

1st General Meeting of COST Action COSMIC WISPers (CA21106) *Bari - September 5-8, 2023*

Birefringence in CMB anisotropies due to cosmological pseudoscalar fields



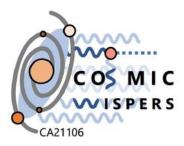
Vatican Observatory



Matteo Galaverni

Based on a work with:

Fabio Finelli and Daniela Paoletti (INAF/OAS Bologna & INFN Bologna) Phys.Rev.D 107 (2023) 8, 083529 (arXiv 2301.07971 [astro-ph.CO])



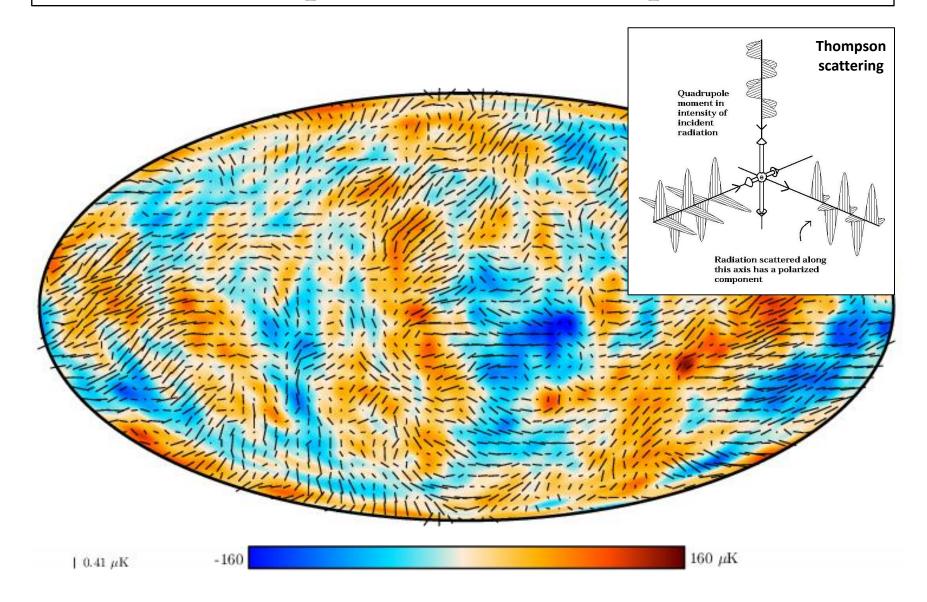
1st General Meeting of COST Action COSMIC WISPers (CA21106) *Bari - September 5-8, 2023*

Birefringence in CMB anisotropies due to cosmological pseudoscalar fields

Plan of the talk

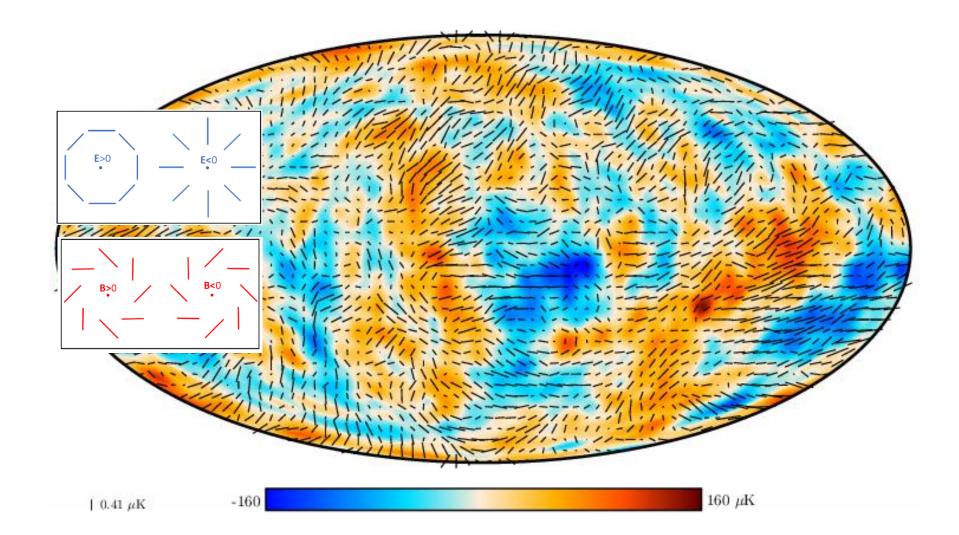
- CMB polarization anisotropies & cosmic birefringence;
- **Redshift/time evolution** of cosmic birefringence and CMB;
- Cosmological pseudoscalar field: Early Dark Energy (EDE),
 Quintessence (DE) or axion-like Dark Matter (DM);
- Conclusions.

CMB polarization anisotropies

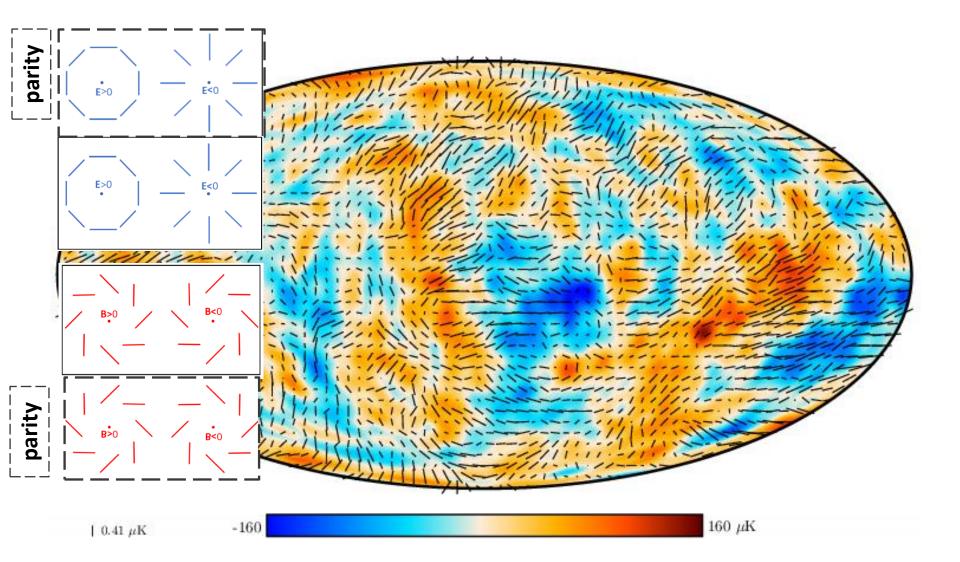


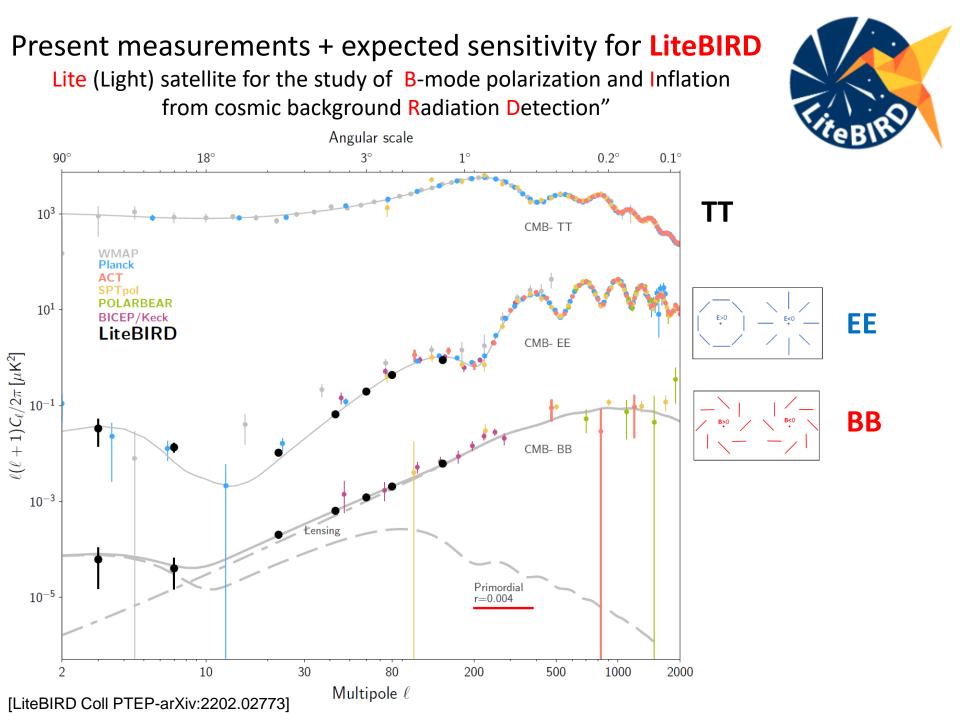
[2018 Planck map of the polarized CMB anisotropies]

