

Unveiling dark fifth forces with Large Scale Structures

Salvatore Bottaro

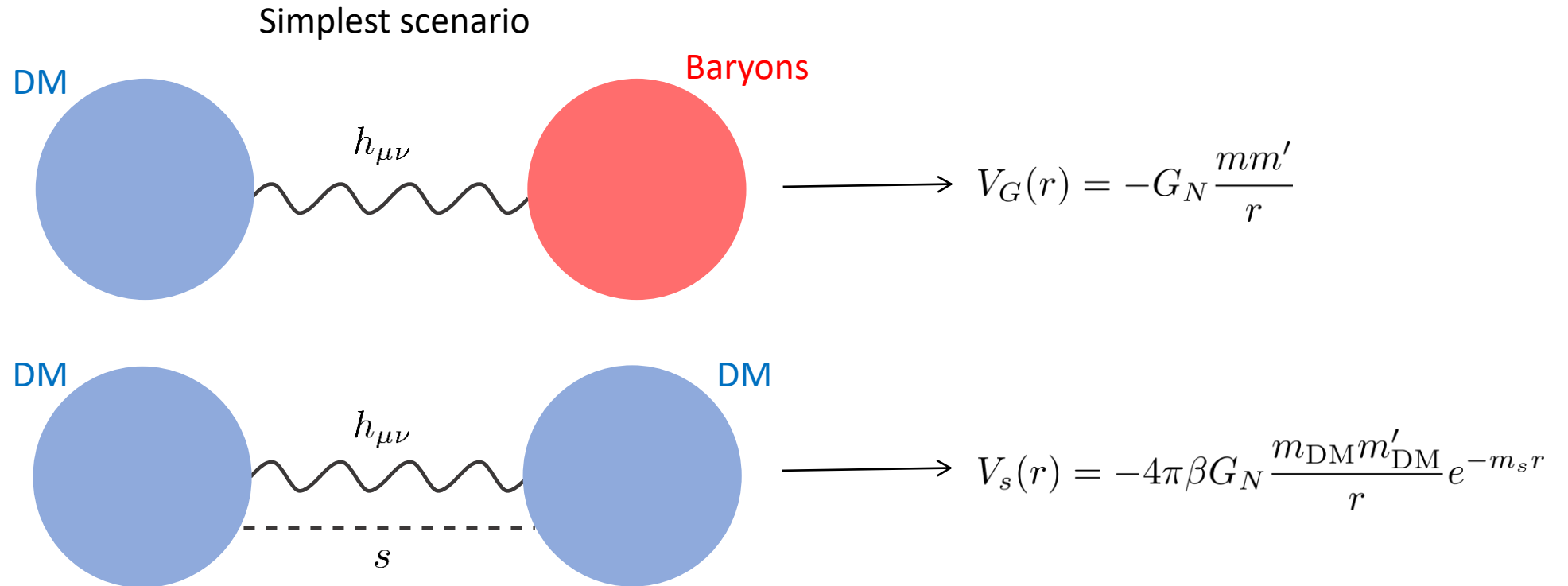
In collaboration with: E. Castorina, M. Costa, D. Redigolo, E. Salvioni



COSMIC WISPerS - September 5, 2023

Fifth forces in the Dark Sector

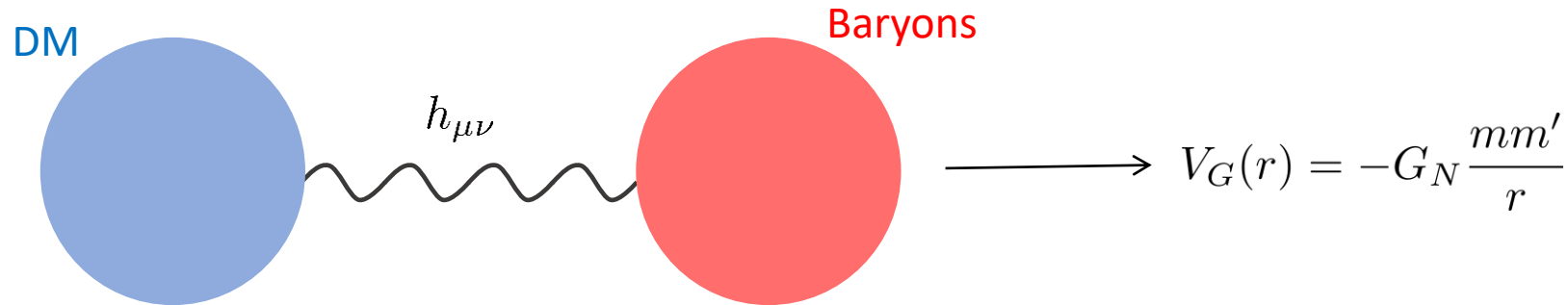
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- Sensibly more precision will be reached with present and future galaxy surveys (DESI, EUCLID...)



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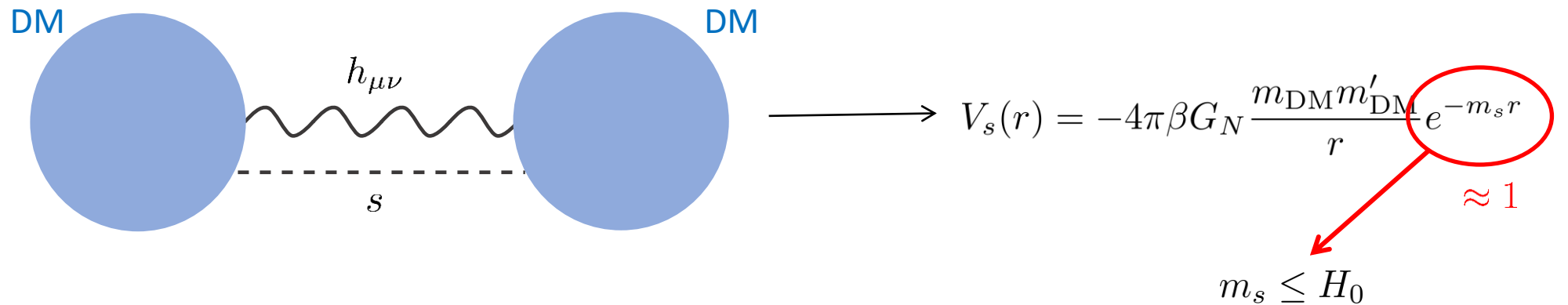
Simplest scenario



Example

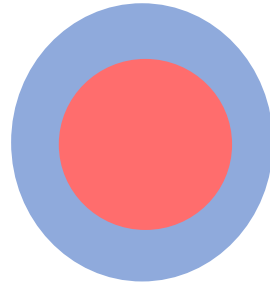
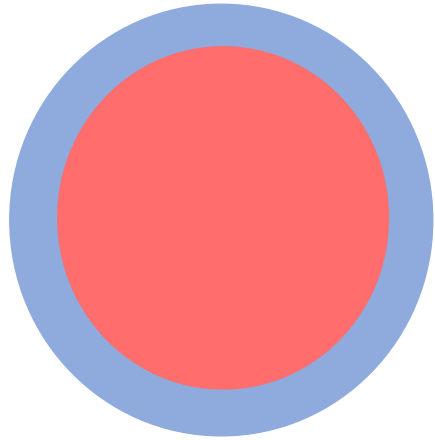
$$\mathcal{L}_{\text{int}} \propto s\chi^2$$

Space-time
dependent
DM mass!



Fifth forces in the Dark Sector

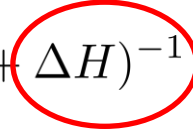
Different bkg evolution: modified distances $d(z) \propto H^{-1} = (H_{\Lambda\text{CDM}} + \Delta H)^{-1}$



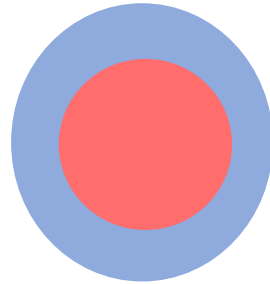
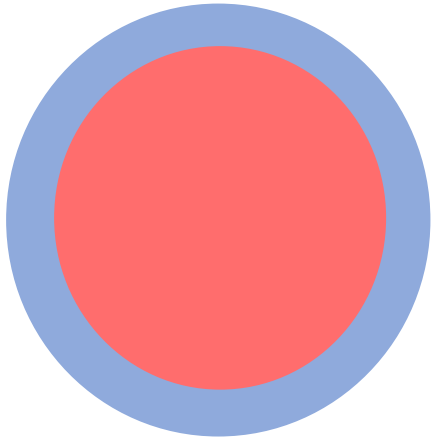
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From $m_{\text{DM}}(s)$
Typically overlooked
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Faster growth of matter fluctuations

