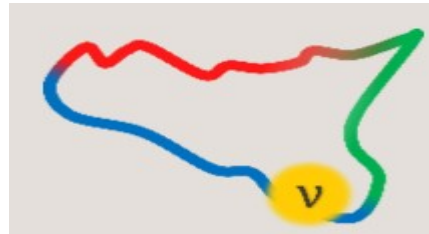


# HIGH-ENERGY ASTROPHYSICAL NEUTRINOS

**Esteban Roulet**  
(CONICET, Bariloche, Argentina)

**MAYORANA School, Modica 2023**

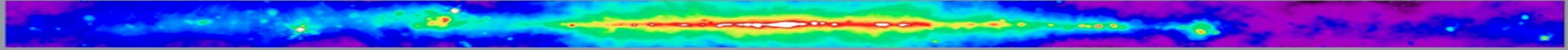


# MULTIWAVELENGTH ELECTROMAGNETIC ASTRONOMY

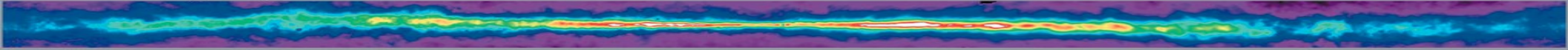
Traditionally, the sky has been observed in electromagnetic radiation  
different wavelengths lead to different insights



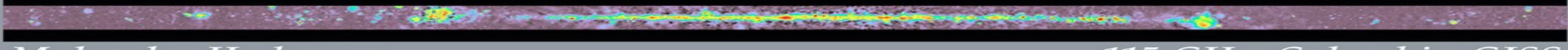
*Radio Continuum* 408 MHz Bonn, Jodrell Banks, & Parkes



*Atomic Hydrogen* 21 cm Leiden-Dwingeloo, Maryland-Parkes



*Radio Continuum* 2.4-2.7 GHz Bonn & Parkes



*Molecular Hydrogen* 115 GHz Columbia-GISS



*Infrared* 12, 60, 100  $\mu\text{m}$  IRAS



*Near Infrared* 1.25, 2.2, 3.5  $\mu\text{m}$  COBE/DIRBE



*Optical* Laustsen et al. Photomosaic



*X-Ray* 0.25, 0.75, 1.5 keV ROSAT/PSPC



*Gamma Ray* >100 MeV CGRO/EGRET



# multiwavelength M87

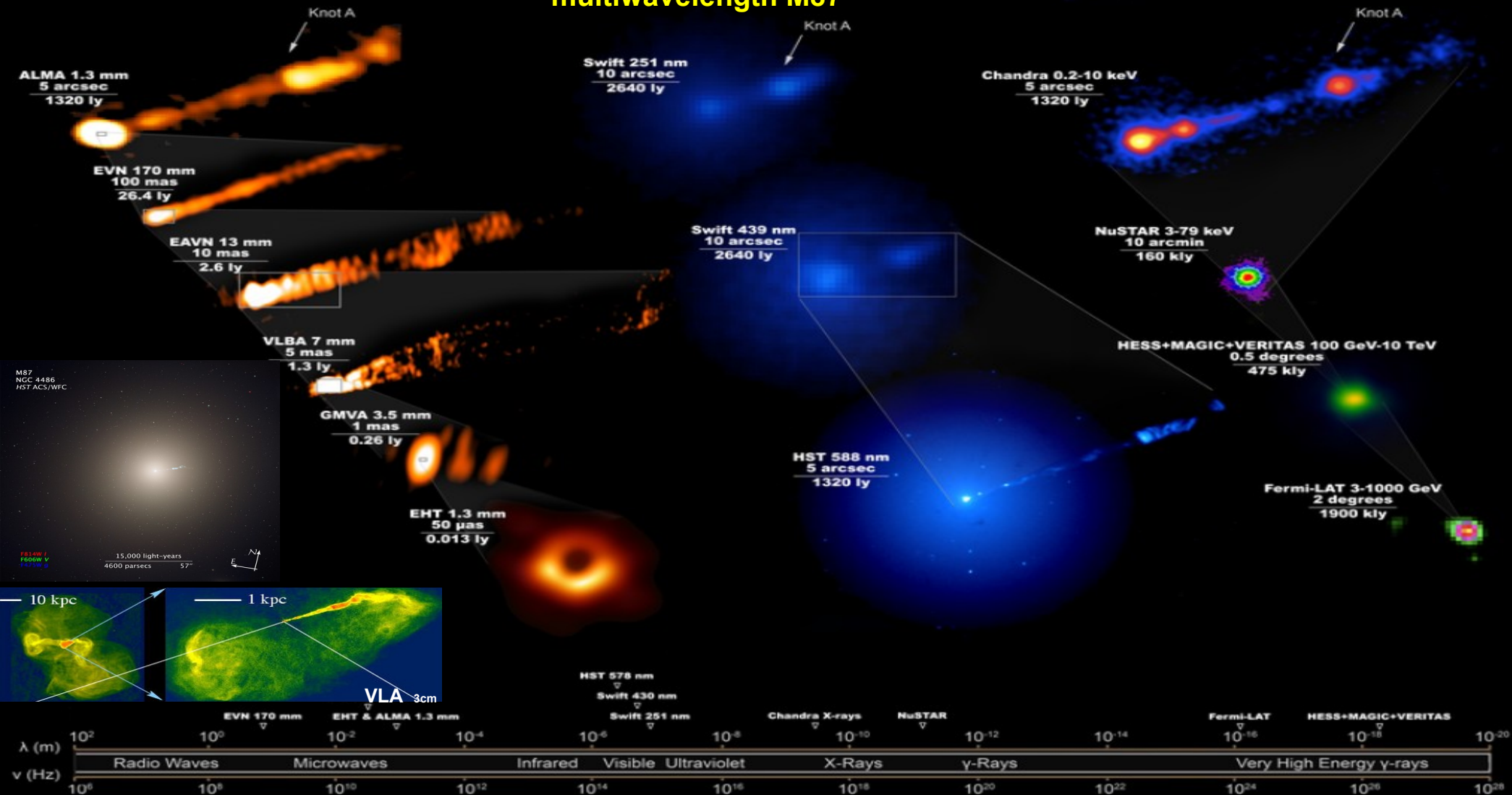
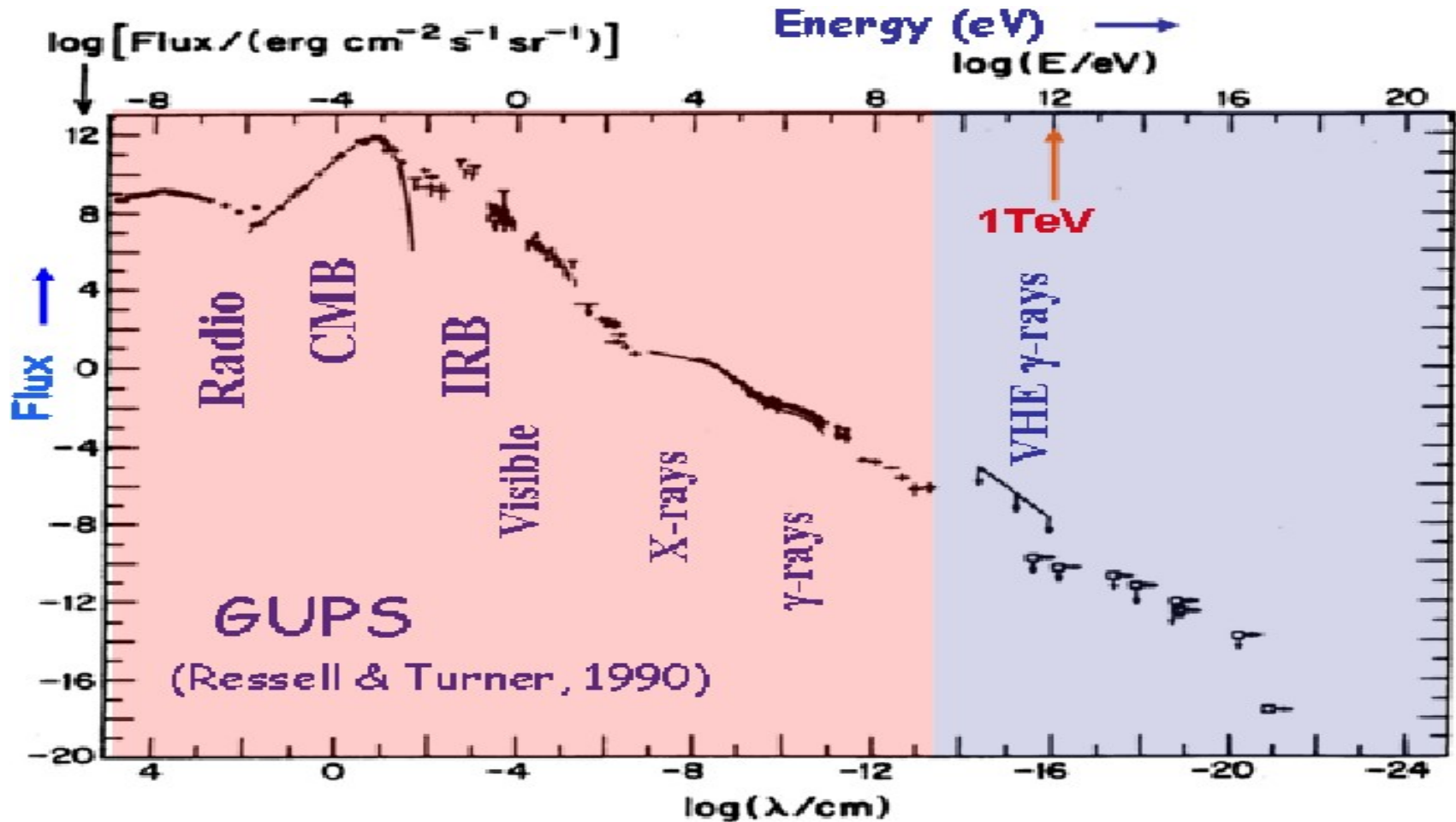


Image Credit: The EHT Multi-wavelength Science Working Group; the EHT Collaboration; ALMA (ESO/NAOJ/NRAO); the EVN; the EAVN Collaboration; VLBA (NRAO); the GMVA; the Hubble Space Telescope; the Neil Gehrels Swift Observatory; the Chandra X-ray Observatory; the Nuclear Spectroscopic Telescope Array; the Fermi-LAT Collaboration; the H.E.S.S. collaboration; the MAGIC collaboration; the VERITAS collaboration; NASA and ESA. Composition by J. C. Alagaba

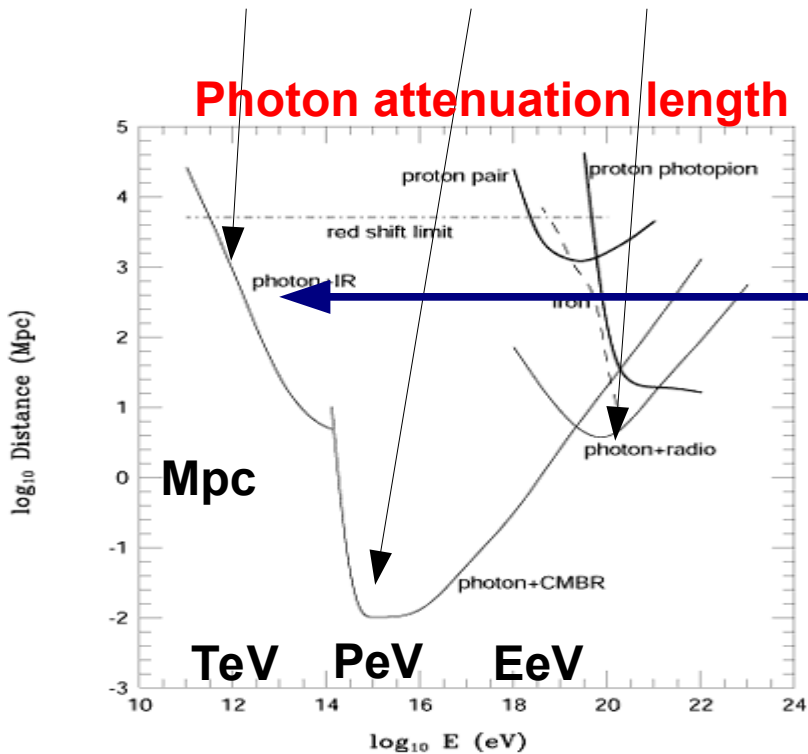
# DIFFUSE PHOTON BACKGROUND



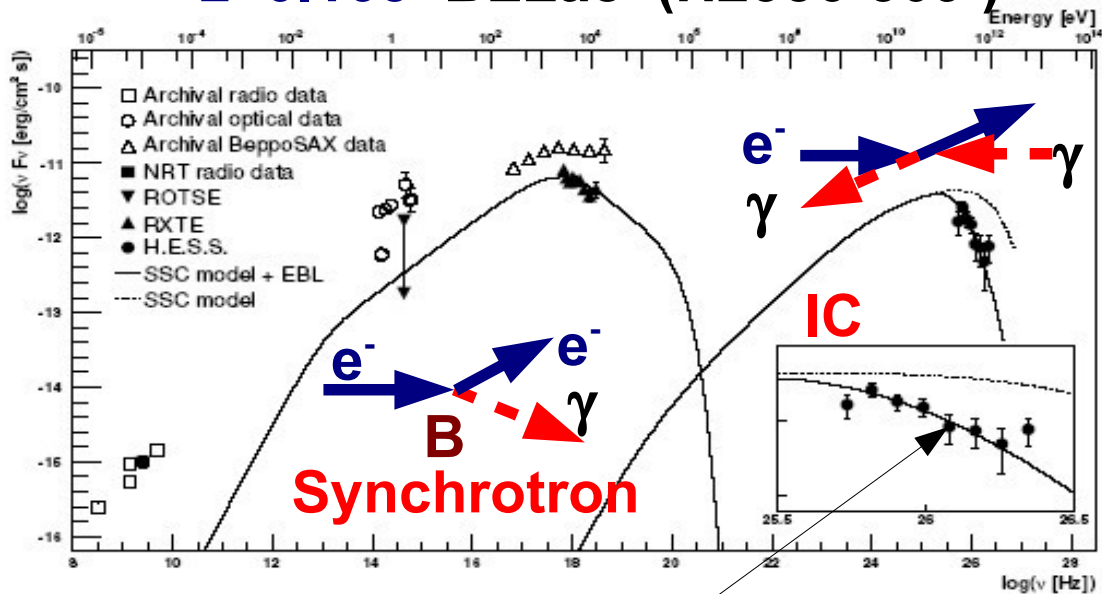
the Universe is filled with background radiation

# Distant $\gamma$ sources are strongly attenuated by background photons (starlight, CMB, radio, ...): $\gamma\gamma \rightarrow e^+e^-$

(  $\text{MeV}^2 = \text{eV TeV} = \text{meV PeV}$  )



**$z=0.165$  BLLac (H2356-309)**



can even measure IR background from observed attenuation

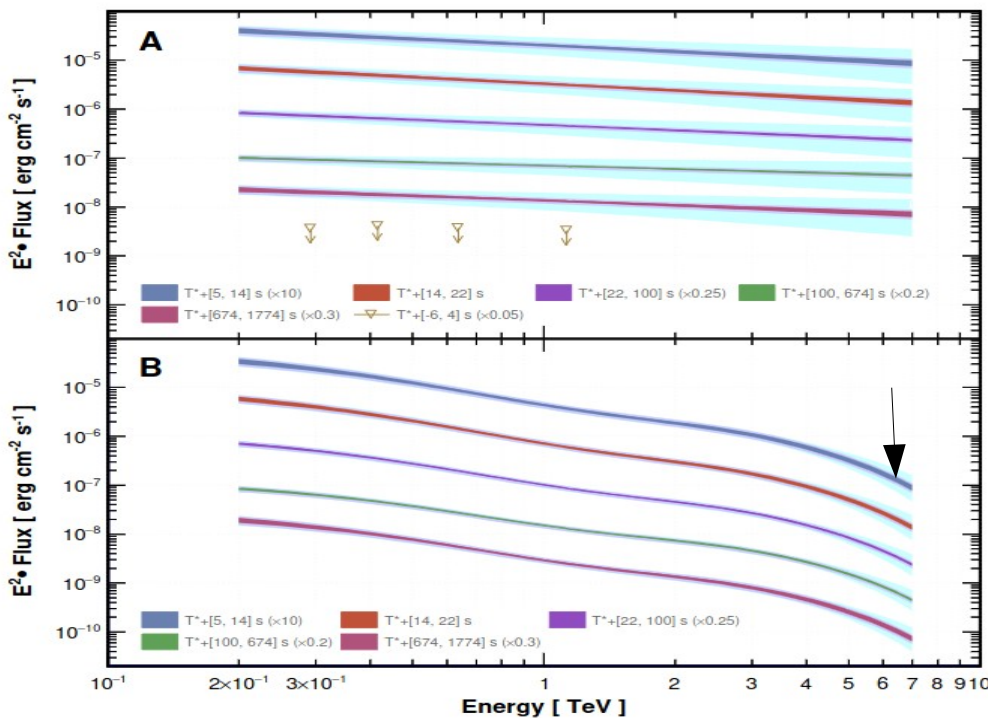
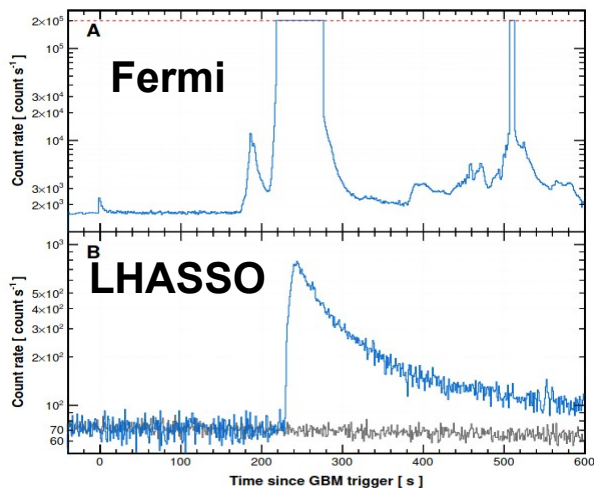
**beyond few TeV, high redshift Universe is unobservable with photons**

# GRB 221009A Observed by LHAASO



redshift  $z=0.151 \rightarrow$  distance  $\sim 740$  Mpc

2306.06372

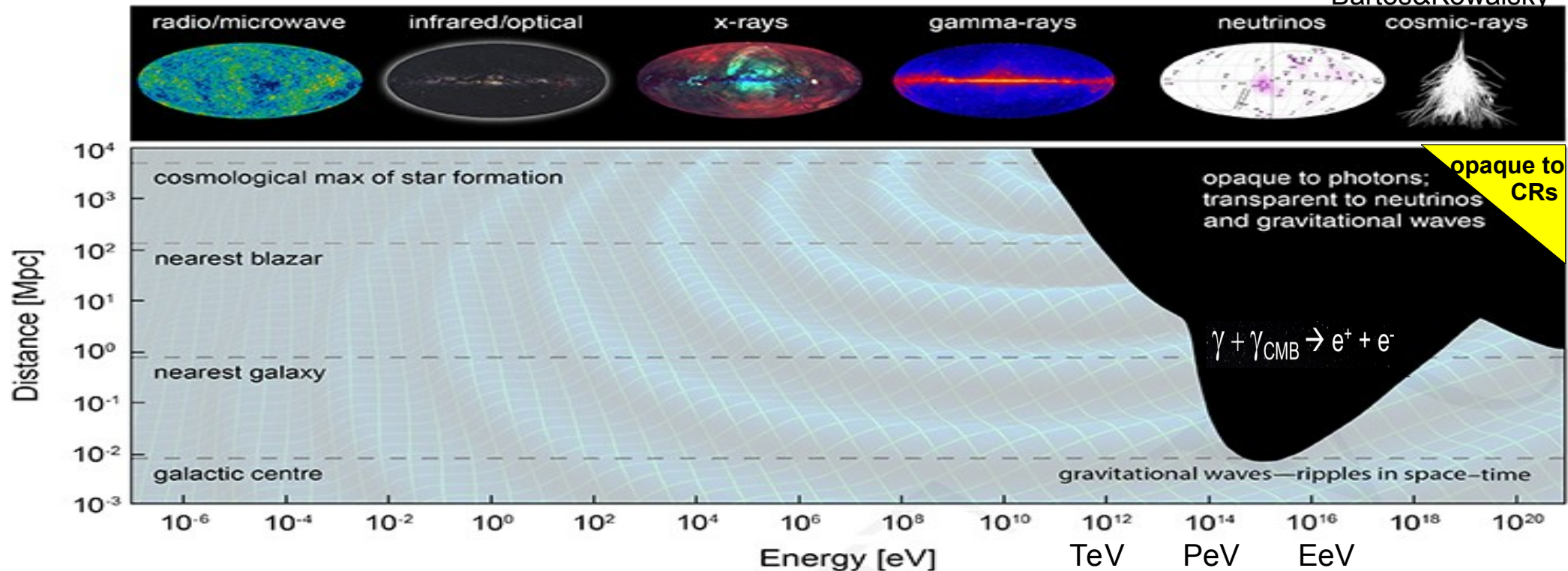


re 2: Intrinsic and observed flux spectra for five time-intervals. (A) the intrinsic spectra corrected

moreover, absence of TeV emission in initial X-ray burst due to  $\gamma$  attenuation in dense inner jet?  
 TeV burst in external afterglow

# beyond TeV, photons get absorbed by radiation backgrounds

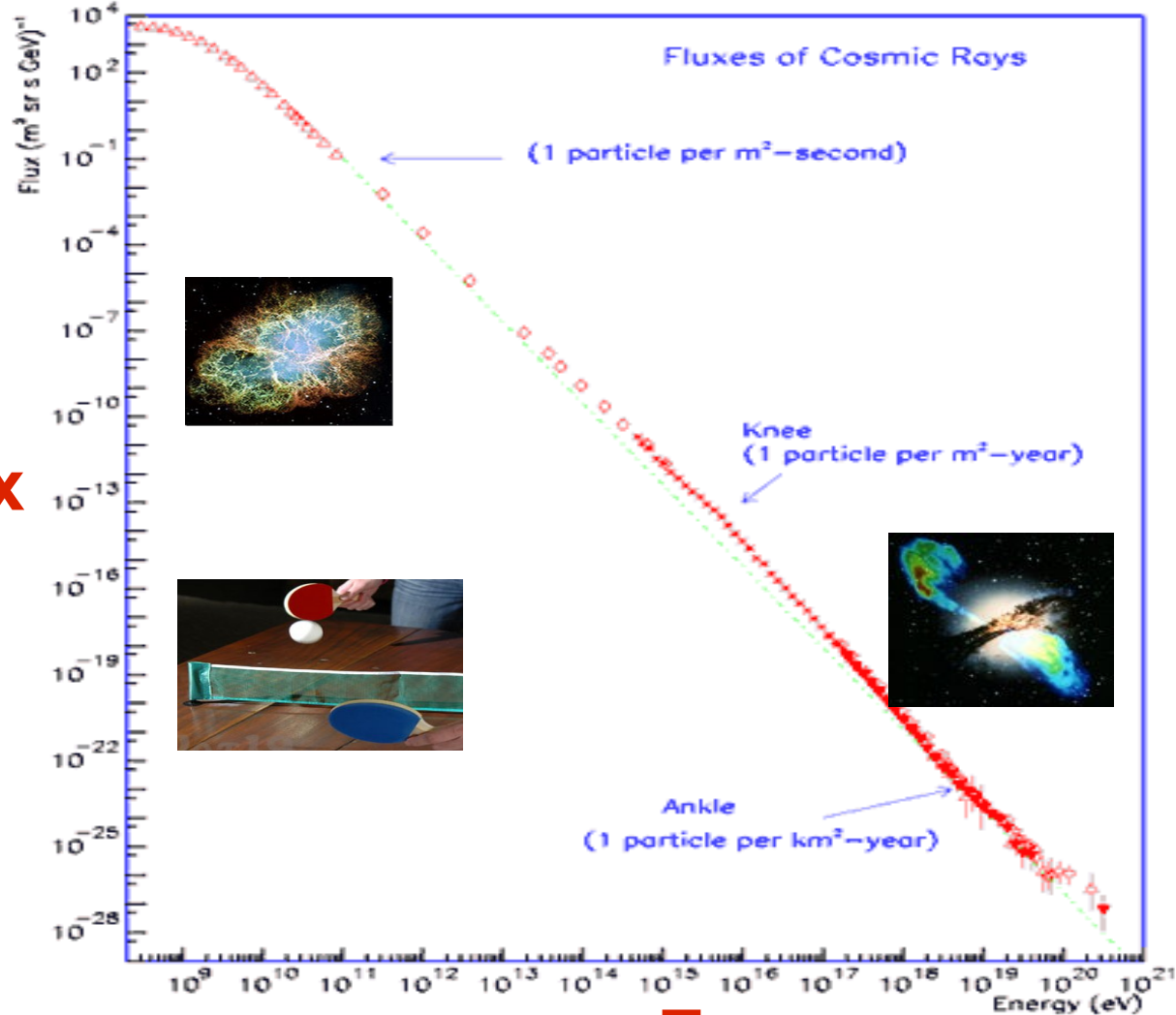
Bartos&Kowalsky



beyond 100 TeV, only Galactic sources visible in photons  
sources at more than 100 Mpc (e.g. blazars) visible up to few TeV  
beyond those energies only neutrinos and cosmic rays can reach us  
moreover, while photons emitted from star photospheres,  
neutrinos come from deep inside stars and supernovae  
multimessenger searches now include also gravitational waves!

# COSMIC RAY SPECTRUM

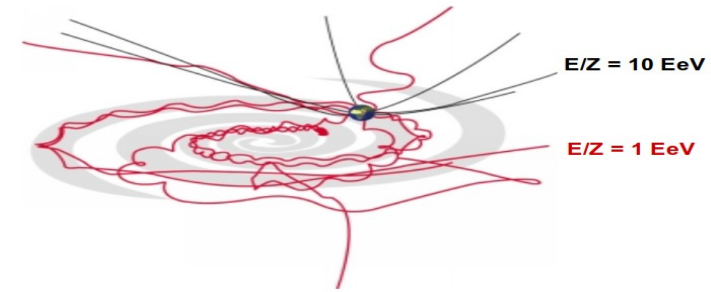
flux



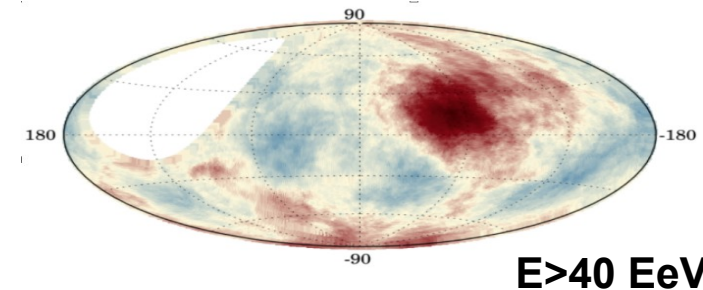
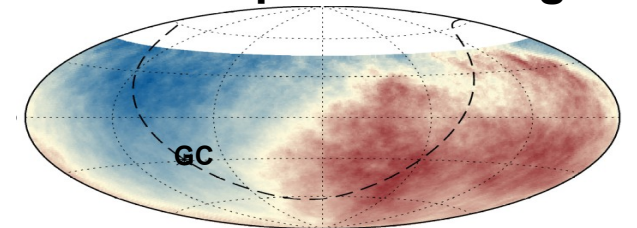
Energy

# CR ANISOTROPIES

but CRs are charged particles

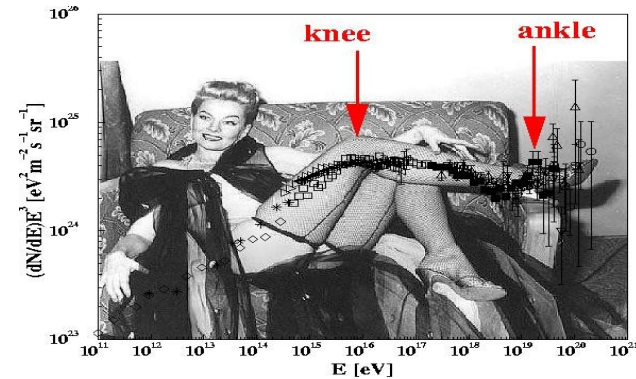
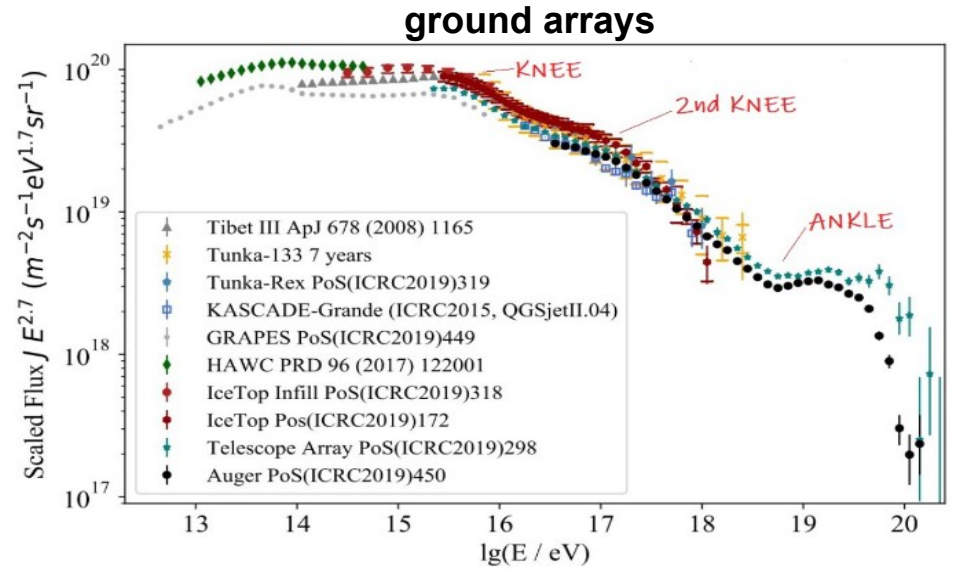
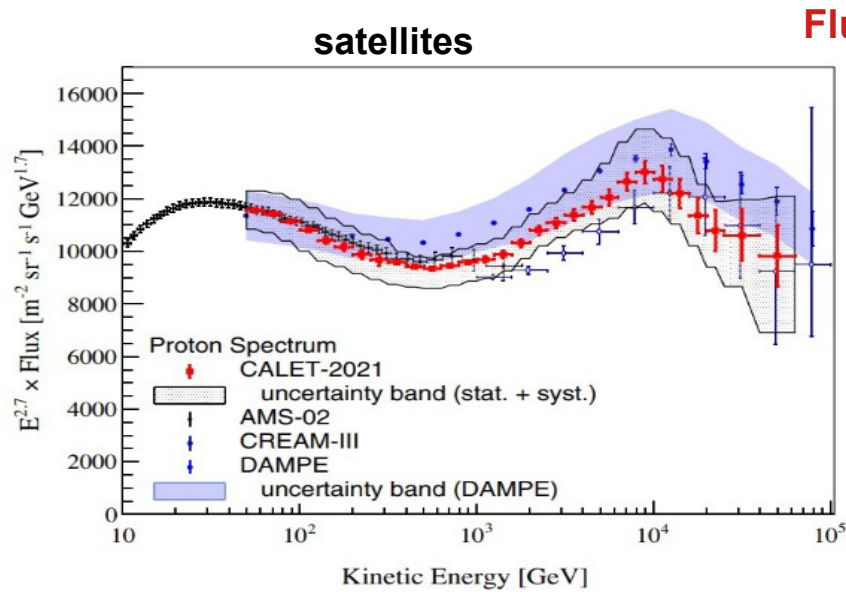


anisotropies with Auger



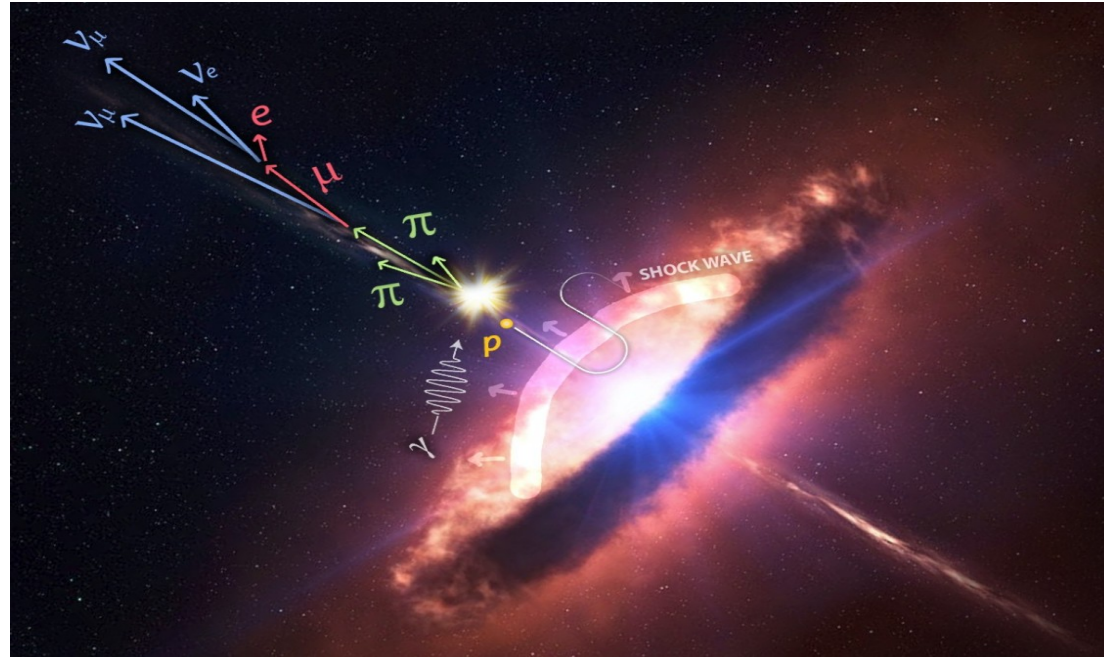


# CR spectrum actually not so featureless



## NEUTRINO ASTRONOMY AT VERY HIGH ENERGIES:

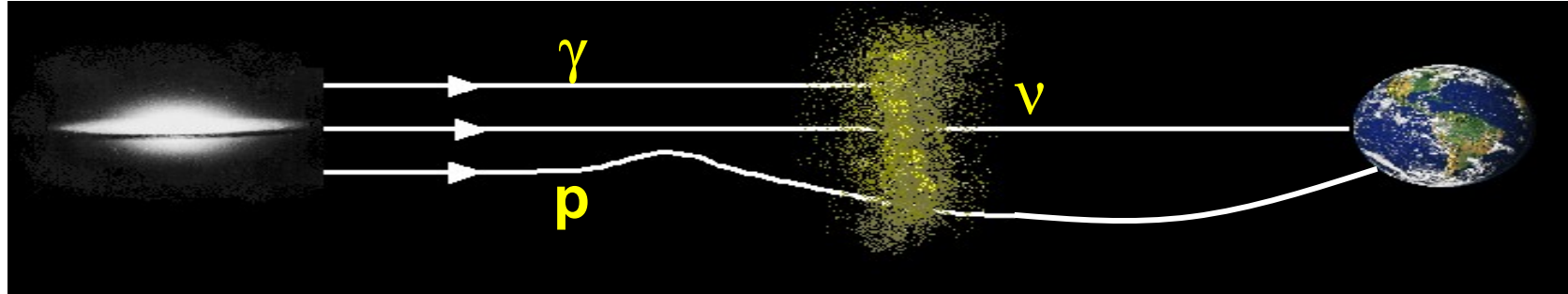
- NEUTRINOS PROPAGATE STRAIGHT AND UNATTENUATED
- NEUTRINOS GET PRODUCED IN CR SOURCES



neutrinos may tell us where CRs get accelerated,  
and they may also help us to learn a lot about their sources

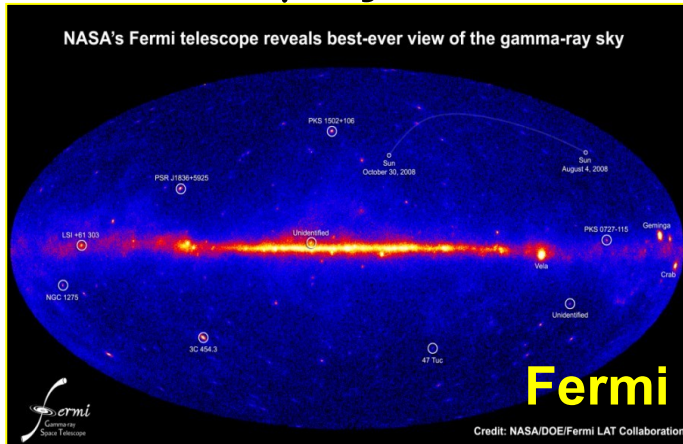
# THE ENERGETIC UNIVERSE

multi-messenger astronomy

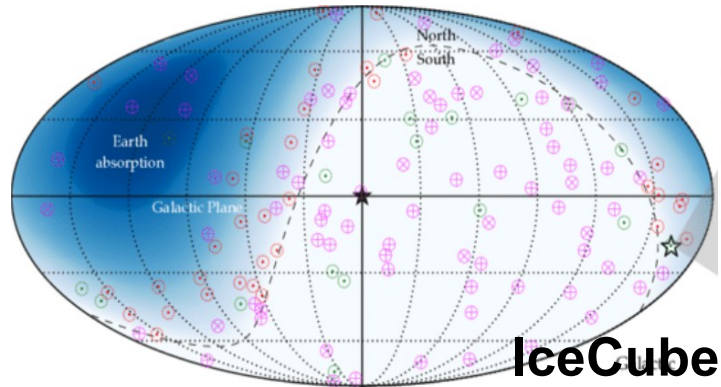


$\gamma$  rays

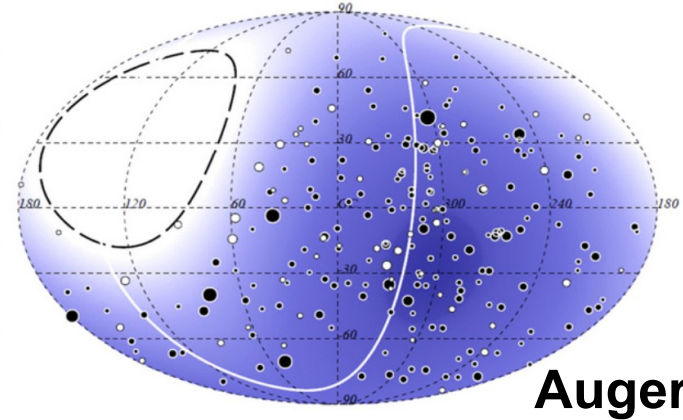
NASA's Fermi telescope reveals best-ever view of the gamma-ray sky



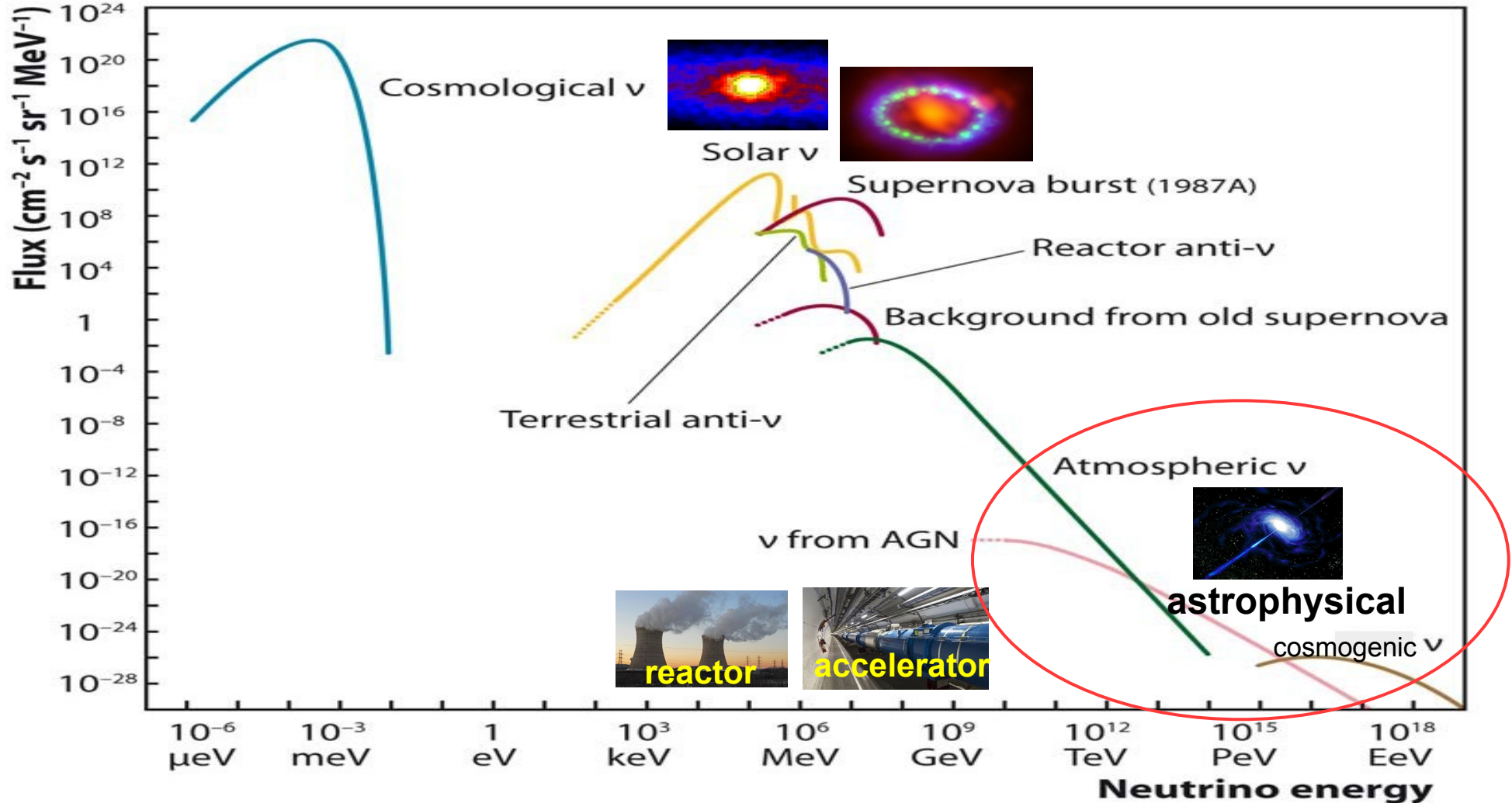
neutrinos



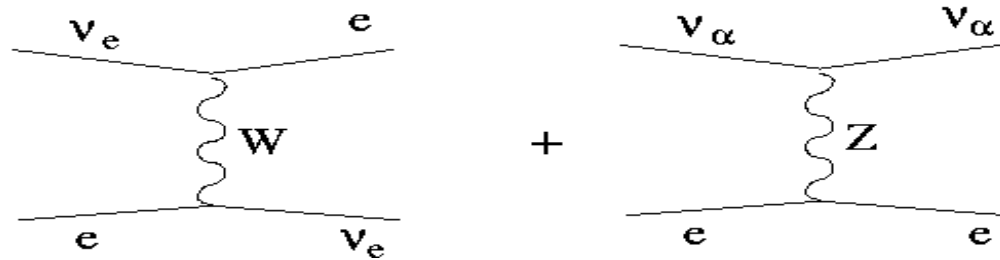
UHE Cosmic rays



# Neutrino spectra from different sources



# Neutrino – electron interactions



$$H^{eff} = \frac{G_F}{\sqrt{2}} [\bar{\nu}_\alpha \gamma_\mu (1 - \gamma_5) \nu_\alpha] [\bar{e} \gamma^\mu (c_L P_L + c_R P_R) e]$$

$$c_L = \delta_{\alpha e} + \sin^2 \theta_W - \frac{1}{2} \qquad c_R = \sin^2 \theta_W$$

**For  $E_\nu \gg m_e$ :**

$$\sigma(\nu l \rightarrow \nu l) = \frac{2G_F^2}{\pi} m_l E_\nu \left[ c_L^2 + \frac{c_R^2}{3} \right]$$

( $\sigma(\bar{\nu} l \rightarrow \bar{\nu} l) : c_L \Leftrightarrow c_R$ )

**for antineutrinos:**

$$\sigma(\nu_e e) \approx 4 \times 10^{-44} \text{ cm}^2 \left( \frac{E_\nu}{10 \text{ MeV}} \right) \approx 2.5 \sigma(\bar{\nu}_e e) \approx 6 \sigma(\nu_{\mu, \tau} e) \approx 7 \sigma(\bar{\nu}_{\mu, \tau} e)$$

(check it)

**[for  $E_\nu \gg m_e$  it is directional  $\rightarrow$  good for neutrino telescopes ]**

# Neutrino nucleon CC interactions

For  $E_\nu < 50$  MeV (Quasi Elastic):

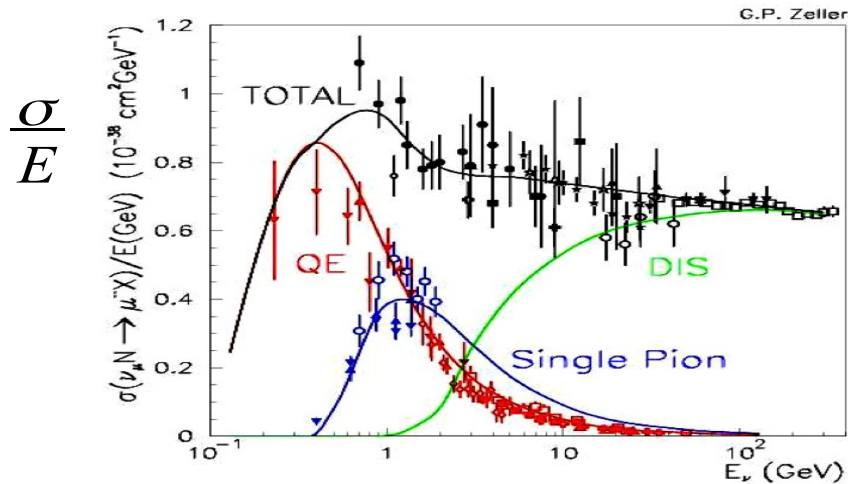
$$\sigma(\nu_e n \rightarrow pe) \approx \sigma(\bar{\nu}_e p \rightarrow n\bar{e}) \simeq \frac{G_F^2}{\pi} \cos^2 \theta_c (g_V^2 + 3g_A^2) E_\nu^2$$

$$g_V = 1$$

$$g_A = 1.27$$

inverse beta decay (IBD) used to discover  $\nu_s$  from reactors  
Reines&Cowan

for  $E_\nu > 50$  MeV the nuclear form factors smooth-out the point-like behavior (QE x-section saturates):



$$F_{V,A}(q^2) \approx F_{V,A}(0) \left[ 1 + \frac{q^2}{m_{V,A}^2} \right]^{-2} \quad (m_{V,A} \simeq \text{GeV})$$

**Quasi Elastic:**

$$\nu_\mu p \rightarrow \mu n$$

**resonant:**

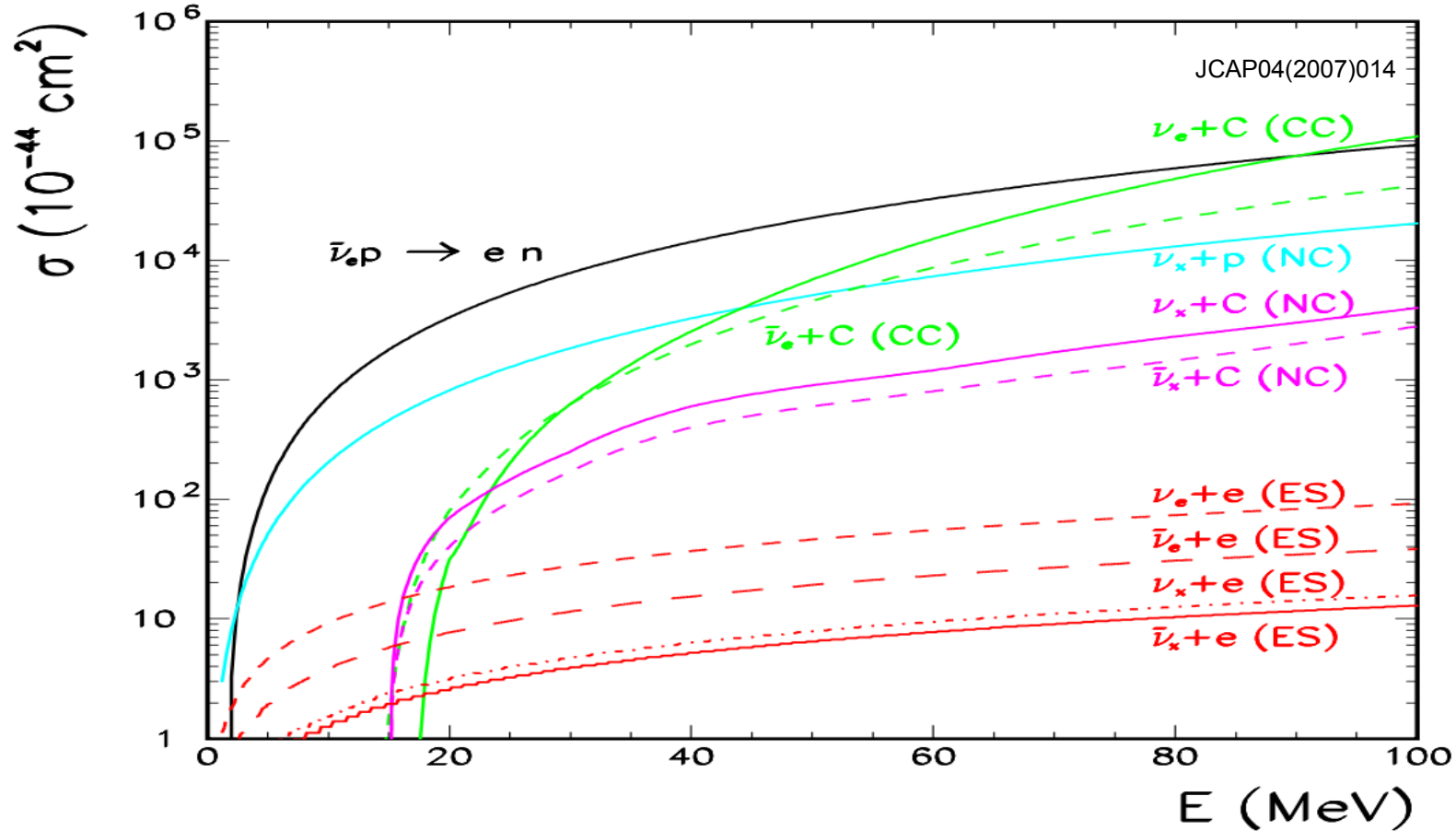
$$\nu_\mu p \rightarrow \mu p \pi$$

**Deep Inelastic:**

$$\nu_\mu p \rightarrow \mu X$$

(see Formaggio&Zeller, RevModPhys 2012)

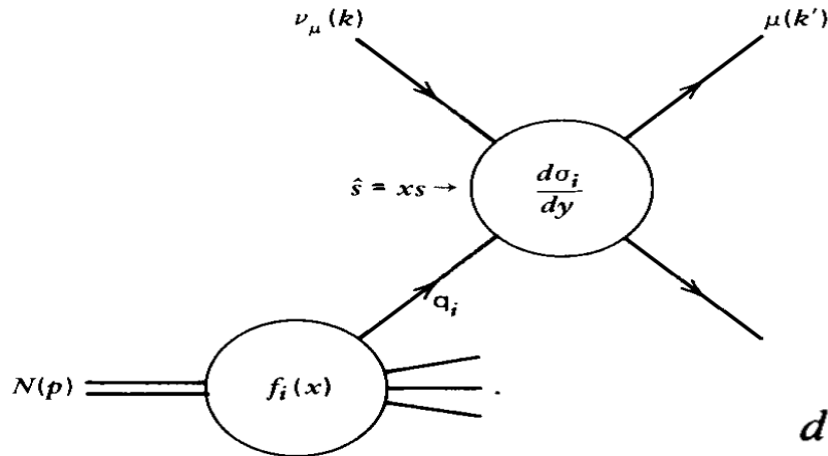
# Neutrino cross sections below 100 MeV



CC cross sections with nuclei have larger thresholds (final state blocking)

