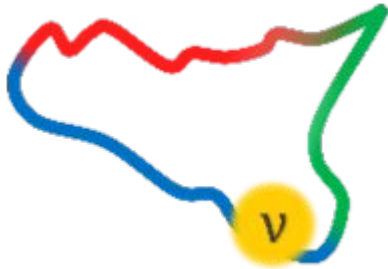




Fundamental Physics with Nuclei

MAYORANA International Workshop

13 July 2023
Saori Pastore



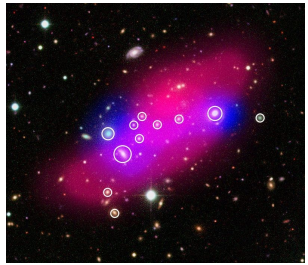
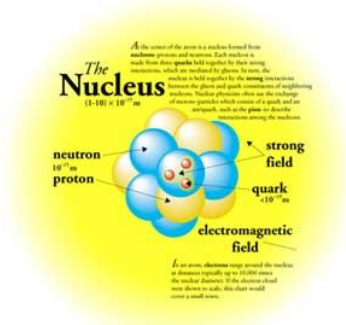
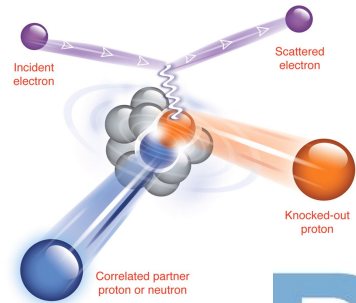
<https://physics.wustl.edu/quantum-monte-carlo-group>

Quantum Monte Carlo Group @ WashU

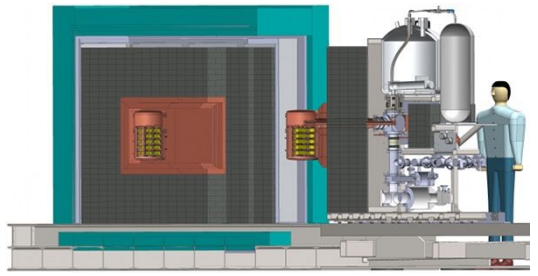
Lorenzo Andreoli (PD) Jason Bub (GS) Graham Chambers-Wall (GS) Garrett King (GS)
Anna McCoy (FRIB TA Fellow)
Maria Piarulli and Saori Pastore

Computational Resources awarded by the DOE ALCC, INCITE and SciDAC programs

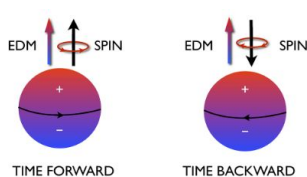
Understand Nuclei to Understand the Cosmos



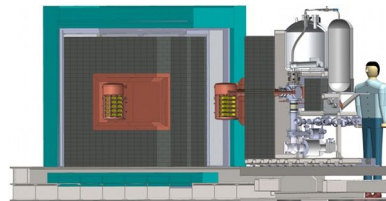
ESA, XMM-Newton, Gaspardello, CFHTL



Ground States'
Electroweak Moments,
Form Factors, Radii



Neutrinoless Double
Beta Decay,
Muon-Capture



Accelerator Neutrino
Experiments,
Lepton-Nucleus XSecs

$(\omega, q) \sim 0$ MeV

$\omega \sim \text{few MeVs}$
 $q \sim 0$ MeV

$\omega \sim \text{few MeVs}$
 $q \sim 10^2$ MeV

$\omega \sim \text{tens of MeVs}$

$\omega \sim 10^2$ MeV



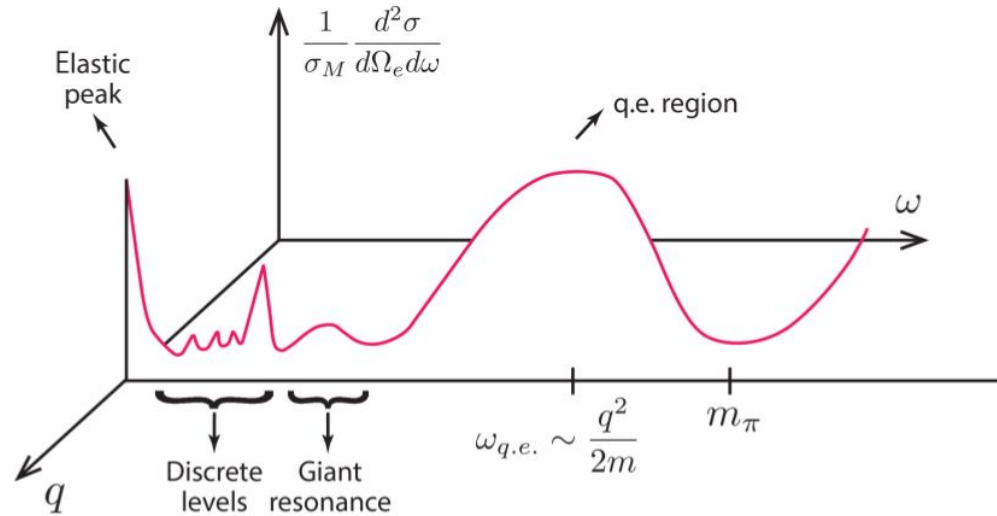
Electromagnetic
Decay, Beta Decay,
Double Beta Decay &
inverse processes



Nuclear Rates for
Astrophysics



Electron-Nucleus Scattering Cross Section



Energy and momentum transferred (ω, q)

Current and planned experimental programs rely on theoretical calculations at different kinematics

Strategy

Validate the Nuclear Model against available data for strong and electroweak observables

- Energy Spectra, Electromagnetic Form Factors, Electromagnetic Moments, ...
- Electromagnetic and Beta decay rates, ...
- Muon Capture Rates, ...
- Electron-Nucleus Scattering Cross Sections, ...

Use attained information to make (accurate) predictions for BSM searches and precision tests

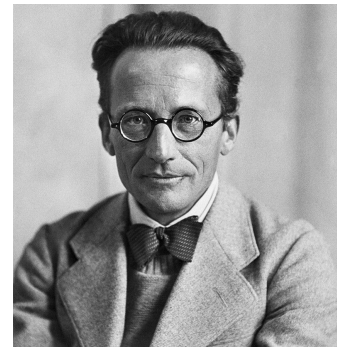
- EDMs, Hadronic PV, ...
- BSM searches with beta decay, ...
- Neutrinoless double beta decay, ...
- Neutrino-Nucleus Scattering Cross Sections, ...
- ...

Microscopic (or *ab initio*) Description of Nuclei

Comprehensive theory that describes quantitatively and predictably nuclear structure and reactions

Requirements:

- Accurate understanding of the interactions/correlations between nucleons in **pairs, triplets, ... (two- and three-nucleon forces)**
- Accurate understanding of the electroweak interactions of external probes (electrons, neutrinos, photons) with nucleons, correlated nucleon-pairs, ... (**one- and two-body electroweak currents**)
- **Computational methods** to solve the many-body nuclear problem of strongly interacting particles



Erwin Schrödinger

$$H\Psi = E\Psi$$

Many-body Nuclear Problem

Nuclear Many-body Hamiltonian

$$H = T + V = \sum_{i=1}^A t_i + \sum_{i<j} v_{ij} + \sum_{i<j<k} V_{ijk} + \dots$$

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A, s_1, s_2, \dots, s_A, t_1, t_2, \dots, t_A)$$

Ψ are **spin-isospin** vectors in **3A** dimensions with $2^A \times \frac{A!}{Z!(A-Z)!}$ components

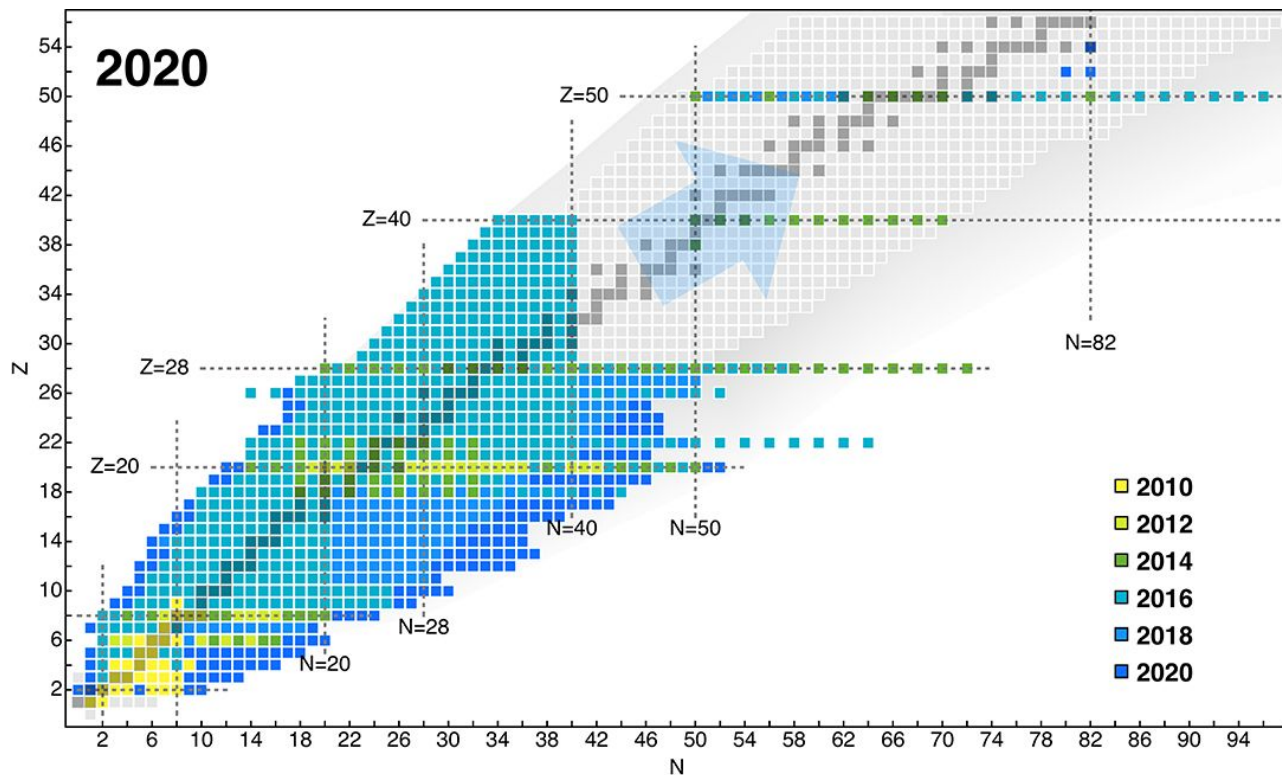
Develop Computational Methods to solve (numerically) exactly or within approximations that are under control the many-body nuclear problem