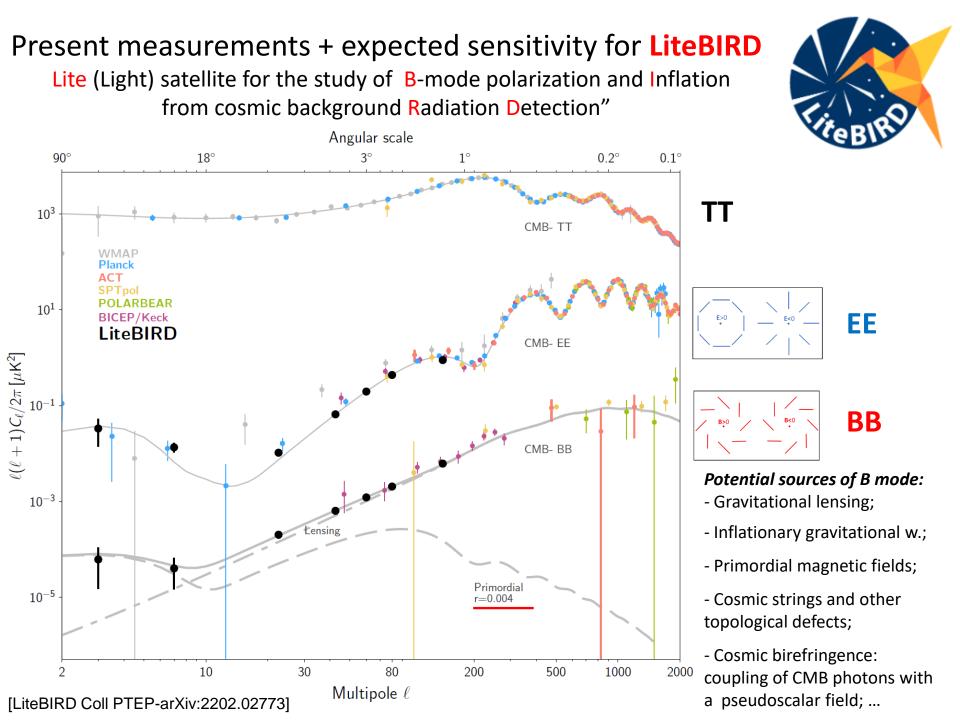
CMB polarization anisotropies

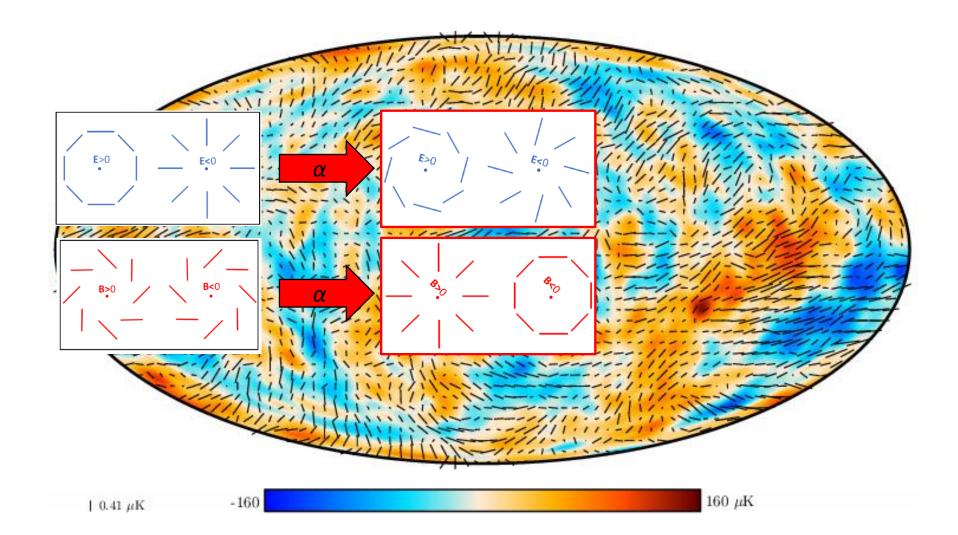


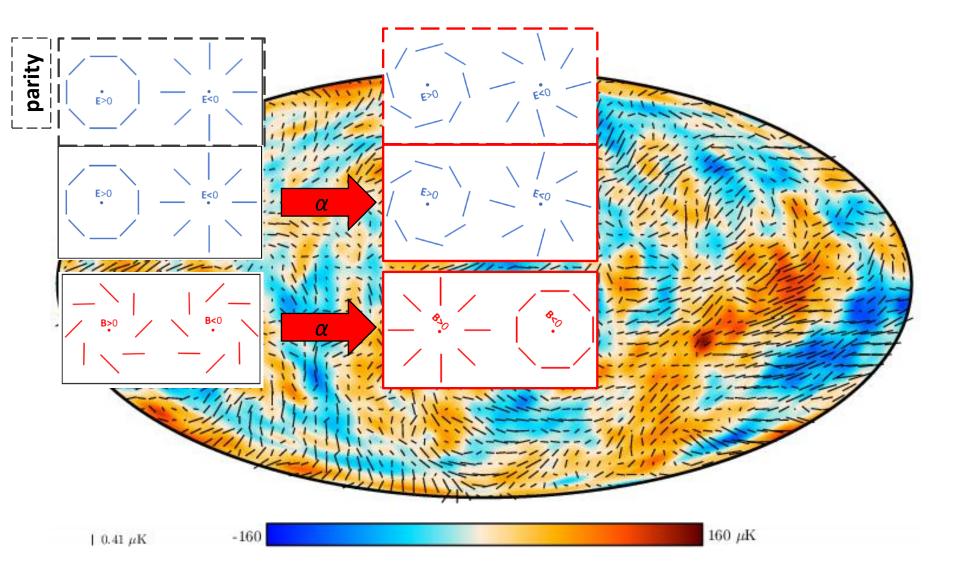
CMB polarization anisotropies

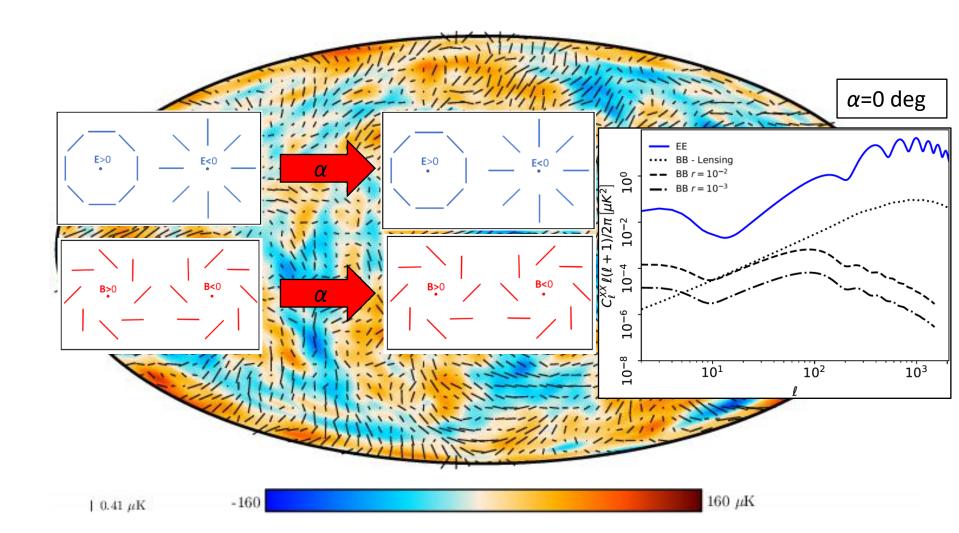


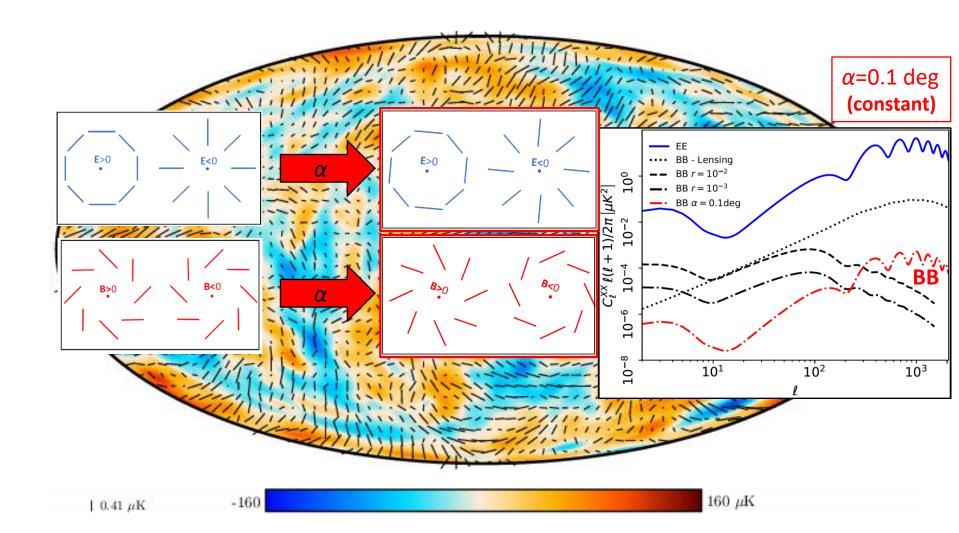


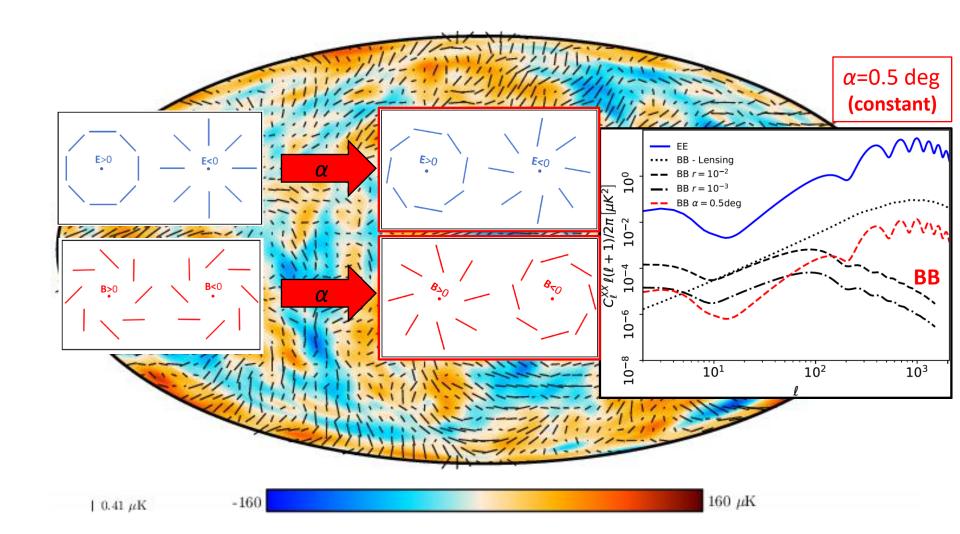


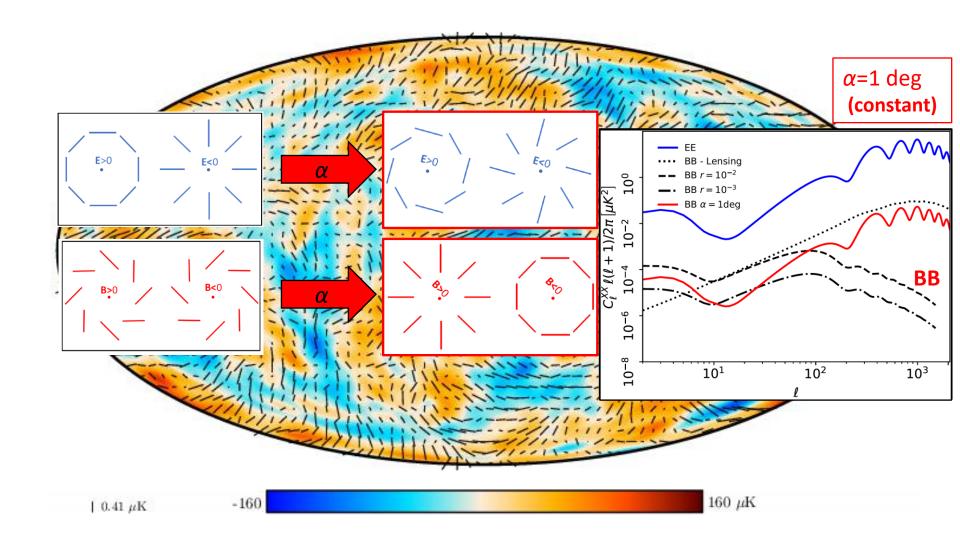




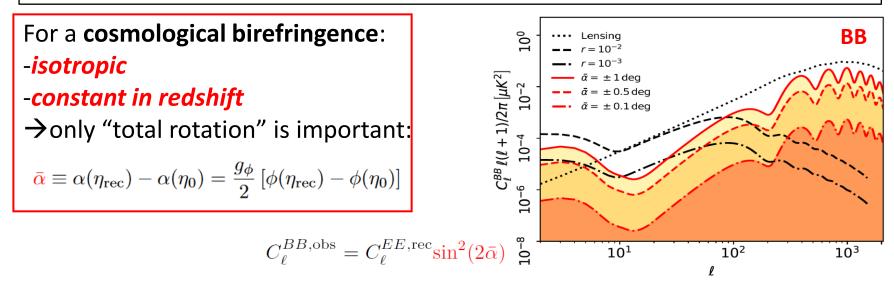




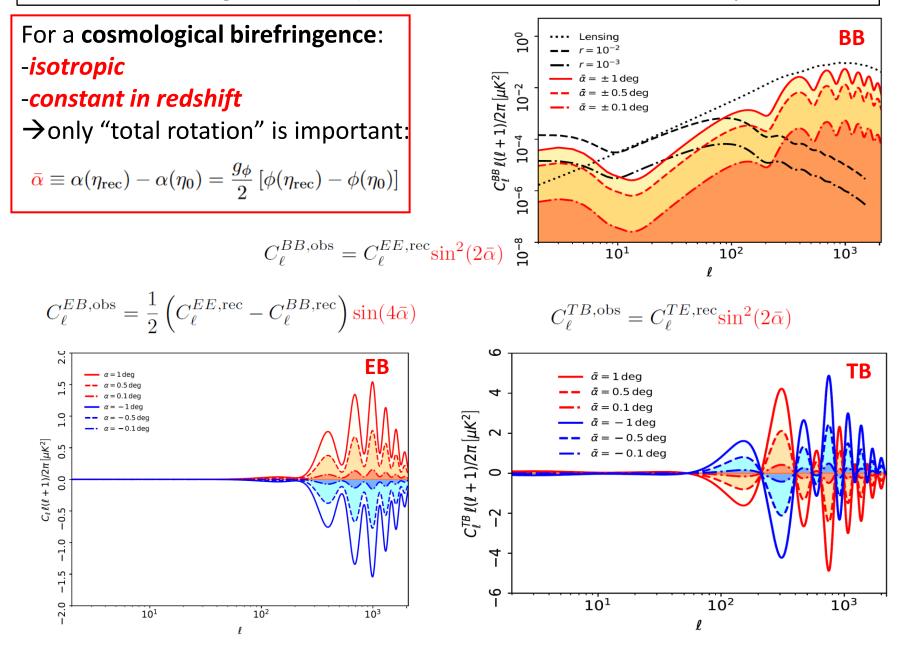




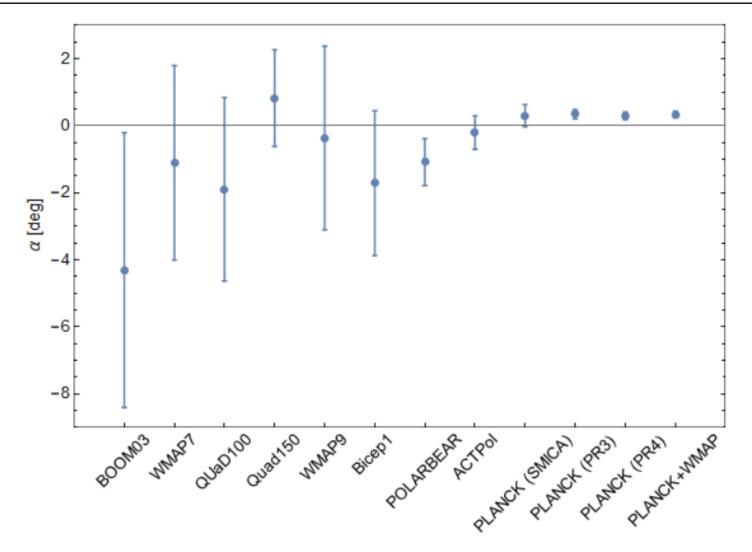
α [deg]: *isotropic* and *constant in redshift*



α [deg]: *isotropic* and *constant in redshift*

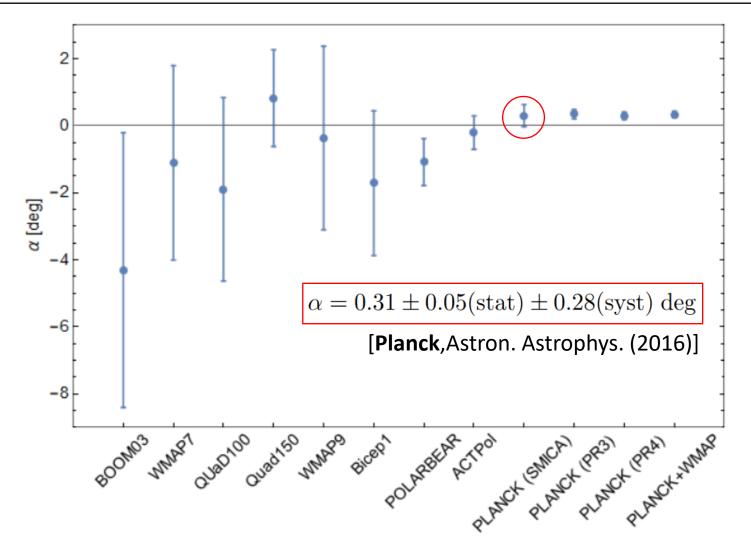


Constraints for α (*isotropic* and *constant in z*)



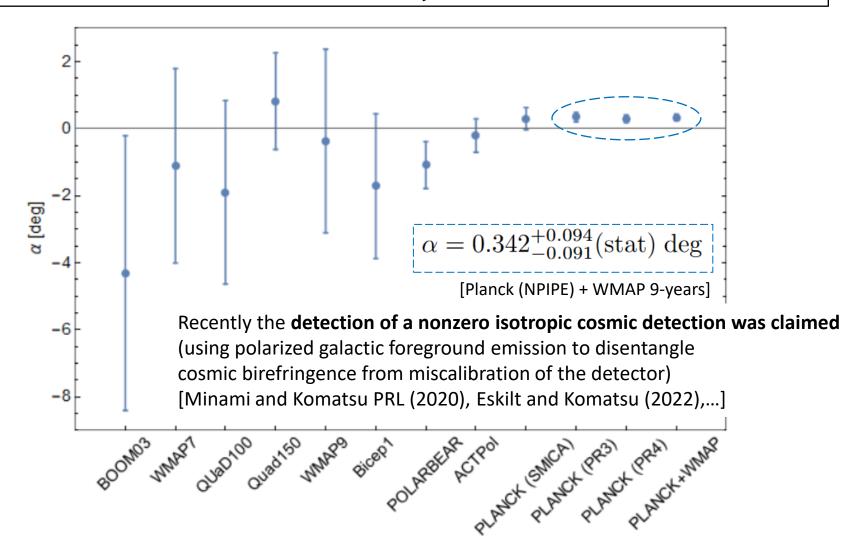
CMB polarization experiments constraints on isotropic cosmological birefringence as reviewed in Kaufman et al. (2016), Planck (2016), Greco et al. (2023), Williams (2023).

Constraints for α (*isotropic* and *constant in z*)



CMB polarization experiments constraints on isotropic cosmological birefringence as reviewed in Kaufman et al. (2016), Planck (2016), Greco et al. (2023), Williams (2023).

