



Observation of Entanglement in ______ Top Quark Pairs at ATLAS

Firenze, 05/11/2023

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 Testing quantum entanglement at hadron colliders is a brand new idea! Isn't it?

> CORRELATION EXPERIMENTS AND THE NONVALIDITY OF ORDINARY IDEAS ABOUT THE PHYSICAL WORLD**

> > Henry P. Stapp

Lawrence Berkeley Laboratory University of California Berkeley, California 94720

July 9, 1976

- In 1968, <u>Henry Stapp</u> proposed using spin correlations in proton-proton scattering.
- Add in top quarks and you have the ATLAS paper title.



• Fair to say that the 'buzzword' effect is in full swing:



Shameless Advertising



• Fair to say that the 'buzzword' effect is in full swing:



Breaking news: ATLAS ...

(one of Clara's worst performing videos...)

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• Interest appears to be growing (plot from summer 2023)



Interesting increase in non-LHC paper!



• What is quantum entanglement?

"An entangled state is one that cannot be written as a convex combination of product states of density matrices"

 If two particles are entangled, the quantum state of one particle cannot be described independently from the other:

$$|\psi\rangle = |a\rangle_A \otimes |b\rangle_B$$
 separable

 $|\psi\rangle = |a_1\rangle_A \otimes |b_1\rangle_B + |a_2\rangle_A \otimes |b_2\rangle_B$ non-separable

• Spin/polarisation is the canonical example of an observable to use to test entanglement.

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- Why do we care about doing this at the LHC?
- Entanglement has been measured before in composite systems:
 Diamonds, Mesons, Tardigrades, electrons in atoms).
- And in free particle systems (photons).
- top quarks are the first time it has been measured in an unbound fundamental fermion.
- top quarks, even when produced near threshold, are genuinely relativistic ($\beta \sim 0.4$) and Entanglement has never been explored under such conditions.

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- How does QE manifest in tt events?
- Key feature of top quarks is that it is VERY heavy! This leads to some unique properties:



 QCD has no time to dilute the top's quantum numbers, and they are transferred directly to its decay particles (where we can access them).



- QCD is P-conserving and T invariant → tops have no preferential polarisaton in tt production.
- But spins are correlated!

$$C = \frac{N(\uparrow\uparrow\uparrow) + N(\downarrow\downarrow) - N(\uparrow\downarrow) - N(\downarrow\uparrow)}{N(\uparrow\uparrow) + N(\downarrow\downarrow) + N(\uparrow\downarrow) + N(\downarrow\uparrow)}$$

• More formally:

PolarisationSpin Correlation $\frac{1}{\sigma} \frac{d^2 \sigma}{d \cos \theta^a_+ d \cos \theta^b_-} = \frac{1}{4} (1 + B^a_+ \cos \theta^a_+ + B^b_- \cos \theta^b_- - C(a, b) \cos \theta^a_+ \cos \theta^b_-)$

 Measuring B and C in tt essentially means constructing angles with the decay particles (usually charged leptons).



• It matters how you measure these angles!





It matters how you measure these angles!





• It matters how you measure these angles!





• The goal of the ATLAS measurement is to measure:

$$D = \frac{tr[C]}{3} = -3 \cdot < \cos(\phi) >$$

- Where φ is the angle between the top spin analysers in their parent top rest frames.
- An observation of D < -1/3 is a sufficient condition to claim entanglement in tt pairs (equivalently, that their density matrices are not factorable).

Getting the D

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- If D is observed to be less than -1/3, then the tops can be said to be entangled.



 This occurs in tt production when the tops are close to threshold (in gg fusion) or very boosted (in qqbar). We focus on the former.

Getting the small D

• The primary experimental challenges in this result are to reconstruct the tops with sufficient sensitivity to isolate the threshold region where tops are entangled.

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Signal/Validation Regions



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• Why is it hard to reconstruct top quarks?





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 Charged leptons are the perfect spin analyser!



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 Charged leptons are the perfect spin analyser!

Neutrinos notdetected (directly)by ATLAS

 ATLAS selects events with two charged leptons in the final state (+ 1 or more b-tagged jets).

Reconstructing Tops

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- In order to measure D, we need to fully reconstruct both tops (we need measure cos(Φ) in parent top rest frames).
 This means somehow dealing with two neutrinos
- There are a number of methods to achieve this, but this measurements relies heavily on the "Ellipse method".



Nucl.Instrum.Meth.A 736 (2014) 169-178

- Employs a geometry approach to analytically solve the system using linear algebra.
- Some other numerical methods used in small number of events.

Selection



 Events are selected with exactly 1 electron and 1 muon (standard p_T, η cuts).

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 Require 1 or more b-tagged jets (85% W.P):
 Ioose working point to ensure high stats in signal region.

• Three regions in m(tt) are defined:

SR: 340 < m(tt̄) < 380 GeV [High degree of entanglement]
 VR1: 380 < m(tt̄) < 500 GeV [some entanglement]
 VR2: m(tt̄) > 500 GeV [no entanglement]

Selection



• This selection is a very robust one (similar selection used in dozens of analyses).



