



# Dalitz decays $D_{sJ}^{(*)} \rightarrow D_s^{(*)} \ell^+ \ell^-$

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Based on: P. Colangelo, F. De Fazio, F. Loporco, N.L., Dalitz decays  $D_{sJ}^{(*)} \rightarrow D_s^{(*)} \ell^+ \ell^-$ , [arXiv:2308.03453 [hep-ph]]

# The nature of the positive parity $c\bar{s}$ P-wave states

$D_{s1}(1^+)$  and  $D_{s2}^*(2^+)$   
well established states

$D_{s0}^*(0^+)$ ,  $D_{s1}'(1^+)$  first discovered in  
2003 by BaBar and CLEO Collaborations

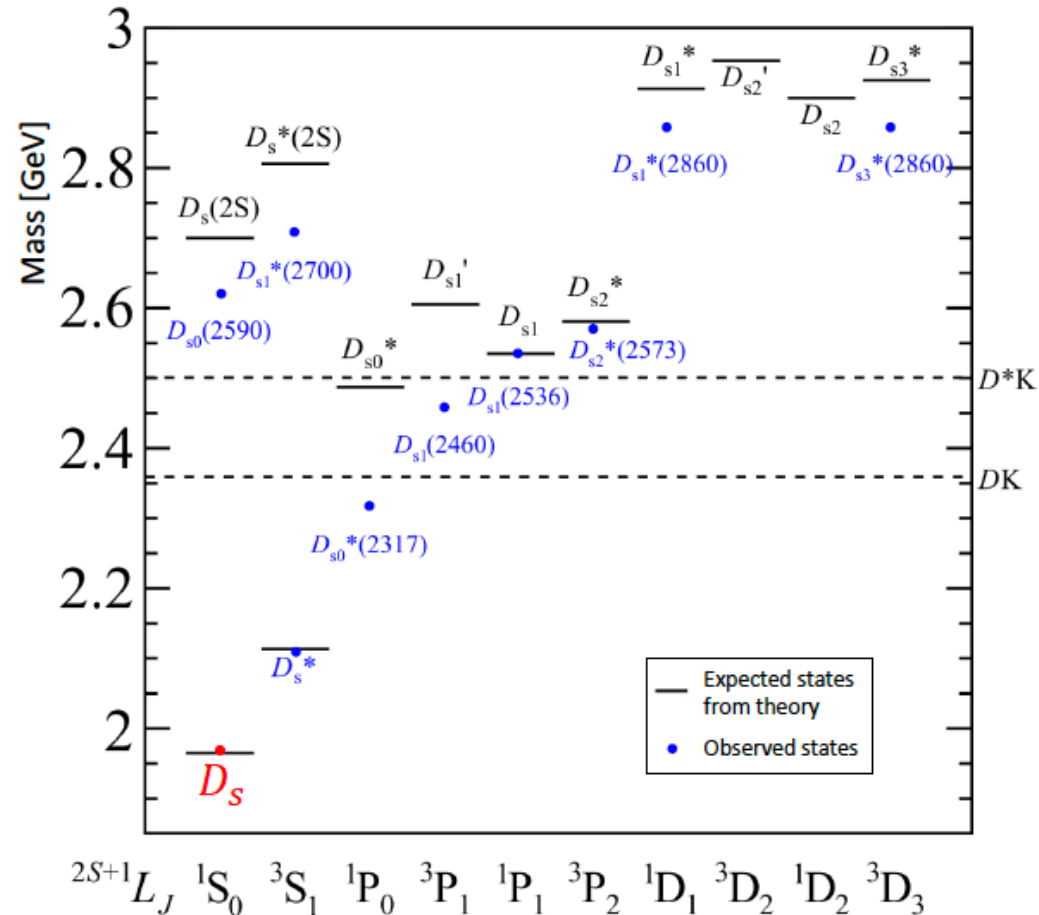
Phys. Rev. Lett. 90 (2003) 242001, Phys. Rev. D 68 (2003) 032002



Possible  $c\bar{s}$  P-wave states together  
with  $D_{s1}(1^+)$  and  $D_{s2}^*(2^+)$

MASS	
EXPECTED	OBSERVED
2.45 – 2.5 GeV	2.3178(5) GeV
2.6 – 2.65 GeV	2.4595(6) GeV

Mass near  $DK$  and  $D^*K$  threshold



Possible interpretations:

- ordinary  $c\bar{s}$  P-wave states
- molecular states
- multiquark states

# Method of investigation

Use Dalitz decays  $D_{sJ}^{(*)} \rightarrow D_s^{(*)} \ell^+ \ell^-$   
to probe the nature of  $D_{s0}^*$  and  $D_{s1}'$



complement the information from  
the electric dipole **radiative decays**

$$D_{s0}^* \rightarrow D_s^* \gamma, D_{s1}' \rightarrow D_s^{(*)} \gamma$$

$c\bar{s}$  system composed  
of **heavy-light** quarks

Heavy degrees of freedom decouple



Heavy quark spin  $\vec{s}_Q$  and total angular  
momentum of the light degrees of  
freedom  $\vec{s}_\ell$  separately conserved

**heavy quark spin symmetry**



States classified in doublets

$$H_a = \frac{1 + \not{v}}{2} [P_{a\mu}^* \gamma^\mu - P_a \gamma_5]$$

$$(s_\ell^P = \frac{1^-}{2}) \quad \mathbf{D}_s, \mathbf{D}_s^*$$

$$S_a = \frac{1 + \not{v}}{2} [P_{1a}^{\prime\mu} \gamma_\mu \gamma_5 - P_{0a}^*]$$