Constraints for α (*isotropic* and *constant in z*)



CMB polarization experiments constraints on isotropic cosmological birefringence as reviewed in Kaufman *et al.* (2016), Planck (2016), Greco *et al.* (2023), Williams (2023).

Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] \mp i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$

Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] \mp i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$

line-of-sight $C_{\ell}^{EE/BB/EB} = (4\pi)^2 \frac{9(\ell+2)!}{16(\ell-2)!} \int k^2 dk \left[\Delta_{E/B/E}(k,\eta_0) \Delta_{E/B/B}(k,\eta_0) \right]$ strategy [Seljak and Zaldarriaga (1996)]

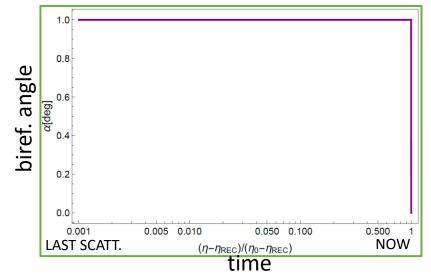
$$\Delta_{E}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \cos 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$
$$\Delta_{B}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \sin 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$

Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] \mp i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$

we follow the **line-of-sight** $C_{\ell}^{EE/BB/EB} = (4\pi)^2 \frac{9(\ell+2)!}{16(\ell-2)!} \int k^2 dk \left[\Delta_{E/B/E}(k,\eta_0) \Delta_{E/B/B}(k,\eta_0) \right]$ strategy [Seljak and Zaldarriaga (1996)]

$$\Delta_{E}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \cos 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$
$$\Delta_{B}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \sin 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$

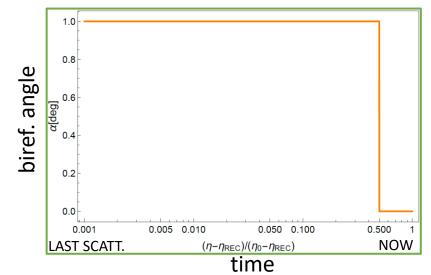


Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] \mp i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$

we follow the **line-of-sight** $C_{\ell}^{EE/BB/EB} = (4\pi)^2 \frac{9(\ell+2)!}{16(\ell-2)!} \int k^2 dk \left[\Delta_{E/B/E}(k,\eta_0) \Delta_{E/B/B}(k,\eta_0) \right]$ strategy [Seljak and Zaldarriaga (1996)]

$$\Delta_{E}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \cos 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$
$$\Delta_{B}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \sin 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$

