

$U(1)$ Entanglement asymmetry in the Ising CFT

How much is a defect non-topological?

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Outline

- Introduce entanglement asymmetry
- Provide interpretation in terms of (non) topological defects
- Show explicit calculation for Ising CFT

Entanglement asymmetry in general

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Ingredients: a **state** ρ and a **group action** on the Hilbert space $U_g, g \in G$.

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$$\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_{\bar{A}} \quad \rho_A = \text{tr}_{\bar{A}} \rho \quad U_g = U_{A,g} \otimes U_{\bar{A},g} \quad \rho_A \mapsto U_{A,g}^{-1} \rho_A U_{A,g}$$

Symmetrized state:

$$\rho_{A,G} = \int_{g \in G} U_{A,g} \rho_A U_{A,g}^{-1}$$

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Entanglement asymmetry = difference in Renyi entropy between $\rho_{A,G}$ and ρ_A

$$\Delta S_n = \frac{1}{1-n} (\log \text{Tr} \rho_{A,G}^n - \log \text{Tr} \rho_A^n)$$

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Properties ¹

- $\Delta S_n \geq 0$
- $\Delta S_n = 0 \Leftrightarrow \rho_A = \rho_{A,G}$
- $\Delta S_n \leq \log |G|$ for finite groups, unbounded for G continuous

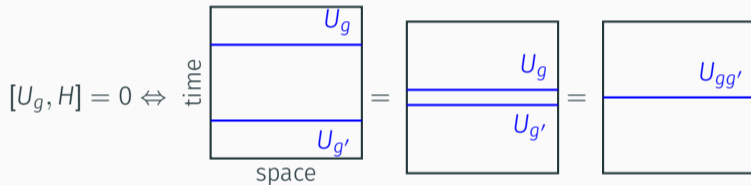
¹Ares, Murciano, Calabrese *Nat. Comm.* (2023)

Ferro, Ares, Calabrese (2023)

Capizzi, Mazzoni (2023)

Interpretation with defects

“Modern language”: symmetries \leftrightarrow topological operators/defects²

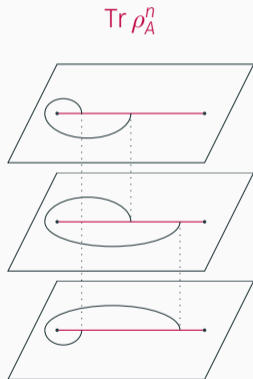


On the contrary, if the symmetry is broken (explicitly or spontaneously), the defects are not topological

²Gaiotto, Seiberg, Kapustin, Willett *JHEP* (2015)

Entanglement asymmetry with defects

$$\Delta S_n = \frac{1}{1-n} \log \frac{\text{Tr} \rho_{A,G}^n}{\text{Tr} \rho_A^n}$$

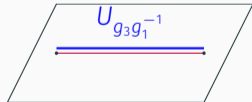
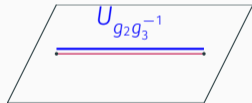
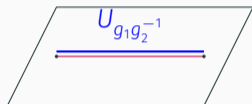
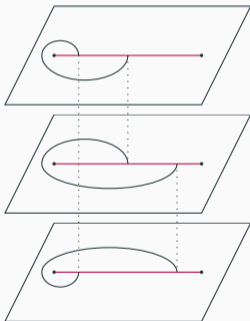


Entanglement asymmetry with defects

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$$\text{Tr} \rho_{A,G}^n = \int_{g_1, \dots, g_n} \text{Tr} U_{A,g_1}^{-1} \rho_A U_{A,g_1} U_{A,g_2}^{-1} \dots U_{A,g_n}^{-1} \rho_A U_{A,g_n}$$

$\text{Tr} \rho_A^n$

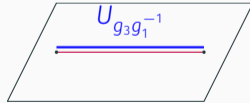
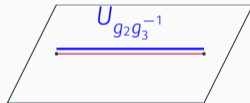
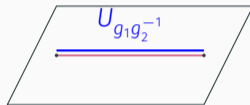
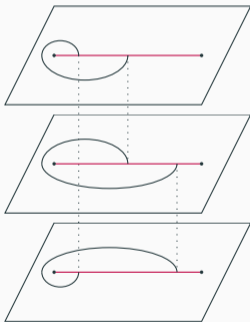


Entanglement asymmetry with defects

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$\text{Tr} \rho_A^n$



Entanglement asymmetry = how much the defects are not topological

Our case

Our system

LATTICE

State: $|gs\rangle$ of critical Ising chain: $H = -\sum_{j \in \mathbb{Z}} (\sigma_j^x \sigma_{j+1}^x + \sigma_j^z)$



Group action: $U_{A,\alpha} = \exp\left(i\frac{\alpha}{2} \sum_{j \in A} \sigma_j^z\right)$ rotation around \hat{z} .

