



# Probing New Physics through Entanglement in Diboson Production

Rafael Aoude, E.M., Fabio Maltoni, Luca Mantani – arXiv:2307.09675 [hep-ph]

Eric Madge

GGI – November 8, 2023

# Introduction

so far in the workshop ...

## top pairs

$$h \rightarrow ZZ^*, WW^*$$

$$WW, ZZ, WZ$$

- Standard Model ✓  
see talks by  
J.R.M. de Nova,  
J.W. Howarth and  
D. Gonçalves
- Standard Model ✓  
see talks by A. Barr,  
T. Maurin and J. Moreno
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see talk by A. Barr  
(also R.A. Morales)

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- Standard Model ✓  
see talk by A. Barr  
(also R.A. Morales)

- BSM ✓  
see talks by L. Mantani  
and C. Severi

- BSM ✓  
see talks by A. Bernal  
and L. Marzola

- BSM ✗  
⇒ this talk!

also: beyond QM (M. Eckstein)

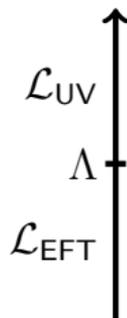
# Standard Model Effective Field Theory

○ **new physics** at high scale  $\Lambda$

→ at low energies: EFT with SM fields only

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \geq 5} \sum_n \frac{C_n \mathcal{O}_n^{(d)}}{\Lambda^{d-4}}$$

⇒ model independent



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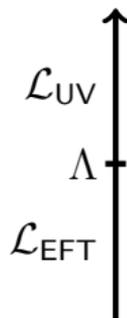
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dim- $d$  operator composed  
of SM fields (respecting  
SM symmetries)

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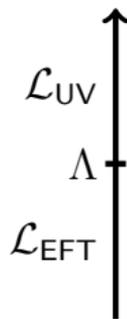
$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \geq 5} \sum_n \frac{C_n \mathcal{O}_n^{(d)}}{\Lambda^{d-4}}$$

Wilson coefficient  $C_n$  (indicated by a red arrow)

$\mathcal{O}_n^{(d)}$  (indicated by a grey arrow)

dim- $d$  operator composed of SM fields (respecting SM symmetries)

⇒ model independent



# Standard Model Effective Field Theory

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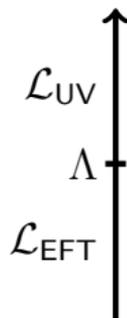
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$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \geq 5} \sum_n \frac{C_n \mathcal{O}_n^{(d)}}{\Lambda^{d-4}}$$

Wilson coefficient  $C_n$       dim- $d$  operator composed of SM fields (respecting SM symmetries)  $\mathcal{O}_n^{(d)}$

NP scale  $\Lambda$

⇒ model independent



# Standard Model Effective Field Theory

○ **new physics** at high scale  $\Lambda$

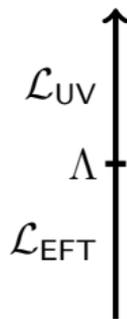
→ at low energies: EFT with SM fields only

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \geq 5} \sum_n \frac{C_n \mathcal{O}_n^{(d)}}{\Lambda^{d-4}} \quad \rightarrow \quad \mathcal{L}_{\text{SM}} + \sum_n c_n \mathcal{O}_n^{(6)} + \dots$$

Annotations:

- Wilson coefficient  $C_n$  (points to  $C_n$ )
- NP scale  $\Lambda$  (points to  $\Lambda^{d-4}$ )
- dim- $d$  operator composed of SM fields (respecting SM symmetries)  $\mathcal{O}_n^{(d)}$  (points to  $\mathcal{O}_n^{(d)}$ )

⇒ model independent



# Standard Model Effective Field Theory

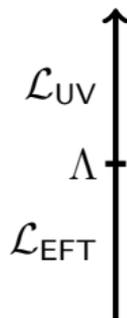
○ new physics at high scale  $\Lambda$

→ at low energies: EFT with SM fields only

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \geq 5} \sum_n \frac{C_n \mathcal{O}_n^{(d)}}{\Lambda^{d-4}} \quad \xrightarrow{\text{dim-}d \text{ operator composed of SM fields (respecting SM symmetries)}} \quad \mathcal{L}_{\text{SM}} + \sum_n c_n \mathcal{O}_n^{(6)} + \dots$$

Annotations:  
- "Wilson coefficient" points to  $C_n$   
- "NP scale" points to  $\Lambda$   
- "dim- $d$  operator composed of SM fields (respecting SM symmetries)" points to  $\mathcal{O}_n^{(d)}$   
- Red arrow points to  $c_n = \frac{C_n}{\Lambda^2}$

⇒ model independent



# Standard Model Effective Field Theory

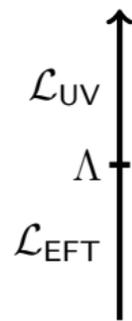
- **new physics** at high scale  $\Lambda$

→ at low energies: EFT with SM fields only

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \geq 5} \sum_n \frac{C_n \mathcal{O}_n^{(d)}}{\Lambda^{d-4}} \quad \rightarrow \quad \mathcal{L}_{\text{SM}} + \sum_n c_n \mathcal{O}_n^{(6)} + \dots$$

Wilson coefficient  $\swarrow$   $C_n$   $\searrow$   $\mathcal{O}_n^{(d)}$   
 NP scale  $\swarrow$   $\Lambda^{d-4}$   $\searrow$   $\Lambda$   
 dim- $d$  operator composed of SM fields (respecting SM symmetries)

$\uparrow$   
 $c_i = \frac{C_i}{\Lambda^2}$



⇒ model independent

- **Warsaw basis:** (flavor universal)

[Grzadkowski, Iskrzyński, Misiak, Rosiek (JHEP 2010)]

- 1 dim-5 operator
- 59 (non- $\mathcal{B}$ ) + 4 ( $\mathcal{B}$ ) dim-6 operators

- **here:** 13 relevant CP-even dim-6 flavor-universal operators

# Dimension-6 Operators

Definition		95 % CL	Definition		95 % CL
two-fermion operators			bosonic operators		
$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$	$[-0.17, 0.14]$	$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$	$[-0.18, 0.22]$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$	$[-0.07, 0.09]$	$c_{\varphi W}$	$\left(\varphi^\dagger \varphi - \frac{v^2}{2}\right) W_I^{\mu\nu} W_{\mu\nu}^I$	$[-0.15, 0.30]$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$	$[-0.06, 0.22]$	$c_{\varphi B}$	$\left(\varphi^\dagger \varphi - \frac{v^2}{2}\right) B_{\mu\nu} B^{\mu\nu}$	$[-0.11, 0.11]$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$	$[-0.21, 0.05]$	$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$	$[-0.17, 0.27]$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$	$[-0.21, 0.26]$	$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$	$[-0.52, 0.43]$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$	$[-0.11, 0.13]$	four-fermion operator		
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$	$[-0.21, 0.05]$	$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$	$[-0.16, 0.02]$

95 % CL bounds in  $\text{TeV}^{-2}$  from [SMFIT (JHEP 2021)]

# Coupling Modification

- universal shifts in  $G_f$ ,  $\sin \theta_W$  and  $g_Z$

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$
$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

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$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
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$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$
$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

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$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

# Coupling Modification

- universal shifts in  $G_f$ ,  $\sin\theta_W$  and  $g_Z$
- direct modification of **quark couplings to  $Z$**

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{u} \gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{d} \gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{q} \gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi) (\bar{q} \gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{e} \gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{l} \gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi) (\bar{l} \gamma^\mu \tau^I l)$
$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l} \gamma_\mu l) (\bar{l} \gamma^\mu l)$

# Coupling Modification

- universal shifts in  $G_f$ ,  $\sin\theta_W$  and  $g_Z$
- direct modification of quark couplings to  $Z$
- direct modification of (LH) quark couplings to  $W$

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$
$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

# Coupling Modification

- universal shifts in  $G_f$ ,  $\sin \theta_W$  and  $g_Z$
- direct modification of quark couplings to  $Z$
- direct modification of (LH) quark couplings to  $W$
- direct modification of **electron couplings to  $Z$**

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
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$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

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- universal shifts in  $G_f$ ,  $\sin \theta_W$  and  $g_Z$
- direct modification of quark couplings to  $Z$
- direct modification of (LH) quark couplings to  $W$
- direct modification of electron couplings to  $Z$
- direct modification of (LH) electron couplings to  $W$

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
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$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
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# Coupling Modification

- universal shifts in  $G_f$ ,  $\sin \theta_W$  and  $g_Z$
- direct modification of quark couplings to  $Z$
- direct modification of (LH) quark couplings to  $W$
- direct modification of electron couplings to  $Z$
- direct modification of (LH) electron couplings to  $W$
- direct modification of **triple gauge couplings**

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$
$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

# Coupling Modification

- universal shifts in  $G_f$ ,  $\sin \theta_W$  and  $g_Z$
- direct modification of quark couplings to  $Z$
- direct modification of (LH) quark couplings to  $W$
- direct modification of electron couplings to  $Z$
- direct modification of (LH) electron couplings to  $W$
- direct modification of triple gauge couplings  
(including new Lorentz structures)

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$
$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

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- universal shifts in  $G_f$ ,  $\sin \theta_W$  and  $g_Z$
- direct modification of quark couplings to  $Z$
- direct modification of (LH) quark couplings to  $W$
- direct modification of electron couplings to  $Z$
- direct modification of (LH) electron couplings to  $W$
- direct modification of triple gauge couplings (including new Lorentz structures)
- direct modification of **Higgs couplings**

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$
$c_W$	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi B}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) B_{\mu\nu} B^{\mu\nu}$
$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$c_{ll}$	$(\bar{l}\gamma_\mu l)(\bar{l}\gamma^\mu l)$

# Coupling Modification

- universal shifts in  $G_f$ ,  $\sin \theta_W$  and  $g_Z$
- direct modification of quark couplings to  $Z$
- direct modification of (LH) quark couplings to  $W$
- direct modification of electron couplings to  $Z$
- direct modification of (LH) electron couplings to  $W$
- direct modification of triple gauge couplings (including new Lorentz structures)
- direct modification of Higgs couplings

$c_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}\gamma^\mu u)$
$c_{\varphi d}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}\gamma^\mu d)$
$c_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}\gamma^\mu q)$
$c_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}\gamma^\mu \tau^I q)$
$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e}\gamma^\mu e)$
$c_{\varphi l}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}\gamma^\mu l)$
$c_{\varphi l}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}\gamma^\mu \tau^I l)$
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$c_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$
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$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$
$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
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# SMEFT corrections

- Calculation at LO, including dim-6 and (dim-6)<sup>2</sup> SMEFT corrections

$$R = \sum_{\text{IS}} L_{\text{IS}}(\hat{s}) \overline{\sum}_{\substack{\text{DOFs (excl.} \\ \text{FS spin)}}} \mathcal{M}^* \mathcal{M}$$

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- expand  $R$ -matrix to order  $\mathcal{O}(\Lambda^{-4})$

$$\rho = \frac{R}{\text{tr } R} = \frac{1}{9} \left\{ \frac{R_{\text{SM}}}{\tilde{A}_{\text{SM}}} + \sum_n \frac{C_n}{\Lambda^2} \left[ \frac{R_n}{\tilde{A}_{\text{SM}}} - \frac{\tilde{A}_n R_{\text{SM}}}{\tilde{A}_{\text{SM}}^2} \right] + \sum_{n,m} \frac{C_n C_m}{\Lambda^4} \left[ \frac{R_{nm}}{\tilde{A}_{\text{SM}}} - \frac{\tilde{A}_n R_m + \tilde{A}_{nm} R_{\text{SM}}}{\tilde{A}_{\text{SM}}^2} + \frac{\tilde{A}_n \tilde{A}_m R_{\text{SM}}}{\tilde{A}_{\text{SM}}^3} \right] \right\}$$

- similar: concurrence, purity, Bell violation, ...

$$\mathcal{C}(\rho) = \inf_{\{|\Psi\rangle\}} \left[ \sum_i p_i \mathcal{C}(|\Psi_i\rangle) \right], \quad \mathcal{C}(|\Psi\rangle) = \sqrt{2 \left( 1 - \text{tr}_A \left[ \left( \text{tr}_B |\Psi\rangle\langle\Psi| \right)^2 \right] \right)}$$

○  $0 \leq \mathcal{C}(\rho) \leq \frac{2}{\sqrt{3}}, \quad \mathcal{C}(\rho) > 0 \implies \text{entangled}$

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○ for qutrits: analytically calculable only for pure states

$\implies$  provide lower and upper bound:

$$\rho_{A/B} = \text{tr}_{B/A} \rho$$

$$(\mathcal{C}(\rho))^2 \geq (\mathcal{C}_{\text{LB}}(\rho))^2 = 2 \max \left[ \text{tr} \rho^2 - \text{tr} \rho_A^2, \text{tr} \rho^2 - \text{tr} \rho_B^2 \right]$$

$$(\mathcal{C}(\rho))^2 \leq (\mathcal{C}_{\text{UB}}(\rho))^2 = 2 \min \left[ 1 - \text{tr} \rho_A^2, 1 - \text{tr} \rho_B^2 \right]$$

