

Crash course on CATS

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23.01.2023



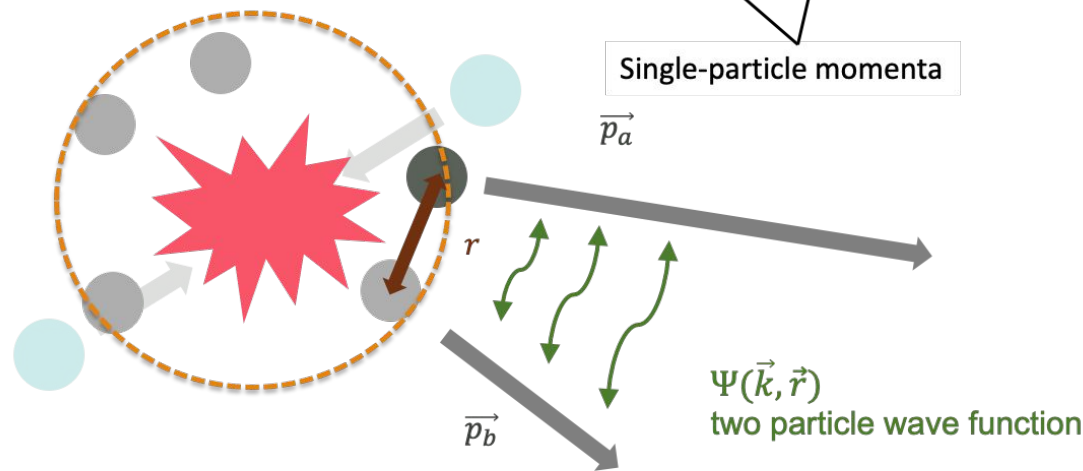
1. Femtoscopy in a nutshell
2. Basics of CATS
 - a. main features
 - b. a deeper understanding with few examples

Useful links:

- [CATS paper](#)
- [Dr. D. Mihaylov PhD thesis](#)
- [GitHub CATS repo](#)
- [\(Old\) GitHub tutorial](#)



Source function $S(\vec{r})$



Statistical definition

Experimental definition

Theoretical definition

$$C(k^*) = \frac{\mathcal{P}(\vec{p}_a, \vec{p}_b)}{\mathcal{P}(\vec{p}_a)\mathcal{P}(\vec{p}_b)} = \mathcal{N} \frac{N_{\text{Same}}(k^*)}{N_{\text{Mixed}}(k^*)} = \int S(r) |\Psi(\vec{k}^*, \vec{r})|^2 d^3r \xrightarrow{k^* \rightarrow \infty} 1$$

