

# Simulation of Tapered Co-propagating Structures for Dielectric Laser Accelerators

S. Quevedo Díaz<sup>1</sup>, A. Leiva Genre<sup>1</sup>, G. Torrisi<sup>1</sup>, D. Mascali<sup>1</sup>, G. S. Mauro<sup>1</sup>, G. Sorbello<sup>1,2</sup>, A. Bacci<sup>3</sup>, M. Rossetti Conti<sup>3</sup>, R. Palmeri<sup>4</sup>



<sup>1</sup>Istituto Nazionale di Fisica Nucleare – Laboratori Nazionali del Sud, Via S. Sofia 62, 95123 Catania, Italy.

<sup>2</sup>Dipartimento di Ingegneria Elettrica, Elettronica e Informatica (DIEEI), Università degli Studi di Catania, Via S. Sofia 64, 95125 Catania, Italy.

<sup>3</sup>Istituto Nazionale di Fisica Nucleare – Sezione di Milano, Via Celoria 16, 20133 Milano, Italy.

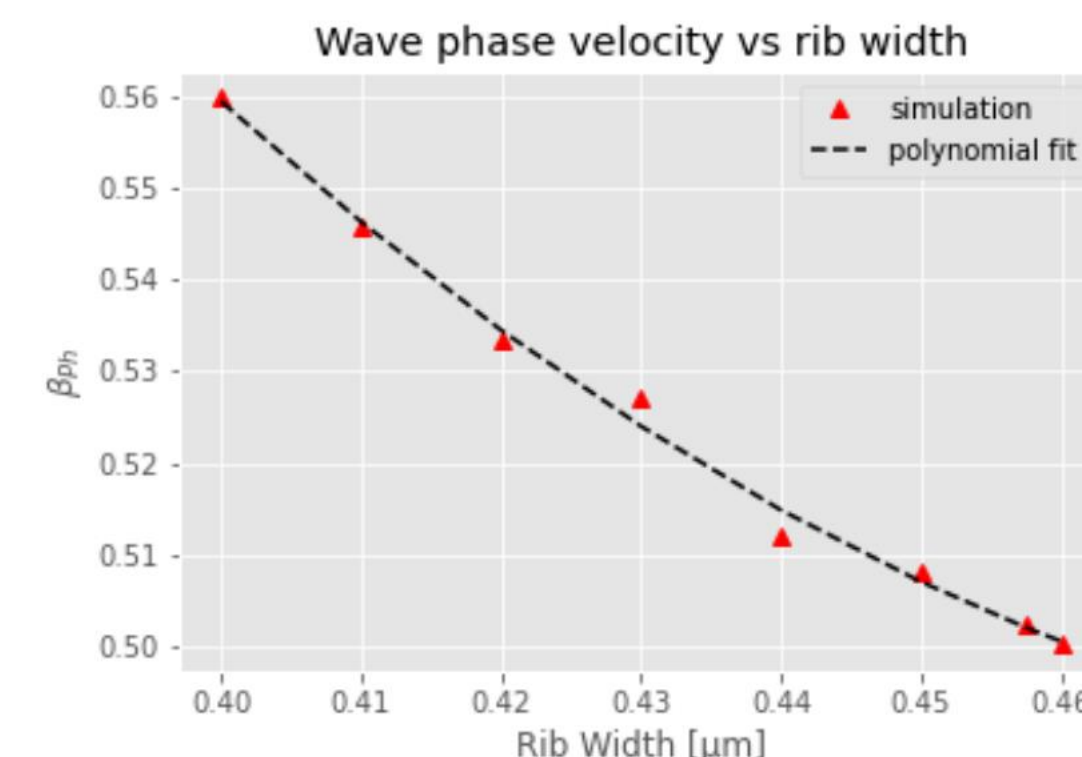
<sup>4</sup>Consiglio Nazionale delle Ricerche, Istituto per il rilevamento elettromagnetico dell'ambiente, Via Diocleziano 328 - 80124 Napoli, Italy