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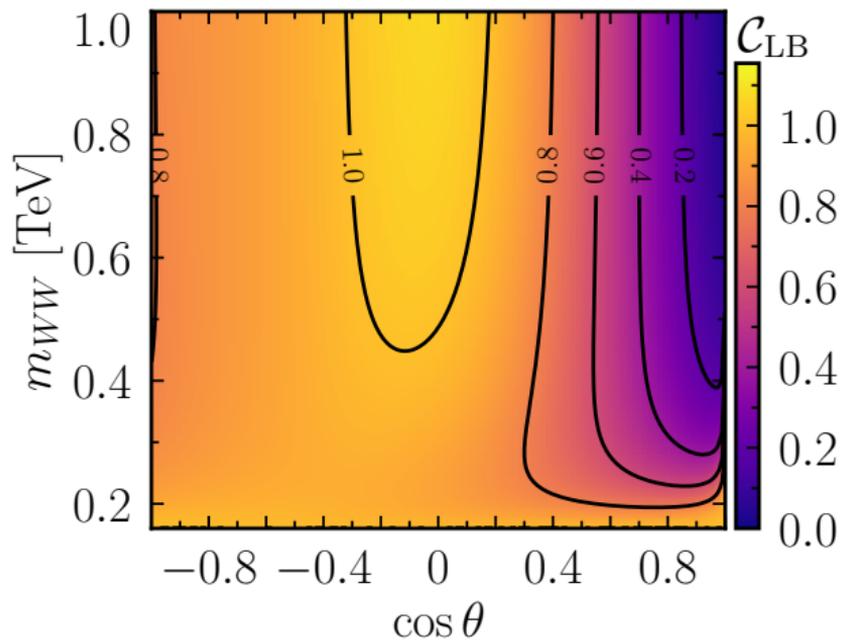
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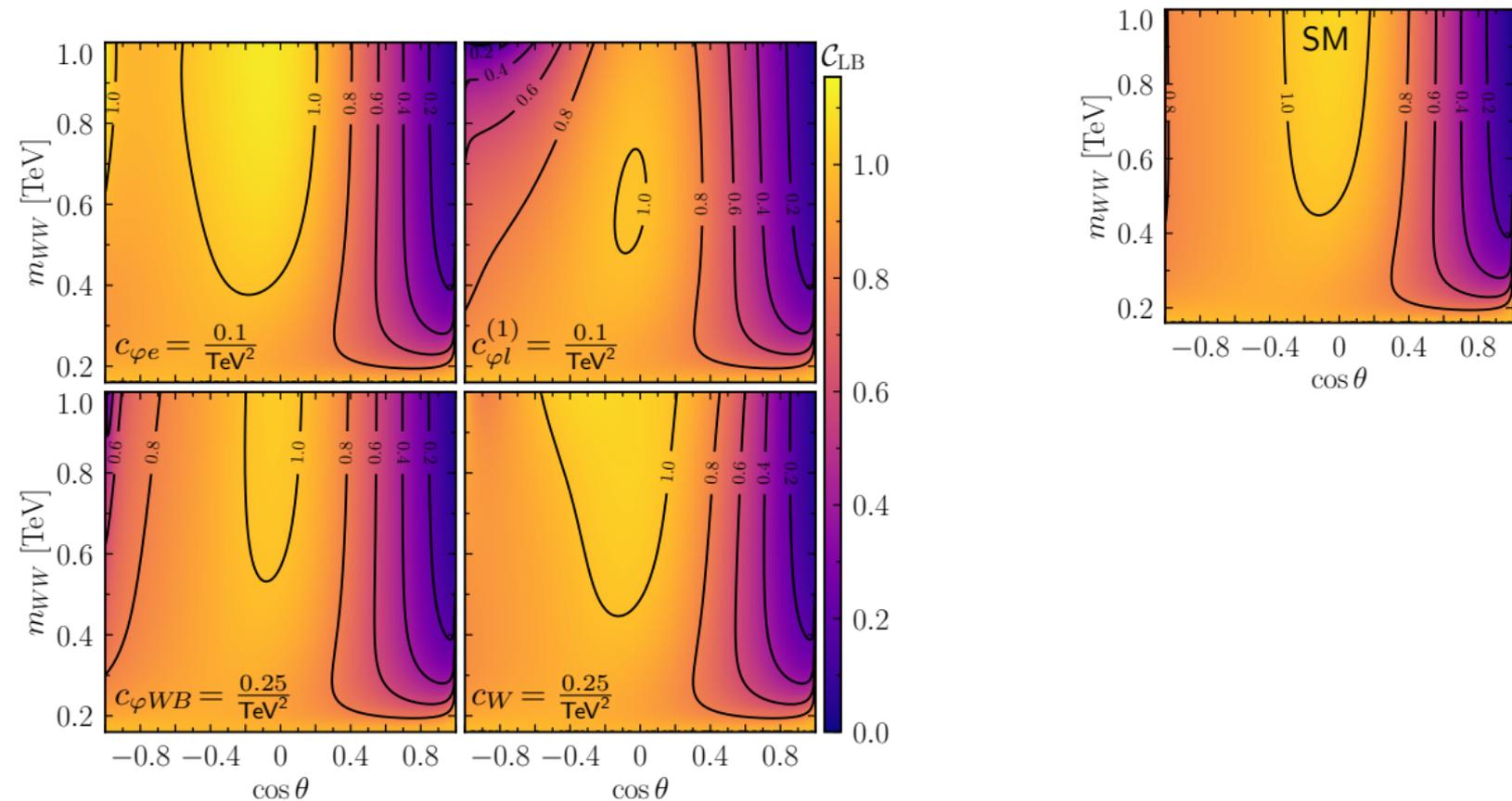
$$(\mathcal{C}(\rho))^2 \leq (\mathcal{C}_{\text{UB}}(\rho))^2 = 2 \min \left[ 1 - \text{tr} \rho_A^2, 1 - \text{tr} \rho_B^2 \right]$$

○ for pure state:  $P(\rho) = \text{tr} \rho^2 = 1 \implies \mathcal{C}_{\text{LB}}(\rho) = \mathcal{C}(\rho) = \mathcal{C}_{\text{UB}}(\rho)$

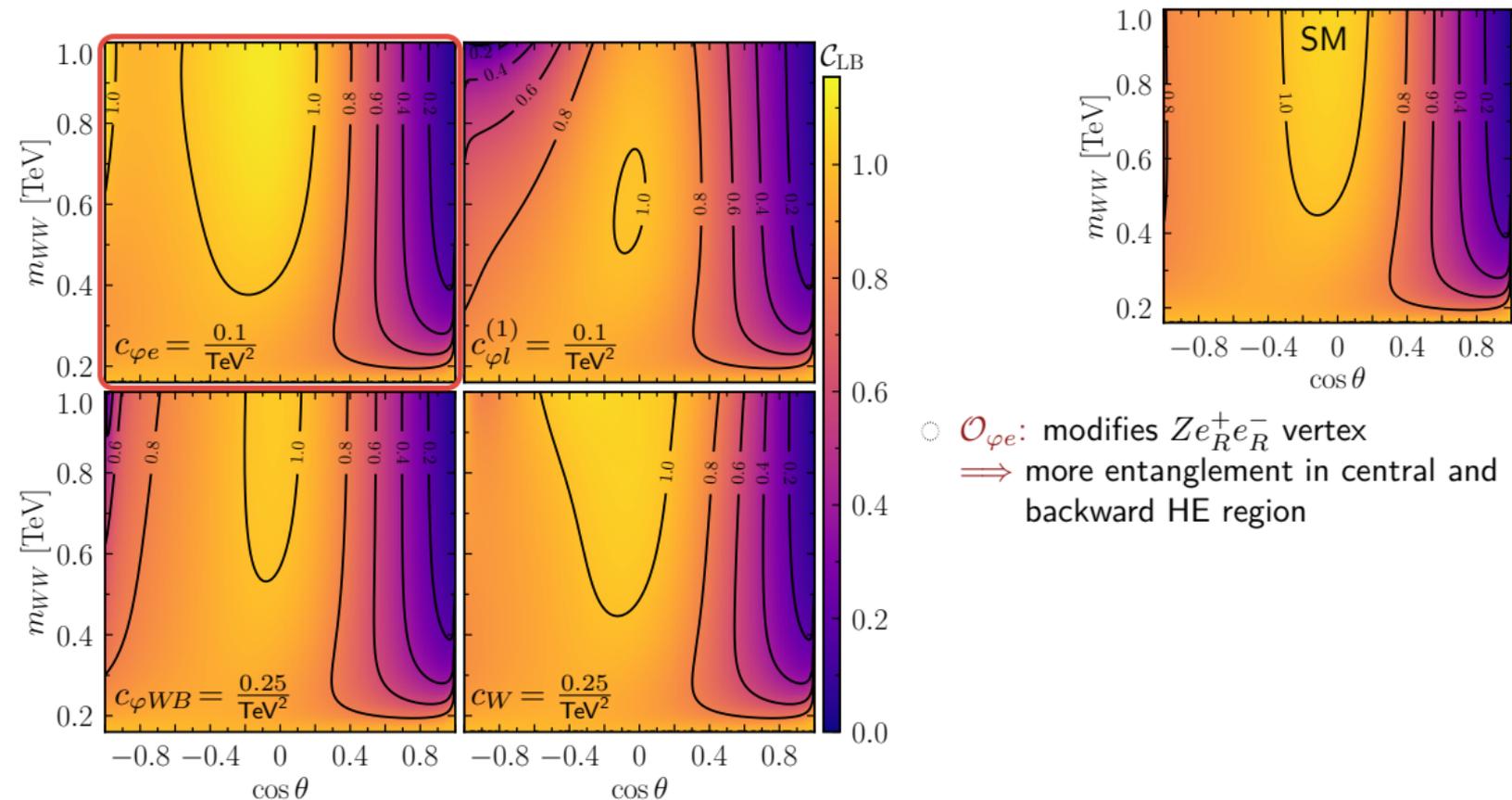
# SMEFT effects in $e^+e^- \rightarrow W^+W^-$



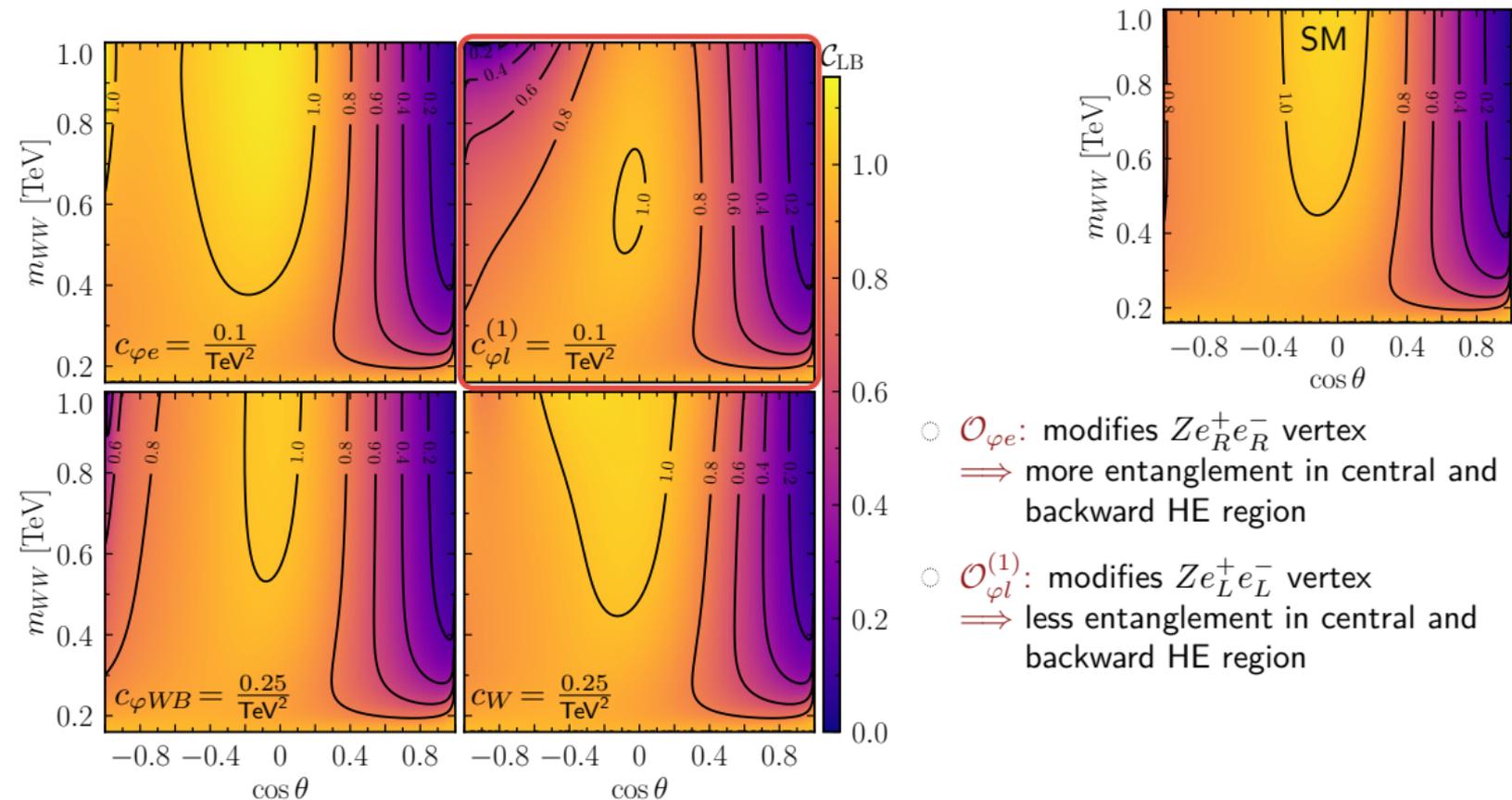
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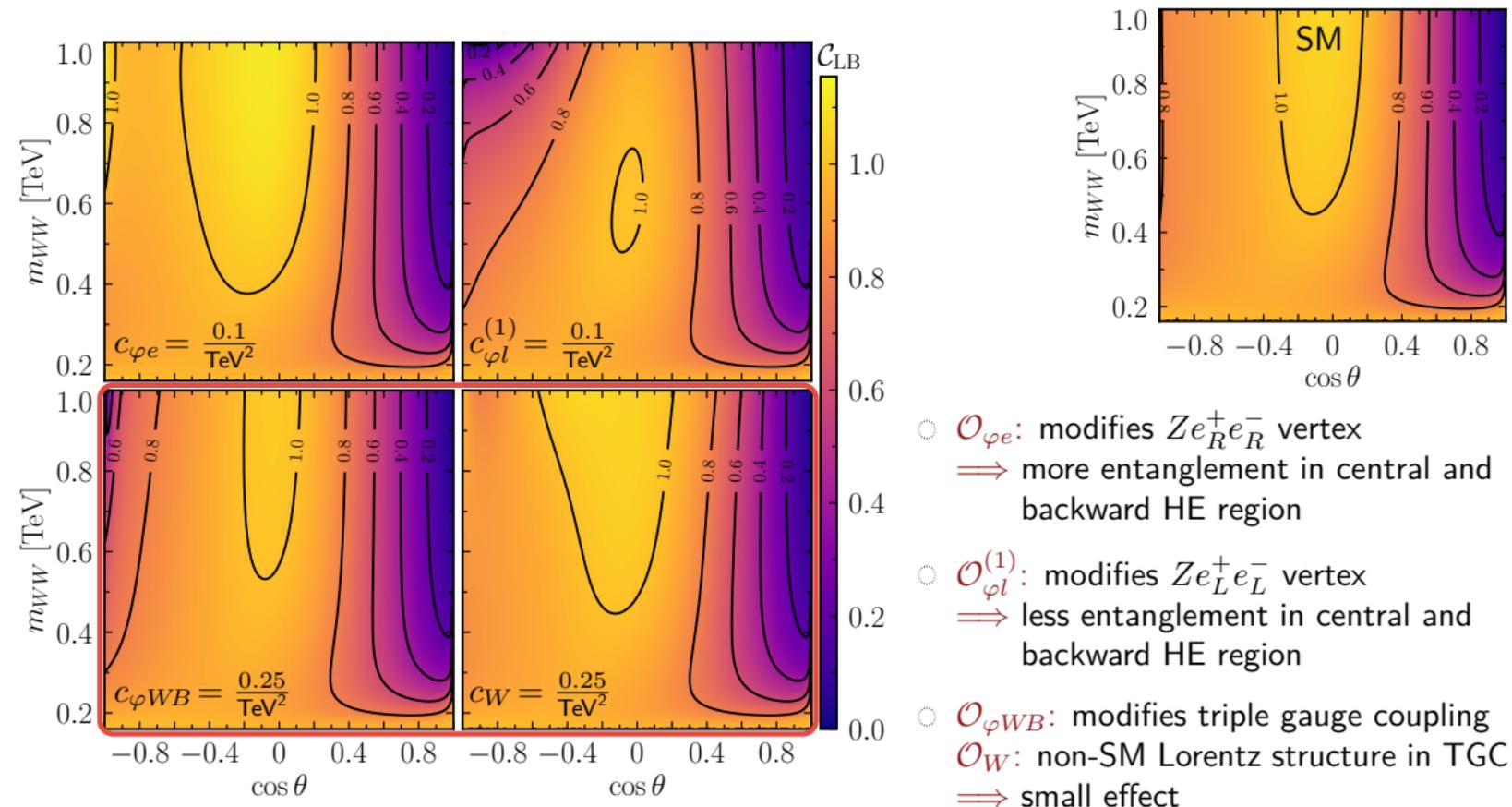
# SMEFT effects in $e^+e^- \rightarrow W^+W^-$



