Francesco Sanfilippo, INFN Roma Tre XXXVII Convegno Nazionale di Fisica Teorica 27-29 Sept 2023

## 

Parametrization of the interaction of a Muon with an external magnetic field

$$
\mu=g \frac{e}{2 m} S
$$



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Giromagnetic factor $g$ : relation to the particle spin $S$
Proton $g_{p}=5.5856946893$
Neutron $\quad g_{n}=-3.82608545$
Electron $g_{e}=-2.00231930436256$
Muon $g_{\mu}=-2.0023318418$

## 

Parametrization of the interaction of a Muon with an external magnetic field

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Giromagnetic factor $g$ : relation to the particle spin $S$


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\end{aligned}
$$

Pointlike particles: $g=2$ Dirac equation 1928

## Loop corsection to g=?

Vacuum polarization renormalizes $g$

$$
a=\frac{g-2}{2}
$$

Anomalous magnetic moment


## Loop cossection to g=?

Vacuum polarization renormalizes $g$

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Anomalous magnetic moment

... and possibly, any sort of unknown particle

## Loop cossection to g=2

Vacuum polarization renormalizes $g$

$$
a=\frac{g-2}{2}
$$

Anomalous magnetic moment

... and possibly, any sort of unknown particle
Measure precisely $a \rightarrow$ probe completeness of the Standard Model

## 

 electron : $\quad a_{e}=0.00115965218073$ muon : $\quad a_{\mu}=0.00116592089$tau: $\quad a_{\tau}=0$

## Thes andunalobs magnetic monent of leptons

$$
\begin{aligned}
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Electron is stable: 1000 times more precise than muon

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It also need small-size apparatus for measurement


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Electron is stable: 1000 times more precise than muon
It also need small-size apparatus for measurement
Why not studying the electron?

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Simple dimensional analysis: $a_{\ell}^{N P} \sim \kappa m_{\ell}^{2} / m_{N P}^{2}$

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Muon wins over electrons by a factor $m_{\mu}^{2} / m_{e}^{2} \sim 43000$

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Electron is stable: 1000 times more precise than muon
It also need small-size apparatus for measurement
Why not studying the electron?
Simple dimensional analysis: $a_{\ell}^{N P} \sim \kappa m_{\ell}^{2} / m_{N P}^{2}$
Muon wins over electrons by a factor $m_{\mu}^{2} / m_{e}^{2} \sim 43000$
Tau would be even better, but decays too fast to measure (but there are ideas)

## T'se anconsalobs nsagnecic noment of rnuon

## 

BNL E821 exp. up to 2006

## 



BNL E821 exp. up to 2006


## 



BNL E821 exp. up to 2006



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BNL E821 exp. up to 2006



## 



BNL E821 exp. up to 2006




Transfer of the ring to Fermilab

## 



BNL E821 exp. up to 2006


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$g_{\mu}-2$ experiment @ Fermilab

## 



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$g_{\mu}-2$ experiment @ Fermilab


2023: measurement confirmed

## 



BNL E821 exp. up to 2006


Transfer of the ring to Fermilab



$g_{\mu}-2$ experiment @ Fermilab


2023: measurement confirmed What about the theory?!?

## 

 https://muon-gm2-theory.illinois.edu/Target: match the theory precision \& accuracy with the upcoming g-2 experiment White paper: Physics Reports 887 (2020) 1-166 [arXiv:2006.04822]