${}^4\text{He}$: 96
 ${}^6\text{Li}$: 1280
 ${}^8\text{Li}$: 14336
 ${}^{12}\text{C}$: 540572



<http://exascale.org/np/>

Current Status



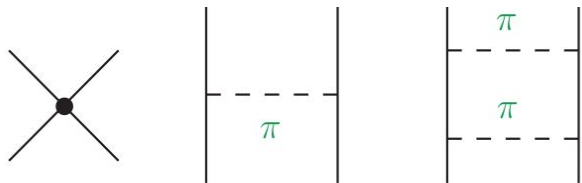
H. Hergert
Front. Phys.
07 October 2020

Many-body Nuclear Interactions

Many-body Nuclear Hamiltonian

$$H = T + V = \sum_{i=1}^A t_i + \sum_{i<j} v_{ij} + \sum_{i<j<k} V_{ijk} + \dots$$

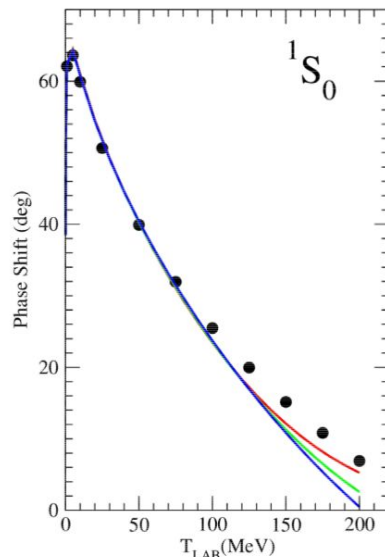
v_{ij} and V_{ijk} are **two-** and **three-**nucleon operators based on experimental data fitting; fitted parameters subsume underlying QCD dynamics



Contact term: short-range

Two-pion range: intermediate-range $r \propto (2m_\pi)^{-1}$

One-pion range: long-range $r \propto m_\pi^{-1}$



SP et al. PRC80(2009)034004



Hideki Yukawa

AV18+UIX; **AV18+IL7**

Wiringa, Schiavilla, Pieper
et al.

chiral $\pi N\Delta$

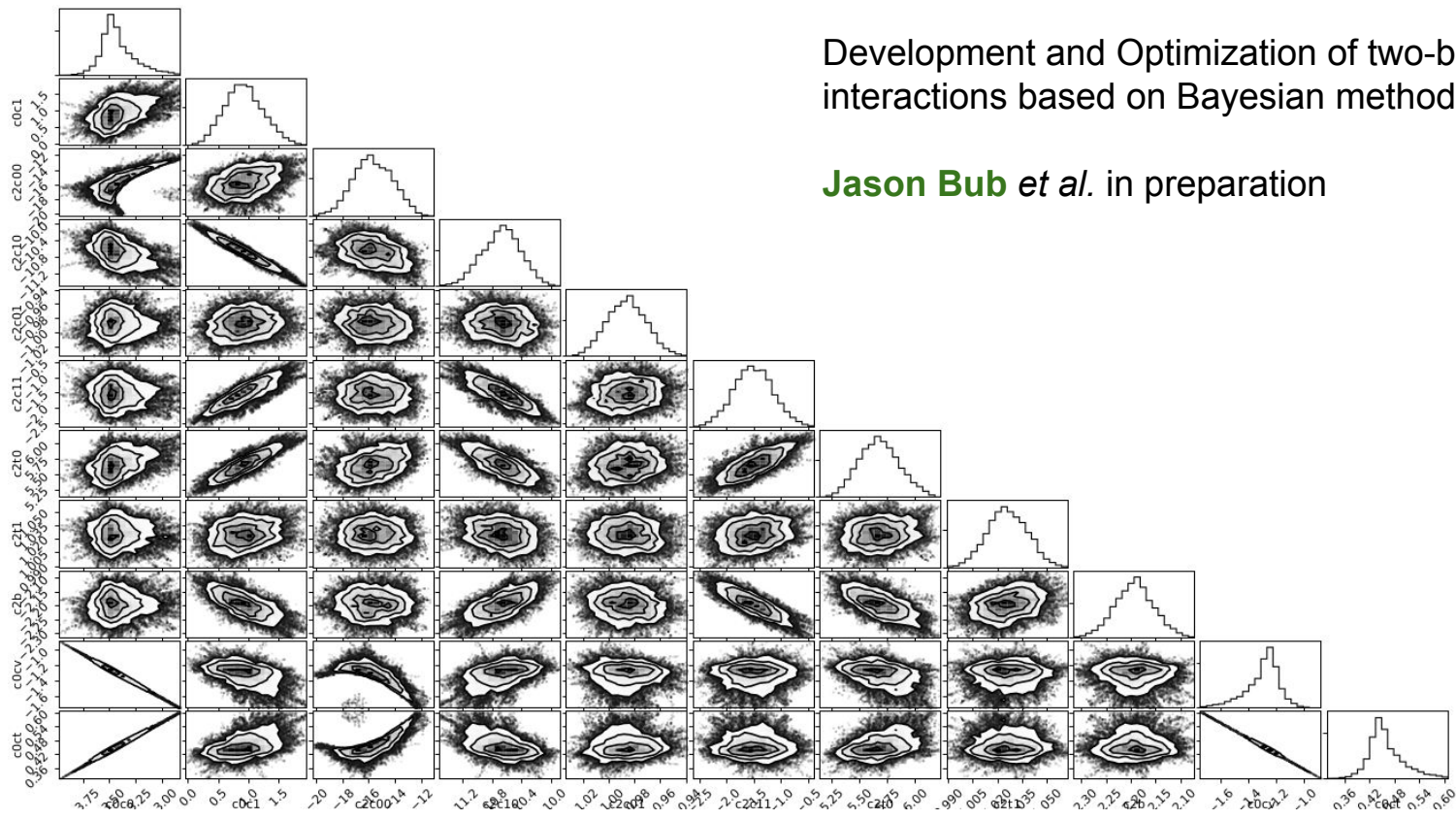
N3LO+N2LO Piarulli *et al.*

et al. **Norfolk Models**

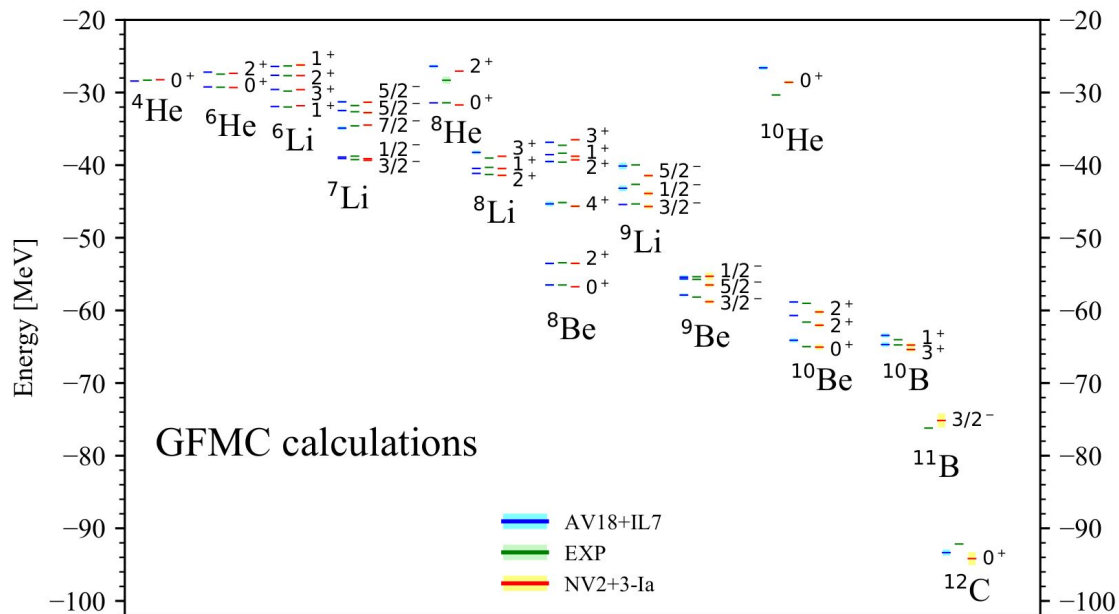
Optimization of Nuclear Two-body Interactions

Development and Optimization of two-body interactions based on Bayesian methods

Jason Bub *et al.* in preparation

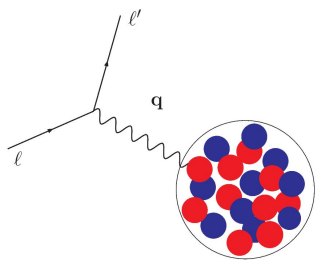


Energies

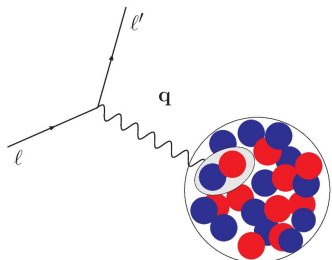


Piarulli *et al.* PRL120(2018)052503

Many-body Nuclear Electroweak Currents



one-body



two-body

- Two-body currents are a manifestation of two-nucleon correlations
- Electromagnetic two-body currents are required to satisfy current conservation

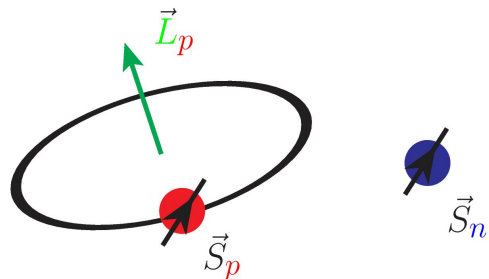
$$\mathbf{q} \cdot \mathbf{j} = [H, \rho] = [t_i + v_{ij} + V_{ijk}, \rho]$$