## Introduction

- **Enhancing Compactness and Efficiency:** By utilizing laser pulses and advanced dielectric materials, this emerging technology holds the promise of delivering particle acceleration systems that are not only more compact but also significantly more efficient and cost-effective than conventional accelerator technologies.
- **Exploring Innovative Approaches:** Within the realm of accelerator research, one particularly promising avenue is the investigation of hollow-core dielectric Electromagnetic Band Gap (EBG) microstructures powered by lasers. These structures represent an exciting and novel frontier in accelerator science.
- **Comparative Advantages:** In contrast to traditional metallic accelerators, dielectric accelerators offer several key advantages. They possess higher damage thresholds and can support greater accelerating gradients, particularly in the GV/m regime, outperforming their metallic counterparts.
- **Foundation in Photonic Crystals:** Dielectric Laser Accelerators are fundamentally grounded in the principles of Photonic Crystals (PhCs). PhCs are structures engineered to either permit or obstruct the propagation of electromagnetic waves within specific frequency bands [1].
- **Diverse Accelerating Structures:** The versatility of dielectric laser accelerators is reflected in the variety of structures that can be employed, including 2D or 3D accelerating waveguides based on different PhC configurations, as well as dielectric slot waveguides.
- **Numerical Insights:** Valuable insights into the performance of these accelerators have been gained through electromagnetic mapping and Particle-in-Cell (PIC) simulations, carried out using advanced software like CST Studio Suite. These simulations have enabled the extraction of several critical figures of merit.
- **Expansive Applications:** Beyond the realm of particle physics, dielectric laser accelerators have the potential to revolutionize a multitude of fields, ranging from medicine and industry to scientific research and space exploration. As research in this exciting field continues to advance, we can anticipate the emergence of even more compelling applications.

## Tapering: Phase velocity vs Rib width

- Phase velocity depends on rib width. The corresponding pair values were calculated by performing a parametric sweep in time domains simulations [2].
- Electrons traveling along the accelerator increase their velocity as they gain energy.



Phase velocity corresponding to different values of the rib width. It is calculated by fixing all the structural fields and extracting their phase velocity for each value of  $a$ .

- Particle phase dictates the interaction between the electron and the EM field.
- Synchronous condition is given by  $v_{ph} = v_e$ .
- With adequate tapering of the slot waveguide, phase velocity can match the electron velocity, providing maximum energy gain during the acceleration stage as a result

$$E_z(z, t) = E_0 \cos(\varphi)$$

where  $\varphi$  is given by:

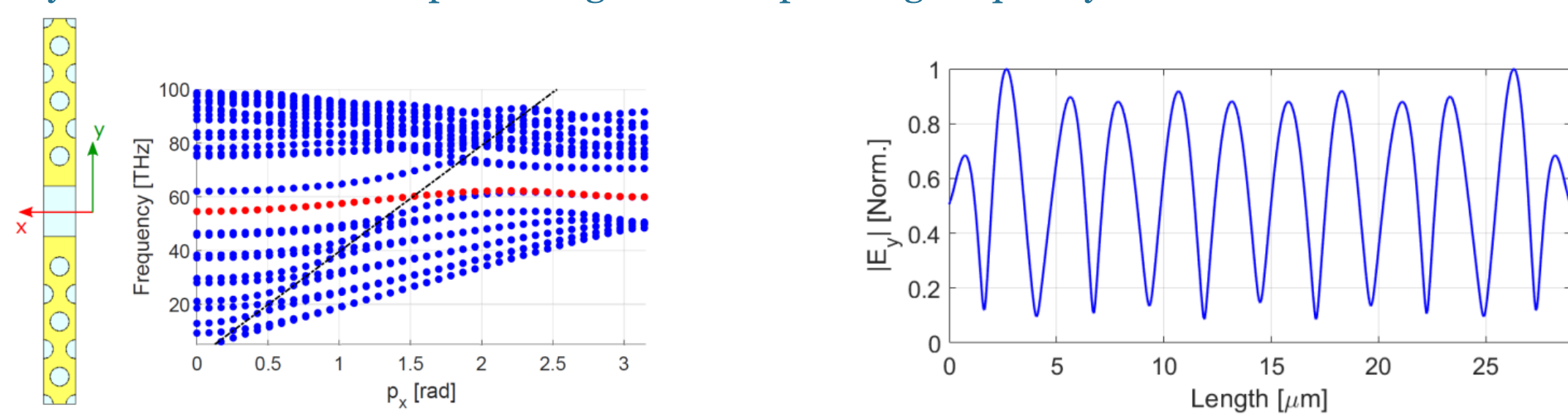
$$\varphi = \omega \left( t_0 + \int_0^z \frac{dz'}{v_e(z')} - \int_0^z \frac{dz'}{v_{ph}(z')} \right) + \varphi_0$$

