

Lorentz Violation and Sagnac Gyroscopes

Jay D. Tasson

Carleton College

jtasson@carleton.edu

June 14, 2023

Outline

- introduction
 - Lorentz violation & the Standard-Model Extension (SME)
 - gravity
- tests: recent and proposed tests
 - gravimeters
 - **Sagnac gyroscopes**
 - gravitational waves

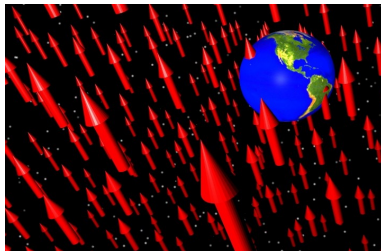
1

Lorentz violation

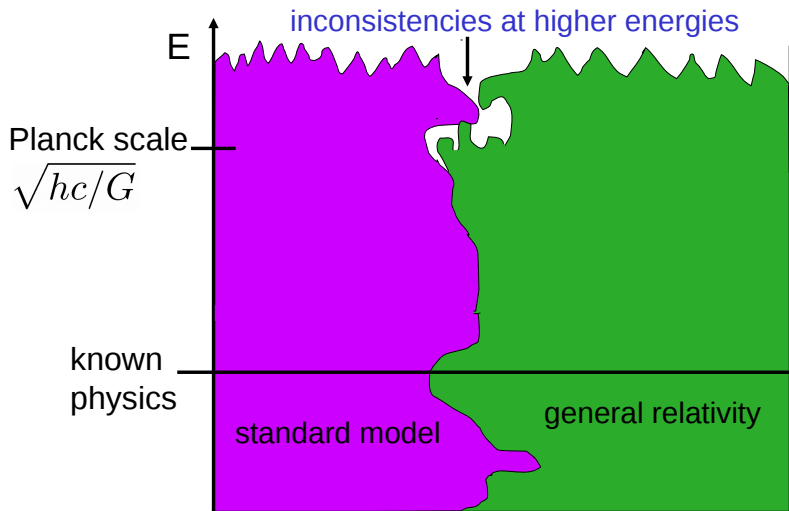
- Lorentz symmetry



- Lorentz violation



motivation



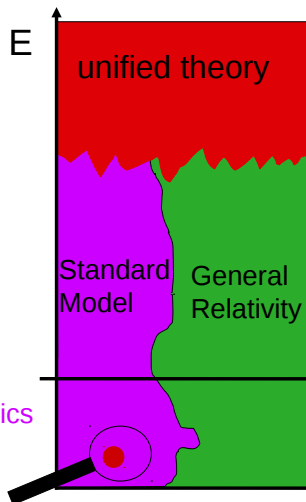
motivation

options for probing experimentally

- galaxy-sized accelerator



- suppressed effects in sensitive experiments
- Lorentz and CPT violation
- can arise in theories of new physics
 - difficult to mimic with conventional effects



the Standard-Model Extension (SME)

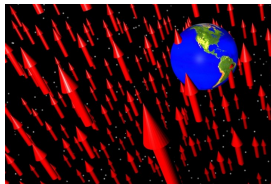
effective field theory which contains:

- General Relativity (GR)
 - Standard Model (SM)
 - arbitrary coordinate-independent CPT & Lorentz violation
- $$L_{\text{SME}} = L_{\text{GR}} + L_{\text{SM}} + L_{\text{LV}}$$
- CPT violation comes with Lorentz violation

CPT & Lorentz-violating terms

- constructed from GR and SM fields
- parameterized
by coefficients for Lorentz violation
- samples

$$\bar{\psi} a_{\mu} \gamma^{\mu} \psi$$



the Standard-Model Extension (SME)

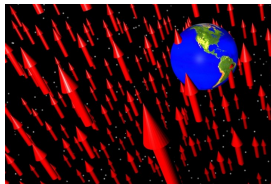
effective field theory which contains:

- General Relativity (GR)
 - Standard Model (SM)
 - arbitrary coordinate-independent CPT & Lorentz violation
- $$L_{\text{SME}} = L_{\text{GR}} + L_{\text{SM}} + L_{\text{SME}}$$
- CPT violation comes with Lorentz violation