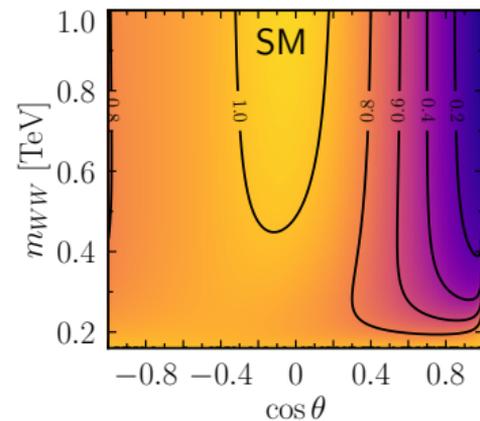
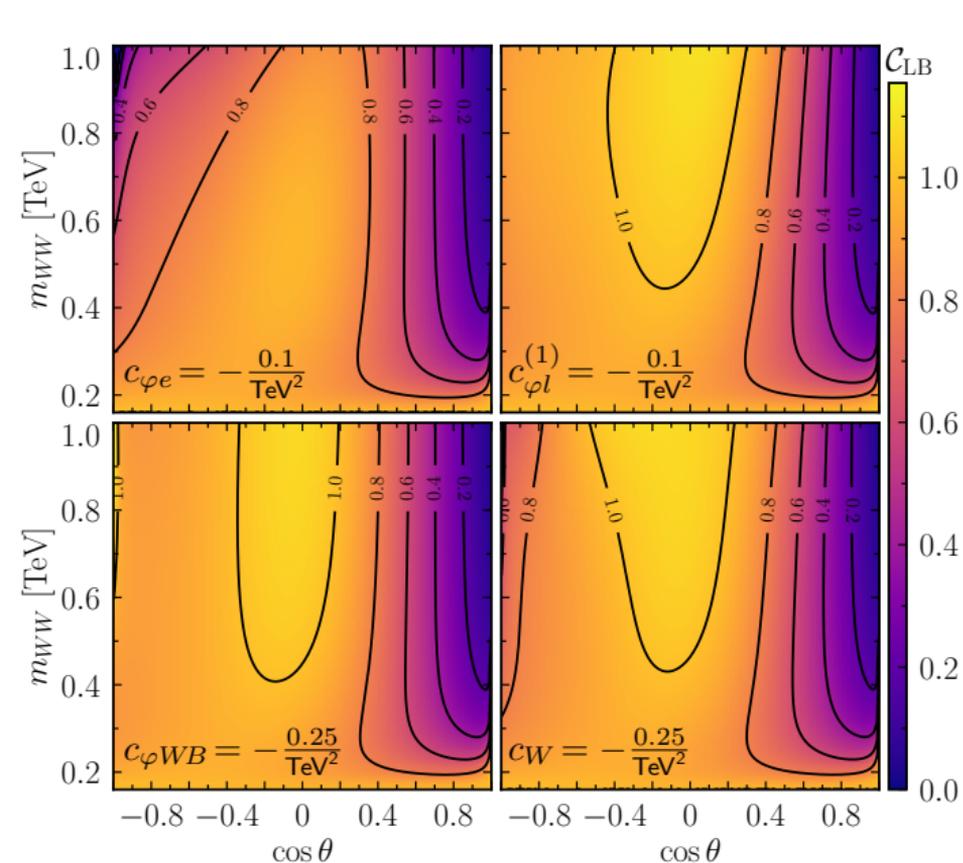
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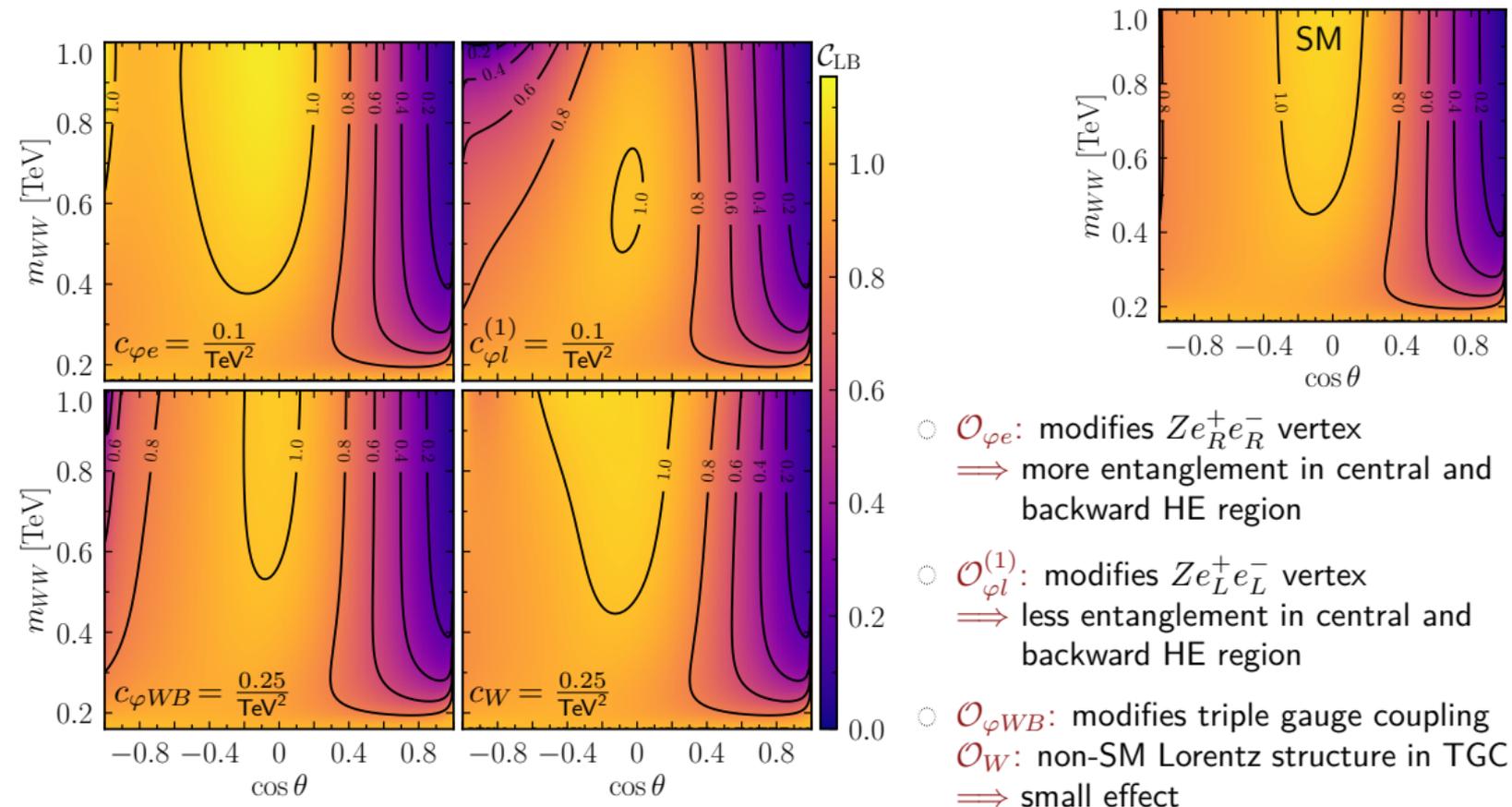


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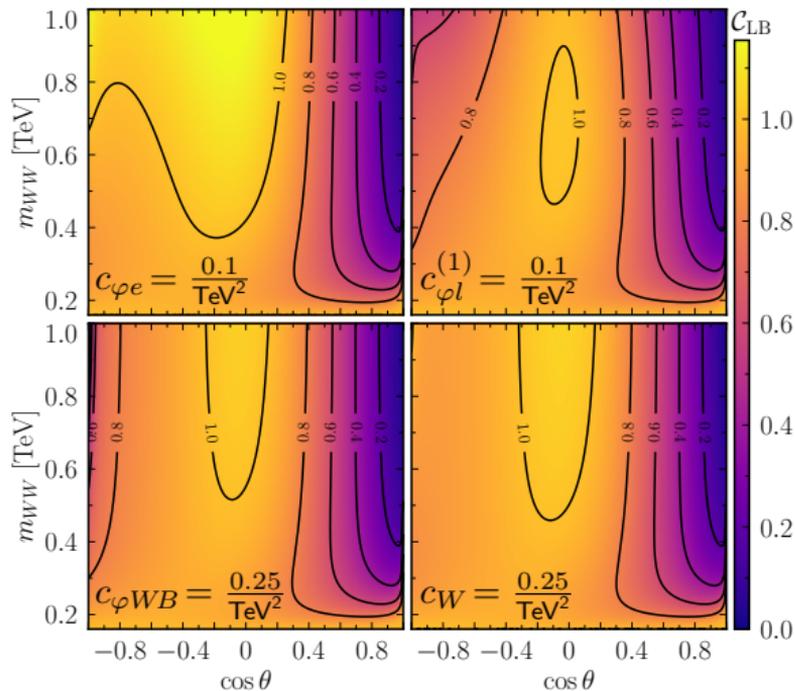
- $\odot \mathcal{O}_{\varphi e}$ : modifies  $Ze_R^+e_R^-$  vertex  
 $\Rightarrow$  more entanglement in central and backward HE region
- $\odot \mathcal{O}_{\varphi l}^{(1)}$ : modifies  $Ze_R^+e_R^-$  vertex  
 $\Rightarrow$  less entanglement in central and backward HE region
- $\odot \mathcal{O}_{\varphi WB}$ : modifies triple gauge coupling  
 $\Rightarrow$  opposite effect for opposite sign
- $\odot \mathcal{O}_W$ : non-SM Lorentz structure in TGC  
 $\Rightarrow$  small effect