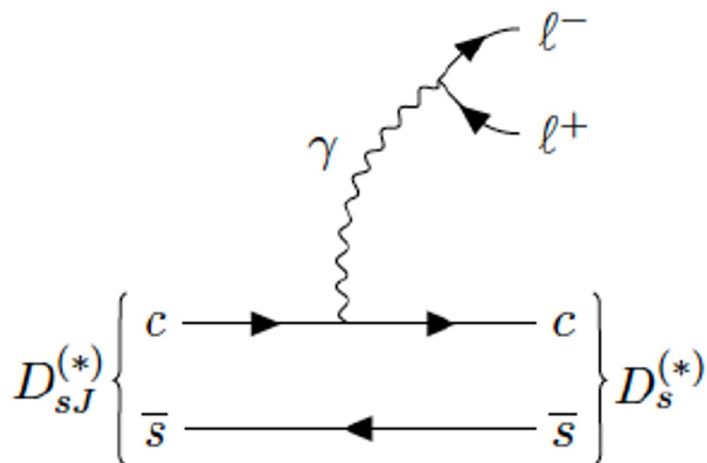
$$(s_\ell^P = \frac{1^+}{2}) \quad \mathbf{D}_{s0}^*, \mathbf{D}'_{s1}$$

$$T_a^\mu = \frac{1 + \not{v}}{2} \left\{ P_{2a}^{\mu\nu} \gamma_\nu - P_{1a\nu} \sqrt{\frac{3}{2}} \gamma_5 \left[ g^{\mu\nu} - \frac{1}{3} \gamma^\nu (\gamma^\mu - v^\mu) \right] \right\}$$

$$(s_\ell^P = \frac{3^+}{2}) \quad \mathbf{D}_{s1}, \mathbf{D}_{s2}^*$$

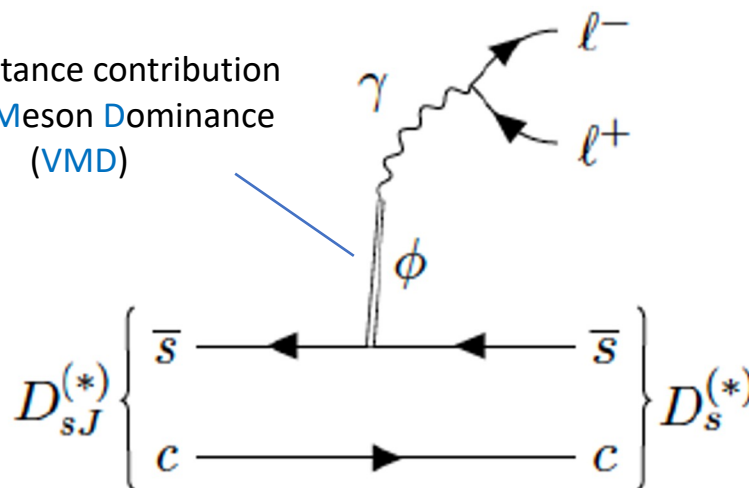
# Contributions to the decays

Charm



Strange

Long distance contribution  
Vector Meson Dominance  
(VMD)



$$\mathcal{A}(D_{sJ}^{(*)}(p') \rightarrow D_s^{(*)}(p)l^-(p_1)l^+(p_2)) = \langle D_s^{(*)}(p, \epsilon) | iJ_\mu^{\text{em}} | D_{sJ}^{(*)}(p', \epsilon') \rangle \frac{-ig^{\mu\nu}}{q^2} (-ie)\bar{u}(p_1)\gamma_\nu v(p_2)$$

$$J_\mu^{\text{em}} = e(e_c \bar{c}\gamma_\mu c + e_s \bar{s}\gamma_\mu s)$$

Computed using heavy quark  
spin symmetry

Effective Lagrangian constructed using

HQS + CS + HGI

Heavy Quark Symmetry

Chiral Symmetry

Hidden Gauge Invariance

# Strange quark contribution

$$\langle D_s^{(*)}(p, \epsilon) | \bar{s} \gamma_\mu s | D_{sJ}^{(*)}(p', \epsilon') \rangle$$

↓ VMD

$$\langle 0 | \bar{s} \gamma_\mu s | \phi(q, \eta) \rangle = m_\phi f_\phi \eta_\mu$$

$$\langle D_s^{(*)}(p, \epsilon) \phi(q, \eta) | D_{sJ}^{(*)}(p', \epsilon') \rangle \frac{i}{q^2 - m_\phi^2} \langle 0 | \bar{s} \gamma_\mu s | \phi(q, \eta) \rangle$$

Interaction with light vector meson through the field

$$\rho_\mu = i \frac{g_V}{\sqrt{2}} \hat{\rho}_\mu \quad g_V = 5.8 \text{ chosen to satisfy low energy relations (KRSF)}$$

$$\hat{\rho}_\mu = \begin{pmatrix} \sqrt{\frac{1}{2}}\rho^0 + \sqrt{\frac{1}{6}}\phi^{(8)} & \rho^+ & K^{*+} \\ \rho^- & -\sqrt{\frac{1}{2}}\rho^0 + \sqrt{\frac{1}{6}}\phi^{(8)} & K^{*0} \\ K^{*-} & \bar{K}^{*0} & -\sqrt{\frac{2}{3}}\phi^{(8)} \end{pmatrix}_\mu$$

Interaction Lagrangian terms for negative and positive parity doublets

$$\mathcal{L}_1^S = -g_1^S \text{Tr} [\bar{H} S \gamma^\alpha (\mathcal{V}_\alpha - \rho_\alpha)] + \text{h.c.}$$

$$\mathcal{F}_{\mu\nu} = \partial_\mu \rho_\nu - \partial_\nu \rho_\mu + [\rho_\mu, \rho_\nu] \quad \mathcal{L}_2^S = g_2^S \frac{1}{\Lambda} \text{Tr} [\bar{H} S \sigma^{\alpha\beta} \mathcal{F}_{\alpha\beta}] + \text{h.c.}$$

$$\mathcal{L}_2^T = i h^T \frac{1}{\Lambda^2} \text{Tr} [\bar{H} T_\mu \sigma^{\alpha\beta} \mathcal{D}^\mu \mathcal{F}_{\alpha\beta}] + \text{h.c.}$$

$g_1^S, g_2^S, h^T$ : strong couplings that relate several channels

# Charm quark contribution

$$\langle D_s^{(*)}(p, \epsilon) | \bar{c} \gamma_\mu c | D_{sJ}^{(*)}(p', \epsilon') \rangle$$

Trace formalism to parametrize charm current matrix elements

$$\langle H(v) | \bar{c} \Gamma c | S(v') \rangle = -\tau_{1/2}(w) \text{Tr}[\bar{H}(v) \Gamma S(v')]$$

$$\langle H(v) | \bar{c} \Gamma c | T(v') \rangle = -\tau_{3/2}(w) \text{Tr}[\bar{H}(v) \Gamma v_\mu T^\mu(v')]$$

Universal functions

$$\tau_i(w) = \tau_i(1) [1 - (w - 1) \rho_i^2]$$

$$w = v' \cdot v = \frac{m_{D_{sJ}^{(*)}}^2 + m_{D_s^{(*)}}^2 - q^2}{2m_{D_{sJ}^{(*)}} m_{D_s^{(*)}}}$$

