# MAGO: A SRF Cavity for Gravitational Wave Observation

**Detection Theory and MAGO Prototype** 

Lars Fischer, Erice 2. September 2023 (Lars.fischer@desy.de) in collaboration with T. Krokotsch, R. Löwenberg, G. Moortgat-Pick, M. Paulsen, K. Peters, M. Wenskat

#### MAGO Prototype Cavity

- Built in 2005 at INFN Genoa
- FNAL, INFN Genoa, UHH and DESY collaborate on the project
- Prototype designed for gravitational wave (GW) frequencies 5-10 kHz
- Radio frequency mode pumped by external oscillator
- Indirect coupling via mechanical modes
- **Direct coupling** of GW to RF mode



#### Lorentz Force Detuning<sup>1</sup>

• Deriving the equations of motion from

$$\mathcal{L}_{mech} = \frac{1}{2} \sum_{n} 2U_n (e'_n{}^2 - b'_n{}^2) + \sum_{l} \left( \frac{1}{2} M \dot{q}_l^2 - \frac{1}{2} M \omega^2 q_l^2 + q \right)$$

• Additional term appears in the mechanical EoMs

$$\ddot{q}_l(t) + \frac{\omega_l}{Q_l} \dot{q}_l(t) + \omega_l^2 q_l(t) = \frac{1}{M} \left( f_l(t) + f_l^{ba}(t) \right)$$

 Back-action of the EM fields on the cavity shell damps the signal strength

 $f_l^{ba}(t) = 2V_{cav}^{-1/3} \sqrt{U_0 U_1} C_{01}^l b_0(t) b_1(t)$ 

• Measure RF power in higher mode



## **Direct coupling (Gertsenshtein effect)**

### Indirect coupling

- GW deform the cavity and excite mechanical modes
- Tidal force density from equation of geodesic deviation

 $\vec{f}(t,\vec{x}) = -\rho(\vec{x})R_{0i0j}(t)x_j\vec{e_i}$ 

- Generalized force density on one mode  $f_l(t, k_z) = -\frac{\omega_g^2}{2} M V_{cav}^{1/3} [h_+ \Gamma_+^l - h_\times \Gamma_\times^l] e^{i(\omega_g t - k_z z)}$
- Separation of displacement  $\vec{u}(t, \vec{x}) = \sum_l \vec{\xi}_l(\vec{x})q_l(t)$

$$\Gamma_{+}^{l} = \frac{V_{cav}^{-1/3}}{M} \int_{V_{cav}} d^{3}x \rho(\vec{x}) (x\xi_{l,x}(\vec{x}) - y\xi_{l,y}(\vec{x}))$$

$$\Gamma_{\times}^{l} = \frac{V_{cav}^{-1/3}}{M} \int_{V_{cav}} d^{3}x \rho(\vec{x}) (x\xi_{l,y}(\vec{x}) + y\xi_{l,x}(\vec{x}))$$

Coupling coefficient for GW from arbitrary directions

 $h_{ij}^{TT}(\alpha,\beta) = \mathcal{R}(\alpha,\beta)h_{ij,z}^{TT}\mathcal{R}^{T}(\alpha,\beta)$ 

 $|\Gamma_+|$  for  $f_{32} = 3.9642$  kHz

## Signal power

• Signal power with damping of the EM field back-action

 $P_{sig} = \frac{\omega_1}{Q_{cpl}} \omega_g^4 U_0 \left| \frac{1}{2} \frac{\omega_1^2 C_{01}^l (h_+ \Gamma_+^l + h_\times \Gamma_\times^l)}{\beta_1 \beta_l - \gamma_1 \gamma_l} - \frac{\beta_l H(\kappa_1 \eta_{01}^E + \lambda_1 \eta_{01}^B)}{\beta_1 \beta_l - \gamma_1 \gamma_l} \right|^2$ 

• With the damping term

$$\gamma_1 \gamma_l = \frac{1}{M} V_{cav}^{-2} U_0 \left( \omega_1 C_{01}^l \right)^2$$



- GW induces current in Maxwells equations
- The Lagrangian for the EM field

 $\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} - j^{\mu}_{eff} A_{\mu}$ 

• Current for GW strain  $h_{\mu\nu}$ 

 $j^{\mu}_{eff} = \partial_{\nu} (\frac{h}{2} F^{\mu\nu} + h^{\nu}_{\ \alpha} F^{\alpha\mu} - h^{\mu}_{\ \alpha} F^{\alpha\nu})$ 

• For monochromatic GW propagating in z direction the overlap factors are

 $\eta_{0n}^E \coloneqq \frac{1}{H\sqrt{U_0 U_n}} \int_{V_{cav}} d^3 x \, H_0(\vec{x}) \epsilon_0 \vec{E}_0(\vec{x}) \vec{E}_n(\vec{x})$ 

 $\eta_{0n}^{E} \coloneqq \frac{1}{H\sqrt{U_{0}U_{n}}} \int_{V_{cav}} d^{3}x \, H_{0}(\vec{x}) \frac{1}{\mu_{0}} \vec{B}_{0}(\vec{x}) \vec{B}_{n}(\vec{x})$ 

• Where H is the normalization of

 $H_0(\vec{x}) = (h_+(x^2 - y^2) + 2h_\times xy)$ 

• Gertsenshtein effect has small effect at low frequencies





Plus polarized GW coupling to vibrational mode

#### **RF mode overlap factor**

• RF mode overlap in cavity perturbation theory

 $\vec{E}'_n = \vec{E}_n + \sigma \vec{E}_n^{(1)} + \mathcal{O}(\sigma^2)$ 

• Displacement field for isotropic elastic solids

$$\vec{f}(t,\vec{x}) = \rho(\vec{x})\frac{\partial^2 \vec{u}}{\partial t^2} - (\lambda + \mu)\nabla(\nabla \cdot \vec{u}) - \mu\Delta \vec{u}$$

• Overlap factor of RF modes for one mechanical mode

$$C_{01}^{l} = \frac{V_{cav}^{1/3}}{2\sqrt{U_0 U_1}} \int_{\partial V_{cav}} d\vec{S} \cdot \vec{\xi}_l(\vec{x}) \left[ \frac{1}{\mu_0} \vec{B}_0(\vec{x}) \vec{B}_1(\vec{x}) - \epsilon_0 \vec{E}_0(\vec{x}) \vec{E}_1(\vec{x}) \right]$$



Signal power without damping

#### **MAGO Sensitivity**

- Scanning mode, GW is on resonance with RF mode difference  $\omega_g = \omega_1 - \omega_0$
- Broadband mode, GW of all frequencies couple to one mechanical mode  $\omega_l$  with  $\omega_l = \omega_1 - \omega_0$ .
- High sensitivity in high frequency regime  $h_{min} \sim \mathcal{O}(10^{-21})$



#### Purely polarized GW couplings to the E field of the $TE_{011}$ mode



Purely polarized GW couplings to the B field of the  $TE_{011}$  mode

**Detector simulation with COMSOL** 



Absolute E-field distribution of the  $TE_{011}$  mode, proposed by the MAGO collaboration<sup>2</sup>

[1] Löwenberg et al., arXiv/2307.14379 (2023)
[2] Bernard et al., arXiv:gr-qc/0203024v1 (2002)



#### **CLUSTER OF EXCELLENCE**

QUANTUM UNIVERSE