Emitting Source

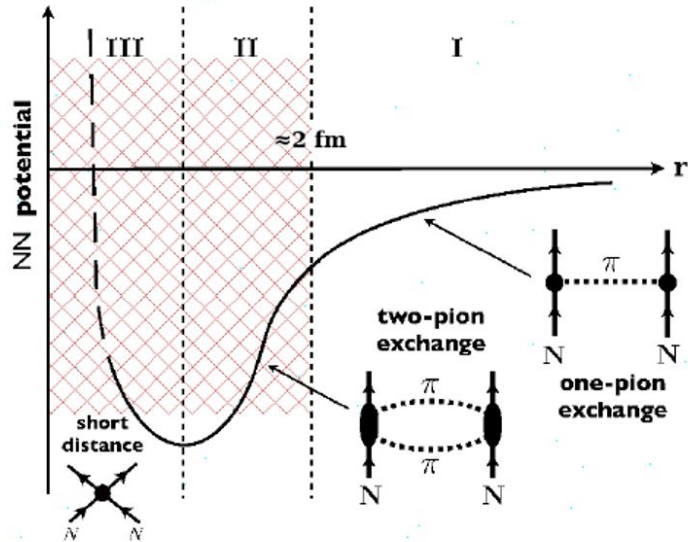
Relative Wavefunction

$C(k^*) \rightarrow$

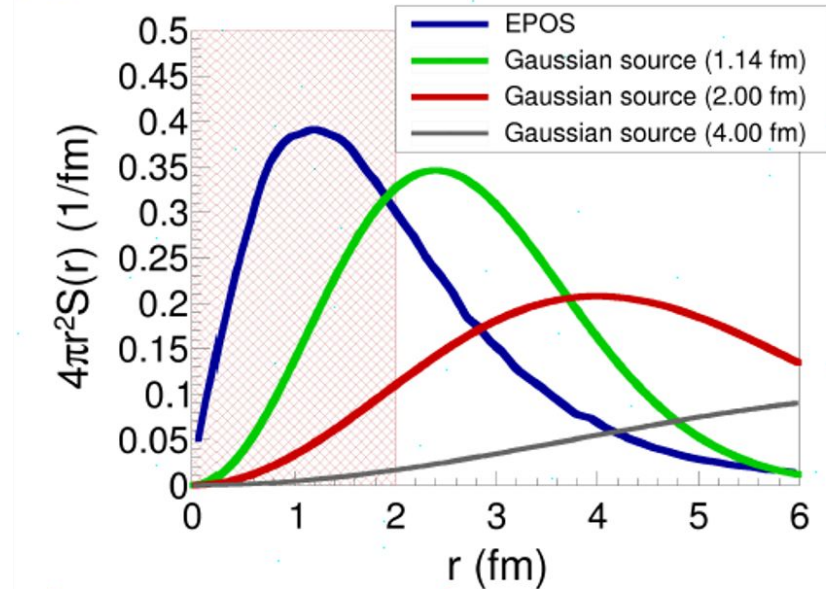
- > 1 attraction
- $= 1$ no interaction
- < 1 repulsion



$$C(k^*) = \int S(r) |\Psi(\vec{k}^*, \vec{r})|^2 d^3r \xrightarrow{k^* \rightarrow \infty} 1$$



Based on Holt, Kaiser and Weise,
Prog.Part.Nucl.Phys. 73 (2013) 35-83



A-A collisions: $R_G \sim 4$ fm

p-A collisions: $R_G \sim 2$ fm

p-p collisions: $R_G \sim 1$ fm



$$C(k) = \frac{\mathcal{P}(\vec{p}_a, \vec{p}_b)}{\mathcal{P}(\vec{p}_a)\mathcal{P}(\vec{p}_b)} = \int \underbrace{S(\vec{r})}_{\text{source}} \underbrace{|\Psi(\vec{r}, k)|^2}_{\text{2-particle wave function}} d\vec{r} \xrightarrow{k \rightarrow \infty} 1$$

Final output

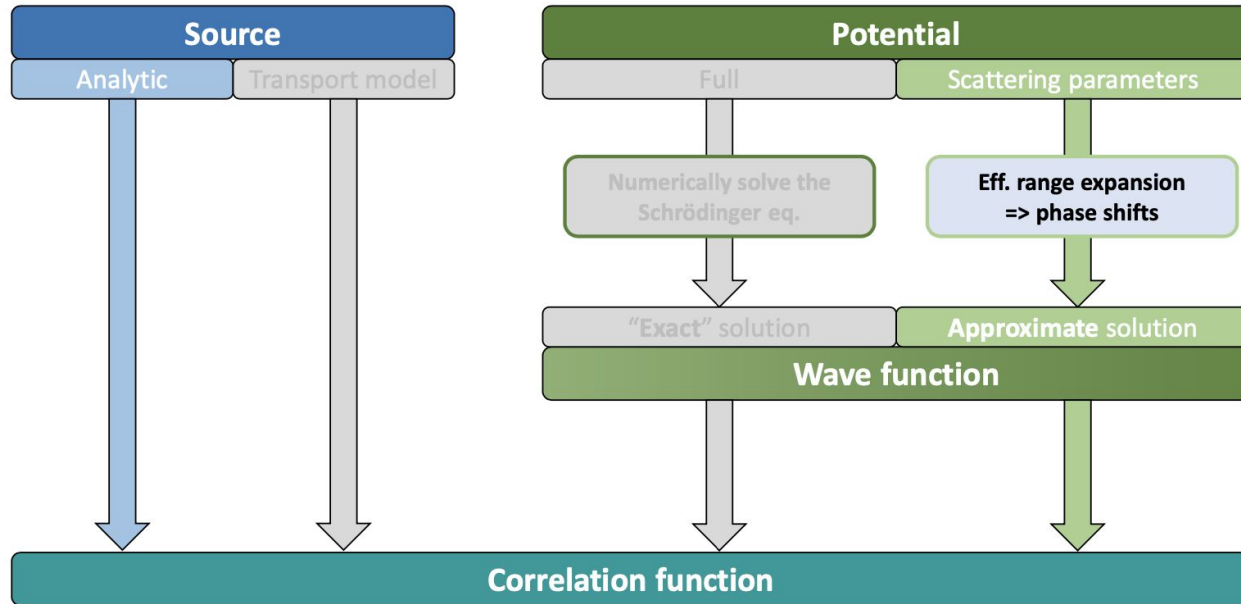
input

eval. for any local potential $V(r)$

1. **Study** the **interaction potential** between different particle species
→ one needs to be capable of computing the wave function
2. To get reliable results we must find ways to realistically **model** the **emission source**
→ going beyond the gaussian assumption if required, input from transport models