 Very good overall agreement between the number of signal+background events and the observed number of events in data.

Calibration Curve



We somehow need to correct our observed D for detector effects: We achieve this with a calibration curve.



- To construct this curve we need to change the amount of entanglement in our MC.
- We create 5 hypothesis points corresponding to the SM and 4 different reweighing points: (+20%, -20%, -40%, -60%)

Reweighting

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- How these alternative hypothesis points are constructed is one of the key points of the measurement.
- We cannot dial entanglement up or down in the MC, so we reweight the cos(Φ) distribution as a function of m(tt̄).



• If this is not done correctly, the relation:

$$D = \frac{tr[C]}{3} = -3 \cdot < \cos(\phi) >$$

does not hold.

 The method we have used ensures that this relationship remains correct.



• The relative size of the systematics is not fixed and changes at each hypothesis point:





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 Ideally, truth and reco shift in a correlated way, and there is no resultant uncertainty.



• The relative size of the systematics is not fixed and changes at each hypothesis point:



• In practice, most uncertainties shift reco but not truth and therefore change the slope (all detector uncertainties do this).



• The relative size of the systematics is not fixed and changes at each hypothesis point:



 In the worst case, systematics shift slope and offset and have a large effect (our dominant uncertainties behave this way).



 The relative size of the systematics is not fixed and changes at each hypothesis point:

Systematic source	$\Delta D_{\rm observed} (D = -0.547)$	ΔD (%)	$\Delta D_{\text{expected}}(D = -0.470)$	$\Delta D (\%)$
Signal Modelling	0.017	3.2	0.015	3.2
Electrons	0.002	0.4	0.002	0.4
Muons	0.001	0.1	0.001	0.1
Jets	0.004	0.7	0.004	0.8
<i>b</i> -tagging	0.002	0.4	0.002	0.4
Pile-up	< 0.001	< 0.1	< 0.001	< 0.1
$E_{\mathrm{T}}^{\mathrm{miss}}$	0.002	0.3	0.002	0.4
Backgrounds	0.010	1.8	0.009	1.8
Total Statistical Uncertainty	0.002	0.3	0.002	0.4
Total Systematic Uncertainty	0.021	3.8	0.018	3.9
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 As with most top measurements, we are limited by signal modelling, though background modelling (Z+jets) matters too due to looser b-tag and shape of the background.



• We have a large suite of MC modelling related systematic uncertainties:

Systematic uncertainty source	Relative size (for SM D value)
Top-quark decay	1.6%
Parton distribution function	1.2%
Recoil scheme	1.1%
Final-state radiation	1.1%
Scale uncertainties	1.1%
NNLO reweighting	1.1%
pThard setting	0.8%
Top-quark mass	0.7%
Initial-state radiation	0.2%
Parton shower and hadronization	0.2%
$h_{\rm damp}$ setting	0.1%

 Colour reconnection, string vs cluster fragmentation, spin correlation in parton shower, EW shower were all tested but found to be negligible effects.





• The observed (expected) results are:

SR $D = -0.547 \pm 0.002$ [stat.] ± 0.021 [syst.] (-0.470 ± 0.002 [stat.] ± 0.018 [syst.]),

VR1 $D = -0.222 \pm 0.001$ [stat.] ± 0.027 [syst.] (-0.258 ± 0.001 [stat.] ± 0.026 [syst.]),

VR2 $D = -0.098 \pm 0.001$ [stat.] ± 0.021 [syst.] (-0.103 ± 0.001 [stat.] ± 0.021 [syst.]),



 The observed results excludes the entanglement limit at more than 5 sigma significance.



• How reliable are the elements of this result?



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• How reliable are the elements of this result?



Jay Howarth

• Corrections to the data: very reliable

 A comprehensive and conservative (even by ATLAS's standards) list
 of systematic
 uncertainties has
 been considered on
 all aspects of the
 analysis.

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• How reliable are the elements of this result?



- Predictions of the SM: Reliable but limited.
- These predictions come from general purpose MC generators:
 - We understand them very well, but they are not designed to model threshold perfectly.

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• How reliable are the elements of this result?



- Entanglement limits: Reliable but limited.
- Same limitations as predictions.
- Two models give different limits, but source is understood and we've taken the most conservative of the two.

What about Topponium?

 Bound state effects are most prevalent in the region that we care about.



• These are not directly included in our MC simulations (but we have attempted to introduce them as a cross-check and other uncertainties cover similar effects).

Kiyo, Kühn, Moch, Steinhauser, Uwer, 2009

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- Bound state effects should be increasing entanglement:
 - Including them only makes result more significant, not less.



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*exaggerated, the effect on the error bars would be too small to see.

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Parton Shower



• Difference seems to come from the ordering of the shower.



 Angular ordered showers have a large effect compared to dipole showers.

Jay Howarth

Doesn't effect detector corrections significantly.



- ATLAS has observed quantum entanglement for the first time in a pair of fundamental quarks, at the highest labmade energies.
- This is the first step in a program to use the LHC as a tool for exploring quantum information.
- Important questions about how entanglement (and spin correlation) is modelled in this threshold region:
 Would be a very profitable area for further study in the theory community!