

COST «Cosmic Wispers» Training School

Lecce, September 14th, 2023

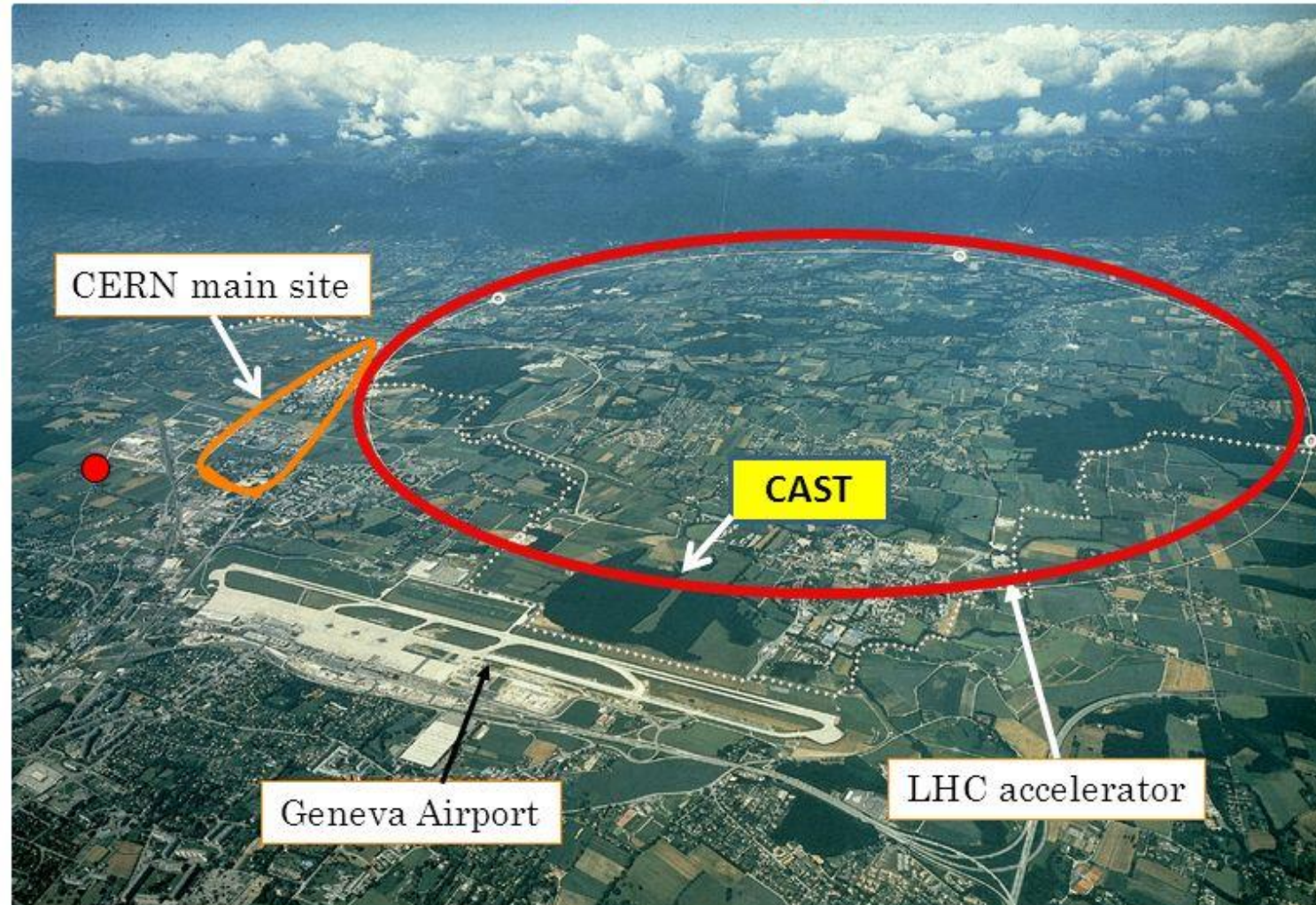


Antonios Gardikiotis
INFN - LNL

OUTLINE

- Axion helioscope : CAST
- Dielectric haloscope: MADMAX
- Cavity haloscope: QUAX

CAST 2002-2022



CAST (Cern Axion Solar Telescope) – 2002 -2022

CERN LIBRARIES, GENEVA
SC00001091

CERN- SPSC -99-21
CERN 99-21
SPSC/P312
August 9, 1999
Sw 199948

Proposal to the SPSC

A solar axion search using a
decommissioned LHC test magnet

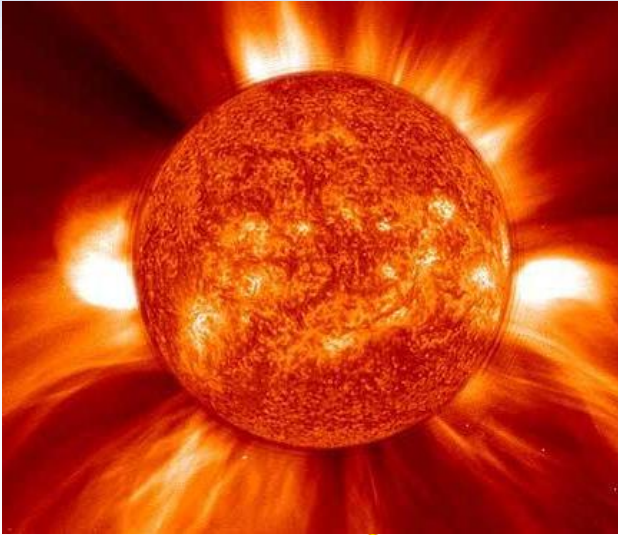


The Solar Axion Telescopic ANTenna

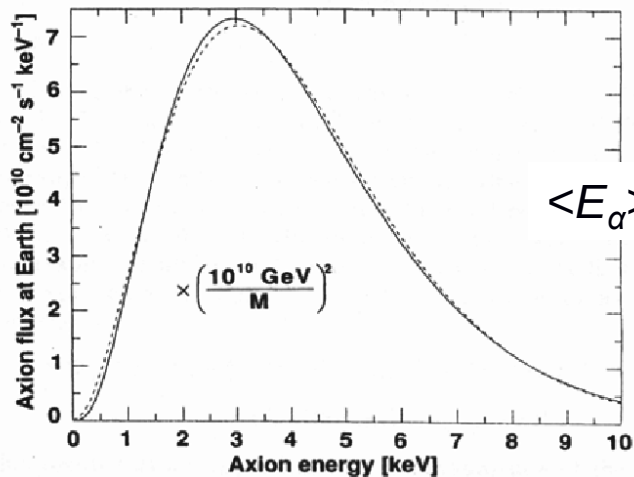


Axion Helioscope Experiment

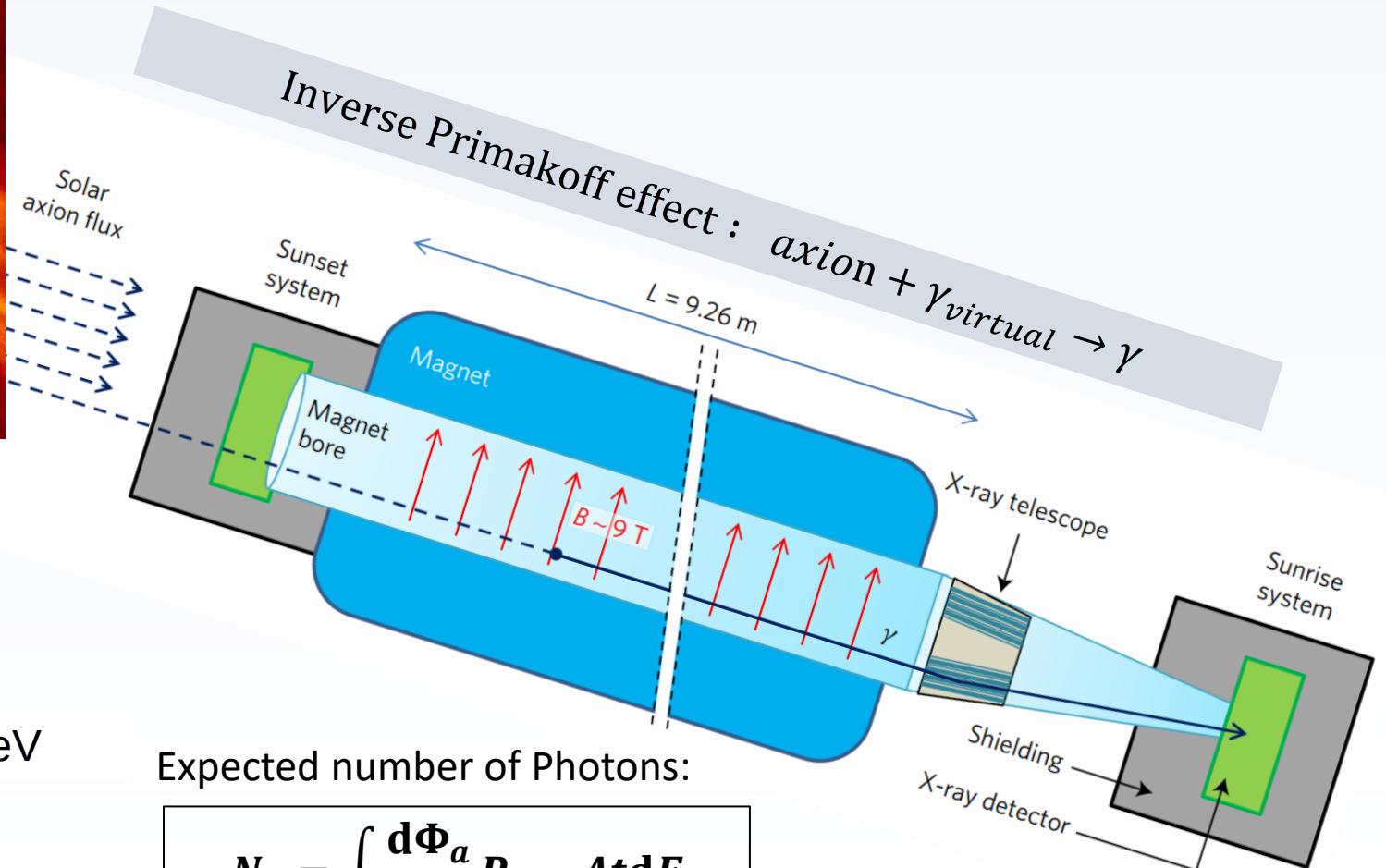
Axion source :



Standard Solar Model



$\langle E_\alpha \rangle = 4.2 \text{ keV}$

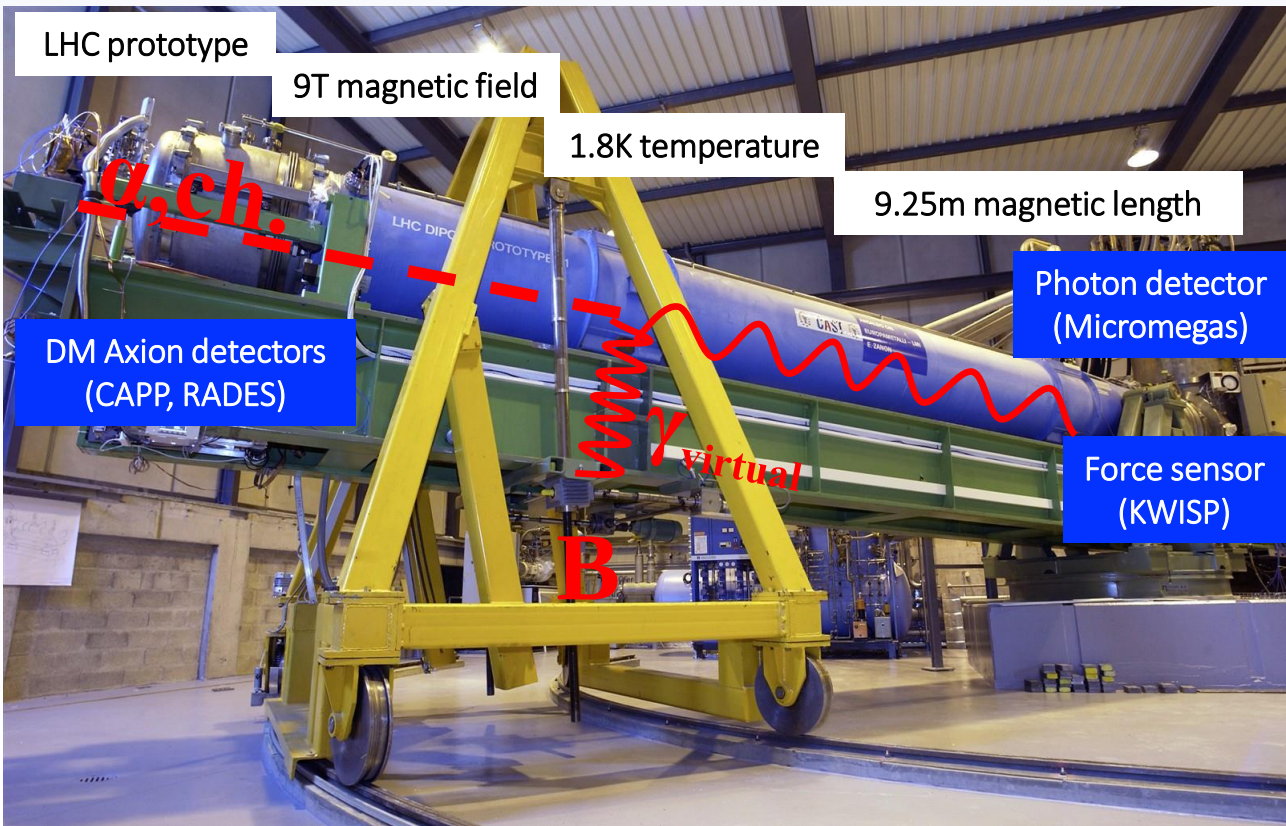


Expected number of Photons:

$$N_\gamma = \int \frac{d\Phi_a}{dE} P_{\alpha \rightarrow \gamma} A t dE$$

$$\frac{d\Phi_\alpha}{dE} = g_{10}^2 6.0 \times 10^{10} \text{ cm}^{-2} \text{ s}^{-1} \text{ keV}^{-1} E^{2.481} e^{-E/1.205}$$

CAST Axion Helioscope



Rotating platform

- (Vertical: $\pm 80^\circ$, Horizontal: $\pm 40^\circ$)
- 2x90 min solar tracking/day
- Sunrise: X-ray Focusing Device coupled to a CCD + 1 Micromegas
- Sunset: 2 Micromegas

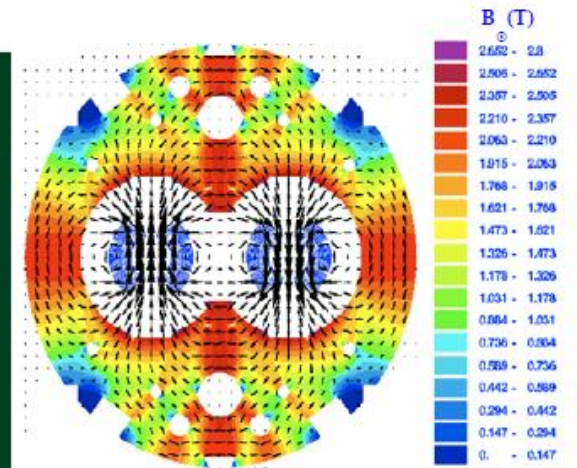
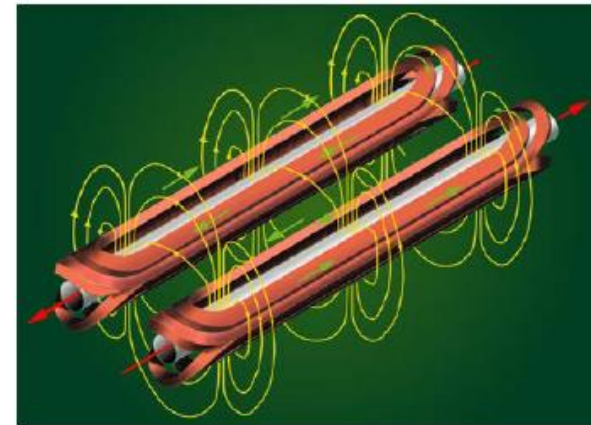
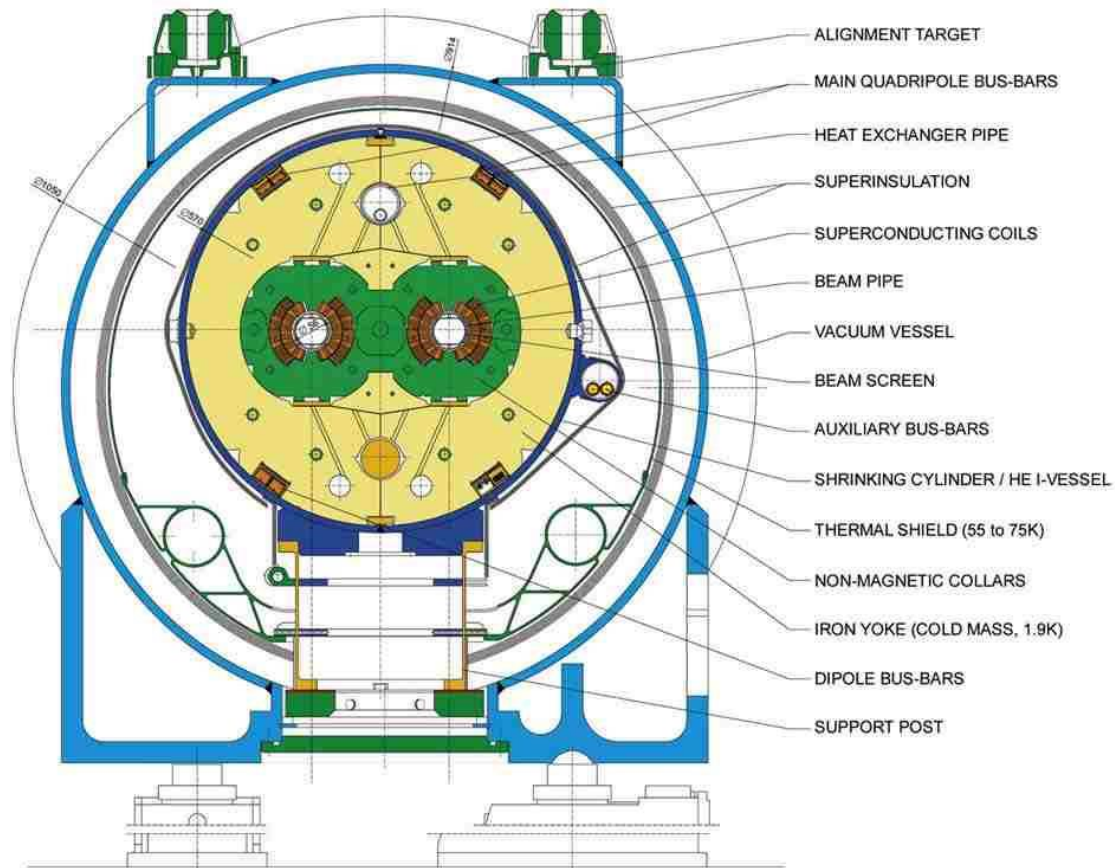
CAST Axion Helioscope



Decommissioned LHC (Large Hadron Collider) superconductive dipole magnet

LHC DIPOLE : STANDARD CROSS-SECTION

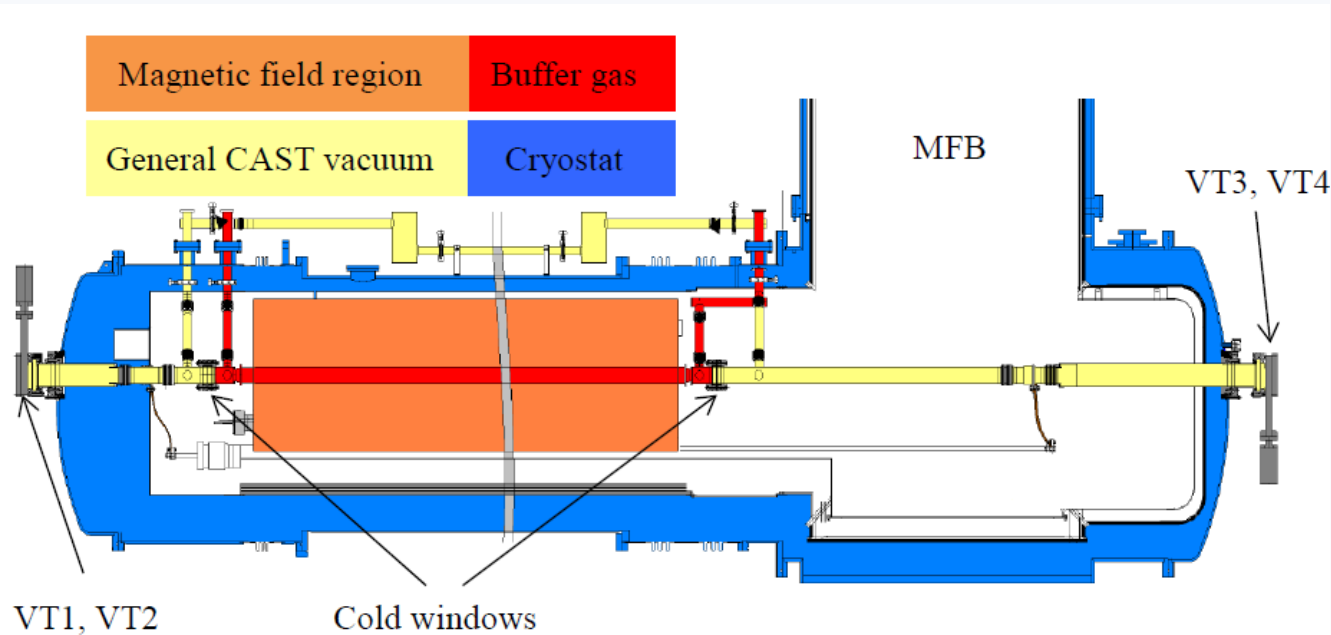
CERN AC/OL/MM - HE107 - 30 04 1999



spare LHC magnet

- 45 mm bore diameter
- cross section area $A = 14.5 \text{ cm}^2$
- ~10 m long
- $B \sim 9 \text{ T}$

CAST Cryogenics

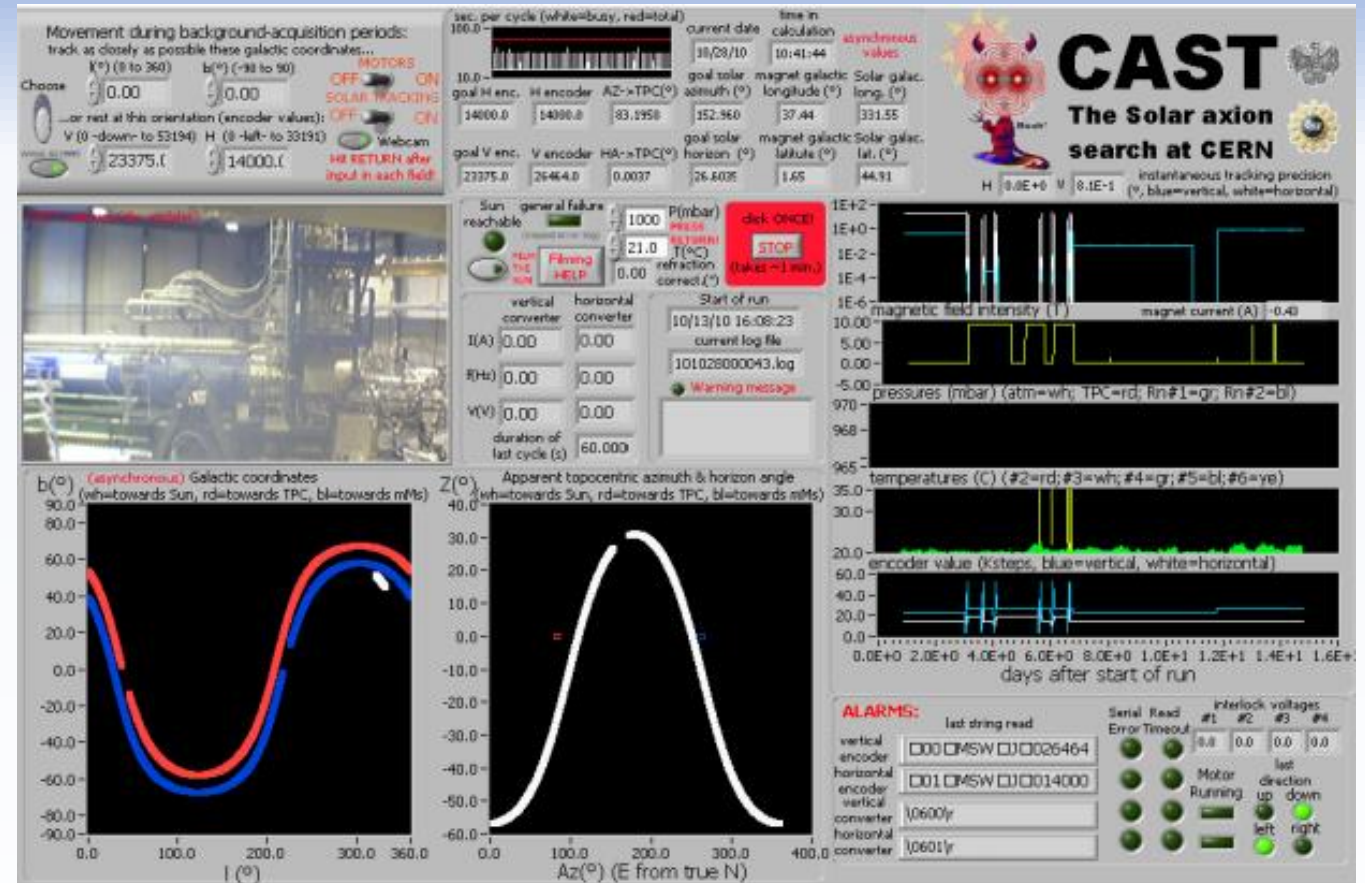
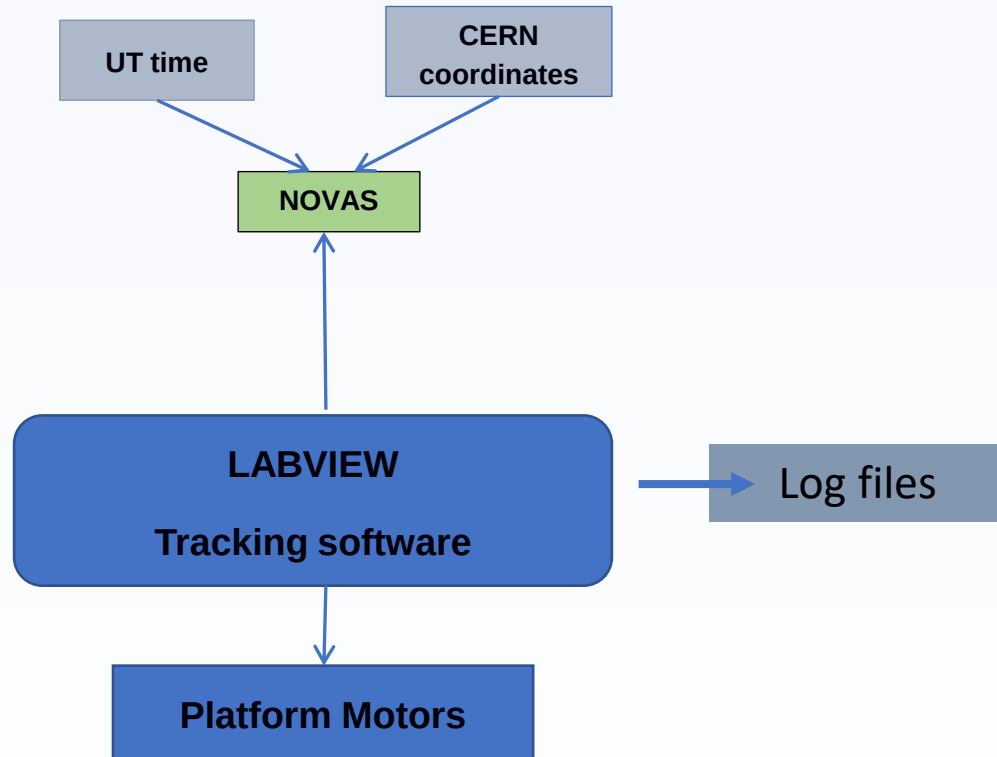


The **interlocks** system (PLC) automatically protects the vacuum setup, the buffer gas, the detectors vacuum side and the cryostat whenever a problem occurs

Quench : a fast discharge of the current is triggered



Solar Tracking



magnet positions to encoder values
 1 vertical encoder unit = 30 μ m of magnet movement

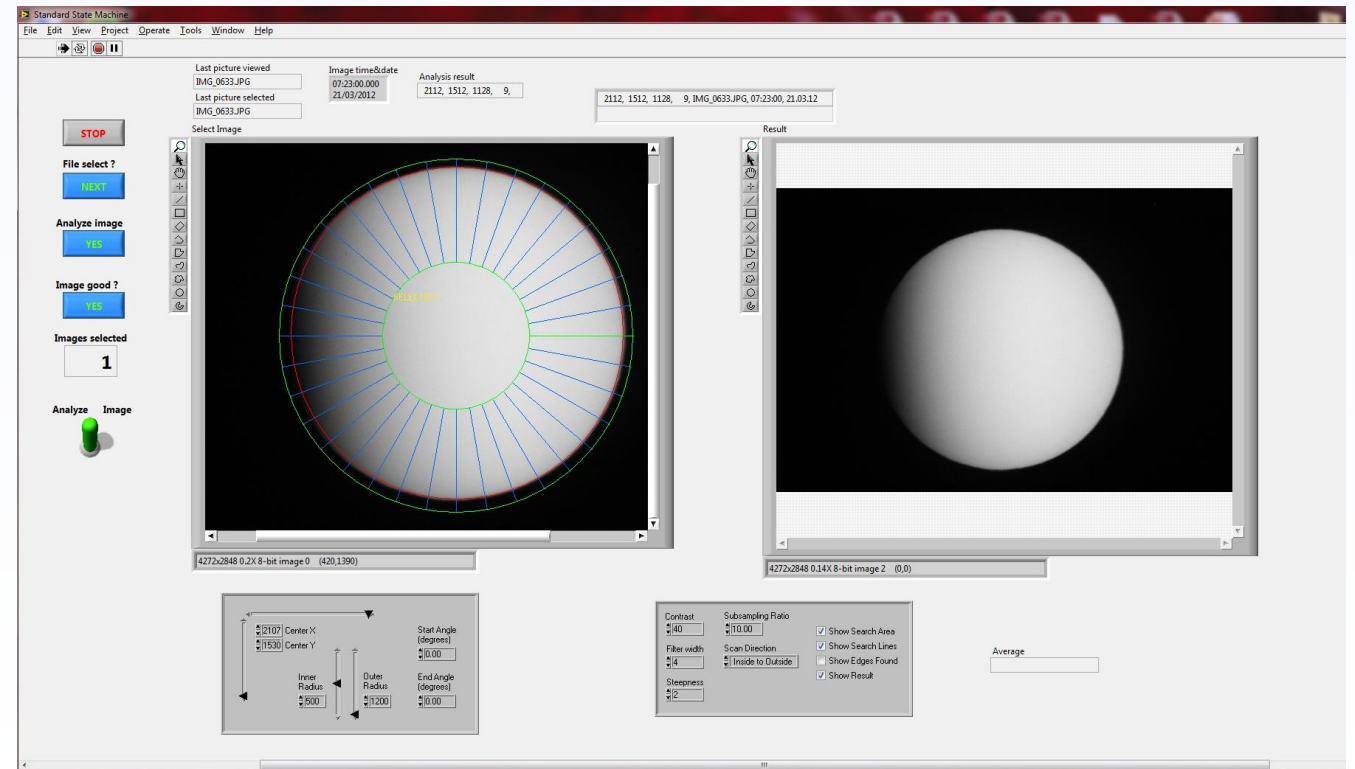
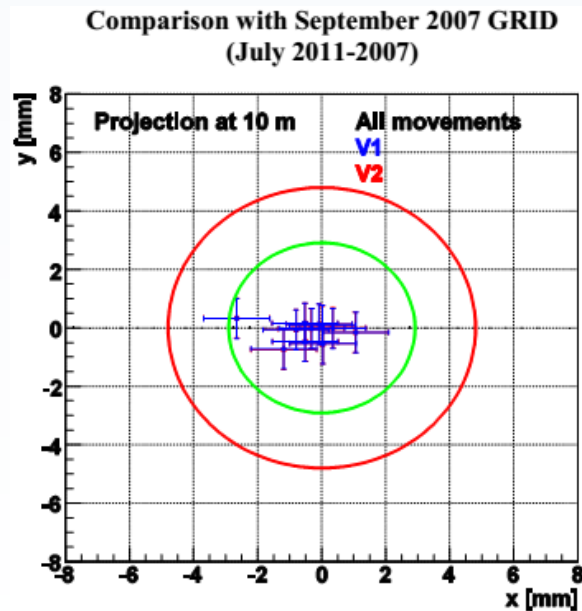


Grid measurements and Sun filming

- independent magnet positions
in a set of reference coordinates (GRID)
- green circle: required precision of 1 arcmin



