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Istituto Nazionale di Fisica Nucleare



Transversely polarized Λ production in e^+e^- annihilation and SIDIS processes within a TMD approach

Marco Zaccheddu – Università degli Studi di Cagliari & INFN

In collaboration with:

Umberto D'Alesio, Leonard Gamberg and Francesco Murgia

Sar Wors 2023

Motivations and Contents

Observation of Transverse $\Lambda/\bar{\Lambda}$ Hyperon Polarization in e^+e^- Annihilation at Belle

- 2 data set @ $\sqrt{s} = 10.58$ GeV

[Y. Guan et al., Phys. Rev. Lett. 122. 042001 (2019)]

Double hadron production:

- $e^+e^- \rightarrow \Lambda\pi/K + X$: 128 points - bins of the energy fractions $z_\Lambda - z_{\pi,K}$

Single-inclusive hadron production:

- $e^+e^- \rightarrow \Lambda(jet) + X$: 32 points - $\Lambda(jet)$, in bins of $z_\Lambda - p_\perp$

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- Callos, Kang, Terry; *Phys.Rev.D* 102 (2020) 9, 096007
- Chen, Liang, Pan, Song, Wei; *Phys.Lett.B* 816 (2021) 136217
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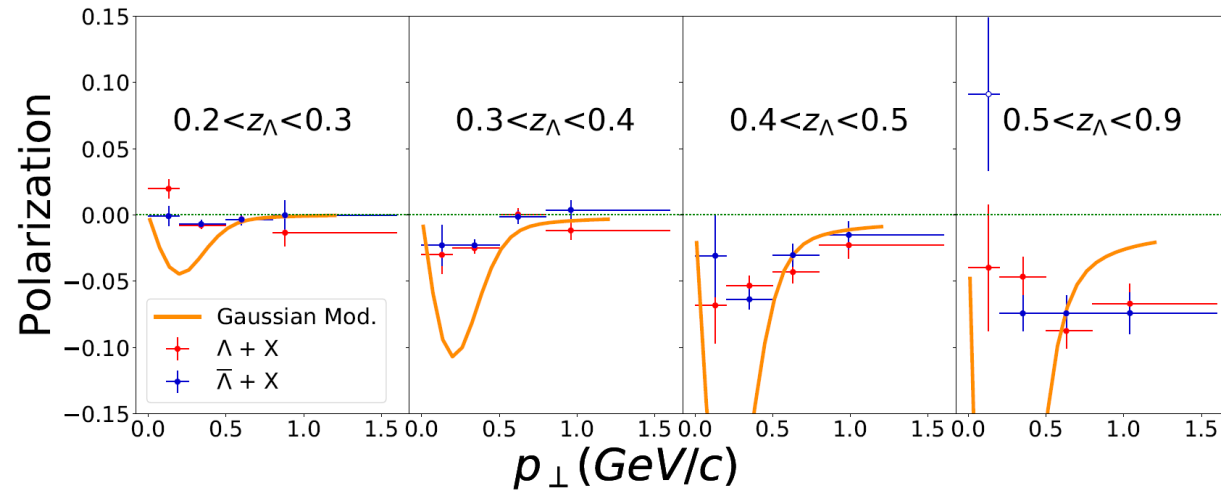
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Two open issues!

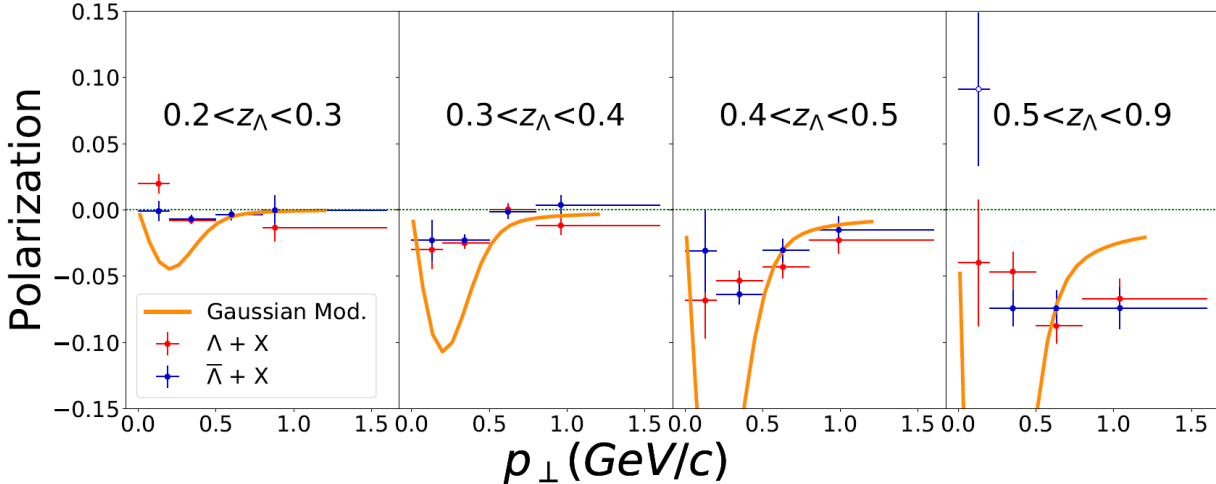
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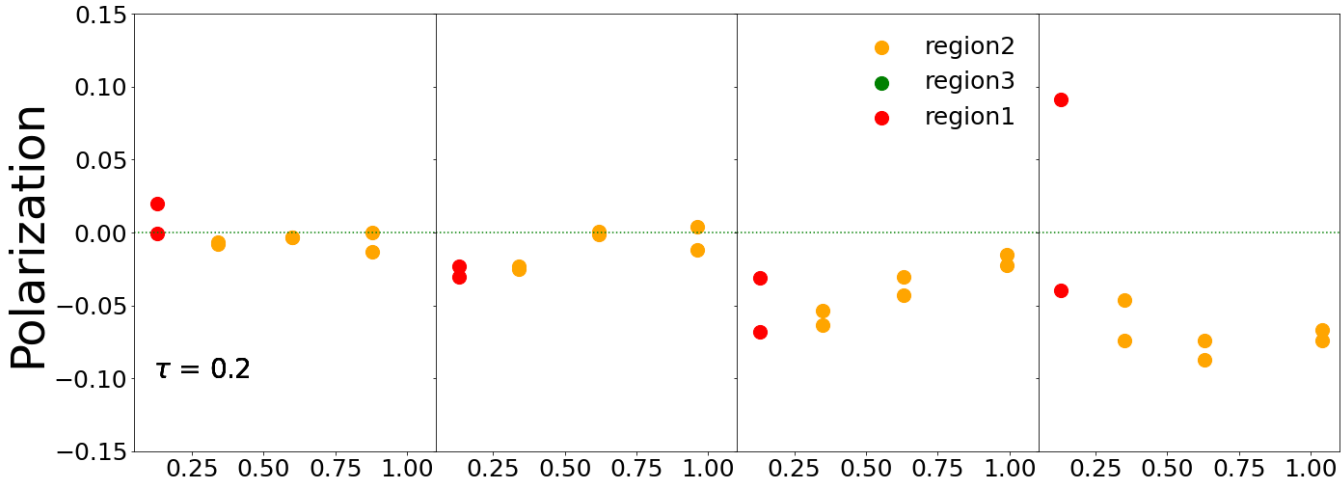
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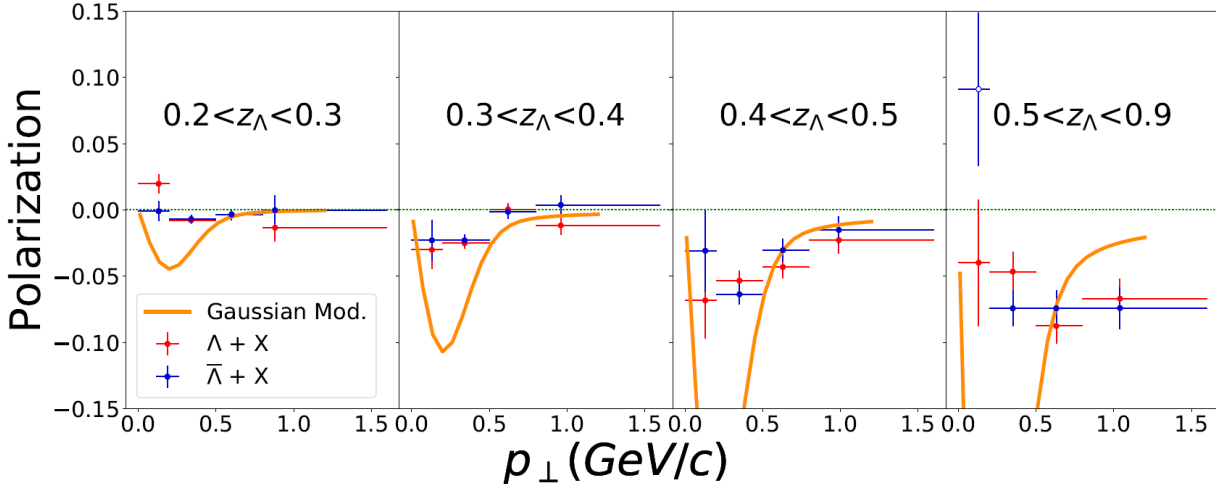
See A. Simonelli's talk!



Data not binned in Thrust!

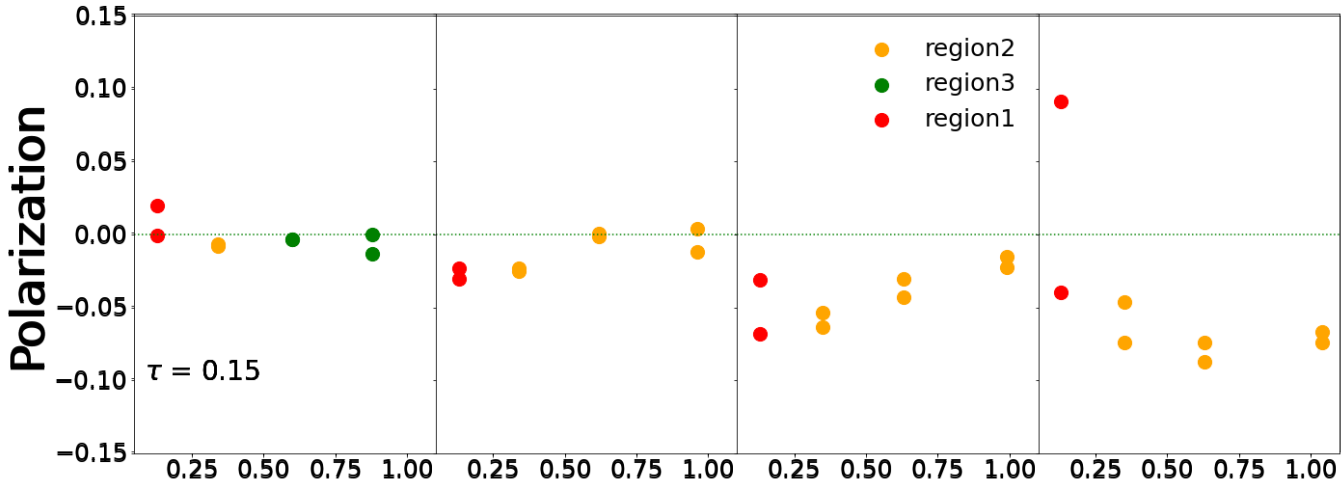
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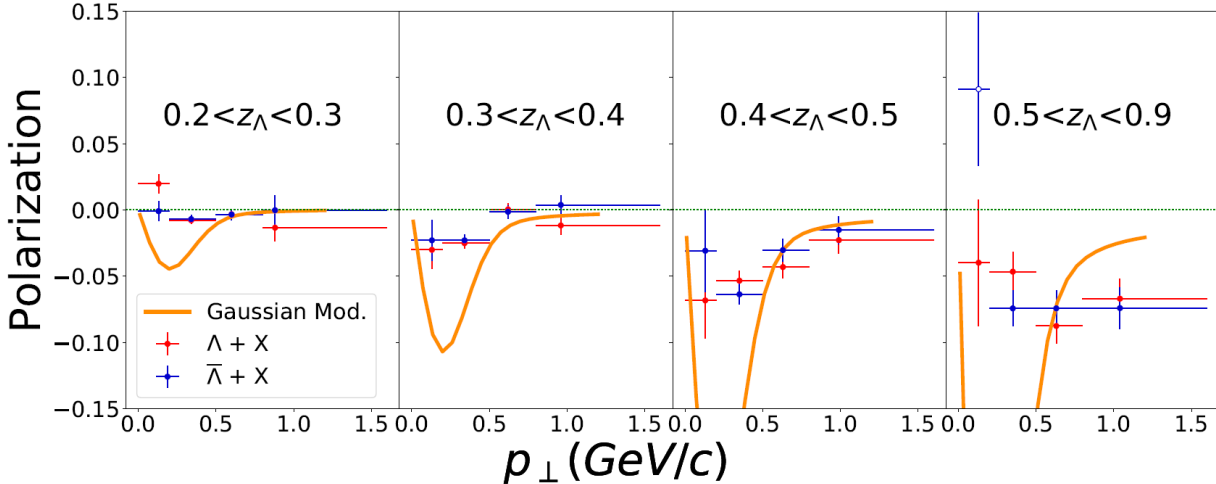
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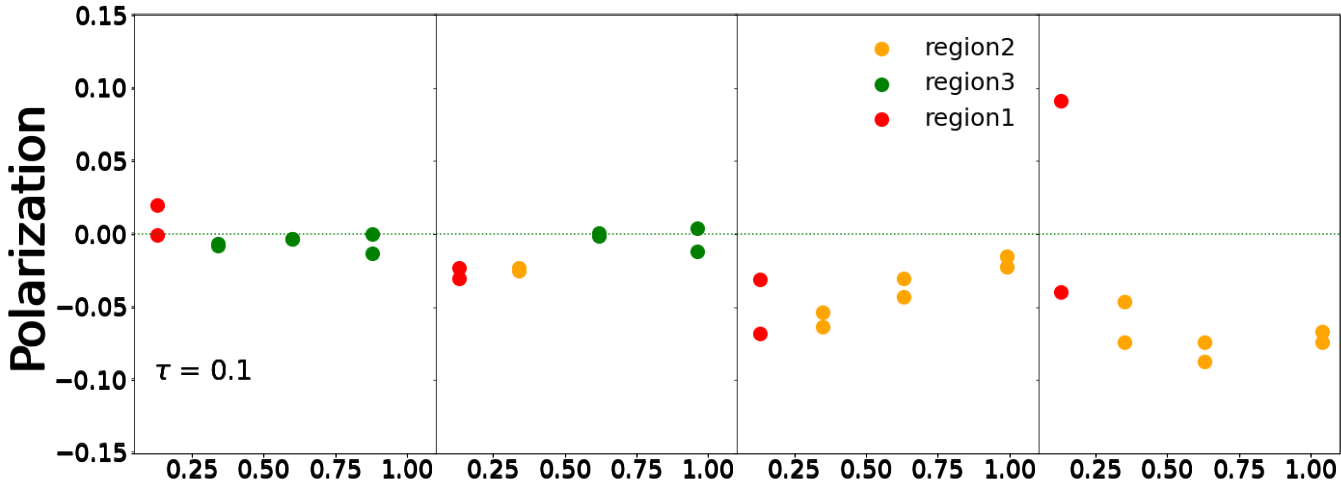
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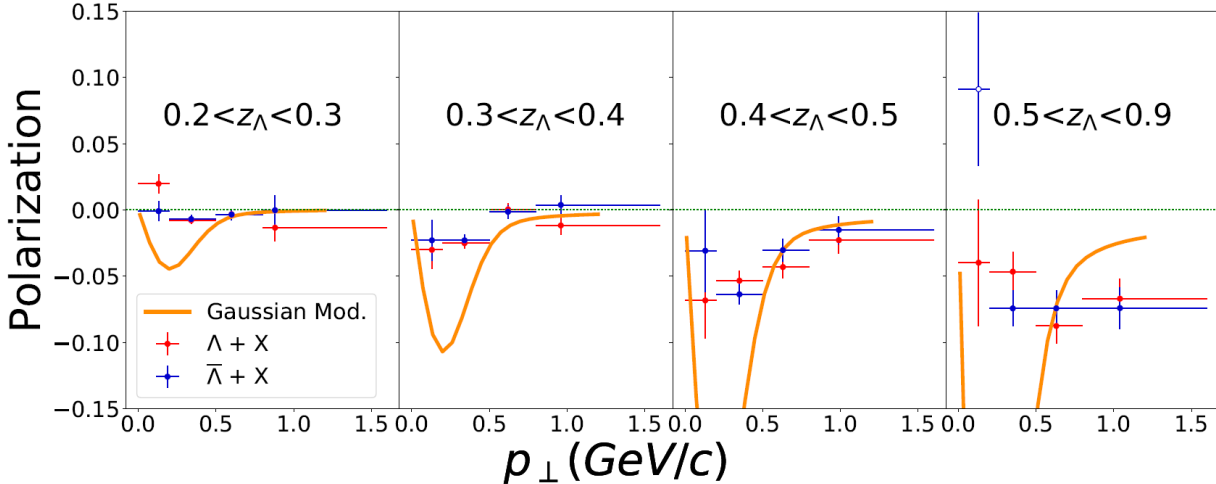
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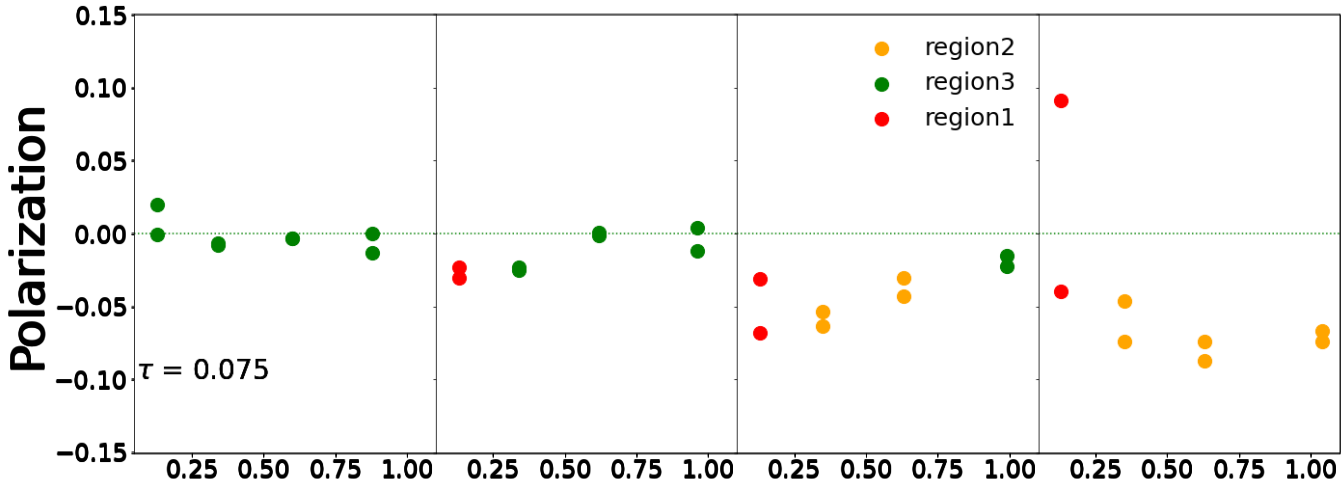
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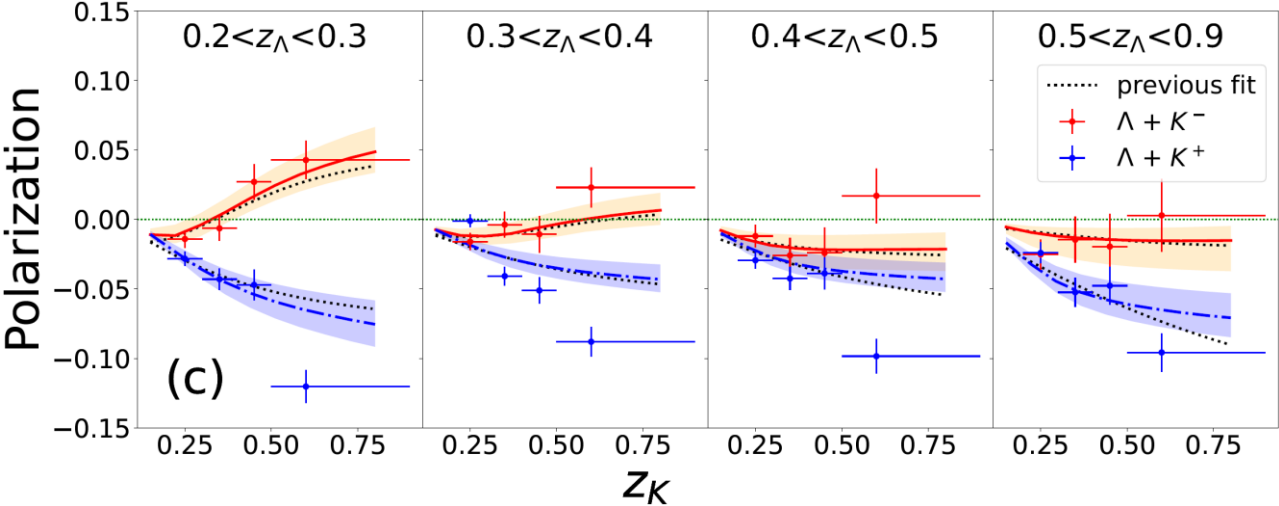
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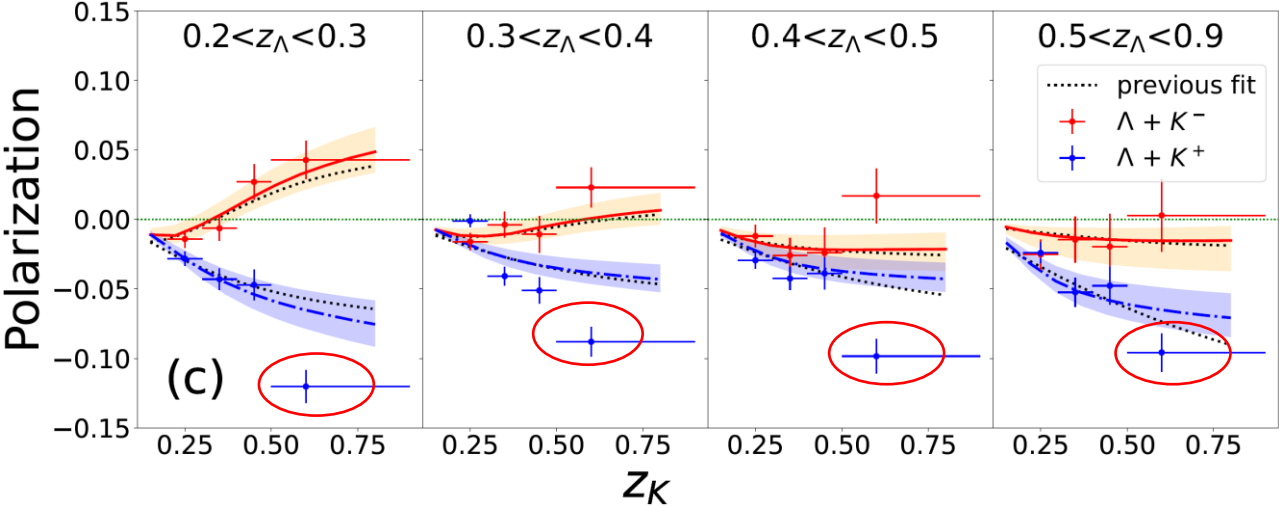
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- Preliminary analysis: necessary to include heavier flavors