Nuclear Charge Operator

$$\rho = \sum_{i=1}^A \rho_i + \sum_{i<j} \rho_{ij} + \dots$$

Nuclear (Vector) Current Operator

$$\mathbf{j} = \sum_{i=1}^A \mathbf{j}_i + \sum_{i<j} \mathbf{j}_{ij} + \dots$$



Magnetic Moment: Single Particle Picture

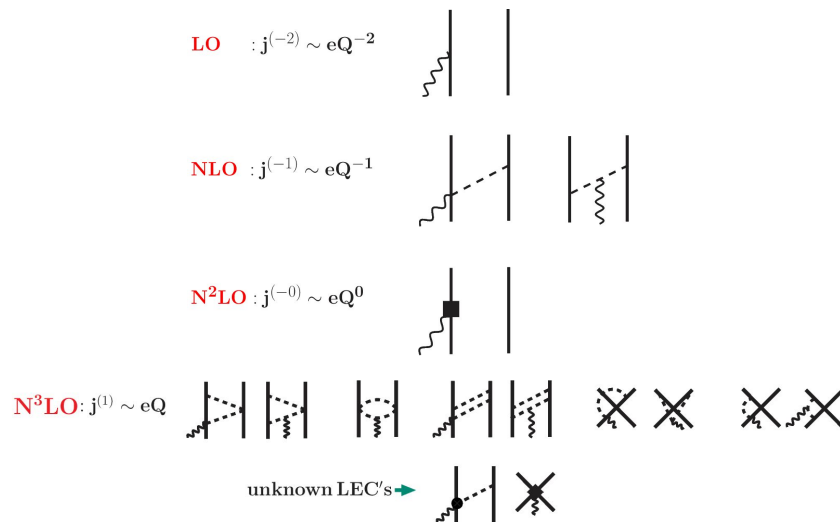
Many-body Currents

- **Meson Exchange Currents (MEC)**

Constrain the MEC current operators by imposing that the current **conservation relation is satisfied with the given two-body potential**

- **Chiral Effective Field Theory Currents**

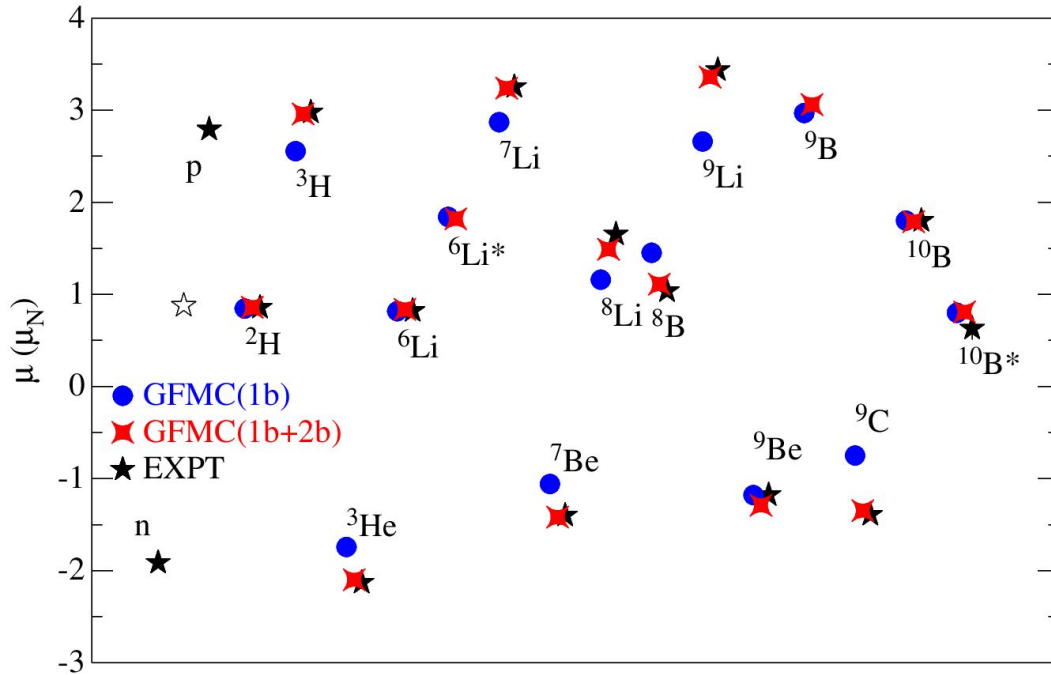
Are constructed consistently with the two-body chiral potential; Unknown parameters, or Low Energy Constants (**LECs**), need to be **determined by either fits to experimental data or by Lattice QCD calculations**



Electromagnetic Current Operator

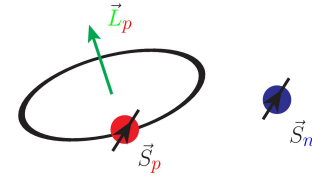
SP *et al.* PRC78(2008)064002, PRC80(2009)034004,
 PRC84(2011)024001, PRC87(2013)014006
 Park *et al.* NPA596(1996)515, Phillips (2005)
 Kölling *et al.* PRC80(2009)045502 & PRC84(2011)054008

Magnetic Moments of Light Nuclei



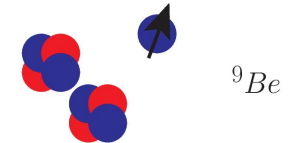
SP *et al.* PRC87(2013)035503

Single particle picture

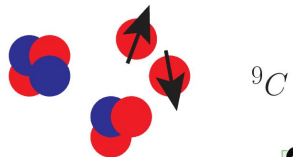


$$\mu_N(1b) = \sum_i [(L_i + g_p S_i)(1 + \tau_{i,z})/2 + g_n S_i(1 - \tau_{i,z})/2]$$

Small two-body current effects

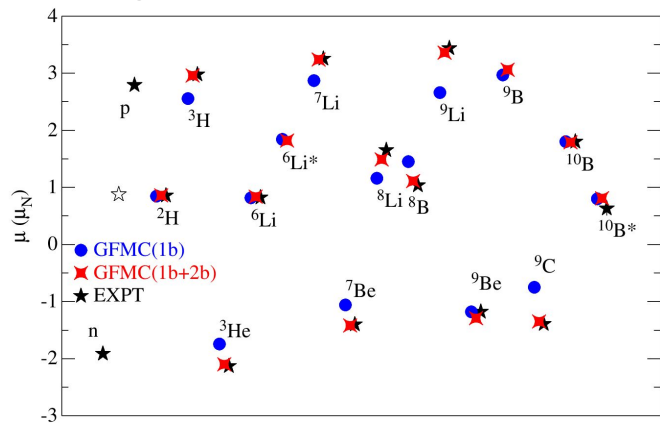


Large two-body current effects



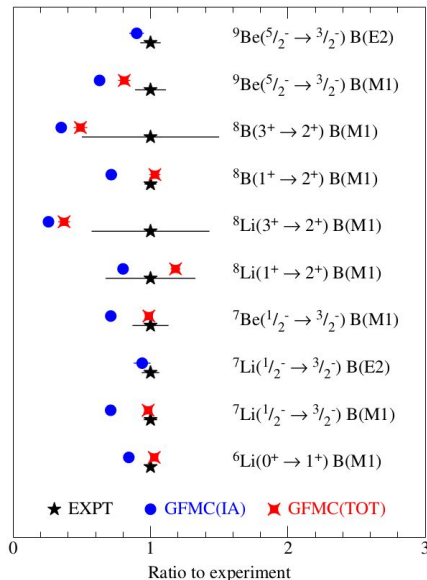
Electromagnetic Observables

Magnetic moments

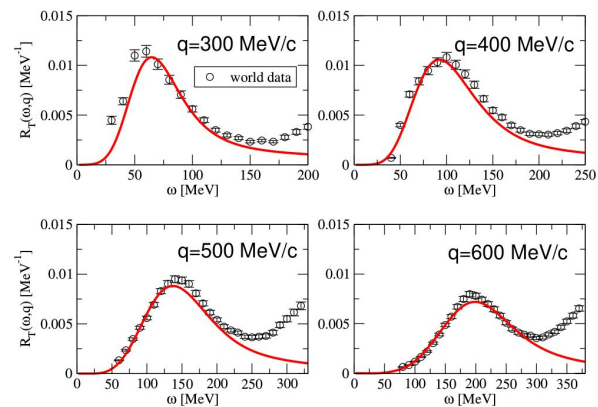


SP *et al.* PRC87(2013)035503,
 PRC101(2020)044612

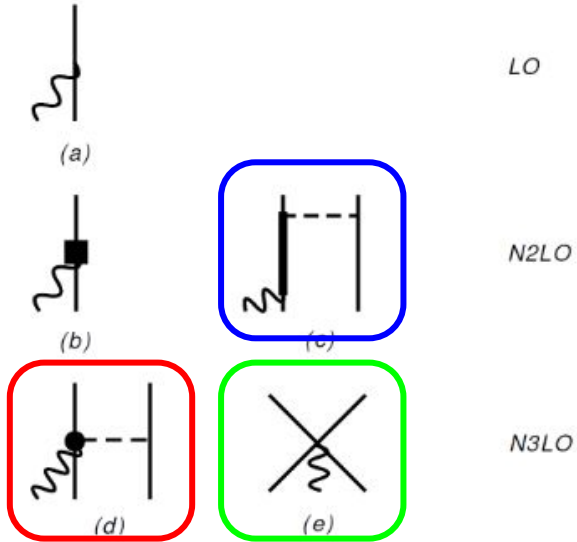
EM decay



e - ^4He particle scattering



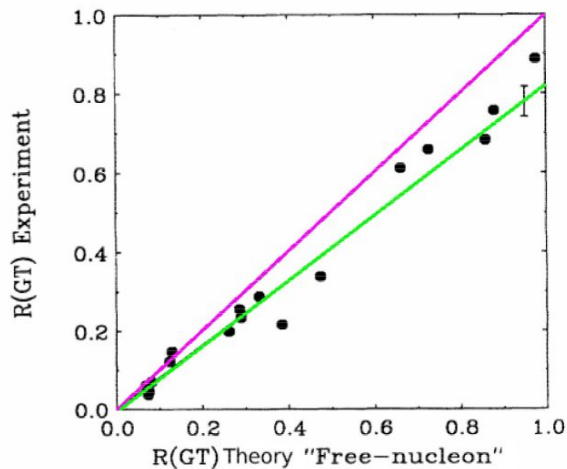
Axial currents with Δ at tree-level



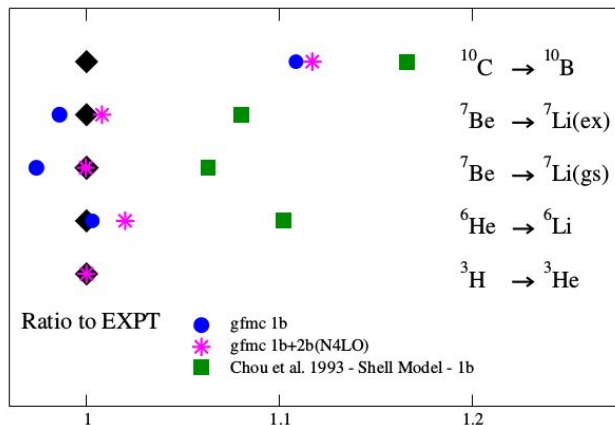
Two body currents of one pion range
(red and blue) with c_3 c_4 from Krebs
et al. Eur.Phys.J.(2007)A32

