

PADME Run III

Theoretical perspective and expected sensitivity



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Outline

Targeting the X_{17} anomaly

PADME Run-3: a “resonant” fixed target experiment

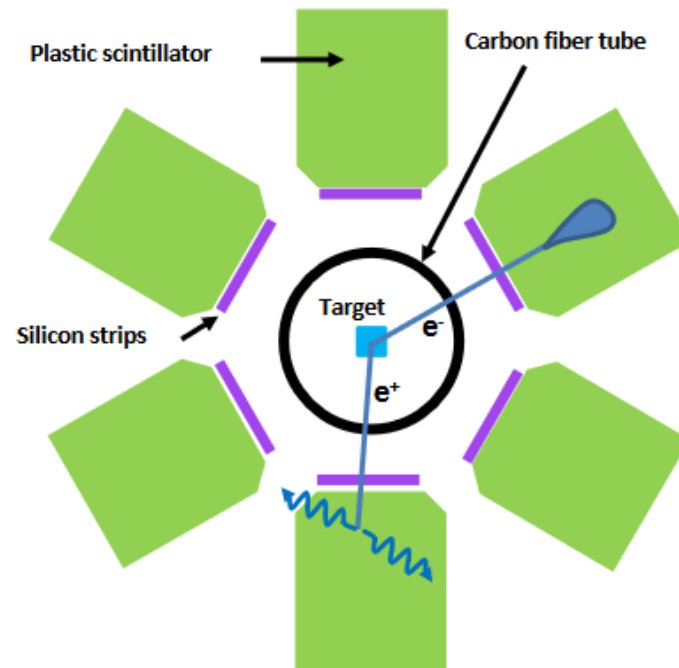
Run III projections

The X17 anomaly

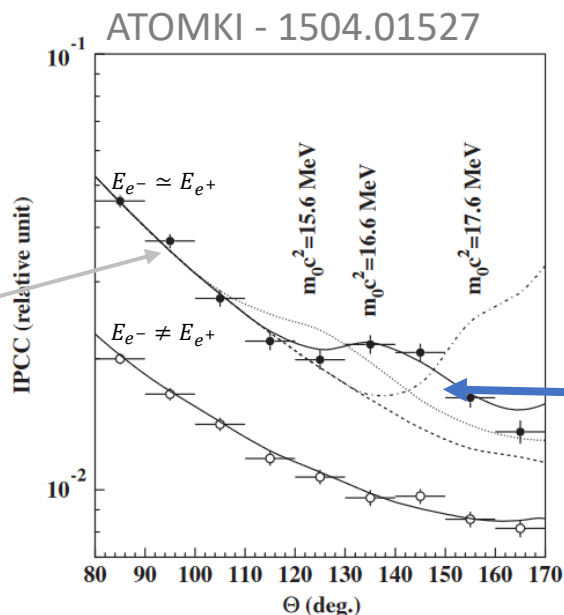
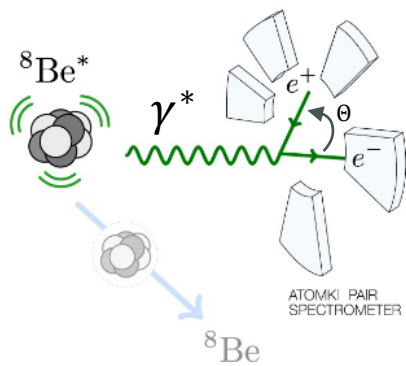
- The signal: a 17 MeV boson in the ATOMKI spectrometer?

2209.10795 1504.01527 2104.10075

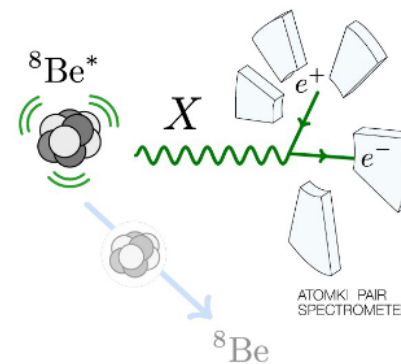
- Production of excited nuclei ^{12}C , ^8Be and ^4He , followed by radiative decay $N^* \rightarrow N \gamma^* \rightarrow N e^+ e^-$
- Very large excess, cannot be statistical. But from one collaboration only.