Uncertainties from the value  
at zero recoil and the slope

Form factors  $\tau_{1/2}$  and  $\tau_{3/2}$  appear in several channels

# Numerical Results

Uncertainties from  $\tau_{1/2}$ ,  $\tau_{3/2}$  and  $g_1^S, g_2^S, h^T$

$g_1^S$ : from the semileptonic  $D \rightarrow K^*$  form factor  
Phys. Rept. 281 (1997) 145{238}

$g_2^S$ : from light-cone QCD sum rule computation of the decay amplitude of the positive parity charmed mesons to real photons  
Phys. Rev. D 72 (2005) 074004

$h^T$ : from strong decay width of excited charmed mesons  
Phys. Rev. D 98 (2018) 114028

$\tau_{1/2}, \tau_{3/2}$ : from semileptonic  $B$  decays to positive parity charmed mesons  
Phys. Rev. D 58 (1998) 116005

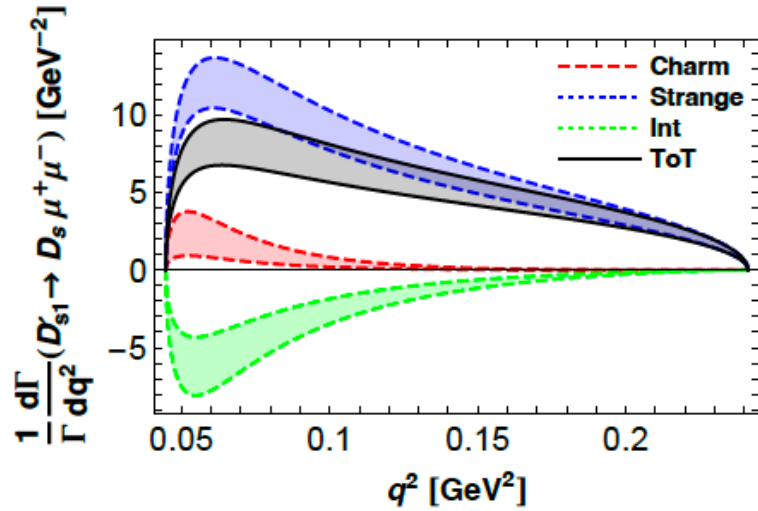
Sign of interference not known  Two extreme cases depending on the product between  $\tau_{1/2}, \tau_{3/2}$  and  $g_1^S, g_2^S, h^T$

Case A  
POSITIVE

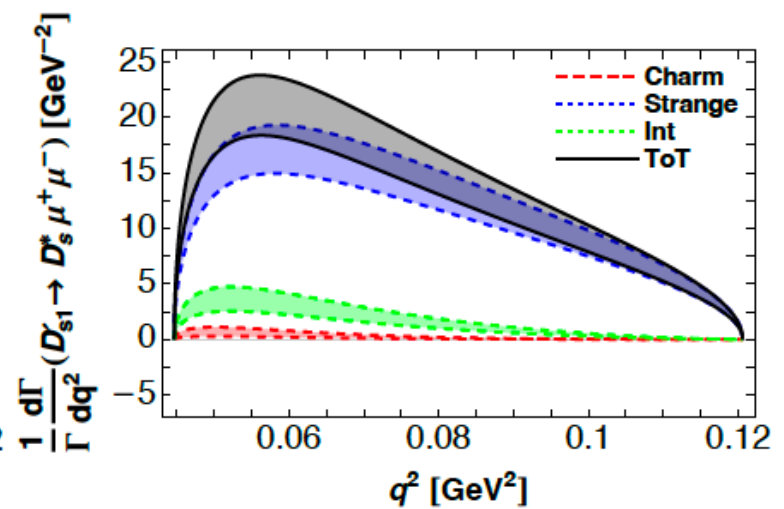
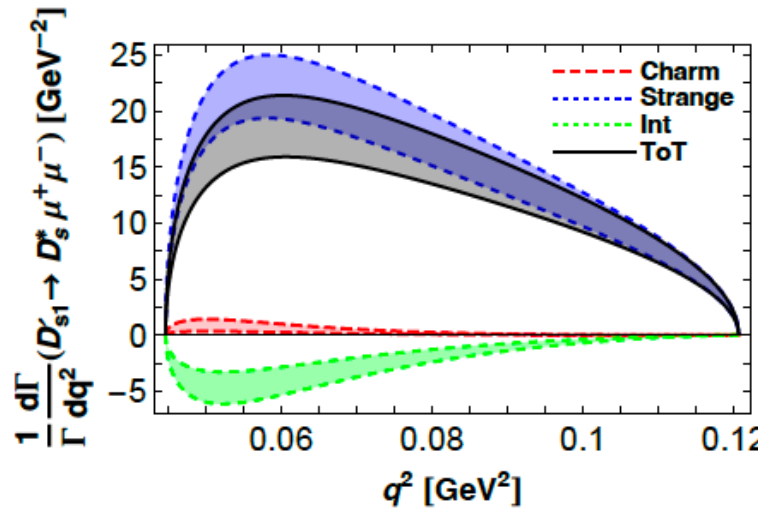
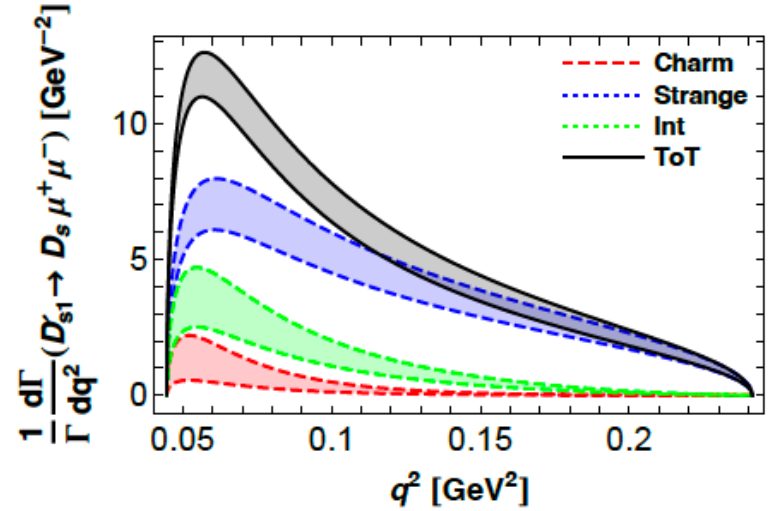
Case B  
NEGATIVE

