

COSMIC WISPer
1st General Meeting

Search for WISPer with gravitational-wave detectors

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on behalf of the LIGO-Virgo-KAGRA collaborations



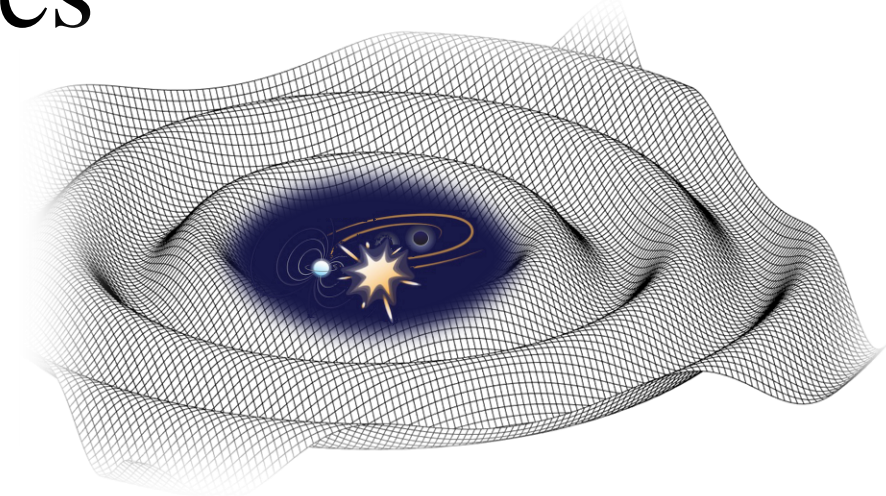
Istituto Nazionale di Fisica Nucleare



Gravitational Waves

- Gravitational Waves (GWs) are solutions of the linearized Einstein field equations in vacuum:

Space-time metric perturbation $\leftarrow \square h_{ij} = -\frac{16\pi G}{c^4} T_{ij} \rightarrow$ Matter



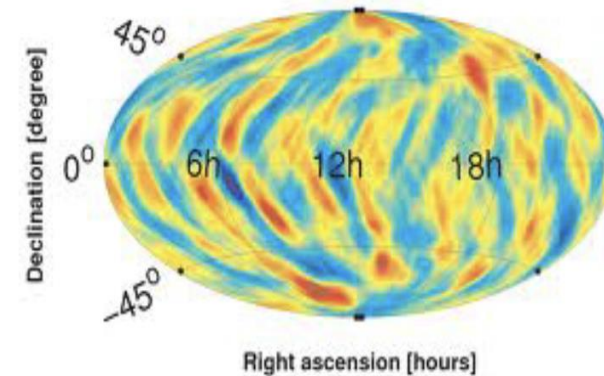
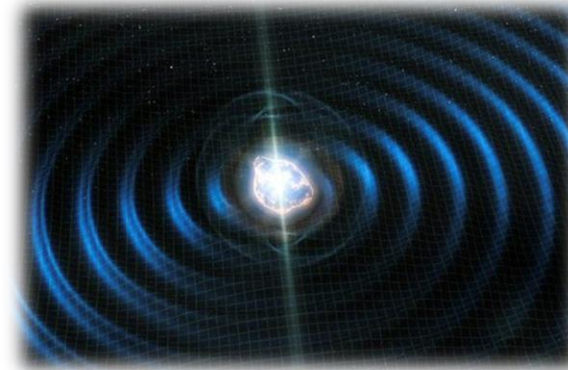
- Produced by the bulk motion of matter. Some sources we search for:

Coalescing compact binaries
(black holes, neutron stars)

Supernova explosions

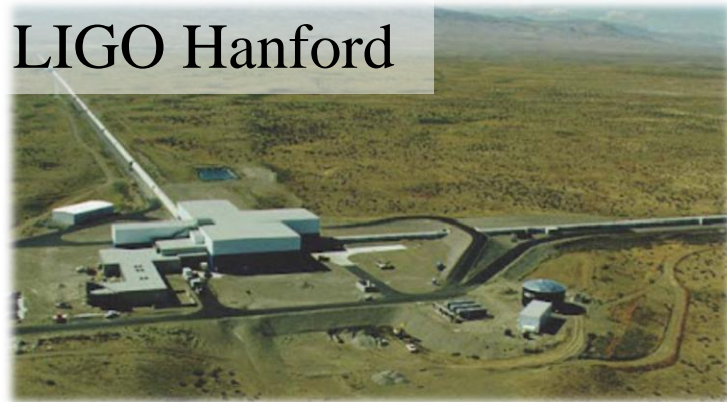
Continuous Waves
(e.g., spinning neutron stars)

Stochastic background
(Astrophysical,
Cosmological)