# Neutrino deep inelastic CC scattering ( $E \gg 100$ MeV):

$$\frac{d\sigma}{dx dy}(\nu N \rightarrow \mu X) = \sum_i$$



$$= \sum_i f_i(x) \left( \frac{d\sigma_i}{dy} \right)_{\hat{s}=xs}$$

$$\frac{d\sigma}{dy}(\nu_\mu d \rightarrow \mu^- u) = \frac{G^2 xs}{\pi},$$

$$\frac{d\sigma}{dy}(\bar{\nu}_\mu u \rightarrow \mu^+ d) = \frac{G^2 xs}{\pi} (1-y)^2$$

## Scaling (Bjorken) variables:

$$x \equiv Q^2 / 2m_N (E_\nu - E_l), \quad y \equiv (E_\nu - E_l) / E_\nu, \quad Q^2 \equiv -(p_\nu - p_l)^2$$

isoscalar  
target

$$d^p(x) + d^n(x) = d(x) + u(x) \equiv \mathbf{q}(x)$$

$$\bar{u}^p(x) + \bar{u}^n(x) = \bar{u}(x) + \bar{d}(x) \equiv \bar{\mathbf{q}}(x)$$



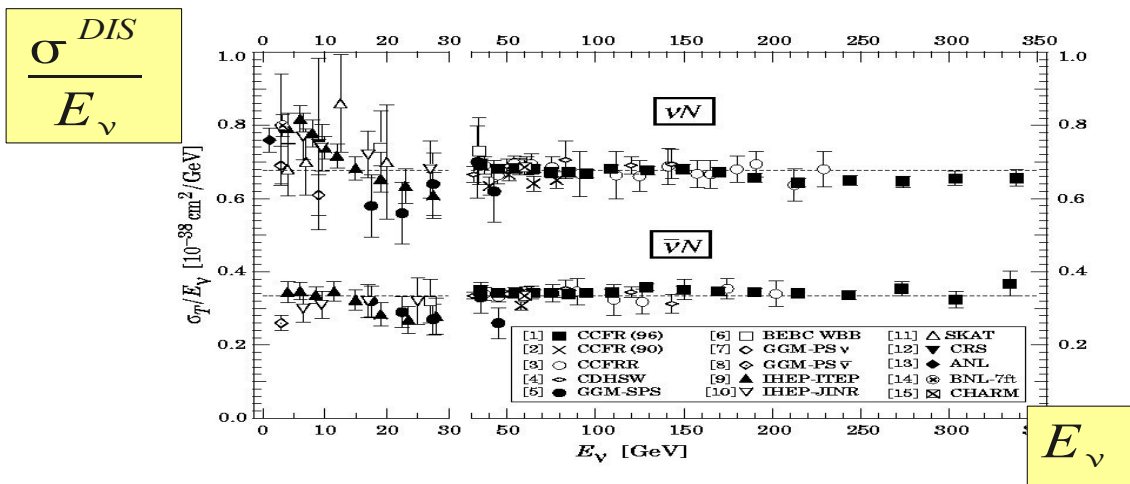
# Deep inelastic CC cross section:

$$\frac{d^2 \sigma^{DIS}}{dx dy} (\nu N \rightarrow l X) = 2 \frac{G_F^2}{\pi} m_N E_\nu \frac{M_W^4}{(Q^2 + M_W^2)^2} [xq(x, Q^2) + x(1-y)^2 \bar{q}(x, Q^2)]$$

for

$$E_\nu < M_W^2 / 2 m_N \simeq 5 \text{ TeV}$$

$$\Rightarrow \sigma^{DIS} \propto E_\nu$$

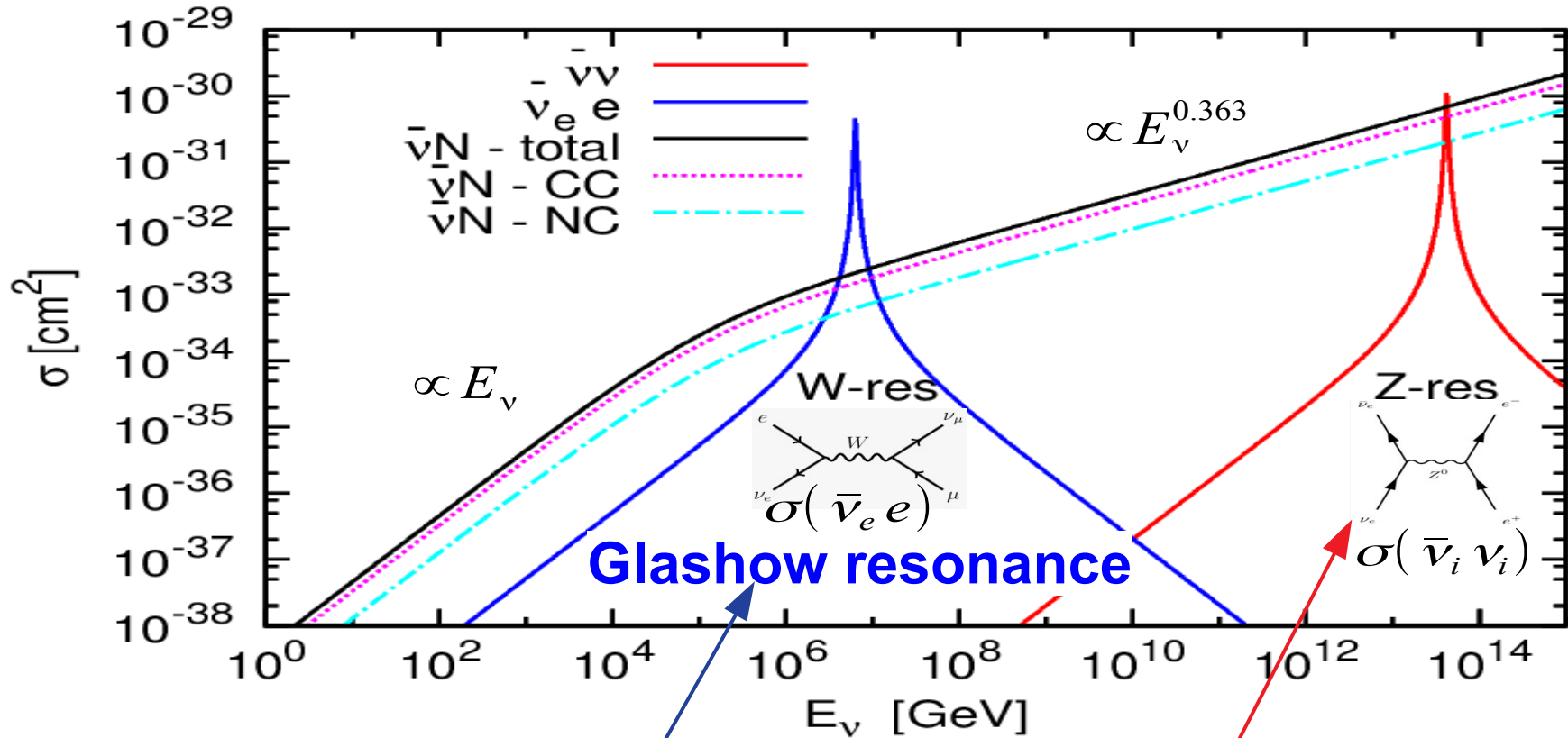


For  $E_\nu \gg 5 \text{ TeV}$ , W propagator  $\Rightarrow$  sensitivity to  $Q^2 < M_W^2$   
 i.e., depends on small values of  $x$  ( $\ll 1$ )

$$\sigma_{CC}(\nu N) \simeq \sigma_{CC}(\bar{\nu} N) \simeq 5.5 \times 10^{-36} \text{ cm}^2 (E_\nu / \text{GeV})^{0.363}$$

$$\sigma_{NC}(\nu N) \simeq \sigma_{NC}(\bar{\nu} N) \simeq 2.3 \times 10^{-36} \text{ cm}^2 (E_\nu / \text{GeV})^{0.363}$$

# The W and Z resonances



**W-resonance:**

$$E_\nu \simeq M_W^2 / 2m_e \approx 6.3 \text{ PeV}$$

**Z-resonance:**

$$E_\nu \simeq M_Z^2 / 2m_\nu \approx 4 \times 10^{22} \text{ eV} (0.1 \text{ eV} / m_\nu)$$

Glashow '60

Weiler '84  
ER '93

# COSMIC NEUTRINO BACKGROUND

- decoupled from  $e^+e^-$  plasma at  $T \sim \text{MeV}$  ( $t \sim 1\text{s}$ ),  
while CMB decoupled in e-p recombination at  $T \sim 0.3\text{eV}$  (370 kyrs)

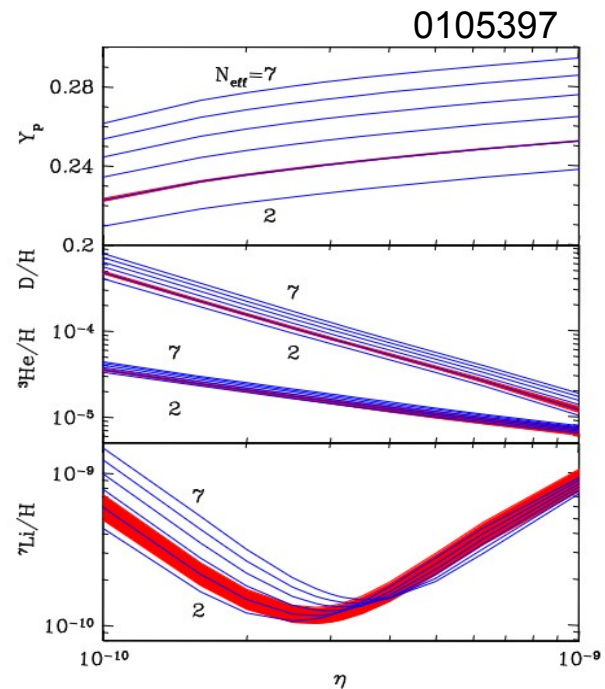
- affect Universe expansion (they contribute to radiation):

→ affect Big Bang Nucleosynthesis (BBN):  
more neutrino families → faster expansion  
→ less neutrons decay → more He →  $N_\nu < 3.4$

→ affect matter-radiation equality → affect CMB  
combining with BAO from galaxy surveys →  $N_{\text{eff}} = 2.99 \pm 0.33$

become non-relativistic at  $z_i = 188$  ( $m_i/0.1\text{eV}$ )

→ values of masses affect structure formation →  $\Sigma m_i < 0.12 \text{ eV}$

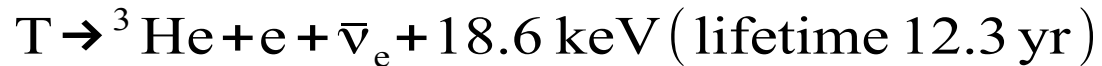


Today  $T_\nu = 1.9\text{K} \rightarrow n_\nu = 56/\text{cm}^3$  per species ( $\nu_\alpha$  or anti  $\nu_\alpha$ )

(while  $T_\gamma = 2.7\text{K}$  due to reheating when  $e^+e^-$  annihilated)

may be detected by capture on beta decaying nuclei such as Tritium:

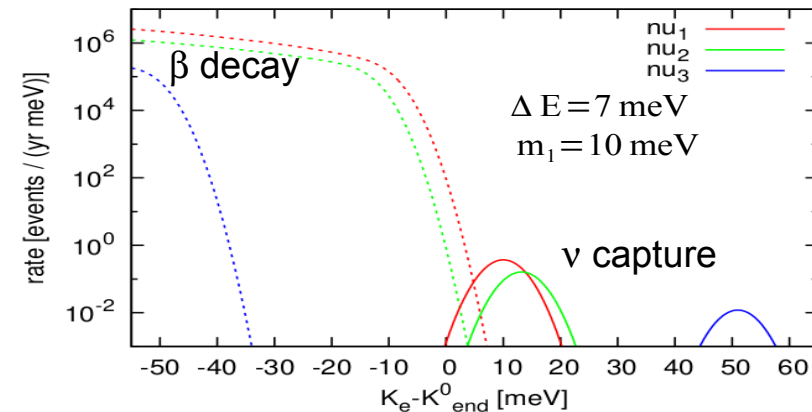
(Weinberg 62)



1810.00505

→ PTOLEMY experiment@GS  
requires exquisite E resolution

(aim at  $\Delta E < 50 \text{ meV}$ )

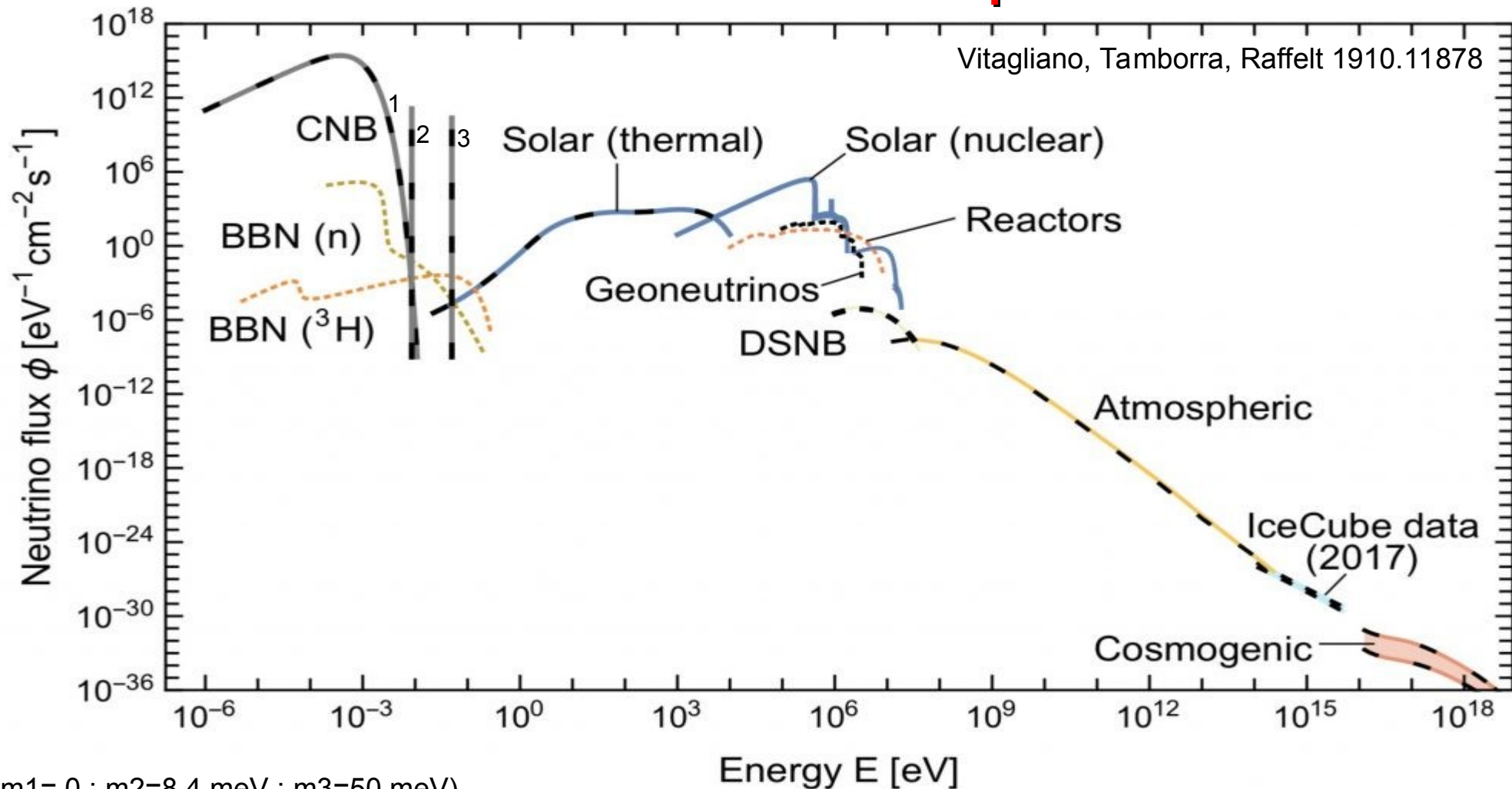


if neutrino non-relativistic now:

rate for Majorana = 2 rate of Dirac (half of states became sterile)

for 100 g of T:  $\sim 8$  events/yr (Majorana),  $\sim 4$  events/yr (Dirac)

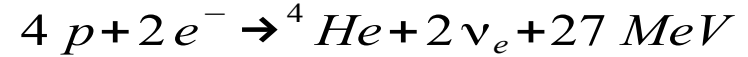
# Grand Unified neutrino spectrum



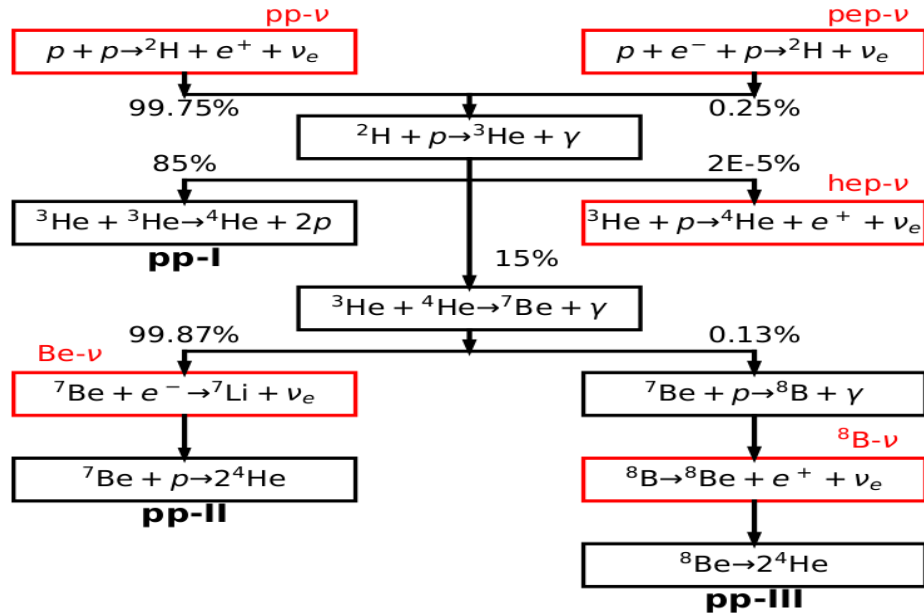
( $m_1 = 0$  ;  $m_2 = 8.4$  meV ;  $m_3 = 50$  meV)

# SOLAR NEUTRINOS

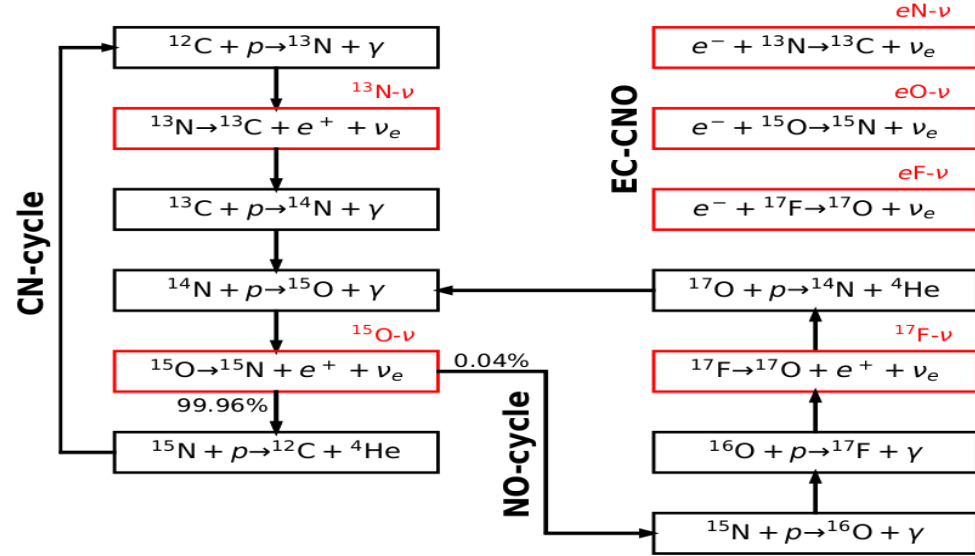
Solar Energy produced by fusion:



## pp-chain



## CNO-cycle



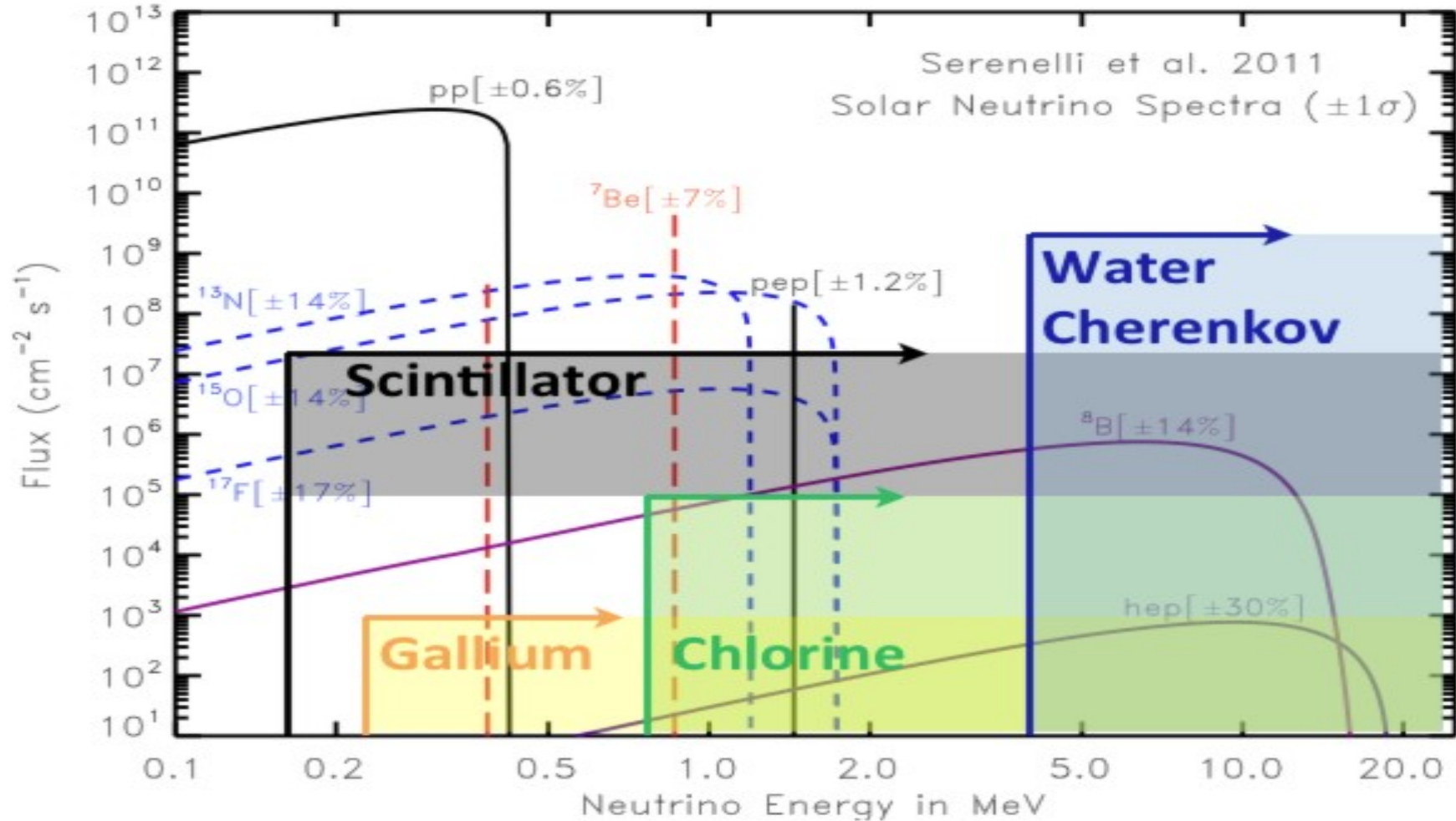
energy flux  
on Earth:

$$\phi_E \simeq 1.5 \text{ kW} / \text{m}^2$$

⇒ neutrino  
flux:

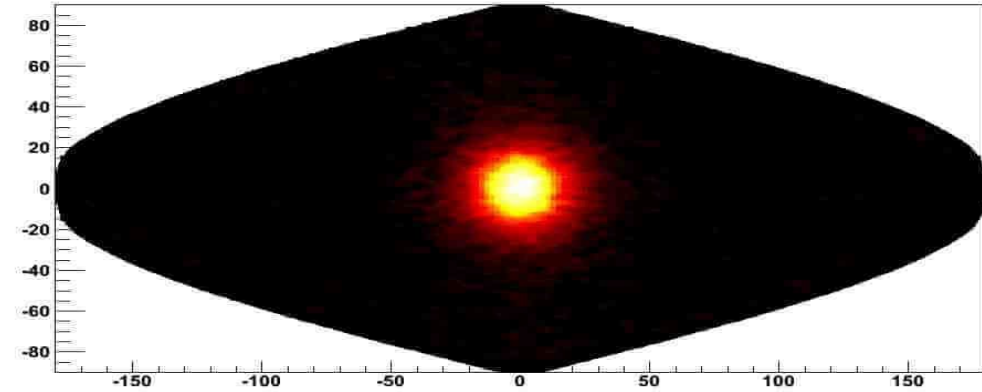
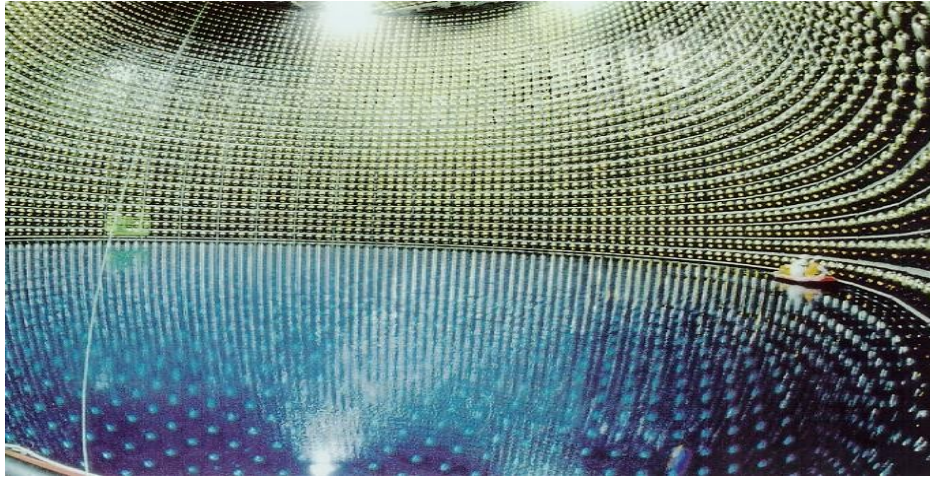
$$\phi_\nu \simeq \frac{2 \times \phi_E}{27 \text{ MeV}} \simeq 6 \times 10^{10} \frac{\nu_e}{\text{cm}^2 \text{ s}}$$

# SOLAR NEUTRINO SPECTRA



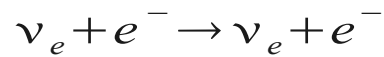
$\sim 3\%$  of solar energy emitted in  $\nu_e$  with  $E_\nu \sim 0.1-14$  MeV

# Super Kamiokande



first image of a neutrino source

## Cherenkov in H<sub>2</sub>O

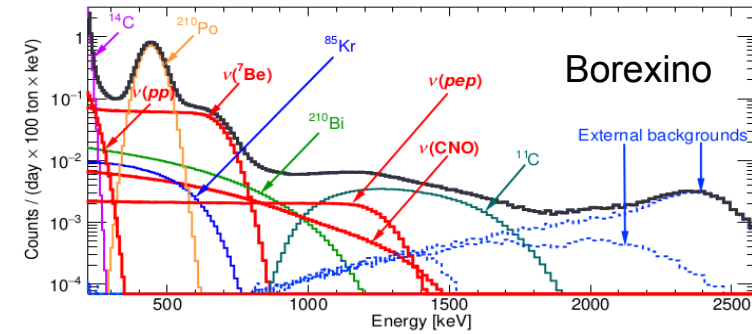


(and some sensitivity to  $\nu_{\mu, \tau}$ )

scintillators (Borexino, KamLAND ..)  
lower threshold but no directionality

SNO heavy water, smaller volume than SK

solar neutrinos also observed radiochemically (with Cl & Ga),  
but just counted a month later → no direction nor original energy





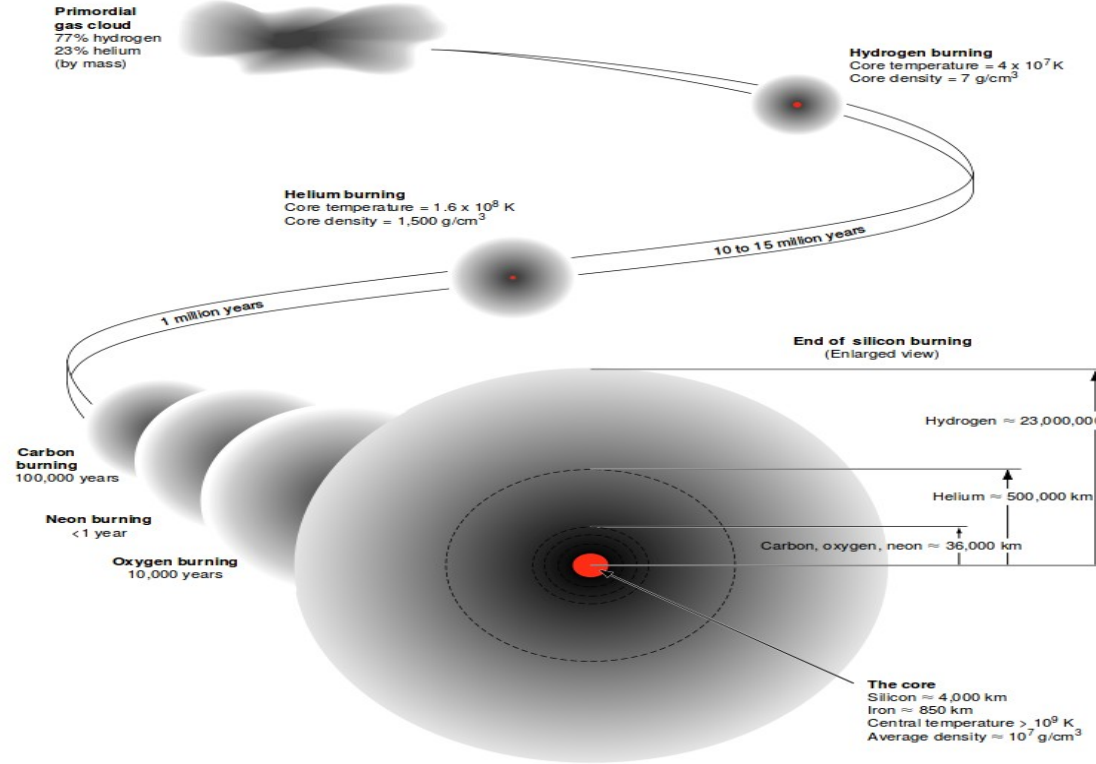
# Supernova Neutrinos

The life of stars:  
gravity against thermal pressure  
→ need to burn fuel

Stars with  $M > 8M_{\text{sun}}$  burn  
H → He → C,O → Ne → Si → Fe

but Fe burning is endothermic  
→ contracts until e degeneracy  
pressure prevents collapse

but when  $M_{\text{Fe}} \approx M_{\text{Ch}} \approx 1.4 M_{\text{sol}}$   
⇒ collapse to n-star ( $R = 10 \text{ km}$ )  
releasing  $E_b \approx GM_{\text{sun}}^2/R \approx 3 \times 10^{53} \text{ erg}$



**SUPERNOVA explosion → 99% emitted as MeV neutrinos in ~ 10 s !**

1% in kinetic energy of expanding envelope, 0.01% in light over few months

hypernovas may produce long GRBs, accelerating CRs in the jets → UHE  $\nu$   
short GRBs (< 2 sec) from neutron star binaries can also produce UHE  $\nu$

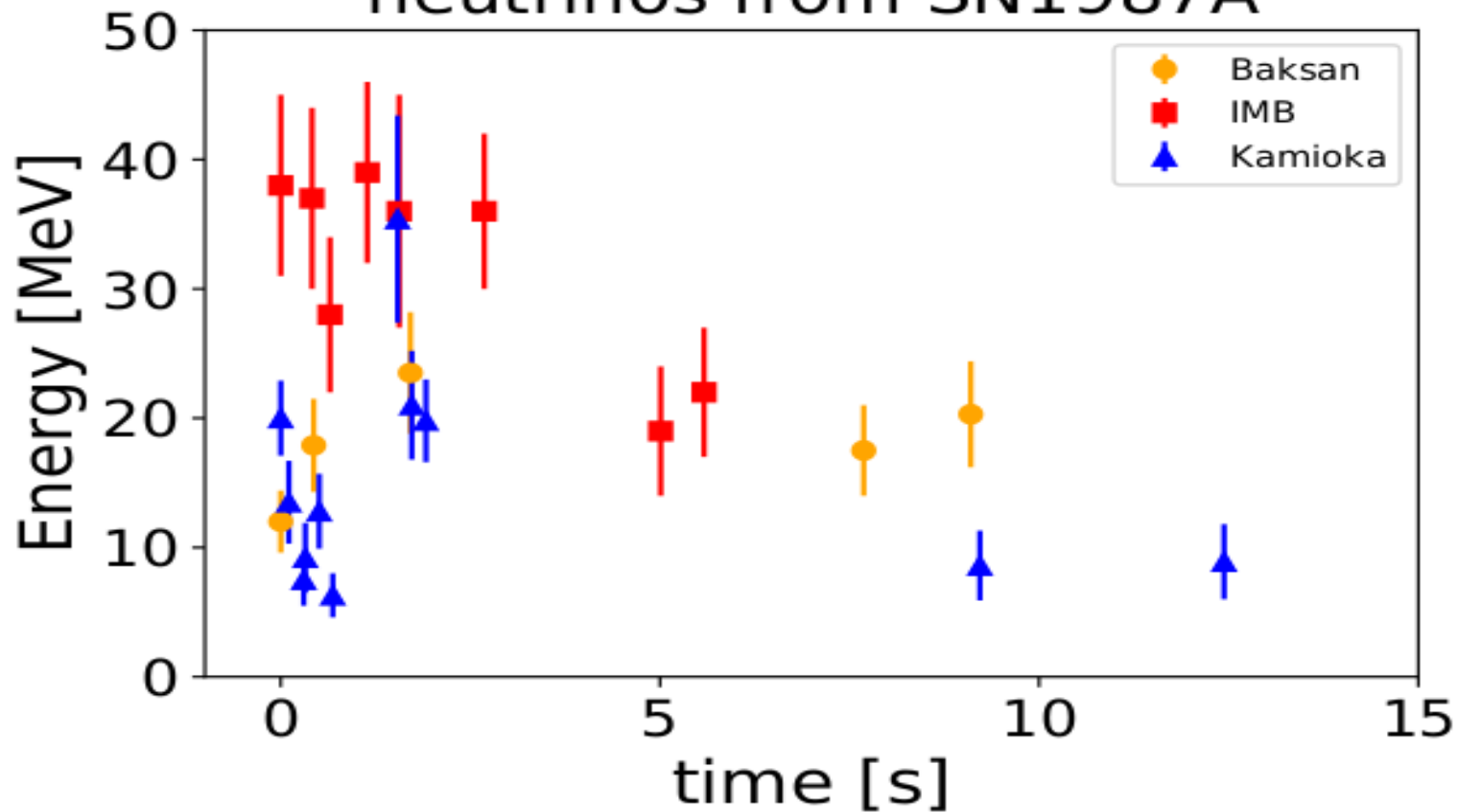
**No Galactic SN seen since the invention of the telescope :(  
although about 1 to 3 per century expected, but obscured by gas**



**SN 1987A was seen in the nearby LMC in light and in neutrinos !**



# neutrinos from SN1987A

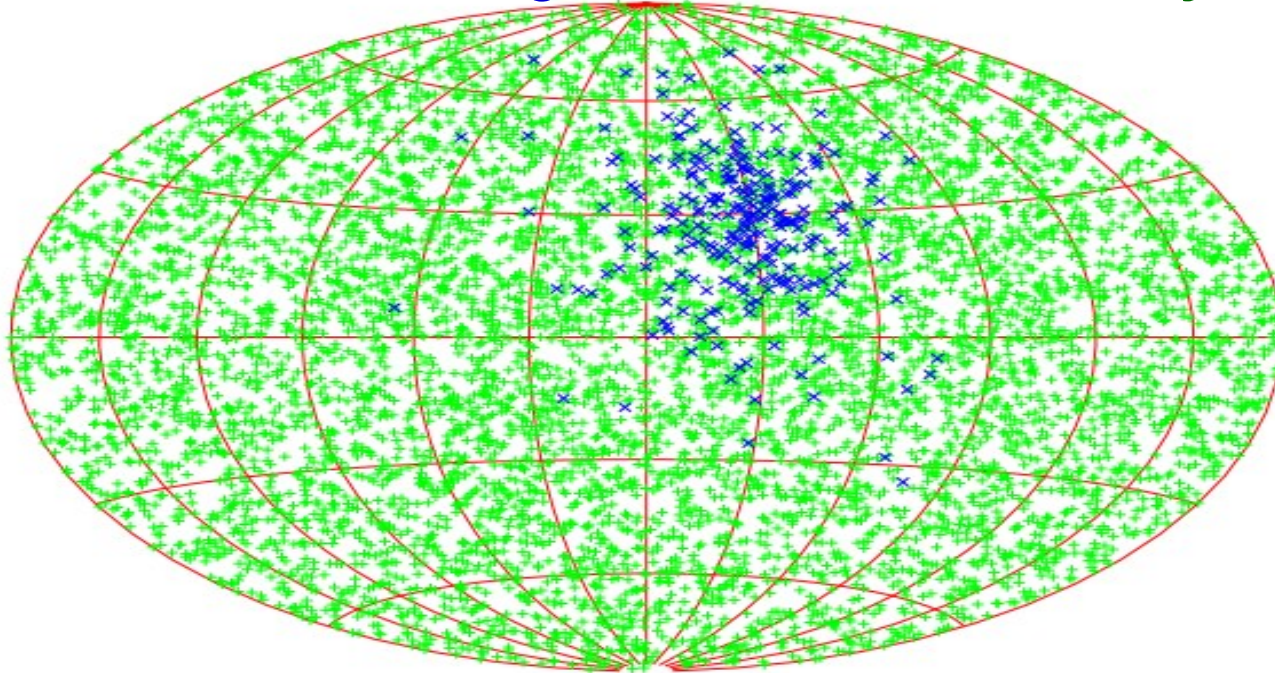


most (all?) IBD of  $\bar{\nu}_e$  → isotropic, no picture :(

# future Galactic supernova (d~10 kpc) at SuperKamiokande:

Electron Scattering

Inverse Beta Decay

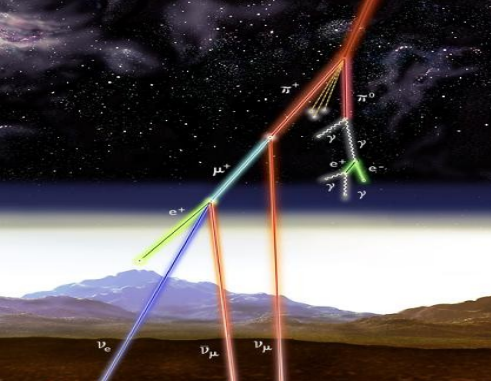


~6000 events  
expected in SK



pointing may reach 2-3 deg with SK Gd  
it happens few hours before light is emitted → early warning

<http://snews.bnl.gov>



# High energy atmospheric neutrinos

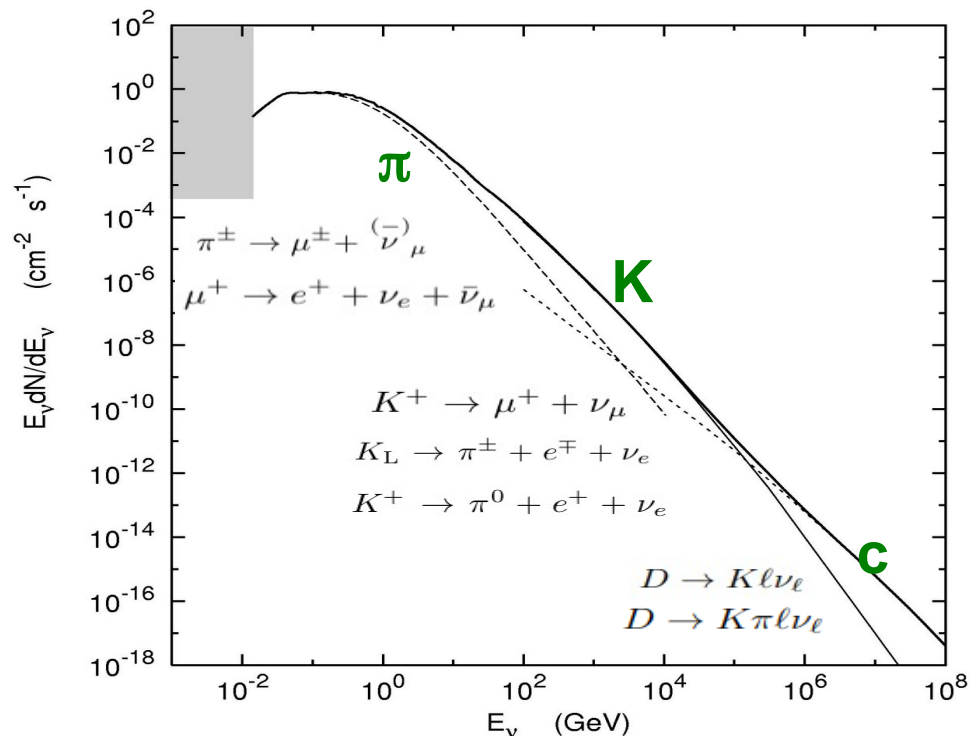
(from weak decays of hadrons in showers produced by CR-air interactions)

meson decay length:  $L = \gamma c \tau$  ,  $\gamma = E/m$

$$L_{\pi} \simeq 6 \text{ km} (E_{\pi} / 100 \text{ GeV})$$

$$L_K \simeq 7.5 \text{ km} (E_K / \text{TeV})$$

$$L_D \simeq 2 \text{ km} (E_D / 10 \text{ PeV})$$



interaction length  $\lambda_{\pi} = 120 \text{ g/cm}^2$

attenuation length:  $L^{att} = \frac{1.2 \text{ km}}{\rho / 10^{-3} \text{ g/cc}} \simeq 5 \text{ km}$  (h~10 km)

at low energies, atmospheric  $\nu$ s mainly from pion decays

but above 100 GeV pions are stopped before decay  $\rightarrow$  kaons become the main source

and above ~100 TeV kaons suppressed  $\rightarrow$  prompt charm decays dominate

# Prompt charm production

Enberg et al. /0808.0418

For  $E > 200 \text{ TeV} \rightarrow \nu$  mostly from  $c$  decays

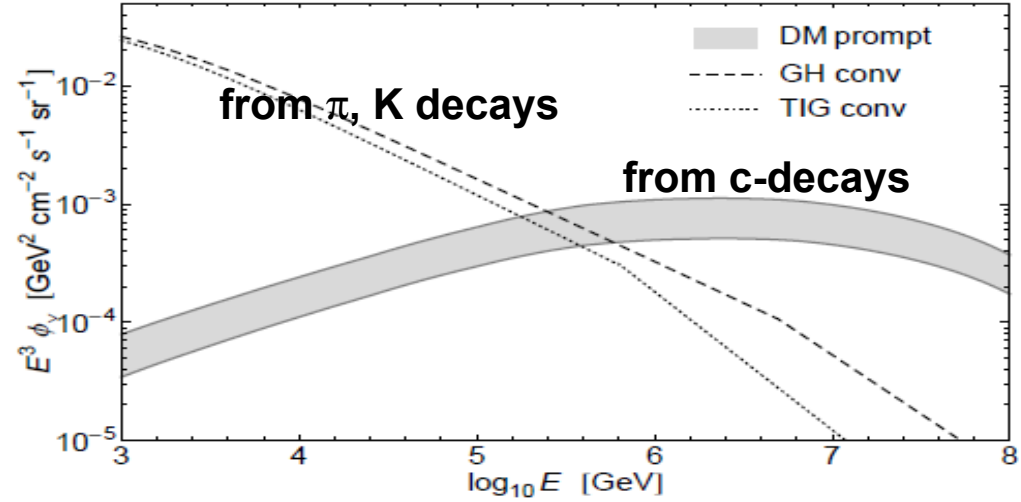
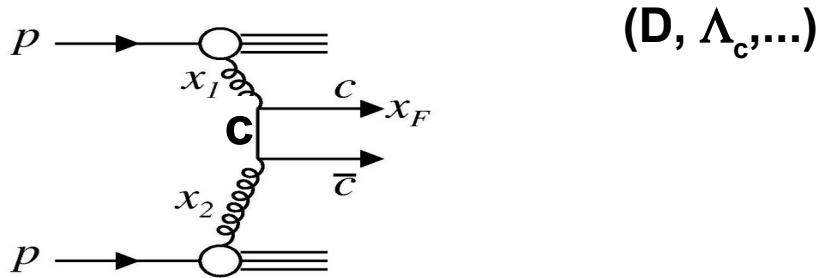
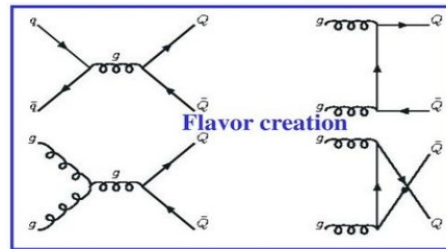


FIG. 5: Prompt and conventional  $\nu_\mu + \bar{\nu}_\mu$  fluxes in the vertical

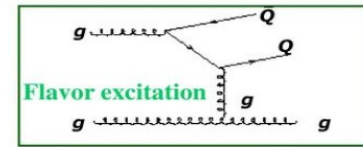
sample gluon density distribution at tiny momentum fractions  
 ( $x_2 < 10^{-5}$  for  $E > 10^{15} \text{ eV}$ )  $\rightarrow$  need to extrapolate from measured values

also requires to include large NLO processes

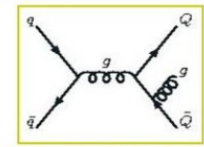
$\rightarrow$  significant uncertainties



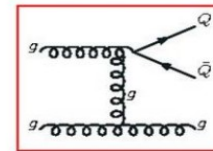
Leading Order and Next to Leading Order



Flavor excitation

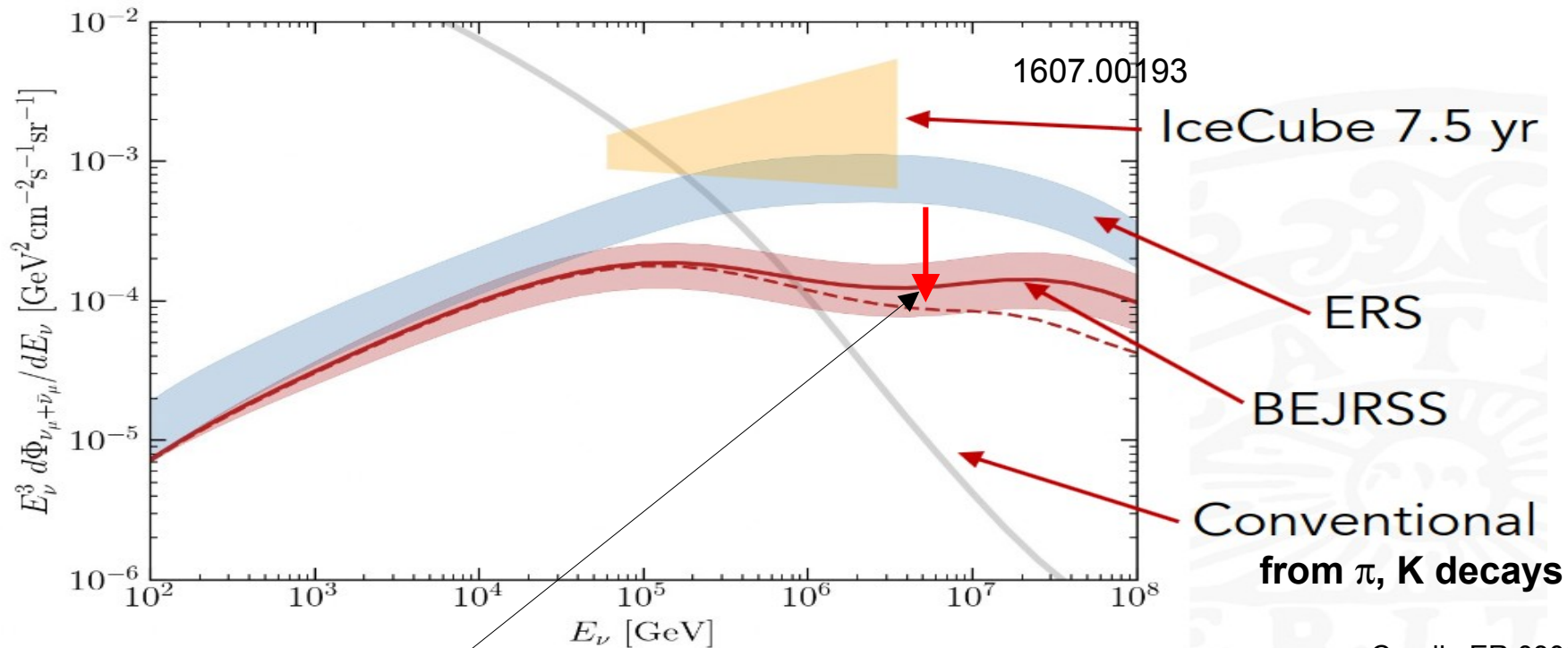


radiative corrections

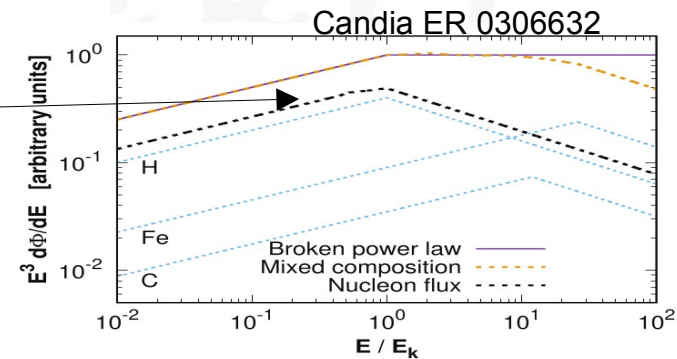


Gluon splitting

# Prompt neutrino fluxes



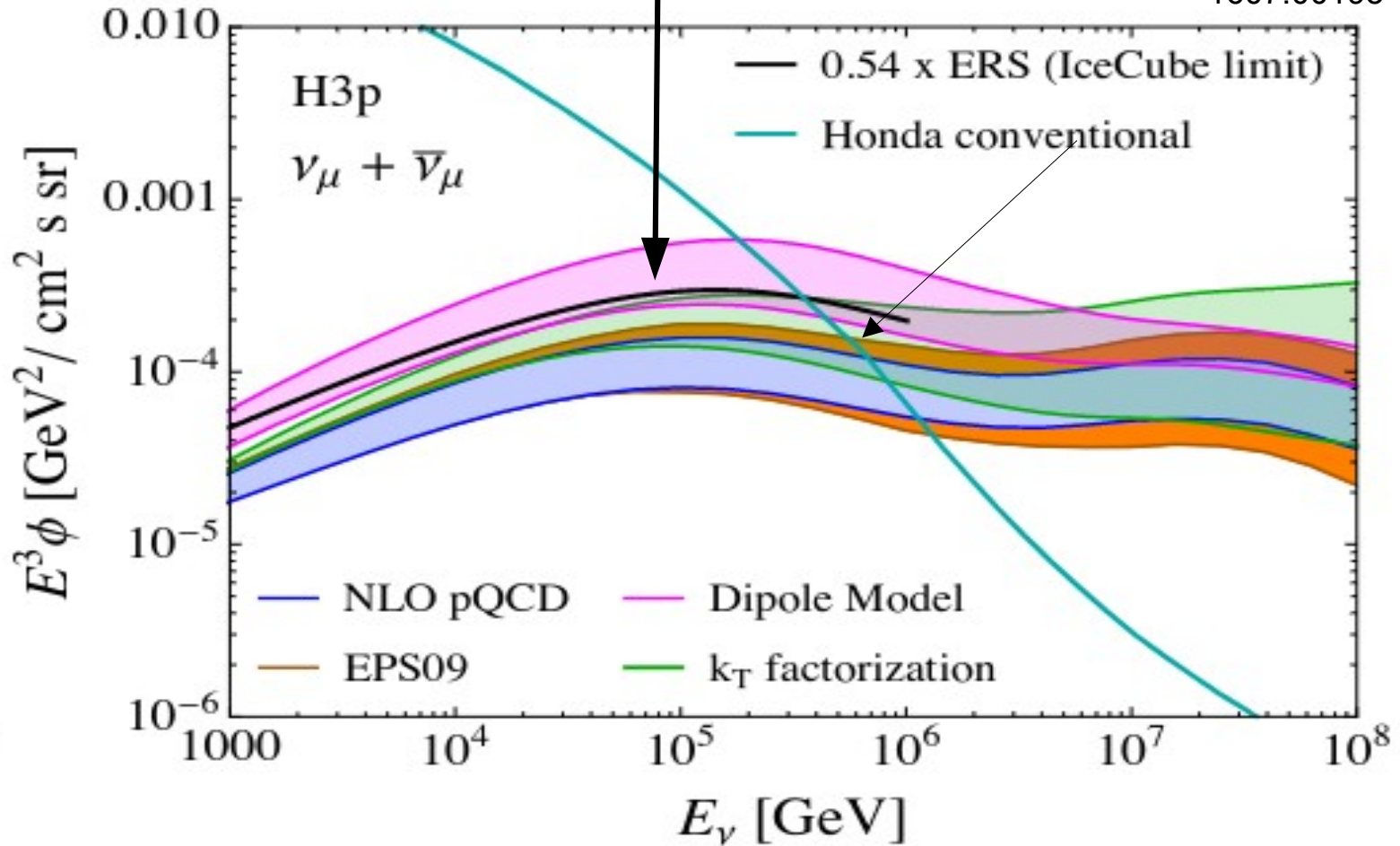
what is relevant is the flux of nucleons  
 not just the total CR fluxes  
 (ERS assumed all CRs where H)