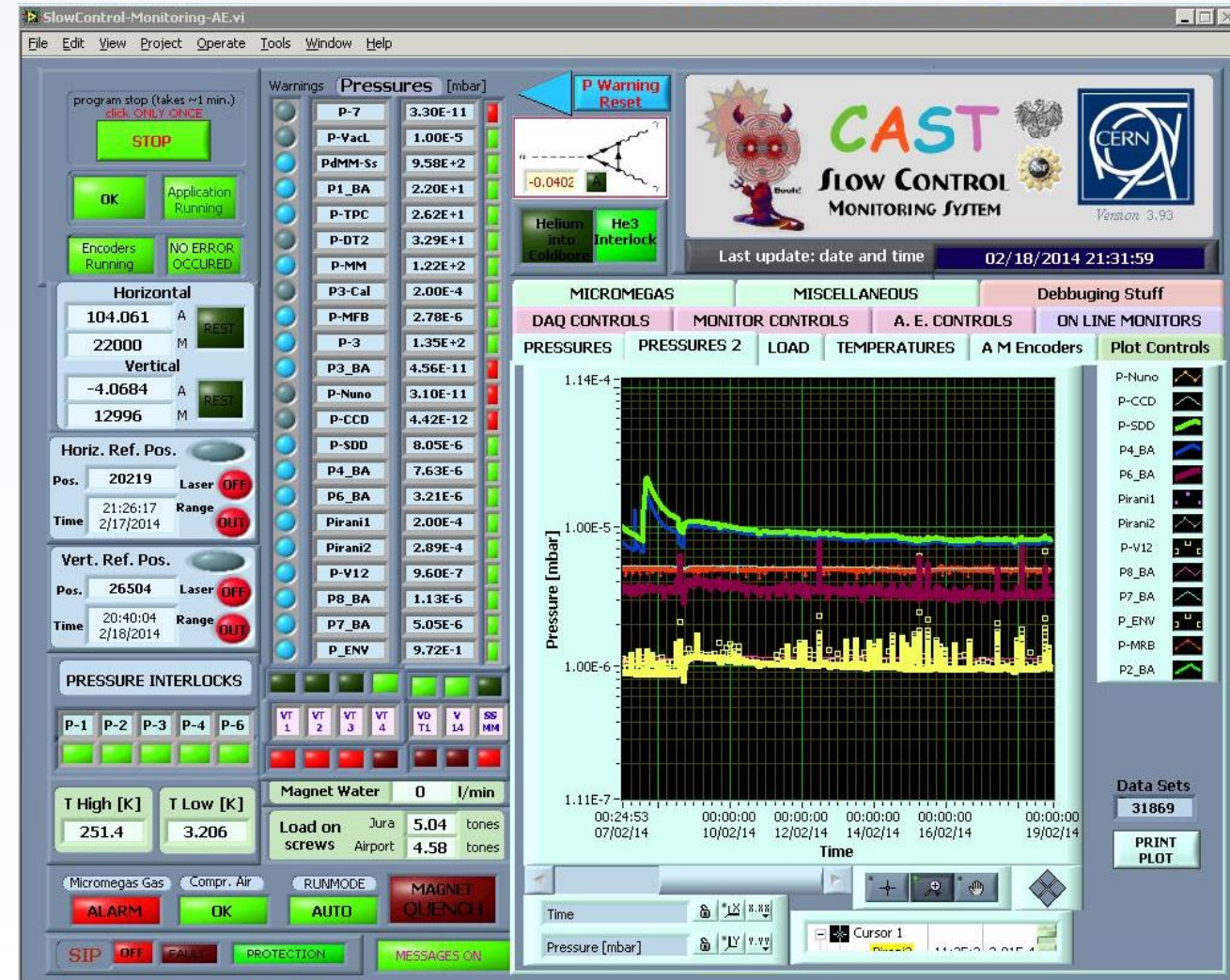
cross-check the tracking system precision: the Sun-filming



total estimated CAST tracking precision was better than 0.01°

Slow control of the CAST experiment

- Vacuum and cryostat pressure
- Pressures in the cold bores – detectors
- Various temperatures in different parts of the magnet
- Magnet movement parameters
- Magnet valves status (open or closed)
- Safety of the system and alarms



Typical values recorded every minute

Buffer Gas into the CAST coldbores

conversion probability for a uniform optical medium

$$P_{\alpha \rightarrow \gamma} = \left(\frac{g_{\alpha\gamma} B}{2} \right)^2 \frac{1}{q^2 + \Gamma^2/4} \left[1 + e^{-\Gamma L} - 2e^{-\Gamma L/2} \cos qL \right]$$

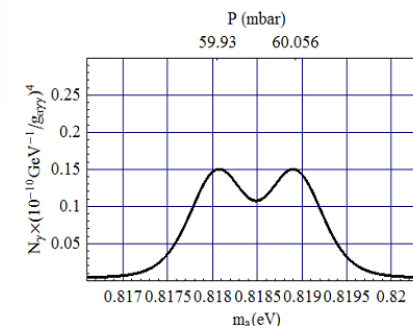
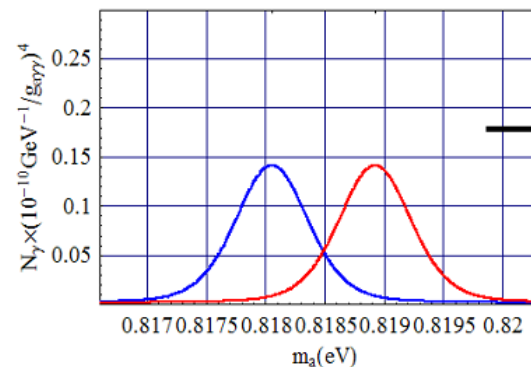
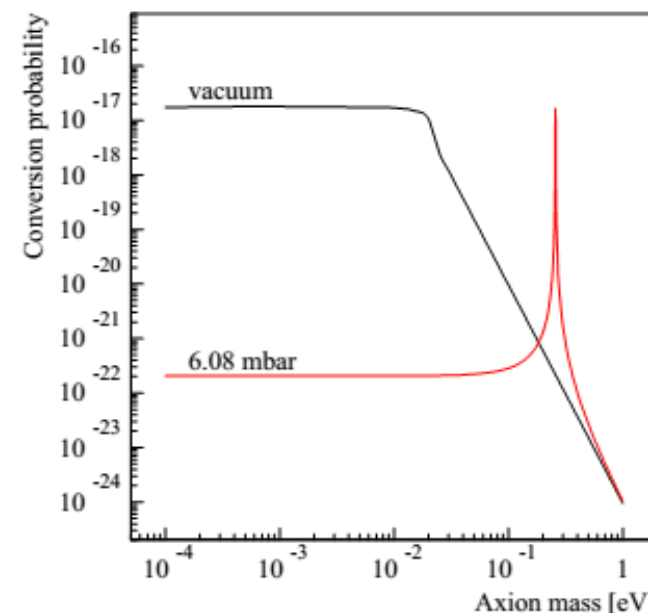
L = magnet length, Γ = absorption coefficient

$$q = \left| \frac{m_\gamma^2 - m_\alpha^2}{2E_\alpha} \right|$$

In CAST phase I (vacuum), coherence was lost for $m_a > 0.02$ eV

By using a buffer gas such as helium, the photon velocity can be decreased to match the incoming axion velocity

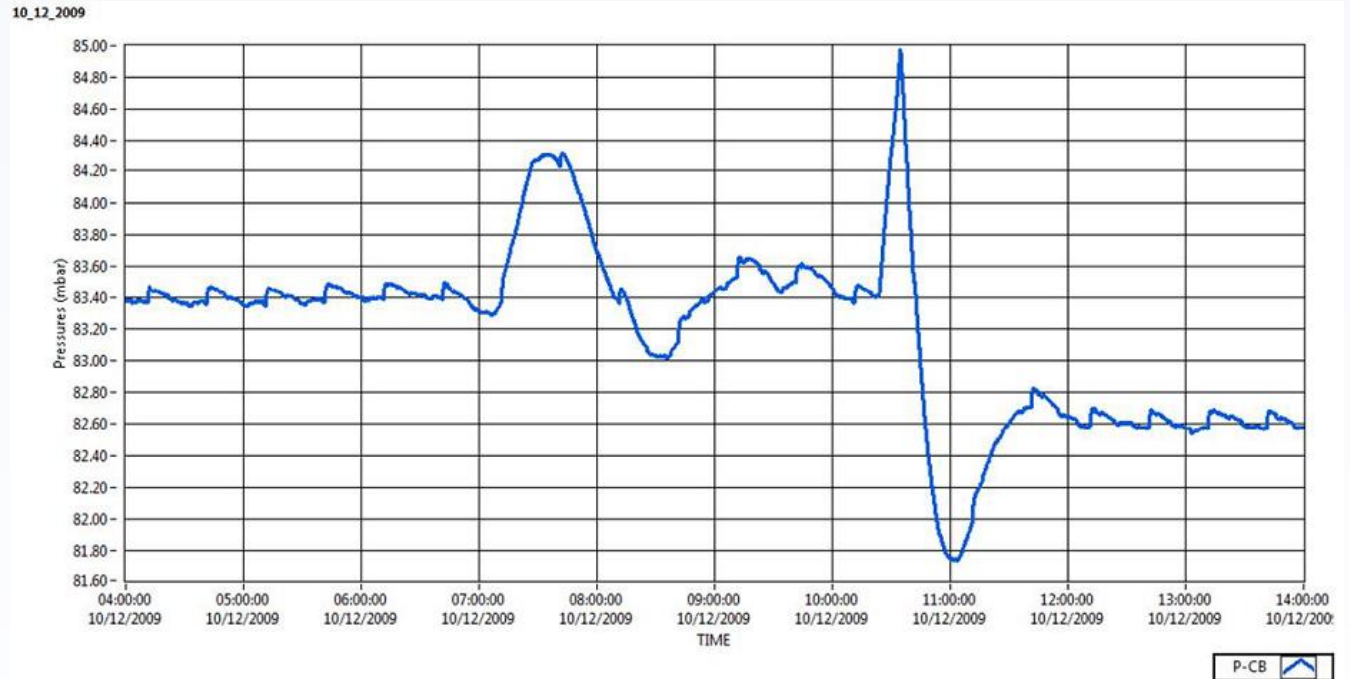
conversion probability inside a 10 m long pipe of magnetic field at 9T



Density profile in the CB

The effective photon mass from Primakoff conversion is related directly to the square root of the gas density

$$\frac{m_\gamma}{c^2} \cong 23.4267 \sqrt{\rho \left[\frac{g}{cm^3} \right] \left[\frac{eV}{c^2} \right]}$$



A vertical movement of the magnet of +6 degrees at 37mbars causes a shift in P_{cb} of about +1.05 mbar

The actual **density profile** which is needed to directly calculate the coherence length cannot be measured directly

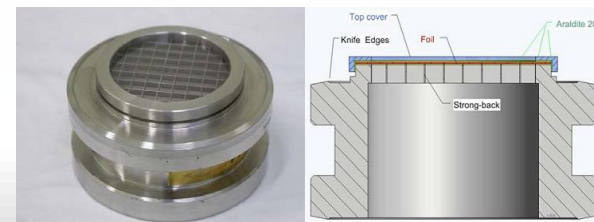
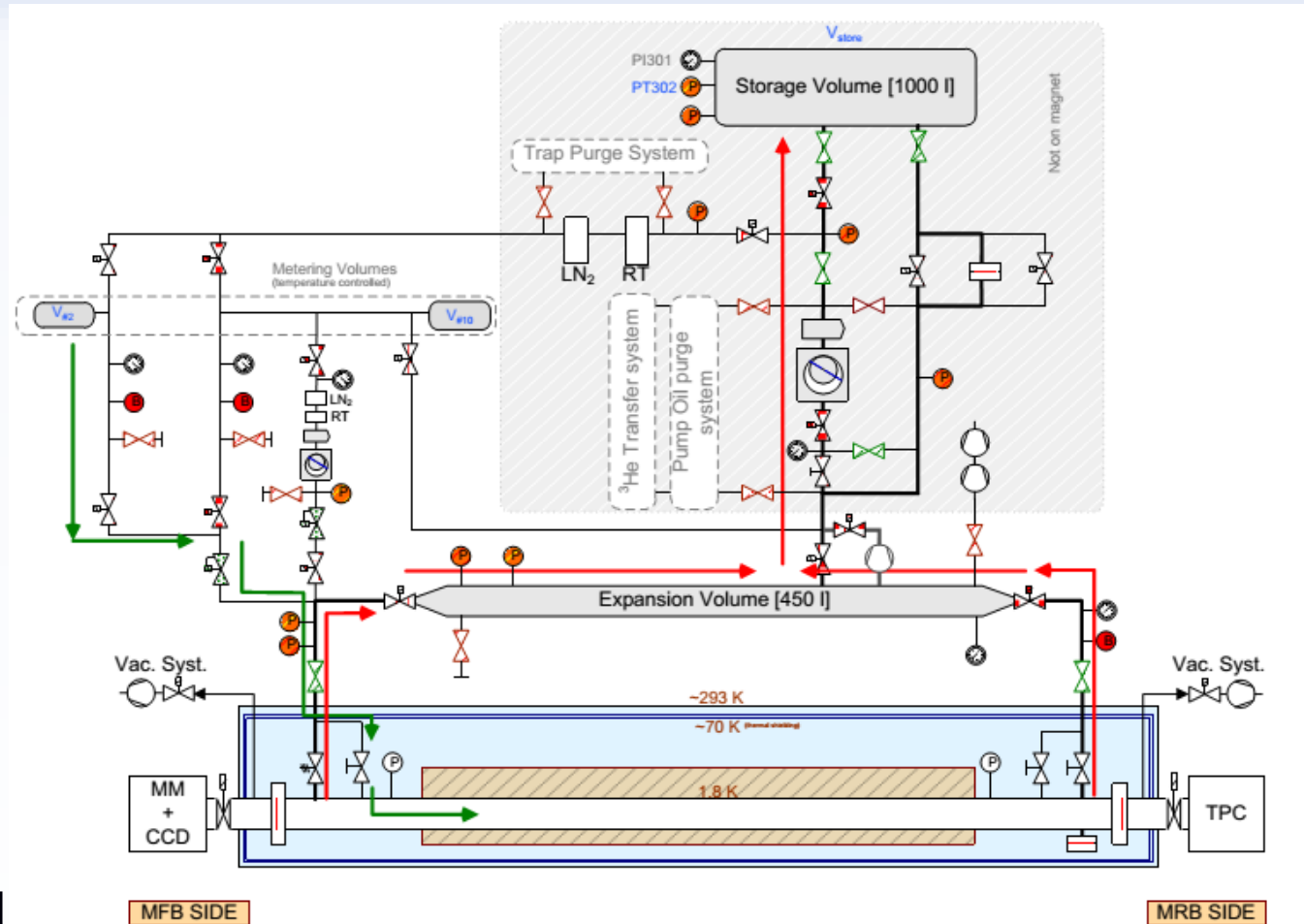
Computational Fluid Dynamics (CFD) simulations

The ^3He gas system

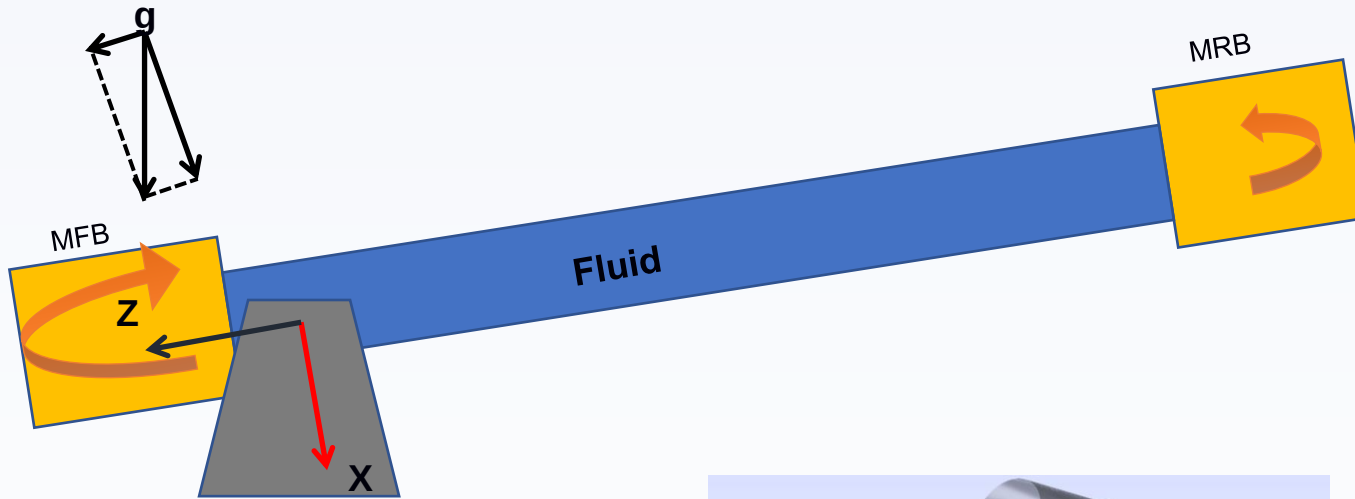


^3He gas system had to be accurate, stable and capable to precisely measuring small quantities of gas inserted into the cold bore volume in a reproducible way.

Pressure setting ~ -0.08 mbar

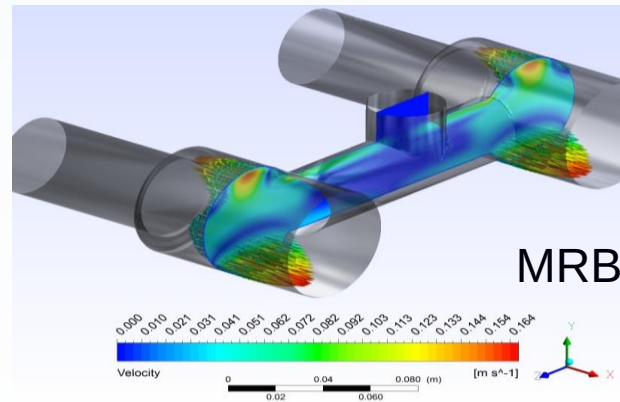


^3He CFD simulations for CAST (heat transfer and turbulence)

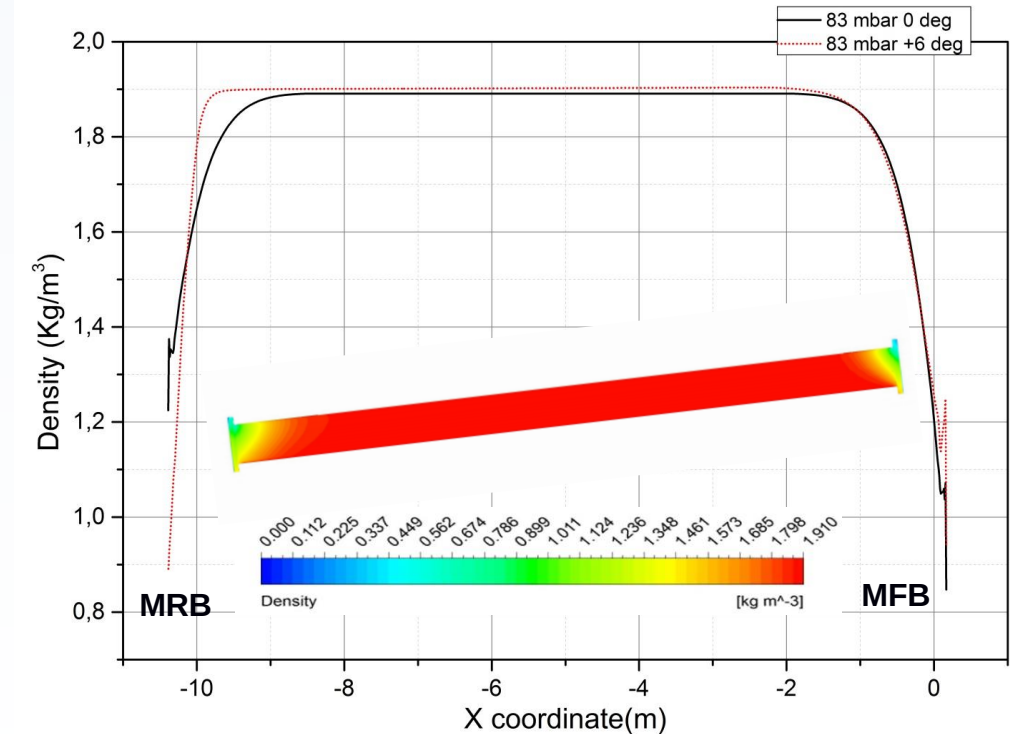
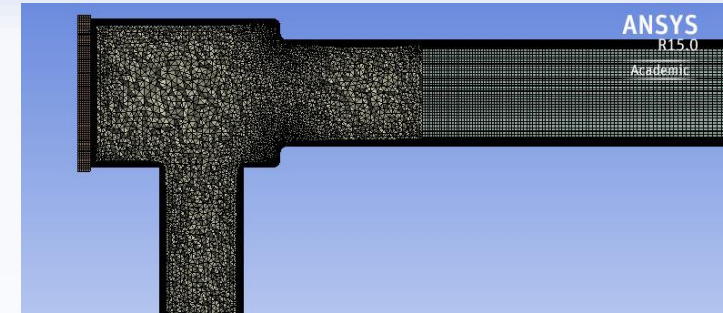


Peng-Robinson equation of state:

$$P = \frac{RT}{V_m - b} - \frac{a(T)}{V_m^2 + 2V_m b - b^2}$$



Extensive simulations for a most detailed model of the system under the different configuration



Density distributions along the center longitudinal line of the model

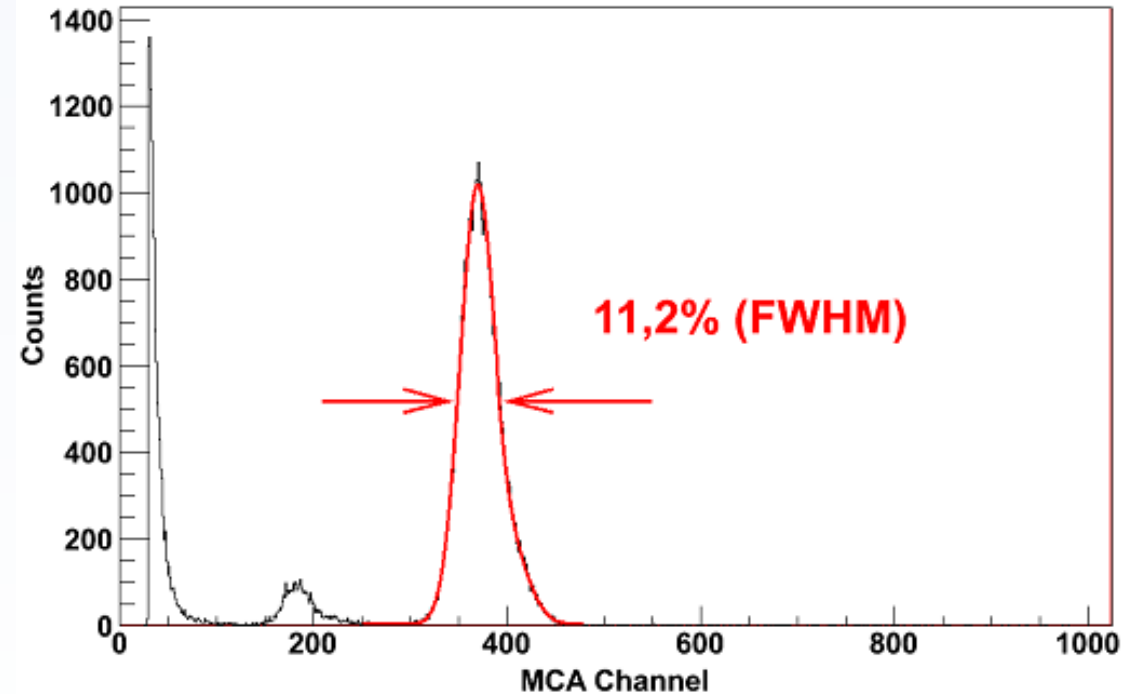
CAST Detectors

axion signal:

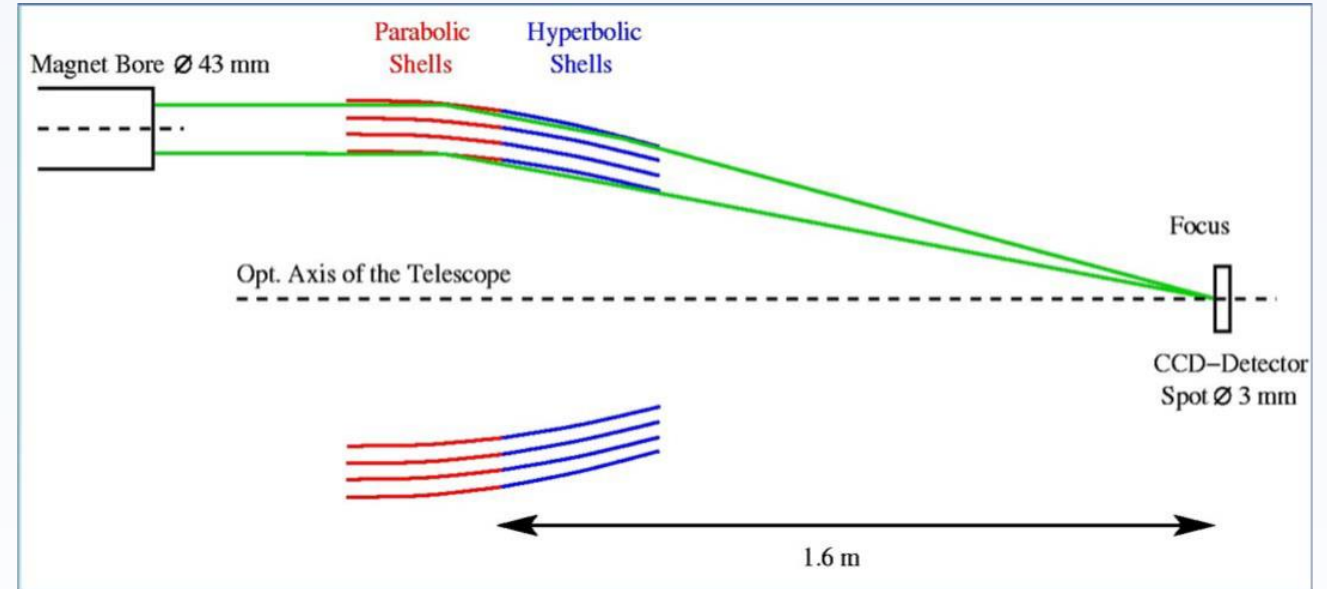
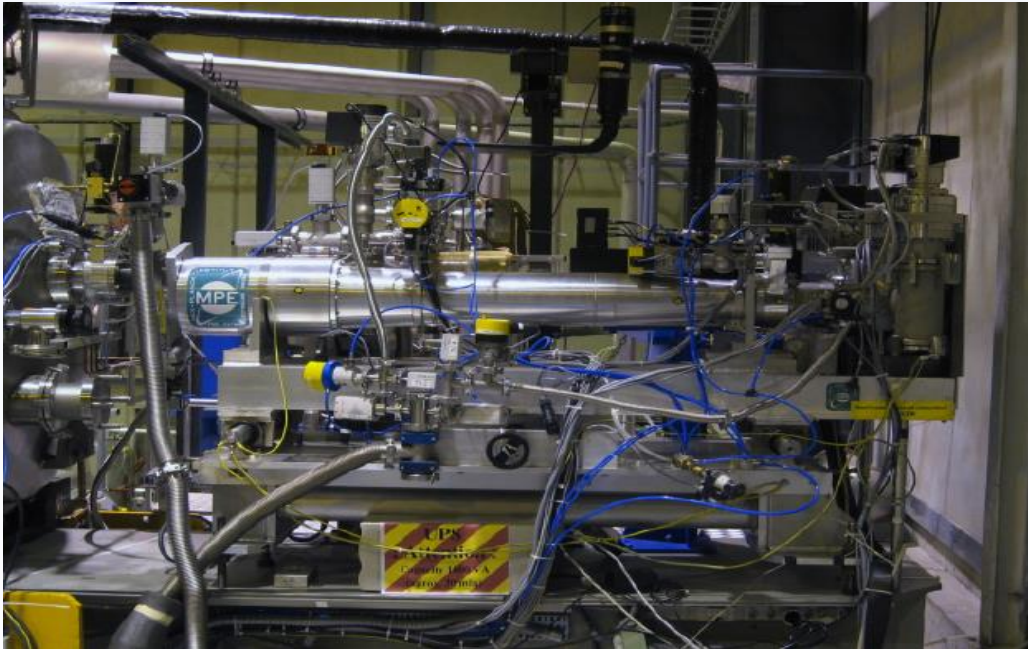
Excess of X-ray in tracking over background

$$N_\gamma = \int \frac{d\Phi_a}{dE} P_{\alpha \rightarrow \gamma} A t dE$$

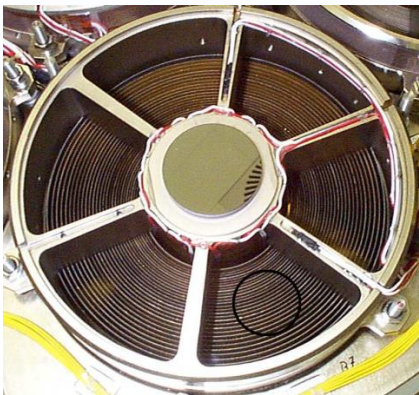
- ❑ time projection chamber (TPC)
- ❑ gaseous chamber Micromegas
- ❑ charge coupled device (CCD)



X-ray telescope of the CAST experiment



S/B improvement of ~150!



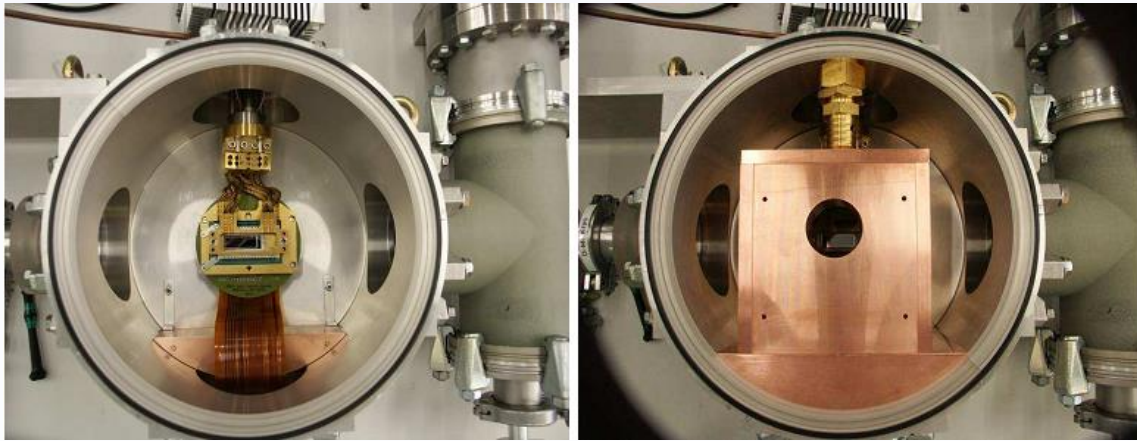
Wolter-I-type telescope (Prototype of ABRIXAS mission)

- 27 nested, gold-coated mirror shells
- Only one sector of telescope illuminated at CAST

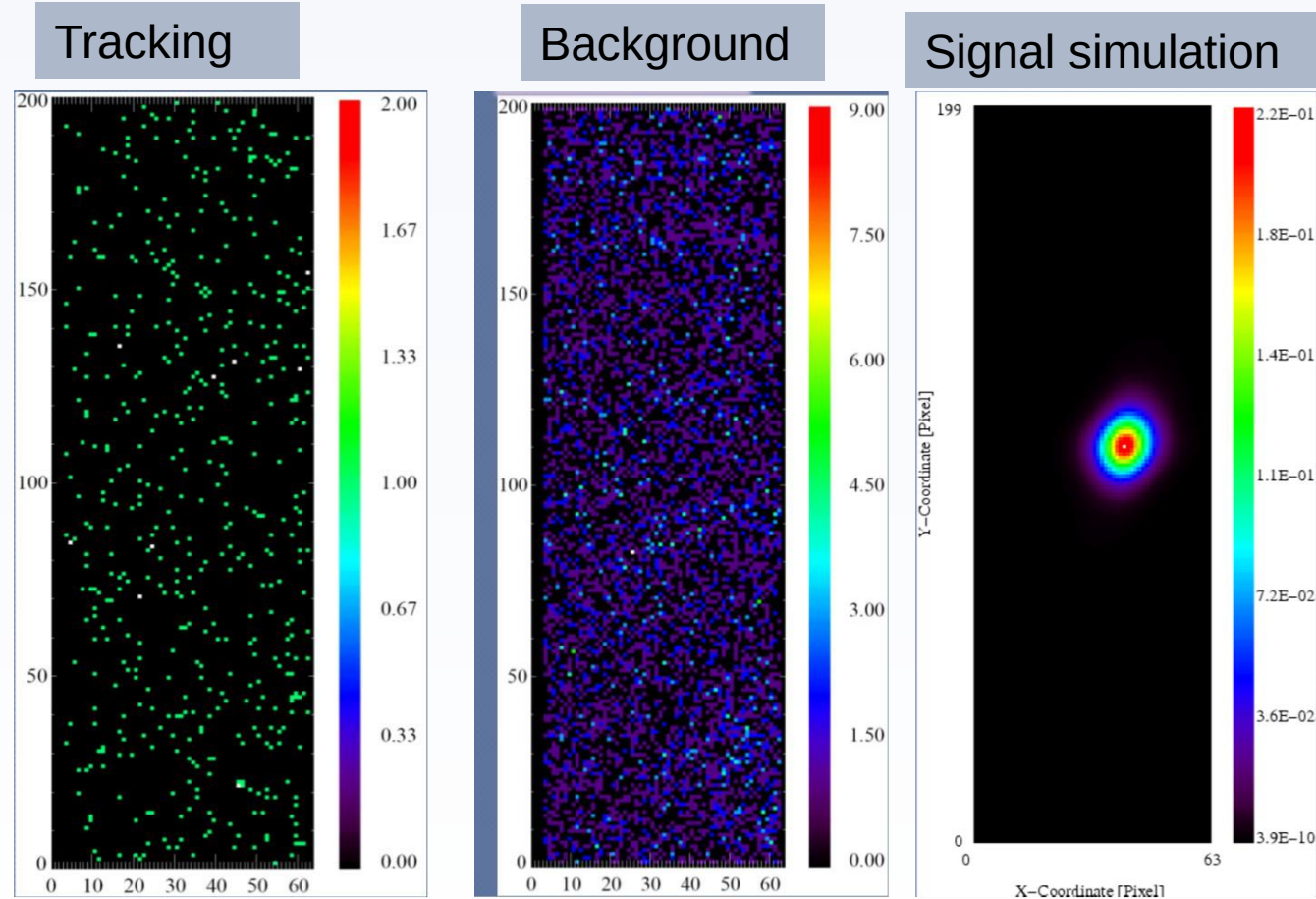
pn-CCD (Prototype of XMM-Newton mission)

- Very good spatial and energy resolution
- Simultaneous measurement of signal and background

Charged Coupled Device (CCD) and the X-ray telescope

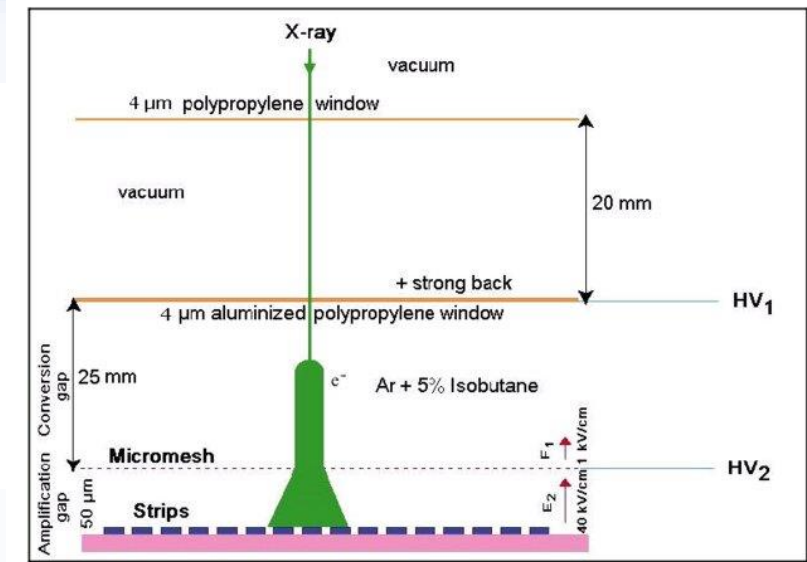
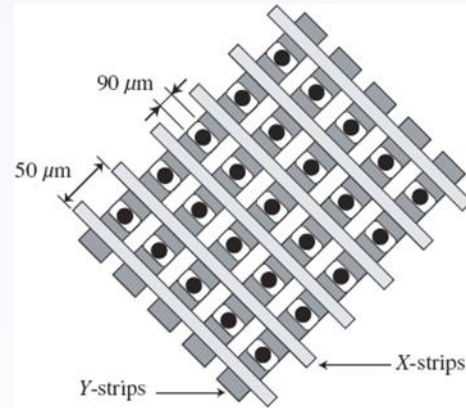
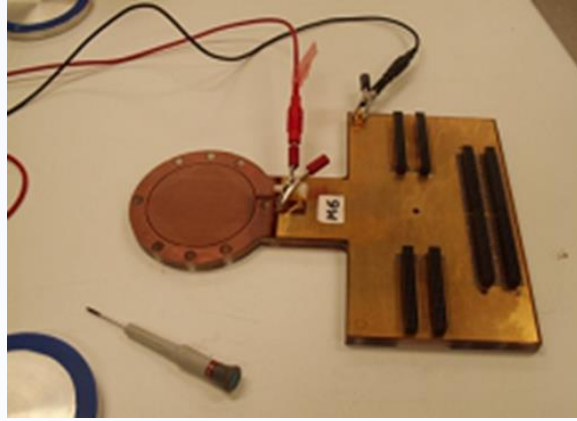