CPT & Lorentz-violating terms

- constructed from GR and SM fields
- parameterized by coefficients for Lorentz violation
- samples

$$\bar{\psi} a_{\mu} \gamma^{\mu} \psi$$



A broad and general search for new physics

Recent Gravitational Tests

System	E.g.	
MICROSCOPE	Pihan-le Bars <i>et al.</i>	PRL 123, 231102 (2019)
gravimeters	Flowers <i>et al.</i>	PRL 119, 201101 (2017)
solar-system data	Hees <i>et al.</i>	PRD 92, 064049 (2015)
lunar laser ranging	Bourgoin <i>et al.</i>	PRL 117, 241301 (2016)
gravitational waves	LIGO, Virgo, FermiGBM, Integral	ApJL 848, L13 (2017)
WEP experiments	Hohensee <i>et al.</i>	PRL 111, 151102 (2013)
pulsar timing	L. Shao	PRL 112, 111103 (2014)
short-range gravity	C.-G. Shao <i>et al.</i>	PRL 122, 011102 (2019)
VLBI	Le Poncin-Lafitte <i>et al.</i>	PRD 94, 125030 (2016)

...

For a full list, see,
Kostelecký & Russell, *Data tables for Lorentz and CPT violation*,
arXiv:0801.0287

...updated annually

Complete linearized pure gravity

known physics
SM + GR

+ ○ + ● + ◉ + ... =

quantum
gravity

$$\mathcal{L} = \frac{1}{4} h_{\mu\nu} \hat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} h_{\rho\sigma}$$

$$\hat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} = \mathcal{K}^{(d)\mu\nu\rho\sigma\epsilon_1\epsilon_2\dots\epsilon_{d-2}} \partial_{\epsilon_1} \partial_{\epsilon_2} \dots \partial_{\epsilon_{d-2}}$$

Complete linearized pure gravity

known physics
SM + GR

+ ○ + ● + ◉ + ... =

quantum
gravity

$$\mathcal{L} = \frac{1}{4} h_{\mu\nu} \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} h_{\rho\sigma} \quad \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} = \mathcal{K}^{(d)\mu\nu\rho\sigma\epsilon_1\epsilon_2\dots\epsilon_{d-2}} \partial_{\epsilon_1} \partial_{\epsilon_2} \dots \partial_{\epsilon_{d-2}}$$

0) gauge-structure preserving operators

$$\widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} \xrightarrow{\text{tensor decomposition}} \widehat{\mathcal{S}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{Q}}^{(d)\mu\nu\rho\sigma} + \dots$$

Complete linearized pure gravity

known physics
SM + GR

+ ○ + ● + ◉ + ... =

quantum
gravity

$$\mathcal{L} = \frac{1}{4} h_{\mu\nu} \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} h_{\rho\sigma} \quad \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} = \mathcal{K}^{(d)\mu\nu\rho\sigma\epsilon_1\epsilon_2\dots\epsilon_{d-2}} \partial_{\epsilon_1} \partial_{\epsilon_2} \dots \partial_{\epsilon_{d-2}}$$

0) gauge-structure preserving operators

$$\widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} \xrightarrow{\text{tensor decomposition}} \widehat{\mathcal{S}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{Q}}^{(d)\mu\nu\rho\sigma} + \dots$$

1) linearized limit of Einstein-Hilbert

$$\widehat{\mathcal{S}}^{(4)\mu\nu\rho\sigma} \rightarrow \epsilon^{\mu\rho\alpha\kappa} \epsilon^{\nu\sigma\beta\lambda} \eta_{\kappa\lambda} \partial_\alpha \partial_\beta$$

Complete linearized pure gravity

known physics
SM + GR

+ ○ + ● + ◉ + ... =

quantum
gravity

$$\mathcal{L} = \frac{1}{4} h_{\mu\nu} \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} h_{\rho\sigma} \quad \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} = \mathcal{K}^{(d)\mu\nu\rho\sigma\epsilon_1\epsilon_2\dots\epsilon_{d-2}} \partial_{\epsilon_1} \partial_{\epsilon_2} \dots \partial_{\epsilon_{d-2}}$$