Gravitational-wave detectors network

LIGO Hanford



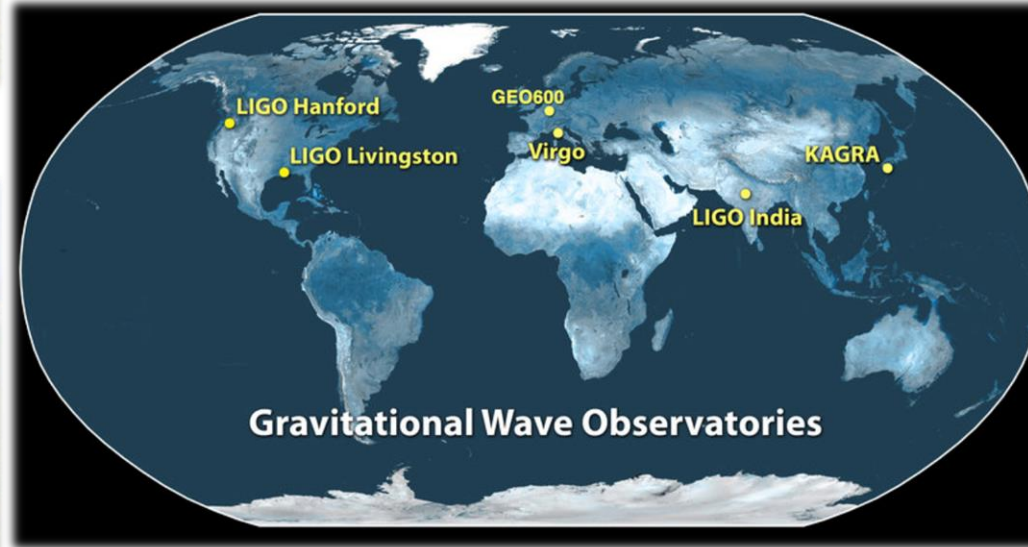
GEO600



Virgo



LIGO Livingston

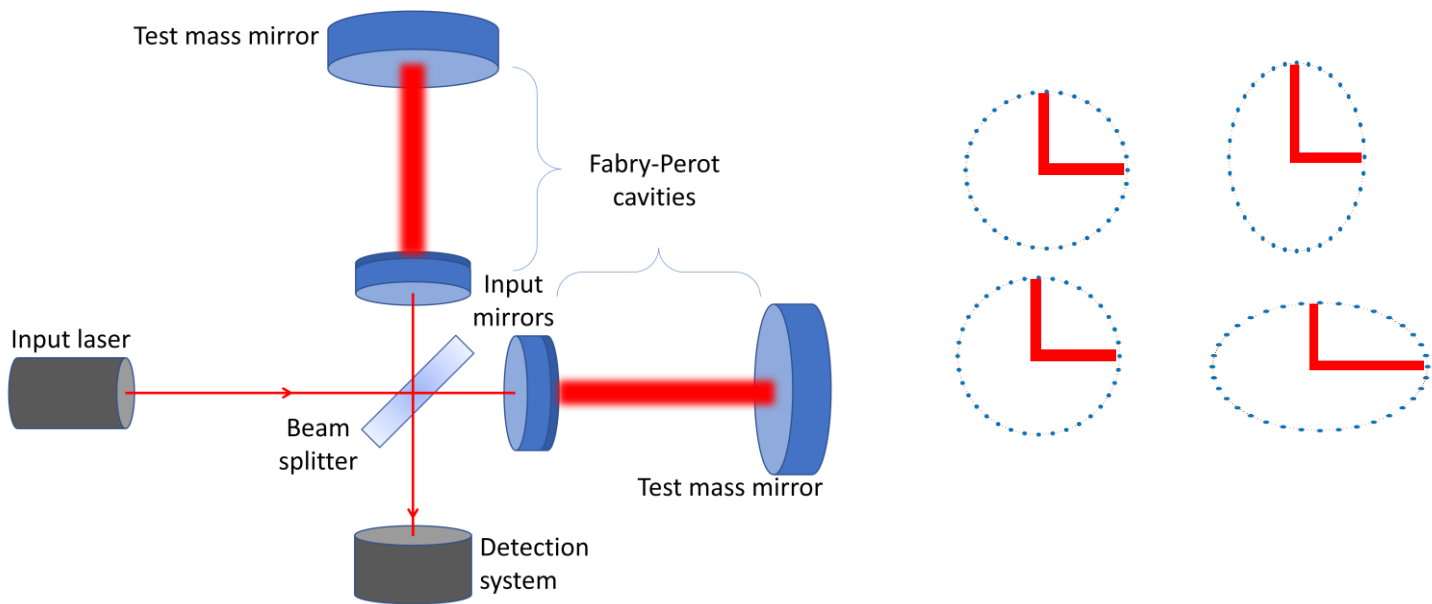


KAGRA

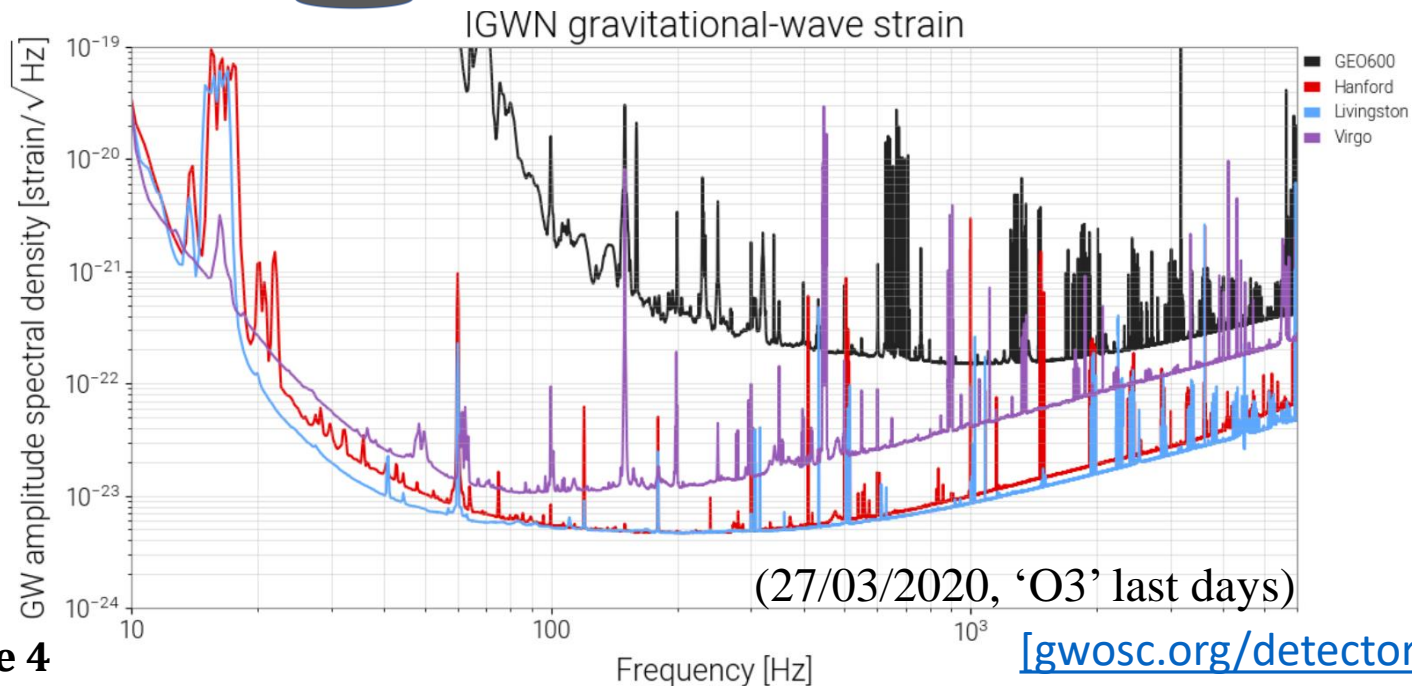


- Three observing runs completed in the Advanced configuration from 2015 to 2020 (O1-O3).
- Run O4 started on 24th May 2023, expected duration 1.5 years.
LIGO has started O4 and is already observing. Virgo has postponed to face extra noises.
slide 3 KAGRA has joined the run and then stopped to continue commissioning and join again later.

Gravitational-wave detectors basics



- Michelson interferometers with Fabry-Perot cavities.
- GW effect: perturbation of the laser optical path in the arms.
- Output → phase shift between two recombined laser beams.
- → $h(t)$ time series.



Limiting noises:

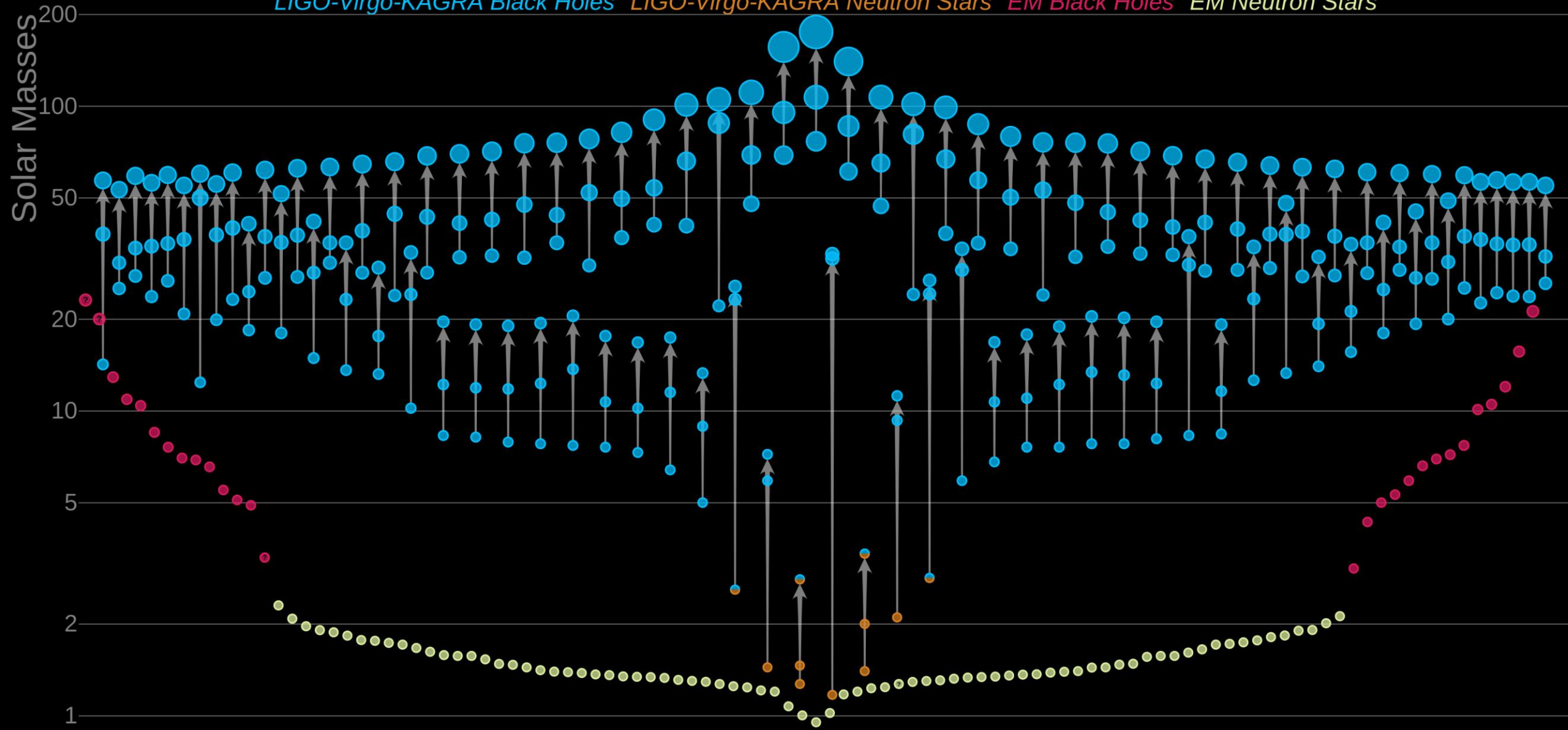
(< **80 Hz**) Newtonian, seismic, thermal, quantum radiation pressure, technical noises.

(**80 – 300 Hz**) Thermal noise.

(> **300 Hz**) Quantum shot noise.

90 coalescing binaries detected in the first 3 runs

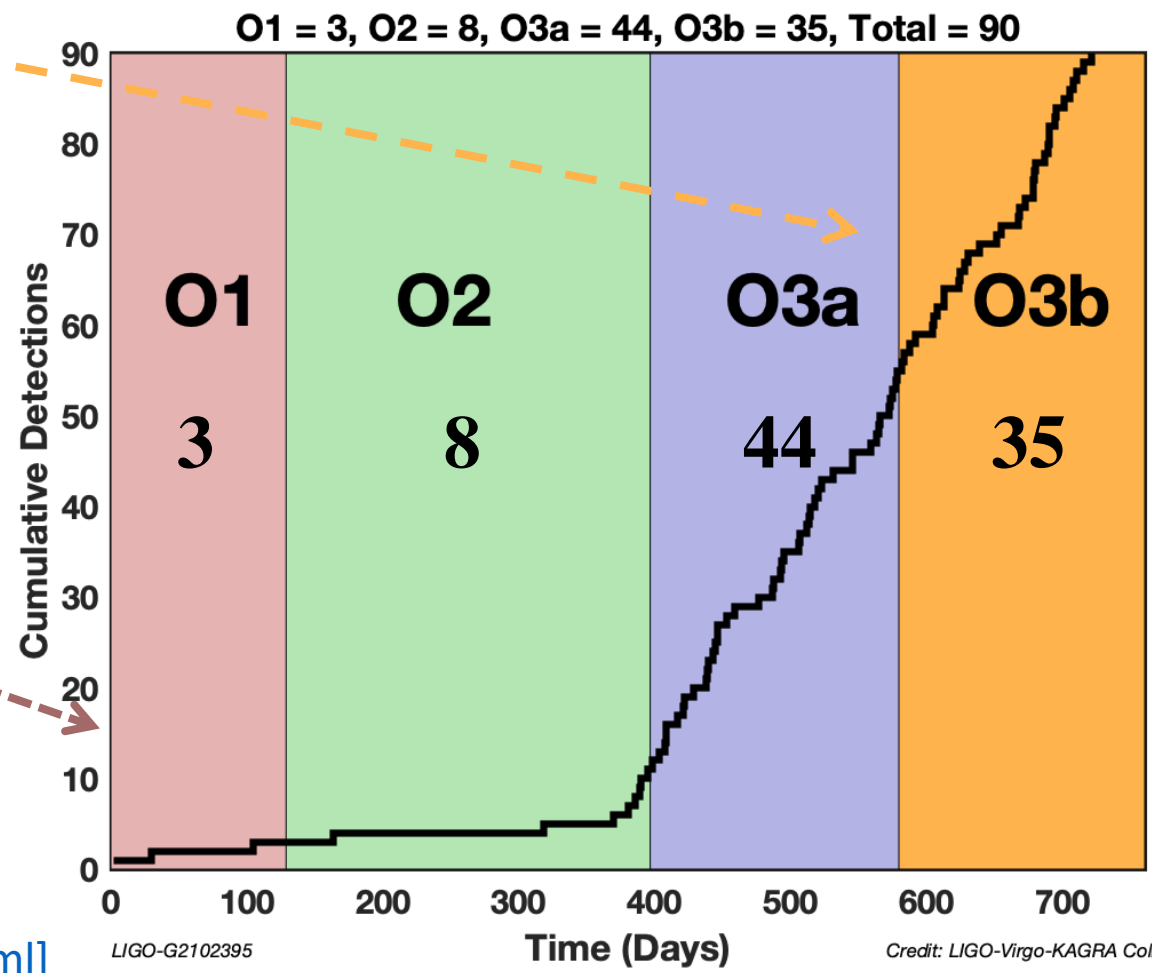
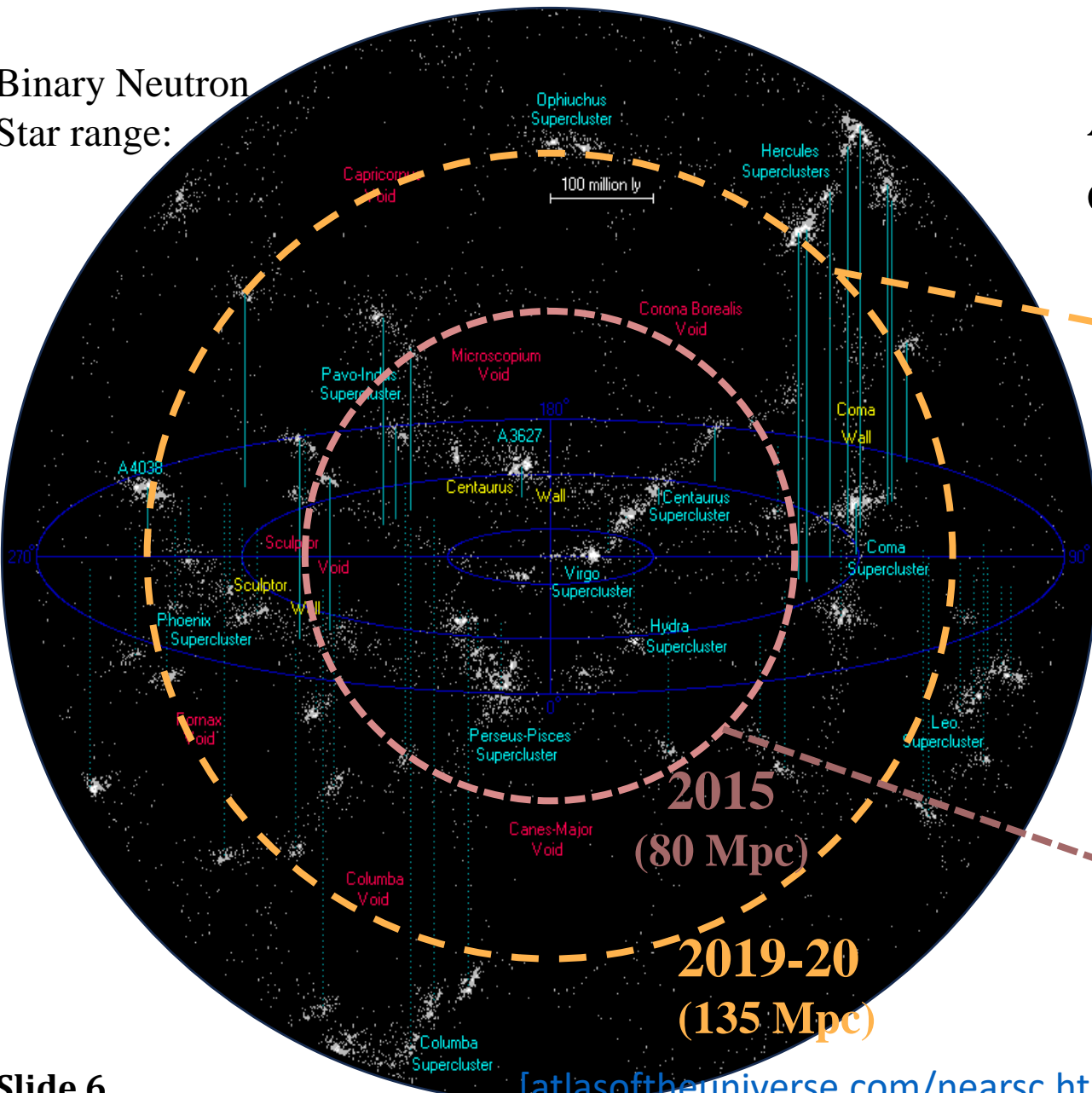
LIGO-Virgo-KAGRA Black Holes *LIGO-Virgo-KAGRA Neutron Stars* *EM Black Holes* *EM Neutron Stars*



