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Isovector and isoscalar spin-multipole giant resonances in the parent and daughter nuclei of double-beta decay triplets

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MAYORANA School

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Motivation

- Neutrinoless double-beta decay is one of the most interesting subjects in nuclear and particle physics.
- If it were detected, it would mean new physics beyond the standard model.
- By studying the spin-multipole giant resonances in mother and daughter double-beta decay nuclei, one could obtain more information about the wave functions of the virtual state of the intermediate nuclei.



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- The poster will present a (semi) recent publication.

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Isvector and isoscalar spin-multipole giant resonances in the parent and daughter nuclei of double- β -decay triplets

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Intro

- The study of double-beta decay nuclei and their daughter nuclei.
- The strength functions of isovector and isoscalar spin-multipole giant resonances were studied.
- Used theory was quasiparticle random-phase approximation (QRPA).
- Transitions from the ground state to the excited states were calculated.
 - $J^\pi = 0^-, 1^-, 2^-$ ($L = 1$, spin-dipole)
 - $J^\pi = 1^+, 2^+, 3^+$ ($L = 2$, spin-quadrupole)

Studied nuclei:

(^{76}Ge , ^{76}Se)

(^{82}Se , ^{82}Kr)

(^{96}Zr , ^{96}Mo)

(^{100}Mo , ^{100}Ru)

(^{116}Cd , ^{116}Sn)

(^{128}Te , ^{128}Xe)

(^{130}Te , ^{130}Xe)

(^{136}Xe , ^{136}Ba)



Theory (QRPA)

- Nucleons pair with each other
 - Protons with protons and neutrons with neutrons
- Nucleon pairs are handled as quasiparticles that behave like bosons
- QRPA excitation operator can be written as

$$Q^\dagger = \sum_{a \leq b} [X_{ab}^\omega A_{ab}^\dagger(JM) - Y_{ab}^\omega \tilde{A}_{ab}(JM)]$$

- Two-quasiparticle operators:

$$A_{ab}^\dagger(JM) = \mathcal{N}_{ab}(J)[a_a^\dagger a_b^\dagger]_{JM} \quad \tilde{A}_{ab}(JM) = -\mathcal{N}_{ab}(J)[\tilde{a}_a \tilde{a}_b]_{JM}$$



Theory

- Transition operators for isovector and isoscalar excitations:

$$\mathcal{O}_{L,JM}^{0,v} = i^l r^l [Y_L \sigma]_{JM} t_0 \quad \mathcal{O}_{L,JM}^{0,s} = i^l r^l [Y_L \sigma]_{JM}$$

- Are used to obtain reduced transition nuclear matrix elements
- Transition strengths are calculated as the square of the transition NME:

$$S_{nJ^\pi}^L = |(nJ^\pi || \mathcal{O}_{L,J}^0 || \text{QRPA})|^2$$



Example of the results

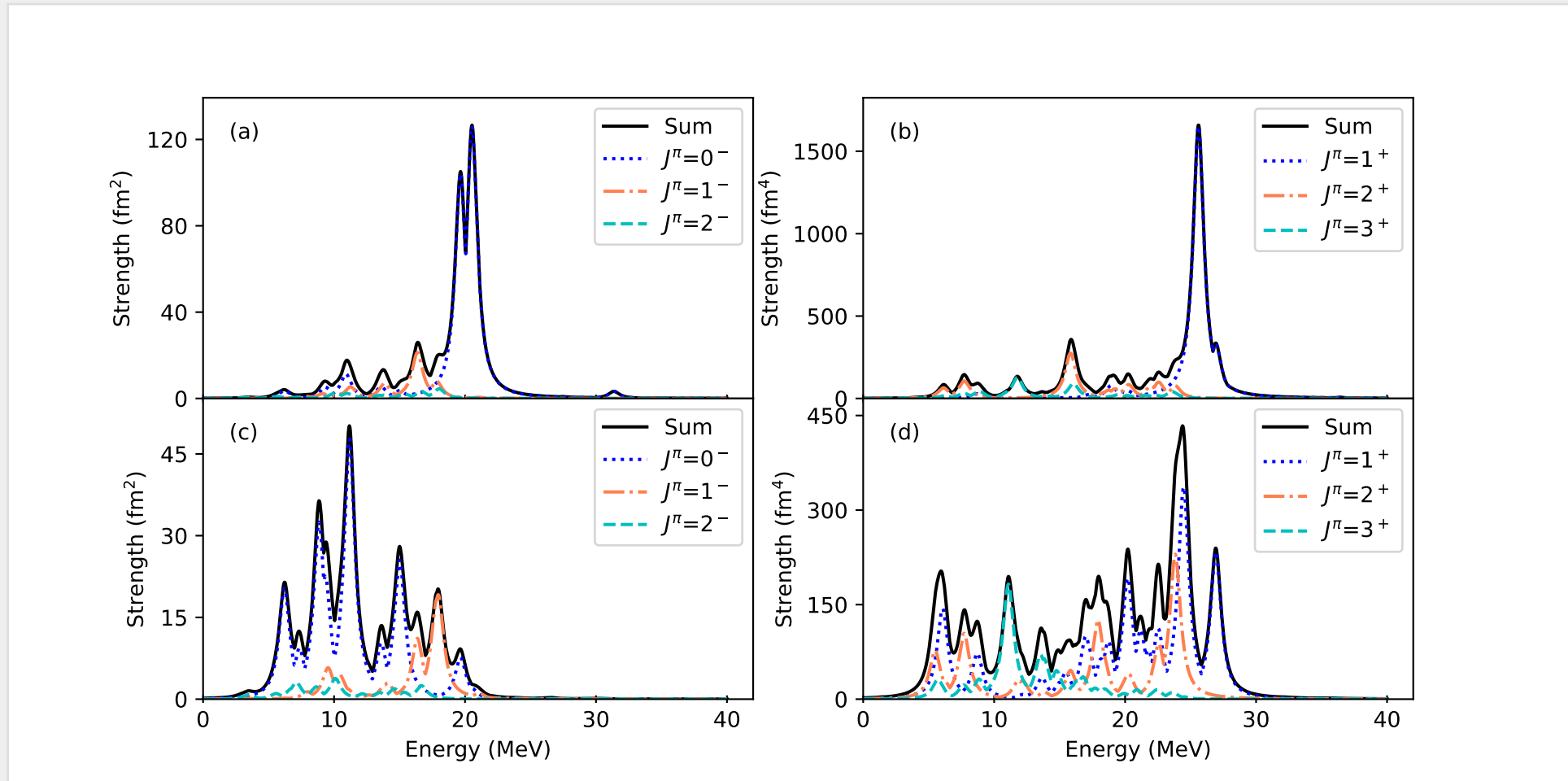
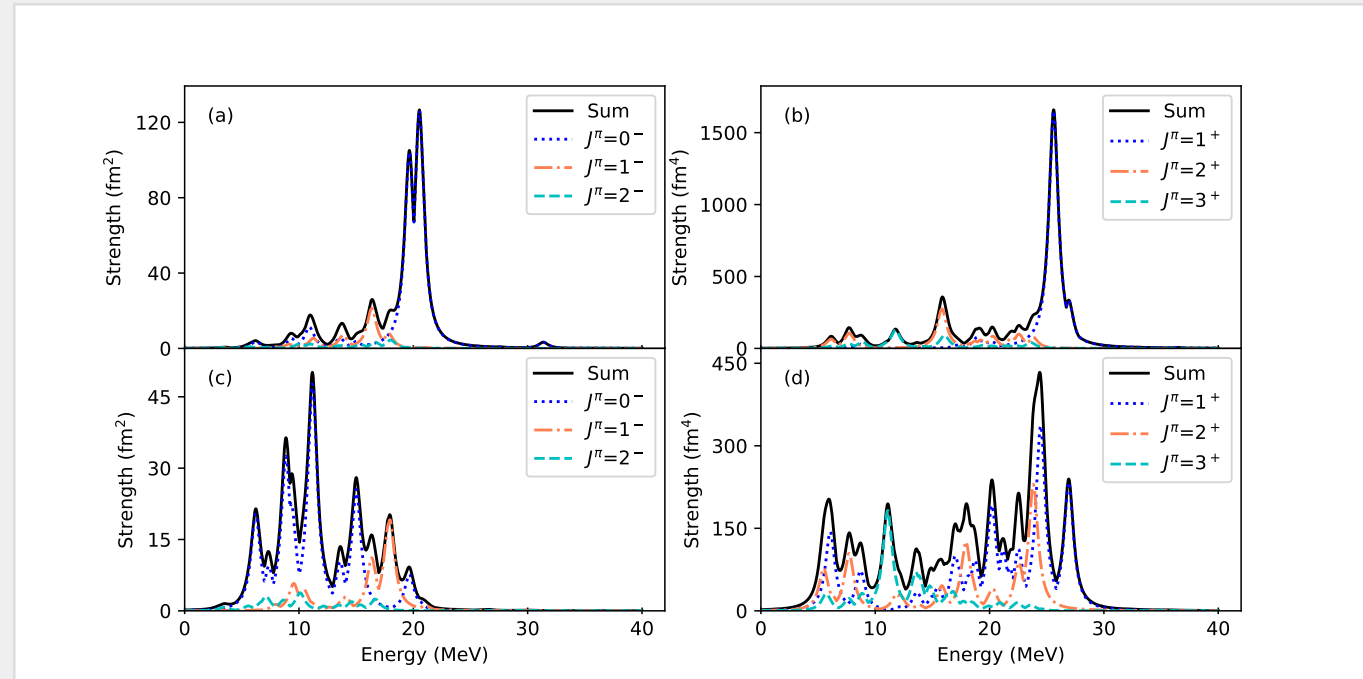