Regular meetings: latest in Bern, 4-8 September 2023

The anomalous magnetic moment of the muon in the Standard Model
T. Aoyama ${ }^{1,2,3}$, N. Asmussen ${ }^{4}$, M. Benayoun ${ }^{5}$, J. Bijnens ${ }^{6}$, T. Blum ${ }^{7,8}$, M. Bruno ${ }^{9}$, I. Caprini ${ }^{10}$ C. M. Carloni Calame ${ }^{11}$, M. Cè ${ }^{9,12,13}$, G. Colangelo ${ }^{\dagger 14}$, F. Curciarello ${ }^{15,16}$, H. Czyż ${ }^{17}$, I. Danilkin ${ }^{12}$, M. Davier ${ }^{\dagger 18}$ C. T. H. Davies ${ }^{19}$, M. Della Morte ${ }^{20}$, S. I. Eidelman ${ }^{\dagger 21,22}$, A. X. El-Khadra ${ }^{\dagger 23,24}$, A. Gérardin ${ }^{25}$, D. Giusti ${ }^{26,27}$ M. Golterman ${ }^{28}$, Steven Gottlieb ${ }^{29}$, V. Gülpers ${ }^{30}$, F. Hagelstein ${ }^{14}$, M. Hayakawa ${ }^{31,2}$, G. Herdoíza ${ }^{32}$, D. W. Hertzog ${ }^{33}$ A. Hoecker ${ }^{34}$, M. Hoferichter ${ }^{\dagger 14,35}$, B.-L. Hoid ${ }^{36}$, R. J. Hudspith ${ }^{12,13}$, F. Ignatov ${ }^{21}$, T. Izubuchi ${ }^{37,8}$, F. Jegerlehner ${ }^{38}$ . L. Lellouch $^{25}$ I. Logashenko ${ }^{21}$, B. Malaescu ${ }^{5}$, K. Maltman ${ }^{44,45}$, M. Kupsc Marinković ${ }^{46,47}$, P. Masjuanne ${ }^{48,49}$ L. Lellouch ${ }^{25}$, I. Logashenko ${ }^{21}$, B. Malaescu ${ }^{5}$, K. Maltman ${ }^{44,45}$, M. K. Marinkovi chen $^{46,47}$, P. Masjuan ${ }^{48,49}$
A. S. Meyer ${ }^{37}$, H. B. Meyer ${ }^{12,13}$, T. Mibe ${ }^{\dagger 1}$, K. Miura ${ }^{12,13,3}$, S. E. Müller ${ }^{50}$, M. Nio ${ }^{2,51}$, D. Nomura ${ }^{52,53}$,
A. Nyffeler ${ }^{\dagger 12}$, V. Pascalutsa ${ }^{12}$, M. Passera ${ }^{54}$, E. Perez del Rio ${ }^{55}$, S. Peris ${ }^{48,49}$, A. Portelli ${ }^{30}$, M. Procura ${ }^{56}$,
C. F. Redmer ${ }^{12}$, B. L. Roberts ${ }^{\dagger 57}$, P. Sánchez-Puertas ${ }^{49}$, S. Serednyakov ${ }^{21}$, B. Shwartz ${ }^{21}$, S. Simula ${ }^{27}$, D. Stöckinger ${ }^{58}$, H. Stöckinger-Kim ${ }^{58}$, P. Stoffer ${ }^{59}$, T. Teubner ${ }^{\dagger 60}$, R. Van de Water ${ }^{24}$, M. Vanderhaeghen ${ }^{12,13}$ G. Venanzoni ${ }^{61}$, G. von Hippel ${ }^{12}$, H. Wittig ${ }^{12,13}$, Z. Zhang ${ }^{18}$
M. N. Achasov ${ }^{21}$, A. Bashir ${ }^{62}$, N. Cardoso ${ }^{47}$, B. Chakraborty ${ }^{63}$, E.-H. Chao ${ }^{12}$, J. Charles ${ }^{25}$, A. Crivellin ${ }^{64,6 .}$. O. Deineka ${ }^{12}$, A. Denig ${ }^{12,13}$, C. DeTar ${ }^{66}$, C. A. Dominguez ${ }^{67}$, A. E. Dorokhov ${ }^{68}$, V. P. Druzhinin ${ }^{21}$, G. Eichmann ${ }^{69,47}$,
M. Fael ${ }^{70}$, C. S. Fischer ${ }^{71}$, E. Gámiz ${ }^{72}$, Z. Gelzer ${ }^{23}$, J. R. Green ${ }^{9}$, S. Guellati-Khelifa ${ }^{73}$, D. Hatton ${ }^{19}$
N. Hermansson-Truedsson ${ }^{14}$, S. Holz ${ }^{36}$, B. Hörz ${ }^{74}$, M. Knecht ${ }^{25}$, J. Koponen ${ }^{1}$, A. S. Kronfeld ${ }^{24}$, J. Laiho ${ }^{75}$,
S. Leupold ${ }^{42}$, P. B. Mackenzie ${ }^{24}$, W. J. Marciano ${ }^{37}$, C. McNeile ${ }^{76}$, D. Mohler ${ }^{12,13}$, J. Monnard ${ }^{14}$, E. T. Neil ${ }^{17}$
A. V. Nesterenko ${ }^{68}$, K. Ottnad ${ }^{12}$, V. Pauk ${ }^{12}$, A. E. Radzhabov ${ }^{78}$, E. de Rafael ${ }^{25}$, K. Raya ${ }^{79}$, A. Risch ${ }^{12}$, A. Rodríguez-Sánchez ${ }^{6}$, P. Roig ${ }^{80}$, T. San José ${ }^{12,13}$, E. P. Solodov ${ }^{21}$, R. Sugar ${ }^{81}$, K. Yu. Todyshev ${ }^{21}$, A. Vainshtein ${ }^{82}$ A. Vaquero Avilés-Casco ${ }^{66}$, E. Weil ${ }^{71}$, J. Wilhelm ${ }^{12}$, R. Williams ${ }^{11}$, A. S. Zhevlakov ${ }^{88}$


