

# New foundational experiments with quantum process tomography

*Proc. R. Soc. A.* **478**:20210806 (2022), arXiv:2103.12000

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JAGIELLONIAN UNIVERSITY  
IN KRAKÓW



University  
of Gdańsk



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OF TECHNOLOGY

Firenze, 6 November 2023

# Routes towards New Physics

$$\text{Standard Model} \subset \text{QFT} = \text{Quantum Mechanics} + \text{Special Relativity}$$

Routes towards New Physics:

- ① Beyond Standard Model, but still in QFT
  - SUSY, composite Higgs, dark sector, inflation, ...
- ② Beyond Special Relativity, but assuming QM
  - QFT in curved spacetimes – ‘semi-classical’ (Unruh effect, Hawking radiation ...)
  - quantum gravity
- ③ Beyond Quantum Mechanics, but assuming relativity
  - Super-quantum correlations, GPTs, ...
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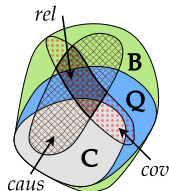
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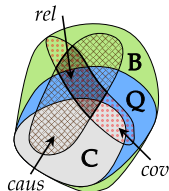


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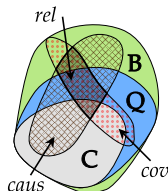


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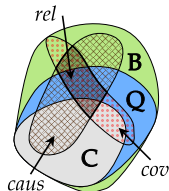


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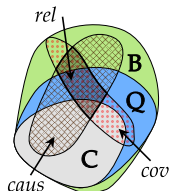


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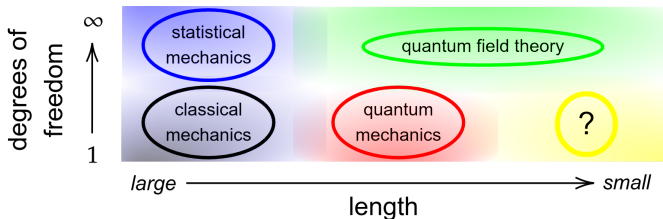
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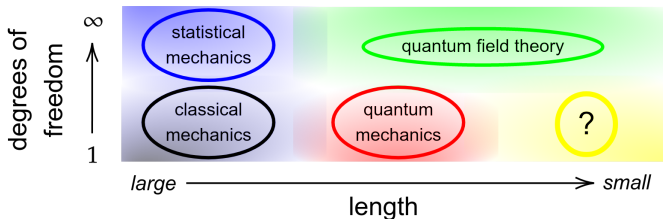


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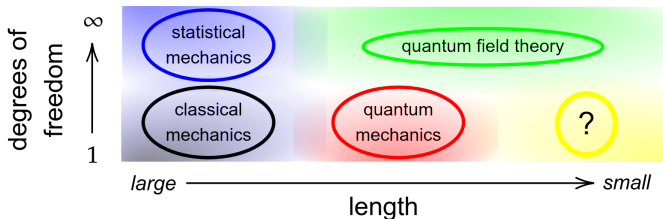
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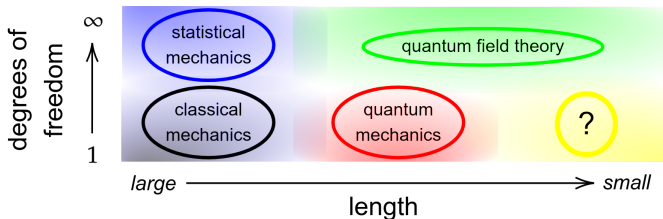
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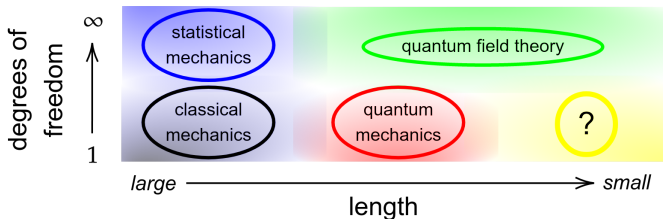
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# The theory independent **black box** methodology

