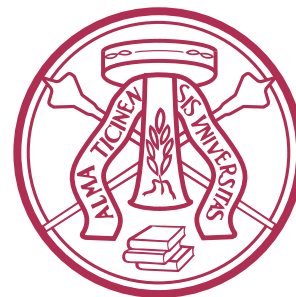




Istituto Nazionale di Fisica Nucleare



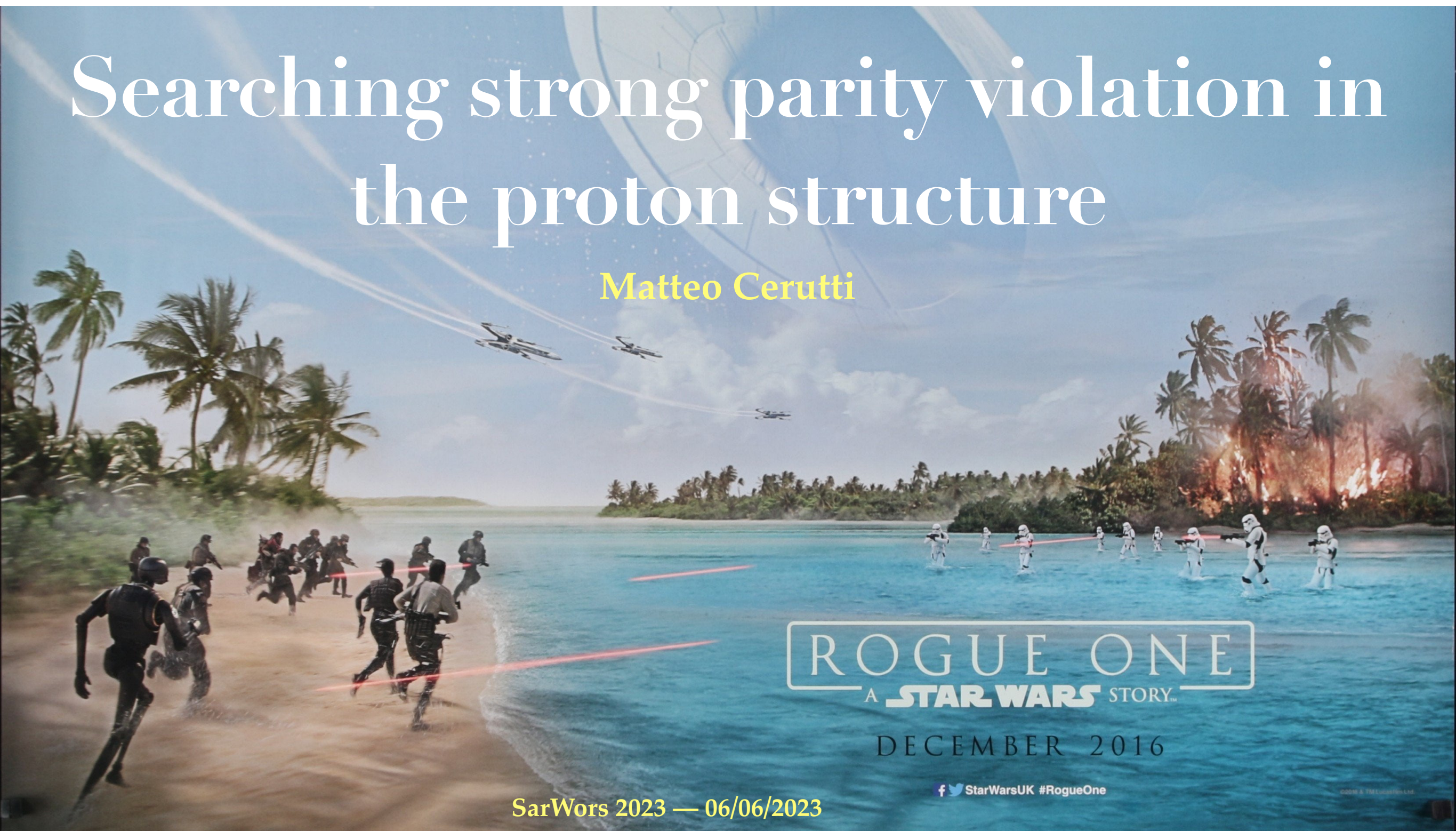
**HAS QCD**  
HADRONIC STRUCTURE AND  
QUANTUM CHROMODYNAMICS



**UNIVERSITÀ  
DI PAVIA**

# Searching strong parity violation in the proton structure

Matteo Cerutti



**ROGUE ONE**  
A **STAR WARS** STORY™

DECEMBER 2016

f StarWarsUK #RogueOne

SarWors 2023 — 06/06/2023

©2016 & TM Lucasfilm Ltd.





Istituto Nazionale di Fisica Nucleare



**HAS QCD**  
HADRONIC STRUCTURE AND  
QUANTUM CHROMODYNAMICS



**UNIVERSITÀ  
DI PAVIA**

# Searching strong parity violation in the proton structure

Matteo Cerutti







Istituto Nazionale di Fisica Nucleare



**HAS QCD**  
HADRONIC STRUCTURE AND  
QUANTUM CHROMODYNAMICS



**UNIVERSITÀ  
DI PAVIA**

# Searching strong parity violation in the proton structure

Matteo Cerutti

**SAR WORS**



**2023**

3<sup>rd</sup> SARDINIAN WORKSHOP ON SPIN

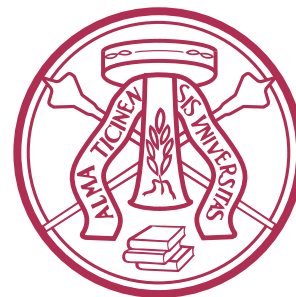




Istituto Nazionale di Fisica Nucleare



**HAS QCD**  
HADRONIC STRUCTURE AND  
QUANTUM CHROMODYNAMICS



**UNIVERSITÀ  
DI PAVIA**

# Searching strong parity violation in the proton structure

**Matteo Cerutti**

in collaboration with A. Bacchetta, L. Manna,  
M. Radici and X. Zheng

**SAR WORS**



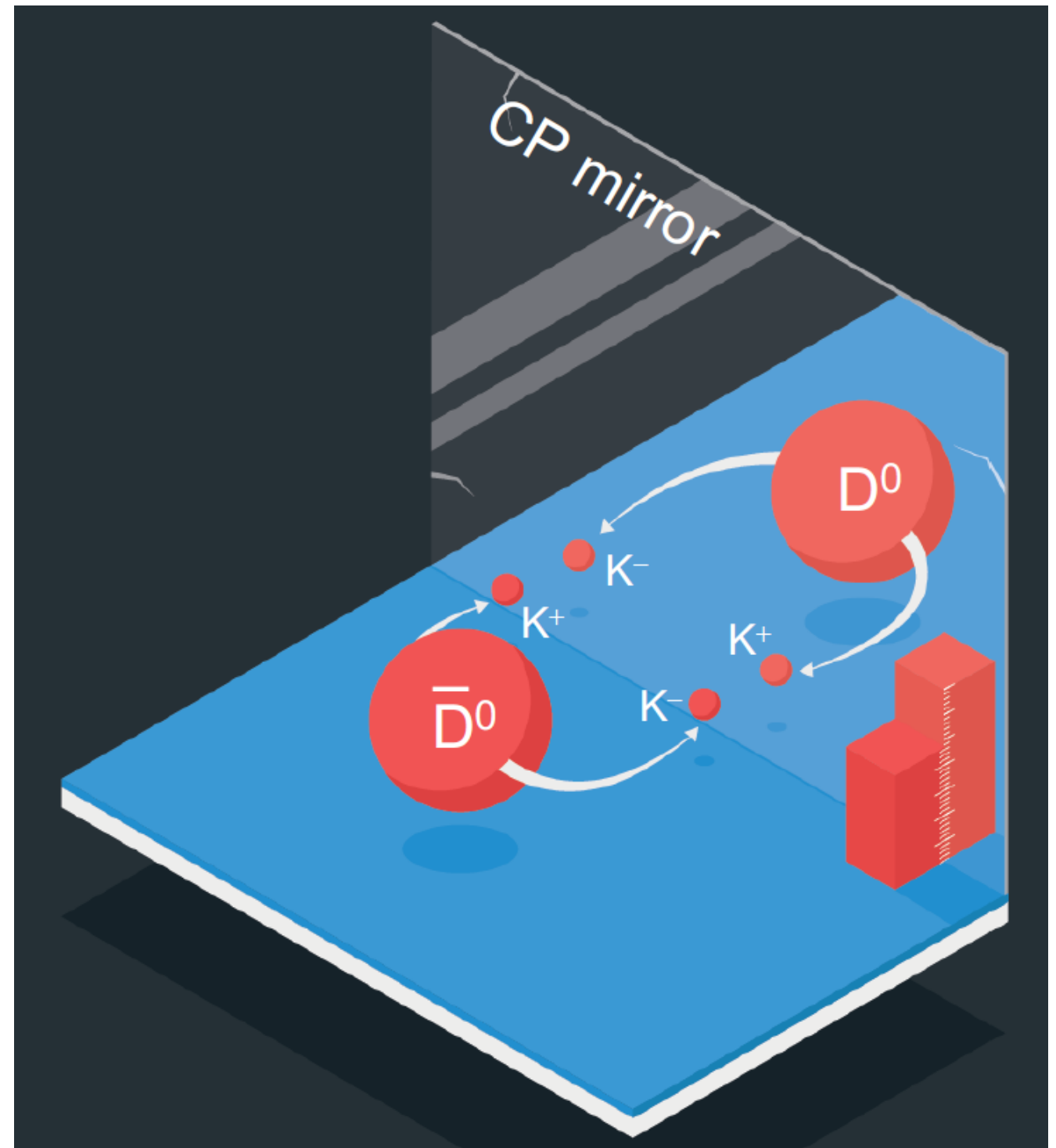
**2023**

3<sup>rd</sup> SARDINIAN WORKSHOP ON SPIN



# Motivations

Investigation of the  
“Strong CP problem”



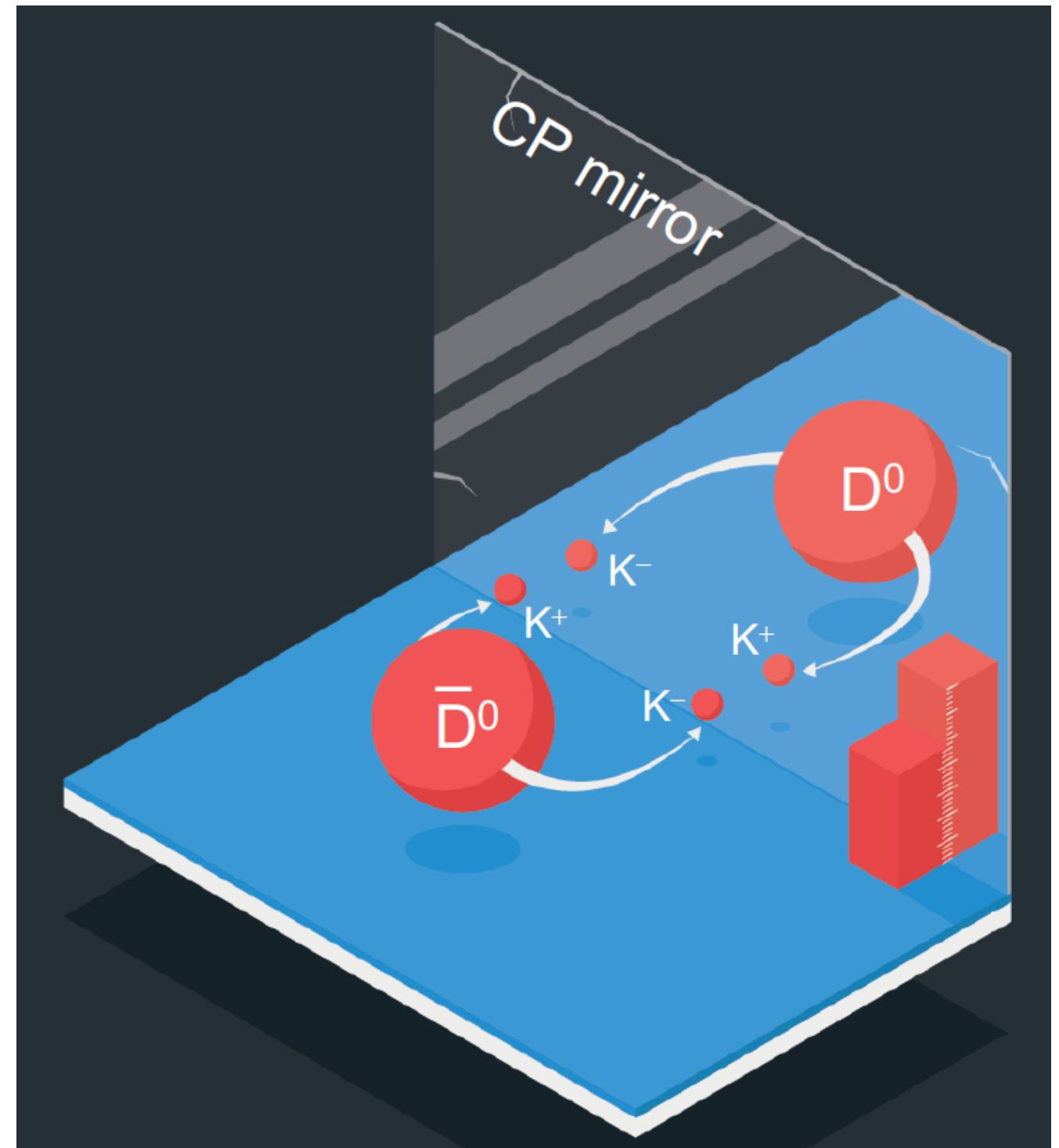


# Motivations

Investigation of the  
“Strong CP problem”



Matter-Antimatter  
imbalance





# Motivations

EW sector

CP violation is included



# Motivations

EW sector

Weak CP

CP violation is included





# Motivations

EW sector

Weak CP

CP violation is included

*too small...*





# Motivations

EW sector

Weak CP

QCD sector

CP violation is included

*too small...*





# Motivations

EW sector

Weak CP

CP violation is included

*too small...*

QCD sector

Strong CP





# Motivations

EW sector

Weak CP

CP violation is included

*too small...*

QCD sector

Strong CP

$$\mathcal{L}'_{\text{QCD}} = \mathcal{L}_{\text{QCD}} + \mathcal{L}^{\text{CP}}$$





# Motivations

## EW sector

Weak CP

CP violation is included

*too small...*

## QCD sector

Strong CP

$$\mathcal{L}'_{\text{QCD}} = \mathcal{L}_{\text{QCD}} + \mathcal{L}^{\text{CP}}$$

$\theta$ -term

SMEFT operators



# Motivations

## EW sector

Weak CP

CP violation is included

*too small...*

## QCD sector

Strong CP

$$\mathcal{L}'_{\text{QCD}} = \mathcal{L}_{\text{QCD}} + \mathcal{L}^{\text{CP}}$$

$\theta$ -term

SMEFT operators



Nucleon electric dipole moment





# Motivations

## EW sector

Weak CP

CP violation is included

*too small...*

## QCD sector

Strong CP

$$\mathcal{L}'_{\text{QCD}} = \mathcal{L}_{\text{QCD}} + \mathcal{L}^{\text{CP}}$$

$\theta$ -term

SMEFT operators



Nucleon electric dipole moment

*never measured...*



# Motivations

P-symmetry



# Motivations

P-symmetry

QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

# Motivations

P-symmetry

QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

*Are there any effects of QCD  
P-violation on the internal  
structure of nucleons?*



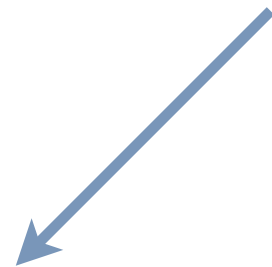
# Motivations

P-symmetry

QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

*Are there any effects of QCD  
P-violation on the internal  
structure of nucleons?*



Terms from EW sector

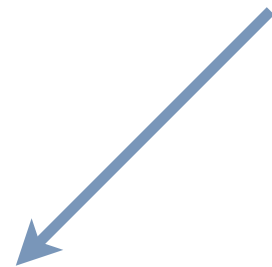
# Motivations

P-symmetry

QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

*Are there any effects of QCD  
P-violation on the internal  
structure of nucleons?*



Terms from EW sector

Weak P-violation





# Motivations

P-symmetry

QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

*Are there any effects of QCD  
P-violation on the internal  
structure of nucleons?*

Terms from EW sector

Weak P-violation



Terms from QCD sector

# Motivations

P-symmetry

QCD sector

QCD Lagrangian is assumed to be invariant under parity transformations

*Are there any effects of QCD  
P-violation on the internal  
structure of nucleons?*

Terms from EW sector

Weak P-violation



Terms from QCD sector

Strong P-violation

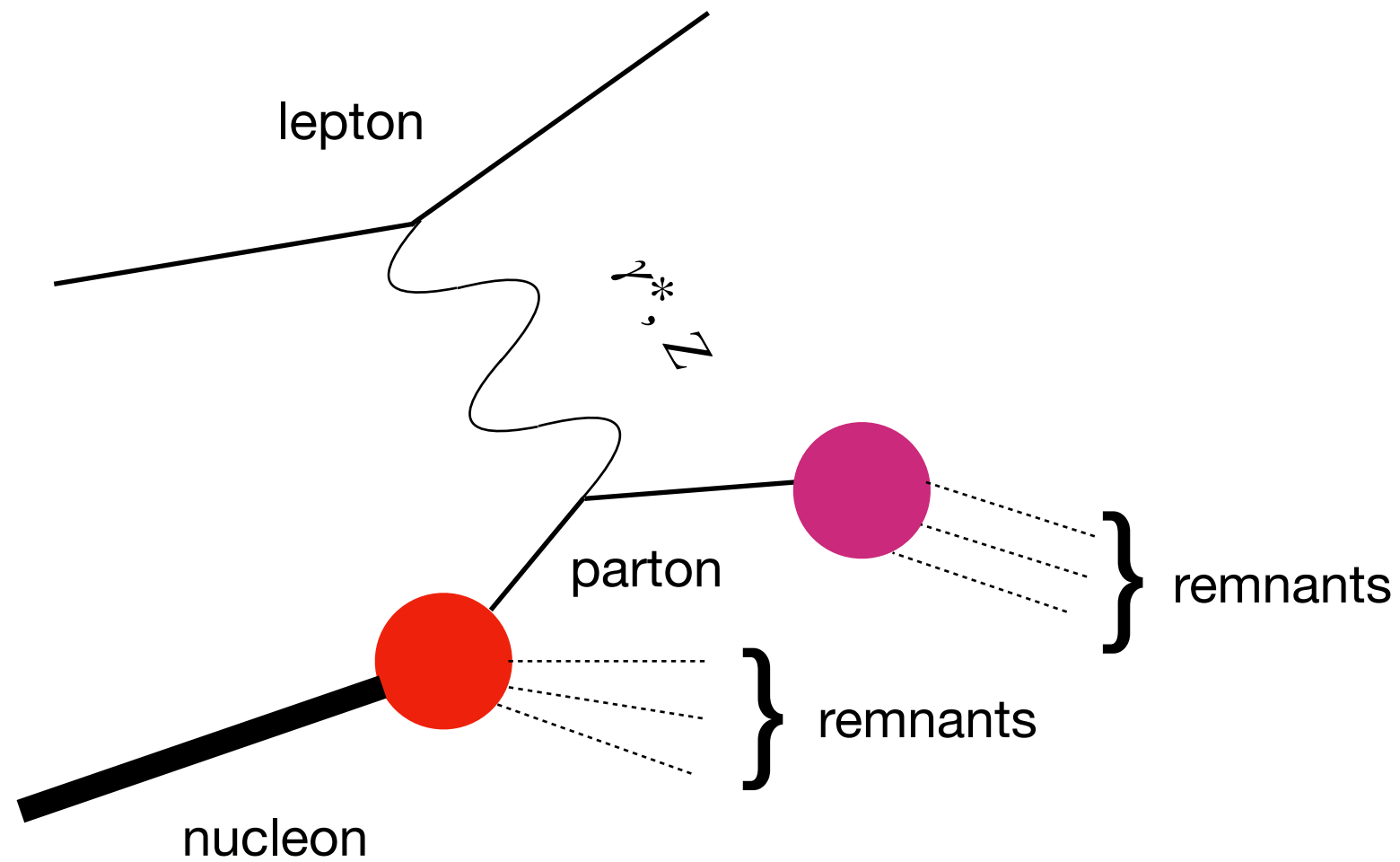




Which implications could the  
presence of strong P-violation cause  
to inclusive DIS?