# IceCube bounds

1607.00193



Note that:

Probability for pion to decay before interacting  $\sim \lambda_{\text{int}}/\lambda_{\text{dec}} \sim E^{-1}$

→ conventional atmospheric flux  $\sim \text{CR flux} / E \sim E^{-3.7}$

Prompt flux  $\sim \text{CR flux} \sim E^{-2.7}$

(and steeper above 100 TeV due to knee)

$\nu_e$  conventional fluxes from K decays

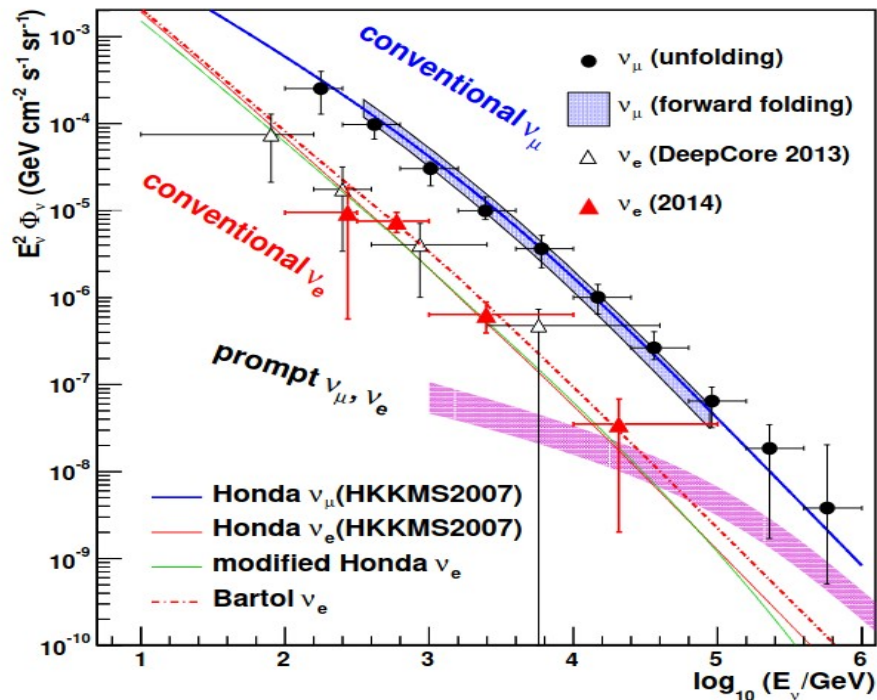
$\sim \nu_\mu$  fluxes/30

while  $\nu_e \sim \nu_\mu$  for prompt

→  $\nu_e$  fluxes cross at  $\sim 30$  TeV

→  $\nu_\mu$  fluxes cross at  $> 500$  TeV

1504.03753

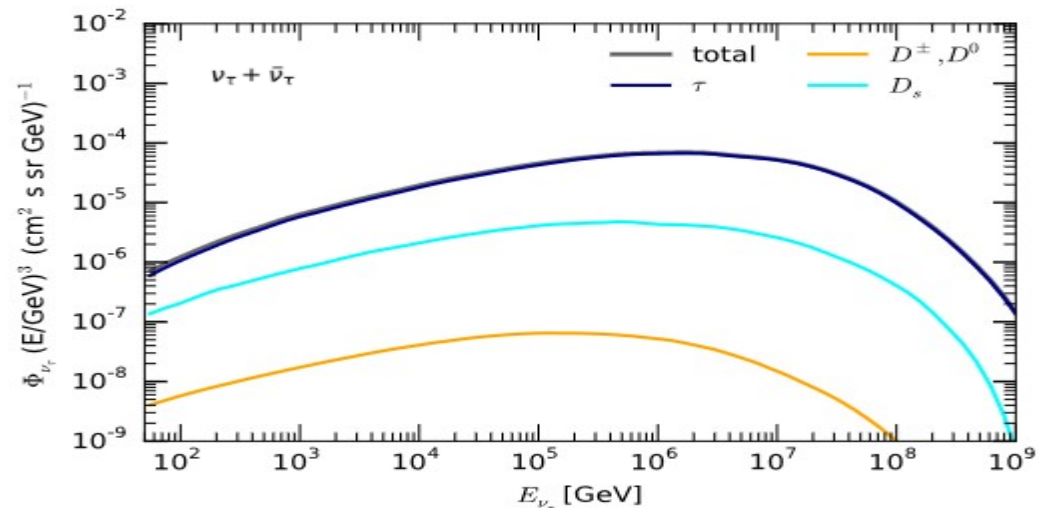
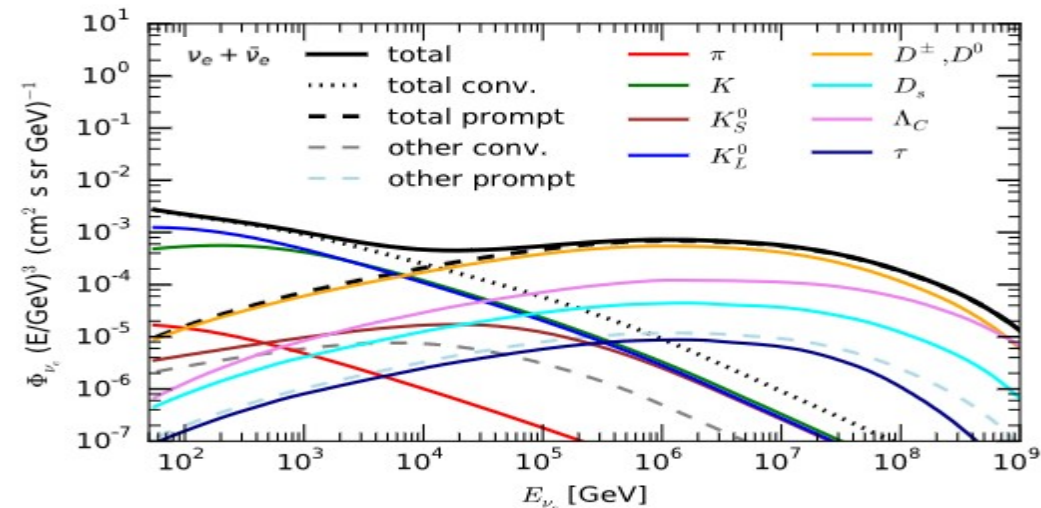
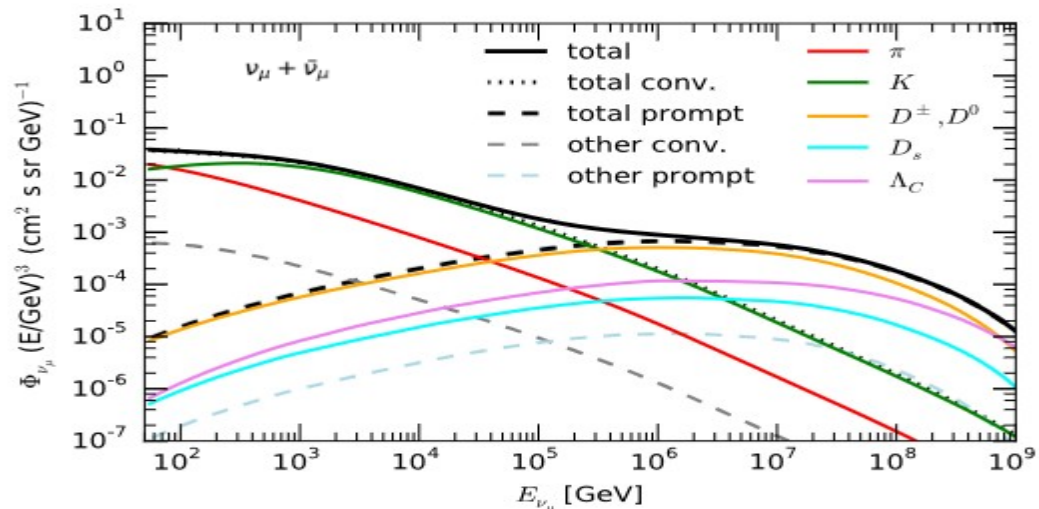
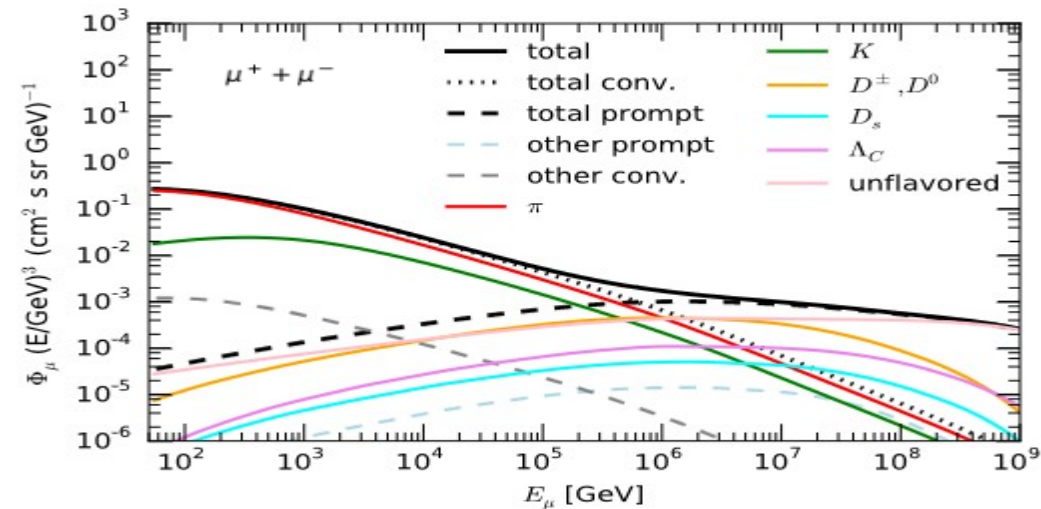


astrophysical sources have less dense targets →  $\pi$  decay without interaction

→ astrophysical fluxes  $\sim E^{-\gamma}$ , with  $\gamma \sim 2$  to 2.4 ?

They may become dominant above few tens or hundreds of TeV

# ATMOSPHERIC MUON & NEUTRINO FLUXES

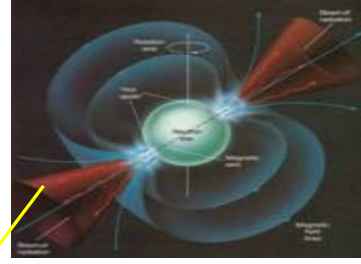


# Examples of powerful astrophysical objects [potential CR accelerators]

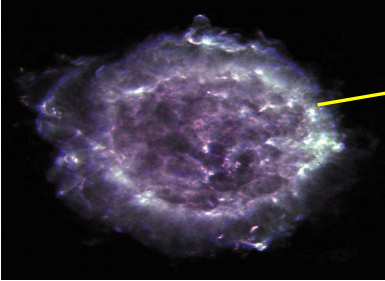
AGN



Pulsar



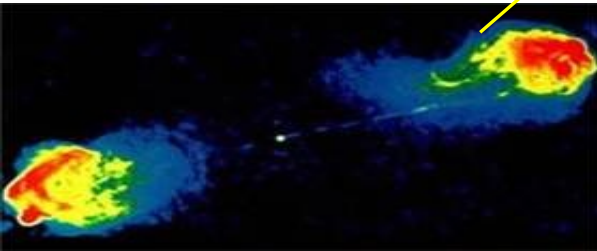
SNR



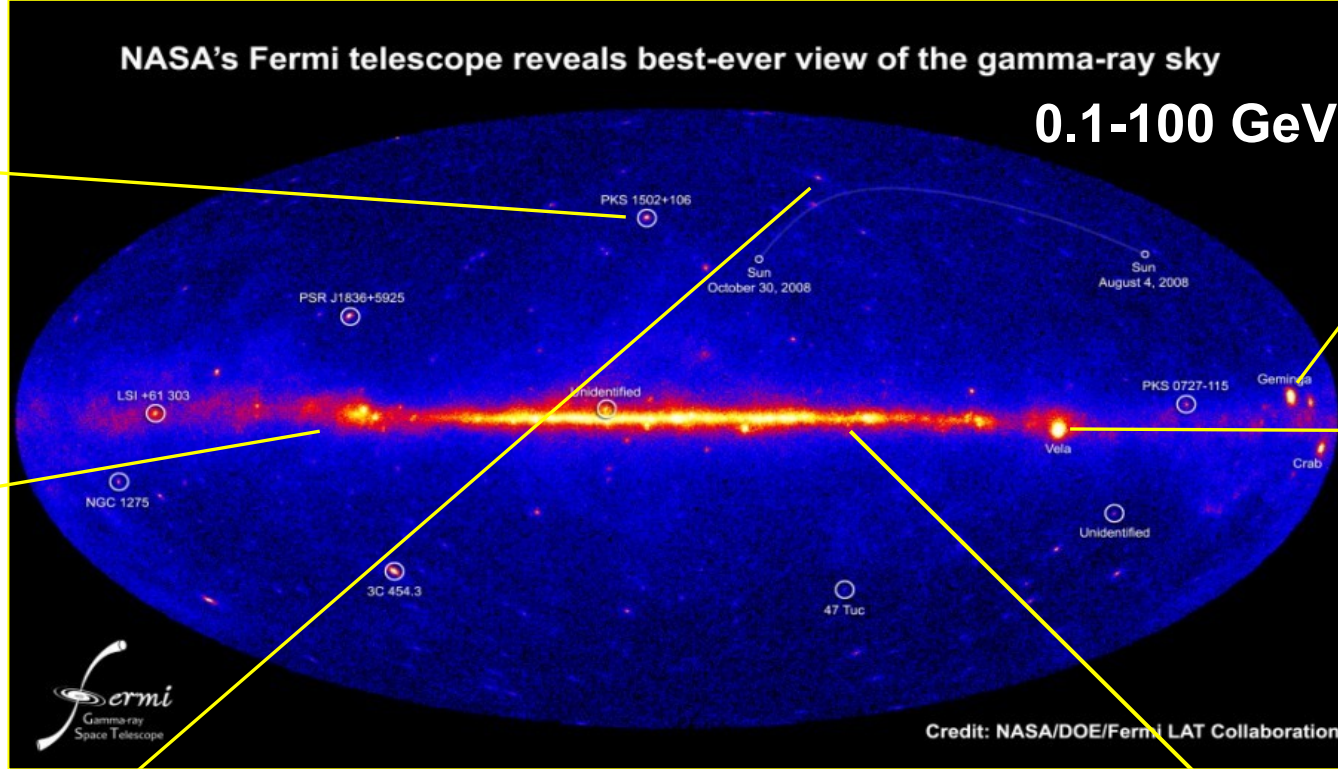
PWN



radio Galaxy



GRB

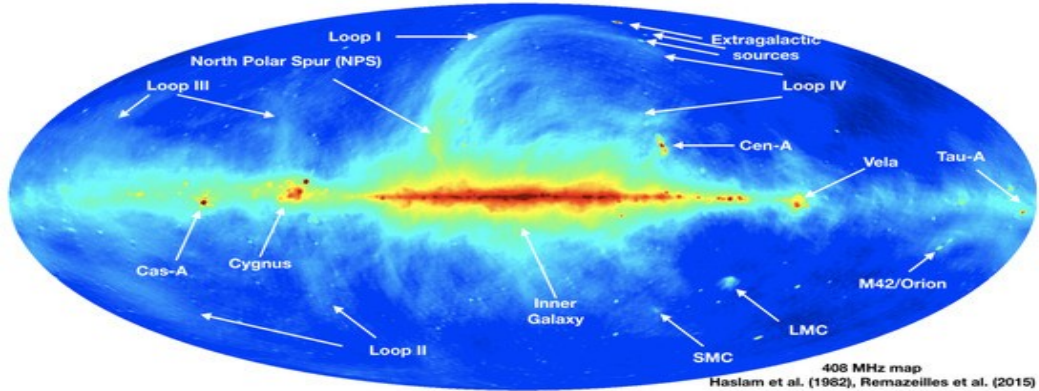


colliding galaxies

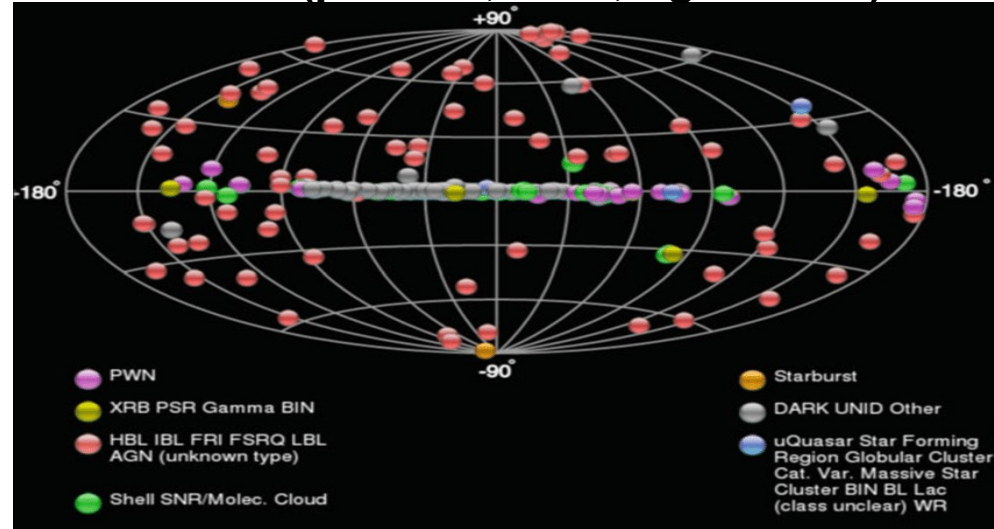


diffuse emission

Radio maps:  
 Synchrotron emission  
 (accelerated e, SNR, pulsars,  
 radiogalaxies, ...)



TeVCAT (IACT, HAWC, LHAASO)  
 $E > 100$  GeV (pulsars, SNR, xg AGNs..)



Synchrotron  
 from e in jet

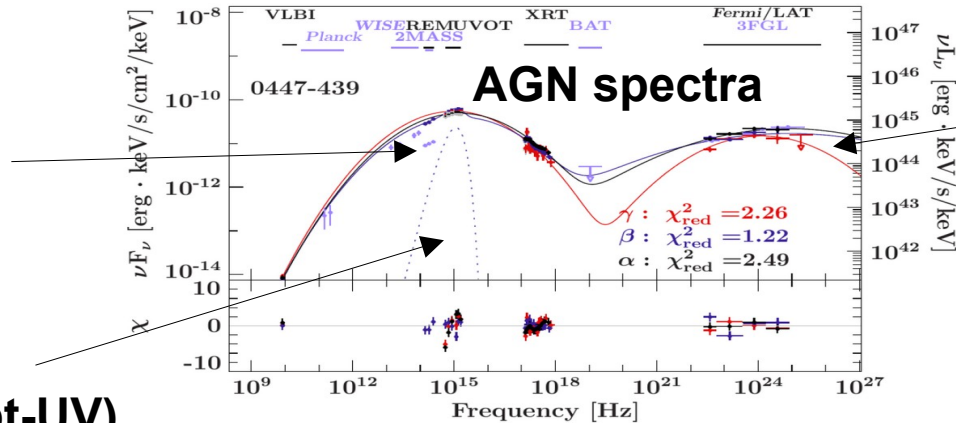


Figure 1.5: Broadband spectral energy distribution of PKS 0447-439. The model includes both peaks and an additional blackbody for present thermal excess. (Krauß et al., 2016)

thermal from  
 accretion disk (opt-UV)

Compton on synchrotron  $\gamma$   
 or on external radiation

and/or

hadronic  $\gamma$  from  $\pi^0$  decays  
 produced by interactions  
 of accelerated CRs  
 $\rightarrow \nu$  from associated  
 charged pion decays

# Neutrino production mechanisms

## pp scenarios (gas targets)

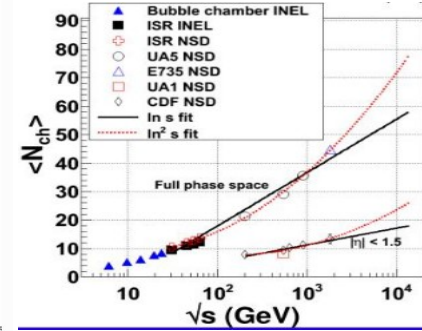
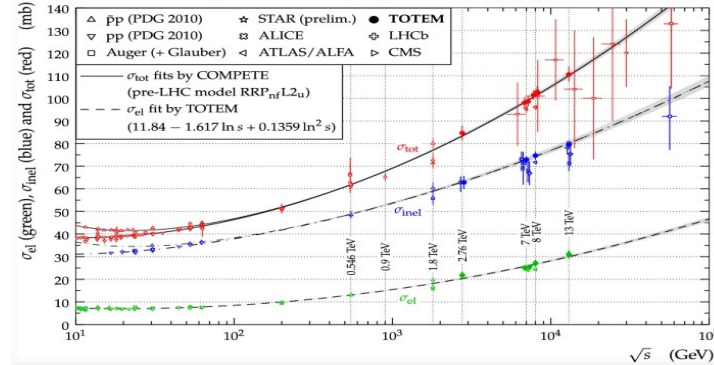
$$p+p \rightarrow \frac{N_\pi}{3} (\pi^+ + \pi^0 + \pi^-) + X$$

$$\pi^\pm \rightarrow \mu^\pm + \bar{\nu}_\mu$$

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

$$\pi^0 \rightarrow \gamma + \gamma$$

$$E_\pi \approx 2 E_\gamma \approx 4 E_\nu$$



Large  $\pi$  multiplicities  $\rightarrow \langle E_\nu \rangle \sim E_p / 50$

low threshold, smooth rise, large multiplicities

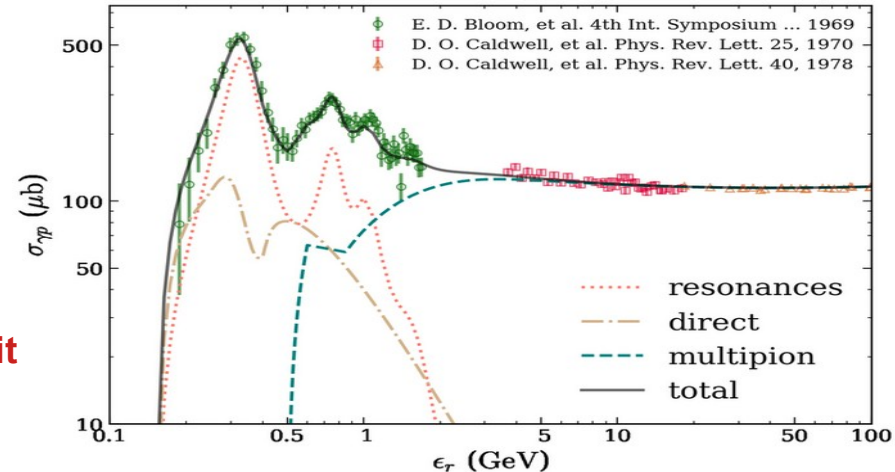
## py scenarios (radiation targets)

e.g.

$$p + \gamma \rightarrow \Delta^+ \rightarrow n + \pi^+$$

$$E_\nu \approx E_\pi / 4 \approx E_p / 20$$

threshold at:  $E'_p > \frac{m_\Delta^2 - m_p^2}{4\varepsilon'_\gamma} = \left( \frac{100 \text{ eV}}{\varepsilon'_\gamma} \right) \times 1.6 \text{ PeV}$   
**prove it**



In blazars or GRB jets, observed E is boosted

$$E_\nu \approx \Gamma E'_\nu$$

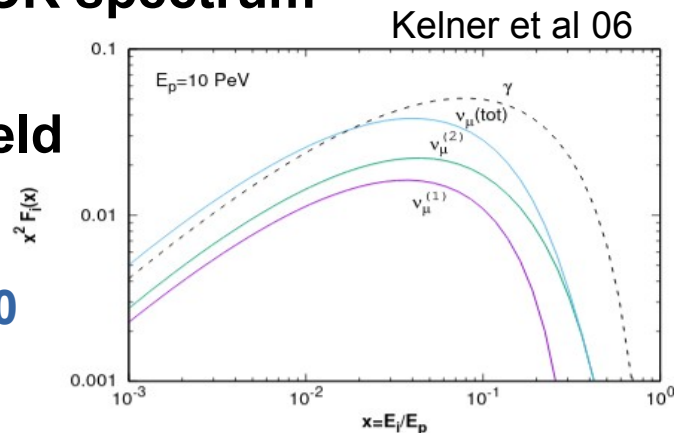
pp scenarios (SNR, AGN lobes & cores,  
CR reservoirs: SBG, galaxy clusters, ...)

**$\nu$  emissivity:**

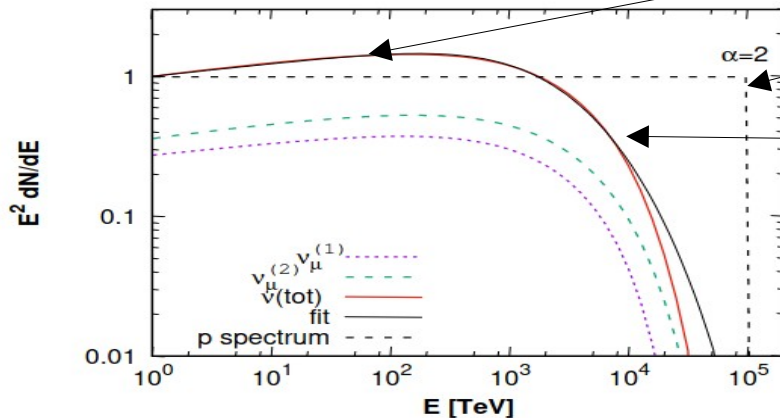
$$\frac{dq_{\nu_i}}{dE_\nu} = N_N \int_{E_\nu}^{\infty} \frac{dE_p}{E_p} \sigma_{pp} F_{\nu_i}(x_\nu, E_p) \frac{dq_p}{dE_p}$$

$N_N$  ← column density  
 $F_{\nu_i}(x_\nu, E_p)$  ←  $\nu$  yield  
 $\frac{dq_p}{dE_p}$  ← CR spectrum

Neutrino energy distribution much broader than just  $E_p/20$   
(corresponding to  $E_\nu \sim E_\pi/4 \sim 0.2 E_p/4$ )



The neutrino spectrum approximately follows CR spectrum (scaling)  
→ almost power-law down to low energies (just slightly harder)



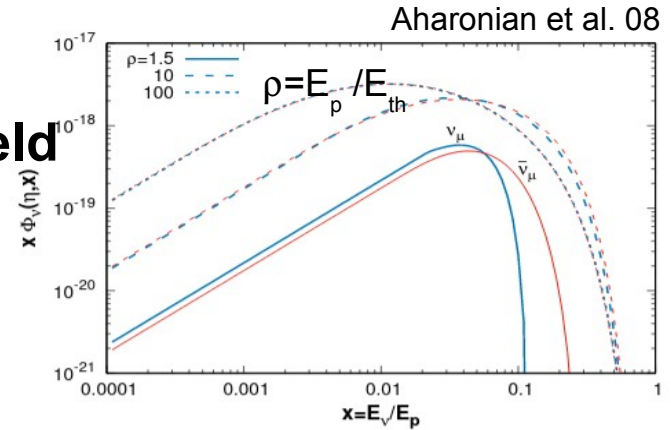
assuming CR proton spectrum  $E^{-2}$   
with sharp cutoff  
neutrino cutoff shape much softer  
(actually depends on CR cutoff shape)

# $p\gamma$ scenarios (AGN jets&cores, GRBs, ...)

$$\frac{dq_i}{dE'_i}(E'_i) = \int \frac{dE'_p}{E'_p} d\varepsilon' \frac{dq_p}{dE'_p}(E'_p) \frac{dn_{ph}}{d\varepsilon'}(\varepsilon') \Phi_i(\eta, x_i)$$

photon density

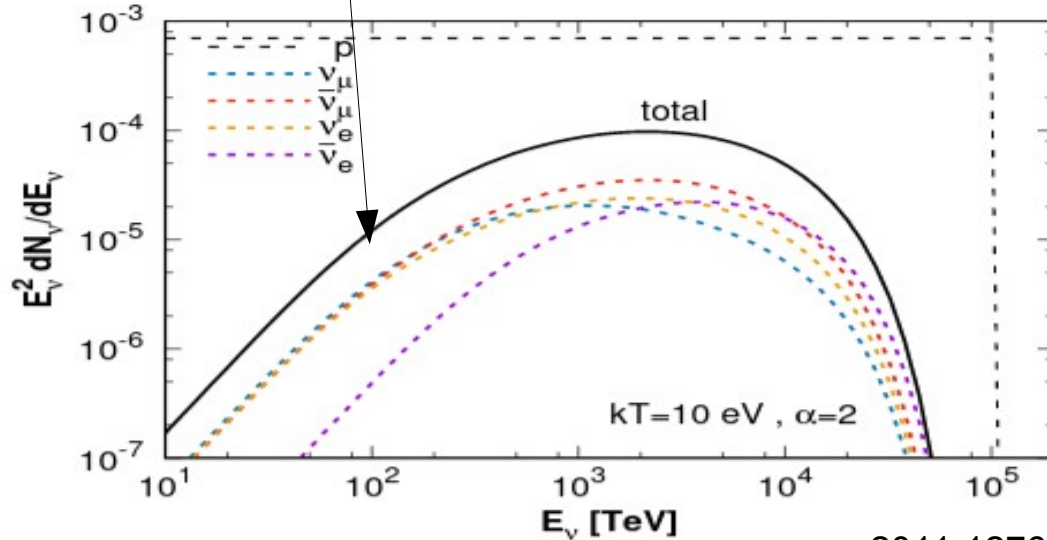
$\nu$  yield



Aharonian et al. 08

the neutrino spectrum depends on both the photon and CR spectra  
strongly suppressed below threshold of  $\Delta$  production

$$(\eta \equiv 4 \epsilon' E'_p / m_p^2)$$





# relation between neutrino and photon fluxes

$$\pi^\pm \rightarrow \mu^\pm + \bar{\nu}_\mu$$

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

$$\pi^0 \rightarrow \gamma + \gamma$$

$$E_\pi \simeq 2 E_\gamma \simeq 4 E_\nu$$

in hadronic processes, neutrino and photon fluxes intimately related

$$\frac{1}{3} \sum_\alpha q_{\nu_\alpha}(E_\nu) \simeq q_{\pi^\pm}(4 E_\nu) \frac{dE_\pi}{dE_\nu} \simeq 4 K_\pi q_{\pi^0}(4 E_\nu) \simeq 2 K_\pi q_\gamma(2 E_\nu) \frac{dE_\gamma}{dE_\pi} \simeq K_\pi q_\gamma(2 E_\nu)$$

$$K_\pi \equiv \frac{q_{\pi^\pm}}{q_{\pi^0}}$$

$q_i(E)$  is differential source generation rate of species  $i$

similarly:

$$\frac{1}{3} \sum_\alpha [E_\nu^2 q_{\nu_\alpha}(E_\nu)] \simeq \frac{K_\pi}{4} [E_\gamma^2 q_\gamma(E_\gamma)]_{E_\gamma=2 E_\nu}$$

prove it

pp production:

$$p + p \rightarrow \frac{N_\pi}{3} (\pi^+ + \pi^0 + \pi^-) + X \quad K_\pi \simeq 2$$

py production:

$$p + \gamma \rightarrow \Delta \rightarrow \pi^+ + n$$

$$p + \gamma \rightarrow \Delta \rightarrow \pi^0 + p$$

$$BR(\Delta \rightarrow \pi^0 + p) = 2/3$$

$$K_\pi = \frac{1}{2} \text{ (if } \Delta \text{ resonance dominates)}$$

→ prove it using that  
 $(\Delta^+, \Delta^0, \Delta^-)$  isospin=3/2  
 $(\pi^+, \pi^0, \pi^-)$   $I=1$   
 $(p, n)$   $I=1/2$

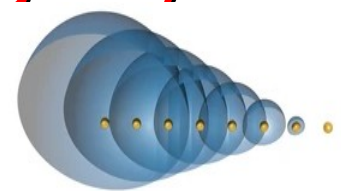
but if multipion production dominates,  $K_\pi \rightarrow 2$

if  $\gamma$  hadronic and source transparent,  $\gamma/\nu$  depends on production mechanism

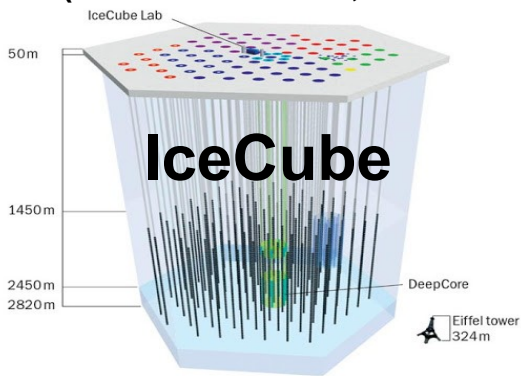


# NEUTRINO TELESCOPES (5 GeV to 10 PeV and beyond)

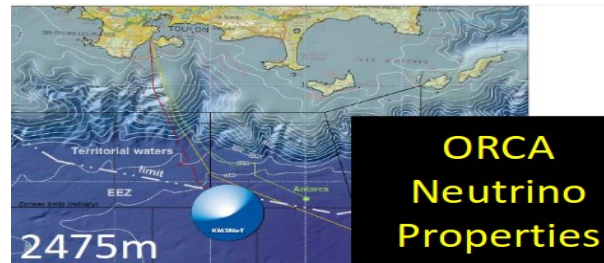
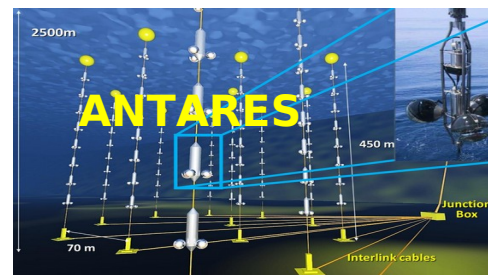
observe Cherenkov light from relativistic charged particles



(DUMAND †1995, AMANDA 1996-2005)

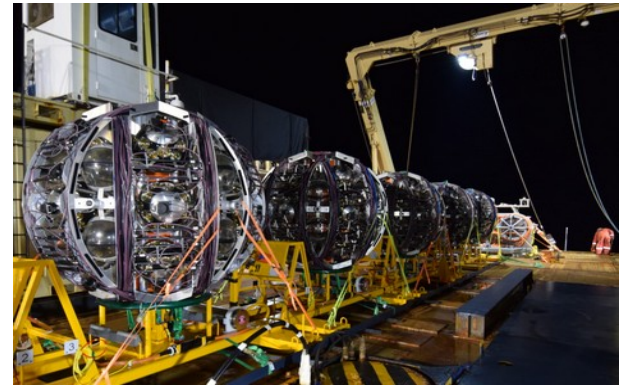
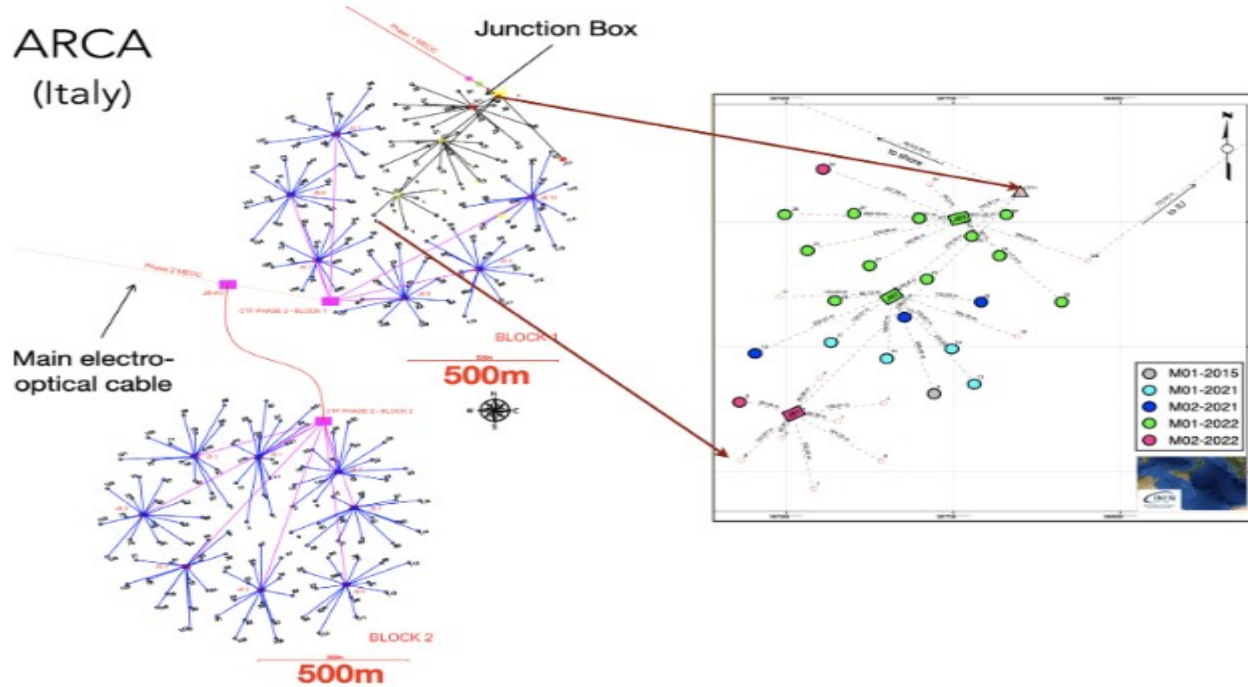


km<sup>3</sup> detector at South Pole,  
completed by 2011,  
looking at northern  $\nu$  sky  
and to southern sky above 100 TeV

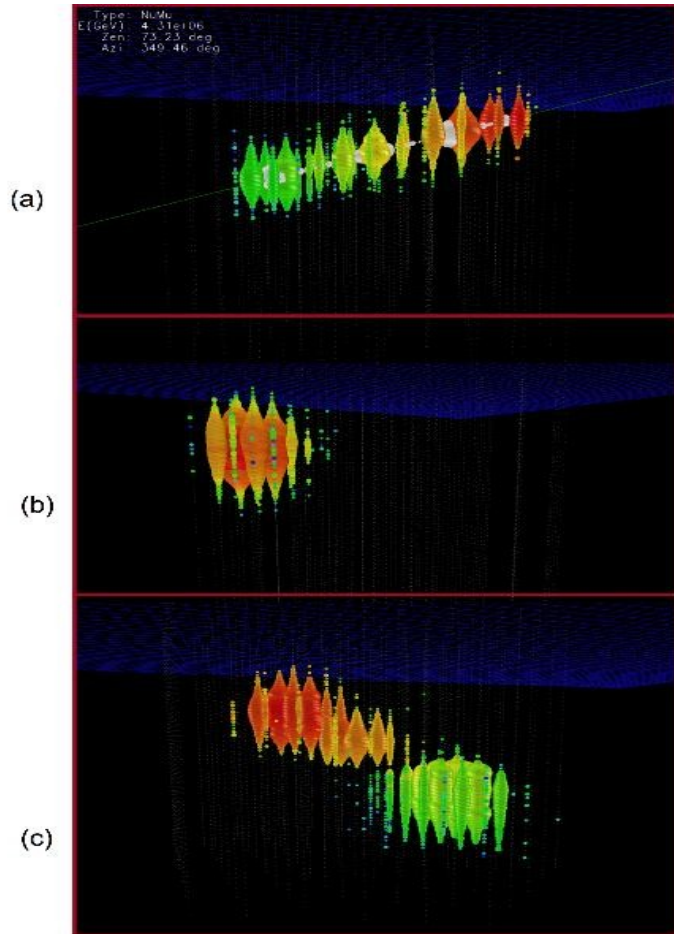


km<sup>3</sup> detector in Mediterranean  
looking at southern neutrino sky  
and northern sky above 100 TeV  
KM3NeT: ORCA & ARCA  
also GVD in lake Baikal

# ARCA deployment



# One may even distinguish neutrino flavors



## muon neutrino (track)

good angular resolution: <1-2 deg

poor E resolution (factor ~2)

muon can be produced outside detector

## electron neutrino (cascade, also from NC)

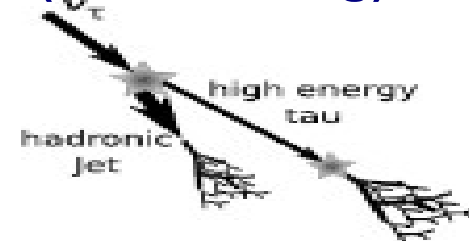
7-15 deg angular resolution (at IC)

15% E resolution

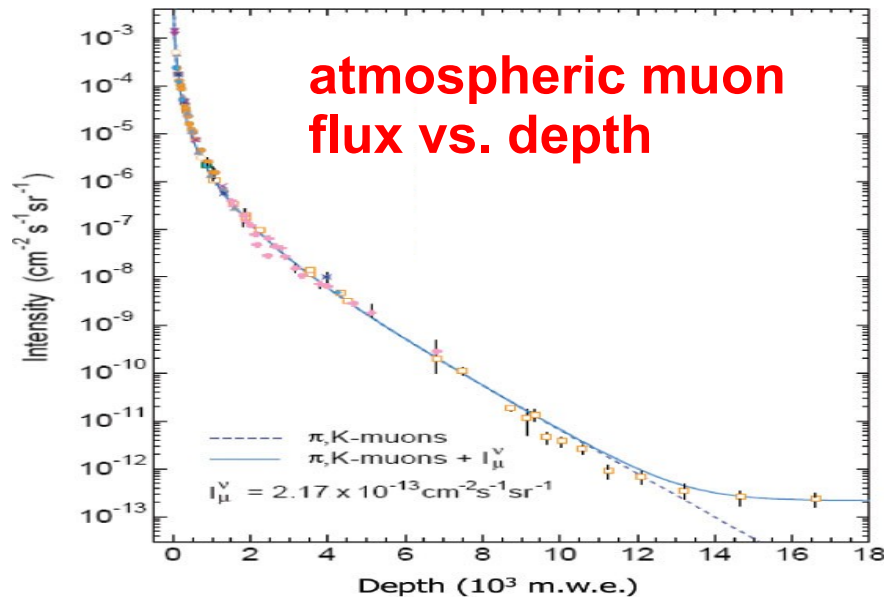
NC showers have only fraction of  $E_\nu$

## tau neutrino (double bang)

$$\gamma c \tau \simeq \frac{E}{\text{PeV}} 50 \text{ m}$$



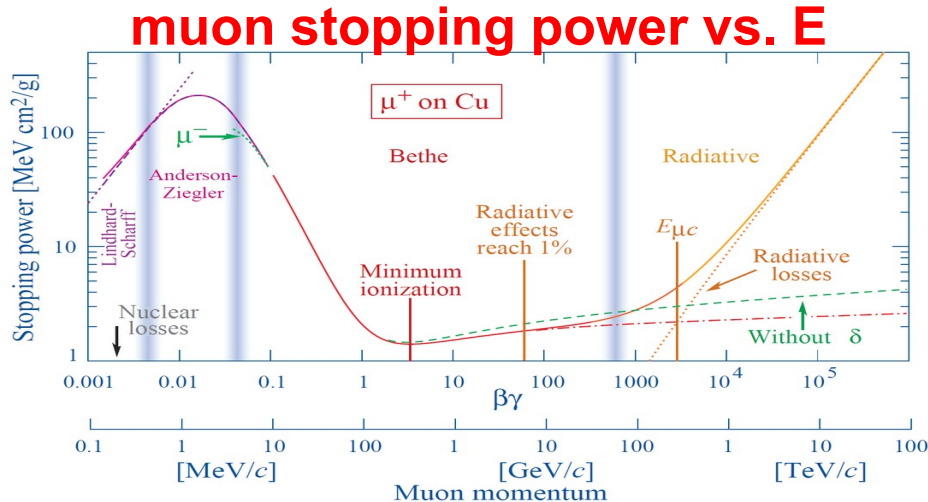
(seen as double pulse cascade for  $E < \text{PeV}$ )



**need to be underground and better look downward**

(at 2km depth, Rate at IC~10<sup>3</sup> Hz)

check it



for  $E_\mu > 0.1 \text{ GeV}$ : 
$$\frac{dE_\mu}{dX} \simeq -\alpha - \beta E_\mu$$

$\alpha \simeq 2 \text{ MeV cm}^2/\text{g}$        $\beta \simeq 4 \times 10^{-6} \text{ cm}^2/\text{g}$

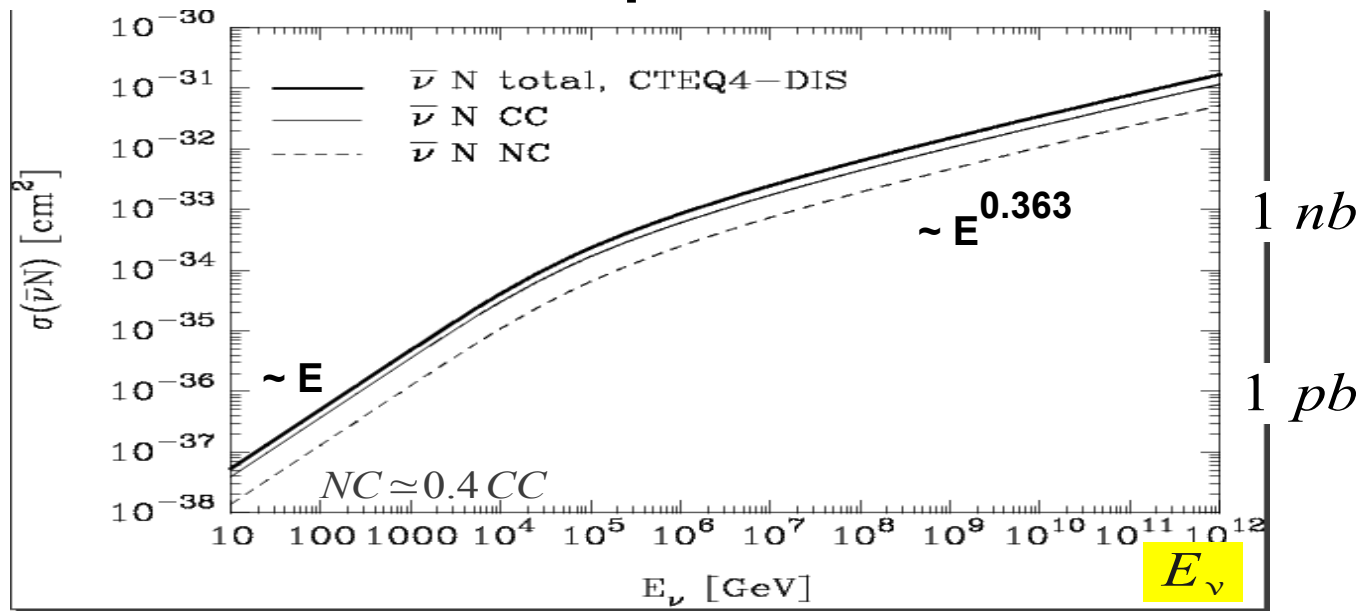
**muon range from  $E_\mu$  to  $E_{th}$ :**

$$d(E_\mu, E_{th}) \simeq D \ln \frac{(1 + E_\mu/\epsilon)}{(1 + E_{th}/\epsilon)}$$
 prove it

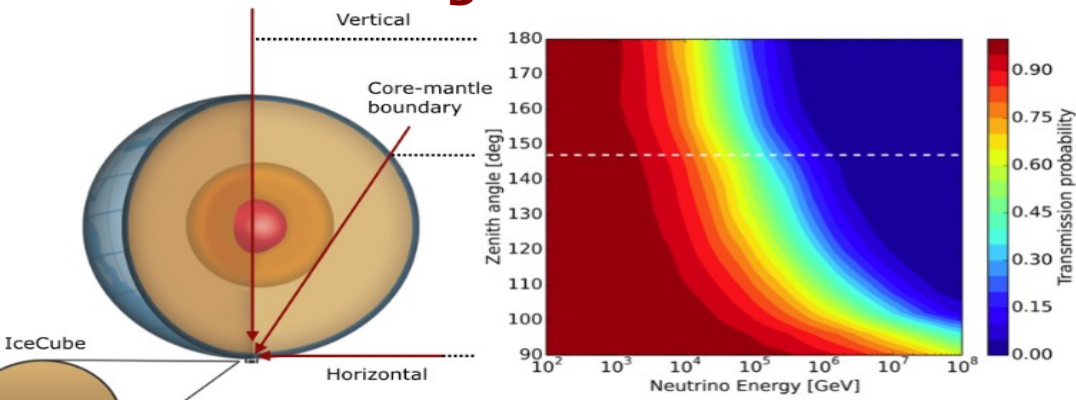
$$D \equiv \frac{1}{(\rho\beta)} \simeq \frac{2.5 \text{ km}}{\rho/\rho_{ice}} \quad \epsilon \equiv \frac{\alpha}{\beta} \simeq 0.5 \text{ TeV}$$

**for  $E > \text{TeV}$ , muons reach detector from several km distance (while e are not MIP)**

# neutrinos detected via deep inelastic $\nu$ nucleon interactions



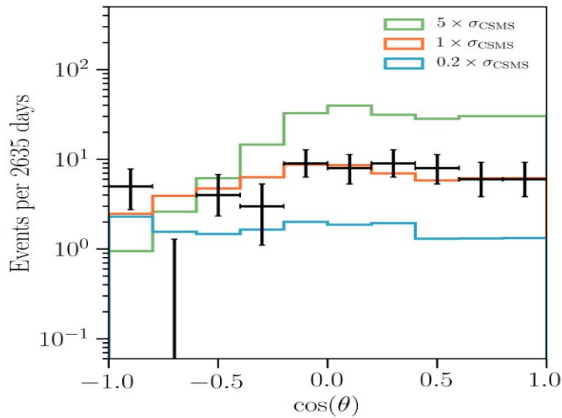
**cross section grows with E ! but Earth becomes opaque for  $E > 100$  TeV**



$$\text{attenuation} = \exp(-\bar{n} \sigma L)$$

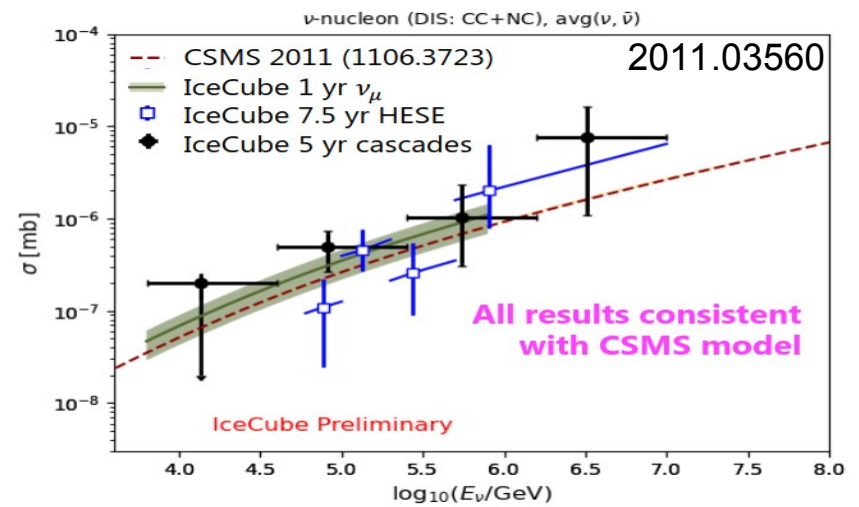
$$\bar{n} \sigma L \simeq \frac{\bar{\rho}}{5 \text{ g/cm}^3} \frac{\sigma}{300 \text{ pb}} \frac{L}{10^4 \text{ km}}$$

