

# Finite Feynman Integrals in Momentum and Parameter Space

Pavel Novichkov

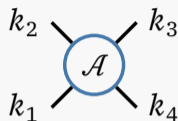
IPhT CEA/Saclay

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[2311.16907] Giulio Gambuti, David Kosower, PN, Lorenzo Tancredi

# Motivation


$$= c_1 \text{Master}_1 + \dots + c_N \text{Master}_N$$

1. IR factorization:  $\mathcal{A}^{(L)} = \mathcal{F}^{(L)} + \mathcal{A}_{\text{IR}}^{(<L)}$

☞ organize integrals according to their divergence properties

[Arkani-Hamed, Bourjaily, Cachazo, Caron-Huot, Trnka 2011; Badger, Mogull, Peraro 2016; Henn, Peraro, Stahlhofen, Wasser 2019]

2. Finite integrals are simpler to evaluate

☞ (quasi-)finite basis [von Manteuffel, Panzer, Schabinger 2015]

# Questions

## 1. Classification

- » locally finite
- » evanescently finite:  $1/\epsilon$  divergence  $\times \mathcal{O}(\epsilon)$  integrand
- » evanescent:  $\mathcal{O}(\epsilon)$

## 2. Construction

- » momentum space
- » parameter space

## 3. Evaluation

- » numerical
- » IBPs, differential equations, ...

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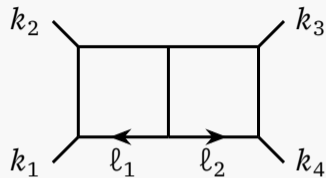
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- » numerical
- » IBPs, differential equations, ...

Constructing finite integrals in momentum space

# Setup

Find **Num** such that  $\int \prod_{i=1}^L d^4 \ell \frac{\text{Num}}{\text{Den}_1 \cdots \text{Den}_E} < \infty$

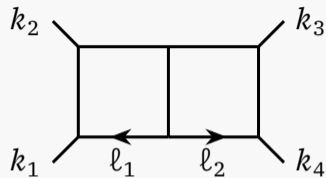


**Num** = Poly( $\ell_i \cdot \ell_j, \ell_i \cdot k_j$ ) whose coefficients are Rational ( $k_i \cdot k_j, m_i$ )

$$\text{Den} = \left( \sum \pm \ell_i \pm k_j \right)^2 - m^2 + i\delta$$

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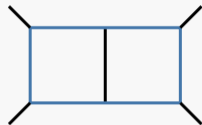
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locally finite  $\sim$  absolutely convergent: 
$$\int dx f(x) = \int_{f(x)>0} dx |f(x)| - \int_{f(x)<0} dx |f(x)|$$

# Strategy

UV



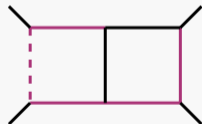
divergent surface

$$\left[ \begin{array}{l} \ell_1 = \infty \\ \ell_1 + \ell_2 = C \end{array} \right]$$

power-counting rule

$$\left[ \begin{array}{l} \ell_1 \rightarrow \lambda \ell \\ \ell_2 \rightarrow C - \lambda \ell \end{array} \right]$$

IR



$$\left[ \begin{array}{l} \ell_1 = k_1 \\ \ell_2 = x k_4 \end{array} \right]$$

$$\left[ \begin{array}{l} \ell_1 \rightarrow k_1 + \lambda^2 \ell_s \\ \ell_2 \rightarrow x k_4 + \lambda^2 \eta_4 + \lambda \ell_{\perp} \end{array} \right]$$

Constrain the ansatz:  $\text{Num} = c_1 + c_2 (\ell_1 \cdot k_1) + \dots + c_n (\ell_2^2)^2 (\ell_2 \cdot k_3)$



## IR divergences from Landau equations

$$\forall l_i : \frac{\partial}{\partial l_i} \sum_e \alpha_e \text{Den}_e = 0$$

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## IR divergences from Landau equations

$$\forall l_i : \frac{\partial}{\partial l_i} \sum_e \alpha_e \text{Den}_e = 0 \xrightarrow{u(\alpha) \neq 0} \begin{cases} l_1 = \frac{\alpha_1}{\alpha_1 + \alpha_2} k_1 \\ l_2 = \dots \end{cases}$$

