

Normal Forms for non linear betatronic motion

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Many thanks to Prof. G. Turchetti & Dr. M. Giovannozzi

Outline

- 1 Non linear charged particles motion
- 2 Normal Forms
- 3 Main application to betatronic motion

The equation of motion

4-D phase space: $\mathbf{x} = (x, p_x, y, p_y)$. The Hamiltonian function

$$\mathcal{H}(\mathbf{x}; s) = \frac{p_x^2 + p_y^2}{2} + \underbrace{\left(\frac{1}{\rho(s)^2} - k_1(s) \right) \frac{x^2}{2} + k_1(s) \frac{y^2}{2}}_{\text{Linear Motion}}$$

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$$\underbrace{-\text{Re} \left[\sum_{n=2}^M \frac{k_n(s) + i j_n(s)}{(n+1)!} (x + iy)^{n+1} \right]}_{\text{Non Linear Motion}}$$

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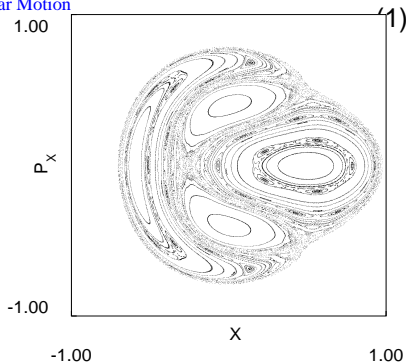
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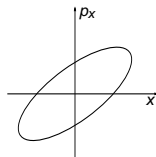
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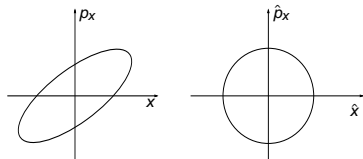
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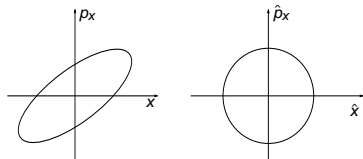
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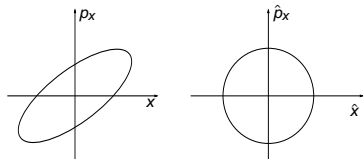
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$$\begin{pmatrix} \hat{x}' \\ \hat{p}'_x \\ \hat{y}' \\ \hat{p}'_y \end{pmatrix} = \mathbf{R}(\omega) \begin{pmatrix} \hat{x} \\ \hat{p}_x + \frac{K_2}{2} \beta_x^{3/2} (\hat{x}^2 - \beta \hat{y}^2) \\ \hat{y} \\ \hat{p}_y - K_2 \beta_x^{3/2} \beta \hat{x} \hat{y} \end{pmatrix}$$

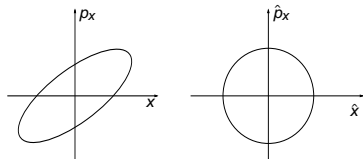
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The theory 1/2...the construction...

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$$\mathbf{z}' = \mathbf{F}(\mathbf{z}_1, \mathbf{z}_1^*, \mathbf{z}_2, \mathbf{z}_2^*) \Rightarrow \begin{cases} z_1' = e^{i\omega_1} \left(z_1 - \frac{i}{4} \left[(z_1 + z_1^*)^2 - \beta (z_2 + z_2^*)^2 \right] \right) \\ z_2' = e^{i\omega_2} \left(z_2 + \frac{i}{2} \beta (z_1 + z_1^*) (z_2 + z_2^*) \right) \end{cases} \quad (4)$$

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Normal forms coordinate $\xi = (\xi_1, \xi_1^*, \xi_2, \xi_2^*)$ and **NORMAL FORM** $\mathbf{U}(\xi)$ such that $\Delta_\alpha \mathbf{U} \equiv \mathbf{U}(e^{i\alpha}\xi, e^{-i\alpha}\xi^*) - e^{i\alpha}\mathbf{U}(\xi, \xi^*) = 0$

