

Indirect constraints on top quark operators from a global SMEFT analysis

F. Garosi, D. Marzocca, A. Rodriguez-Sanchez, A. Stanzione

Assumption: only operators involving **top quarks** are generated at tree level (19 dim-6 operators)

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$

E.g. :

$$\mathcal{O}_{Hq}^{(1)} = (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{q}^3 \gamma^\mu q^3)$$

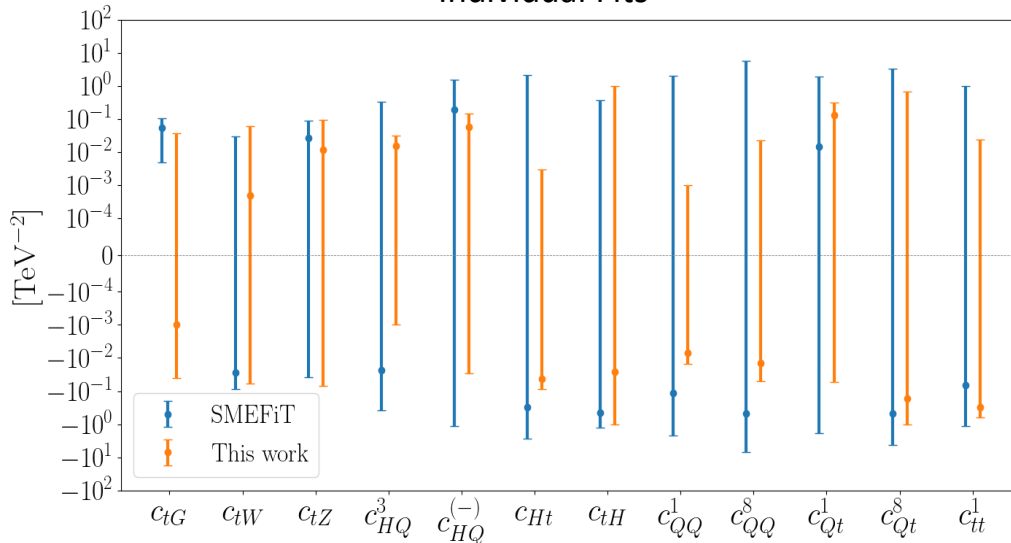
$$\mathcal{O}_{\ell q}^{(1), \alpha\beta} = (\bar{\ell}^\alpha \gamma_\mu \ell^\beta) (\bar{q}^3 \gamma^\mu q^3)$$

- B and K physics
- Higgs measurements
- Z and W decays
- Cabibbo angle
- g-2
- Lepton decays
- LFU tests
- LFV decay channels

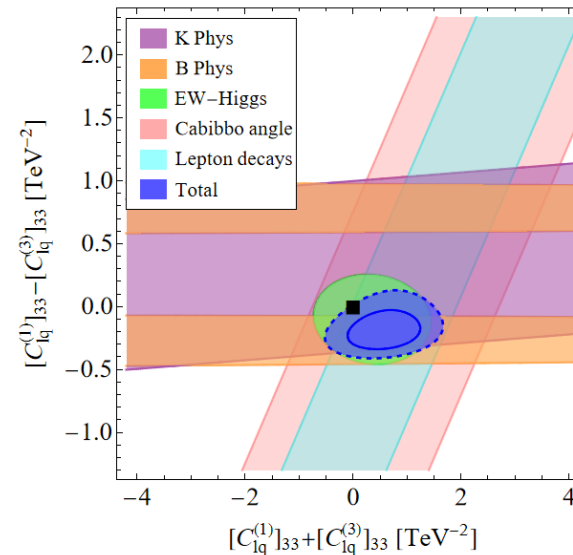
We study the impact of these operators on a large set of **low energy observables**
RGE and **1-loop matching** procedures are considered

We build a **global likelihood** and study some applications:

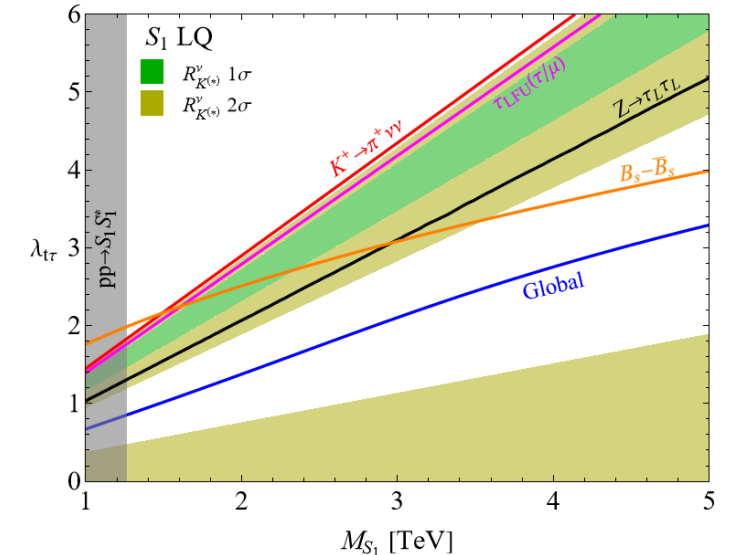
Individual Fits



Pairwise Fits



Consequences for UV models



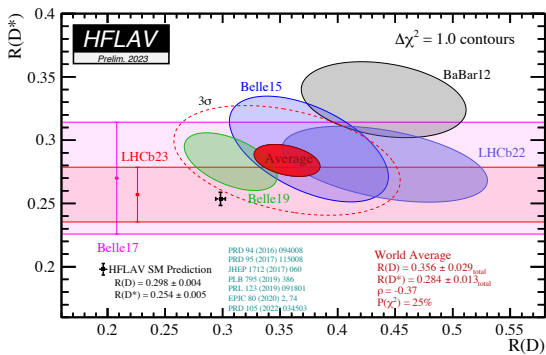
Relevance of B_c for the Standard Model and beyond