## Electroyeals contributions

T. Aoyama, M. Hayakawa,
T. Kinoshita, M. Nio
[PRLs, 2012]


## Electroyeals contributions

T. Aoyama, M. Hayakawa,
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[PRLs, 2012]

Computed up to $5^{\text {th }}$ order!!!
~10000 diagrams


## Lignts-by-lignt contribution



- Related to two-photons scattering
- Nasty hadronic contribution
- Long distance effects hard to compute
- Nonperturbative contribution



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Dispersive approach

[several contributions put together]

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Dispersive approach

[several contributions put together]

- Tiny, but model-dependent


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Dispersive approach


(a)

(b)

(c)

$$
a_{\mu}^{l b l, h a d}=92(19) \cdot 10^{-11}
$$

[several contributions put together]

- Tiny, but model-dependent
? Size comparable to the full HVP error...


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Lattice calculations

[RBC/UKQCD coll, PRL 124, 2020]

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[several contributions put together]

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Lattice calculations

[RBC/UKQCD coll, PRL 124, 2020]

- First principle calculation
$\square$ Larger error, but validates model
faclrontic vactulas polarizetion



## fanclrontic vacudus polarization

$a_{\mu}^{H V P}=\left(\frac{\alpha}{\pi}\right)^{2} \int_{0}^{\infty} d Q^{2} f\left(Q^{2}\right)\left[\Pi\left(Q^{2}\right)-\Pi(0)\right]$
analytic kernel vectorial polarization

Vector Correlation function

$$
C_{\mu \nu}(x)=\left\langle j_{\mu}(x) j_{\nu}(0)\right\rangle
$$

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Polarization Tensor

$$
\Pi_{\mu \nu}\left(Q^{2}\right)=\int d^{4} x e^{i Q x} C_{\mu \nu}(x)
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## fanclrontic vacudus polarization

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analytic kernel vectorial polarization

Vector Correlation function

$$
\operatorname{Cuv}_{\mu}(\mathfrak{C})=\left\langle\dot{J}_{\mu}(\underset{\mathscr{C}}{ }) \dot{J}_{\nu}(0)\right\rangle
$$



Polarization Tensor

- Long distance contributions

Vectorial Polarization

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Polarization Tensor

- Long distance contributions
- Nonperturbative QCD effects

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## fagaronic vactulas polarization

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$$

- Long distance contributions
- Nonperturbative QCD effects
- How to evaluate?!!?


## How to evaluate the flyp?

## 

...replace it with another, unrelated experimental measurement!

## flow io evaluate the flyp?

...replace it with another, unrelated experimental measurement!

## Optical theorem

Elastic scattering amplitude
Total $\mathrm{e}^{+} \mathrm{e}^{-}$cross section

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## Optical theorem

Elastic scattering amplitude

$=$
$=$

Total $\mathrm{e}^{+} \mathrm{e}^{-}$cross section

hadrons

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## Optical theorem

Elastic scattering amplitude


$$
a_{\mu}^{H V P}=\int_{0}^{\infty} d Q^{2} K\left(Q^{2}\right) \hat{\Pi}\left(Q^{2}\right)=\int_{0}^{\infty} R(E) K(E) d E
$$

$=$

hadrons

## flow io evaluate the flyp?

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## Optical theorem

Elastic scattering amplitude

$=$


Total $\mathrm{e}^{+} \mathrm{e}^{-}$cross section $e^{+}$e

Can we call this "theoretical prediction"...?