0) gauge-structure preserving operators

$$\widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} \xrightarrow{\text{tensor decomposition}} \widehat{\mathcal{S}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{Q}}^{(d)\mu\nu\rho\sigma} + \dots$$

1) linearized limit of Einstein-Hilbert

$$\widehat{\mathcal{S}}^{(4)\mu\nu\rho\sigma} \rightarrow \epsilon^{\mu\rho\alpha\kappa} \epsilon^{\nu\sigma\beta\lambda} \eta_{\kappa\lambda} \partial_\alpha \partial_\beta$$

2) $d = 4$ (minimal) Lorentz violation

$$\widehat{\mathcal{S}}^{(4)\mu\nu\rho\sigma} \rightarrow \mathcal{S}^{(4)\mu\nu\rho\sigma\alpha\beta} \partial_\alpha \partial_\beta = \epsilon^{\mu\rho\alpha\kappa} \epsilon^{\nu\sigma\beta\lambda} \bar{\mathcal{S}}_{\kappa\lambda} \partial_\alpha \partial_\beta$$

Complete linearized pure gravity

known physics
SM + GR

+ ○ + ● + ◉ + ... =

quantum
gravity

$$\mathcal{L} = \frac{1}{4} h_{\mu\nu} \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} h_{\rho\sigma} \quad \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} = \mathcal{K}^{(d)\mu\nu\rho\sigma\epsilon_1\epsilon_2\dots\epsilon_{d-2}} \partial_{\epsilon_1} \partial_{\epsilon_2} \dots \partial_{\epsilon_{d-2}}$$

0) gauge-structure preserving operators

$$\widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} \xrightarrow{\text{tensor decomposition}} \widehat{\mathcal{S}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{K}}^{(d)\mu\nu\rho\sigma} + \widehat{\mathcal{Q}}^{(d)\mu\nu\rho\sigma} + \dots$$

1) linearized limit of Einstein-Hilbert

$$\widehat{\mathcal{S}}^{(4)\mu\nu\rho\sigma} \rightarrow \epsilon^{\mu\rho\alpha\kappa} \epsilon^{\nu\sigma\beta\lambda} \eta_{\kappa\lambda} \partial_\alpha \partial_\beta$$

2) $d = 4$ (minimal) Lorentz violation

$$\widehat{\mathcal{S}}^{(4)\mu\nu\rho\sigma} \rightarrow \mathcal{S}^{(4)\mu\nu\rho\sigma\alpha\beta} \partial_\alpha \partial_\beta = \epsilon^{\mu\rho\alpha\kappa} \epsilon^{\nu\sigma\beta\lambda} \overline{\mathcal{S}}_{\kappa\lambda} \partial_\alpha \partial_\beta$$

3) $d = 5$ example

$$\widehat{\mathcal{Q}}^{(5)\mu\nu\rho\sigma} \rightarrow \mathcal{Q}^{(5)\mu\rho\alpha\nu\beta\sigma\gamma} \partial_\alpha \partial_\beta \partial_\gamma$$

action to metric via spontaneous breaking

$$\begin{aligned}g_{00} &= -1 + 2U + 3\bar{s}^{00}U + \bar{s}^{jk}U^{jk} \\g_{0j} &= -\bar{s}^{0j}U - \bar{s}^{0k}U^{jk} + \frac{1}{2}\hat{Q}^j\chi \\g_{jk} &= \delta^{jk} + (2 - \bar{s}^{00})\delta^{jk}U \\&\quad + (\bar{s}^{lm}\delta^{jk} - \bar{s}^{jl}\delta^{mk} - \bar{s}^{kl}\delta^{jm} + 2\bar{s}^{00}\delta^{jl}\delta^{km})U^{lm}\end{aligned}$$