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$$\rightarrow \forall \alpha_e : \alpha_e \frac{\partial}{\partial \alpha_e} \mathcal{F}(\alpha) = 0$$

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↓

$$\forall \alpha_e : \alpha_e \text{Den}_e = 0 \longrightarrow \forall \alpha_e : \alpha_e \frac{\partial}{\partial \alpha_e} \mathcal{F}(\alpha) = 0$$

in many cases [\[Stephen's talk\]](#)

can form a subtraction-free combination:

$$\sum_e w_e \left( \alpha_e \frac{\partial}{\partial \alpha_e} \mathcal{F}(\alpha) \right) = 0$$

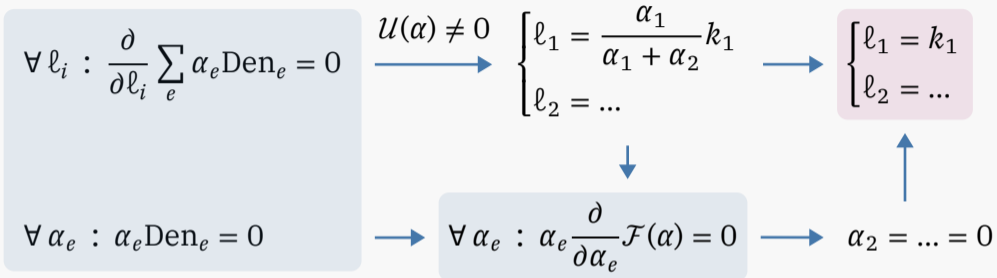
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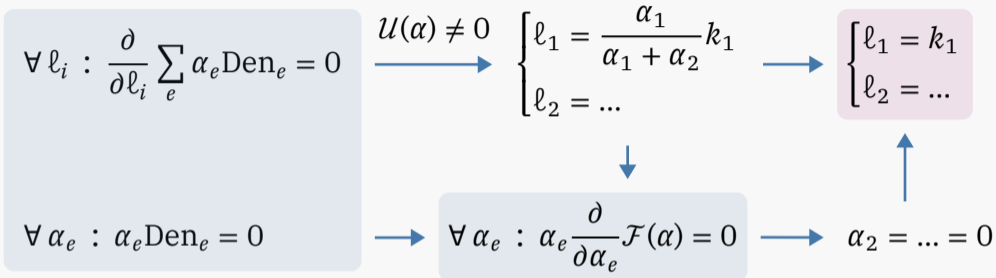
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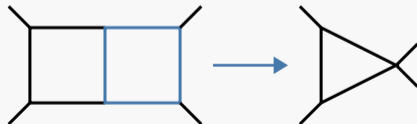
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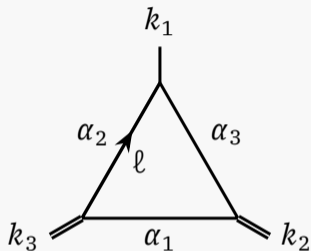
# IR divergences from Landau equations



degenerate case:  $U(\alpha) = 0$



## Example



$$\frac{\partial}{\partial l} (\alpha_1 (\ell + k_3)^2 + \alpha_2 \ell^2 + \alpha_3 (\ell - k_1)^2) = 0$$

$$\Rightarrow \ell = \frac{\alpha_3 k_1 - \alpha_1 k_3}{\alpha_1 + \alpha_2 + \alpha_3}$$

$$\forall \alpha_e : \quad \alpha_e \frac{\partial}{\partial \alpha_e} (m_3^2 \alpha_1 \alpha_2 + m_2^2 \alpha_1 \alpha_3) = 0$$

$$\Rightarrow \alpha_1 = 0$$

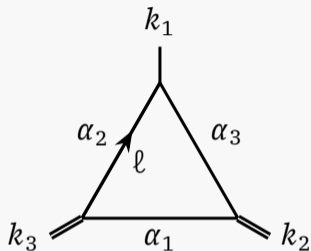
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$$\ell = -k_3$$