The rate of detections is continuously growing

Binary Neutron
Star range:

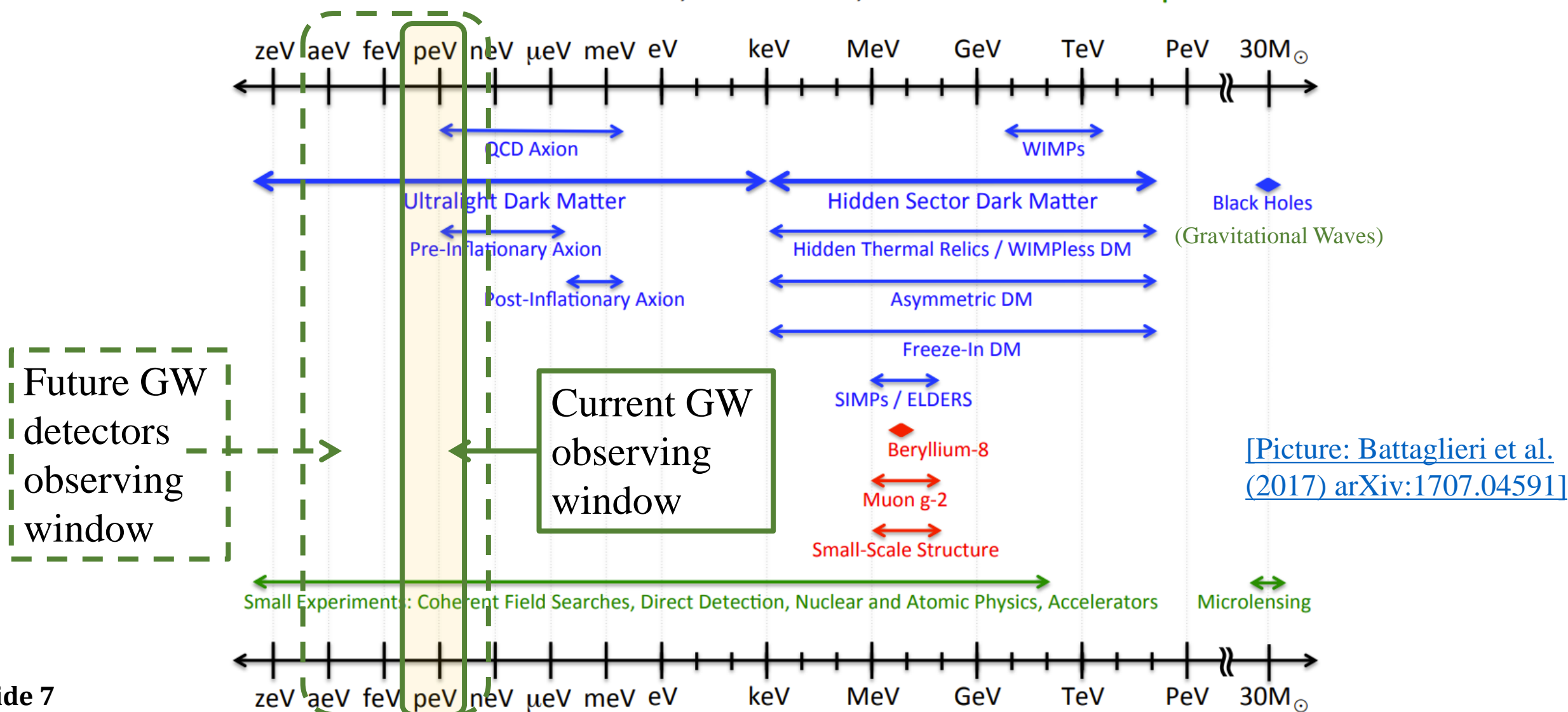
As the sensitivity increases, the volume of inspected universe scales as r^3



The Dark Matter energy scale we can probe

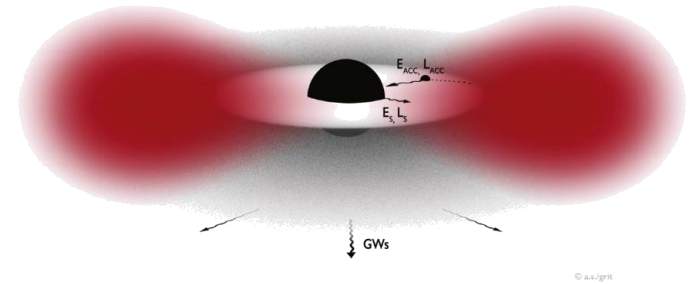
Dark Matter (DM) candidates cover ~90 orders of magnitude in mass

Dark Sector Candidates, Anomalies, and Search Techniques

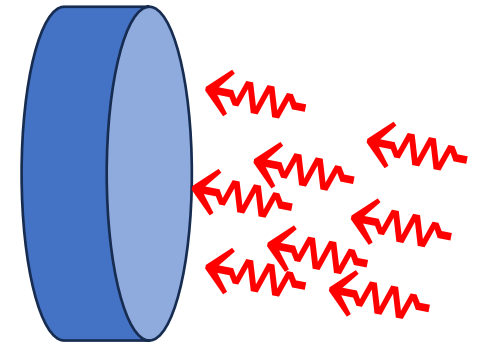
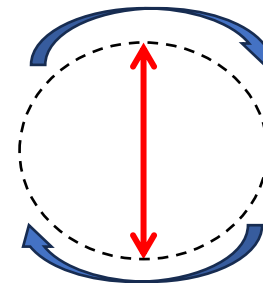


WISPs searches with Earth-based GW detectors

- Ultra-light boson condensates around compact objects. These condensates could directly emit GW, but can also interact with the object's dynamics.



- Direct interaction of DM fields with detectors. They can produce a differential strain on the detectors data without being GWs.
 - Interaction with optics
 - Rotation of laser polarization



Ultra-light *boson clouds*

- Typical subjects: dark photon, QCD axion, axion-like particles from string theory.
Considered masses in the range $[10^{-20}, 10^{-11}]eV$
- If they exist in that mass range, they can trigger superradiant instability around spinning black holes.
- Condition: Compton wavelength (λ) comparable with the black hole size (r_s).

- The bosonic field can bind to the BH, forming a “gravitational atom”.

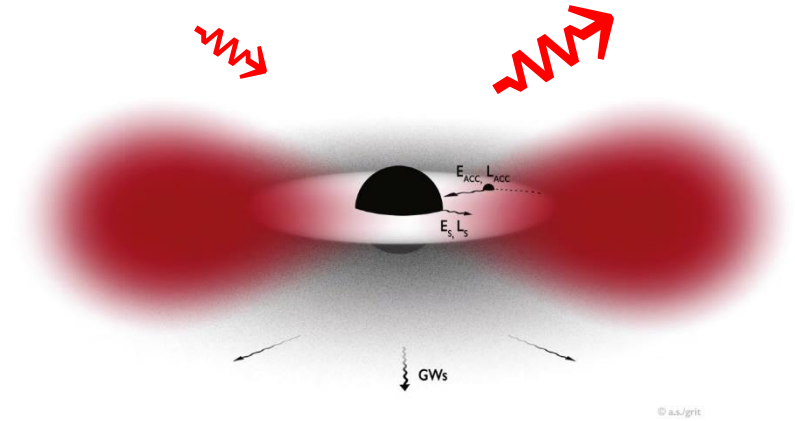
Gravitational fine-structure constant:
$$\alpha = \frac{r_s}{\lambda} = \frac{GM_{\text{bh}}}{c^3} \frac{m_b}{\hbar}$$

- Impact on BH spins and dynamics. → Searches affordable with future detectors

[\[Ng *et al.* \(2021\) PRD 103, 063010.\]](#) [\[De Luca, Pani \(2021\) JCAP, 032\(08\)\]](#) [\[Qing Yang *et al.* \(2018\) RAA 18 065\]](#)

- Emission of GW from the depletion of the boson cloud. → Searches already started!