Contact current involves the LEC c_D

Beta decay

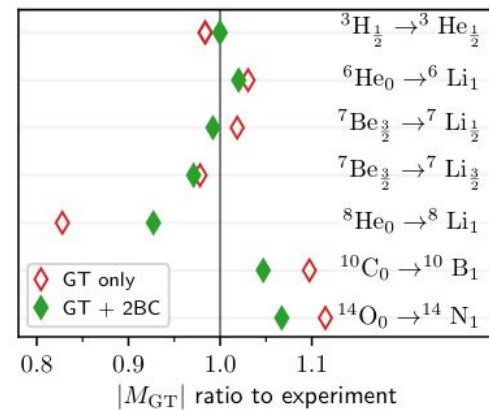


Chou et al. PRC47(1993)163



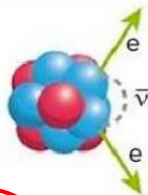
gfm (1b) and gfm (1b+2b); shell model (1b)

SP et al. PRC97(2018)022501

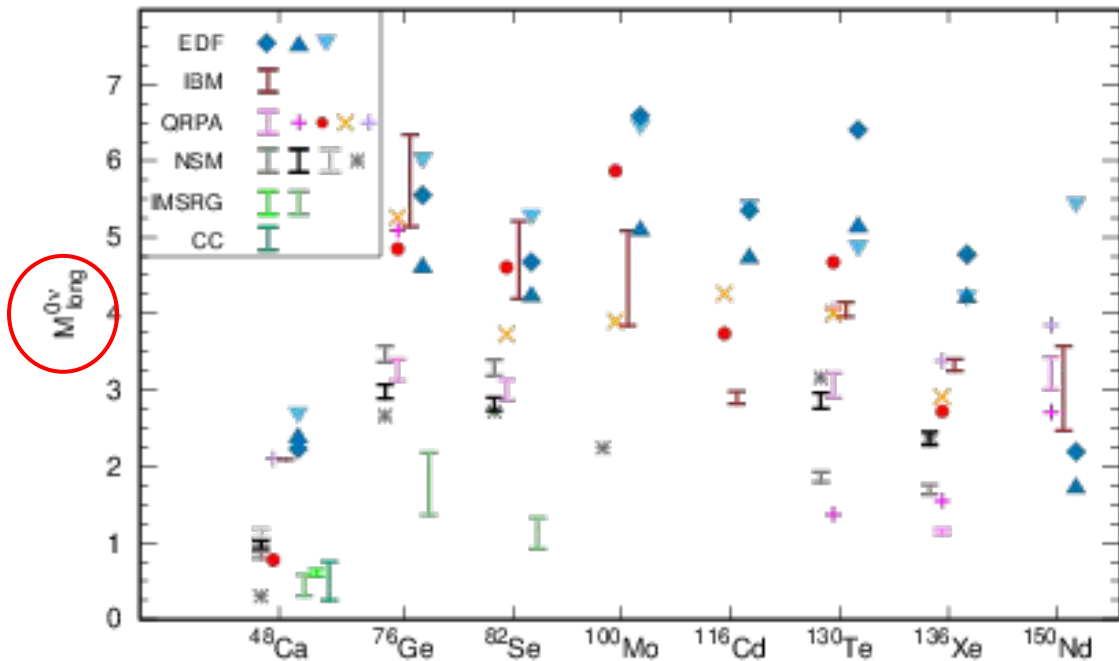


P. Gysbers *Nature Phys.* 15 (2019)

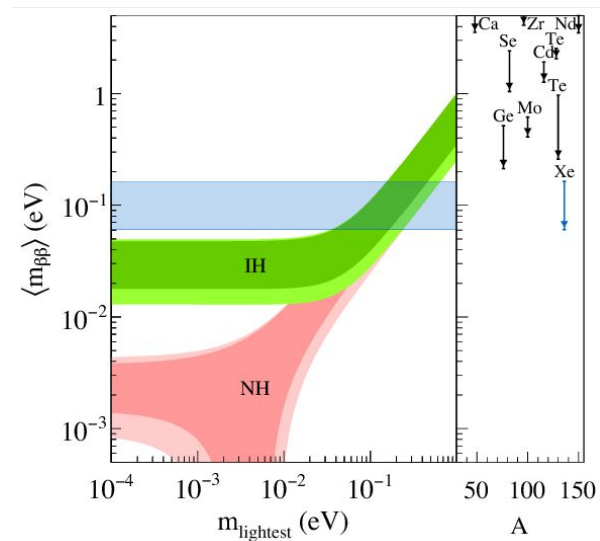
Neutrinoless Double Beta Decay



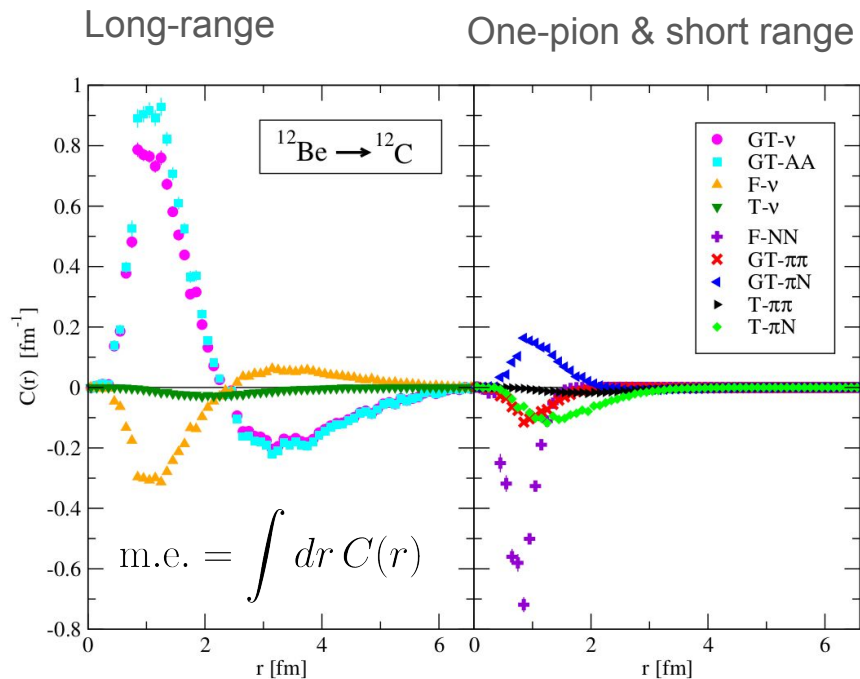
$$[T_{1/2}^{0\nu}]^{-1} = G_{0\nu}(Q, Z) |M_{0\nu}|^2 m_{\beta\beta}^2$$



Menéndez Vissani et al. *Rev.Mod.Phys.* 95 (2023) 2, 025002



Neutrinoless Double Beta Decay Matrix Elements



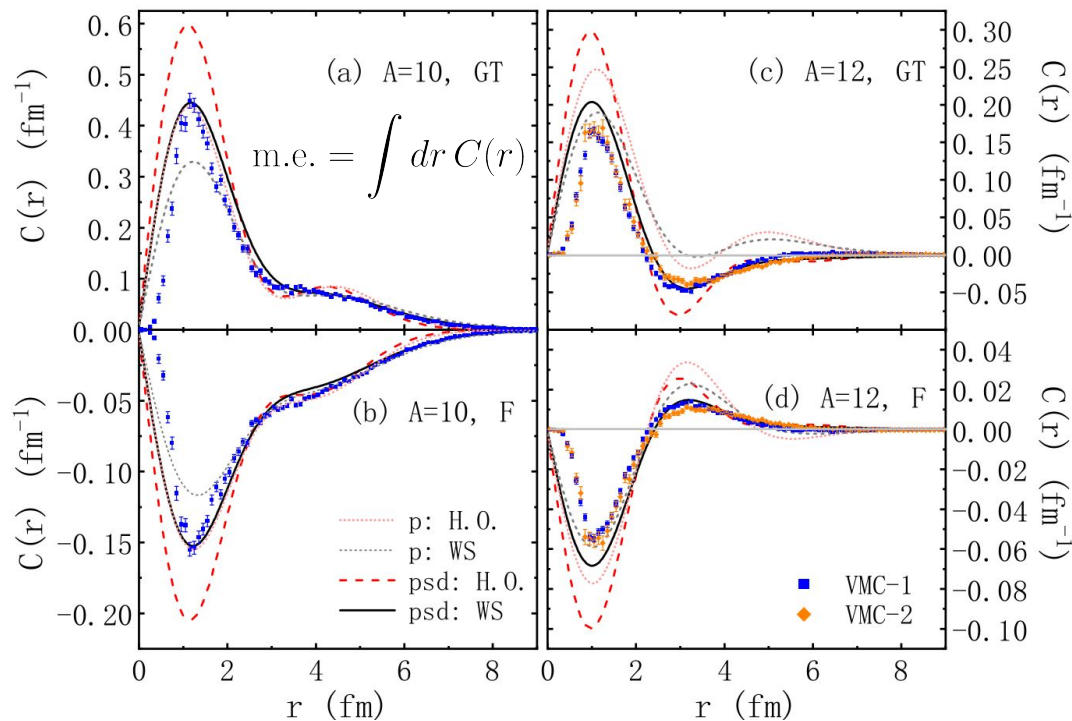
SP *et al.* PRC97(2018)014606



Cirigliano Dekens DeVries Graesser Mereghetti *et al.*
 PLB769(2017)460, JHEP12(2017)082, PRC97(2018)065501

