## **Precision in EFT predictions**

### Eleni Vryonidou University of Manchester









European Research Council Established by the European Commission

HP2, Torino 12/09/24

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### A model independent probe of heavy New Physics



## measurements at low energy.



neory 
$$\mathcal{L}_{SM}(\phi) + |\mathcal{L}_{dim6}(\phi)| + \dots$$

Effective Field Theory reveals high energy physics through precise



## SMEFT The global aspect



- 50 coefficients fitted: bounds varying between operators
- Most Wilson coefficient bounds below 1 for  $\Lambda = 1$  TeV
- No sign of significant deviations from the SM
- Least constrained coefficients are 4-top operators

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SMEFiT3.0 Celada, Giani, Mantani, Rojo, Rossia, Thomas, EV, ter Hoeve arXiv:2404.12809 HP2, 12/9/24



## SMEFT Not just a theorists' tool

### **ATLAS CONF-2023-052**



Higgs+EW

### **CMS-PAS-TOP-22-006**

### Top sector $t\bar{t}l\nu, t\bar{t}l\bar{l}, tl\bar{l}q, tHq$





## **EFT pathway to New Physics**







 $\frac{1}{2}\sum c_i^6(\mu)a_{n,i}^6(\mu) + 0$ Precise EFT predictions

### Precise SM predictions

Precise experimental measurements













### Constraints

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Precise experimental measurements

$$\frac{1}{2}c_{i}^{6}(\mu)$$



## **EFT pathway to New Physics**

 $\Delta Obs_n$ 



EXP

### Constraints $\frac{1}{\Lambda 2}c_i^6(\mu)$

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 $Obs_n^{SM} = \frac{1}{\Lambda^2} \sum c_i^6(\mu) a_{n,i}^6(\mu) + O\left(\frac{1}{\Lambda^4}\right)$ Precise EFT predictions Precise SM predictions

Precise experimental measurements



### **Aspects of EFT predictions** And how to improve them

- \* Higher Orders in  $1/\Lambda^4$ 
  - \* squared dim-6 contributions
  - \* double insertions of dim-6
  - \* dim-8/10... contributions

\* Higher Orders in QCD and EW \* EFT is a QFT, renormalisable order-by-order  $1/\Lambda^2$  $\mathcal{O}(\alpha_s, \alpha_{ew}) + \mathcal{O}\left(\frac{1}{\Lambda^2}\right) + \mathcal{O}\left(\frac{\alpha_s}{\Lambda^2}\right) + \mathcal{O}\left(\frac{\alpha_{ew}}{\Lambda^2}\right)$ 





## Why bother with higher orders?

Higher orders in SMEFT bring:

- **\*** Accuracy
- \* Precision
- Improved sensitivity \*
  - \* Accurate knowledge of the deviations (distribution shapes, correlations between observables, etc.) can be the key to disentangle them from the SM.

\* Loop-induced new sensitivity: operators entering at one-loop



## Accuracy and precision: QCD **Example 1: k-factors and shape**



Different shapes at NLO

Degrande, Maltoni, Mimasu, EV, Zhang arXiv:1804.07773

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	13 TeV	$\sigma$ NLO	К
	$\sigma_{SM}$	$0.507^{+0.030+0.000+0.007}_{-0.048-0.000-0.008}$	1.09
	$\sigma_{t\phi}$	$-0.062\substack{+0.006+0.001+0.001\\-0.004-0.001-0.001}$	1.13
	$\sigma_{\phi G}$	$0.872\substack{+0.131+0.037+0.013\\-0.123-0.035-0.016}$	1.39
	$\sigma_{tG}$	$0.503^{+0.025+0.001+0.007}_{-0.046-0.003-0.008}$	1.07
	$\sigma_{t\phi,t\phi}$	$0.0019\substack{+0.0001+0.0001+0.0000\\-0.0002-0.0000-0.0000}$	1.17
	$\sigma_{\phi G,\phi G}$	$1.021\substack{+0.204+0.096+0.024\\-0.178-0.085-0.029}$	1.58
	$\sigma_{tG,tG}$	$0.674\substack{+0.036+0.004+0.016\\-0.067-0.007-0.019}$	1.04
	$\sigma_{t\phi,\phi G}$	$-0.053\substack{+0.008+0.003+0.001\\-0.008-0.004-0.001}$	1.42
	$\sigma_{t\phi,tG}$	$-0.031\substack{+0.003+0.000+0.000\\-0.002-0.000-0.000}$	1.10
	$\sigma_{\phi G,tG}$	$0.859^{+0.127+0.021+0.017}_{-0.126-0.020-0.022}$	1.37
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Maltoni, EV, Zhang arXiv:1607.05330



