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Three-loop amplitudes

for $V+jet$ and $H+jet$ production

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with

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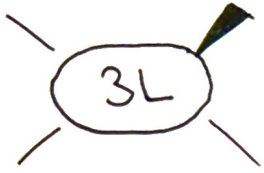
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William Torres-Bobadilla

based on 2301.10849, 2306.10170, 2307.15405, 2410.xxxxx

MOTIVATION

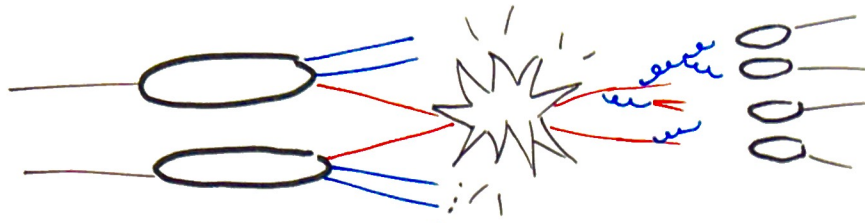


amplitudes vital to collider programme
of the next 2 decades

incarnations:

- $e^+e^- \rightarrow 3 \text{ jets @ } N^3\text{LO}$: clean, event shapes, α_s extraction
- $pp \rightarrow V + \text{jet @ } N^3\text{LO}$: large xsection, leptonic decay - clear signature
→ PDF fitting
- $pp \rightarrow H + \text{jet @ } N^3\text{LO}$: large p_T : resolve loop, light Yukawa couplings,
new physics coupling^{to} H
- $gg \rightarrow H @ N^4\text{LO}$ differential : main production mechanism @ LHC

Progress in all parts of particle theory workflow required

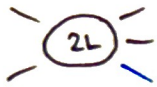


AMPLITUDES

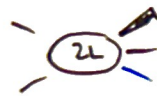


next order =

legs ↑



done [many; review: 2207.12255]



[Badger, Hartanto, Zoia, Krysi 2021-22]

[Abreu, F. Cordero, Ita, Kliewert, Page, Sotnikov '21]



PL integrals [Abreu, Chicherin, Sotnikov,
Zoia 2024]

&

loops ↑



[Caola, Chakraborty, Gambuti, von Manteuffel,
Tauscher; Bargiata 2021-22]



this talk!

(... and others: 4L FF; Gpt; with internal masses ...)

THEMES

- comp. cost (IBPs)
- size of expressions
- function spaces
- efficient evaluation
- ⋮

COMMON WORKFLOW

- ① Identify master integrals & derive DEs
- ② Solve in terms of a class of transcendental functions $\left\{ \begin{array}{l} \text{HPLs} \\ \text{MPLs} \\ \text{elliptic polylogs} \\ \dots \end{array} \right.$
- ③ IBP-reduce amplitude (get rational coefficients) & insert masters

⚡ integrals have repeated, spurious, evanescent structures
- contribute to intermediate swell
- cancel only at the end

traditional:

integrals $\xrightarrow{\text{insert}}$ amplitude $\xrightarrow{\text{observe}}$ physics

"bootstrap":

eg. [Dixon, Gürdogan, McLeod, Wilhelm '22]

~~integrals~~ amplitude $\xleftarrow{\text{enforce}}$ physics

this talk:

integrals $\xrightarrow{\text{tie}}$ amplitude $\xleftarrow{\text{impose}}$ physics

TOOL - Chen iterated integrals [Chen 1973]

$$I(l_1, l_2, \dots, l_n) = \int_0^1 dt \partial_t l_1 I(l_2, \dots, l_n) \quad \text{(path)}$$

- "alphabet" - independent kernels $l_i =$ functions of kinematics
- "length/weight" - number of integrations

ENCODES
ANALYTIC STRUCTURE

For ϵ -factorised DE system $d\vec{J}(x_i; \epsilon) = \epsilon dA(x_i) \vec{J}(x_i; \epsilon)$

formal solution $\vec{J} = \sum_{n=0}^{\infty} \epsilon^n \sum_{i=0}^n \tilde{A}^i \vec{V}_{n-i}$

connection matrix

"EASY"
TO OBTAIN

boundary vector (weight $n-i$)