## fow io evaluate the flyp?

...replace it with another, unrelated experimental measurement!

## Optical theorem

Elastic scattering amplitude =

$=$


Total $\mathrm{e}^{+} \mathrm{e}^{-}$cross section $e^{+}$e $\gamma$ ~n hadrons

Can we call this "theoretical prediction"...?
NO! We are plugging a substantial experimental input

## Electron-posjiton cross section o



Probability of electron-positrons to annihilate into hadrons

$$
R(E)=\frac{\sigma\left(e^{+} e^{-} \rightarrow \text { hadrons }\right)}{\sigma\left(e^{+} e^{-} \rightarrow \mu^{+} \mu^{-}\right)}
$$

normalizing each energy $E$ with the annihilation into muons

## 玉゙ecirosj-posjirosu cross secijon o



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normalizing each energy $E$ with the annihilation into muons

A number of worldwide experiments since the early ' 60


KLOE @ DAФNE FRASCATI


BABAR @ SLAC STANFORD


CMD3 @ VEPP-2000 NOVOSIBIRSK

## 




## 

|  | $\vdots$ |  |  |
| :---: | :---: | :---: | :--- |
|  |  |  |  |

## fyp ifors phenomenological R-ratio

|  |  |  |  |
| :---: | :---: | :---: | :--- |
|  |  | GR 1969 |  |
|  |  |  |  |


M. Davier, A. Hoecker, B. Malaescu, Z. Zhang,
[Eur. Phys. J. C 80 (2020) 241]

## 「'se restowssed g-2 puzzle



August 2023: release of Run III results by Fermilab g-2 experiment

## Scientists may be on brink of discovering fifth force of nature

Experts closing in on potentially identifying new force after surprise wobble of subatomic particle

© The muon g-2 ring sits in its detector hall at the Fermilab in Illinois. Photograph: Ryan
Postel/Fermi national accelerator laboratory/Reuters
The tantalising theory that a fifth force of nature could exist has been given a boost thanks to unexpected wobbling by a subatomic particle, physicists have revealed.

## folel oss, fijitis fiorce!

## Do we really control the theory uncertainties?

After all, we are replacing HVP with a combination of other experiments
Let us look back at the R-ratio...

## fold on, filith force!

## Do we really control the theory uncertainties?

After all, we are replacing HVP with a combination of other experiments
Let us look back at the R-ratio...

Hints of tension in the two-pions final state


Disagreement of 2023 CMD3 measurement


Leís folt inte cjuestions back on track

## Leさt's put infe cjuestions back on track

- "Computing" HVP via dispersive method is the weakest part of the story


## Lér's put inse questions back on track

- "Computing" HVP via dispersive method is the weakest part of the story
- Can we compute for real HVP from the first principle of the theory?


## Lér's put inse questions back on track



- "Computing" HVP via dispersive method is the weakest part of the story
- Can we compute for real HVP from the first principle of the theory?
$\rightarrow$ Lattice QCD comes to the rescue!


## Consoptijstg finge firons the first principles

## Original proposal

"Lattice Calculation of the Lowest-Order Hadronic Contribution to the Muon Anomalous Magnetic Moment"
[T. Blum, PRL 91 (2003)]

## Consoputing finla firons the fiost prinaciples

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Fourier transform of lattice-computed correlation function

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\Pi_{\mu \nu}\left(Q^{2}\right)=\int d^{4} x e^{i Q x} C_{\mu \nu}(x)
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"simple" two points correlation function


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$$

"simple" two points correlation function
Issue: Convolution kernel enhances $Q^{2} \sim m_{\mu}^{2} \sim 0.01 \mathrm{GeV}^{2}$

- Momenta on the lattice are quantized
- Lowest momenta are very noisy



## Larijce OdCD sissumation

First principle simulation of strong interactions
Quantum modynamics on a Lattice
Euclidean spacetime with $\mathrm{O}\left(10^{10}\right)$ degrees of freedom