▪  $D'_{s1} \rightarrow D_s^{(*)} \mu^+ \mu^-$

Case A



Case B

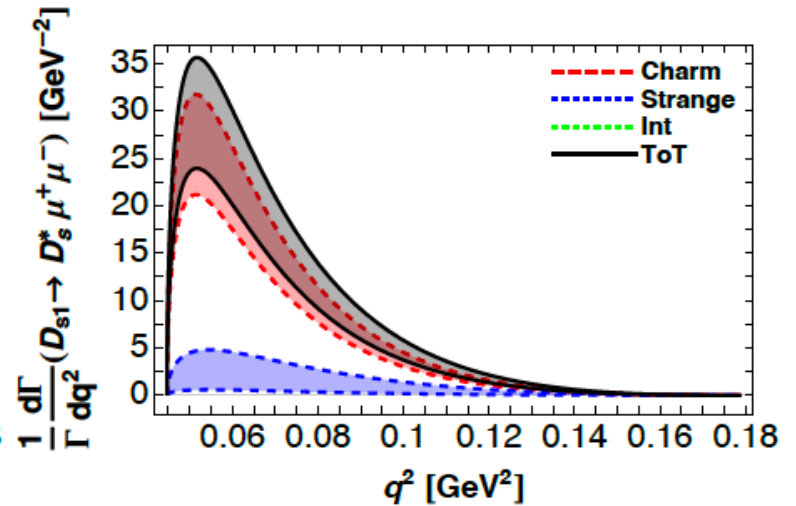
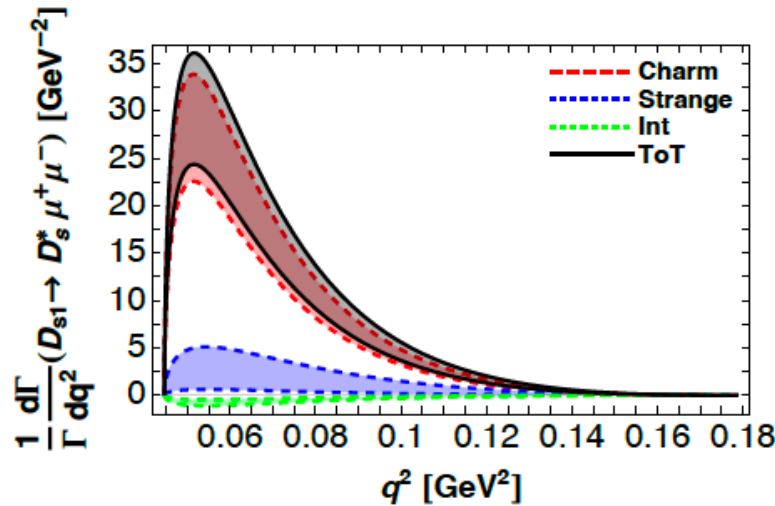
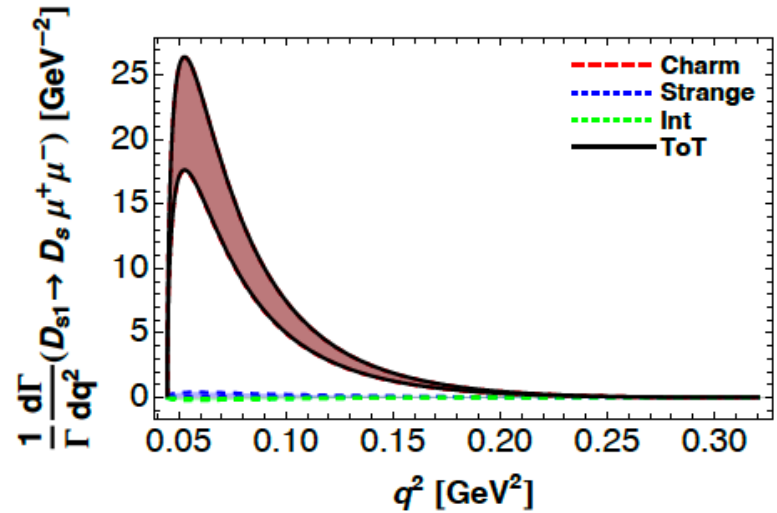
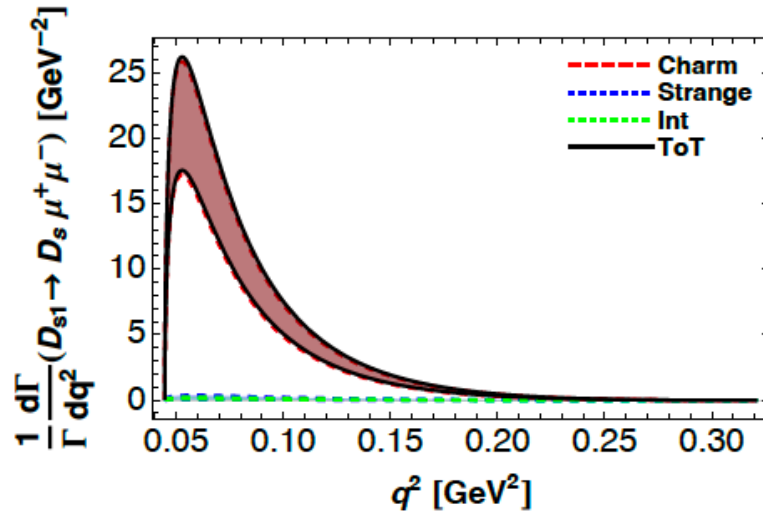