1×3 cm² sensitive area
64 × 200 pixels
280 mm thick
14 frames/sec_



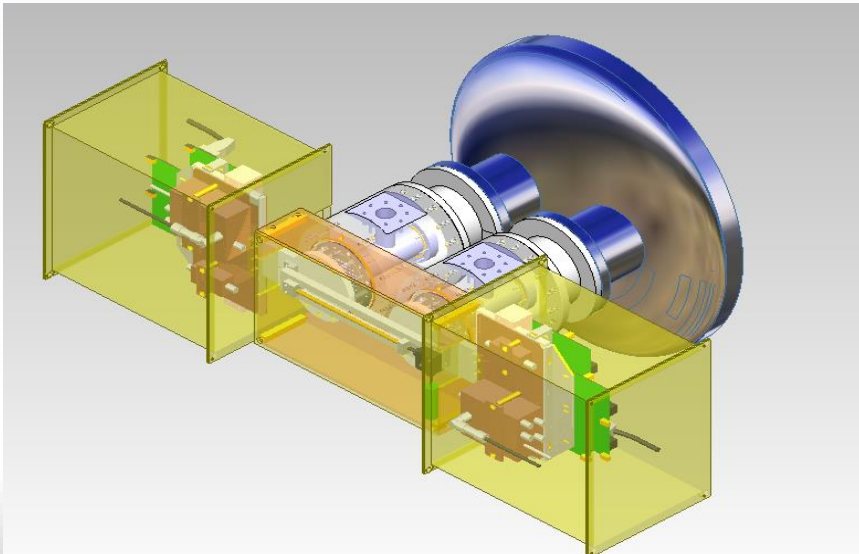
- Lower background (better selection of materials)
- Lower threshold (~ 200 eV)
- Better E resolution (< 160 eV (FWHM) @ 6 keV)
- State of the art technology

Micromegas detectors (MICRO MESH Gaseous Structure)



Micromegas area : 33.9 cm²

Magnet cold-bore area : 14.55 cm²



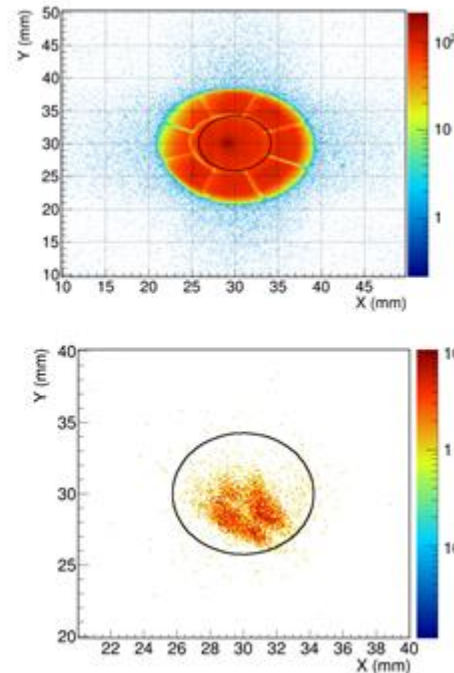
- The track time resolution of a Micromegas detector is ~ 0.60 ns
- Very stable operation during long periods
- The space resolution is less than 100 μm .
- High counting rate capabilities.
- sparking probability is now reduced

Typical rate: <2 c/h

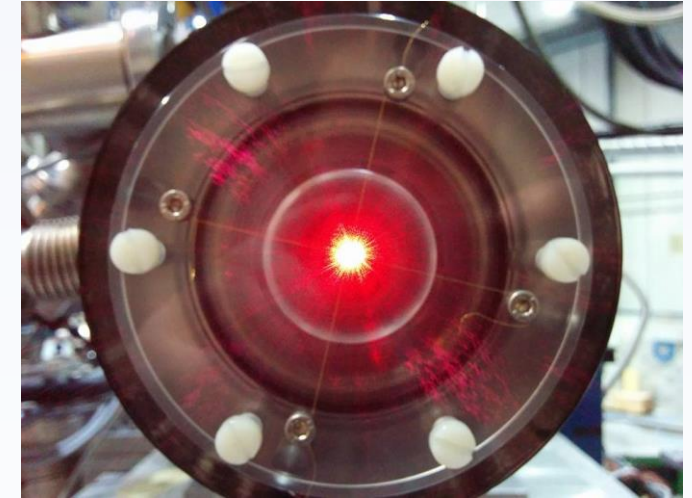
Nobel Prize, for the invention of the MultiWire Proportional Chamber (MWPC) by G. Charpak in 1968

Micromegas detectors -- Sunrise

^{55}Fe calibration

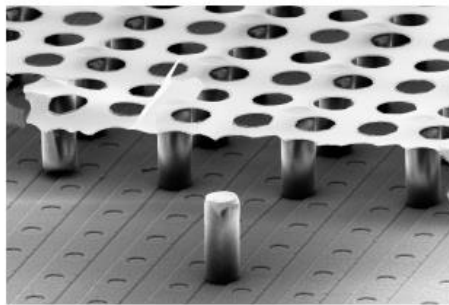


^{55}Fe calibration @13 m

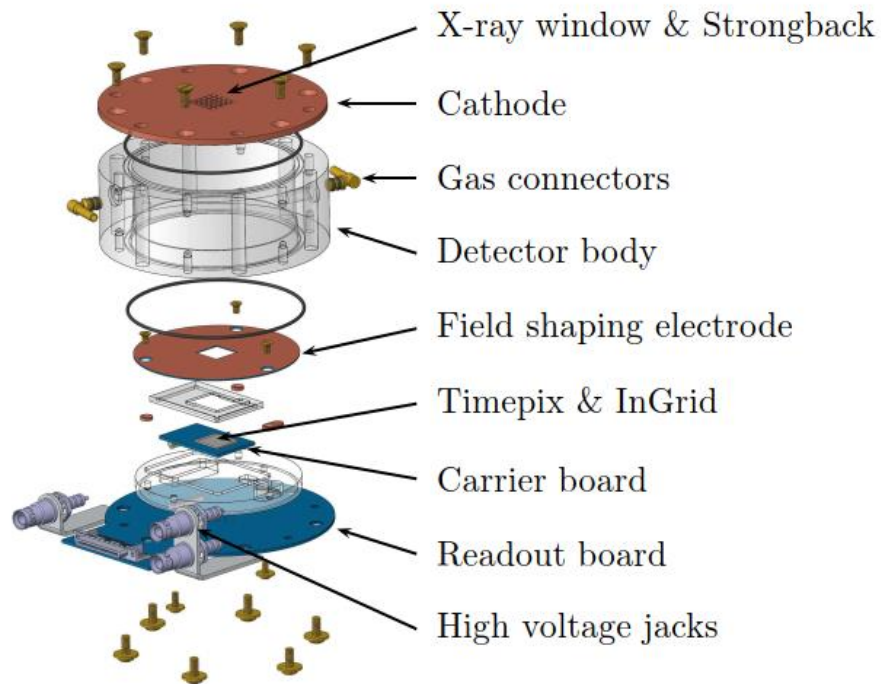
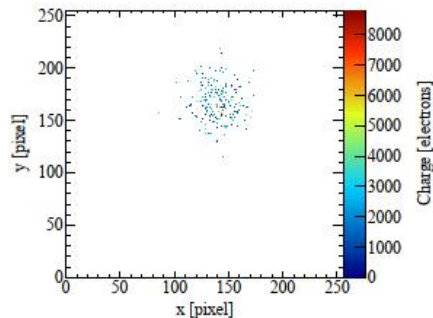


Improving Micromegas: InGrid X-ray Detector

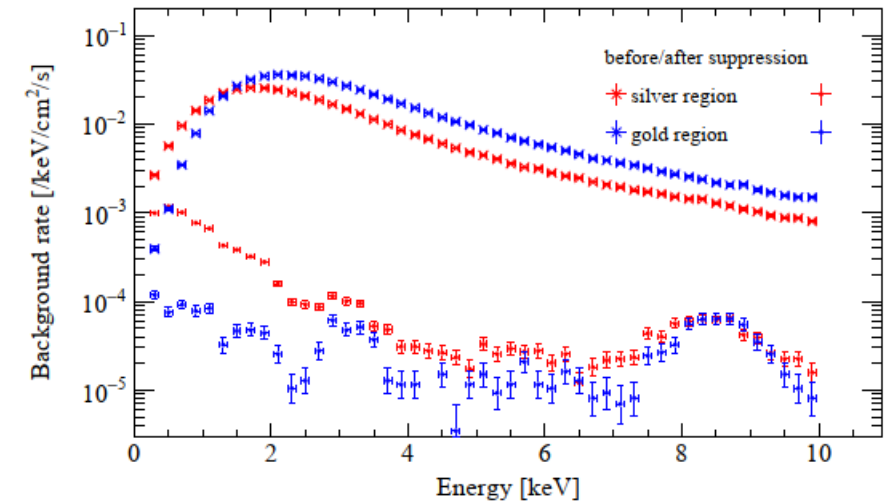
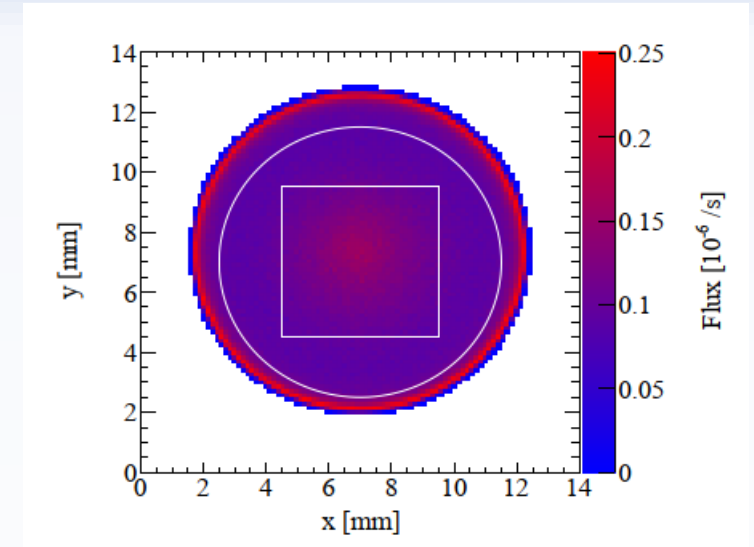
X-ray detector based on the combination of an integrated Micromegas stage with a pixel chip



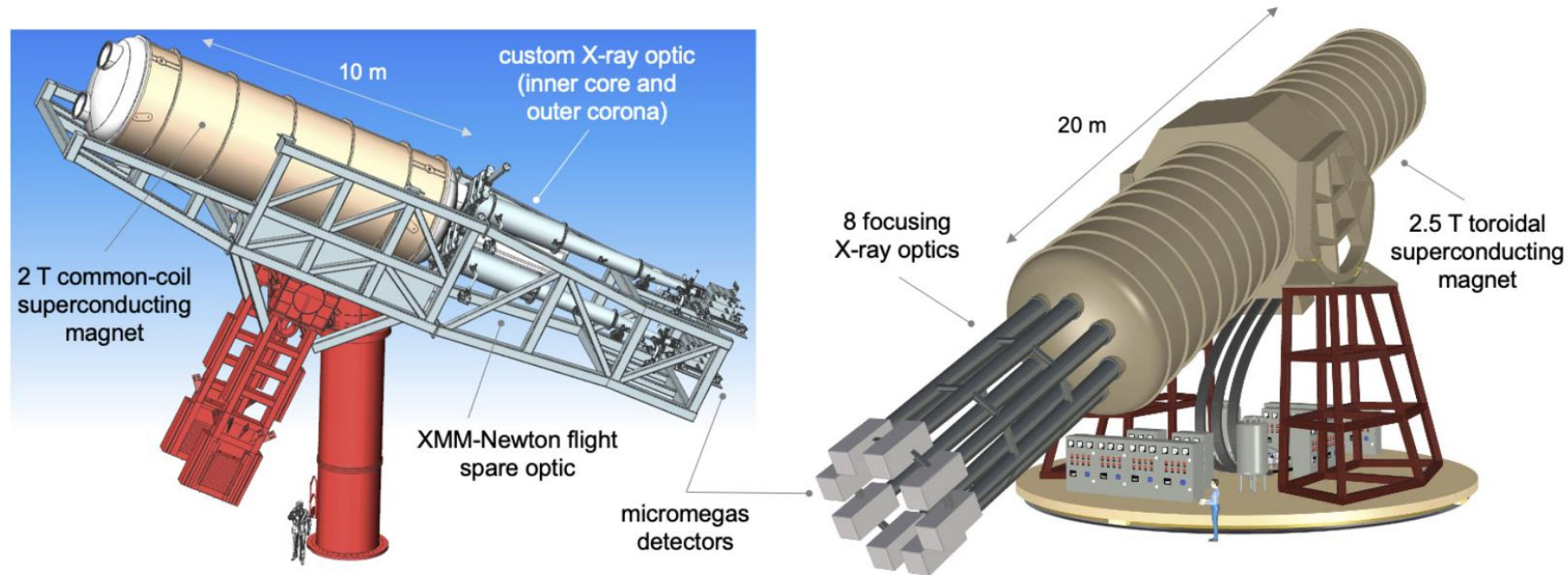
(a)



1×2 cm² sensitive area
256 × 256 pixels



CAST $\rightarrow$ IAXO (International Axion Observatory)

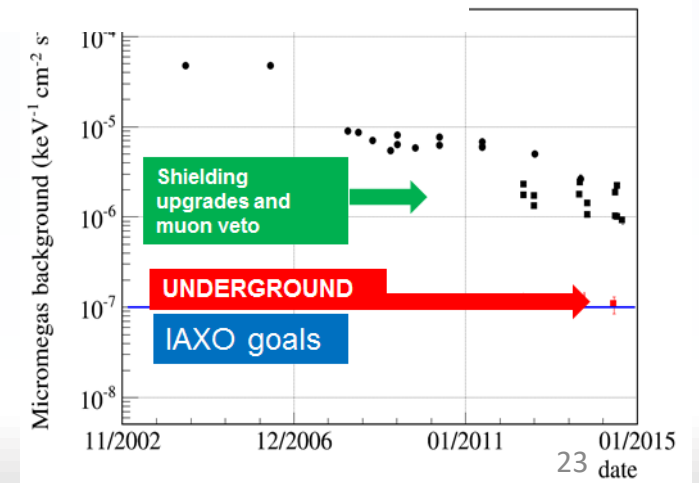


BabyIAXO (left) and IAXO (right)

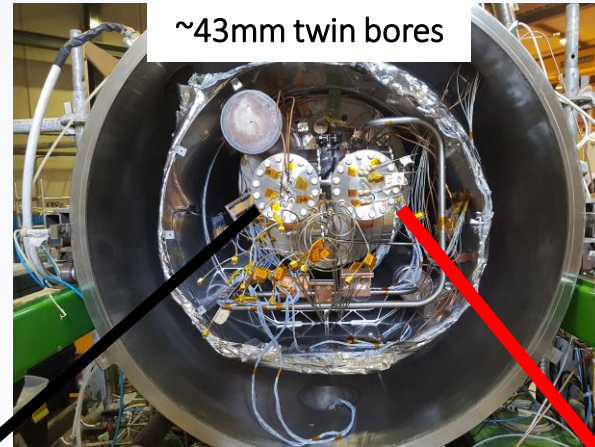
A next-generation axion helioscope, designed to search for axions from the Sun that are reconverted into X-ray photons via a strong laboratory magnetic field

IAXO: factor of 20 improvement in sensitivity

BabyIAXO : factor of 5 improvement in sensitivity



CAST 2019 → axion Haloscope



CAST-RADES

CAST-CAPP

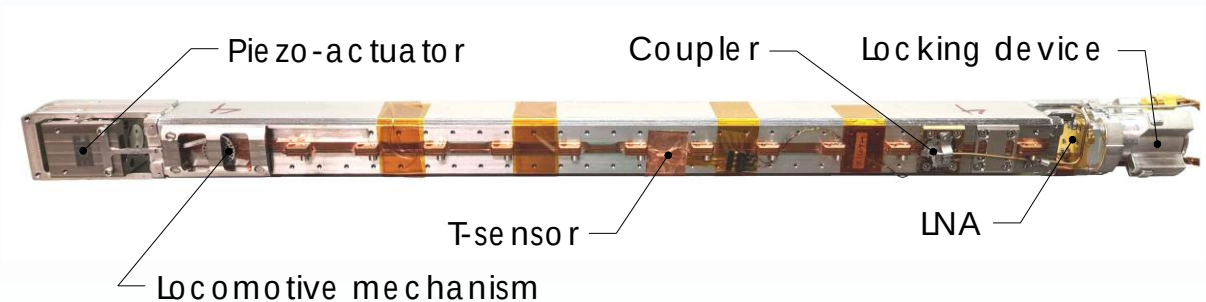
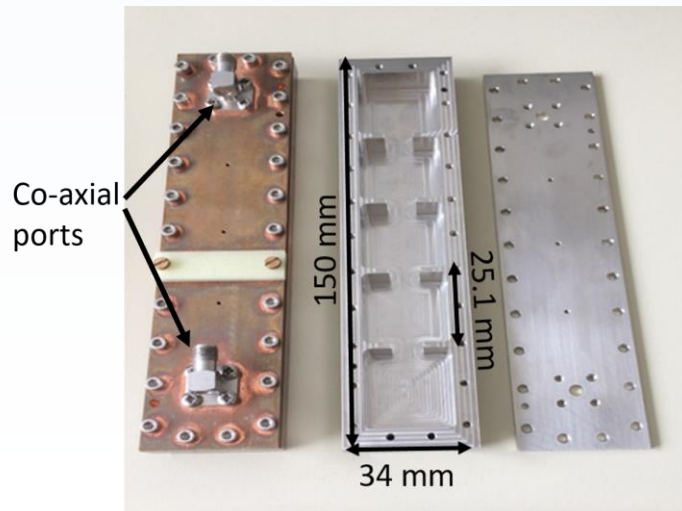
23 x 25 x 390 mm

$V = 224 \cdot 10^{-6} \text{ m}^3$

$Q_L = 20000$

$T_{\text{SYS}} = 10 \text{ K}$

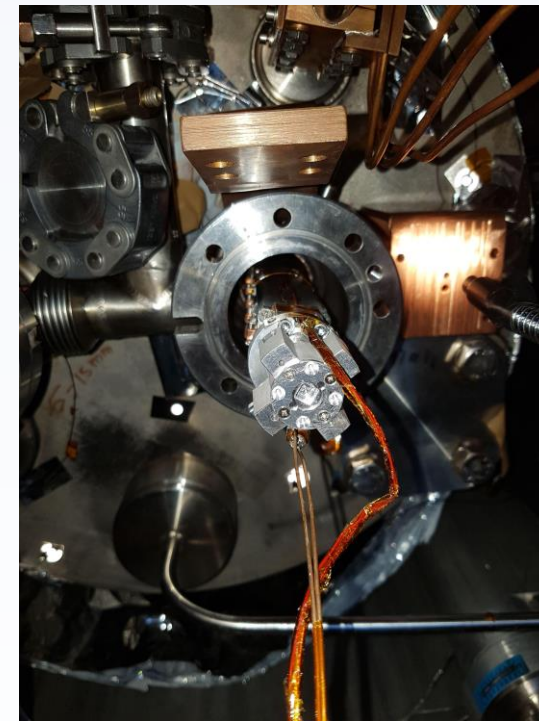
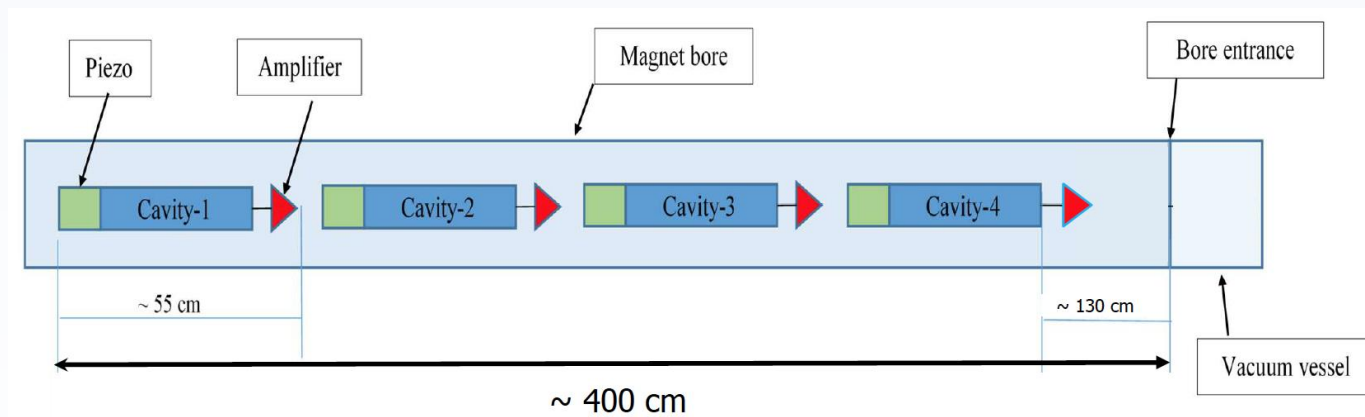
$C = 0.53$



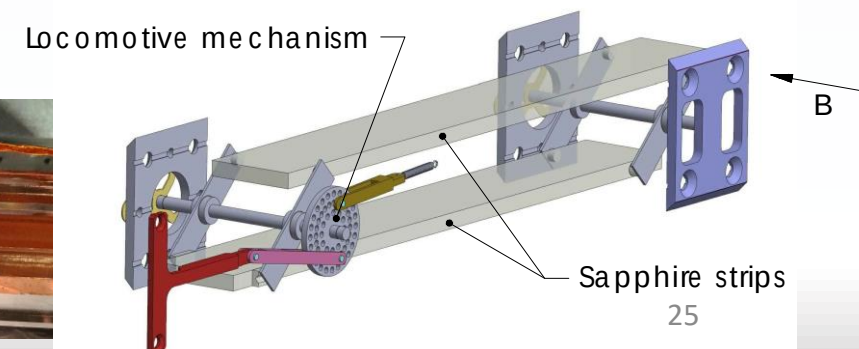
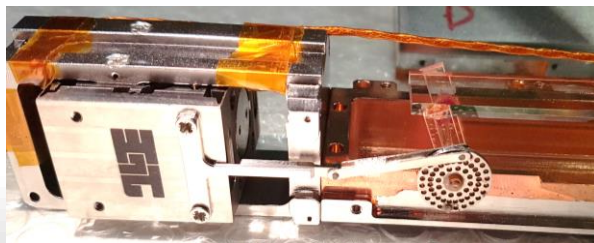
CAST-CAPP haloscope

- ✓ Phase-matching of all four cavities.
- ✓ Fast resonance scanning.
- ✓ Unexplored parameter space.

“phase matching” technique by combining coherently the power outputs of each frequency-matched cavity after individual signal amplification

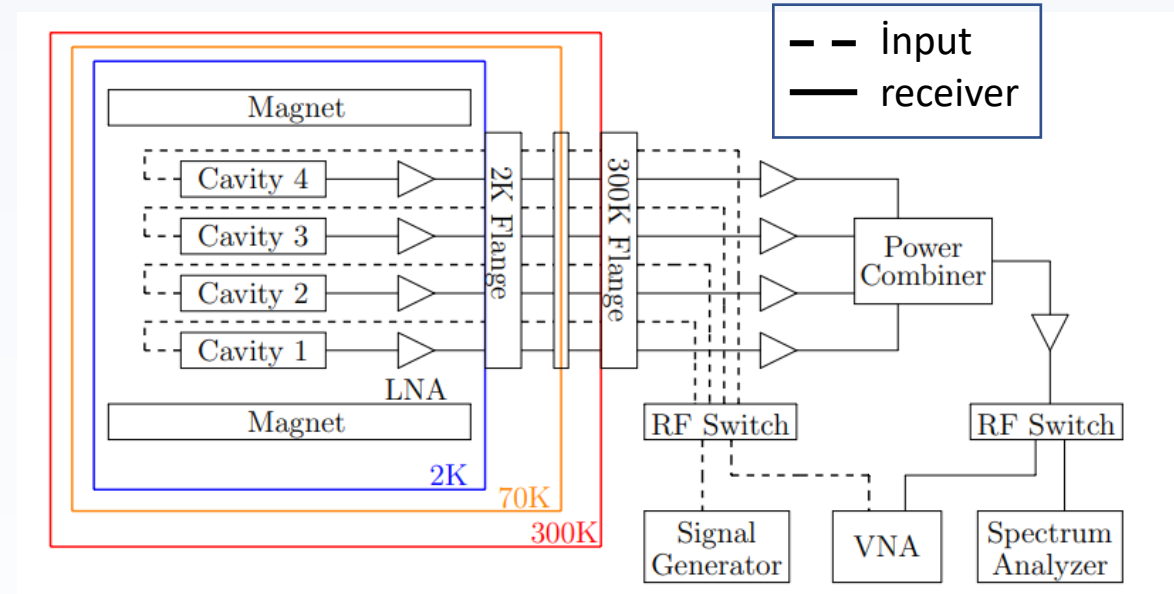


- Electromagnetic mode of interest : TE_{101}
- Frequency range: $\sim 4.8 - 5.4$ GHz (660 MHz)
- Axion mass range: $\sim 19.7 - 22.4$ μeV
- Resolution: 100 Hz
- Max piezo speed: 10 MHz / min

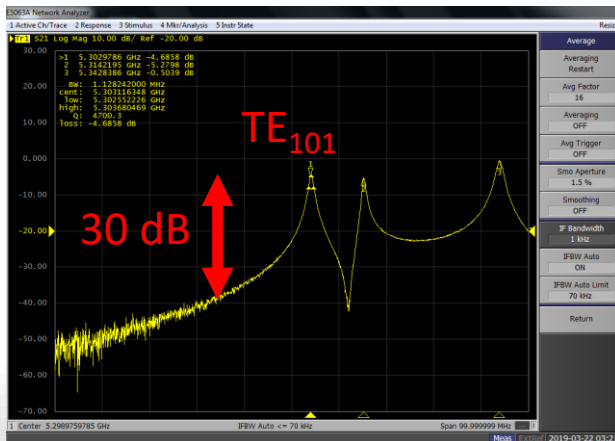


Detection scheme of CAST-CAPP

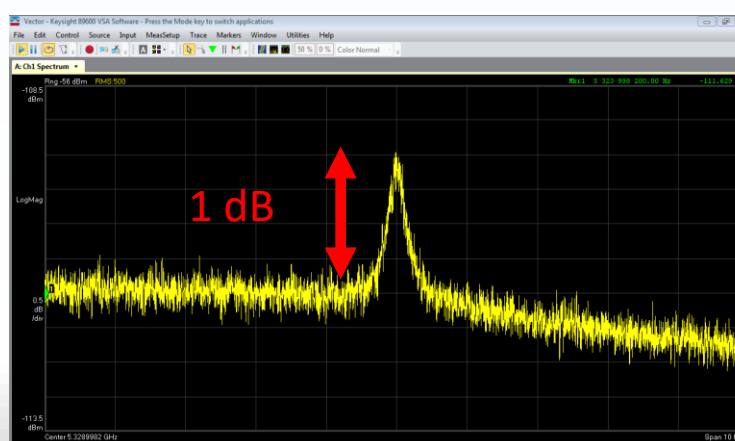
- CAST-CAPP diagram:
 - CAST cryostat layers (colors)
 - Cavities aligned in magnet bore
 - Cold and warm amplifiers
 - Power combiner
- 2nd receiver channel setup for environmental EMI



Vector Network Analyzer

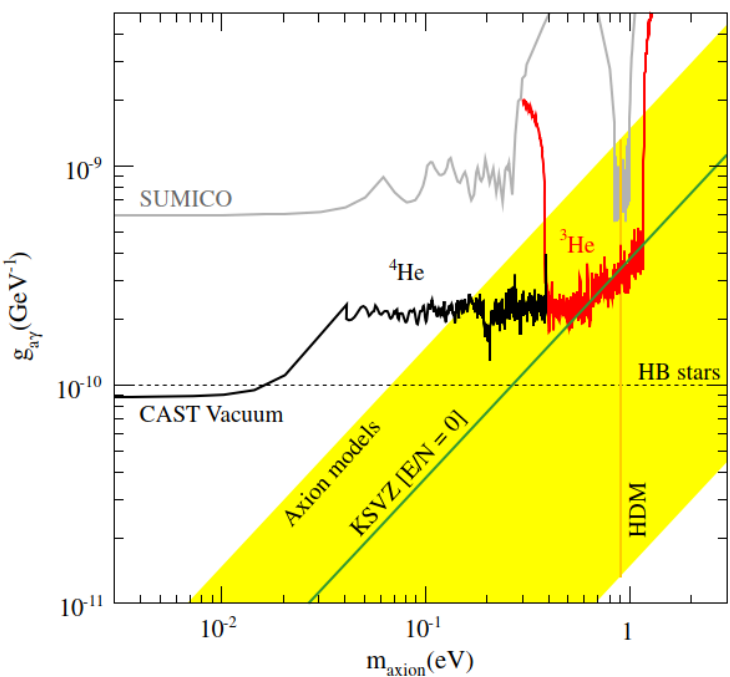


Spectrum Analyzer



CAST axion limits and major publications

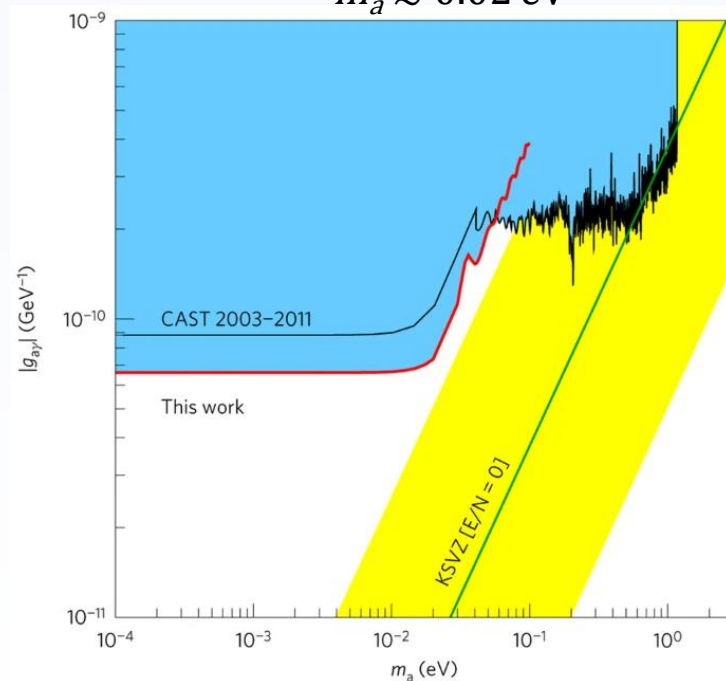
$g_{\alpha\gamma} \leq 3.3 \times 10^{-10} \text{ GeV}^{-1}$ at 95% CL.
0.64–1.17 eV



PRL 112, 091302 (2014)

2013–2015 run

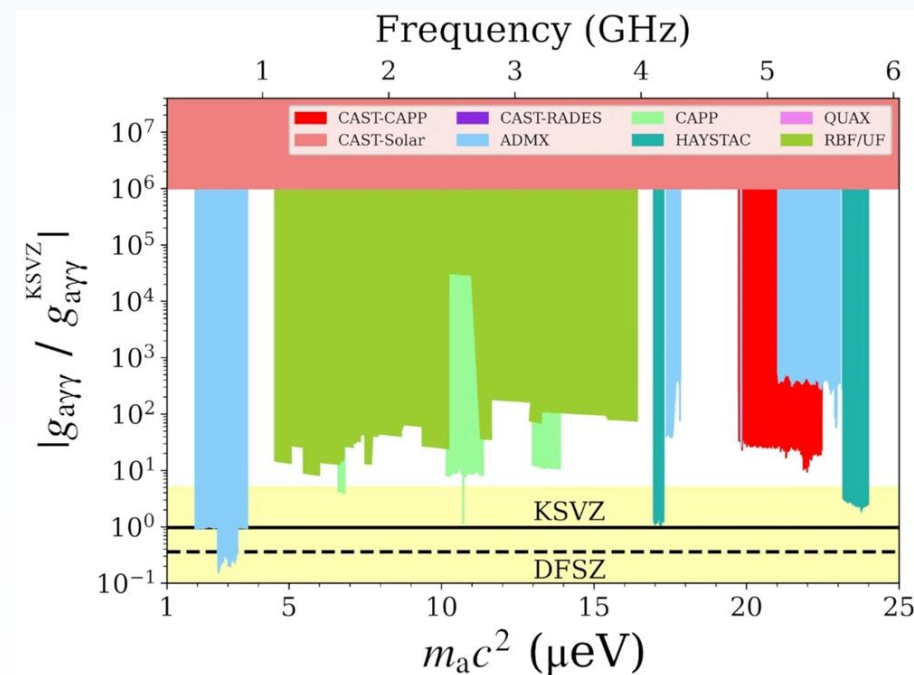
$g_{\alpha\gamma} \leq 0.66 \times 10^{-10} \text{ GeV}^{-1}$ at 95% CL.
 $m_a \lesssim 0.02 \text{ eV}$