After shuffling, $I(\vec{l}_1) I(\vec{l}_2) = \sum_{\vec{l} = \vec{l}_1 \cup \vec{l}_2} I(\vec{l})$

no hidden zeros

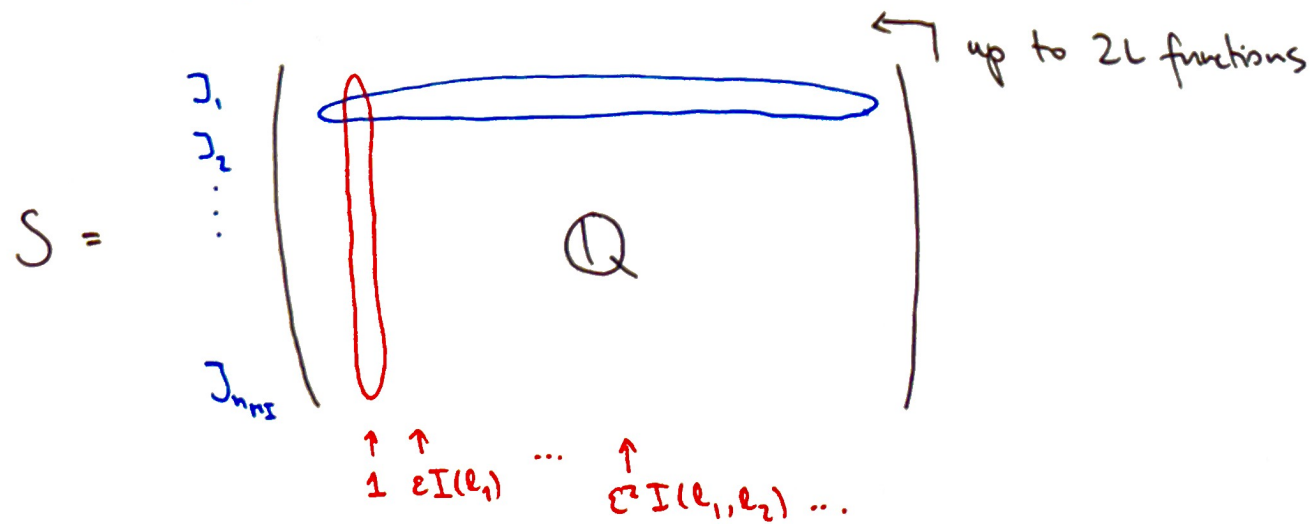
INDEPENDENT

→ VECTOR SPACE \mathcal{I}_w of shuffled iterated integrals up to weight w
over \mathbb{Q}

with basis $b_{\mathcal{I}_w} = \left\{ \varepsilon^{-a} \sum_n^b I(l_1, \dots, l_n; x) \right\}$ $a + bn + c \leq 2L$

@ any rat. kinematic point : scattering amplitudes $\in \mathcal{I}_w$
master integrals $\in \mathcal{I}_w$

cast master integrals as a \mathbb{Q} matrix (truncate @ $w=2L$)



coefficients of Chen integrals in MI

Master Integral Subspace $\mathcal{M}_w \subseteq \mathcal{I}_w$

spanned by the MI solution-vectors
= row space of matrix S

$$\# \text{ columns } (S) = \dim(\mathcal{I}_w) \Rightarrow n_{\text{MI}} \geq \dim(\mathcal{M}_w) = \text{rank}(S) \quad \text{!}$$

typically: thousands / millions
of iterated integrals

but!

of indep. combinations
bounded by # MIs

generous bound!

Sources of dependence:

- symmetries
- iteration
- truncation

All taken into account in Chen solutions

used for pentagon functions!
[Chicherin, Sotnikov '20]
[Chicherin, Sotnikov, Zoia '21]