- Physical systems are treated (merely!) as information-processing devices (“**black boxes**”) and probed by free agents.
- The conclusions are drawn from the **output–input correlations**.

$$P(\text{outputs} \mid \text{inputs})$$

Bell test: 2 agents (Alice and Bob) — 2 inputs  $(x, y)$  — 2 outputs  $(a, b)$

The *experimental* (frequency)  
correlation function:

$$C_e(x, y) = P(a = b \mid x, y) - P(a \neq b \mid x, y)$$

The key assumption of *freedom of choice* (“measurement independence”):

$$P(x, y \mid \lambda) = P(x) \cdot P(y)$$

- No pre-correlations between the inputs  $(x, y)$  and the box  $(\lambda)$ .



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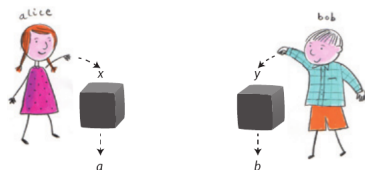
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[Sandu Popescu, *Nature Physics* 10, 264 (2014)]

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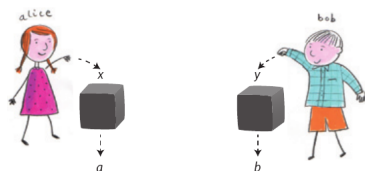
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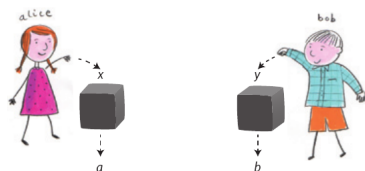
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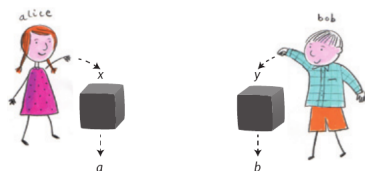
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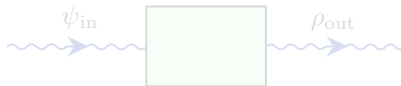
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# Quantum-data boxes

- We treat physical systems as **Q-data boxes**, i.e. *quantum-information* processing devices.
- A Q-data box is probed *locally* with quantum information.



- $p$  are classical parameters (e.g. scattering kinematics)
- The *pure input* state is **prepared**,  $P : x \rightarrow \psi_{\text{in}}$ .
- The *output state* is reconstructed via **quantum tomography** from the outcomes of projective measurements  $M : \rho_{\text{out}} \rightarrow a$ .

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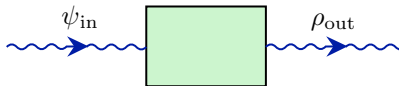