The electric field seen by the electron is given by a sinusoidal dependence. The phase plays a fundamental role in the acceleration.

## Dielectric waveguides

### 2D Accelerating waveguides based on triangular lattice PhC ( $\beta=1$ )

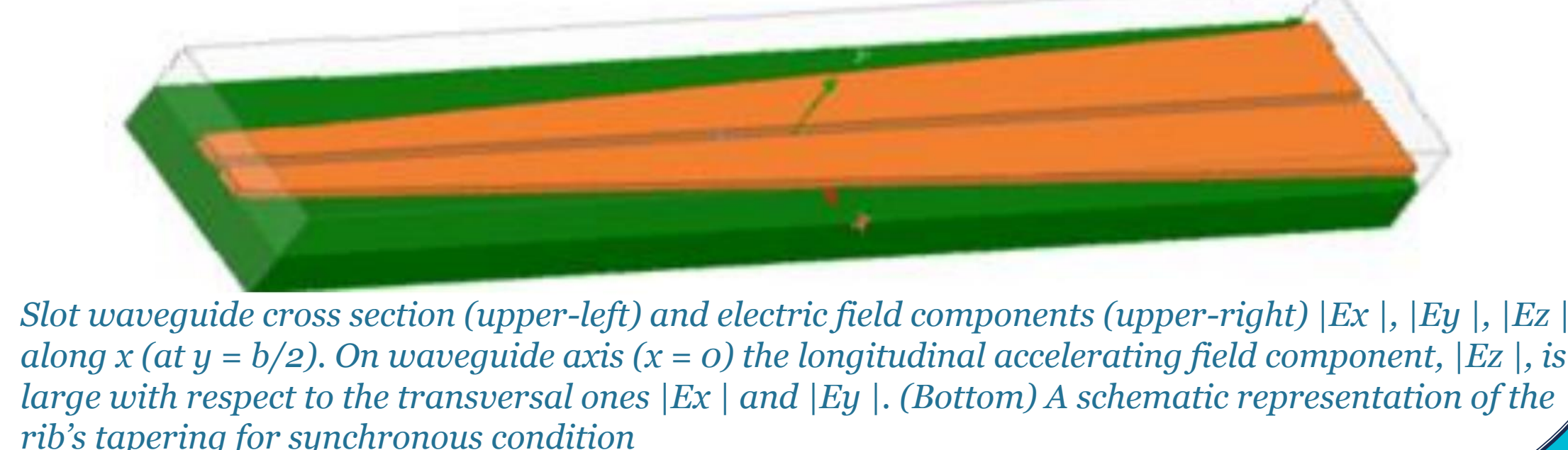
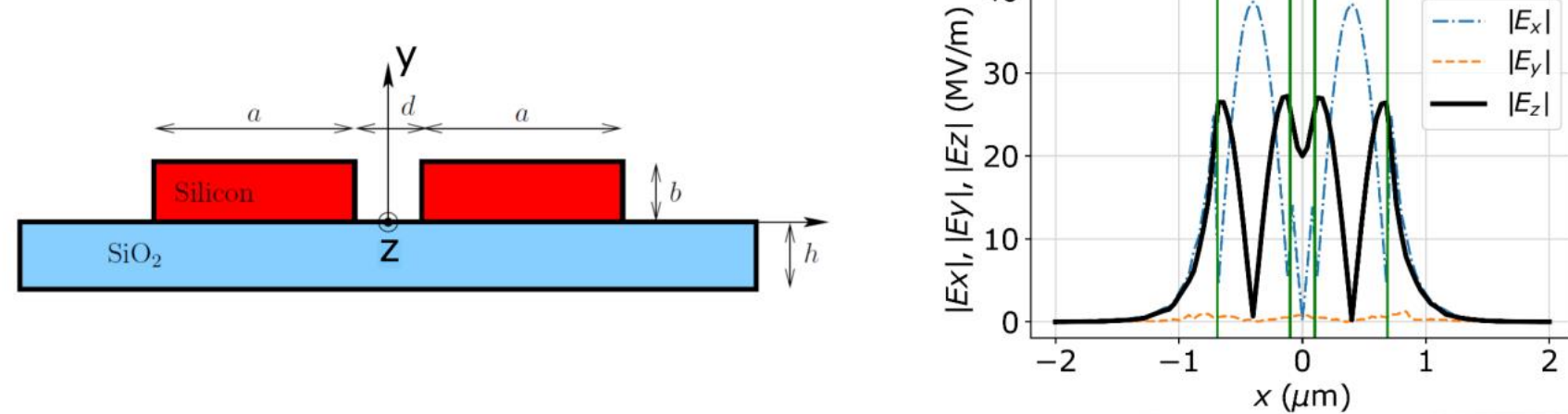
- Based on a 2D triangular lattice PhC, composed of periodically arranged vacuum holes with radius  $r = 0.3d$ , where  $d$  is the lattice constant (distance between two adjacent vacuum holes center) equal to  $1.207 \mu\text{m}$ , practiced on a slab ( $\epsilon_r = 9.7$ ). Dimensions depends on operating wavelength.
- The central vacuum channel, called “defect”, has been tuned to support an accelerating mode, synchronous with the speed of light at the operating frequency of 60 THz.