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$$\begin{array}{ccc} \mathbf{z} & \xrightarrow{\mathbf{F}} & \mathbf{z}' \\ \Phi \uparrow & & \uparrow \Phi \\ \xi & \xrightarrow{\mathbf{U}} & \xi' \end{array} \Rightarrow \begin{cases} \mathbf{z} = \Phi(\xi) = \underbrace{\xi}_{\text{Start with Identity}} + [\Phi(\xi)]_{n \geq 2} \\ \xi' = \underbrace{\mathbf{U}(\xi, \xi^*)}_{\mathbf{U} \text{ same linear motion of } \mathbf{F}} = \begin{cases} \xi'_1 = e^{i\omega_1} \xi_1 + [U_1(\xi)]_{n \geq 2} \\ \xi'_2 = e^{i\omega_2} \xi_2 + [U_2(\xi)]_{n \geq 2} \end{cases} \end{cases} \quad (5)$$

A. Bazzani, G. Servizi, E. Todesco and G. Turchetti, *CERN Yellow Report*, 94-02, 1994.

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$$\omega \cdot \mathbf{k} = 0 \leftrightarrow \mathbf{k} = l_1\mathbf{e}_1 + l_2\mathbf{e}_2$$

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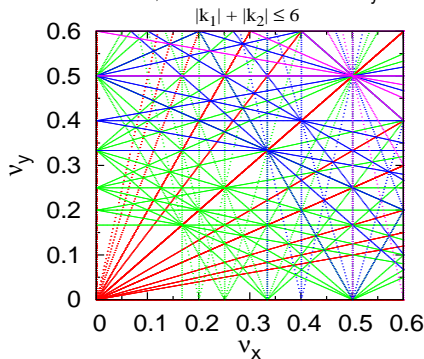
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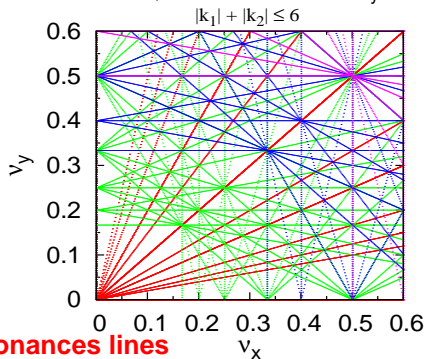
$$\omega \cdot \mathbf{k} = 0 \leftrightarrow \mathbf{k} = l\mathbf{e}$$

$$l \in \mathbb{N}, \mathbf{e} \in \mathbb{N}^3$$

DOUBLE-RESONANCE CASE

$$\omega \cdot \mathbf{k} = 0 \leftrightarrow \mathbf{k} = l_1\mathbf{e}_1 + l_2\mathbf{e}_2$$

$$l_j \in \mathbb{N}, \mathbf{e}_j \in \mathbb{N}^3$$



Intricate net of resonances lines

4D resonances and Henón map 2/2

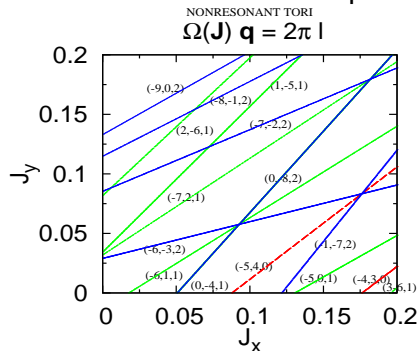
If we care about **non linear stuff**...

4D resonances and Henón map 2/2

If we care about **non linear stuff**... $\Omega = \Omega(j)$

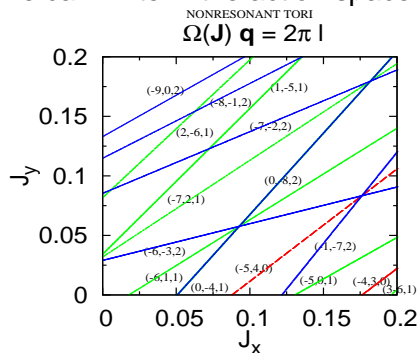
4D resonances and Henón map 2/2

If we care about **non linear stuff**... $\Omega = \Omega(j) \Rightarrow$ **amplitude detuning**
 We can write in the action space the resonance lines...