Physical subspace $\mathcal{A}_w^{\parallel}$

\mathcal{M}_w encompasses any amplitude with this topology

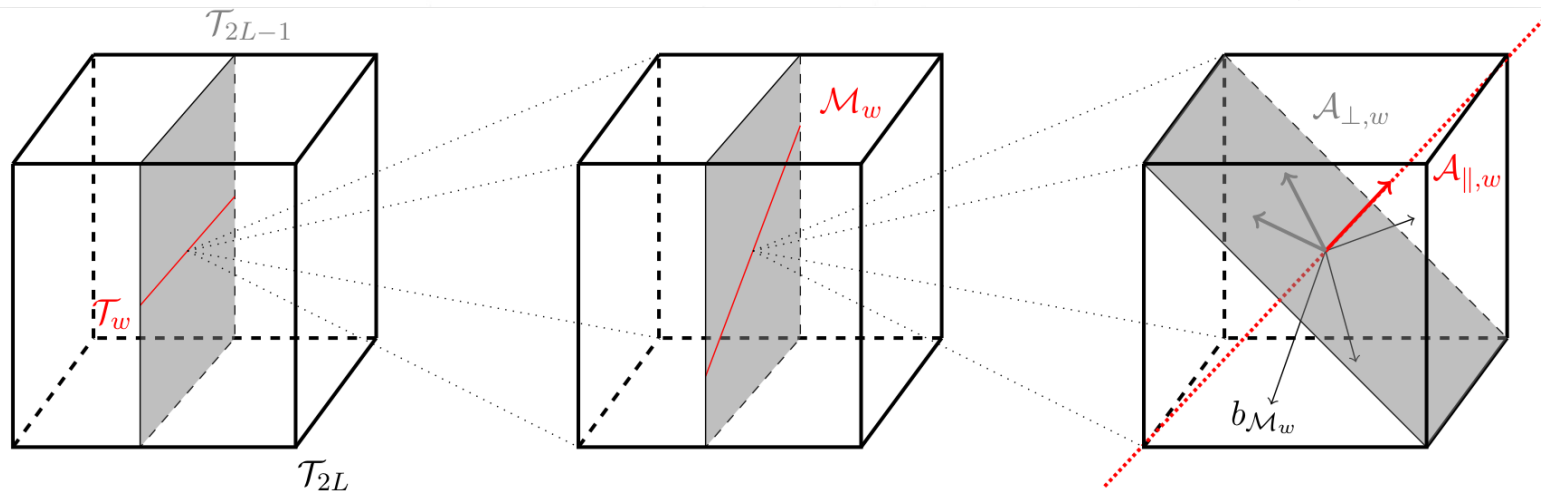
but for particular objects still too big cf. Won Lim's talk

$$\mathcal{M}_w = \mathcal{A}_w^{\parallel} \oplus \mathcal{A}_w^{\perp}$$

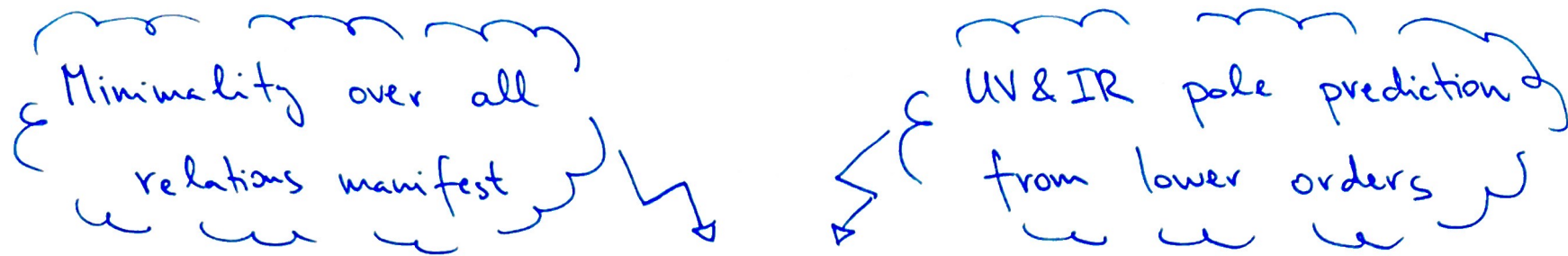
masters: 21 alphabet letters

FF of $\text{Tr}(\phi_R^2)$: 6 letters, restrictions...

physically favoured: alphabet, disc's, adjacencies, poles...



CONSEQUENCES



restrictions on finite
part of amp. without knowledge of integrand

Ex 1 I_1 and I_2 only appear in one minimal function

$$\Psi_i = 1 + \dots + \varepsilon^n I_1 + \dots + \varepsilon^{n+1} I_2 + \dots \rightarrow \begin{array}{l} \text{coeff. of } I_1 \text{ in } \varepsilon^1 \text{ of amp.} \\ \text{(known!)} \\ \doteq \text{coeff. of } I_2 \text{ in } \varepsilon^0 \text{ of amp.} \end{array}$$

Ex 2 Letter l_x appears for the first time @ this loop order

\rightarrow entire function Ψ_i must be suppressed by ε^{2L-n}

(no cancellation only vanishing due to minimality)

