

Open bottom production at hadron colliders at NNLO+NNLL

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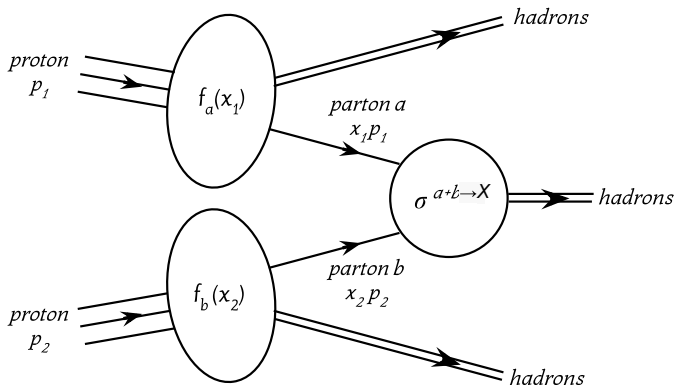
The goal

- Want to describe the process $pp \rightarrow b/\bar{b} + X$ with high precision
- At high energies, the b -quark is effectively almost massless
- \Rightarrow Quasi-collinear singularities ($b \rightarrow b g$, $g \rightarrow b \bar{b}$, etc.)
- These introduce large logarithms $\sim \ln p_{T,b}/m_b$
- So: want to perform resummation (through NNLL)
- Many details already covered by Matteo Cacciari on Tuesday!
- Note: also logs from initial state and UV
- Those already resummed by PDF collaboration (4-flavour vs. 5-flavour PDF+ α_s sets)

Reminder:

fragmentation functions

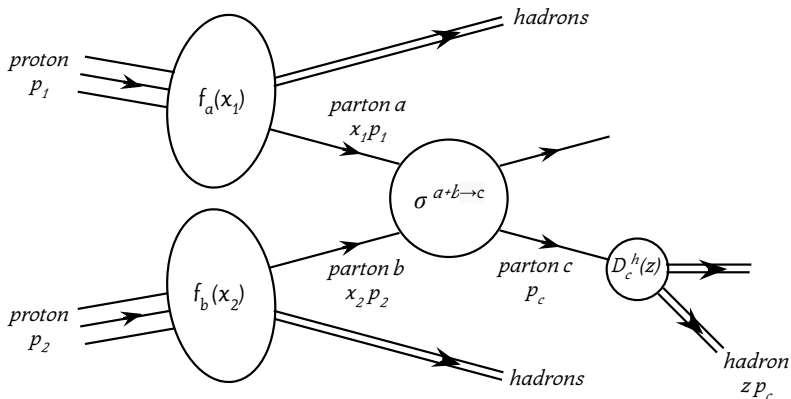
Standard QCD collinear factorisation



$$\sigma^{pp \rightarrow X}(p_1, p_2) = \sum_{a,b} \int_0^1 \int_0^1 f_a(x_1) f_b(x_2)$$

$$\sigma^{a+b \rightarrow X}(p_a = x_1 p_1, p_b = x_2 p_2) dx_1 dx_2 + \mathcal{O}(m_p^n / Q^n)$$

Final-state factorisation



$$\sigma^{pp \rightarrow h+X}(p_1, p_2, p_h) = \sum_{a,b,c} \iiint_0^1 f_a(x_1) f_b(x_2) D_{c \rightarrow h}^{\text{NP}}(z) \delta(p_h - z p_c)$$

$$\sigma^{a+b \rightarrow c}(p_a = x_1 p_1, p_b = x_2 p_2, p_c) dx_1 dx_2 dz + \mathcal{O}(m_p^n / Q^n, m_h^m / E_h^m)$$

Perturbative fragmentation functions

- So far FFs of non-perturbative nature (parton \rightarrow hadron)
- Can also define perturbative FFs as follows: *Mele, Nason (1991)*

$$\sigma_b(m_b) = \sum_i \sigma_i(m_b = 0) \otimes D_{i\rightarrow b} + \mathcal{O}(m_b^n/E_b^n)$$

- Convolution with $D_{i\rightarrow b}$ restores correct behaviour of mass-independent and log-enhanced terms ($\ln^m p_{T,b}/m_b$)
- Works for any massive particle, multiple masses, multiple particles, etc.
- As long as $m_b \gg \Lambda_{\text{QCD}}$: perturbatively calculable!