# DIS process

$$l(\ell) + N(P) \rightarrow \gamma^*(q) \rightarrow l(\ell') + X$$



# Cross Section

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

In general



# Cross Section

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

In general

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} \sum_{j=\gamma, \gamma Z, Z} \eta^j L_{\mu\nu}^{(j)}(l, l'; \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

# Cross Section

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(l, l', \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

In general

$$\frac{d^3\sigma}{dx_B dy d\phi_S} = \frac{\alpha^2 y}{2Q^4} \sum_{j=\gamma, \gamma Z, Z} \eta^j L_{\mu\nu}^{(j)}(l, l'; \lambda_e) 2MW^{\mu\nu}(q, P, S)$$

$$\eta^\gamma = 1 \quad \eta^{\gamma Z} = \left( \frac{G_F M_Z^2}{2\sqrt{2}\pi\alpha} \right) \frac{Q^2}{Q^2 + M_Z^2} \quad \eta^Z = (\eta^{\gamma Z})^2$$

# Hadronic Tensor (unpolarized)

$$2MW_{\mu\nu}(q, P) = \sum_X \int \frac{d^3 P_X}{2E_X} \delta^4(P + q - P_X) \langle P | J_\mu^\dagger(0) | P_X \rangle \langle P_X | J_\nu(0) | P \rangle$$



# Hadronic Tensor (unpolarized)

$$2MW_{\mu\nu}(q, P) = \sum_X \int \frac{d^3 P_X}{2E_X} \delta^4(P + q - P_X) \langle P | J_\mu^\dagger(0) | P_X \rangle \langle P_X | J_\nu(0) | P \rangle$$

Dominant contribution on the Light-Cone

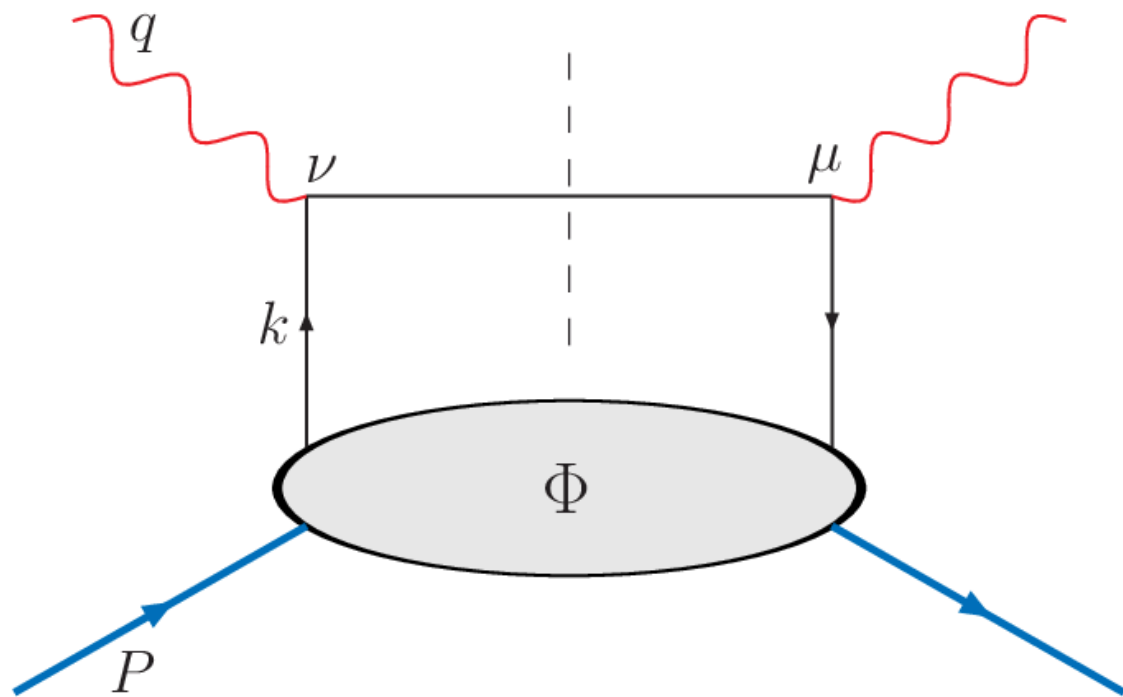
# Hadronic Tensor (unpolarized)

$$2MW_{\mu\nu}(q, P) = \sum_X \int \frac{d^3 P_X}{2E_X} \delta^4(P + q - P_X) \langle P | J_\mu^\dagger(0) | P_X \rangle \langle P_X | J_\nu(0) | P \rangle$$

Dominant contribution on the Light-Cone

$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(q, P, S) \Gamma^\mu \gamma^+ \Gamma^\nu]$$

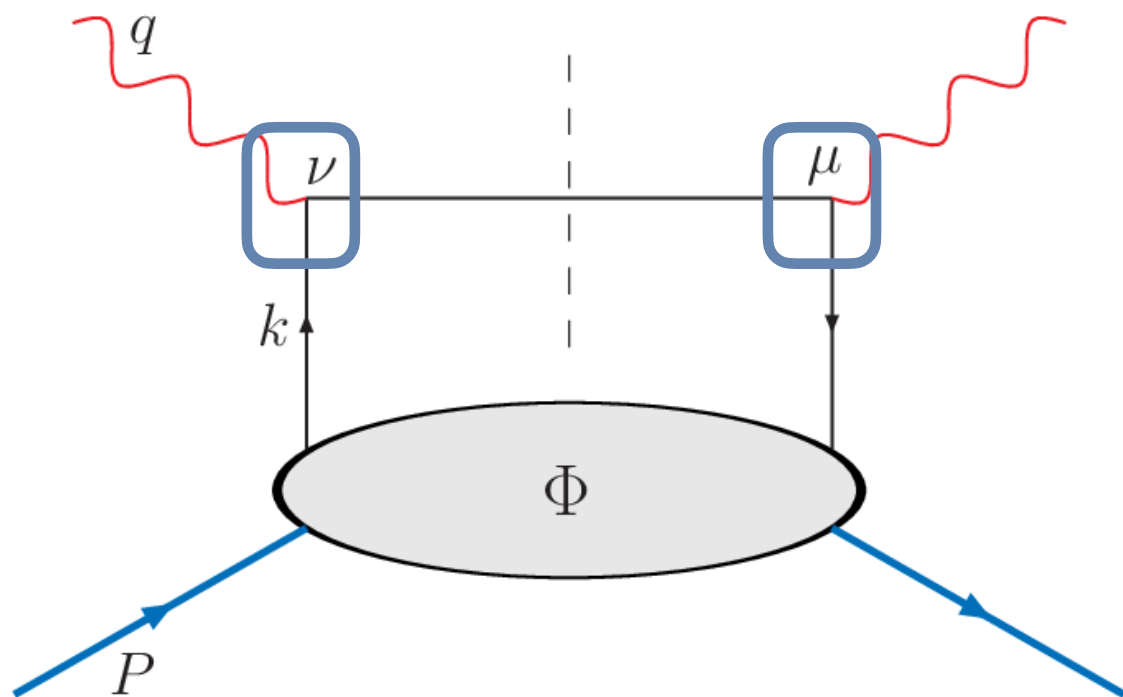
# Hadronic Tensor (unpolarized)



$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(q, P, S) \Gamma^\mu \gamma^+ \Gamma^\nu]$$



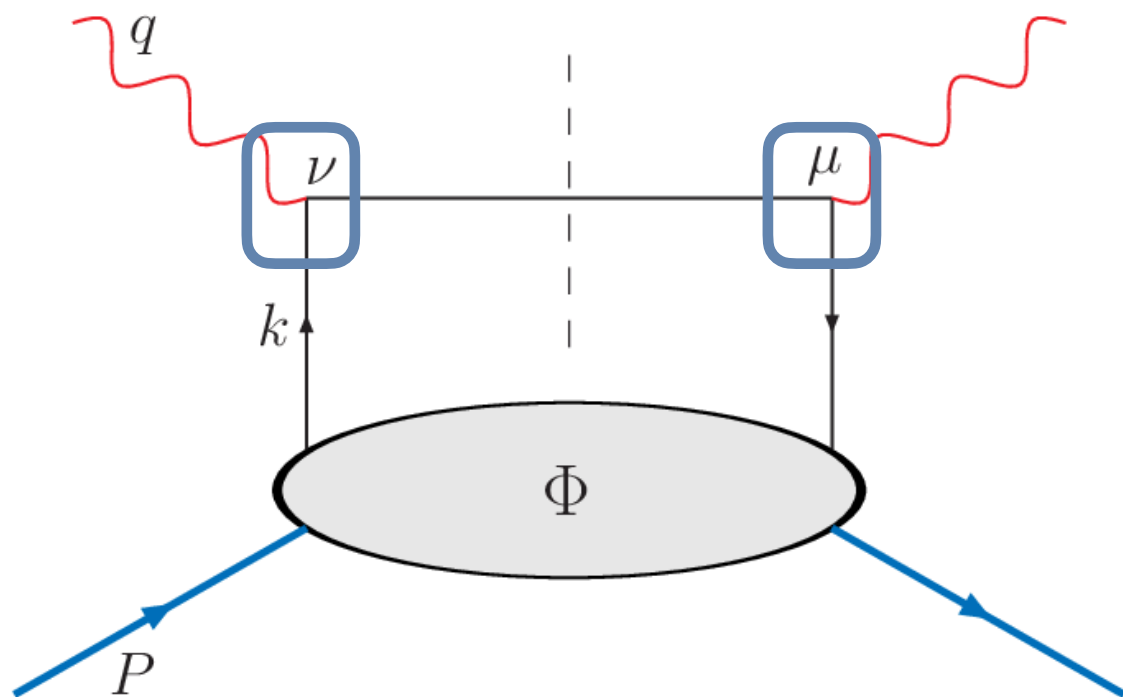
# Hadronic Tensor (unpolarized)



Vertices of the interactions

$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(q, P, S) \Gamma^\mu \gamma^+ \Gamma^\nu]$$

# Hadronic Tensor (unpolarized)

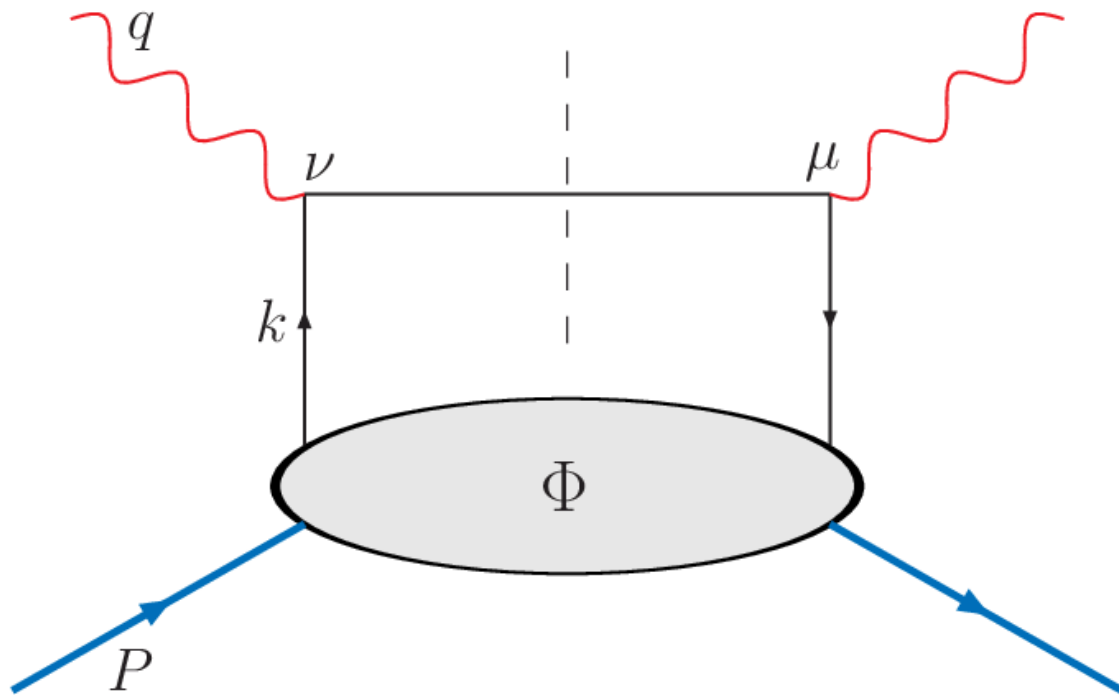


Vertices of the interactions

***P-odd structures  
already present in the  
hadronic tensor!***

$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(q, P, S) \Gamma^\mu \gamma^+ \Gamma^\nu]$$

# Hadronic Tensor (unpolarized)

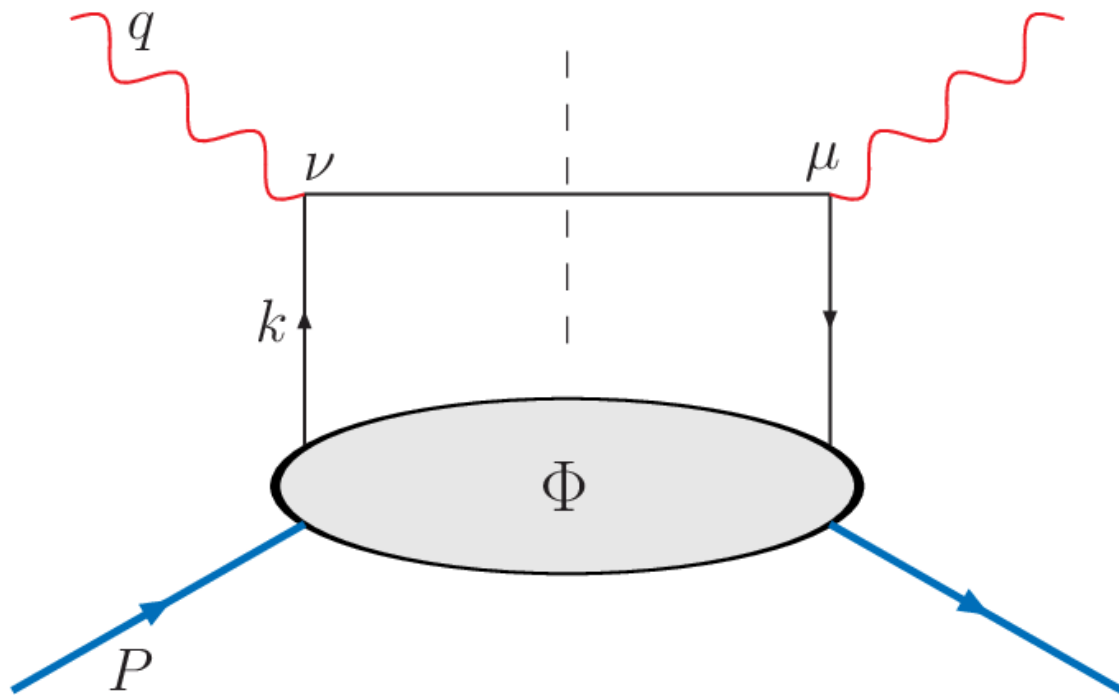


***P-odd structures  
already present in the  
hadronic tensor!***

$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(q, P, S) \Gamma^\mu \gamma^+ \Gamma^\nu]$$

Correlation distribution function

# Hadronic Tensor (unpolarized)



***P-odd structures  
already present in the  
hadronic tensor!***

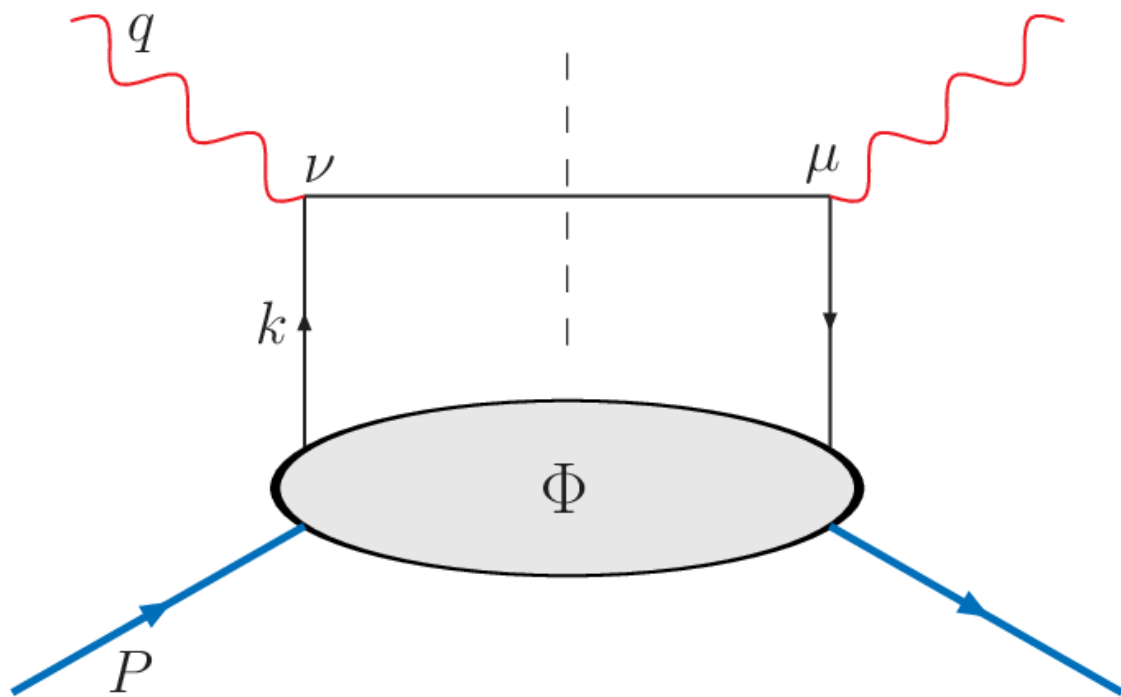
$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(q, P, S) \Gamma^\mu \gamma^+ \Gamma^\nu]$$