$$U = G \int d^3x' \frac{\rho(\vec{x}', t)}{R}$$

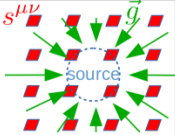
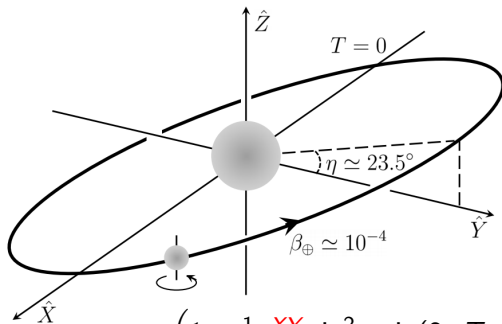
$$\chi = -G \int d^3x' \rho(\vec{x}', t)R$$

$$U^{jk} = \partial_j \partial_k \chi + \delta_{jk}U$$

in the lab ... gravimeters

Lorentz violation

$$\mathcal{L}_{\text{grav}} = \mathcal{L}_{\text{GR}} + \mathcal{L}_{\text{LV,grav}}$$

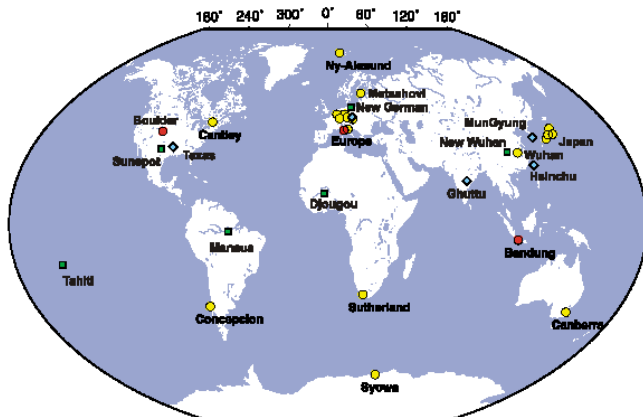
$$\mathcal{L}_{\text{LV,grav}} = e s^{\mu\nu} R_{\mu\nu} \dots$$



$$a = g \left(1 + \frac{1}{2} \bar{s}^{XY} \sin^2 \chi \sin(2\omega T + \phi) + \dots \right)$$

$\chi = \text{colatitude}$ $\omega = \text{sidereal frequency}$

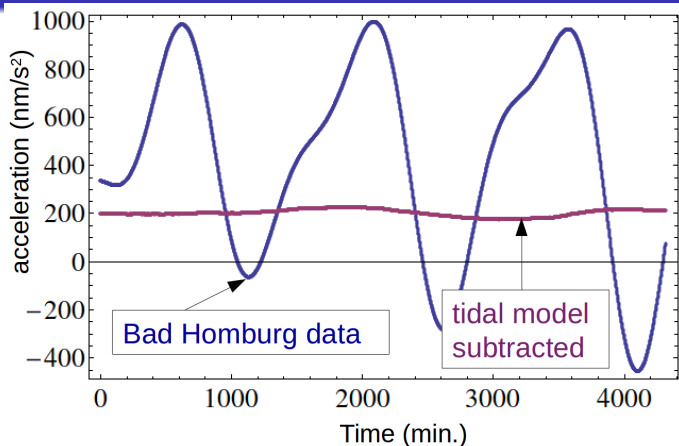
in the lab ... gravimeters

global geodynamics project



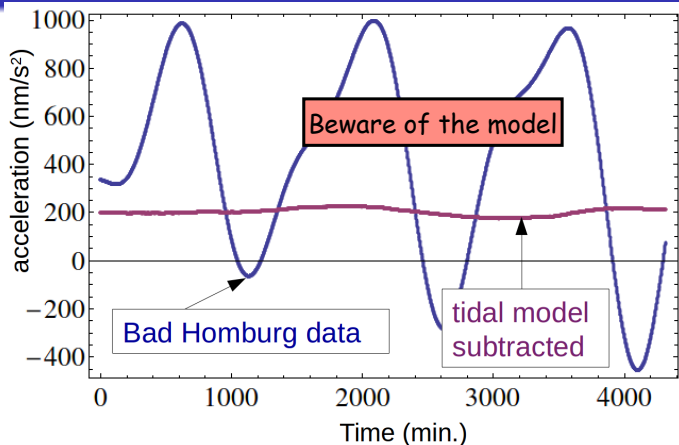
$$a = g \left(1 + \frac{1}{2} \bar{s}^{XY} \sin^2 \chi \sin(2\omega T + \phi) + \dots \right)$$