The SM signal: $N^* \rightarrow N \gamma^* \rightarrow N e^+ e^-$



NP sigma: $N^* \rightarrow N V \rightarrow N e^+ e^-$



Which NP bosons for the X17 ?

- No current nuclear physics explanation
- Use spin-parity analysis for hints of possible NP candidate
 - Scalar (0+) excluded by ^8Be data
 - ^4He data mixes 0+ and 0- excited nuclei

$J_* = L \oplus J_X$
 $P_* = (-1)^L P_X$

$N^* \rightarrow N V$	V			
N_*	0^+	0^-	1^-	1^+
$^4\text{He } 0^+$	S	—	—	P
$^4\text{He } 0^-$	—	S	P	—
$^8\text{Be } 1^+$	—	P	P	S, D

T. Tait from Feng et al. 2016-2020

- The most obvious requirements is on the mass !

$m_{X_{17}} =$	$16.70 \pm 0.35 \pm 0.50 \text{ MeV}$	(^8Be ATOMKI 2015)
	$17.01 \pm 0.16 \text{ MeV}$?	(^8Be proceedings 2018)
	$16.94 \pm 0.12 \pm 0.21 \text{ MeV}$	(^4He ATOMKI 2021)
	$17.03 \pm 0.11 \pm 0.20 \text{ MeV}$	(^{12}C ATOMKI 2022)

Re-analysis but no systematics !

Exp1	Exp2	Average
16.86(6)	17.17(7)	17.01(16)

Keeping the 2015 systematics for ^8Be + naive combination of both ^8Be measurements and of ^4He and ^{12}C

$$m_{X_{17}} \approx 16.99 \pm 0.04 \text{ (stat)} \pm 0.21 \text{ (syst)} \text{ MeV}$$

Fixing notations: explicit Lagrangian

- X17 as a light vector V^μ , potentially with both vector and axial couplings

$$\mathcal{L} \supset -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + \frac{1}{2}M_V^2 V_\mu V^\mu + \sum_{f=\ell,q} V_\mu \bar{f} (g_{Vf} + \gamma^5 g_{Af}) f \longrightarrow g_{Vf} \text{ corresponds to } eq_f \varepsilon \text{ in dark photon models}$$

- X17 as an axion-like particle (ALP) a , interacting via $\bar{f}\gamma^\mu\gamma^5 f$

$$\mathcal{L} \subset \frac{1}{2}(\partial_\mu a)(\partial^\mu a) - \frac{1}{2}m_a^2 a^2 + \sum_{f=\ell,q} \frac{g_{af}}{2}(\partial_\mu a) \bar{f} \gamma^\mu \gamma^5 f \longrightarrow \frac{g_{af}}{f_a} \text{ corresponds to } \frac{Q_{af}}{f_a} \text{ in the literature}$$

Most of the e^+/e^- -driven production rates shown in the rest of the talk satisfy approximately:

$$m_\ell g_{al} \longleftrightarrow g_{V\ell}$$



e^+/e^- -driven production rates are pretty agnostic concerning the X17 nature/couplings

The X17 couplings

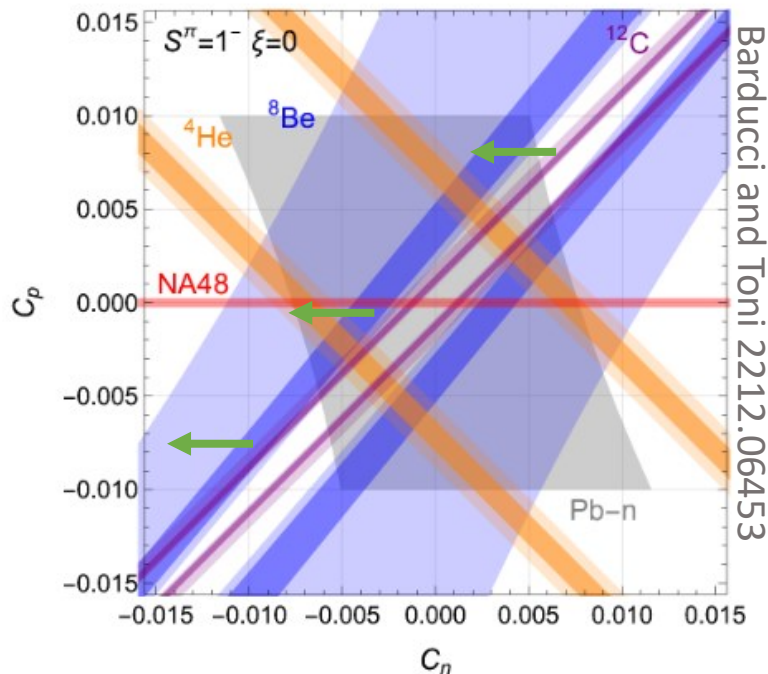
- Need coupling to e^+e^- AND a large couplings to quarks to fit the excess