Nature Phys. 13 (2017) 584-590

2019 to 2021 run

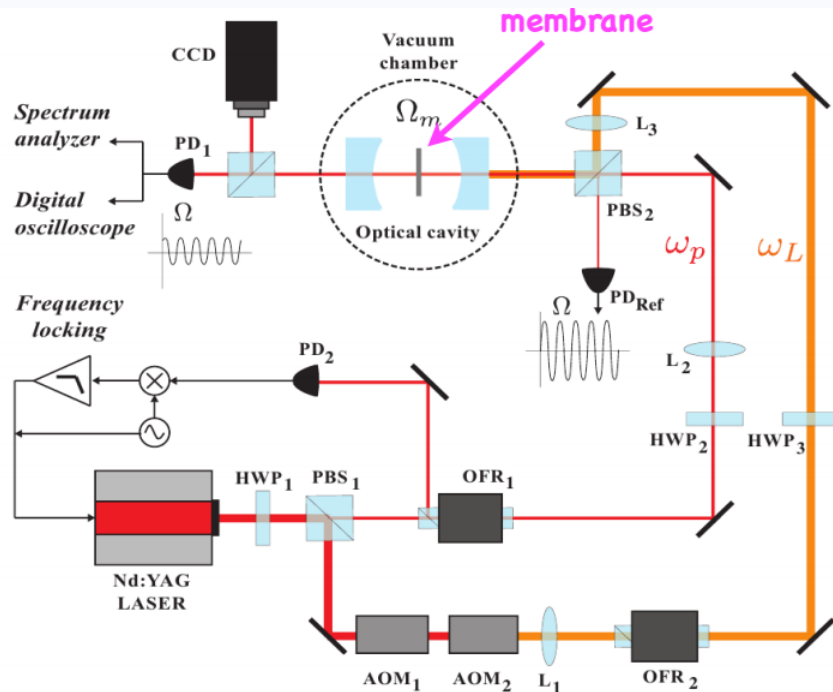
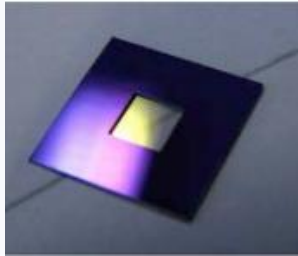
$g_{\alpha\gamma} \leq 8 \times 10^{-14} \text{ GeV}^{-1}$ at 90% CL.
19.74 μeV and 22.47 μeV



Nat Commun 13, 6180 (2022)

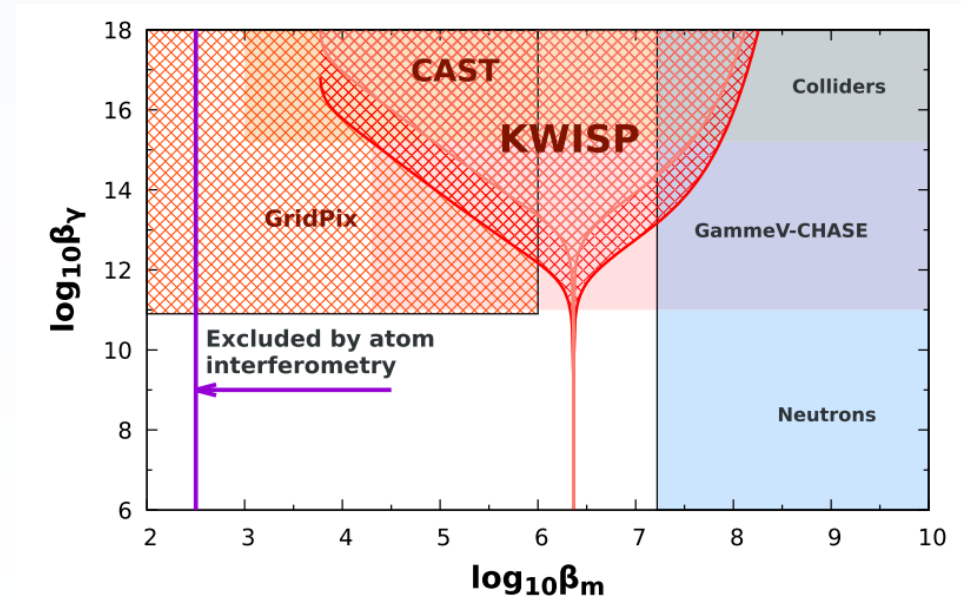
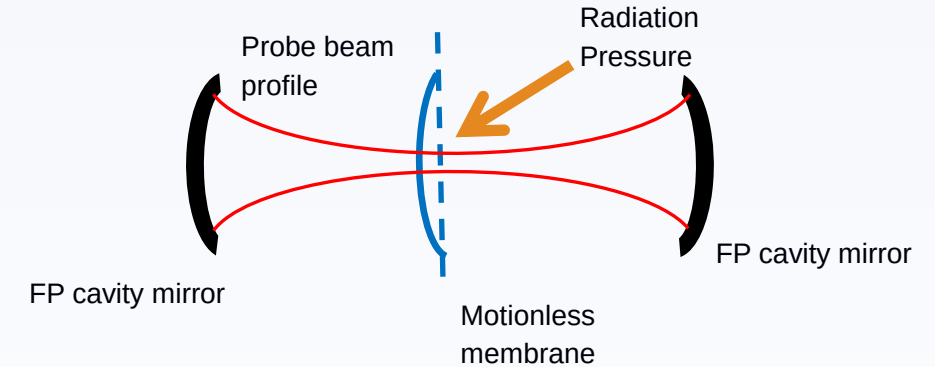
Detection of Solar Chameleons through Radiation Pressure

A force sensor that is called KWISP for “Kinetic WISP detection” (Weakly Interactive slim Particles)



Commercially available

- "large" area, up to 5x5 mm
- Transparent @1064 nm
- Density $\sim 3 \text{ g/cm}^3$
- Can be coated with metal
- High stress
- High resonant frequency
- High mechanical $Q \sim 10$





02/2016

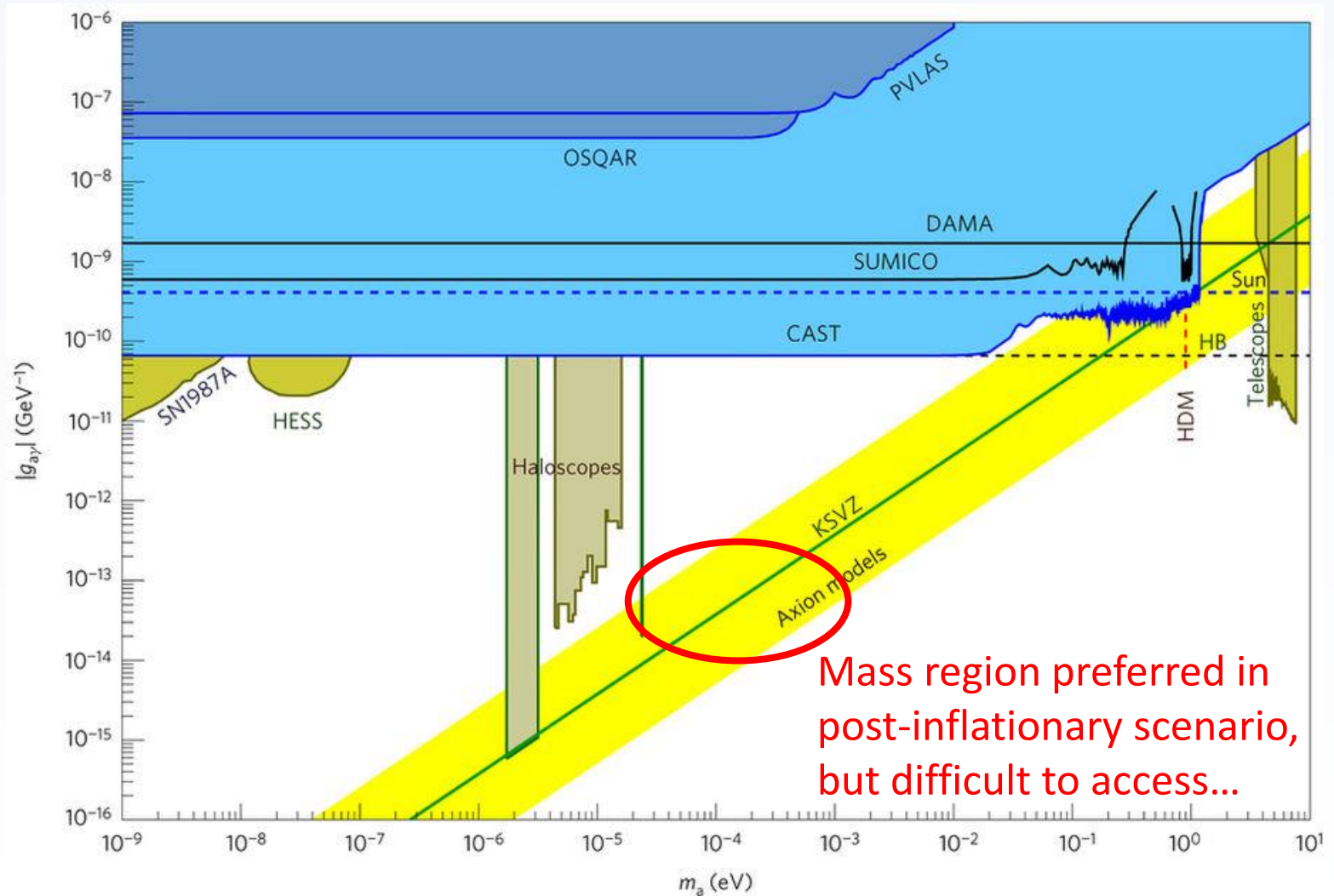
Dielectric Haloscopes

(post-inflationary scenario) $N_{\text{DW}}=1$
40 to 180 μeV

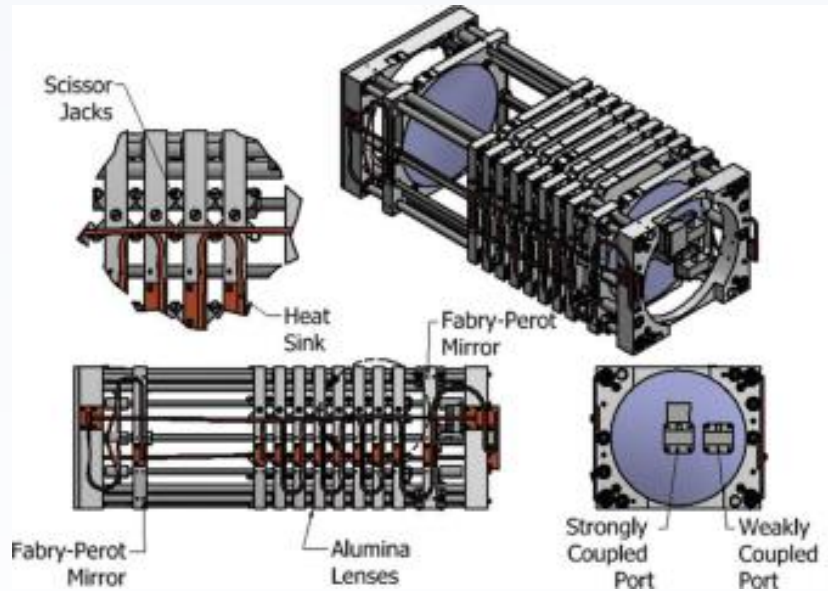
Nat. Commun. **13**, 1049 (2022)

For domain wall number $N_{\text{DW}}>1$
28 μeV to meV

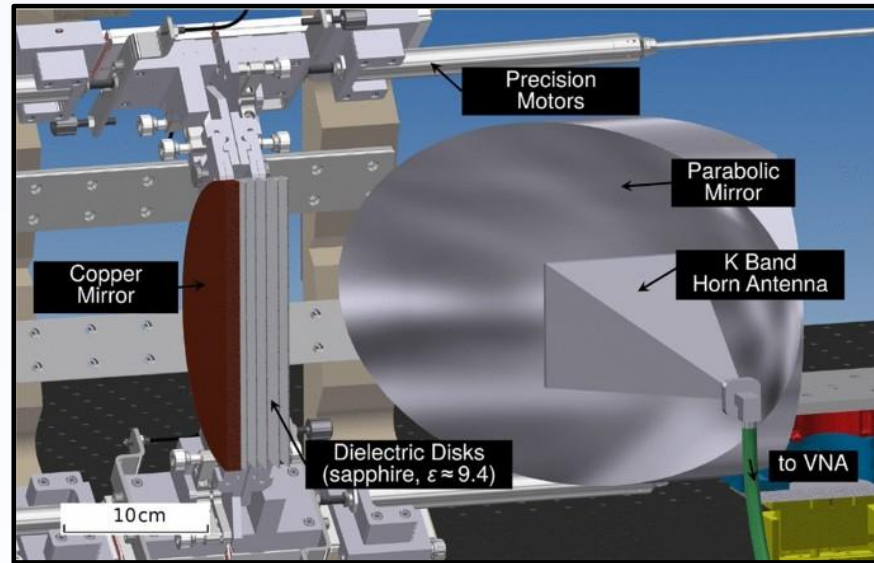
Andreas Ringwald and Ken'ichi Saikawa
Phys. Rev. D 93, 085031



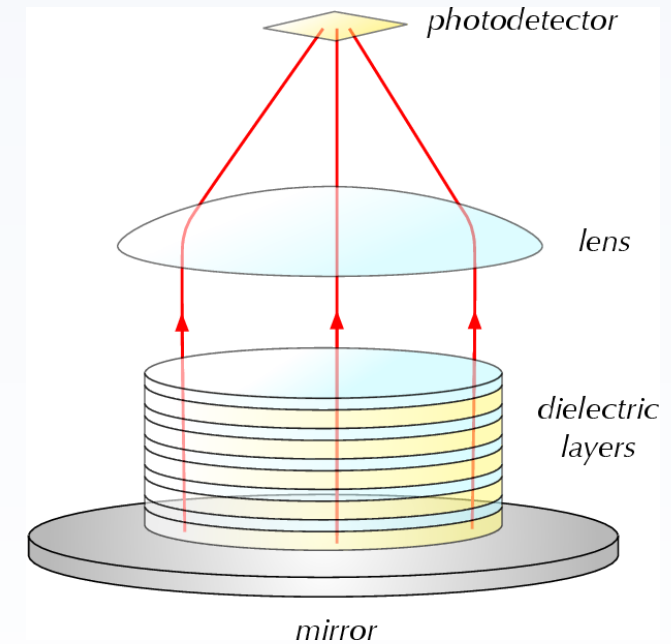
Dielectric Haloscopes



ADMX-Orpheus

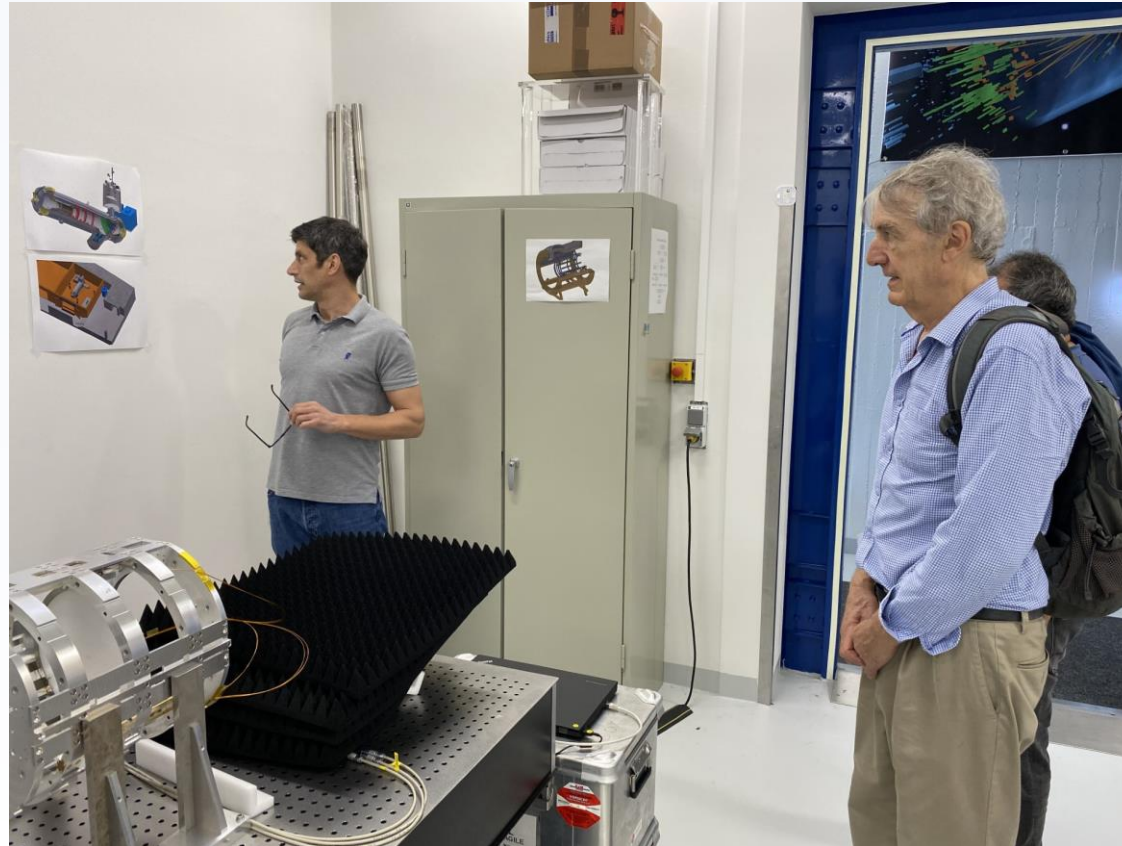


MADMAX



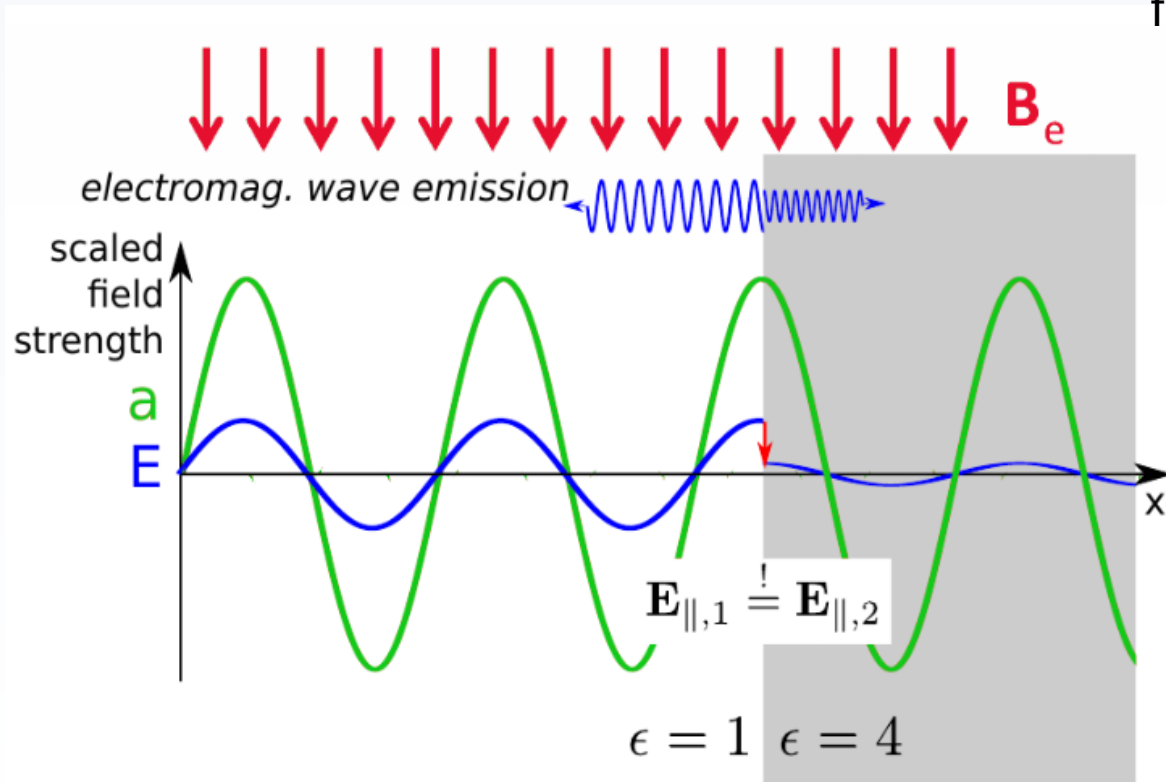
MuDHI

Dielectric Haloscopes - MADMAX



MADMAX Dielectric Haloscope

The cosmological axion field $\mathbf{a}(\mathbf{t})$ inside an external magnetic field $\mathbf{B}_e$ sources a tiny axion induced electric field $\mathbf{E}_\alpha$



$$E_\alpha \sim 10^{-12} \text{ V/m for } 10 \text{ T in vacuum}$$

E_α is different in materials with different ϵ

$E_\parallel$ must be continuous at the **boundaries**

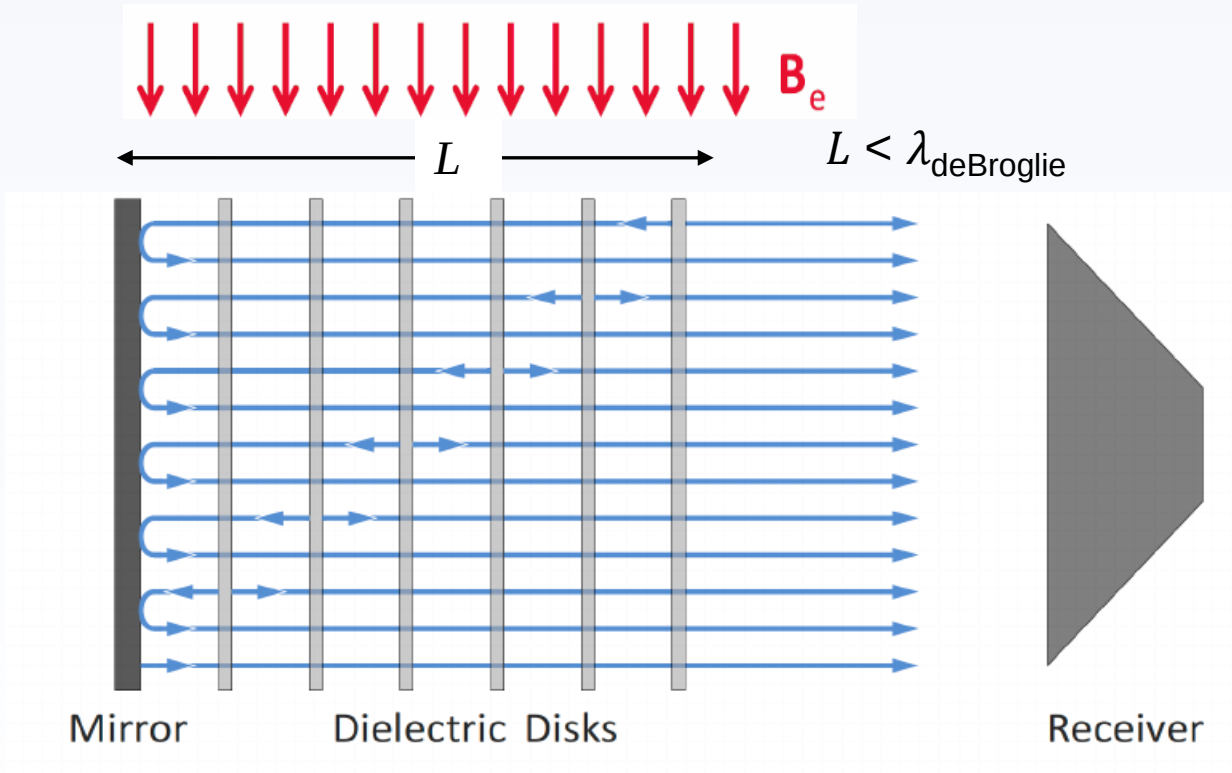
$$E_\alpha = -\frac{g_{\alpha\gamma} B_e}{\epsilon} \alpha \quad \alpha = \alpha_0 \cos(m_\alpha t)$$

Power emitted from a single surface:

$$P_{sig} = 2.2 \cdot 10^{-27} \text{ W} \left(\frac{A}{1 \text{ m}^2} \right) \left(\frac{B_e}{10 \text{ T}} \right)^2$$

or 2 photons/hour @ 25 GHz

Detection principle



- EM radiation escapes at open end
- Detect Traveling wave instead of standing wave modes

Boosted emitted power :

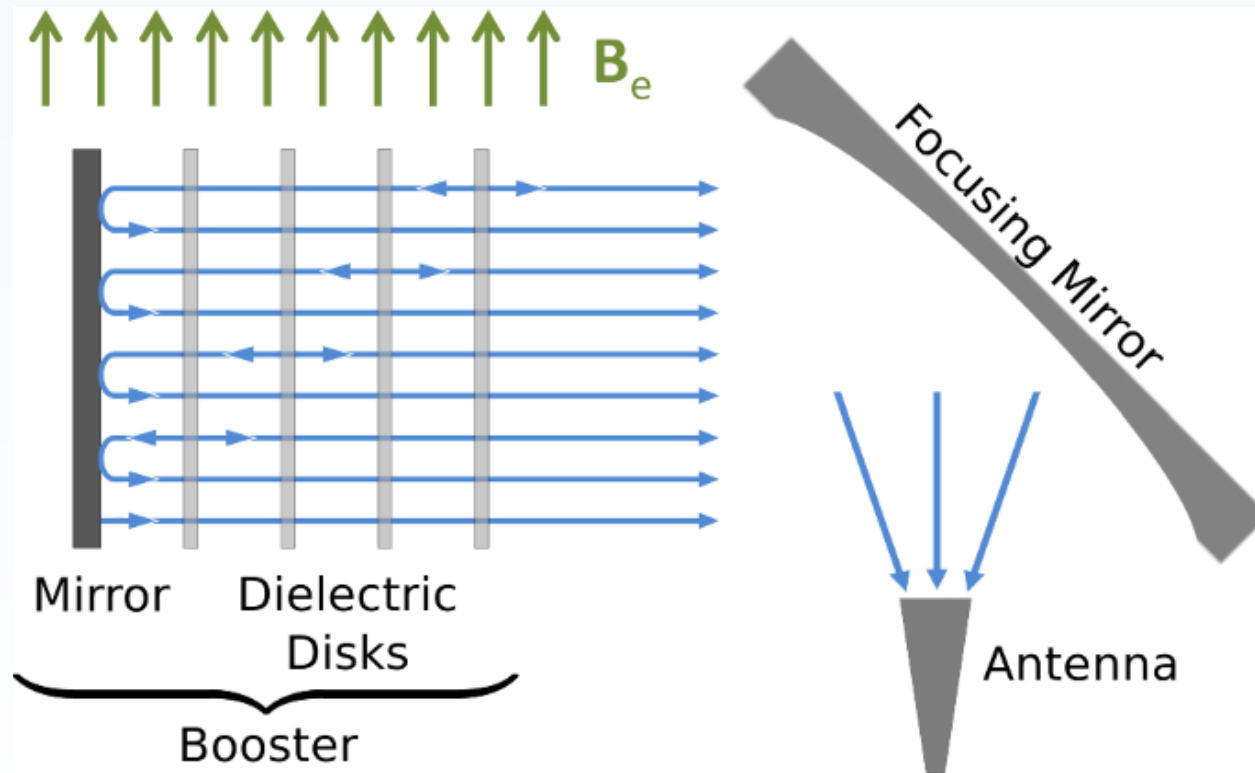
- coherent emission from multiple interfaces by controlling the discs separations and
- constructive interference

$$P_{sig} = 2.2 \cdot 10^{-27} W \left(\frac{A}{1m^2} \right) \left(\frac{B_e}{10T} \right)^2 \cdot \beta^2$$

power “Boost factor” β^2

$$\beta^2 = \frac{P_{sig}}{P_{mirror}}$$

Detection principle



- EM radiation escapes at open end
- Detect Traveling wave instead of standing wave modes (?)

power “Boost factor” β^2

$$\beta^2 = \frac{P_{sig}}{P_{mirror}}$$

Simulations indicate:

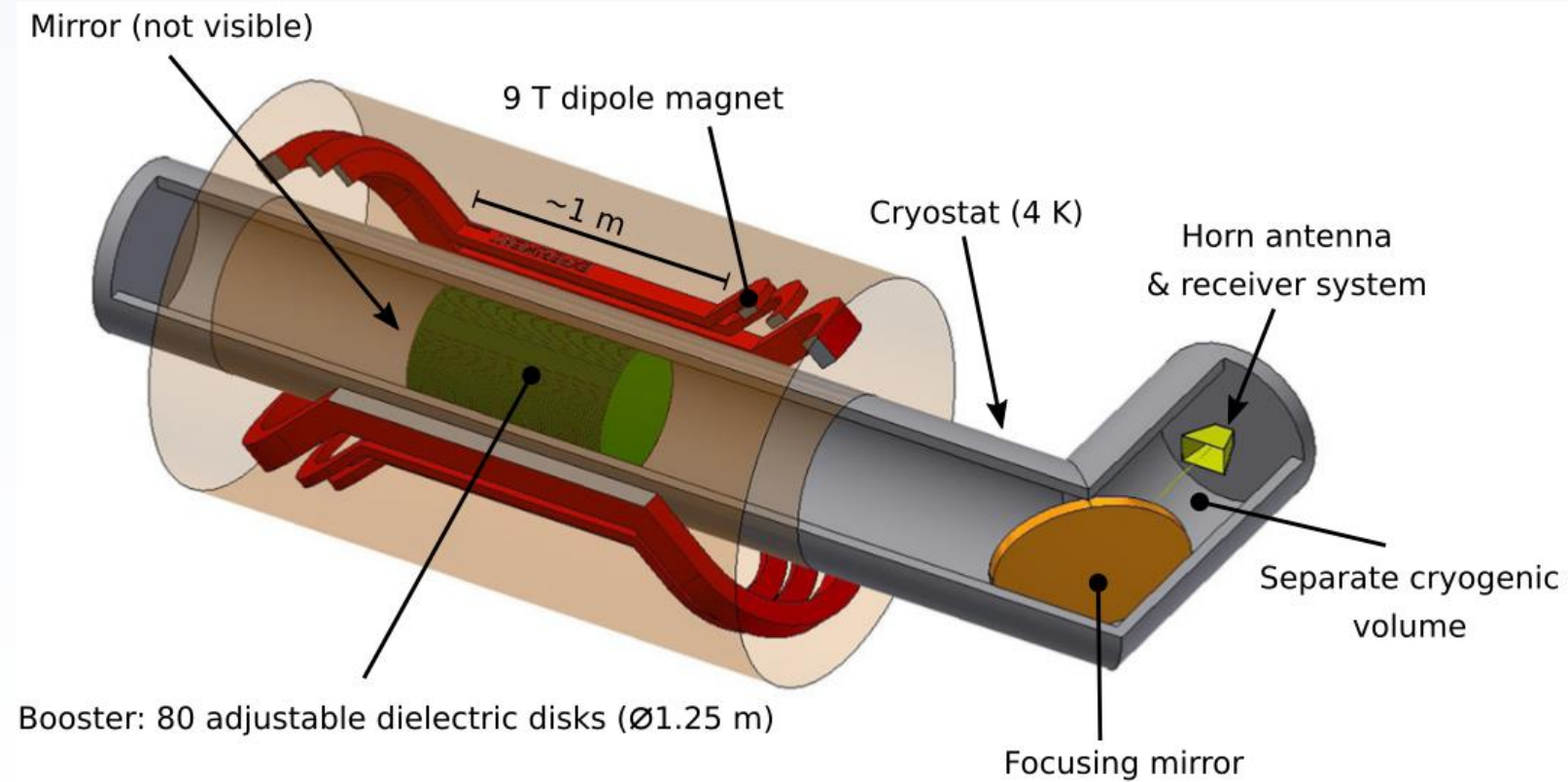
$|\beta^2| > 10^4$ achievable with 80 discs
with dielectric constant $\epsilon \sim 24$ (LaAlO₃)

β^2 is affected by BC’s inaccuracies of the discs:
Disc mis-positioning, tilting, thickness variations

MADMAX full scale experiment

Signal power

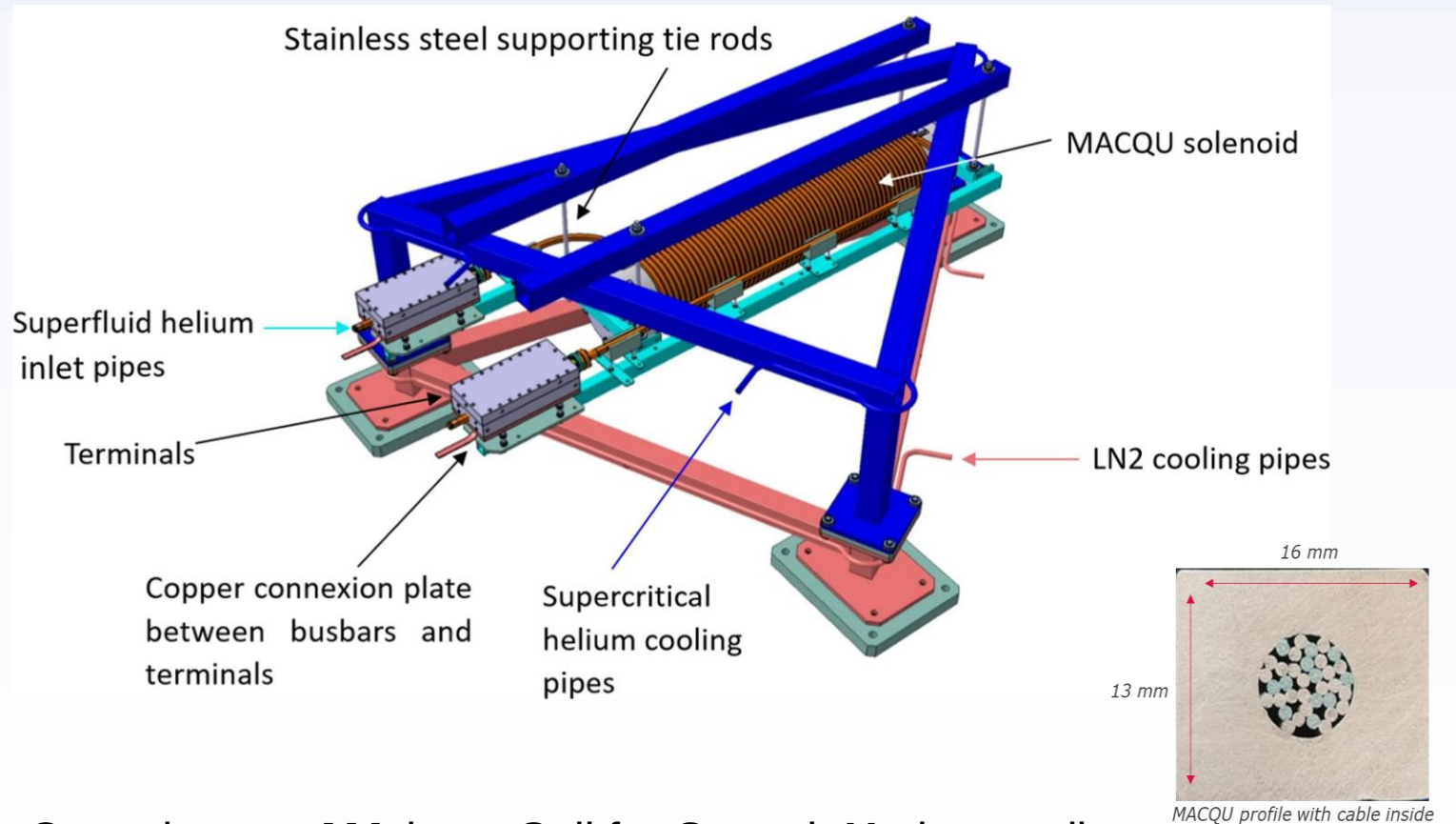
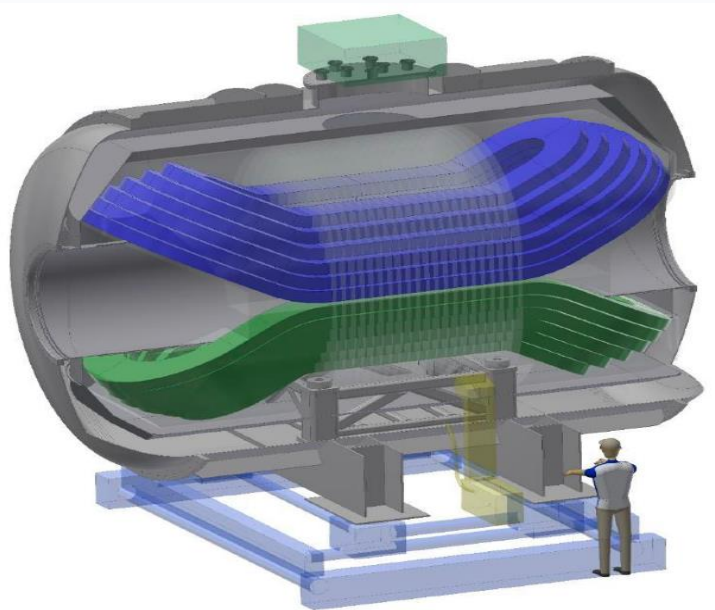
$$P_{sig} = 2.2 \cdot 10^{-27} W \left(\frac{A}{1m^2} \right) \left(\frac{B_e}{10T} \right)^2 \cdot \beta^2$$