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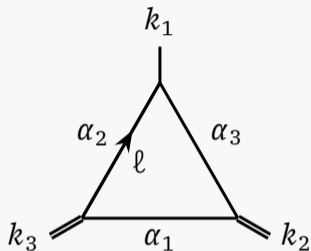
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by power counting

## Example



$$\text{Num} = c_1 + c_2 (\ell \cdot k_1) + c_3 (\ell \cdot k_2)$$

$$\text{Num} (\ell = x k_1) = 0$$

$$c_1 + x (k_1 \cdot k_2) c_3 = 0$$

$$\Rightarrow c_1 = c_3 = 0$$

$$\Rightarrow \text{Num} = (\ell \cdot k_1)$$

$$\Rightarrow \int d^4 \ell \frac{(\ell \cdot k_1)}{\ell^2 (\ell - k_1)^2 (\ell + k_3)^2} < \infty$$

Constructing finite integrals in parameter space

## From momentum to parameter space

$$\int \prod_{i=1}^L d^D \ell_i \frac{\text{Num}(\ell, k)}{\text{Den}_1 \cdots \text{Den}_E}$$



$$\Gamma(E - LD/2 - \lfloor r/2 \rfloor) \int_{\mathbb{R}^+} d^{E-1} \alpha \text{Num}(\alpha; s, D) \mathcal{U}(\alpha)^{E-(L+1)D/2-r} \mathcal{F}(\alpha; s)^{LD/2-E}$$

$r$  is rank of  $\text{Num}(\ell, k)$

$$s = \{(k_i \cdot k_j), m_i\}$$

- »  $\text{Num}(\alpha; s, D)$  is a polynomial in  $\alpha$ s of degree  $rL$
- » not every  $\text{Num}(\alpha; s, D)$  can be obtained from  $\text{Num}(\ell, k)$
- » for some  $\text{Num}(\ell, k) \neq 0$  we have  $\text{Num}(\alpha; s, D) = 0$

## Convergence in parameter space

$$\int_{\mathbb{R}^+} d^n \alpha \frac{\text{Num}(\alpha)}{\text{Den}(\alpha)} = \sum_x c_x \int_{\mathbb{R}^+} \frac{d\alpha_1}{\alpha_1} \dots \frac{d\alpha_n}{\alpha_n} \frac{\alpha_1^{x_1} \dots \alpha_n^{x_n}}{\text{Den}(\alpha)}$$

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Euler–Mellin integral [[Pierpaolo's talk](#)]

Newton polytope [[Leonardo's talk](#)]

Theorem:  $\int$  converges  $\Leftrightarrow (x_1, \dots, x_n) \in \text{Newt}(\text{Den})$

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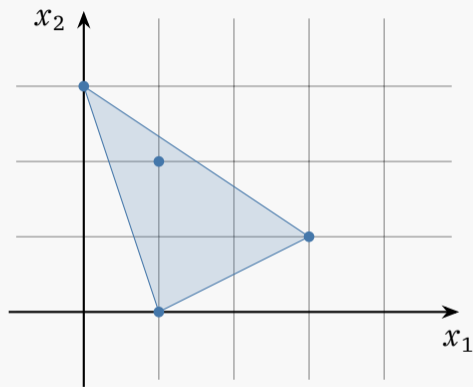
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- ✗ theorem applies to single terms of  $\text{Num}(\alpha)$ 
  - ✓ can show there are no cancellations between terms
- ✗ theorem probes only end-point divergences
  - ✓ map inner divergences to the boundary [\[Stephen's talk\]](#)