$$|g_{Vu} + 2g_{Vd}| \sim [0.6 \cdot 10^{-3}, 3 \cdot 10^{-3}] \quad \longrightarrow \quad |2g_{Vu} + g_{Vd}| \lesssim 0.4 \cdot 10^{-3}$$

Protophobia needed to escape NA48 $\pi^0 \rightarrow \gamma V$

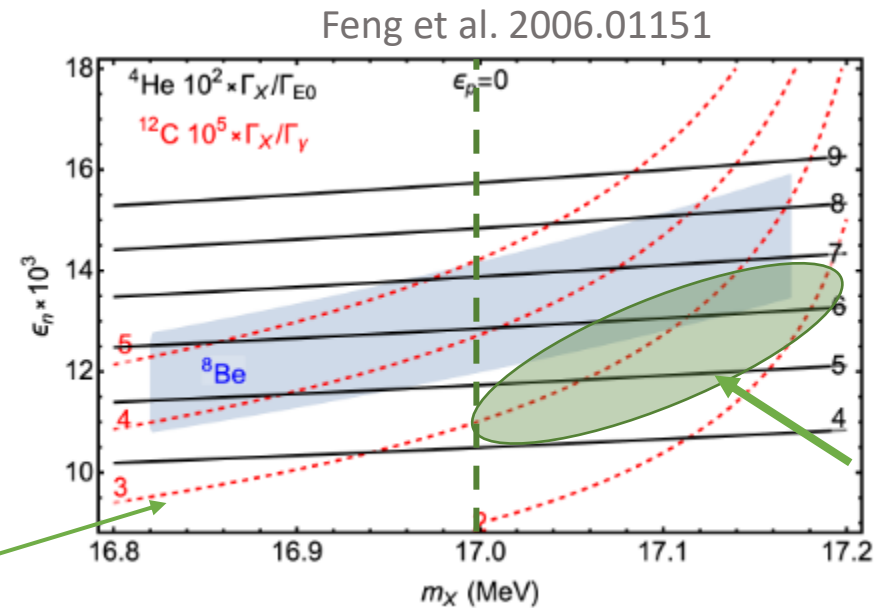
See e.g. 1608.03591, assumed BR_{ee} at 1.

- The couplings and mass both enters in fitting the excess \rightarrow simultaneous fit of ^{12}C , ^8Be and ^4He required, not available yet with the latest data



It may be that the central value for the mass will move to accomodate the fit in couplings!