in the lab ... gravimeters



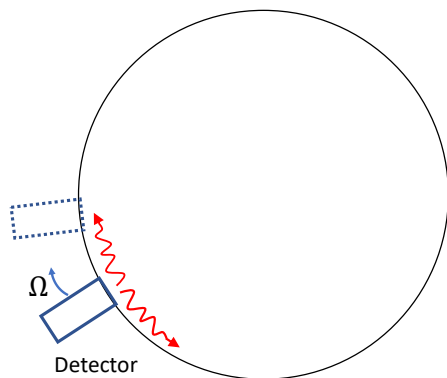
- among best lab bounds $\bar{s}^{JK} < 10^{-10}$
(c.f. solar system 10^{-12} , astrophysics 10^{-14})

in the lab ... gravimeters



- among best lab bounds $\bar{s}^{JK} < 10^{-10}$
(c.f. solar system 10^{-12} , astrophysics 10^{-14})

Sagnac gyros – Minkowski spacetime



Arrival time difference

$$\begin{aligned}\Delta t &\approx \frac{2\Omega r t}{c} \\ &= \frac{2\Omega r 2\pi r}{c^2} \\ &= \frac{4\Omega A}{c^2} \\ &\rightarrow 4 \int \vec{\Omega} \cdot d\vec{A}\end{aligned}$$

Phase difference per orbit

$$\Delta\psi = 2\pi \frac{c\Delta t}{\lambda}$$

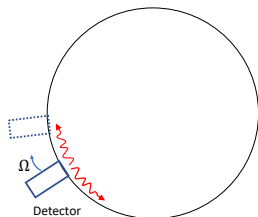
N – orbits to be in phase

$$2\pi = N\Delta\psi \quad N = \frac{\lambda}{c\Delta t}$$

beat period/frequency
via perimeter P

$$T = N\frac{P}{c} \Rightarrow f_b = \frac{\Delta t}{\lambda P}$$

Sagnac via the metric



...in rotating detector frame

light-like geodesics

$$0 = g_{00}dt^2 + \underline{2g_{0j}dtdx^j} + g_{jk}dx^jdx^k$$

quadratic in dt with 2 solutions

$$\Delta t = 2 \oint \frac{g_{0j}}{g_{00}} dx^j$$

in the cases of interest

$$\Delta \tau \approx 2 \oint g_{0j} dx^j$$

note that

$$2\Omega_j = \epsilon_{jkl} \partial_k g_{0l}$$

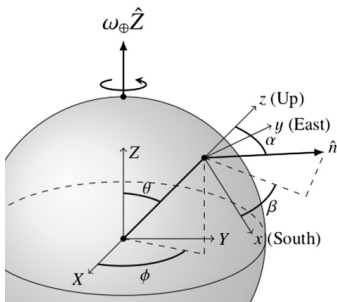
Ampere's law analogy

$$\Delta \tau \approx 4 \int \vec{\Omega} \cdot d\vec{A}$$

Conclusion: take the same approach for other contributions to g_{0j} in the detector frame

- gravitomagnetism
- Lorentz violation

SME results



beat frequency³

$$\frac{4AGM_{\oplus}}{\lambda PR_{\oplus}^2} \sin \alpha \left[\cos \beta (\bar{s}^{TX} \sin \phi - \dots) + \sin \beta (\cos \theta (\bar{s}^{TX} \cos \phi + \bar{s}^{TY} \sin \phi) + \dots) \right]$$

post-Newtonian metric²

$$g_{0j} = -\bar{s}^{0j} U - \bar{s}^{0k} U^{jk} + \frac{1}{2} \hat{Q}^j \chi$$

(U = Newton potential; U^{jk} , χ = post-Newton potentials)