- Leading operators in neutrinoless double beta decay are two-body operators
- These observables are particularly sensitive to short-range and two-body physics
- Transition densities calculated in momentum space indicate that the momentum transfer in this process is of the order of $\mathbf{q} \sim 200 \text{ MeV}$

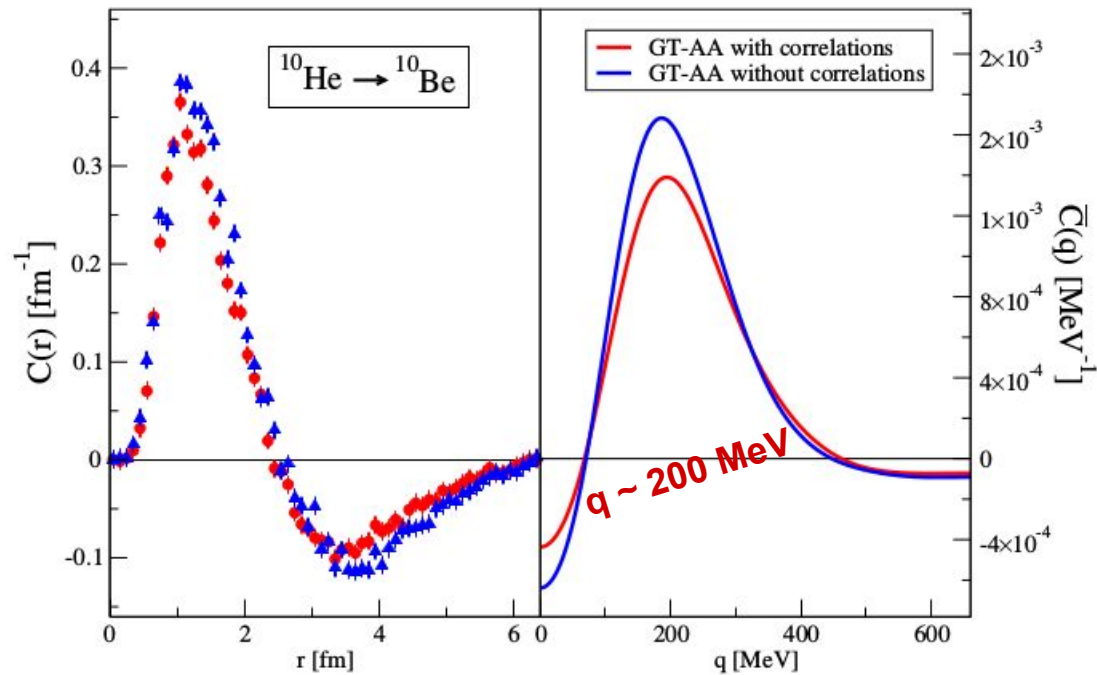
Comparison with Shell-Model Calculations



Closer agreement between Shell-Model calculations with Variational Monte Carlo results is reached by

- Increasing the size of the model space
- Wood-Saxon single particle wave functions are superior in describing the tails of the densities wrt harmonic oscillator wave functions
- Phenomenological Short-Range-Correlations functions further improve the agreement

Correlations in neutrinoless double beta decay ME



Partial muon capture rates: VMC calculations

$$\Gamma_{\text{VMC}}(\text{avg.}) = 1495 \text{ s}^{-1} \pm 19 \text{ s}^{-1}$$

$$\Gamma_{\text{expt}} = 1496.0 \text{ s}^{-1} \pm 4.0 \text{ s}^{-1}$$

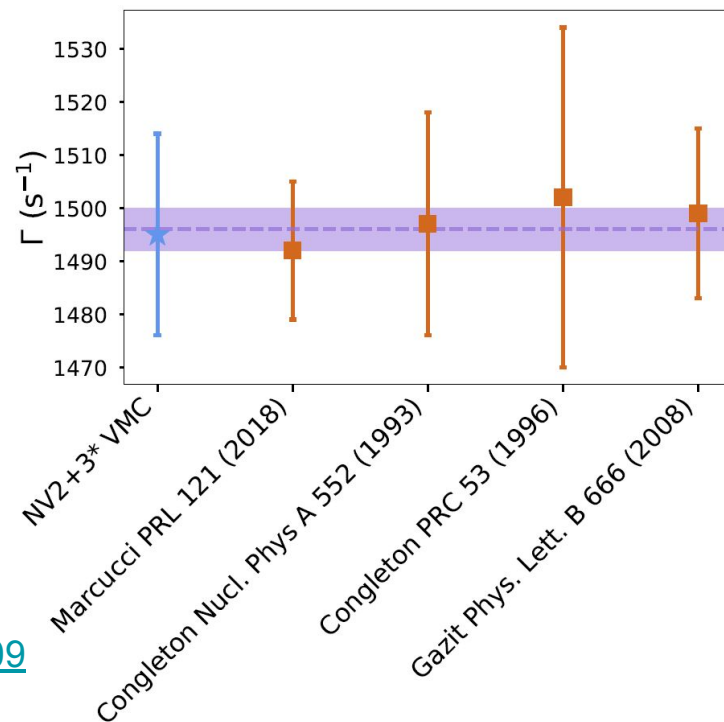
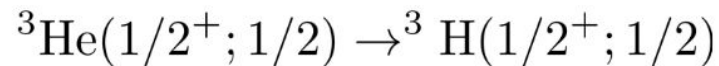
Ackerbauer *et al.* PLB417, 224(1998)

Momentum transfer $q \sim 100 \text{ MeV}$

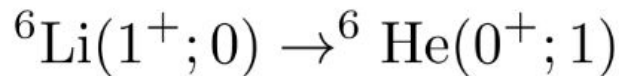
Two-body correction is $\sim 8\%$ of total rate on average for $A=3$

Garrett King *et al.* PRC2022

Review by Measday [Physics Reports 354 \(2001\) 243–409](#)



Partial muon capture rates: VMC calculations



$$\Gamma_{\text{VMC}}(\text{avg.}) = 1235 \text{ s}^{-1} \pm 101 \text{ s}^{-1}$$

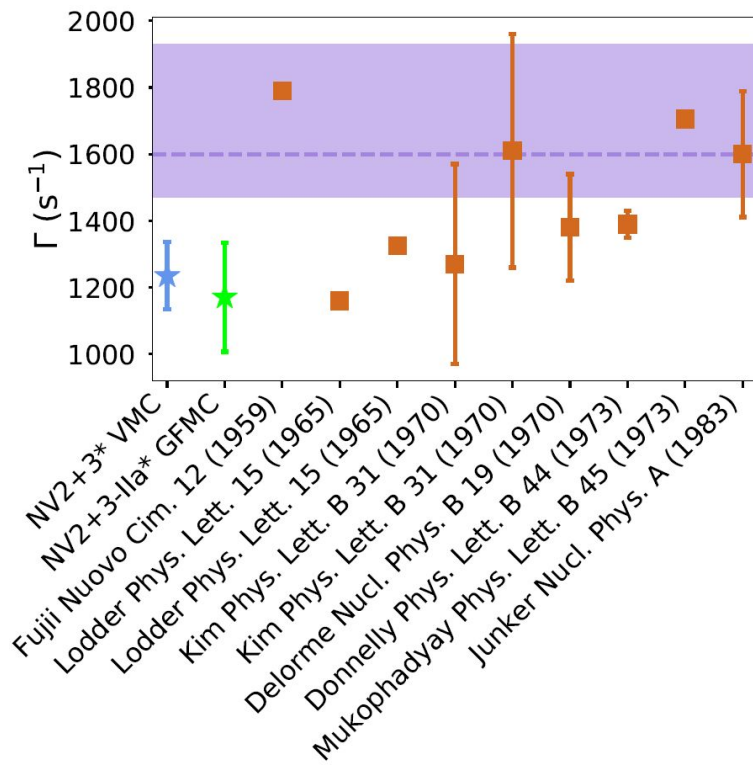
$$\Gamma_{\text{GFMC}}(\text{IIa}^*) = 1171 \text{ s}^{-1} \pm 164 \text{ s}^{-1}$$

$$\Gamma_{\text{expt}} = 1600 \text{ s}^{-1} +330/-129 \text{ s}^{-1}$$

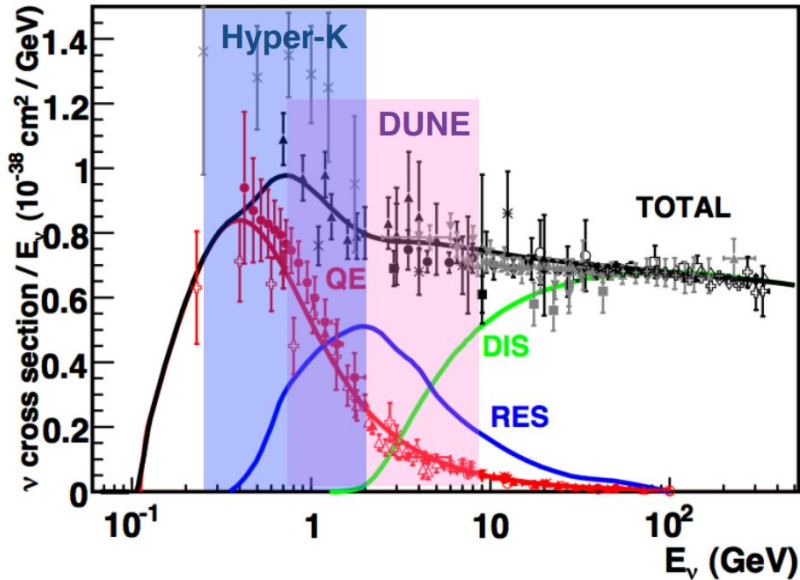
Deutsch *et al.* PLB26(1968)315

Garrett King *et al.* PRC2022

With FRIB experimentalist colleagues:
Gamow-Teller strength in A=11;
Schmitt *et al.* PRC106(2022)