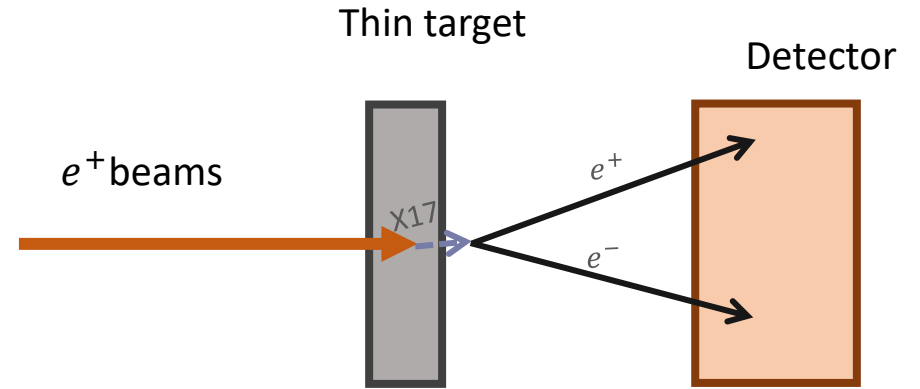
Measured ^{12}C is just below this line



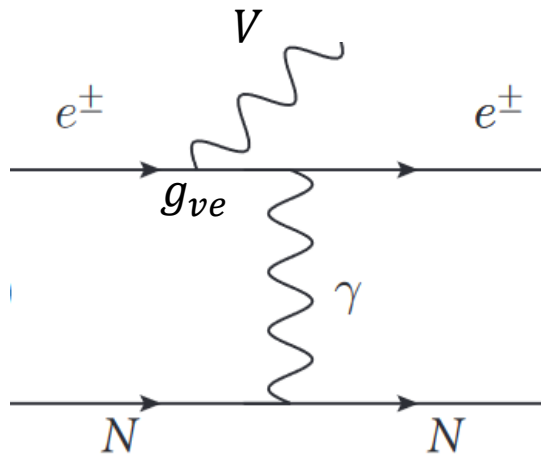
A full mass + couplings fit would put us around here?

Resonant X17 production on a thin target

- Main idea: use resonant production and search for visible X17 decay in $e^+ e^-$
- This is more or less the PADME setup, but no missing energy in the final state

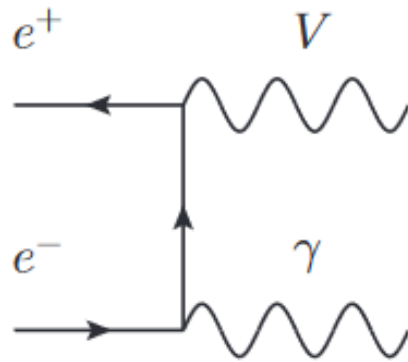


Bremsstrahlung process

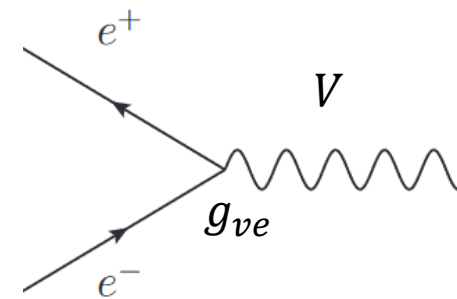


$$\sigma_{brem} \sim \frac{g_{ve}^2}{M_V^2} \alpha^2 Z^2$$

Associated annihilation process



$$\sigma_{assoc} \sim \frac{g_{ve}^2}{2m_e E_+} \alpha Z$$



$$\sigma_{res} \sim \frac{g_{ve}^2}{2m_e} \pi Z \delta(E_+ - E_{res})$$

Resonant process

→ Cross-section x100 times larger

→ Scales as Z only

Scanning on energies

- To use resonant production, you need to start with positrons AND get to the resonant energies

$$E_{res} = \frac{M_V^2}{2 m_e} \quad \rightarrow \quad \text{A “scanning” procedure is required}$$

- Analysis strategy: vary the beam energy, fit the background, calibrate the luminosity and look for resonance
 - The resonance shape is exactly the one of the beam energy distribution !