**→ need to look above horizon at the highest energies**



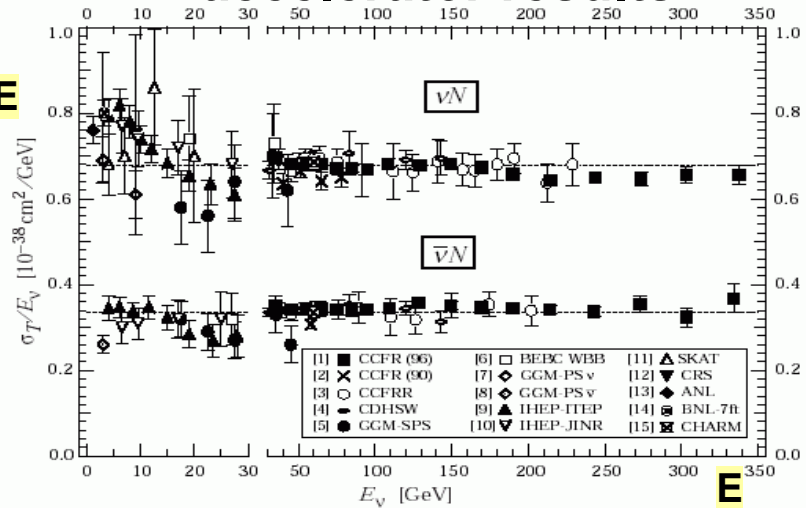
60 HESE events above 60 TeV

attenuation has been used to measure  $\nu N$  cross section beyond reach of accelerators

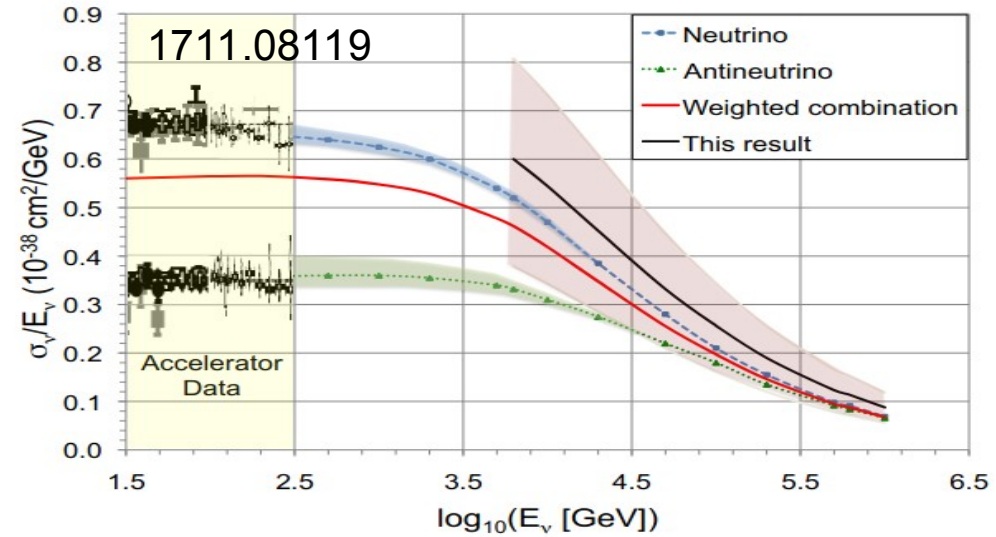


### accelerator results

$\sigma / E$



### IceCube results

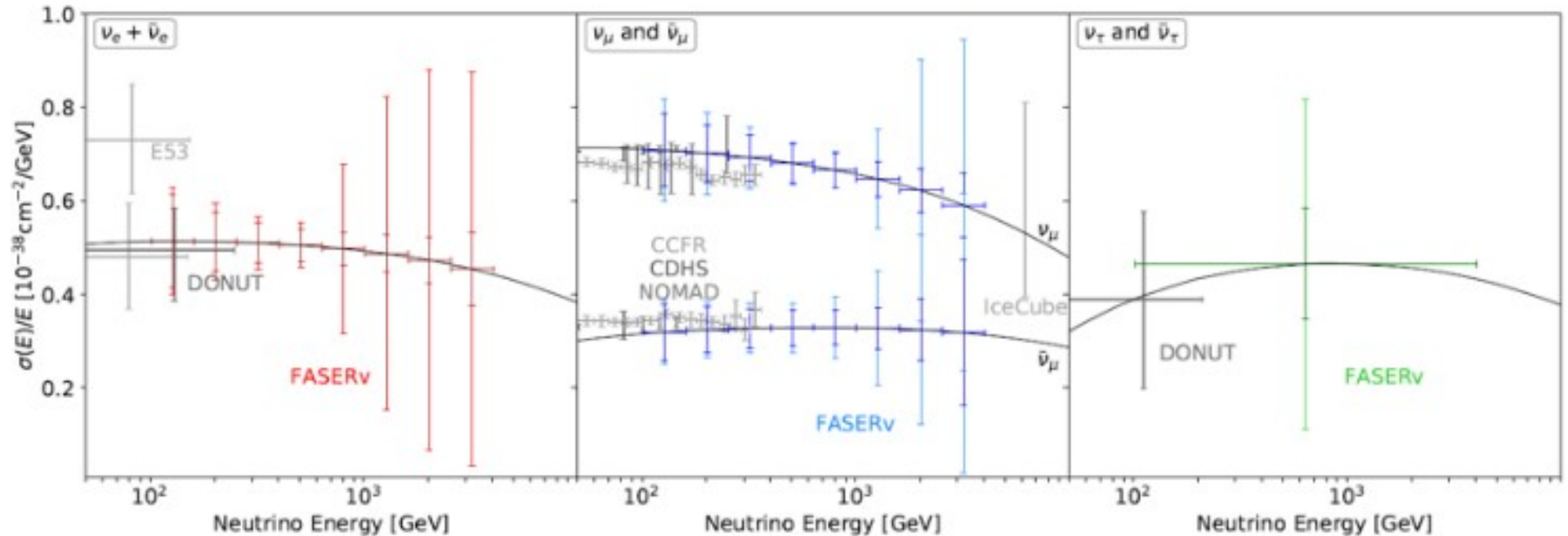


(10784 up-going mus above 6 TeV)

Figure 39.10:  $\sigma_T/E_\nu$ , for the muon neutrino and anti-neutrino charged-current total cross section as a function of neutrino energy. The error bars include both statistical and systematic errors. The straight lines are the averaged values over all energies as measured by the experiments in Refs. [1–4]:  $= 0.677 \pm 0.014$  ( $0.334 \pm 0.008$ )  $\times 10^{-38}$   $\text{cm}^2/\text{GeV}$ . Note the change in the energy scale at 30 GeV. (Courtesy W. Seligman and M.H. Shapov, Columbia University, 2001.)



# LHC will fill the gap (2022-2025)

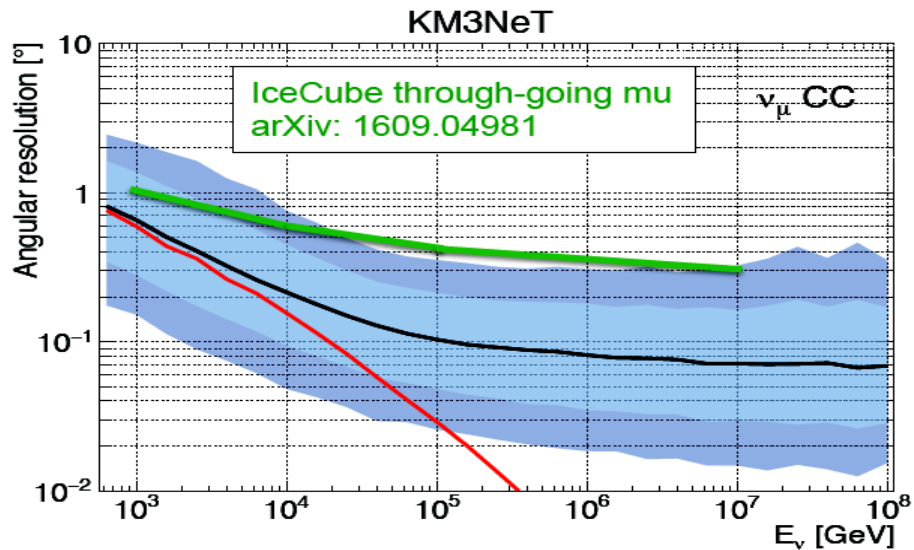


**Figure 3:** The expected sensitivity to  $\nu$ -nucleon CC cross-section for  $\nu_e$  (left),  $\nu_\mu$  (middle) and  $\nu_\tau$  (right) with  $150 \text{ fb}^{-1}$  of data-taking at FASER $\nu$  in LHC Run 3.

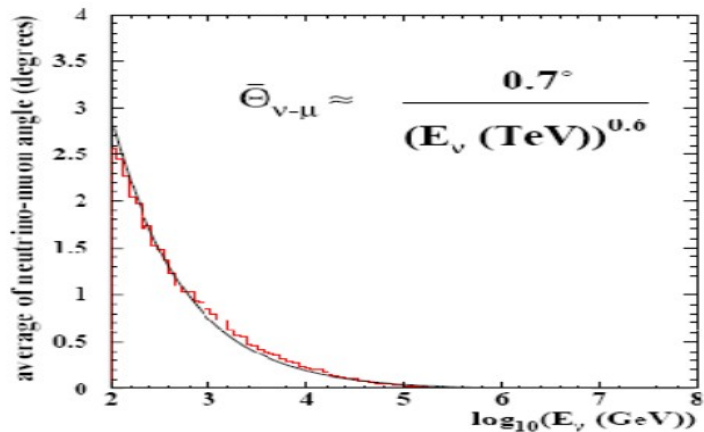
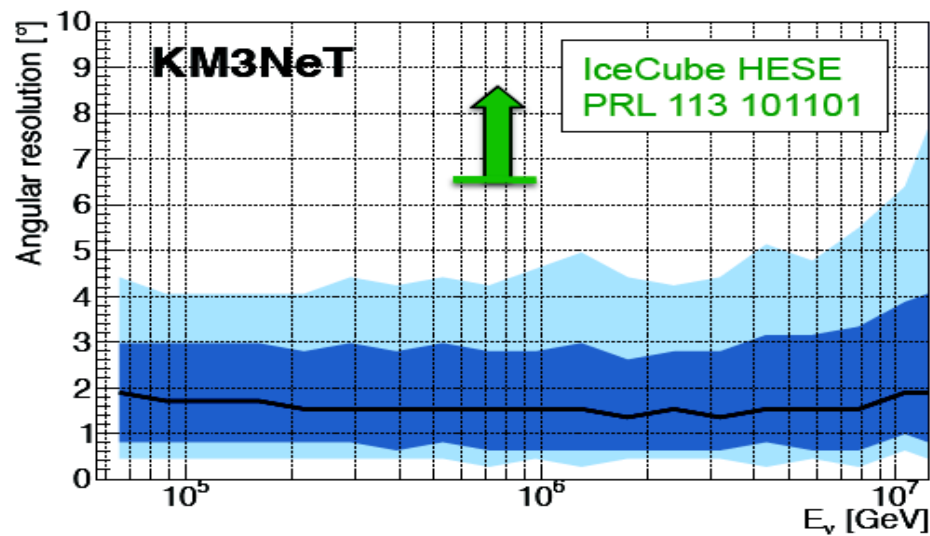


# angular resolution

tracks



showers

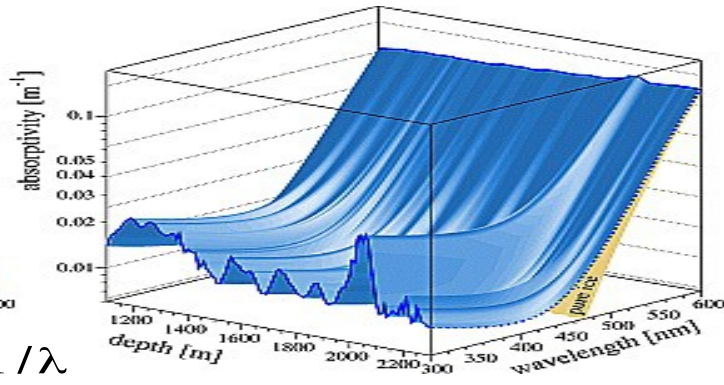
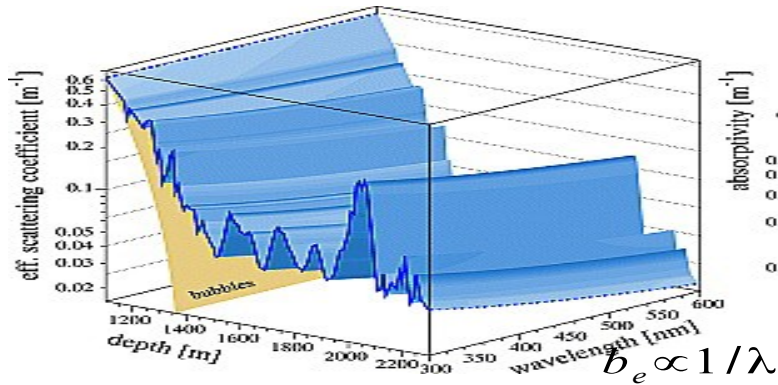
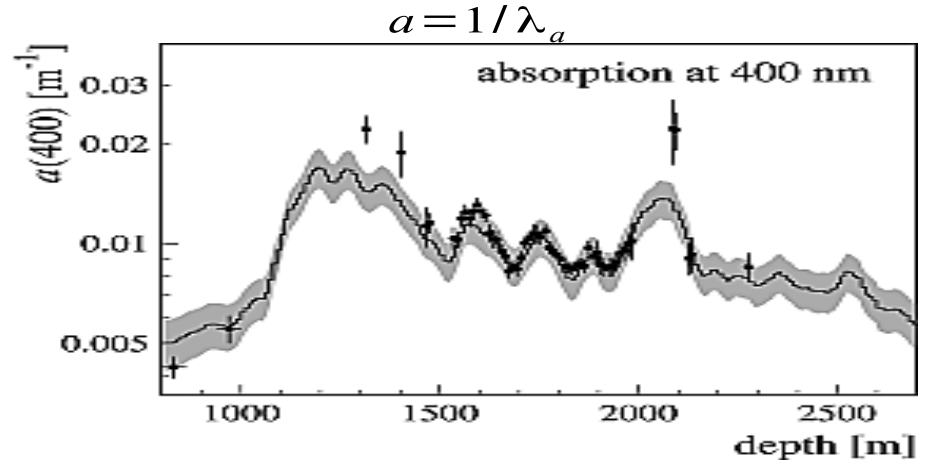
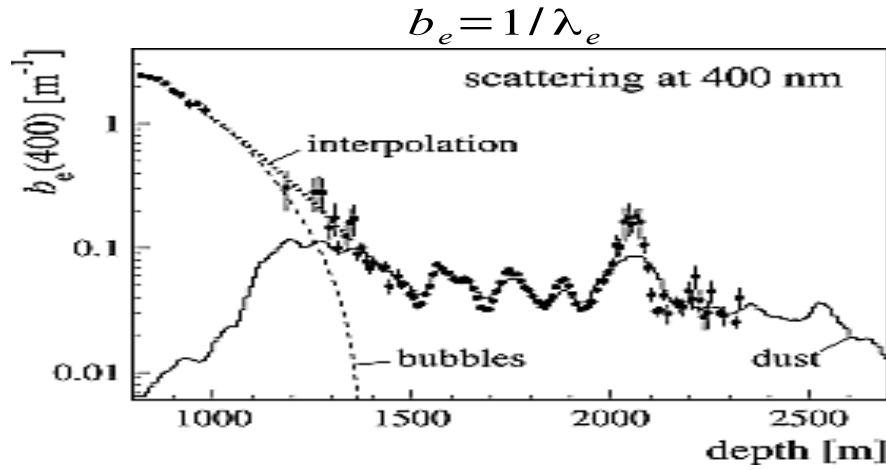


resolution in sea-water (in ice)

tracks ~ 0.1° (> 0.5°)

showers ~ 2° (> 6°)

# ICE SCATTERING AND ATTENUATION



Effective scattering length for isotropization at 400 nm  $\lambda_e \sim 20 - 30 \text{ m}$   
 while  $\lambda_a \sim 100 \text{ m}$

# absorption/scattering in sea water

blue

Toulon

Date	Effective attenuation length (m)	Absorption length (m)	Scattering length (m)
July 1998	60.6±0.4± 5	68.6±1.3±5	265±4±28
Mar. 1999	51.9±0.7±1	61.2±0.7±1	228±11±24
June 2000	46.4±1.9±2	49.3±0.3±2	301±3±27

Table 6-1: Summary of results obtained for the attenuation lengths at the Toulon site using blue light. The first error is statistical and the second one is systematic.

UV

Date	Effective attenuation length (m)	Absorption length (m)	Scattering length (m)
July 1999	21.9±0.8±2	23.5±0.1±2	119±2±10
Sept. 1999	22.8±0.3±2	25.6±0.2±2	113±3±10
June 2000	26.0±0.5±1	28.9±0.1±1	133±3±12

Table 6-2: Summary of results obtained for the attenuation lengths at the Toulon site using UV light. The first error is statistical and the second one is systematic.

CapoPassero

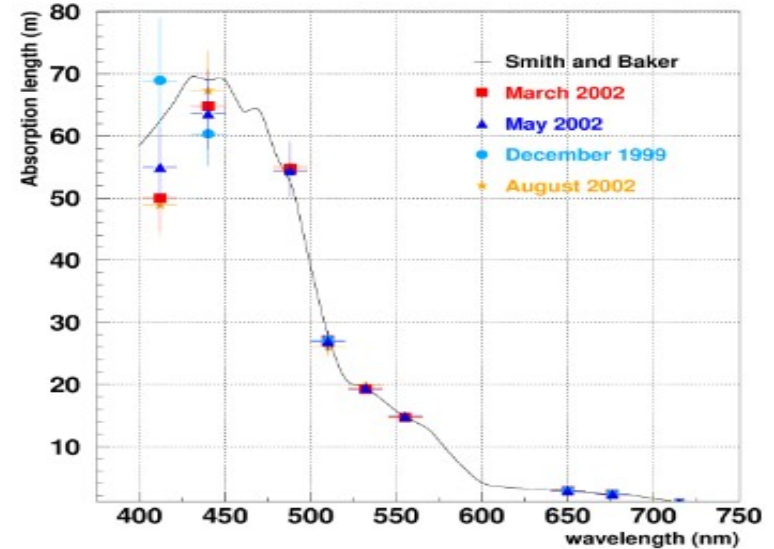


Figure 6-2: Average absorption length as a function of wavelength, for four seasons at the Capo Passero site.

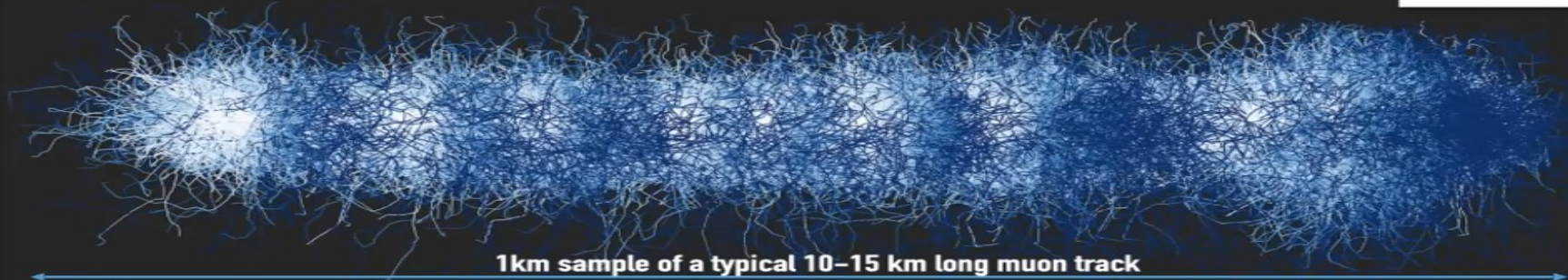
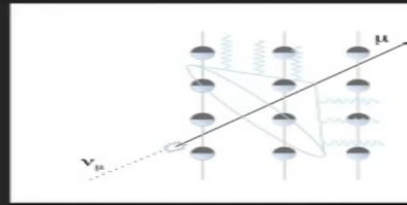
Km3NeT CDR

Scattering is not an issue → better angular resolution in sea water  
but absorption strong → cannot separate PMTs too much

# HORIZONTAL HIGH ENERGY MUONS: THE SIGNATURE

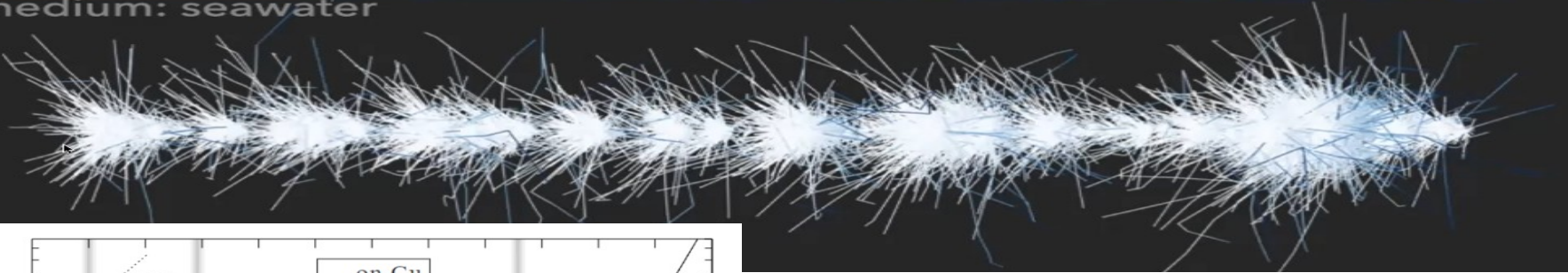
1 PeV horizontal muon

medium: IceCube ice

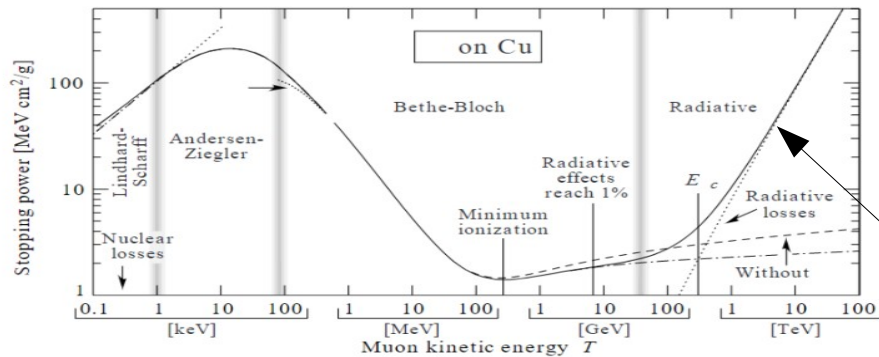


1 km sample of a typical 10-15 km long muon track

medium: seawater



K. Krings



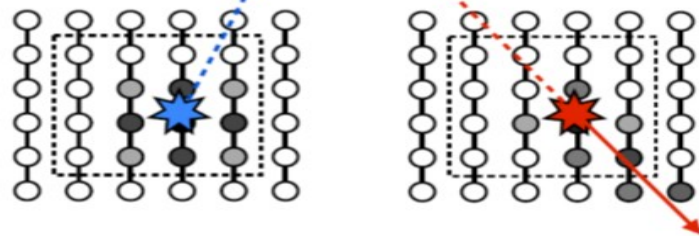
In regime of catastrophic stochastic E losses

At IceCube: downgoing muons  $\sim 10^3$  Hz, atmo  $\nu \sim 10^{-3}$  Hz, astrophysical  $\sim 10^{-5}$  Hz

( $\rightarrow$  need to look for upgoing, or veto downgoing muons, and focus on high E)

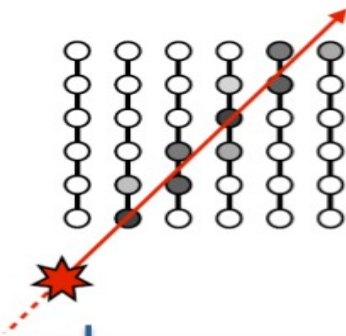
"starting" events

using edge to VETO



upgoing tracks

using Earth as filter

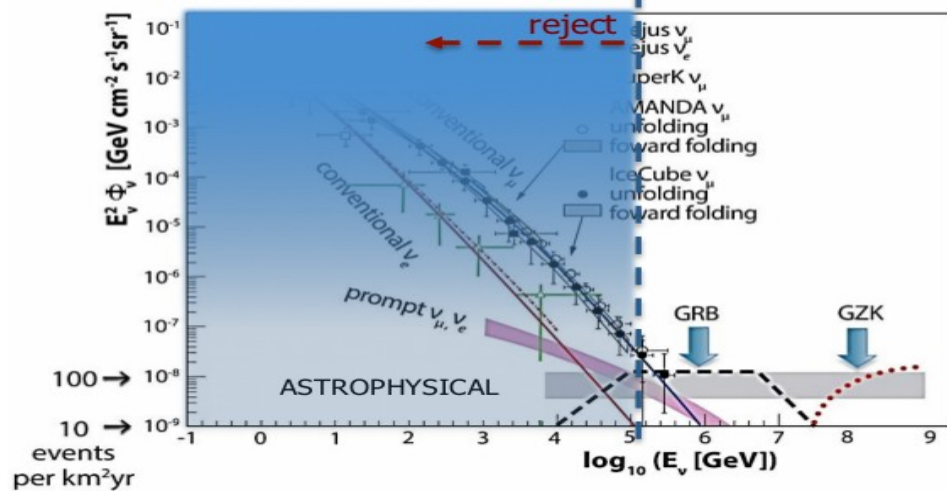


suppress atm.  $\mu$   
background

suppress atm.  $\nu$   
background

using ENERGY  
diffuse searches

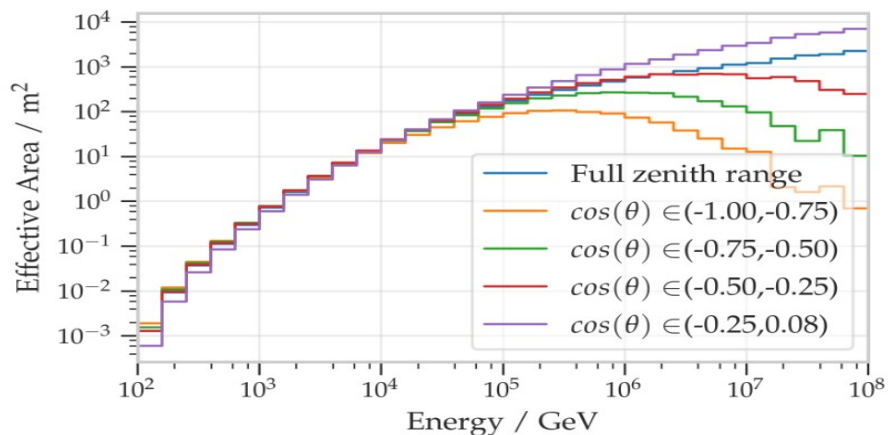
SPATIAL  
TEMPORAL CORRELATIONS  
point sources



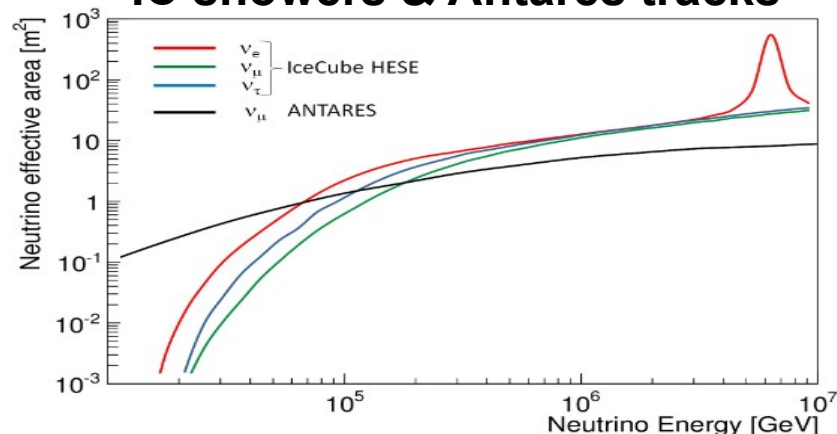
# Effective area of detector:

$$N_{events} = A_{eff} Fluence \quad ; \quad Fluence = \int dt Flux$$

## IC tracks



## IC showers & Antares tracks



depends on cross section, Earth attenuation, fiducial volume, ...

$A_{eff} \sim 1 \text{ m}^2$  at TeV,  $100 \text{ m}^2$  at 100 TeV (tracks)

$\sim 10 \text{ m}^2$  for cascades at PeV

it is like catching water with a strainer

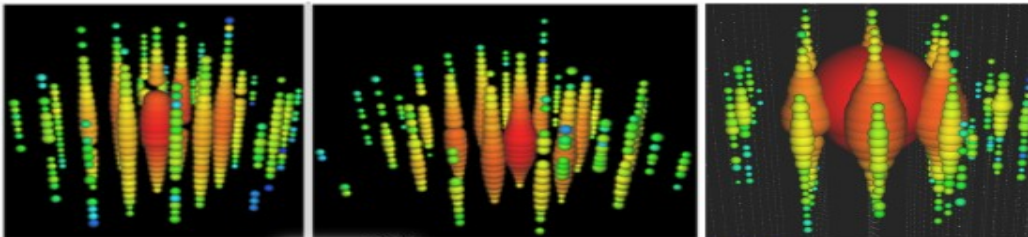
Note that typical astro flux is:

$$E_{\nu}^2 \Phi_{\nu} \sim 10^{-8} \text{ GeV} / (\text{cm}^2 \text{ s sr}) \rightarrow d\Phi / d \ln E \sim 3 (\text{TeV} / E_{\nu}) / (\text{m}^2 \text{ yr sr})$$





# Using the High Energy Starting Events (HESE), IceCube observed PeV neutrinos measuring ~ a dozen per year of astrophysical neutrinos with $E > 60$ TeV



"Bert"  
1.04 PeV  
Aug. 2011



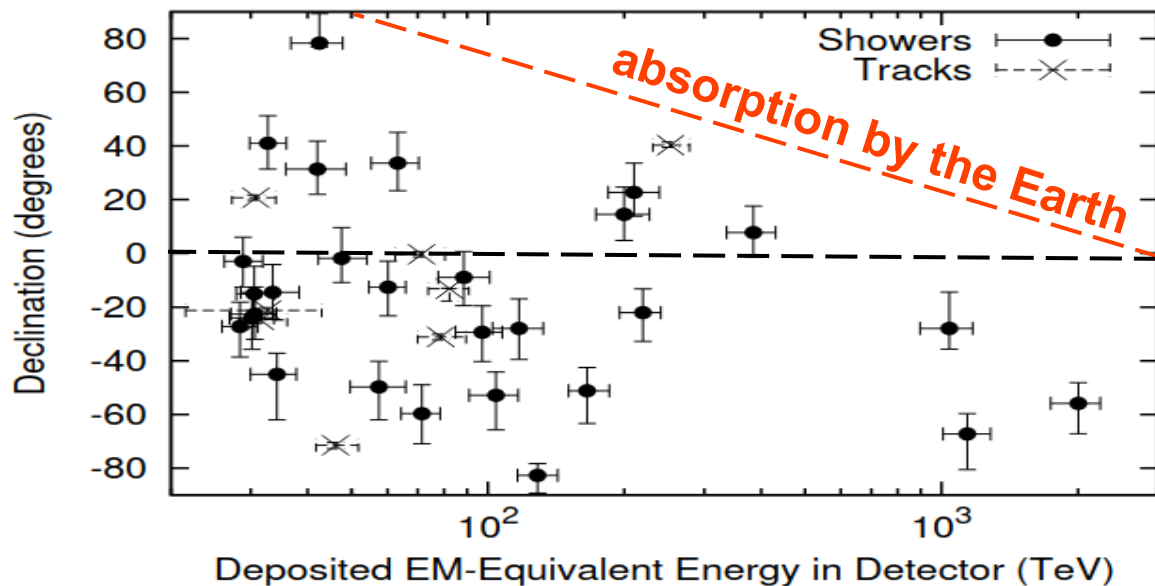
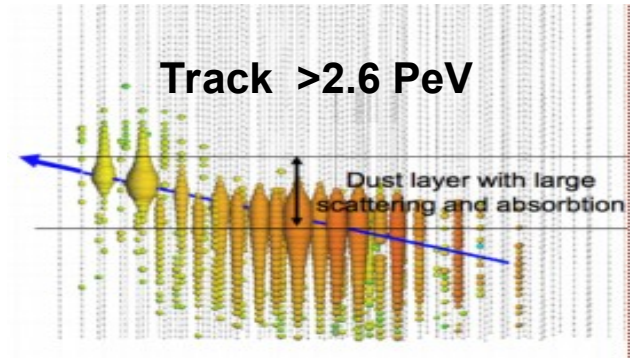
"Ernie"  
1.14 PeV  
Jan. 2012



"Big Bird"  
2 PeV  
Dec. 2012

(Science 2013)  
1405.5303

color: time  
size: E deposit

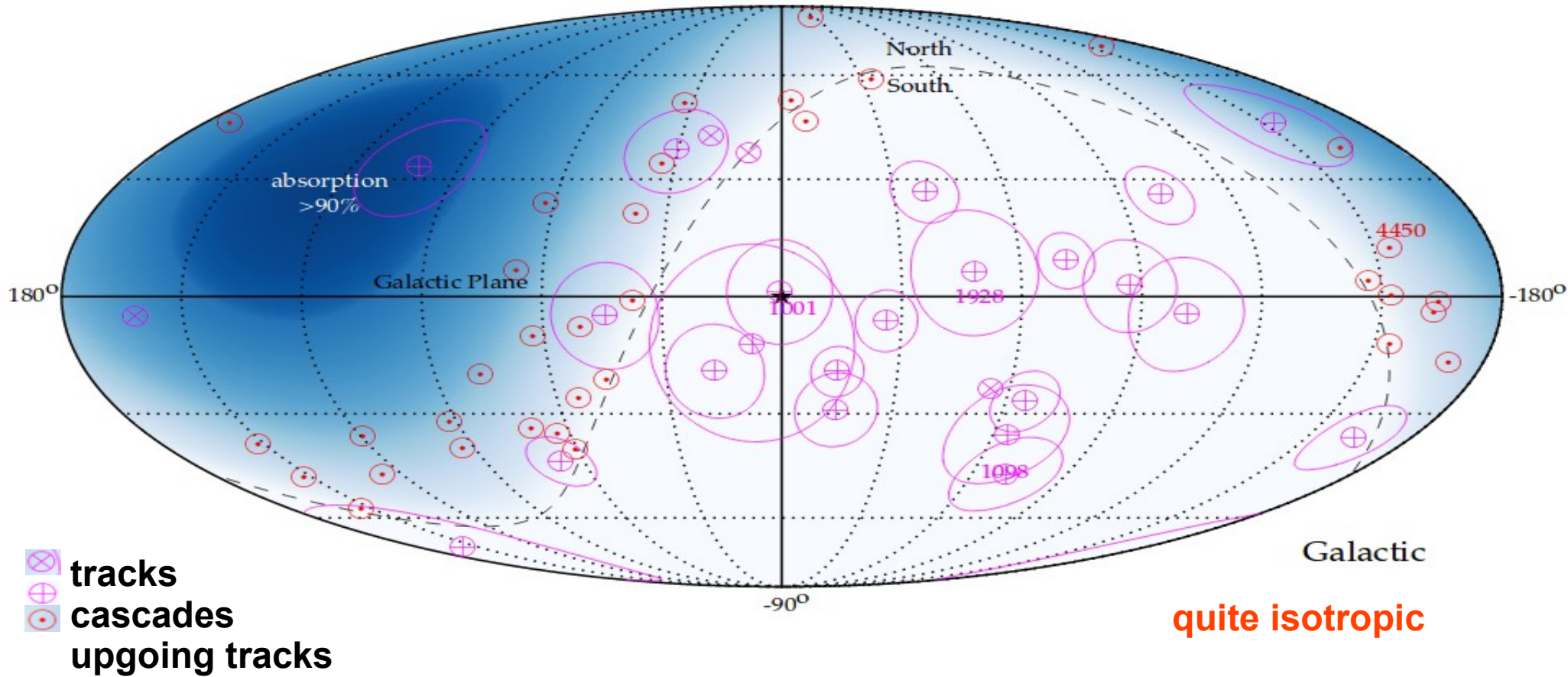


upgoing

downgoing

# Map of astrophysical neutrino arrival directions

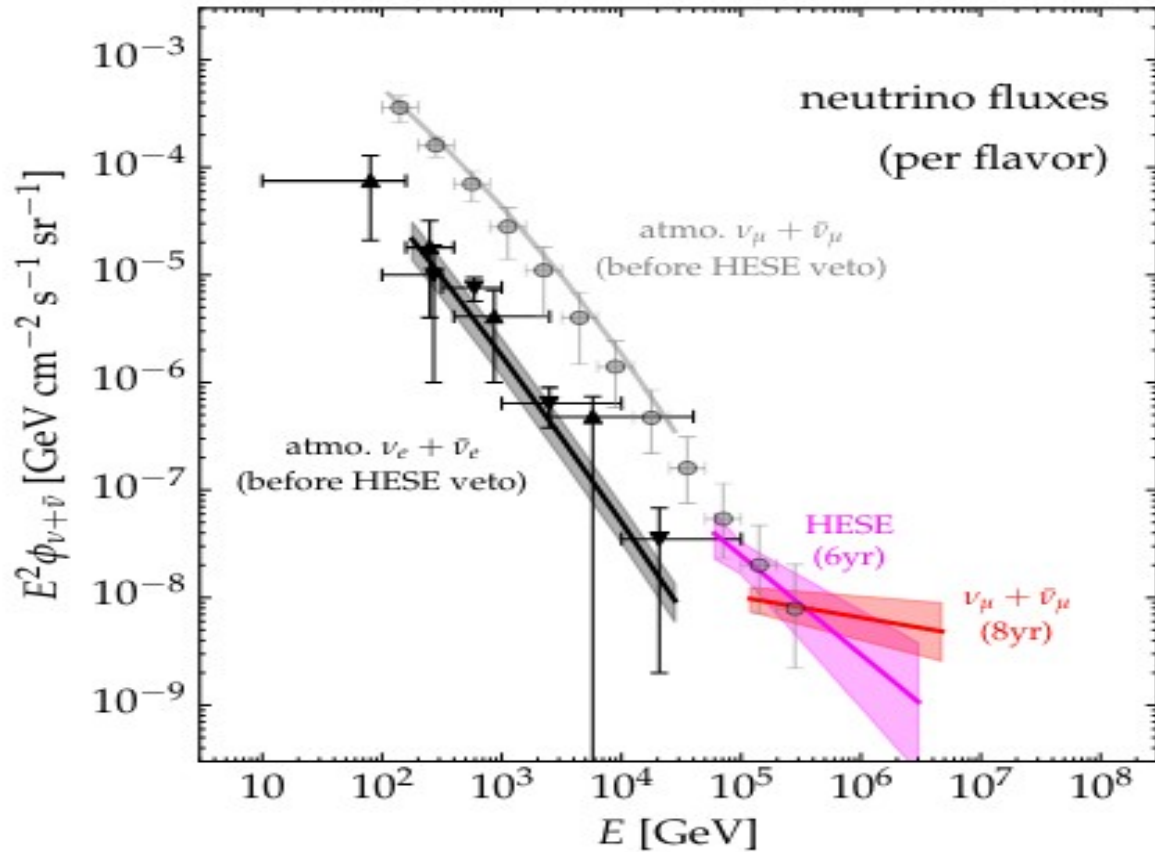
Arrival directions of most energetic neutrino events (HESE 6yr (magenta) &  $\nu_\mu + \bar{\nu}_\mu$  8yr (red))



upgoing tracks with  $E_\mu > 200$  TeV, HESE with  $E > 100$  TeV

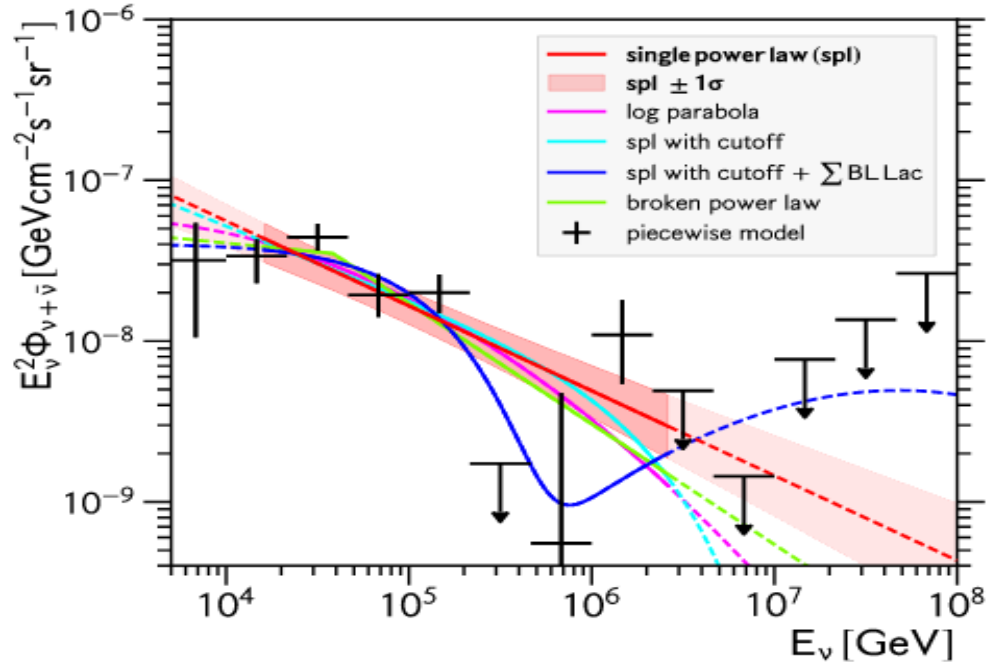
1805.11112

# Spectrum of HESE & throughgoing muons from $\nu_\mu$ CC



# CASCADE EVENTS

From  $\nu_e$  and  $\nu_\tau$  CC + NC from all flavors



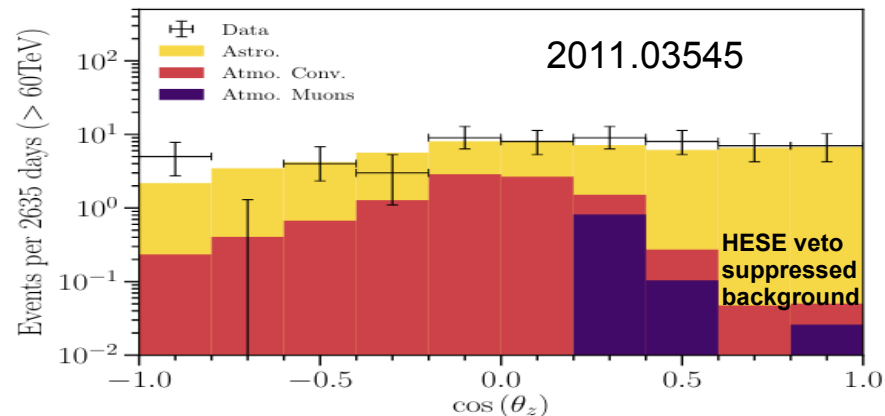
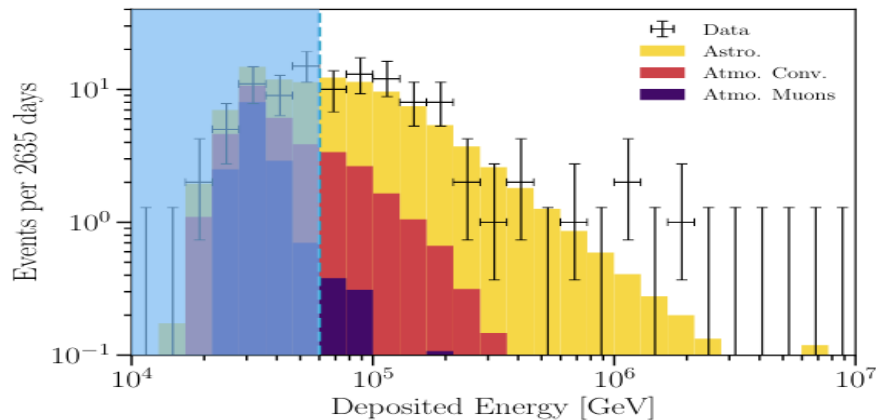
sensitive energy range  
from 16 TeV to 2.6 PeV

small  
atmosph  
backg

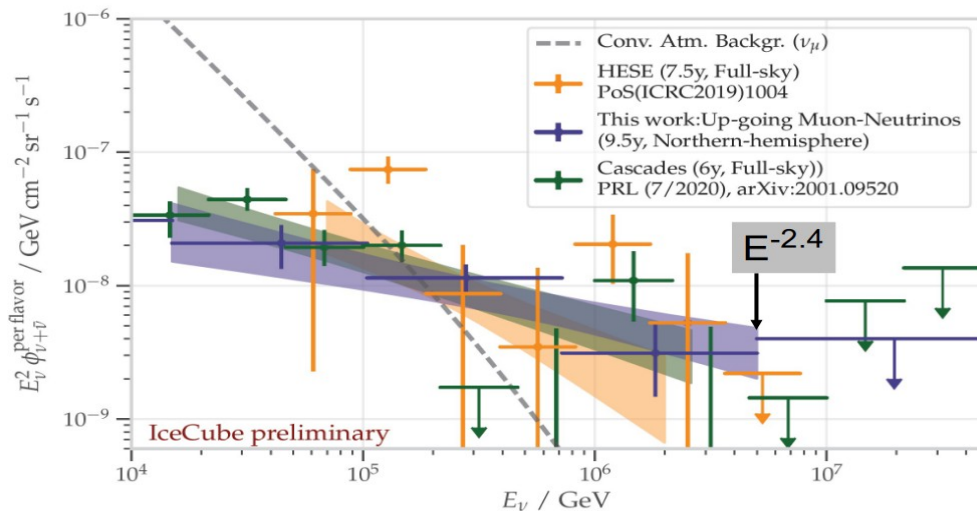
Number of Events	$\nu_e + \bar{\nu}_e$	$\nu_\mu + \bar{\nu}_\mu$	$\nu_\tau + \bar{\nu}_\tau$
astro.	$303^{+46}_{-45}$	$59^{+8}_{-7}$	$204^{+28}_{-27}$
	$(127^{+12}_{-12})$	$(22^{+2}_{-2})$	$(80^{+7}_{-7})$
astro. GR	$0.73^{+0.31}_{-0.22}$	-	-
atmo. conv.	$851^{+23}_{-23}$	$2901^{+64}_{-65}$	-
	$(50^{+3}_{-3})$	$(143^{+8}_{-8})$	-
atmo. prompt	< 192	< 32	-
	(< 57)	(< 7)	-

TABLE III. Number of events for the six years cascade data. The number of astrophysical neutrinos results from the single power law best fit. Numbers of events given in brackets refer to neutrinos with reconstructed energies above 10 TeV. The number of atmospheric tau neutrinos is negligible. Number of Glashow Resonance (astro. GR) events are evaluated assuming  $pp$  type sources in the 4 – 8 PeV energy range.

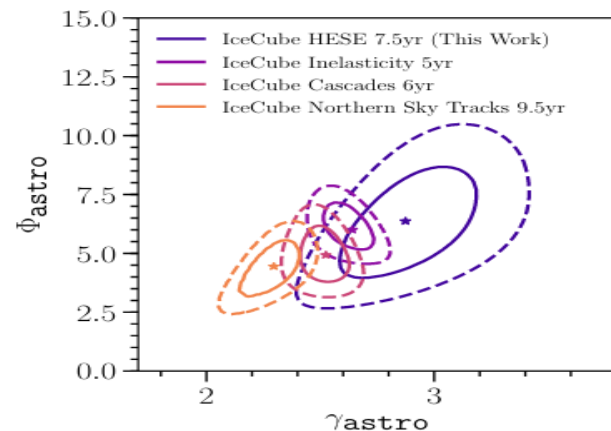
# Latest update



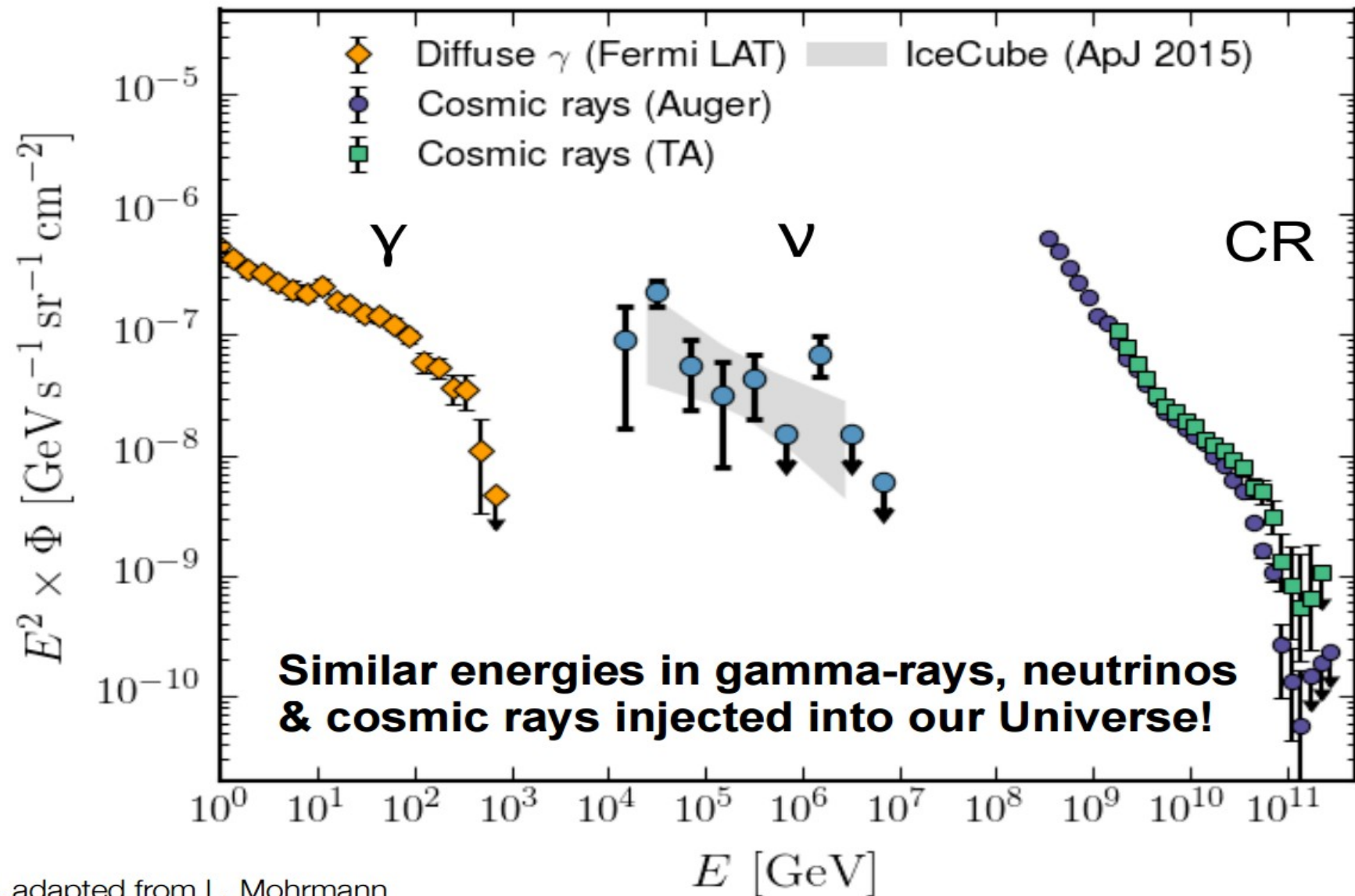
harder than atmospheric background & isotropic



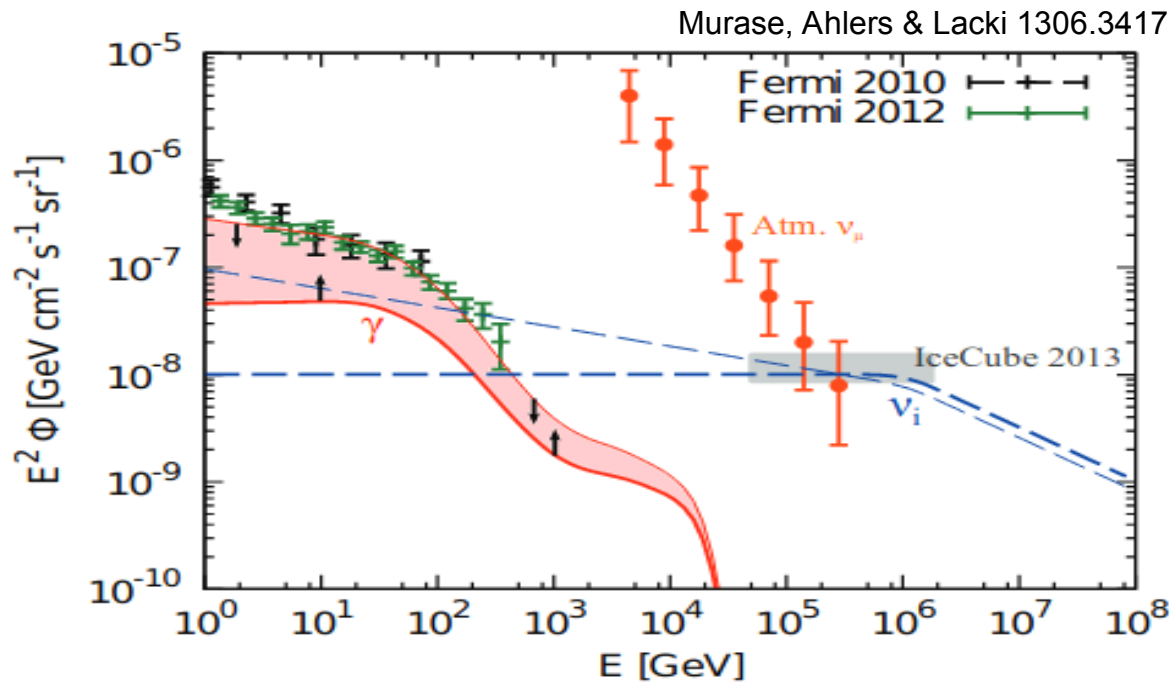
$$\frac{d\Phi_{6\nu}}{dE} = \Phi_{\text{astro}} \left( \frac{E_\nu}{100 \text{ TeV}} \right)^{-\gamma_{\text{astro}}} \cdot 10^{-18} \text{ GeV}^{-1} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$



some tension, but is it a single power-law?  
are fluxes in the north (tracks) equal to the south (seen with cascades)?