## Convergence in parameter space

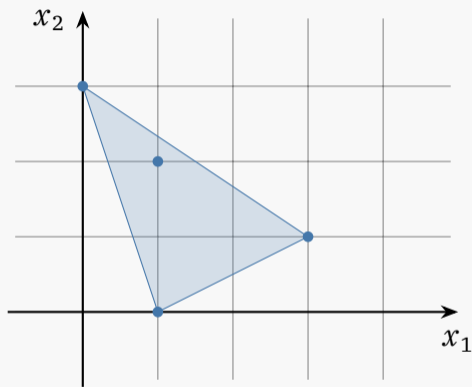


$$\text{Den}(\alpha) = \sum_{x \in X} c_x \alpha^x = \alpha_1 + \alpha_1^3 \alpha_2 + \alpha_1 \alpha_2^2 + \alpha_2^3$$

$$X = \{(1, 0), (3, 1), (1, 2), (0, 3)\}$$

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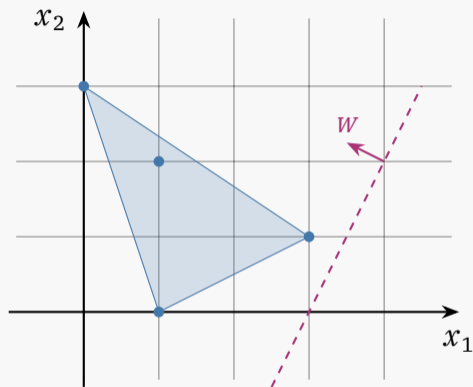
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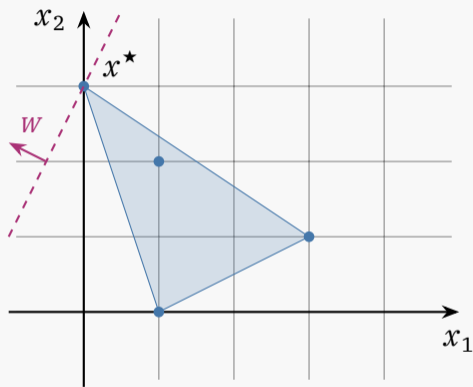
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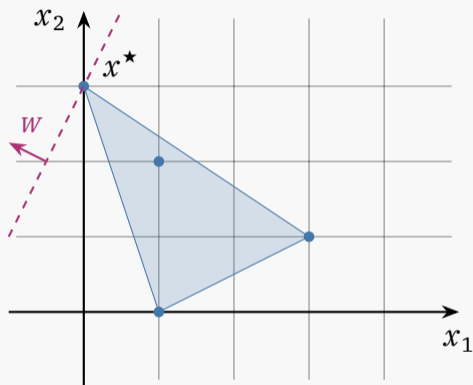
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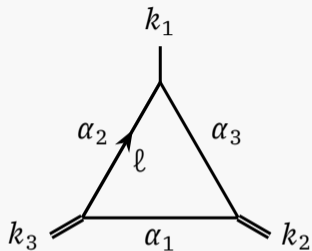
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$\alpha^x \subset \text{Num}$  cannot become dominant (for any choice of  $w$ ) if  $x$  lies inside  $\text{Newt}$

## Example

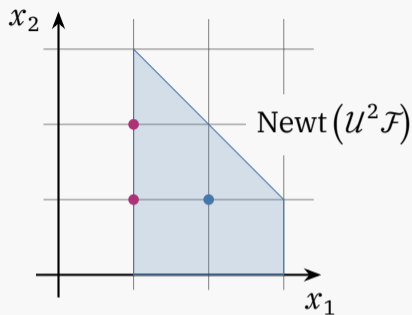
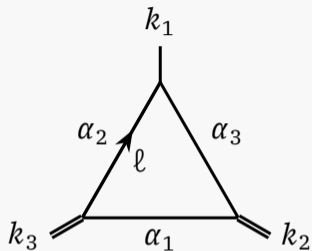


$$\text{Num}(\ell) = c_1 + c_2 (\ell \cdot k_1) + c_3 (\ell \cdot k_2) \rightarrow \int_{\mathbb{R}^+} \frac{d\alpha_1}{\alpha_1} \frac{d\alpha_2}{\alpha_2} \frac{\text{Num}(\alpha)}{U^2(\alpha) \mathcal{F}(\alpha)}$$

$$\begin{aligned} \text{Num}(\alpha) &= 2c_1 \alpha_1 \alpha_2^2 + [2c_1 + (m_3^2 - m_2^2) c_3] \alpha_1 \alpha_2 \\ &\quad + [2c_1 + (m_3^2 - m_2^2) c_2 + (m_2^2 + m_3^2) c_3] \alpha_1^2 \alpha_2 \end{aligned}$$