Magnet

Main challenges

- Booster mechanics
- **Magnet**
- Receiver at cold, B-field environment
- Funding

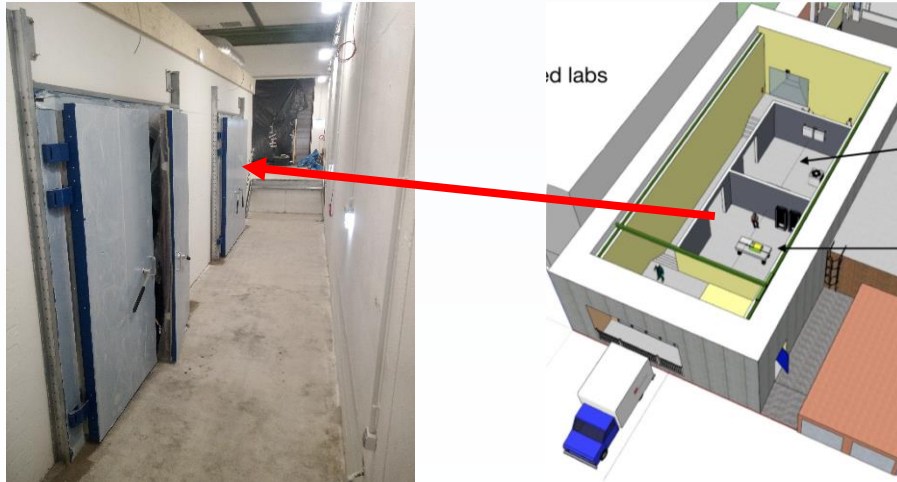


- Quench tests: **M**Admax **C**oil for **Q**uench **U**nderstanding
- Reliable quench protection
- Feasibility of conductor production

$$\text{FoM} = B^2 A = 100 \text{ T}^2 \text{ m}^2$$

Timeline

- Commissioning of full-scale MADMAX to start around 2028
- Medium scale prototype:
 - 20 discs with $\varnothing 300$ mm
 - Frequency range 18 to 25 GHz
 - To be commissioned in 2024 at Hamburg in new RF lab
 - Possible physics measurements in 2022 and 2023
(currently under discussion: Morpurgo magnet at CERN (1.6 T))



RF lab

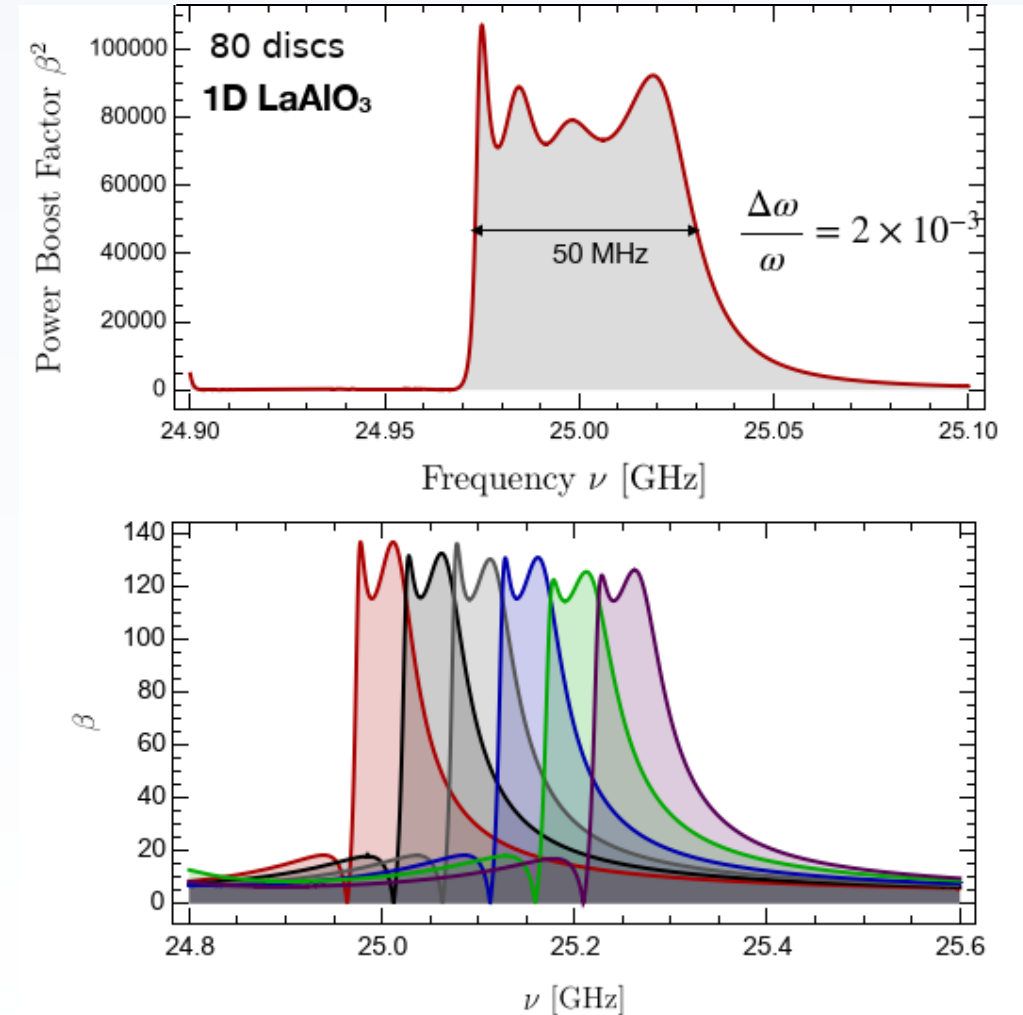


Morpurgo magnet

Properties Of The “Boost Factor”

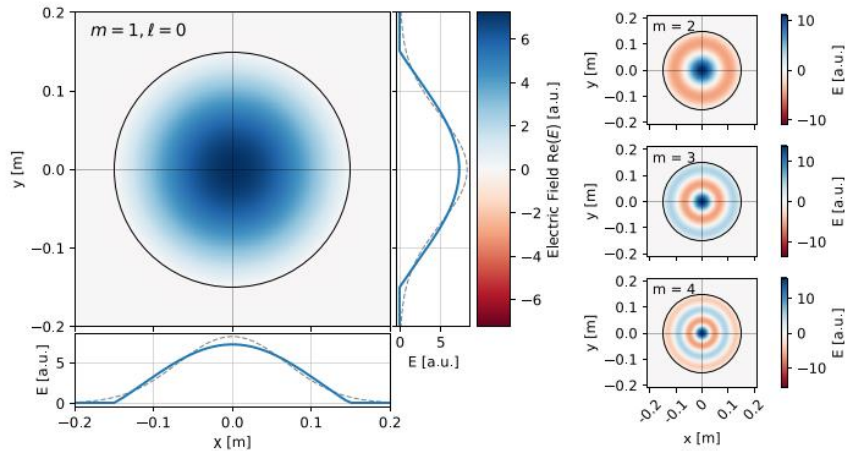
Options for broadband and narrowband scans

- Large single volume
- Approach QCD sensitivity: $\beta^2 > 2 \cdot 10^3$ possible
- Frequency tuning: Disk spacings control $\beta(f)$

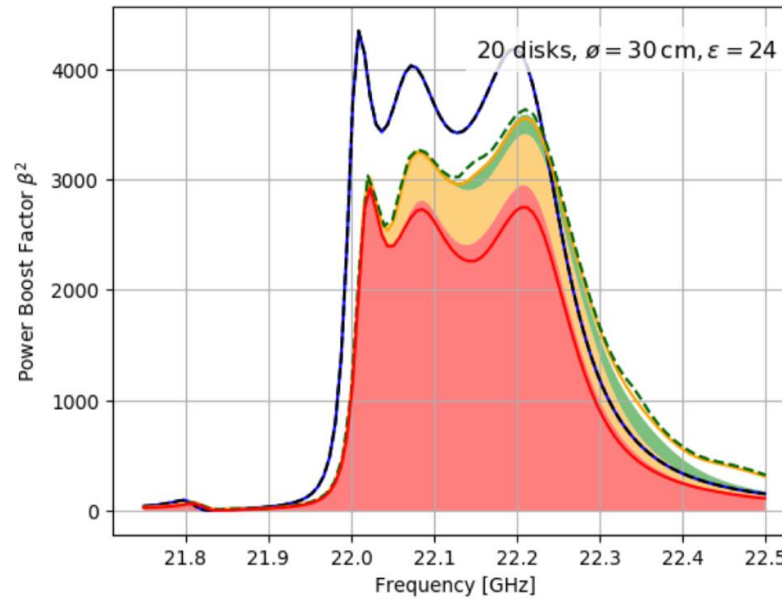


Frequency can be tuned by changing the disc positions

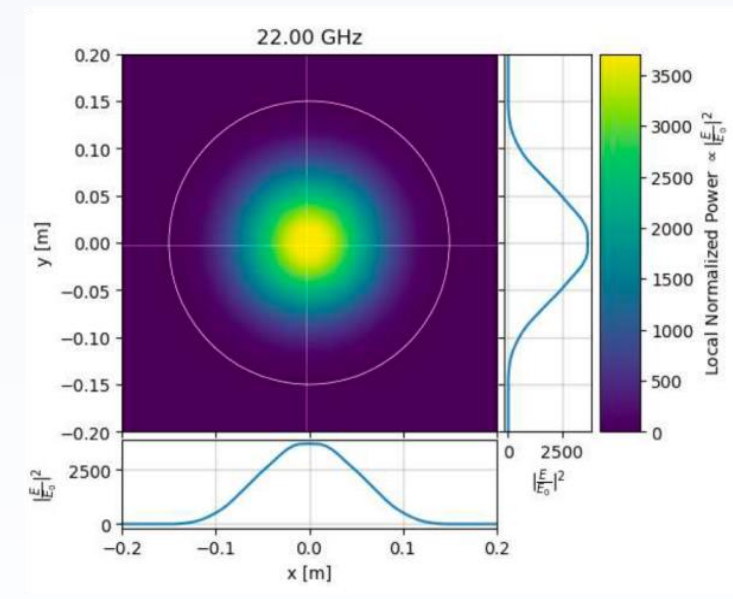
Ideal Booster – simulations



Modes of the circular dielectric haloscope.



Power boost factor β^2 considering dielectric disks of finite size with spacings tuned to cover a bandwidth of ≈ 50 MHz



EXPERIMENTAL CHALLENGES

1. estimating the boost factor in Open Booster

- ▶ The boost factor is not a measurable quantity (?)

2. *3D effects can have significant impact on the boost factor*

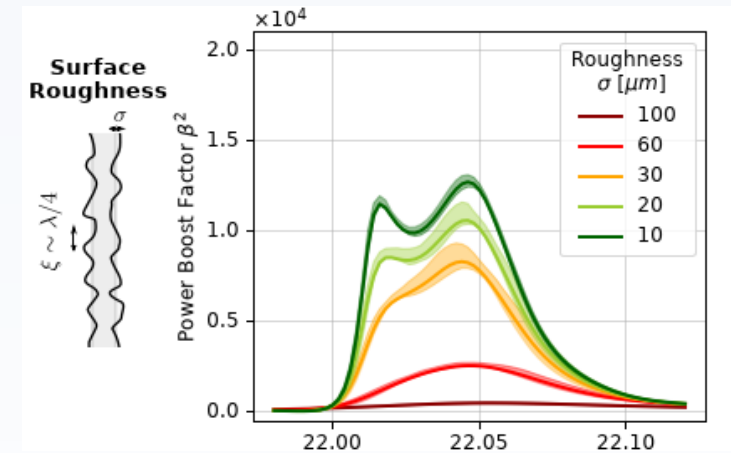
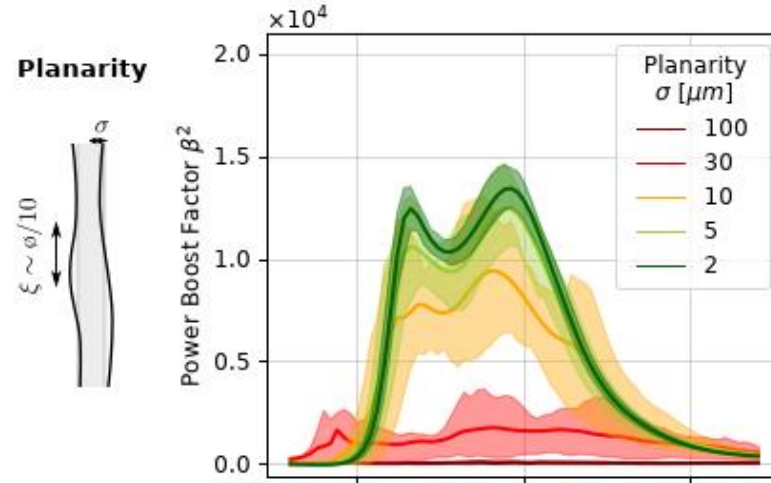
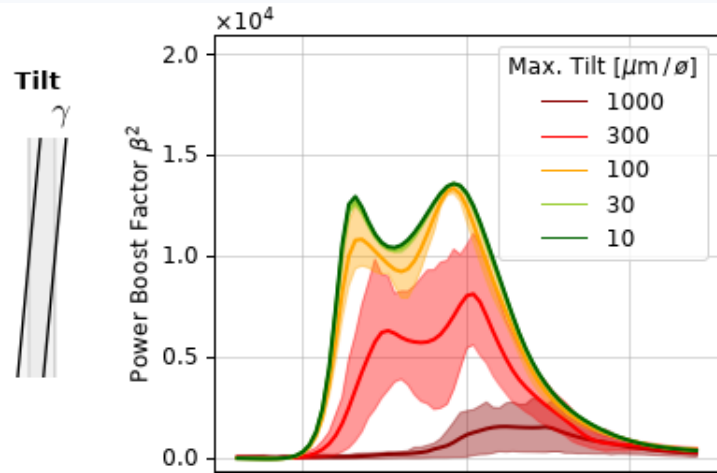
3. RF receiver chain to detect a signal of 10^{-22} W at 4 K

4. Engineering challenges

- ▶ E.g. high field (~ 9 T), large bore ($\sim \text{m}^2$) magnet

BOOST FACTOR β^2 : 3D EFFECTS

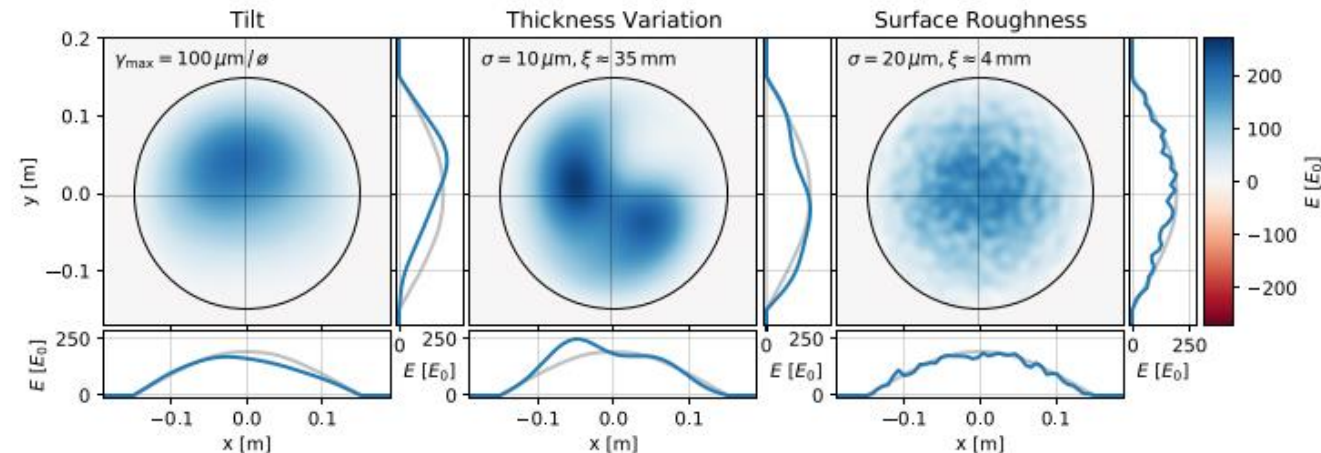
MADMAX Prototype 20 Disks, $\varnothing = 30$ cm



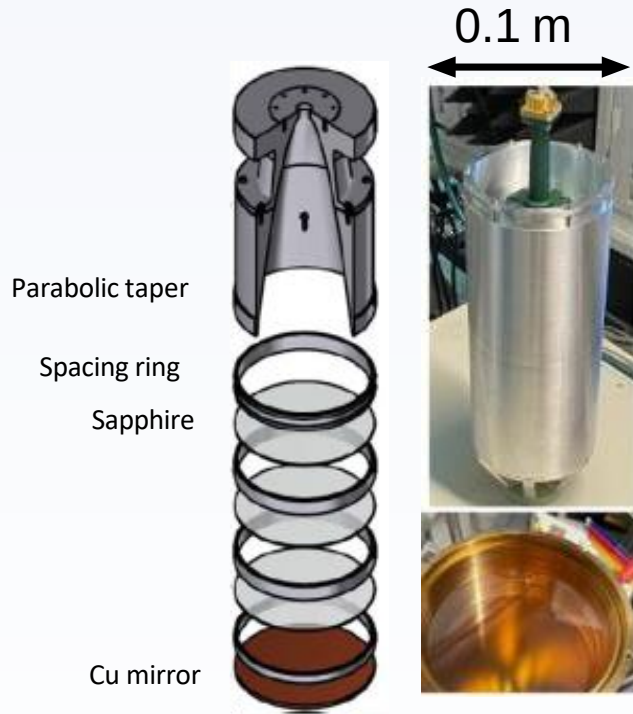
3D effects such as disc tilting, surface roughness and axion velocity have been studied.

Requirements :

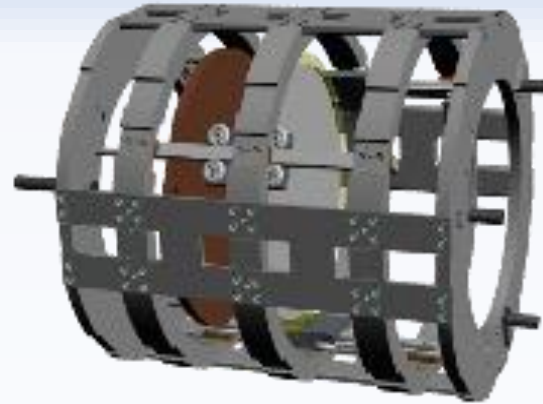
(tiling < 0.1 miliradian, surface roughness < 10 μm)



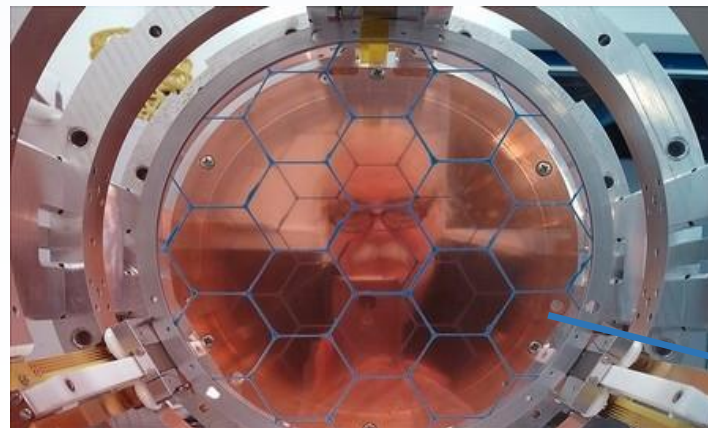
Open and closed Boosters



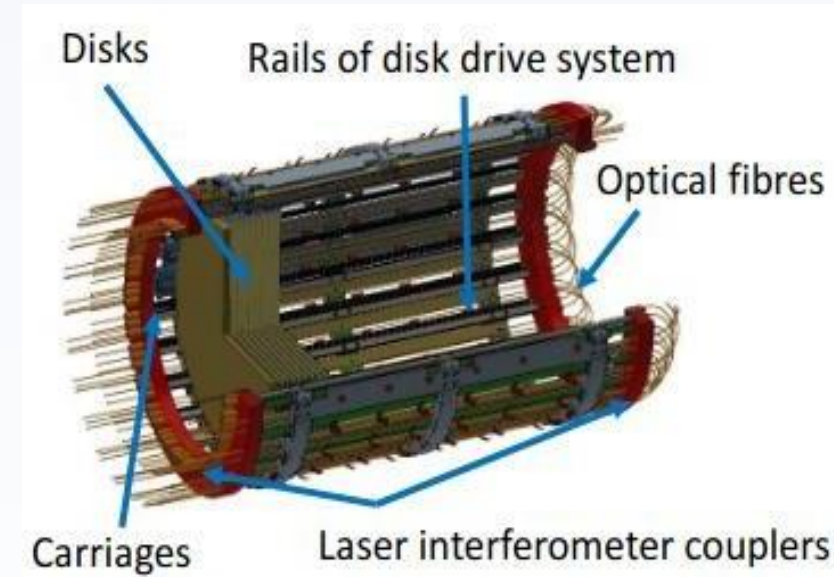
Closed Boosters (CB):
 $\varnothing = 100$ mm (CB100), 3 Al_2O_3 disks
 $\varnothing = 200$ mm (CB200), 3 Al_2O_3 disks



Open Boosters (OB):
 $\varnothing = 200$ mm (OB200), 2 Al_2O_3 disks
 $\varnothing = 300$ mm (OB300), 3 disks (Al_2O_3 & LaAlO_3)



Prototype booster



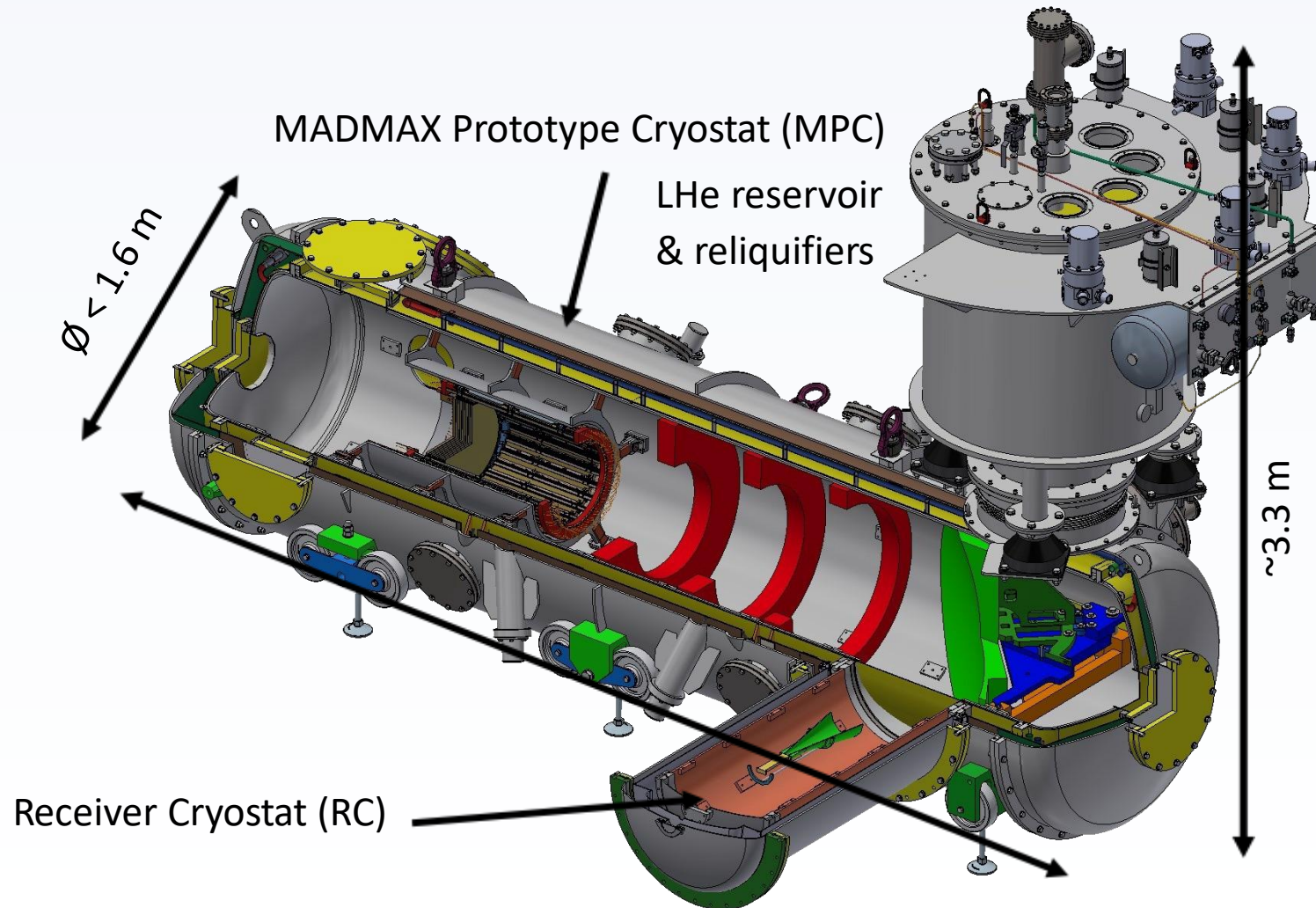
Operating conditions:

- Cryogenic temperatures: 4 K
- High magnetic field: up to ~ 10 T
- Vacuum or cold He exchange gas

Single Sapphire $\varnothing 300$ mm discs
or

Tiled discs made of LaAlO_3 3" wafers

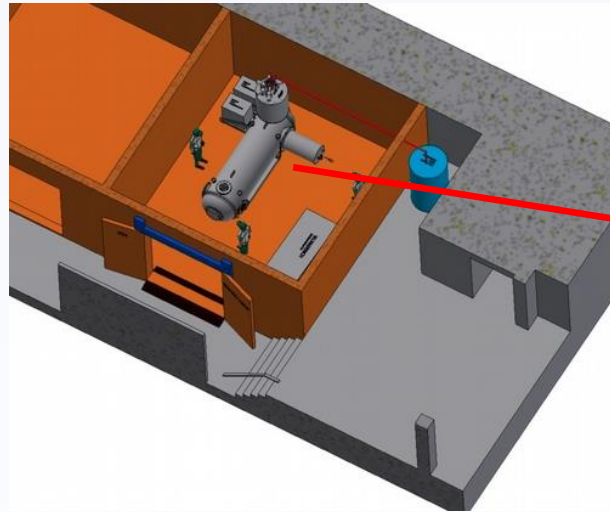
MADMAX Prototype Cryostat



Large bore ($\varnothing = 760 \text{ mm}$) cryostat allows operation of all prototypes
Fits into the 1600 mm warm bore of **MORPURGO** magnet at CERN

- Delivery expected beginning of 2024
- Commissioning with OB300 in Hamburg

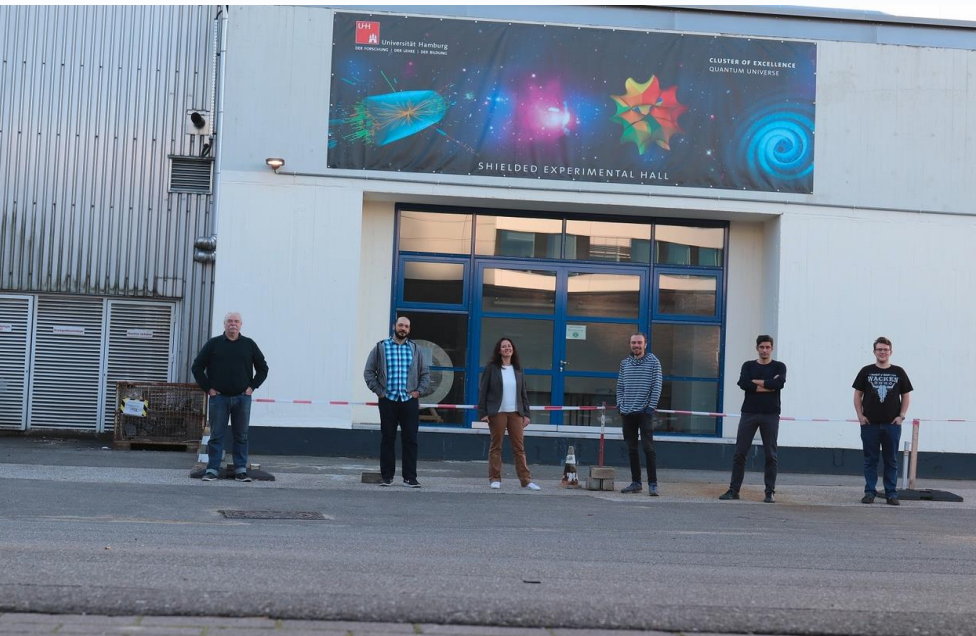
Use of the MORPURGO magnet @ CERN



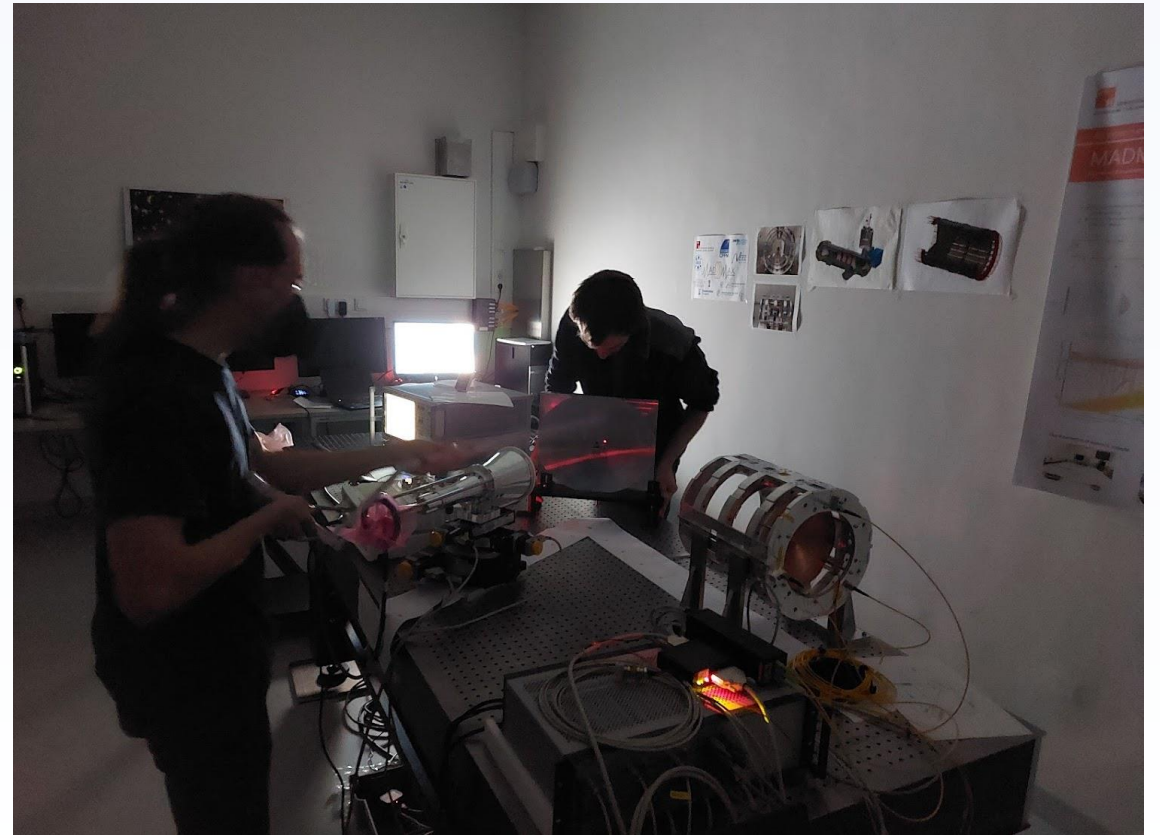
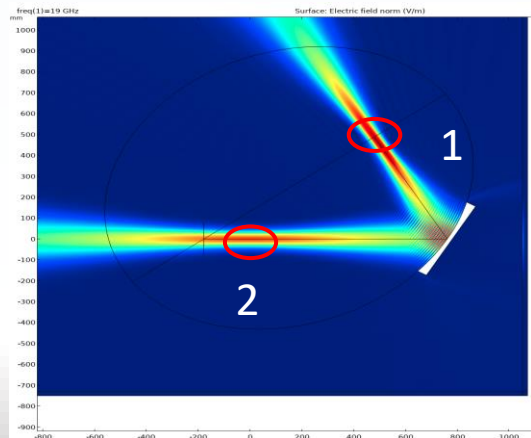
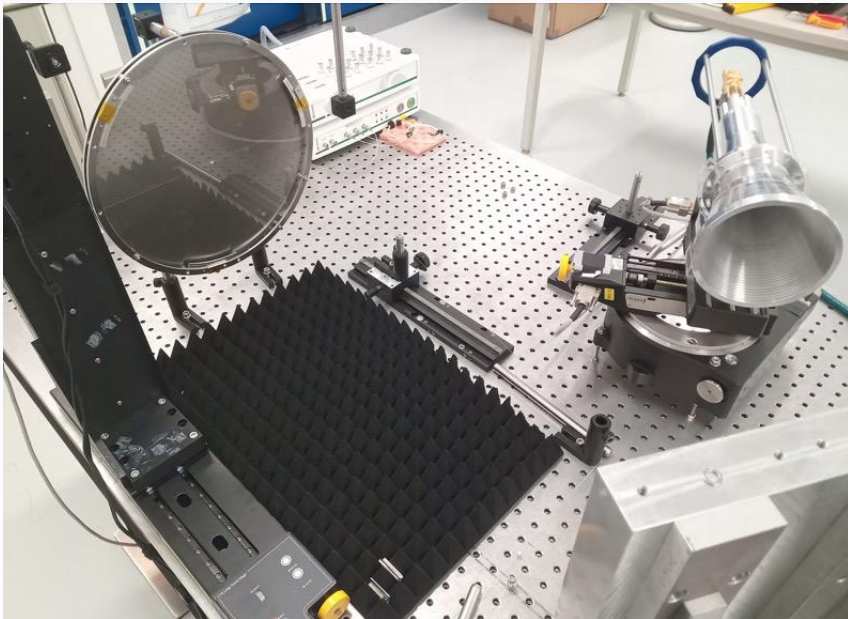
MORPURGO magnet @ CERN