4D resonances and Henón map 2/2

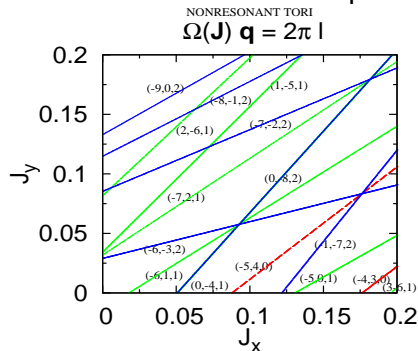
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lowest order

4D resonances and Henón map 2/2

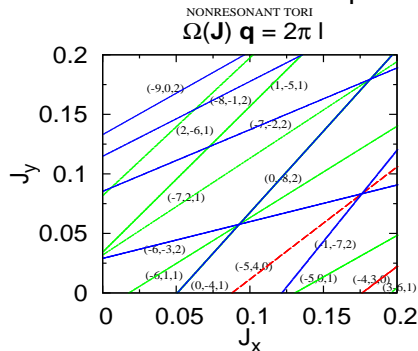
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lowest order \rightarrow **H quadratic** in \vec{j}

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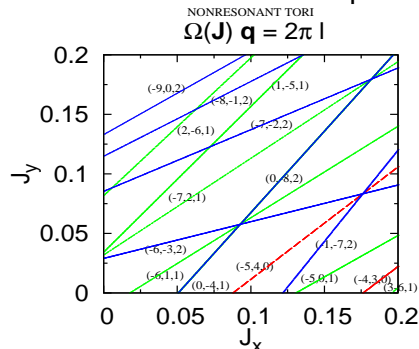


lowest order \rightarrow **H quadratic** in \vec{j}

$$\begin{cases} \Omega_x = \omega_x + 2\alpha_{20}j_1 + \alpha_{11}j_2 \\ \Omega_y = \omega_y + \alpha_{11}j_1 + 2\alpha_{02}j_2 \end{cases} \quad (9)$$

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Resonance conditions as function of emittances $\rightarrow \epsilon_j = 2j_j$

Details: A. Bazzani, L. Bongini, G. Turchetti "Analysis of resonances in action space for symplectic maps" Phys. Rev. E 57, 1178

2D Henón map and Normal Forms application 1/3

x-plane **sextupolar** dynamics from Eq. (3) if we let $\beta \rightarrow 0$

2D Henón map and Normal Forms application 1/3

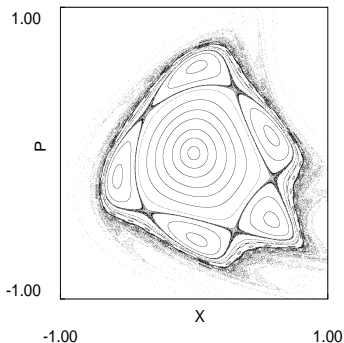
x-plane **sextupolar** dynamics from Eq. (3) if we let $\beta \rightarrow 0$

$$z' = F(z, z^*) = e^{i\omega} \left(z - \underbrace{\frac{i}{4} (z + z^*)^2}_{\text{Nonlinear term}} \right) \quad (10)$$

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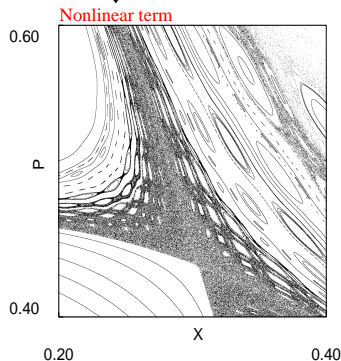
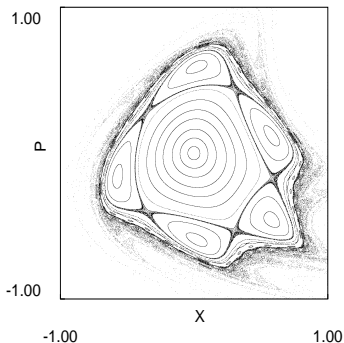
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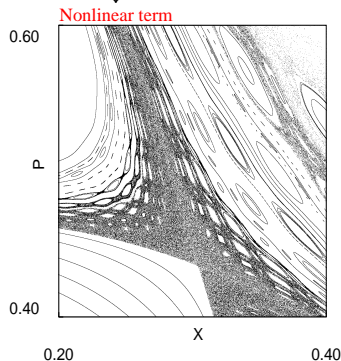
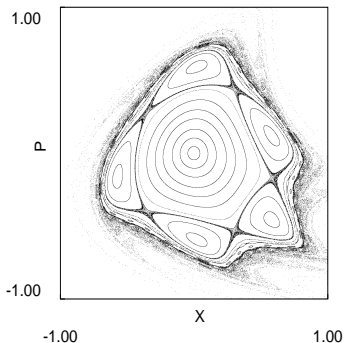


..phase portrait with $\nu = \omega/2\pi = 0.212\dots$

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..phase portrait with $\nu = \omega/2\pi = 0.212... \mathbf{1/5 \text{ resonance!}}$

2D Henón map and Normal Forms application 2/3

Analytical **nonlinear tune** value...