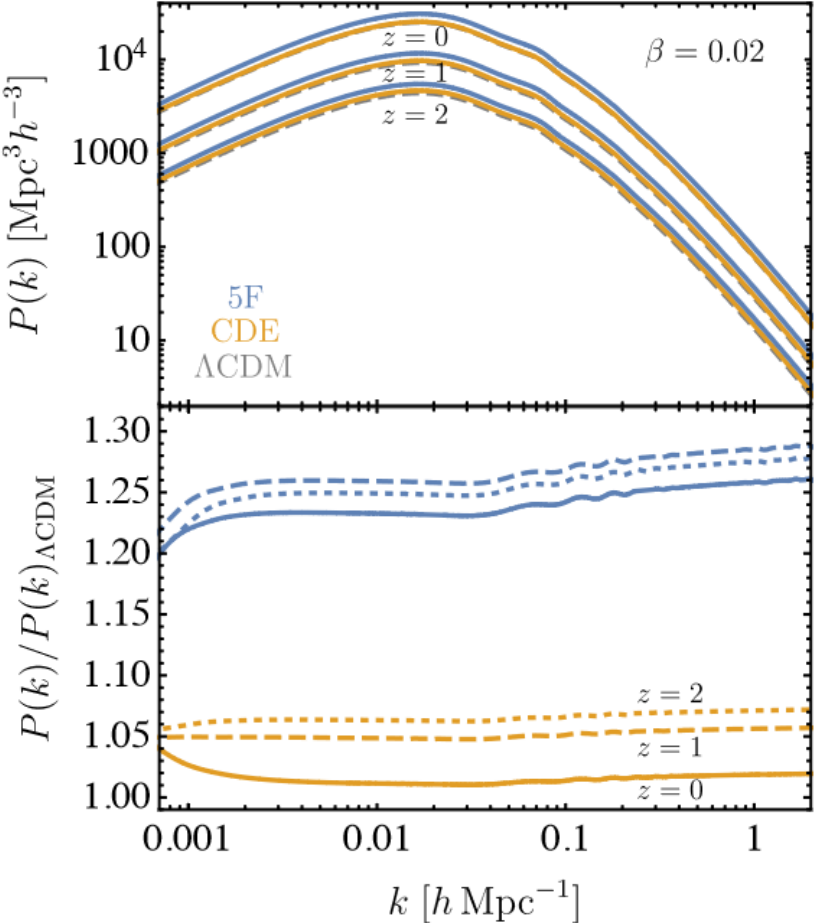
$$\delta_m(a) = D_{m,\Lambda\text{CDM}}(a) \left(1 + \frac{6}{5} \beta \tilde{m}^2 f_{\text{DM}}^2 \log \frac{a}{a_{\text{eq}}} \right) \delta_m(a_{\text{eq}})$$

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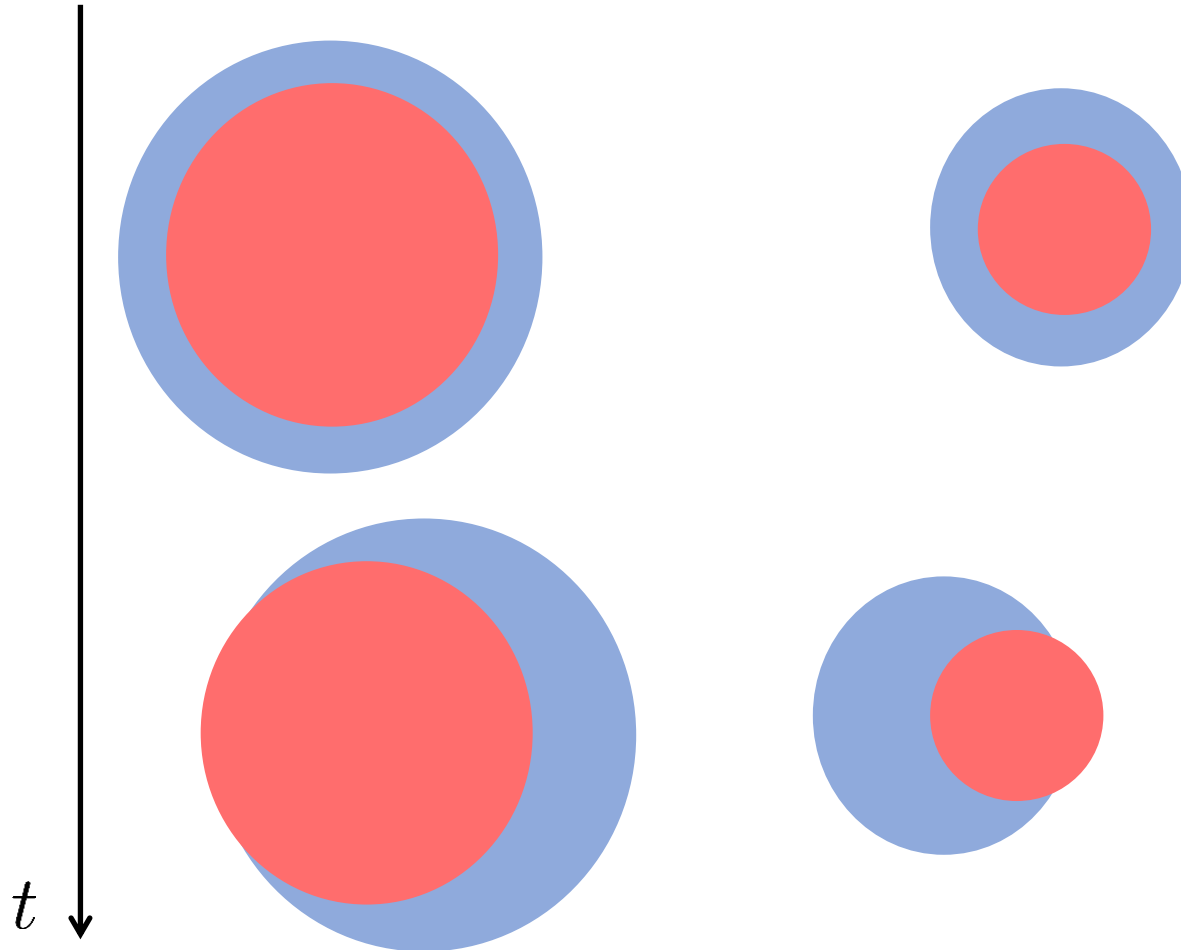
Log-enhanced growth

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EP violation: non-trivial evolution of
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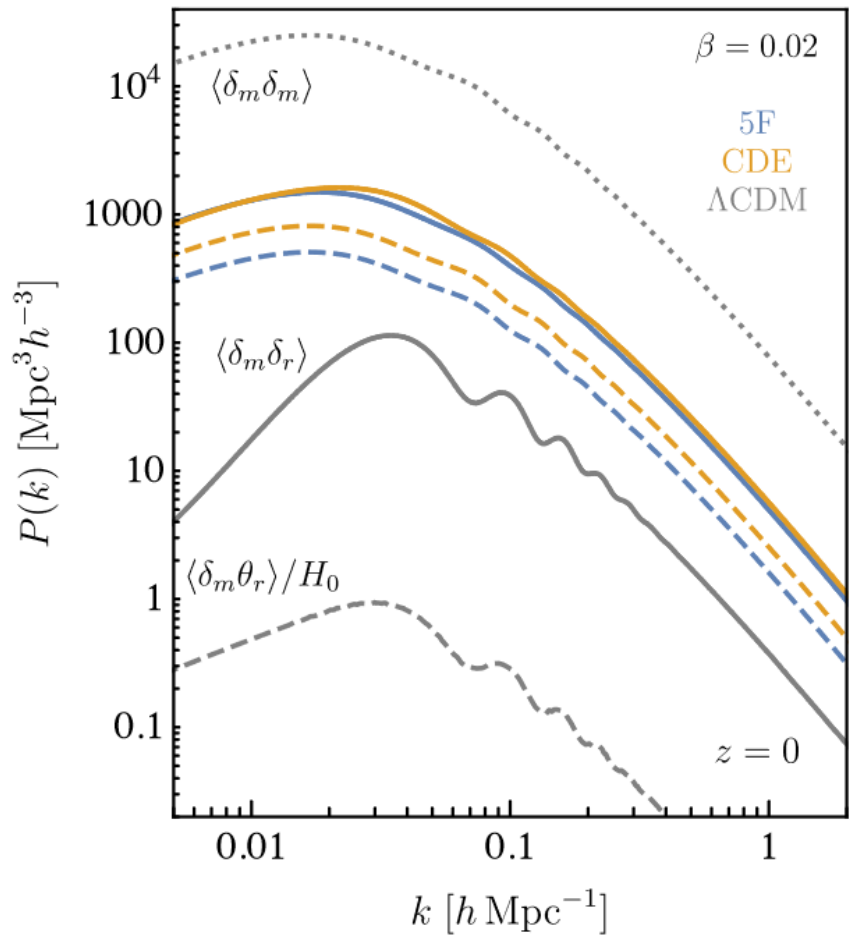
$$\delta_r(a) = \frac{5}{3} \beta \tilde{m}^2 f_{\text{DM}} \delta_{m,\Lambda\text{CDM}}(a)$$