[\[Arvanitaki *et al.* \(2010\) PRD 81, 123530\]](#) [\[Brito *et al.* \(2017\) PRD 96, 064050\]](#)



[Picture credit: [Ana Sousa Carvalho](#)]

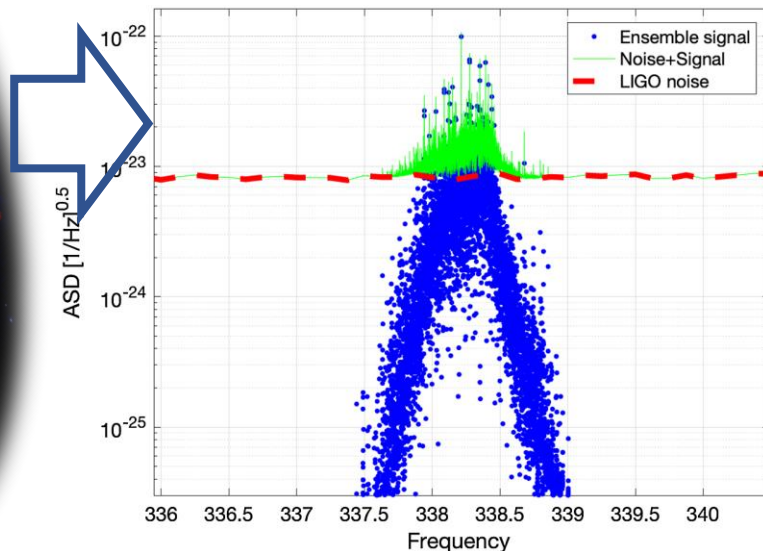
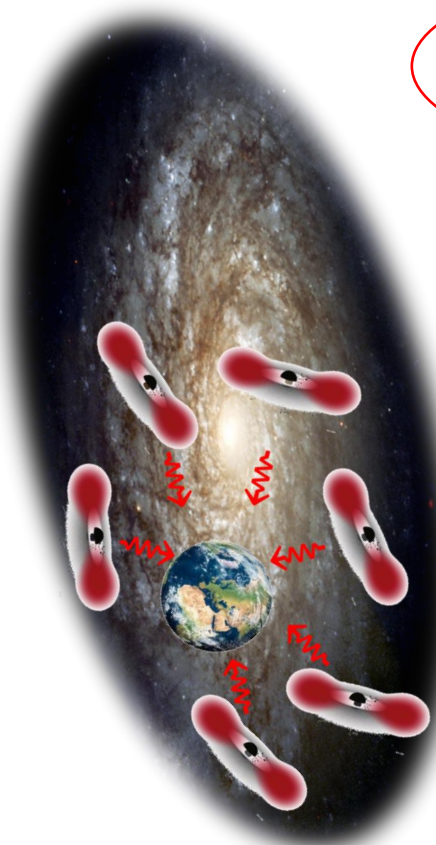
GWs from ultra-light scalar boson clouds

In the case of scalar bosons:

- Cloud forms in \sim day timescale
- Annihilation of bosons into gravitons (GWs).
- Quasi-monochromatic GWs with $\sim 10^4$ years timescale.

$$f_{\text{gw}} \simeq 483 \text{ Hz} \frac{m_b}{10^{-12} \text{ eV}} \left[1 - 7 \cdot 10^{-4} \left(\frac{M_{\text{BH}}}{10 M_{\odot}} \cdot \frac{m_b}{10^{-12} \text{ eV}} \right)^2 \right]$$

2nd order correction



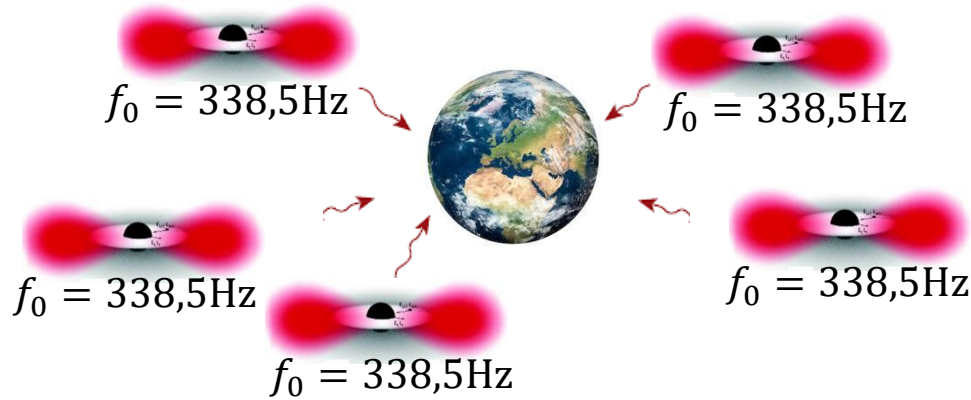
If the superradiance affects a significant fraction of galactic black holes, we would expect to have several GWs all around a same frequency!

- Search them as an ensemble of signals.
- Search them as multiple individual signals.

[\[Zhu et al. \(2020\) PRD 102, 063020\]](#)

[\[D'Antonio et al. \(2018\) PRD 98\(10\), 103017\]](#)

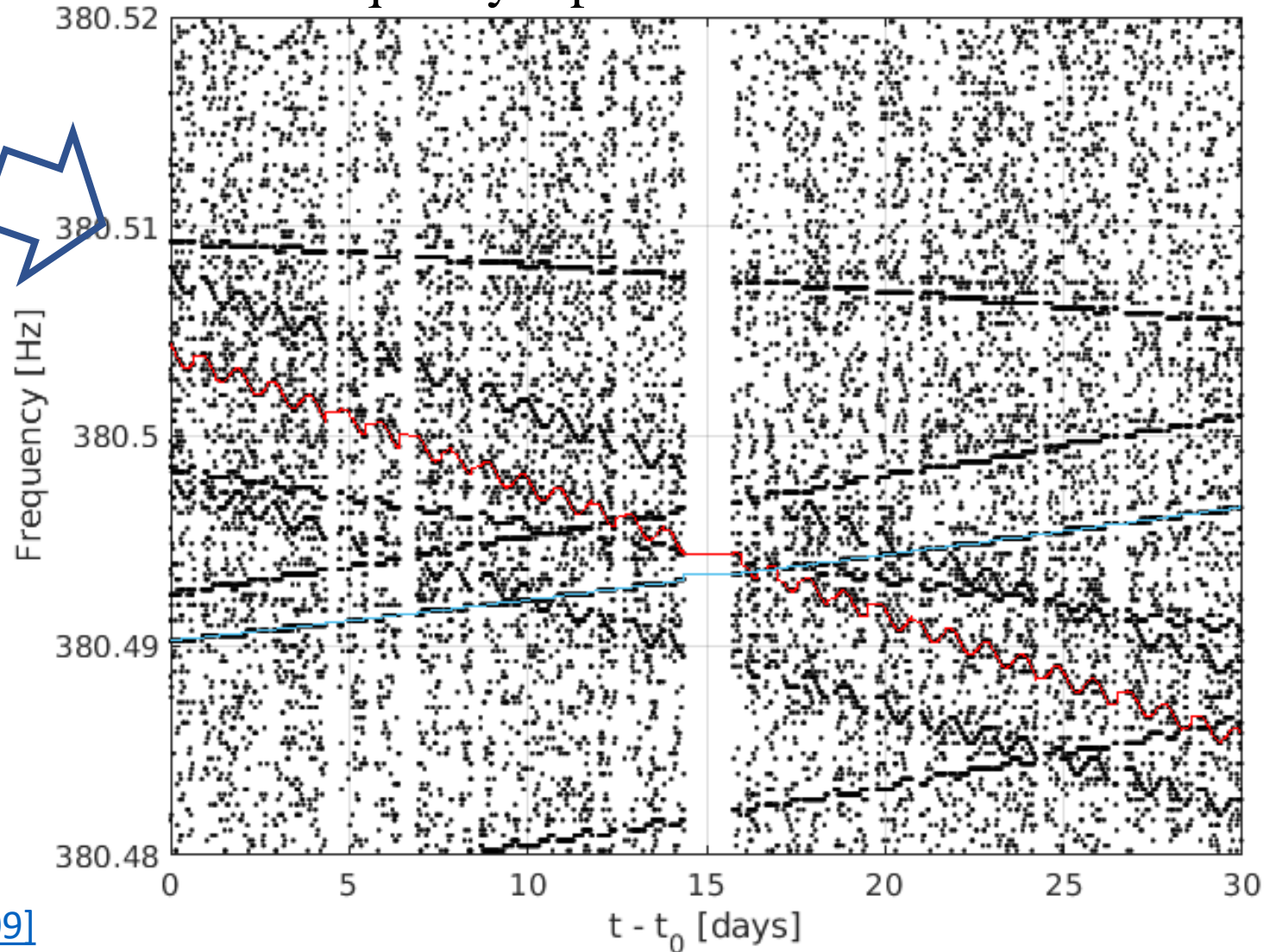
Multiple signals can be resolved!



- The daily and annual Doppler effect due to the Earth motion modulates the observed signal.
- Exploiting the Doppler, we can resolve individual signals up to the sky resolution.

