# Quantum metrology of displacements



um information

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Metrology: estimation of a parameter, through measurements.

The estimation is always performed by averaging over N measurements, so that (central limit theorem), the error of the average goes as  $1/\sqrt{N}$ 



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Quantum Metrology: estimation of a parameter with increased precision (thanks to quantum effects, e.g. entanglement)

Usually:  $\sqrt{N}$  enhancement: the error goes as 1/N



# Want to estimate a parameter $\varphi$ written onto a probe by a transformation $U\varphi$



Want to estimate a parameter  $\varphi$  written onto a probe by a transformation  $U_{\varphi}$ 









# Optimize fidelity (overlap) $\rightarrow$ DISCRIMINATION " $U_{\varphi}$ is present or not"





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#### Optimize **Q Fisher Information** → ESTIMATION

"what's the value of  ${\cal Q}$  ?"



Optimize fidelity (overlap)  $\rightarrow$  DISCRIMINATION " $U_{\varphi}$  is present or not"

Optimize **Q Fisher Information** → ESTIMATION

, "what's the value of arphi ?"

the metric in Hilbert space (to measure distances)



# Compare quantum strategies to the corresponding classical strategy





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#### Same number of uses of $~U_{arphi}$

Same employed energy





use 
$$-U_{\varphi}$$
 in parallel:







Classical strategies:

use

-*U*\_\_\_\_



in parallel:











Quantum strategies:



the N transformations act on an entangled state





Quantum strategies:



(Heisenberg bound)

the N transformations act on an entangled state



the N transformations act on an entangled state

Note: entanglement at the measurement u stage is useless!

Classical strategies:

Quantum



the N transformations act on an entangled state

#### Squeezing vs entanglement arXiv:1901.07482



### Squeezing

arXiv:1901.07482



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 Take the energy used by N coherent (classical) probes



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## Squeezing

- •Take the energy used by N coherent (classical) probes
- •Use it to squeeze one probe



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Noiseless case is simple and that's where the  $\sqrt{N}$  quantum advantage is

# The noisy case is a **mess**, but the noiseless case is an upper bound.





 $U_{lpha,arphi} = U_{arphi}^{\dagger} e^{ilpha G} U_{arphi}$ 

#### with random rotation $\varphi$



$$U_{lpha, \varphi} = U_{\varphi}^{\dagger} e^{i lpha G} U_{\varphi}$$

#### with random rotation $\varphi$

# Want to estimate $\alpha$ , don't care about $\varphi$ (it's just noise)



# $U_{\alpha,\varphi} = U_{\varphi}^{\dagger} e^{i\alpha G} U_{\varphi}$ with random rotation $\varphi$ Want to estimate $\alpha$ , don't care about $\varphi$

NOTE: the channel rotates, not the state

(it's just noise)



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## NOTE: the channel rotates, not the state

Example: displacement (nonzero  $\alpha$ if the axion and em field interact)  $D(\alpha, \varphi) = e^{\alpha(e^{i\varphi}a^{\dagger} - e^{-i\varphi}a)}$ 



#### Our results:

## 1) A squeezed state is an optimal input

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# A squeezed state is an optimal input Squeezing+photodetection is an optimal measurement





## Our results:

A squeezed state is an optimal input
Squeezing+photodetection is an optimal measurement
we retain the√n advantage over the classical strategies! Just as the noiseless case!

### Homodyne detection does not work!

#### A surprising connection to Rayleigh's curse in imaging, which is a particular case of the class of channels we consider



#### Results



### Results



Large Fisher info=good estimation Small fidelity between initial and final state=good discrimination





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- Wait for the axion (displacement of the state)
- 3) Anti-squeeze+photodetection





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- 3) use convexity of the QFI to show that, in the limit  $\alpha \rightarrow 0$ , the noise doesn't matter: the QFI is equal to the QFI averaged by the noise.
- 4) for optical displacements show that the QFI averaged over noise is bounded by the average energy of the state: both Fock and sq. vacuum saturate the bound



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#### Open problems! (see Quntao's paper)

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Adapt protocols to what we can do

Not all required transformations can be easily implemented in the lab, the single mode analysis may not be

What's the best figure of merit?



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Typically we compare quantum and classical strategies that use the **same resources** (energy or uses of the channel)

In practice  $\rightarrow$  other figures of merit may be more relevant (scan rate!)



#### What did I say?



- 1. Quantum metrology noiseless case  $\sqrt{N}$
- 2. QFI vs fidelity: overlap between initial and final state
- 3. Spreading channels
- 4. Squeezing input and antisqueezing+photodetection optimal.

Take home message

Squeezing is optima to estimate displacements with random (and irrelevant) phase in the noiseless case Lorenzo Maccone maccone@unipv.it

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Squeezing metrology: arXiv:1901.0748

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