Correlation distribution function

$$\Phi_{ij}(k, P, S) = \int \frac{d^4\xi}{(2\pi)^4} e^{ik \cdot \xi} \langle P | \bar{\psi}_i(0) U(0, \xi) \psi_i(\xi) | P \rangle$$



# Hadronic Tensor (unpolarized)



***P-odd structures  
already present in the  
hadronic tensor!***

$$2MW^{\mu\nu}(q, P, S) = \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(q, P, S) \Gamma^\mu \gamma^+ \Gamma^\nu]$$

Correlation distribution function

$$\Phi_{ij}(k, P, S) = \int \frac{d^4\xi}{(2\pi)^4} e^{ik \cdot \xi} \langle P | \bar{\psi}_i(0) U(0, \xi) \psi_i(\xi) | P \rangle$$

Decomposition in partonic densities

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Lorenz scalar

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Lorenz scalar

Hermiticity



# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

Lorenz scalar

Hermiticity

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

Lorenz scalar

Hermiticity

$$\mathbb{1}, \gamma^\mu, \sigma^{\mu\nu}$$

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

Lorenz scalar

Hermiticity

~~Parity invariance~~

$$\mathbb{1}, \gamma^\mu, \sigma^{\mu\nu}$$

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

$$\mathbb{1}, \gamma^\mu, \sigma^{\mu\nu}$$

Lorenz scalar

Hermiticity

~~Parity invariance~~

$$i\gamma^5, \gamma^\mu \gamma^5, i\gamma^5 \sigma^{\mu\nu}$$

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

Lorenz scalar

Hermiticity

~~Parity invariance~~

$$\mathbb{1}, \gamma^\mu, \sigma^{\mu\nu}$$

$$i\gamma^5, \gamma^\mu \gamma^5, i\gamma^5 \sigma^{\mu\nu}$$

Leading twist contributions



# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

Lorenz scalar

Hermiticity

~~Parity invariance~~

$$\mathbb{1}, \gamma^\mu, \sigma^{\mu\nu}$$

$$i\gamma^5, \gamma^\mu \gamma^5, i\gamma^5 \sigma^{\mu\nu}$$

Leading twist contributions

$$\Phi_{\text{PE}}(x) \simeq \frac{1}{2} f_1(x) \gamma^-$$

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

Lorenz scalar

Hermiticity

~~Parity invariance~~

$$\mathbb{1}, \gamma^\mu, \sigma^{\mu\nu}$$

$$i\gamma^5, \gamma^\mu \gamma^5, i\gamma^5 \sigma^{\mu\nu}$$

Leading twist contributions

$$\Phi_{\text{PE}}(x) \simeq \frac{1}{2} f_1(x) \gamma^-$$

$$\Phi_{\text{PV}}(x) \simeq \frac{1}{2} g_1^{\text{PV}}(x) \gamma^5 \gamma^-$$

# Partonic correlator (unpolarized)

Integrated correlator

$$\Phi_{ij}(x_B) = \int \frac{d\xi^-}{2\pi} e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle_{\xi^+ = \xi_T = 0}$$

Lorenz scalar

Hermiticity

Parity invariance

Lorenz scalar

Hermiticity

~~Parity invariance~~

$$\mathbb{1}, \gamma^\mu, \sigma^{\mu\nu}$$

$$i\gamma^5, \gamma^\mu \gamma^5, i\gamma^5 \sigma^{\mu\nu}$$

Leading twist contributions

$$\Phi_{\text{PE}}(x) \simeq \frac{1}{2} f_1(x) \gamma^-$$

$$\Phi_{\text{PV}}(x) \simeq \frac{1}{2} g_1^{\text{PV}}(x) \gamma^5 \gamma^-$$

$$\Phi(x) = \Phi_{\text{PE}}(x) + \Phi_{\text{PV}}(x)$$

# Neutral-current DIS

$$\frac{d\sigma^\pm}{dxdy} = \frac{2\pi\alpha^2}{xyQ^2} \left[ \begin{aligned} & \left( Y_+ + \gamma^2 y^2 / 2 \right) \left( F_{2UU} + \lambda F_{2LU}^\pm \right) \\ & - y^2 \left( F_{L,UU} + \lambda F_{L,LU}^\pm \right) \\ & - \frac{Y_-}{\sqrt{1 + \gamma^2}} \left( x F_{3UU}^\pm + \lambda x F_{3LU} \right) \end{aligned} \right]$$

$$\frac{d\sigma^\pm}{dxdy} = \frac{2\pi\alpha^2}{xyQ^2} \left[ Y_+ F_2^\pm - y^2 F_L^\pm \mp Y_- x F_3^\pm \right]$$

PDG 2023

# Focus: structure function $xF_3(x, Q^2)$

$$xF_{3LU}(x, Q^2) = xF_3^{(\gamma)} - g_V^e \eta_{\gamma Z} xF_3^{(\gamma Z)} + (g_V^e{}^2 + g_A^e{}^2) \eta_Z xF_3^{(Z)}$$

# Focus: structure function $xF_3(x, Q^2)$

$$xF_{3LU}(x, Q^2) = xF_3^{(\gamma)} - g_V^e \eta_{\gamma Z} xF_3^{(\gamma Z)} + (g_V^e{}^2 + g_A^e{}^2) \eta_Z xF_3^{(Z)}$$

$$xF_3^{(\gamma)}(x, Q^2) = 0$$

$$xF_3^{(\gamma Z)}(x, Q^2) = \sum_q 2e_q g_A^q x f_1^{(q-\bar{q})}$$

$$xF_3^{(Z)}(x, Q^2) = \sum_q 2g_V^q g_A^q x f_1^{(q-\bar{q})}$$



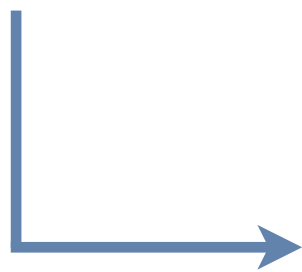
# Focus: structure function $x F_3(x, Q^2)$

$$x F_{3LU}(x, Q^2) = x F_3^{(\gamma)} - g_V^e \eta_{\gamma Z} x F_3^{(\gamma Z)} + (g_V^e{}^2 + g_A^e{}^2) \eta_Z x F_3^{(Z)}$$

$$x F_3^{(\gamma)}(x, Q^2) = 0$$

$$x F_3^{(\gamma Z)}(x, Q^2) = \sum_q 2e_q g_A^q x f_1^{(q-\bar{q})}$$

$$x F_3^{(Z)}(x, Q^2) = \sum_q 2g_V^q g_A^q x f_1^{(q-\bar{q})}$$



Additional contributions  
due to the new PV parton  
distribution

$$x \Delta F_3^{(\gamma)}(x, Q^2) = - \sum_q e_q^2 x g_1^{\text{PV}(q+\bar{q})}$$

$$x \Delta F_3^{(\gamma Z)}(x, Q^2) = - \sum_q 2e_q g_V^q x g_1^{\text{PV}(q+\bar{q})}$$

$$x \Delta F_3^{(Z)}(x, Q^2) = - \sum_q (g_V^{q2} + g_A^{q2}) x g_1^{\text{PV}(q+\bar{q})}$$

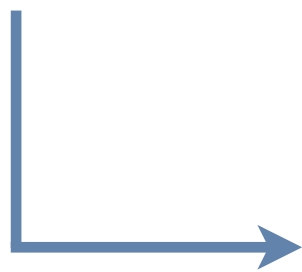
# Focus: structure function $x F_3(x, Q^2)$

$$x F_{3LU}(x, Q^2) = x F_3^{(\gamma)} - g_V^e \eta_{\gamma Z} x F_3^{(\gamma Z)} + (g_V^e{}^2 + g_A^e{}^2) \eta_Z x F_3^{(Z)}$$

$$x F_3^{(\gamma)}(x, Q^2) = 0$$

$$x F_3^{(\gamma Z)}(x, Q^2) = \sum_q 2e_q g_A^q x f_1^{(q-\bar{q})}$$

$$x F_3^{(Z)}(x, Q^2) = \sum_q 2g_V^q g_A^q x f_1^{(q-\bar{q})}$$



**Additional contributions  
due to the new PV parton  
distribution**

$$x \Delta F_3^{(\gamma)}(x, Q^2) = - \sum_q e_q^2 x g_1^{\text{PV}(q+\bar{q})}$$

$$x \Delta F_3^{(\gamma Z)}(x, Q^2) = - \sum_q 2e_q g_V^q x g_1^{\text{PV}(q+\bar{q})}$$

$$x \Delta F_3^{(Z)}(x, Q^2) = - \sum_q (g_V^{q2} + g_A^{q2}) x g_1^{\text{PV}(q+\bar{q})}$$

# Focus: structure function $x F_3(x, Q^2)$

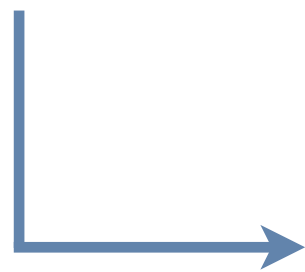
$$x F_{3LU}(x, Q^2) = x F_3^{(\gamma)} - g_V^e \eta_{\gamma Z} x F_3^{(\gamma Z)} + (g_V^e{}^2 + g_A^e{}^2) \eta_Z x F_3^{(Z)}$$

$$x F_3^{(\gamma)}(x, Q^2) = 0$$

$$x F_3^{(\gamma Z)}(x, Q^2) = \sum_q 2e_q g_A^q x f_1^{(q-\bar{q})}$$

$$x F_3^{(Z)}(x, Q^2) = \sum_q 2g_V^q g_A^q x f_1^{(q-\bar{q})}$$

**MAIN INNOVATION  
OF PV-HYPOTHESIS**



**Additional contributions  
due to the new PV parton  
distribution**

$$x \Delta F_3^{(\gamma)}(x, Q^2) = - \sum_q e_q^2 x g_1^{\text{PV}(q+\bar{q})}$$

$$x \Delta F_3^{(\gamma Z)}(x, Q^2) = - \sum_q 2e_q g_V^q x g_1^{\text{PV}(q+\bar{q})}$$

$$x \Delta F_3^{(Z)}(x, Q^2) = - \sum_q (g_V^{q2} + g_A^{q2}) x g_1^{\text{PV}(q+\bar{q})}$$

# Neutral-current DIS

$$\frac{d\sigma^\pm}{dxdy} = \frac{2\pi\alpha^2}{xyQ^2} \left[ \begin{aligned} & \left( Y_+ + \gamma^2 y^2 / 2 \right) \left( F_{2UU} + \lambda F_{2LU}^\pm \right) \\ & - y^2 \left( F_{L,UU} + \lambda F_{L,LU}^\pm \right) \\ & - \frac{Y_-}{\sqrt{1 + \gamma^2}} \left( x F_{3UU}^\pm + \lambda x F_{3LU} \right) \end{aligned} \right]$$

# Neutral-current DIS

$$\frac{d\sigma^\pm}{dxdy} = \frac{2\pi\alpha^2}{xyQ^2} \left[ \begin{aligned} & \left( Y_+ + \gamma^2 y^2 / 2 \right) \left( F_{2UU} + \lambda F_{2LU}^\pm \right) \\ & - y^2 \left( F_{L,UU} + \lambda F_{L,LU}^\pm \right) \\ & - \frac{Y_-}{\sqrt{1 + \gamma^2}} \left( x F_{3UU}^\pm + \lambda x F_{3LU} \right) \end{aligned} \right]$$

Standard DIS structure functions

# Neutral-current DIS

$$\frac{d\sigma^\pm}{dxdy} = \frac{2\pi\alpha^2}{xyQ^2} \left[ \begin{aligned} & \left( Y_+ + \gamma^2 y^2 / 2 \right) (F_{2UU} + \lambda F_{2LU}^\pm) \\ & - y^2 (F_{L,UU} + \lambda F_{L,LU}^\pm) \\ & - \frac{Y_-}{\sqrt{1 + \gamma^2}} (x F_{3UU}^\pm + \lambda x F_{3LU}) \end{aligned} \right]$$

## Standard DIS structure functions

$$F_{2UU}(x, Q^2) = F_2^{(\gamma)} - g_V^e \eta_{\gamma Z} F_2^{(\gamma Z)} + (g_V^e{}^2 + g_A^e{}^2) \eta_Z F_2^{(Z)},$$

$$F_{2LU}^\pm(x, Q^2) = \mp g_A^e \eta_{\gamma Z} F_2^{(\gamma Z)} \pm 2g_V^e g_A^e \eta_Z F_2^{(Z)},$$

$$x F_{3UU}^\pm(x, Q^2) = \mp g_A^e \eta_{\gamma Z} x F_3^{(\gamma Z)} \pm 2g_V^e g_A^e \eta_Z x F_3^{(Z)},$$

$$x F_{3LU}(x, Q^2) = x F_3^{(\gamma)} - g_V^e \eta_{\gamma Z} x F_3^{(\gamma Z)} + (g_V^e{}^2 + g_A^e{}^2) \eta_Z x F_3^{(Z)},$$

# Phenomenology



# Experimental observable

PVDIS Asymmetry

$$A_{PV} \equiv \frac{d\sigma(\lambda = 1) - d\sigma(\lambda = -1)}{d\sigma(\lambda = 1) + d\sigma(\lambda = -1)}$$

PVDIS Collaboration, *Nature* 506 (2014)  
D. Wang et al., *Phys.Rev.C* 91 (2015)

# Experimental observable

PVDIS Asymmetry

$$A_{\text{PV}} \equiv \frac{d\sigma(\lambda = 1) - d\sigma(\lambda = -1)}{d\sigma(\lambda = 1) + d\sigma(\lambda = -1)}$$
$$= \frac{Y_+ F_{2LU} - y^2 F_{L,LU} - Y_- x F_{3LU}}{Y_+ F_{2UU} - y^2 F_{L,UU} - Y_- x F_{3UU}}$$

PVDIS Collaboration, *Nature* 506 (2014)  
D. Wang et al., *Phys.Rev.C* 91 (2015)

$$Y_{\pm} = 1 \pm (1 - y)^2$$

# Experimental observable

PVDIS Asymmetry

$$A_{PV} \equiv \frac{d\sigma(\lambda = 1) - d\sigma(\lambda = -1)}{d\sigma(\lambda = 1) + d\sigma(\lambda = -1)}$$

PVDIS Collaboration, *Nature* 506 (2014)  
D. Wang et al., *Phys.Rev.C* 91 (2015)

$$= \frac{Y_+ \boxed{F_{2LU}} - y^2 \boxed{F_{L,LU}} - Y_- \boxed{x F_{3LU}}}{Y_+ \boxed{F_{2UU}} - y^2 \boxed{F_{L,UU}} - Y_- \boxed{x F_{3UU}}}$$