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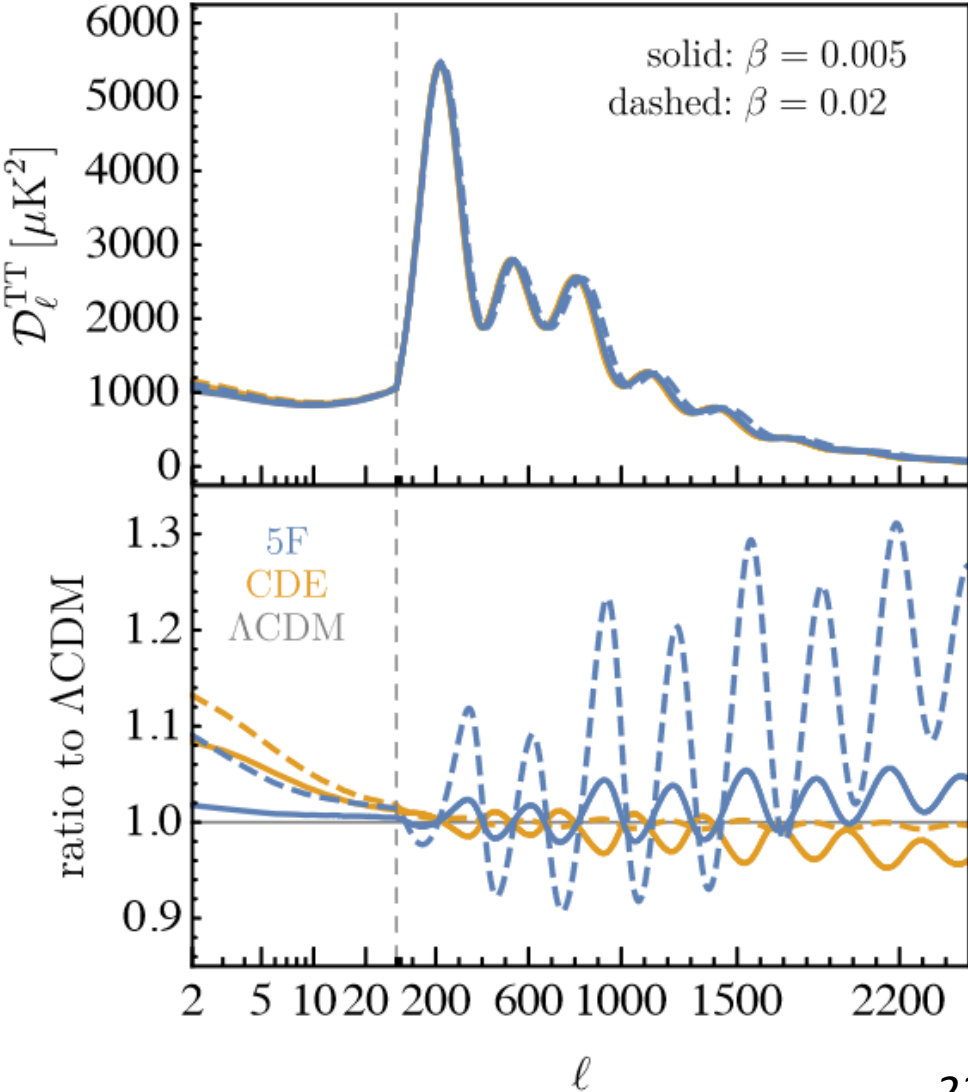
Effects on linear cosmology

CMB power spectrum mostly affected by bkg

$$\beta \tilde{m}^2 f_{\text{DM}}^2 \log \frac{a_{\text{rec}}}{a_{\text{eq}}} \approx \beta \ll 1$$

Shift in the peaks from modified angular diameter distance

$$l_n \approx \frac{n\pi}{c_s t_{\text{rec}}} D_A(z_{\text{rec}}) \propto \int_0^{z_{\text{rec}}} \frac{dz}{H_{\Lambda\text{CDM}}(z) + \Delta H(z)}$$



Effects on LSS

Problem: we observe galaxies, which track dark matter fluctuations \longrightarrow *Bias expansion*

Effects on LSS

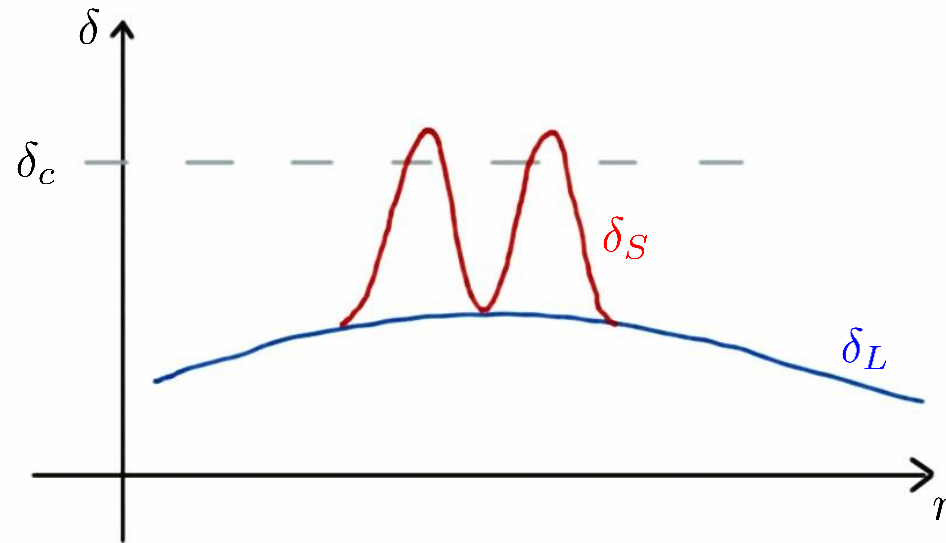
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$$\delta_g(\vec{k}) = b_1 \delta_m(\vec{k}) + b_r \delta_r(\vec{k}) + \dots$$

Fluctuation of the galaxy
number density

Bias parameters

$$b_i = \frac{1}{\bar{n}_g} \frac{d\bar{n}_g}{d\delta_i}$$



Effects on LSS

Power spectrum $\langle \delta_g(\vec{k}) \delta_g(\vec{k}') \rangle = (2\pi)^3 \delta^{(3)}(\vec{k} + \vec{k}') P(k)$

$$P(k) = \left(1 + \frac{6}{5} f_{\text{DM}}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\text{eq}}}\right)^2 P_{\Lambda\text{CDM}}^{\text{lin}}(k) + \left(1 + \frac{6}{5} f_{\text{DM}}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\text{eq}}}\right)^4 P_{\Lambda\text{CDM}}^{1\text{-loop}}(k) + f_{\text{DM}} \tilde{m}^2 \beta \Delta P(k)$$

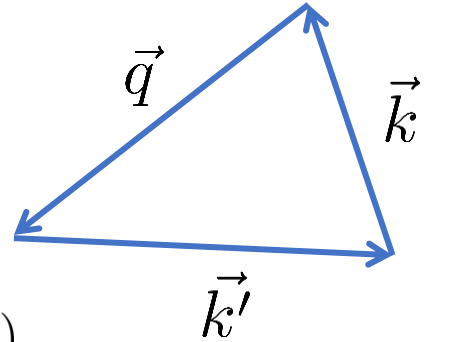
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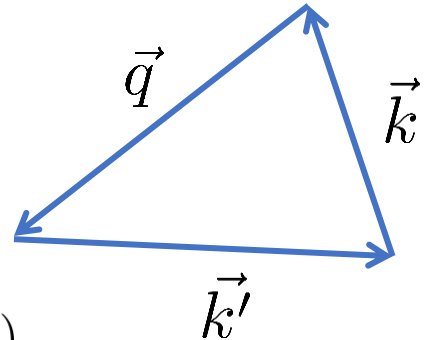
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- From $\delta_m \subset \delta_g$
- Log-enhancement from bkg effects
- Typically overlooked in literature