[\[Pierini et al. \(2022\) PRD 106\(4\), 042009\]](#)

Time-frequency representation of one-month data

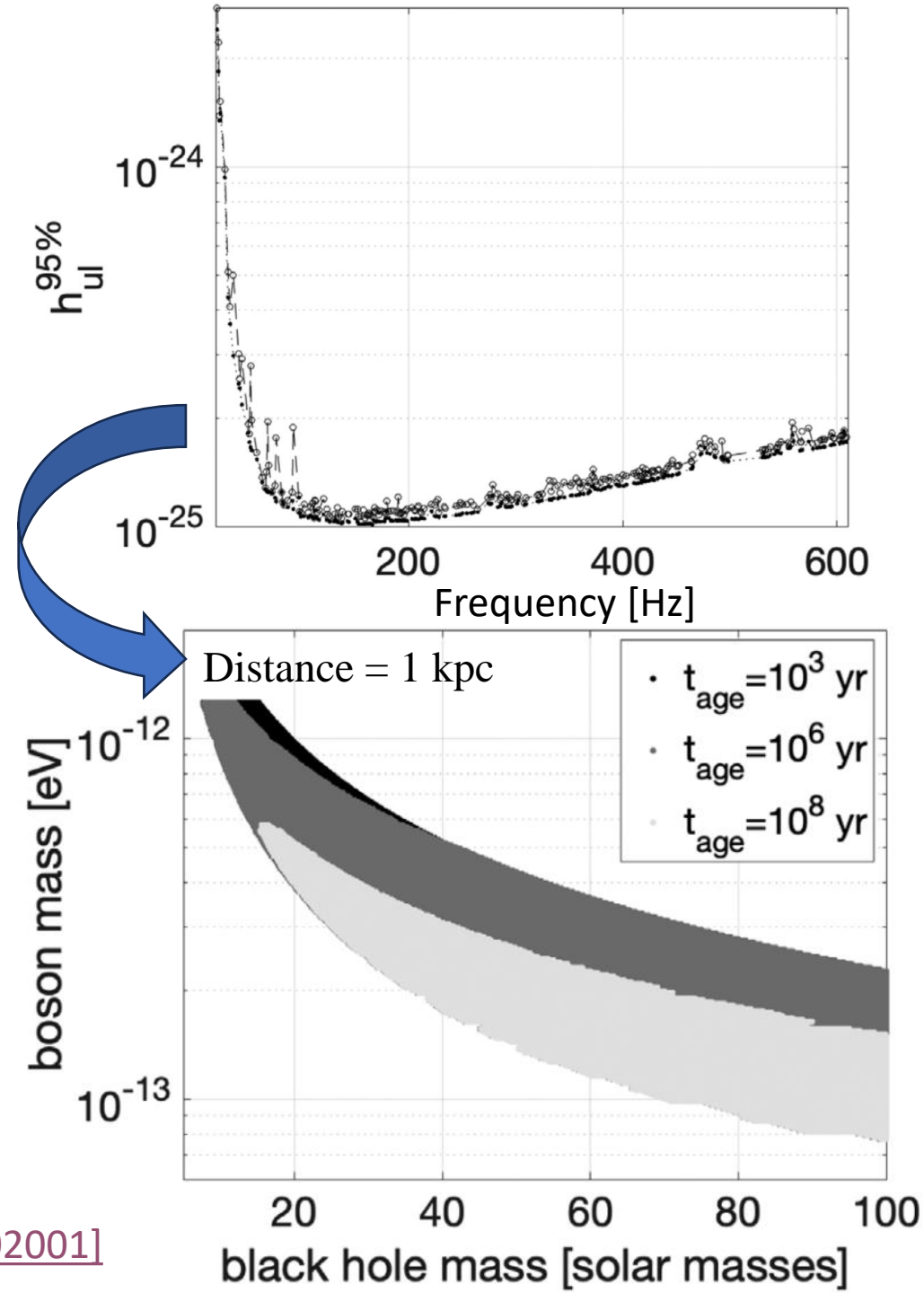


Exclusion regions from the last observing run O3

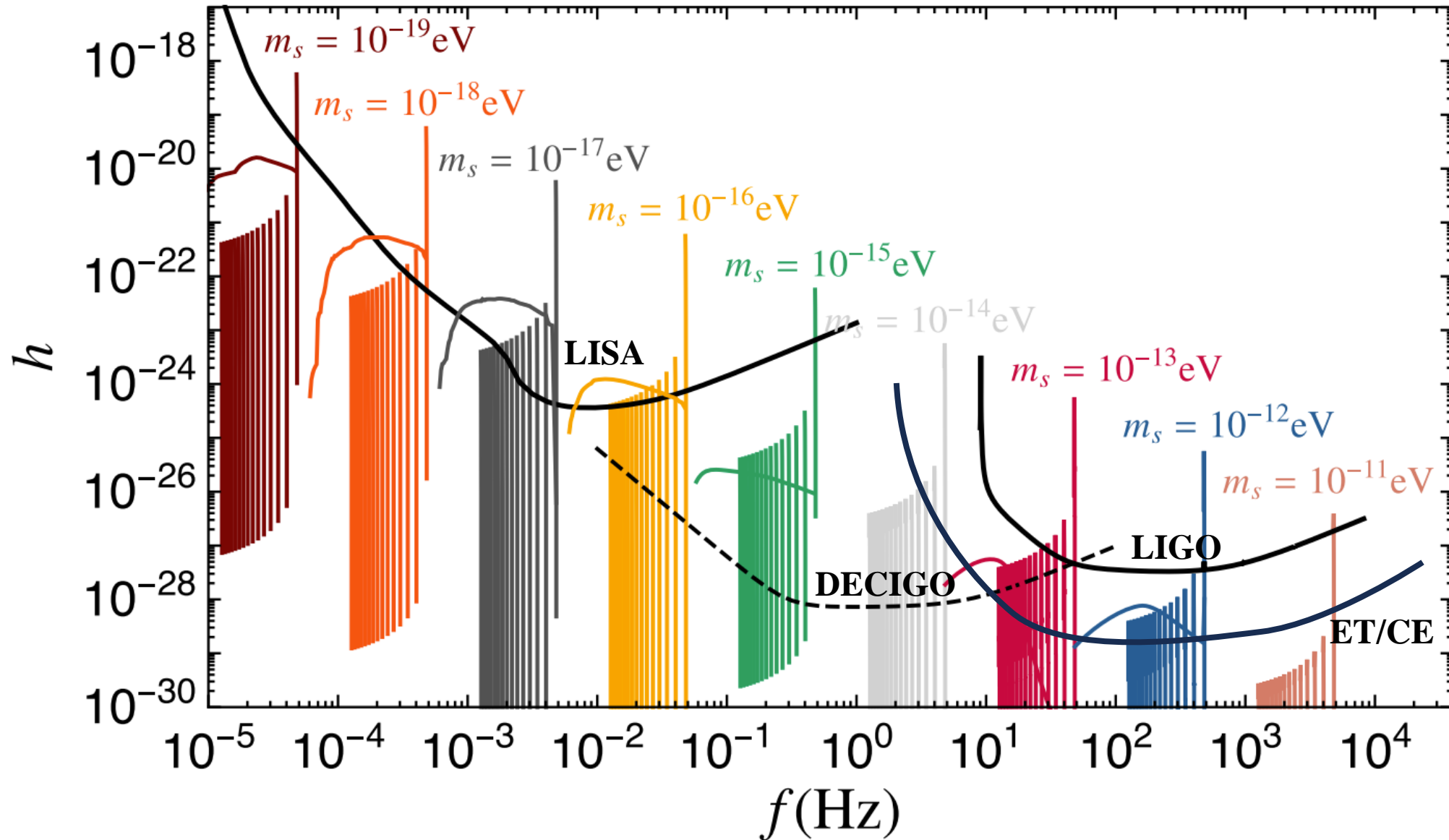
- The LVK collaboration makes all-sky searches for boson clouds around spinning BHs in the Galaxy.
- No significant candidates have been found
 - upper limits in the GW amplitude
 - exclusion regions for the boson mass.

$$h_0 \sim 6 \cdot 10^{-24} \left(\frac{M_{\text{bh}}}{10 M_{\odot}} \right) \left(\frac{\alpha}{0.1} \right)^7 \left(\frac{1 \text{ kpc}}{r} \right) (\chi_i - \chi_c) \left(1 + \frac{t}{\tau_{\text{age}}} \right)^{-1}$$

- At fixed distance, spin and age of black hole, we can exclude there are black holes with boson clouds with masses such to produce a detectable GW.
- At growing distance or age, the exclusion region becomes smaller.

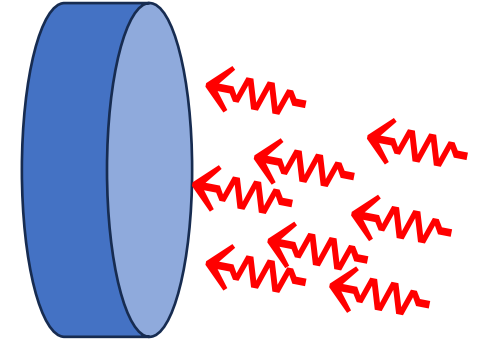


Future observations



Direct detection of dark matter

- Ultra-light DM can directly interact with interferometer optical components producing a potentially detectable signal



→ It is not a GW signal, but nevertheless the interaction can cause a differential strain.

- The mass scale to which detectors are sensitive is set by the particle field frequency $f_0 = m_A c^2 / h$

Earth-based detectors sensitivity window: 10 – 1000 Hz.

➤ Probed particle masses: $10^{-14} - 10^{-11} \text{ eV}$

Dark photon

- Dark Photon (DP) couples to baryon ($U(1)_B$) or neutron number ($U(1)_{B-L}$)

[\[Nelson & Scholz \(2011\) PRD 84, 103501\]](#)

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_A^2 A^\mu A_\mu - \epsilon_D e J_D^\mu A_\mu$$