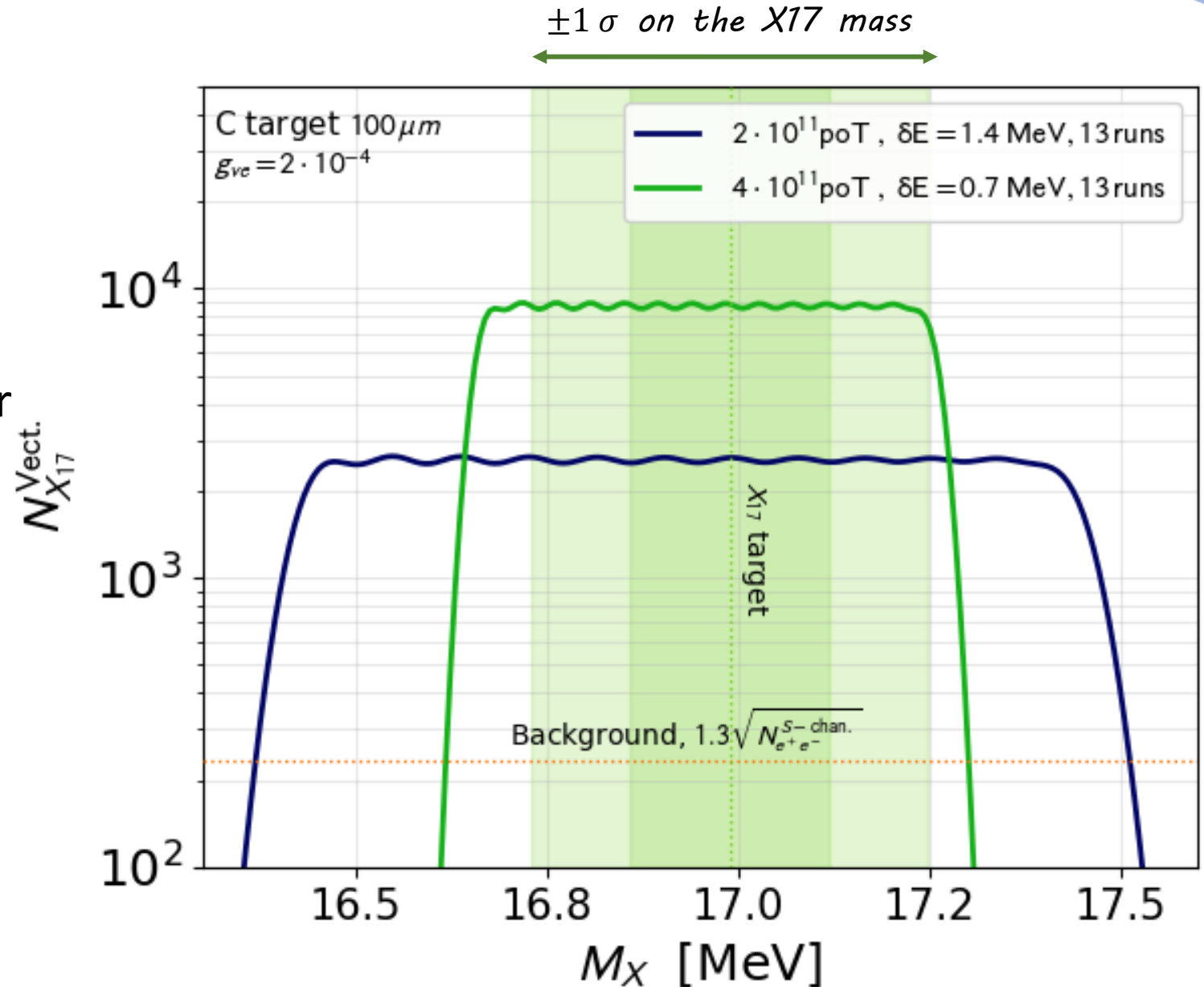
$$\mathcal{N}_{X_{17}}^{\text{per poT}}(E) = \frac{\mathcal{N}_A Z \rho}{A} \ell_{\text{tar}} \frac{g_{ve}^2}{2m_e} f(E_{res}, E)$$

With f the beam spread, typically modelised by a Gaussian distribution with spread δE

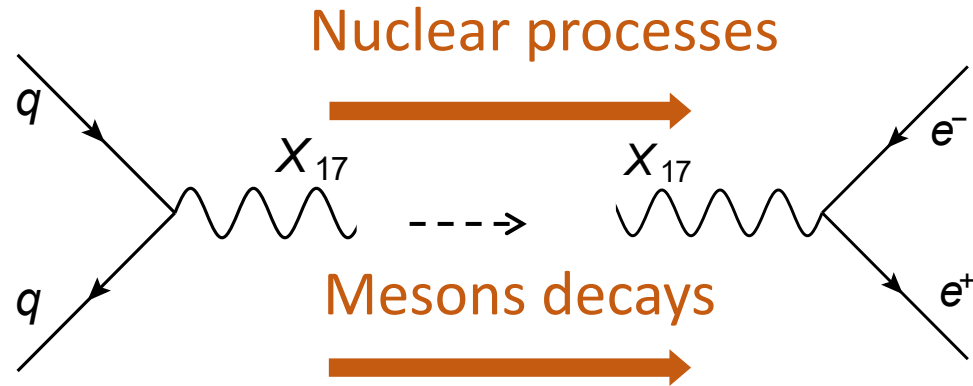
- Main background is from Bhabha scattering, fitted directly from the data
 - “Large angular acceptance” detector important to reduce the t-channel contribution

Production rates

- The beam energies used must be adapted depending on the beam spread δE
 - Smaller spread implies lower background as the signal a “bump” with spread δE
- In PADME, **better resolution on the beam energy steps than on e^+e^- invariant mass**
 - Each individual run at a given energy is used as a data point



PADME and the X17 boson, the perfect target



- e^+ / e^- beam dump and e^+ / e^- collider
- $e^{(*)} \rightarrow e X$ emission