(left) 2D triangular lattice waveguide supercell and projected band diagram along the accelerating  $x$  axis. The  $TM_{01}$ -like mode (red curve) is synchronous with the speed of light at  $f_0 = 60$  THz. (right) Accelerating  $E_y$  field component along waveguide axis, for a structure of  $29 \mu\text{m}$  total length.

### Slot waveguides ( $0.4 < \beta < 0.7$ )

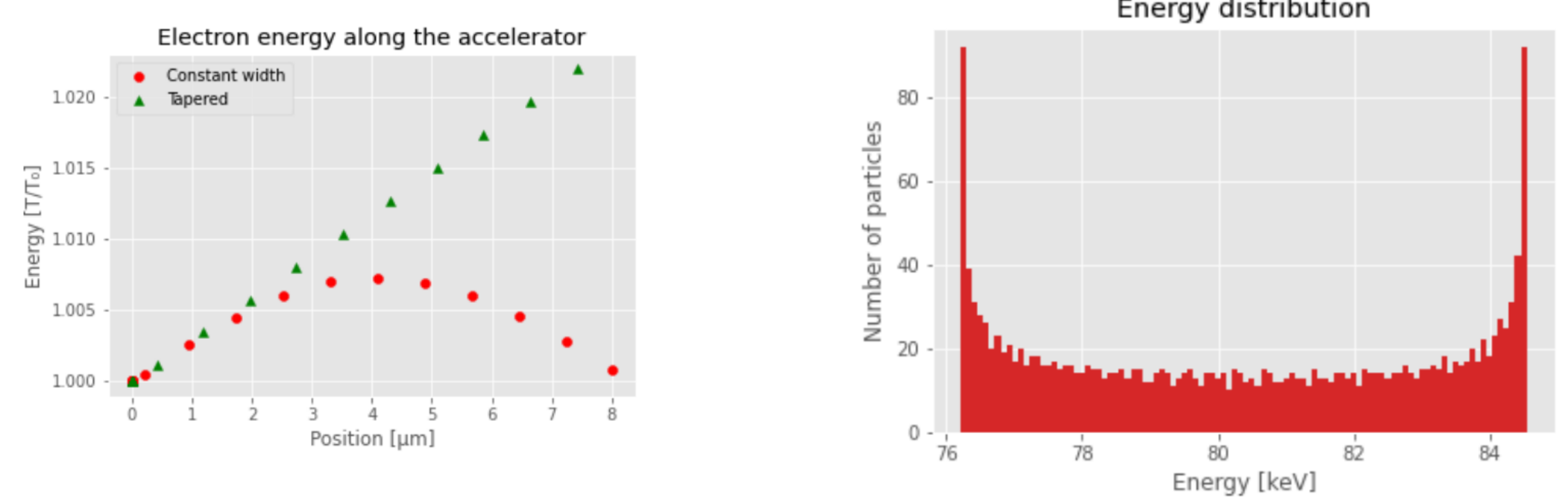
- The slot waveguide, is a dielectric structure employed for sub-relativistic electron acceleration.
- It supports a longitudinal accelerating field whose phase velocity can be tuned to maintain the synchronism with the accelerated electrons as they gain velocity [2].
- It consists of two silicon tapered “rib” waveguides of width  $a = (0.4575-0.44345) \mu\text{m}$  and height  $b = 0.172 \mu\text{m}$  supported by a  $\text{SiO}_2$  substrate, separated by a hollow-core channel of width  $d = 0.155 \mu\text{m}$ . Accelerator length is  $L = 20 \mu\text{m}$ .
- The accelerating mode is excited by a  $1.55 \mu\text{m}$  laser with input power  $P_{inj} = 250$  W.



Slot waveguide cross section (upper-left) and electric field components (upper-right)  $|E_x|$ ,  $|E_y|$ ,  $|E_z|$  along  $x$  (at  $y = b/2$ ). On waveguide axis ( $x = 0$ ) the longitudinal accelerating field component,  $|E_z|$ , is large with respect to the transversal ones  $|E_x|$  and  $|E_y|$ . (Bottom) A schematic representation of the rib's tapering for synchronous condition

## Beam dynamics

- An electron source with initial energy  $T_0 = 80$  keV is placed at the beginning of the DLA
