TMD DISTRIBUTIONS AT N4LLL

arXiv:2305.07473 [hep-ph] in collaboration with V. Moos, A. Vladimirov and P. Zurita And JHEP 10 (2022) 118 M. Bury et al.



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TMD FACTORIZATION IP!

$$\frac{d\sigma}{dQ^2 dy dq_T^2} = \sigma_0 \sum_{f_1, f_2} \int \frac{d^2 \boldsymbol{b}}{4\pi} e^{i(\boldsymbol{b} \cdot \boldsymbol{q}_T)} H_{f_1 f_2}(Q, Q) \{ R[\boldsymbol{b}; (Q, Q^2)] \}^2 F_{f_1 \leftarrow h_1}(x_1, \boldsymbol{b}) F_{f_2 \leftarrow h_2}(x_2, \boldsymbol{b}) \}$$

In recent years we have learnt a lot about this formula. For instance: Its range of applicability is provided by $\delta = \frac{q_T}{Q} \ll 1$, fixed- q_T , $\delta \sim 0.25$ We have a non-perturbative evolution kernel (whose perturbative part is known at N3LO!!). We can work with different schemes (CSS, ζ -prescription, ..). We have a re-factorization of TMD at large transverse momentum in Wilson coefficients (now at N3LO!!) and PDF (now at NNLO!!) PDF are just part of a model



TMD FACTORIZATION

We can : Perform an extraction of TMD at N4LL (higher order than PDF..) Analyze the source of errors Be ready for NLP corrections

> In this talk I will consider the first two points We call the new Artemide code extraction



In SV19 we tried with several PDF sets

PDF set	χ^2_{DY}/N_{pt}
CT14	1 50
C114	1,37
HERAPDF2.0	0,97
MMHT14	1,34
NNPDF3.1	1,14
PDF4LHC15	1,53

SV19

2.5

u-quark

Also, in SV19, for $b \rightarrow 0$, the uncertainty bands $\rightarrow 0$.

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THE PDF BIAS

- So we have some questions to answer:
 - 1. Can we get good TMD fits for different collinear PDFs?
 - 2. Would they have sensible uncertainty bands?
 - 3. Would they the consistent with each other?

FLAVOR INDEPENDENT 7NP

Most of replicas (64%) have $\chi^2/N>2$. Each replica has a peculiar shape



PDF UNCERTAINTIES AND FLAVOUR DEPENDENCE

M. Bury, F. Hautmann, S. Leal-Gomez, I. Scimemi, A. Vladimirov, PZ, JHEP 10 (2022) 118

Flavor separation make fits more PDF set independent and modeling simpler

$$f_{NP}^{f}(x,b) = \exp\left(-\frac{\lambda_{1}^{f}(1-x) + \lambda_{2}^{f}x}{\sqrt{1+\lambda_{0}x^{2}\mathbf{b}^{2}}}\mathbf{b}^{2}\right) \qquad f = u, \,\bar{u}, \,d, \,\bar{d}, \,sea$$



FINAL χ^2 : MSHT20 (1.12), HERA20 (0.91), NNPDF31(1.21), CT18 (1.08)

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We include the PDF uncertainties while keeping f_{NP} fixed.

We re-fit TMD, for each PDF replica.

We get reasonable uncertainty bands.



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The TMD obtained from different sets agree reasonably





- Flavor dependence.
- All the latest datasets!
- W-boson production!
- **Increased perturbative accuracy! (** N^4LL **)**
- Includes collinear PDF uncertainties!
- A full new fit to Drell-Yan data.

ART23: DETAILS

Evolution:

- **We use the** ζ **prescription** (I.S., A. Vladimirov *JHEP* 08 (2018) 003)
- We use the integral form of the evolution kernel to introduce a scale dependence similar to CSS for direct comparison

$$\mathscr{D}(b,\mu) = \mathscr{D}_{\text{small-b}}(b^*,\mu^*) + \int_{\mu^*}^{\mu} \frac{d\mu'}{\mu'} \Gamma_{\text{cusp}}(\mu') + \mathscr{D}_{\text{NP}}(b) \qquad b^*(b) = \frac{b}{\sqrt{1 + \frac{b^2}{B_{\text{NP}}^2}}} = \frac{2e^{-\gamma_E}}{\mu^*}$$

We discover that we are sensitive to log corrections to the NP part of the evolution kernel $\mathscr{D}_{NP}(b) = bb^* \left[c_0 + c_1 \ln \left(\frac{b^*}{B_{NP}} \right) \right]$

ART23: DETAILS

Simple Parameterization:
$$f_{NP}^{f}(x,b) = \frac{1}{\cosh\left(\left(\lambda_{1}^{f}(1-x)+\lambda_{2}^{f}x\right)b\right)}$$

 $f = u, \bar{u}, d, \bar{d}, sea$

- In total, 13 parameters
- Reference PDFs: MSHT20



New in!

- PHENIX: DY data at $\sqrt{s} = 200 \text{ GeV}$ STAR: Z/ γ -boson production at $\sqrt{s} = 510 \text{ GeV}$ (preliminary). CMS and LHCb: ydifferential Z-boson production at $\sqrt{s} = 13$ TeV. ATLAS: high precision differential Z-boson cross
 - section.
- CMS: high-Q neutralboson production.
- Tevatron: W-boson production.

ART23: DETAILS

- Fitting procedure: construct simultaneous replicas of the data
 AND the PDFs. Then fit.
- The number of replicas needed to have a faithful representation of the TMDPDF distribution was deemed to be 1000.



- overall improvement w.r.t. SV19. Specially for the LHC data.
 Higher precision plays a key role here.
- more realistic uncertainty bands than in SV19.



Scale variations $\left\{ \mu \to s_2 \mu, \mu^* \to s_3 \mu^*, \mu_{OPE} \to s_4 \frac{2e^{-\gamma_E}}{b} + 2 \text{GeV} \right\}.$













ART23:	RESULTS
$\lambda_1^u = 0.87^{+0.10}_{-0.10},$	$\lambda_2^u = 0.91_{-0.29}^{+0.33},$
$\lambda_1^d = 0.99^{+0.09}_{-0.12},$	$\lambda_2^d = 6.06^{+1.36}_{-1.34},$
$\lambda_1^{\bar{u}} = 0.35_{-0.22}^{+0.23},$	$\lambda_2^{\bar{u}} = 46.6^{+7.9}_{-8.1},$
$\lambda_1^{\bar{d}} = 0.12^{+0.13}_{-0.11},$	$\lambda_2^{\bar{d}} = 1.53^{+0.54}_{-0.17},$
$\lambda_1^{sea} = 1.32_{-0.24}^{+0.23},$	$\lambda_2^{sea} = 0.46^{+0.13}_{-0.45},$





green: ART23 grey: SV19

SUMMARY

- We have performed a novel TMDPDF extraction: *ART23*.
- We used all the newest measurements and also W-boson production data, finding a good description.
- For the first time, the PDF uncertainties are systematically included. And we have realistic uncertainty bands.
- The flavor dependence in the NP ansatz is crucial to reduce the PDF bias.
- The global fit (including SIDIS data) is ... closer



NNPDF3.1: COMPARISON