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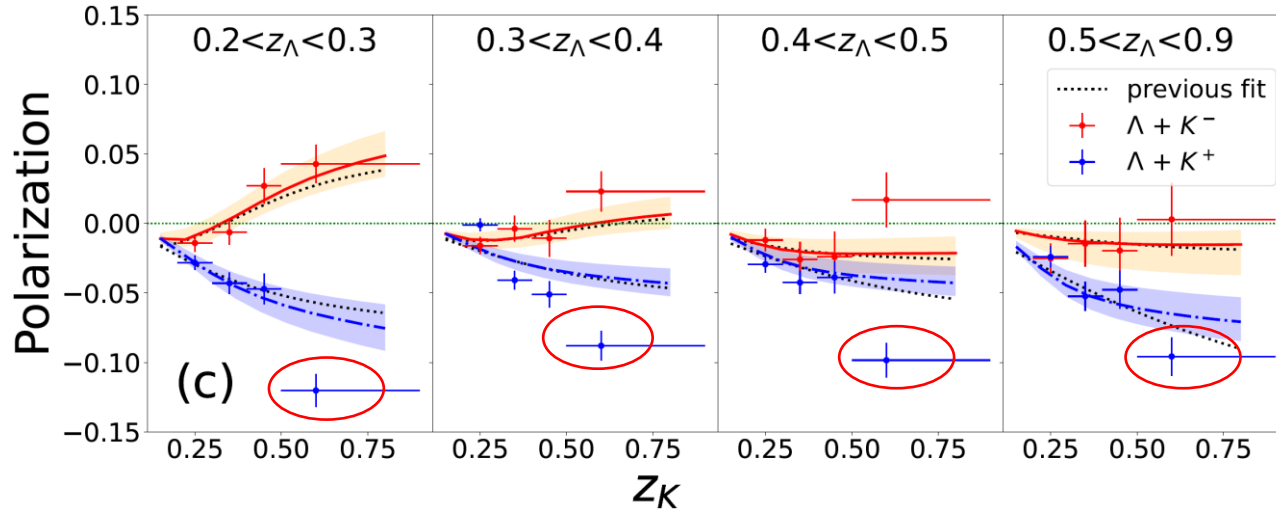
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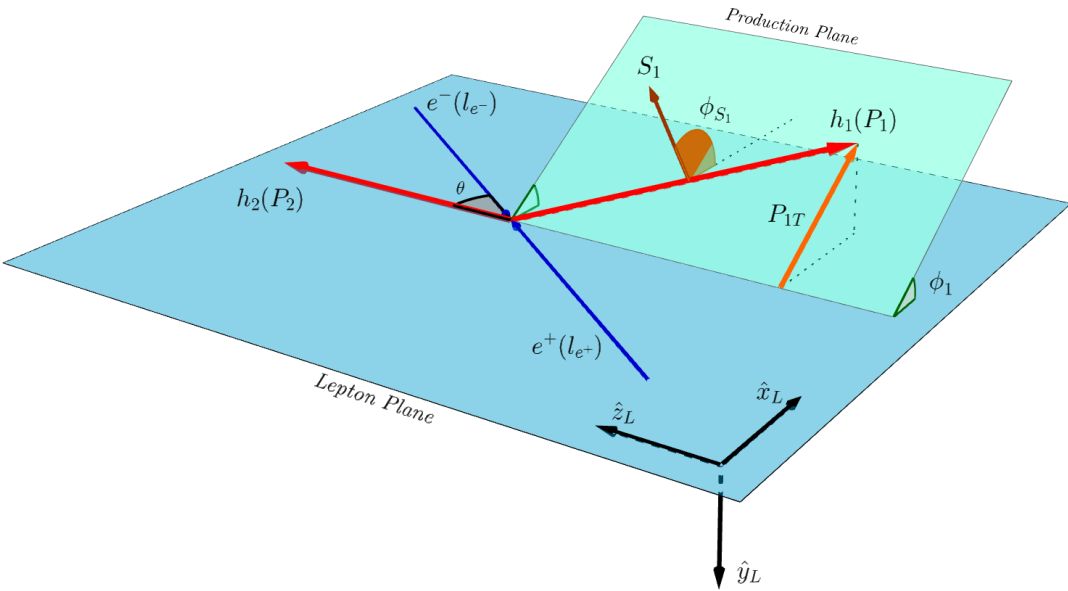


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Contents:

- Convolutions and Polarization: $e^+e^- \rightarrow h_1^\uparrow h_2 X$
- Fit results: Belle 2-h
- SIDIS predictions and Intrinsic Charm
- Conclusions

Double hadron production in e^+e^- processes



Polarization:

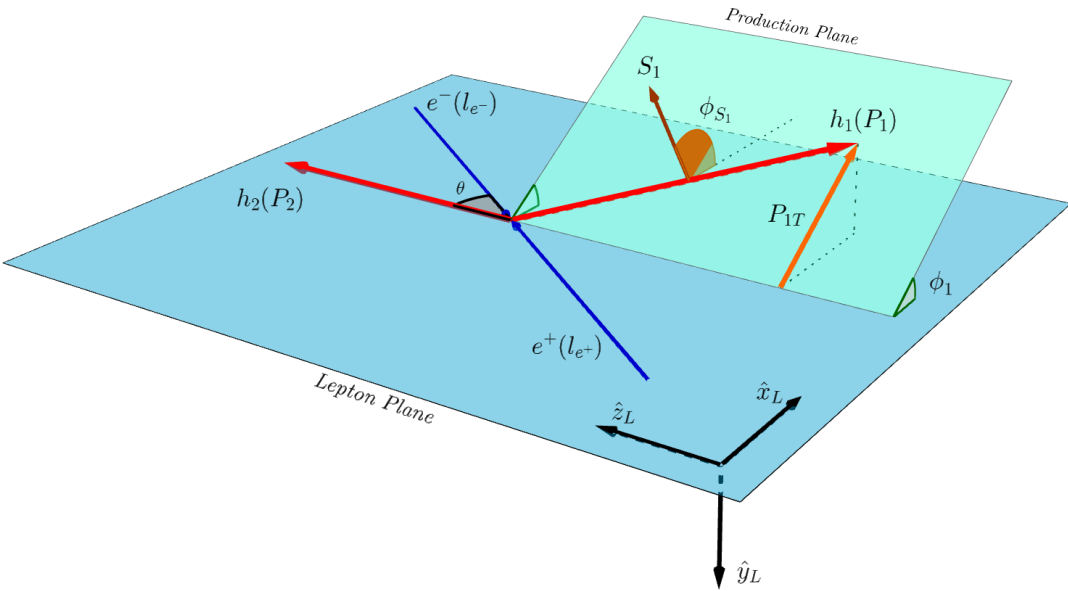
$$P_n^{h_1}(z_{h_1}, z_{h_2}) = \frac{M_1 \int dq_T q_T d\phi_1 \mathcal{B}_1 \left[\tilde{D}_{1T}^{\perp(1)} \tilde{D}_1 \right]}{\int dq_T q_T d\phi_1 \mathcal{B}_0 \left[\tilde{D}_1 \tilde{D}_1 \right]}$$

Convolutions:

$$\mathcal{B}_0 \left[\tilde{D} \tilde{D} \right] = \frac{1}{z_1^2 z_2^2} \sum_q e_q^2 \int \frac{db_T}{2\pi} b_T J_0(b_T q_T) d_{q/h_1}(z_1; \bar{\mu}_b) d_{\bar{q}/h_2}(z_2; \bar{\mu}_b) \\ \times M_{D_1}(b_c(b_T), z_1) M_{D_2}(b_c(b_T), z_2) e^{-g_K(b_c(b_T); b_{\max}) \ln \left(\frac{Q^2 z_1 z_2}{M_1 M_2} \right) - S_{\text{pert}}(b_*; \bar{\mu}_b)}$$

$$\mathcal{B}_1 \left[\tilde{D}_{1T}^{\perp(1)} \tilde{D}_1 \right] = \frac{1}{z_1^2 z_2^2} \sum_q e_q^2 \int \frac{db_T}{(2\pi)} b_T^2 J_1(b_T q_T) D_{1T}^{\perp(1)}(z_1; \bar{\mu}_b) d_{\bar{q}/h_2}(z_2; \bar{\mu}_b) \\ \times M_{D_1}^{\perp}(b_c(b_T), z_1) M_{D_2}(b_c(b_T), z_2) e^{-g_K(b_c(b_T); b_{\max}) \ln \left(\frac{Q^2 z_1 z_2}{M_1 M_2} \right) - S_{\text{pert}}(b_*; \bar{\mu}_b)},$$

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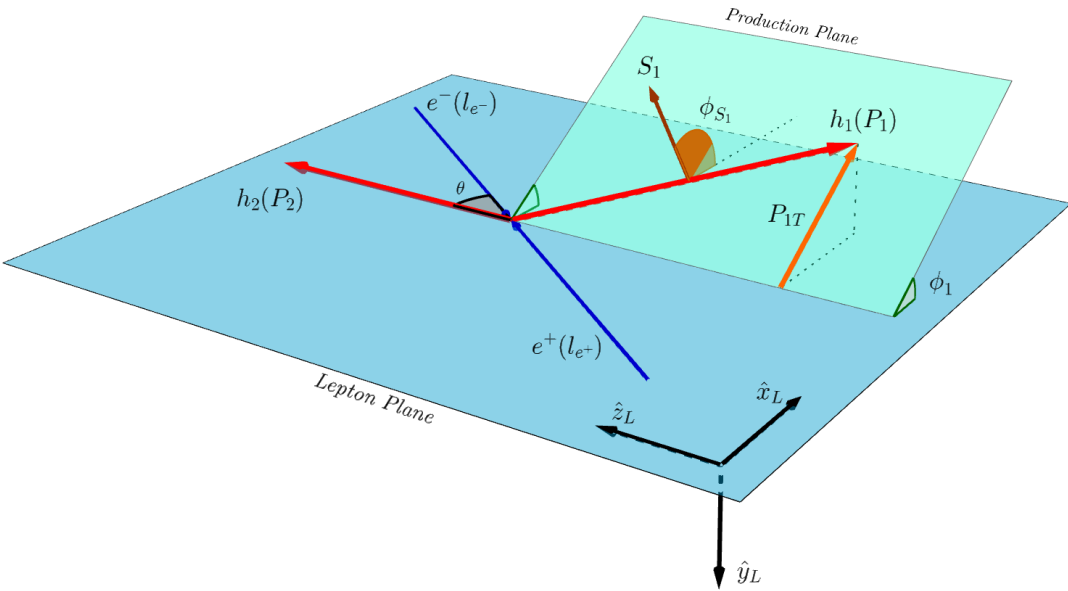
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Unpolarized FFs:
DSS set for π/K
AKK set for Λ

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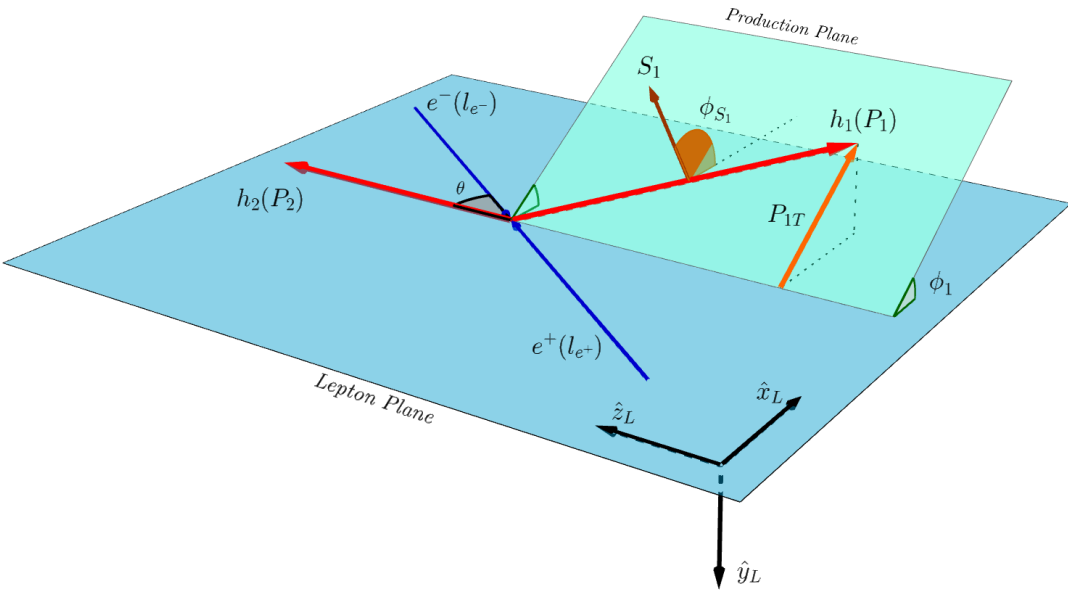
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Polarizing FF first moment

$$\tilde{D}_{1T, \Lambda/q}^{\perp(1)}(z; \mu_b) = \mathcal{N}_q^p(z) d_{q/\Lambda}(z; \mu_b)$$

$$\mathcal{N}_q^p(z) = N_q z^{a_q} (1-z)^{b_q} \frac{(a_q + b_q)^{(a_q + b_q)}}{a_q^{a_q} b_q^{b_q}}$$