Booster in cryostat operation and commissioning in SHELL @ UHH

First axion physics measurements in MORPURGO magnet @ CERN



SHELL RF measurements



OB 200 (booster) alignment

Signal Power

Change in
reflection
coefficient

$$\Delta\Gamma = \frac{\alpha_e \omega}{P_{\text{in}}} \mathbf{E}_R^2$$

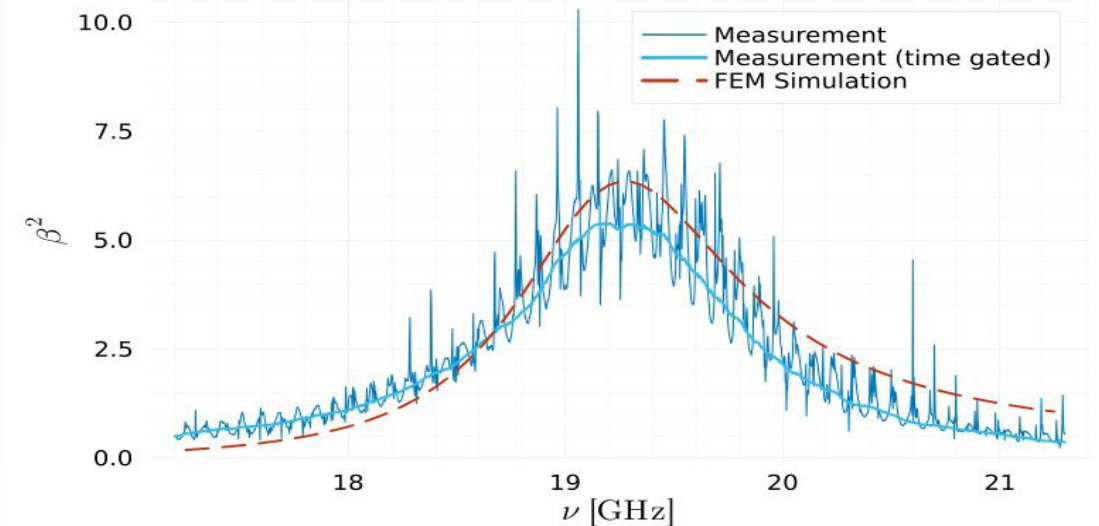
Electric
field

$$\beta^2 = \frac{P_{\text{sig}}}{P_0}$$

$$P_{\text{sig}} = \frac{g_{a\gamma}^2}{16P_{\text{in}}} \left| \int_{V_a} dV \mathbf{E}_R \cdot \dot{\mathbf{A}} \mathbf{B}_e \right|^2$$



- Measure max. E-field between disk and mirror
- Calculate signal power
- Includes effects currently not simulated:
 - Antenna coupling
 - Transverse field perturbations

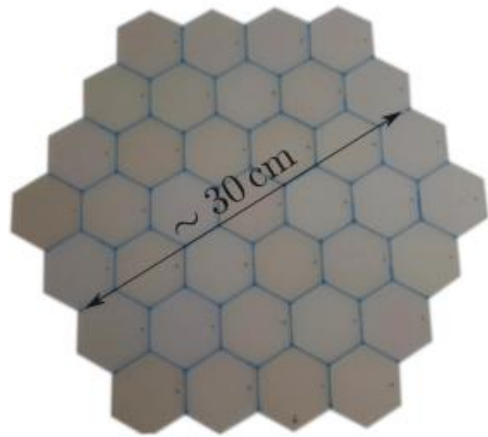


Single disk “low” boost factor

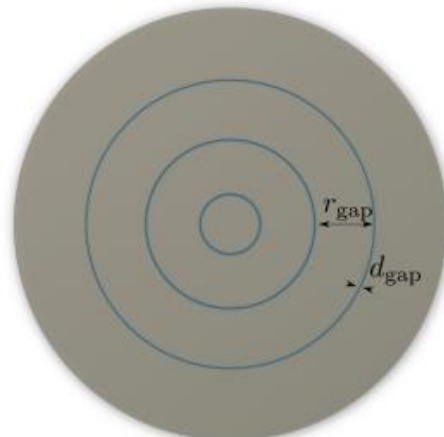
Disk tiling

Tiled Dielectric Disks

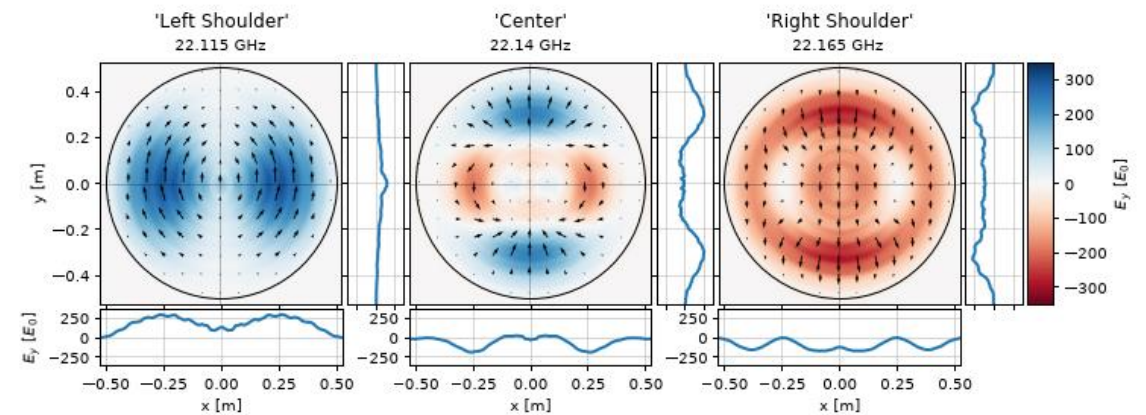
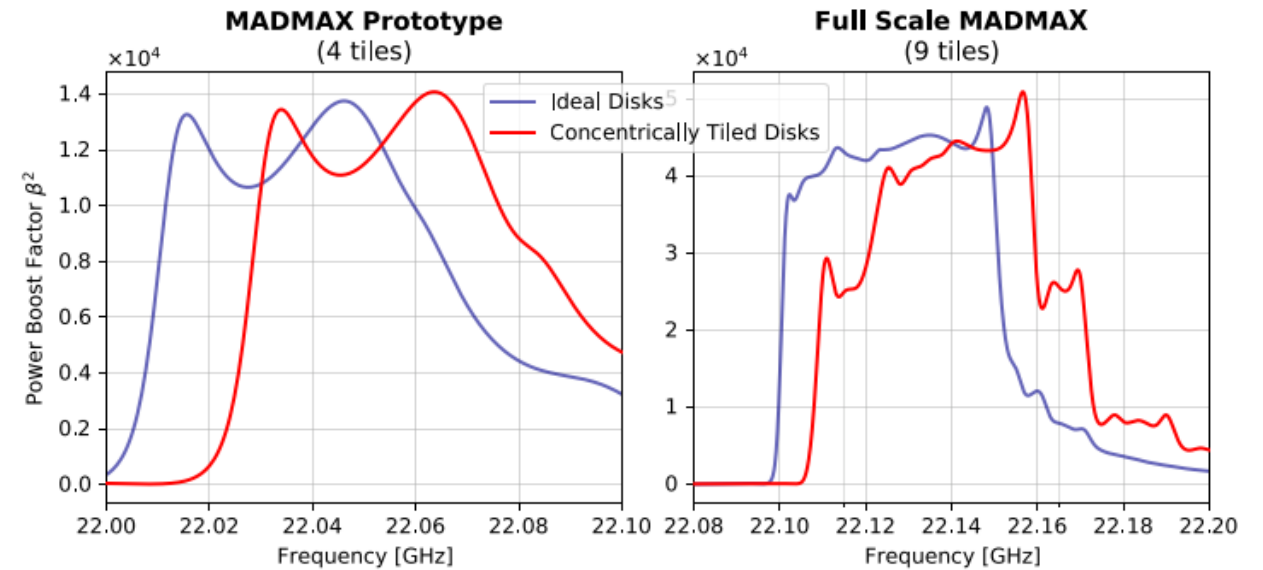
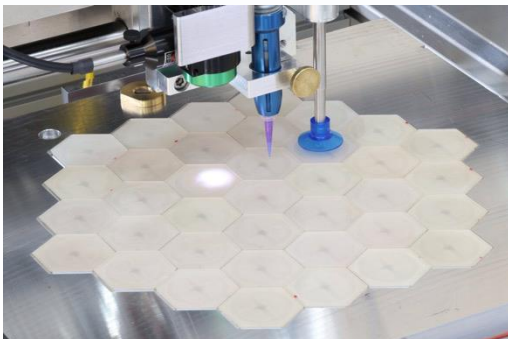
Currently Available Disk



Simulated Disk



- Tiling of $\varnothing = 300$ mm LaAlO_3 disks

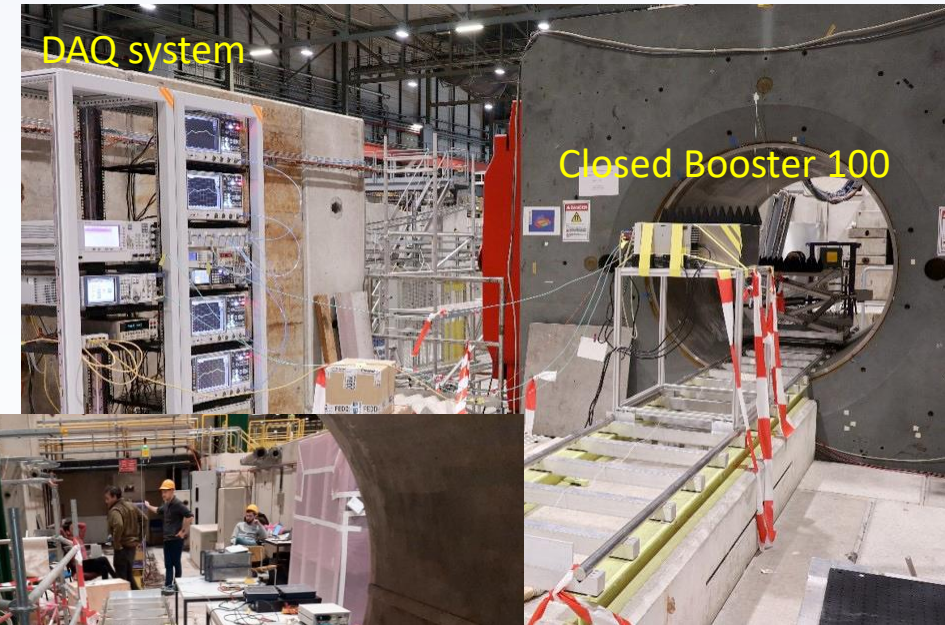
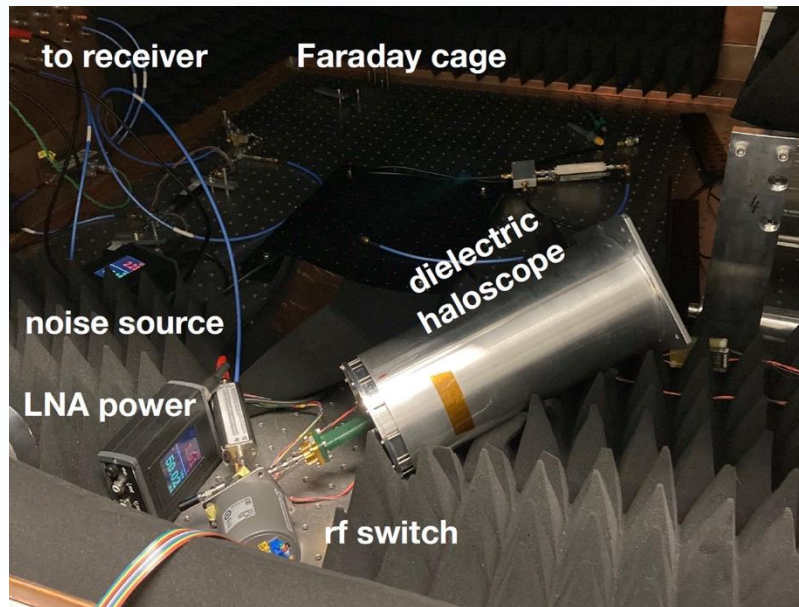


for the full scale MADMAX the beam shape is significantly altered due to polarization effects caused by the tiling

CB measurement at CERN

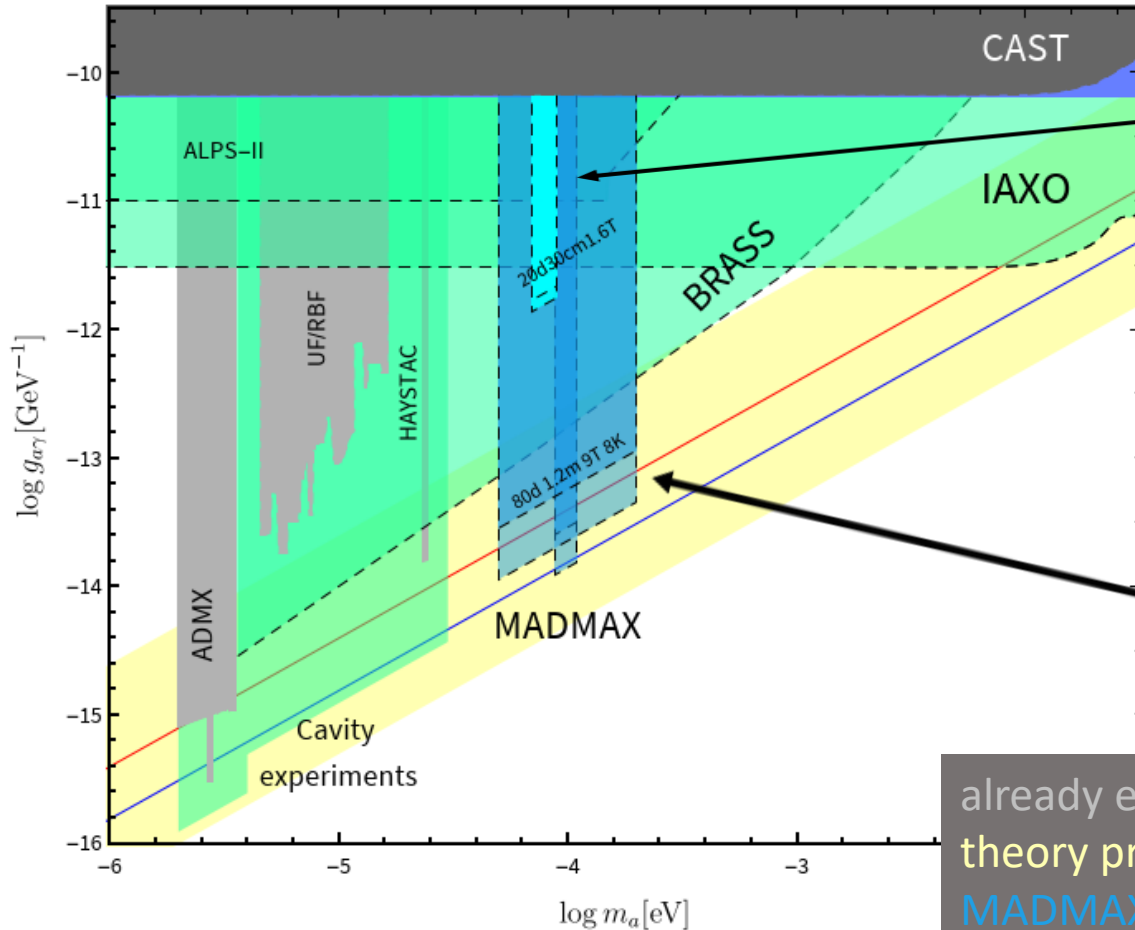
Simple closed system to understand RF behaviour

- Receiver
- Parabolic taper
- 3x Ø100 mm disks (fixed distances)
- Copper mirror



Tested in
MORPURGO
@ CERN
in 2022 and 2023

MADMAX limits



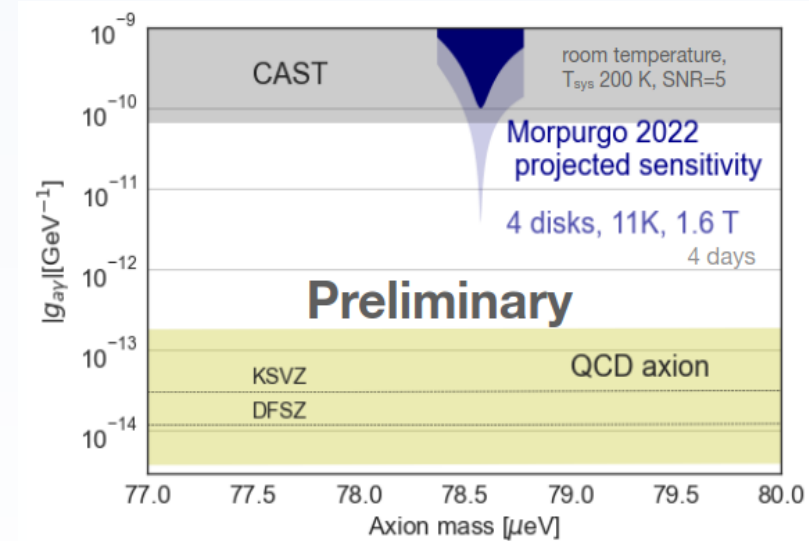
MADMAX Prototype:

$$\begin{aligned} N_{\text{disk}} &= 20 \\ A_{\text{disk}} &= 0.07 \text{ m}^2 \\ B_{\parallel} &= 1.6 \text{ T} \\ T_{\text{sys}} &= 8 \text{ K} \end{aligned}$$

MADMAX:

$$\begin{aligned} N_{\text{disk}} &= 80 \\ A_{\text{disk}} &= 1.2 \text{ m}^2 \\ B_{\parallel} &= 9 \text{ T} \\ T_{\text{sys}} &= 8 \text{ K} \end{aligned}$$

already explored/excluded
theory prediction
MADMAX (Prototype)
other axion experiments



QUAX (QUest for AXions) @ INFN

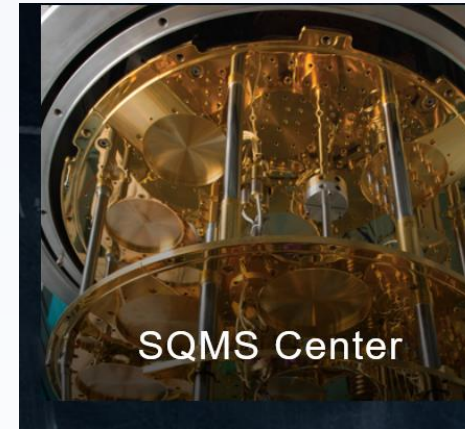


QUAX a- γ

constructing two haloscopes to search for dark matter axions with the standard Sikivie type technique based on the axion photon coupling.

QUAX a-e

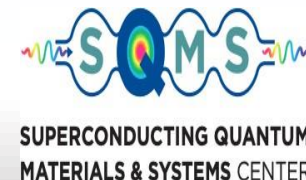
ferrimagnetic haloscope experiment, located at the Laboratori Nazionali di Legnaro, is developing a detector for dark matter axion exploiting the coupling to electrons.



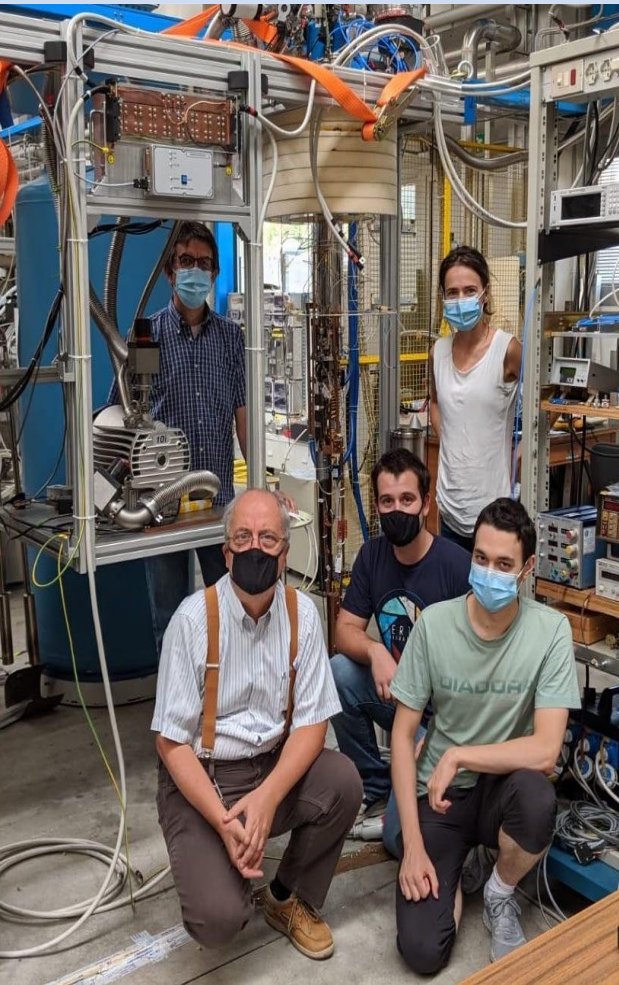
- EU and US collaborations for the integration of:
1. **high-Q cavities** (SQMS, Fermilab)
 2. state-of-the-art **itinerant microwave photon counters** (Quantronics group, Saclay)
 3. **TWPA** (N. Roch group, Neel Institute in Grenoble)



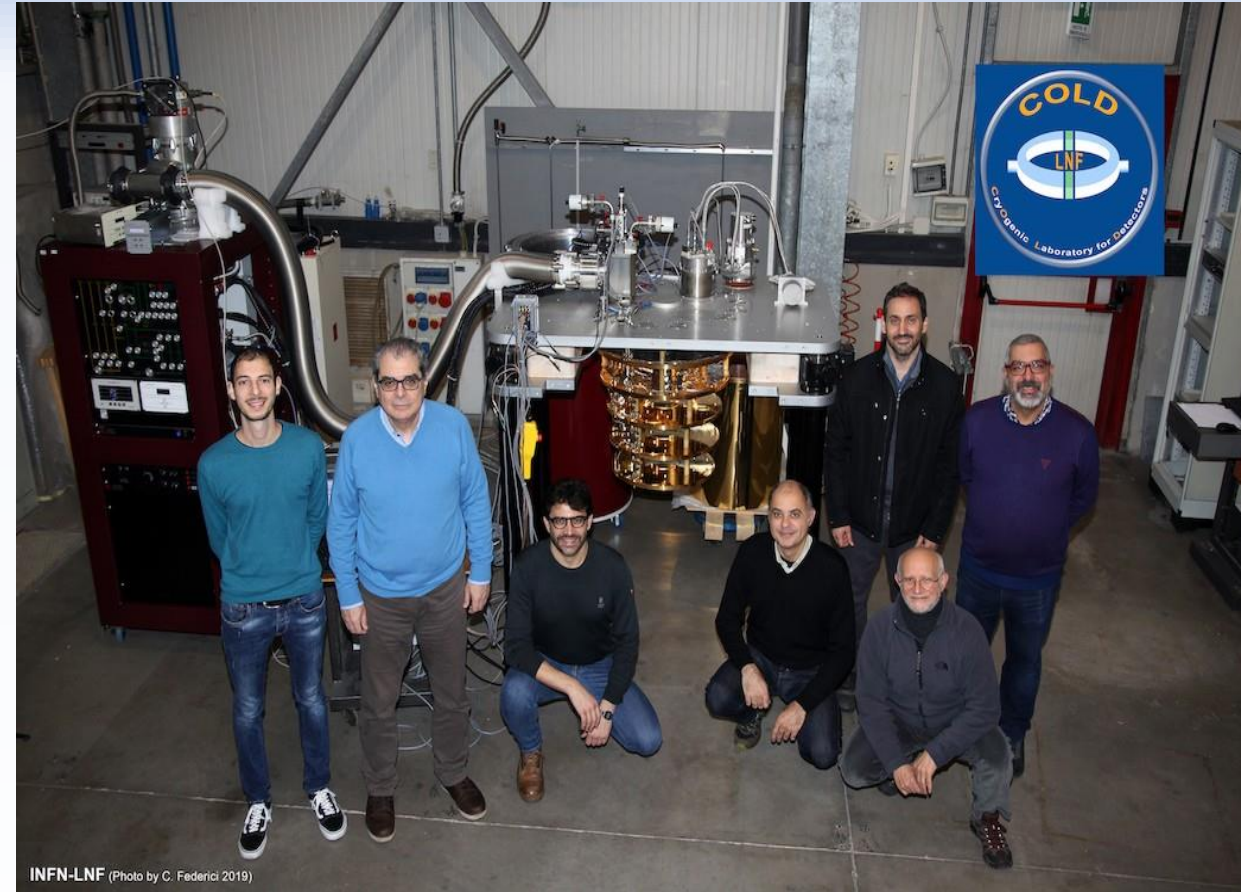
Quantronics Group
Research Group in Quantum
Electronics, CEA-Saclay, France



QUAX Collaboration



LNL

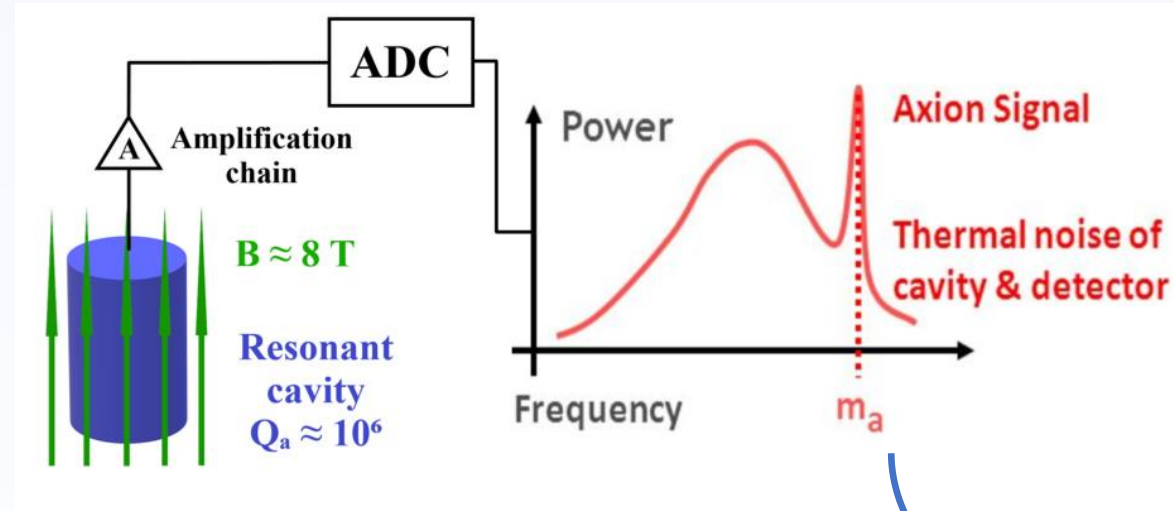


LNF

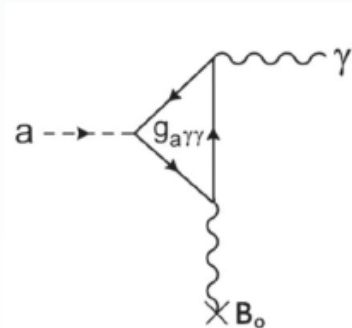


Funded program to run an axion observatory with two Haloscopes :
One in **LNL** & one in **LNF**

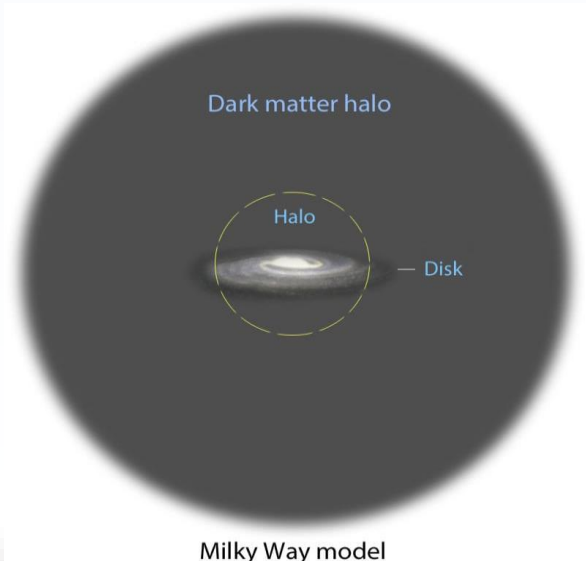
Haloscope detection principle



$$P_{\alpha \rightarrow \gamma} = \left(g_{\alpha \gamma \gamma}^2 \frac{\rho_\alpha}{m_\alpha} \right) \left(\frac{\beta}{1 + \beta} B_0^2 V C Q_L \right)$$



- axions converted to photons inside **intense magnetic field**
- **high-Q** microwave cavity collects and integrates signal
- signal amplified by **low noise receiver chain**, then digitized by ADC
- axion signal **revealed as excess power over noise fluctuations**



Haloscopes in the parameter space

- Excluded parameter space for axions, zoomed into microwave cavity haloscopes.
- Theoretical predictions for post-inflationary axions.
- How to reach the cosmological relevant region?