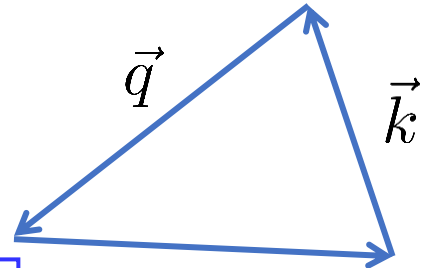
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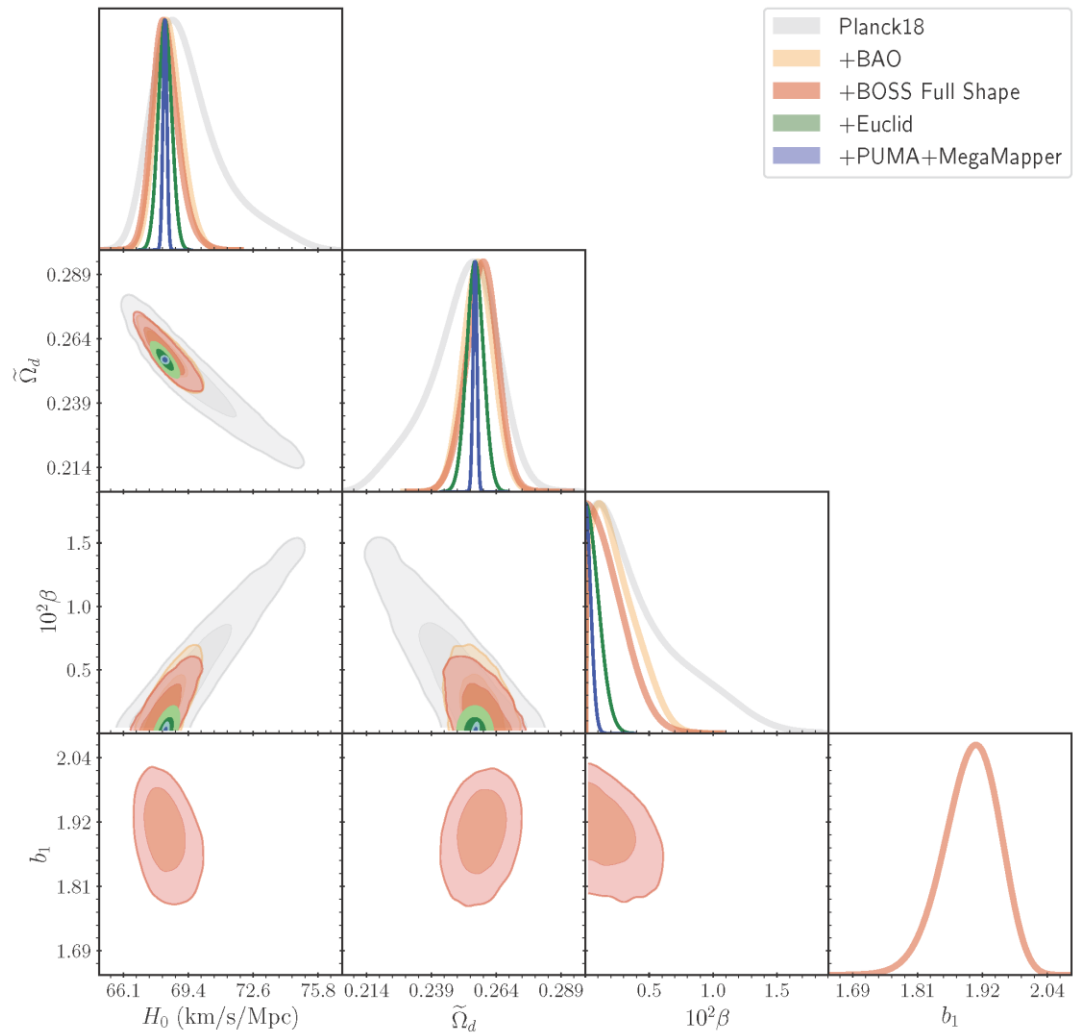
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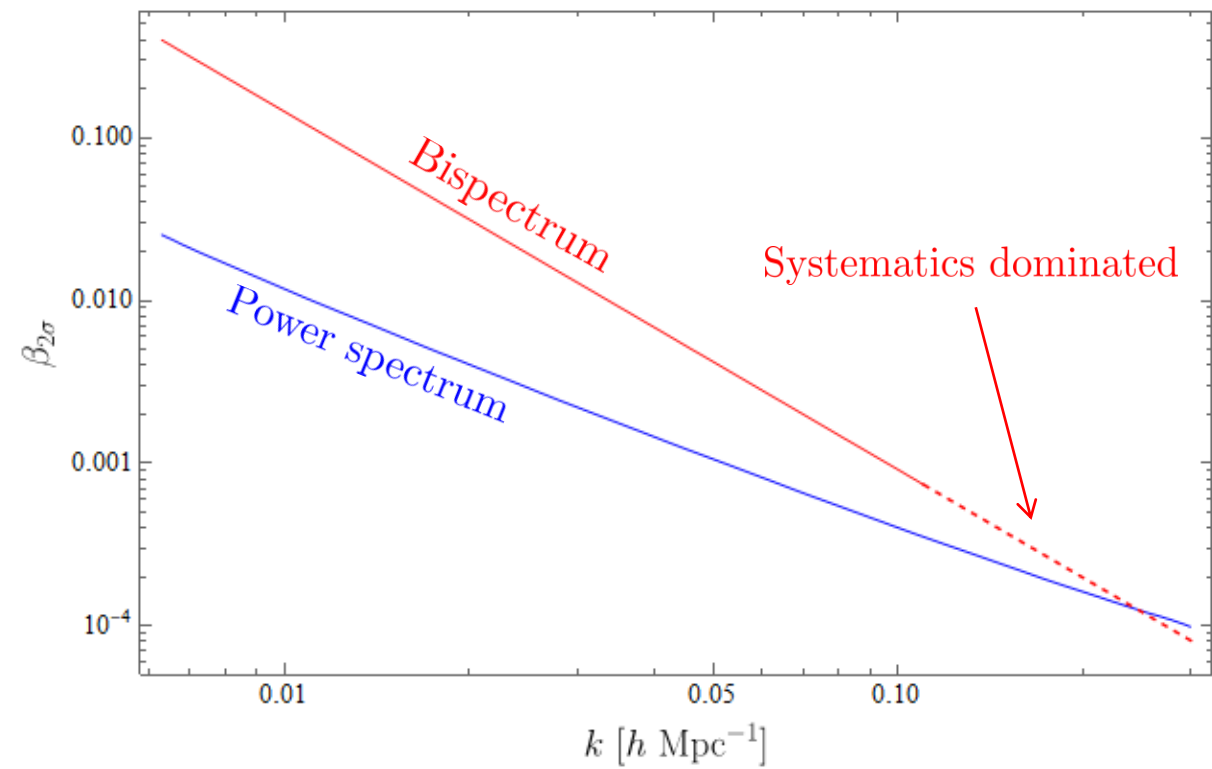
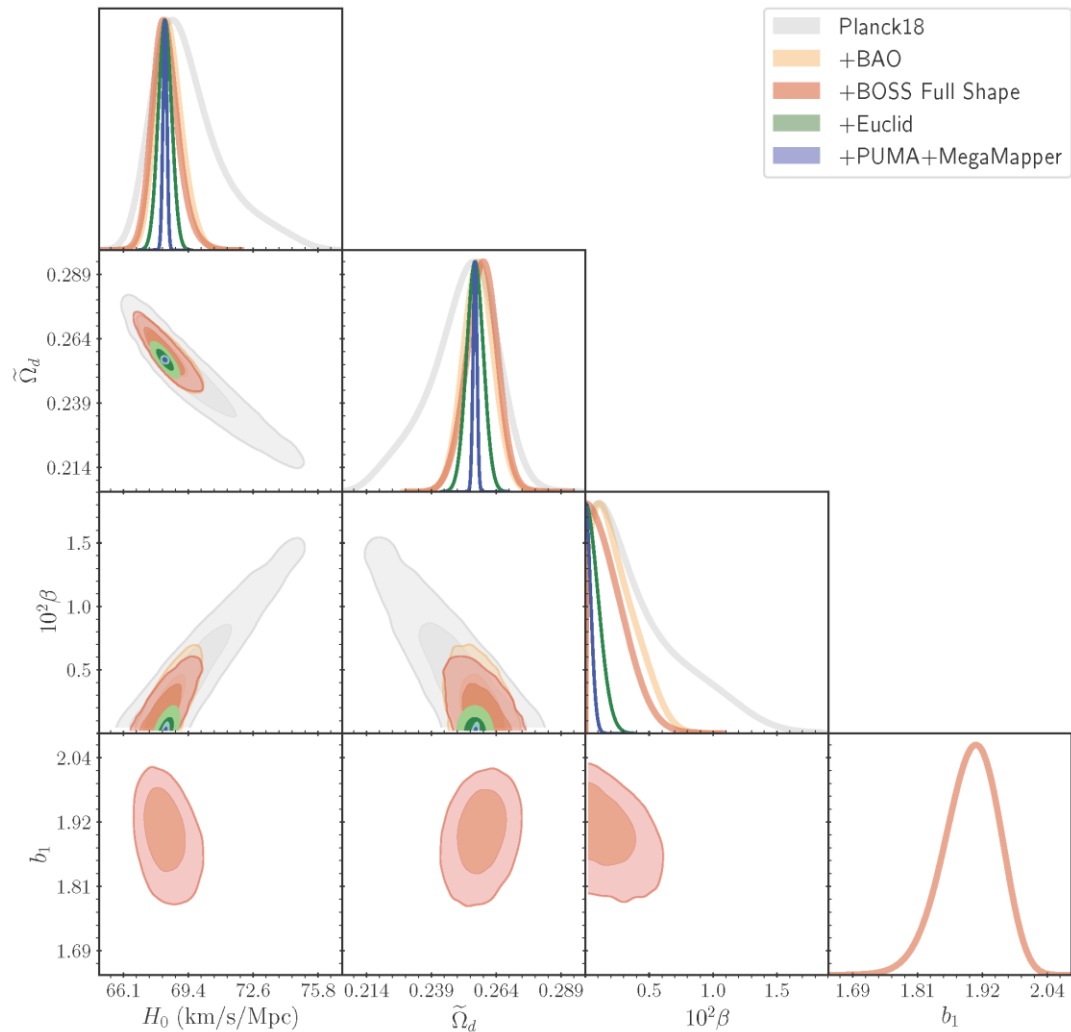
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- From $\delta_r \subset \delta_g$
- Different spatial structure
- Not log-enhanced
- Possible poles in the squeezed bispectrum