## Laticice ? CD simsulation

First principle simulation of strong interactions

## Quantum modynamics on a Lattice

Euclidean spacetime with $\mathrm{O}\left(10^{10}\right)$ degrees of freedom

Hybrid Monte Carlo + Molecular Dynamics simulations


Numerical solution of the discrete Dirac Equation (partial derivative equation $\rightarrow$ large sparse matrix)

## Laricice ? OCD sissumation

First principle simulation of strong interactions

## Quantum modynamics on a Lattice

Euclidean spacetime with $\mathrm{O}\left(10^{10}\right)$ degrees of freedom

Hybrid Monte Carlo + Molecular Dynamics simulations


Numerical solution of the discrete Dirac Equation (partial derivative equation $\rightarrow$ large sparse matrix)

A long list of scientific achievements:

- \% reconstruction of the hadron spectrum, $\square$



## Cosselatios fusuction iss laticice QCD

Task \#1: Producing O(100-1000) "configurations" of gluonic fields.

- 1 configuration: $O(1-50 G B$ data) $\sim 1$ day of simulation on $O(5000)$ cores.
- 100 s MCoreHours in national, European \& worldwide supercomputers
- Similar in spirit to storing collision events at particle accelerators $\rightarrow$ a handful of collaboration worldwide



## Cosselatios fusuction ins latice QCD

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Task \#2: Propagate $\mathrm{O}(100)$ quark on the gluon backgrounds \& take algebraic combinations

- 100 propagator $\sim 1$ hour of simulation on $\mathrm{O}(5000)$ cores/few GPUS.
- Similar in spirit to data analysis of collision events.

- "Smaller" national, European calls.


## Corselarion function in latcice QCD

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Key point: Lattice is a real (Euclidean) space method So let's stay in real space!

## hyy firons real space

By simply taking Laplace transform:

$$
a_{\mu}^{H V P}=\int_{0}^{\infty} d Q^{2} K\left(Q^{2}\right) \hat{\Pi}\left(Q^{2}\right) \quad \rightarrow \quad \int_{0}^{\infty} d t \tilde{K}(t) C(t)
$$

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Integration kernel enhances long euclidean times: $\tilde{K}(t) \xrightarrow{t \rightarrow \infty} t^{2}$

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$$

Integration kernel enhances long euclidean times: $\tilde{K}(t) \xrightarrow{t \rightarrow \infty} t^{2}$


Integral converges
because $C(t)$ falls
exponentially in time

$$
C(t) \rightarrow e^{-A t}
$$

## fyp firons real space

By simply taking Laplace transform:

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a_{\mu}^{H V P}=\int_{0}^{\infty} d Q^{2} K\left(Q^{2}\right) \hat{\Pi}\left(Q^{2}\right) \quad \rightarrow \quad \int_{0}^{\infty} d t \tilde{K}(t) C(t)
$$

Integration kernel enhances long euclidean times: $\tilde{K}(t) \xrightarrow{t \rightarrow \infty} t^{2}$


Integral converges
because $C(t)$ falls exponentially in time
$C(t) \rightarrow e^{-A t}$


Issue
Exponentially large noise at large time $\frac{S}{N}(t) \rightarrow e^{-B t}$ [G.Parisi, 1984]
Djificule tasts：farcd work，special tools SUPERCOMPUTERS GOOD USAGE
＂La potenza è nulla senza il controllo＂


|  | $\begin{aligned} & \stackrel{\rightharpoonup}{w} \\ & \stackrel{y}{0} \\ & 0 \end{aligned}$ |  |  |
| :---: | :---: | :---: | :---: |
|  | $\begin{aligned} & \text { ou } \\ & \stackrel{y ⿸ ⿻ 一 丿 口}{1} \end{aligned}$ |  |  |
|  |  | NODE 1 |  |



MODERN ALGORITHMS \&
NEW METHODS


Adaptative solvers

DD- ААMG


Multigrid

Eigendeflation


MODERN ALGORITHMS \&
NEW METHODS


Adaptative solvers

DD- $\AA \mathrm{AMG}$


Multigrid

Eigendeflation



L-sitice 〔ejulis ys, clispersive resultos

Dispersive results incompatible
with the experiments


KNT 18
Jegerlehner 17
DHMZ 17
DHMZ 11
HLMNT 11

From 2020: Lattice results are compatible with the experiments!!!