Nicola Losacco Dipartimento di Fisica, Università di Bari, INFN Bari

arXiv: 2205.08933

arXiv: 2208.13398

$b \rightarrow c \ell \bar{\nu}$ anomalies



need to improve theoretical
hadronic uncertainties

$$B_c \rightarrow M_{c\bar{c}} \ell \bar{\nu}$$



symmetry of the problem:
heavy quark spin symmetry



both negative and positive parity states

$J/\psi, \eta_c$ and χ_{cJ}, h_c

4-plet organization helps to get
hints on the structure of
debated states: $\chi_{c1}(3872)$

$X(3872)$ $\chi_{c1}(2P)$

$$R(D^{(*)}) = \frac{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell^- \bar{\nu}_\ell)}$$

doublet

4-plet

$$H(v) = \frac{1+\not{v}}{2} [B_c^{*\mu} \gamma_\mu - B_c \gamma_5] \frac{1-\not{v}}{2}$$

$$H'(v') = \frac{1+\not{v}'}{2} [\Psi_c^{*\mu} \gamma_\mu - \eta_c \gamma_5] \frac{1-\not{v}'}{2}$$

$$\mathcal{M}^\mu(v') = \frac{1+\not{v}'}{2} \left[\chi_{c2}^{\mu\nu} \gamma_\nu + \frac{1}{\sqrt{2}} \chi_{c1,\gamma} \epsilon^{\mu\alpha\beta\gamma} v'_\alpha \gamma_\beta + \frac{1}{\sqrt{3}} \chi_{c0} (\gamma^\mu - v'^\mu) + h_c^\mu \gamma_5 \right] \frac{1-\not{v}'}{2}$$

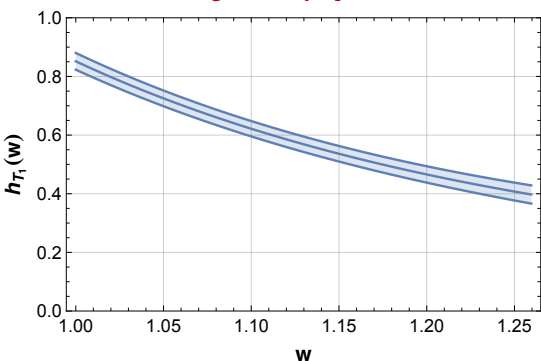
$$\langle M'(v') | J_0 | M(v) \rangle = -\Delta(w) \text{Tr} [\bar{H}'(v') \Gamma H(v)]$$

heavy quark expansion

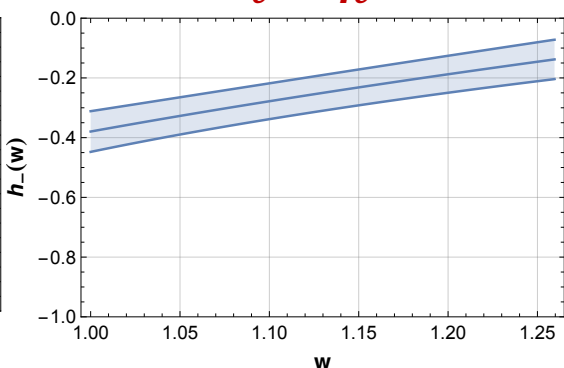
universal function describing the form
factors at LO

few results:

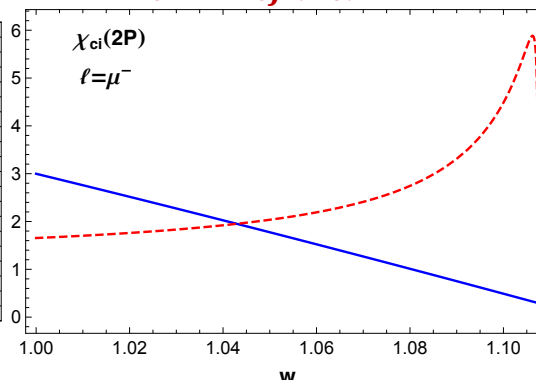
$B_c \rightarrow J/\psi \ell \bar{\nu}$



$B_c \rightarrow \eta_c \ell \bar{\nu}$



$B_c \rightarrow \chi_{cJ}(h_c) \ell \bar{\nu}$



$$\frac{d\Gamma(B_c \rightarrow \chi_{c1} \ell \bar{\nu})/dw}{d\Gamma(B_c \rightarrow \chi_{c0} \ell \bar{\nu})/dw}$$

$$\frac{d\Gamma(B_c \rightarrow \chi_{c2} \ell \bar{\nu})/dw}{d\Gamma(B_c \rightarrow \chi_{c1} \ell \bar{\nu})/dw}$$