2D Henón map and Normal Forms application 2/3

Analytical **nonlinear tune** value... $\Omega = \omega + \Omega_2 j + \mathcal{O}(j^2)$

2D Henón map and Normal Forms application 2/3

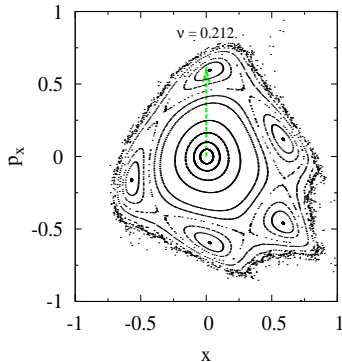
Analytical **nonlinear tune** value... $\Omega = \omega + \Omega_2 j + \mathcal{O}(j^2)$

The Average Phase Advance $\nu^{\text{APA}}(z; M) = \frac{1}{2\pi M} \sum_{j=1}^M (\vartheta^{\circ j} - \vartheta^{\circ(j-1)})$
gives us the **nonlinear tune**

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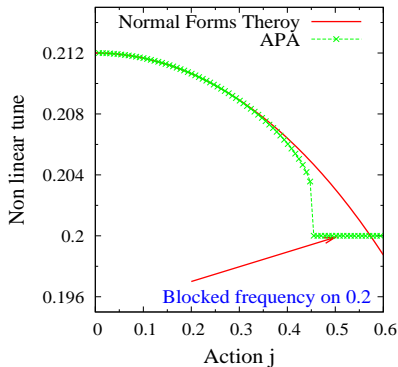
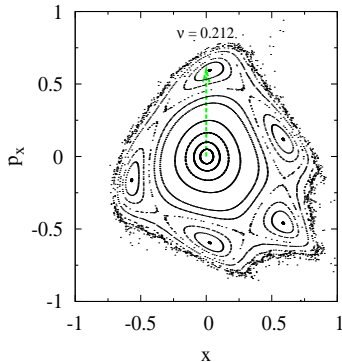
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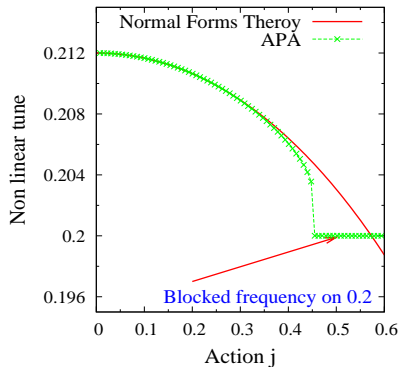
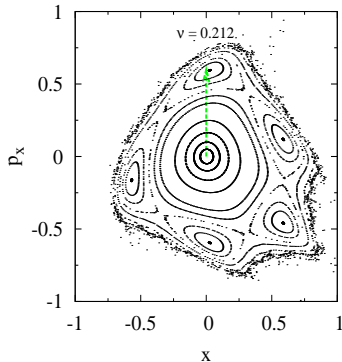
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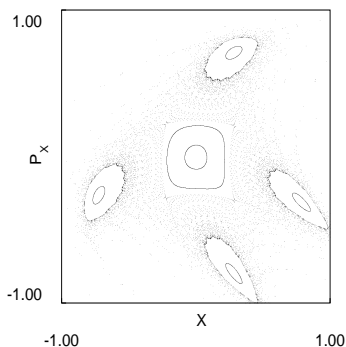
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Good agreement between numerical and theoretical values!!!

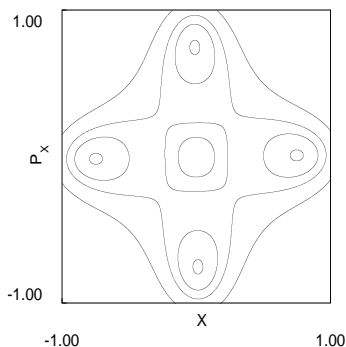
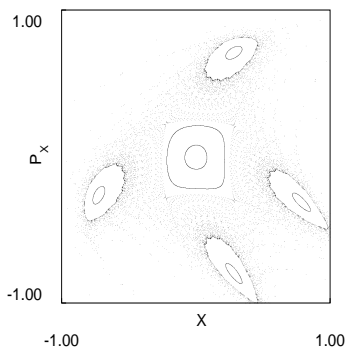
2D Henón map and Normal Forms application 3/3