Generally: hundreds/thousands of restrictions on coefficients in finite part! ₉

RECAP

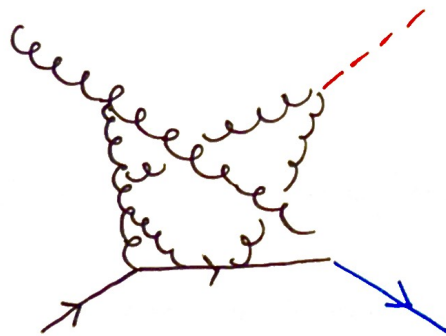
- Chen iterated integrals :
- formal solution to DEs
 - all (even unknown!) relations manifest
 - independent (if indep. letters & shuffled)

⇒ suitable for building basis of functions: minimal, graded, tailored
to an amplitude

Before amp. is known :- insights about analytic structure,
- information poles → finite part

After amp. is known: - no cancellations,
- DEs only for relevant functions → solve (choose fibration) & analytic continuation & evaluation
- finite-field: fewer coefficients with more restrictions

Application to $X + \text{jet}$
amplitudes



COMPUTATIONAL SETUP

- ① Lagrangian
- Z boson: vector (g_V) & axial-vector ($g_A \gamma_5$) coupling (Larin)
suppressed but in diff. observables NNLO (AV) \approx N³LO (V)
Singlet contribution due to $m_b \ll m_t$; 2L amps
 - H boson: infinite m_t limit [Gehrmann, Tauscher, Peraro: 2306.10170]
 $\mathcal{L}_{\text{HEFT}} = -\frac{\sum}{4} H G_a^{\mu\nu} G_{a,\mu\nu}$ ← coupling directly to gluons, no int. masses

② Tensor structure following [Tauscher, Peraro '19-20]

- avoid evanescent terms from d-dimensional Dirac algebra
- zero (even) or one (odd) instance of $\epsilon^{\mu\nu\sigma}$ in each basis elem. → parity manifest
- # tensor structures = # helicity amplitudes → minimal

$$Vg\bar{g}: 6+6 \text{ hel. amp.}$$

$$Vggg: 12+12 \text{ hel. amp.}$$

$$Hg\bar{g}: 1 \text{ hel. amp.}$$

$$Hggg: 2 \text{ hel. amp.}$$

(3) Colour decomposition

Assume $N \sim N_c \sim N_f$
 (3) (5)

and $N^2 = \mathcal{O}(10)$

$$\cdot N_c^3 + \cdot N_c^2 N_f + \cdot N_c N_f^2 + \cdot N_f^3$$

leading colour

$$\mathcal{A}^{(3L)} = + \cdot N_c + \cdot N_f + \cdot N_f^2 / N_c$$

$$+ \cdot \frac{1}{N_c} + \cdot \frac{N_f}{N_c^2}$$

$$+ \cdot \frac{1}{N_c^3}$$

V+jet planar

H+jet also non-planar

(4) Kinematics

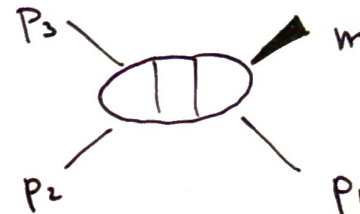
2 indep. dimensionless variables

$$(x+y+z=1)$$

$$x = \frac{s_{12}}{s^2}$$

$$y = \frac{s_{23}}{s^2}$$

$$z = \frac{s_{13}}{s^2}$$



(5) Symbol alphabet

• 2L integrals [Gehrmann, Remiddi '01]

• 3L planar integrals [Di Vita, Mastrolia, Schubert, Yundin '14; Canko, Syrakos '21]

$$\{x, y, z, 1-x, 1-y, 1-z\}$$

VS.

• 3L NPL integrals [Heun, Lim, T.-Bobadilla '23; in prep.]

3 new types:

$$\left\{ \frac{xy - \sqrt{xyz}}{xy + \sqrt{xyz}}, x^2 + xy + y, x^2 - x + y \right\}$$

(x6 perm.)

REDUCTION

- tensor rank up to 6
 - $\Theta(1M)$ integrals at full colour
 - top sectors ($t=10$) with up to 18 masters (8 @ L.C.)
- } prohibitively expensive



Fully determine analytic structure without complicated reduction!