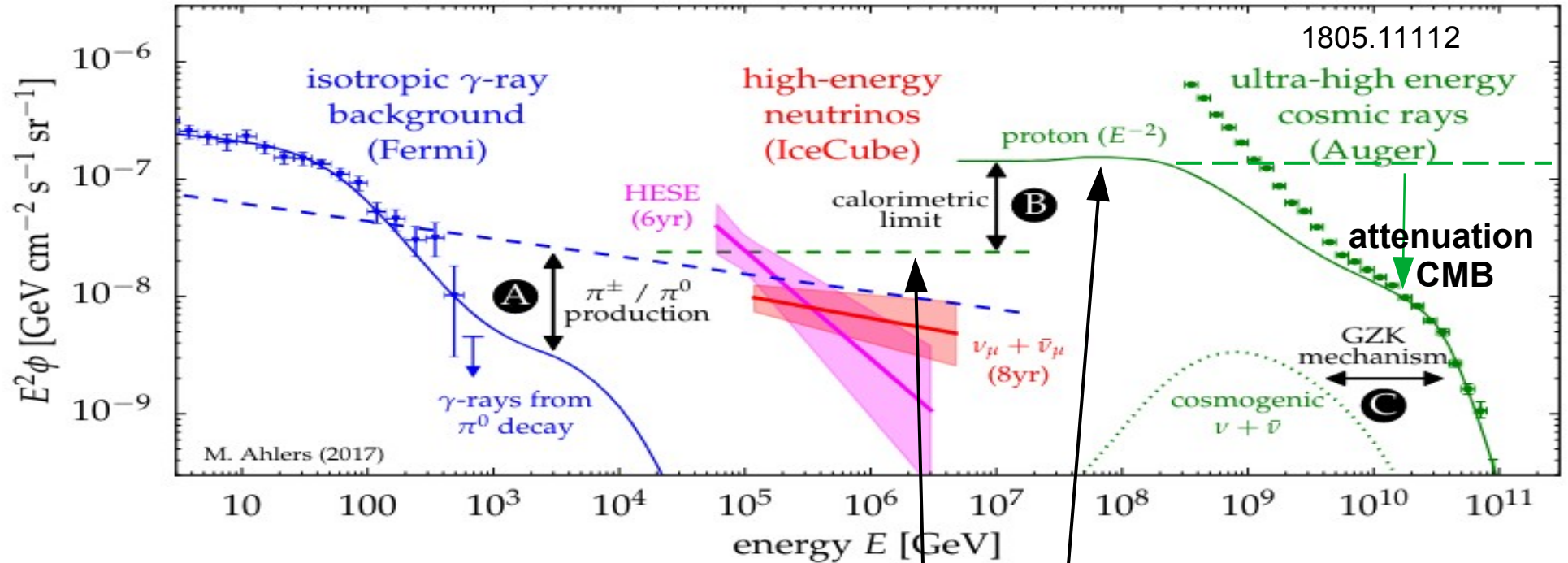
# photons from associated neutral pion decays cascade down below PeV by interactions with CMB and EBL radiation backgrounds



typically the neutrino spectrum below 500 TeV cannot be softer than  $\sim E^{-2.2}$   
to not overshoot gamma background bounds

→ suggest sources opaque to VHE photons (hidden sources)

and/or astrophysical fluxes strongly suppressed below 100 TeV (py scenarios?)



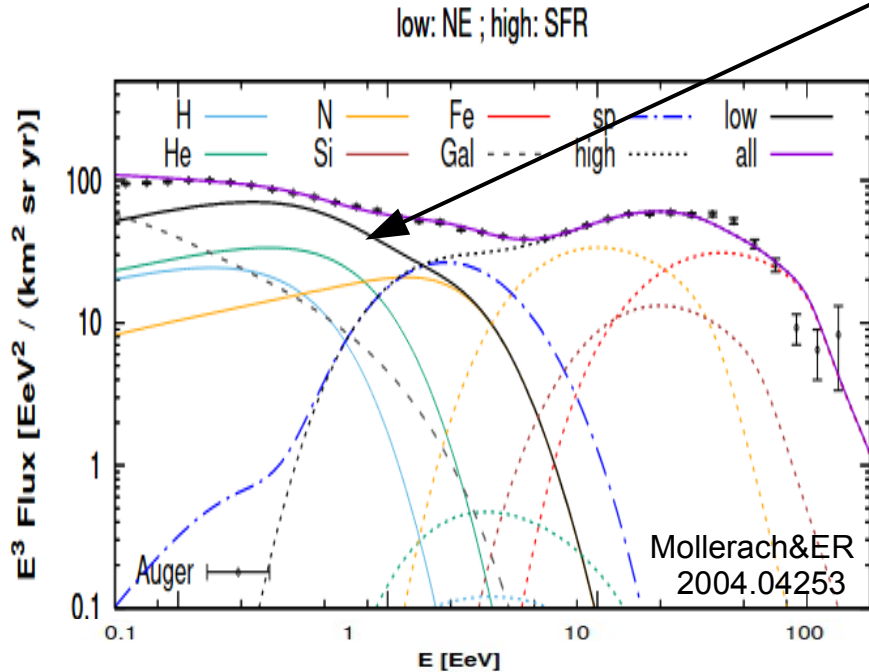
**Waxman-Bahcall 'bound', or reference flux (1999):**  
 if extragalactic UHECRs were protons with  $E^{-2}$  spectrum  
 and significantly converted to pions at PeV energies,  
 the maximum  $\nu$  flux would be  $\sim E^2 \Phi \sim 3 \cdot 10^{-8}$  GeV/cm<sup>2</sup> s sr (all flavors)

observed flux near WB would suggest sources almost opaque to CRs

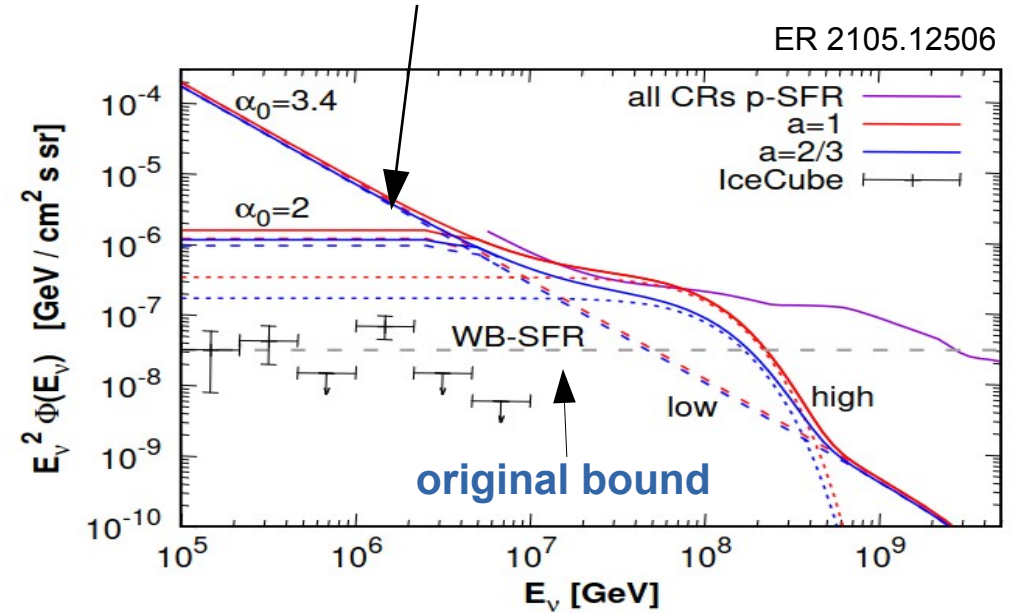


## Waxman-Bahcall flux revisited:

- now we know that UHECRs are not dominantly protons above 5 EeV
- below 5 EeV CRs are likely extragalactic and from different sources



## Upper-bounds in realistic scenarios



realistic scenarios would only require that less than few% of CRs interact at the sources  
the PeV sources can be thin to CRs and still produce observed  $\nu$  fluxes

# IN 2018 ICECUBE OBSERVED 300 TeV NEUTRINO FROM BLAZAR TXS 0506+056 (first neutrino clearly associated to a high-energy extragalactic source)

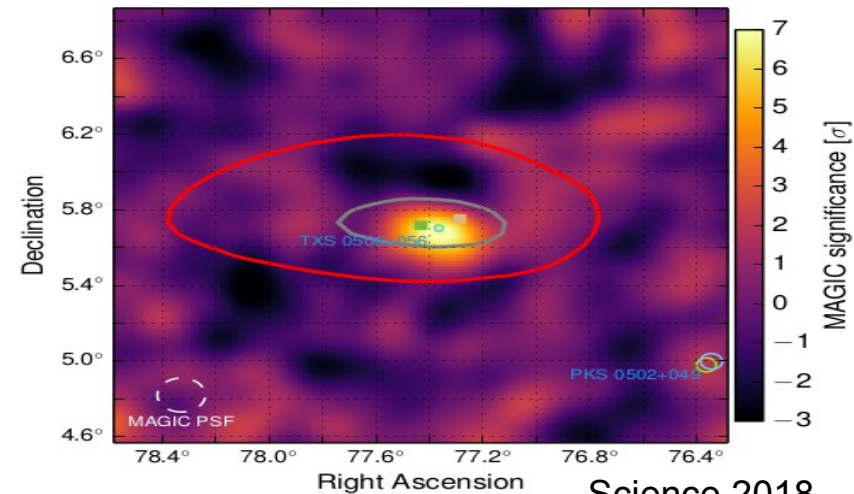
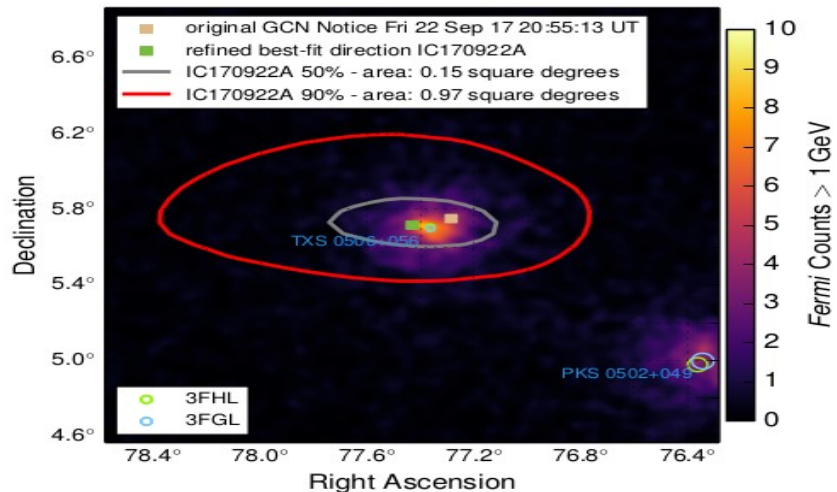
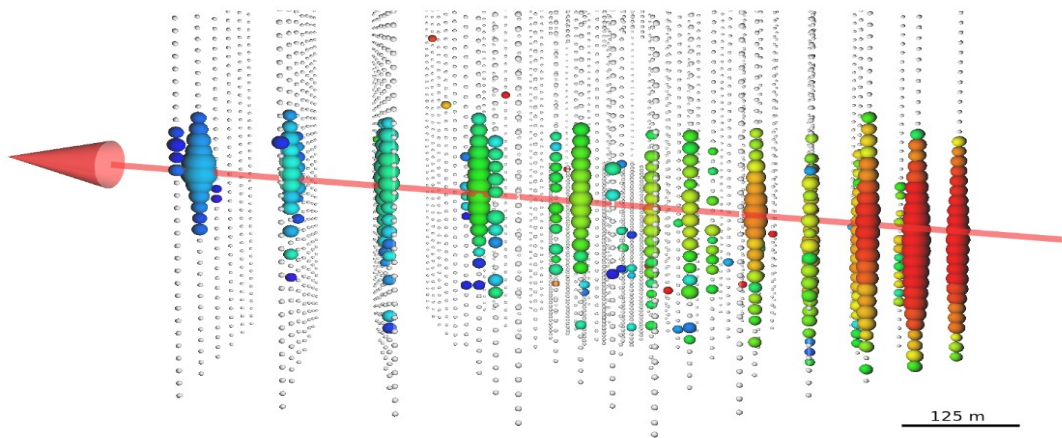
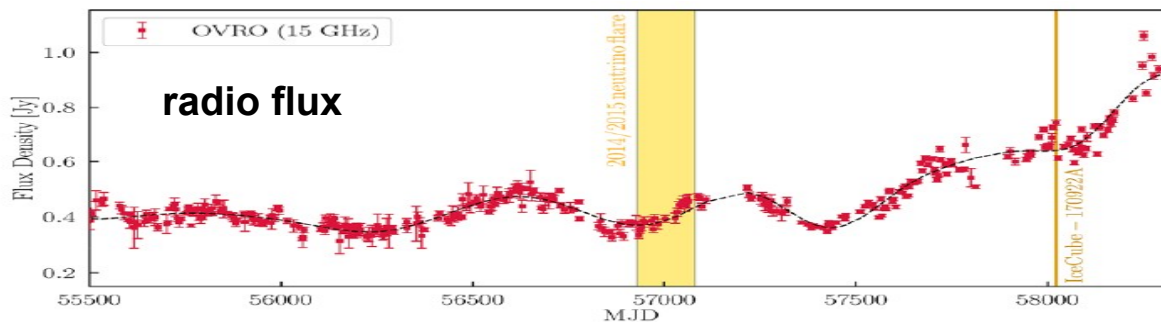
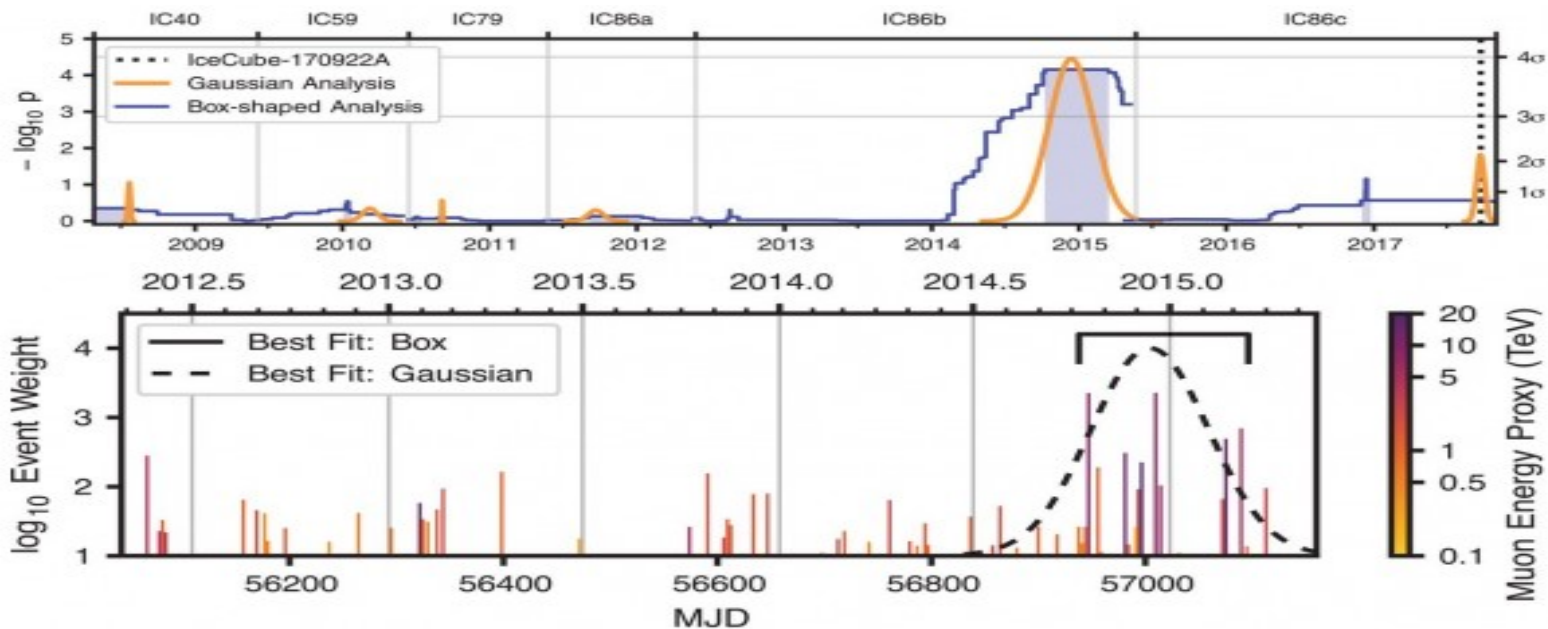


Figure 1: Event display for neutrino event IceCube-170922A. The time at which a DOM

**IceCube event direction coincides  
with Blazar TXS 0506+056,  
that was in a state of increased activity (Fermi, MAGIC)**

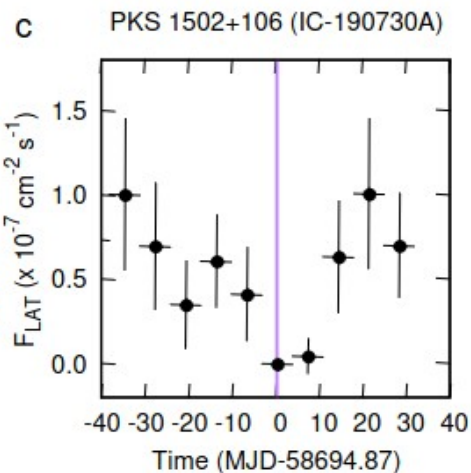
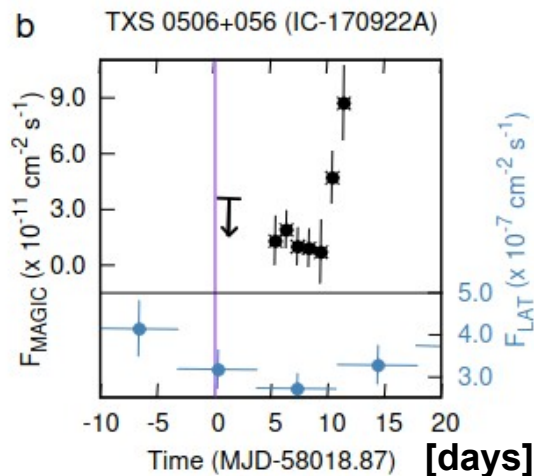
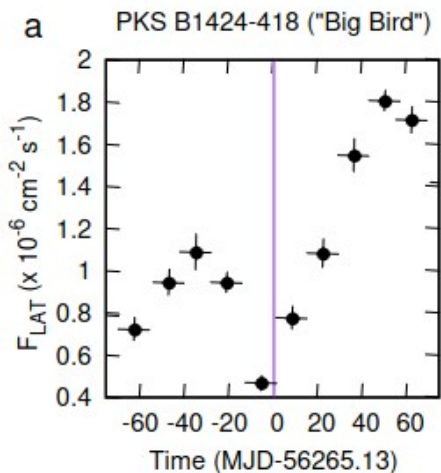
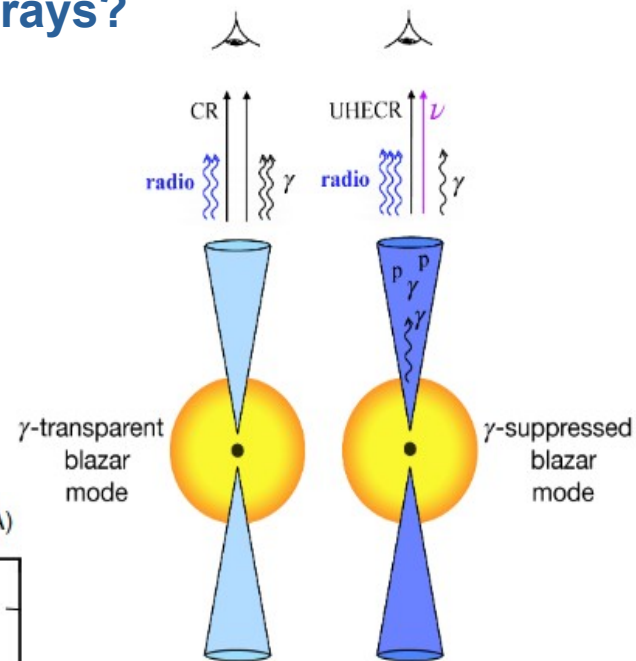
**source redshift  $z = 0.33$  ( $> 1.7$  Gpc!)**

# searching in archival data in the TXS direction: → ~1-20 TeV flare lasting 160 days near end 2014 (13+/-5 events)

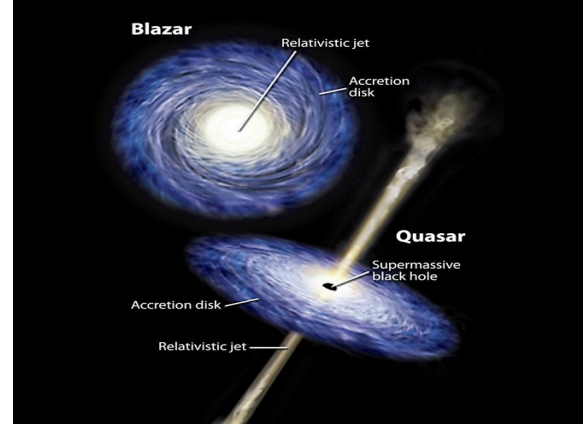
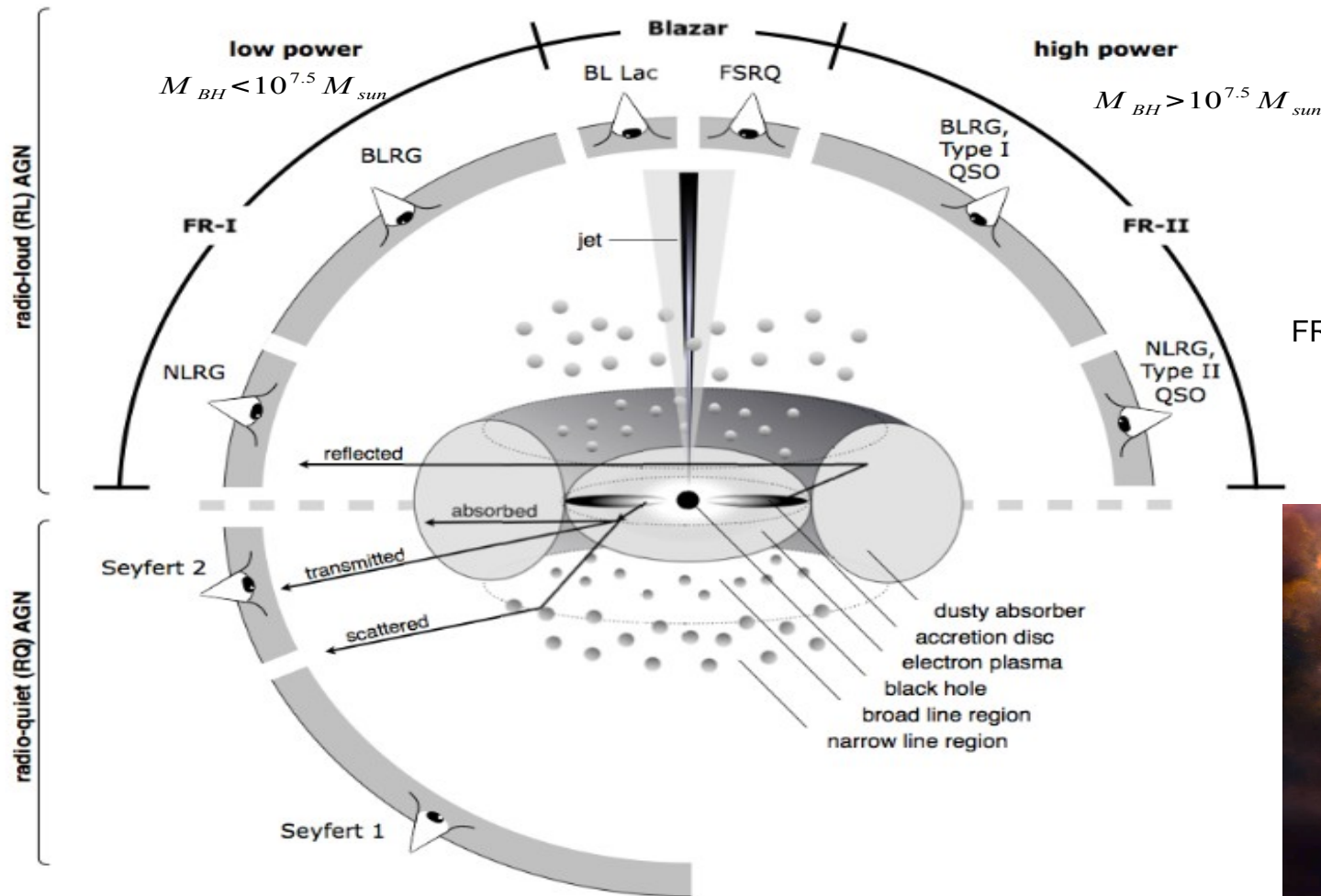


remarkably, the energetic neutrino in TXS (and other less significant blazar candidates) appeared during brief periods of reduced gamma flux: the enhanced target that produced the neutrinos also attenuated the gamma rays?

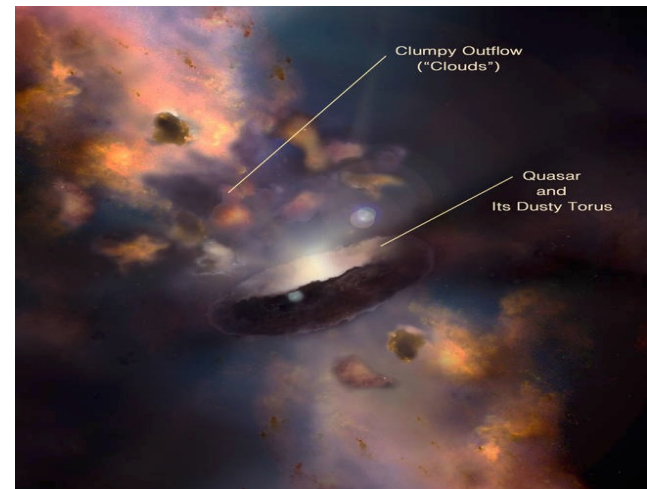
$$\tau_{\gamma\gamma} \approx \frac{\sigma_{\gamma\gamma}}{k_{p\gamma}\sigma_{p\gamma}} \tau_{p\gamma} \approx 10^3 \tau_{p\gamma}$$



# AGN Grand Unification



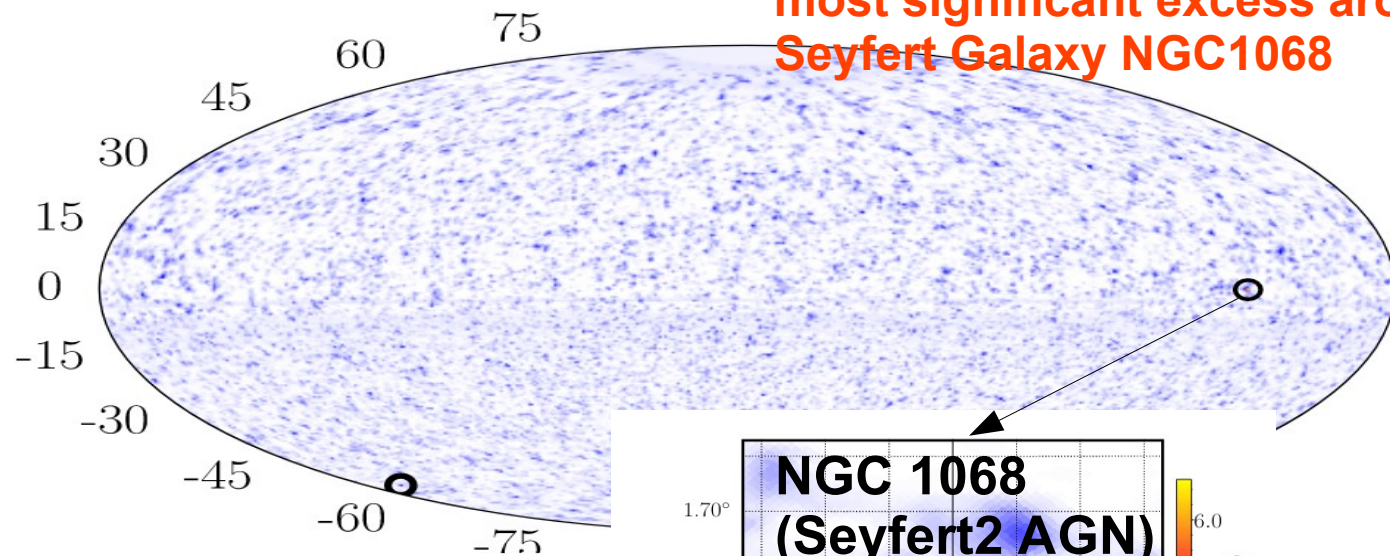
FR: Fanaroff Riley



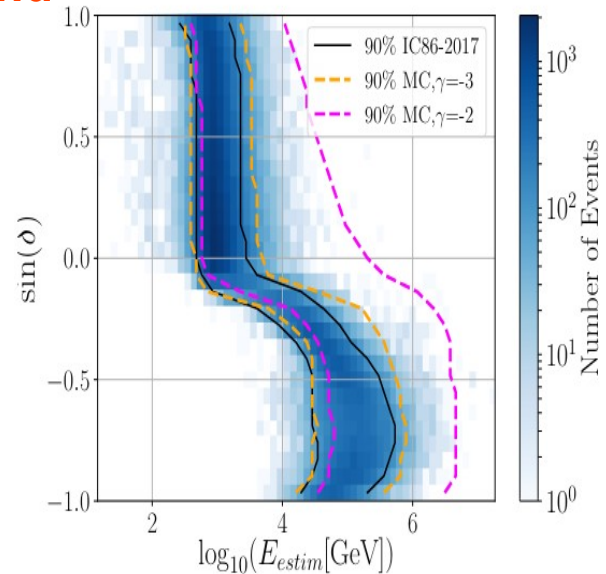
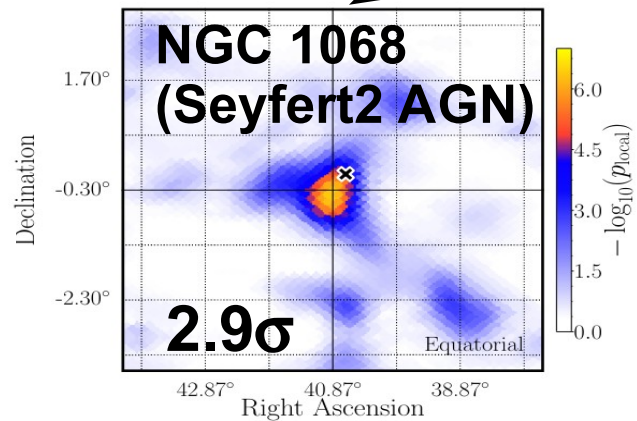
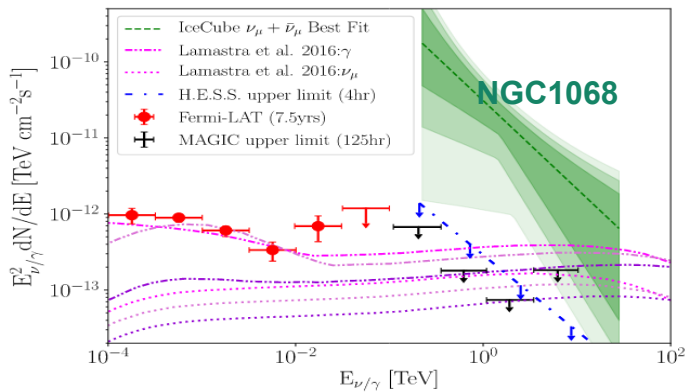
# IceCube SEARCH FOR POINT SOURCES WITH MUON TRACKS

(angular resolution  $< 1^\circ$ )

most significant excess around  
Seyfert Galaxy NGC1068



(about 130k events per yr)



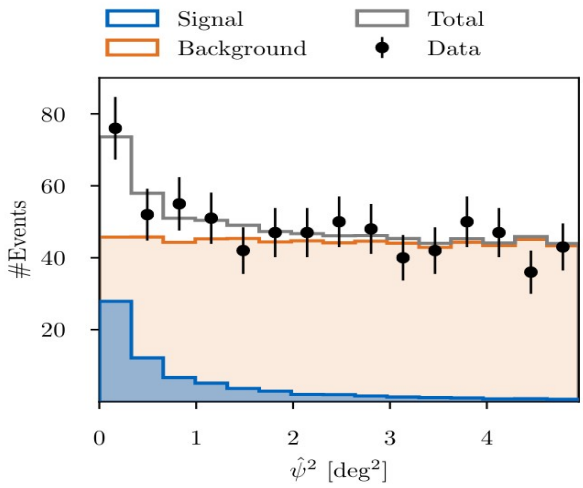
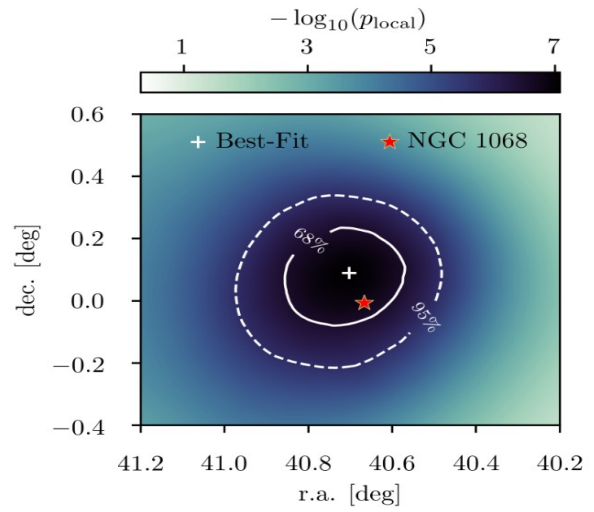
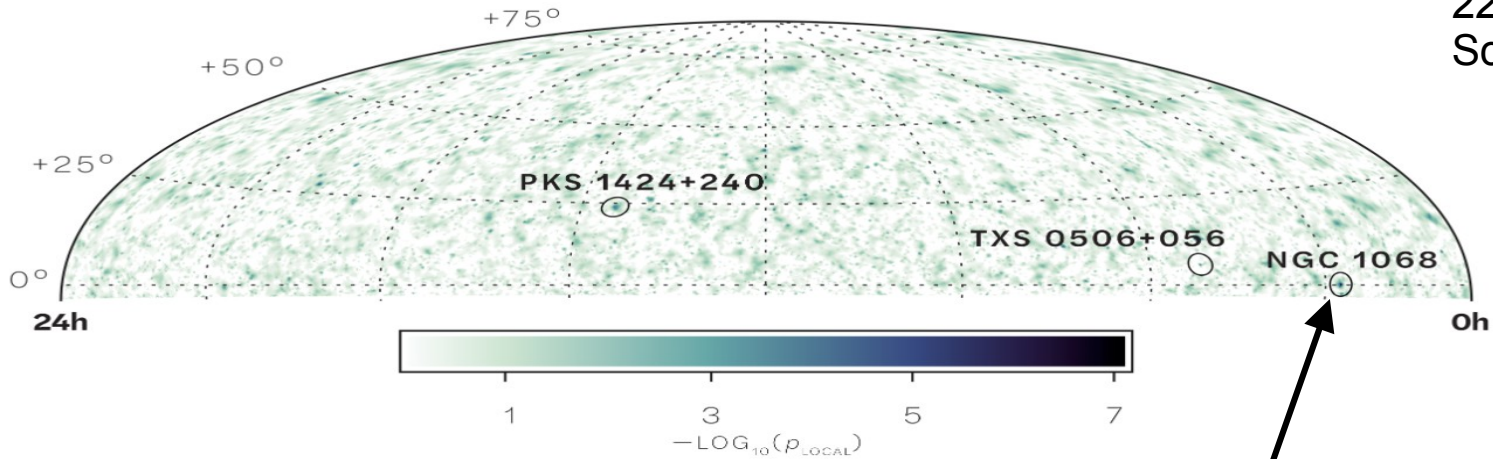
for upgoing,  $E > 1$  TeV,  
backg from atmo neutrinos

for downgoing,  $E > 100$  TeV,  
backg from atmo muons

1910.08488,

# NGC 1068 latest results

2211.09972  
Science 22

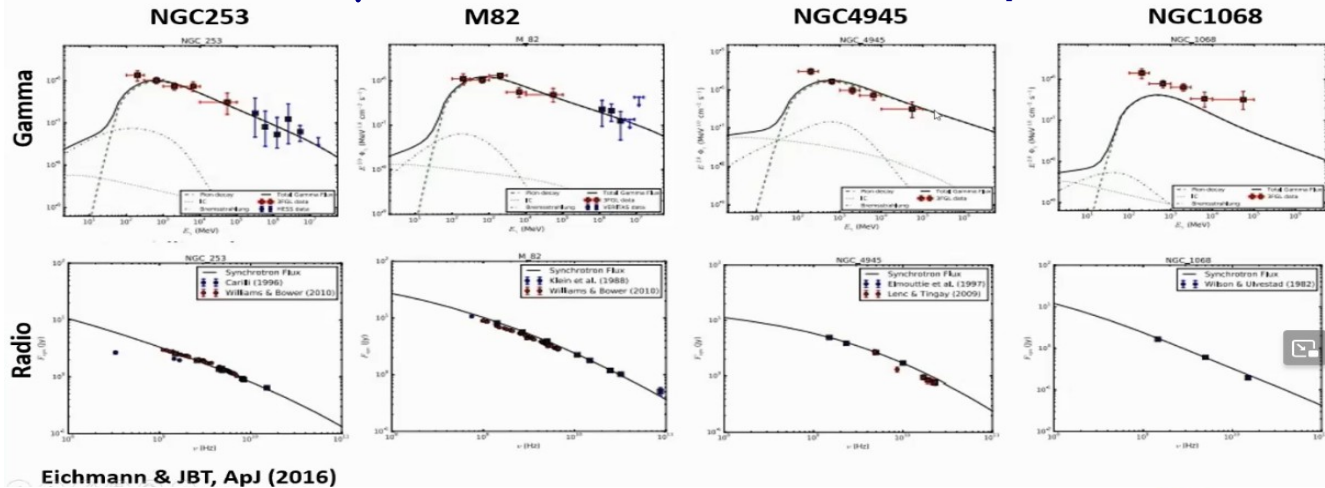


**79+20 excess neutrinos,  
4.2 $\sigma$  in catalog search**



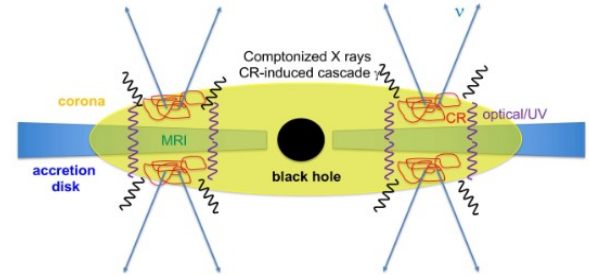
**emission consistent with being steady**

# SBG $\gamma$ emission from radio inferred e & p



Eichmann & JBT, ApJ (2016)

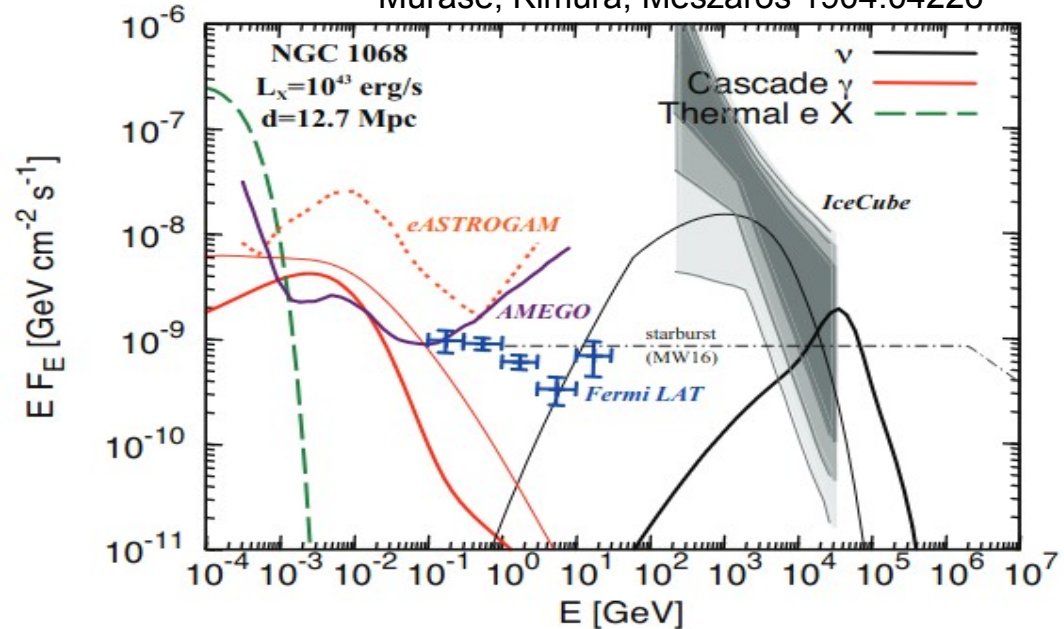
falls short for NGC 1068  
extra emission from corona?



In the case of NGC 1068, the observed gamma-ray flux and the SN rate cannot be described by our model. Here, an SN rate that is about an order of magnitude bigger than the observed one is need in order to accurately describe the gamma-ray data.

gamma-ray data. Thus, there needs to be another relativistic proton source that dominates the acceleration of CRs within this starburst. Here, the active galactic nucleus and the related jet-driven particle acceleration are promising source candidates. In the following, we

Murase, Kimura, Meszaros 1904.04226

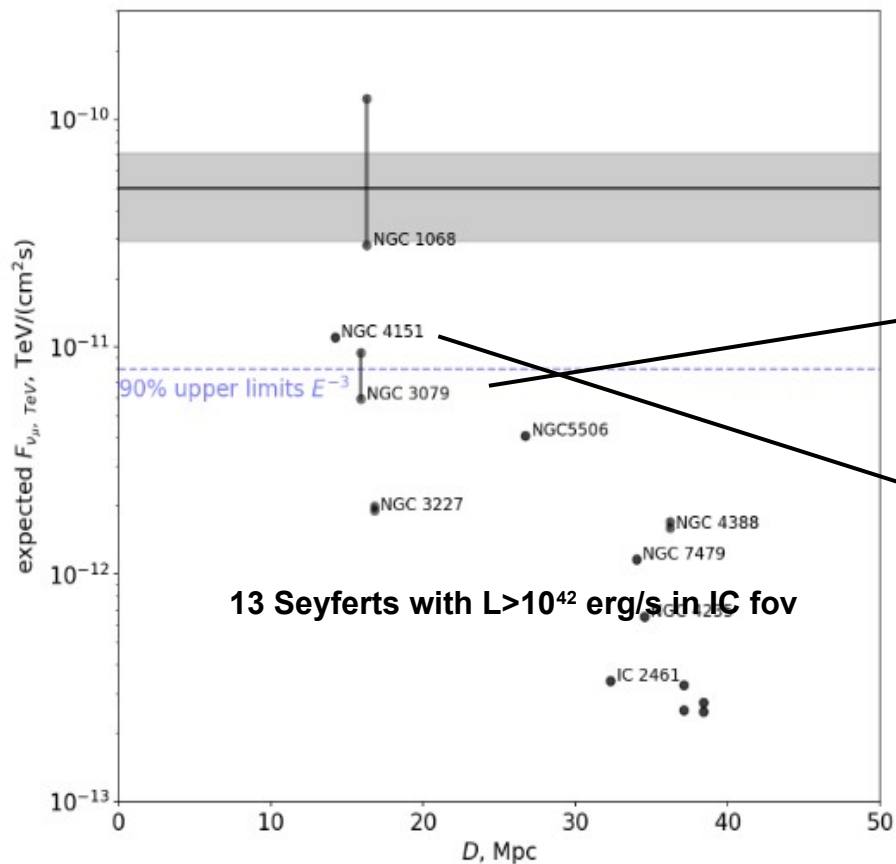




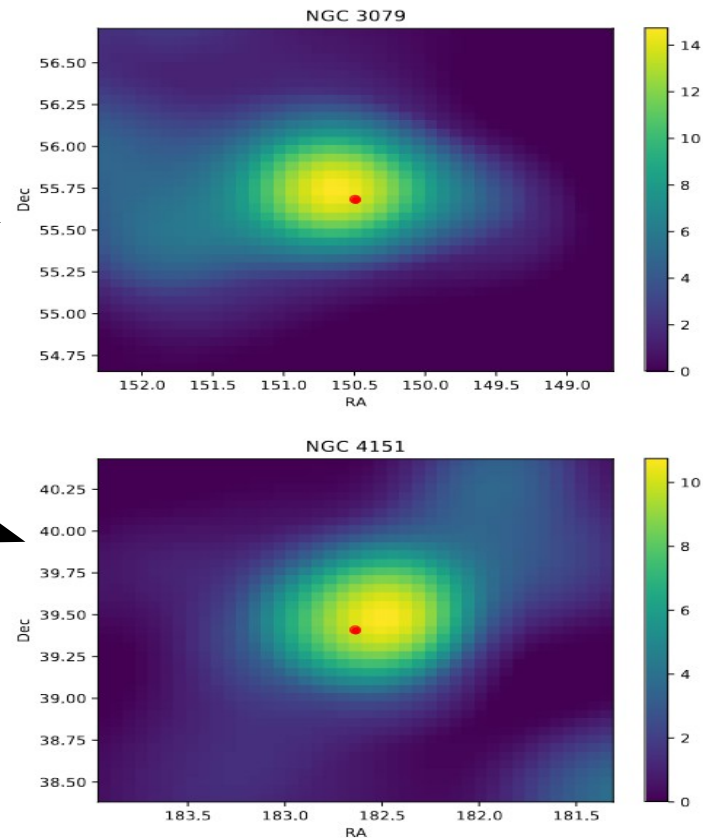
# Are all Seyfert TeV neutrino emitters?