# SMEFT effects in $e^+e^- \rightarrow W^+W^-$

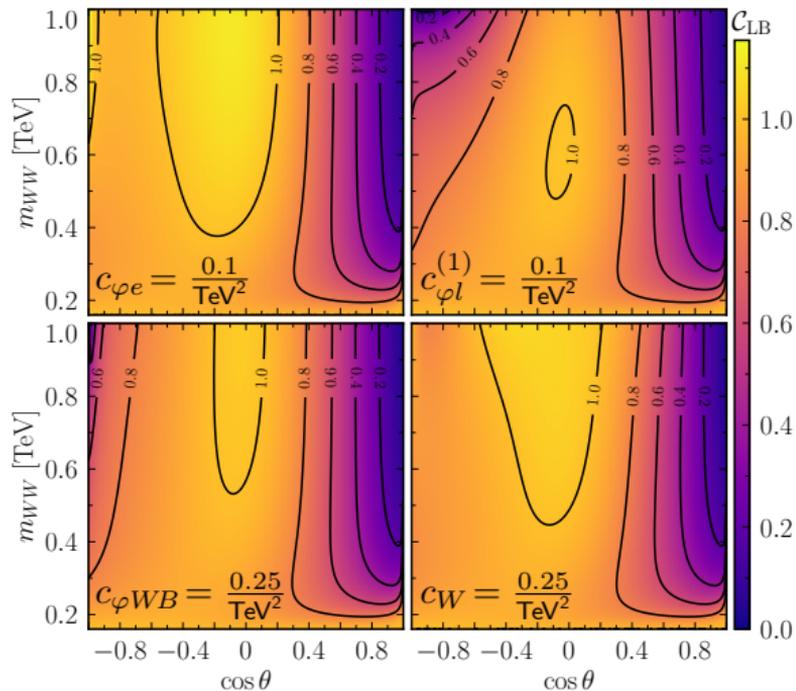


# SMEFT effects in $e^+e^- \rightarrow W^+W^-$ at linear order

SM + dim-6



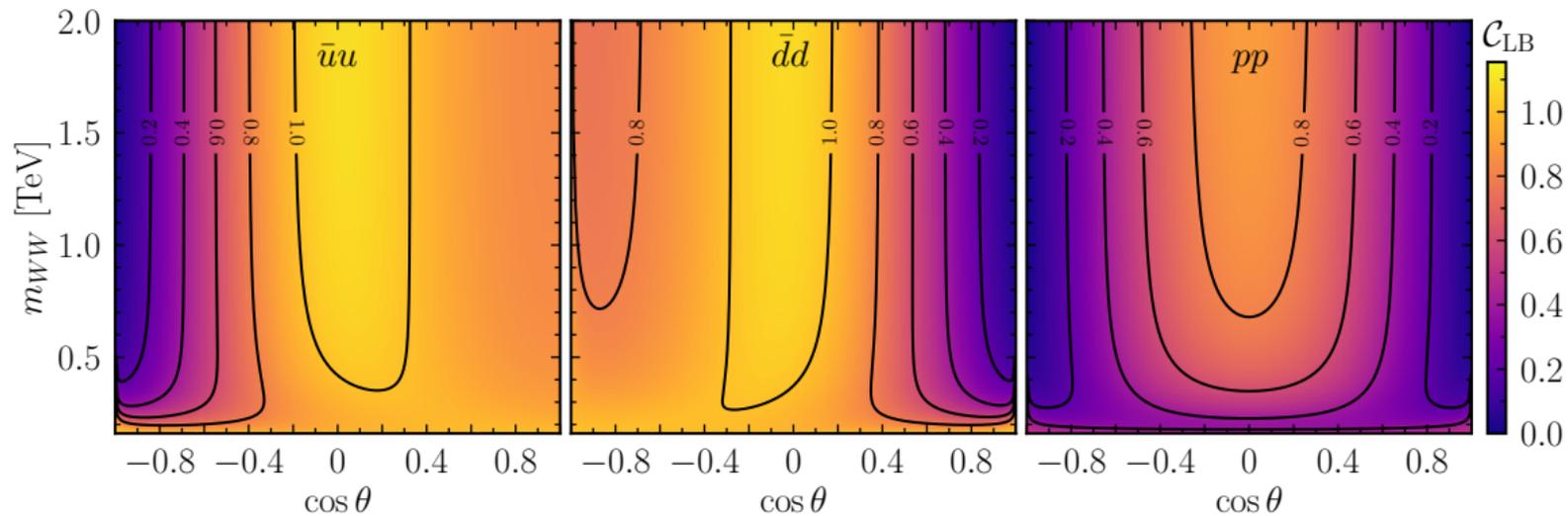
SM + dim-6 + (dim-6)<sup>2</sup>



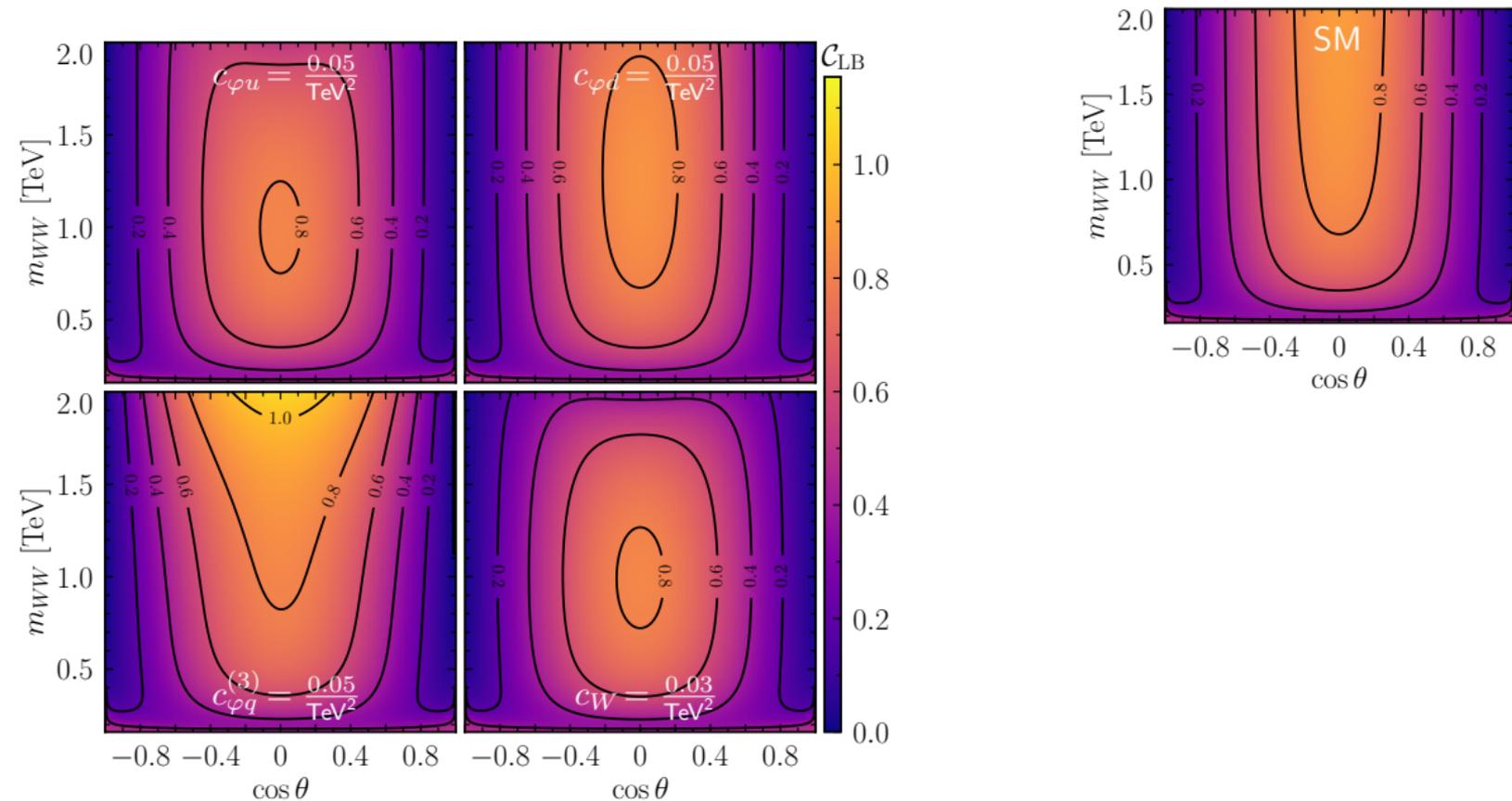




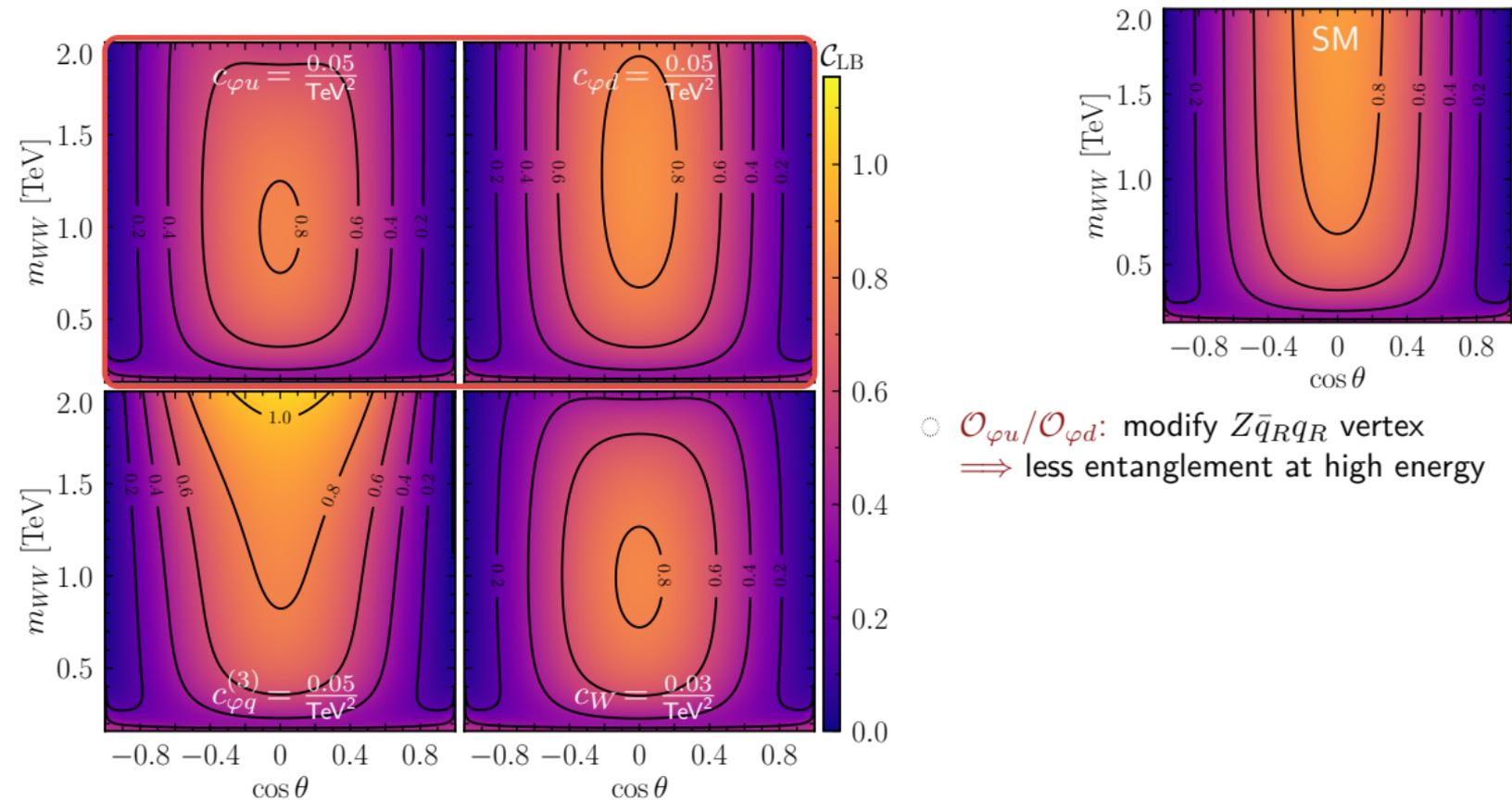
# SMEFT effects in $pp \rightarrow W^+W^-$



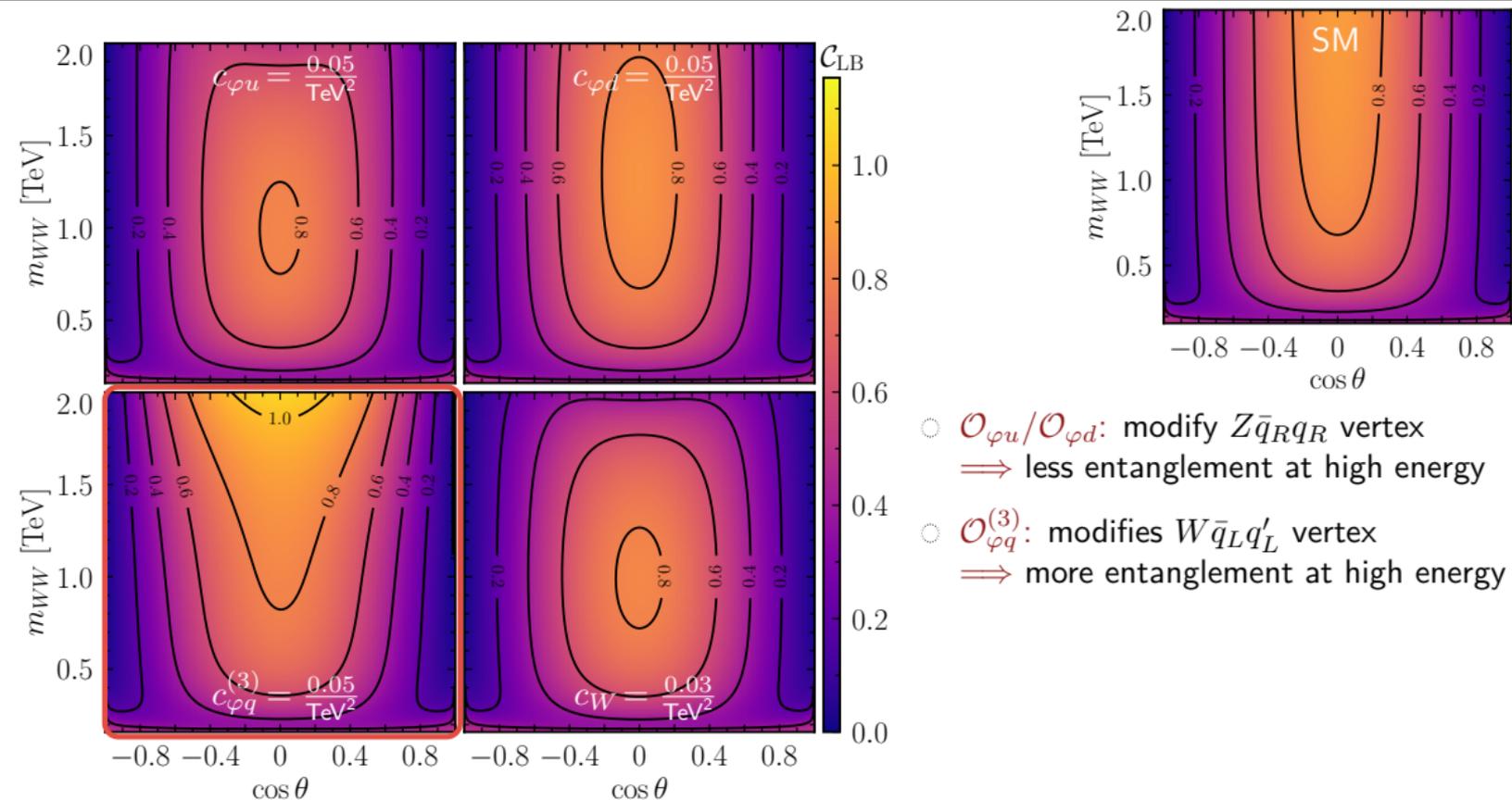
# SMEFT effects in $pp \rightarrow W^+W^-$