Contribution of  $g_1^{PV}$  in each of  
the structure functions due to  
 $\gamma Z$  and  $Z$  channels

$$Y_{\pm} = 1 \pm (1 - y)^2$$

# Available experimental data

HERA dataset  
(Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

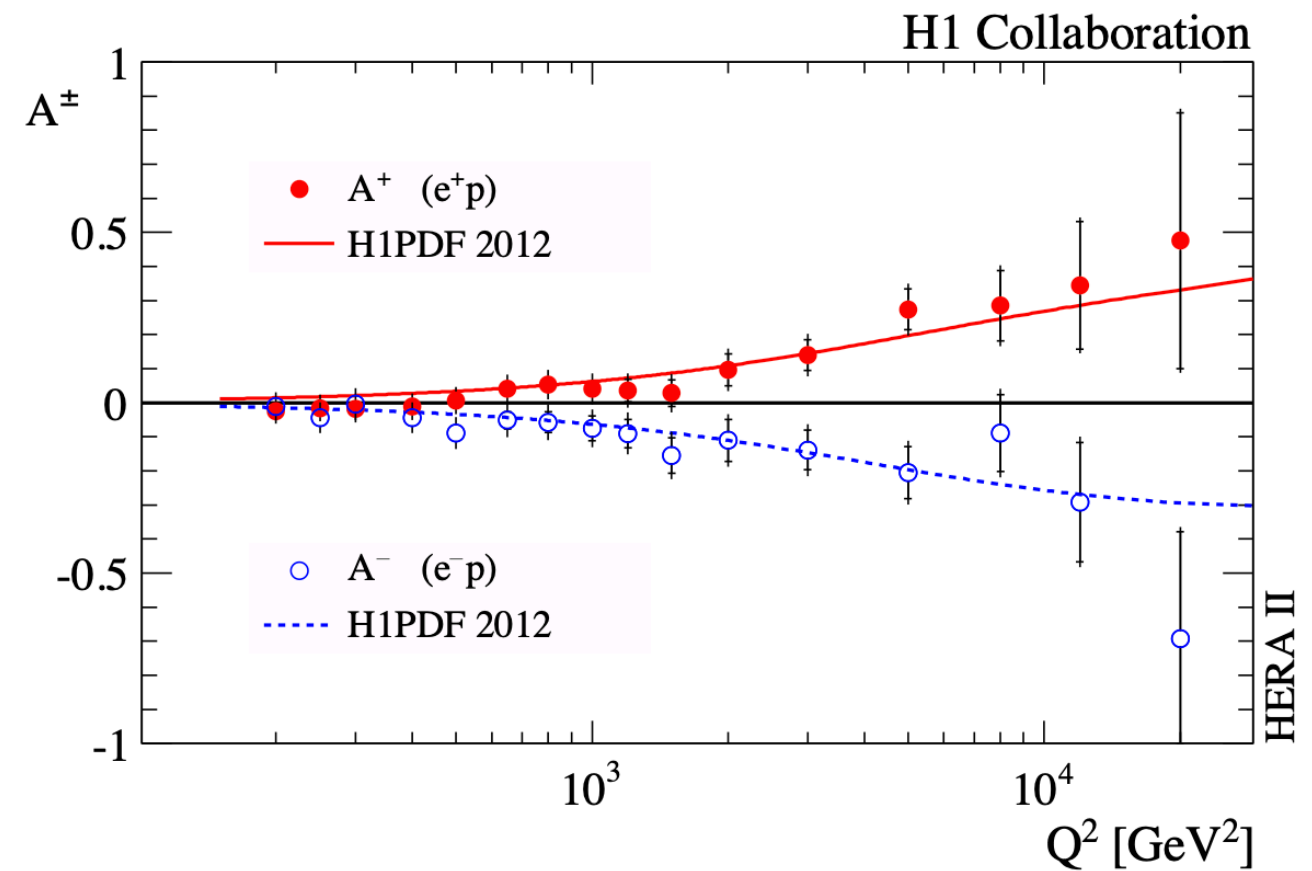
# Available experimental data

HERA dataset  
(Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

$e^+$  asymmetry: 136 data

$e^-$  asymmetry: 138 data



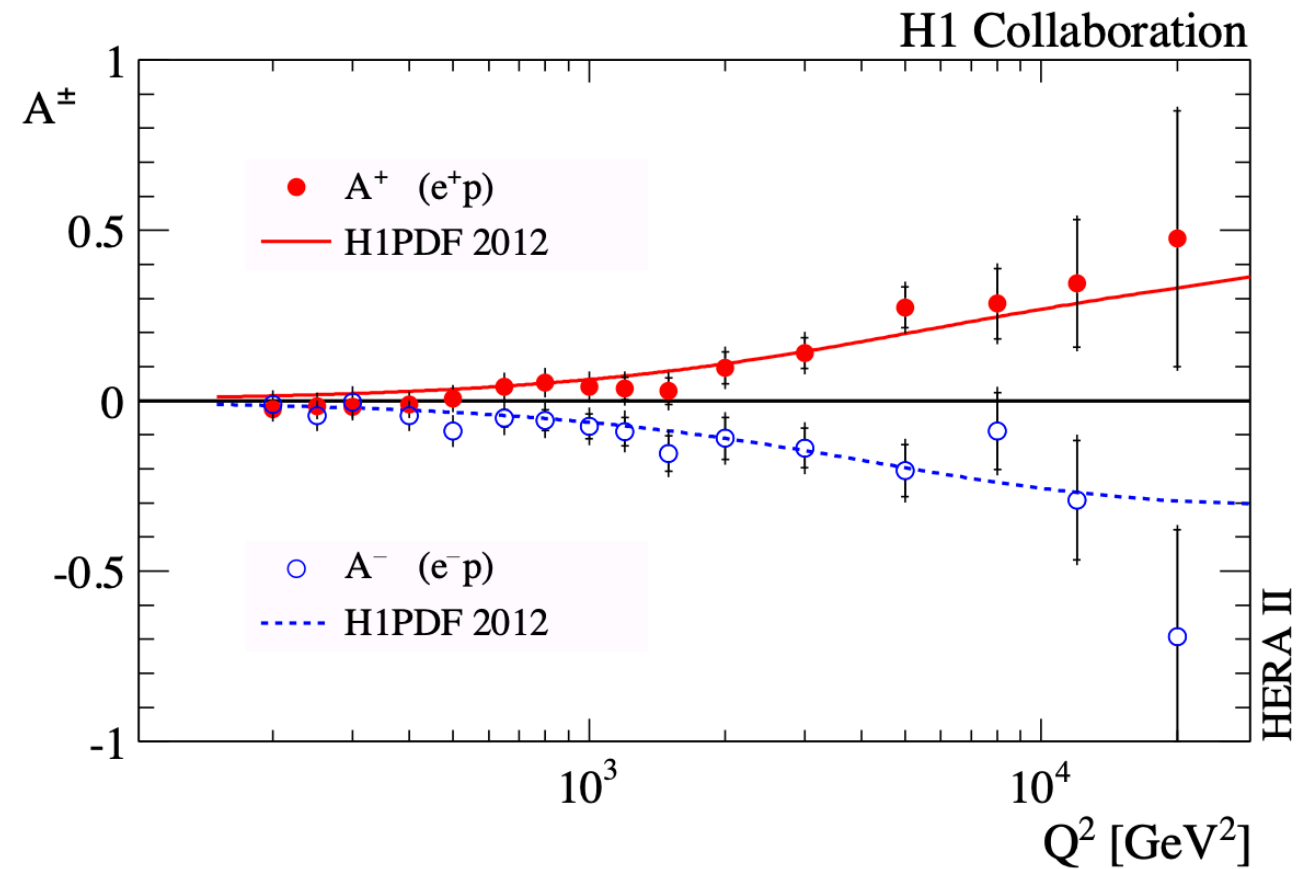
# Available experimental data

## HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

$e^+$  asymmetry: 136 data

$e^-$  asymmetry: 138 data



## JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014)

D. Wang et al., Phys.Rev.C 91 (2015)

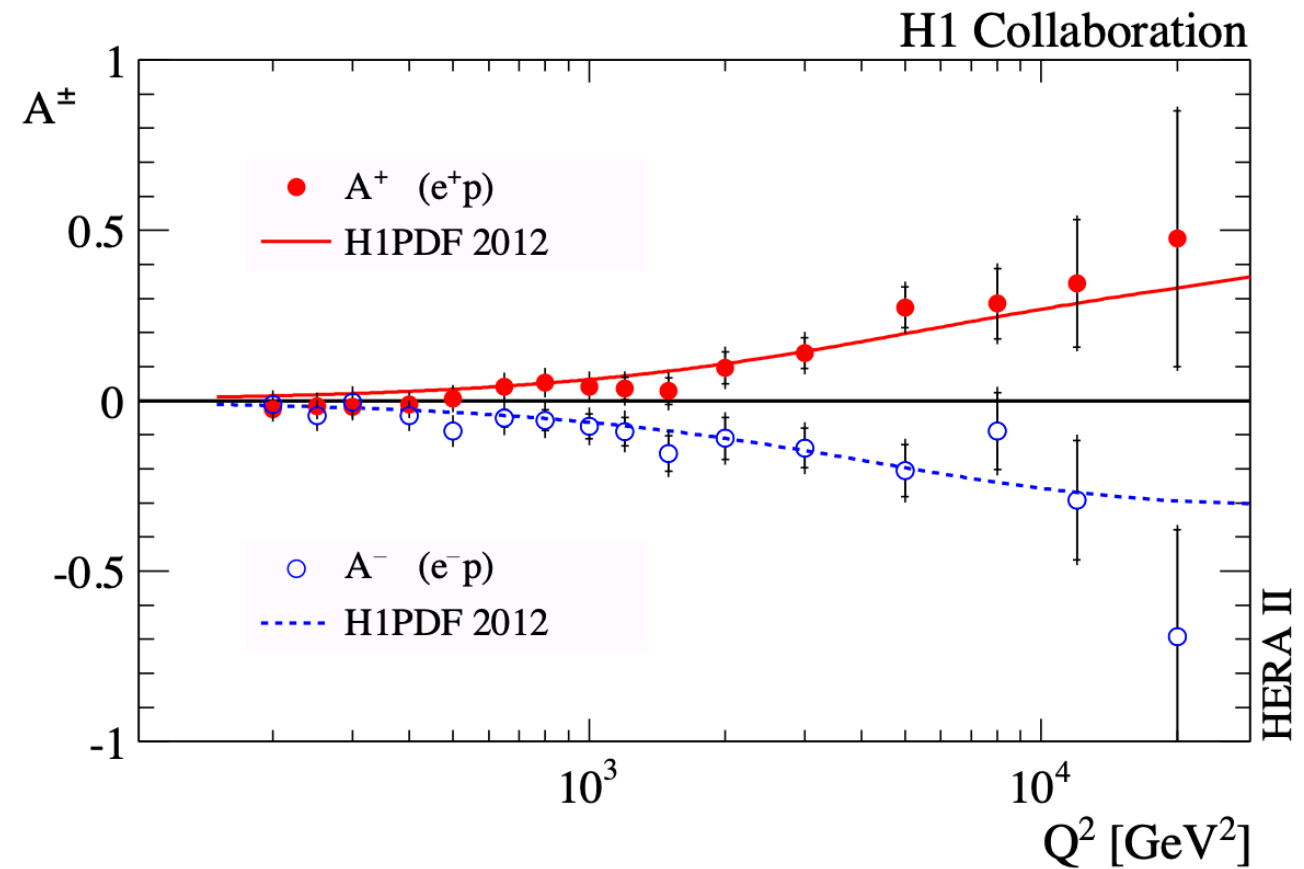
# Available experimental data

## HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

$e^+$  asymmetry: 136 data

$e^-$  asymmetry: 138 data



## JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014)

D. Wang et al., Phys.Rev.C 91 (2015)

$e^-$  asymmetry: 2 data



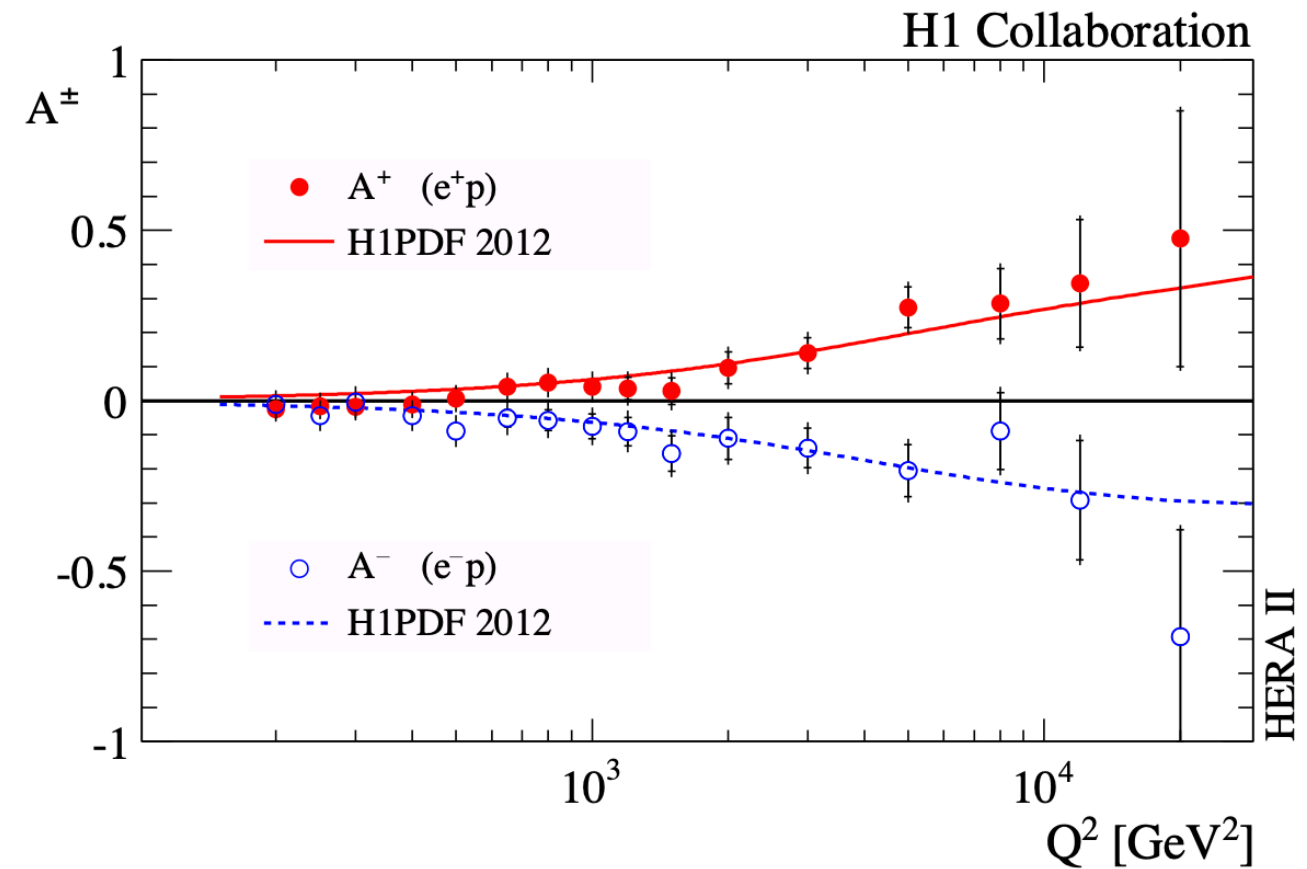
# Available experimental data

## HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

$e^+$  asymmetry: 136 data

$e^-$  asymmetry: 138 data



## JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014)  
D. Wang et al., *Phys.Rev.C* 91 (2015)

$e^-$  asymmetry: 2 data

## SLAC-E122 dataset

C.Y. Prescott et al., *Phys. Lett. B* (1979)

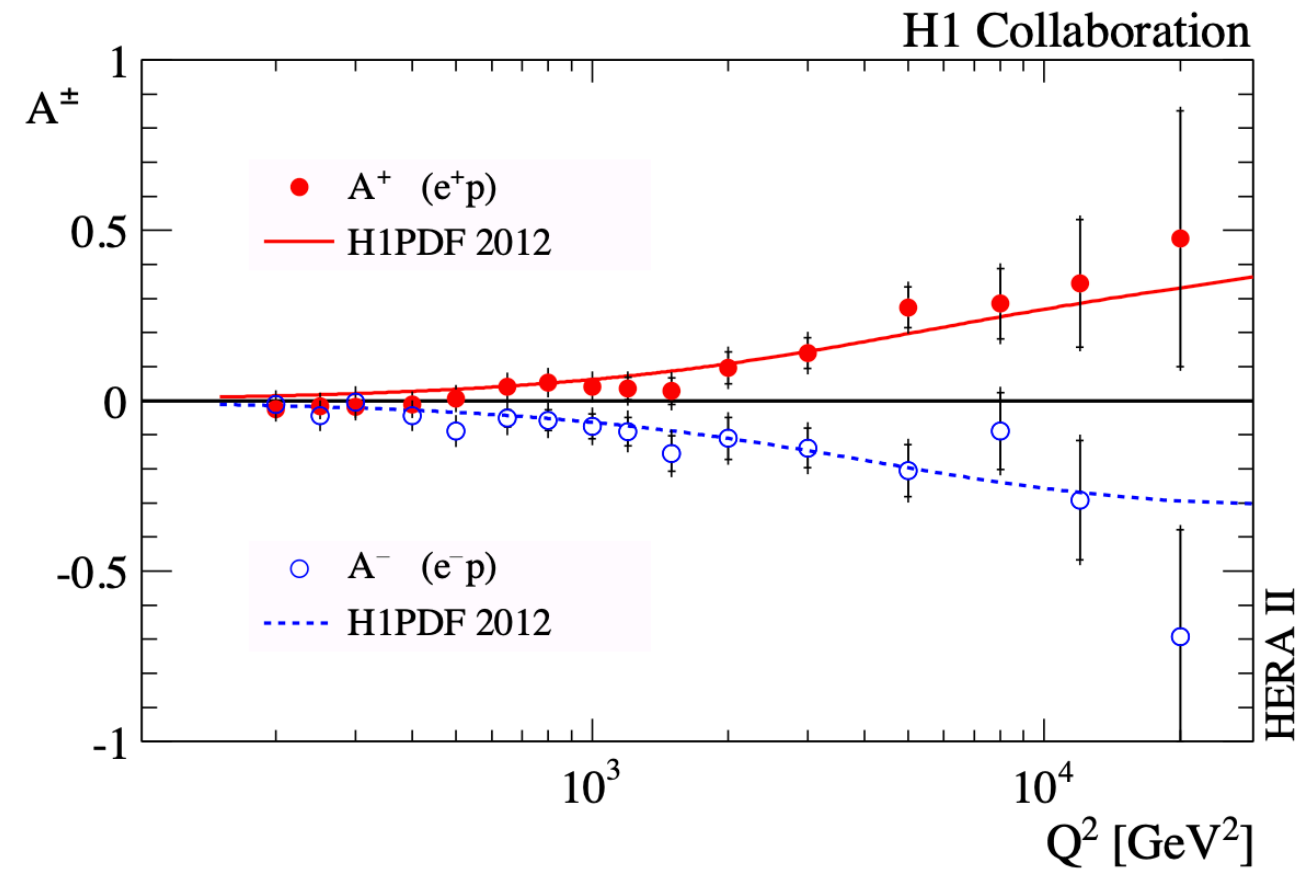
# Available experimental data

## HERA dataset (Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

$e^+$  asymmetry: 136 data

$e^-$  asymmetry: 138 data



## JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014)

D. Wang et al., *Phys.Rev.C* 91 (2015)

$e^-$  asymmetry: 2 data

## SLAC-E122 dataset

C.Y. Prescott et al., *Phys. Lett. B* (1979)