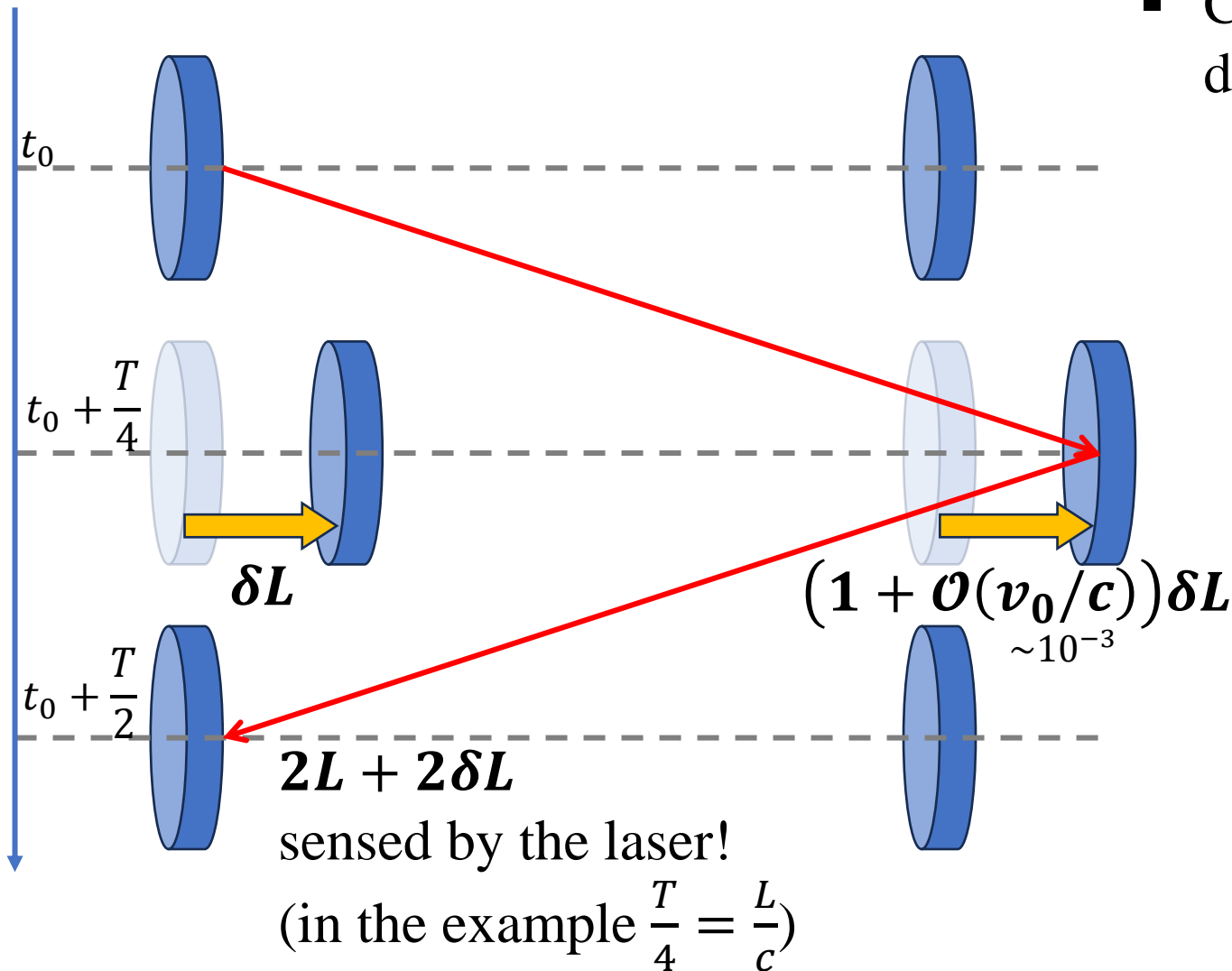
$F_{\mu\nu}$: DP field strength tensor
 A_μ : DP potential
 m_A : DP mass
 ϵ_D : DP normalized coupling strength

- For DP masses $10^{-14} - 10^{-11}$ eV and a local DM density $\rho_{DM} \approx 0.4 \text{ GeV/cm}^3$ the resulting occupation number is $O(10^{54})$
 - The DP field can be described as a superposition of plane waves

$$\vec{A}(\vec{x}) = \sum_i \vec{A}_i \cos(2\pi f_i t - \vec{k}_i \cdot \vec{x} + \phi_i)$$

- Frequency spread due to the Maxwell-Boltzmann velocity distribution of DPs around the virial velocity $v_0 \sim 220 \text{ km/s} \rightarrow \Delta f / f \sim 10^{-6}$

Dark photon interaction with detectors



T: oscillation period

- Coupling with mirrors \rightarrow differential strain due to the spatial gradient of the DP field:

$$\sqrt{\langle h_D^2 \rangle} \sim 6.56 \cdot 10^{-27} \left(\frac{\epsilon_D}{10^{-23}} \right) \left(\frac{100\text{Hz}}{f_0} \right)$$

[Pierce *et al.* (2018) *PRL* **121**, 061102]

- Equivalent differential strain due to finite light-traveling time in the arms:

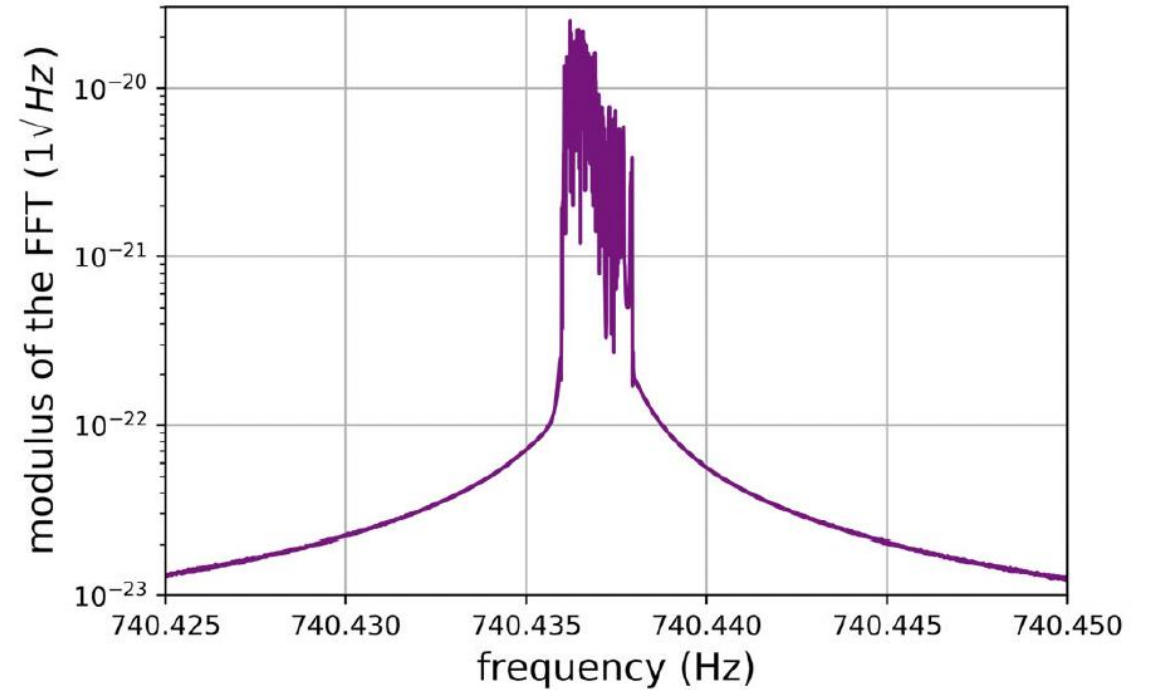
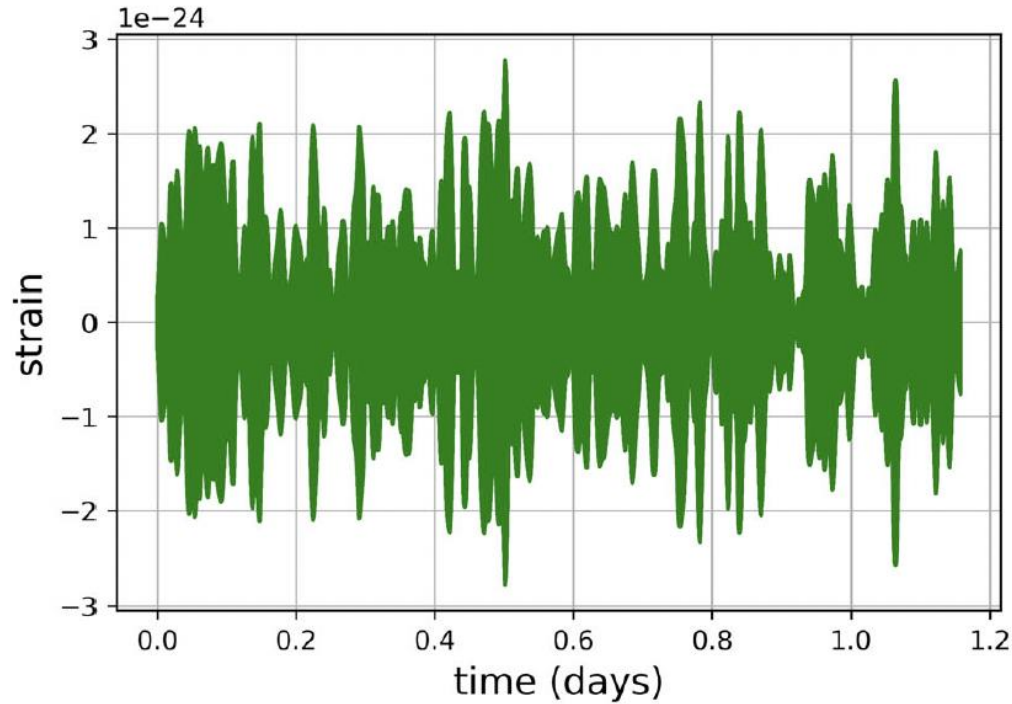
$$\sqrt{\langle h_C^2 \rangle} = \frac{\sqrt{3}}{2} \frac{2\pi f_0 L}{v_0} \sqrt{\langle h_D^2 \rangle}$$

$$\sim 6.58 \cdot 10^{-26} \left(\frac{\epsilon_D}{10^{-23}} \right)$$

[Morisaki *et al.* (2021) *PRD* **103**, L051702]