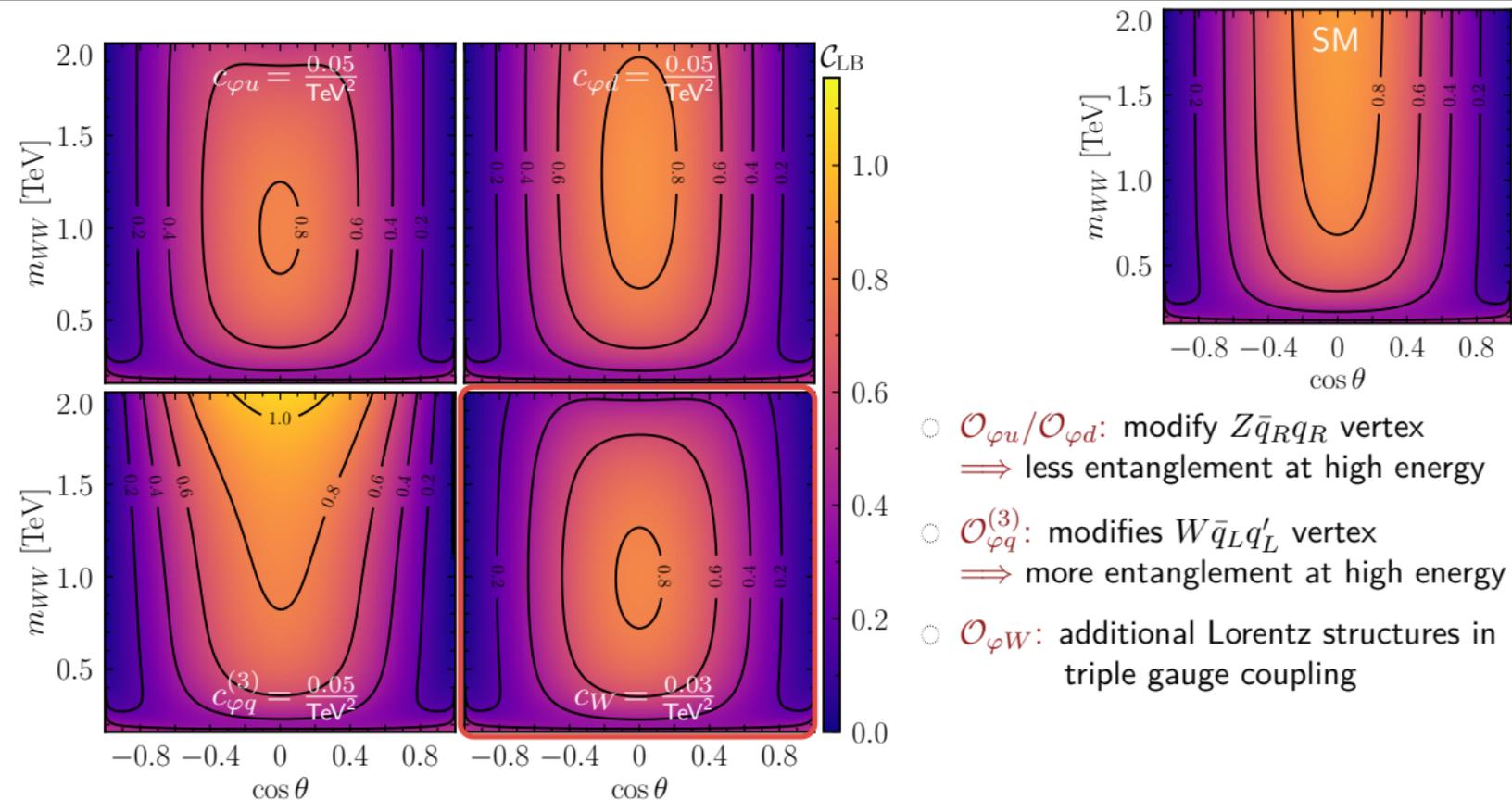
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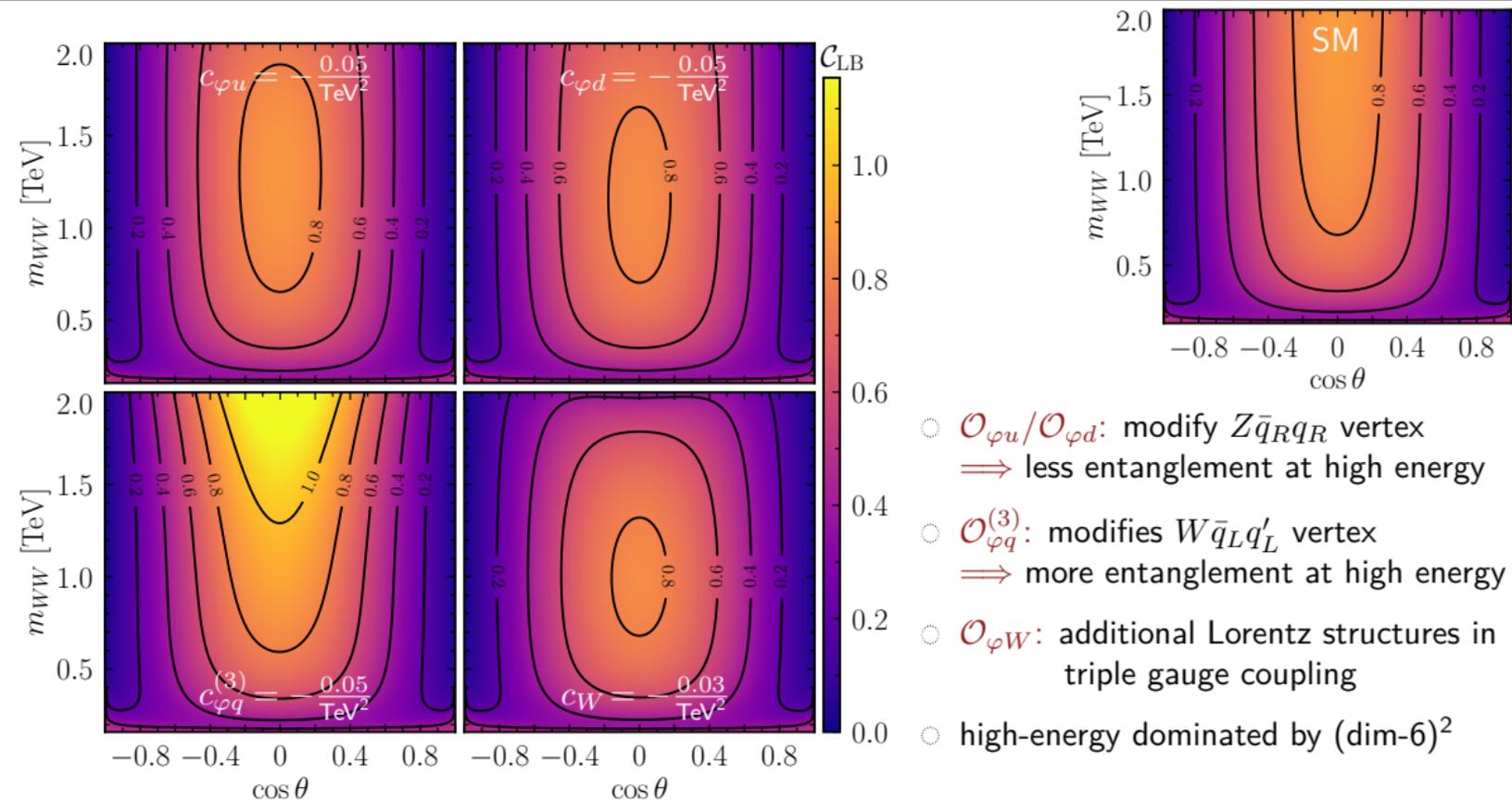


# SMEFT effects in $pp \rightarrow W^+W^-$

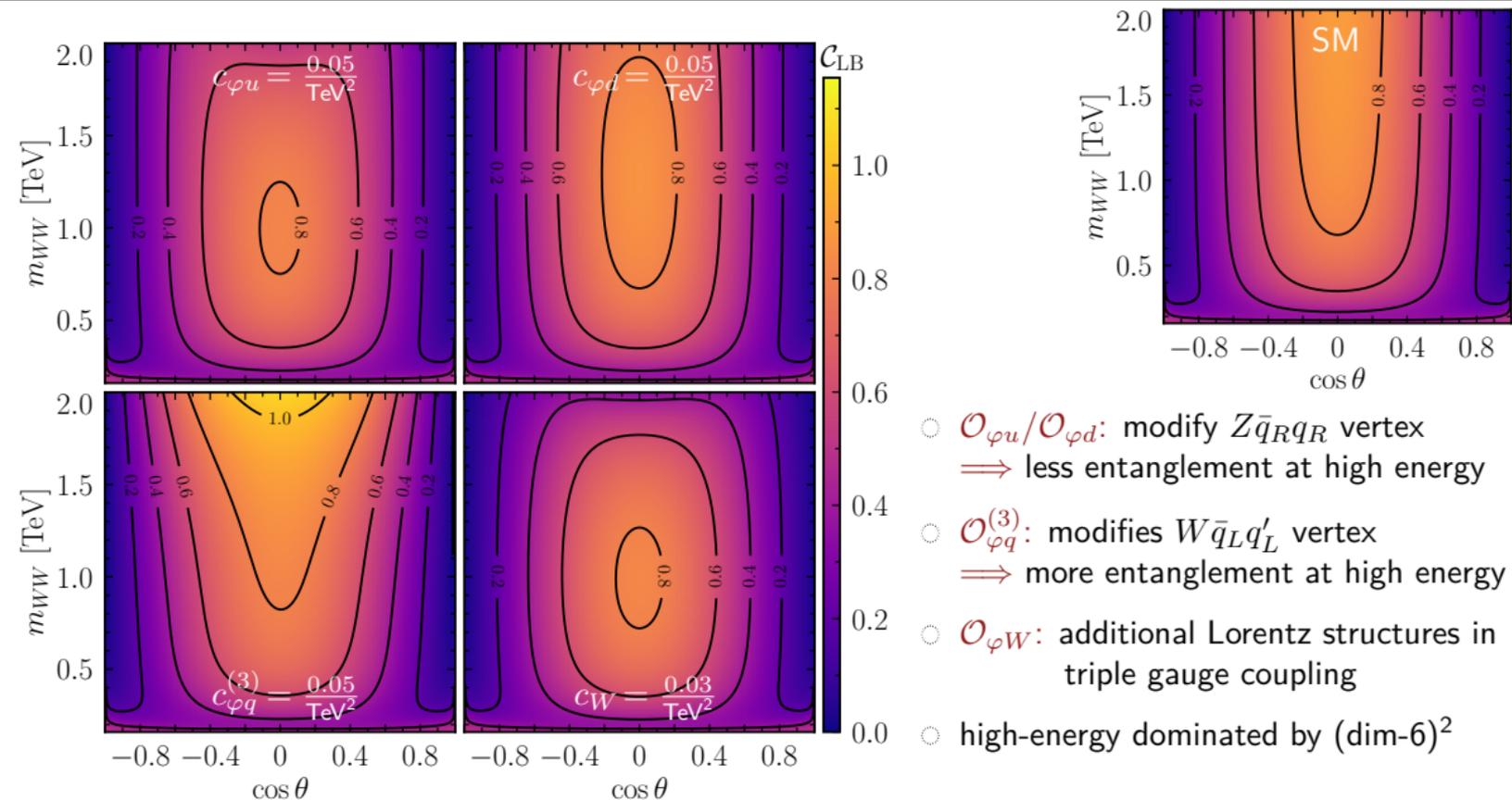


- $\odot$   $\mathcal{O}_{\varphi u}/\mathcal{O}_{\varphi d}$ : modify  $Z\bar{q}_R q_R$  vertex  
 $\Rightarrow$  less entanglement at high energy
- $\odot$   $\mathcal{O}_{\varphi q}^{(3)}$ : modifies  $W\bar{q}_L q'_L$  vertex  
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- $\odot$   $\mathcal{O}_{\varphi W}$ : additional Lorentz structures in triple gauge coupling

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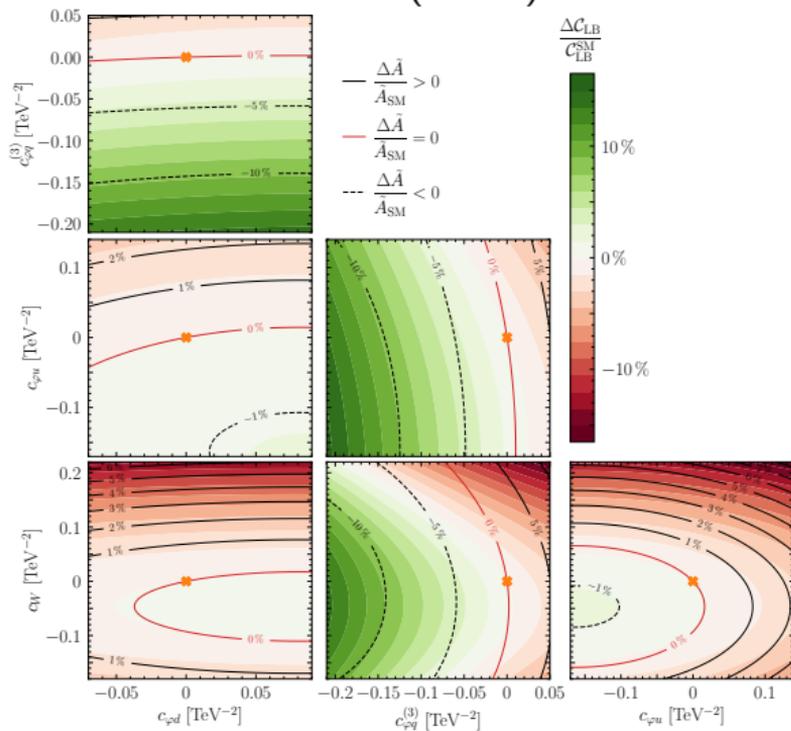


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- $\mathcal{O}_{\varphi W}$ : additional Lorentz structures in triple gauge coupling
- high-energy dominated by (dim-6)<sup>2</sup>