$e^-$  asymmetry: 11 data

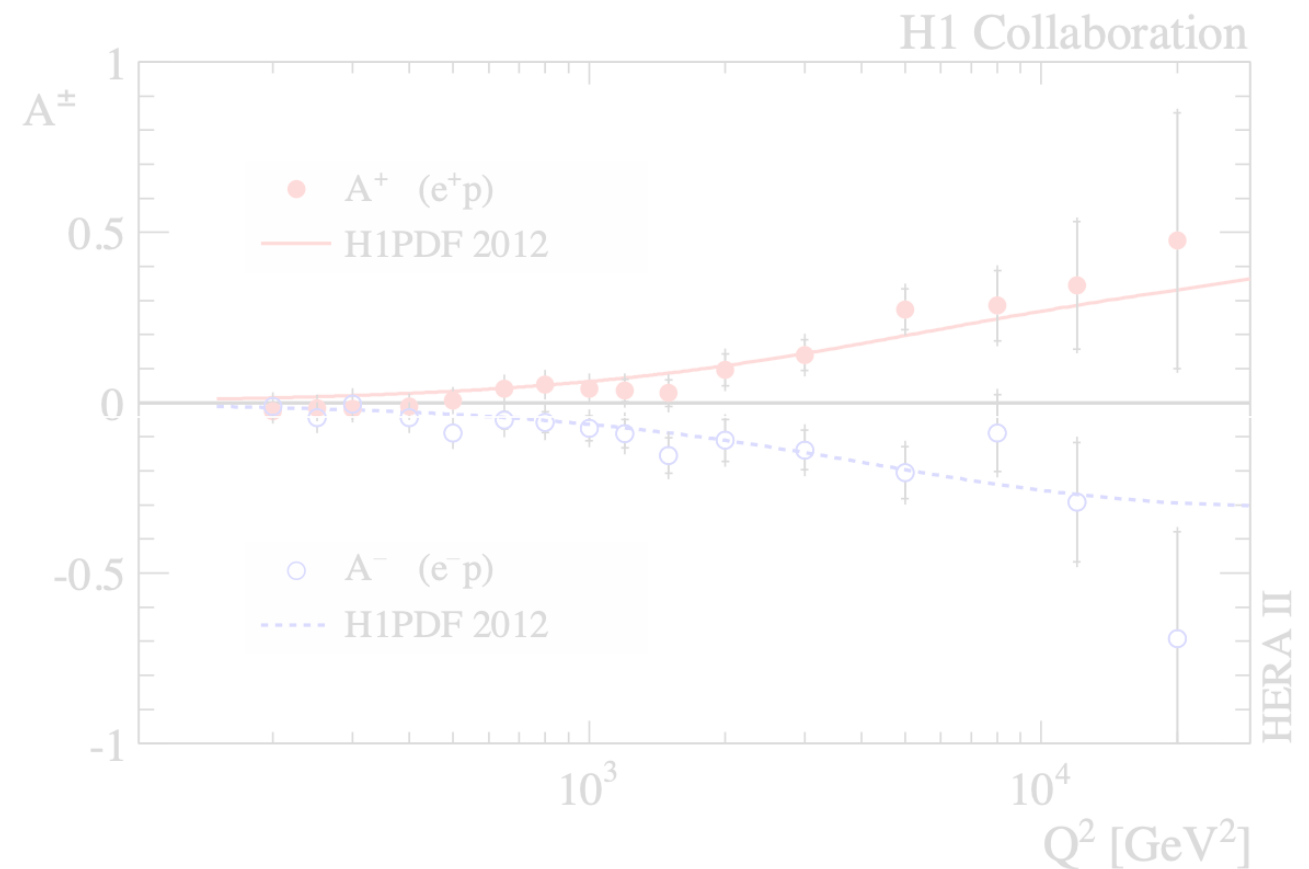
# Available experimental data

HERA dataset  
(Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

$e^+$  asymmetry: 136 data

$e^-$  asymmetry: 138 data



JLab6 PVDIS dataset

PVDIS Collaboration, *Nature* 506 (2014)

D. Wang et al., *Phys.Rev.C* 91 (2015)

$e^-$  asymmetry: 2 data

SLAC-E122 dataset

C.Y. Prescott et al., *Phys. Lett. B* (1979)

$e^-$  asymmetry: 11 data

# Available experimental data

HERA dataset  
(Run I + II combined)

H1 Collaboration, Eur. Phys. J. C 78 (2018)

$e^+$  asymmetry: 136 data

$e^-$  asymmetry: 138 data

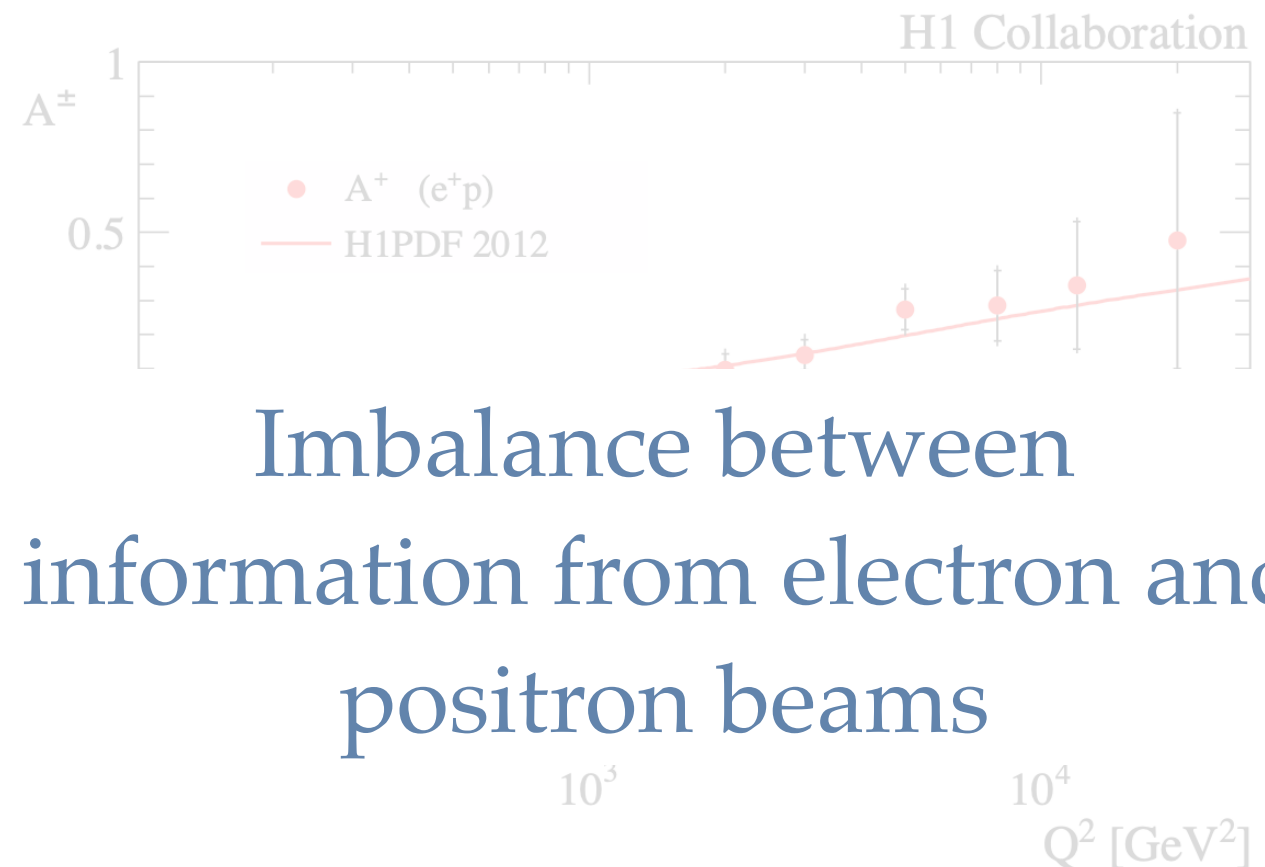
JLab6 PVDIS dataset

PVDIS Collaboration, Nature 506 (2014)

D. Wang et al., Phys.Rev.C 91 (2015)

SLAC-E122 dataset

C.Y. Prescott et al., Phys. Lett. B (1979)



Imbalance between  
information from electron and  
positron beams

$e^-$  asymmetry: 2 data

$e^-$  asymmetry: 11 data

# Experimental data: energy range

HERA dataset

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

JLab6 + SLAC-E122 datasets



# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

JLab6 + SLAC-E122 datasets

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

low-energy

$$Q^2 \simeq M_N^2$$

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

JLab6 + SLAC-E122 datasets

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

low-energy

$$Q^2 \simeq M_N^2$$

applicability of the theory?

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

JLab6 + SLAC-E122 datasets

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

low-energy

$$Q^2 \simeq M_N^2$$

applicability of the theory?

**Target-Mass Corrections**

e.g., A. Bacchetta et al., JHEP 02 (2007)

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

JLab6 + SLAC-E122 datasets

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

low-energy

$$Q^2 \simeq M_N^2$$

applicability of the theory?

**Target-Mass Corrections**

e.g., A. Bacchetta et al., JHEP 02 (2007)

**EW radiative corrections**

J. Erler, S. Su, Prog.Part.Nucl.Phys. 71 (2013)

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

JLab6 + SLAC-E122 datasets

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

low-energy

$$Q^2 \simeq M_N^2$$

applicability of the theory?

**Target-Mass Corrections**

e.g., A. Bacchetta et al., JHEP 02 (2007)

**EW radiative corrections**

J. Erler, S. Su, Prog.Part.Nucl.Phys. 71 (2013)

$$C_{1u} = 2g_A^e g_V^u = 2 \left( -\frac{1}{2} \right) \left( \frac{1}{2} - \frac{4}{3} \sin^2 \theta_W \right)$$

$$C_{2u} = 2g_V^e g_A^u = 2 \left( -\frac{1}{2} + 2 \sin^2 \theta_W \right) \left( \frac{1}{2} \right)$$

$$C_{1d} = 2g_A^e g_V^d = 2 \left( -\frac{1}{2} \right) \left( -\frac{1}{2} + \frac{2}{3} \sin^2 \theta_W \right)$$

$$C_{2d} = 2g_V^e g_A^d = 2 \left( -\frac{1}{2} + 2 \sin^2 \theta_W \right) \left( -\frac{1}{2} \right)$$

# Experimental data: energy range

HERA dataset

$$Q^2 \in (200, 30000) \text{ GeV}^2$$

high-energy

$$Q^2 \gg M_N^2$$

no need of modification of the theory

JLab6 + SLAC-E122 datasets

$$Q^2 \in (0.9, 1.9) \text{ GeV}^2$$

low-energy

$$Q^2 \simeq M_N^2$$

applicability of the theory?

$$\begin{aligned} C_{1u} &= 2g_A^e g_V^u = 2 \left( -\frac{1}{2} \right) \left( \frac{1}{2} - \frac{4}{3} \sin^2 \theta_W \right) \\ C_{2u} &= 2g_V^e g_A^u = 2 \left( -\frac{1}{2} + 2 \sin^2 \theta_W \right) \left( \frac{1}{2} \right) \\ C_{1d} &= 2g_A^e g_V^d = 2 \left( -\frac{1}{2} \right) \left( -\frac{1}{2} + \frac{2}{3} \sin^2 \theta_W \right) \\ C_{2d} &= 2g_V^e g_A^d = 2 \left( -\frac{1}{2} + 2 \sin^2 \theta_W \right) \left( -\frac{1}{2} \right) \end{aligned}$$

## **Target-Mass Corrections**

e.g., A. Bacchetta et al., JHEP 02 (2007)

## **EW radiative corrections**

J. Erler, S. Su, Prog.Part.Nucl.Phys. 71 (2013)

$$\begin{aligned} C_{1u}^{\text{SM}} &= -0.1887 - 0.0011 \times \frac{2}{3} \ln(\langle Q^2 \rangle / 0.14 \text{ GeV}^2) \\ C_{1d}^{\text{SM}} &= 0.3419 - 0.0011 \times \frac{-1}{3} \ln(\langle Q^2 \rangle / 0.14 \text{ GeV}^2) \\ C_{2u}^{\text{SM}} &= -0.0351 - 0.0009 \ln(\langle Q^2 \rangle / 0.078 \text{ GeV}^2) \\ C_{2d}^{\text{SM}} &= 0.0248 + 0.0007 \ln(\langle Q^2 \rangle / 0.021 \text{ GeV}^2) \end{aligned}$$

# Parameterization of $g_1^{PV}(x, Q^2)$

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$$\gamma^5 \gamma^\mu$$



# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF  $g_1(x, Q^2)$**

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF  $g_1(x, Q^2)$**   
 $\longrightarrow$  **C-odd**

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF  $g_1(x, Q^2)$**

$\longrightarrow$  **C-odd**

$$xF_3^j(x, Q^2) = \sum_q C_q^j x f_1^{(q-\bar{q})}$$

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF**  $g_1(x, Q^2)$

$\longrightarrow$  **C-odd**

$$xF_3^j(x, Q^2) = \sum_q C_q^j x f_1^{(q-\bar{q})}$$

$$\Delta x F_3^j(x, Q^2) = - \sum_q C_q^{\prime j} x \alpha g_1$$

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF**  $g_1(x, Q^2)$

$\longrightarrow$  **C-odd**

$$xF_3^j(x, Q^2) = \sum_q C_q^j x f_1^{(q-\bar{q})}$$

$$\Delta x F_3^j(x, Q^2) = - \sum_q C_q^{\prime j} x \alpha g_1^{(q+\bar{q})}$$

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF**  $g_1(x, Q^2)$

$\longrightarrow$  **C-odd**

$$xF_3^j(x, Q^2) = \sum_q C_q^j x f_1^{(q-\bar{q})}$$

$$\Delta x F_3^j(x, Q^2) = - \sum_q C_q^{\prime j} x \alpha g_1^{(q+\bar{q})}$$

$$F_2^j(x, Q^2) = \sum_q \hat{C}_q^j x f_1^{(q+\bar{q})}$$

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF**  $g_1(x, Q^2)$

$\longrightarrow$  **C-odd**

$$xF_3^j(x, Q^2) = \sum_q C_q^j x f_1^{(q-\bar{q})}$$

$$\Delta x F_3^j(x, Q^2) = - \sum_q C_q^{\prime j} x \alpha g_1^{(q+\bar{q})}$$

$$F_2^j(x, Q^2) = \sum_q \hat{C}_q^j x f_1^{(q+\bar{q})}$$

$$\Delta F_2^j(x, Q^2) = - \sum_q \hat{C}_q^{\prime j} x \alpha g_1^{(q-\bar{q})}$$

# Parameterization of $g_1^{PV}(x, Q^2)$

PV parton density comes from the structure

$\gamma^5 \gamma^\mu$   $\longrightarrow$  **Same evolution as helicity PDF**  $g_1(x, Q^2)$

$\longrightarrow$  **C-odd**

$$xF_3^j(x, Q^2) = \sum_q C_q^j x f_1^{(q-\bar{q})}$$

$$\Delta x F_3^j(x, Q^2) = - \sum_q C_q^{\prime j} x \alpha g_1^{(q+\bar{q})}$$

$$F_2^j(x, Q^2) = \sum_q \hat{C}_q^j x f_1^{(q+\bar{q})}$$

$$\Delta F_2^j(x, Q^2) = - \sum_q \hat{C}_q^{\prime j} x \alpha g_1^{(q-\bar{q})}$$

**1 parameter to be fitted**



# Error propagation

PDF set for

# Error propagation

PDF set for

$$f_1(x, Q^2)$$

*NNPDF4.0*

Ball et al. (NNPDF), EPJ C 82 (2022)

# Error propagation

PDF set for

$$f_1(x, Q^2)$$

*NNPDF4.0*

Ball et al. (NNPDF), EPJ C 82 (2022)

$$g_1(x, Q^2)$$

*NNPDFpol1.1*

Nocera et al. (NNPDF), Nucl. Phys. B 887 (2014)

# Error propagation

PDF set for

$$f_1(x, Q^2)$$

*NNPDF4.0*

Ball et al. (NNPDF), EPJ C 82 (2022)

$$g_1(x, Q^2)$$

*NNPDFpol1.1*

Nocera et al. (NNPDF), Nucl. Phys. B 887 (2014)

100 MC replicas of unpolarized PDF

# Error propagation

PDF set for

$$f_1(x, Q^2)$$

*NNPDF4.0*

Ball et al. (NNPDF), EPJ C 82 (2022)

$$g_1(x, Q^2)$$

*NNPDFpol1.1*

Nocera et al. (NNPDF), Nucl. Phys. B 887 (2014)

100 MC replicas of unpolarized PDF

100 MC replicas of helicity PDF

# Error propagation

PDF set for

$$f_1(x, Q^2)$$

*NNPDF4.0*

Ball et al. (NNPDF), EPJ C 82 (2022)

$$g_1(x, Q^2)$$

*NNPDFpol1.1*

Nocera et al. (NNPDF), Nucl. Phys. B 887 (2014)

100 MC replicas of unpolarized PDF

100 MC replicas of helicity PDF

100 MC replicas experimental data

# Error propagation

PDF set for

$$f_1(x, Q^2)$$

*NNPDF4.0*

Ball et al. (NNPDF), EPJ C 82 (2022)

$$g_1(x, Q^2)$$

*NNPDFpol1.1*

Nocera et al. (NNPDF), Nucl. Phys. B 887 (2014)

100 MC replicas of unpolarized PDF

100 MC replicas of helicity PDF

100 MC replicas experimental data



# Error propagation

PDF set for

$$f_1(x, Q^2)$$

*NNPDF4.0*

Ball et al. (NNPDF), EPJ C 82 (2022)

$$g_1(x, Q^2)$$

*NNPDFpol1.1*

Nocera et al. (NNPDF), Nucl. Phys. B 887 (2014)