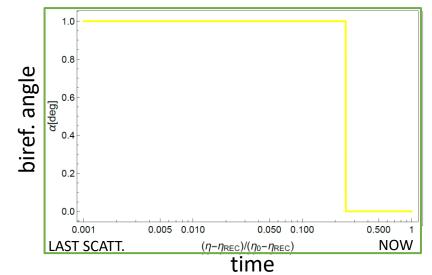


Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] = i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$
we follow the

line-of-sight $C_{\ell}^{EE/BB/EB} = (4\pi)^2 \frac{9(\ell+2)!}{16(\ell-2)!} \int k^2 dk \left[\Delta_{E/B/E}(k,\eta_0) \Delta_{E/B/B}(k,\eta_0) \right]$ strategy [Seljak and Zaldarriaga (1996)]

$$\Delta_{E}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \cos 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$
$$\Delta_{B}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \sin 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$

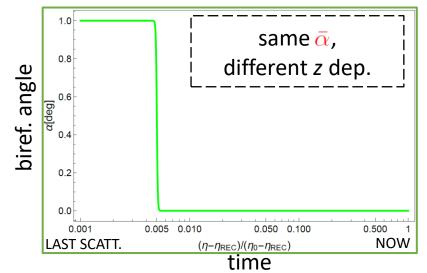


Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] \mp i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$

we follow the **line-of-sight** $C_{\ell}^{EE/BB/EB} = (4\pi)^2 \frac{9(\ell+2)!}{16(\ell-2)!} \int k^2 dk \left[\Delta_{E/B/E}(k,\eta_0) \Delta_{E/B/B}(k,\eta_0) \right]$ strategy [Seljak and Zaldarriaga (1996)]

$$\Delta_{E}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \frac{\cos 2 \left[\alpha(\eta)-\alpha(\eta_{0})\right]}{\cos 2 \left[\alpha(\eta)-\alpha(\eta_{0})\right]}$$
$$\Delta_{B}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta \, g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \frac{\sin 2 \left[\alpha(\eta)-\alpha(\eta_{0})\right]}{\sin 2 \left[\alpha(\eta)-\alpha(\eta_{0})\right]}$$

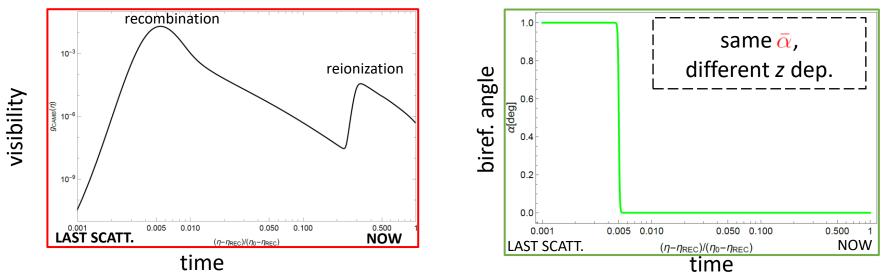


Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] \mp i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$

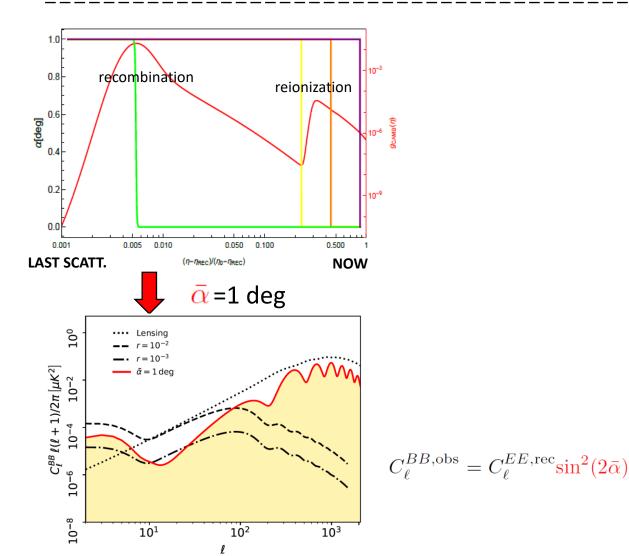
we follow the line-of-sight $C_{\ell}^{EE/BB/EB} = (4\pi)^2 \frac{9(\ell+2)!}{16(\ell-2)!} \int k^2 dk \left[\Delta_{E/B/E}(k,\eta_0) \Delta_{E/B/B}(k,\eta_0) \right]$ strategy [Seljak and Zaldarriaga (1996)]

$$\Delta_{E}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \cos 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$
$$\Delta_{B}(k,\eta_{0}) = \int_{0}^{\eta_{0}} d\eta g(\eta) S_{P}^{(0)}(k,\eta) \frac{j_{\ell}(k\eta_{0}-k\eta)}{(k\eta_{0}-k\eta)^{2}} \sin 2\left[\alpha(\eta)-\alpha(\eta_{0})\right]$$

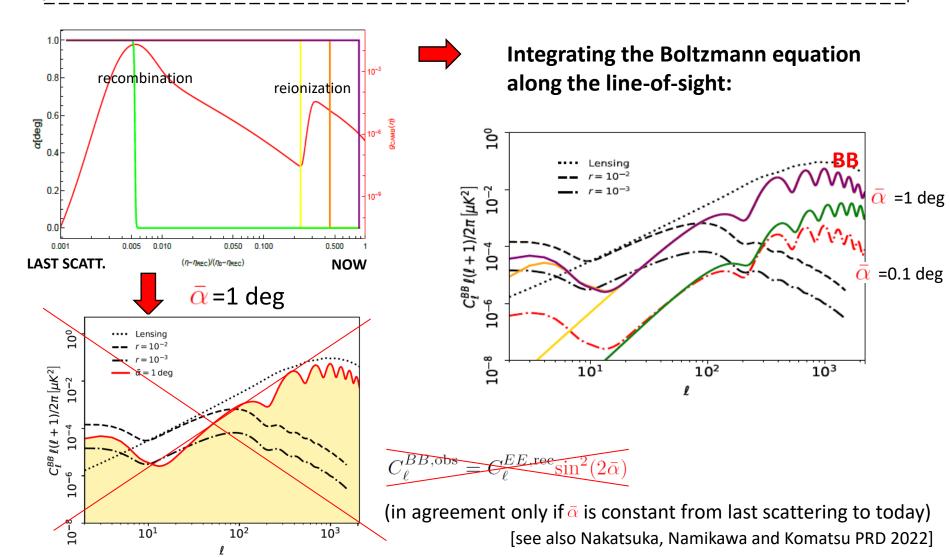


α [deg]: *isotropic* and *<u><i>z*</u> *dependent*

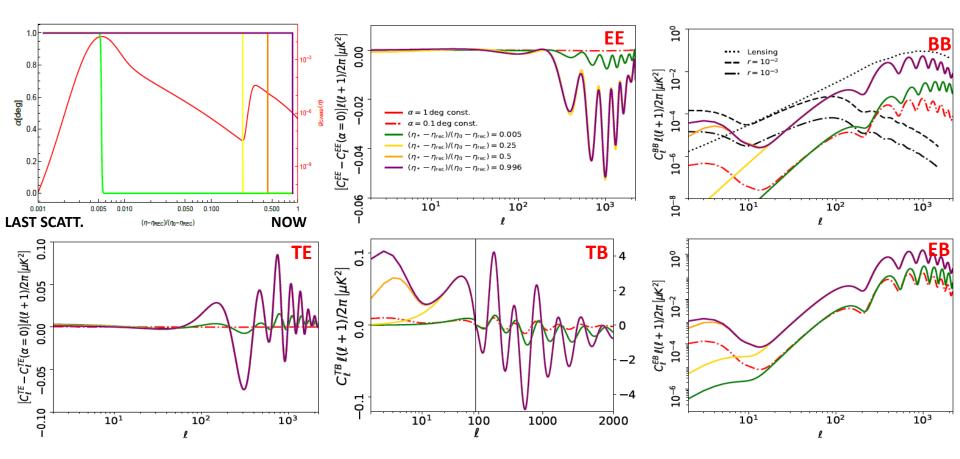
- same $\bar{\alpha} = 1 \deg$
- but linear polarization rotation happens at different epochs!