Effects on LSS



Effects on LSS



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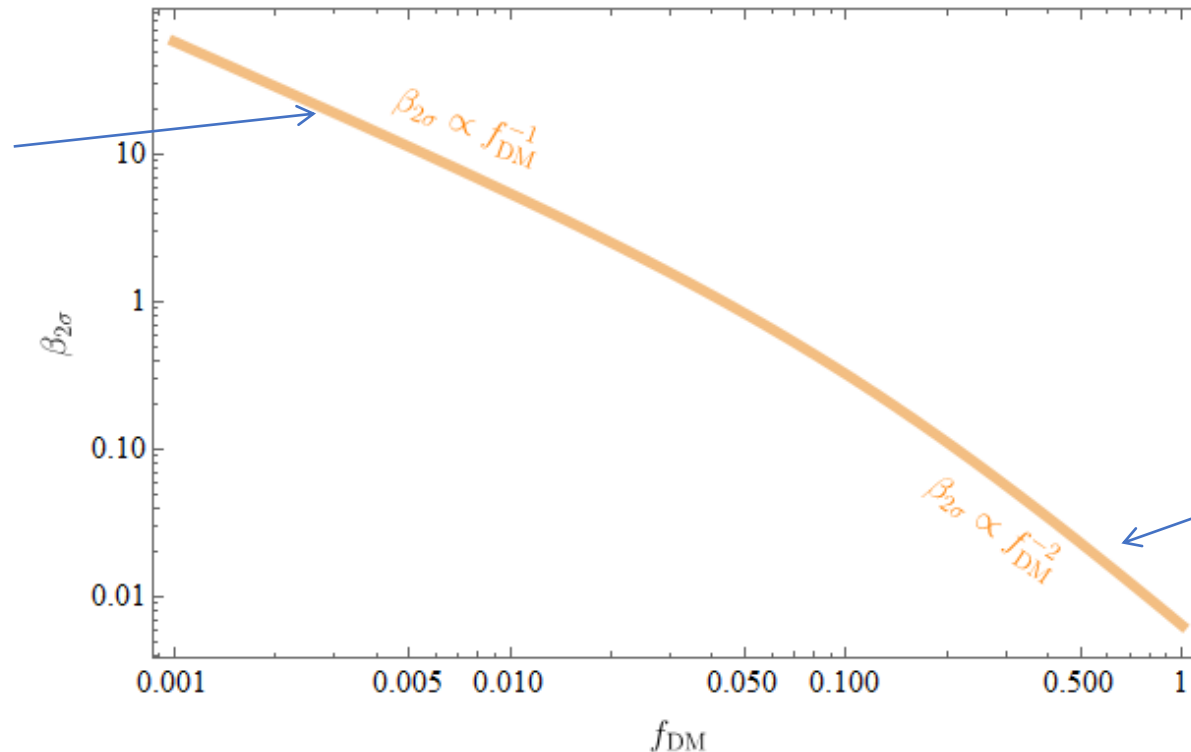
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Different scaling with the DM fraction

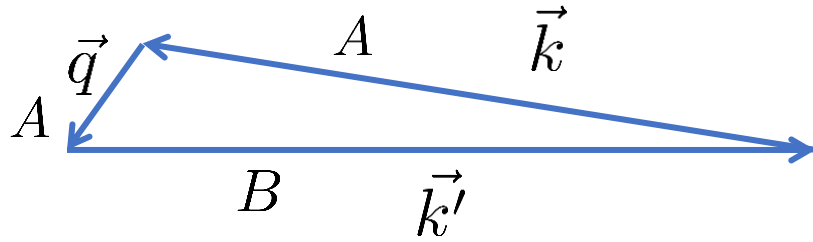
Signal dominated by new spatial features



Signal dominated by log-enhanced corrections

Effects on LSS - Pole of the Bispectrum

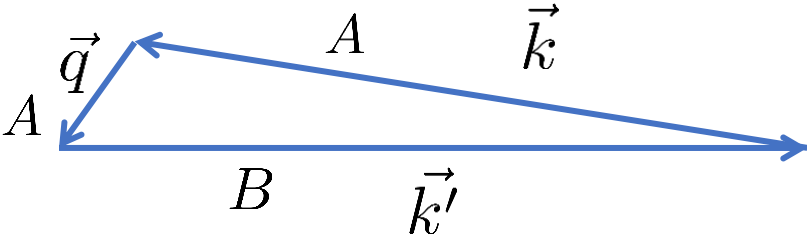
In the squeezed limit with two *different* tracers, the bispectrum has a pole



$$\lim_{q \rightarrow 0} \Delta \mathcal{B}_{AAB}(q, k, k') \propto \frac{\vec{q} \cdot \vec{k}}{q^2} P_{\text{lin}}(q) P_{\text{lin}}(k)$$

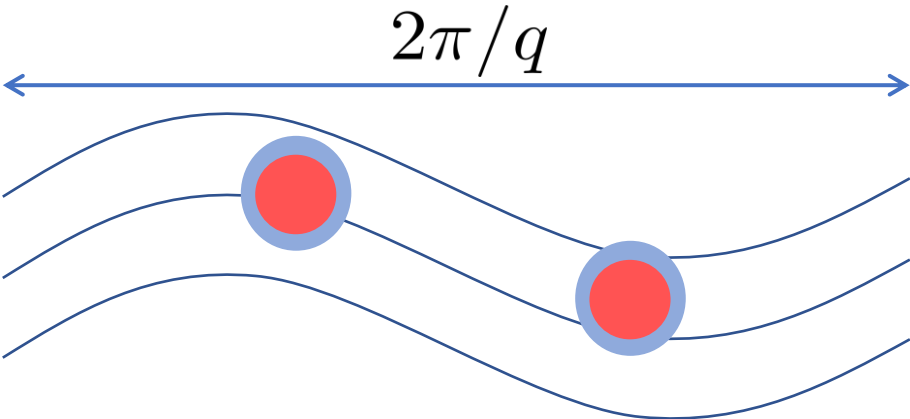
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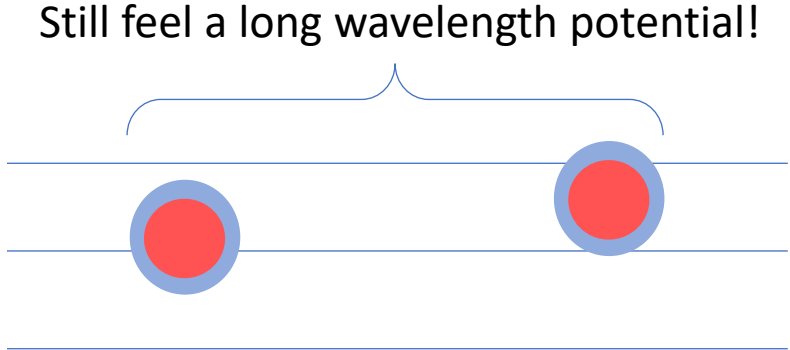