Polarizing FF non-perturbative function

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u, d, s + charm

Isospin symmetry

- No SU(2): N_u, N_d, N_s, N_{sea}
- SU(2): $N_u = N_d, N_{\bar{u}} = N_{\bar{d}}, N_s, N_{\bar{s}}$

See also:

Chen, Liang, Pan, Song, Wei;
Phys.Lett.B 816 (2021) 136217

Double hadron production in e^+e^- processes

Scenarios considered:

1. No Charm, No SU(2) sym.

pFFs for: *up, down strange and sea*;

2. Charm, No SU(2) sym.

pFFs for: *up, down strange and sea*;

3. Charm, SU(2) sym.

pFFs for: *up/down, $\overline{up}/\overline{down}$, strange, $\overline{strange}$*

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χ_{dof}^2
96 points
1,174
1,259
1,361

Double hadron production in e^+e^- processes

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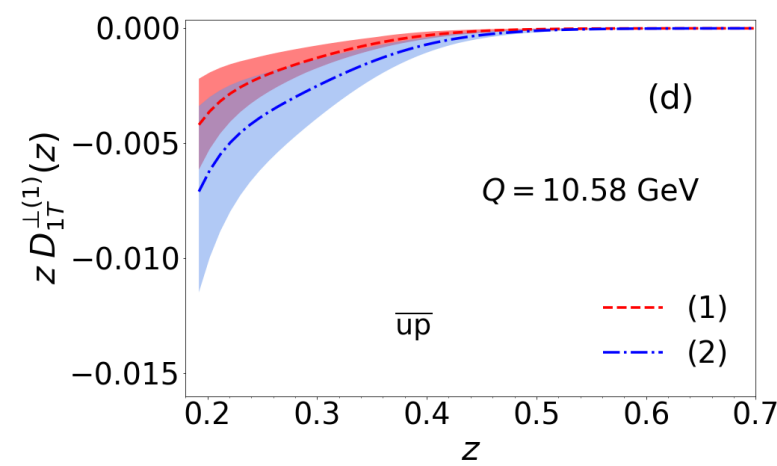
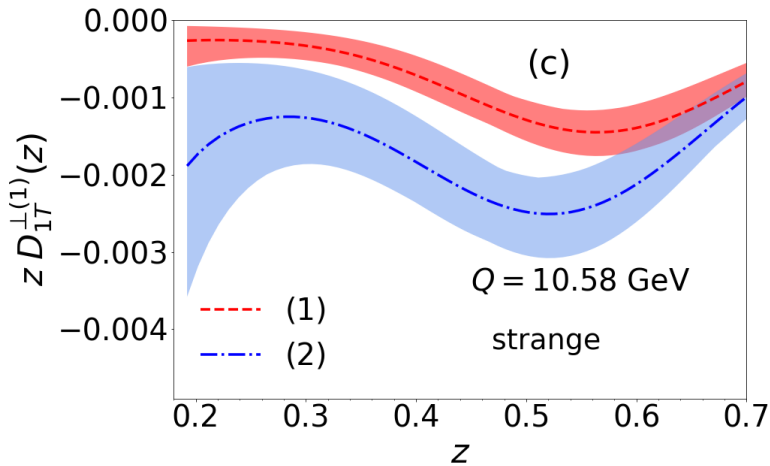
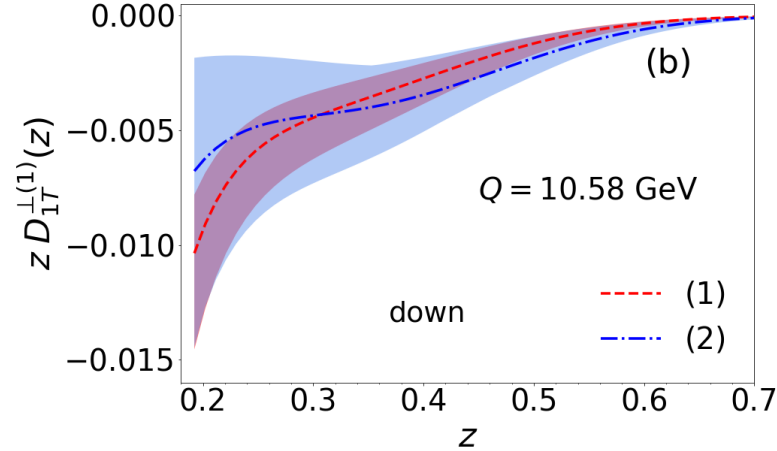
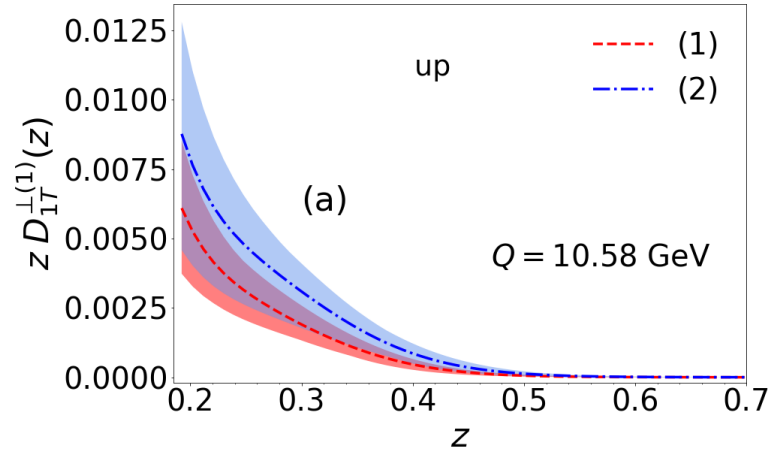
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χ^2_{dof}	χ^2_{dof}
96 points	128 points
1,174	1,903
1,259	1,622
1,361	1,645

Double hadron production in e^+e^- processes

First moments: (1) & (2) scenarios

- (1) No Charm, No SU2
- (2) Charm, No SU2

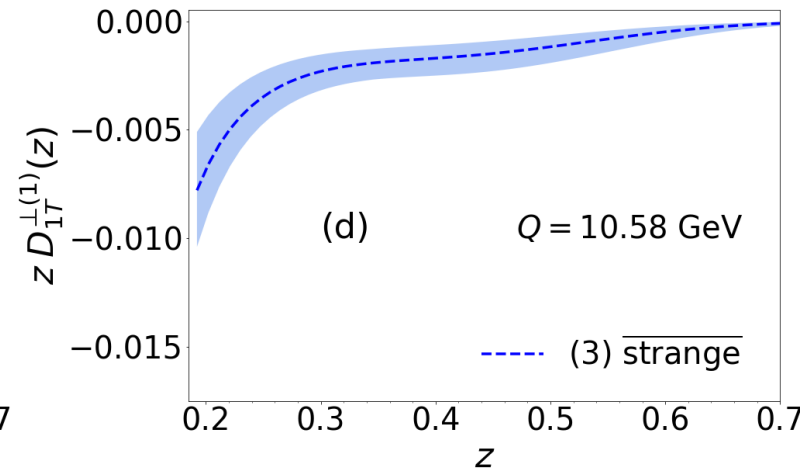
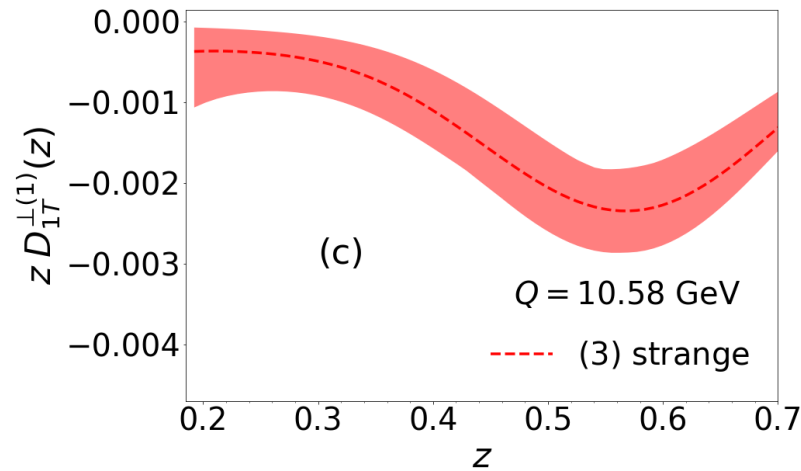
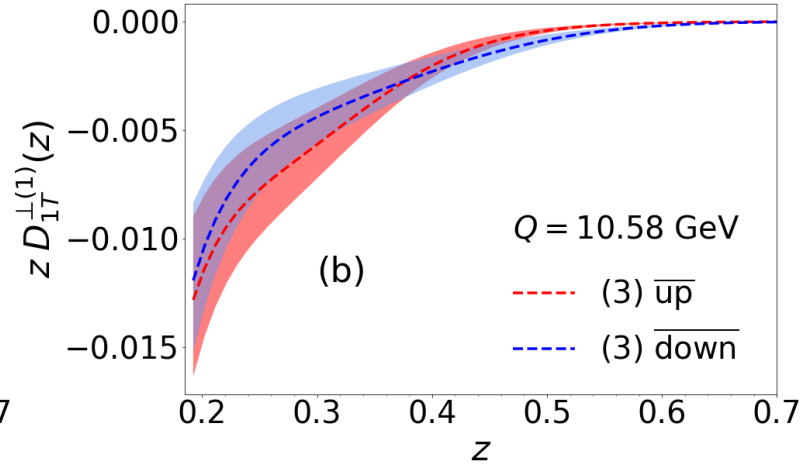
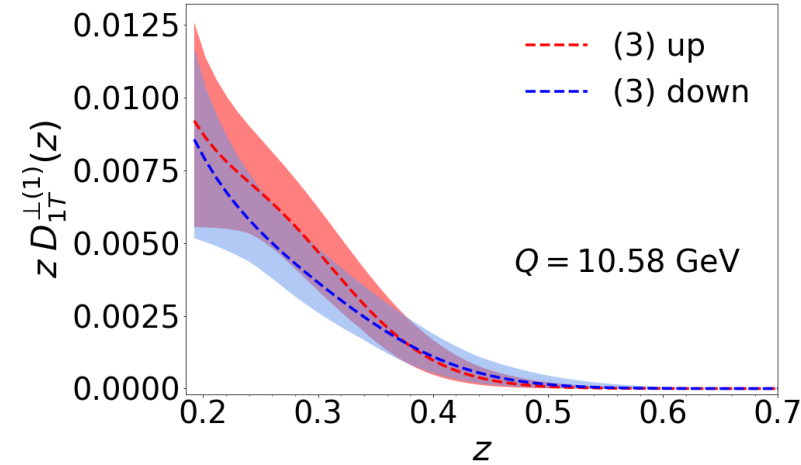


- pFFs are different in magnitude due to the charm contribution;
- up pFF is positive;
- First moments are compatible, except for the strange f.m.
- Similar size for the Gaussian width.

Double hadron production in e^+e^- processes

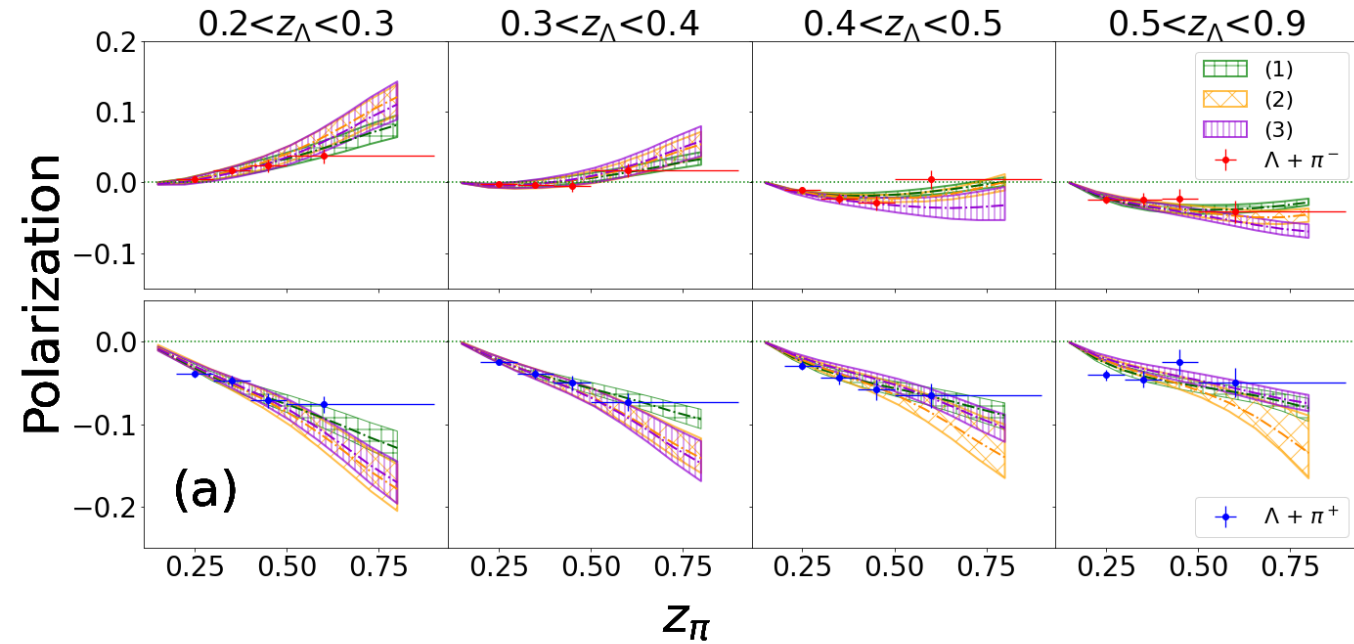
First moments: (3) scenario

(3) Charm, SU2



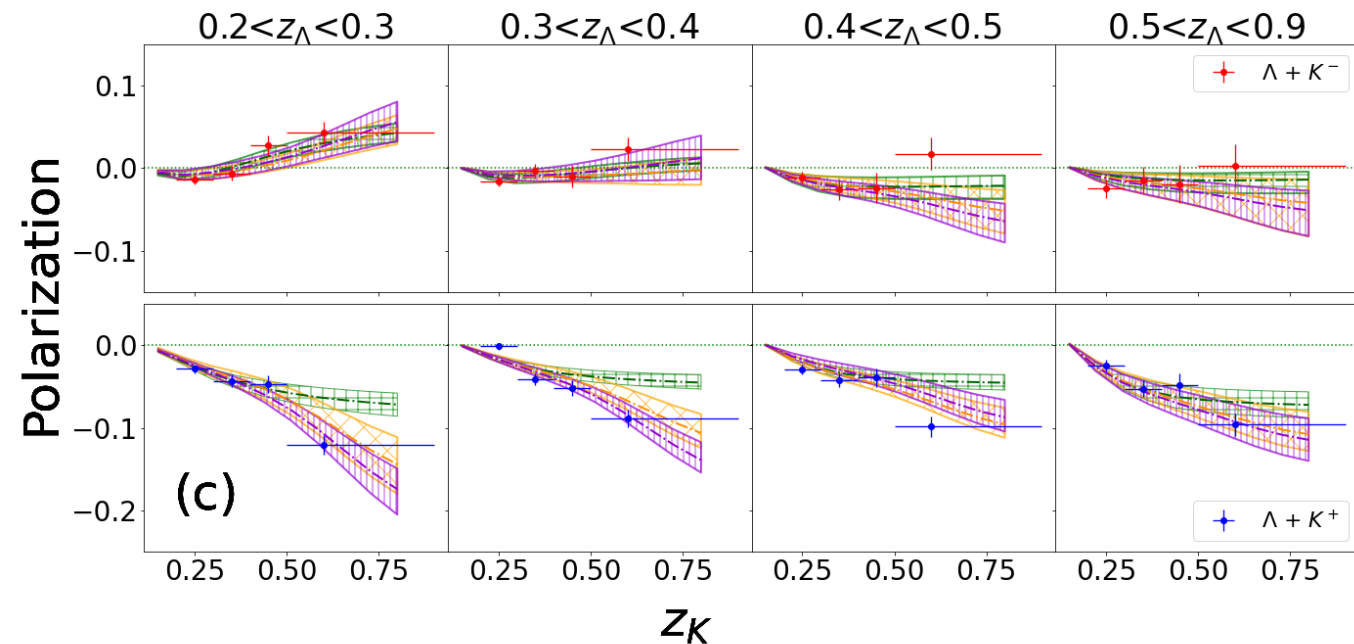
- $up/down$ pFFs are positive;
- $\overline{up}/\overline{down}$ pFFs are negative;
- $strange/\overline{strange}$ pFFs are negative;
- up & $\overline{strange}$ compatible with (1,2) scn.
- The negative sea contribution is larger in size;
- Similar size for the Gaussian width.