Model independent + test NP origin of the signal

- We look for a light boson decaying to mostly to $e^+ e^-$ with mass:

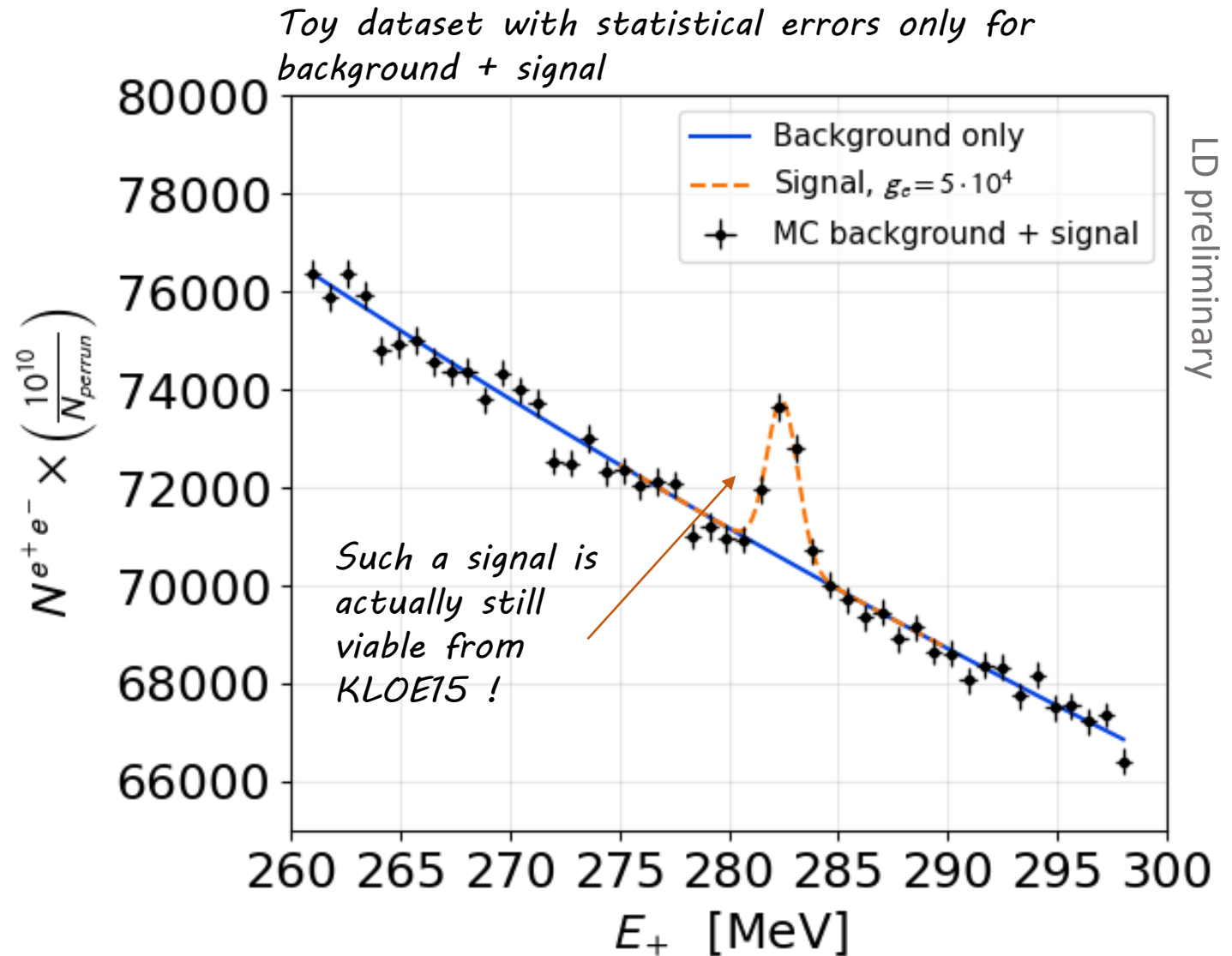
$$m_{X_{17}} \simeq 16.99 \pm 0.04 (stat) \pm 0.21(syst) \text{ MeV}$$