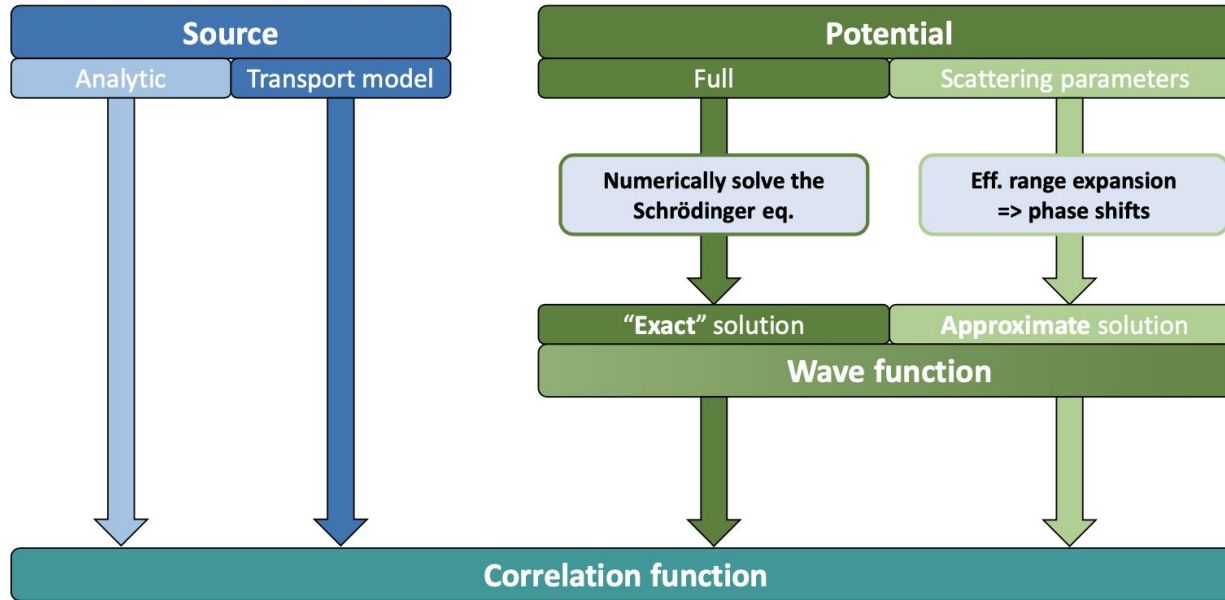
$$C(k) = \frac{\mathcal{P}(\vec{p}_a, \vec{p}_b)}{\mathcal{P}(\vec{p}_a)\mathcal{P}(\vec{p}_b)} = \int \underbrace{S(\vec{r})}_{\text{source}} \underbrace{|\Psi(\vec{r}, k)|^2}_{\text{2-particle wave function}} d\vec{r} \xrightarrow{k \rightarrow \infty} 1 \quad (1)$$



- asymptotic behaviour (far from scatt. centre)
- ERE expansion
- not suited (always) for small sources



$$C(k) = \frac{\mathcal{P}(\vec{p}_a, \vec{p}_b)}{\mathcal{P}(\vec{p}_a)\mathcal{P}(\vec{p}_b)} = \int \underbrace{S(\vec{r})}_{\text{source}} \underbrace{|\Psi(\vec{r}, k)|^2}_{\text{2-particle wave function}} d\vec{r} \xrightarrow{k \rightarrow \infty} 1 \quad (1)$$



Numerically solve SE



$$\Psi_k(\vec{r}) = R_k(k, r)Y(\theta) = \sum_{l=0}^{\infty} R_{k,l}(r)Y_l(\theta) = \sum_{l=0}^{\infty} i^l(2l+1) \frac{\mathbf{u}_{k,l}(r)}{r} P_l(\cos \theta)$$

↑
Separation of variables
↑
Expansion in partial waves
↑
Legendre polynomials

With $\mathbf{u}_{k,l}(r) = r \cdot R_{k,l}(r)$

$$\frac{d^2 \mathbf{u}_{k,l}}{dr^2} = \left[2mV(r) + \frac{l(l+1)}{r^2} - k^2 \right] \mathbf{u}_{k,l} \quad \text{radial SE}$$