Perturbative fragmentation functions

$$\sigma_b(m_b) = \sum_i \sigma_i(m_b = 0) \otimes D_{i \rightarrow b} + \mathcal{O}(m_b^n/E_b^n)$$

- The $\sigma_i(m_b = 0)$ contain additional collinear singularities w.r.t. $\sigma_b(m_b)$
- \Rightarrow The PFF must contain collinear poles which cancel these
- \Rightarrow Predictable poles, which can be renormalised away
- Introduces new and arbitrary factorisation scale
+ corresponding RGEs (time-like DGLAP equations)
- Plan for resummation:
 - ① Evaluate PFF at $\mu \sim m_b$ (no large log of m_b)
 - ② Evolve to $\mu \sim Q$ (hard scale of process)
 - ③ Combine with massless cross sections evaluated at $\mu \sim Q$

Soft-gluon resummation

- PFF contains large threshold logs \Rightarrow soft-gluon resummation
- Discussed in detail in Matteo Cacciari's talk
- We perform soft-gluon resummation with NNLL accuracy
- But: care must be taken at very large x (see Andrea Ghira's talk)
- Fortunately not a problem here (large x not probed)
- Note: partonic cross section also contains threshold logarithms
- Those are not resummed here

The final resummation formula

- Factorising the mass dependence into FFs allows resummation
- But only accurate up to power corrections $\mathcal{O}(m_b/E_b)$ (leading-power, LP)
- At low p_T , resummation not needed
- \Rightarrow Can use fixed-order result with full mass dependence
- ‘Best prediction’ across full spectrum: combine the two!
- Naïve matching:

$$\text{Best prediction} = (\text{FO massive}) + (\text{resummed LP}) - \underbrace{(\text{FO LP})}_{\text{double counting}}$$

The final resummation formula

- Resummation of PFF performed in previous work
Czakon, TG, Mitov, Poncelet (2022)
- ‘Double counting’ is FO expansion of resummed calculation
- Combination of partonic cross sections with
 - 1 NNLO perturbative fragmentation functions
Melnikov, Mitov (2004, 2005)
 - 2 NNLO PDF threshold matching conditions
Buza, Matiounine, Smith, Migneron, van Neerven (1996)
 - 3 NNLO α_s threshold matching condition
Bernreuther, Wetzel (1982/1983)
- These extra contributions now also implemented in STRIPPER
- STRIPPER: fully general NNLO subtraction scheme code

The final resummation formula

- Naïve matching:

$$\text{Best prediction} = (\text{FO massive}) + (\text{resummed LP}) - (\text{FO LP})$$

- But: LP cross sections blow up at low p_T
- Solution: change $E_T(m_b = 0) \rightarrow E_T(m_b) \Leftrightarrow p_T^2 \rightarrow p_T^2 - m_b^2$ for LP
- In practice still not enough: large resummation effects at low p_T
- Artificially suppress resummation effects at low p_T
- FONLL approach: [Cacciari, Greco, Nason \(1998\)](#)

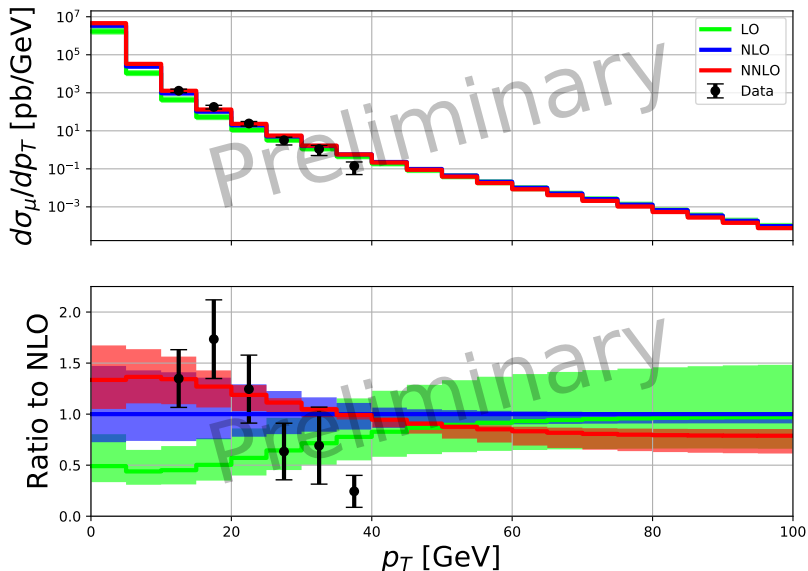
$$\text{FONLL} = (\text{FO massive}) + G(p_{T,b}) \cdot [(\text{resummed LP}) - (\text{FO LP})]$$