- ① Source of new quadratic & sqrt letters are NPL graphs with ≥ 8 propagators
easy reductions!
→ 7-particle cuts discard only functions with 2L letters
& span the coefficients of all new functions
- ② Expressed in minimal set of graded transc. functions (combinations of Chen integrals), all expected / conjectural properties are manifest

RESULTS - example

→ reductions exact in ϵ at several numerical points with FIRE 6 [Smirnov, '15]

Chukharev

and KIRA 2 [Klappert, Lange, Maierhöfer, Usavitsch '20]

- 1282 minimal function basis
- only 93 contain new letters (but 100s of MIs!)

Ex $H \rightarrow g^+ g^+ g^+$, N_c^3 colour layer, point $x = \frac{3}{13}$, $y = \frac{11}{17}$

new letters
starting @ w4

$$\begin{aligned} \psi_1 &\rightarrow \mathcal{O}(\epsilon^3) \\ \psi_2 &\rightarrow \mathcal{O}(\epsilon^3) \\ \psi_3 &\rightarrow \mathcal{O}(\epsilon^3) \\ \psi_4 &\rightarrow \frac{1746}{48641} \epsilon^2 + \mathcal{O}(\epsilon^3) \\ \psi_5 &\rightarrow \mathcal{O}(\epsilon^3) \\ \psi_6 &\rightarrow \frac{22}{289} \epsilon^2 + \mathcal{O}(\epsilon^3) \\ \psi_7 &\rightarrow \frac{16}{169} \epsilon^2 + \mathcal{O}(\epsilon^3) \\ \psi_8 &\rightarrow \mathcal{O}(\epsilon^3) \\ \psi_9 &\rightarrow \mathcal{O}(\epsilon^3) \end{aligned}$$

starting
@ w5

$$\begin{aligned} \psi_{10} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{25} &\rightarrow -\frac{121}{63} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{11} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{26} &\rightarrow \frac{11}{6} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{12} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{27} &\rightarrow -\frac{121}{6} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{13} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{28} &\rightarrow \frac{685}{63} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{14} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{29} &\rightarrow \frac{3179}{126} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{15} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{30} &\rightarrow \frac{55}{18} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{16} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{31} &\rightarrow -\frac{11}{2} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{17} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{32} &\rightarrow \frac{187}{6} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{18} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{33} &\rightarrow \frac{85}{63} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{19} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{34} &\rightarrow \frac{11}{6} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{20} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{35} &\rightarrow -\frac{55}{126} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{21} &\rightarrow \mathcal{O}(\epsilon^2) & \psi_{36} &\rightarrow \frac{11}{9} \epsilon + \mathcal{O}(\epsilon^2) \\ \psi_{22} &\rightarrow \mathcal{O}(\epsilon^2) \\ \psi_{23} &\rightarrow \mathcal{O}(\epsilon^2) \\ \psi_{24} &\rightarrow \mathcal{O}(\epsilon^2) \end{aligned}$$

starting
@ w6

$$\begin{aligned} \psi_{37-63} &\rightarrow \mathcal{O}(\epsilon) \\ \psi_{64-93} &\rightarrow \mathcal{O}(\epsilon) \end{aligned}$$

❖ digradatic, sqrt letters

❖ only unigradatic

RESULTS - summary

H+jet amplitudes do not exhibit complexity of integrals!

① sqrt and diguadratic letters drop out of L.C. amplitudes

② maximum transcendentality principle [Kotikov, Lipatov '04 etc.]

Seemingly obeyed = no new letters in weight 6 functions in Hggg

↳ cf. Sudakov FF @ 3L
[Gehrmann, Henn, Huber '12]

↳ same as $\text{Tr}(\phi^2)$ FF in planar $\mathcal{N}=4$
[see Won Lim's talk]

Hggg amp.s @ 2L
[Brandhuber, Travaglini, Yang '12]

③ remarkable simplicity & patterns, size of rational coefficients;
at most one new uniguadratic letter / function

→ New avenues for bootstrapping approaches

SUMMARY

computed:

- 1L and 2L amplitudes extended to $\mathcal{O}(\epsilon^2)$ [2301.10845]
and with axial-vector coupling of Z [2306.10170]
- $\gamma^*/Z/W^\pm + \text{jet}$ 3L amplitudes @ L.C. [2307.15405]
- NPL master integrals for $H + \text{jet}$ @ L.C. [2302.12776; in prep.]
- all $X + \text{jet}$ integrands

minimal graded transc. functions

for simpler amplitude & evaluation of int.s & early insight into structure

coming:

- full analytic expressions for $H + \text{jet}$ @ L.C.
- remaining NPL topologies