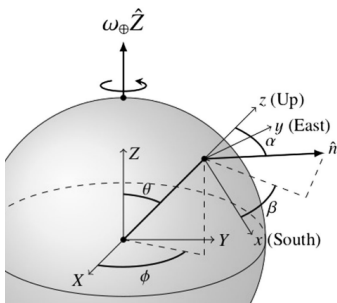
$$\text{where } \hat{Q}^j = [q^{(5)0jk0l0m} + q^{(5)n0knljm} + q^{(5)njknl0m}] \partial_k \partial_l \partial_m$$

look gravitomagnetic but the source is at rest

²Bailey&Havert PRD'17

³Moseley *et al.* PRD'19

SME results



beat frequency³

$$\frac{4AGM_{\oplus}}{\lambda PR_{\oplus}^2} \sin \alpha \left[\cos \beta (\bar{s}^{TX} \sin \phi - \dots) + \sin \beta (\cos \theta (\bar{s}^{TX} \cos \phi + \bar{s}^{TY} \sin \phi) + \dots) \right]$$

- sidereal dependence $\phi = \omega_{\oplus} t + \phi_0$
- orientation dependence
- $q^5 \Rightarrow$ additional harmonics
- lab-competitive \bar{s} sensitivities
- best q^5 sensitivities

post-Newtonian metric²

$$g_{0j} = -\bar{s}^{0j} U - \bar{s}^{0k} U^{jk} + \frac{1}{2} \hat{Q}^j \chi$$

(U = Newton potential; U^{jk} , χ = post-Newton potentials)

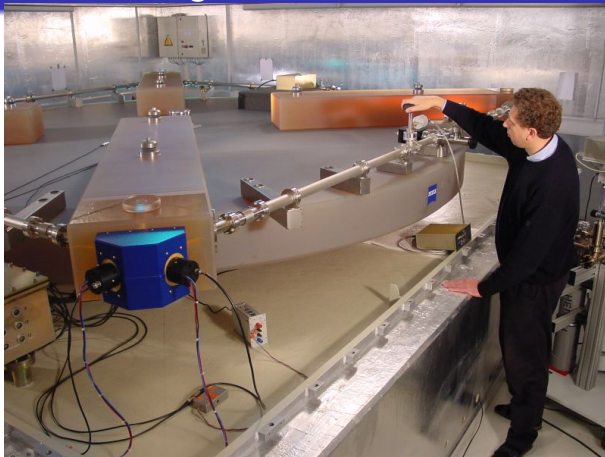
$$\text{where } \hat{Q}^j = [q^{(5)0jk0l0m} + q^{(5)n0knljm} + q^{(5)njknl0m}] \partial_k \partial_l \partial_m$$