Double hadron production in e^+e^- processes



- (1) No Charm, No SU2
- (2) Charm, No SU2
- (3) Charm, SU2

- All scenarios can describe $\Lambda\pi^\pm, \bar{\Lambda}\pi^\pm, \Lambda K^-, \bar{\Lambda}K^+$ data;
- Scenario (1) cannot describe $\Lambda K^+, \bar{\Lambda}K^-$ data with $z_K > 0,5$;
- With the Charm contribution we obtain similar good fits and description



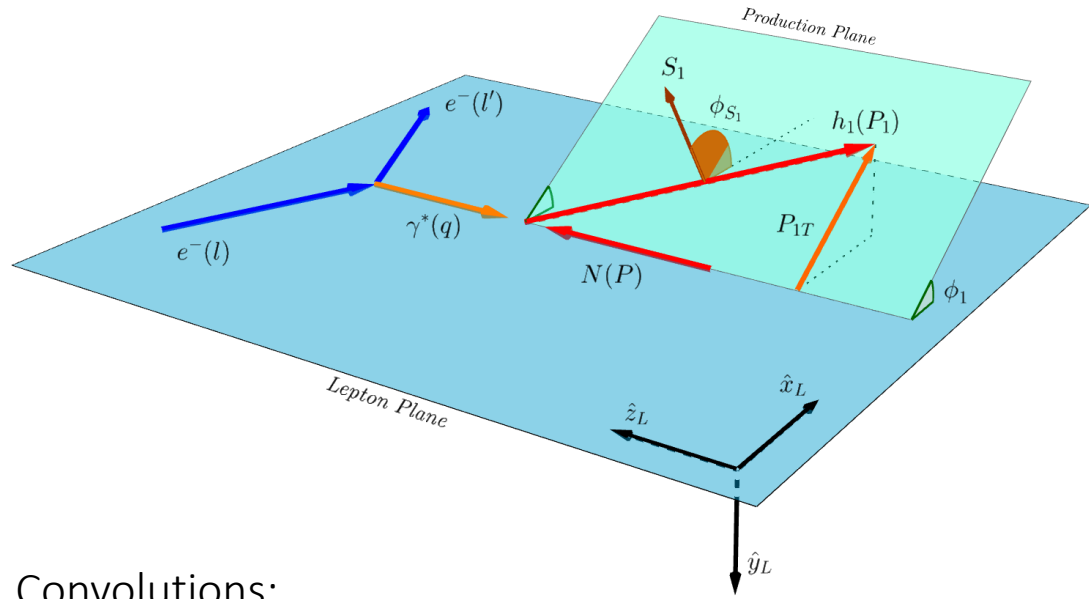
Double hadron production in e^+e^- processes

Some remarks:

- Charm contribution in the unpolarized C.S. is necessary;
Attempts were made to include this contribution also in the polarized c.s.
→ No improvements!
- We cannot distinguish between the (2) and (3) scenarios.
If Normalization factors are free,
up & *down* come out opposite, violating the SU(2) symmetry.
- Investigate the polarization in:
 e^+e^- at different energies → we cannot distinguish between (2) & (3) scenarios

SIDIS

Semi-inclusive Deep Inelastic Scattering



Polarization:

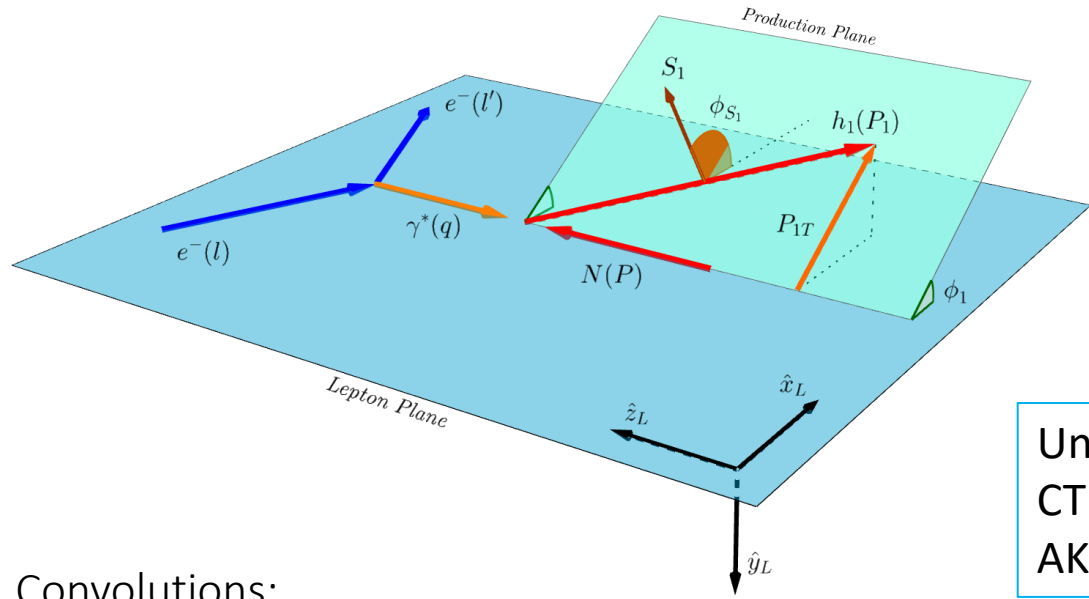
$$P_n^{h_1}(x_B, z_h) = \frac{M_1 \int dq_T q_T d\phi_1 \mathcal{B}_1 \left[\tilde{f}_1 \tilde{D}_{1T}^{\perp(1)} \right]}{\int dq_T q_T d\phi_1 \mathcal{B}_0 \left[\tilde{f}_1 \tilde{D}_1 \right]}$$

Convolutions:

$$\mathcal{B}_0 \left[\tilde{f}_1 \tilde{D}_1 \right] = \frac{1}{z^2} \sum_q e_q^2 \int \frac{db_T}{(2\pi)} b_T J_0(b_T q_T) f_{N/q}(x; \bar{\mu}_b) d_{q/h}(z; \bar{\mu}_b) \\ \times M_{f_1}(b_c(b_T), x) M_{D_h}(b_c(b_T), z) e^{-g_K(b_c(b_T); b_{\max}) \ln \left(\frac{Q^2 z}{x M_P M_h} \right) - S_{\text{pert}}(b_*; \bar{\mu}_b)}$$

$$\mathcal{B}_1 \left[\tilde{f}_1 \tilde{D}_{1T}^{\perp(1)} \right] = \frac{1}{z^2} \sum_q e_q^2 \int \frac{db_T}{(2\pi)} b_T^2 J_1(b_T q_T) f_{N/q}(x; \bar{\mu}_b) D_{1T,q}^{\perp(1)}(z; \bar{\mu}_b) \\ \times M_{f_1}(b_c(b_T), x) M_{D_1}^{\perp}(b_c(b_T), z) e^{-g_K(b_c(b_T); b_{\max}) \ln \left(\frac{Q^2 z}{x M_P M_h} \right) - S_{\text{pert}}(b_*; \bar{\mu}_b)}$$

Semi-inclusive Deep Inelastic Scattering



Polarization:

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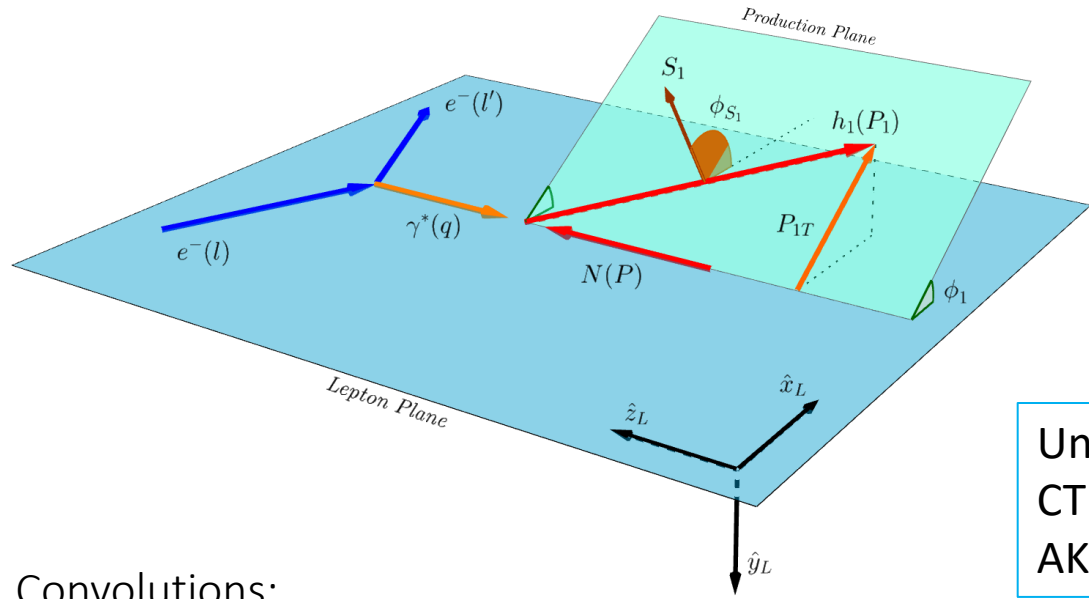
Unpolarized functions:
CT14nnlo set for proton
AKK set for Λ

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Semi-inclusive Deep Inelastic Scattering



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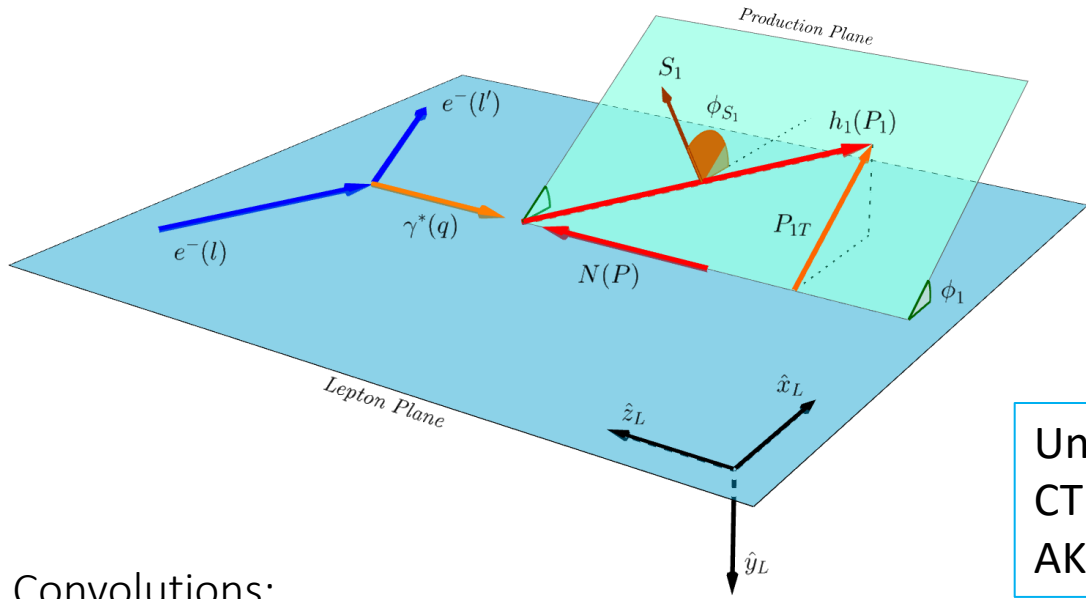
Non-perturbative functions from
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Semi-inclusive Deep Inelastic Scattering



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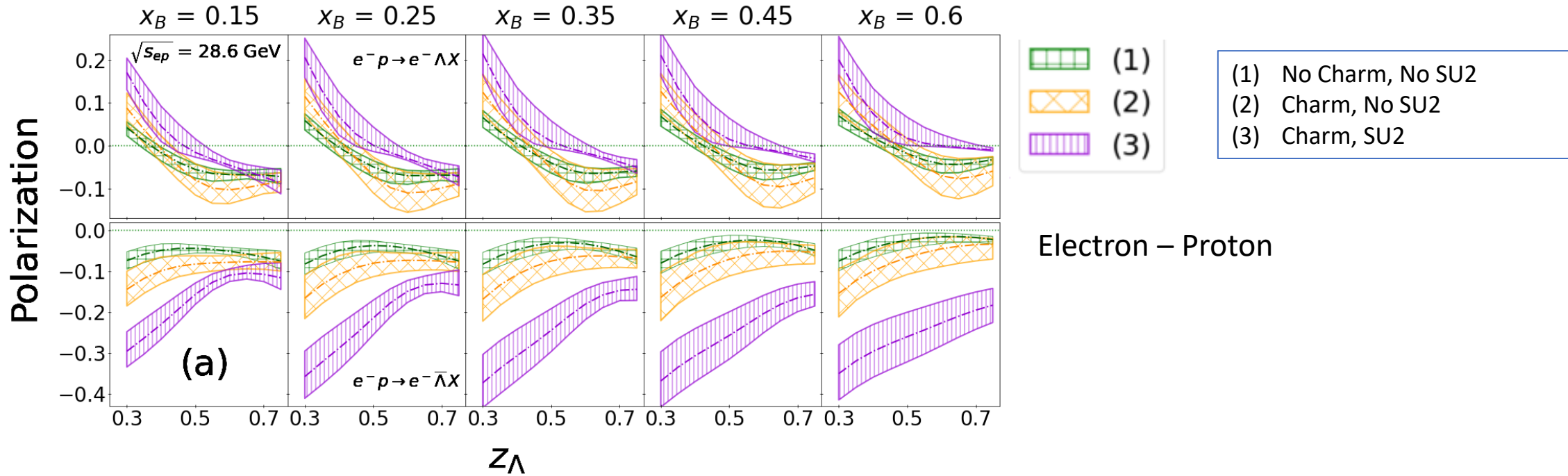
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Predictions are given at different energies:

E_N (GeV)	E_{e^-} (GeV)	$\sqrt{s_{eN}}$ (GeV)
41	5	28.6
100	10	63.2

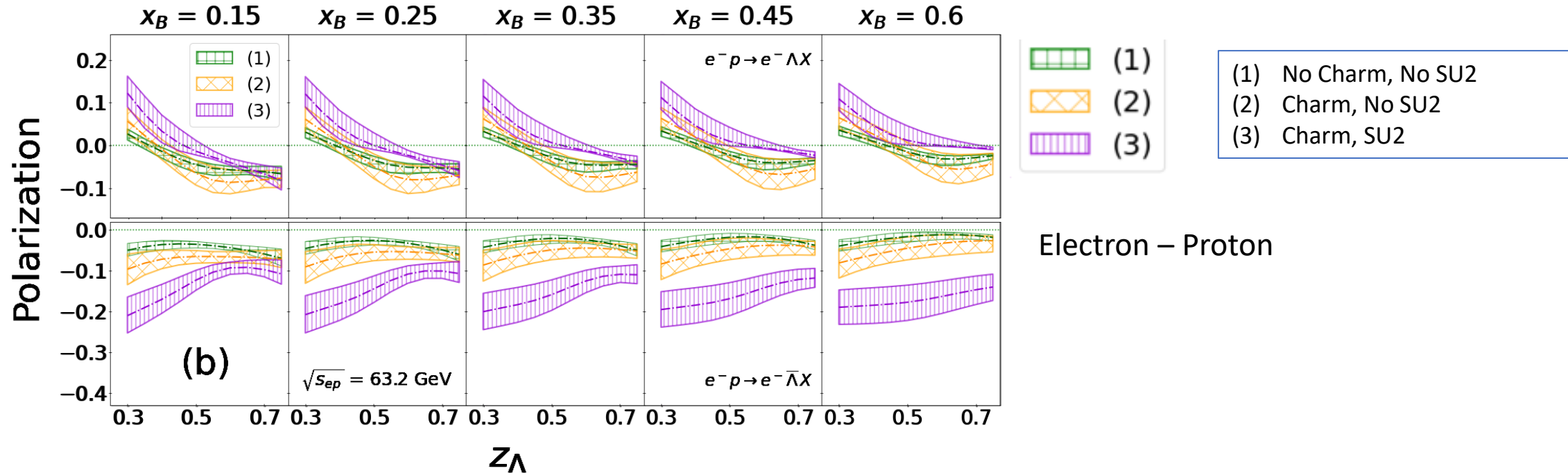
$$s = 4E_N E_e, \quad Q^2 = x_B y s$$

Semi-inclusive Deep Inelastic Scattering



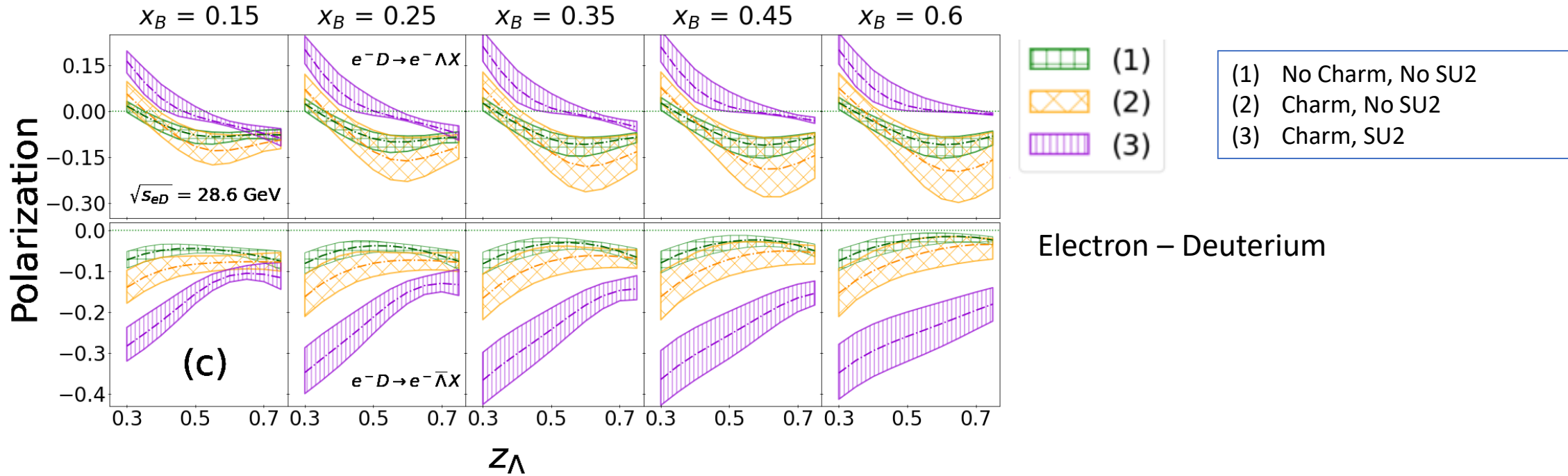
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Semi-inclusive Deep Inelastic Scattering