Figure: Isovector and isoscalar strength functions for ^{136}Xe . (a) isovector, $L = 1$, (b) isovector $L = 2$, (c) isoscalar $L = 1$, and (d) isoscalar $L = 2$.



Observations of the obtained results

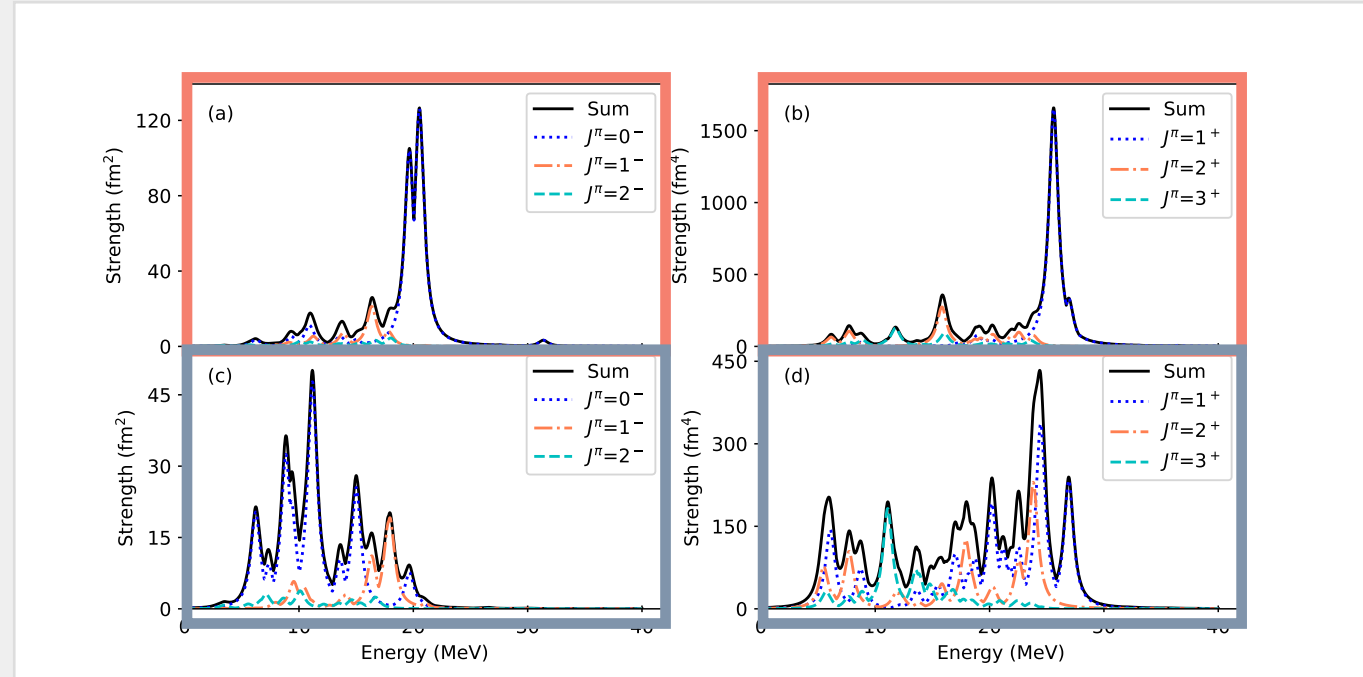
- How widely the peaks are spread depends on the type of the excitation.
- Isoscalar are more spread in energy.
- Isovector strength tend to locate in one or few peaks.