We can also build up the **interpolating Hamiltonian** to control the system stability



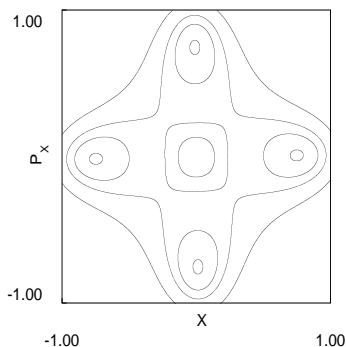
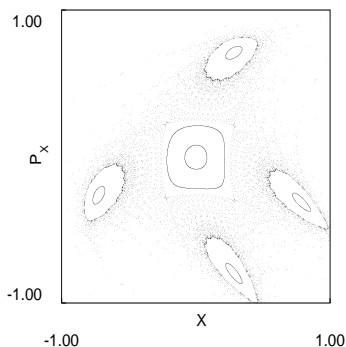
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Map iteration (left) flow of $H(\xi, \xi^*)$ (right) ν being close to $1/4$ resonance

Outline

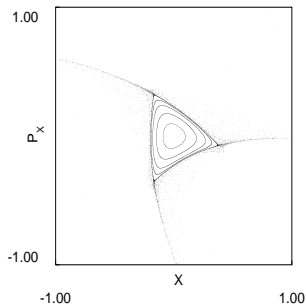
- 1 Non linear charged particles motion
- 2 Normal Forms
- 3 Main application to betatronic motion

Dynamic aperture 1/2

Studies on **slow extraction dynamic aperture**... exciting the 1/3 resonance

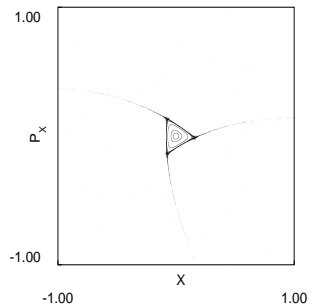
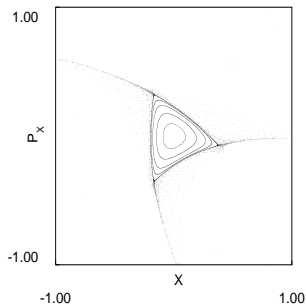
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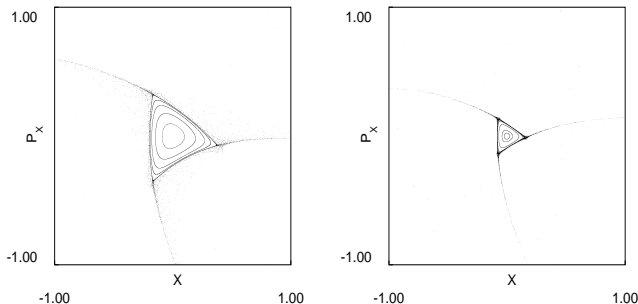
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Dynamic aperture 1/2

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Approaching the (**unstable!**) $\nu = 1/3$ sextupolar resonance...the dynamic aperture is shrinking... particles get extracted. Studies on **hyperbolic NF**

L. Bongini, A. Bazzani, G. Turchetti, I. Hofmann *Analysis of a model of resonant extraction of intense beams by normal forms and frequency map* Phys. Rev. Special Topics - Accelerators and Beams **4**, 114201 (2001)

Dynamic aperture 2/2

Nekhoroshev like estimates on the time stability.

Dynamic aperture 2/2

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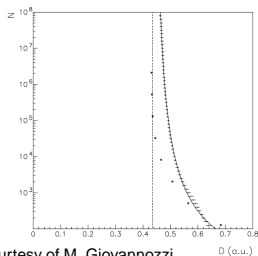
A particle in $\mathcal{B}(0; r/2)$ remains bounded in $\mathcal{B}'(0; r)$ for a time

$$\tau(r) \geq \tau_0 \exp \left[\left(\frac{r^*}{r} \right)^{\frac{2}{1+d}} \right] \quad \mathcal{D}(N) = \mathcal{D}_\infty \left(1 + \frac{b}{\log^k N} \right) \quad (11)$$