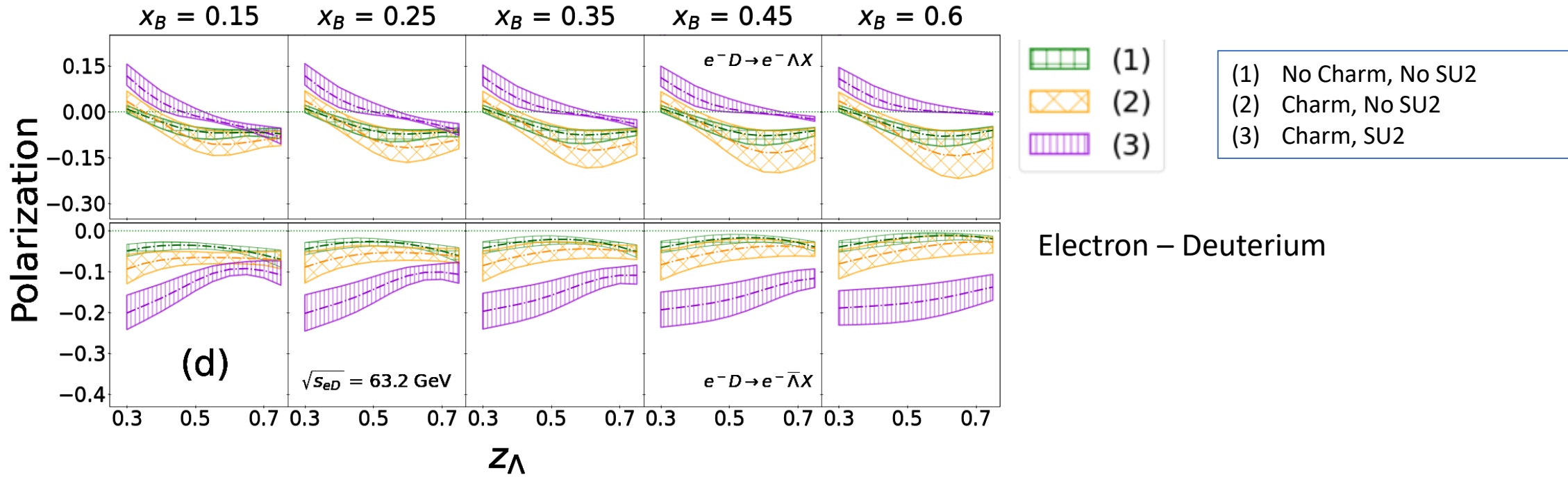
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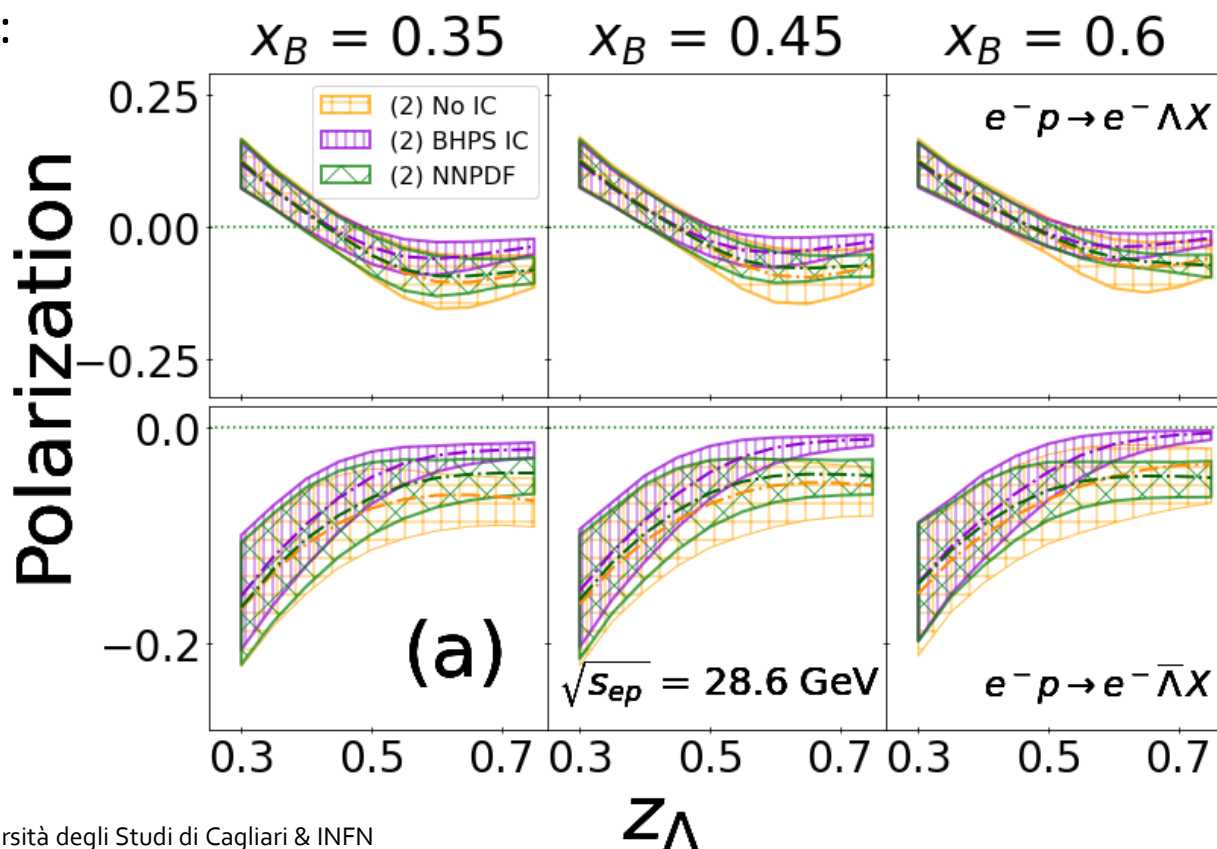
Semi-inclusive Deep Inelastic Scattering

The charm contribution in the fragmentation process is relevant

Intrinsic Charm (IC) component in the proton:

- CT14nnloIC set with BHPS model [T.-J. Hou et al., *JHEP* 02 (2018) 059]
- NNPDF4.0nnlo set [NNPDF Coll., *Eur.Phys.J.C* 82 (2022) 5, 428]

(2) Scenario:

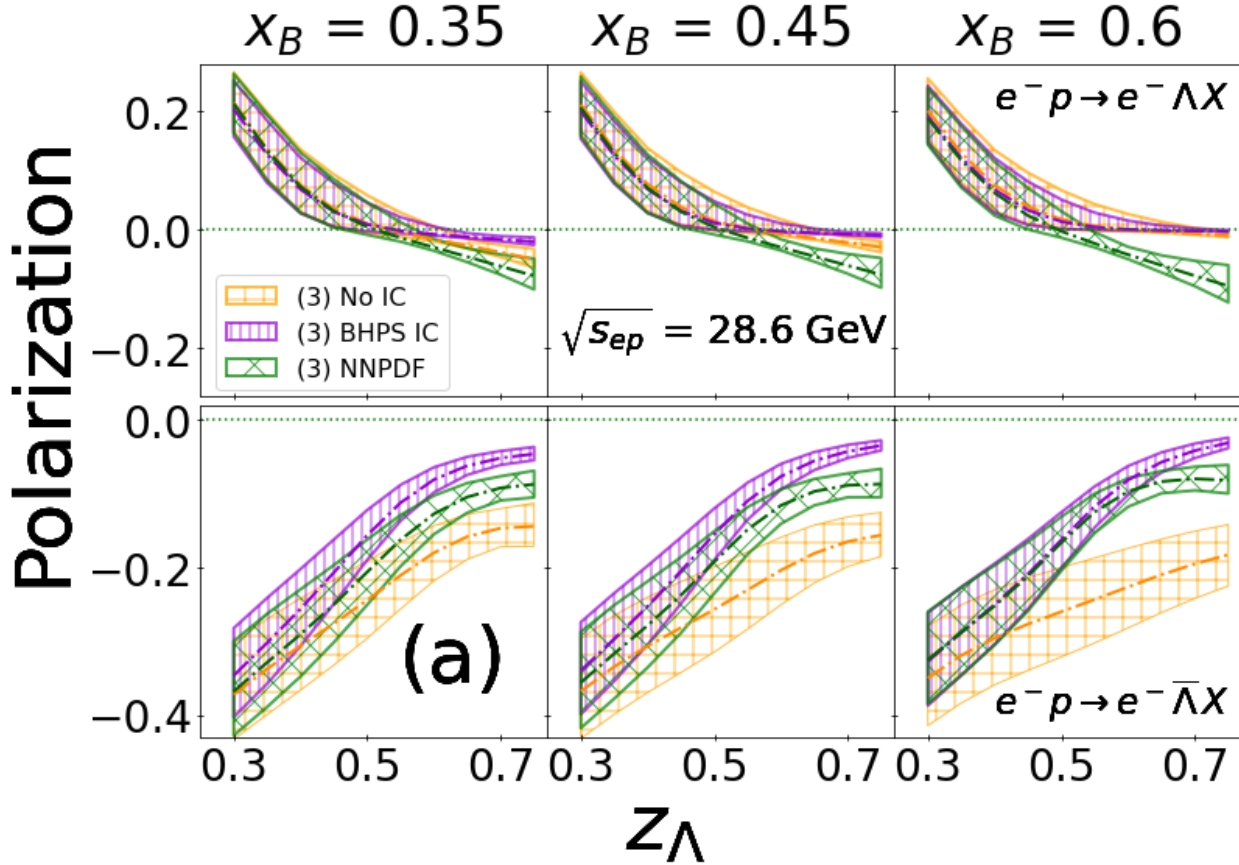


(2) Charm, No SU2

- BHPS and NNPDF: similar polarization of previous predictions
- Same behavior is present for greater values of the c.m. energy.

Semi-inclusive Deep Inelastic Scattering

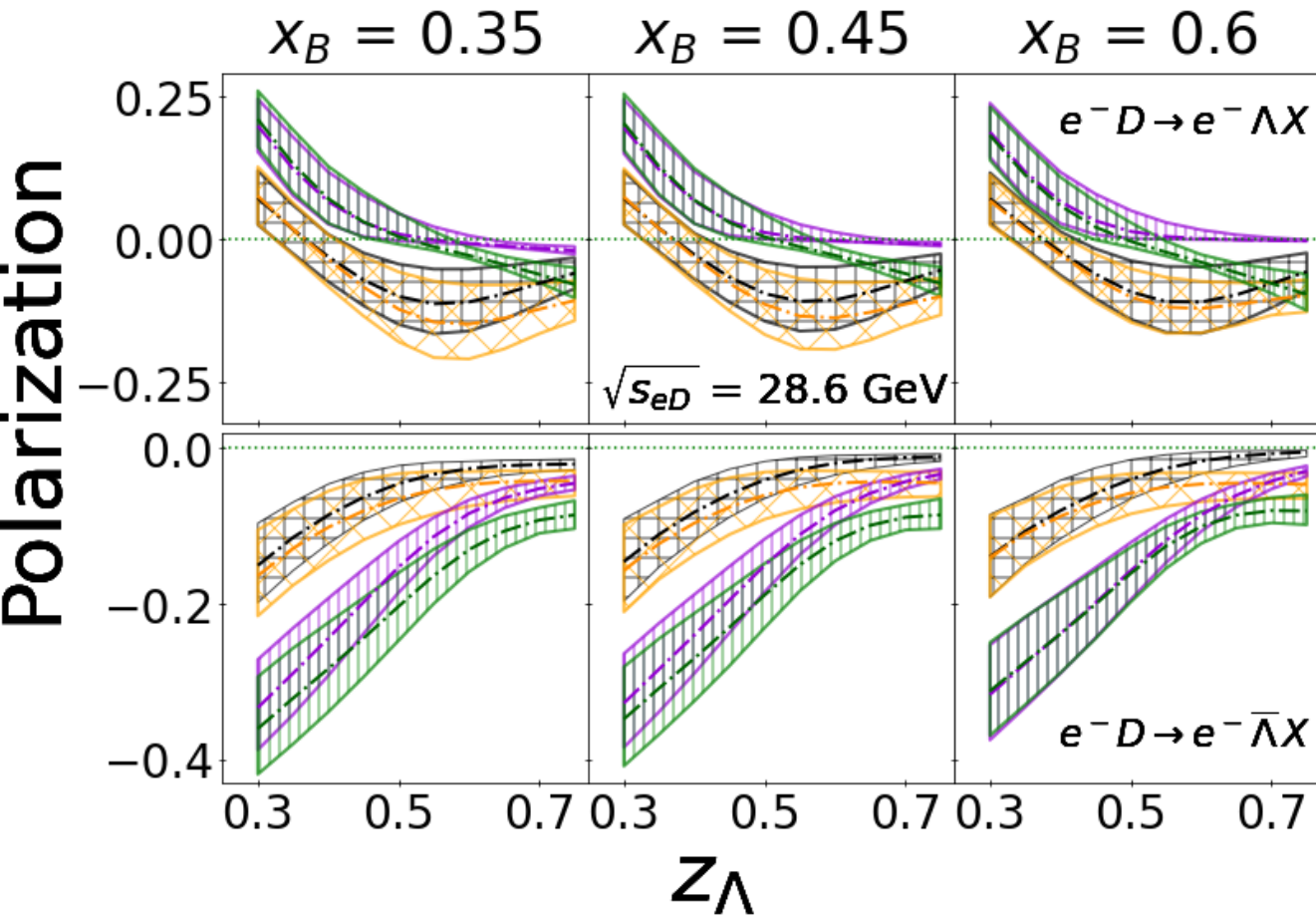
(3) Scenario:



(3) Charm, SU2

- Estimates vary significantly as x_B increases;
- $\bar{\Lambda}$ estimates with BHPS and NNPDF different from the previous ones;
- Λ :decreases to zero
- Λ : NNPDF become negative

Semi-inclusive Deep Inelastic Scattering



- (2) BHPS IC
- (2) NNPDF
- (3) BHPS IC
- (3) NNPDF

- (2) Charm, No SU2
- (3) Charm, SU2

- Λ : the predictions are compatible
- $\bar{\Lambda}$: within (2) and (3) are still different

Conclusions

- Double hadron production in e^+e^- :
Fit results

Charm is necessary

(2) and (3) scenarios cannot be distinguished → open issue!

- SIDIS:
predictions for the transverse Lambda polarization

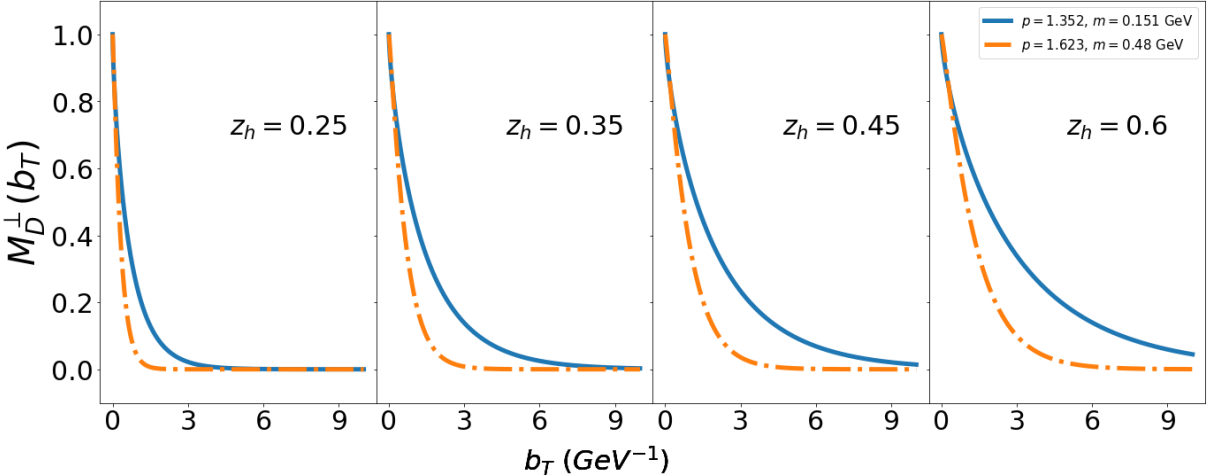
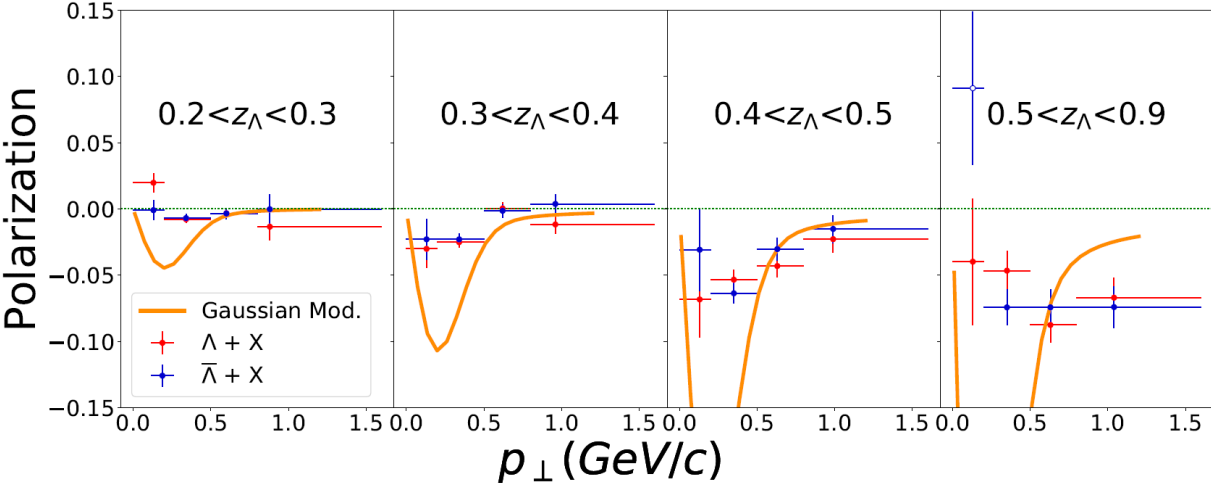
(2) and (3) scenarios predictions are different and can be distinguished



Motivations and Contents

2-h and 1-h data cannot be described by the same non-perturbative model!