- The narrow mass range plus model-independent e^\pm couplings makes this anomaly a perfect target for a resonant search !
- The target energy range is [260 - 300] MeV \rightarrow perfectly adapted to e^+ beam at LNF

$$\mathcal{N}_{X_{17}}^{\text{per poT}} \simeq 3.8 \cdot 10^{-7} \times \left(\frac{g_e}{3 \cdot 10^{-4}} \right)^2 \left(\frac{1 \text{ MeV}}{\delta E} \right)$$

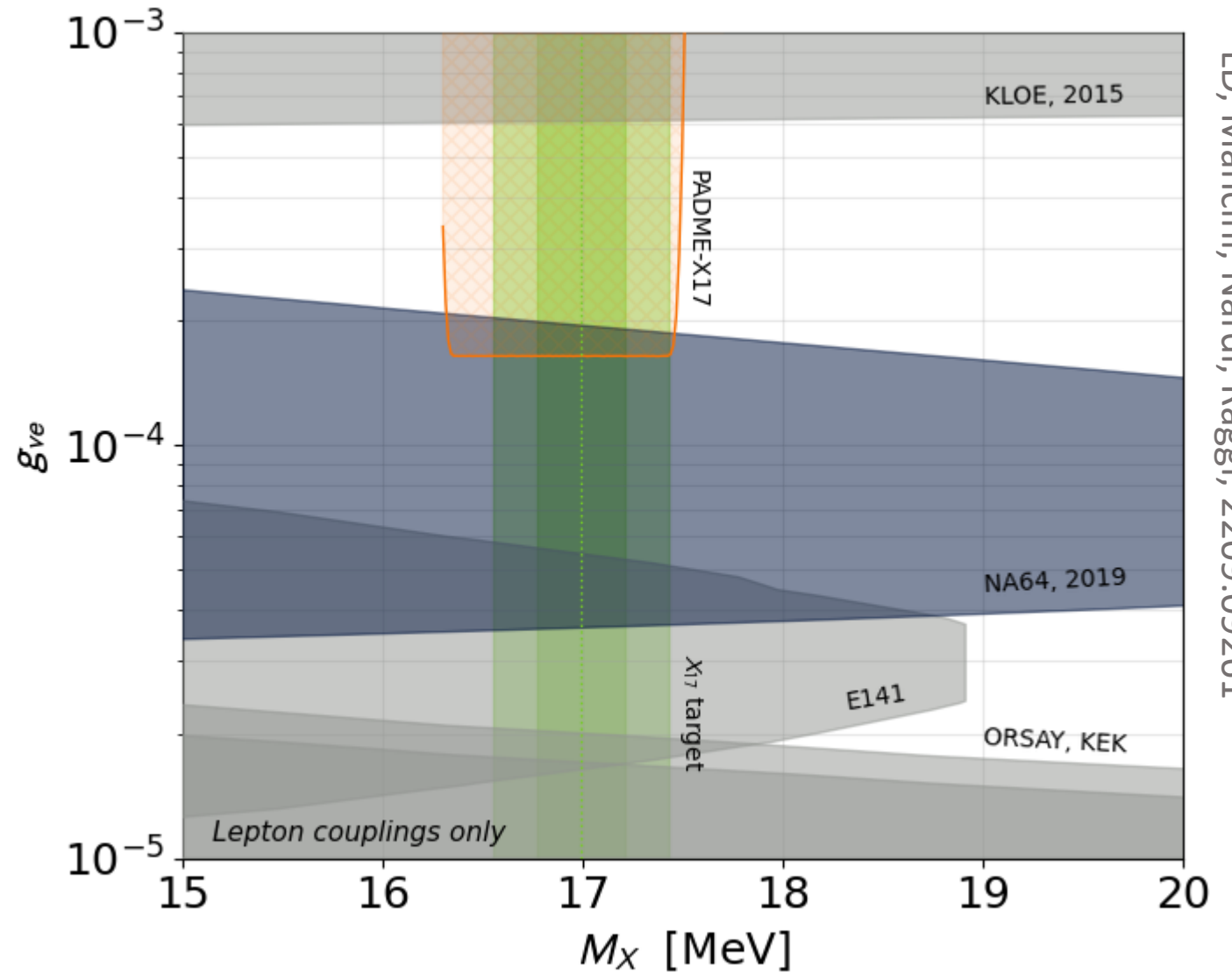
What would a perfect signal look like ...