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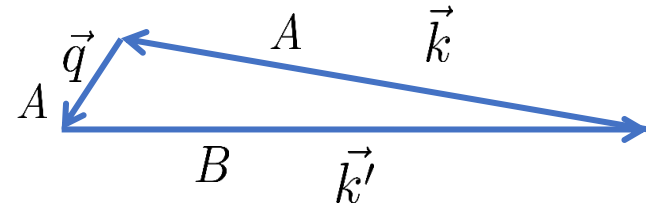
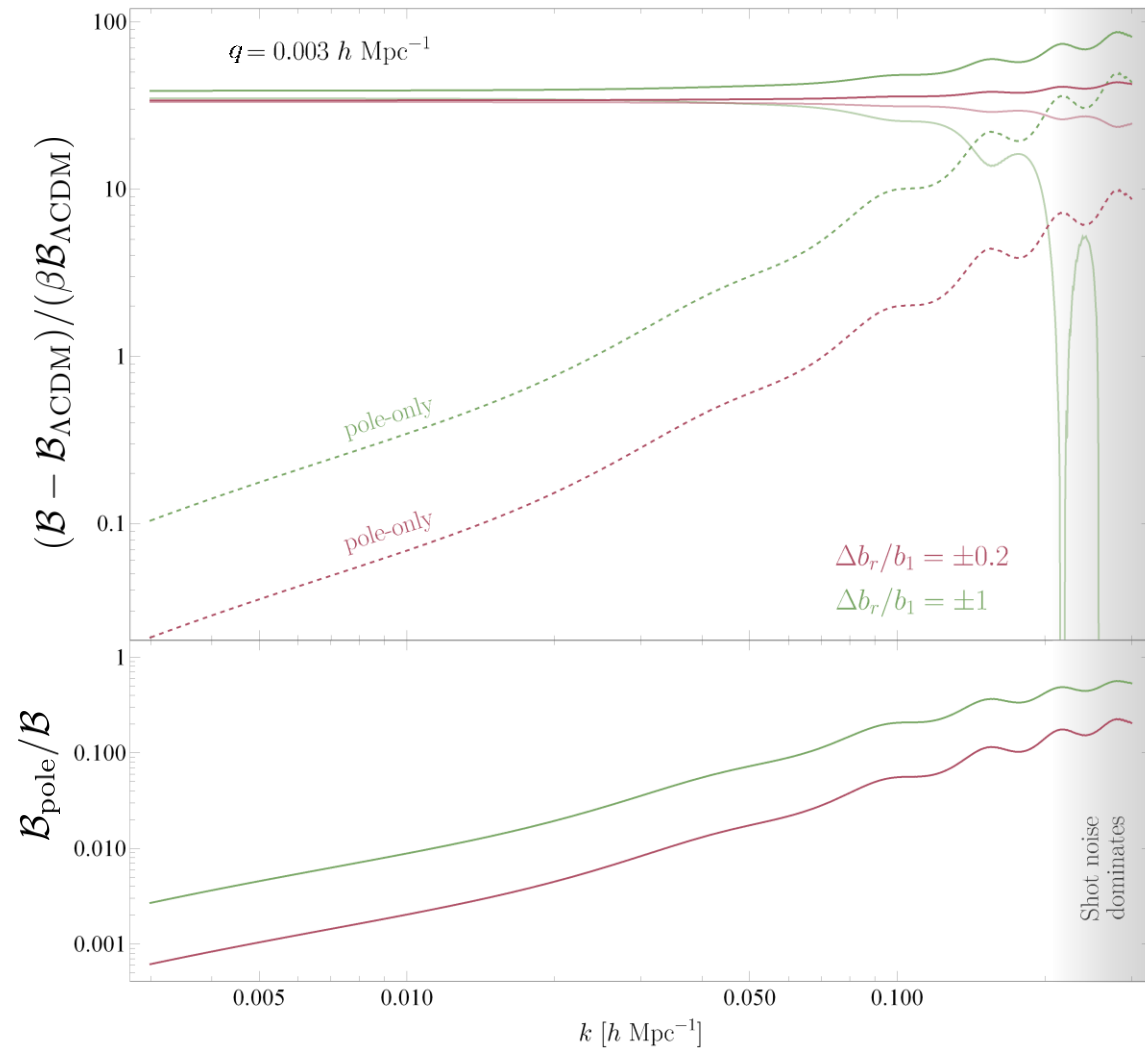
EP violation!



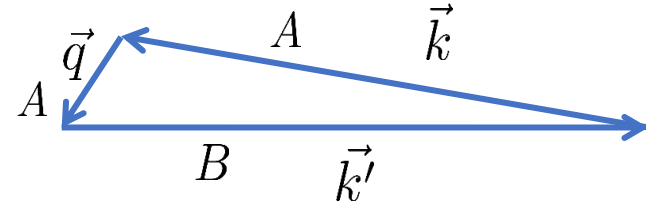
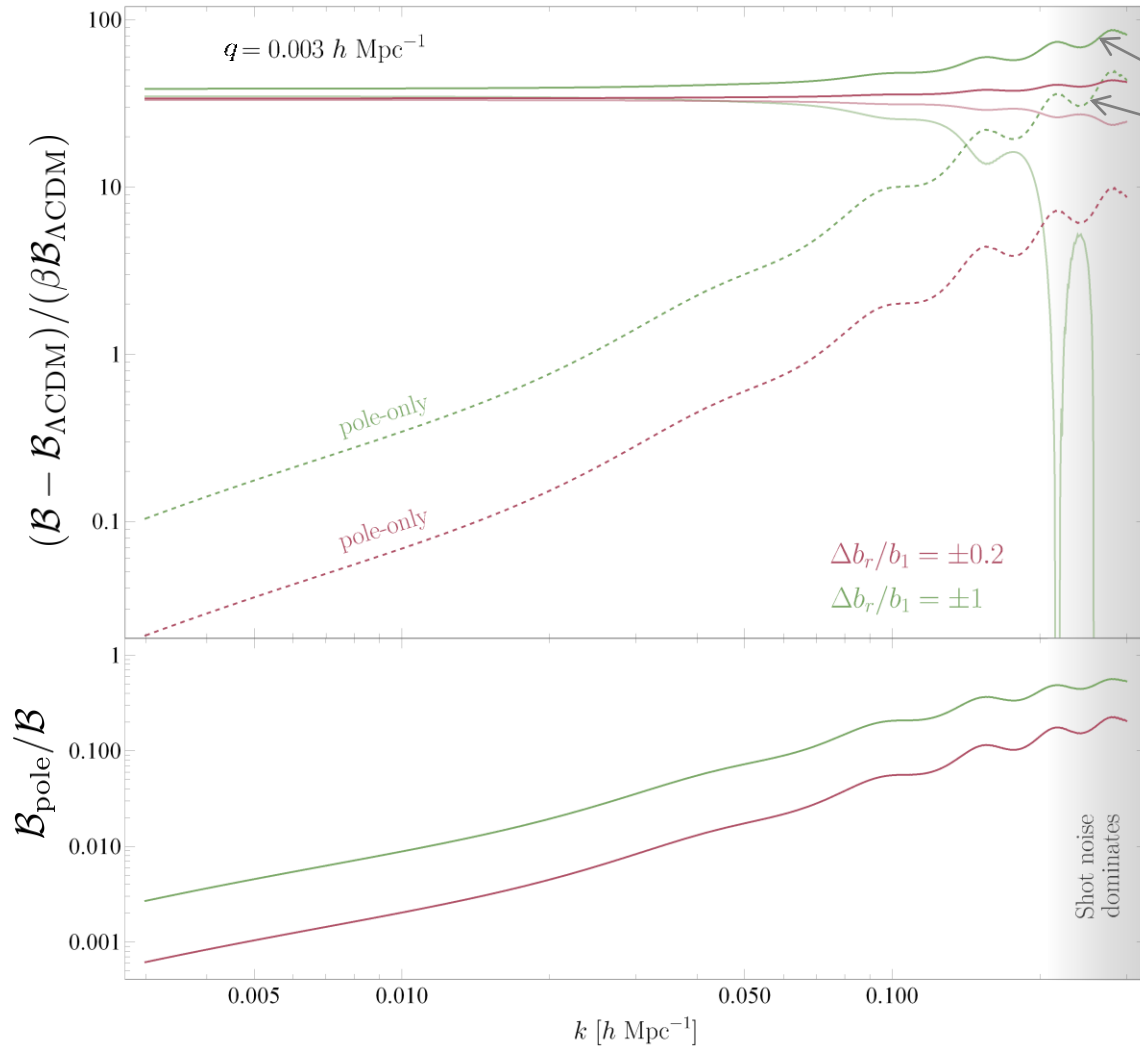
Boost to free-fall system



Effects on LSS - Pole of the Bispectrum

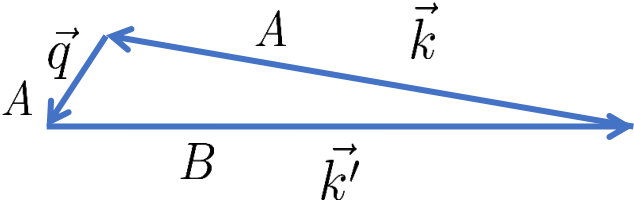
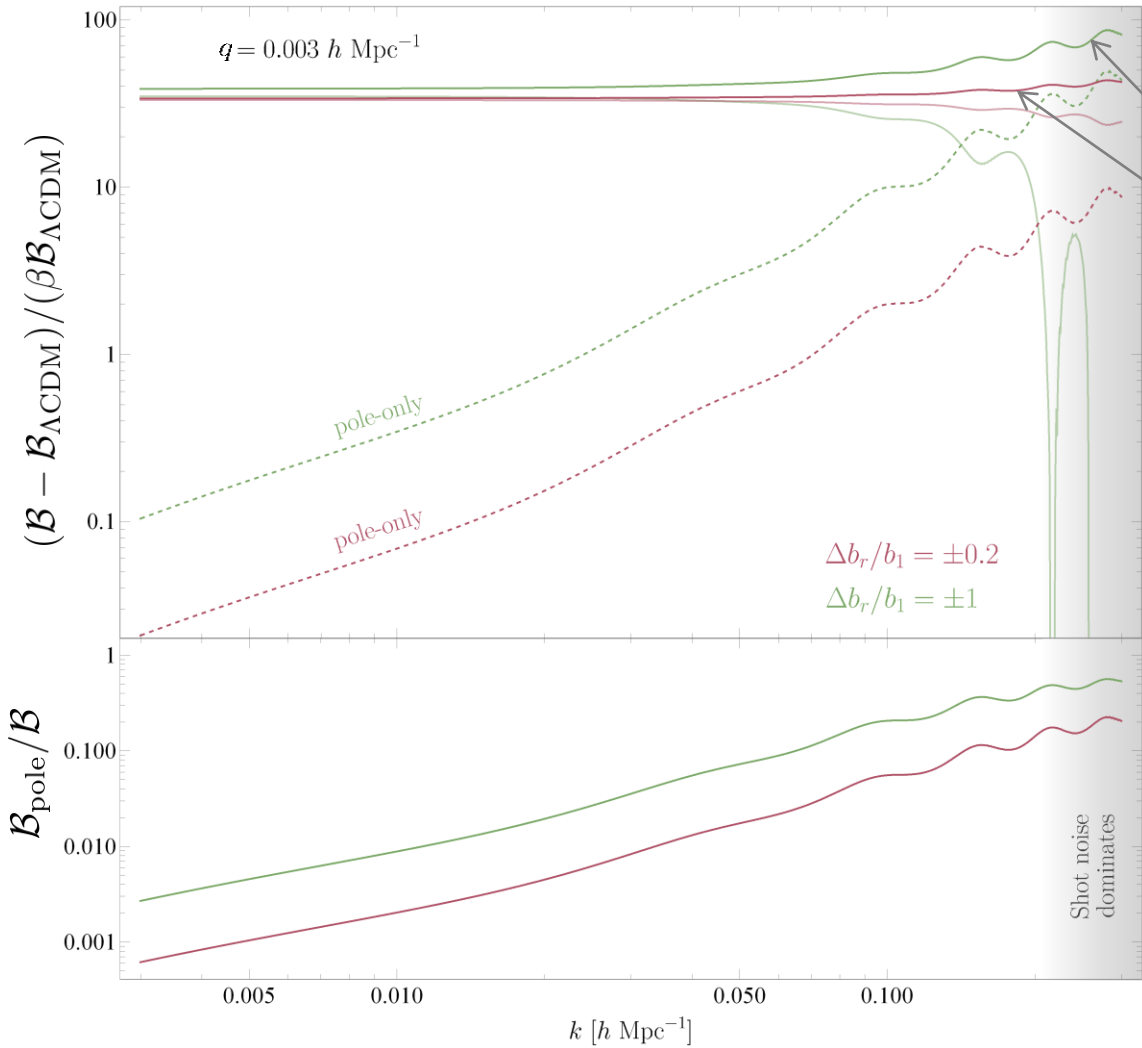