# Central High-Energy Region ( $pp \rightarrow W^+W^-$ )

$$m_{WW} = 500 \text{ GeV}, \cos \theta = 0$$

dim-6 + (dim-6)<sup>2</sup>

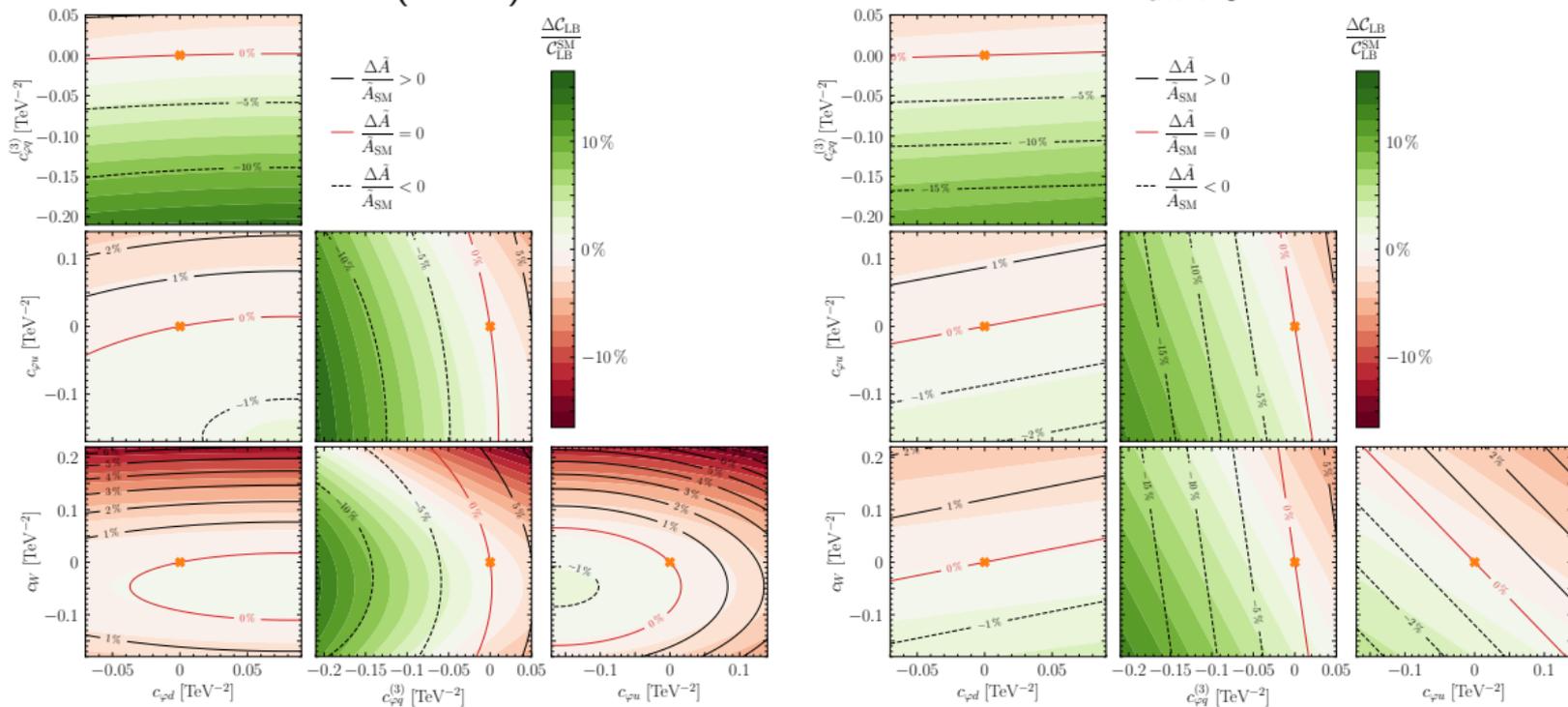


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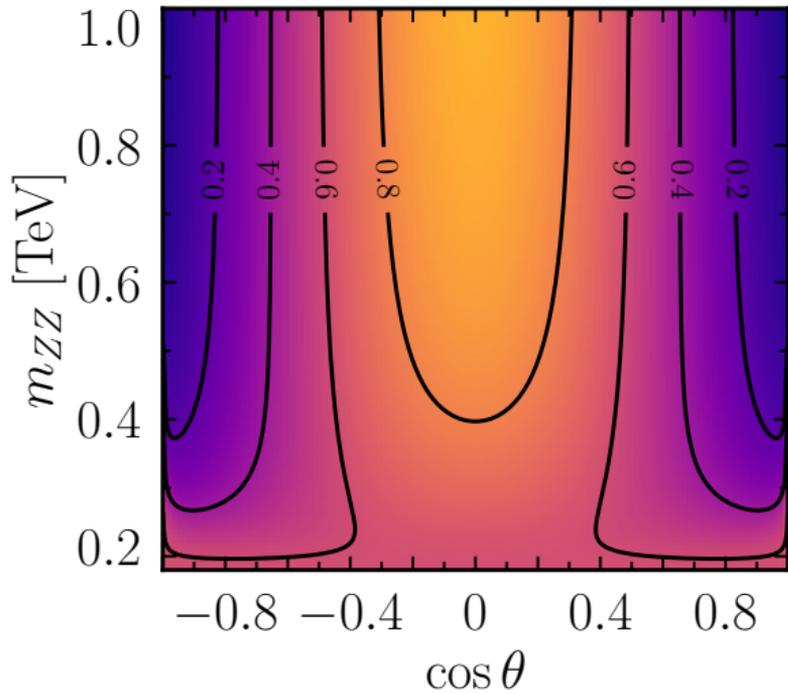
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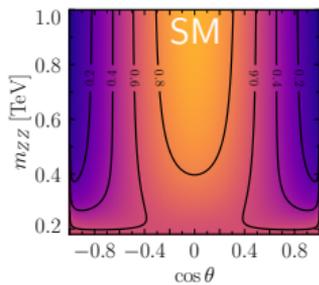
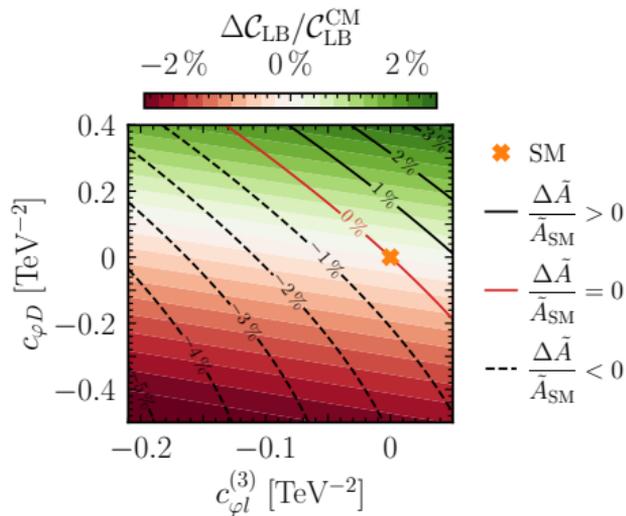
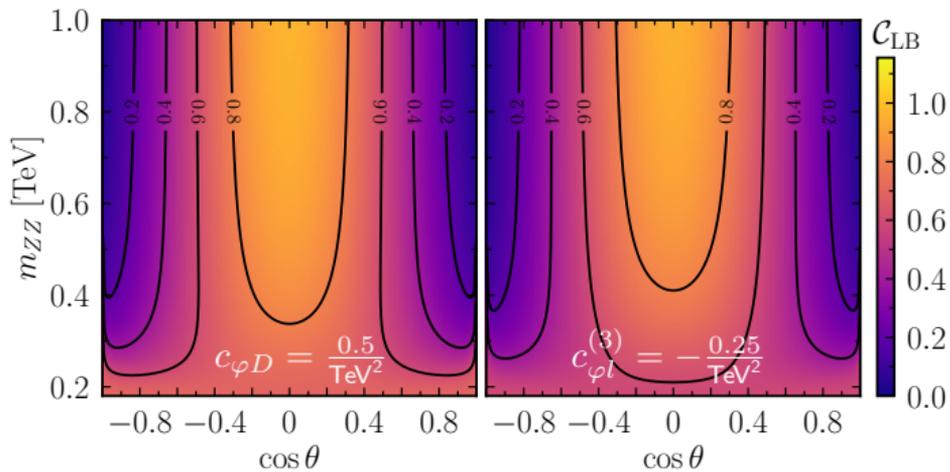
dim-6



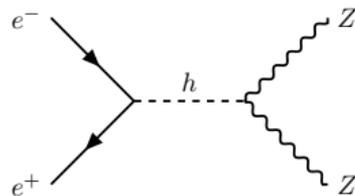
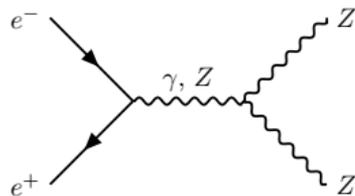
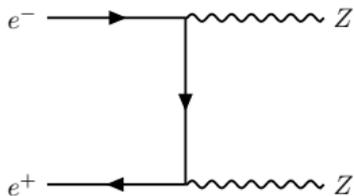
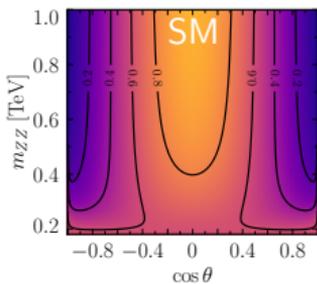
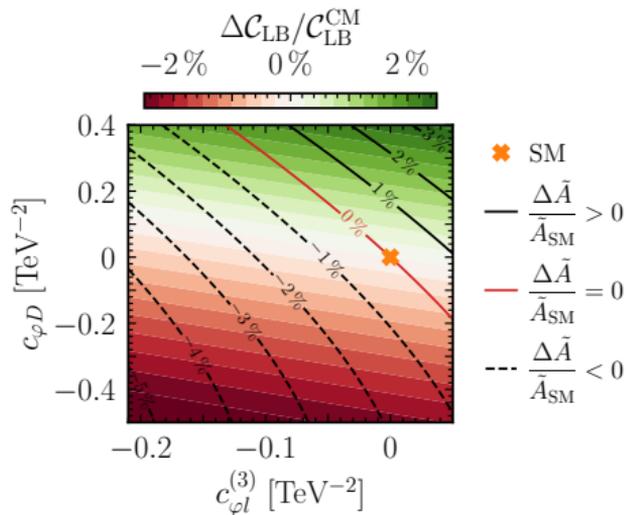
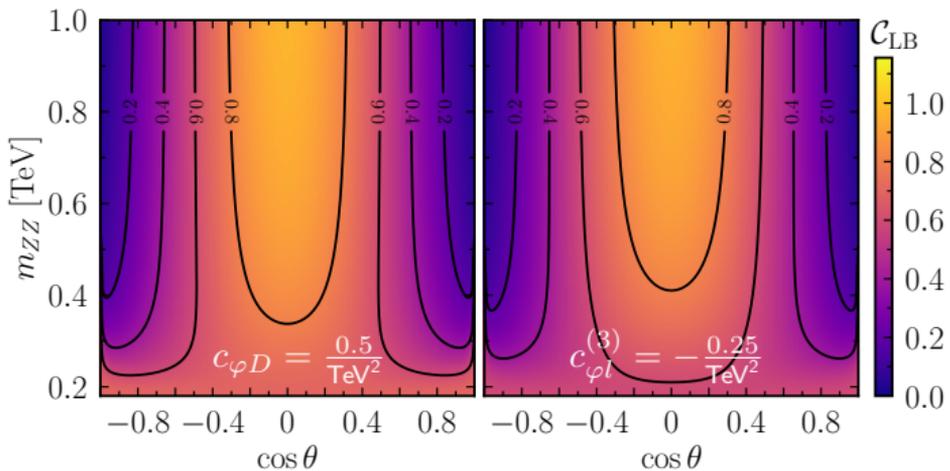
$$e^+e^- \rightarrow ZZ$$



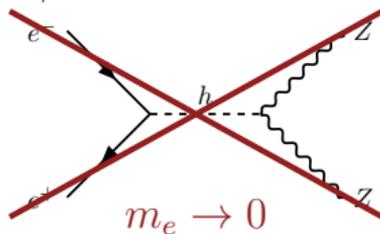
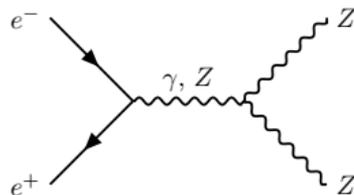
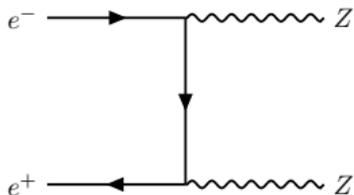
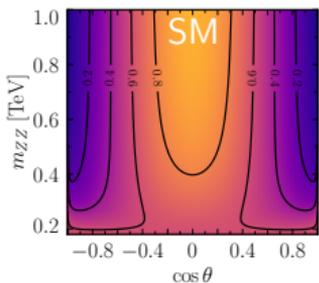
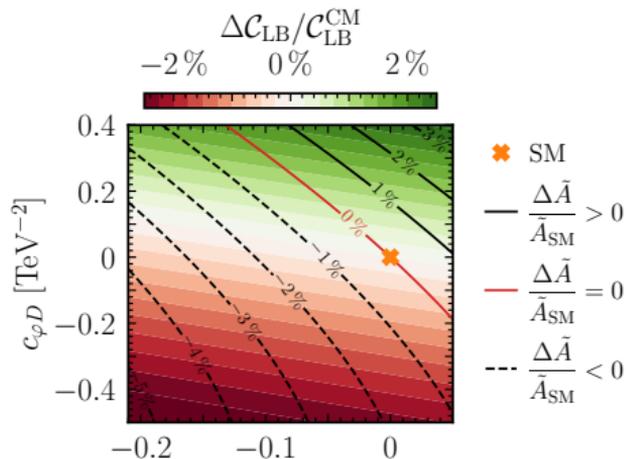
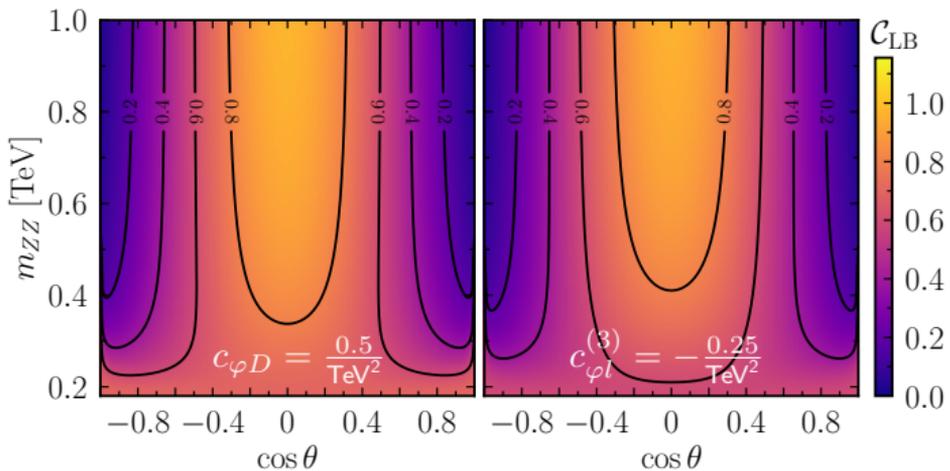
# $e^+e^- \rightarrow ZZ$