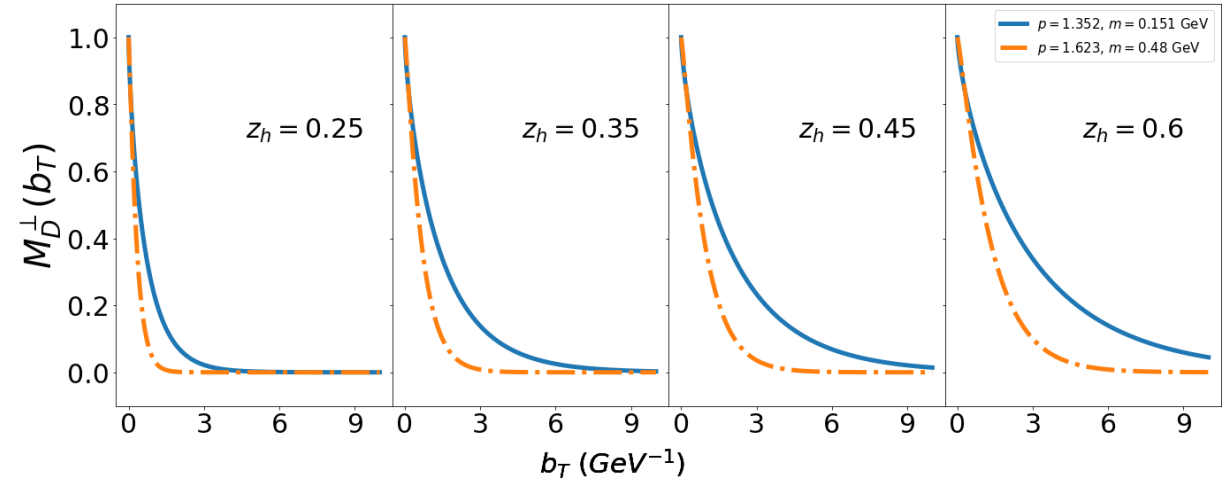
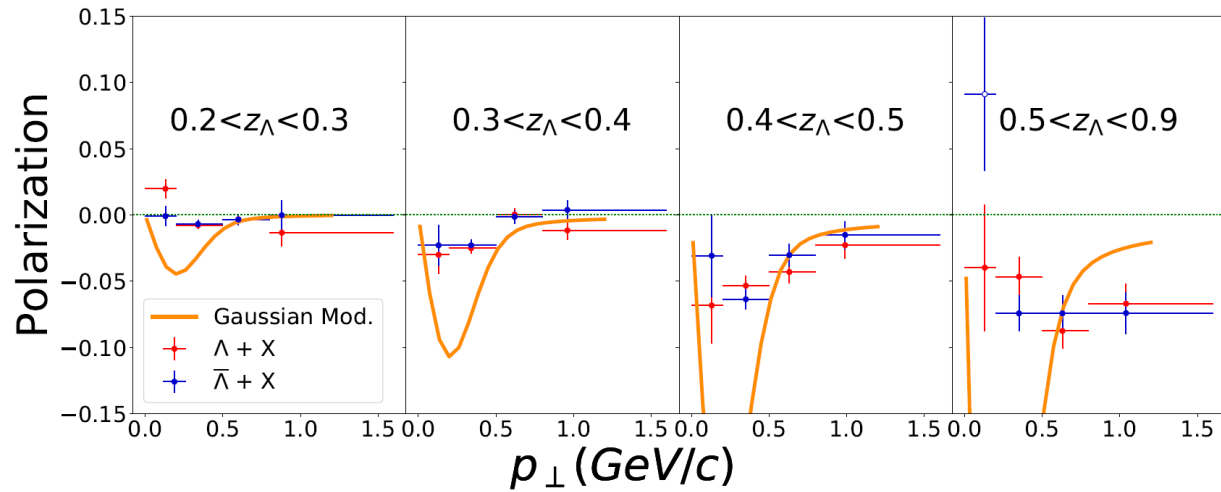
- 2-h Power-Law model
- 1-h Power-Law model



Motivations and Contents

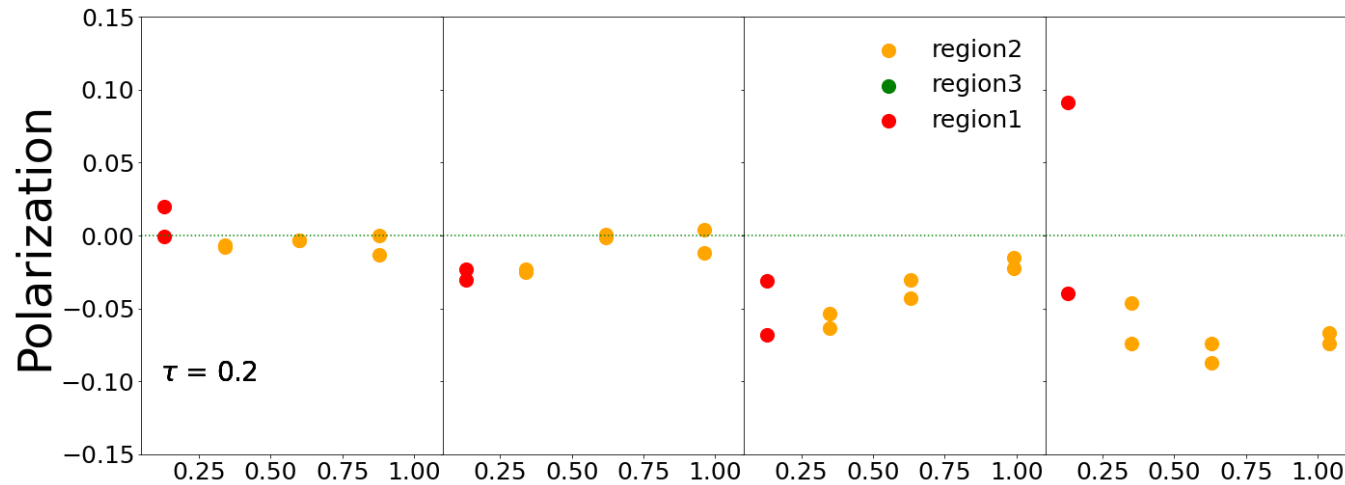
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Data can be analyzed within a different scheme.

See A. Simonelli's talk!

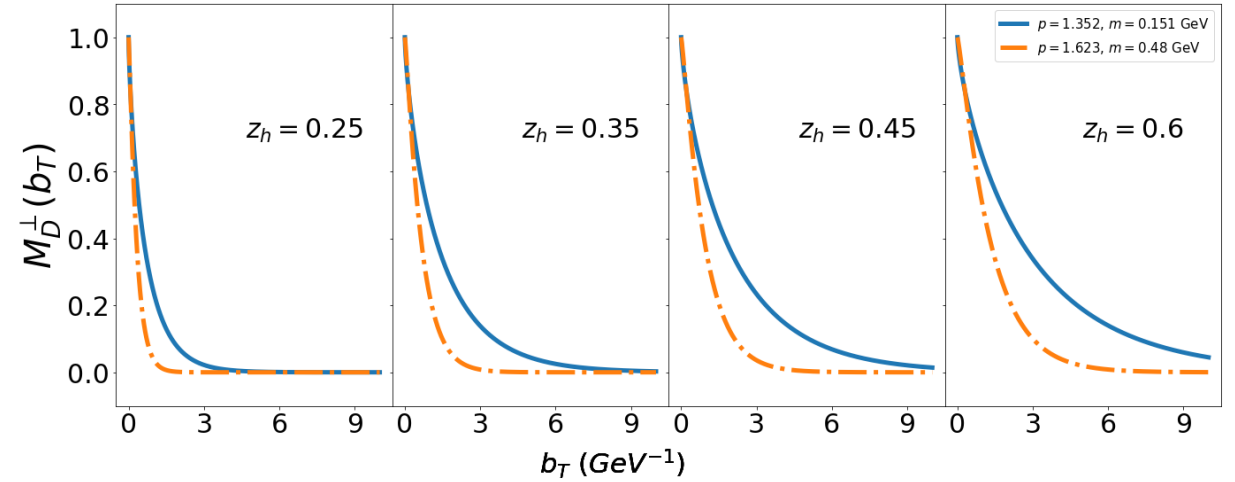
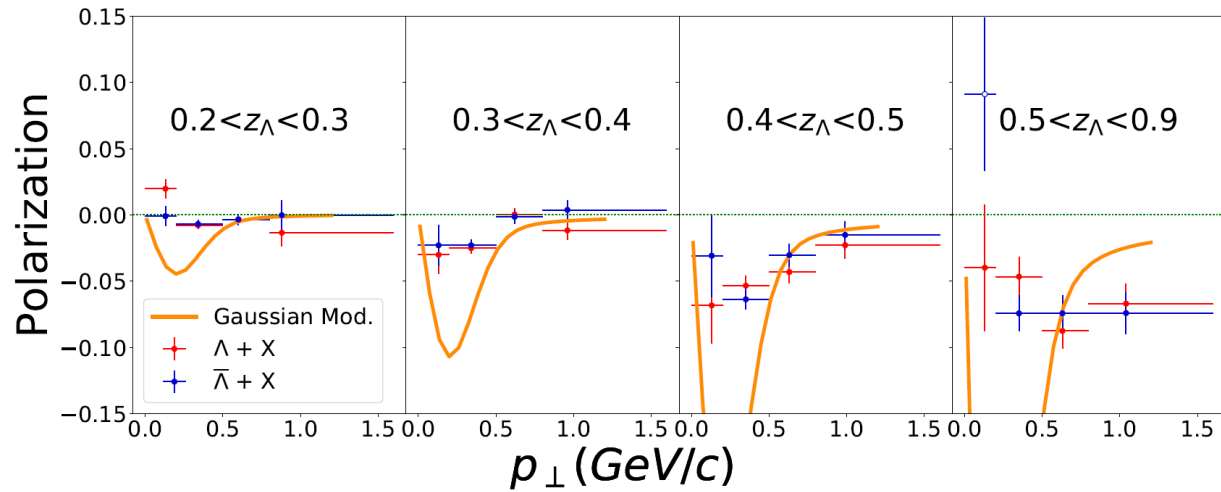


Data not binned in Thrust!

Motivations and Contents

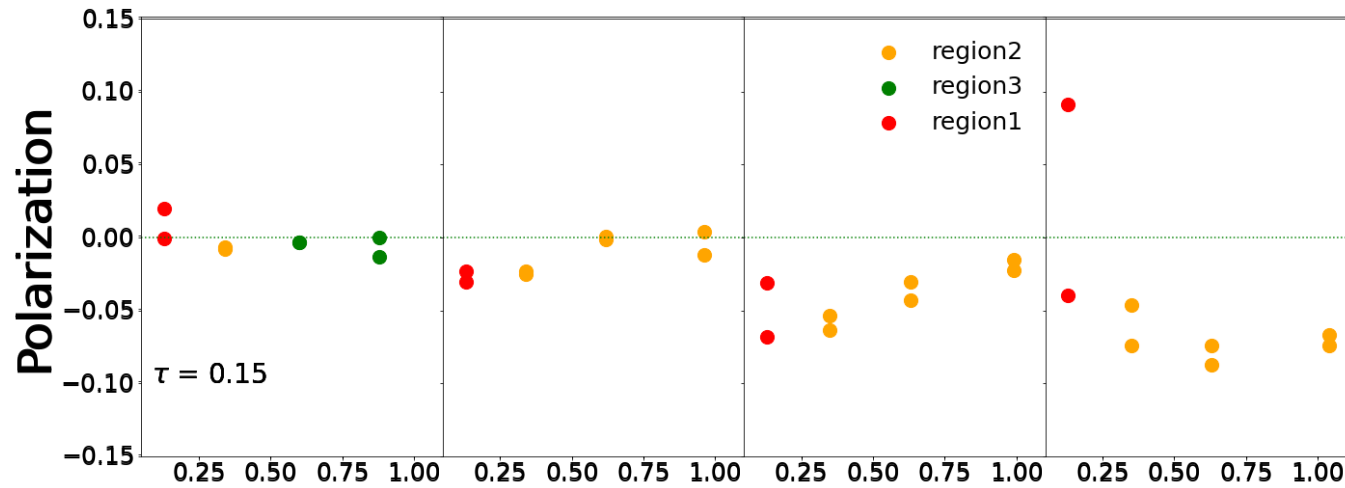
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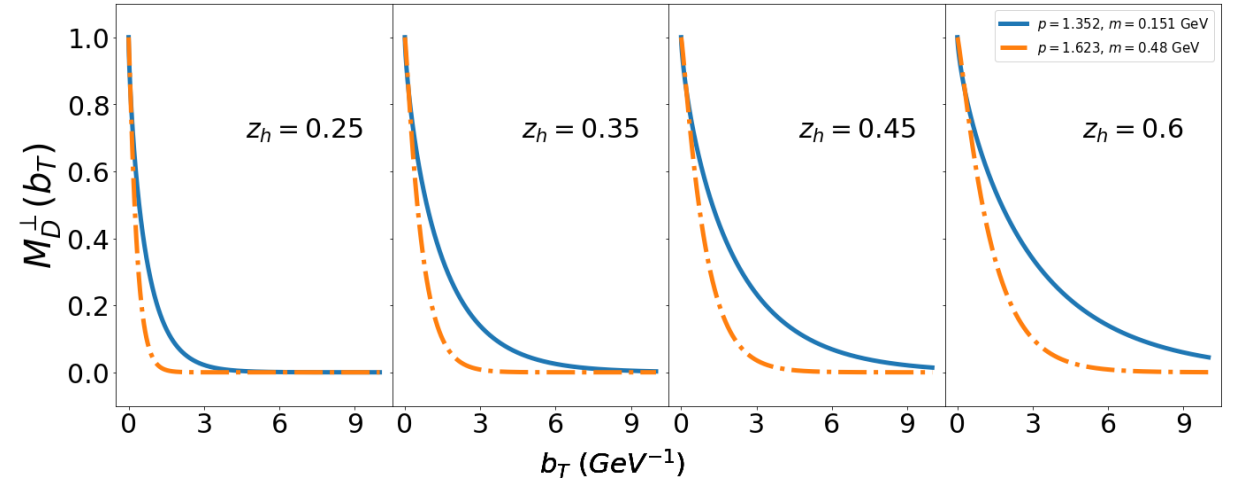
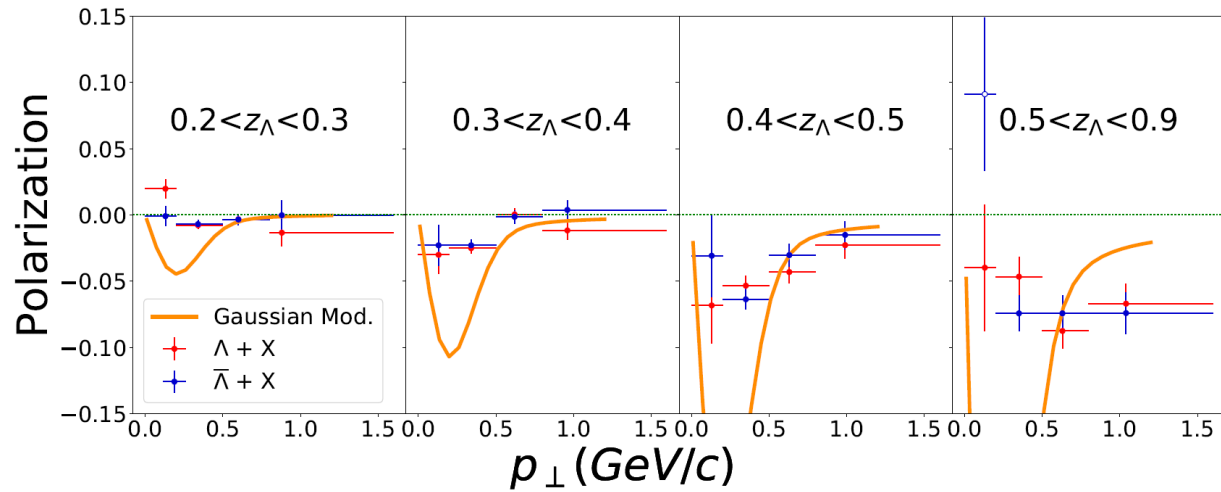