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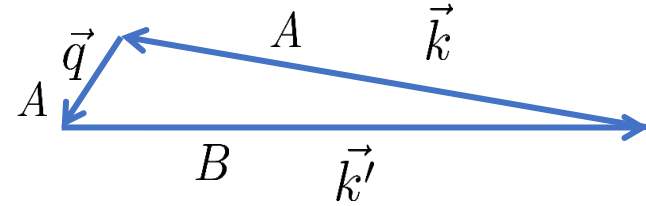
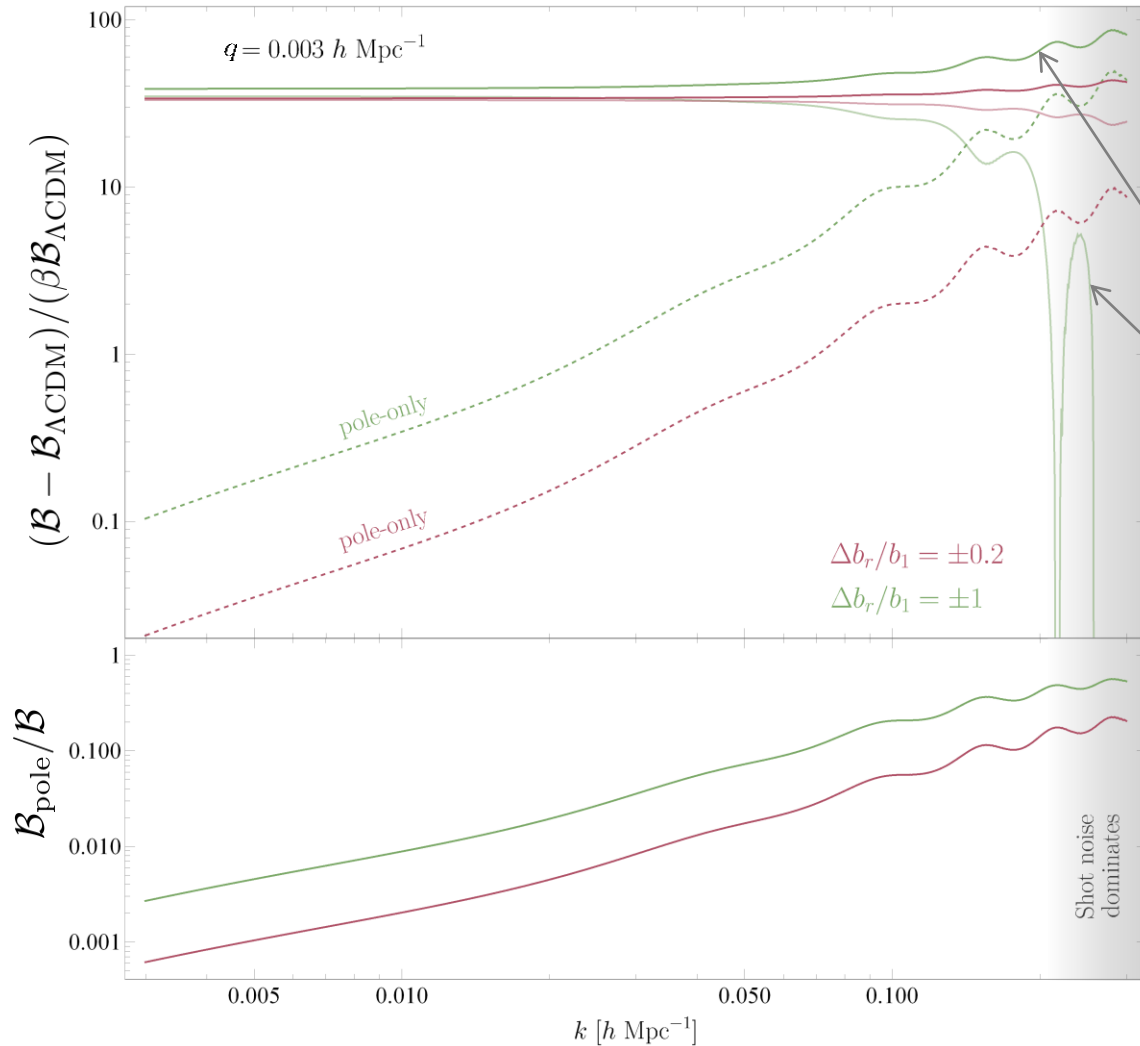
- The log enhanced growth factor “covers” the pole

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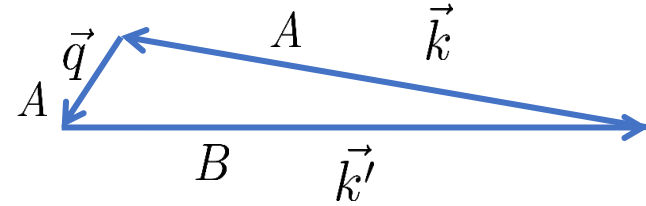
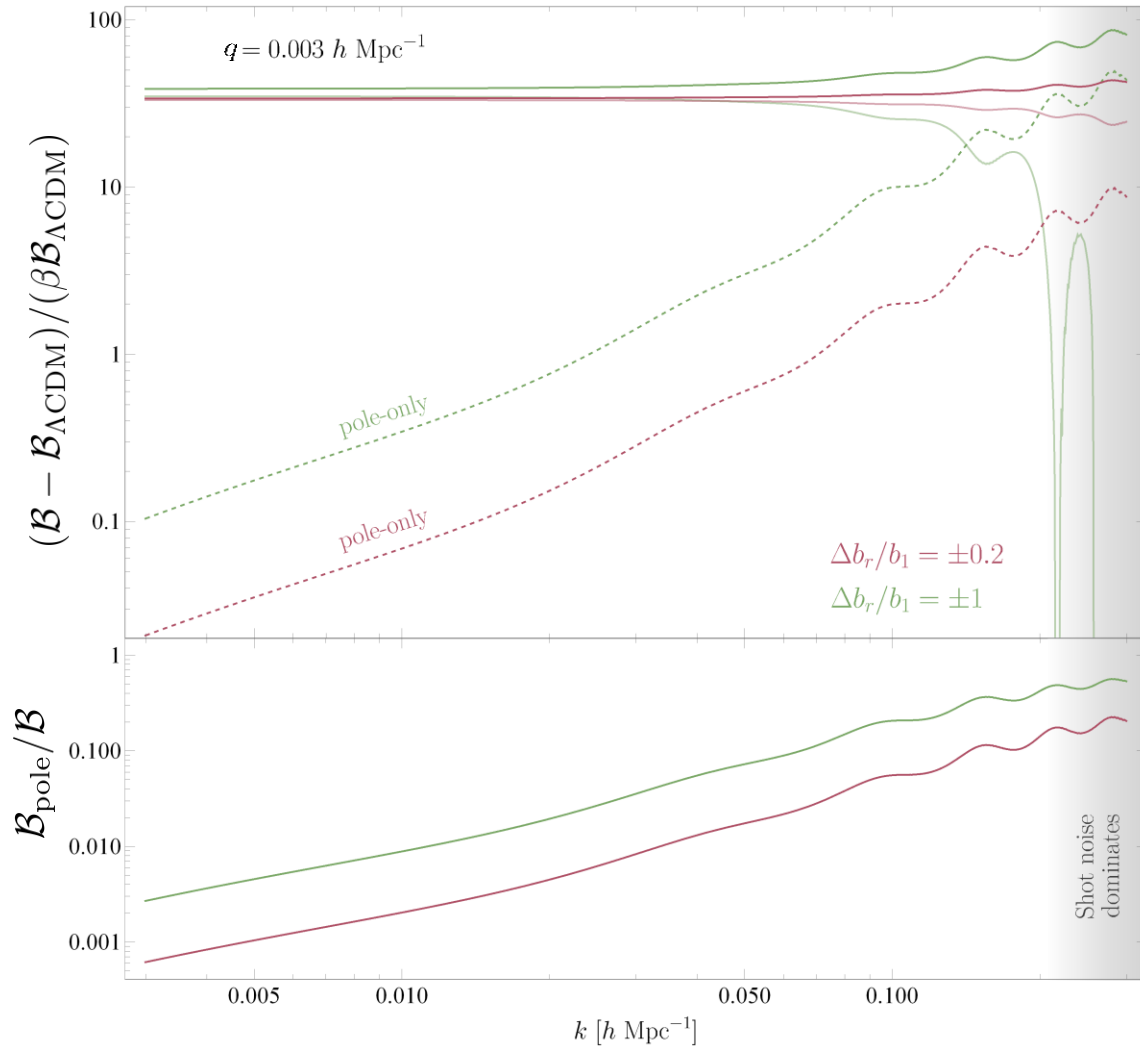
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Effects on LSS - Pole of the Bispectrum