- Take a (slightly) idealised Run-3 dataset
 - $4.9 \cdot 10^{11}$ PoT, a 0.25% beam spread
 - [261, 298] MeV range of energy with a 0.75 MeV energy step
- Assume the luminosity uncertainty has been controlled below the statistical error (i.e. $\sim 0.5\%$ level or so)
 - Normalise each point to 10^{10}
- The beam spread implies that the signal will turn up in several data-point
 - Full bump-search approach required



Projections for PADME – X17

- Complete simulation based on the current PADME setup
 - If all goes to plan the vector X17 hypothesis will be fully covered !)
 - Not enough for the scalar case due to reduced NA64 limit
- Serve as a first “test run” for these kind of analysis
 - Will also provide best constraints on dark photon in this (narrow) mass range, beating NA48



Conclusion

Conclusion

- Positron-based facilities allow to leverage resonant production to significantly increase signal rates
 - With the cost of having to scan over a large energy range
- PADME-Run III is (also) a proof-of-concept for a $e^-e^+ \rightarrow X, X \rightarrow e^+e^-$ experimental strategy for a feebly coupled new particle
- The X17 vector explanation (the most probable one given ^{12}C data) could be fully covered!
 - (provided experimental errors stay nicely where expected)

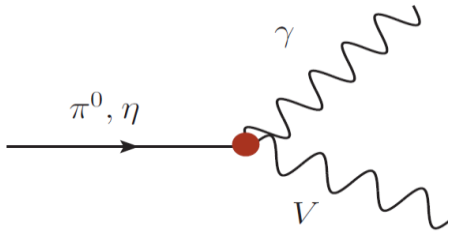
Backup

Accelerator facilities (currently) available

Intensity beam dumps

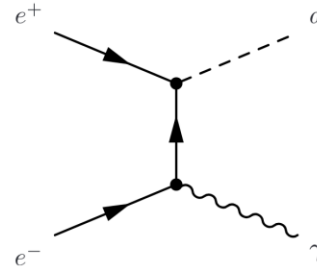
→ p machines (beam neutrinos exp, SHiP).

- Large backgrounds + far away detectors



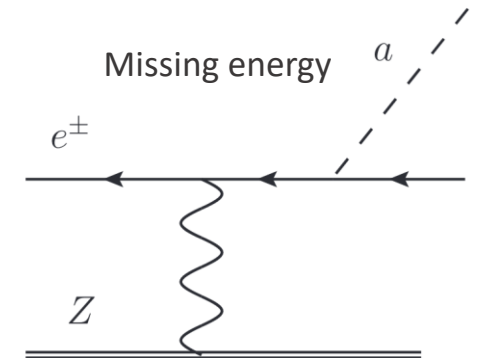
e^+e^- colliders

→ Higher energy but smaller luminosity



e^- “precision” fixed targets (NA64 ...)

→ Small luminosity but no background



- PADME Run-3 is also the first try at a new concept : e^+ “precision” beam dumps → Rely on resonant production for x10-100 improvement in production rates

Rare decays searches

- Rare decays probes are both extremely effective in probing X17, often at the price of a large model dependence

- Mesons decay probes (example from mostly last year)

- | | | | | |
|--------------|---|------------------------------------------------------------------------------------------------------------|-----------------------------------------|------------------------------------------------------|
| Vector state | { | $\pi^0 \rightarrow \gamma V_{17}$, for vector states: NA48 bounds implies proto-phobic | hep-ex/0610072 | Feng et al.
(1604.07411,1608.03591)
2006.01151 |
| | | J/Ψ decays, charm couplings only | Ban et al. 2012.04190 | |
| | | $B^* \rightarrow B V_{17}, D^* \rightarrow D V_{17}$ for vector states | Castro and Quintero 2101.01865 | |
| Axion | { | $\pi^0 \rightarrow a_{17} \rightarrow e^+ e^-, K \rightarrow \pi(\pi) a_{17}, K \rightarrow \mu\nu a_{17}$ | e.g Alves et al. 1710.03764, 2009.05578 | |
| | | $\pi^0 \rightarrow a_{17} a_{17} a_{17}$ and other multi-leptons final states | Hostert and Pospelov 2012.02142 | |

- If flavour-violation, many more available channels both in lepton decays and in “standard” flavoured meson decay.

- Also radiative emission from μ decay

ALP case

- The « tip » of the NA64 search does not cover the relevant range
- Not much is required, should be much better with new data