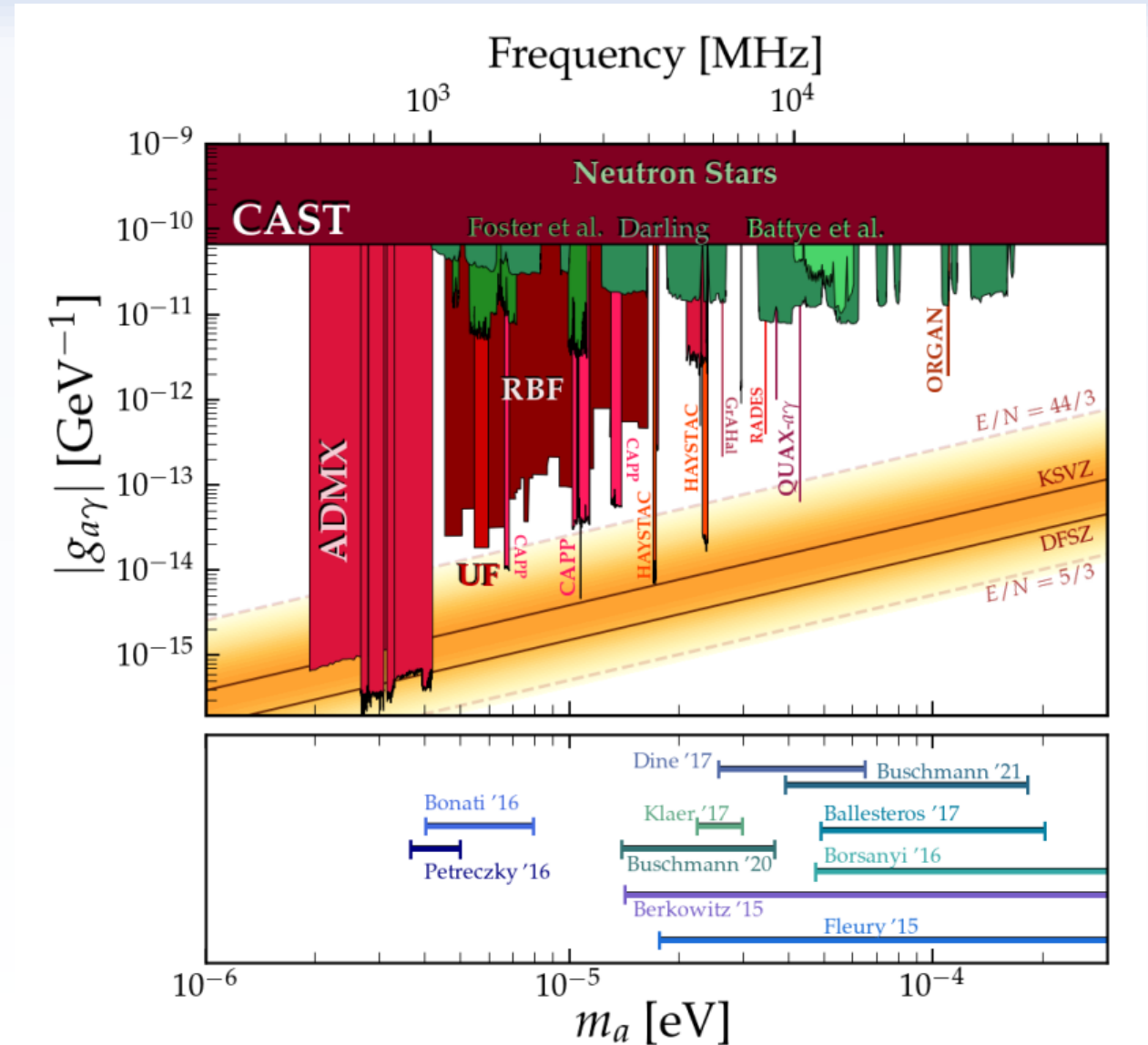
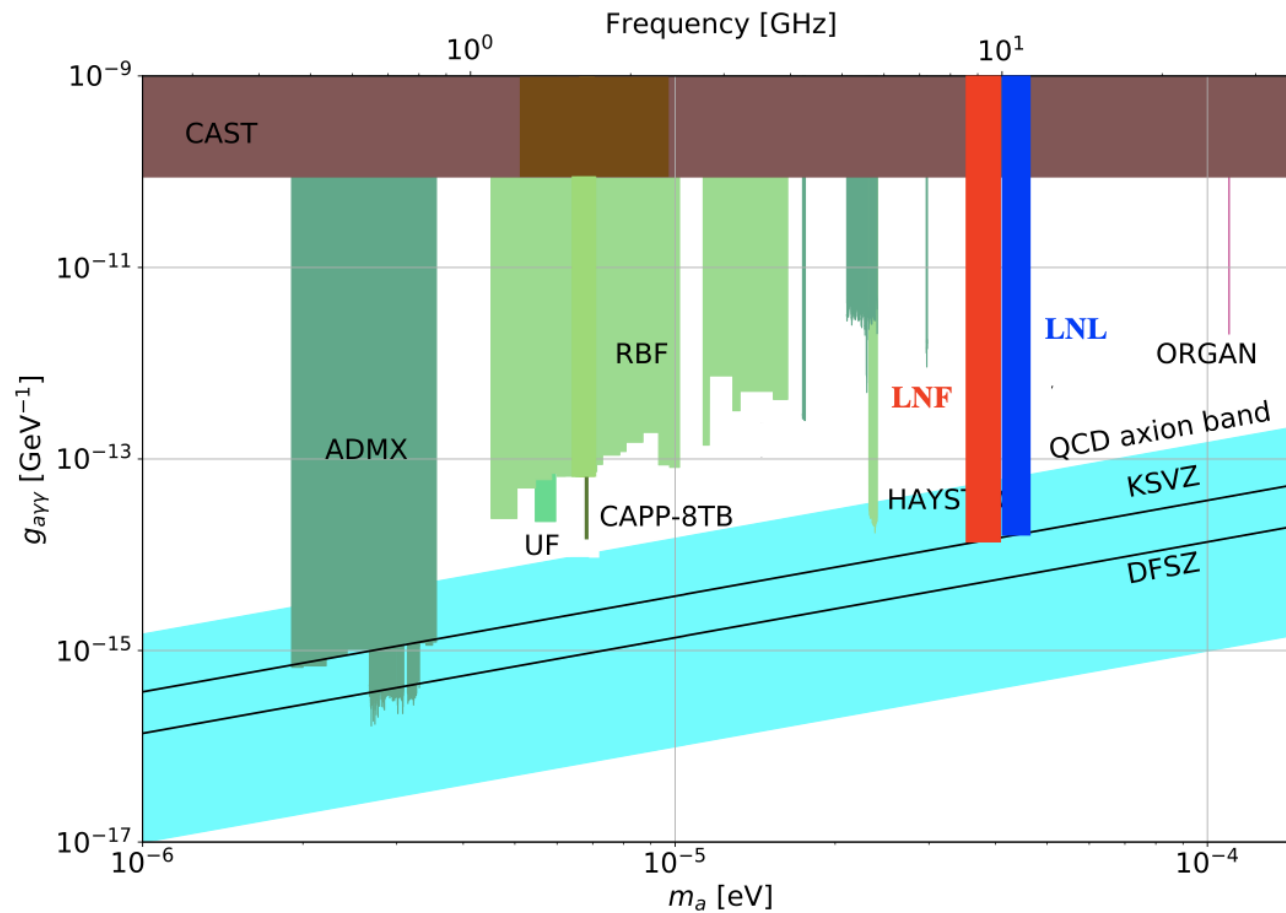


Figure courtesy of Ciaran O'Hare
<https://doi.org/10.5281/zenodo.3932430>

QUAX scientific program



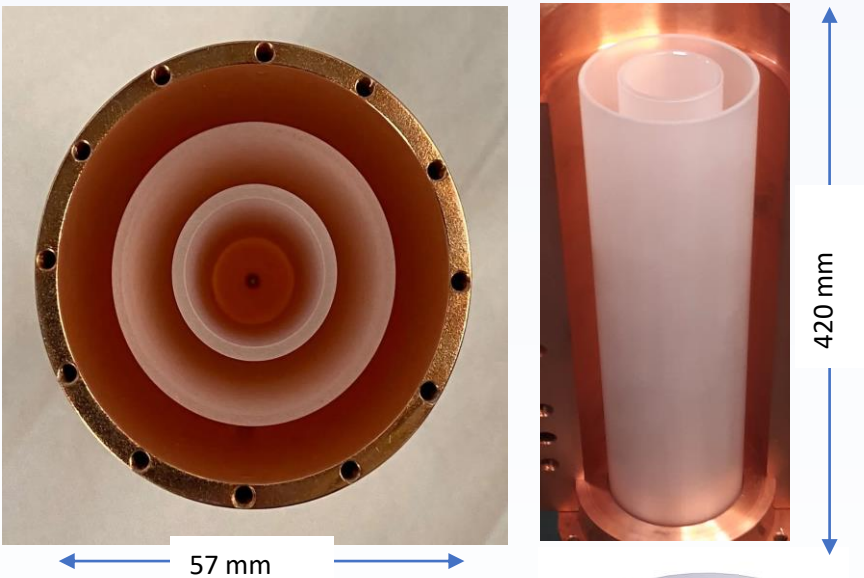
	LNF	LNL
Magnetic field	9 T	14 T
Magnet length	40 cm	50 cm
Magnet inner diameter	9 cm	12 cm
Frequency range	8.5 - 10 GHz	9.5 - 11 GHz
Cavity type	Hybrid SC	Dielectric
Scanning type	Inserted rod	Mobile cylinder
Number of cavities	7	1
Cavity length	0.3 m	0.4 m
Cavity diameter	25.5 mm	58 mm
Cavity mode	TM010	pseudoTM030
Single volume	$1.5 \cdot 10^{-4} \text{ m}^3$	$1.5 \cdot 10^{-4} \text{ m}^3$
Total volume	$7 \otimes 0.15$ liters	0.15 liters
Q_0	300 000	1 000 000
Single scan bandwidth	630 kHz	30 kHz
Axion power	$7 \otimes 1.2 \cdot 10^{-23} \text{ W}$	$0.99 \cdot 10^{-22} \text{ W}$
Preamplifier	TWJPA/INRIM	DJJAA/Grenoble
Operating temperature	30 mK	30 mK

Outline

- ❑ **High Q** dielectric resonators / novel tuning mechanisms and geometries
- ❑ **TWPA** – based amplification chain characterization and measurements
- ❑ **Single Microwave Photon Detectors (SMPD)** for "itinerant " photons

Dielectric resonators

Tunable cavity with dielectric shells
“dielectric boosted” resonator

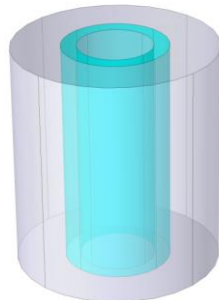


TM₀₃₀ higher order mode

for axion detection

2 shells → $C_{030} = 0.03$

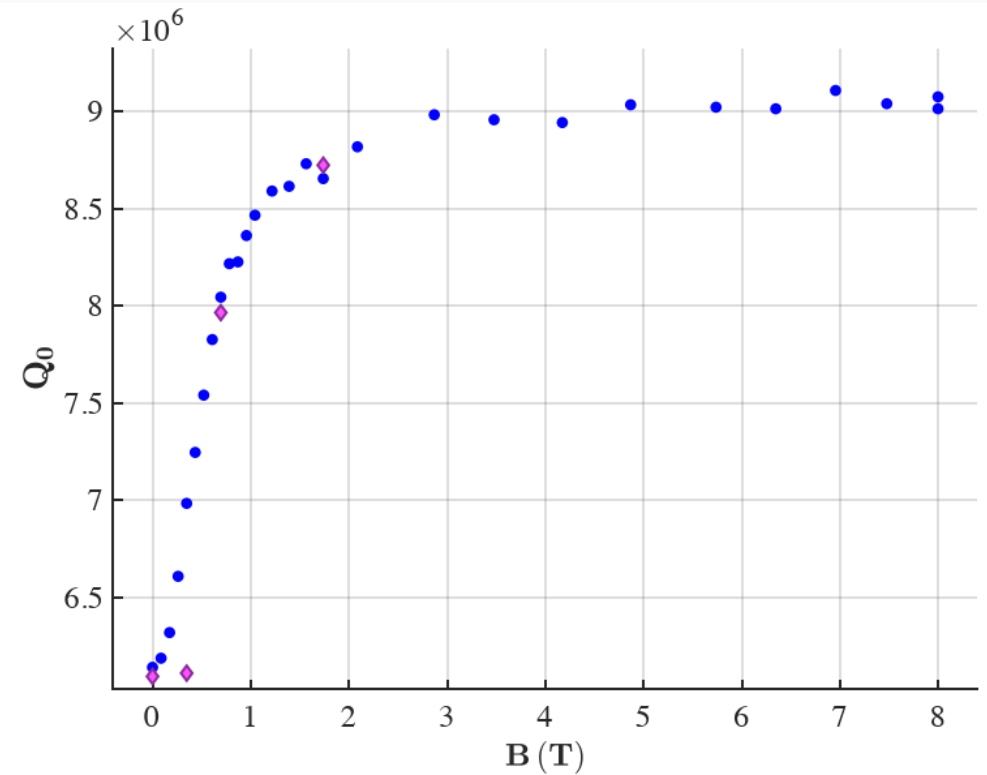
1 shell → $C_{030} = 0.49$



[Phys. Rev. Appl. 17, 054013 \(2022\)](#)

[Nucl. Instrum. Methods A **985**, 164641 \(2021\)](#)

$$Q_0 \sim 9 \cdot 10^6 \text{ @ } 10.3 \text{ GHz} - 4.2 \text{ K}$$



paramagnetic impurities in sapphire :

@ high B-fields they are completely swept away

Wet DU Cryostat

- $\text{He}^3\text{-He}^4$ “wet” dilution refrigerator (refurbished) ! recovery system + compressor at LNL
- Cooling power of 1 mW at 120 mK



8 T field magnet



base temperature of 50 mK after ~ 9 hours

Magnet and DU



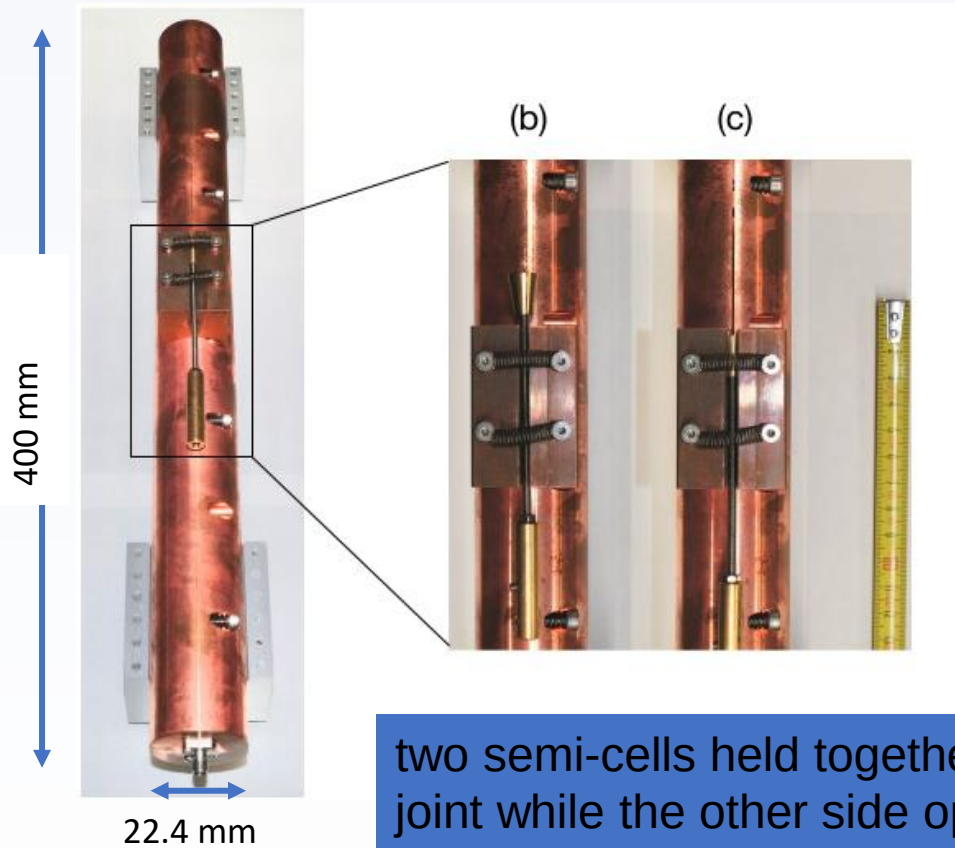
Nb-Ti SC magnet
surrounds the cavity and
provides the 8T magnetic field

Compensation coil shields
the cryogenic electronics
and TWPA
from ~ 0.3 T residual field

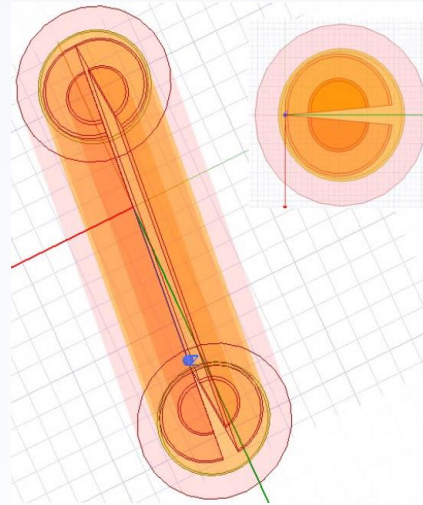


^3He – ^4He wet dilution refrigerator: refurbished unit previously
installed in the gravitational wave bar antenna Auriga test facility

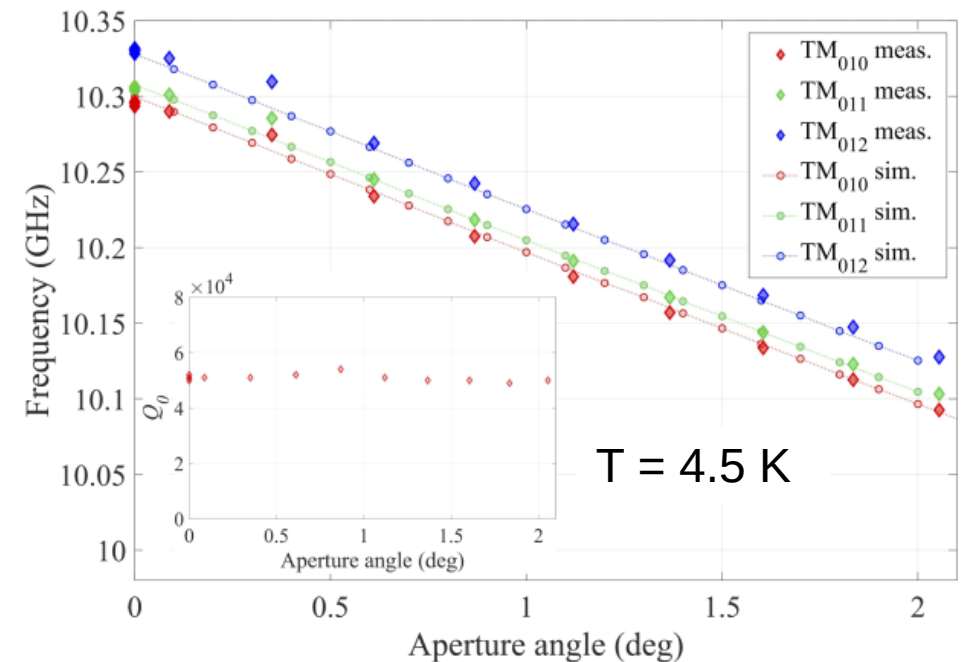
Tuning mechanisms



two semi-cells held together at a fixed joint while the other side opens to tune the frequency of the resonant mode



$V = 0.158$ liters
 $Q_0 \approx 50.000$
 freq. ~ 10 GHz
 $T = 4$ K



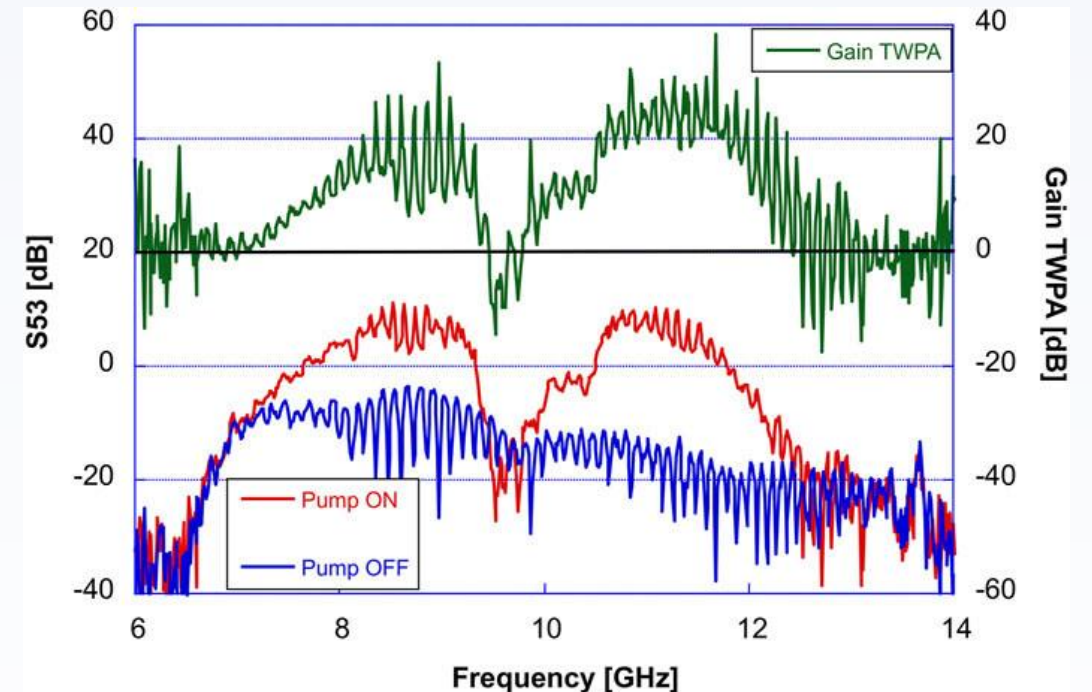
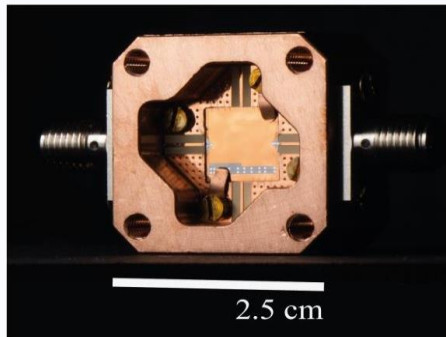
maximum aperture of ~ 2 degrees, is equivalent to a 2% increase in effective volume

Tuning a range of at least 200 MHz for the **fundamental mode TM₀₁₀**

TWPA characterization

In axion DM research we need:

- quantum-limited noise performance
- GHz amplification bandwidth



$$T_{\text{sys}} = 2.1 \pm 0.1 \text{ K @ } 10.353 \text{ GHz}$$

Novel, reliable calibration scheme to measure T_{sys} exactly at the cavity output and without the need for switches nor heated load.

Review of
Scientific Instruments

ARTICLE

scitation.org/journal/rsi

A haloscope amplification chain based on a traveling wave parametric amplifier

Cite as: Rev. Sci. Instrum. 93, 094701 (2022); doi: [10.1063/5.0098039](https://doi.org/10.1063/5.0098039)

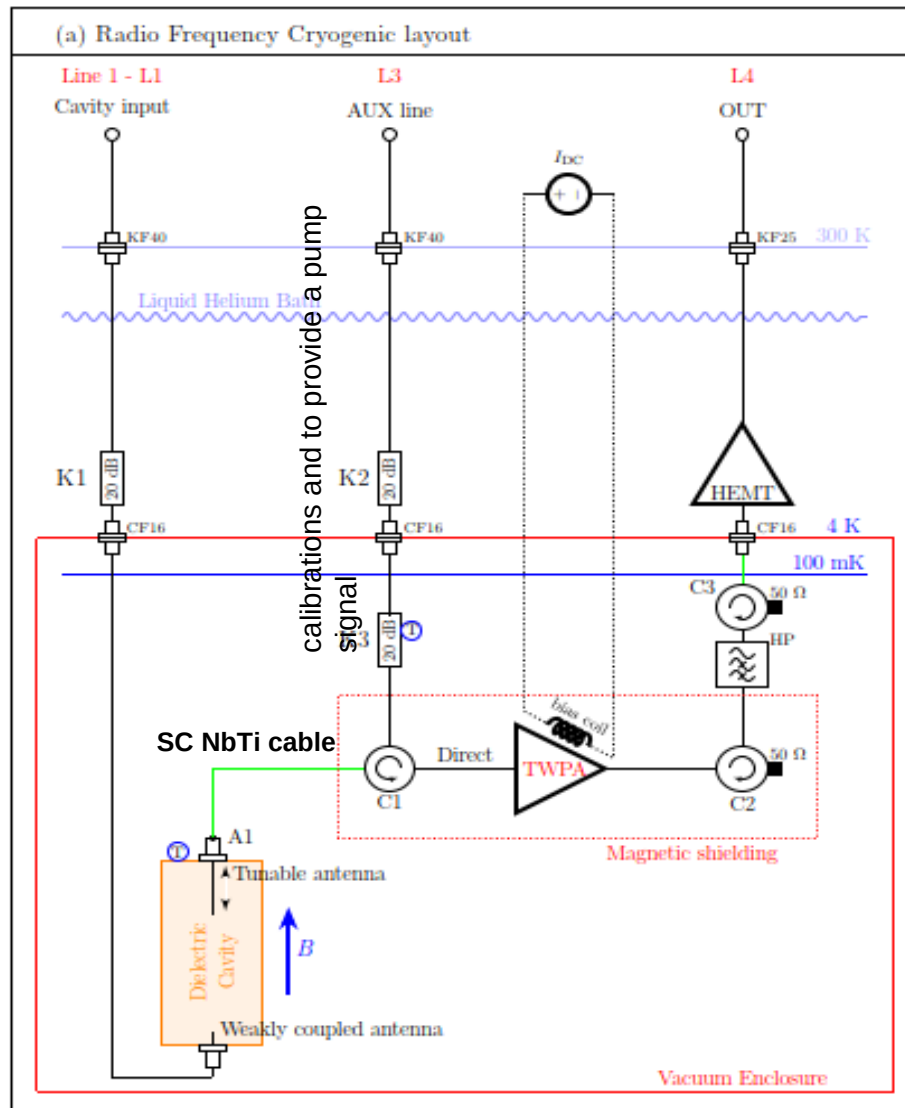
Submitted: 4 May 2022 • Accepted: 28 July 2022 •

Published Online: 6 September 2022



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Schematics of the experimental apparatus



TM030 mode

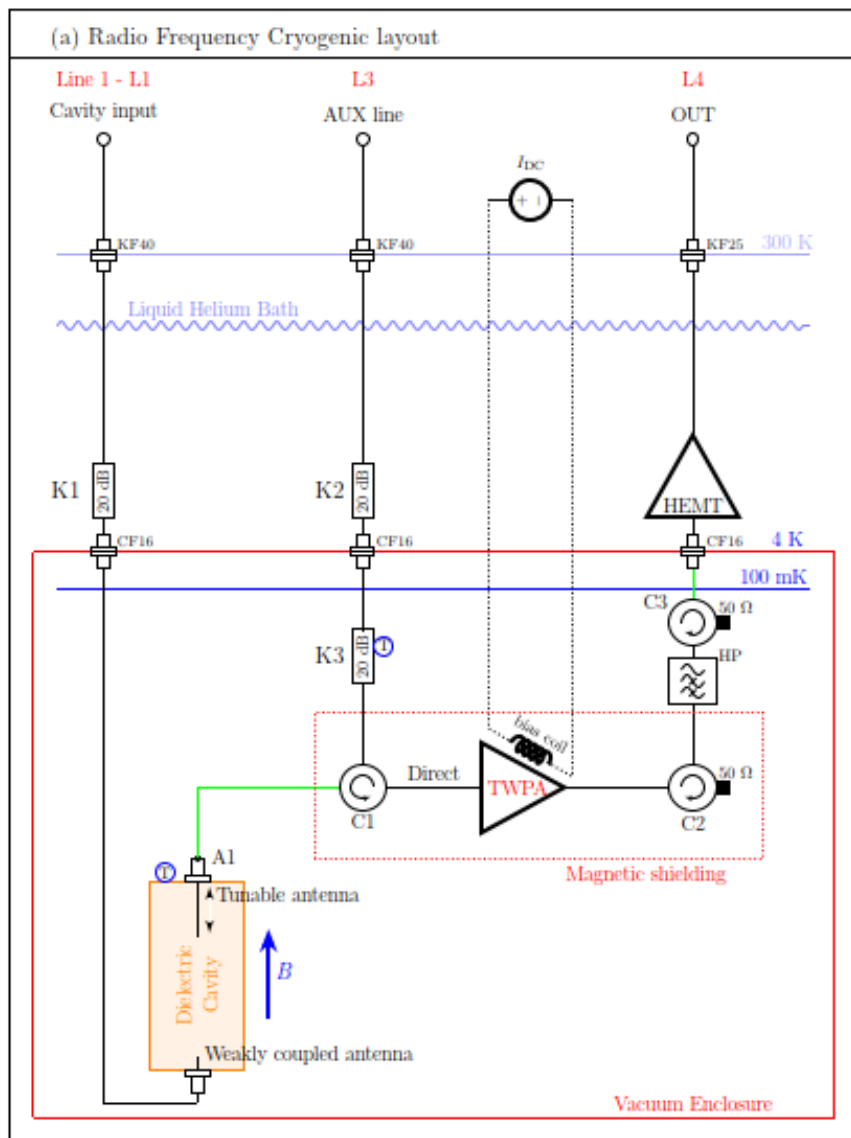
Effective volume $3.4 \times 10^{-2} \text{ lt}$

Freq. = 10.353 GHz

Tuned by 3 2-mm sapphire rods

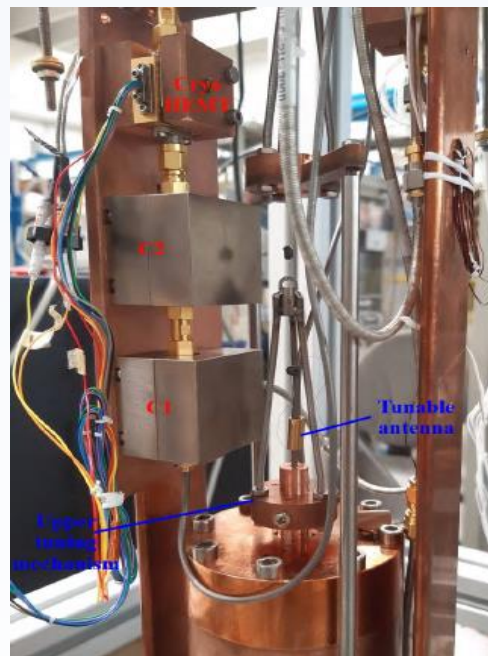
- pre-amplifier TWPA
- HEMT amplifier
- C1, C2 and C3 are circulators,
- HP is an 8 GHz high pass filter
- K1, K2, and K3 are attenuators

high-Q microwave resonant cavity:
internal quality factor of $\sim 9 \times 10^6$



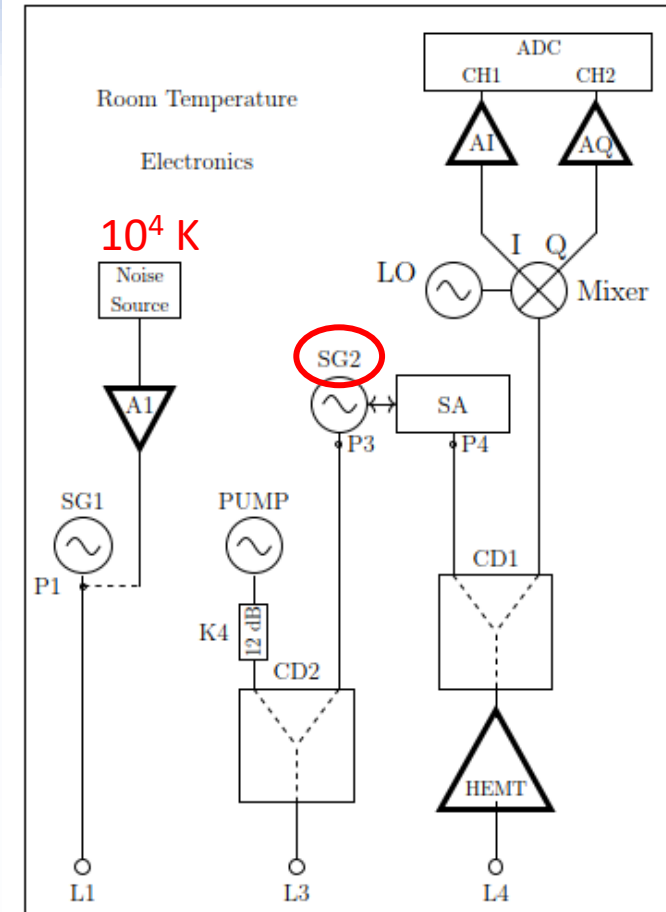
high-Q microwave resonant cavity:
internal quality factor of more than 9×10^6

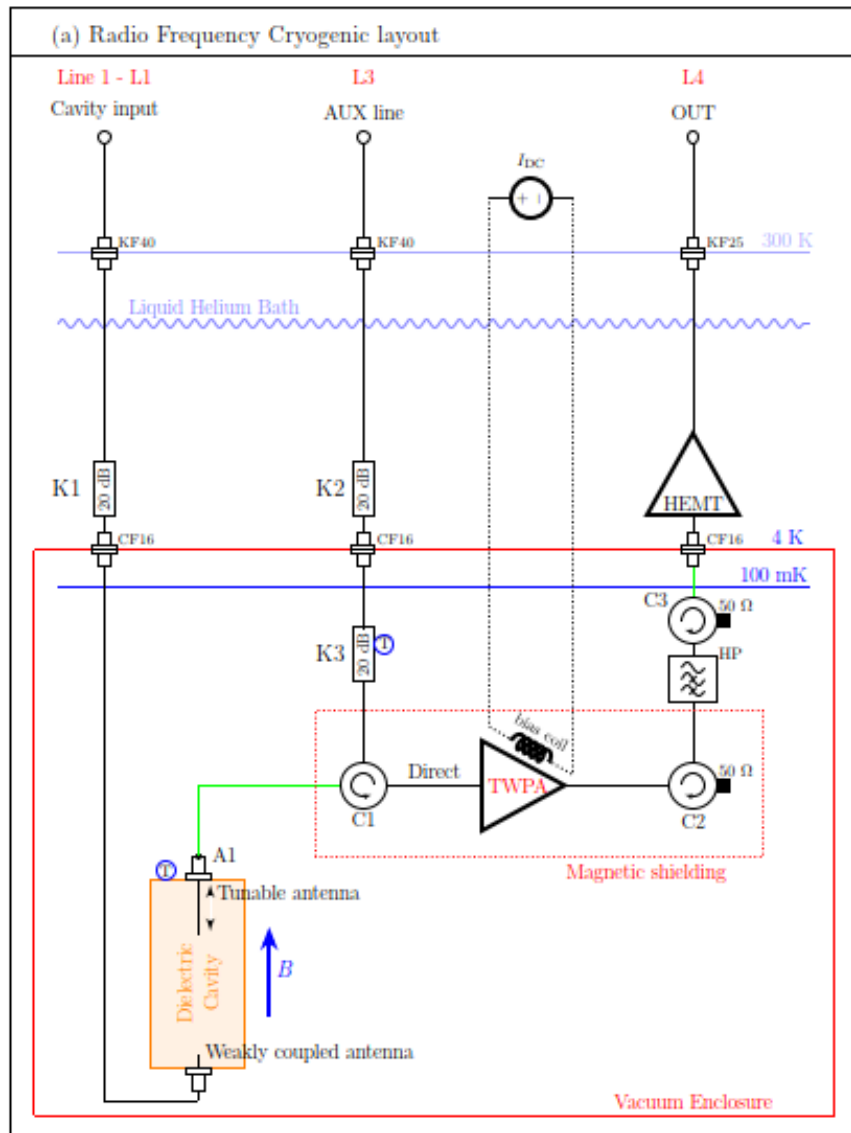
Tuning by 3 2-mm sapphire rods



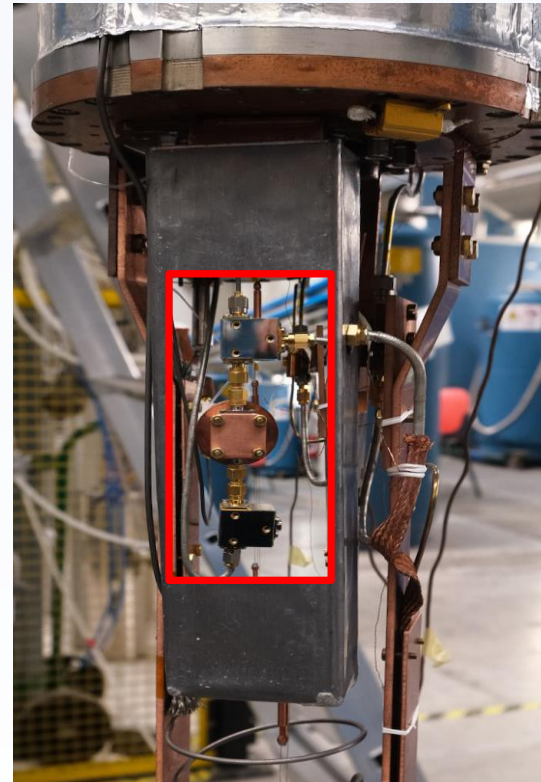
-cavity reflection spectrum S43
(input from line L3 - output from line L4)

-information on the loaded quality
 Q_L , resonance frequency f_c and coupling β of the
tunable antenna