▪  $D_{s1} \rightarrow D_s^{(*)} \mu^+ \mu^-$

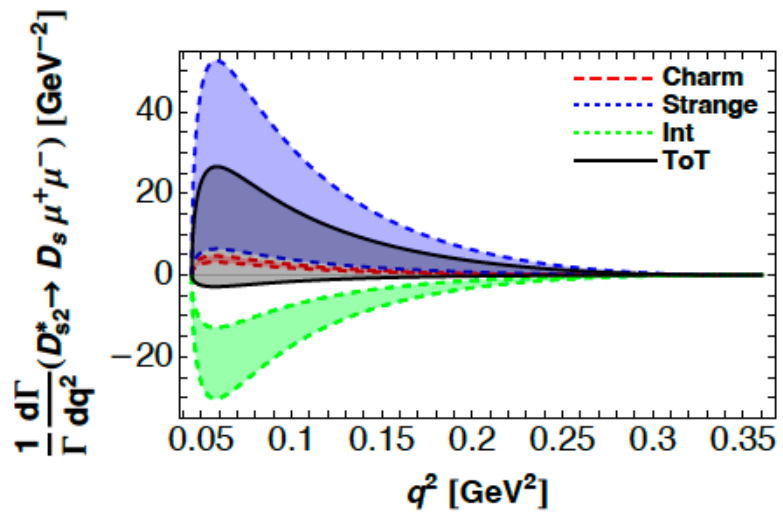
Case A

Case B

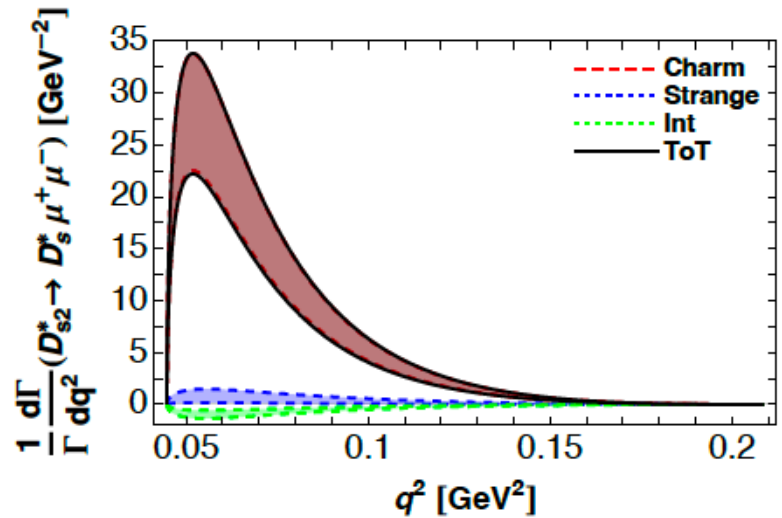
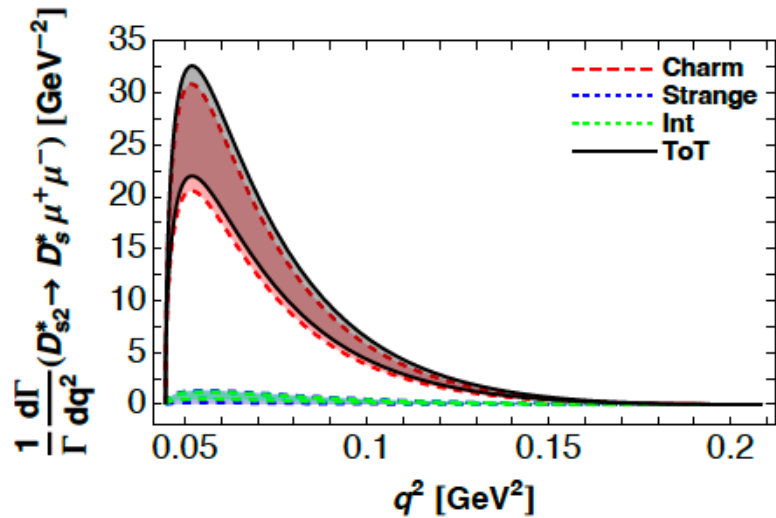
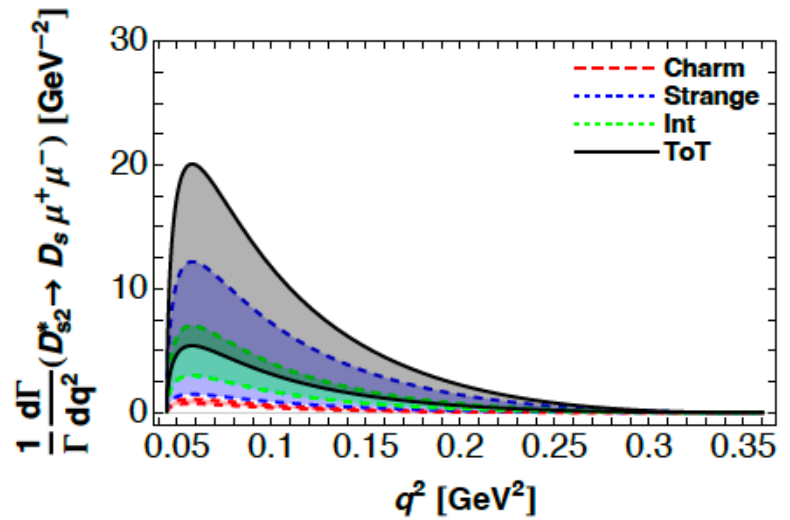