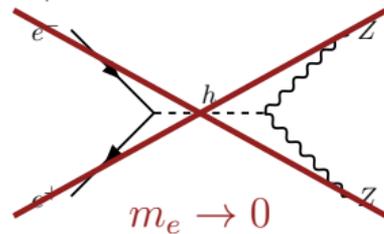
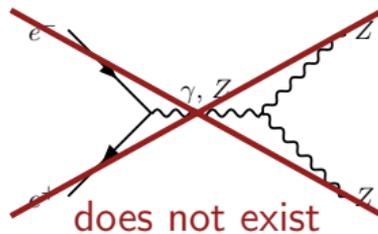
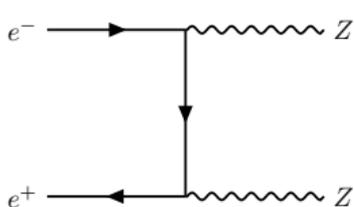
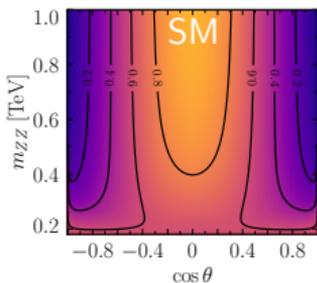
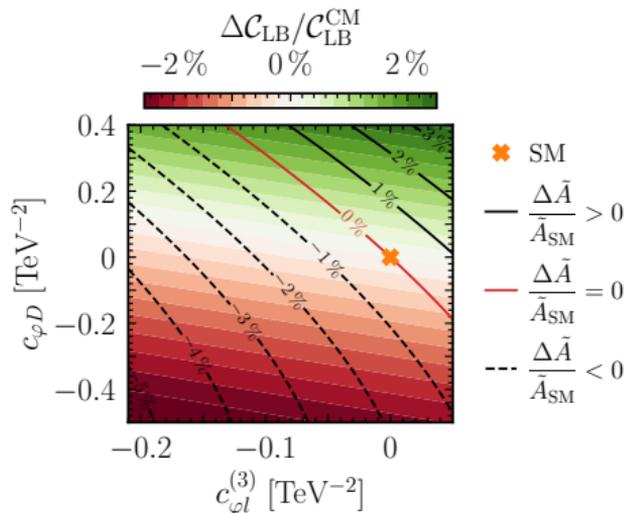
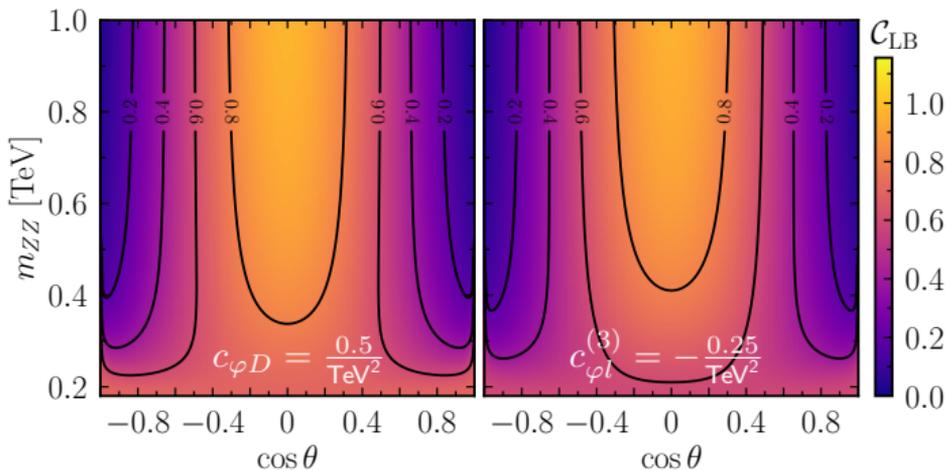
# $e^+e^- \rightarrow ZZ$



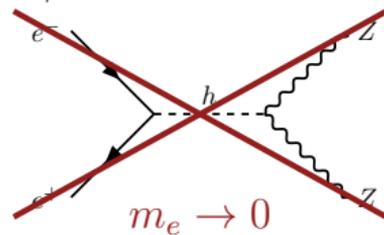
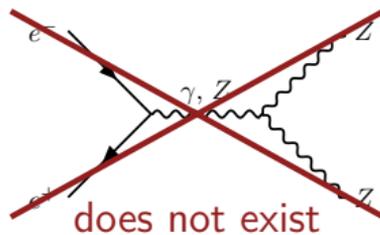
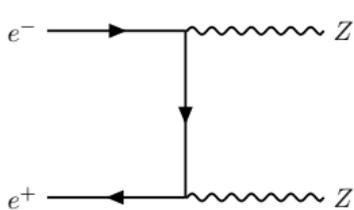
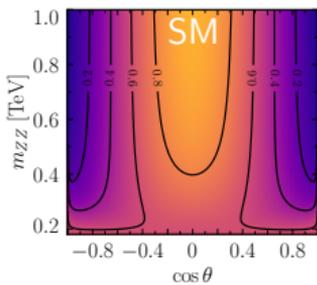
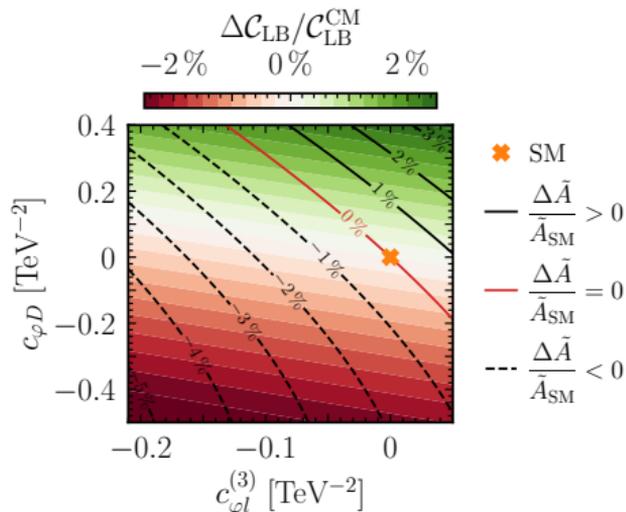
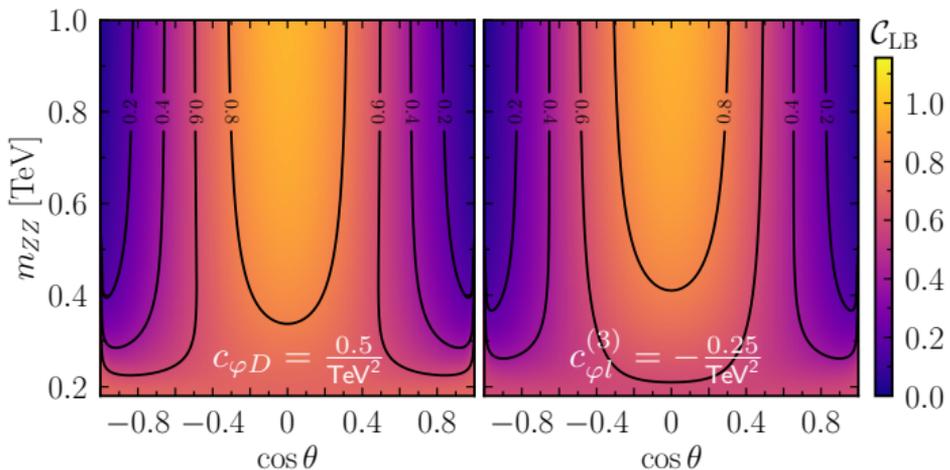
# $e^+e^- \rightarrow ZZ$



# $e^+e^- \rightarrow ZZ$



# $e^+e^- \rightarrow ZZ$



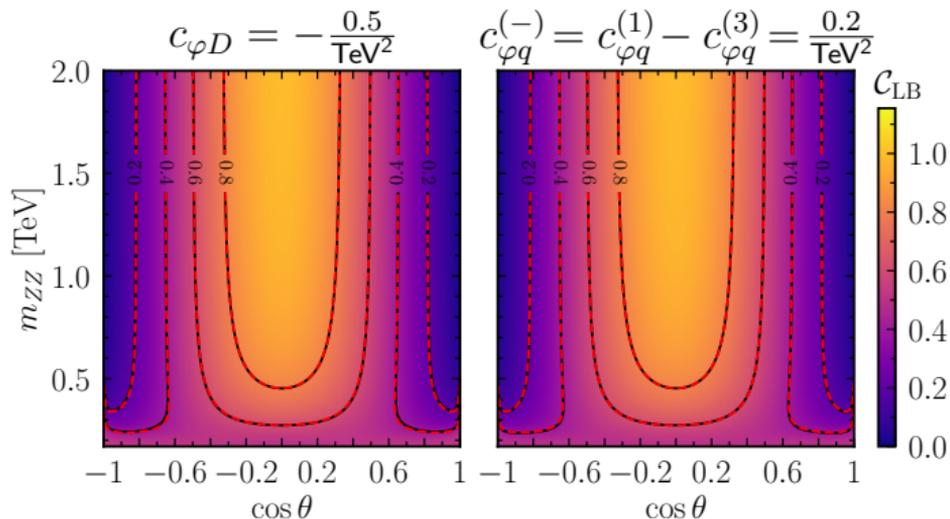
$\Rightarrow$  SMEFT effects only through balance of RH and LH couplings

# $pp \rightarrow ZZ$

- additional effect:

summation of  $\bar{u}u$  and  $\bar{d}d$  initial states

$\Rightarrow$  entanglement reduced  
(compared to partonic channels)



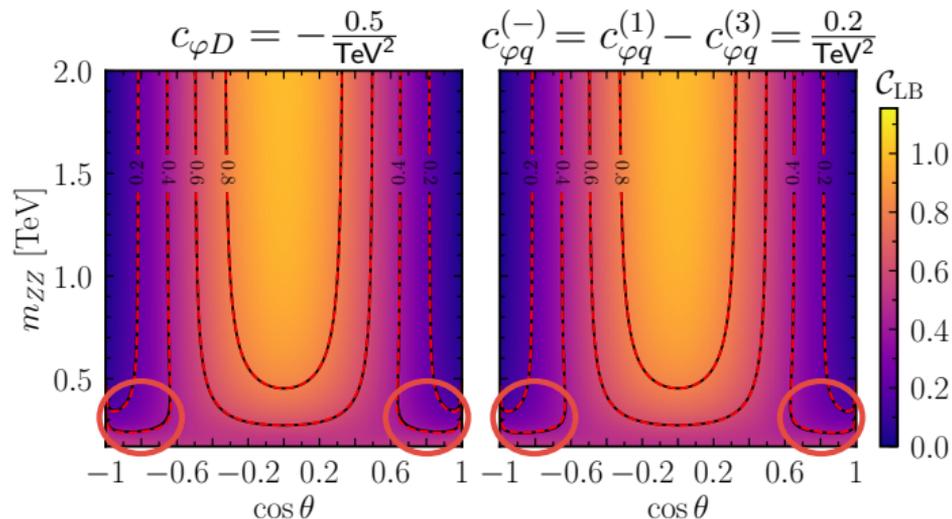
# $pp \rightarrow ZZ$

- additional effect:

summation of  $\bar{u}u$  and  $\bar{d}d$  initial states

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- only small effects in low-energy collinear region



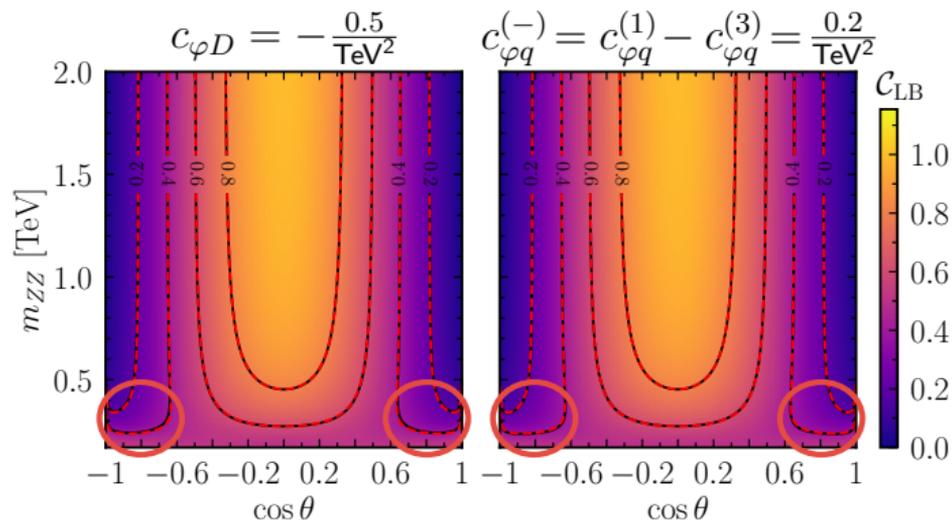
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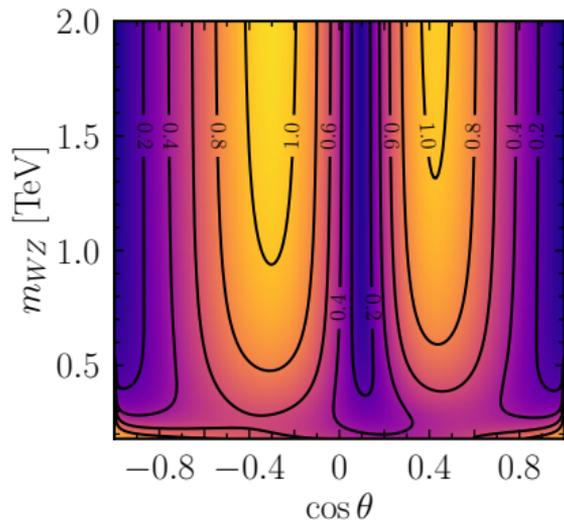