$$G(p_T) = \frac{p_T^2}{p_T^2 + c^2 m_b^2}, \text{ with } c = 5$$

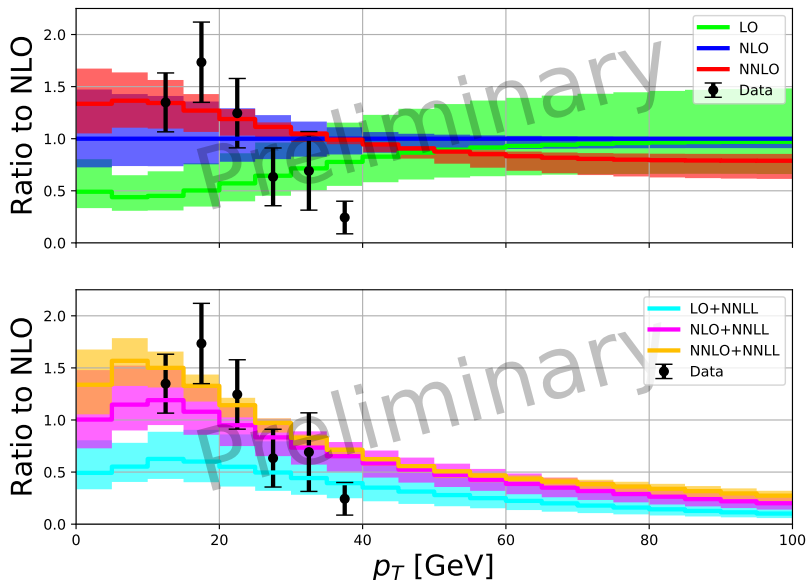
Preliminary results at NNLO+NNLL

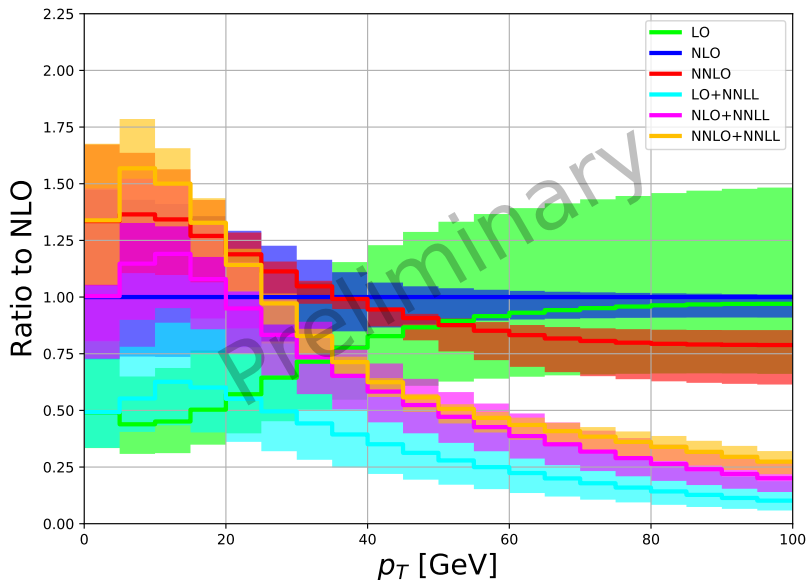
The muon p_T spectrum through NNLO+NNLL

- Study muons from B -hadron decays at 630 GeV $Spp\bar{p}S$
Phys.Lett.B 256 (1991) 121
- From b to B : standard (non-perturbative) fragmentation
 \Rightarrow nothing new
- From B to μ : data-driven implementation
- Data probes muons up to 40 GeV \leftrightarrow b-quarks up to ~ 80 GeV
- $\Rightarrow \ln(p_{T,b}^2/m_b^2) \sim 6$
- Resummation effects expected to be large

The muon p_T spectrum through NNLO+NNLL

The muon p_T spectrum through NNLO+NNLL



The muon p_T spectrum through NNLO+NNLL

Conclusion & outlook

- First NNLO+NNLL calculation of open bottom at hadron colliders
- Resummed $\ln(p_{T,b}/m_b)$ + some soft-gluon resummation
- Including full mass dependence at FO
- Straightforward extension of FONLL to higher order
- Improved precision and agreement with data w.r.t. NLO+NLL
- Prediction shown just an illustrative example
- Also have predictions for LHC & Tevatron
- Stay tuned for the publication!