Neronov, Savchenko, Semikoz, 2306.09018

neutrino flux expected from observed Xrays



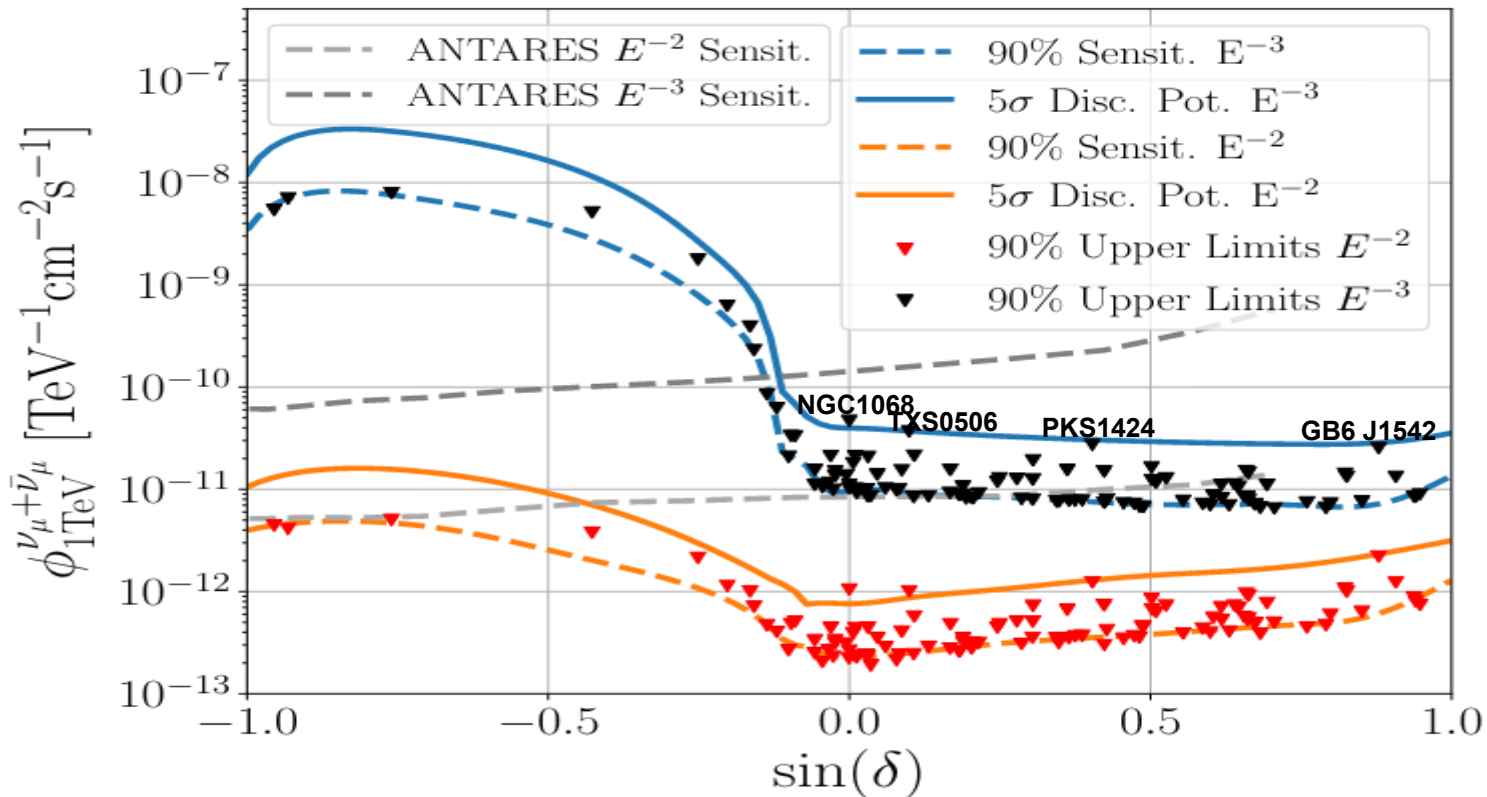
observed in IceCube muon tracks



probably yes

# FLUX BOUNDS FOR LIST OF 110 PRESELECTED GAMMA SOURCES

(from 4FGL: BLLac, FSRQ, starburst gal + few galactic sources from TeVCAT)



1910.08488

(1 Seyfert+3 blazars)

TXS 0506 blazar at 1.7Gpc distance, NGC 1068 Seyfert at ~14 Mpc  
 → populations with very different luminosities and source densities  
 some hints of correlations with TDE? 2005.05340  
 there are also stringent bounds on contribution from GRBs (<1%)

2205.11410,  
 2302.05459

# THE GLASHOW RESONANCE

Interaction of  $\bar{\nu}_e$  with electrons in ice

$$\bar{\nu}_e e \rightarrow W \rightarrow \bar{f} f'$$

resonant for:

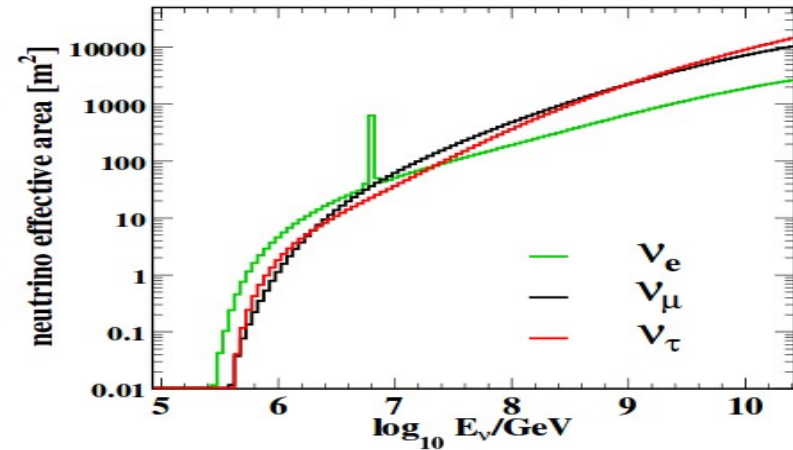
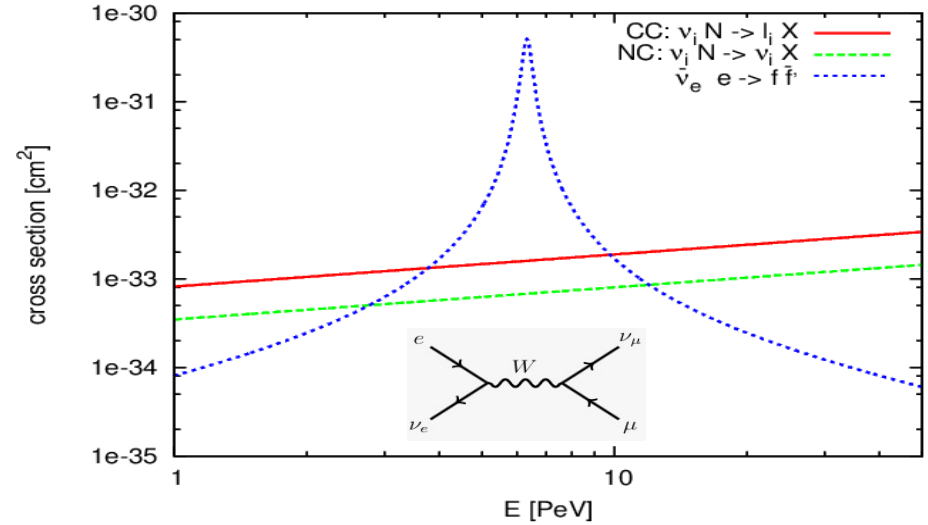
$$E = \frac{M_W^2}{2m_e} = 6.3 \text{ PeV}$$

at the peak,

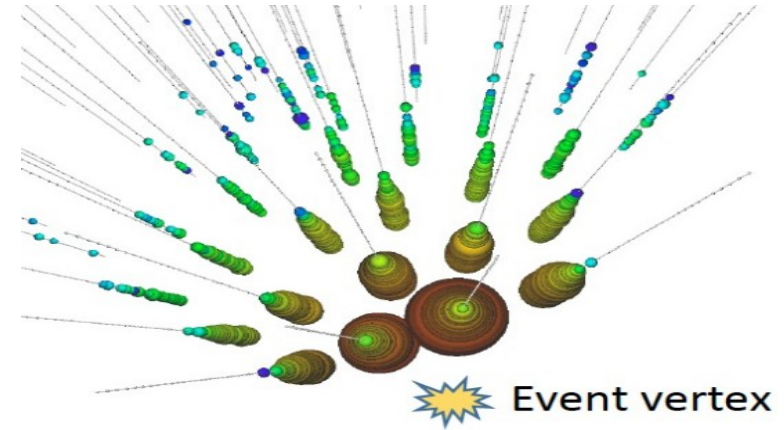
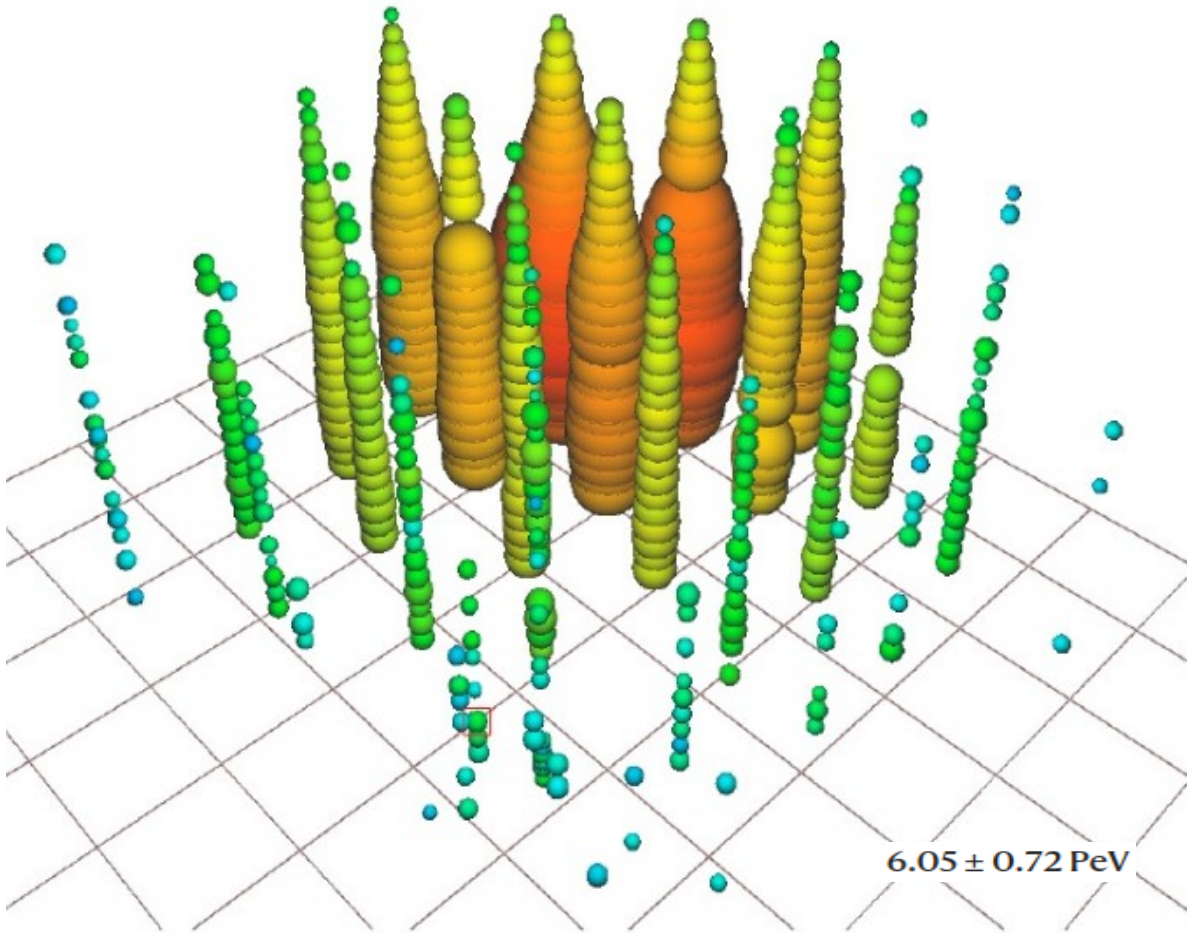
$$\sigma(\bar{\nu}_e e \rightarrow \text{all}) \simeq 350 \sigma^{CC}(\nu_i N \rightarrow l_i X)$$

but peak narrow (0.17 PeV),  
electron antineutrino flavor not dominant,  $n_e/n_N = 5/9$

→ overall contribution to the IceCube rates  
from W resonance is similar to the  
CC+NC ones within 2.5 PeV of the resonance



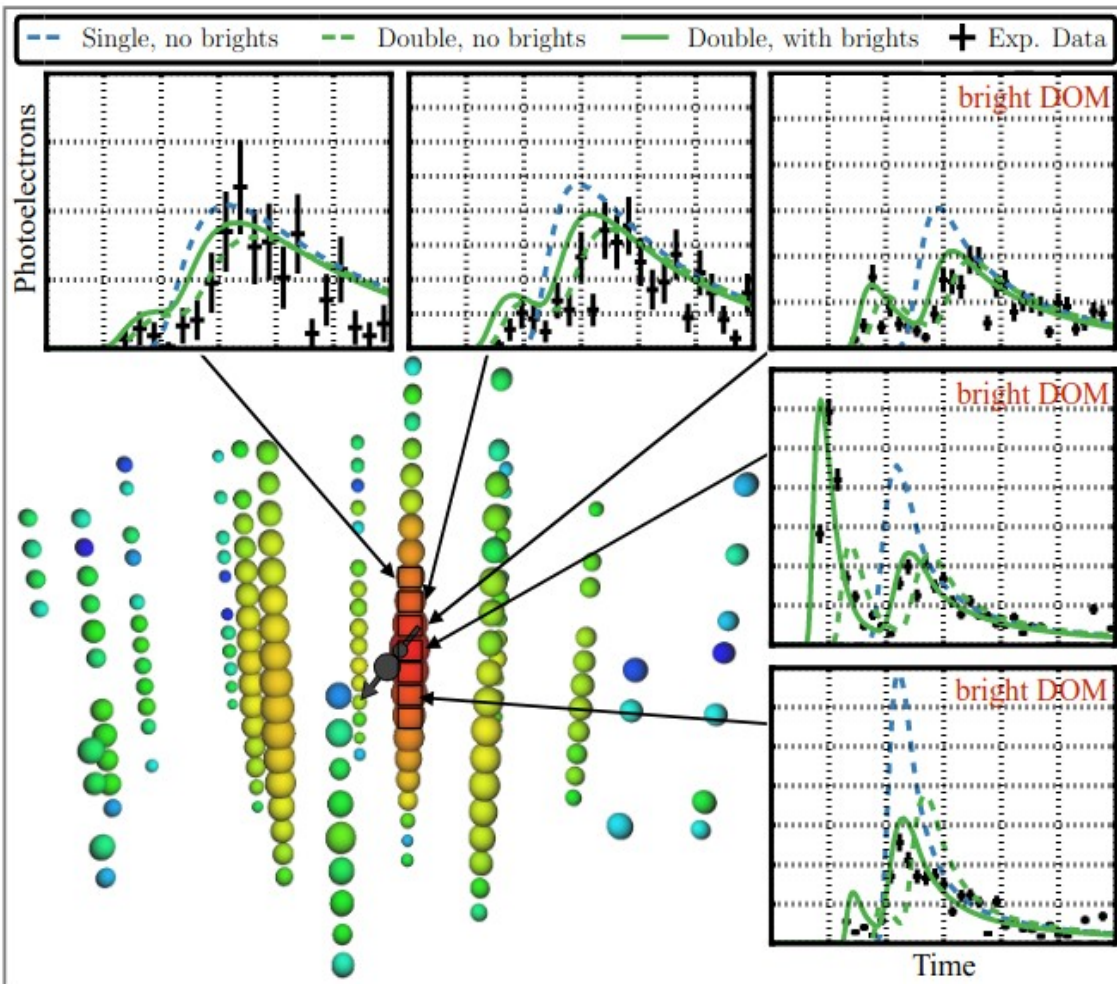
**one W-resonance candidate has been observed!**



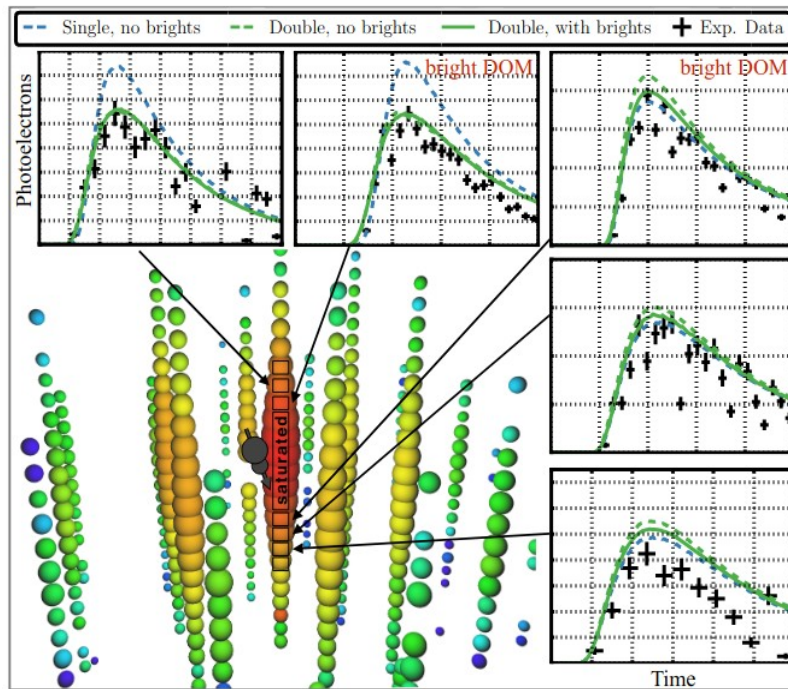
Nature 591 (2021) 7849

# two tau neutrino candidates have been observed !

2011.03561

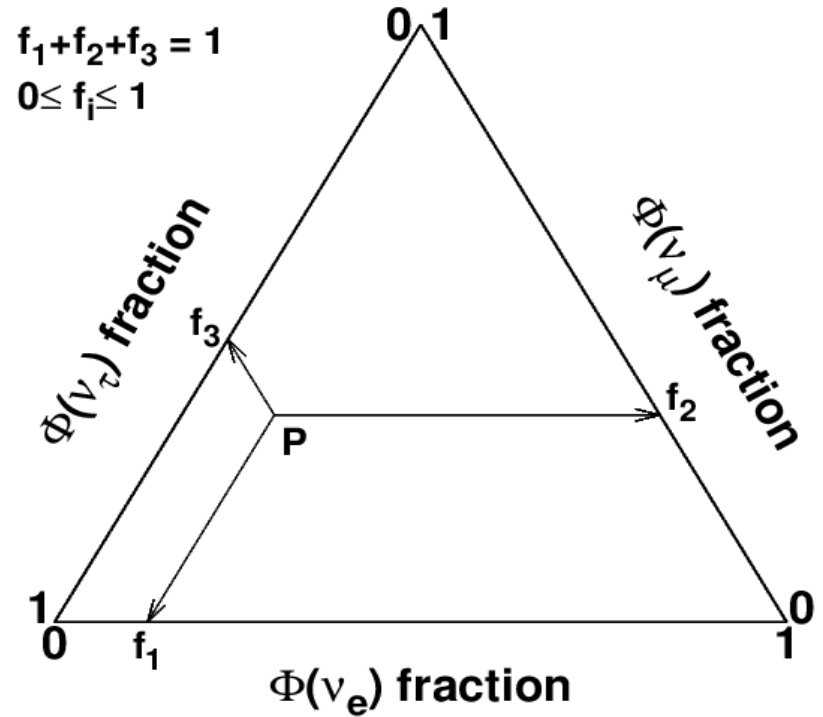
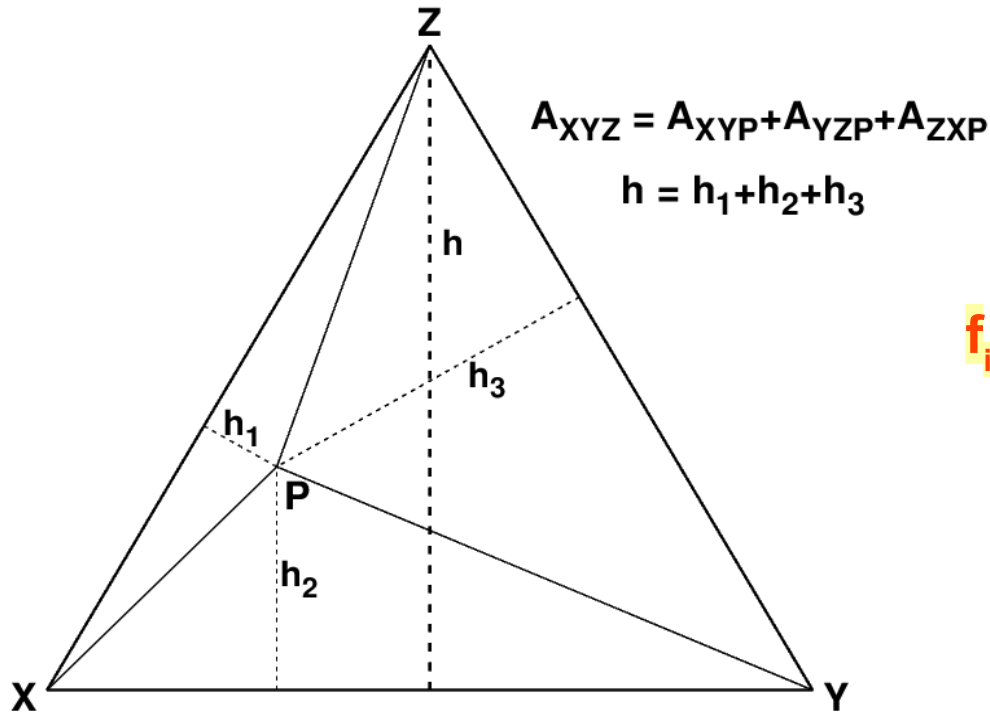


Event #1   Event #2  
 $> 1.5 \text{ PeV}$     $> 65 \text{ TeV}$



double pulse

# FLAVOR TRIANGLE & VIVIANI'S THEOREM



# Flavor oscillations of astrophysical neutrinos

## Incoherent flavor conversions

(Pakvasa et al 2008)

$$P_{\alpha\beta} = \sum_i |U_{\alpha i}|^2 |U_{\beta i}|^2$$

(after traveling from cosmological distances  
oscillations get averaged)

$$\langle \sin^2(\Delta m^2 x / 4E) \rangle \simeq 1/2$$

what arrives is an incoherent superposition of mass eigenstates

**$\pi^-$ -decays:**  $(\nu_e : \nu_\mu : \nu_\tau) = (1 : 1 : 0) \rightarrow 2 \times (0.40 : 0.31 : 0.29)$

$$(\bar{\nu}_e : \bar{\nu}_\mu : \bar{\nu}_\tau) = (0 : 1 : 0) \rightarrow (0.24 : 0.39 : 0.37)$$

[similarly for  $\pi^+$  decays]

**$\mu$  dumps:**  $(\nu_e : \nu_\mu : \nu_\tau) = (0 : 1 : 0) \rightarrow (0.24 : 0.39 : 0.37)$

[if  $\mu$  get damped by synchrotron emission before decay]

(Kashti&Waxman 05)

**n-decays:**  $(\nu_e : \nu_\mu : \nu_\tau) = (0 : 0 : 0) \rightarrow (0 : 0 : 0)$

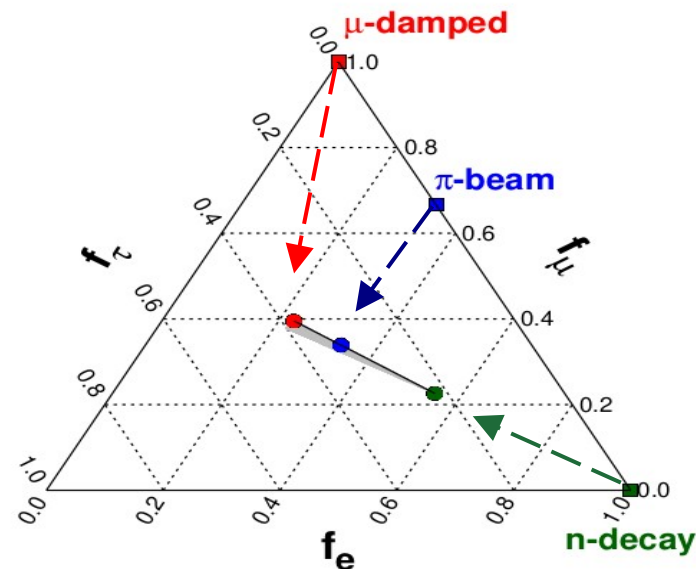
$$(\bar{\nu}_e : \bar{\nu}_\mu : \bar{\nu}_\tau) = (1 : 0 : 0) \rightarrow (0.55 : 0.24 : 0.21)$$

(strong nuclear photodisintegrations)

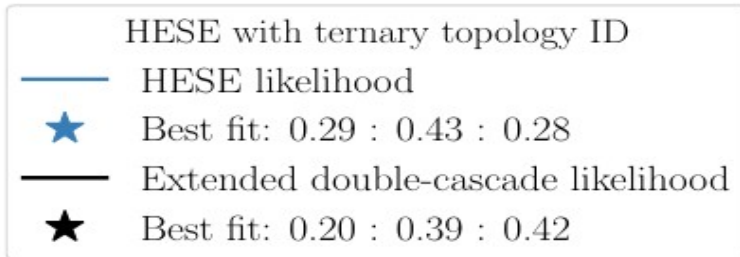
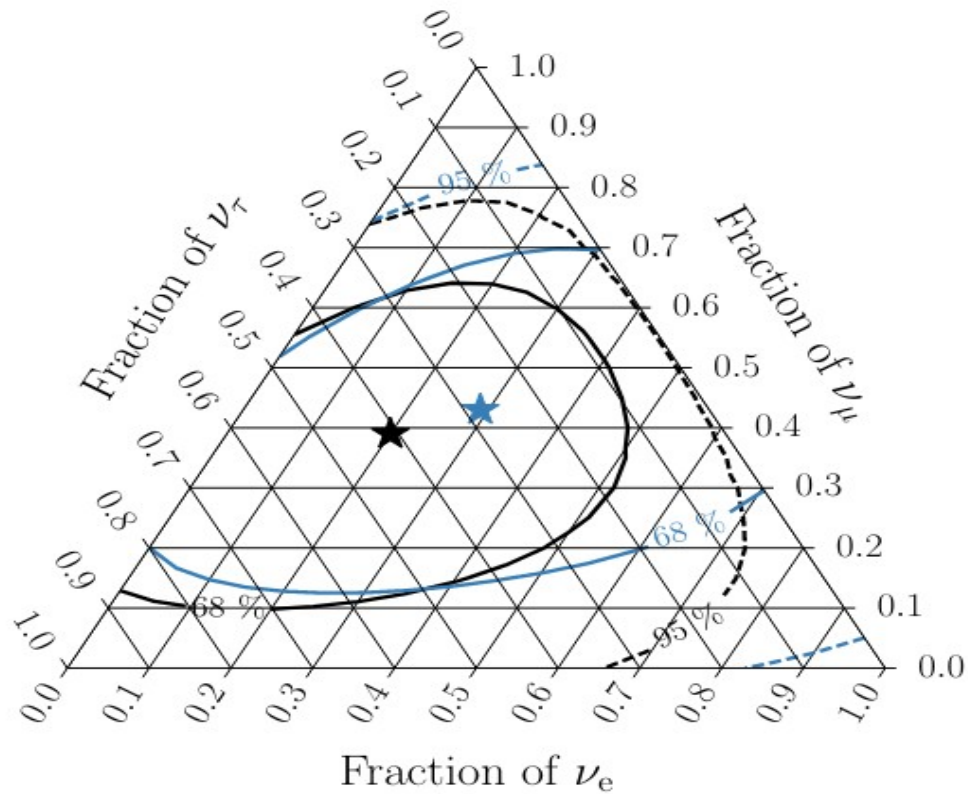
(for NH)

(IceCube measures  
neutrinos+antineutrinos)

$$\nu + \bar{\nu} \rightarrow (0.335 : 0.339 : 0.326)$$



# Results from IceCube



cascade / track ratio

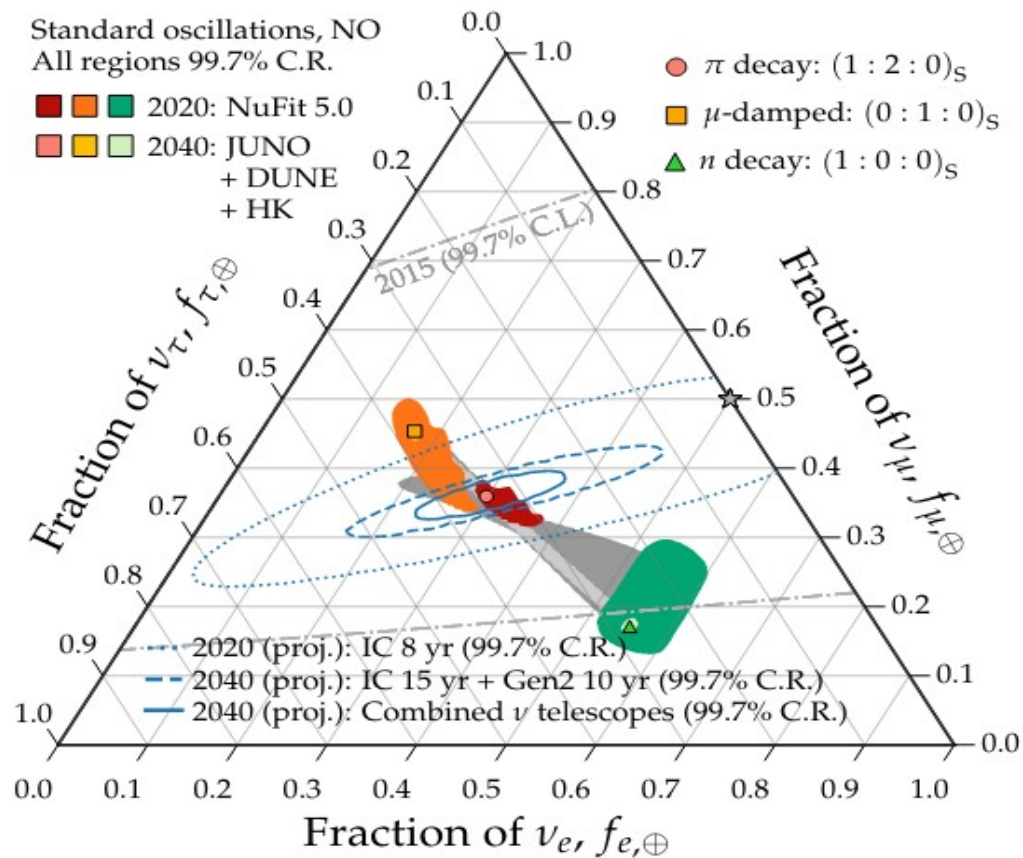
tau neutrinos, ...

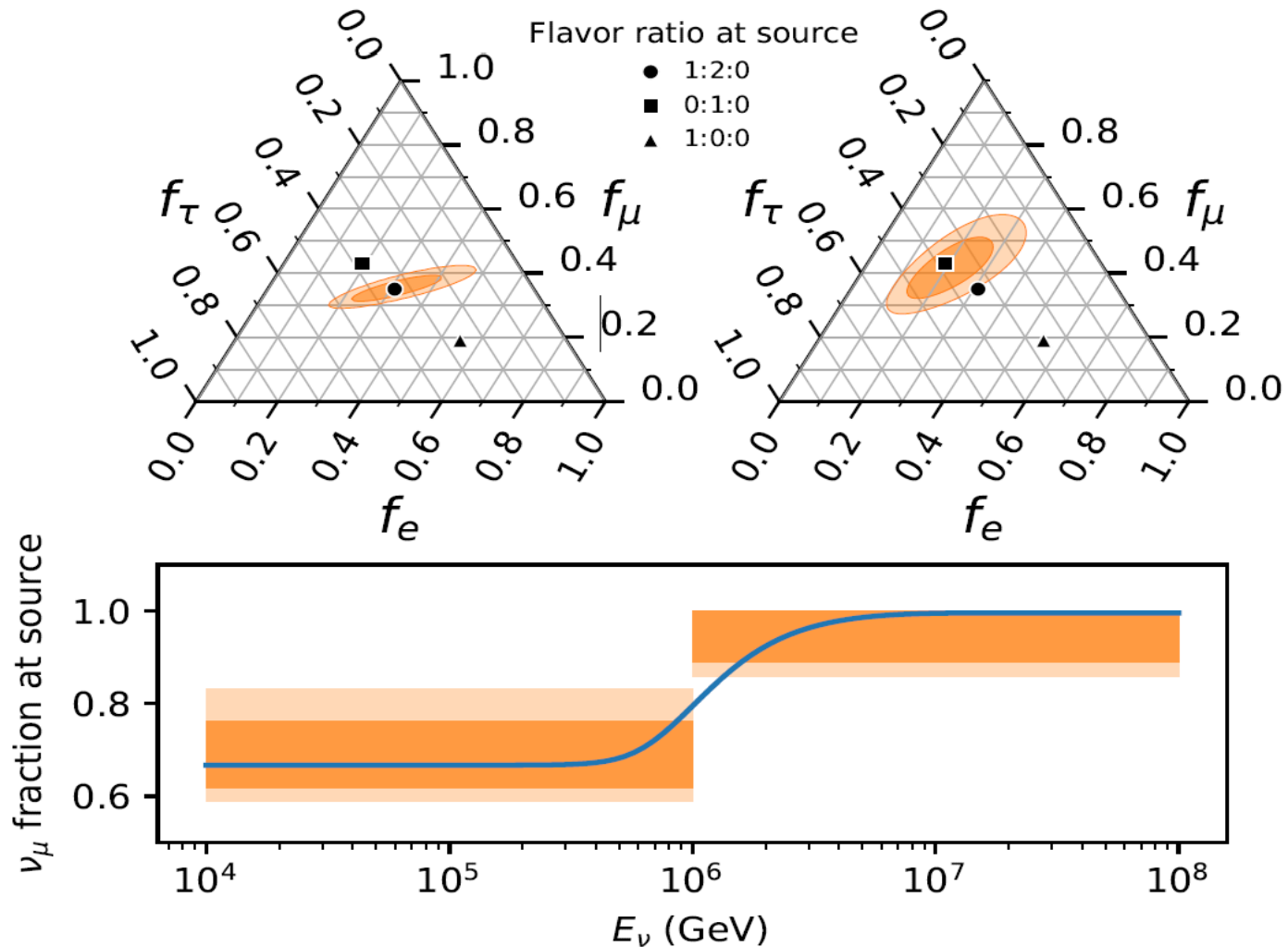


2011.03561



# IN 20 YEARS FROM NOW...



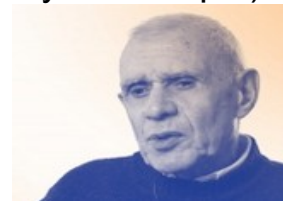


transition to muon damped sources observable with Gen2

# Ultra-high energy neutrinos from interactions with CMB

(Berezinsky & Zatsepin)

UHECRs can be attenuated by background photons as they travel (cosmogenic  $\nu$ )  
besides eventually interacting at the sources (astrophysical  $\nu$ )



**Threshold:**  $p \gamma \rightarrow \pi^+ n$

$$s = (p_p + p_\gamma)^2 > (m_p + m_\pi)^2 \Rightarrow E_p > \frac{m_\pi (2m_p + m_\pi)}{4 E_\gamma} \simeq \frac{70 \text{ EeV}}{E_\gamma / 10^{-3} \text{ eV}}$$

**$10^{20}$  eV for CMB photons,  
 $10^{17}$  eV for optical photons**

**$\nu$  energies:**

$$\begin{array}{l}
 p \gamma \rightarrow \pi^+ n \begin{cases} \rightarrow \pi^+ \rightarrow \mu^+ \nu_\mu \rightarrow e^+ \bar{\nu}_\mu \nu_\mu \nu_e \longrightarrow E_{\bar{\nu}_\mu} \simeq E_{\nu_\mu} \simeq E_{\nu_e} \simeq E_\pi / 4 \simeq E_p / 20 \\
 \rightarrow n \rightarrow p e \bar{\nu}_e \longrightarrow E_{\bar{\nu}_e} \simeq \frac{m_n - m_p - m_e}{2 m_n} E_n \simeq 4 \times 10^{-4} E_n \end{cases}
 \end{array}$$

**Redshift (production at  $0 < z < 4$ ):**

**$T_{\text{CMB}} = (1+z) 2.7 \text{ K} \rightarrow$  redshifted threshold**

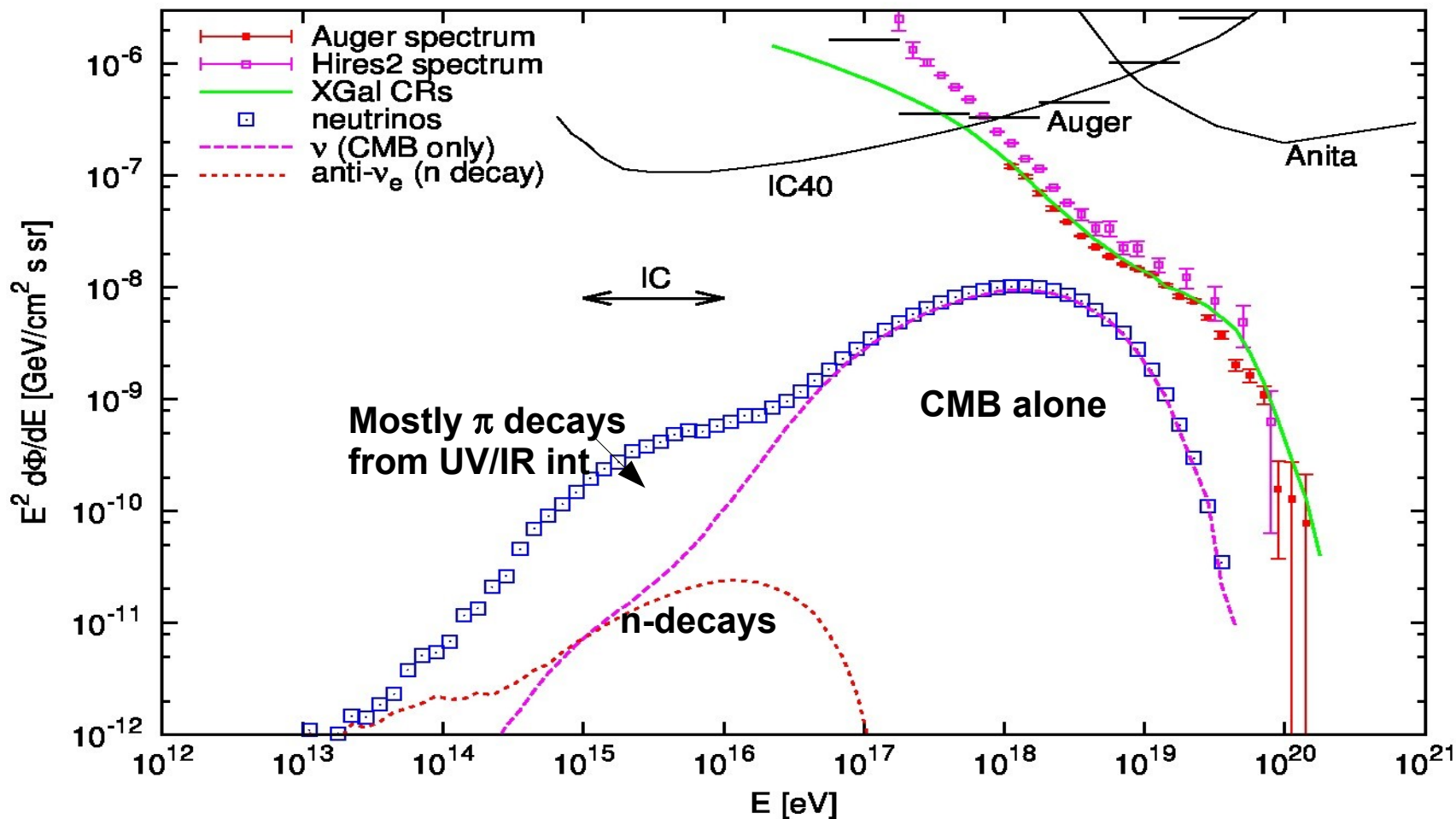
**Redshifted  $\nu$  energy**  $E_{\nu}^{\pi\text{-dec}} \simeq \frac{E_p}{20(1+z)}$

$$E_{\nu}^{\pi\text{-dec}} \simeq \frac{5 \text{ EeV}}{(1+z)(E_\gamma / 10^{-3} \text{ eV})}$$

**EeV  $\nu$  from interactions with CMB photons  
PeV  $\nu$  from interactions with UV/O/IR photons**

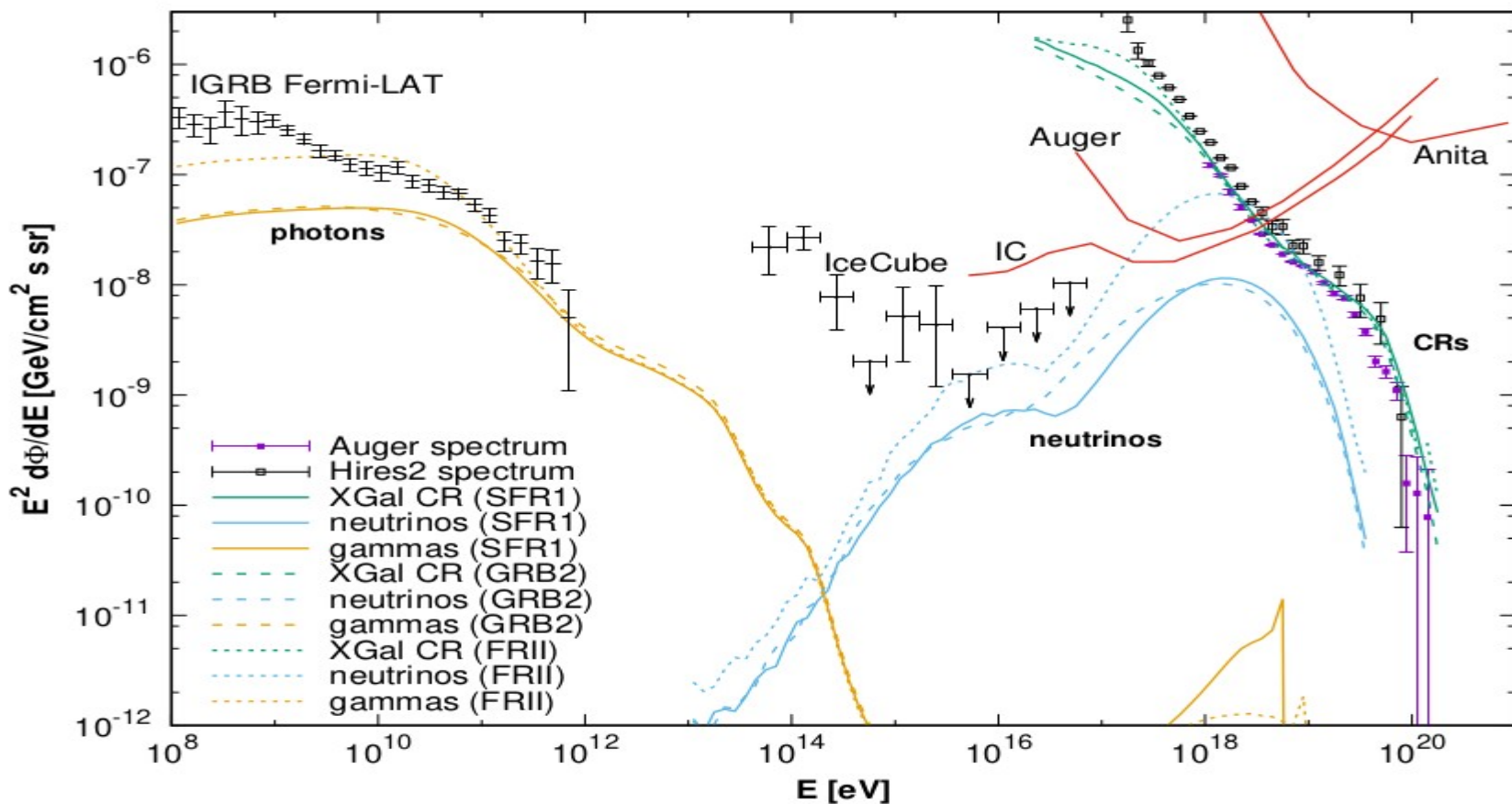
# Cosmogenic neutrinos from UHECR proton sources

proton sources,  $\alpha=2.4$ ,  $E_{\max}=200$  EeV, GRB2



# $\nu$ and $\gamma$ for different CR source evolutions & cascade bound

proton sources,  $E_{\text{max}}=200 \text{ EeV}$

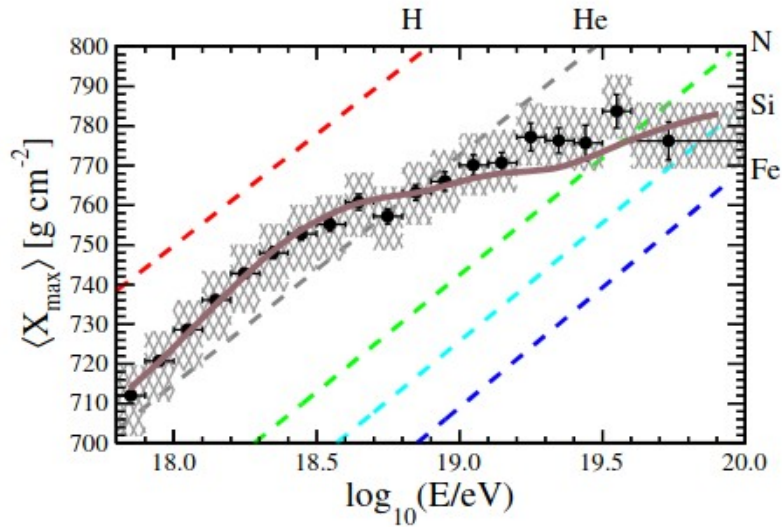


very strong evolution in tension with  $\gamma$  flux

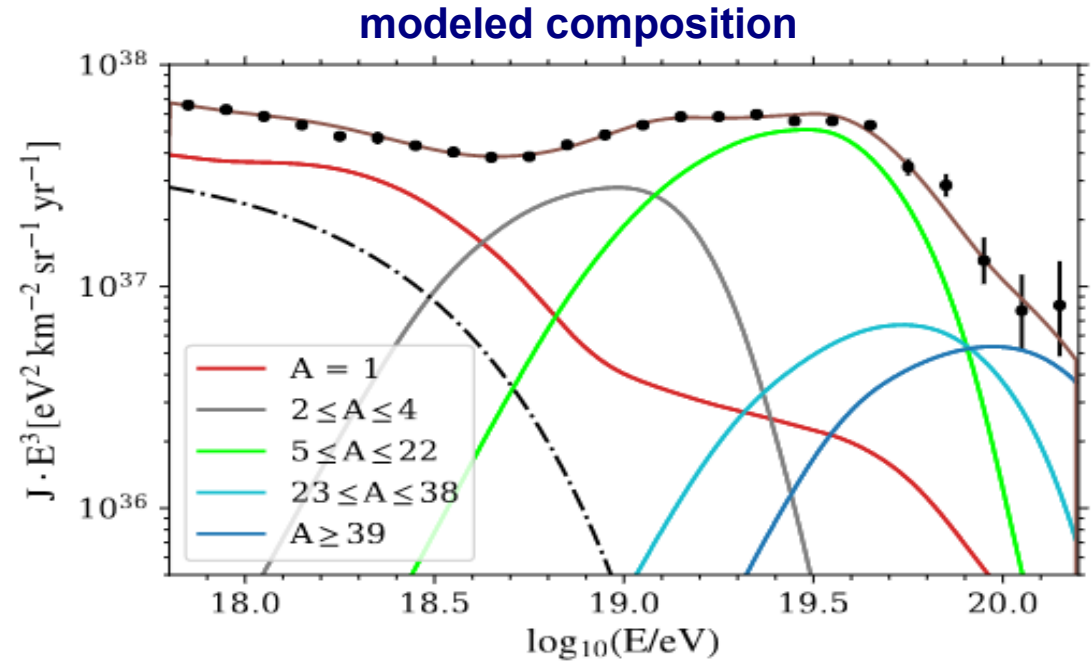
ER, Sigl, vVliet & Mollerach 1209.4033

but the Auger Observatory established that above the ankle (>5 EeV)  
CRs become increasingly heavier

2211.02857



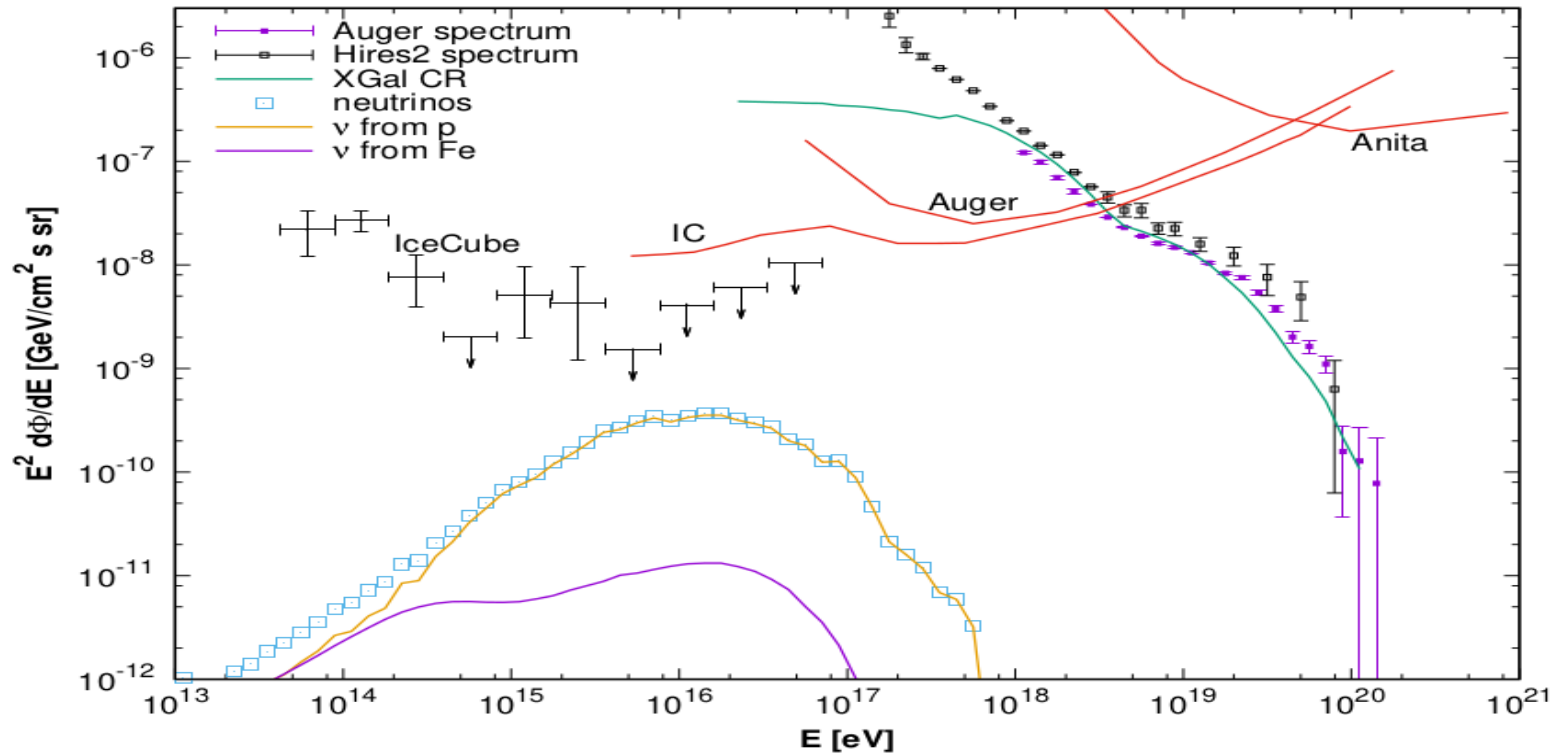
depth of shower maxima vs. E



# scenario with mixed composition with low rigidity cutoff ( $E/Z < 4$ EeV)

mixed composition scenario

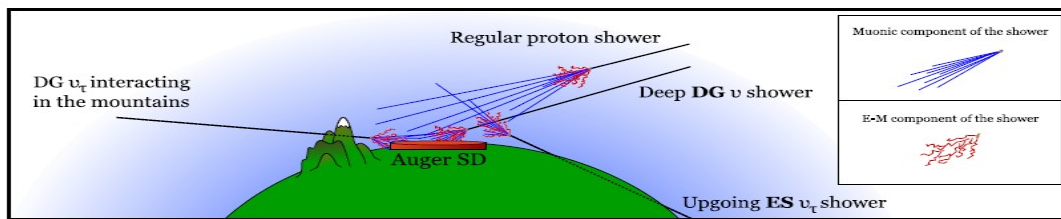
$p/Fe=10$ ,  $\alpha=2.0$ ,  $R_{max}=4$  EV, GRB2



$p$  component below ankle interacting with EBL leads to PeV  $\nu$  fluxes ( $\sim 10\%$  of IC)