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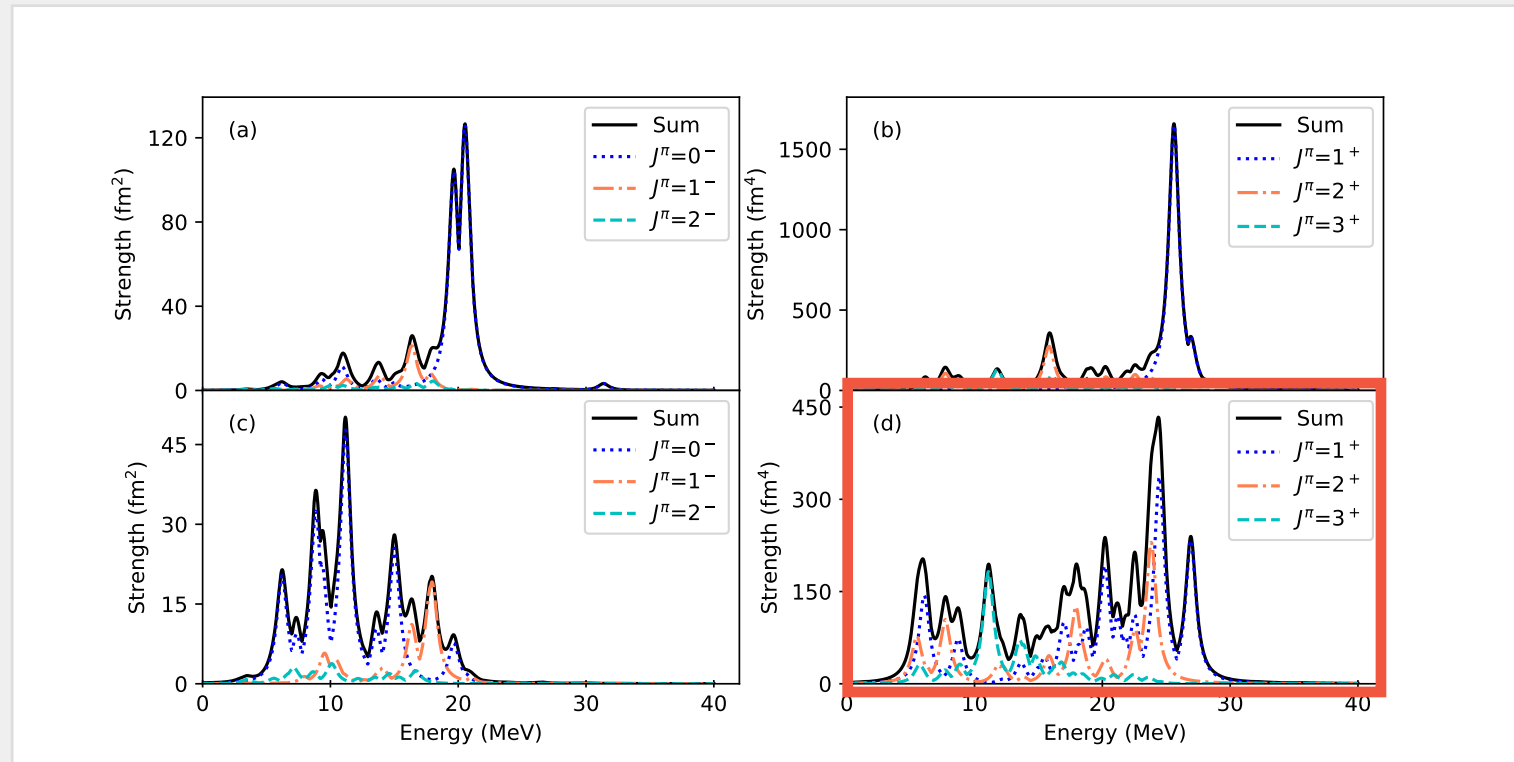
- How widely the peaks are spread depends on the type of the excitation.
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- **Isvector** strength tend to locate in one or few peaks.





Observations of the obtained results

- Only isoscalar strength functions show notable contributions for all the states.





Conclusions

- Comparison of the results obtained in this study with potential future experimental data may shed light on the reliability of the QRPA-based framework describing the wave functions of nuclear states relevant to two-neutrino and neutrinoless double-beta decay.
- More on the poster!



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Thank you!