▪  $D_{s2}^* \rightarrow D_s^{(*)} \mu^+ \mu^-$

Case A



Case B



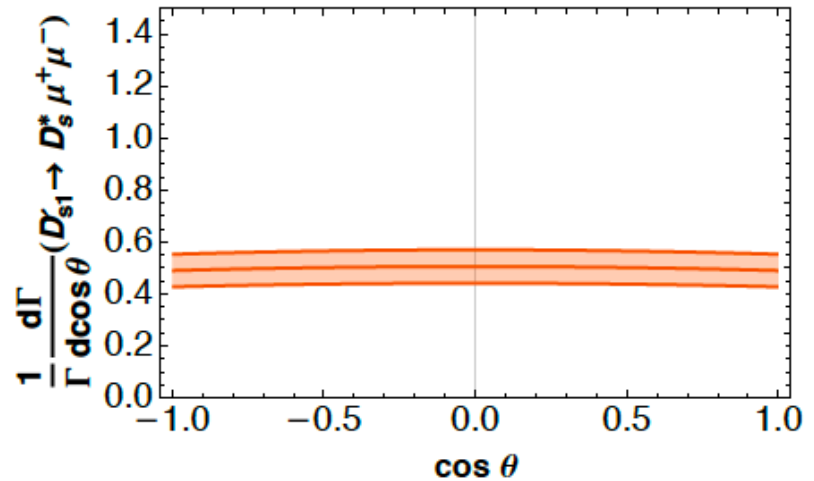
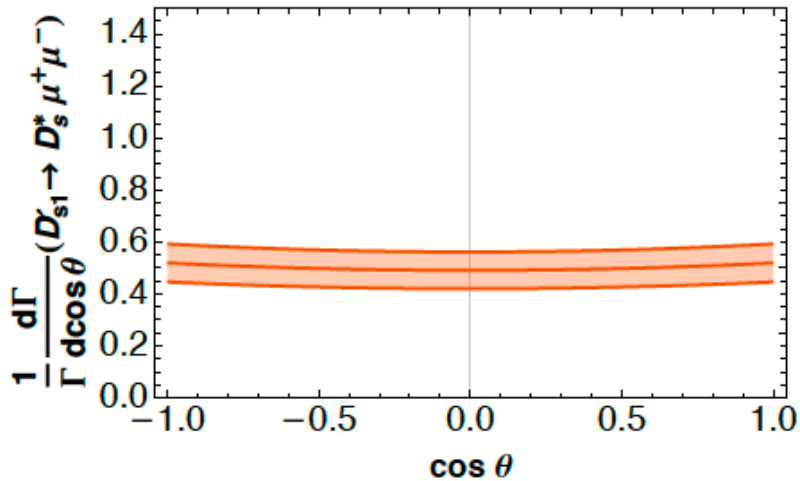
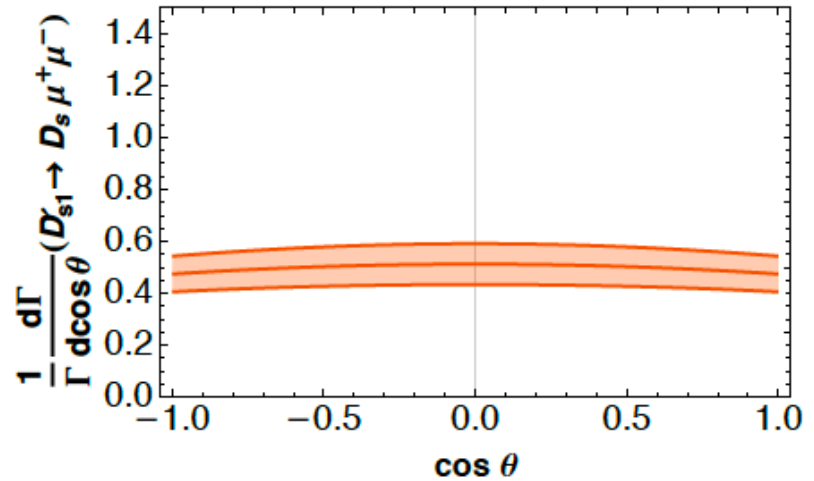
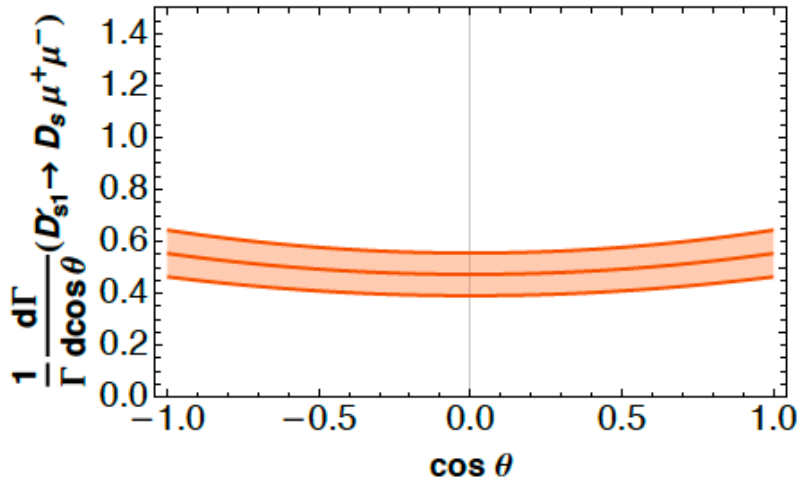
▪  $D'_{s1} \rightarrow D_s^{(*)} \mu^+ \mu^-$  Angular Distribution

$$\cos \theta = \frac{\vec{p}_1 \cdot \vec{p}}{|\vec{p}_1| |\vec{p}|} \quad \vec{p}_1 = \text{lepton momentum}$$

$$\vec{p} = D_s^{(*)} \text{ momentum}$$

Case A

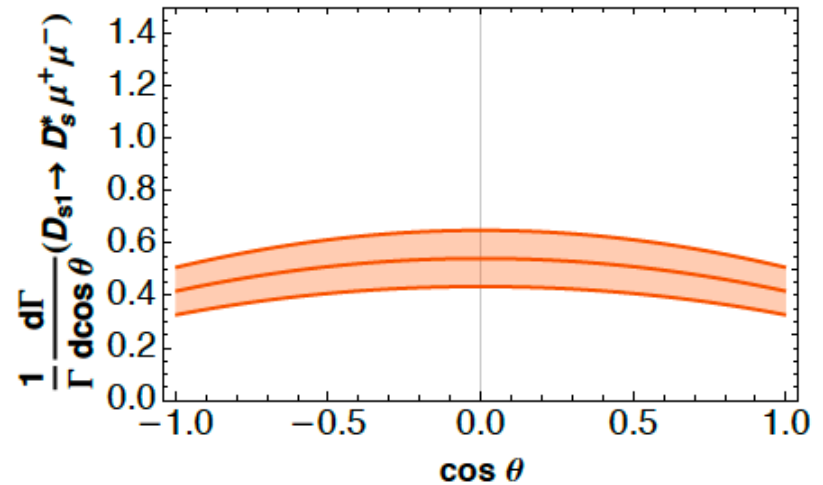
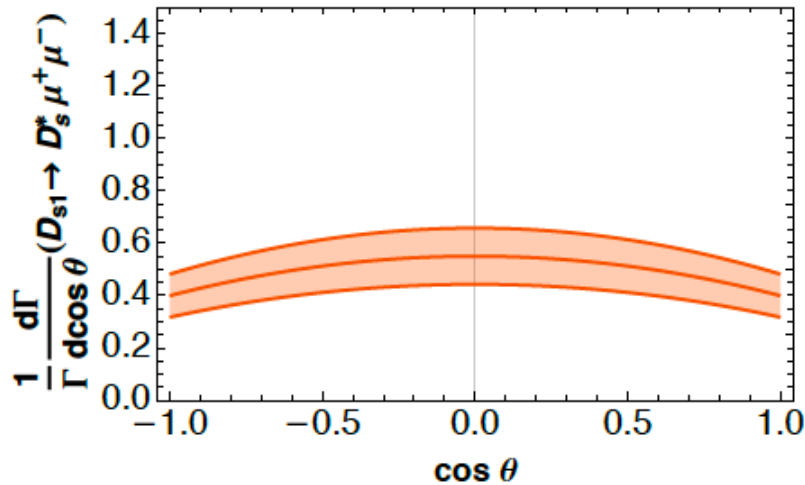
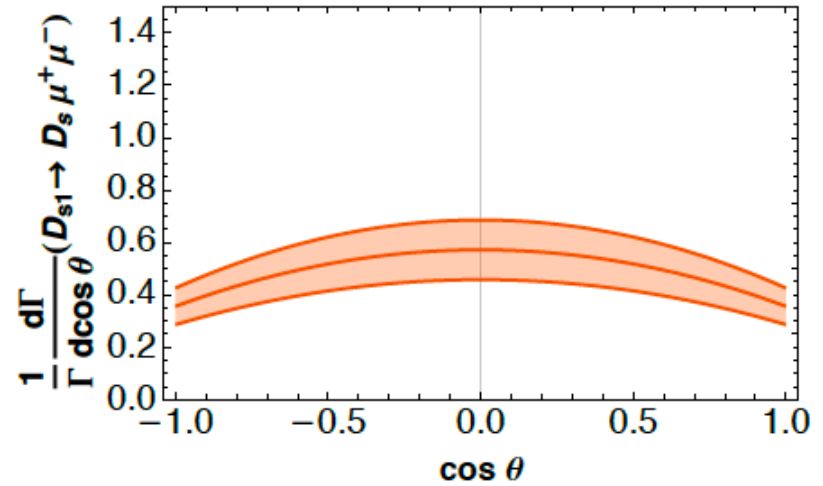
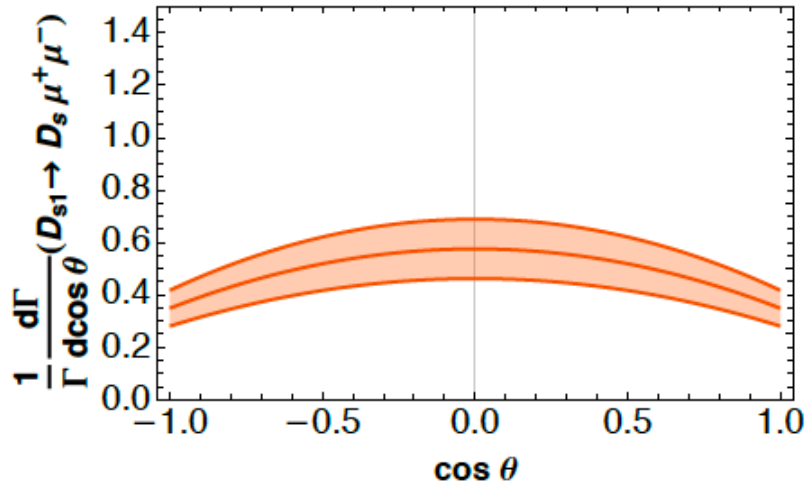
Case B