ttH

Linear

Quadratic



## Accuracy and precision: QCD Example 2: more k-factors and shapes



El Faham, Pelliccioli, EV arXiv:2405.19083

Different K-factors for different operators, different from the SM HP2, 12/9/24

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Celada, Durieux, Mimasu, EV arXiv:2407.09600



## Accuracy and precision: QCD Example 2: more k-factors and shapes



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## **Does NLO QCD change global fits?** Global top fits



SMEFiT3.0 Celada, Giani, Mantani, Rojo, Rossia, Thomas, EV, ter Hoeve arXiv:2404.12809 E. Vryonidou HP2, 12/9/24

Posterior distributions for Wilson coefficients

Blue: LO

Orange: NLO

Significant impact of NLO for some operators

Linear level: NLO resolves non-interference problem for colour singlet 4-fermion operators

NLO is not necessarily more constraining but it is more reliable

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## Where can higher orders help? **Poorly constrained operators**







95% Confidence Level Bounds (1/TeV<sup>2</sup>)

SMEFiT3.0 Celada, Giani, Mantani, Rojo, Rossia, Thomas, EV, ter Hoeve arXiv:2404.12809

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### Only accessible through 4-top and *ttbb* measurements



### **Can precision help** improve this picture?

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Aoude, El Faham, Maltoni, EV arXiv:2208.04962





Aoude, El Faham, Maltoni, EV arXiv:2208.04962



## New probes 4-heavy operators in top pair production



Degrande, Durieux, Maltoni, Mimasu, EV, Zhang arXiv:2008.11743

### Complementary information to ttbb and 4top production HP2, 12/9/24

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Degrande, Rosenfeld, Vasquez arXiv:2402.06528



## Improved sensitivity due to EW loops





95% CL limits on 3<sup>rd</sup> generation 4-fermion operators



Dawson and Giardino arXiv: 2201.09887

New loop-induced sensitivities: EWPO and Higgs production

Bounds are competitive to 4top production

### 4-heavy operators in Higgs production

Alasfar, de Blas, Gröber arXiv:2202.02333

- A combination of all probes can pin these coefficients down!
  - HP2, 12/9/24

## **Towards NLO EW corrections**

EFT effects often pronounced in the tails of distributions

This is where EW corrections are important In the SM NLO EW can be O(10%) at 1/2 TeV

Exact NLO EW corrections are not available in the SMEFT in general (some simple processes studied) At high energy these are described by Sudakov logarithms ~log<sup>2</sup> (s/m<sub>W</sub><sup>2</sup>) and ~log (s/m<sub>W</sub><sup>2</sup>)

Can Sudakov logs be computed and be a good approximation for EW corrections in the SMEFT?



## EW corrections at high energy Formalism

Denner-Pozzorini algorithm

 $\lim_{M_W^2/s\to 0} (\mathcal{M}_{1,\mathrm{EW}}^{\mathrm{SM}})^{i_1\dots i_n}(p_1,\dots,p_n) = \sum_{i_1'\dots i_l'} (\mathcal{M}_0^{\mathrm{SM}})^{i_1'\dots i_n'}(p_1,\dots,p_n) \delta_{\mathrm{EW},\,i_1'i_1\dots i_n'i_n}^{\mathrm{SM}}(p_1,\dots,p_n)$ 



Denner and Pozzorini hep-ph/0010201 and hep-ph/0104127

How about the SMEFT?