- same $\bar{\alpha} = 1 \deg$
- but linear polarization rotation happens at different epochs!



- same $\bar{\alpha}$ =1 deg
- but linear polarization rotation happens at different epochs!



Cosmological pseudoscalar field

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} \nabla_{\mu} \phi \nabla^{\mu} \phi - V(\phi) - \frac{g_{\phi}}{4} \phi F_{\mu\nu} \tilde{F}^{\mu\nu},$$

V(

Cosmological **pseudoscalar field** ϕ acting as:

- Dark Matter (e.g. axion-like particles);
- **Dark Energy** (e.g. ultralight pseudo Nambu-Goldstone bosons);
- Early Dark Energy.

Cosmological pseudoscalar field

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{2}\nabla_{\mu}\phi\nabla^{\mu}\phi - V(\phi) - \frac{g_{\phi}}{4}\phi F_{\mu\nu}\tilde{F}^{\mu\nu}$$

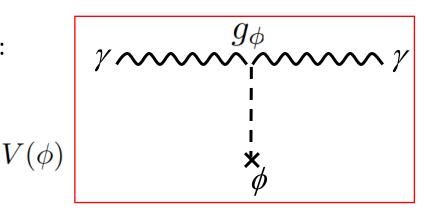
Cosmological **pseudoscalar field** ϕ acting as:

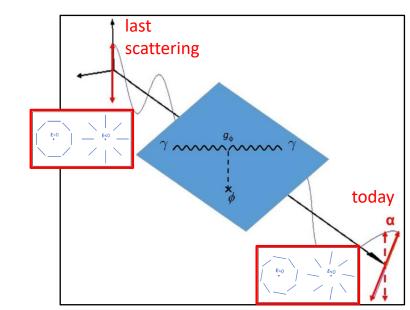
- **Dark Matter** (e.g. axion-like particles);
- Dark Energy (e.g. ultralight pseudo Nambu-Goldstone bosons);
- Early Dark Energy.

$$\alpha(x) = \frac{g_{\phi}}{2} \left[\phi(x) - \phi(x_{\rm em}) \right] \,,$$

rotation of the polarization plane (single photon)

Carrol, Field and Jackiw [PRD 1990], Harari and Sikivie [Phys. Lett. B 1992],...

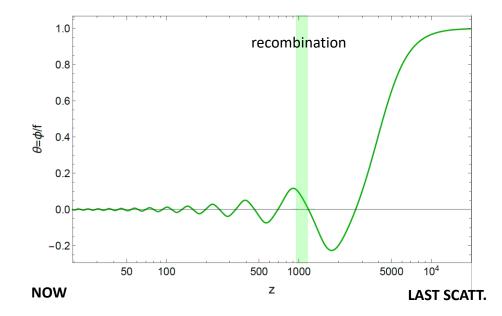


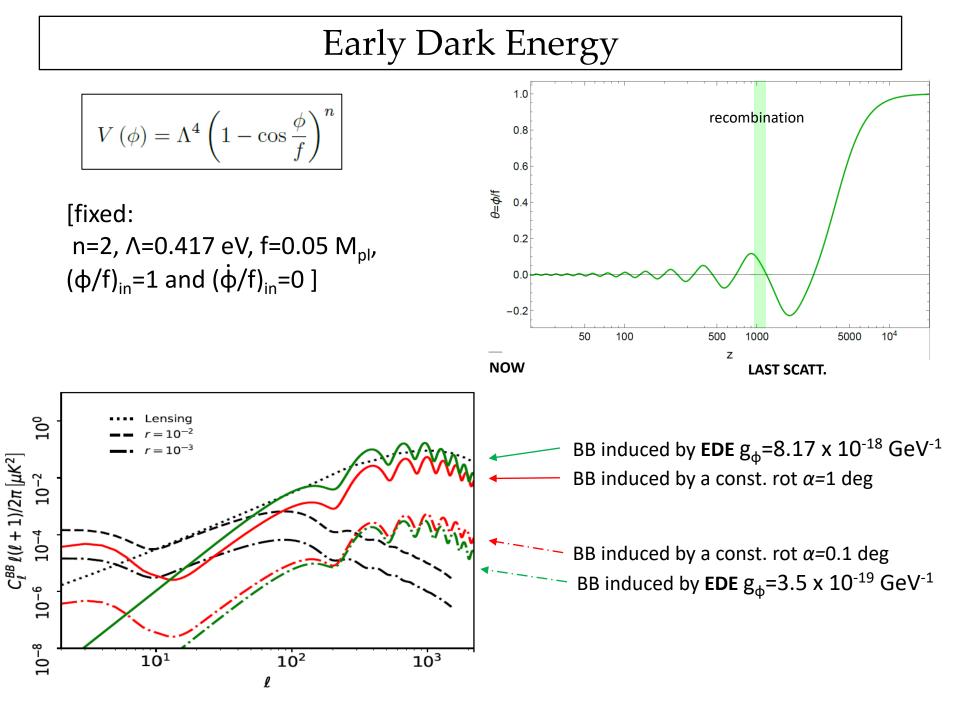


Early Dark Energy

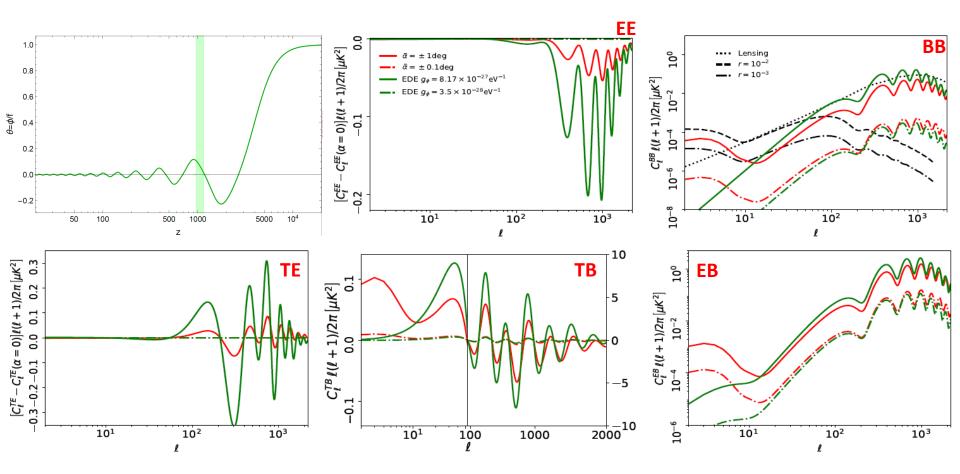
$$V\left(\phi\right) = \Lambda^4 \left(1 - \cos\frac{\phi}{f}\right)^n$$

[fixed: n=2, Λ =0.417 eV, f=0.05 M_{pl}, (ϕ/f)_{in}=1 and ($\dot{\phi}/f$)_{in}=0]





Early Dark Energy



Constraints for a LiteBIRD-like mission for EDE

The parity violating nature of the interaction generates **nonzero parity odd correlators (TB and EB)**, therefore we consider the **full theoretical covariance matrix**:

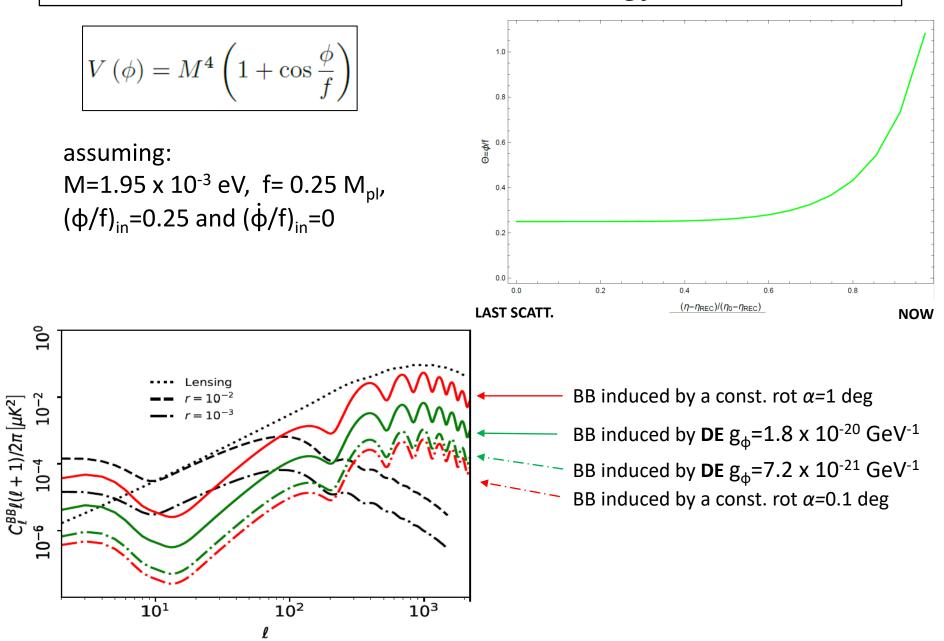
$$\bar{\mathbf{C}}_{l} = \begin{pmatrix} \bar{C}_{\ell}^{TT} & \bar{C}_{\ell}^{TE} & \bar{C}_{\ell}^{TB} \\ \bar{C}_{\ell}^{TE} & \bar{C}_{\ell}^{EE} & \bar{C}_{\ell}^{EB} \\ \bar{C}_{\ell}^{TB} & \bar{C}_{\ell}^{EB} & \bar{C}_{\ell}^{BB} \end{pmatrix}$$

Following Xia *et al.* [Astron. Astrophys. 2008] we introduce the effective χ^2

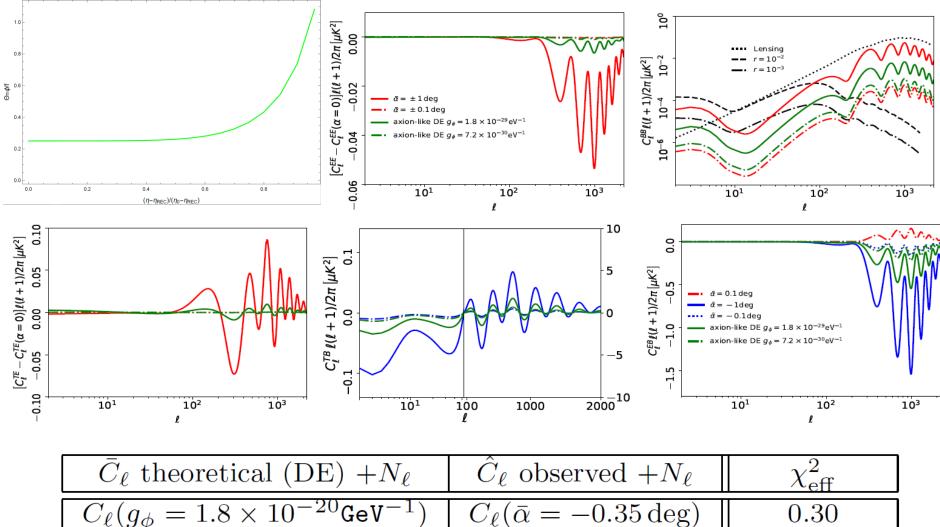
For Early Dark Energy:

\bar{C}_{ℓ} theoretical (EDE) $+N_{\ell}$	\hat{C}_{ℓ} observed $+N_{\ell}$	$\chi^2_{ m eff}$
$C_{\ell}(g_{\phi} = 1.65 \times 10^{-18} \text{GeV}^{-1})$	$C_{\ell}(\bar{\alpha} = 0.35 \mathrm{deg})$	67.3
$C_{\ell}(g_{\phi} = 6.0 \times 10^{-20} {\rm GeV}^{-1})$	$C_{\ell}(\alpha = 0 \deg)$	10.5

Axion-like Dark Energy

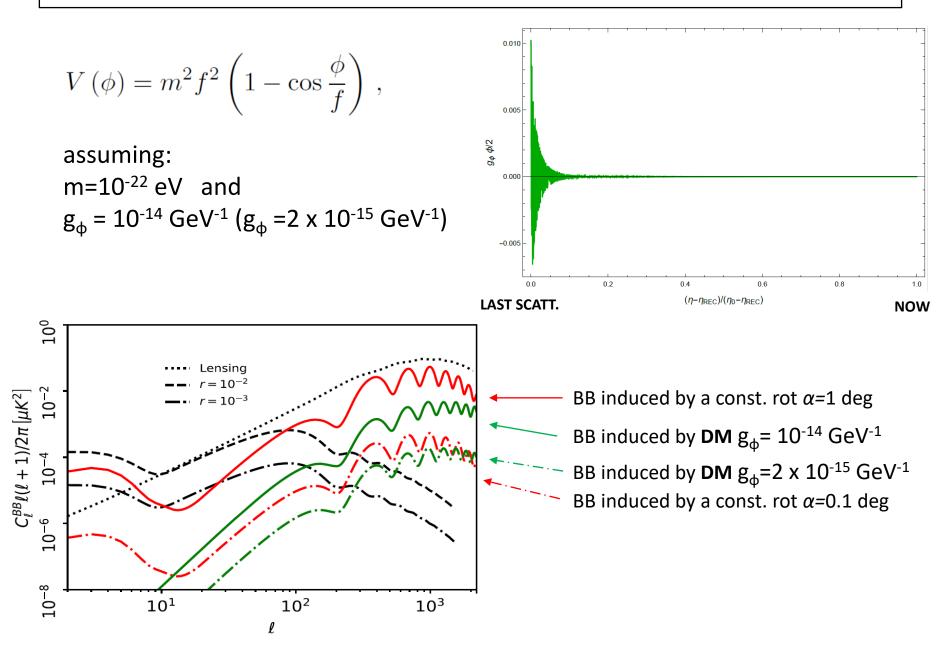


Constraints for a LiteBIRD-like mission for DE

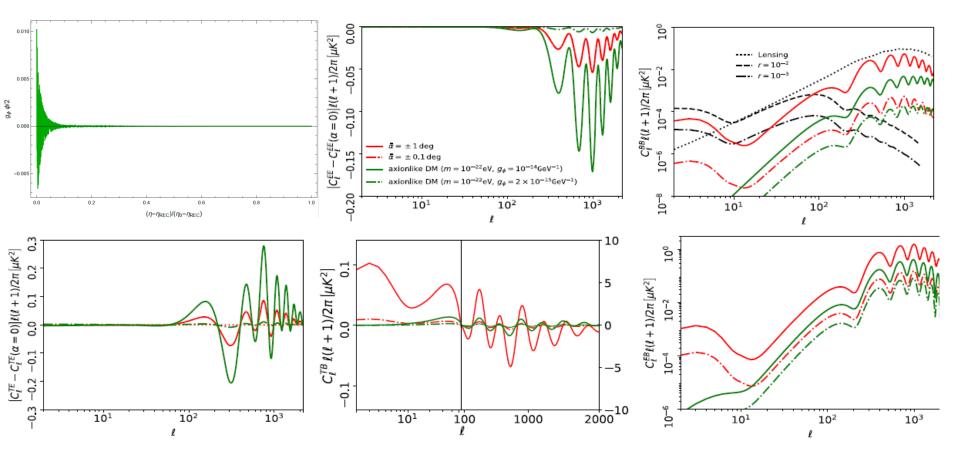


$C_{\ell}(g_{\phi} = 1.8 \times 10^{-5})$		$C_{\ell}(\bar{\alpha} = -0.35 \deg)$	0.30
$C_{\ell}(g_{\phi} = 1.8 \times 10^{-5})$	$^{-20}$ GeV $^{-1})$	$C_{\ell}(\alpha = 0 \deg)$	3.78×10^{3}
$C_{\ell}(g_{\phi} = 9.0 \times 10^{-1})$	$^{-22}$ GeV $^{-1})$	$C_{\ell}(\alpha = 0 \deg)$	9.4