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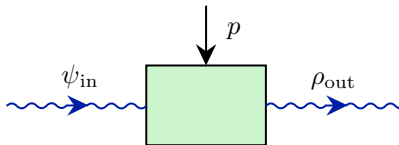
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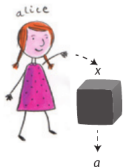
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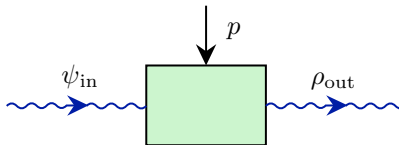
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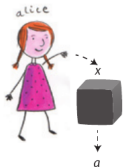
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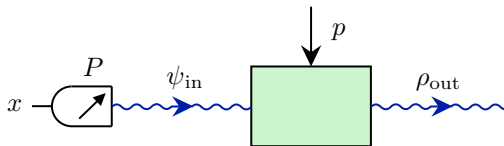
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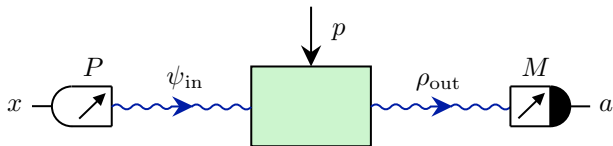
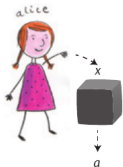
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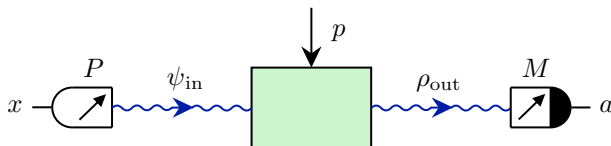
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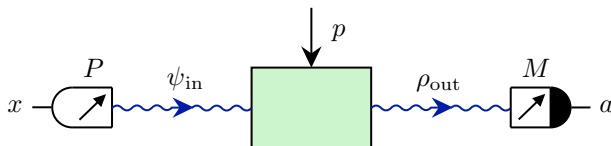
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- $\psi_{\text{in}}$  is pure, initially **uncorrelated** with the box — **freedom of choice**.
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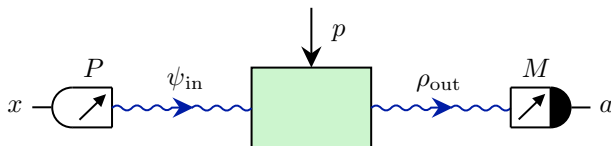
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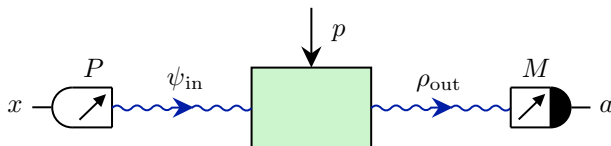


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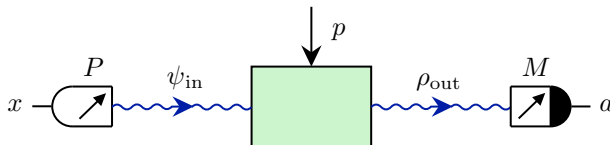
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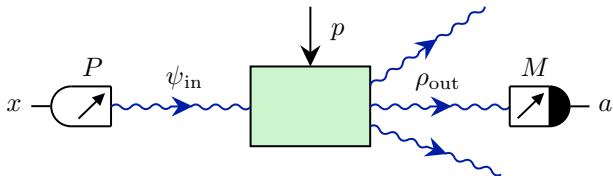
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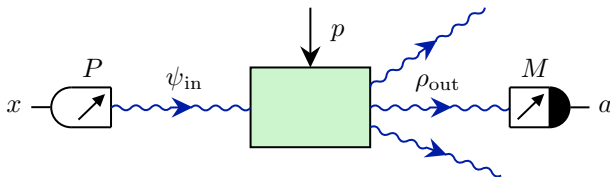
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- Don’t need to gather all outgoing quantum information.

# Quantum-data tests



A **Q-data test** consists in probing a Q-data box with *prepared* input states.

- For every input state  $\psi_{\text{in}}$  one performs the full tomography of  $\rho_{\text{out}}$ .
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- Suppose that we have two available inputs  $\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}$ .
- We choose randomly the input (with probability  $1/2$ ).
- The task is to guess, which of the two states was input.
- Define the **success rate**:  $P_{\text{succ}}(\psi_{\text{in}}^{(1)}, \psi_{\text{in}}^{(2)}) := \frac{1}{2} \sum_{k=1}^2 P(a = k | \psi_{\text{in}}^{(k)})$ .
- In quantum theory  $P_{\text{succ}}$  cannot exceed the **Helstrom bound**

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- Make a Q-data test with  $\{\psi_{\text{in}}^{(k)}; \rho_{\text{out}}^{(k)}\}_{k=1,2}$ .
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# Quantum process tomography