100 MC replicas of unpolarized PDF

100 MC replicas of helicity PDF

100 MC replicas experimental data

Statistical distribution of  
100 values of parameter  $\alpha$

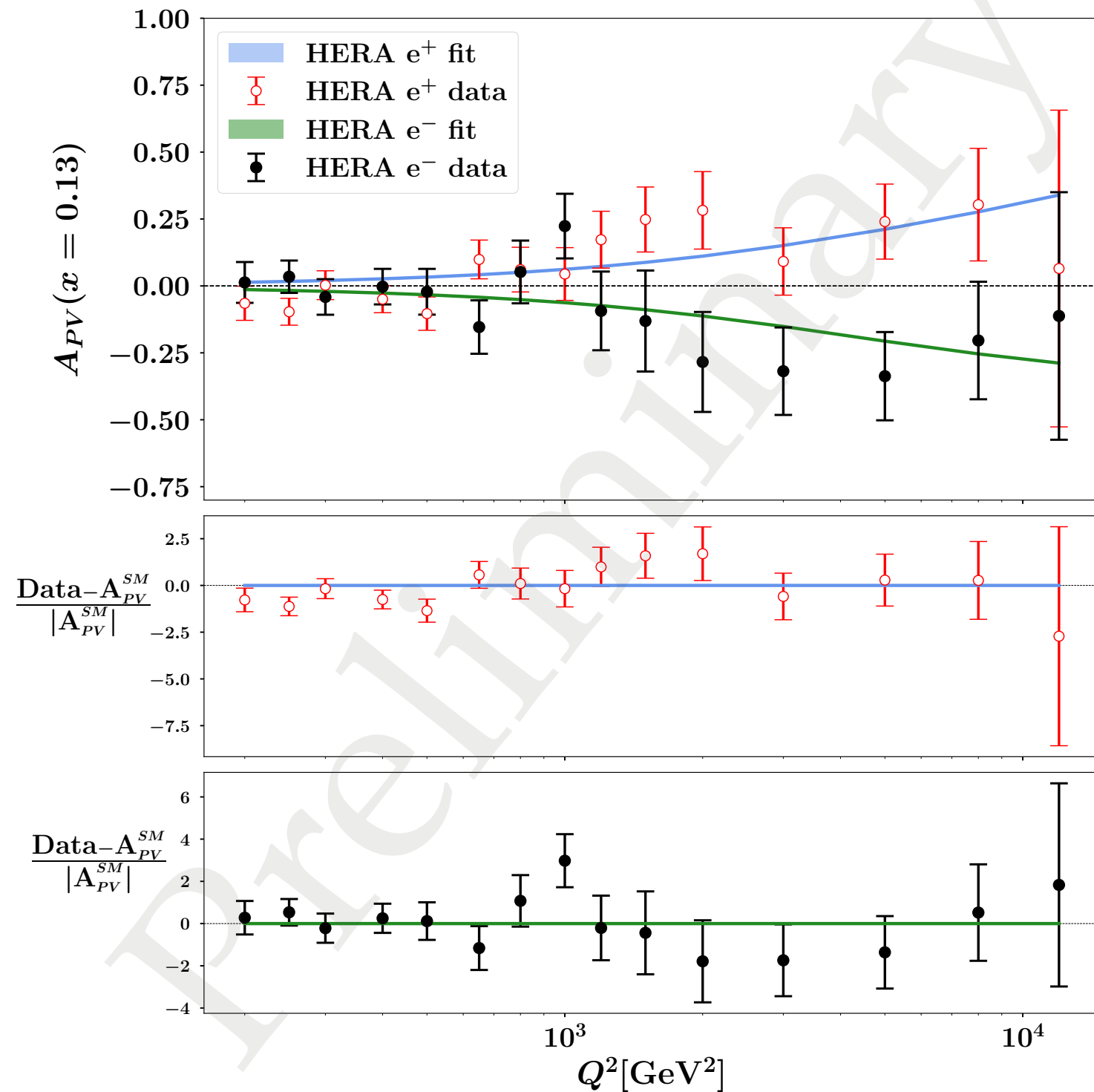


# Results of the fit: $\chi^2$ values

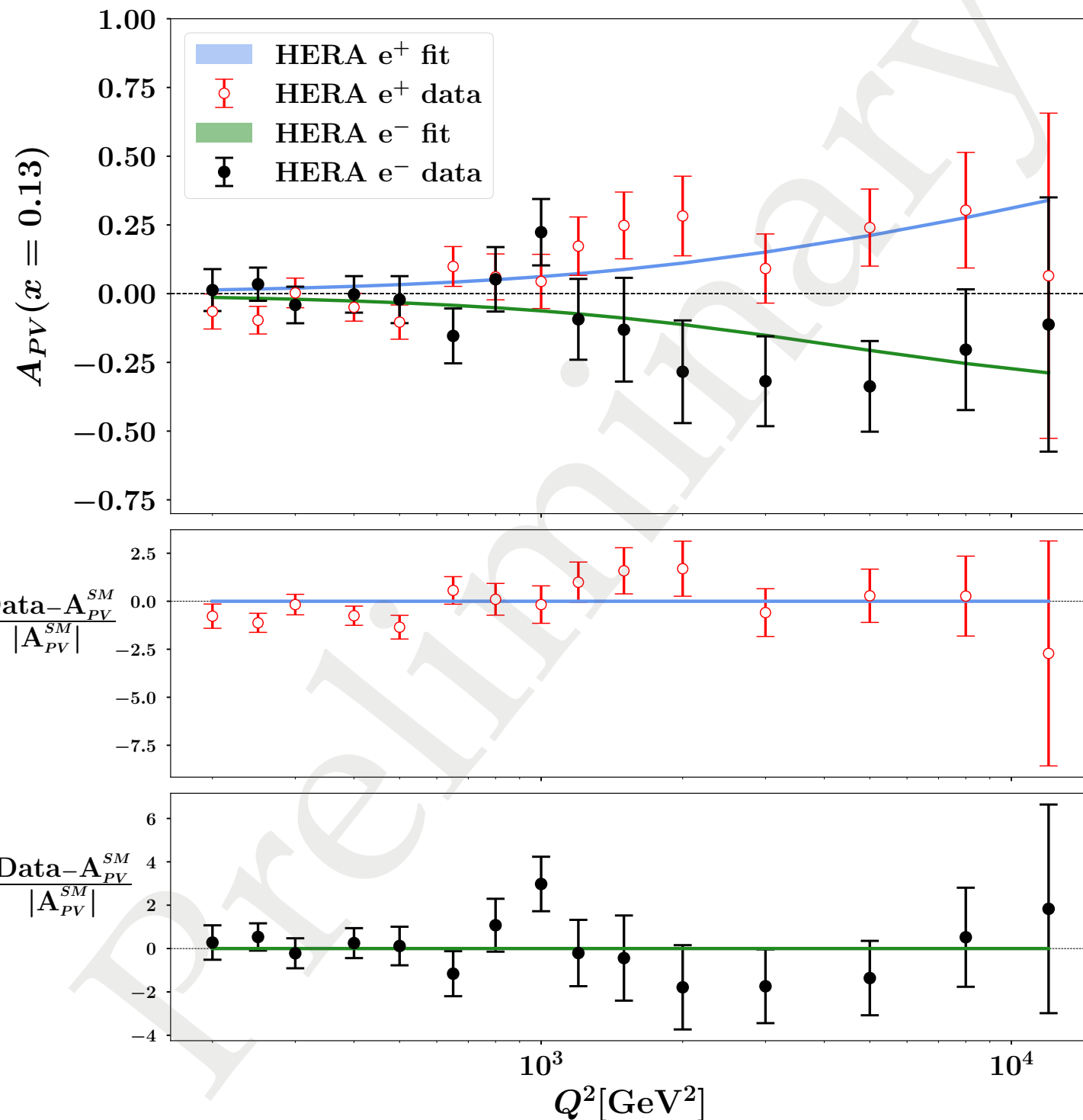
Fit **WITH** EW radiative corrections

	N of points	$\chi^2/N_{\text{data}}$ (SM)	$\chi^2/N_{\text{data}}$ ( <b>Fit</b> )
HERA $A^+$	136	1.12	1.12
HERA $A^-$	138	0.98	0.98
JLab6 $A^-$	2	0.67	0.42
SLAC-E122 $A^-$	11	0.97	0.94
<b><i>TOTAL</i></b>	<b><i>287</i></b>	<b><i>1.042</i></b>	<b><i>1.037</i></b>

# Results of the fit: data-theory comparison

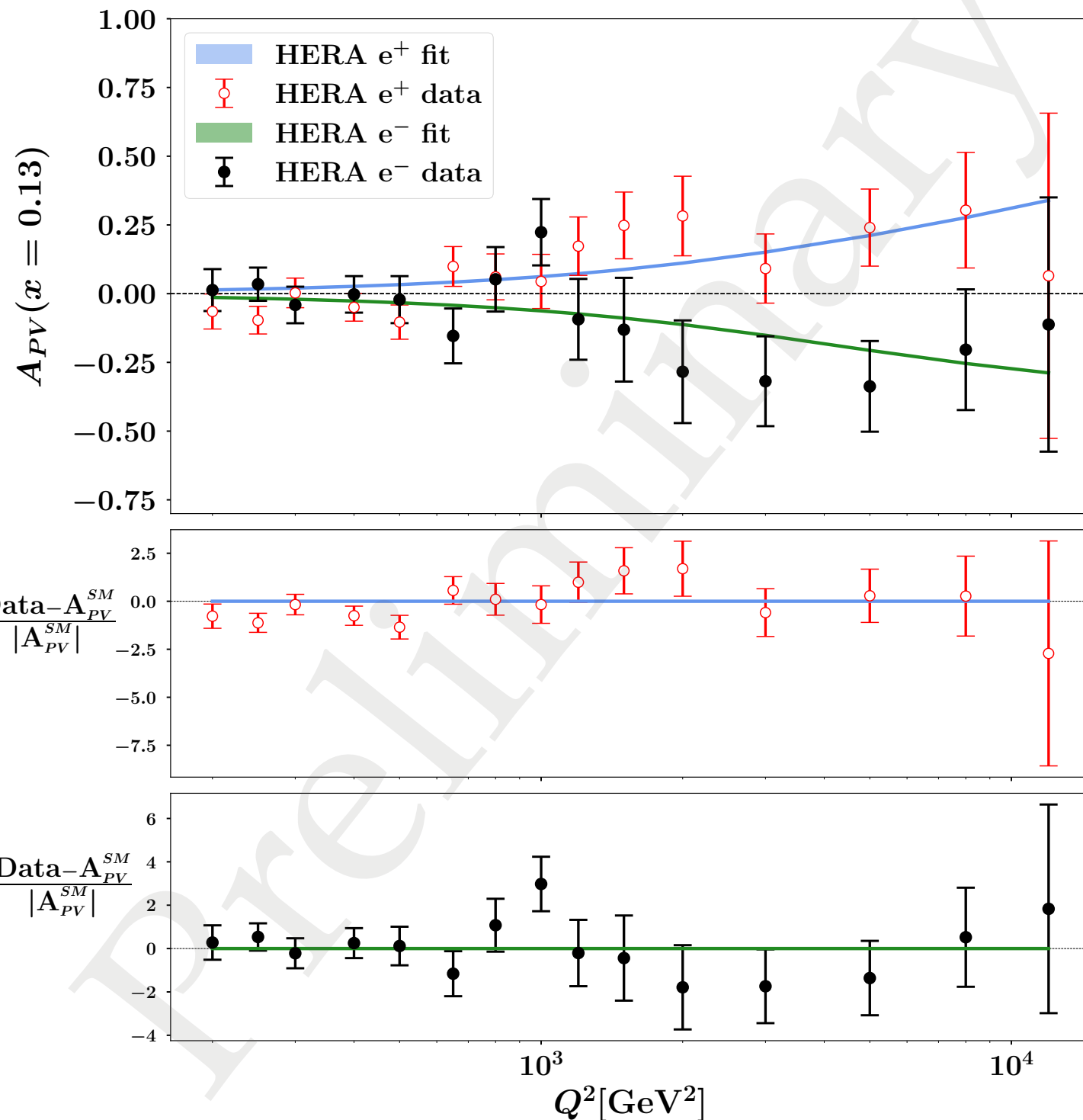


# Results of the fit: data-theory comparison



Very small uncertainties in the predictions because the fit is dominated by data with smaller errors

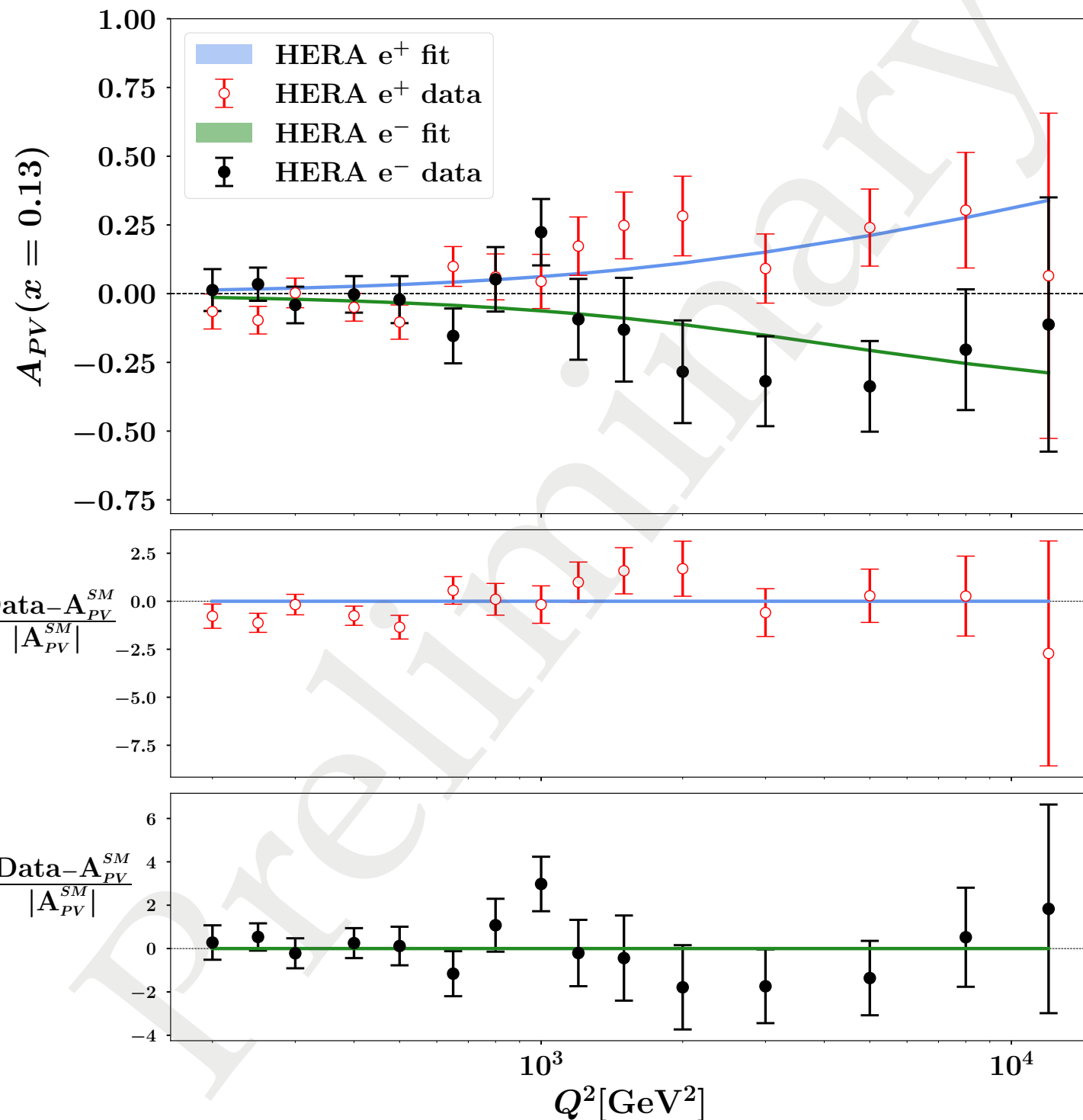
# Results of the fit: data-theory comparison



Very small uncertainties in the predictions because the fit is dominated by data with smaller errors

There's room for a better description for positron asymmetry at low- $Q$

# Results of the fit: data-theory comparison

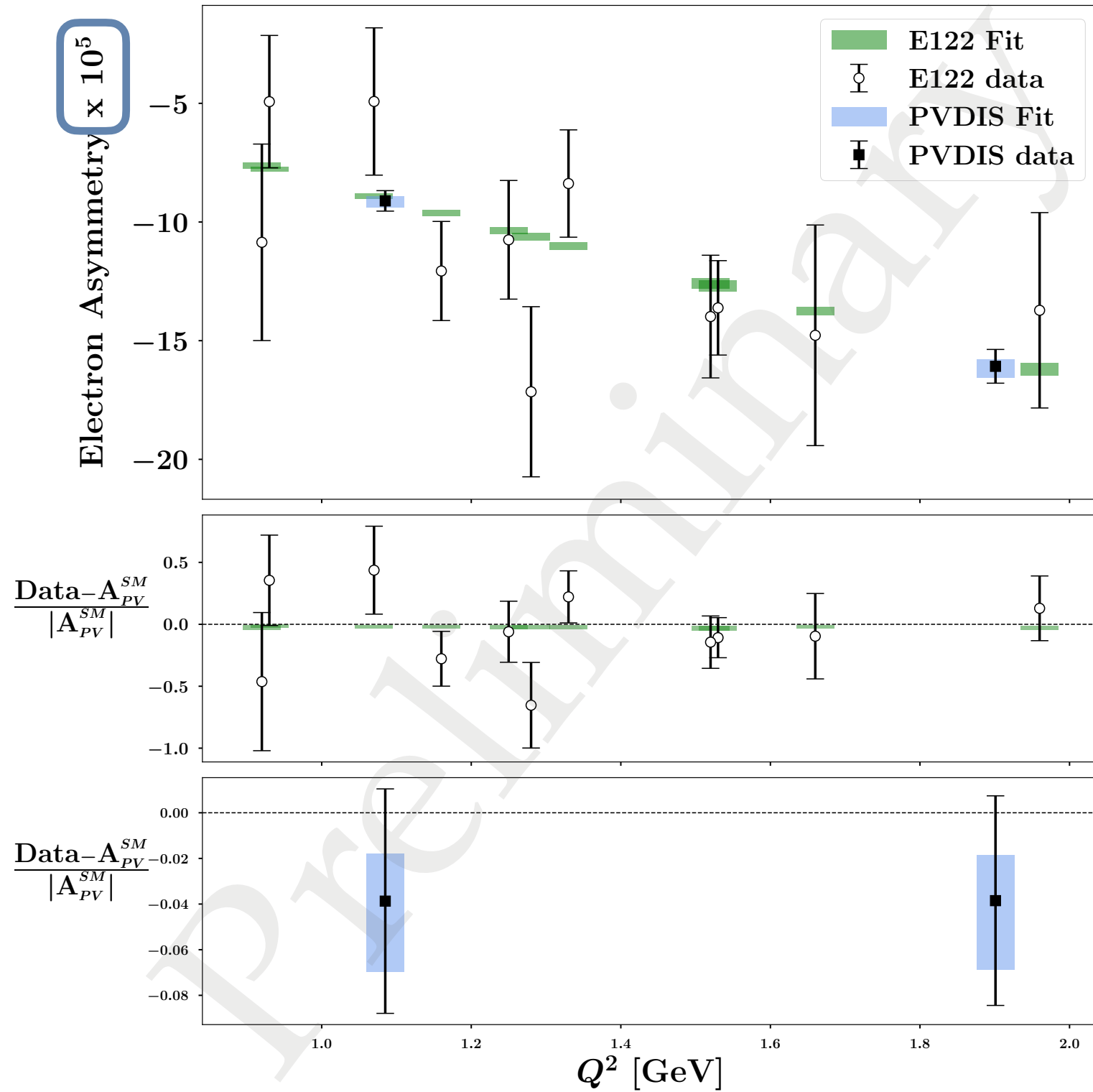


Very small uncertainties in the predictions because the fit is dominated by data with smaller errors