- **Performance comparison:** Tapered vs non-tapered slot waveguide. Tapered delivers continuous acceleration. Its counterpart fails at this task due to its constant phase velocity.
- In a DC source, the particles exiting energy distribution varies according to the initial phase. Particles close to the resonant phase increase their energy up to 5 keV. Dissimilar phases decelerates the electrons, losing some part of their kinetic energy in the process.



Left: Electron energy according to the position inside the accelerator for a tapered and non-tapered (constant width) slot waveguide. Right: Energy distribution at the exit of the accelerator for a continuous monochromatic source.

## Figures of Merit

- The acceleration gradient is around  $227$  MV/m, which is above current conventional accelerators
- The total energy gain is  $4.5$  keV in  $20 \mu\text{m}$ .
- $\Delta W = -eE_0 L \cos(\varphi)$

## Conclusions and Perspectives

- Two different DLA structures has been presented for electron acceleration for both sub-relativistic and relativistic regime.
- Tapered slot waveguide proves to be a potential tool for sub-relativistic electron acceleration. By adapting its structural parameters, a wide range of electron energies can be covered
- Phase acceptance and two-stages acceleration schemes are next step

## References

1. G. Torrisi et al., “Numerical study of photonic-crystal-based dielectric accelerators,” in 10th Int. Particle Acc. Conf. (IPAC'19). JACOW Publishing.
2. Z. Zhao, T. Hughes, S. Tan, H. Deng, N. Saprà, R. England, J. Vuckovic, J. Harris, R. Byer, and S. Fan, “Design of a tapered slot waveguide dielectric laser accelerator for sub-relativistic electrons,” Opt. Express 26, 22801-22815 (2018)