- we need to compute the RSE for all partial waves (PW)
- low-energy (low k^* pairs)
 - interaction dominated by s-wave (residual effects of p-, d- waves)
 - rest of PW follow free particle solution

$$\Psi_k(\vec{r}) = e^{ikz} = \sum_{l=0}^{\infty} i^l (2l+1) j_l(kr) P_l(\cos\theta), \quad \text{spherical Bessel functions}$$

- in PW l there is no interaction → free wave (see next discussion with channels)
- symmetrization/anti-symmetrization wf included in CATS



- Interaction depends on spin/isospin configurations
 - NN interaction ($S=0,1$)
 - $l=1$: pp, nn, $1/\sqrt{2}$ (pn+np)
 - $l=0$: $1/\sqrt{2}$ (pn-np) $\rightarrow$ for $S=1$ we have the deuteron
 - Each configuration contributes to $C(k^*)$ with weights determined by degeneracy.
- $\rightarrow$ p- Λ has $S=0,1$ configurations with weights $1/4$ and $3/4$.

Example for p-p
The shaded PWs are
Pauli-blocked for
identical particles
4 int. channels

Channel 0	Channel 1	Channel 2	Channel 3
$w_0 = \frac{1}{4}$	$w_1 = \frac{3}{4} \cdot \frac{1}{9} = \frac{1}{12}$	$w_2 = \frac{3}{4} \cdot \frac{3}{9} = \frac{1}{4}$	$w_3 = \frac{3}{4} \cdot \frac{5}{9} = \frac{5}{12}$
1P_1	3P_0	3P_1	3P_2
1S_0	3S_1	3S_1	3S_1
Total spin 0	Total spin 1		
Weight 1/4	Weight 3/4		

$$|\Psi_k(r)|^2 = \sum_{\text{channel}} w_{\text{channel}} \sum_l |\Psi_{k,l}(r)|^2.$$



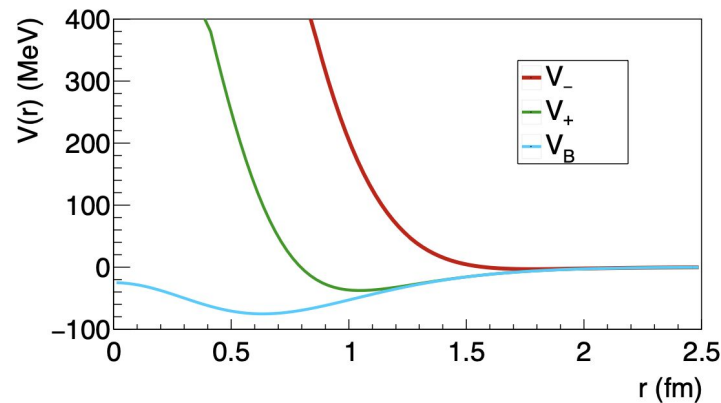
- Toy double-gaussian potential (single channel)

$$V(r) = V_1 \exp(-r^2/\mu_1^2) + V_2 \exp(-r^2/\mu_2^2).$$

- CF as average value of the $|\Psi|^2$ probed by the source acting as pdf

$$C_{\text{th}}(k) = \langle |\Psi|^2 \rangle = \int S(r) |\Psi(\vec{k}, \vec{r})|^2 d^3r.$$

- any depletion/enhancement in wf is transferred to CF, source determines which region in r to probe



1. shifted asymptotic solution determined by δ_0
→ eff. described by scatt. pars

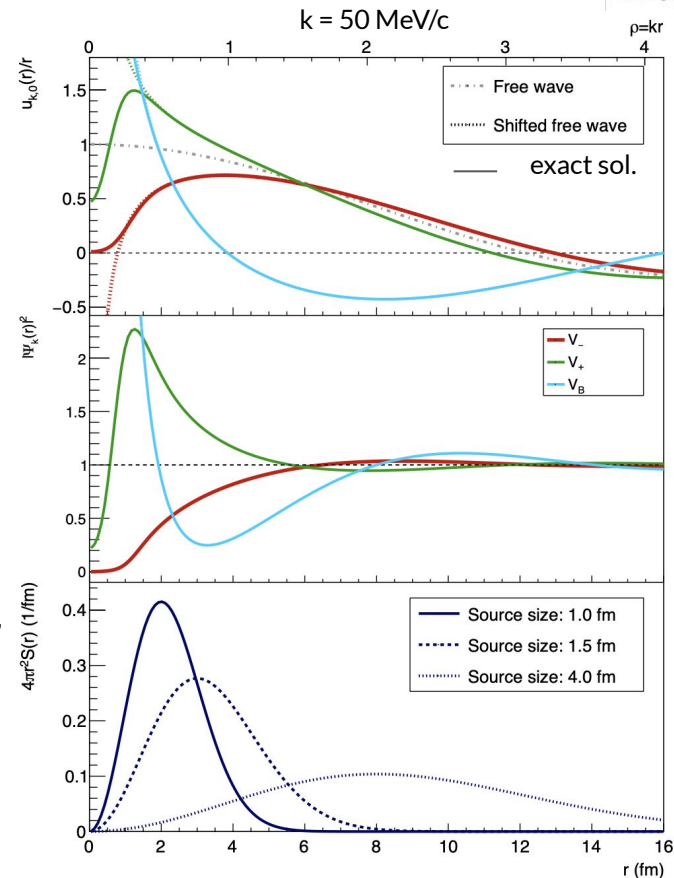
$$k \cot(\delta(k)) \stackrel{k \rightarrow 0}{\approx} \frac{1}{f_0} + \frac{1}{2} d_0 k^2 + \mathcal{O}(k^4).$$