DISPERSIVE PREDICTION


g-2 EXPERIMENT

 DISPERSIVE PREDICTION

## LATTICE PREDICTION



- NEW PHYSICS?
- EXPERIMENTAL ISSUES?
g-2 EXPERIMENT
 DISPERSIVE PREDICTION

LATTICE PREDICTION


## g-2 EXPERIMENT

PROBLEMS IN R-RATIO?

- NEW PHYSICS?
- EXPERIMENTAL ISSUES?


PROBLEMS IN LATTICE?

- BMW IS ALONE
- OTHERS ARE ARRIVING
 DISPERSIVE PREDICTION

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## g-2 EXPERIMENT

PROBLEMS IN R-RATIO?

- NEW PHYSICS?
- EXPERIMENTAL ISSUES?


PROBLEMS IN LATTICE?

- BMW IS ALONE
- OTHERS ARE ARRIVING


## Theneretical unclerstancling of R(E)



$$
R(E)=\frac{\sigma\left(e^{+} e^{-} \rightarrow \text { hadrons }\right)}{\sigma\left(e^{+} e^{-} \rightarrow \mu^{+} \mu^{-}\right)}
$$

A number of important features are qualitatively understood

## T'seoretical buclerstancling of R(E)



$$
R(E)=\frac{\sigma\left(e^{+} e^{-} \rightarrow \text { hadrons }\right)}{\sigma\left(e^{+} e^{-} \rightarrow \mu^{+} \mu^{-}\right)}
$$

A number of important features are qualitatively understood
$R(E) \sim 3$ at large Energies: color gauge theory $J / \Psi$ narrow peak: charm quark (GIM)

## T'seoretical busclerstancling of R(E)



## Thneoretical unclerstanding of $R(E)$



Hadrons are confined states of quarks interacting nonperturbatively
$R(E)$ Semi-quantitatively described only by models/effective theories

Disect sleterssijsacion of̈ is(E) firons V(t)

$$
\underbrace{\frac{\alpha_{e m}^{2}}{3 \pi^{2}} \int_{0}^{\infty} \frac{d E^{2}}{E^{2}} \tilde{K}(E) \boldsymbol{R}(\boldsymbol{E})}_{\text {disperisve, experimental }}=a_{\mu}^{H V P}=\underbrace{2 \alpha_{e m}^{2} \int_{0}^{\infty} d t t^{2} K\left(m_{\mu} t\right) \boldsymbol{V}(\boldsymbol{t})}_{\text {lattice, } S M}
$$



$$
\underbrace{\frac{\alpha_{e m}^{2}}{3 \pi^{2}} \int_{0}^{\infty} \frac{d E^{2}}{E^{2}} \tilde{K}(E) \boldsymbol{R}(\boldsymbol{E})}_{\text {disperisve, experimental }}=a_{\mu}^{H V P}=\underbrace{2 \alpha_{e m}^{2} \int_{0}^{\infty} d t t^{2} K\left(m_{\mu} t\right) \boldsymbol{V}(t)}_{\text {lattice, } S M}
$$

$R(E)$ is the inverse Laplace transform of $V(t)$

- To be computed in presence of noise \& finite sampling
- Notoriously, an ill-posed problem (needs regularization)


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$$

$R(E)$ is the inverse Laplace transform of $V(t)$

- To be computed in presence of noise \& finite sampling
- Notoriously, an ill-posed problem (needs regularization)
...and we are working hard on this!!!
"Probing the Energy-Smeared R-Ratio Using Lattice QCD" [PRL 130 (2023)] see A.De Santis firetalk, today @10:45 am


## Corrsplersentisy o!bservaibles: windows

$$
\underbrace{\frac{\alpha_{e m}^{2}}{3 \pi^{2}} \int_{0}^{\infty} \frac{d E^{2}}{E^{2}} \tilde{K}(E) \boldsymbol{R}(\boldsymbol{E})}_{\text {disperisve, experimental }} \underline{\tilde{\Theta}(E)}=a_{\mu}^{\Theta}=\underbrace{2 \alpha_{e m}^{2} \int_{0}^{\infty} d t t^{2} K\left(m_{\mu} t\right) \boldsymbol{V}(\boldsymbol{t})}_{\text {lattice, } S M} \underline{\Theta(t)}
$$