Not a symmetry, unless $\alpha = \pi$ (\mathbb{Z}_2 symmetry)

³Ares, Murciano, Calabrese *Nat. Comm.* (2023)

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Known results: from lattice calculations³

$$\Delta S_n = \frac{1}{2} \log \ell - \frac{1}{2} \log \frac{4}{\pi} + o(1)$$

Goal: access subleading terms with CFT

³Ares, Murciano, Calabrese *Nat. Comm.* (2023)

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CONTINUUM

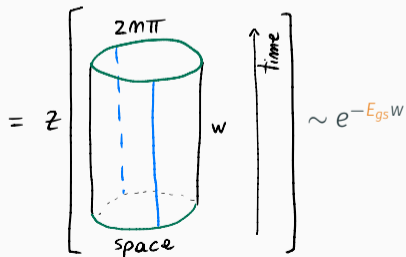
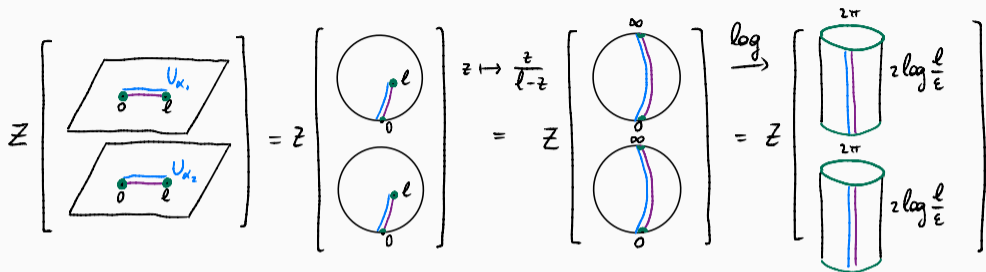
State: ground state of the Majorana CFT: $H = \frac{1}{2i} \int_{\mathbb{R}} dx (\psi \partial_x \psi - \bar{\psi} \partial_x \bar{\psi})$

Group action: $U_{A,\alpha} = \exp(-\alpha \int_A dx \psi \bar{\psi})$

$$U_{\alpha} \begin{pmatrix} \psi \\ \bar{\psi} \end{pmatrix} U_{\alpha}^{-1} = \underbrace{\begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix}}_{\in SO(2)} \begin{pmatrix} \psi \\ \bar{\psi} \end{pmatrix}$$

again, **not a symmetry** unless $\alpha = \pi : \psi \mapsto -\psi, \bar{\psi} \mapsto -\bar{\psi}$

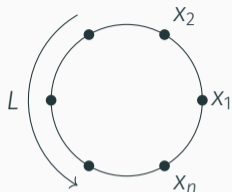
Conformal mappings



$$U_{\alpha} = \exp \left(-\alpha \int dx \psi(x) \overbrace{\bar{\psi}(x)}^{\Delta=1} \right)$$

transforms as a scalar

Hamiltonian with defects



$$H = \frac{1}{2i} \int_x (\psi \partial \psi - \bar{\psi} \partial \bar{\psi}) + \sum_j \overbrace{\alpha_j \psi(x_j) \bar{\psi}(x_j)}^{\text{anti-Hermitian}} = \frac{1}{2i} \int_x \Psi^\dagger D \Psi$$

analytic continuation $\alpha_j = i\lambda_j$

$$\Psi = \begin{pmatrix} \psi \\ \bar{\psi} \end{pmatrix} \quad D = \begin{pmatrix} \partial_x & -\sum_j \lambda_j \delta(x - x_j) \\ \sum_j \lambda_j \delta(x - x_j) & -\partial_x \end{pmatrix}$$

Hamiltonian is quadratic \rightarrow diagonalize kernel $D \rightarrow$ compute E_{gs}