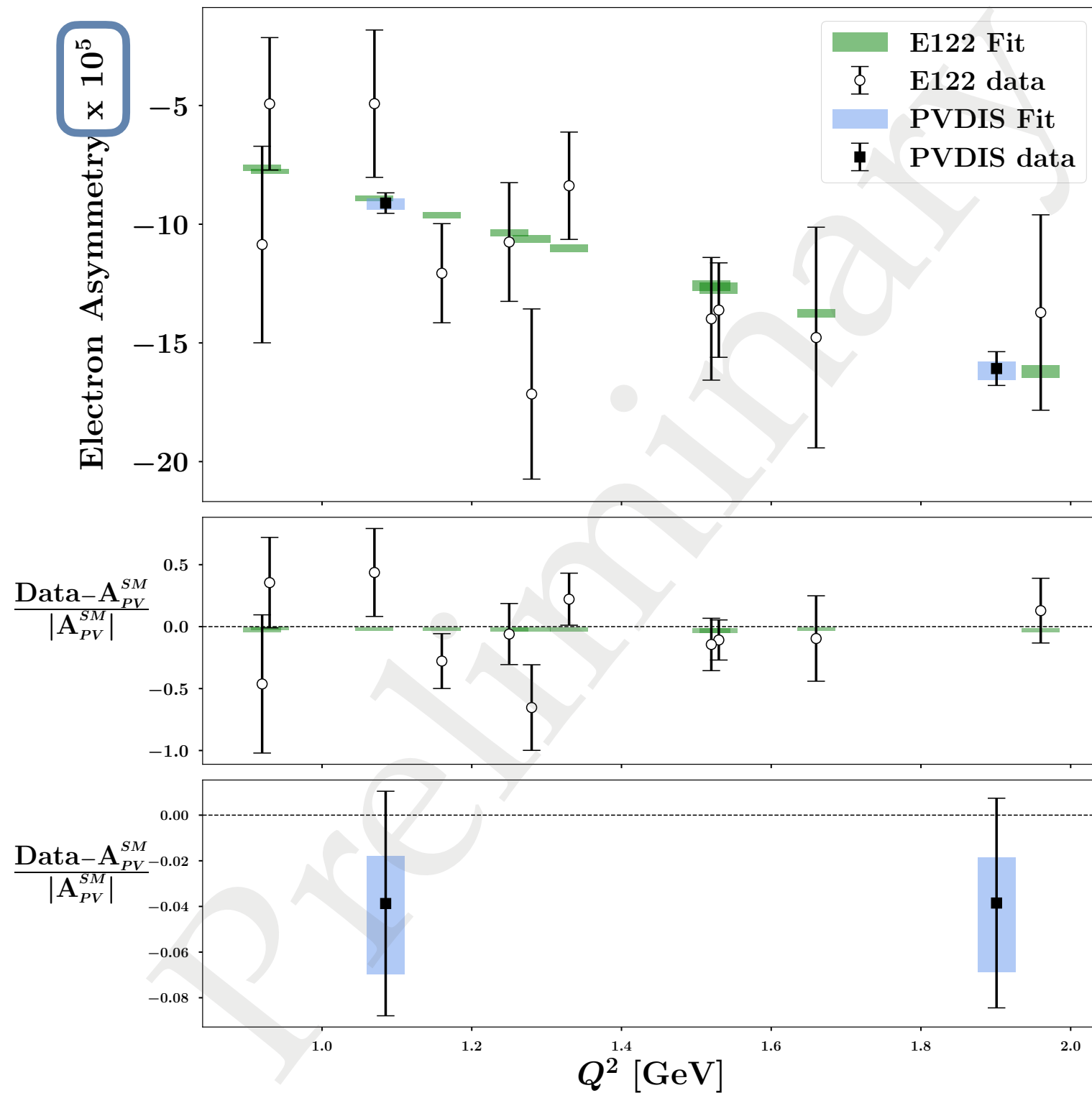
There's room for a better description for positron asymmetry at low- $Q$

Agreement for electron asymmetry, but too large errors at low- $Q$

# Results of the fit: data-theory comparison

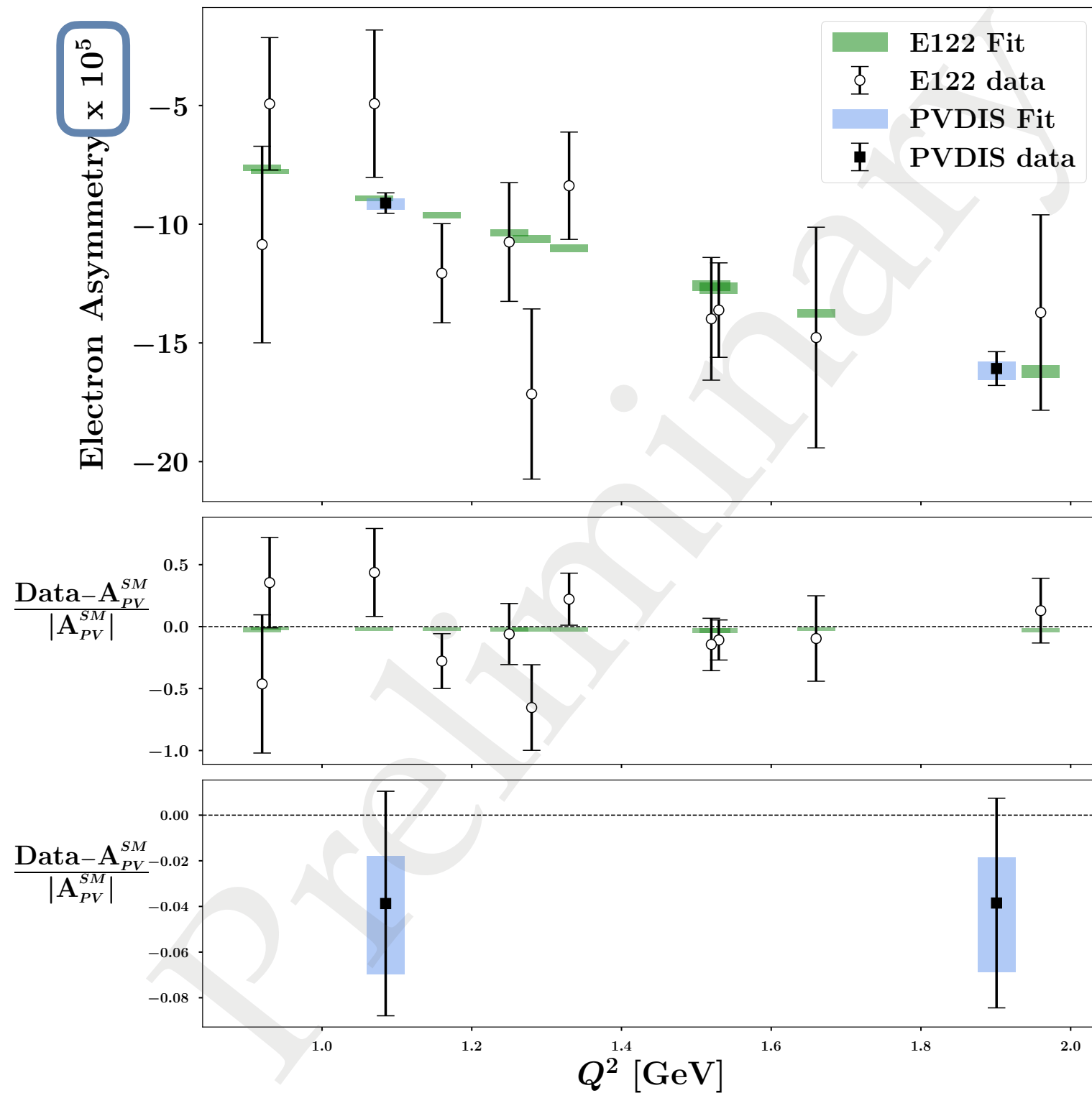


# Results of the fit: data-theory comparison



Sizeable improvement of the fit  
w.r.t. SM predictions

# Results of the fit: data-theory comparison

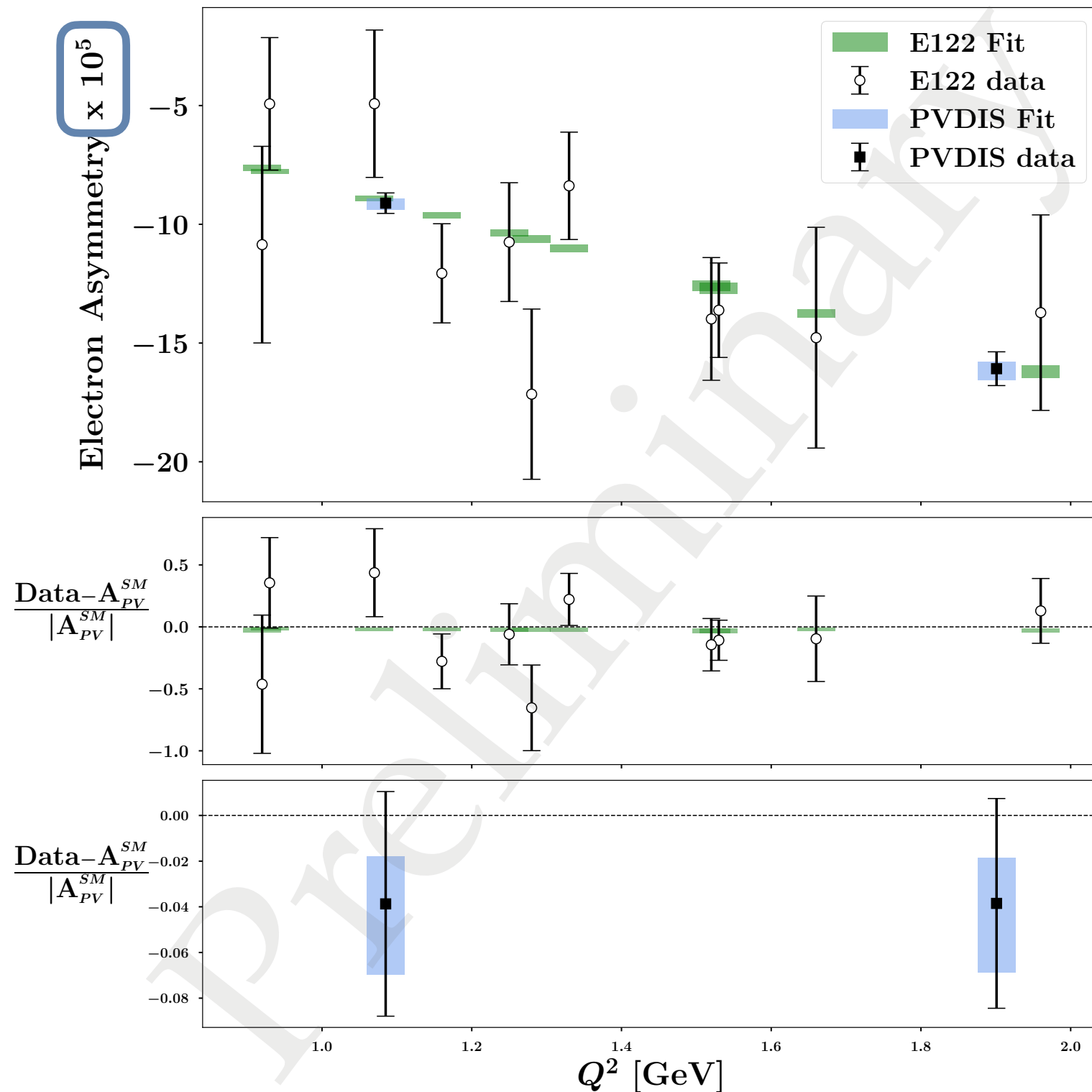


Sizeable improvement of the fit  
w.r.t. SM predictions

Old dataset with still quite large  
experimental errors ( $> 20\%$ )



# Results of the fit: data-theory comparison



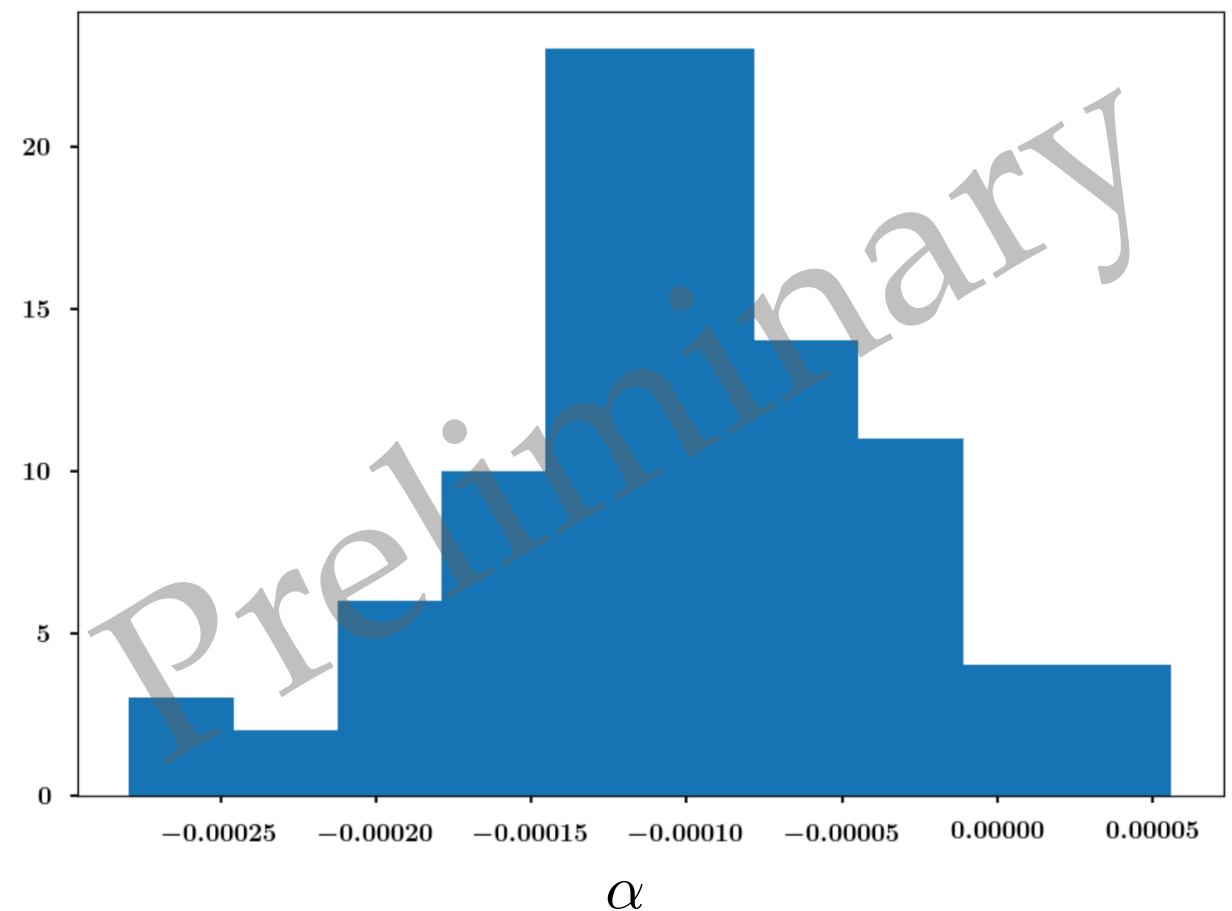
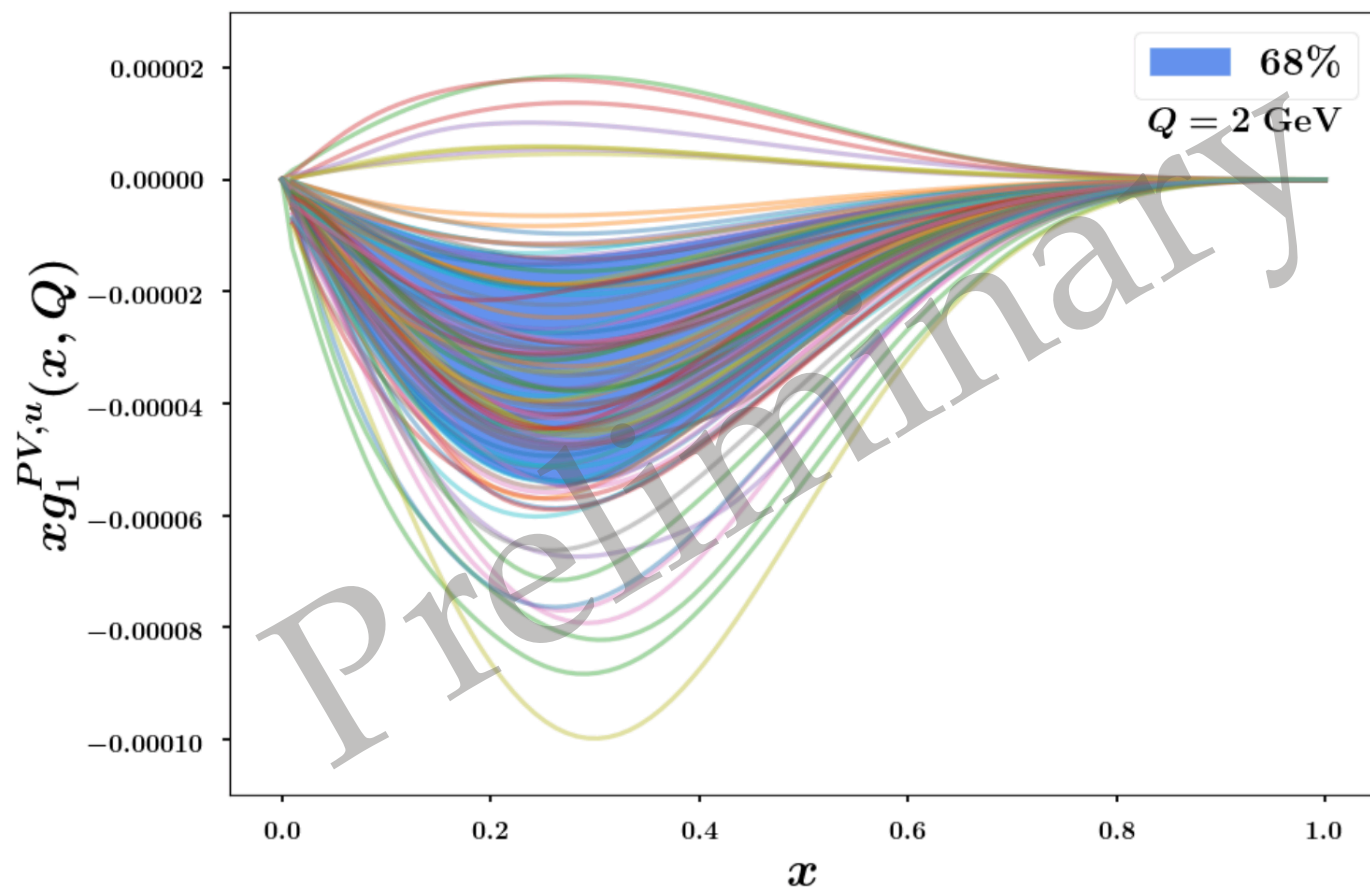
Sizeable improvement of the fit  
w.r.t. SM predictions

Old dataset with still quite large  
experimental errors ( $> 20\%$ )

Data points which actually  
drive the fit due to very small  
experimental errors ( $\sim 1\%$ )

# Results of the fit: $g_1^{PV}(x, Q^2)$ extraction

$$\alpha = (-1.01 \pm 0.66) \cdot 10^{-4}$$



# Conclusions and Outlook

# Conclusions and Outlook

- The strong P- violation can give origin to a new structure function in DIS cross section for one-photon exchange