- overlaps with true solution for $r > 1-2$ fm
- basics of Lednický model relating CF to f_0, d_0
2. $|\Psi|^2$ behaviour tells us about the interaction
- pull/push wf of attr./rep
- interplay for bound-state + localized at low r

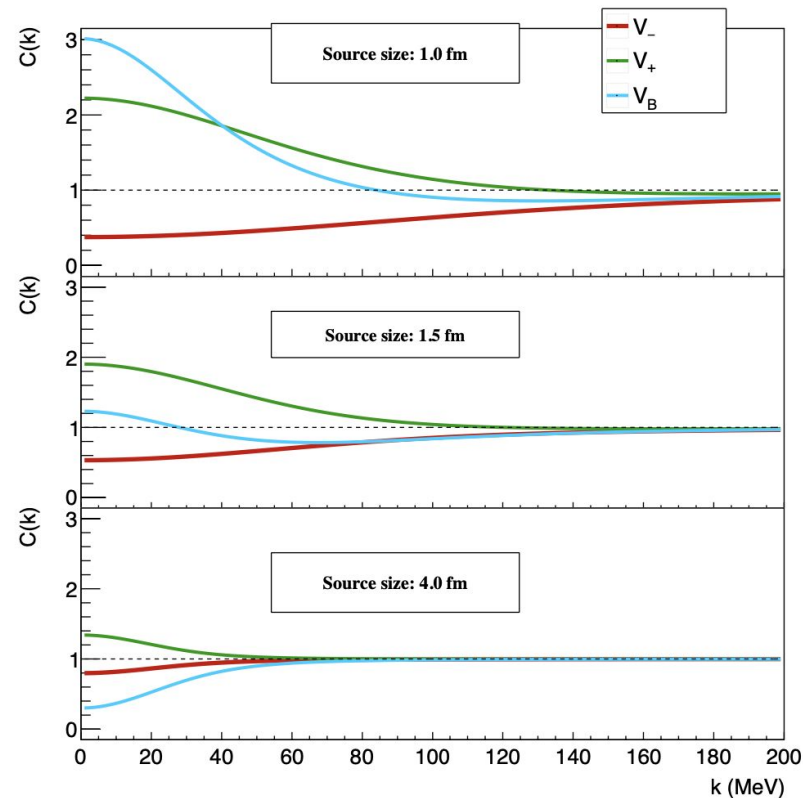
for BS

Potential	f_0 (fm)	d_0 (fm)
V_-	-0.8	0.6
V_+	1.1	5.0
V_B	-4.7	1.3

$$|f_0| \geq 2d_0.$$



1. signal strength at low k^*
 - a. decreases as source size increases
2. For presence of BS
 - a. flip below unity when measured in large colliding systems
 - b. strength of depletion depends on binding energies (e.g. $p\Omega$, DD^*)



Additional slides



- Extract for any potential phase shifts δ_l

$$\sigma = \frac{4\pi}{k^2} \sum_l (2l+1) \sin^2(\delta_l).$$

- p-p as benchmark, Argonne V18 pot.
 - perfect agreement with data
- p- Λ cross section data
 - comparison Usmani/NLO
 - discrepancies between models only at low k^*
→ femtoscopy!!

Parameter	Usmani	χ EFT NLO
$a_0^{(s)}$ (fm)	-2.88	-2.91
$r_{eff}^{(s)}$ (fm)	2.92	2.78
$a_0^{(t)}$ (fm)	-1.66	-1.54
$r_{eff}^{(t)}$ (fm)	3.78	2.72

