Modica, 14th July 2023

Supernovae as factories of neutrinos and other feebly interacting particles

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Core-Collapse Supernovae

For massive stars $(M > 8M_{\odot})$ the nuclear fusion produces heavy elements in an onion structure and a degenerate iron core





Iron in the core cannot be burnt and the star starts to collapse

How many successfull SNe?

H. T. Janka, [arXiv:1702.08825 [astro-ph.HE]].

Very non-trivial to predict if a SN explodes or not, see how much it depends on metallicity



Phases of neutrino emission



Each phase is strongly characterized by a different neutrino signal

SN neutrino emission

T. Fischer et al., Phys. Rev. D 94 (2016) no.8, 085012



Orders of magnitude for SNe

The SN core is an extreme environment



SN1987A: neutrino signal

From the few $\bar{
u}_e p
ightarrow ne^+$ events of SN 1987A we know that...



 $\sim 10^{53}\,{
m erg}$ emitted as neutrinos with energy $\sim {\it O}(15\,{
m MeV})$ in $\sim 10\,{
m s}$

SN1987A: we don't understand the neutrino signal (?) S. W. Li *et al.*, [arXiv:2306.08024 [astro-ph.HE]].



SN models generally agree with each other and disagree with data

Fundamental physics with SNe G. G. Raffelt, 1996, ISBN 978-0-226-70272-8

Constraint on neutrino mass:

$$\Delta t_{
u} = 2.57 \,\mathrm{s} \left(rac{d_{SN}}{50 \,\mathrm{kpc}}
ight) \left(rac{E_{
u}}{10 \,\mathrm{MeV}}
ight)^{-2} \left(rac{m_{
u}}{10 \,\mathrm{eV}}
ight)^2
ightarrow m_{
u} \lesssim 23 \,\mathrm{eV}$$

Constraint on neutrino lifetime:

$$au/m_{
u}\gtrsim 6 imes 10^5$$
 s/eV

Secret interactions with D
u B, speed of neutrinos/light, neutrino charge

Axions: Why? What?

The energy-loss argument

G. Raffelt, Lect. Notes Phys. 741 (2008)

Stars produce feebly interacting particles which escape, draining energy from the core



They strongly affect the SN neutrino burst if

$$\mathcal{L}_{ ext{FIP}} > \mathcal{L}_{
u} = 3 imes 10^{52} \, ext{erg s}^{-1}$$

The strong CP problem

The QCD Lagrangian includes a CP-odd term

$$\mathcal{L}_{ ext{eff}} = \mathcal{L}_{ ext{QCD}} - ar{ heta}_{ ext{QCD}} rac{oldsymbol{g}^2}{32\pi^2} \operatorname{tr} ilde{G}_{\mu
u} G^{\mu
u}$$

where $ilde{G}_{\mu
u} = rac{1}{2} \epsilon_{\mu
ulphaeta} G^{lphaeta}$ and $ar{ heta}_{ ext{QCD}} = heta_{ ext{QCD}} + \operatorname{arg} \det M_{ ext{quark}}$

 $\begin{array}{lll} \mbox{Prediction of neutron electric dipole moment} \\ \mathcal{L} \sim d_n \cdot E & \rightarrow & d_n \approx |\bar{\theta}_{\rm QCD}| \times 10^{-15} e\, {\rm cm} \\ \mbox{Experimental bound: } |\bar{\theta}_{\rm QCD}| < 10^{-10} \\ \mbox{Naturalness problem, why } \bar{\theta}_{\rm QCD} \mbox{ is so small}? \end{array}$

The Peccei-Quinn mechanism

R. D. Peccei et al., Phys. Rev. Lett. 38 (1977)

PQ symmetry

 $U(1)_{PQ}$ is a chiral global symmetry that drives dynamically $ar{ heta}_{PQ}
ightarrow 0$

 $U(1)_{PQ}$ is broken at a scale f_a , the **Peccei-Quinn scale**, and the Goldstone boson is the **axion**

$$\mathcal{L}_{ax} = \frac{1}{2} \partial_{\mu} a \partial^{\mu} a - \xi \frac{a}{f_a} \frac{g^2}{32\pi^2} \tilde{G}^a_{\mu\nu} G^{\mu\nu a} + \frac{g_a}{2m} \bar{\Psi} \gamma^{\mu} \gamma^5 \Psi \partial_{\mu} a - \frac{g_{a\gamma}}{4} a \tilde{F}^{\mu\nu} F_{\mu\nu}$$

The minimum condition removes the CP-odd term: $ar{ heta}_{
m QCD}=0$

Axion-SM interactions

Axion-photon vertex



Axions and Axion-Like Particles

In any axion model

$$m_a \sim rac{1}{f_a} \quad g_{a\gamma} \sim rac{1}{f_a} \quad f_a \gg 246 \ {
m GeV}$$

The typical QCD axion is light and weakly interacting

Axion-Like Particles (ALPs) are a generalization:

- Heavy ALP searches at collider
- Superlight ALPs as fuzzy Dark Matter
- Some ALPs could be the inflaton
- ALPs in flavor-violating processes...

Motivations to study axions and ALPs

Axions and ALPs are a window on high-energy physics



This hot topic is a motivation for interdisciplinary searches

Axion production in nuclear processes

Axion Lagrangian

L. Di Luzio et al., Phys. Rept. 870 (2020), 1-117

$$\mathcal{L}_{
m int} = g_a rac{\partial_\mu a}{2m_N} \Bigg[C_{ap} ar{p} \gamma^\mu \gamma_5 p + C_{an} ar{n} \gamma^\mu \gamma_5 n +$$

$$+ \frac{C_{a\pi N}}{f_{\pi}} (i\pi^+ \bar{p}\gamma^\mu n - i\pi^- \bar{n}\gamma^\mu p) +$$

$$+ C_{aN\Delta} \left(\bar{p} \,\Delta_{\mu}^{+} + \overline{\Delta_{\mu}^{+}} \,p + \bar{n} \,\Delta_{\mu}^{0} + \overline{\Delta_{\mu}^{0}} \,n \right)$$



Axion-nucleon bremsstrahlung in SNe M. S. Turner, Phys. Rev. Lett. **60** (1988)

PC, T. Fischer et al., JCAP 10 (2019) no.10, 016

SN axions are produced by nucleon-axion bremsstrahlung



where we have to include detailed nuclear physics and many body effects

Pion-axion conversion in SNe
PC, B. Fore *et al.*, Phys. Rev. Lett. **126** (2021) no.7, 071102
SN axions are produced by pion-axion conversion



This is the leading axion production process in a SN despite the small density of pions $(\mathcal{O}(1\%))!!$

Flux from pion-axion conversion A. Lella *et al.*, [arXiv:2306.01048 [hep-ph]].

The harder spectrum is due to the pion rest mass



Notice the lower energies as axions become trapped

Consequences on the SN cooling

Clear behavior of free-streaming/trapping regime



SN axion bounds

PC et al., [arXiv:2306.17055 [astro-ph.HE]].



SN axion bound from KII (green): a+ $^{16}O \rightarrow$ $^{16}O^{*}$

Conclusions



"Always the last place you look!"