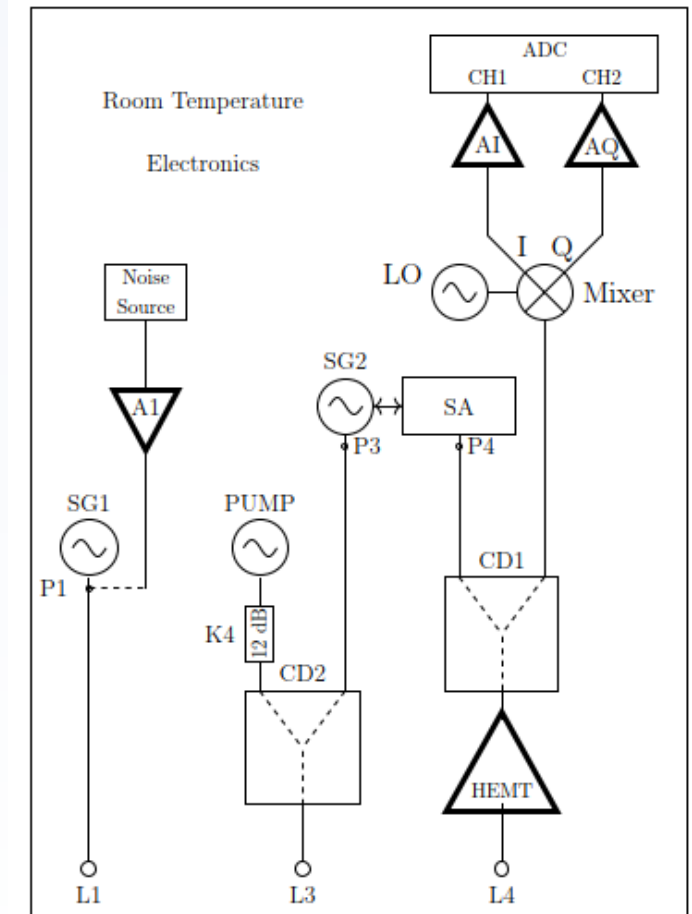




TWPA with circulators inside the shielding box



- Mixing chamber @ 50 mK
- Cavity @ 110 mK



high-Q microwave resonant cavity:
internal quality factor of more than 9×10^6

QUAX_{αγ} TWPA-based amplification measurement

Phys. Rev. D (R. Di Vora et al)

$$N_{\text{sys}} = \frac{k_B}{h\nu_s} T_{\text{sys}} = 4.2 \pm 0.3.$$

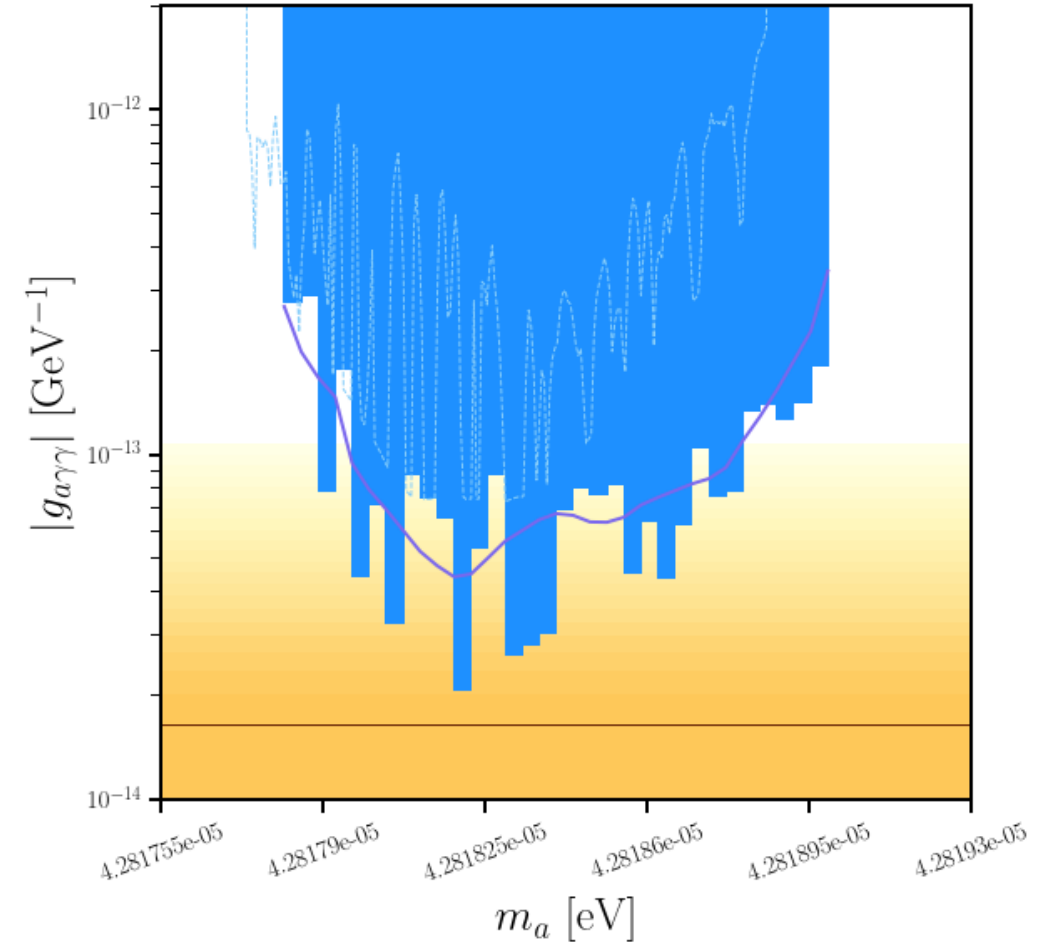
$$T_{\text{sys}} = 2.06 \pm 0.13 \text{ K @ } 10.353 \text{ GHz}$$

Term	Value (K)	N photons
$N(\nu_s, T_s)$	0.27	0.5
$N(\nu_i, T_i)$	0.27	0.7
$N_{\text{HEMT}}/\Lambda_1\Lambda_2G_{\text{TWPA}}$	0.39	0.8
N_{sys}	2.06	4.2
$N_{\text{TWPA}}/\Lambda_1$		2.2

Data taking for 17 hours

- 8 T magnetic field
- $\nu_c = 10.353 \text{ GHz}$
- $Q \sim 2.5 \cdot 10^5$

This is the first time a wide bandwidth quantum limited amplifier has been used in a haloscope working at high frequency



We set a limit for the axion-photon coupling a factor about 4 from the benchmark axion-QCD band

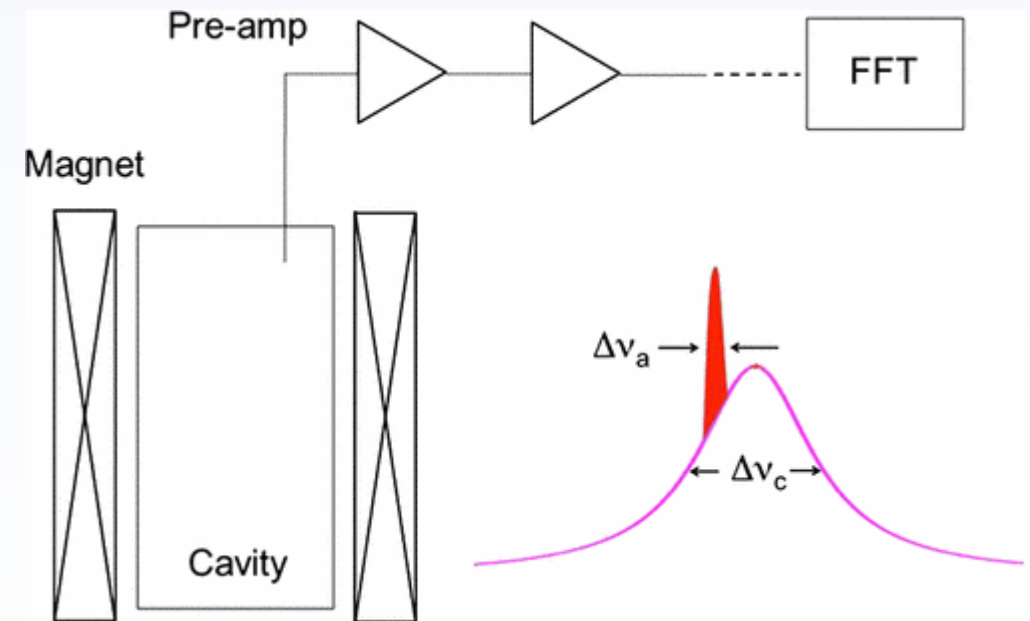
Analysis of single-photon and linear amplifier detectors for microwave cavity dark matter axion searches

at higher frequencies, and thus higher axion masses, single-photon detectors become competitive and ultimately favored, when compared to quantum-limited linear amplifiers, as the detector technology in microwave cavity experimental searches for galactic halo dark matter axions.

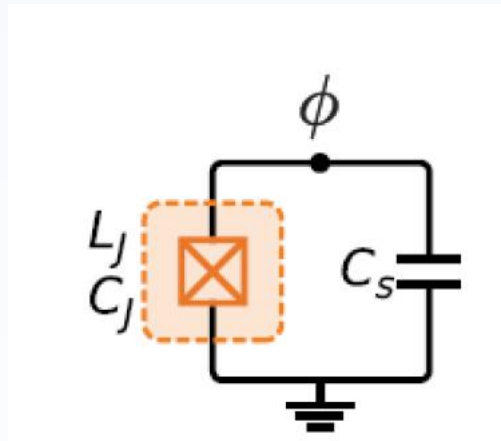
$$\frac{P_\ell}{P_{\text{sp}}} = \left[\sqrt{\bar{n}} + \frac{1}{\sqrt{\bar{n}}} \right] \sqrt{\frac{Q_c}{2\pi\eta Q_a}}.$$

At achievable temperature $T=10$ mK,
 $\bar{n} \approx 3 \times 10^{-11}$ then

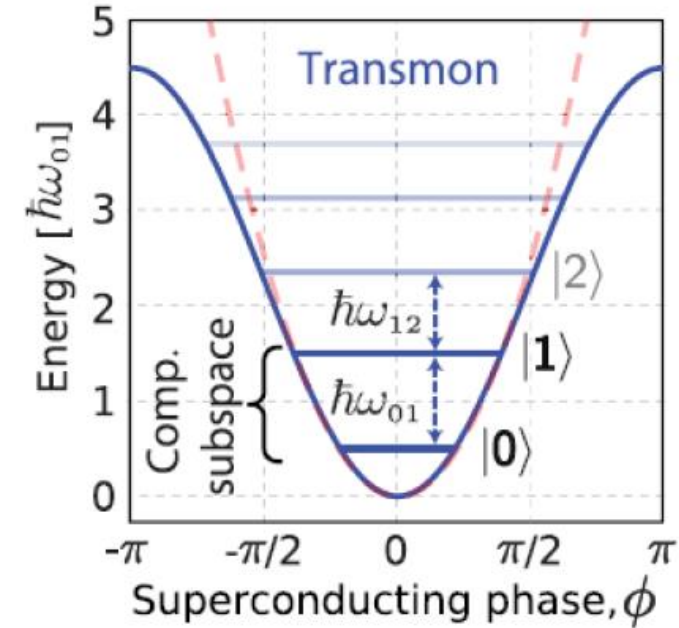
$$P_\ell/P_{\text{sp}} \sim 17,000$$



Single Microwave Photon Detector



Transmon qubit circuit. The orange box represents a Josephson junction with inductance L_J and capacitance C_J

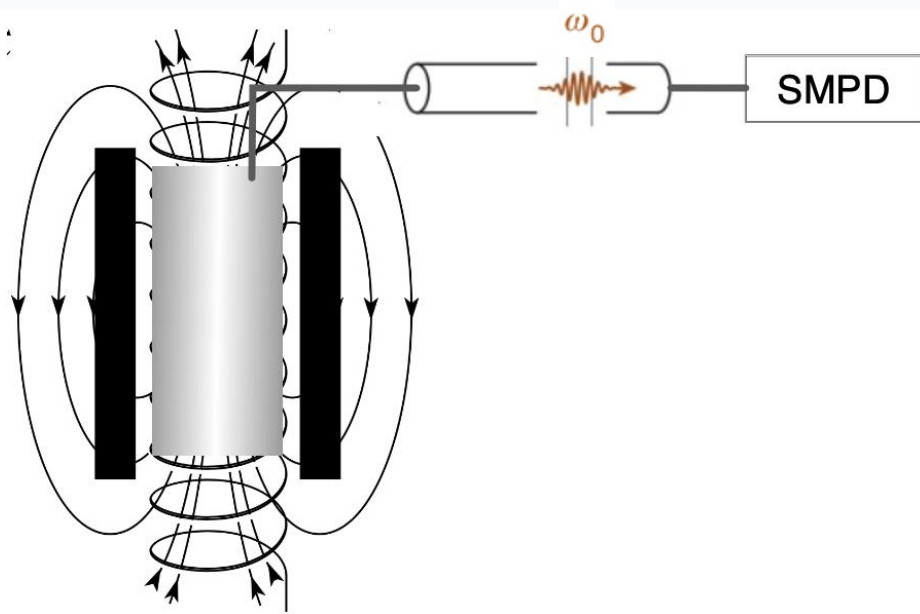


The transmon anharmonic potential yields non-equidistant energy levels. This allows to isolate the two lowest energy levels $|0\rangle$ and $|1\rangle$ that constitute the two-dimensional qubit subspace

A **Single Photon Microwave Counter (SMPD)** architecture is significantly different whether it is meant for **cavity photons** or **itinerant** (traveling) photons.

We are interested in the itinerant version due to the *intense magnetic fields* involved in axion search.

cavity photons $\neq$ itinerant photons



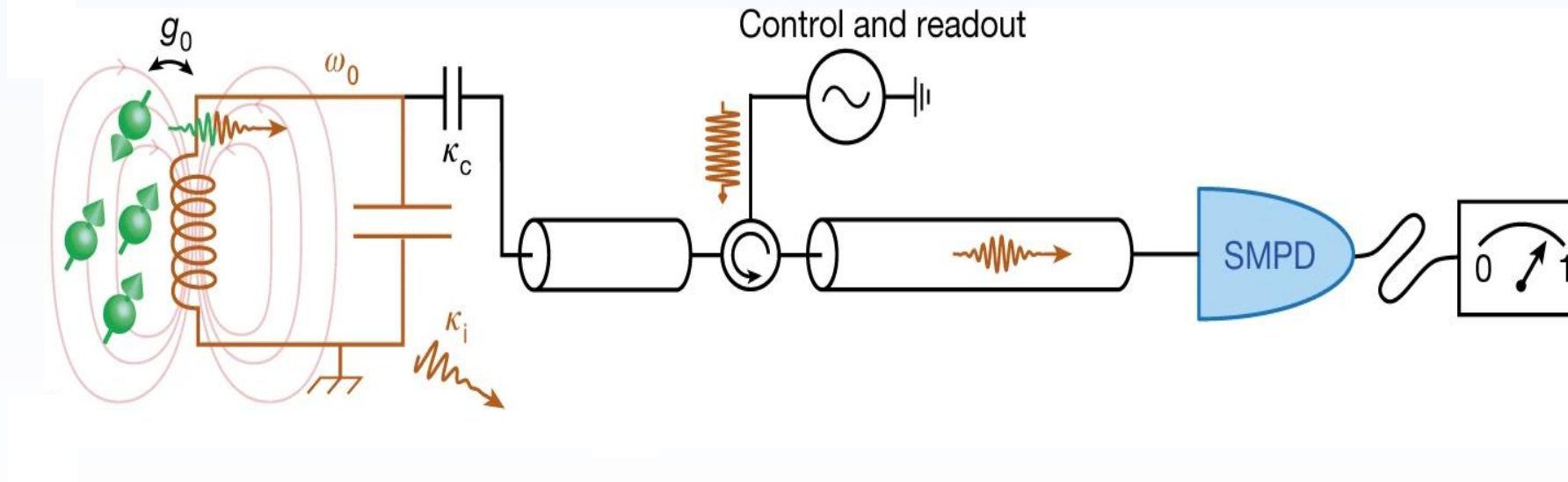
-detection of individual microwave photons is a challenging task because of their low energy $\sim 10^{-5}$ eV

-a solution: use “artificial atoms” introduced in circuit QED, their transition frequencies lie in the $\sim$ GHz range

-or: rely on a single current-biased Josephson junction

Transmon-based SMPD

In the Quantronics group (CEA, Saclay) a transmon-based counter has been developed and used to make **spin fluorescence measurements**, paving the way to single spin flip detection with SMPDs.

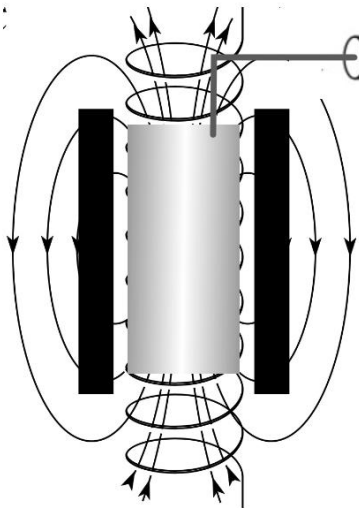


[R. Lescanne et al, Phys. Rev. X 10, 021038 \(2020\)](#)

E. Albertinale , Nature 600, 434 (2021)

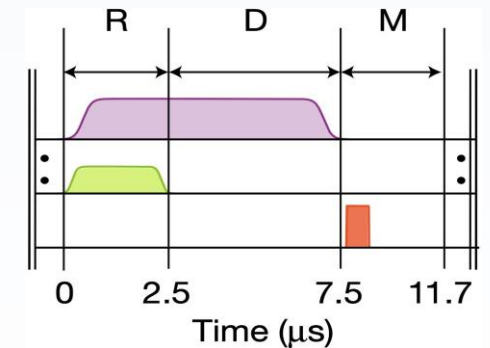
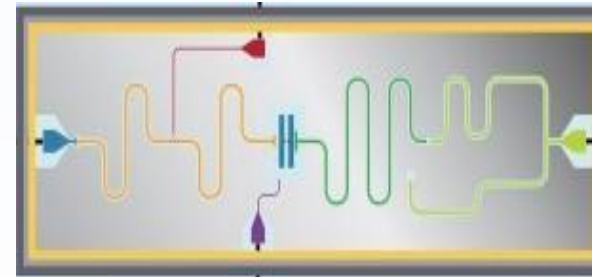
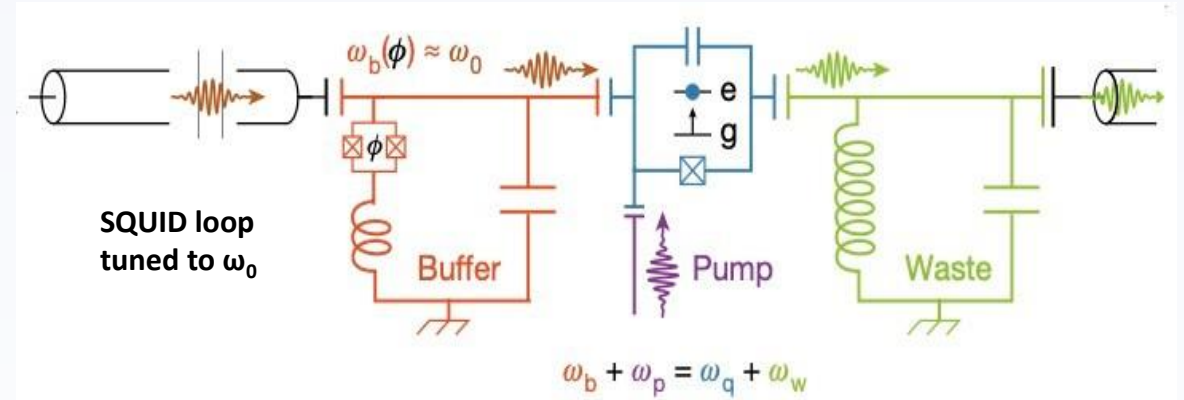
Detecting dark matter through quantum science

cavity photons $\neq$ itinerant photons



SMPD relies on a transmon qubit at frequency ω_q coupled to three ports:

buffer, pump and waste



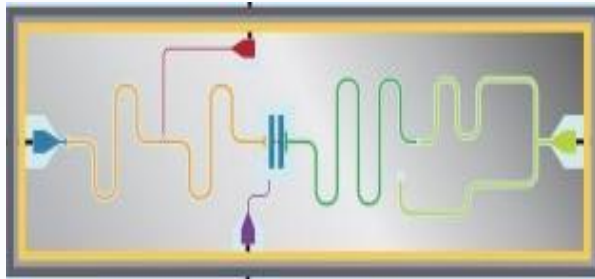
a three-step process repeated continuously

- qubit reset (R) performed by turning on the pump pulse + weak resonant coherent pulse to the waste port
- detection (D) step with the pump pulse on
- measurement (M) step probes the dispersive shift of the buffer resonator to infer the qubit state

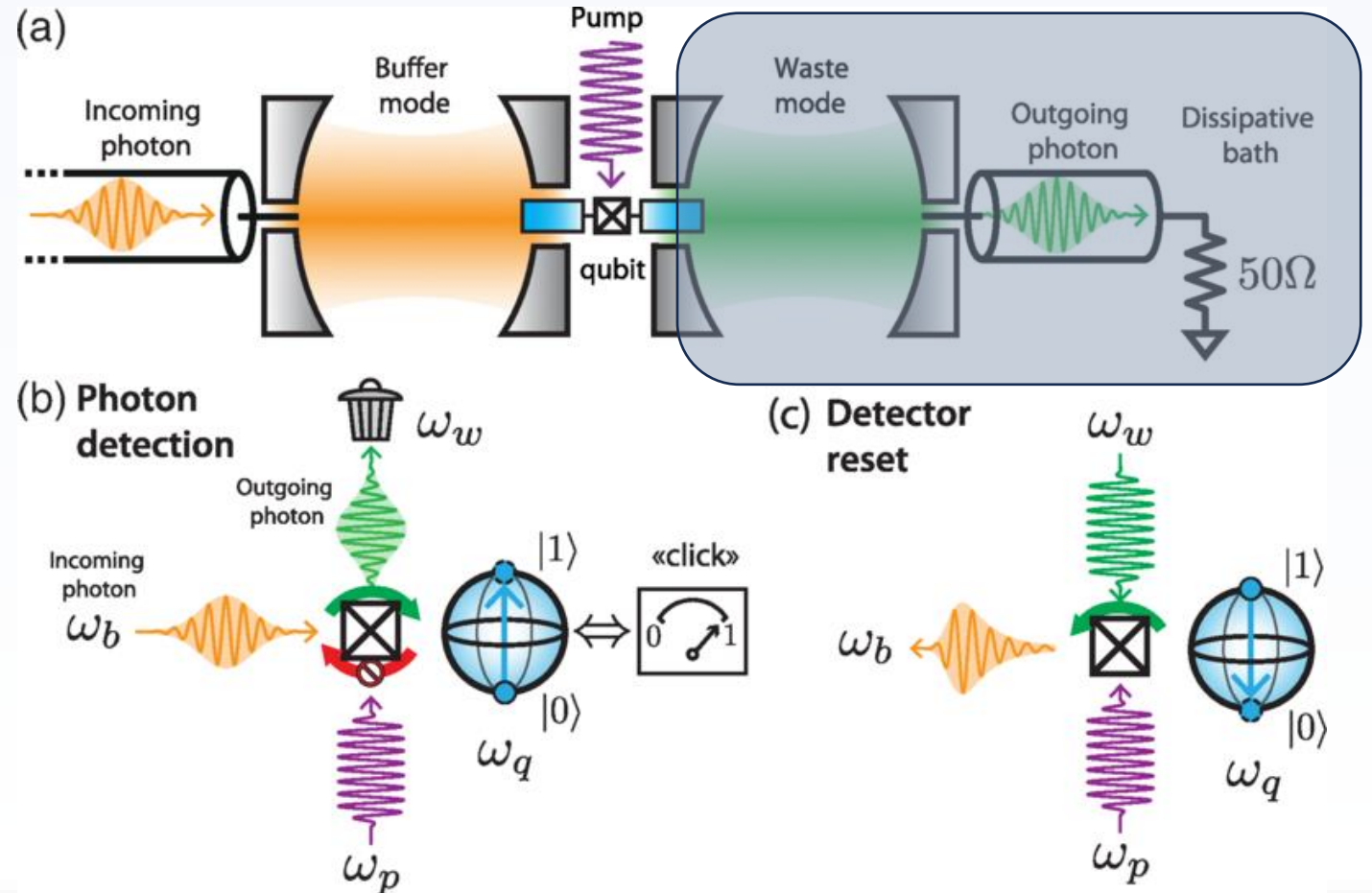
Detecting dark matter through quantum science

irreversible transfer of an incoming photon to an excitation of a transmon qubit

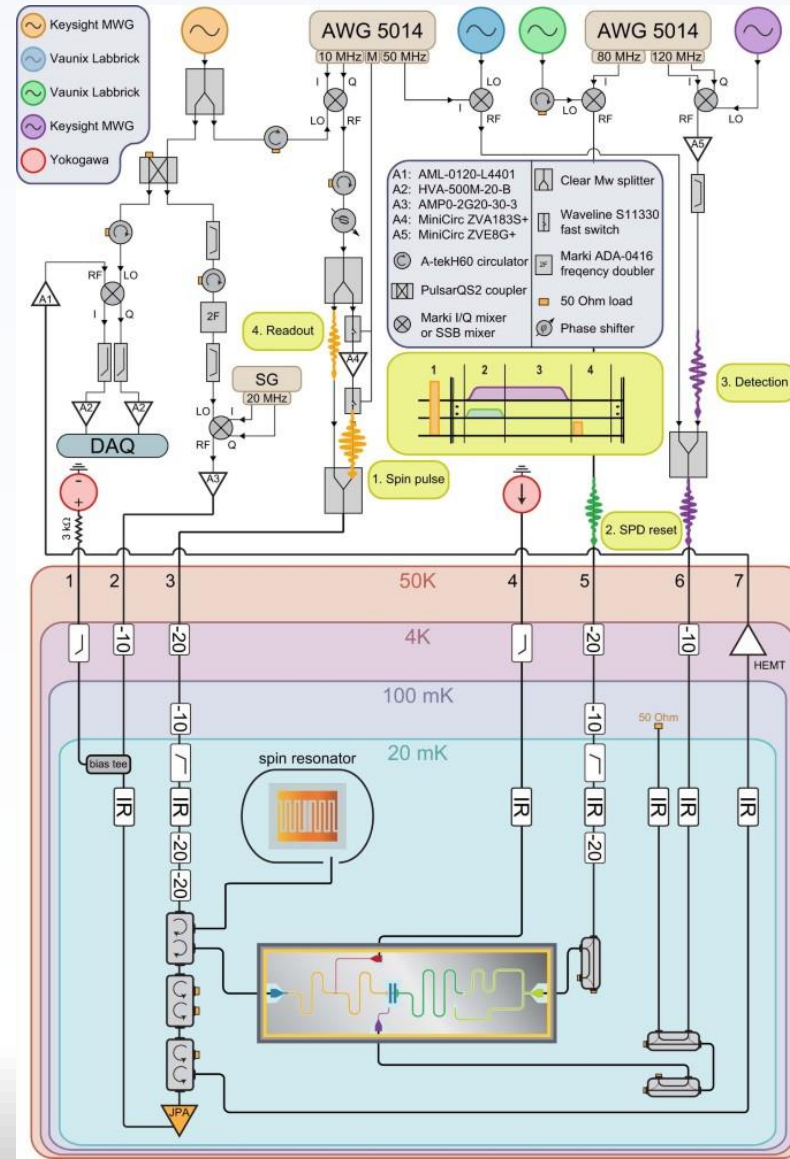
the resonator and qubit do not directly interact but instead dissipate energy to a common cold bath



$$\omega_b + \omega_p = \omega_w + \omega_q$$



Detecting dark matter through quantum science



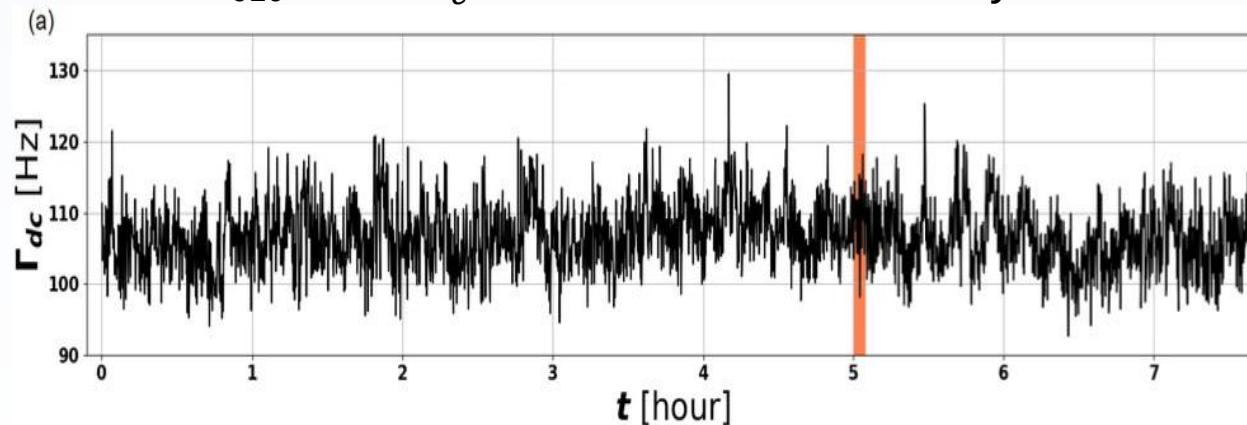
Quantronics Group
Research Group in Quantum
Electronics, CEA-Saclay, France

SMPD in axion dark matter search

single microwave photon detectors (SMPDs) developed in the context of quantum information science have the potential to **greatly improve the search speed at haloscopes**

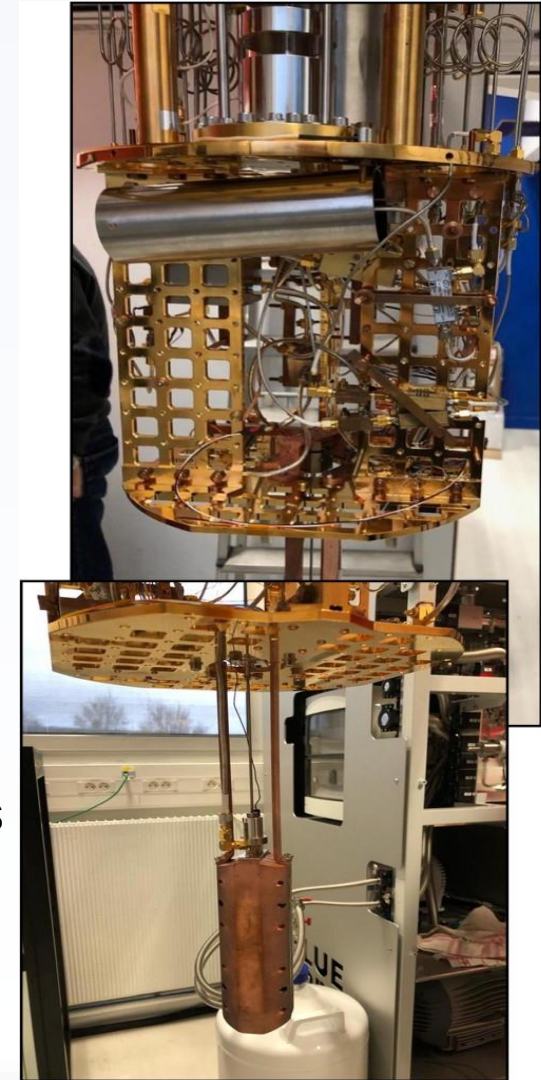
A pilot experiment performed with researchers from **INFN** (Padova, **LNL**) and the **Quantronics group** (CEA Saclay)

- tunable, high quality factor cavity
- TM_{010} mode $\nu_c \rightarrow 7.3$ GHz is readout by a SMPD



Detector bandwidth	~	1 MHz
Tunability range	~	200 MHz
Dark count rate	~	100-120 Hz
Efficiency	~	0.4

Unwanted photons:
false-positive events



Operating temperature = 10 mK

Detecting dark matter through quantum science

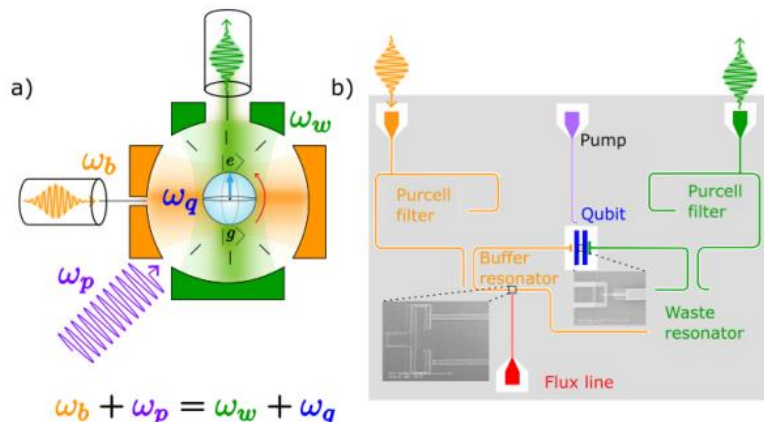
PHOTON DETECTOR PERFORMANCES

The efficiency η of the detector is defined as the probability of detecting the qubit in its excited state assuming a single incoming photon

the dark count rate α defined as the number of false positive detection per unit of time

power sensitivity S of the detector as the noise equivalent power (NEP) for an integration time of 1 s

$$S = \frac{\hbar\omega\sqrt{\alpha}}{\eta}.$$



Bolometer :

$$S = 7 \cdot 10^{-19} \text{ W}/\sqrt{\text{Hz}} \text{ at } 7.9 \text{ GHz},$$

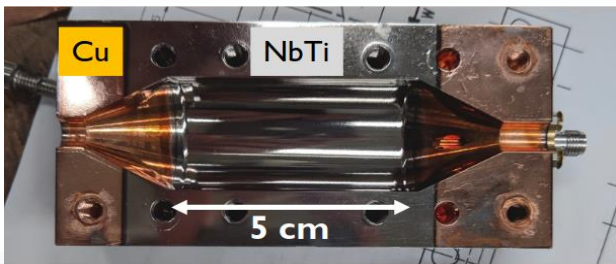
SMPD : $S = 10^{-22} \text{ W}/\sqrt{\text{Hz}}, \quad 7.005 \text{ GHz to } 6.824 \text{ GHz}$

New DU refrigerator + SRF searches

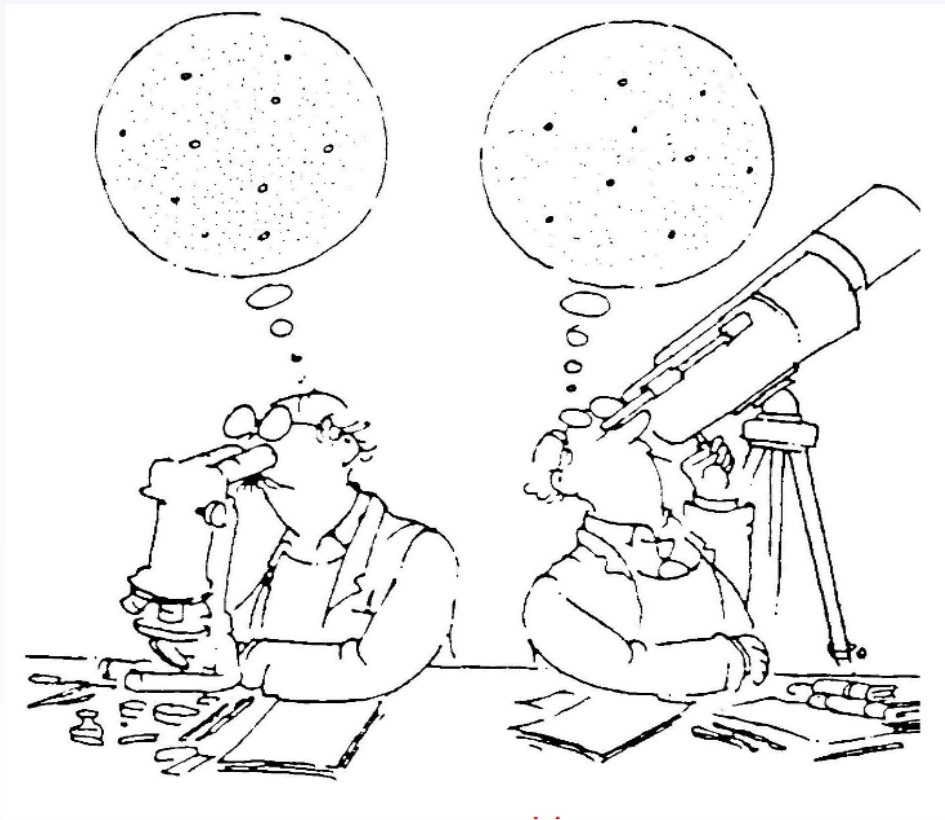
New dilution unit → 7 mK base Temperature

Paris test-run was successful so the SMPD device will be transferred to Italy for haloscope developments in INFN-LNL:

- $\text{QUAX}_{\alpha\gamma\gamma}$
- $\text{QUAX}_{\alpha e}$
- Superconductive cavity (copper cavity sputtered with **NbTi**)



Model CF-CS110-



Thank you