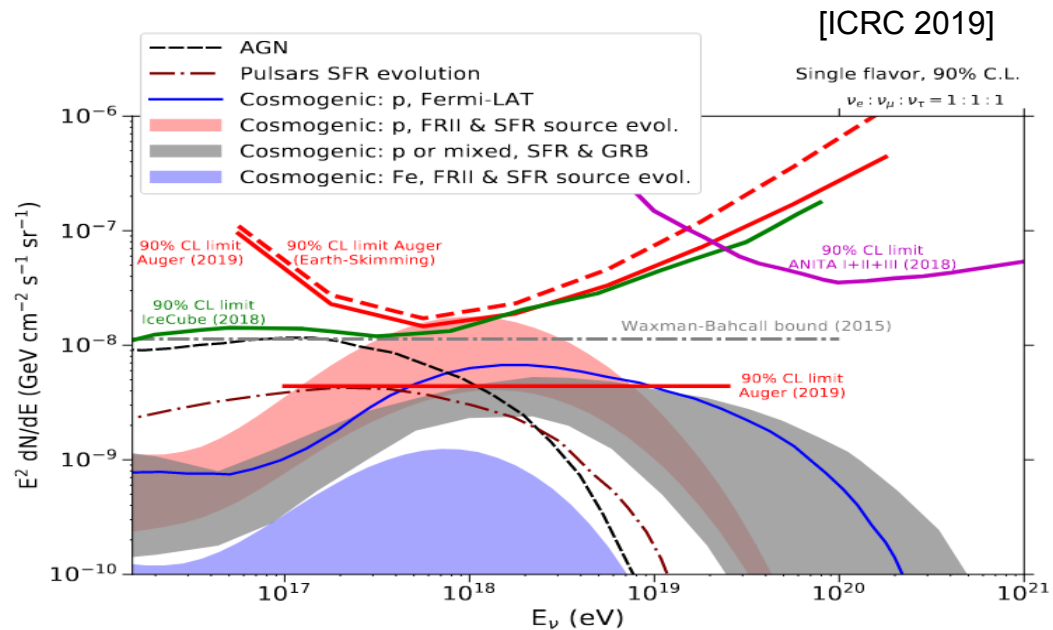
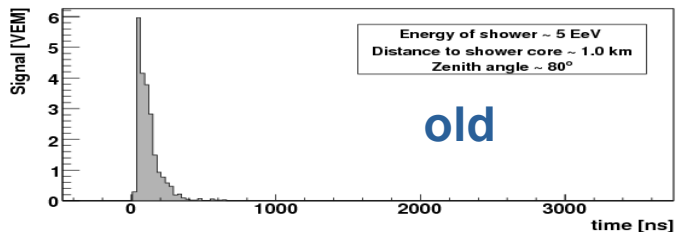
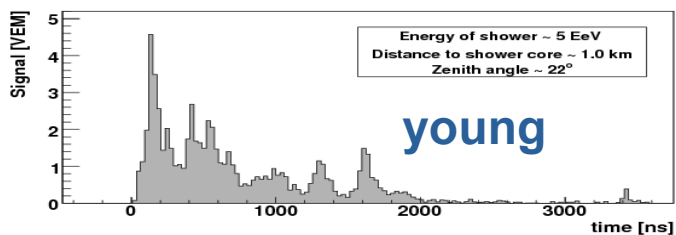
due to low cutoff no pions from CMB  $\rightarrow$  no EeV  $\nu$  (disappointing model)

# NEUTRINO DETECTION IN AUGER



Only neutrinos can produce young horizontal showers

signal vs. time in WCD



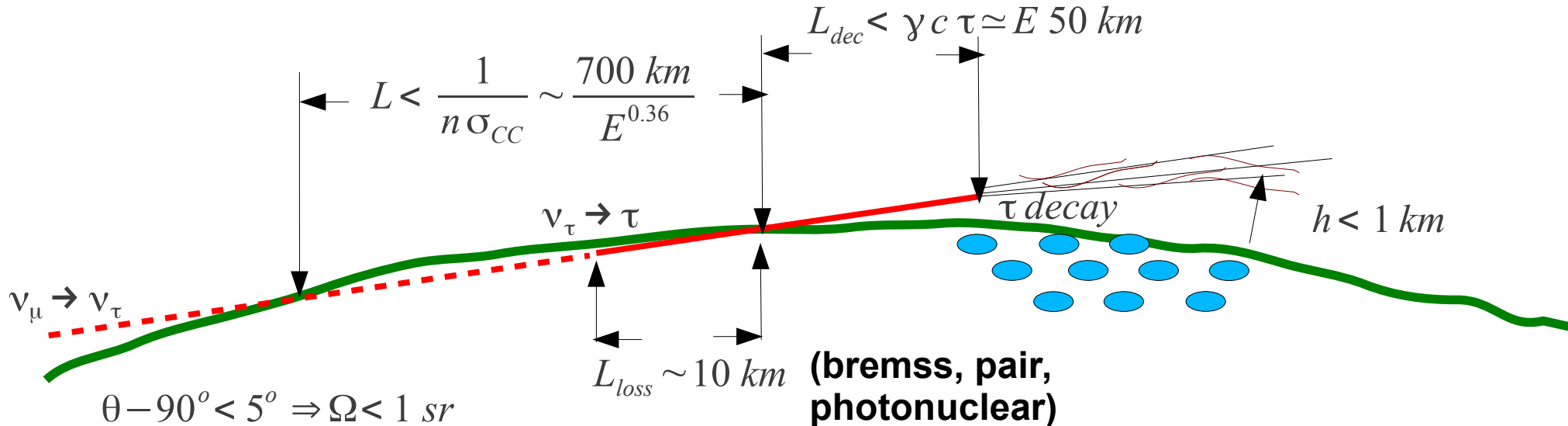
0 events observed → bounds scale linearly with exposure



# Up-going Earth-skimming $\nu_\tau$ showers

Fargion 2000,  
Bertou et al '01  
Feng et al. '02

$$\sigma_{CC} \simeq 10^{-32} \text{ cm}^2 E^{0.36} \quad (E [EeV])$$

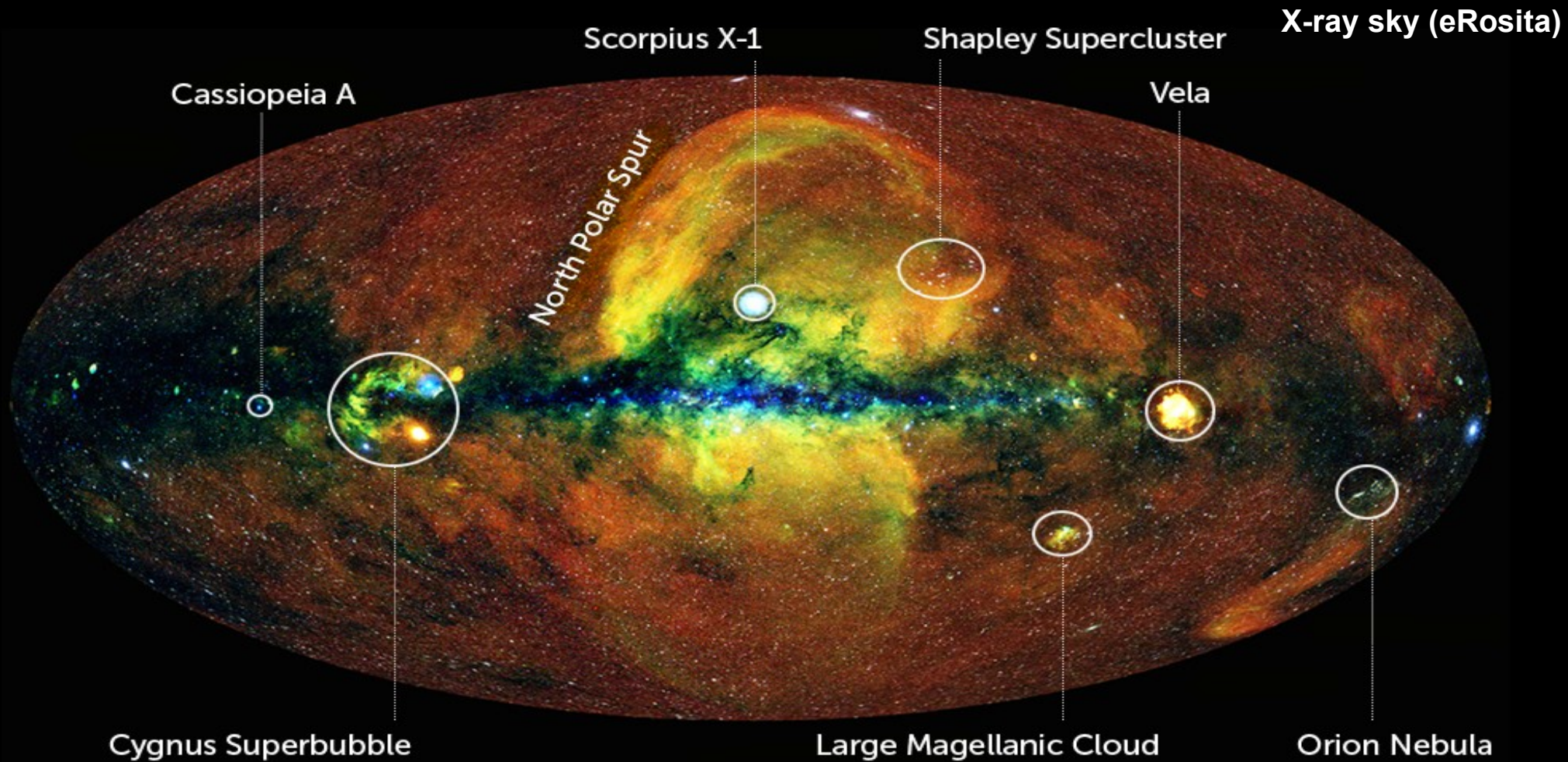


probability of interacting  
in the last 10 km ~ 0.01

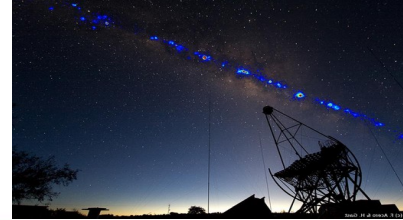
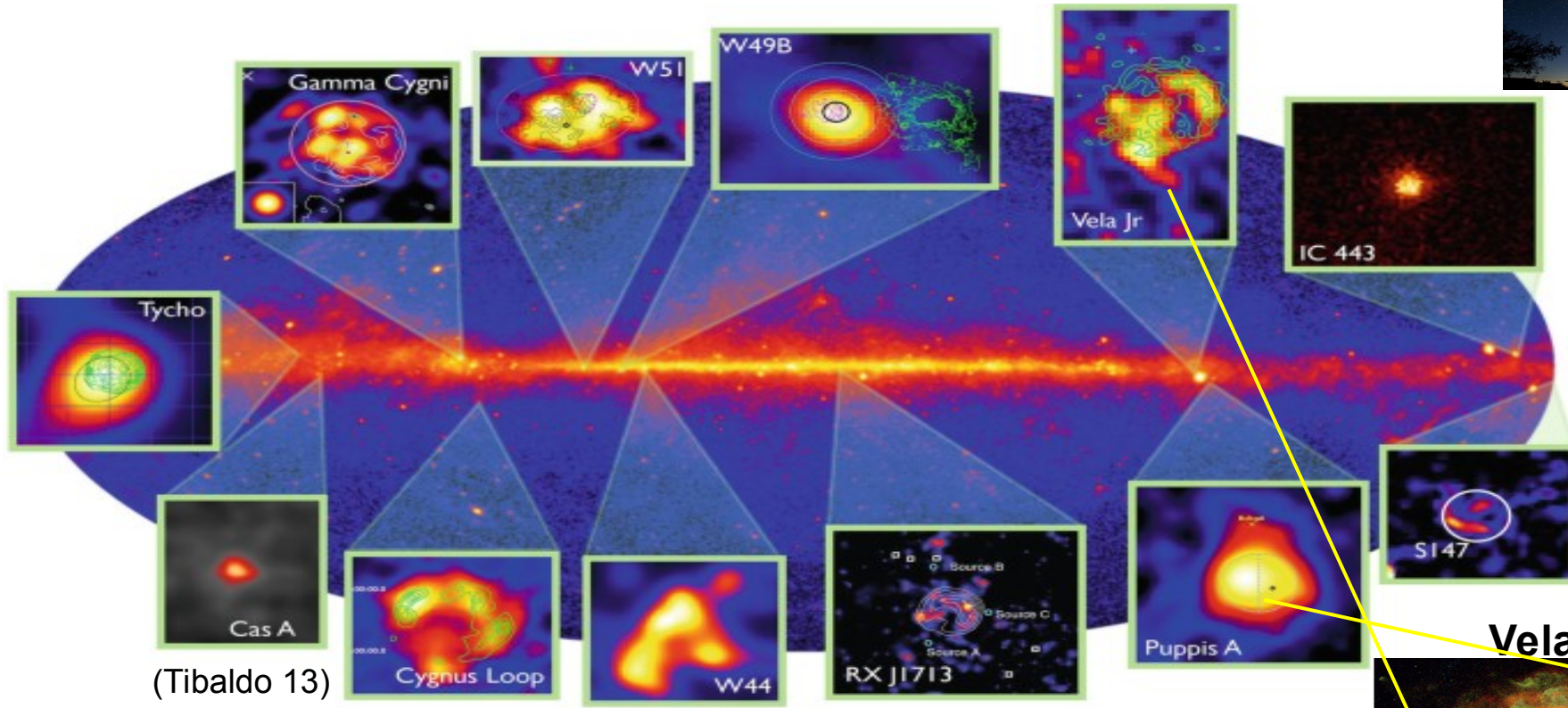
→ effective exposure ~ 0.1 km<sup>2</sup> sr  
(while ~ 10<sup>4</sup> km<sup>2</sup> sr for UHECR)

(bounds also include some downgoing  $\theta > 70^\circ$ )

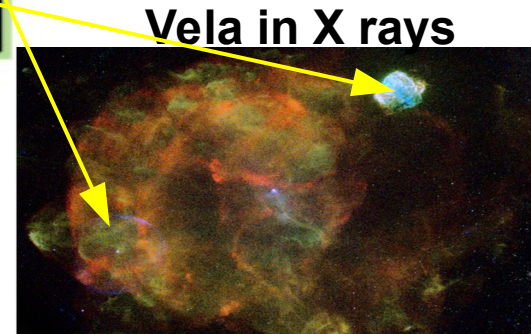
# GALACTIC SOURCES



# Galactic TeV $\gamma$ sources

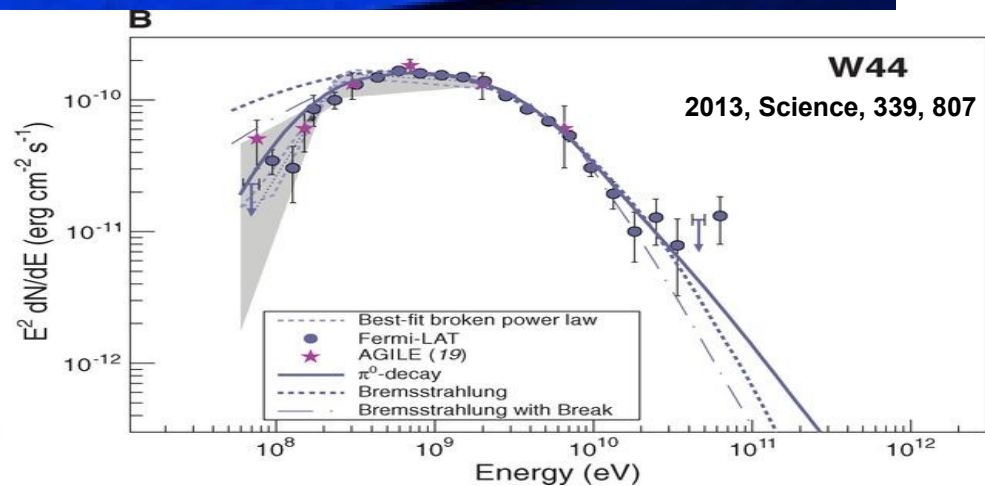
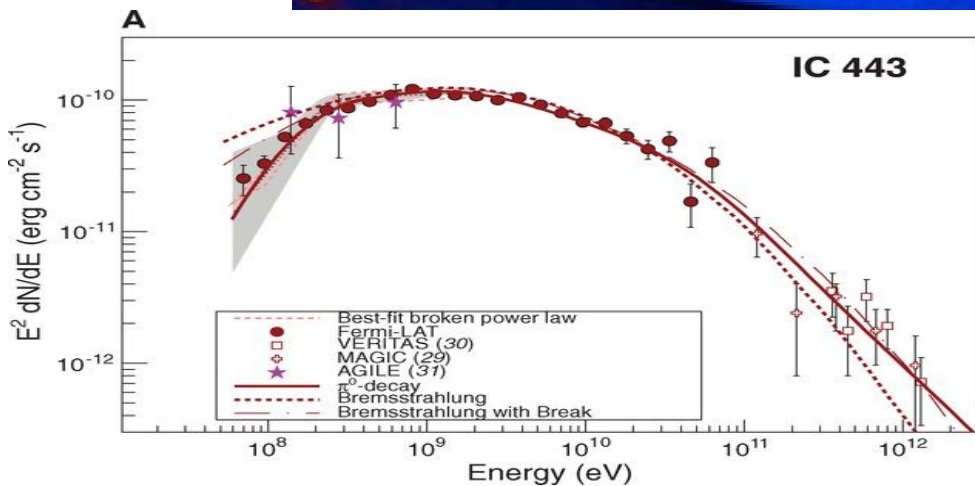
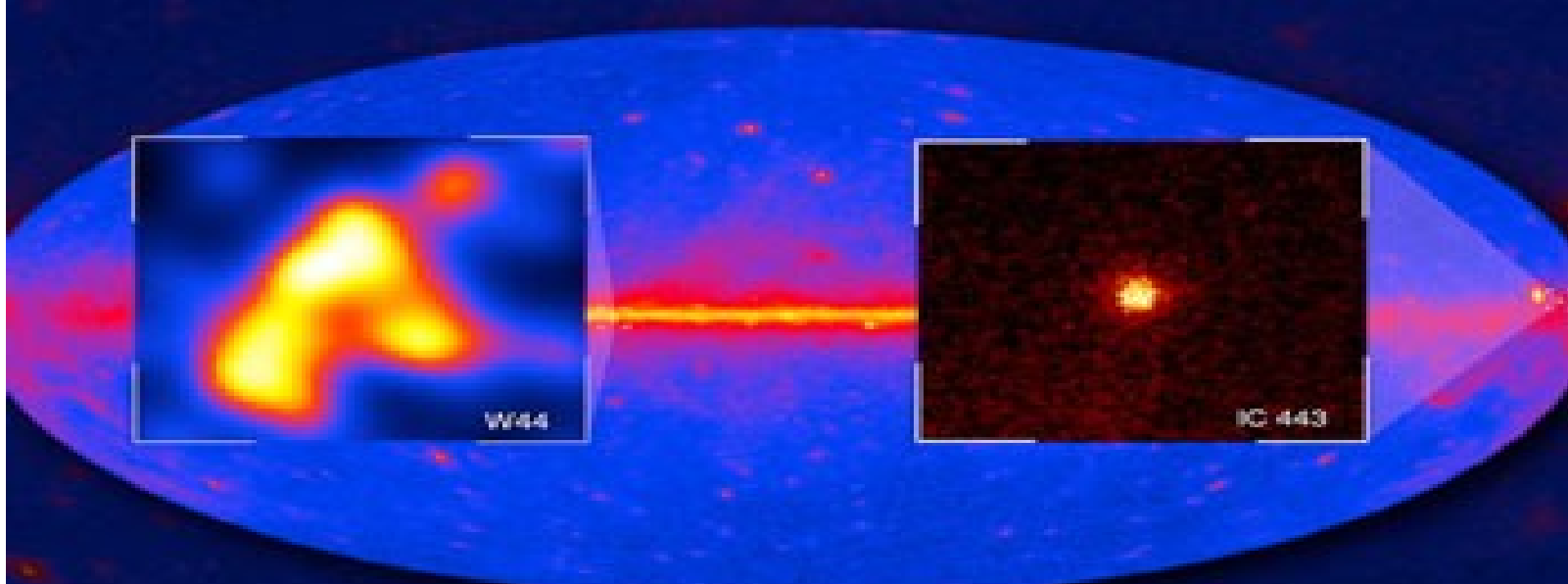


IACT



several show evidence of pion bump  $\rightarrow$  hadronic acceleration (Agile, Fermi)

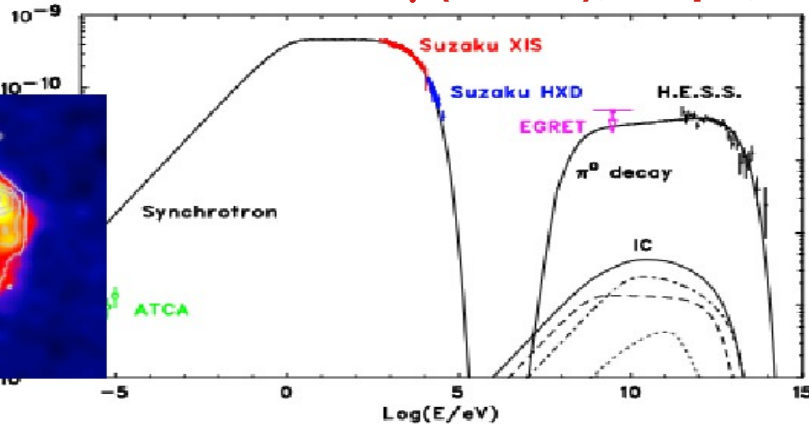
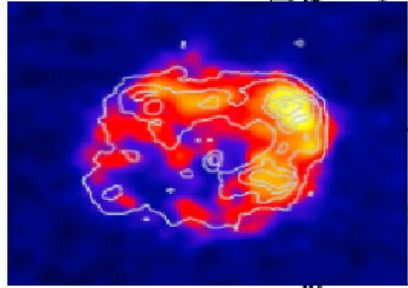
A supernova remnant called Vela (center, reddish green) is one of the most prominent X-ray sources in the sky. The supernova exploded about 12,000 years ago, about 800 light-years away, and overlaps with two other known supernova remnants: Vela Junior (faint purple ring at bottom left) and Puppis A (blue cloud at top right). All three explosions left behind neutron stars, but only the stars at the centers of Vela and Vela Junior are visible to eROSITA.



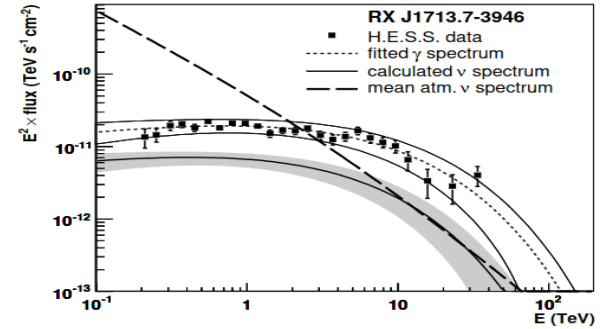
**FERMI found SN remnants with clear signals of gammas from pion decays (low E supp.)**  
**Proton acceleration in Supernovae to beyond 10 TeV proved, associated  $\nu$  flux expected**

# SNR in TeV $\gamma$ (HESS), 1 kpc, 1600yr ?

RX J1713



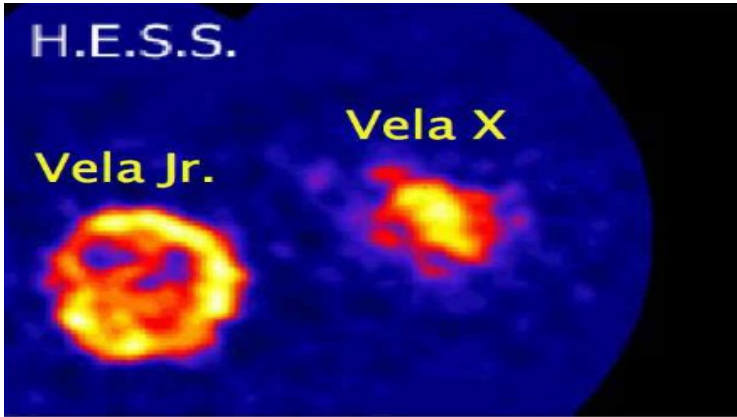
Kappes et al. 0607286



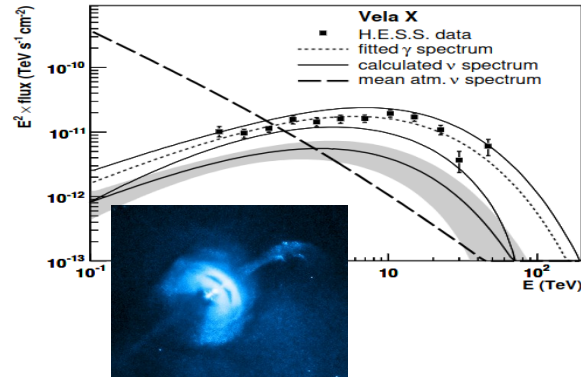
H.E.S.S.

Vela X

Vela Jr.

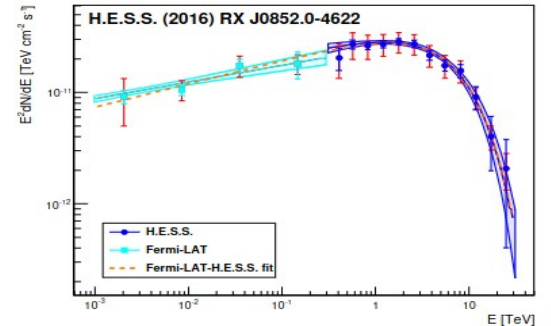


Vela X PWN



Vissani & Aharonian 1112.3911

Vela Jr SNR



expect few events per  $\text{km}^2 \text{ yr}$  above TeV if hadronic

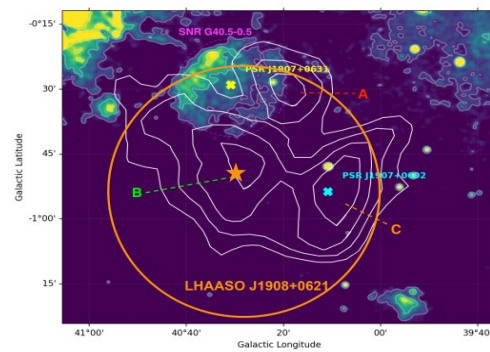
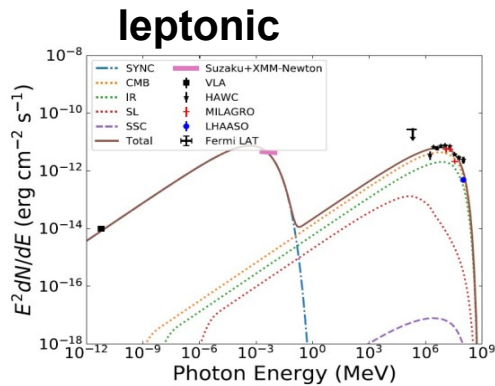
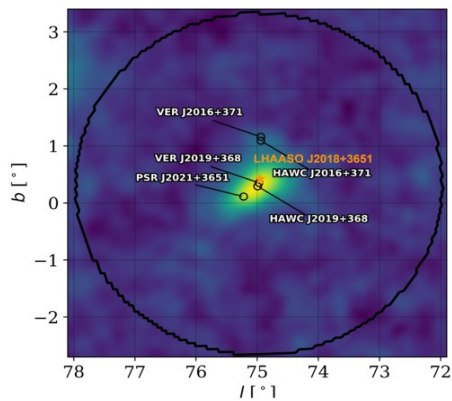
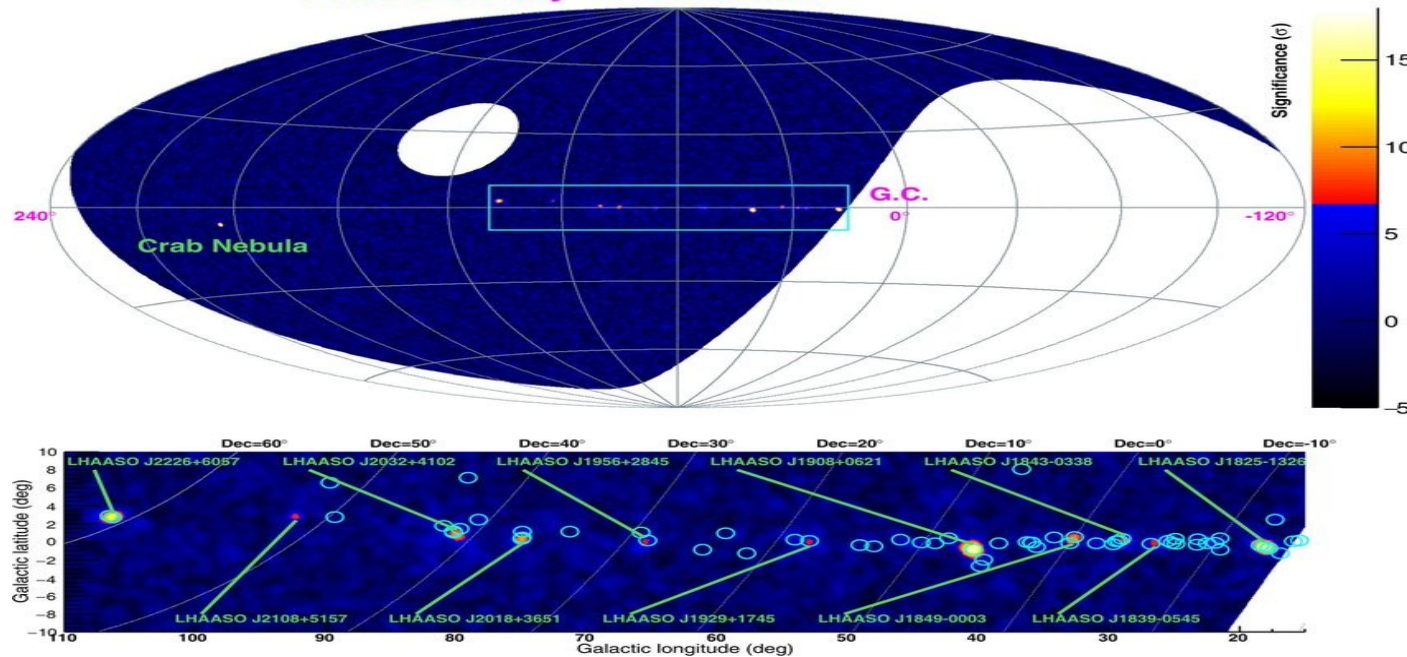
these sources are in southern hemisphere  $\rightarrow$  good targets for KM3NeT

IceCube can see e.g. Cygnus region, Crab, CasA

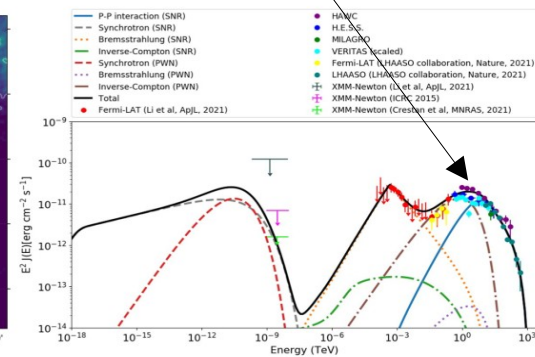
# Ultrahigh-energy photons up to 1.4 petaelectronvolts from 12 $\gamma$ -ray Galactic sources

LHAASO Sky @ >100 TeV

Nature 2021



### hadronic

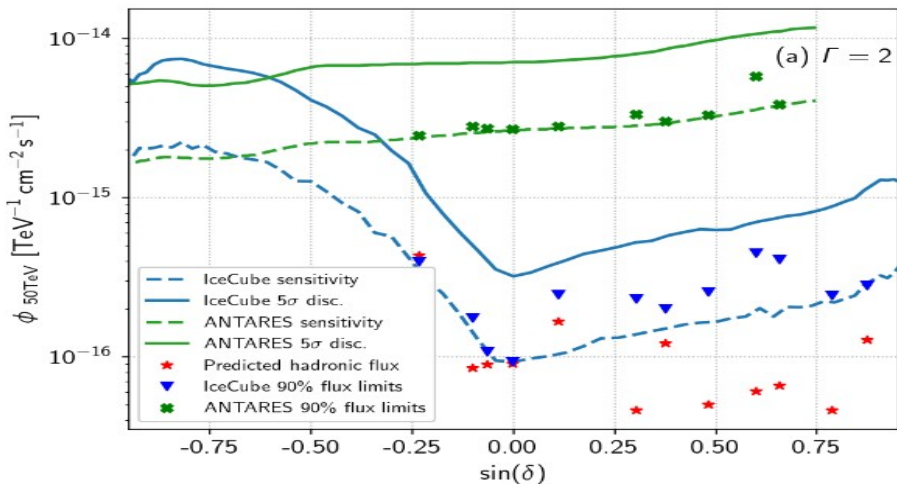


# Searches for Neutrinos from LHAASO ultra-high-energy $\gamma$ -ray sources using the IceCube Observatory

2211.14184

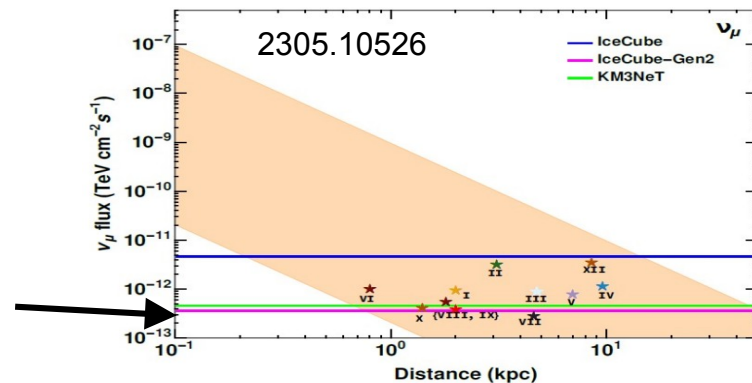
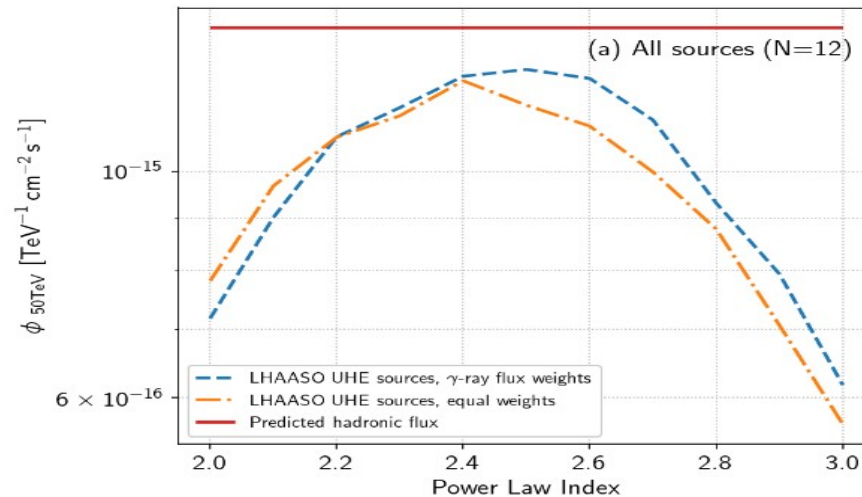
no excess observed  $\rightarrow$  90%CL bound on source fluxes

## individual sources



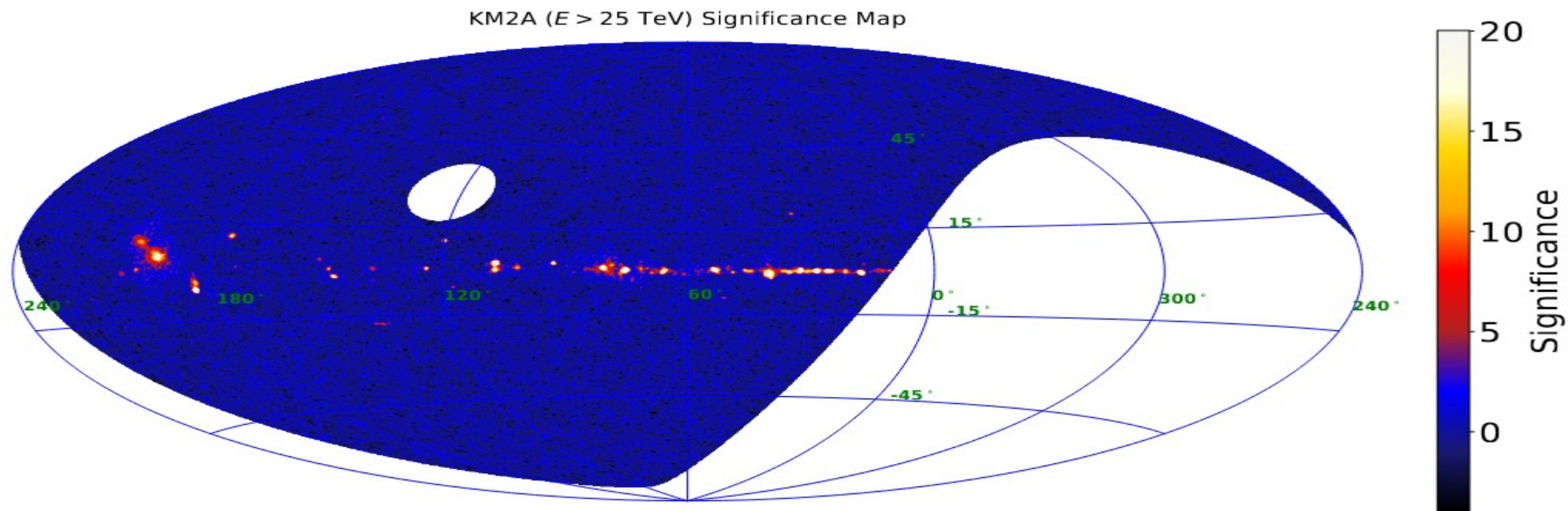
$$\frac{dN_{\nu_{\mu} + \nu_{\bar{\mu}}}}{dE_{\nu}} = \phi_{90\%} \cdot \left(\frac{E_{\nu}}{50\text{TeV}}\right)^{-2} \times 10^{-16} \text{ TeV}^{-1} \text{ cm}^{-2} \text{ s}^{-1}$$

## stacked search



for Crab, hadronic contribution < 59% of total  
 LHAASO J2226+6057 < 47% ,  
 but for others expectations below bounds  
 $\rightarrow$  need larger detectors & better angular resolution

## The First LHAASO Catalog of Gamma-Ray Sources

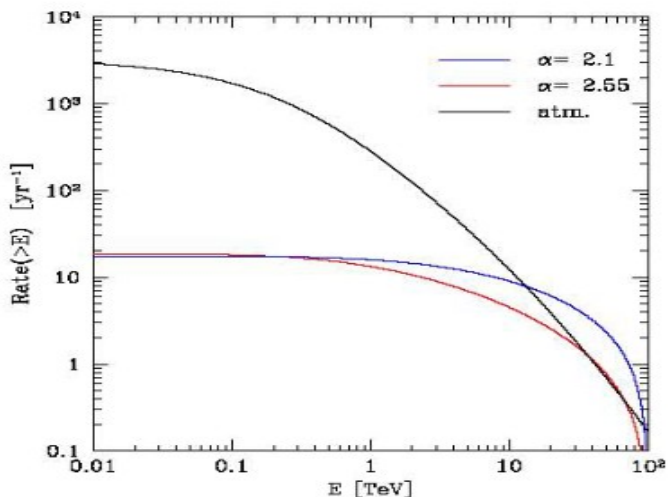


most sensitive  $E > 1$  TeV gamma-ray survey of the sky covering declination from  $-20^\circ$  to  $80^\circ$ . In total, the catalog contains **90 sources** with extended size smaller than  $2^\circ$  and with significance of detection at  $> 5\sigma$ . For each source, we provide its position, extension and spectral characteristics. Furthermore, based on our source association criteria, 32 new TeV sources are proposed in this study. Additionally, 43 sources are detected with ultra-high energy ( $E > 100$  TeV) emission at  $> 4\sigma$  significance level.



# diffuse neutrino flux from Galaxy

observed  $\gamma$  rays are from  $\pi^0$  decays produced in CR - gas interactions  
 → expect neutrinos from associated charged pion decays



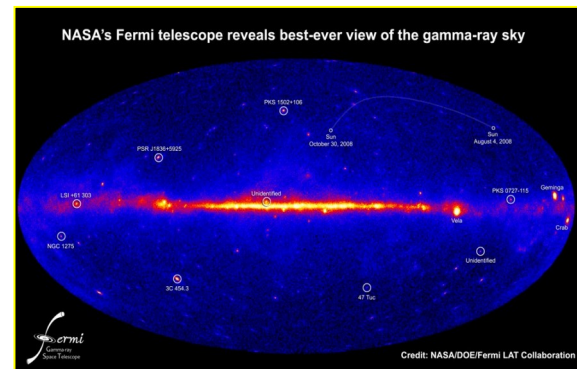
KM3NeT sensitivity

(depends on slope of CR spectrum)

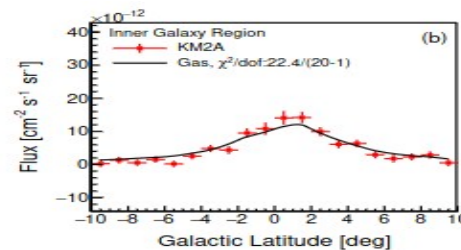
should be observable in the future with muon tracks

recently LHAASO observed diffuse  $\gamma$  emission from the Galaxy in 10 TeV – 1 PeV range 2305.05372

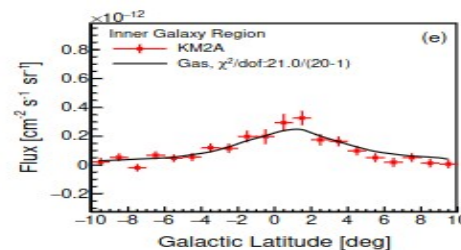
see also Tibet AS array result 2104.05181



a  $2\sigma$  excess observed with Antares



10-63 TeV

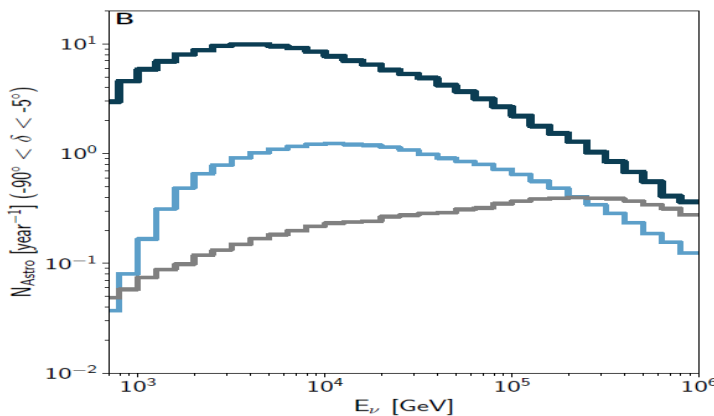
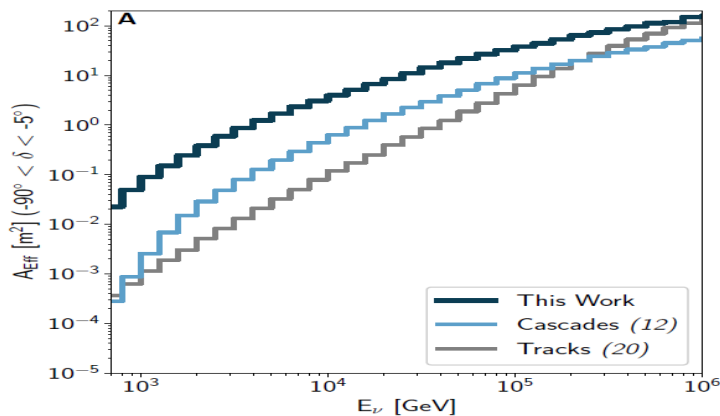


63-1000 TeV

# last week's results from IceCube

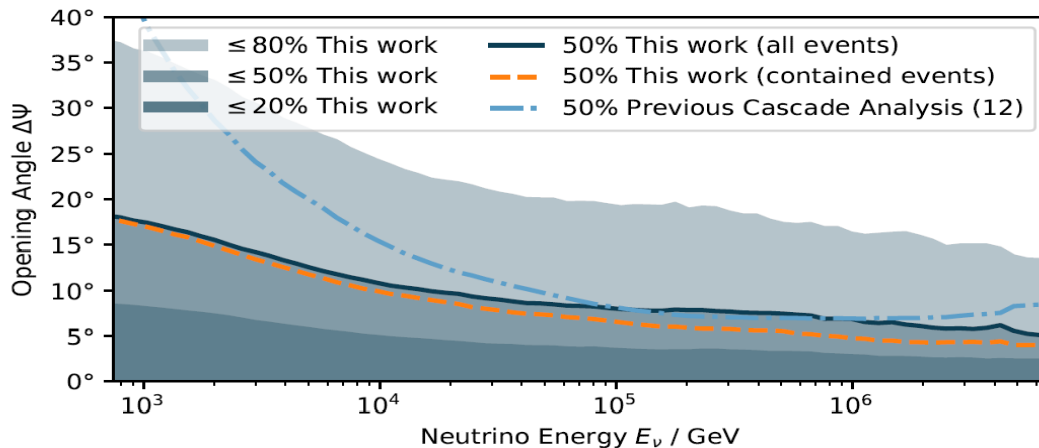
(Science)

new analysis using cascades & improved reconstruction with machine learning



30x more events

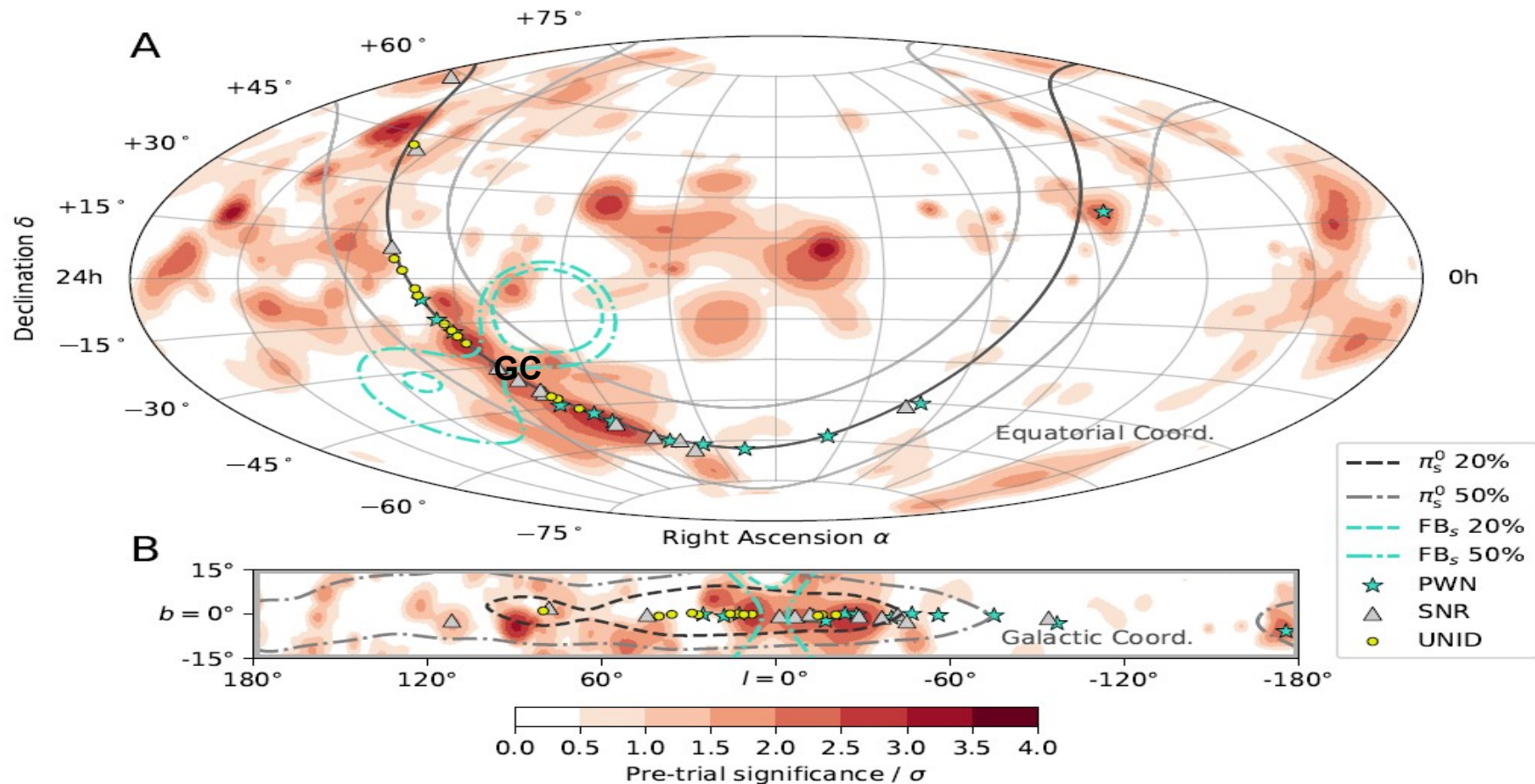
$E > 0.5$  TeV



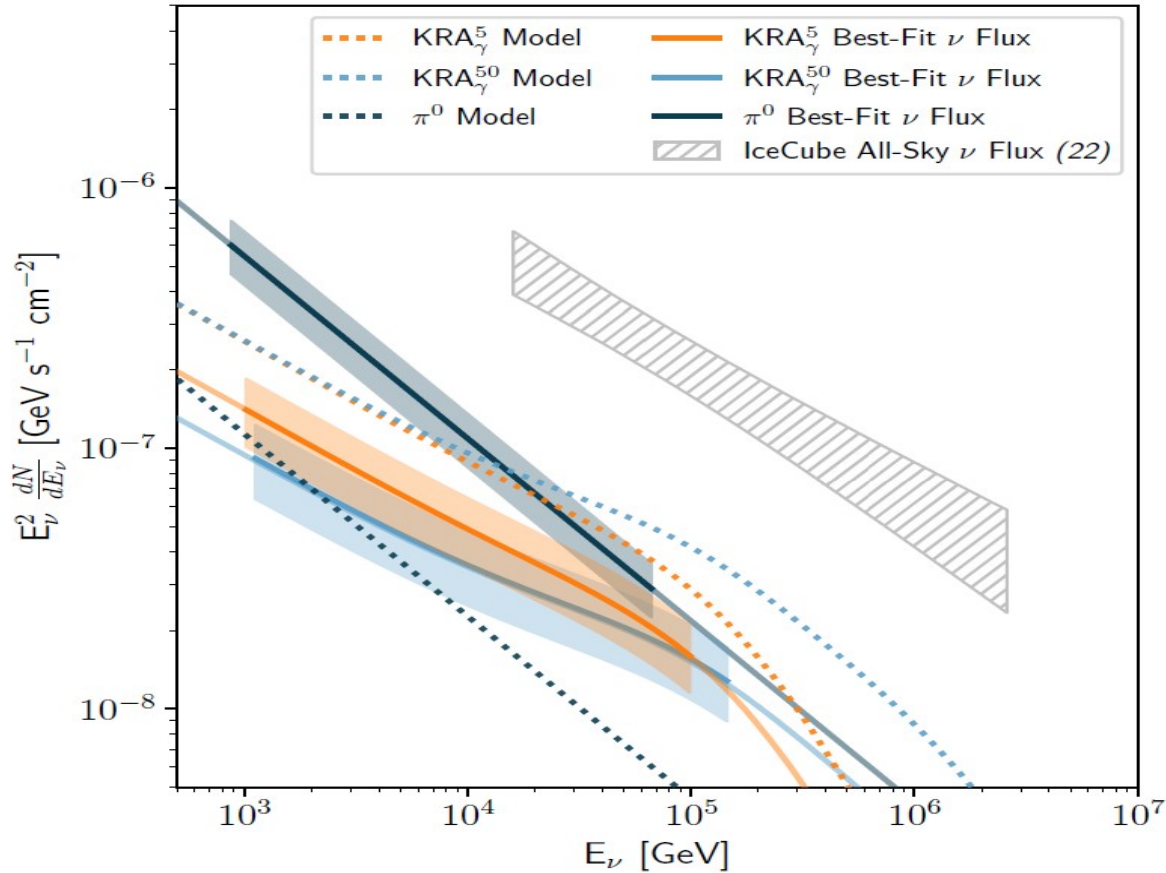
angular resolution  
still reasonable to  
detect diffuse signal  
from Galactic plane

59600 events, with about 7% being astrophysical (87% atmospheric, 6% tracks)