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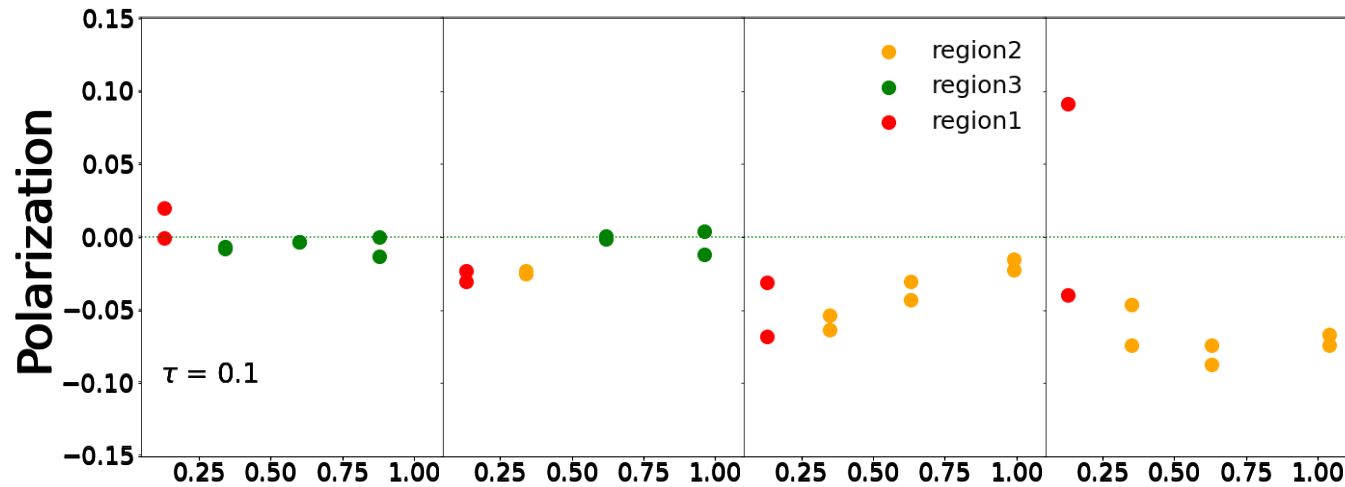
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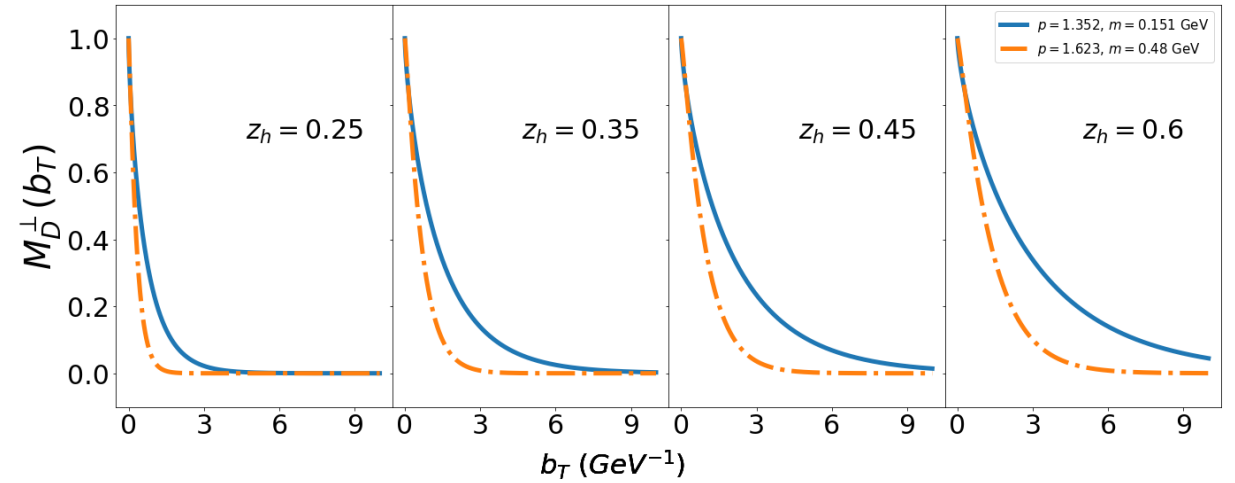
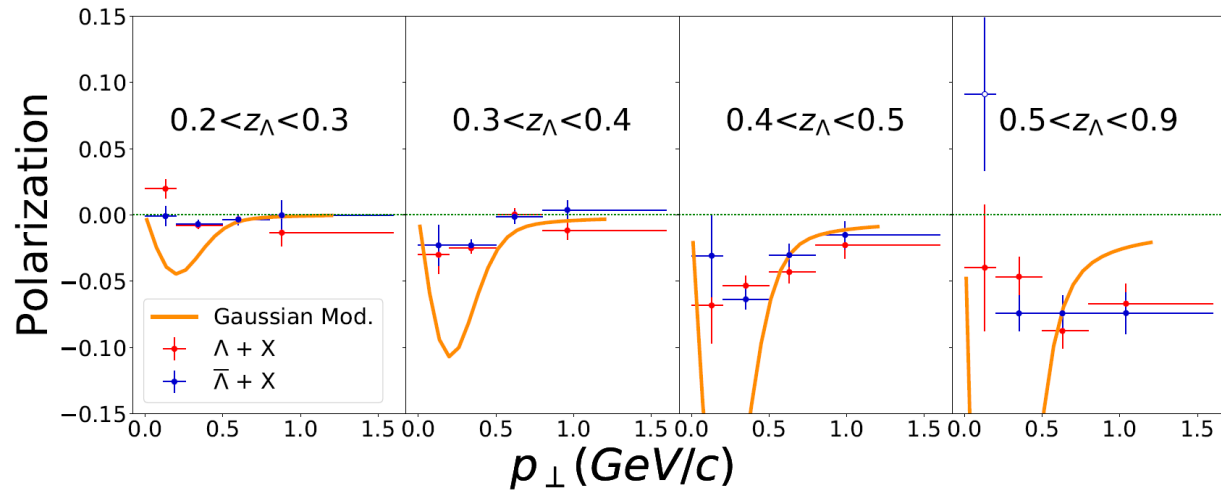


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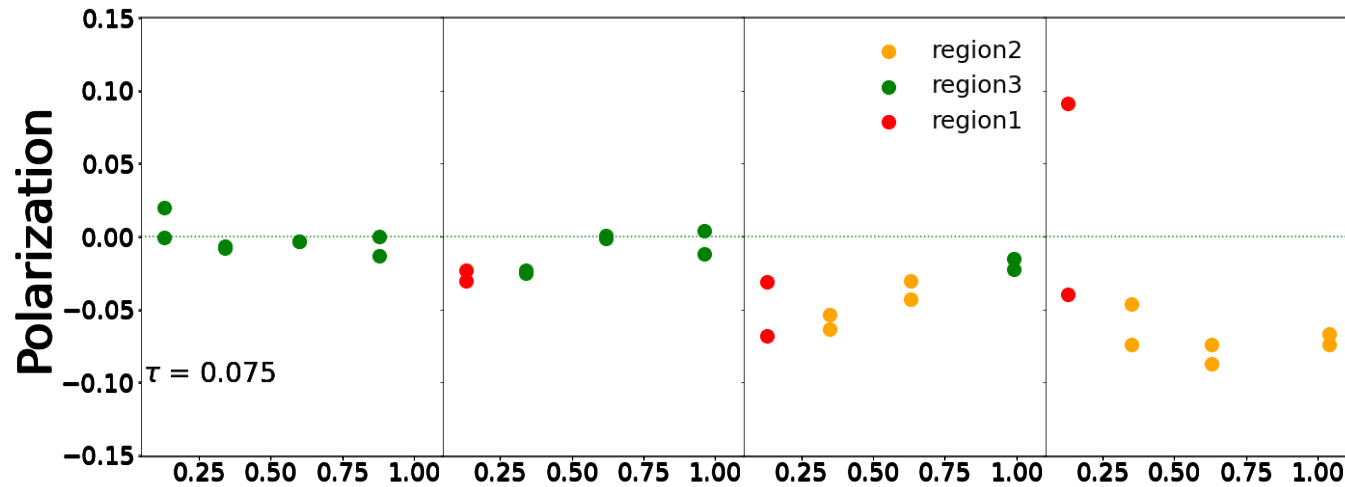
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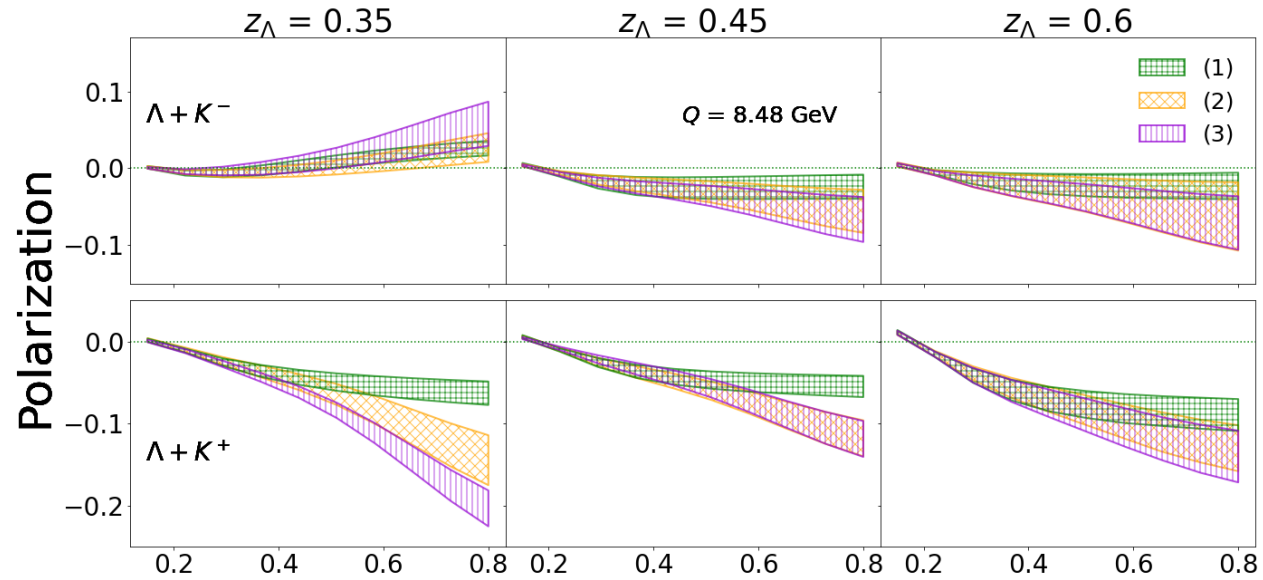
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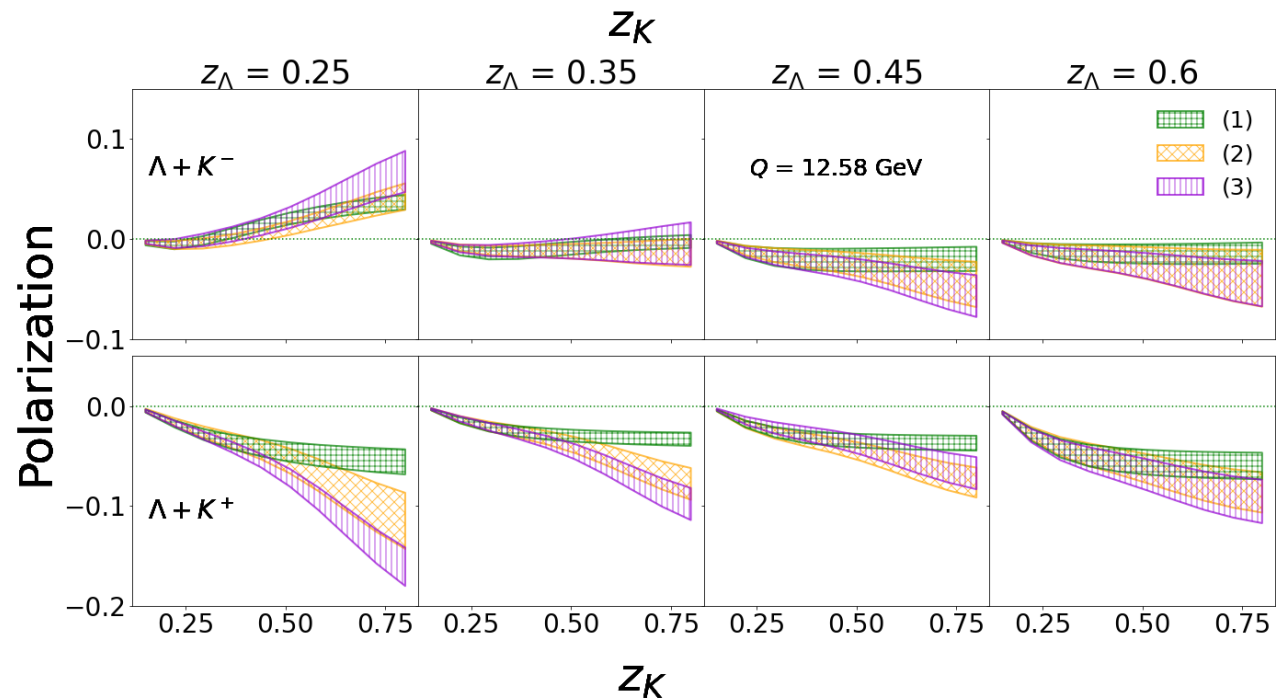
Double hadron production in e^+e^- processes



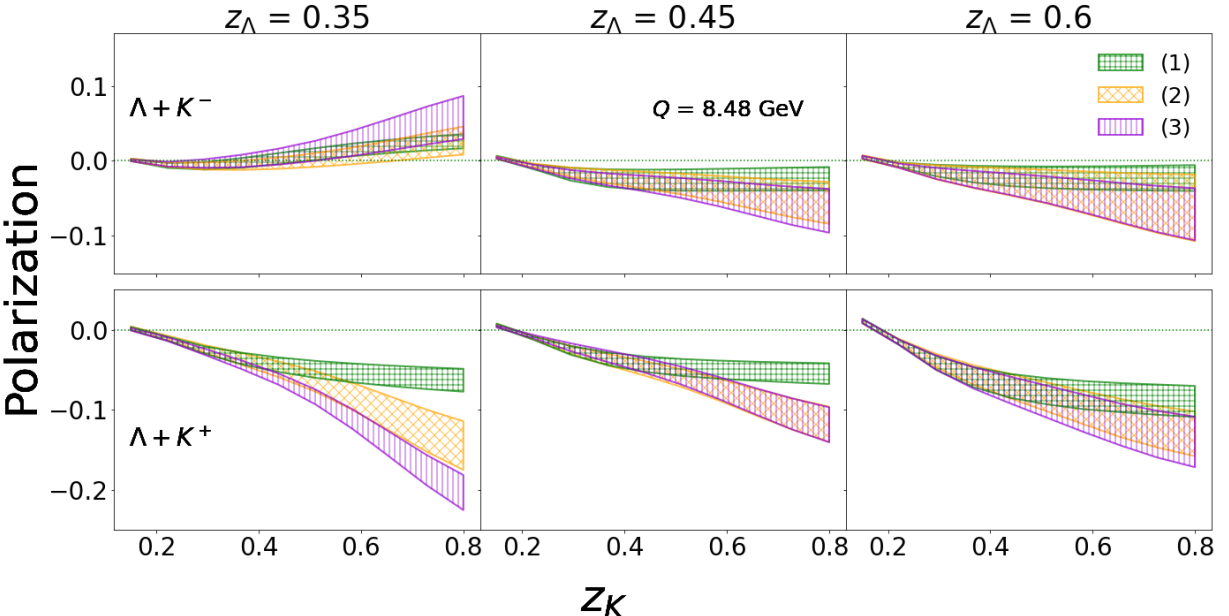
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Predictions at different energies:

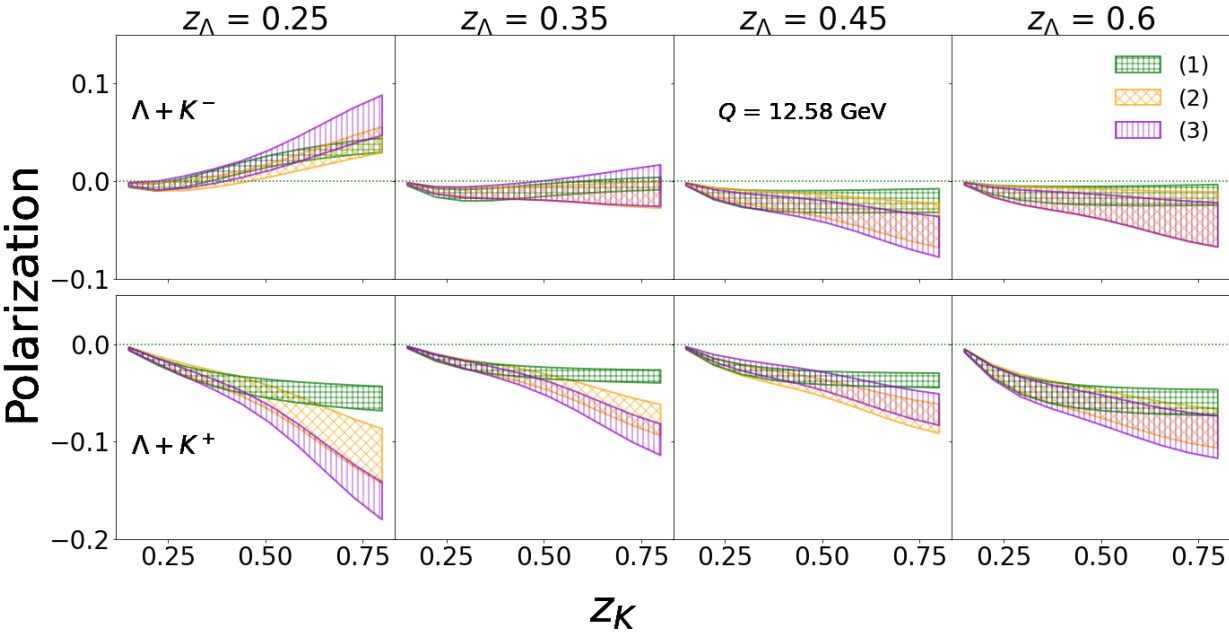
- All scenarios lead to similar predictions for $\Lambda\pi^\pm, \bar{\Lambda}\pi^\pm, \Lambda K^-, \bar{\Lambda}K^+$;
- Scenario (1): different polarization for $\Lambda K^+, \bar{\Lambda}K^-$ with $z_K > 0,5$;
- Scenarios (2) & (3) bands are almost totally overlapped: they cannot be distinguished!