- $D_{s1} \rightarrow D_s^{(*)} \mu^+ \mu^-$  Angular Distribution

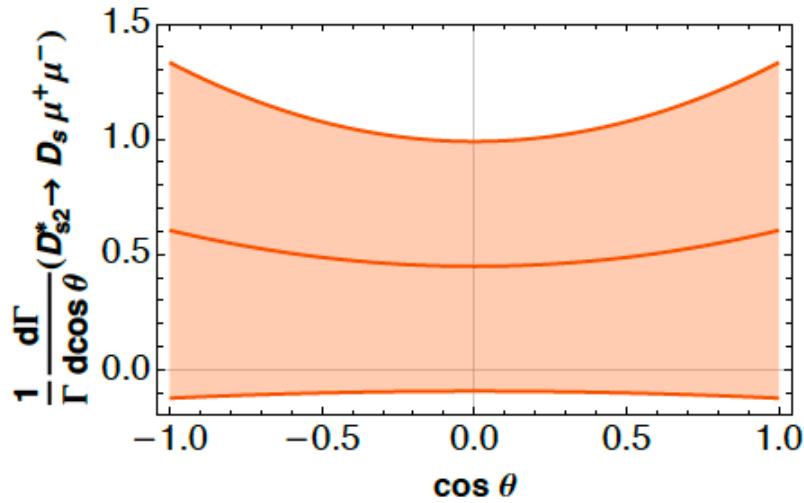
Case A

Case B

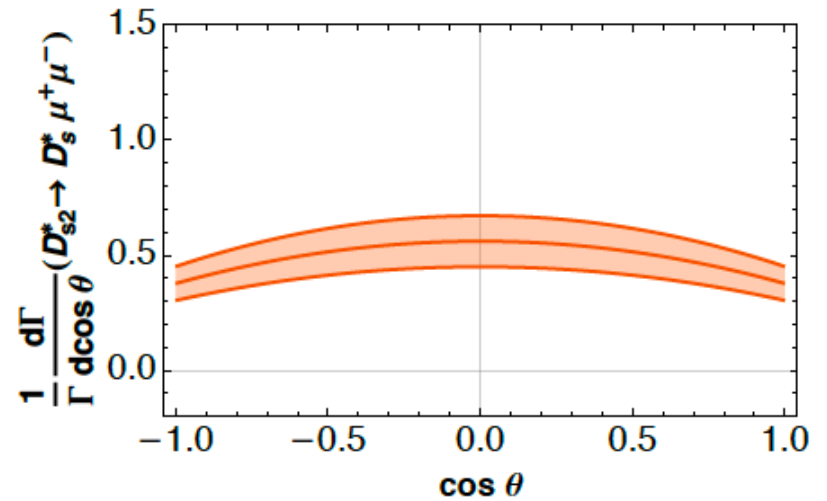
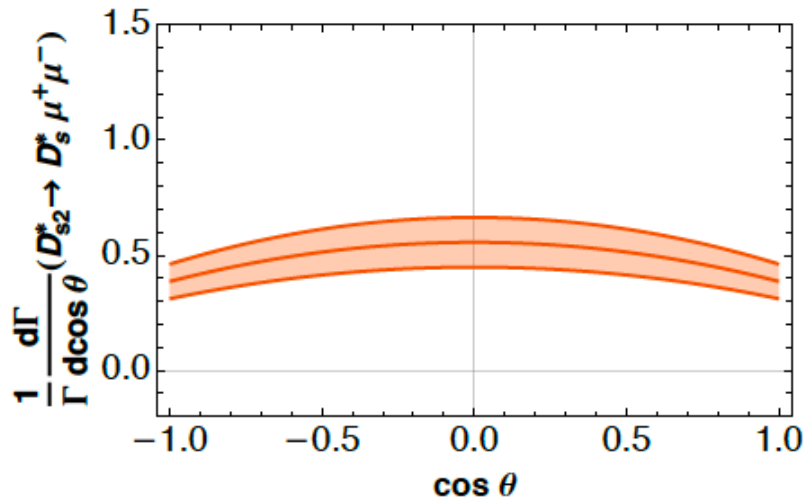
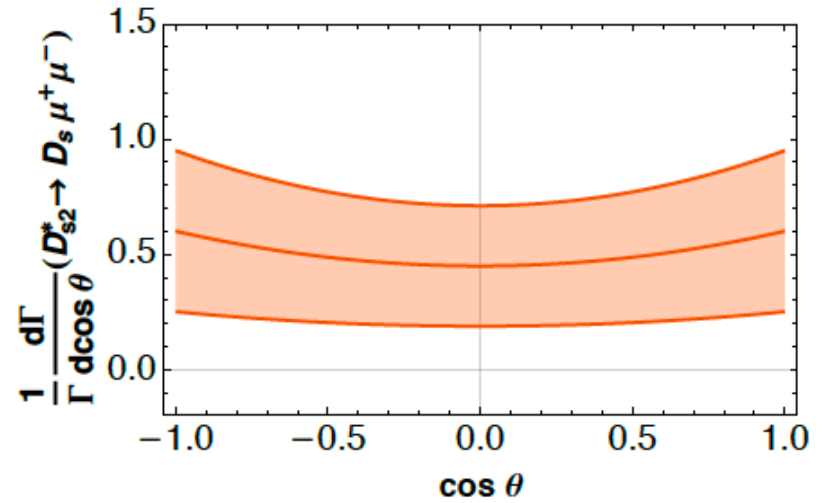