The expected signal from dark photon

Stochastic, narrow-band signal

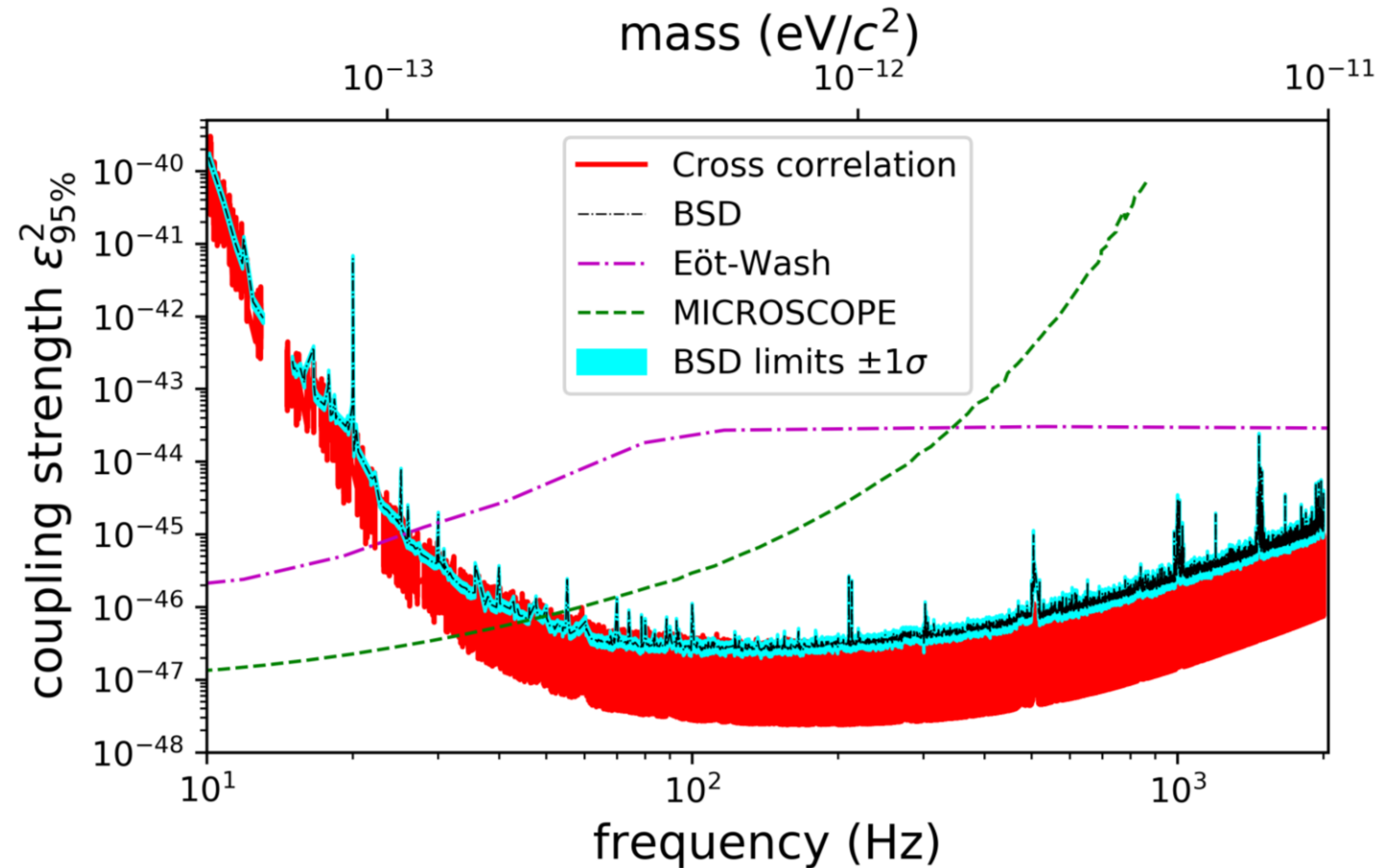


Most recent analysis carried on LIGO-Virgo O3 data, using two different analysis methods:

- Cross-correlation between different detectors. [\[Pierce *et al.* \(2018\) PRL 121, 061102\]](#)
- Excess-of-power statistics. [\[Miller *et al.* \(2021\) PRD 103, 103002\]](#)

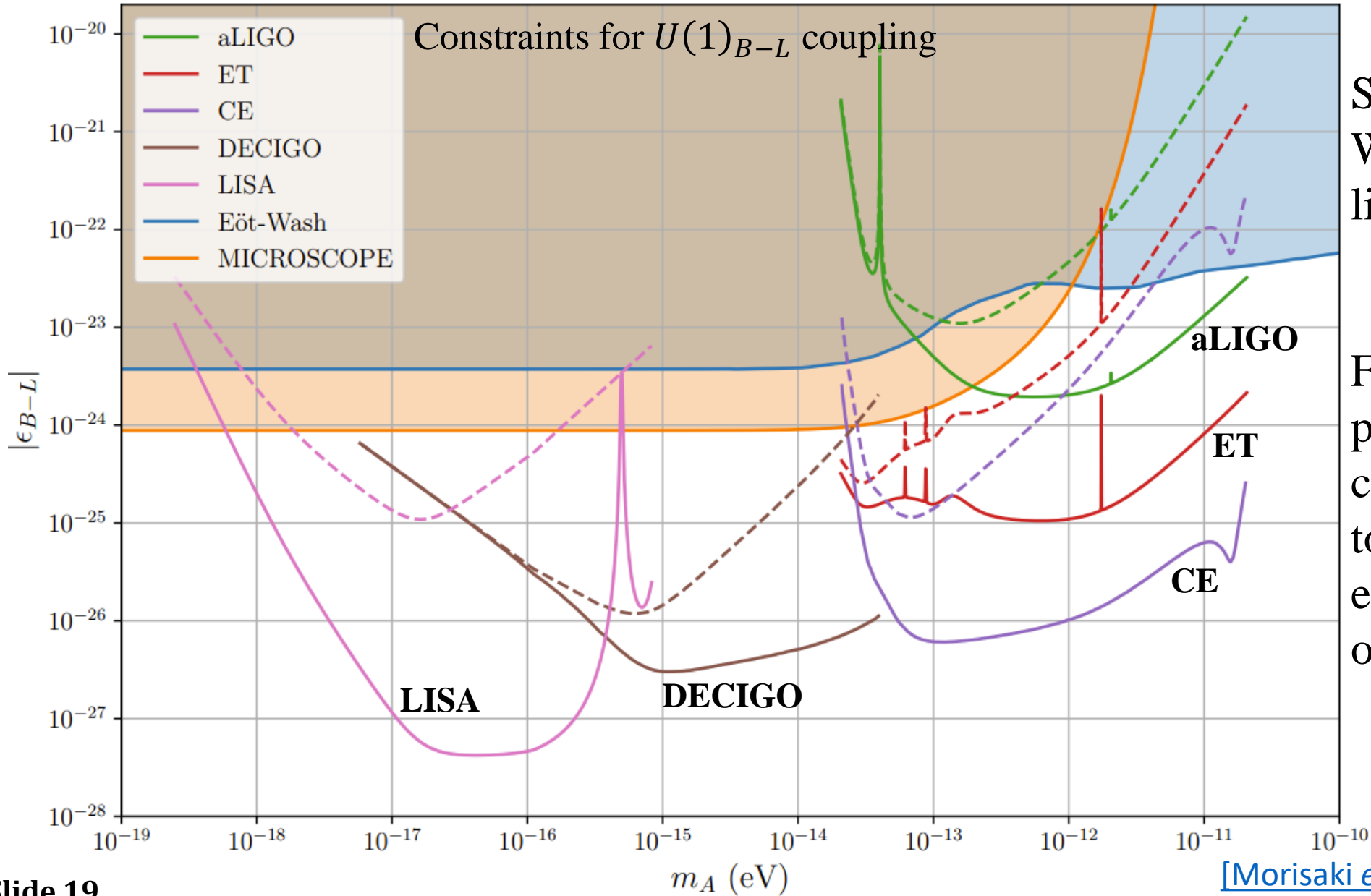
Upper limits on the coupling strength

- No detection, but competitive upper limits are computed.
- Compared to limits from existing torsion balance experiments (Eöt-Wash) and MICROSCOPE satellite.
- Improvement of two order of magnitude w.r.t. direct search experiments, assuming $U(1)_{B-L}$.
- For $U(1)_{B-L}$ upper limits are comparable to direct search experiments.



[Abbott *et al.* (2022) PRD 105, 063030]

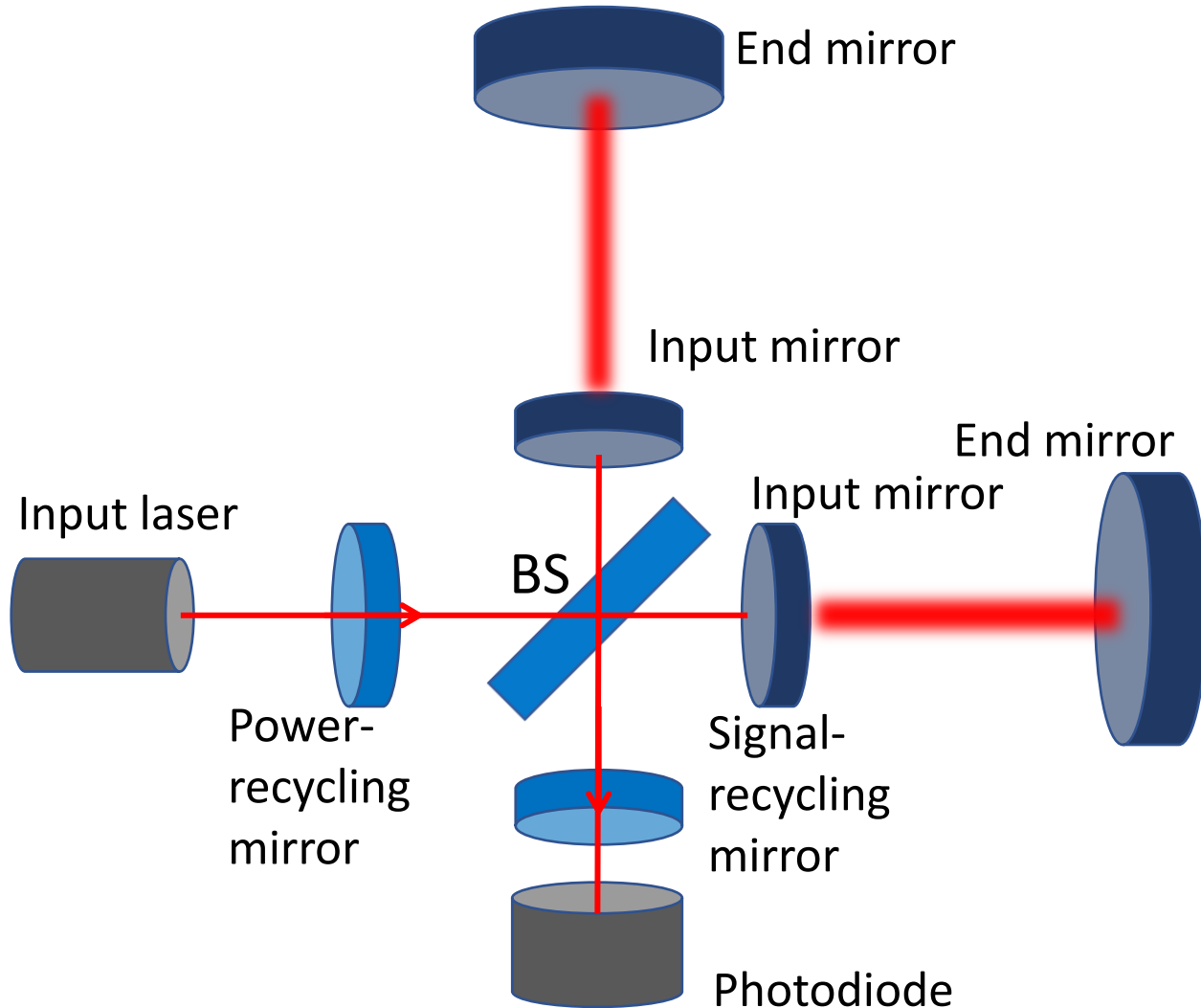
Future observations



Solid/Dashed:
Without/With finite
light-traveling time.