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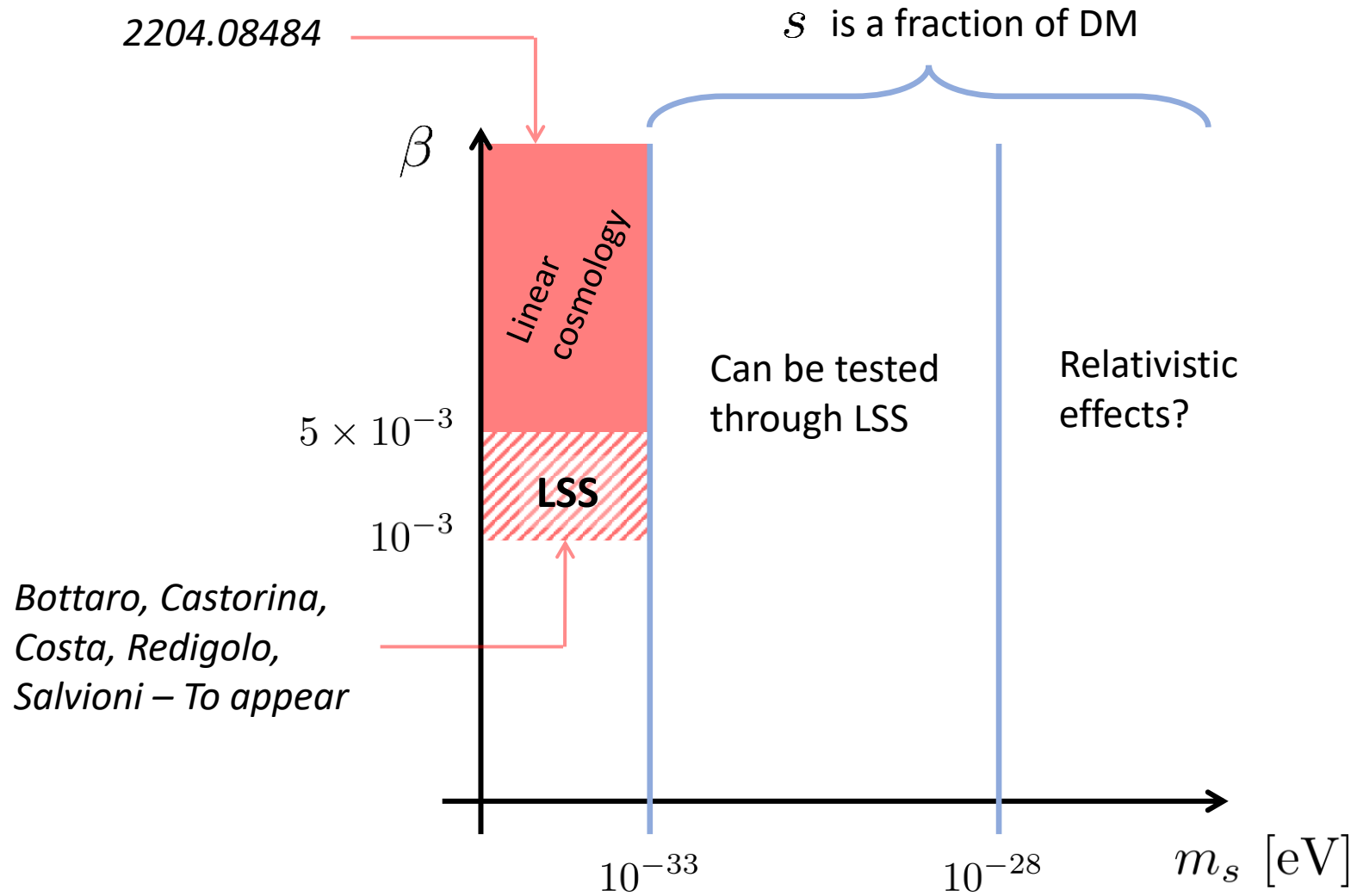
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- Pole observables

$$\mathcal{B}_{\text{pole}}(k_1, k_2, k_3) \equiv \frac{\mathcal{B}_{AAB}(k_1, k_2, k_3) - \mathcal{B}_{ABA}(k_1, k_2, k_3)}{P(k)^2}$$

Conclusions

- Fifth forces in the dark sector modify both the evolution of the background and of the perturbations
- LSS expected to improve bounds from linear theory by a factor 2-5
- Signal dominated by background effects, relative density and velocity effects dominate only for small DM fractions
- Bispectrum limited by systematics, need to improve the theoretical predictions
- Pole observables can make equivalence principle violation manifest in the squeezed limit of the bispectrum

Outlook



Back-up

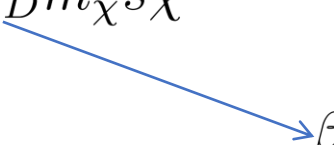
Naturalness of the model

Assuming scalar DM:

$$\mathcal{L} = \frac{1}{2} \partial_\mu \chi \partial^\mu \chi - \frac{1}{2} m_\chi^2 \chi^2 + \frac{1}{2} \partial_\mu s \partial^\mu s - V_s(s) - g_D m_\chi s \chi^2$$

Simplest case: quadratic potential

$$V_s(s) = \frac{1}{2} m_s^2 s^2$$

$$\beta = \frac{g_D^2}{4\pi G_N m_\chi^2}$$


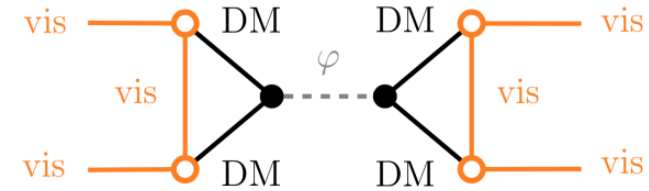
Estimate of the one-loop correction to the scalar mass gives:

$$m_s^2 \geq \frac{\beta}{(4\pi)^2} \frac{m_\chi^4}{M_P^2} \longrightarrow m_\chi \leq 0.02 \text{ eV} \left(\frac{0.01}{\beta} \right)^{\frac{1}{4}} \left(\frac{m_s}{H_0} \right)^{\frac{1}{2}}$$

Relation with other fifth force experiments

The scalar mediator can couple to the SM if DM does, e.g. the axion

$$\mathcal{L} = \frac{\alpha}{8\pi} \frac{E}{N} \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\alpha_3}{8\pi} \frac{a}{f_a} G_{\mu\nu}^a \tilde{G}^{\mu\nu a} - g_D m_a s a^2$$



$$\mathcal{L} = \sqrt{4\pi G_N s} \left(\frac{d_e}{4} F_{\mu\nu} F^{\mu\nu} + \frac{d_g b_3 \alpha_3}{8\pi} G_{\mu\nu}^a G^{\mu\nu a} + \dots \right)$$

$$d_e \simeq \sqrt{\beta} \left(\frac{m_a}{4\pi f_a} \right)^2 \frac{\alpha^2}{16\pi^2} \simeq 2 \times 10^{-10} \sqrt{\frac{\beta}{0.01}} \left(\frac{m_a}{f_a} \right)^2 \leq 2.1 \times 10^{-4}$$

$$d_g \simeq \sqrt{\beta} \left(\frac{m_a}{4\pi f_a} \right)^2 \frac{\alpha_3}{8\pi b_3} \simeq 3 \times 10^{-6} \sqrt{\frac{\beta}{0.01}} \left(\frac{m_a}{f_a} \right)^2 \leq 2.9 \times 10^{-6}$$

MICROSCOPE
(1712.01176)

CMB lensing