- $D_{s2}^* \rightarrow D_s^{(*)} \mu^+ \mu^-$  Angular Distribution

Case A



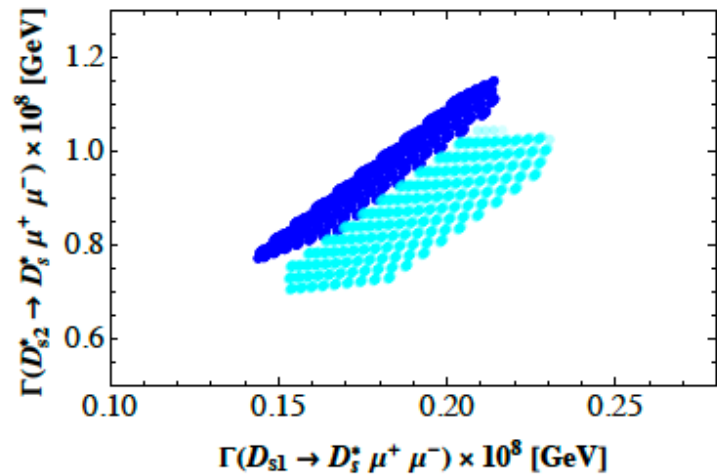
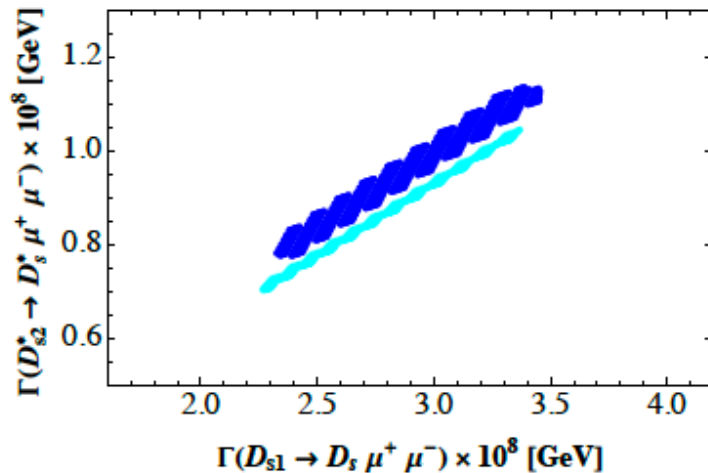
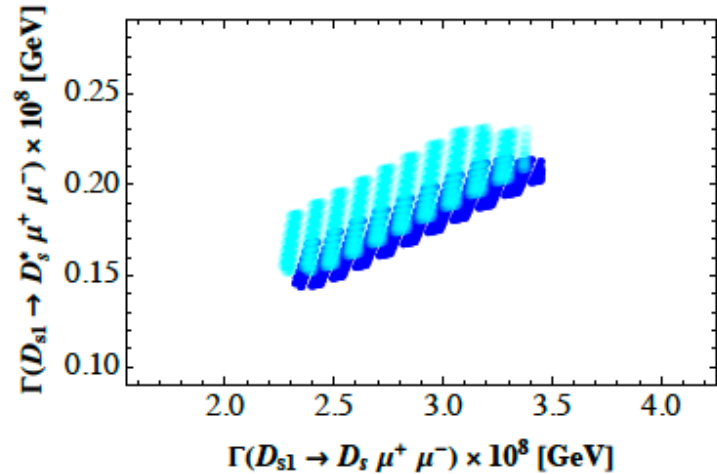
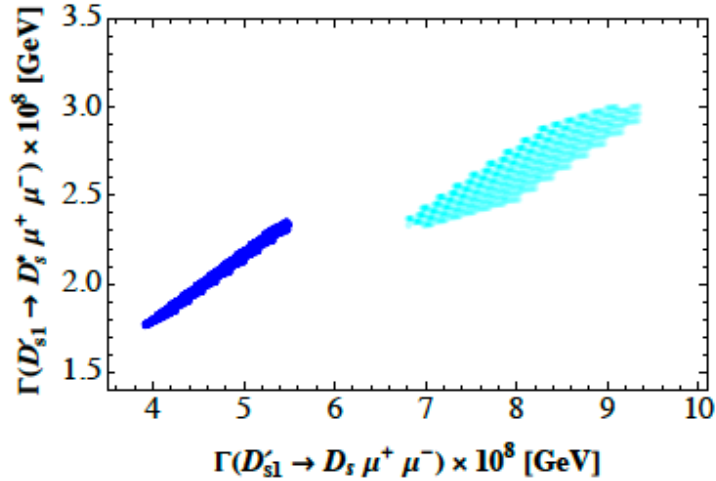
Case B



amplitudes of different modes related through the hadronic parameters



Correlations between decay widths



Although the decay widths are small these measurements are currently under investigation by the LHCb collaboration

# Conclusions

Analysis of the Dalitz decays of the positive parity  $D_{sJ}^{(*)}$  mesons

heavy quark spin symmetry  $\longrightarrow$  Classification in doublets

Observables:

- dilepton invariant mass distributions
- angular distributions
- decay widths correlations  $\longrightarrow$  Same strong couplings and form factors for the members of the same doublets

Experimentally confirmation  $\longrightarrow$  Validation of the classifications scheme  $\longrightarrow$  Hint to an ordinary  $c\bar{s}$  P-wave states interpretation for the  $D_{s0}^*$ ,  $D'_{s1}$  states

THANKS  
FOR YOUR  
ATTENTION