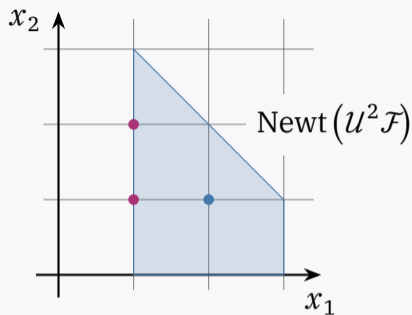
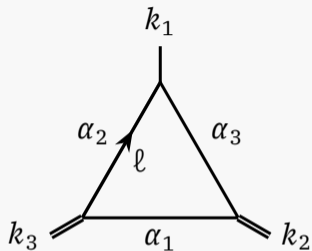
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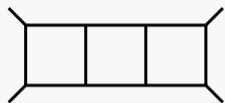
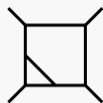
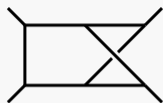
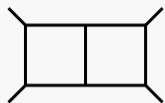
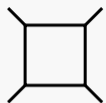


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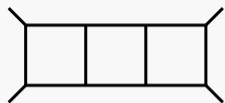
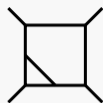
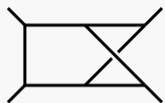
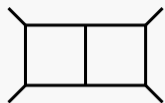
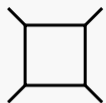
$$\begin{aligned} \text{Num}(\alpha) &= 2c_1 \alpha_1 \alpha_2^2 + [2c_1 + (m_3^2 - m_2^2) c_3] \alpha_1 \alpha_2 \\ &\quad + [2c_1 + (m_3^2 - m_2^2) c_2 + (m_2^2 + m_3^2) c_3] \alpha_1^2 \alpha_2 \end{aligned} \Rightarrow \begin{aligned} c_1 &= c_3 = 0 \\ \text{Num}(\ell) &= (\ell \cdot k_1) \end{aligned}$$

Momentum space vs. parameter space

## Results

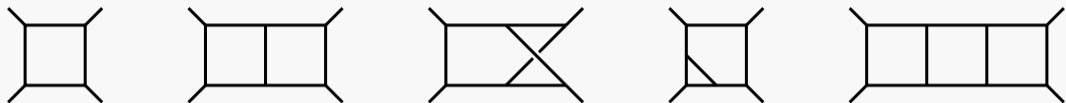


## Results



momentum  
loc. finite  $\subset$  parameter  
loc. finite

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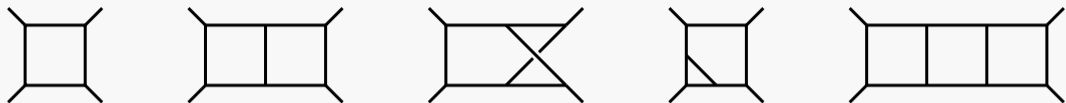


momentum  
loc. finite  $\subset$  parameter  
loc. finite

parameter  
loc. finite = momentum  
loc. finite + weakly  
UV finite + total  
derivatives



## Results



momentum  
loc. finite

$\subset$

parameter  
loc. finite

parameter  
loc. finite

=

momentum  
loc. finite

+

weakly  
UV finite

+

total  
derivatives

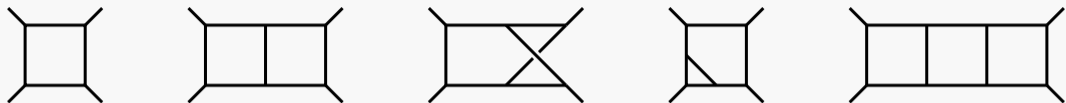
IBP reduce

IBP reduce

compare as linear spaces  
over D-independent coefficients



## Results



momentum  
loc. finite  $\subset$  parameter  
loc. finite

parameter  
loc. finite = momentum  
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## UV power counting (seemingly) doesn't work

$$k \text{ --- } \textcirclearrowleft_{\ell} \text{ --- } [(\ell \cdot k)] = \int d^2\ell \frac{(\ell \cdot k)}{(\ell^2 - m^2 + i\delta)((\ell + k)^2 - m^2 + i\delta)}$$