Neutrino cross section anatomy



Formaggio & Zeller

Quasi-elastic: dominated by single-nucleon knockout

Resonance: excitation to nucleonic resonant states which decay into mesons

Deep-inelastic scattering: where the neutrino resolves the nucleonic quark content

Each of these regimes requires knowledge of both the **nuclear ground state** and the **electroweak coupling and propagation of the struck nucleons, hadrons, or partons**

A challenge for achieving precise neutrino-nucleus cross-section is **reliably bridging the transition regions which use different degrees of freedom**

Lepton-Nucleus scattering: Inclusive Processes

Electromagnetic Nuclear Response Functions

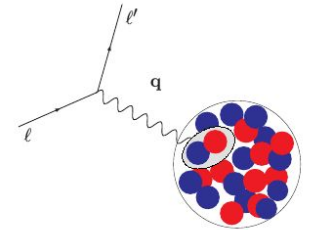
$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

Longitudinal response induced by the charge operator $O_L = \rho$

Transverse response induced by the current operator $O_T = \mathbf{j}$

5 Responses in neutrino-nucleus scattering

$$\frac{d^2 \sigma}{d\omega d\Omega} = \sigma_M [v_L R_L(\mathbf{q}, \omega) + v_T R_T(\mathbf{q}, \omega)]$$



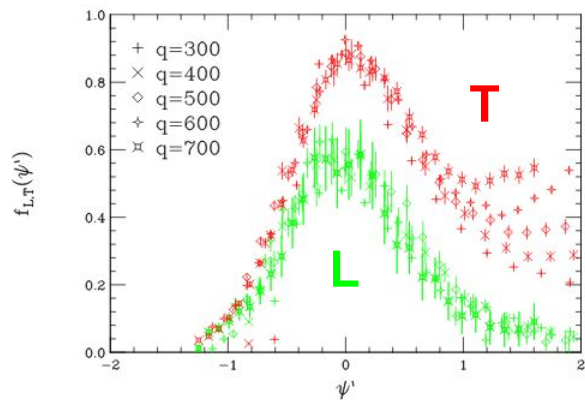
For a recent review on QMC, SF methods see

[Rocco Front. In Phys.8 \(2020\)116](#)

Lepton-Nucleus scattering: Data

Transverse Sum Rule

$$S_T(q) \propto \langle 0 | \mathbf{j}^\dagger \mathbf{j} | 0 \rangle \propto \langle 0 | \mathbf{j}_{1b}^\dagger \mathbf{j}_{1b} | 0 \rangle + \langle 0 | \mathbf{j}_{1b}^\dagger \mathbf{j}_{2b} | 0 \rangle + \dots$$



⁴He Electromagnetic Data
Carlson *et al.* PRC65(2002)024002

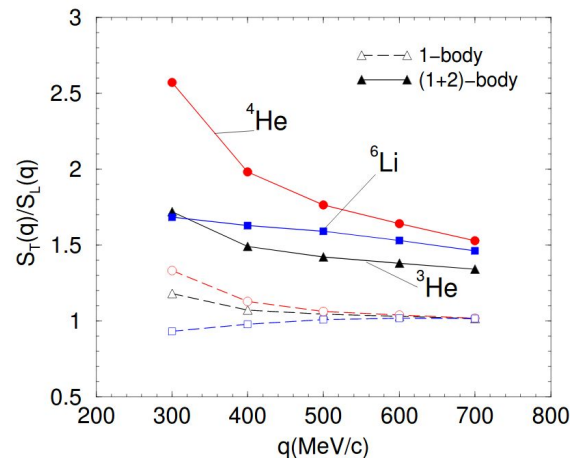
Observed transverse enhancement explained by the combined effect of two-body correlations and currents in the interference term

$$\langle \mathbf{j}_{1b}^\dagger \mathbf{j}_{1b} \rangle > 0$$

Leading one-body term

$$\langle \mathbf{j}_{1b}^\dagger \mathbf{j}_{2b} v_\pi \rangle \propto \langle v_\pi^2 \rangle > 0$$

Interference term

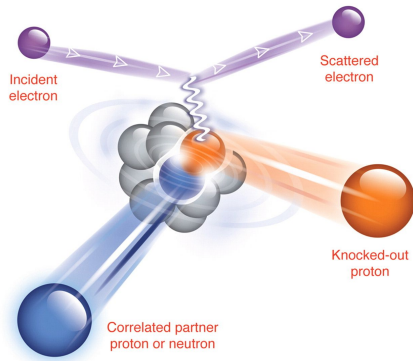


Transverse/Longitudinal Sum Rule
Carlson *et al.* PRC65(2002)024002

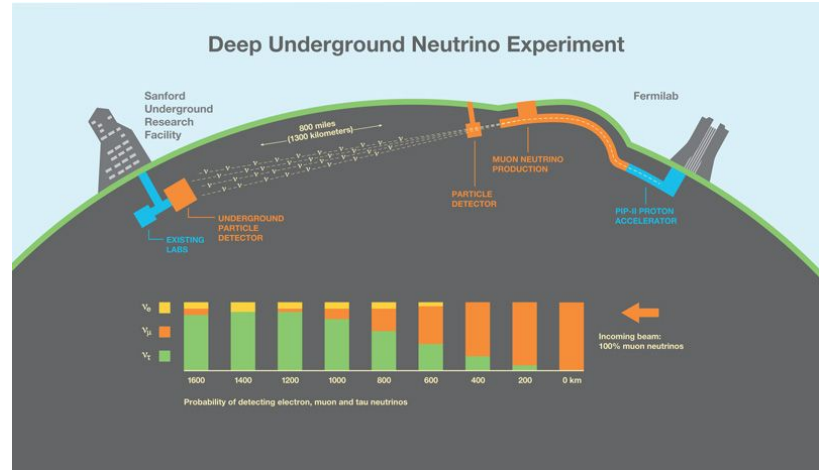
Beyond Inclusive: Short-Time-Approximation

Short-Time-Approximation Goals:

- Describe electroweak scattering from $A > 12$ without losing two-body physics
- Account for exclusive processes
- Incorporate relativistic effects



Subedi et al. Science320(2008)1475



[Stanford Lab article](#)

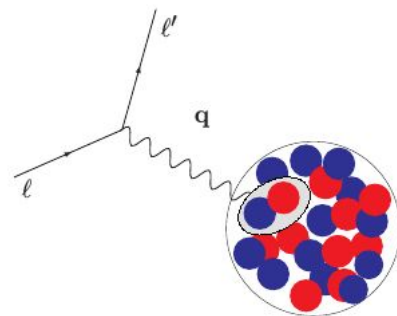
[e4u collaboration](#)



Short-Time-Approximation

Short-Time-Approximation:

- Based on Factorization
- Retains two-body physics
- Correctly accounts for **interference**

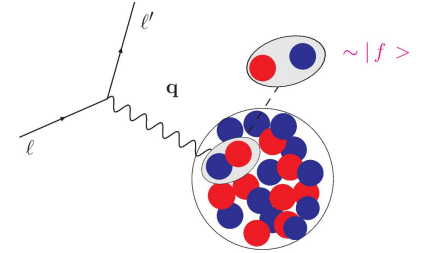


$$R(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_0)t} \langle 0 | O^\dagger e^{-iHt} O | 0 \rangle$$

$$O_i^\dagger e^{-iHt} O_i + O_i^\dagger e^{-iHt} O_j + O_i^\dagger e^{-iHt} O_{ij} + O_{ij}^\dagger e^{-iHt} O_{ij}$$

$$H \sim \sum_i t_i + \sum_{i < j} v_{ij}$$

Short-Time-Approximation



Short-Time-Approximation:

- Based on Factorization
- **Retains two-body physics**
- Response functions are given by the **scattering from pairs of fully interacting nucleons** that propagate into a correlated pair of nucleons
- Allows to retain both two-body correlations and currents at the vertex
- Provides “more” exclusive information in terms of nucleon-pair kinematics via the Response Densities

Response Functions \propto Cross Sections

$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

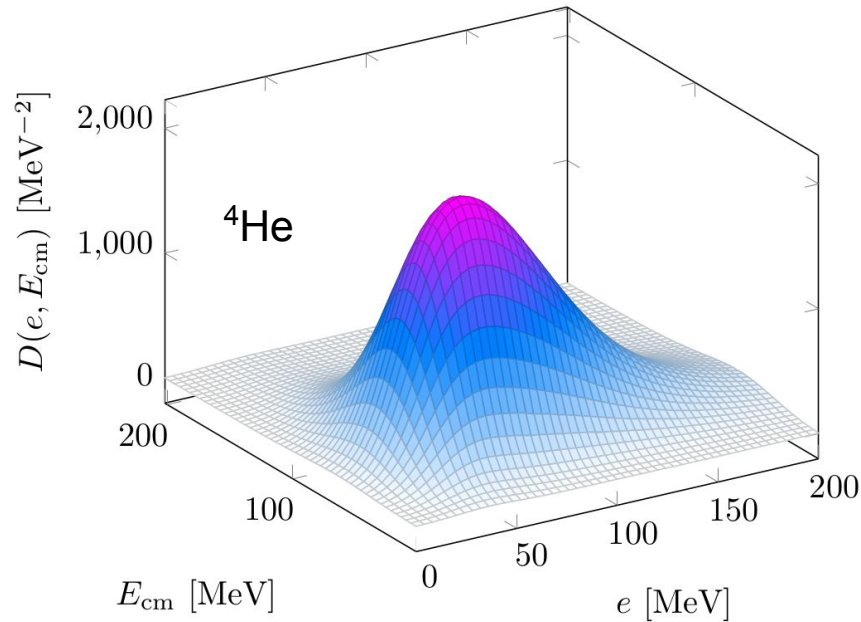
Response **Densities**

$$R(q, \omega) \sim \int \delta(\omega + E_0 - E_f) dP' dp' \mathcal{D}(p', P'; q)$$