Diagonalized Hamiltonian

$$H = \frac{1}{2} \sum_k k (\eta_k^\dagger \eta_k + \bar{\eta}_k^\dagger \bar{\eta}_k) \quad E_{gs} = \sum_{k < 0} k$$

n. defects	momentum space	E_{gs}
0	$\frac{2\pi}{L} (\mathbb{Z} + \frac{1}{2})$	$-\frac{\pi}{12L}$
2	$\frac{2\pi}{L} (\mathbb{Z} + \frac{1}{2} + \frac{\beta_1 - \beta_2}{2\pi})$	$\frac{1}{L} \left[-\frac{\pi}{12} + \frac{(\beta_1 - \beta_2)^2}{\pi} \right]$
3	$\bigcup_{j=1}^3 \frac{6\pi}{L} (\mathbb{Z} + \frac{\theta_j}{2\pi})$	$\frac{6\pi}{L} \left[\frac{1}{4} - \sum_{j=1}^3 \left(\frac{\theta_j}{4\pi} - \frac{1}{2} \left(\frac{\theta_j}{2\pi} \right)^2 \right) \right]$

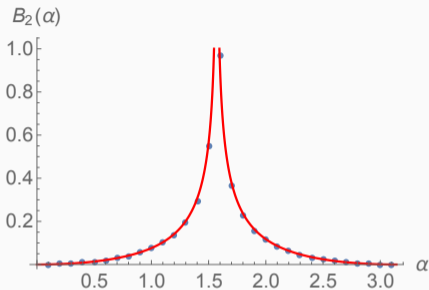
$$\beta_j = 2 \arctan \tanh \frac{\lambda_j}{2}$$

$$\theta_1 = \text{Arg} \left[\frac{1}{2} \left(1 - \sqrt{-s} - \sqrt{-3 - s - 2\sqrt{-s}} \right) \right] \quad s = \sin \beta_1 \sin \beta_2 + \text{cycl. perm.}$$

2-asymmetry ΔS_2

$E_{gs} \rightsquigarrow$ partition function

$$\frac{\text{Tr} [\rho_A e^{i\alpha Q_A} \rho_A e^{-i\alpha Q_A}]}{\text{Tr} \rho_A^2} = \underbrace{e^{A_2(\alpha)\ell}}_{\text{from lattice}} \underbrace{\ell^{B_2(\alpha)}}_{\text{from CFT}}$$



$$B_2(\alpha) = \frac{2}{\pi^2} \text{arctanh}^2 \left(\tan \frac{\alpha}{2} \right)$$

$\Delta S_2 = -\log \int_{\alpha} e^{A_2(\alpha)\ell} \ell^{B_2(\alpha)} \rightarrow$ saddle point around $\alpha = 0$

$$A_2(\alpha) = a_2\alpha^2 + \dots \quad B_2(\alpha) = b_2\alpha^2 + \dots$$

$$\Delta S_2 = \frac{1}{2} \log \ell + \frac{1}{2} \log(-a_2\pi) + \frac{b_2}{2a_2} \frac{\log \ell}{\ell} + o\left(\frac{\log \ell}{\ell}\right)$$

3-asymmetry ΔS_3

$$\frac{\text{Tr} [\rho_A e^{i\alpha_1 Q_A} \rho_A e^{i\alpha_2 Q_A} \rho_A e^{i(-\alpha_1 - \alpha_2) Q_A}]}{\text{Tr} \rho_A^3} = \underbrace{e^{A_3(\boldsymbol{\alpha})\ell}}_{\text{from lattice}} \overbrace{\ell^{B_3(\boldsymbol{\alpha})}}^{\text{from CFT}}$$

$$\Delta S_3 = -\frac{1}{2} \log \int_{\boldsymbol{\alpha}} e^{A_3(\boldsymbol{\alpha})\ell} \ell^{B_3(\boldsymbol{\alpha})} \rightarrow \text{saddle point around } \boldsymbol{\alpha} = 0$$

$$A_3(\boldsymbol{\alpha}) = a_3 \cdot (\alpha_1^2 + \alpha_2^2 + \alpha_1 \alpha_2) + \dots \quad B_3(\boldsymbol{\alpha}) = b_3 \cdot (\alpha_1^2 + \alpha_2^2 + \alpha_1 \alpha_2) + \dots$$

$$\Delta S_3 = \frac{1}{2} \log \ell + \frac{1}{2} \log(-a_3 \pi) + \frac{b_3}{2a_3} \frac{\log \ell}{\ell} + o\left(\frac{\log \ell}{\ell}\right)$$

Same structure as $n = 2$

Conclusions

- Entanglement asymmetry: how much the defects are not topological
- CFT approach $\rightarrow E_{gs}$ of a system with defects
- $\frac{\log l}{l}$ corrections in ΔS_n

Future directions

- Dirac CFT, same group action (again, not a symmetry)
- Extend to $SU(N)$ (e.g. XXZ spin chain)

Thank You!