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# renormalisation

### Only depends on Born amplitudes and the EW charges of the external legs





## **High-energy EW corrections Top pair production@LHC**



El Faham, Mimasu, Pagani, Severi, EV, Zaro in preparation MC implementation based on Pagani, Zaro arXiv:2110.03714, see also Lindert, Mai 2312.07927 HP2, 12/9/24

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## **High-energy EW corrections Top pair production@LHC**



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## **High-energy EW corrections** DY@LHC

 $\mathcal{O}_{qe} = \sum_{\mathrm{f}=1}^{2} (\overline{q}_{\mathrm{f}} \gamma_{\mu} q_{\mathrm{f}}) (\overline{e} \gamma^{\mu} e)$ 





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# High-energy EW corrections

 $\mathcal{O}_{qe} = \sum_{\mathrm{f}=1}^{2} (\overline{q}_{\mathrm{f}} \gamma_{\mu} q_{\mathrm{f}}) (\overline{e} \gamma^{\mu} e)$ 





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## **High-energy EW corrections Breaking degeneracies**



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## **High-energy EW corrections Breaking degeneracies**



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## Future of global fits

More observables: particle level observables new final states

### **Better EFT predictions**

Higher Orders in  $1/\Lambda^4$ • squared dim-6 contributions • double insertions of dim-6 • dim-8 contributions Higher Orders in QCD and EW  $\mathcal{O}(\alpha_s, \alpha_{ew}) + \mathcal{O}$ 

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### **More/less/different operators:** • different flavour assumptions • UV inspired scenarios

EFT is a QFT, renormalisable order-by order in  $1/\Lambda^2$ 

$$-$$
) +  $\mathcal{O}\left(\frac{\alpha_s}{\Lambda^2}\right)$  +  $\mathcal{O}\left(\frac{\alpha_{ew}}{\Lambda^2}\right)$ 



## **SMEFT** computations at dimension-6

$$\Delta Obs_n = Obs_n^{\mathsf{EXP}} - Obs_n^{\mathsf{SM}} = \sum_i \frac{c_i^6(\mu)}{\Lambda^2} a_{n,i}^6(\mu) + 0$$

NLO QCD & loop-induced: Done (SMEFT@NLO) Degrande, Durieux, Maltoni, Mimasu, EV, Zhang arXiv:2008.11743 http://feynrules.irmp.ucl.ac.be/wiki/SMEFTatNLO

NLO EW: Some examples available, progress towards automating these as well.



$$\mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$





## **SMEFT computations at dimension-6**

$$\Delta Obs_n = Obs_n^{\mathsf{EXP}} - Obs_n^{\mathsf{SM}} = \sum_i \frac{c_i^6(\mu)}{\Lambda^2} a_{n,i}^6(\mu) + 0$$

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NLO EW: Some examples available, progress towards automating these as well.

How about this  $\mu$ ?



$$\mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$



## **Running and mixing in SMEFT**

$$\frac{dc_i(\mu)}{d\log\mu} = \boldsymbol{\gamma}_{ij} \, c_j(\mu)$$

One loop anomalous dimension known: (Alonso) Jenkins et al arXiv:1308.2627, 1310.4838, 1312.2014

Example: Turn on 1 operator at high-scale

Compute effect on top pair cross-section





 $c_{Ou}^{1} = 1$  at 2 TeV

### Aoude, Maltoni, Mattelaer, Severi, EV arXiv:2212.05067 HP2, 12/9/24



## Impact of RGE on constraints

### How does running and mixing impacts the constraints?

### Top sector fit:



More important for differential distributions & measurements with very different scales

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### RGE evolution within MC:

PS by PS point computation of coefficients: dynamical scale e.g.  $H_T/2$ 





### Impact of RGE on constraints How does running and mixing impacts the constraints? Higgs sector fit



Maltoni, Ventura, EV arXiv:2406.06670

See also Battaglia, Grazzini, Spira, Wiesemann arXiv: 2109.02987 Di Noi, Grober arXiv:2312.11327 Di Noi, Grober, Mandal arXiv: 2408.03252

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Incl. ATLAS+CMS	ttH CMS	HH Proj
$- \mu = \sqrt{M_H^2 + P_T^2}$	$\mu = \sqrt{M_H^2 + P_T^2}$	μ= <i>Μ</i> <sub>H</sub>
μ= M <sub>H</sub>	μ= M <sub>H</sub>	— μ= M <sub>F</sub>
μ= 1 TeV	μ= 1 TeV	μ= 1 ٦

### Eventually need to be taken into account in a global fit! HP2, 12/9/24





## Summary

Precision computations important to enhance sensitivity (especially for unconstrained operators)

Global fit results affected by the precision of EFT predictions

effects in predictions

Aim to include more and more precise theory predictions in the fits

- Progress in computation of QCD and EW corrections, and inclusion of RGE