A. Bazzani, S. Marmi, G. Turchetti, *Nekhoroshev estimates for non resonant symplectic maps* *Celestial Mechanics* **47**, 333 (1990)

Numerical evidences of this scaling for the Hénon map and for a realistic 4-6D LHC model are given in

M. Giovannozzi, W. Scandale, E. Todesco *Dynamic aperture extrapolation in presence of tune modulation*, *Phys. Rev. E* **57**, 3432 (1998)



Picture is a courtesy of M. Giovannozzi

Multi Turn Extraction

Splitting the beam in phase space by means of **nonlinear elements**

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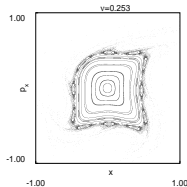
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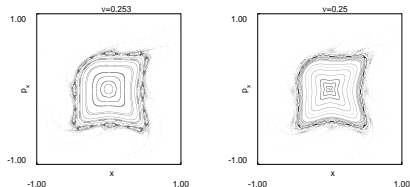


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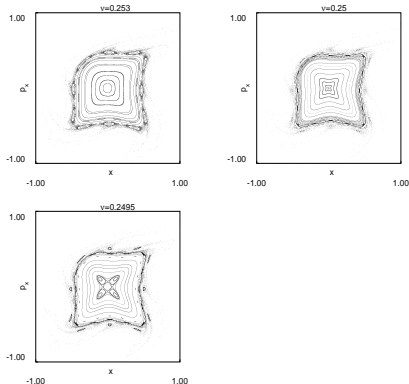


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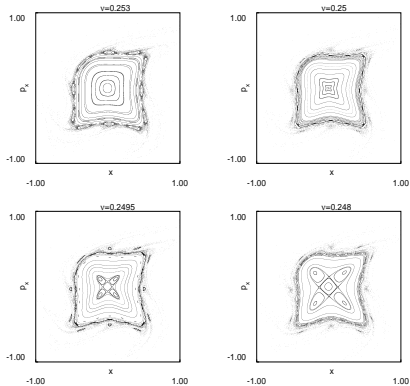


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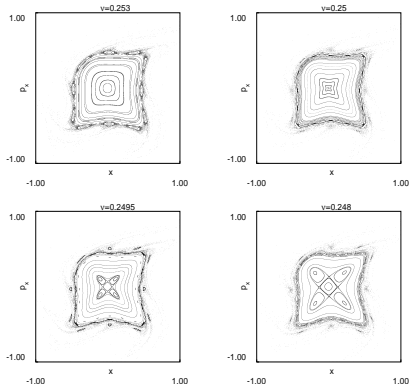


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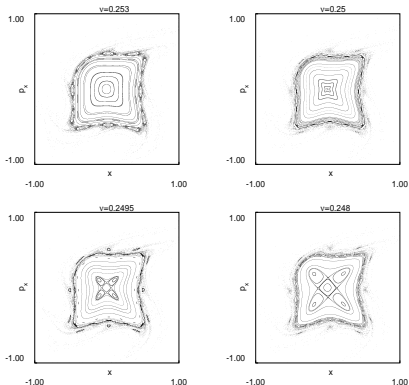
**@ CERN using NF
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Experimental data

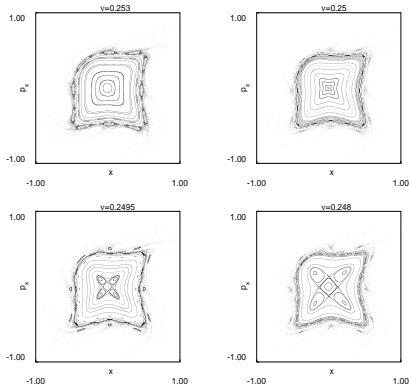
Movie is a courtesy of A. Franchi

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Experimental data

Movie is a courtesy of A. Franchi

Nonlinear parameters from NF analytical calculations!

D. Quatraro, Master Degree Thesis, Bologna University

Conclusion and future perspectives

- NF is an **efficient tool** to deal with nonlinear problems
- Analytical & numerical methods to get informations concerning the **stability** and the **dynamic aperture**
- Several **experiment @ CERN** (MTE) **agreed with NF theory**
- Extensions also to **space charge applications**
C. Benedetti, G. Turchetti *An analytic map for space charge in a nonlinear lattice* Physics Letters **A340**, 461-465 (2006)
- MTE **studies still ongoing @ CERN**