# Conclusions and Outlook

- The strong P- violation can give origin to a new structure function in DIS cross section for one-photon exchange
- A fit of present experimental data is compatible with a non-zero contribution from a new strong PV parton density at more than 1 sigma

# Conclusions and Outlook

- The strong P- violation can give origin to a new structure function in DIS cross section for one-photon exchange
- A fit of present experimental data is compatible with a non-zero contribution from a new strong PV parton density at more than 1 sigma
- To better investigate its behaviour, new data are needed especially at small (medium) values of  $Q$

# Conclusions and Outlook

- The strong P- violation can give origin to a new structure function in DIS cross section for one-photon exchange
- A fit of present experimental data is compatible with a non-zero contribution from a new strong PV parton density at more than 1 sigma
- To better investigate its behaviour, new data are needed especially at small (medium) values of  $Q$
- Experimental data from positron beam are welcome to shed light on the complementarity with electron beam

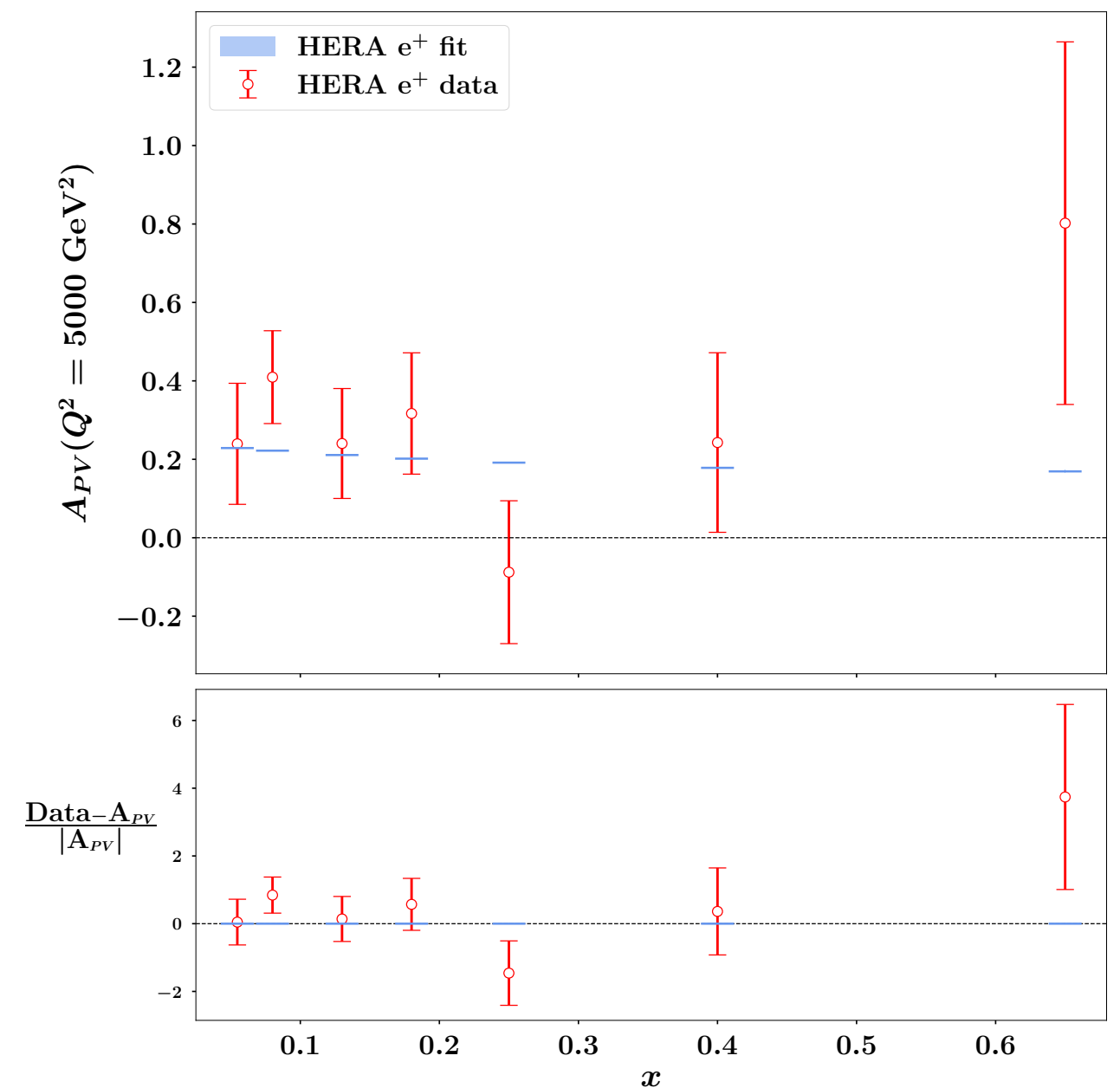
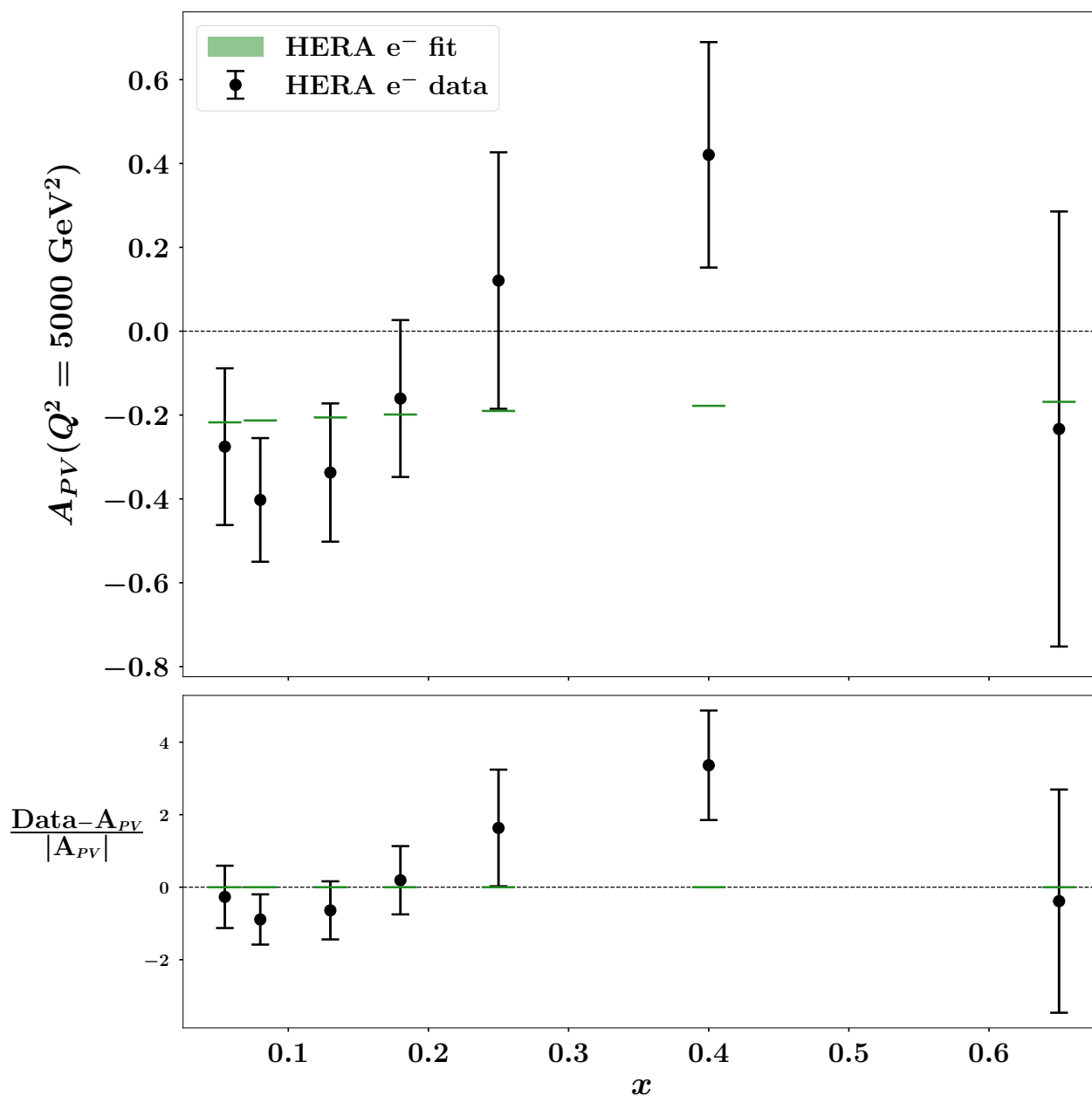
# Conclusions and Outlook

- A different behaviour of the PV parton distribution w.r.t. the variable  $x$  can be investigated



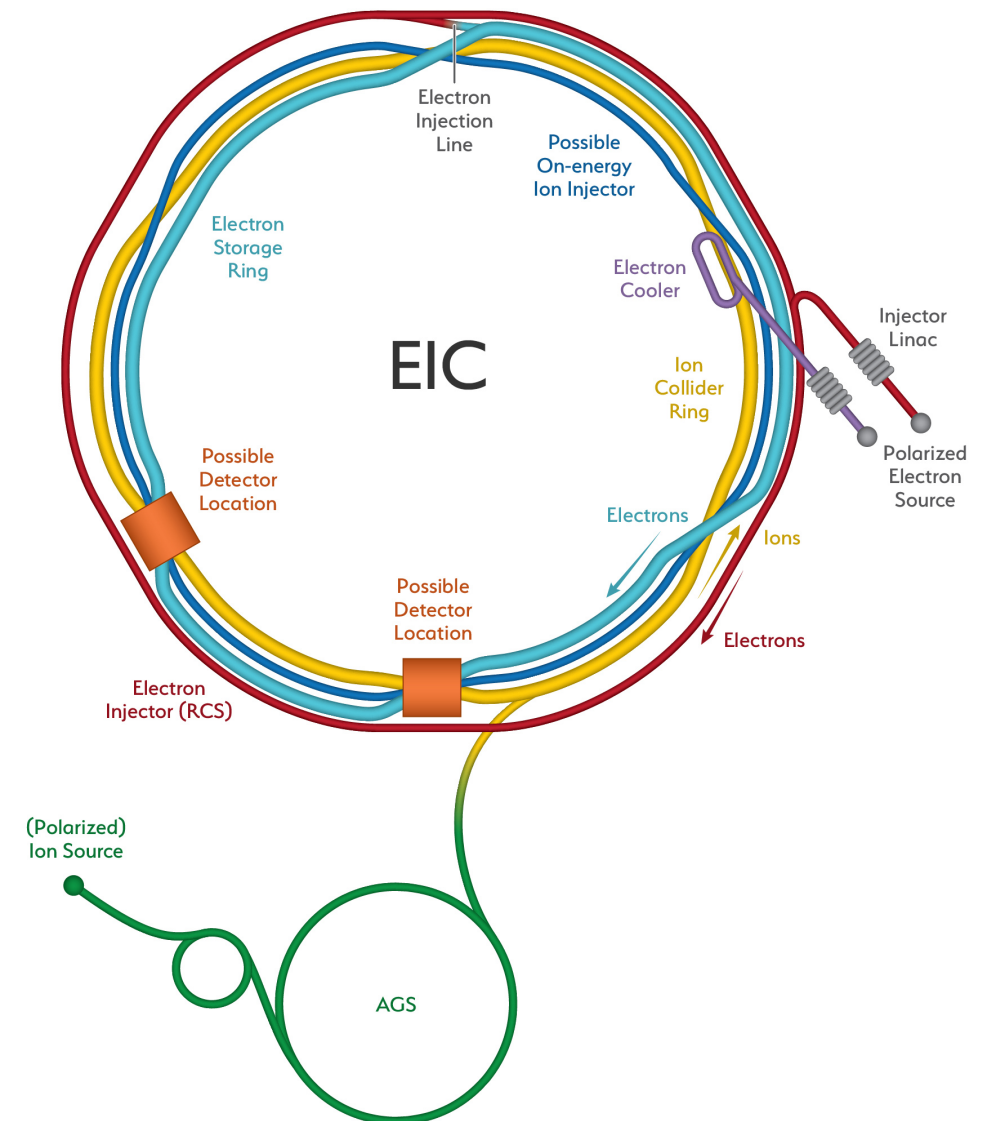
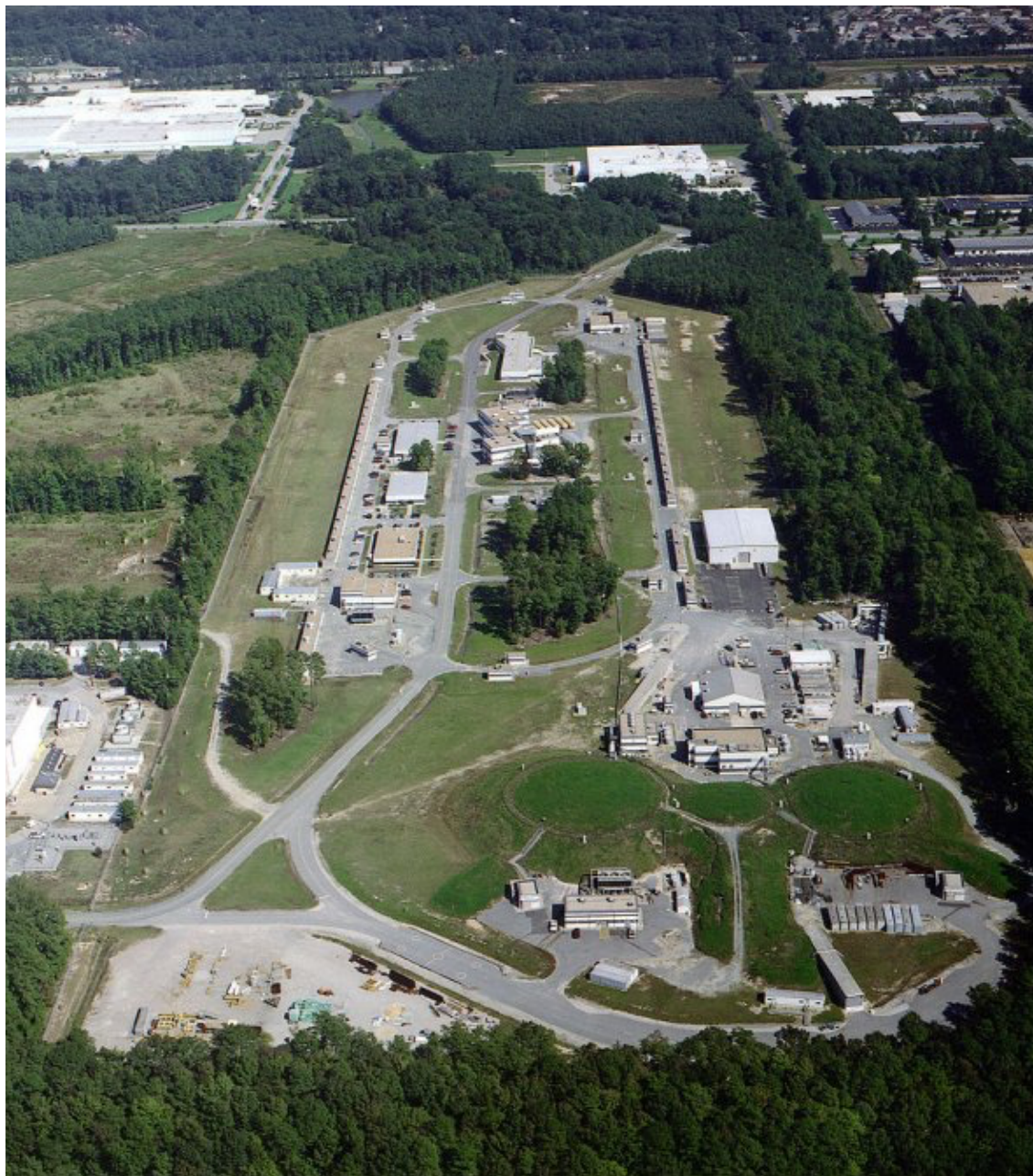
# Conclusions and Outlook

- A different behaviour of the PV parton distribution w.r.t. the variable  $x$  can be investigated



# Conclusions and Outlook

- Predictions of the size of the PV distribution can be made in the kinematic domains of JLab12, JLab20+(?) and EIC



# Conclusions and Outlook

- Further investigations on a new P-odd, CP-odd distribution function arising when considering the polarisation of the target

$$\begin{aligned} \Phi^q(x, Q^2) = & \left\{ f_1^q(x, Q^2) + g_1^{\text{PV}q}(x, Q^2)\gamma_5 \right. \\ & + S_L \left( g_1^q(x, Q^2)\gamma_5 + f_{1L}^{\text{PV}q}(x, Q^2) \right) \\ & \left. - \not{S}_T \left( h_1^q(x, Q^2)\gamma_5 - e_{1T}^{\text{PV}q}(x, Q^2) \right) \right\} \frac{\not{n}_+}{2} \end{aligned}$$

# Conclusions and Outlook

- Further investigations on a new P-odd, CP-odd distribution function arising when considering the polarisation of the target

$$\begin{aligned} \Phi^q(x, Q^2) = & \left\{ f_1^q(x, Q^2) + g_1^{\text{PV}q}(x, Q^2)\gamma_5 \right. \\ & + S_L \left( g_1^q(x, Q^2)\gamma_5 + \boxed{f_{1L}^{\text{PV}q}(x, Q^2)} \right) \\ & \left. - \not{S}_T \left( h_1^q(x, Q^2)\gamma_5 - e_{1T}^{\text{PV}q}(x, Q^2) \right) \right\} \frac{\not{n}_+}{2} \end{aligned}$$

# Conclusions and Outlook

- Further investigations on a new P-odd, CP-odd distribution function arising when considering the polarisation of the target

$$\Phi^q(x, Q^2) = \left\{ f_1^q(x, Q^2) + g_1^{\text{PV}q}(x, Q^2)\gamma_5 \right. \\ \left. + S_L \left( g_1^q(x, Q^2)\gamma_5 + f_{1L}^{\text{PV}q}(x, Q^2) \right) \right. \\ \left. - \not{S}_T \left( h_1^q(x, Q^2)\gamma_5 - e_{1T}^{\text{PV}q}(x, Q^2) \right) \right\} \frac{\not{n}_+}{2}$$

$$\Delta x_B g_5(x_B, Q^2) \approx \Delta x_B g_5^{(\gamma)}(x_B, Q^2) = \frac{1}{2} \sum_q e_q^2 x_B f_{1L}^{\text{PV}(q-\bar{q})}$$