- In QM *any* dynamics  $\mathcal{E} : S(\mathcal{H}_{\text{in}}) \rightarrow S(\mathcal{H}_{\text{out}})$  must be a **CP(TP) map**,  $\mathcal{E} \otimes \mathbb{1}_N \geq 0 \quad \forall N$ .
- $\mathcal{E}$  is completely characterised by  $m^2(n^2 - 1)$  real parameters,  $m = \dim \mathcal{H}_{\text{in}}, n = \dim \mathcal{H}_{\text{out}}$ .
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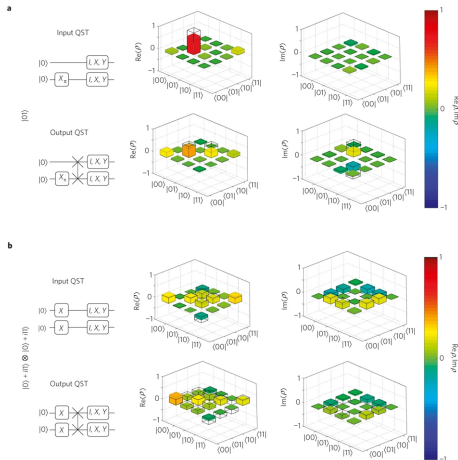
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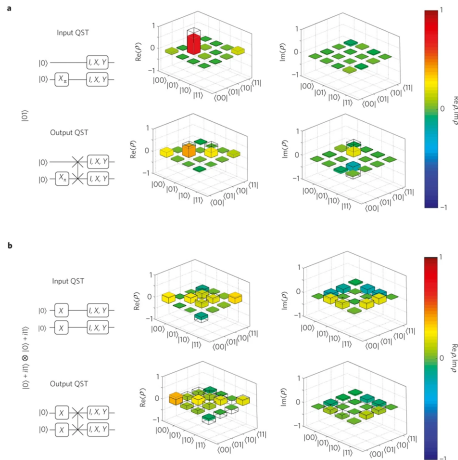


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e.g. electron’s spin or photon’s polarization.  $\rightsquigarrow$  **polarized beams**
  - ➋ Probe a ‘Q-data box’ – that is **collide them!**
  - ➌ Reconstruct the output states  $\rho_{\text{out}}$ .  $\rightsquigarrow$  **weak decays**
- Experimental prospects:
- Explore the spin dynamics with polarised beams,  
 $\mathcal{H}_{\text{in}} = \mathbb{C}^2 \otimes \mathbb{C}^2$ ,  $\mathcal{H}_{\text{out}} = \mathbb{C}^2 \otimes \mathbb{C}^2$  or  $\mathcal{H}_{\text{out}} = \mathbb{C}^3 \otimes \mathbb{C}^3$
  - Direct processes, e.g.  $e^+e^- \rightarrow t\bar{t}$  [C. Altomonte, A. Barr (2022)]
  - Indirect processes, e.g.  $e^-p^+ \rightarrow \gamma g + X \rightarrow t\bar{t} + Y$   
Such an operation  $\mathcal{E} : \mathbb{C}_p^2 \otimes \mathbb{C}_{e^-}^2 \rightarrow S(\mathbb{C}_{t^+}^2 \otimes \mathbb{C}_{t^-}^2)$  is *not* unitary,  
but it must be CP *if* quantum mechanics is valid.
  - Ongoing work with C. Altomonte, A. Barr, P. Horodecki & K. Sakurai.

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*Proc. R. Soc. A.* **478**:20210806 (2022), arXiv:2103.12000

## Take-home messages:

- To make a proper Bell-type test you need the **freedom of choice!**
  - Entanglement detection is not a Bell test!
  - Could we make **direct** projective measurements of spin??
  - A proper Bell test could detect beyond-quantum correlations.
- **Quantum process tomography** offers new opportunities:
  - Seek deviations from QM (unitarity, CP, linearity, ...)
  - Understand quantum *dynamics* in HEP.
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