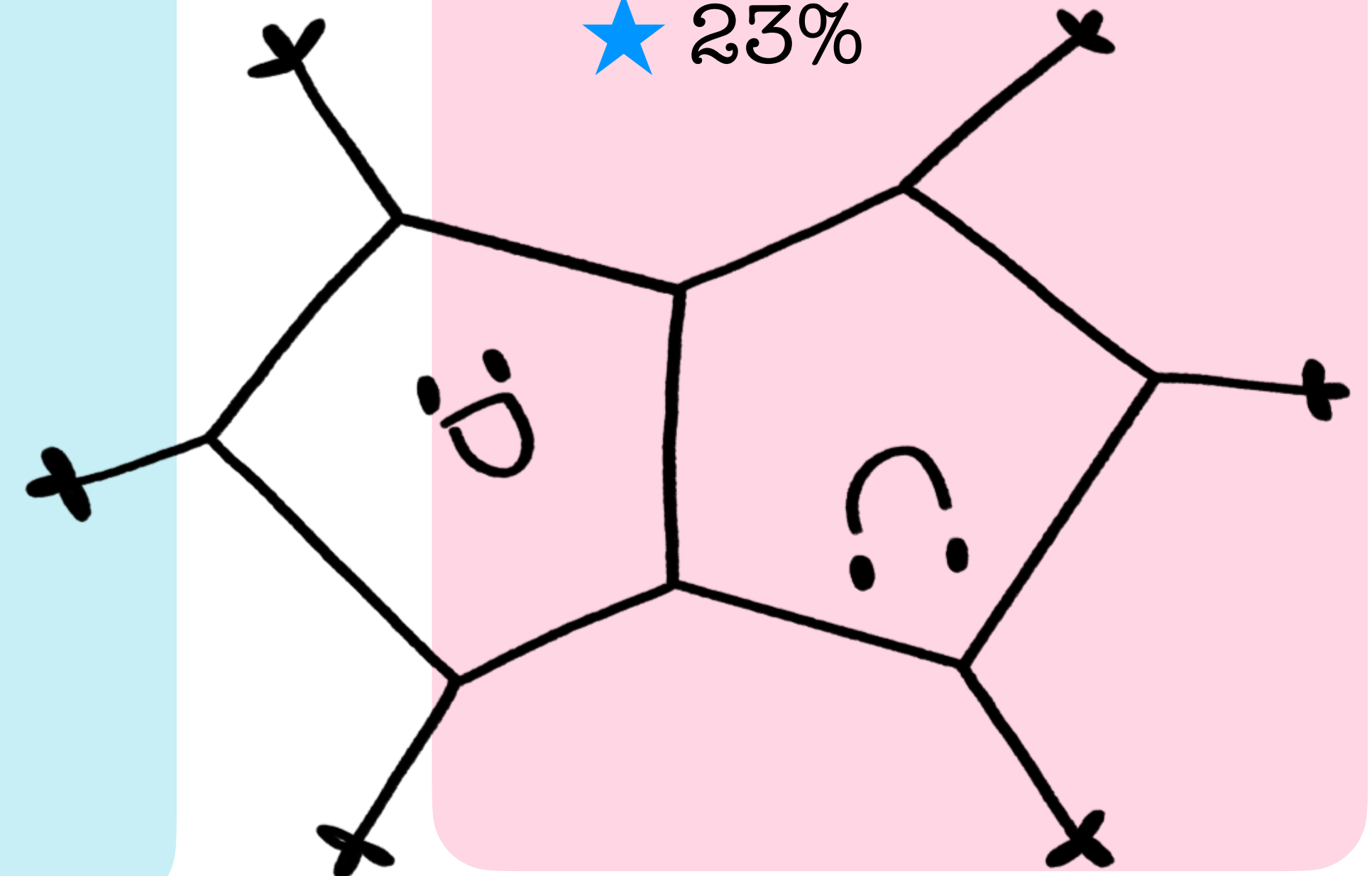
- ★ 73%
- ★ 7.5%
- ★ 12%



Ladybug

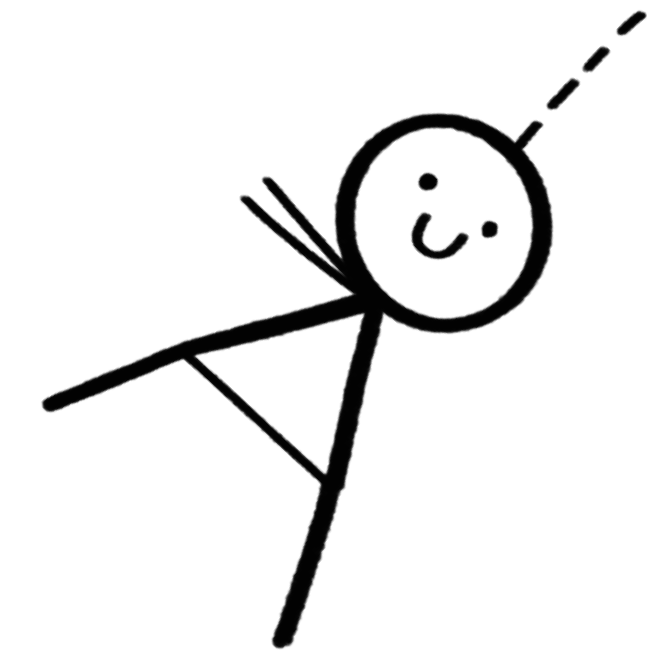
- Speedup: 2.7x

- ★ 48%
- ★ 22%
- ★ 23%

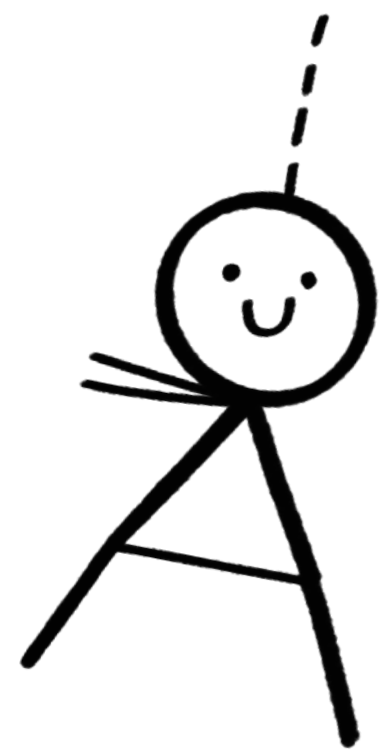


Conclusions...

- IBP reduction: key point ingredient of calculations
- Bottleneck for state-of-the-art precision predictions
- **Transverse integration**

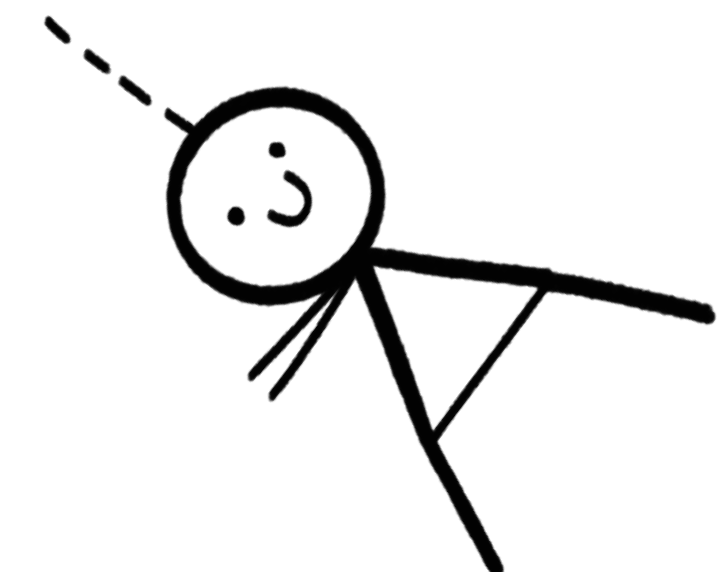


- Build simpler identities to feed into the IBP system
- Map into a new family with fewer invariants and fewer ISPs :
⇒ **easier identities**
- Substantial performance improvements in cutting-edge examples



...& outlook

- Combination with syzygy techniques
- Implementation of factorizable sectors
- Optimizations + release of public package



Thank you for
your attention!