Future detectors will
provide better
constraints with respect
to Equivalence Principle
experiments by orders
of magnitude!

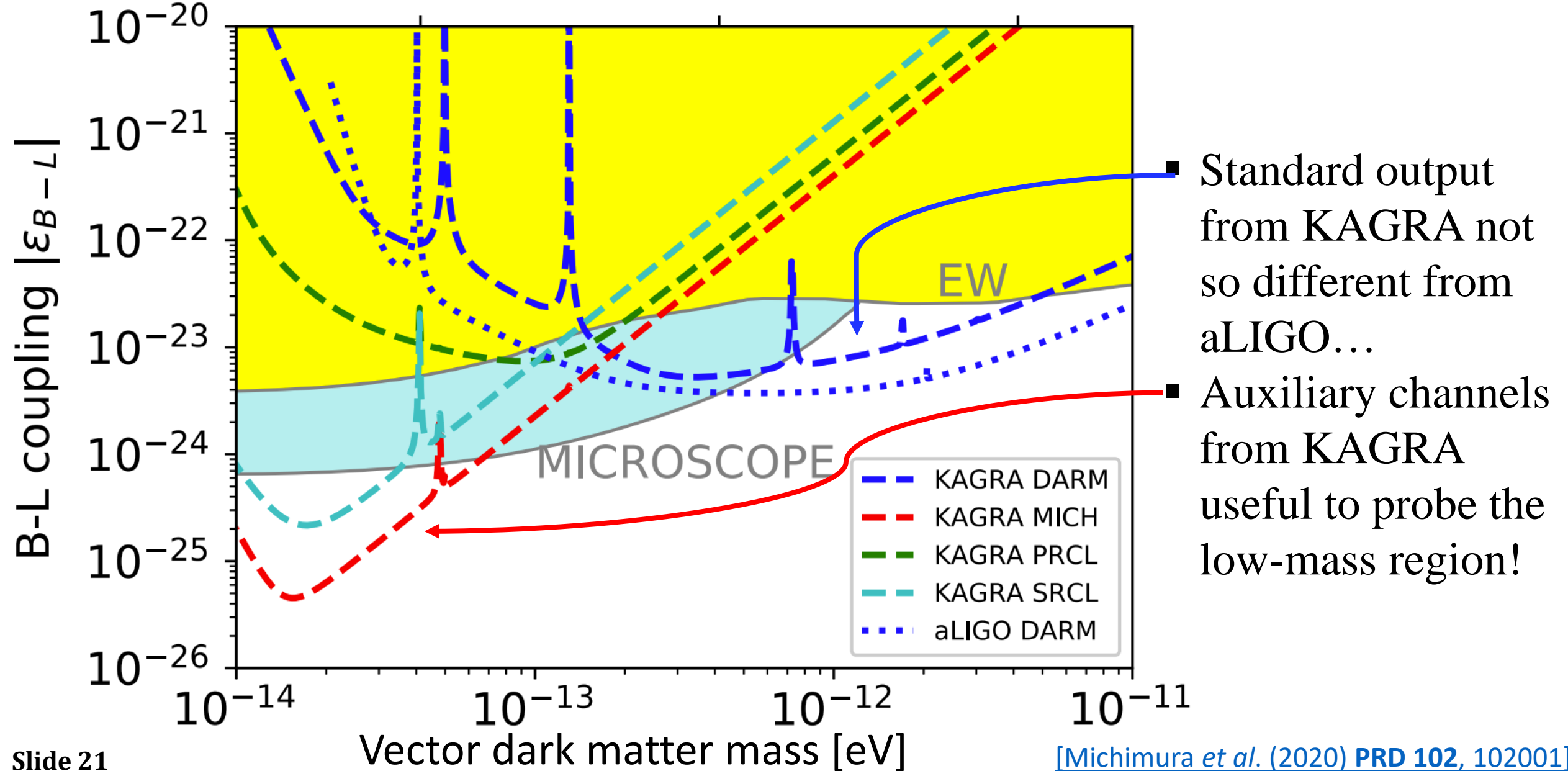
The peculiarity of KAGRA: hardware



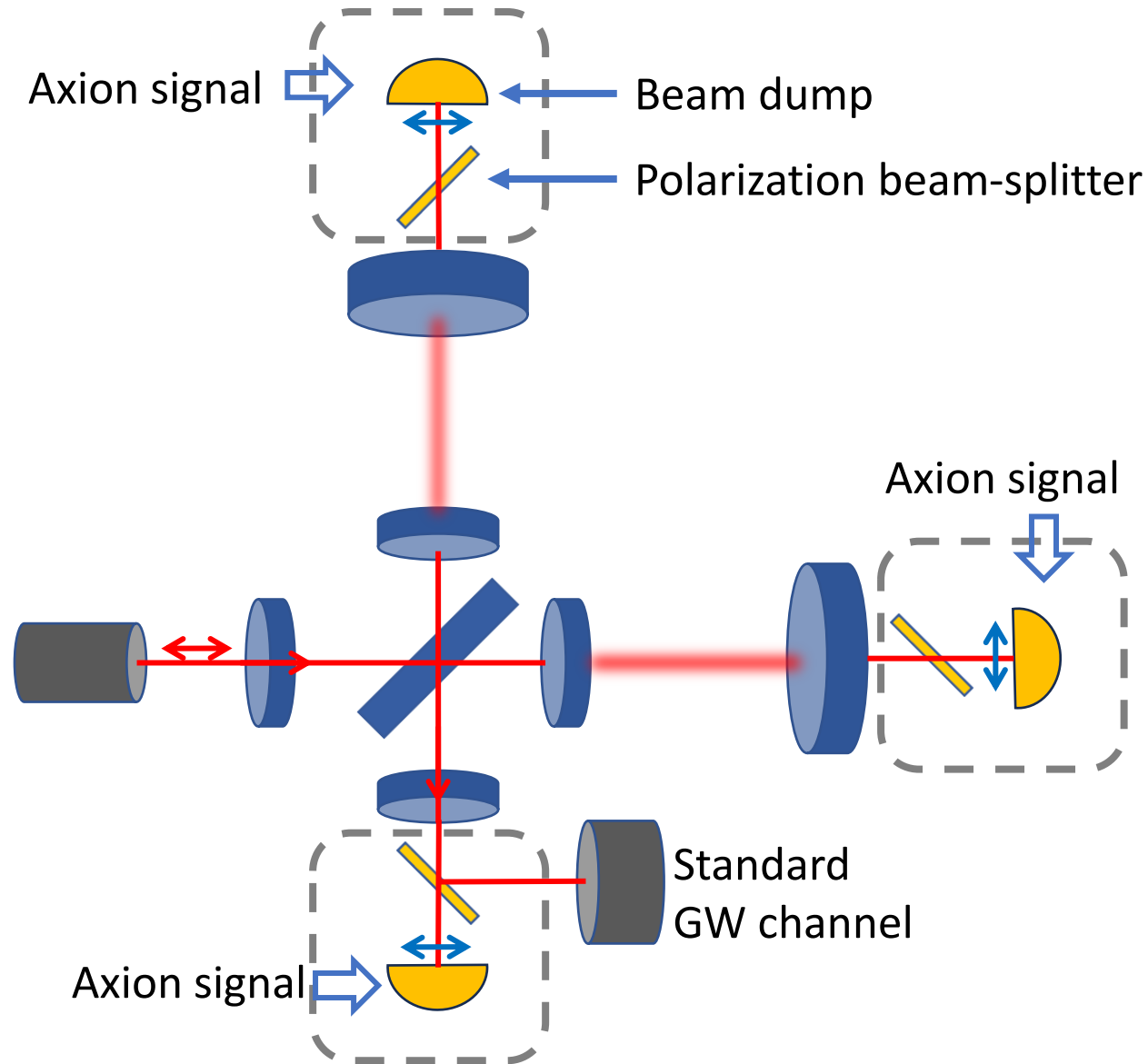
- End and Input mirrors are made by **sapphire** (and are cryogenic).
- Beam-splitter and recycling mirrors are made by **fused-silica**.
- The force on the optics is composition-dependent!
- The effect can be observed by the auxiliary channels, which monitor the intra-optics distances.
- Will provide meaningful results in the next observing runs.

[\[Michimura et al. \(2020\) PRD 102, 102001\]](#)

The peculiarity of KAGRA: future observations

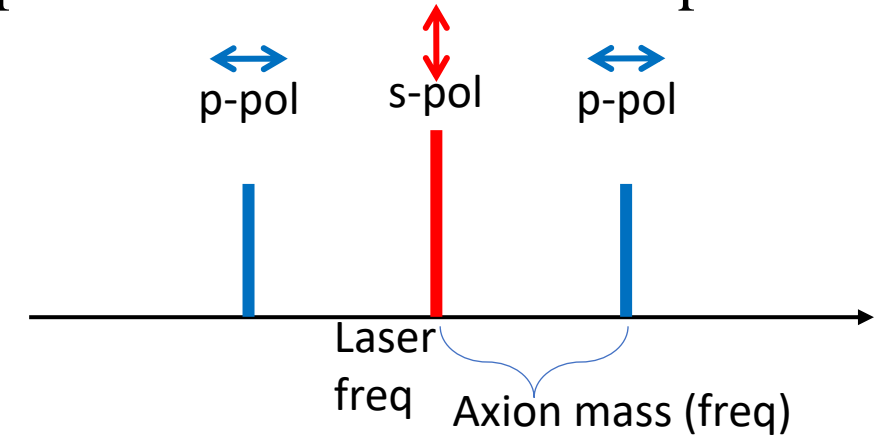


Direct interaction of axions in the detector



ADAM-GD scheme:

- Effect of axions coupling to photons in a cavity: induce modulated linear polarization from s-like to p-like.



- The signal accumulates over multiple, odd round trips.
- The detection would require additional optics at end masses and detection system.

Conclusions

- ❖ Interferometric, Earth-based gravitational-wave detectors have just opened a new window to observe the universe.
- ❖ LIGO, Virgo, KAGRA (and future detectors) have proven to be valid probes for ultra-light dark matter.
- ❖ A new observing run (O4) has just started on 24th May 2023
 - Improved sensitivity.
 - Longer observing time (1.5 years)
- ❖ Stay tuned!

THANK YOU!

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Direct interaction of axions in the detector