look gravitomagnetic but the source is at rest

²Bailey&Havert PRD'17

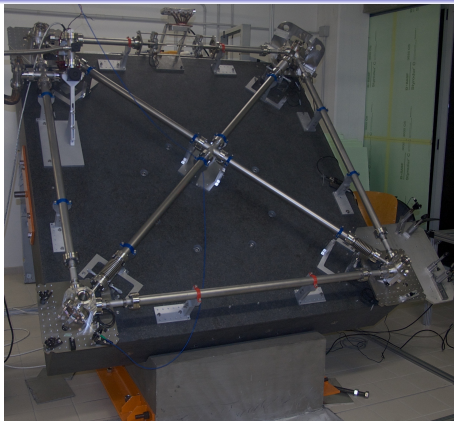
³Moseley *et al.* PRD'19

the experiment...G-ring



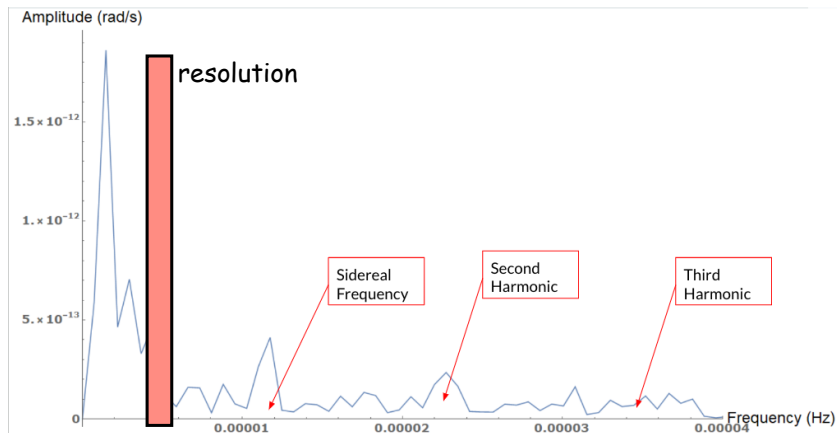
- G-ring, Germany
- insensitive to minimal Lorentz violation

the experiment...GINGER



- GINGERINO, Gyroscopes IN GEneral Relativity, Italy
- sensitive to both minimal and $d = 5$ Lorentz violation

15 days of G-ring



- polar motion corrected

preliminary maximum reach sensitivities

K_{xxxxy}	K_{xxxz}	K_{xxyy}	K_{xxyz}	K_{xxzz}	K_{xyyy}	K_{xyyz}	K_{xyzz}
-16 m	10 m	-113 m	32 m	113 m	48 m	-10 m	-16 m

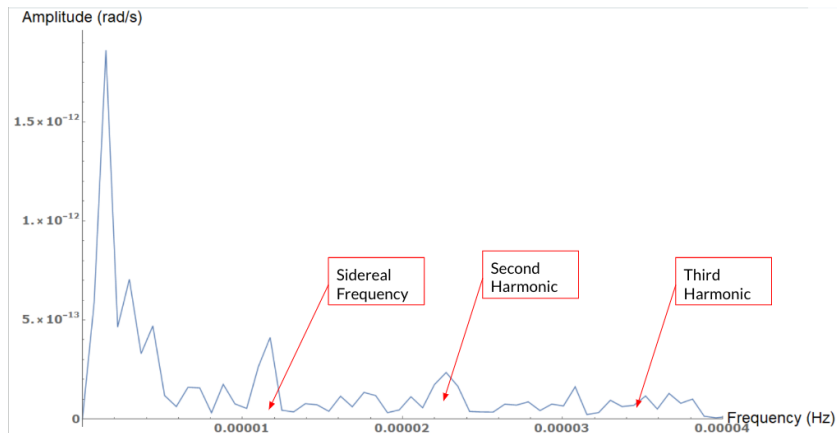
K_{xzzz}	K_{yxxz}	K_{yxyz}	K_{yxzz}	K_{yyyz}	K_{yyzz}	K_{yzzz}
190 m	10 m	-5 m	-16 m	8 m	-113 m	48 m

- compare with pulsar results

$$K_{XXXXY} = \frac{1}{3}(-q^{TXYTXTX} + q^{TXYXYXY} + q^{TXYXZXZ} - q^{XYZXZXT})$$

Coefficient	1- σ limit [m]	Coefficient	1- σ limit [m]
$q^{TXYTXTX}$	22	$q^{TXYTXTY}$	11
$q^{TXYTPTY}$	10	$q^{TXYTPTZ}$	5.7
$q^{TXYXYXY}$	8.0	$q^{TXYXYXZ}$	8.3

15 days of G-ring



- polar motion corrected

what needs doing? Uncertainty!

- We see a signal at some of the frequencies of interest
 - what is the uncertainty in the polar motion correction?
 - is the polar motion modeling trained on data?
 - systematics at the solar day frequency?
- how should we quantify uncertainty?
 - rms of neighboring frequencies → more resolution needed
 - understanding of systematics and modeling uncertainties
- want to collaborate?

summary

- systematic search for new physics
- basic theory in place for many gravitational tests
- expanding phenomenological & experimental breadth & depth
- Sagnac Gyroscopes can be an interesting test
- More data/better understanding of error needed