Probing Ultralight Dark Matter with Binary Pulsars

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Binary pulsars

2 Pulsar timing

3 Ultralight dark matter and its effect on binary systems

4 Results



1 Binary pulsars

- 2 Pulsar timing
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4 Results



Binary pulsar

- pulsar = "pulse" + "star"
 - = highly magnetised rotating neutron star which emits pulses of light, represents a very stable clock





(a) A pulsar

(b) A binary pulsar (it has 1 pulsar only)

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Lighthouse effect



Figure 2: Lighthouse effect - idea

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Keplerian formalism

• 6 orbital parameters, e.g.: $a, e, i, \Omega, \omega, t_0$



Figure 3: Orbital parameters

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Perturbation of Keplerian orbits



Figure 4: Perturbed Keplerian orbit

 \rightarrow Post-Keplerian formalism

$$\dot{a} = \frac{2}{\omega_b} \left(\frac{F_r e}{\sqrt{1 - e^2}} \sin \theta + \frac{F_\theta}{r} a \sqrt{1 - e^2} \right)$$

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• basic quantities:

N ... phase of the pulsar when emitted a pulse t ... pulse's time of arrival N_0 ... reference phase t_0 ... reference time

• "ideal"/trivial case:

$$N = N_0 + \nu(t - t_0)$$

• to describe reality is more difficult, N is a complicated function

Timing model



Figure 5: What makes N complicated in terms of t

Spoiler alert: We want to add one more effect to the picture, the interaction of ULDM with matter (components of the binary system)!

• once we have a model for N and first parameter estimates, we can calculate the time residuals δt :

$$-\nu^{(1)}\delta t = \frac{\partial N}{\partial N_0}|_{(1)}\delta N_0 + \dots,$$

where δN_0 is the difference between the actual and estimated value, etc.

- contributions to δt :
 - 1. inaccuracy in parameter determination
 - 2. oversimplified model
 - 3. noise

Timing model



Figure 6: Example of time residuals calculated with PINT

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(Ultralight) dark matter

- bosonic particles with $m \sim 10^{-22} \div 1 \text{ eV}$
- huge phase space occupanion number
 - \rightarrow wave nature, creates interference patterns, locally described by a classical field (scalar, vector, tensor) $\Phi = \Phi_0 r \cos(m_{\Phi} t + B)$



Figure 7: Dark matter models - range of masses

ULDM-binary system interaction

- if $m \sim 10^{-22} \div 10^{-18}$ eV, negligable backreaction
- for simplificity, scalar field and universal interaction with

$$M_A(\Phi) = \bar{M}_A(1 + \lambda \Phi),$$

• perturbation of binary dynamics by ULDM:

$$\ddot{\vec{r}} = -(1+\lambda\Phi)\frac{G\bar{M}_T\vec{r}}{r^3} - \lambda\dot{\Phi}\,\dot{\vec{r}} = -\frac{G\bar{M}_T\vec{r}}{r^3} + \vec{F},$$
$$\rightarrow \frac{\dot{a}}{a} = -\frac{2\lambda\Phi e\omega_b}{\sqrt{1-e^2}}\frac{a^2}{r^2}\sin\theta - \frac{2\lambda\dot{\Phi}(1+e^2+2e\cos\theta)}{(1-e^2)}$$

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Constraining ULDM with NANOGrav data

- 1. constructing a probability distribution function for using Bayes' theorem
- 2. special assumption: δt comes from UDLM and noise (Gaussian, white) only:

$$\delta t = \delta t^{\rm DM} + \delta t^{\rm PN},$$

3. the present results cannot rule out/confirm the ULDM model, but place a stronger constraint on the value of the coupling constant α

$$\sigma_{\alpha}^{a} = 2.55 \times 10^{11} \frac{\sigma_{a}}{x(t_{0,a})} \frac{m_{\Phi}/\text{eV}}{\sqrt{\rho_{a}/(\text{GeV/cm}^{3})}} \frac{1}{\sqrt{\vec{F}_{a}^{2} + \vec{E}_{a}^{2}}},$$

where \tilde{F}_a^2 and \tilde{E}_a^2 some complicated functions, σ is the noise level and $x \equiv a_{\text{pulsar}} \sin i$.

Low-eccentric binary systems



Figure 8: Blue line - constraint

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High-eccentric binary systems



Figure 9: Blue line - constraint

Binary pulsars

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4 Results



- We probe the particular model of ULDM, which includes interaction with ordinary matter via effective coupling constant λ , with binary pulsars.
- We achieve competing or even stronger constraints than those found in the literature.
- It has not been shown here, but we have developed a procedure that can account for multiple contributions to δt , not just ULDM and noise.
- We will extend the work to vector and tensor ULDM.

Thank you for your attention!

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