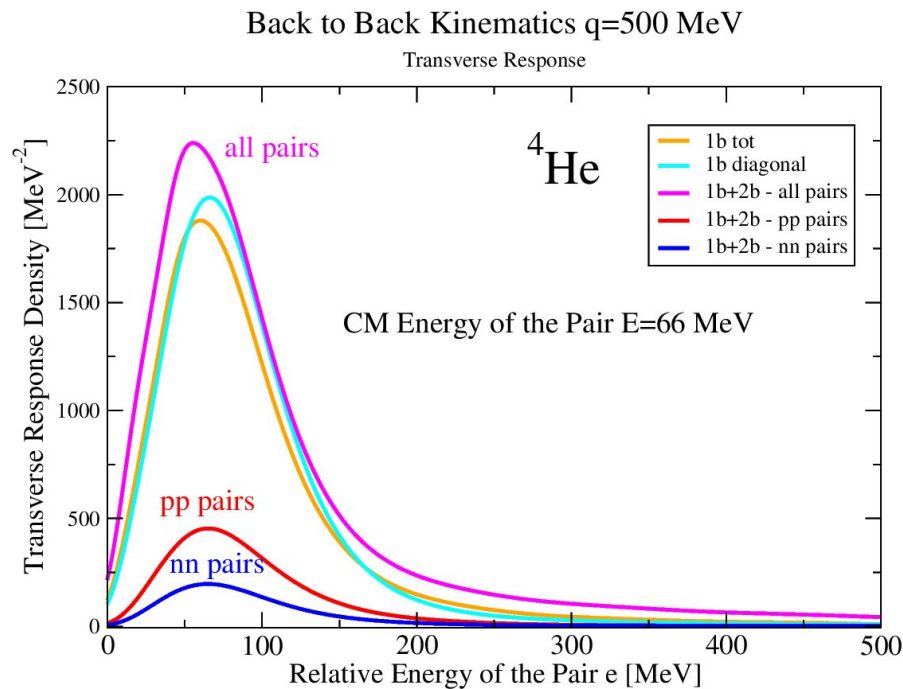
P' and p' are the CM and relative momenta of the struck nucleon pair

Transverse Response Density: e - ${}^4\text{He}$ scattering

Transverse Density $q = 500 \text{ MeV}/c$

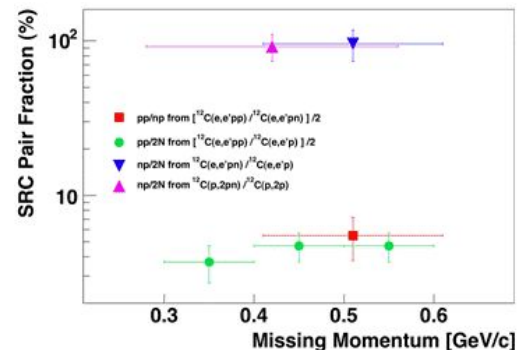


$e^{-4}\text{He}$ scattering in the back-to-back kinematic



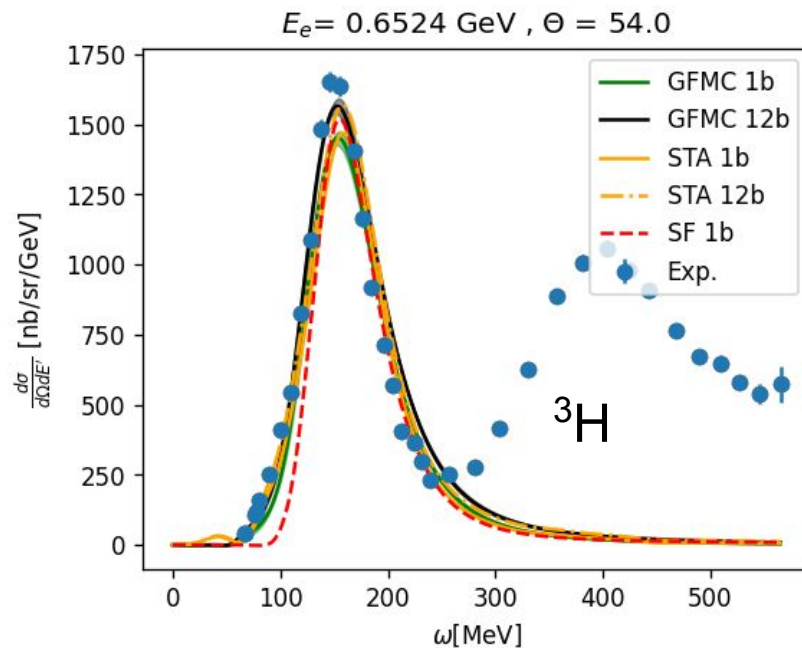
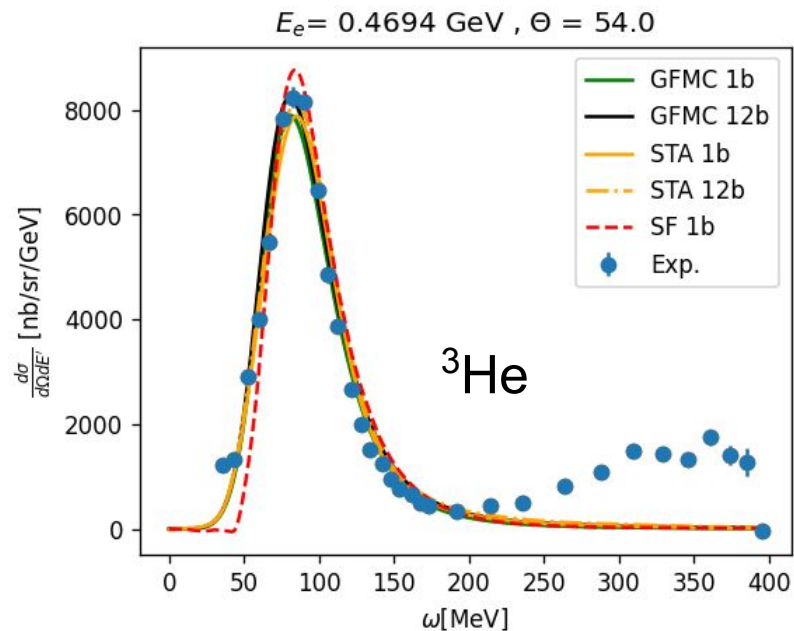
SP *et al.* PRC101(2020)044612