$\Rightarrow$  entanglement in  $pp \rightarrow ZZ$  not very sensitive to dim-6 effects

$$pp \rightarrow W^\pm Z$$

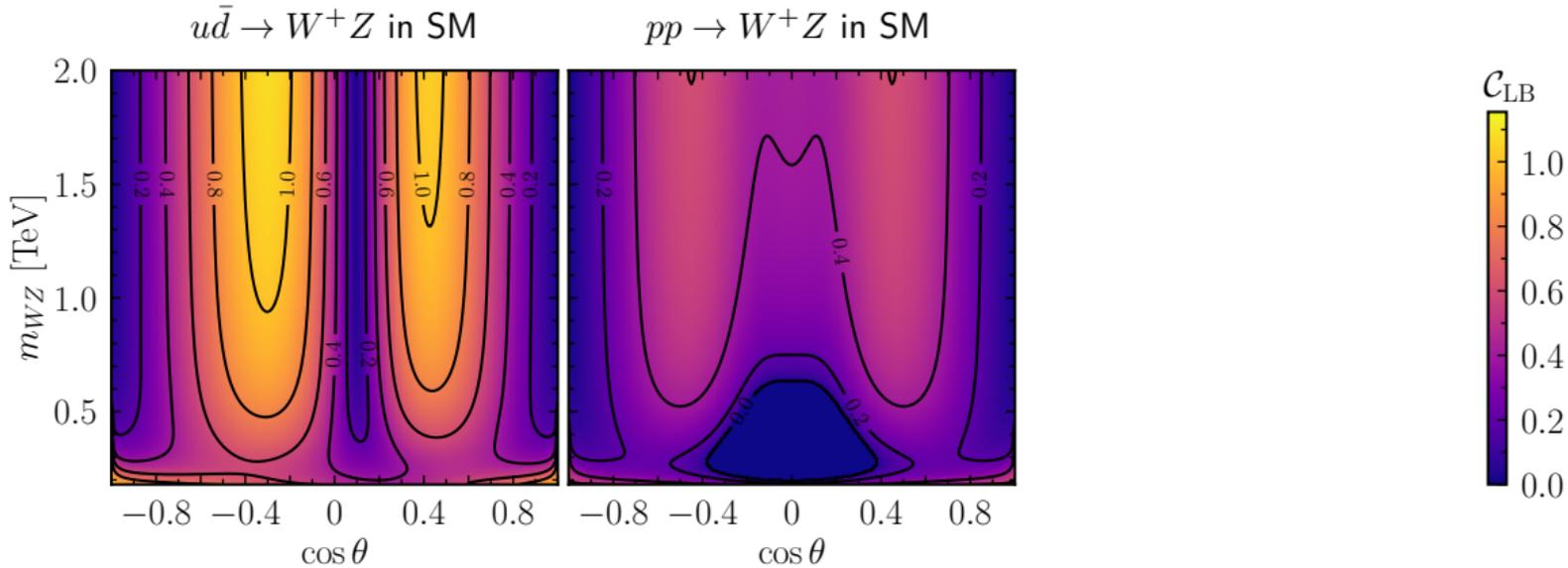
○ only one partonic channel ( $u\bar{d} \rightarrow W^+Z$ ); pure state

$u\bar{d} \rightarrow W^+Z$  in SM



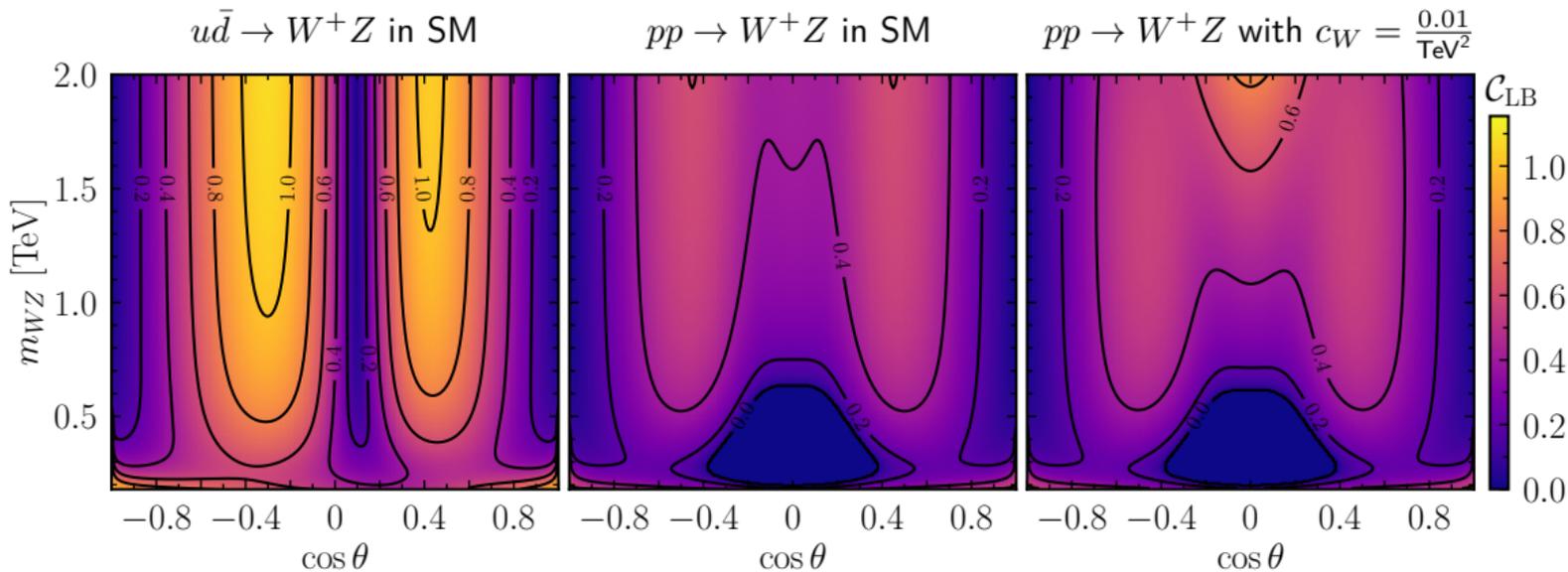
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- only one partonic channel ( $u\bar{d} \rightarrow W^+Z$ ); pure state
- at  $pp$  collider:  $\bar{d}u$  and  $u\bar{d}$   $\implies$  mixed state



# $pp \rightarrow W^\pm Z$

- only one partonic channel ( $u\bar{d} \rightarrow W^+Z$ ); pure state
- at  $pp$  collider:  $\bar{d}u$  and  $u\bar{d}$   $\implies$  mixed state
- quite sensitive to  $\mathcal{O}_W$  (and also  $\mathcal{O}_{\varphi q}^{(3)}$ )



# Summary

- EFT effects can modify the SM entanglement patterns
  - ⇒ entanglement-related observables can be used to probe new physics
- $e^+e^- \rightarrow W^+W^-$ ,  $pp \rightarrow W^+W^-$  and  $pp \rightarrow WZ$  are sensitive to dim-6 modifications
- $pp \rightarrow ZZ$  and  $e^+e^- \rightarrow ZZ$  are less sensitive to dim-6  
(but potentially to dim-8 operators)

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Thank you for your attention!



# Probing New Physics through Entanglement in Diboson Production

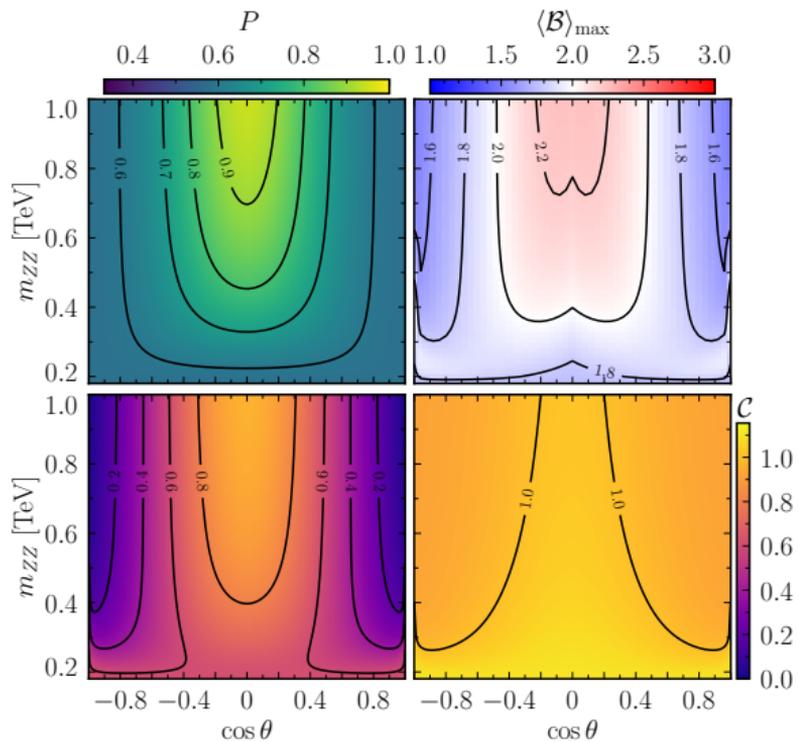
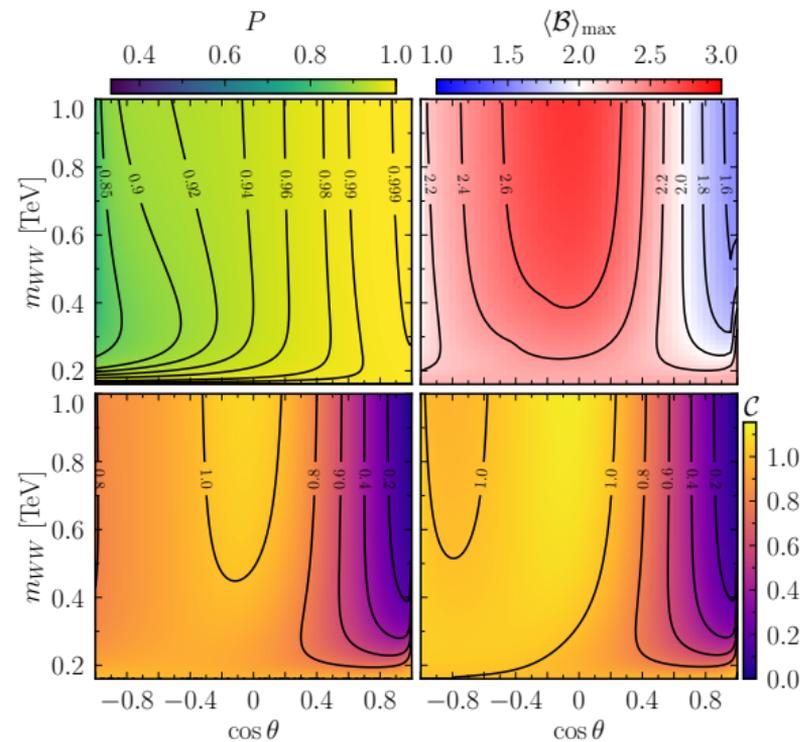
Rafael Aoude, E.M., Fabio Maltoni, Luca Mantani – arXiv:2307.09675 [hep-ph]

backup slides

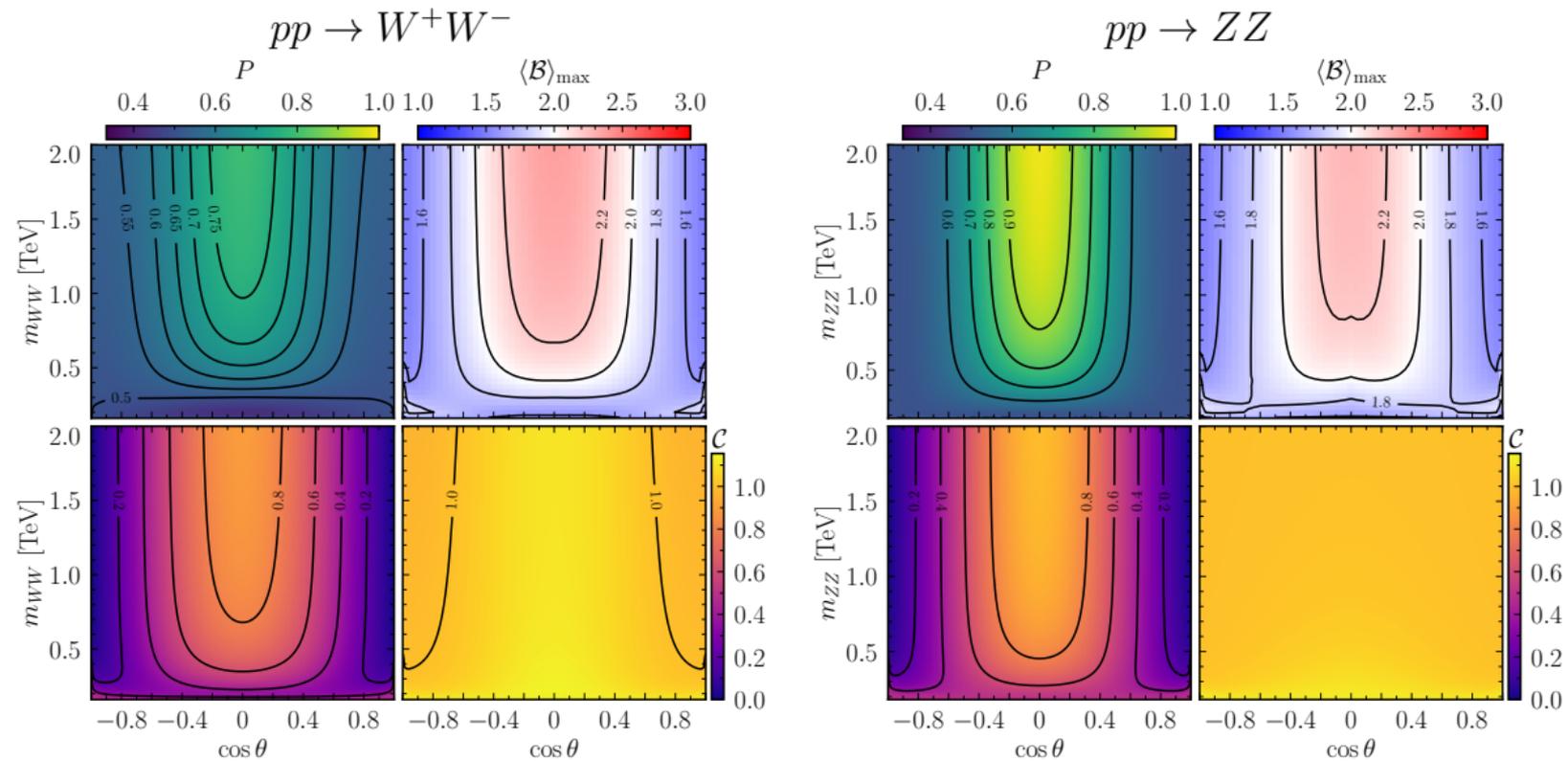
# Weak boson production at electron colliders in the SM

$$e^+e^- \rightarrow W^+W^-$$

$$e^+e^- \rightarrow ZZ$$



# Weak boson production at proton colliders in the SM I



# Weak boson production at proton colliders in the SM II