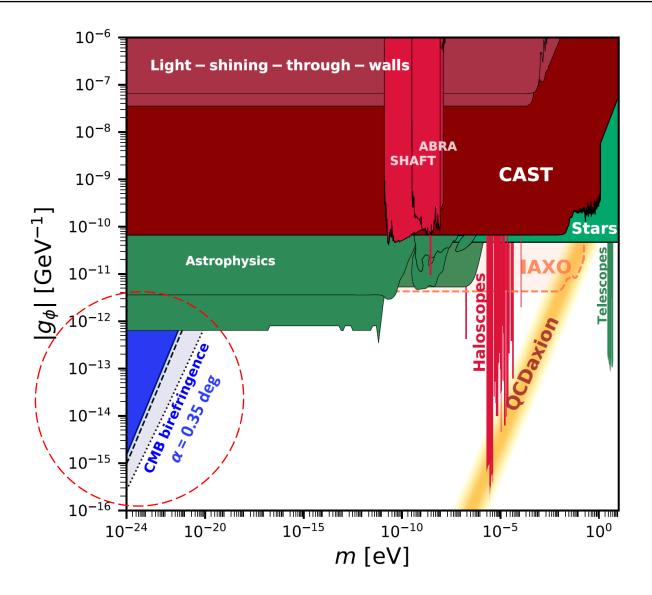
Axion-like Dark Matter



Constraints for a LiteBIRD-like mission for DM



Constraints for axion-like DM



[M.G., F. Finelli and D. Paoletti Phys.Rev.D 107 (2023) 8, 083529; plot created with **AxionLimits** code]

Conclusions

We studied the imprints of an **isotropic** <u>redshift-dependent</u> pseudoscalar field on CMB polarization power spectra for phenomenological and physical motivated models.

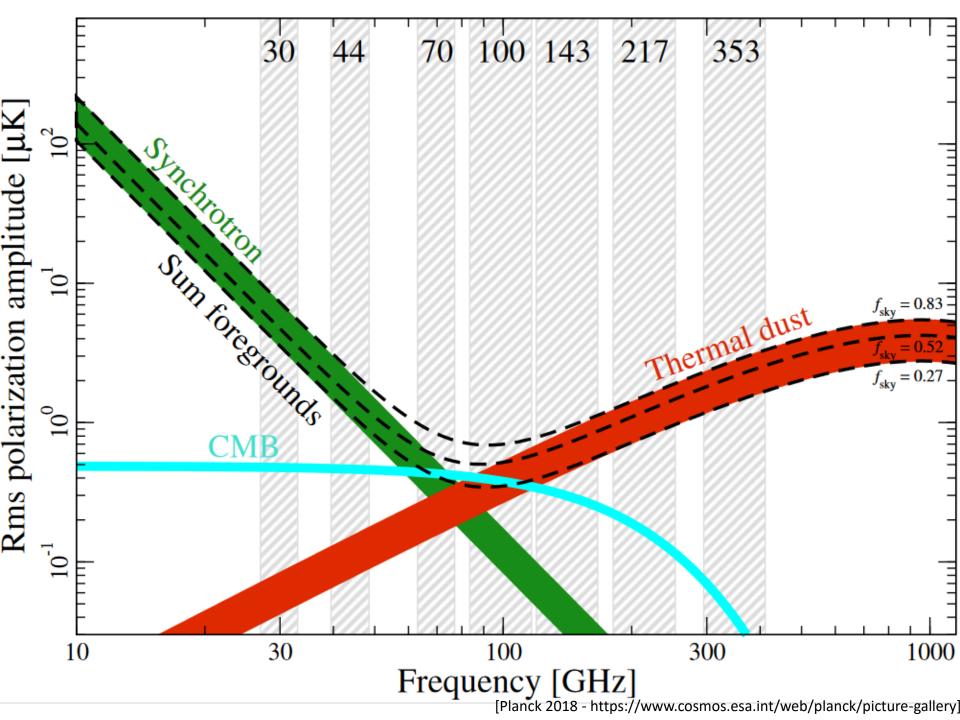
ightarrowRedshift evolution of the birefringence angle

has important effects on CMB polarization power spectra:

- not only $\bar{\alpha} \equiv \alpha(\eta_{rec}) \alpha(\eta_0)$ is important , but also **when** the rotation occurs;
- different theoretical motivated redshift dependencies of the pseudoscalar field (EDE, DE, DM) produce different multipole dependence for the CMB polarization power spectra;
- isotropic birefringence and not only anisotropic one can produce different multipole dependencies of the power spectra, this can be used to break the degeneracy with the miscalibration angle of the detector.

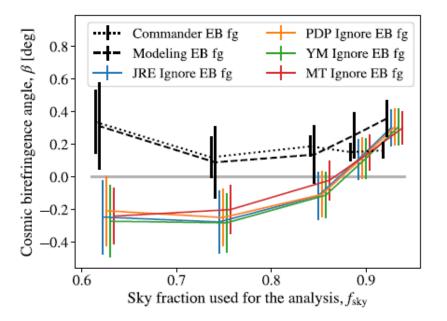
For more details see:

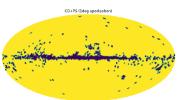
Redshift evolution of cosmic birefringence in CMB anisotropies M.G., F. Finelli and D. Paoletti Phys.Rev.D 107 (2023) 8, 083529 (arXiv 2301.07971 [astro-ph.CO])



Cosmic Birefringence from the *Planck* Data Release 4

P. Diego-Palazuelos,^{1,2,*} J. R. Eskilt,^{3,†} Y. Minami[®],⁴ M. Tristram[®],⁵ R. M. Sullivan[®],⁶ A. J. Banday,^{7,8} R. B. Barreiro[®],¹ H. K. Eriksen[®],³ K. M. Górski,^{9,10} R. Keskitalo[®],^{11,12} E. Komatsu[®],^{13,14} E. Martínez-González,¹ D. Scott[®],⁶ P. Vielva[®],¹ and I. K. Wehus[®]





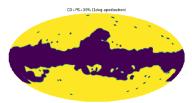


FIG. 1. Constraints on β for various values of f_{sky} with and without accounting for the foreground *EB* correlations. For the former, the dashed and dotted lines show corrections using the filament model [Eq. (2)] and the COMMANDER sky model, respectively. For the latter, the results of four pipelines (JRE, PDP, YM, MT) are shown.

Boltzmann equation for linear polarization with cosmic birefringence (Liu, Lee and Ng [PRL 2006], Finelli and Galaverni [PRD 2009]):

$$\Delta_{Q\pm iU}'(k,\eta) + ik\mu\Delta_{Q\pm iU}(k,\eta) = -n_e\sigma_T a(\eta) \left[\Delta_{Q\pm iU}(k,\eta) + \sum_m \sqrt{\frac{6\pi}{5}} \pm 2Y_2^m S_P^{(m)}(k,\eta)\right] \mp i2\alpha'(\eta)\Delta_{Q\pm iU}(k,\eta)$$
we follow the

line-of-sight $C_{\ell}^{EE/BB/EB} = (4\pi)^2 \frac{9(\ell+2)!}{16(\ell-2)!} \int k^2 dk \left[\Delta_{E/B/E}(k,\eta_0) \Delta_{E/B/B}(k,\eta_0) \right]$ strategy [Seljak and Zaldarriaga (1996)]