# cascades allow IC to look also to southern sky (where GC lies) where tracks are overwhelmed by atmospheric muons



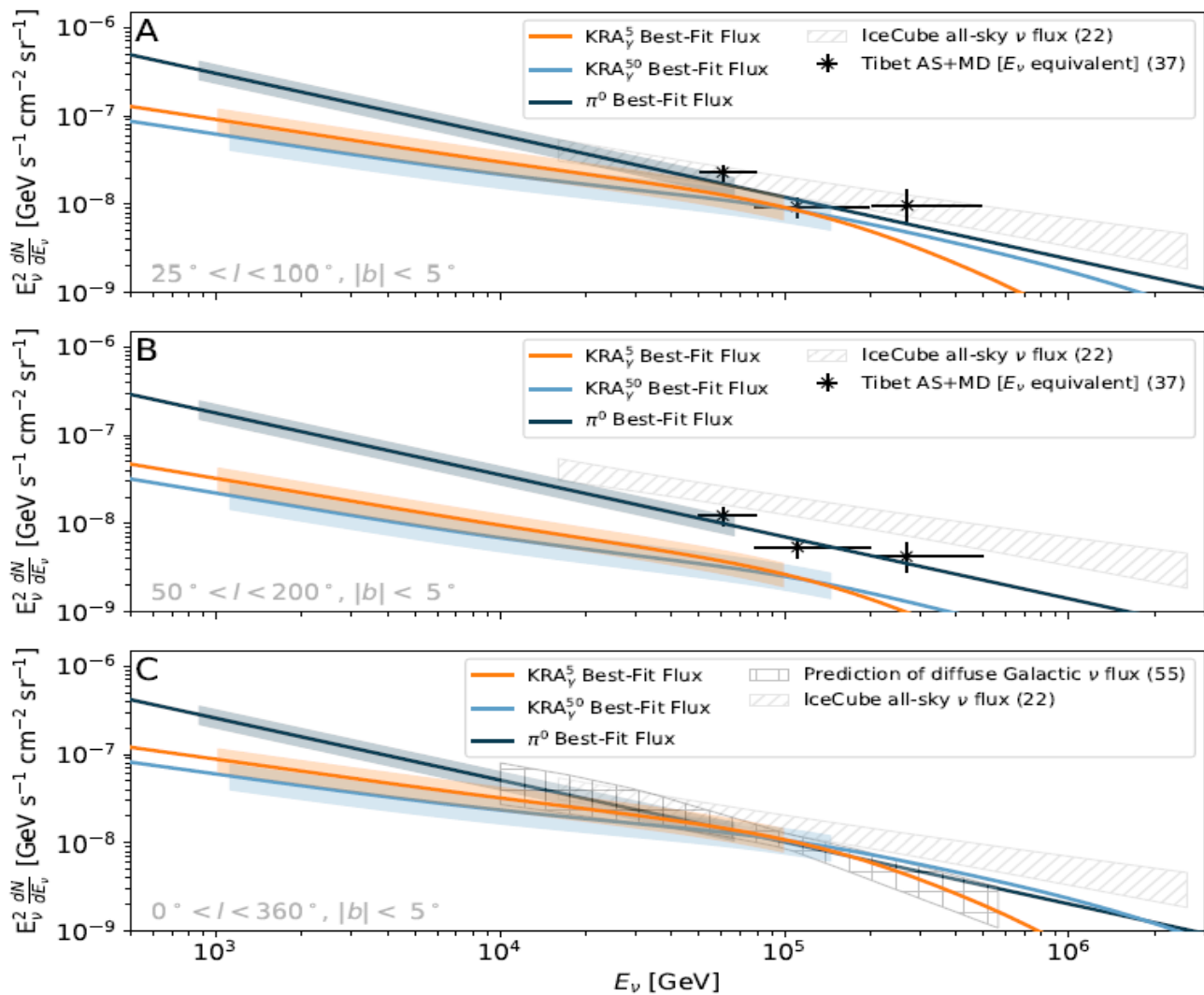
comparing with diffuse emission templates  $\rightarrow$   $4.5\sigma$  excess observed



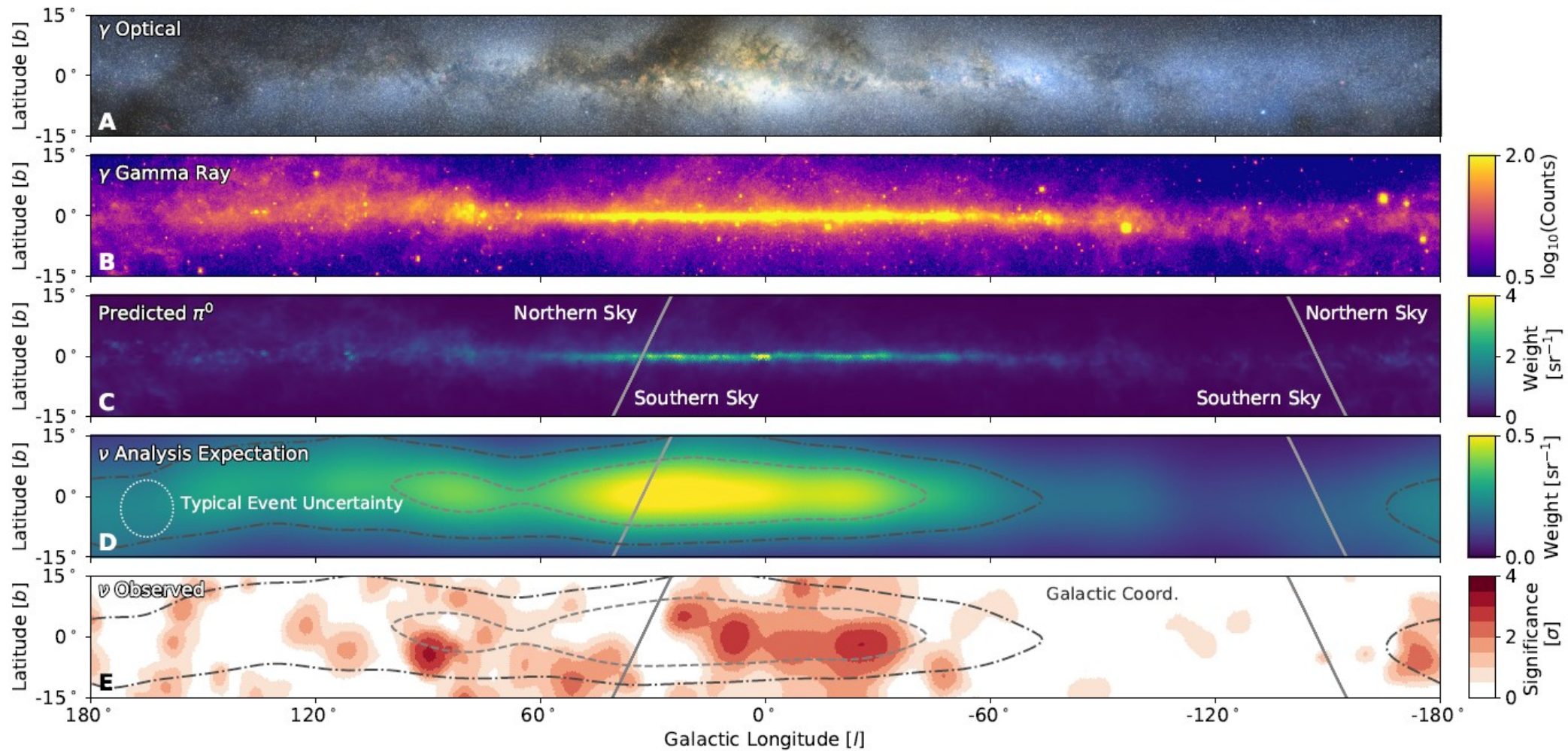
6 to 13% of astrophysical neutrinos come from the Galaxy at 30 TeV

Diffuse Galactic plane analyses	Flux sensitivity $\Phi$	p-value	Best-fitting flux $\Phi$
$\pi^0$	5.98	$1.26 \times 10^{-6}$ ( $4.71\sigma$ )	$21.8^{+5.3}_{-4.9}$
$KRA_\gamma^5$	$0.16 \times \text{MF}$	$6.13 \times 10^{-6}$ ( $4.37\sigma$ )	$0.55^{+0.18}_{-0.15} \times \text{MF}$
$KRA_\gamma^{50}$	$0.11 \times \text{MF}$	$3.72 \times 10^{-5}$ ( $3.96\sigma$ )	$0.37^{+0.13}_{-0.11} \times \text{MF}$
Catalog stacking analyses	p-value		
SNR	$5.90 \times 10^{-4}$ ( $3.24\sigma$ )*		
PWN	$5.93 \times 10^{-4}$ ( $3.24\sigma$ )*		
UNID	$3.39 \times 10^{-4}$ ( $3.40\sigma$ )*		

# flux in the Galactic plane



# the multimessenger picture starting to get completed



# the Galactic contribution may also alleviate the tension between fluxes measured with cascades and tracks

Palladino&Vissani 1601.06678

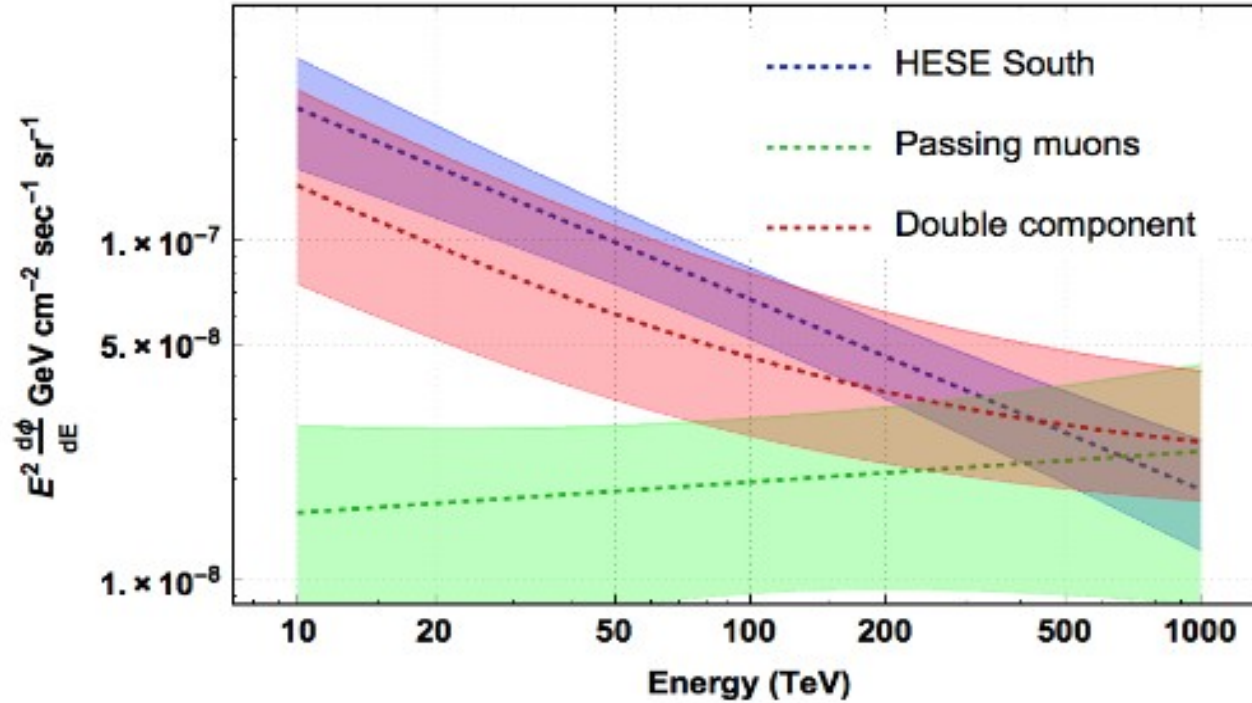
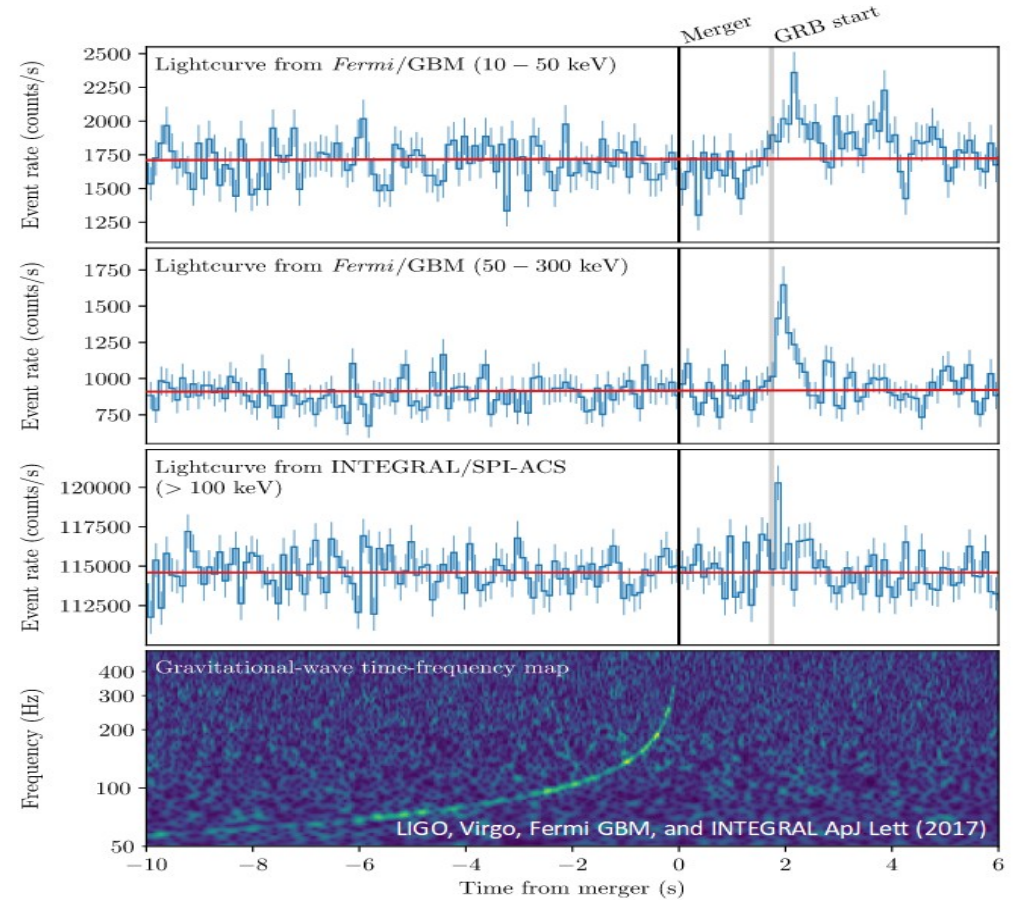
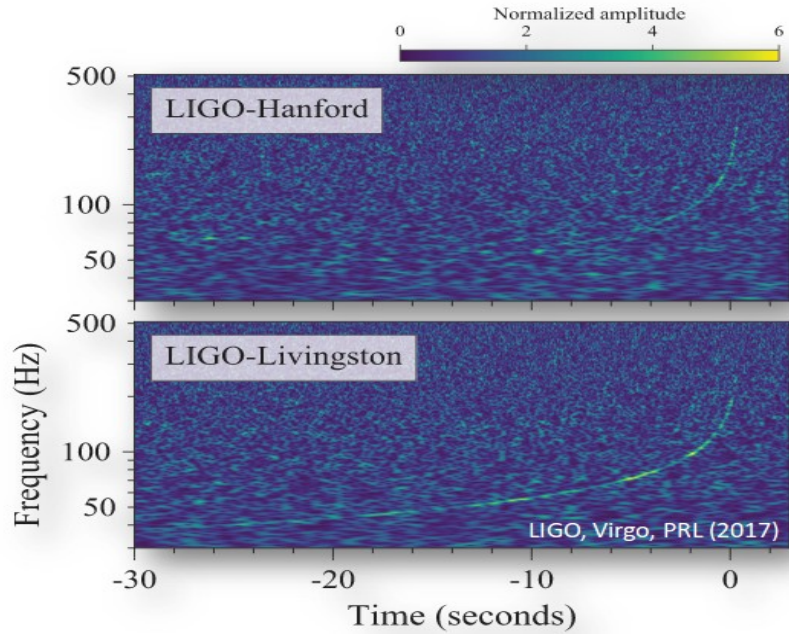
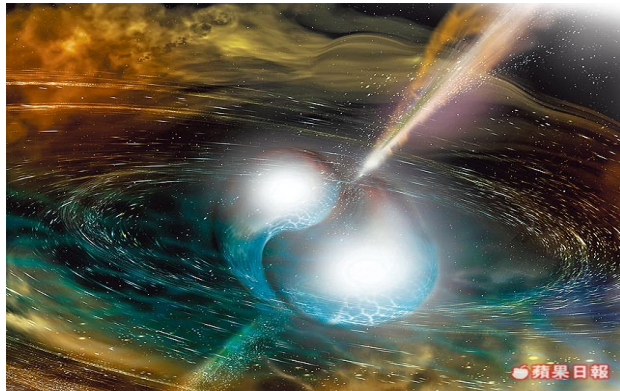


Figure 3: Comparison between the fluxes measured by IceCube from the Northern sky and Southern sky with the two-components flux.

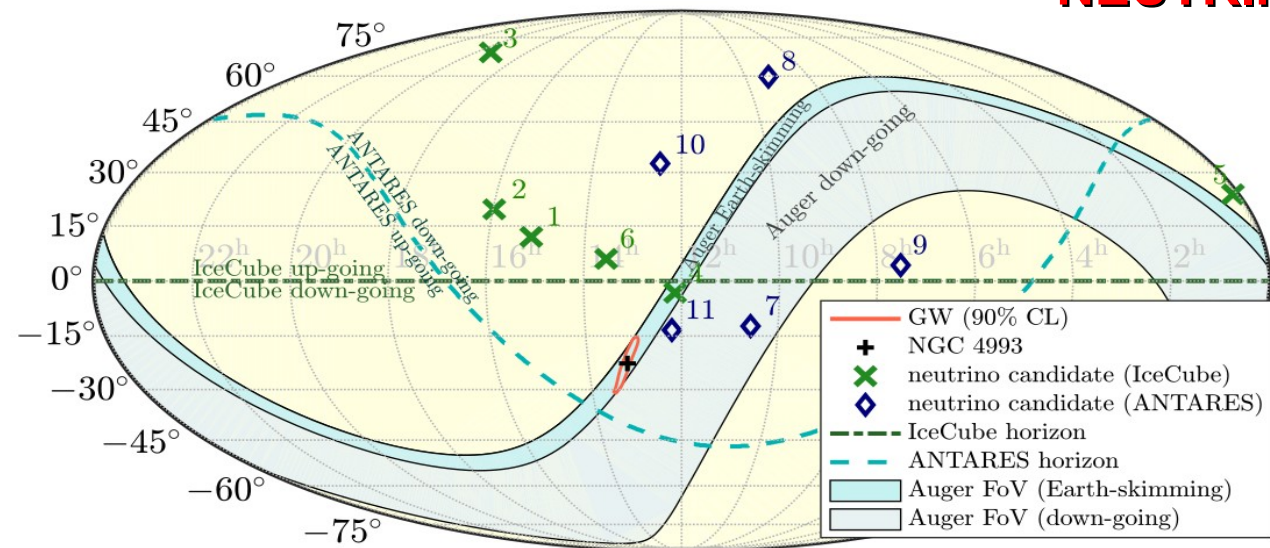
# IN 2017 LIGO/VIRGO OBSERVED BINARY NEUTRON STAR MERGER



also seen in gamma rays, and then optical, radio, X-rays , ...  
Gravitational waves: the new messenger



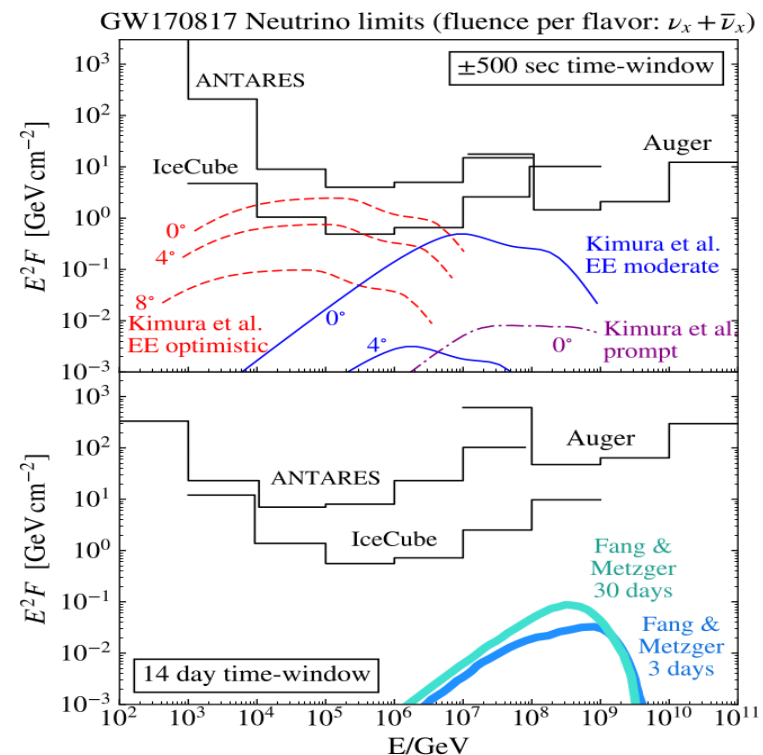




**Neutrinos searched in Auger, IceCube and Antares data, but nothing observed  
→ flux upperbounds**

**still waiting for simultaneous  
GW, gamma and neutrino observation**

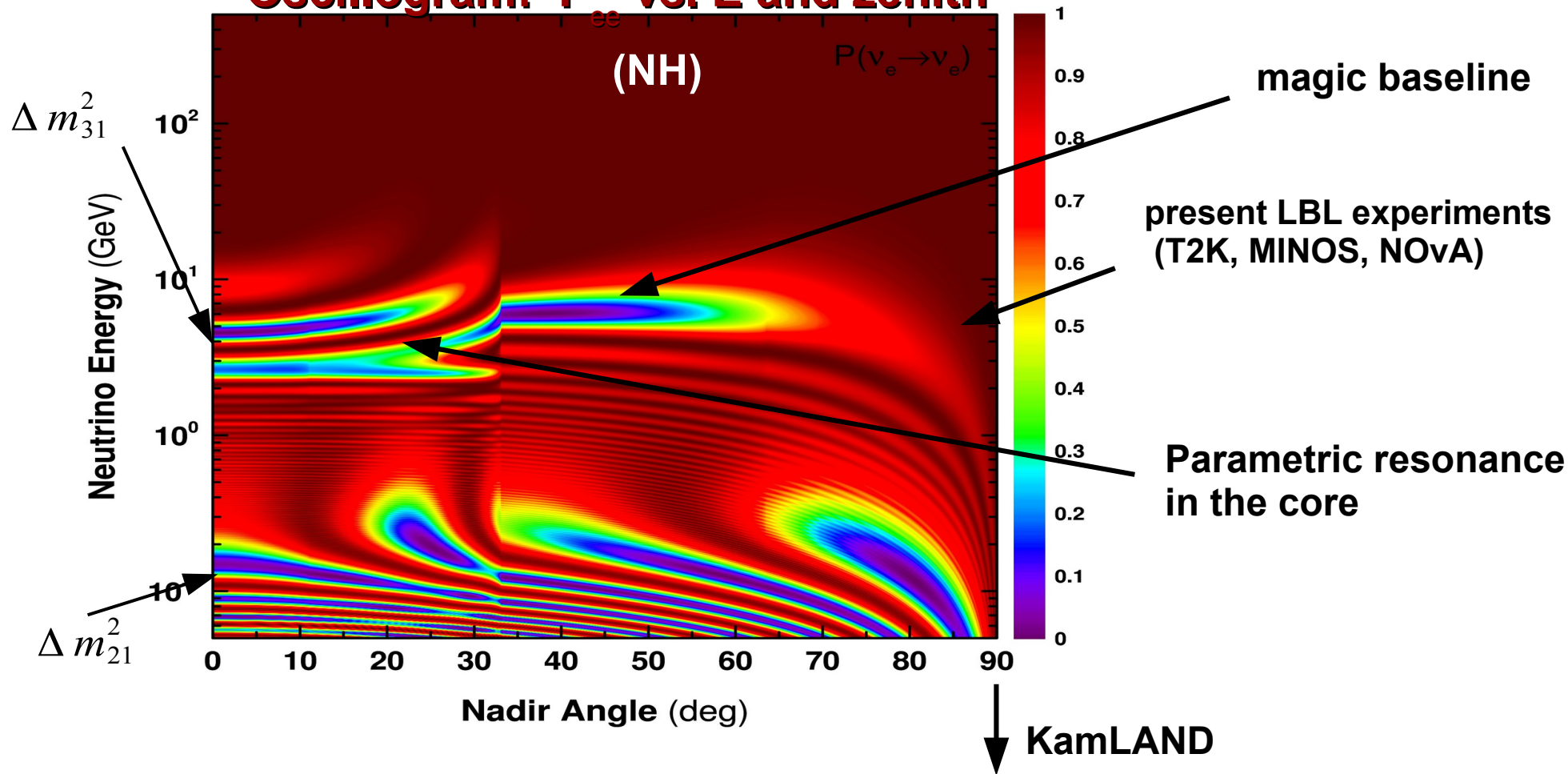
**(CRs, being charged, arrive much later)**



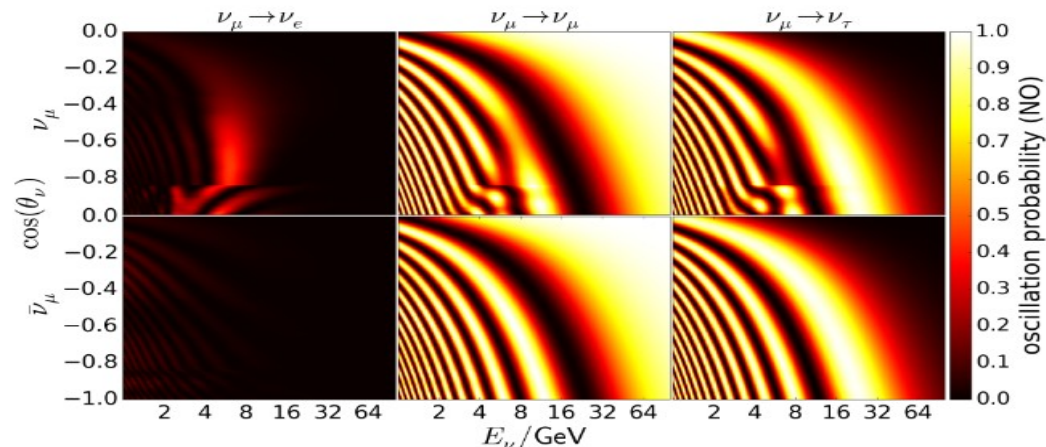
**Figure 2.** Upper limits (at 90% confidence level) on the neutrino spectral fluence from GW170817 during a  $\pm 500$  s window centered on the GW trigger time (top panel), and a 14 day window following the GW trigger (bottom panel).

# oscillations of atmospheric neutrinos (DeepCore, ORCA, HyperK,...)

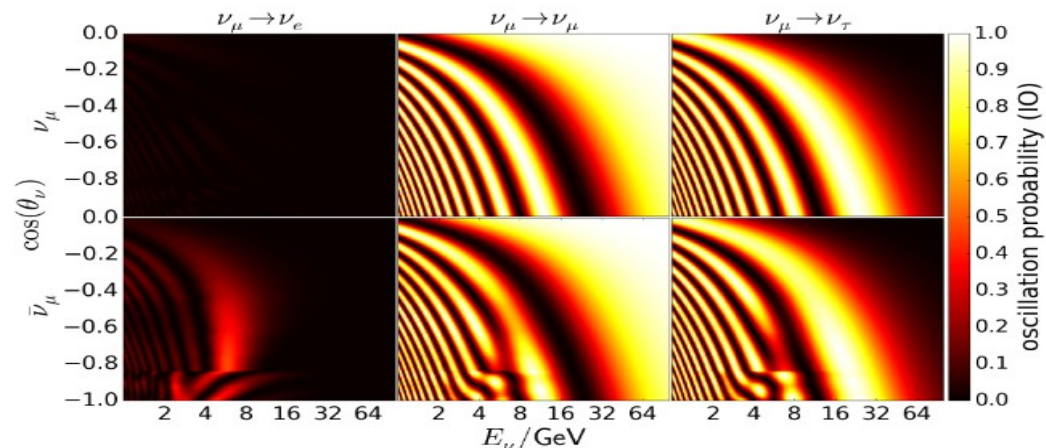
## Oscillogram: $P_{ee}$ vs. E and zenith



# can use atmospheric neutrinos from different directions and energies



(a) Normal Ordering



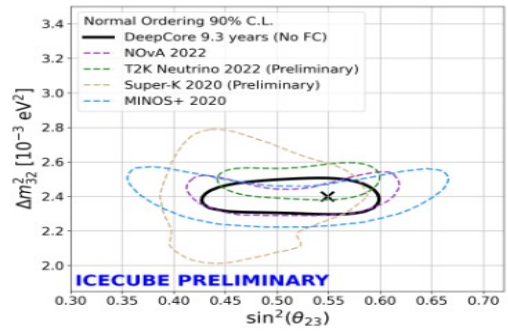
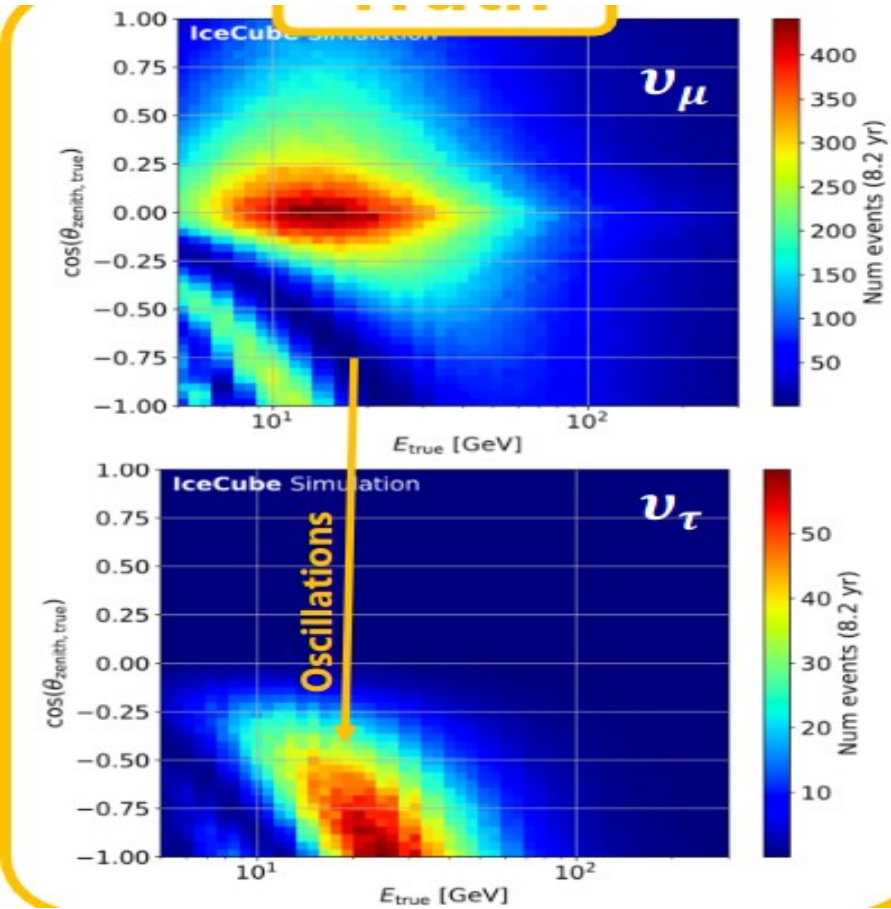
(b) Inverted Ordering

for IO the plots for  $E > \text{GeV}$  are approximately exchanged  
→ could determine hierarchy with atmospheric nus with  $E < 10 \text{ GeV}$  (SK and DeepCore favor NO)

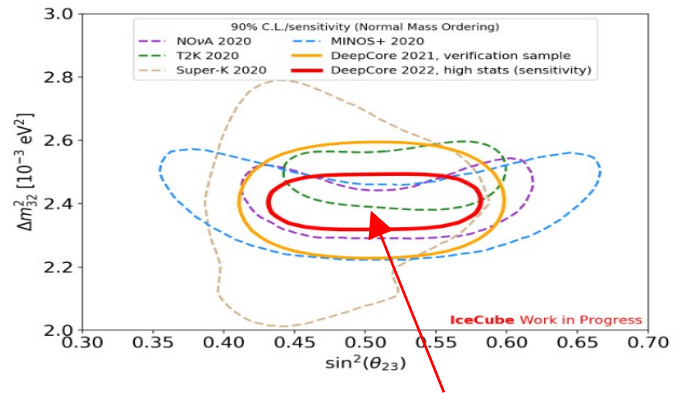
Note that:  $\sigma^{CC}(\nu N) \simeq 2 \sigma^{CC}(\bar{\nu} N)$

→ asymmetry is observable  
(also neutrino flux not identical to antineutrino flux)

future: IC upgrade, ORCA and INO (magnetized →  $\mu^+/\mu^-$ )



### Deep-Core studying $\nu_\mu$ disappearance



Deep-Core studying  $\nu_\tau$  appearance  
 (as excess of upgoing cascades)  
 → will improve mass-mix constraints

similar L/E as LBL experiments probing  $\Delta m^2_{31} = 2.5E-3 \text{ eV}^2$ :  
 L/E  $\sim 10^4 \text{ km}/20 \text{ GeV}$  or  $10^3 \text{ km}/2 \text{ GeV}$   
 but one in DIS regime, the other in the resonant QE regime

$$A = (I_O - N_O) / N_O$$

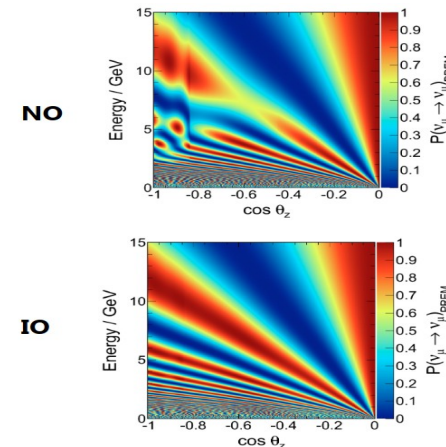
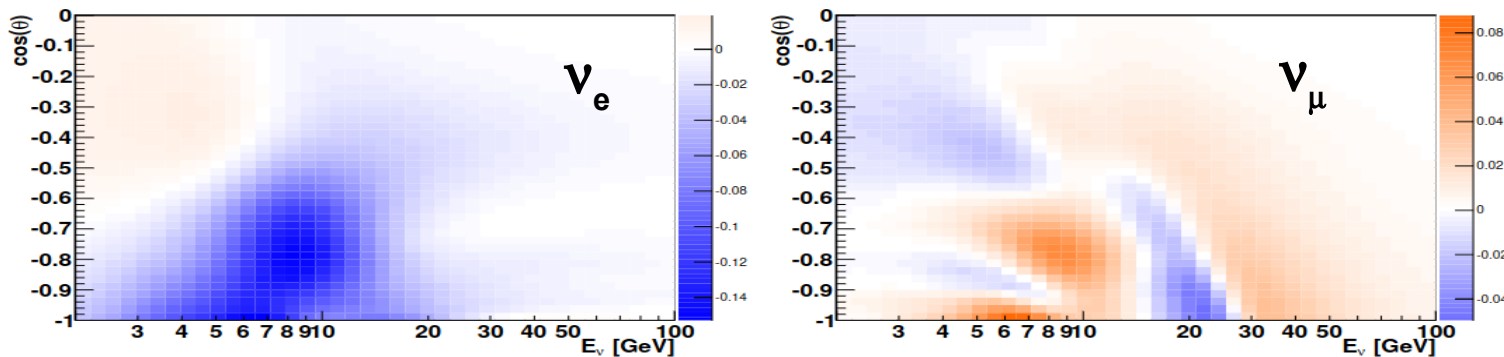
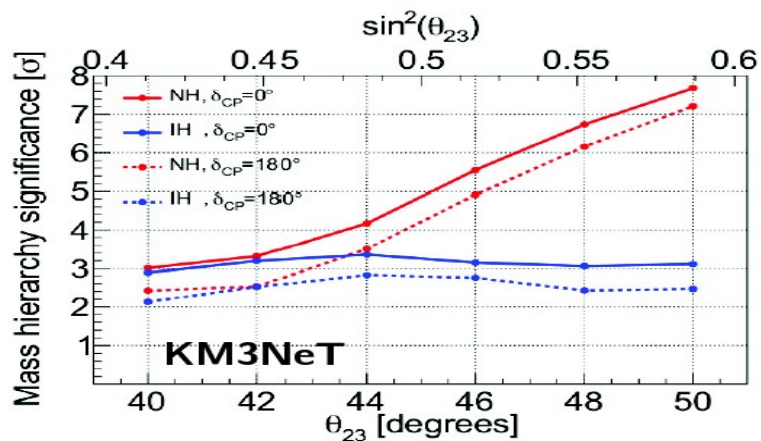


Figure 53: Asymmetry (as defined by Eq. 19) between the number of  $\nu + \bar{\nu}$  CC interactions expected in case of NH and IH, expressed as a function of the energy and the cosine of the zenith angle. The right (left) plot applies to muon (electron) neutrinos. A smearing of 25% is applied on the energy. On the angle,

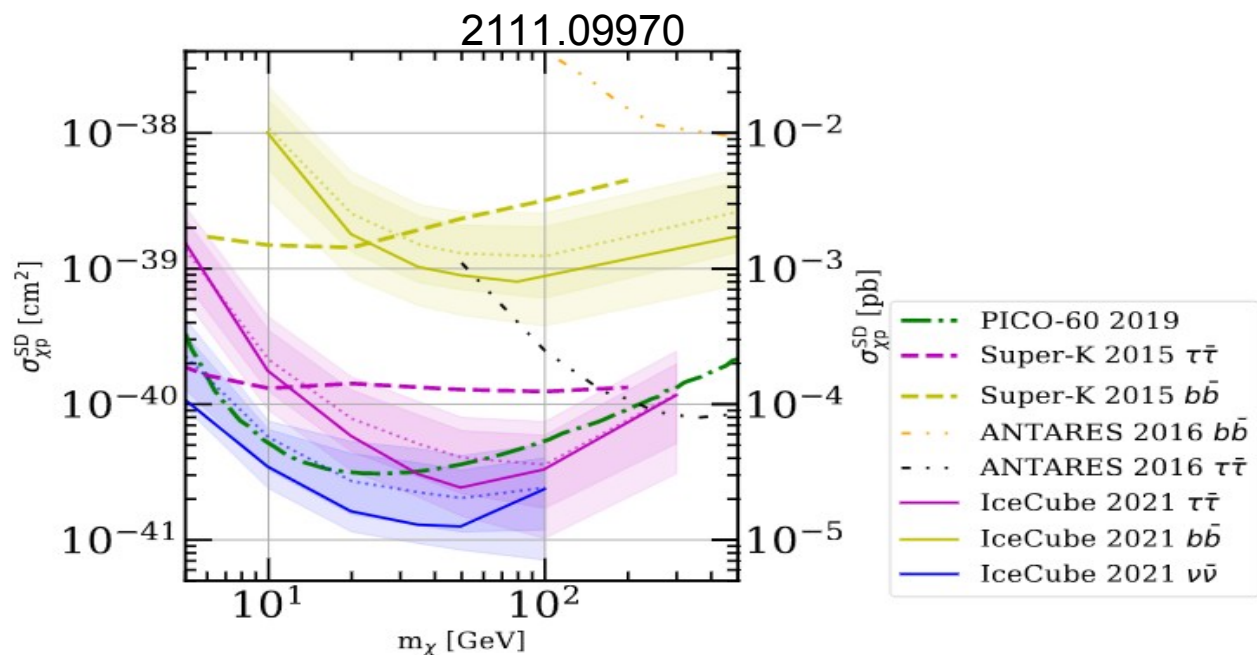
## IC upgrade & ORCA will determine mass hierarchy from matter effects induced nu-antineu differences



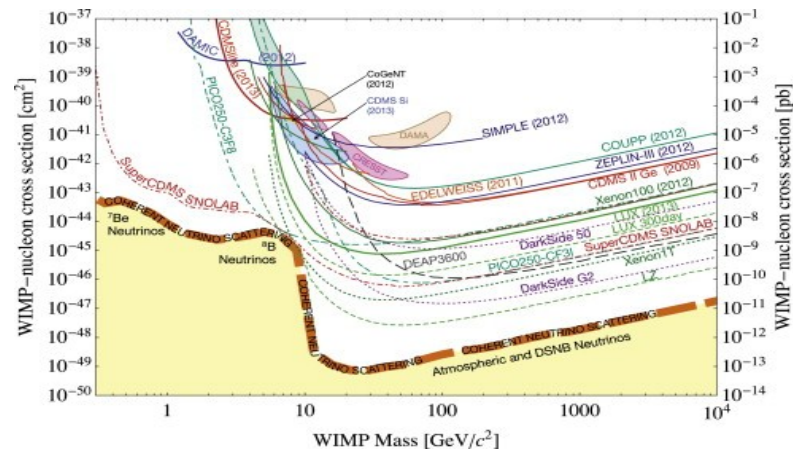
sensitivity on  $\nu$ MO depends on  $\theta_{23}$  octant and CP phase

# GeV neutrinos from WIMP annihilation in the Sun

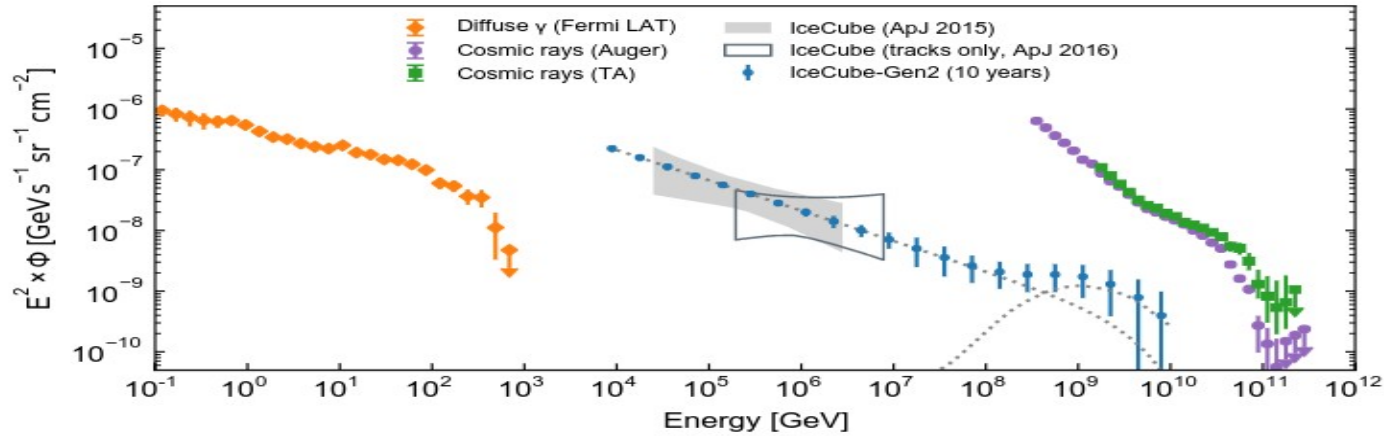
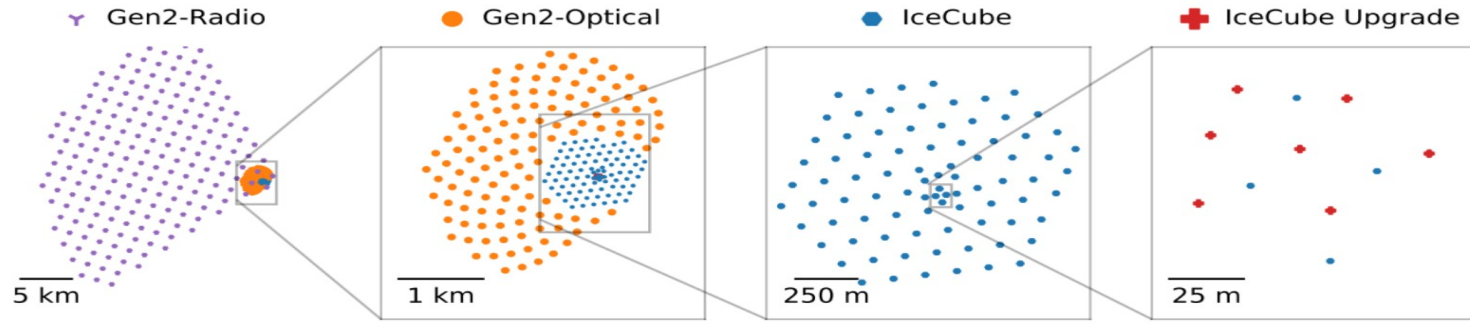
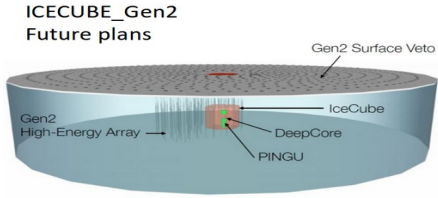
Spin-dependent WIMP nucleon xsect determines WIMP capture in Sun



spin-independent scattering in direct DM search



# the future



## history & future of WCD:

Kamiokande (3 kt; >7 MeV) ; SKamioka (50 kt, >5 MeV) ; HyperK (300 kt)

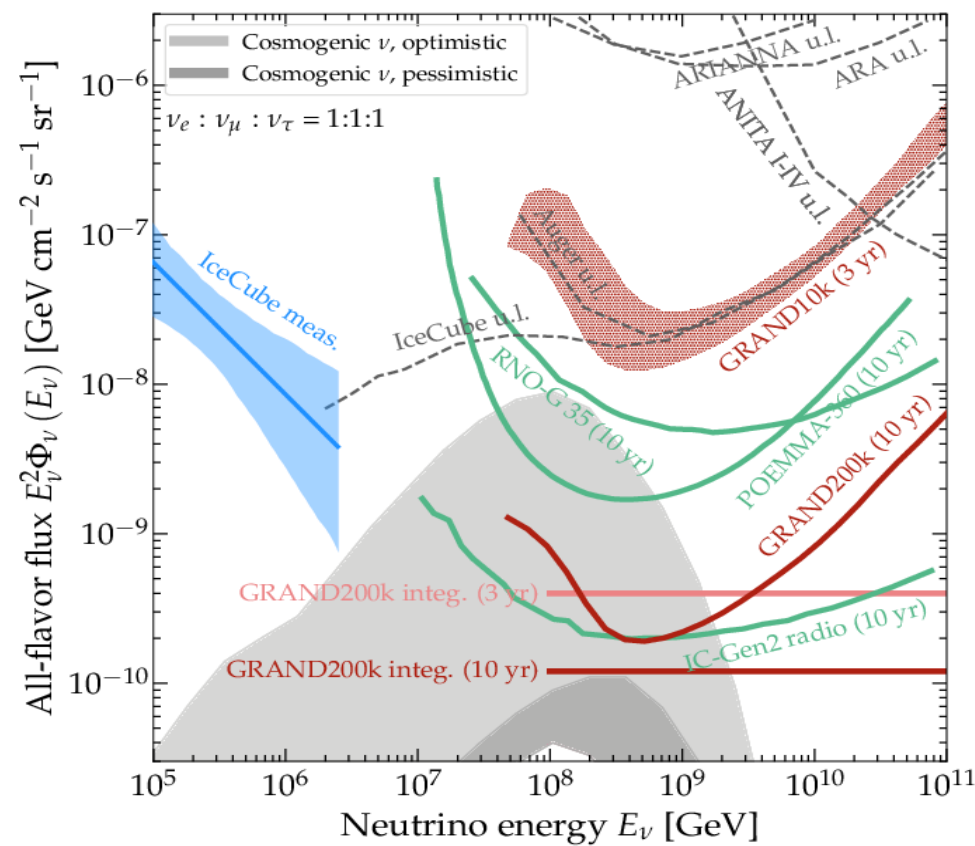
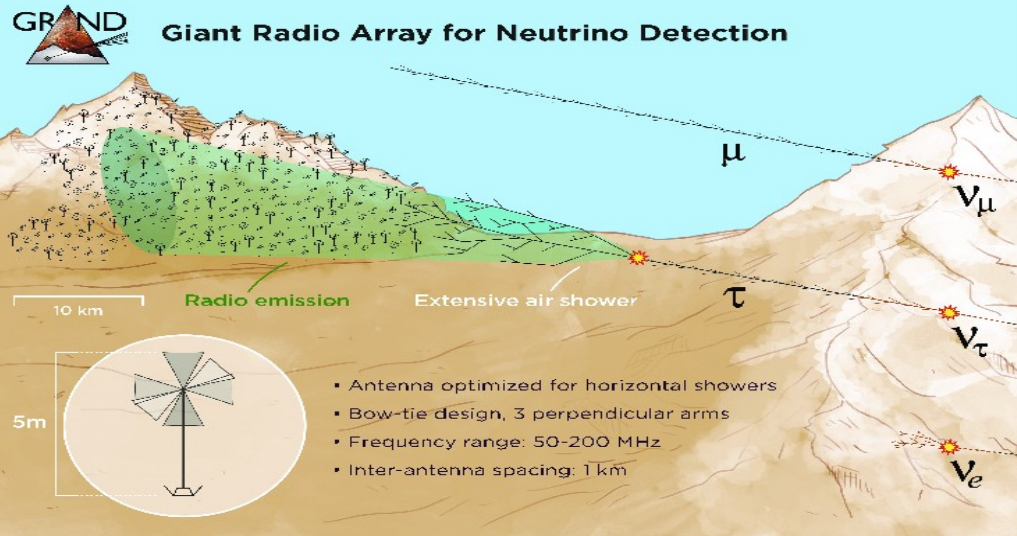
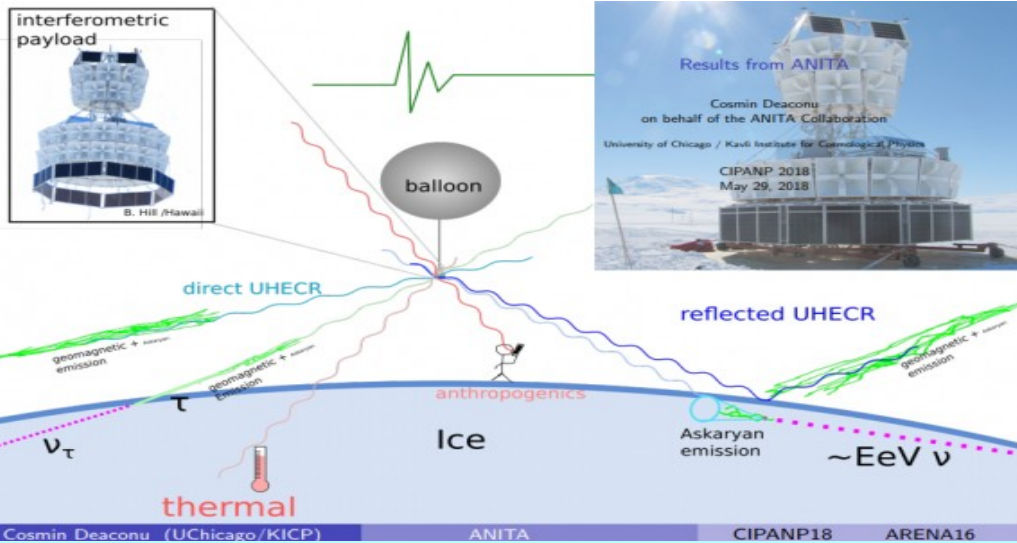
Amanda (Mt) ; Antares (10 Mt) ; IceCube (1 Gt, >100 GeV)

DeepCore (20 Mt, E>10 GeV) ; ORCA (4 Mt, E>5 GeV)

GVD (1Gt) ; ARCA (1.2 Gt, >100 GeV) ; Gen2 (8 Gt, >1 TeV)

P-ONE (3 Gt) ; TRIDENT (>10 Gt)

# radio to chase cosmogenic neutrinos





# CONCLUSIONS

Detection of astrophysical  $\nu$  produced a revolution

We are at the dawn of the era of high-energy neutrino astronomy

- there is now a 300 TeV neutrino from TXS & hints from other blazar sources also the Seyfert NGC1068 has been detected
- will provide clues on the sources of UHECRs, both Xgalactic and Galactic
- contribute to multimessenger astronomy

Also particle physics insights:

- measurement of cross section beyond accelerators & charm production
- $W$  resonance neutrinos , tau neutrino observations
- flavor studies will constrain source mechanism
- searching for cosmogenic neutrinos, neutrinos from WIMP annihilations, neutrino oscillations: mass ordering and mixings, .....

**the coming years will certainly be very exciting**

## Some books about neutrinos:

- **Physics of Neutrinos:**  
Fukugita and Yanagida (2003)
- **Probes of Multimessenger Astrophysics,**  
Maurizio Spurio (2018)
- **Neutrinos in Physics and Astrophysics:**  
Roulet & Vissani (2022)

## Check also:

<https://pdg.lbl.gov/2022/reviews/rpp2022-rev-neutrino-mixing.pdf>

<http://www.nu.to.infn.it>      **Neutrino Unbound page**  
(including links to experiments & recent conferences and schools)