## Consplensentensy obseervaibles: winclows

$$
\underbrace{\frac{\alpha_{e m}^{2}}{3 \pi^{2}} \int_{0}^{\infty} \frac{d E^{2}}{E^{2}} \tilde{K}(E) \boldsymbol{R}(\boldsymbol{E})}_{\text {disperisve, experimental }} \underline{\tilde{\Theta}(E)}=a_{\mu}^{\Theta}=\underbrace{2 \alpha_{e m}^{2} \int_{0}^{\infty} d t t^{2} K\left(m_{\mu} t\right) \boldsymbol{V}(\boldsymbol{t})}_{\text {lattice, } S M} \underline{\Theta(t)}
$$

Modified version of HVP, more localized in energy
$\Theta^{S D}(t)+\Theta^{W}(t)+\Theta^{L D}(t)=1$



## Sinort e Internediate winclows

Short Distance = Large Energies (mostly perturbative)


Lattice compatible with Dispersive

## S'surt : Inserssuediate winadows

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Lattice compatible with Dispersive

Intermediate Distances ~ 1-2 GeV
(two pions final state)


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## Botion line

## Experimental status

- Muon anomalous magnetic is measured since 40 years
- Striking agreement within recent measurements
- Precision will improve in 2024 (Run IV at Fermilab g-2)



## Eoteoss lisse

## Experimental status

- Muon anomalous magnetic is measured since 40 years
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## Theoretical status

- All contributions but HVP are well understood
- Old g-2 puzzle: disagreement with experiments using R-ratio to "compute" HVP
- New g-2 puzzle: ~agreement when using HVP from recent lattice calculation

possitole exolutions to thne puzzles


## Assuming NO NEW PHYSICS

(the so called "everybody go home" scenario)

# posjible solutions to the puzzles 

## Assuming NO NEW PHYSICS

(the so called "everybody go home" scenario)

## Who is wrong? <br> Solves Puzzle 1? Solves Puzzle 2?

Mistake in lattice prediction of HVP
Mistake in g-2 experiments
Mistake in inclusive $\mathrm{e}^{+} \mathrm{e}^{-}$cross section

NO

YES

YES
YES

> Posjible solutions to the puzzles

Allowing for NEW PHYSICS and no experimental/lattice mistake

#  <br> <br> Allowing for NEW PHYSICS and no <br> <br> Allowing for NEW PHYSICS and no experimental/lattice mistake 

 experimental/lattice mistake}

## $\rightarrow$ Difficult to explain at the same time both puzzles

New physics behind the new muon $g$-2 puzzle?

Luca Di Luzio, ${ }^{1,2}$ Antonio Masiero,,${ }^{1,2}$ Paride Paradisi, ${ }^{1,2}$ and Massimo Passera ${ }^{2}$<br>${ }^{1}$ Dipartimento di Fisica e Astronomia 'G. Galilei', Università di Padova, Italy<br>${ }^{2}$ Istituto Nazionale di Fisica Nucleare, Sezione di Padova, Padova, Italy

The recent measurement of the muon $g-2$ at Fermilab confirms the previous Brookhaven result. The leading hadronic vacuum polarization (HVP) contribution to the muon $g$ - 2 represents a crucial ingredient to establish if the Standard Model prediction differs from the experimental value. A recent lattice QCD result by the BMW collaboration shows a tension with the low-energy $e^{+} e^{-} \rightarrow$ hadrons data which are currently used to determine the HVP contribution. We refer to this tension as the new muon $g$-2 puzzle. In this Letter we consider the possibility that new physics contributes to the $e^{+} e^{-} \rightarrow$ hadrons cross-section. This scenario could, in principle, solve the new muon $g$ - 2 puzzle. However, we show that this solution is excluded by a number of experimental constraints.

# Posciolole solutions to the puzzles 

## Allowing for NEW PHYSICS and no experimental/lattice mistake

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...but more complicated scenarios are not ruled out

## In fine f゙ucuse

## Experimental side

- More precise data from g-2 experiment
- Reanalysis of old $\mathrm{e}^{+} e$ experiment KLEO in progress
- Additional measurements from ongoing e+e- experiments
- New experiments MuOnE in the near future


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- More lattice collaborations will compute g-2
- Greater accuracy (better infinite volume \& continuum limits)
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