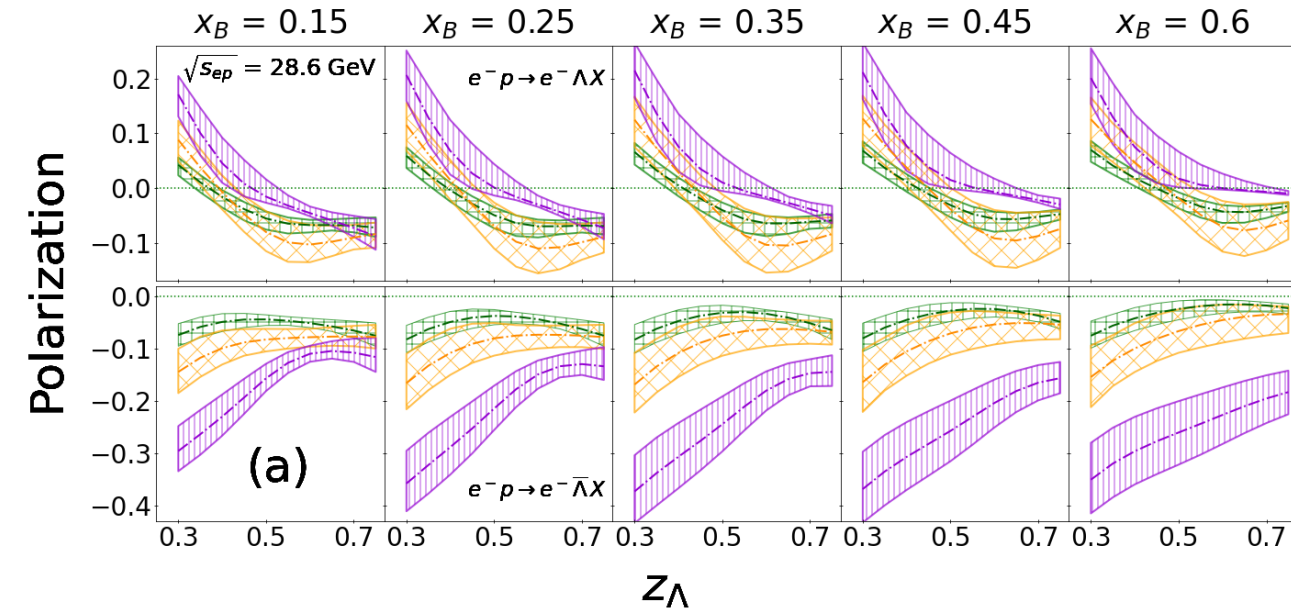
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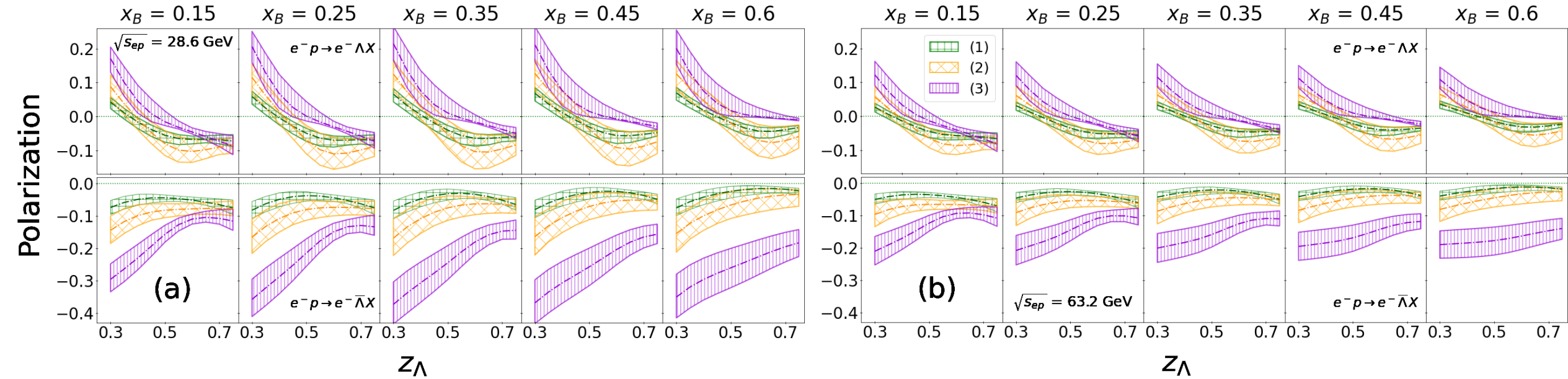


Semi-inclusive Deep Inelastic Scattering



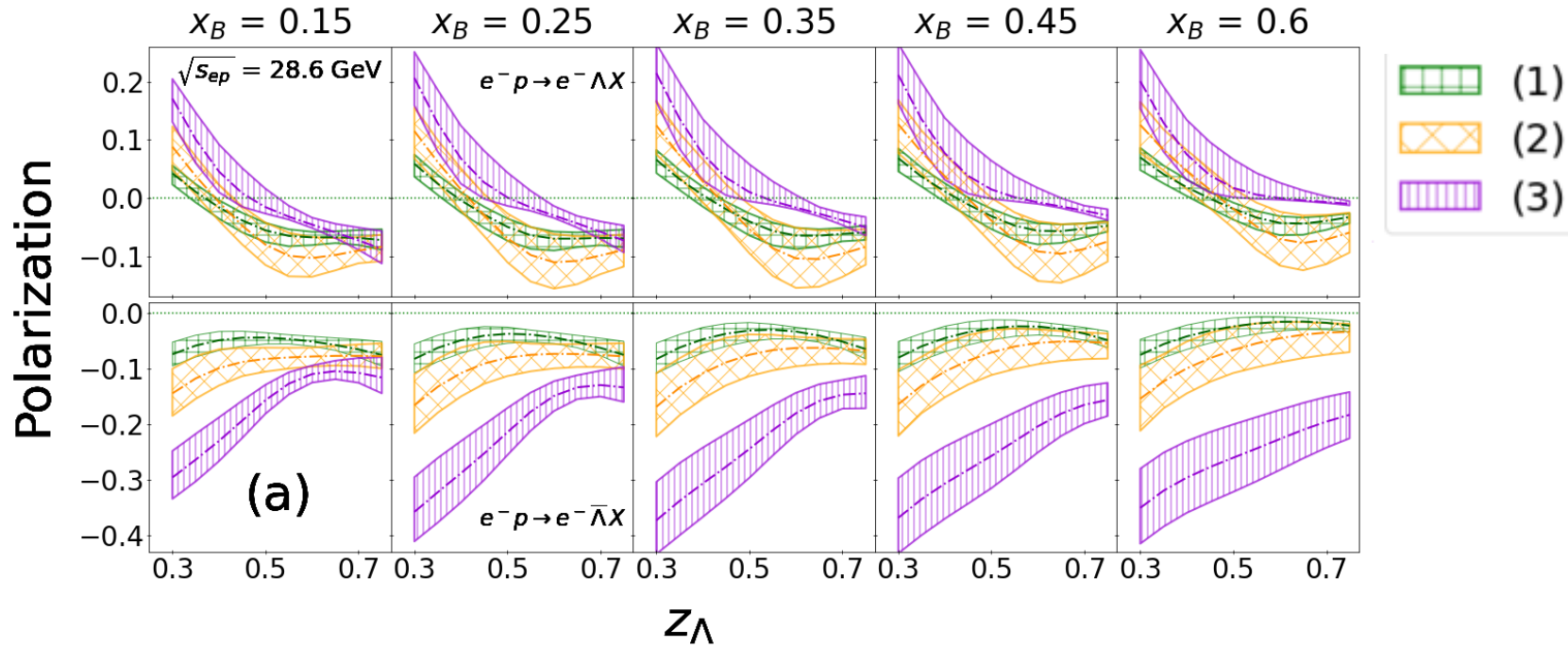
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Semi-inclusive Deep Inelastic Scattering

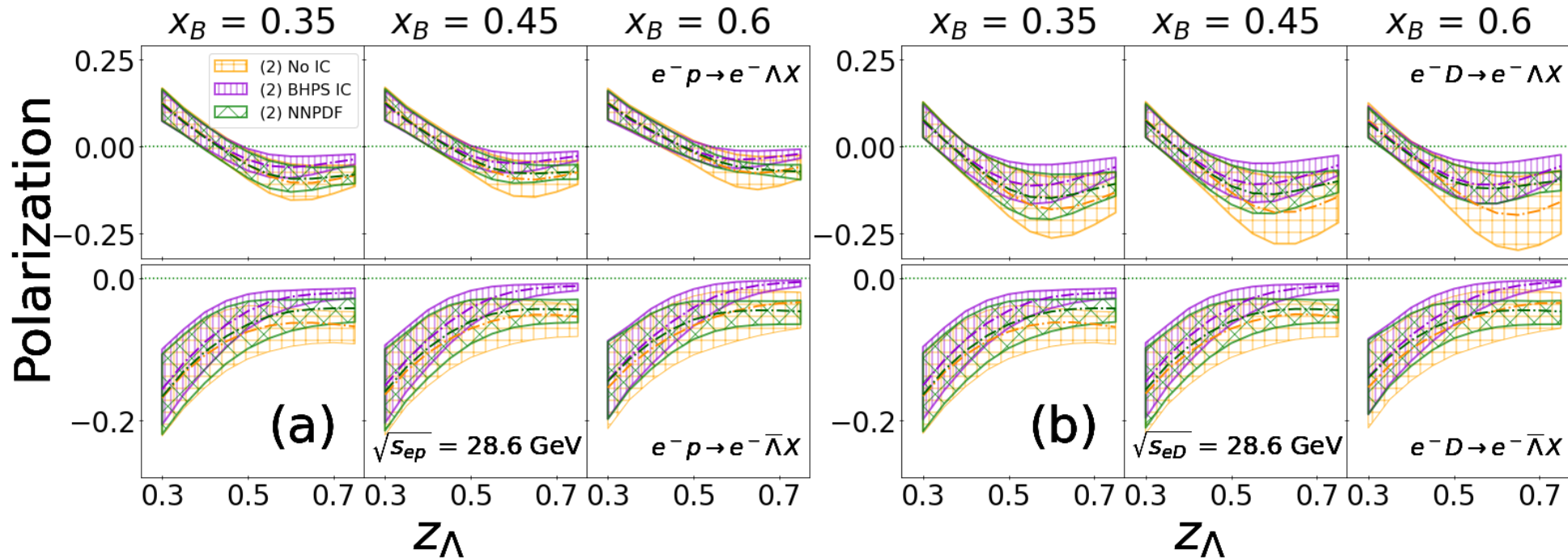
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- CT14nnloIC set with BHPS model [T.-J. Hou et al., *JHEP* 02 (2018) 059]
- NNPDF4.0nnlo set [NNPDF Coll., *Eur.Phys.J.C* 82 (2022) 5, 428]

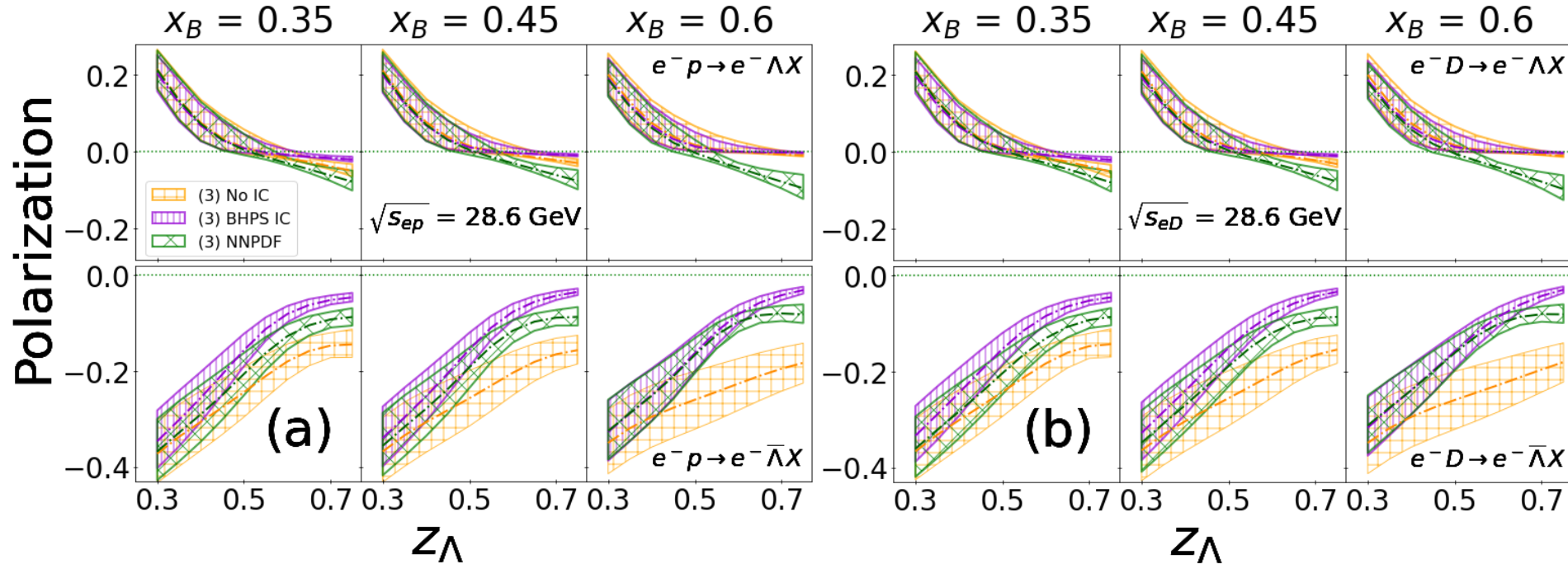
- BHPS and NNPDF: similar polarization of previous predictions
- Same behavior is present for greater values of the c.m. energy.

(2) Scenario:



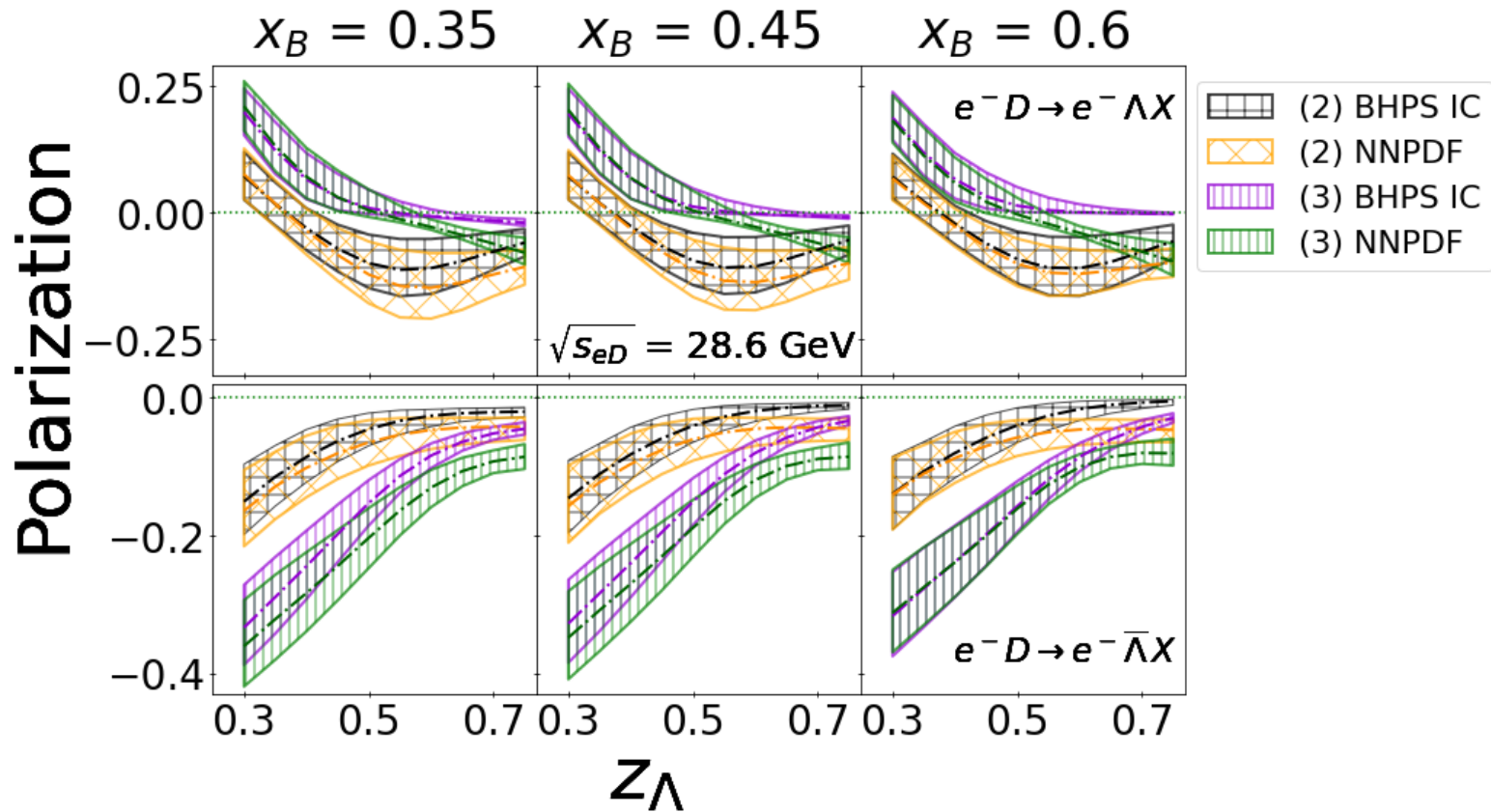
Semi-inclusive Deep Inelastic Scattering

(3) Scenario:



- Estimates vary significantly as x_B increases;
- $\bar{\Lambda}$ estimates with BHPS and NNPDF different from the previous ones;
- Λ :decreases to zero
- Λ : NNPDF become negative

Semi-inclusive Deep Inelastic Scattering



- Λ : the predictions are compatible
- $\bar{\Lambda}$: within (2) and (3) are still different