- pp pairs
- nn pairs
- all pairs 1body
- all pairs tot



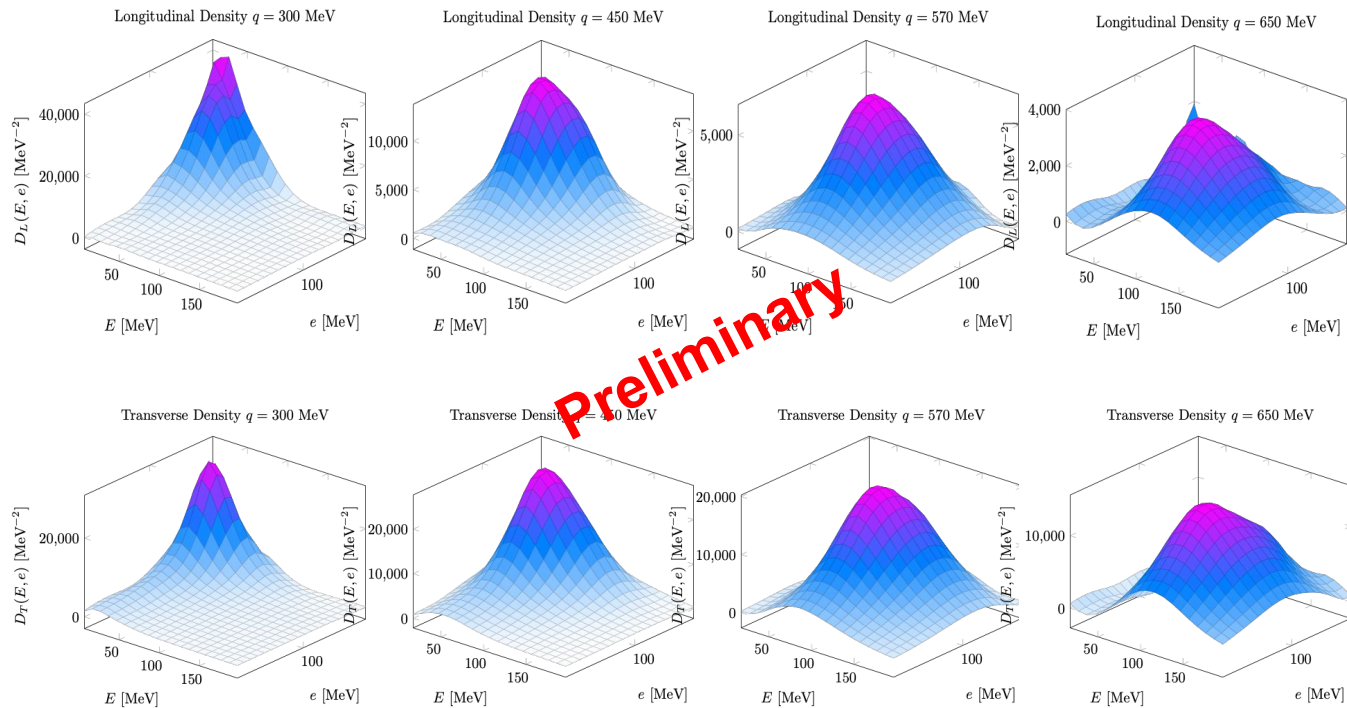
Subedi *et al.* Science320(2008)1475

GFMC SF STA: Benchmark & error estimate

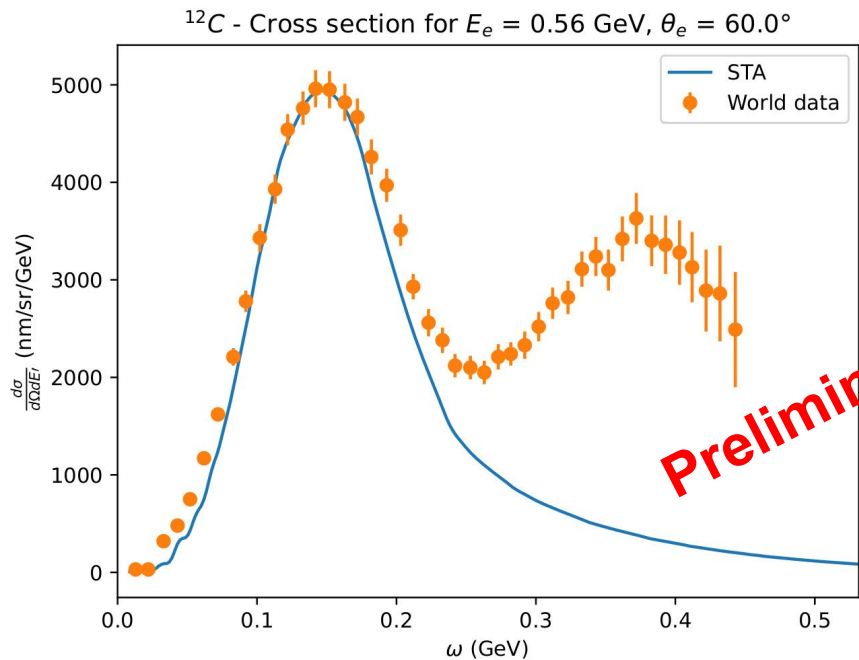


Lorenzo Andreoli, et al. PRC 2021

^{12}C Response Densities



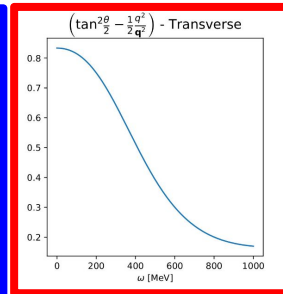
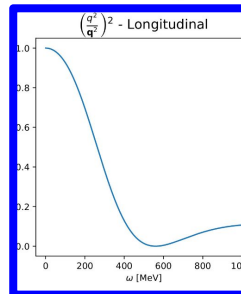
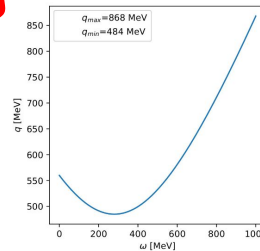
^{12}C cross sections



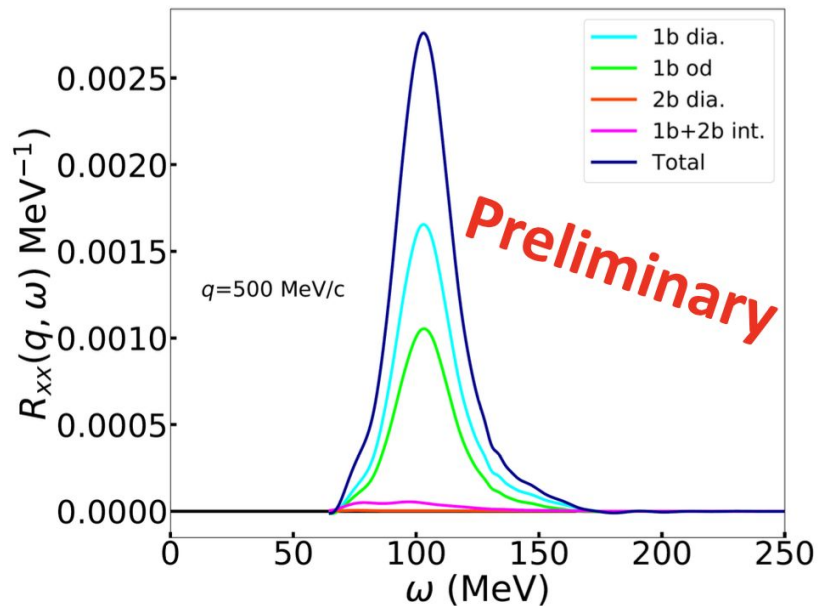
Preliminary

$$\left(\frac{d^2\sigma}{dE'd\Omega'}\right)_e = \left(\frac{d\sigma}{d\Omega'}\right)_M \left[\left(\frac{q^2}{Q^2}\right)^2 F_L(|\mathbf{q}|, \omega) + \left(\tan^2\frac{\theta}{2} - \frac{1}{2}\frac{q^2}{Q^2}\right) R_T(|\mathbf{q}|, \omega)\right]$$

$E_e = 560 \text{ MeV}$, $\theta_e = 60^\circ$



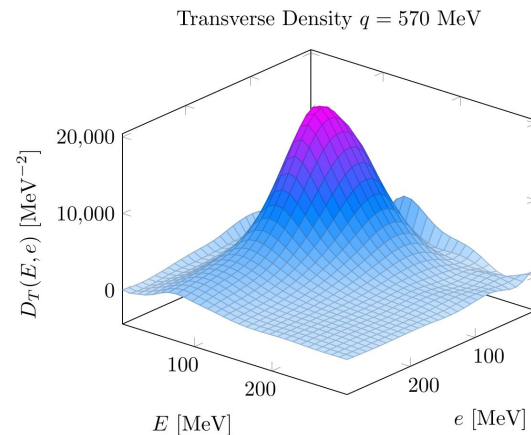
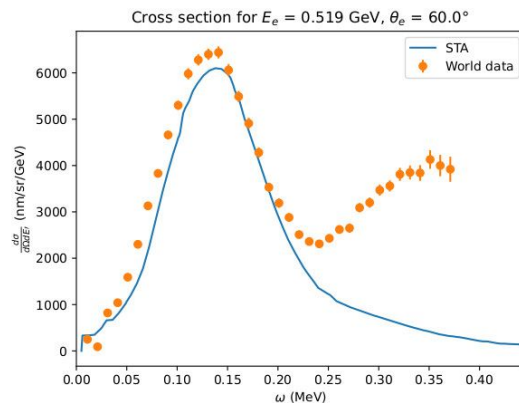
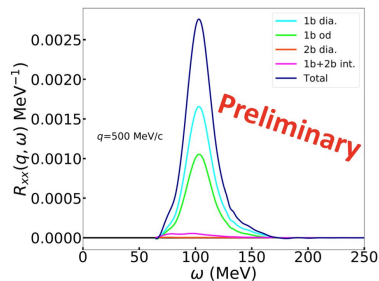
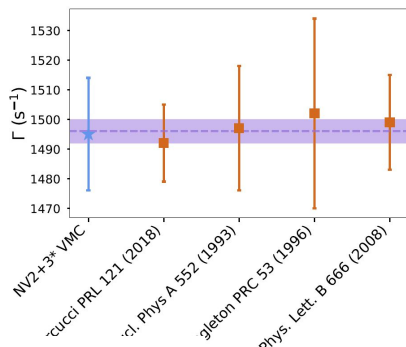
NC processes on deuteron with STA



Garrett King *et al.* in preparation

Summary

Ab initio calculations of light nuclei yield a picture of nuclear structure and dynamics where **many-body effects play an essential role to explain available data.**



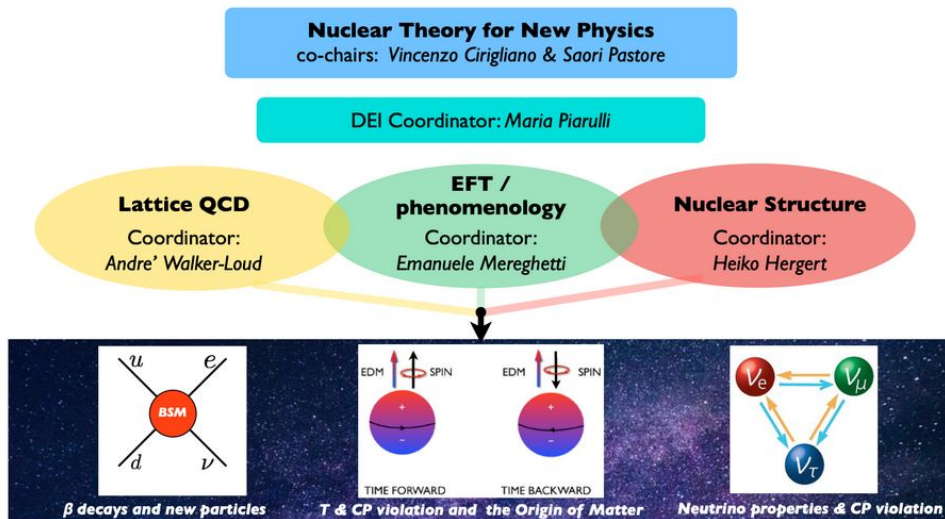
Close **c**ollaborations between **NP, LQCD, Pheno, Hep, Comp, Expt, ...** are required to progress e.g., NP is represented in the Snowmass process

It's a very exciting time!

Nuclear Theory for New Physics NP&HEP TC

Nuclear Theory for New Physics

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- [Funding Acknowledgement](#)



Snowmass:
Topical groups and
Frontier Reports,
Whitepapers, ...

LRP:
White papers,
[2301.03975](#), [FSNN](#),
...

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[NTNP](#)

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Huzhou U: Dong Wang

Fermilab: Gardiner Betancourt

MIT: Barrow

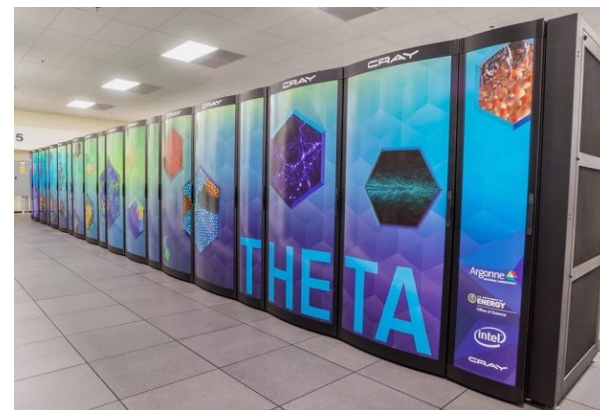


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Quantum Monte Carlo Group for Nuclear Physics

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