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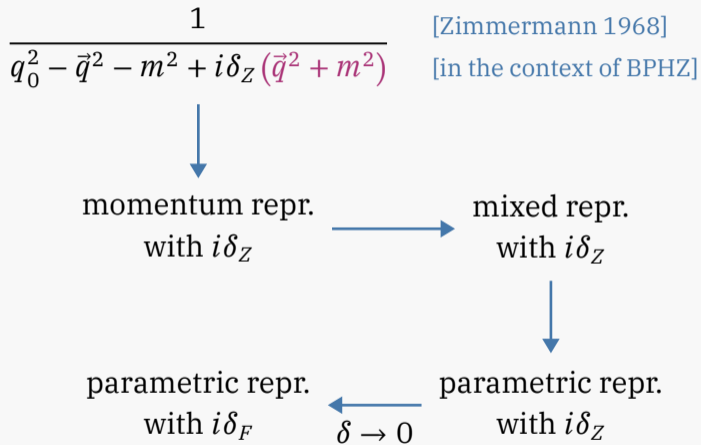
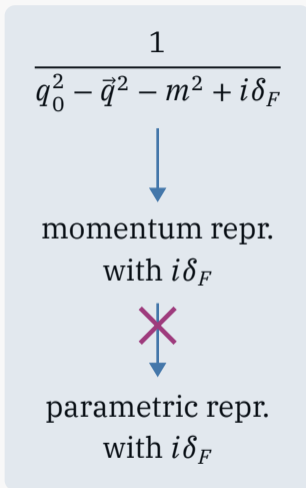
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$$\begin{aligned} k \text{ --- } \bigcirc_{\ell} \text{ --- } [(\ell \cdot k)] &= \int d^2\ell \frac{(\ell \cdot k)}{(\ell^2 - m^2 + i\delta)((\ell + k)^2 - m^2 + i\delta)} \\ &= \int d\ell_+ d\ell_- \frac{\ell_+ k_- + \ell_- k_+}{(\ell_+ \ell_- - m^2 + i\delta)((\ell_+ + k_+)(\ell_- + k_-) - m^2 + i\delta)} \rightarrow \infty \end{aligned}$$

Weinberg's theorem assumes **Euclidean** denominators:

$$\frac{1}{q_0^2 - \vec{q}^2 - m^2 + i\delta} \rightarrow \frac{1}{q_0^2 + \vec{q}^2 + m^2}$$

# Absolute convergence with Zimmermann's $i\delta$



[Zimmermann 1968]

[in the context of BPHZ]

[Lowenstein, Zimmermann 1975; Lowenstein, Speer 1976]

# Absolute convergence with Zimmermann's $i\delta$

$$\frac{1}{q_0^2 - \vec{q}^2 - m^2 + i\delta_F}$$



momentum repr.  
with  $i\delta_F$



parametric repr.  
with  $i\delta_F$

$$\frac{1}{q_0^2 - \vec{q}^2 - m^2 + i\delta_Z (\vec{q}^2 + m^2)}$$



momentum repr.  
with  $i\delta_Z$

parametric repr.  
with  $i\delta_F$



mixed repr.  
with  $i\delta_Z$



parametric repr.  
with  $i\delta_Z$

$\delta \rightarrow 0$

[Zimmermann 1968]

[in the context of BPHZ]

[Lowenstein, Zimmermann 1975; Lowenstein, Speer 1976]

# Summary

1. We can systematically construct locally finite Feynman integrals both in momentum and parameter space
2. Parameter-space locally finite  $\not\Rightarrow$  momentum-space locally finite
3. Momentum-space locally finite  $\Rightarrow$  parameter-space locally finite  $\Rightarrow$  composed of “interior” monomials:
  - » can cross-check momentum-space results
  - » can expand in powers of  $\epsilon$  directly under the integral sign
  - » can compute numerically (term by term  $\Rightarrow$  numerical stability)

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Thank you!