

Polarised physics at the LHC: the $L_{spin}^{\uparrow}C$ project

Pasquale Di Nezza









The LHC beams cannot be polarised

The only possibility to have polarised collisions is through a polarised fixed-target

Collisions provided by a TeV-scale beam (LHC) on fixed target will exploit a unique kinematic region poorly probed. Advanced detectors make available probes never accessed before



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The LHCb detector

- LHCb is a general-purpose forward spectrometer, fully instrumented in $2 < \eta < 5$, and optimised for *c* and *b* hadron detection
- Excellent momentum resolution with VELO + tracking stations:

 $\sigma_p/p = 0.5 - 1.0 \% \ (p \in [2,200] \text{ GeV})$

• Particle identification with RICH+CALO+MUON

 $\epsilon_{\mu} \sim 98 \%$ with $\epsilon_{\pi \to \mu} \lesssim 1 \%$

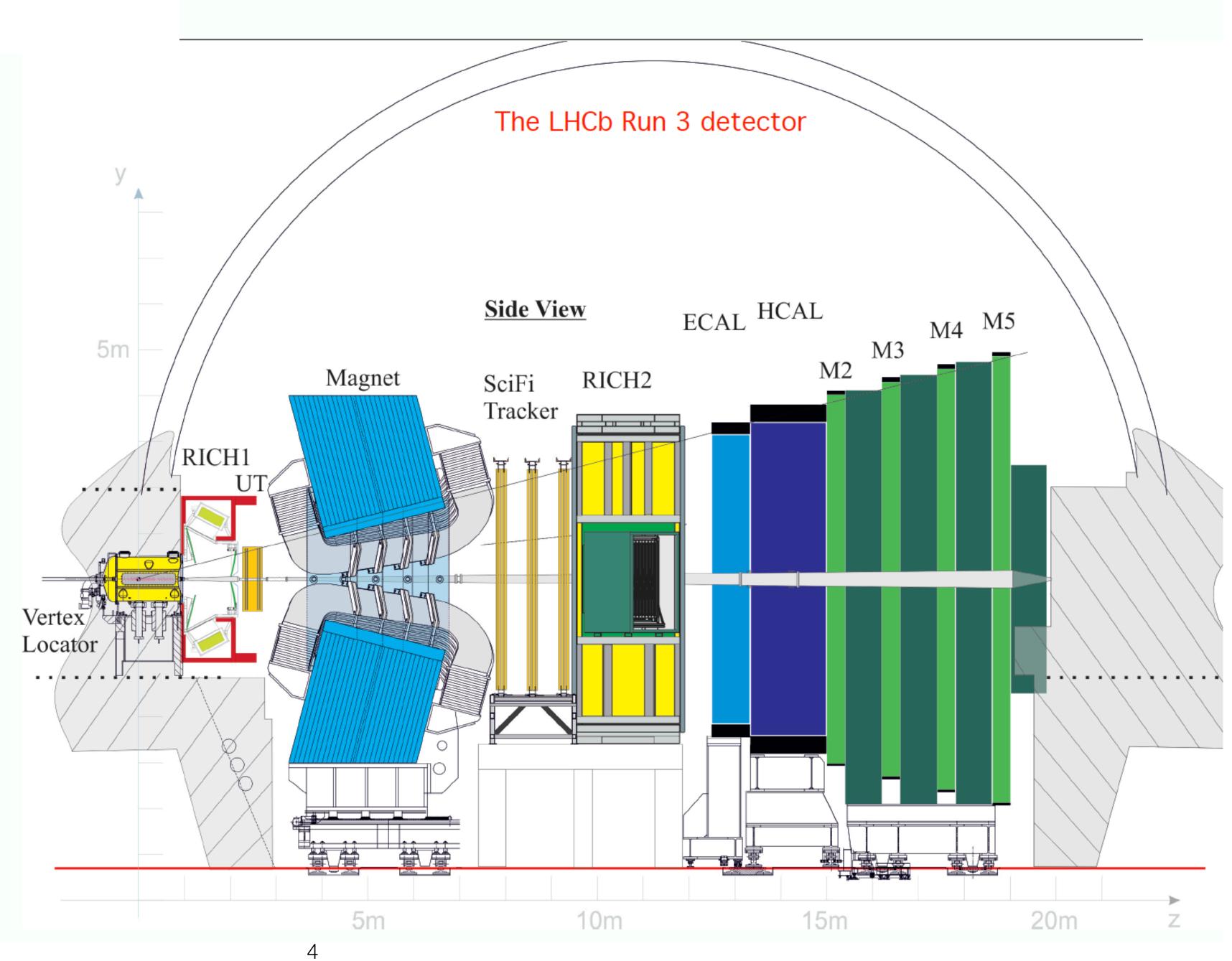
• Low momentum muon trigger:

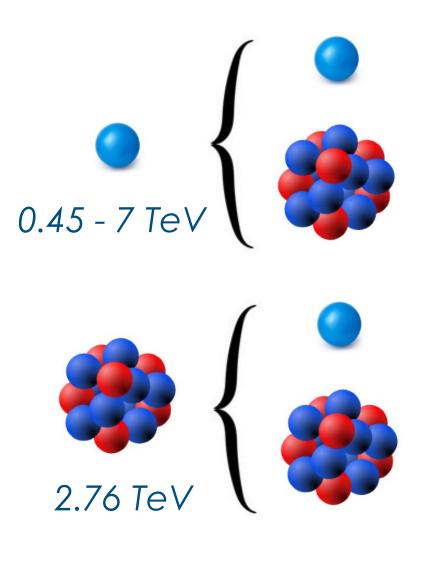
 $p_{T_u} > 1.75 \text{ GeV} (2018)$

will be reduced thanks to the new fullysoftware trigger

• Major detector upgrades performed during LS2 for the Run 3 (5x luminosity)

[<u>JINST 3 (2008) S08005</u>] [<u>IJMP A 30, 1530022 (2015)</u>] [<u>Comput Softw Big Sci 6, 1 (2022)</u>]

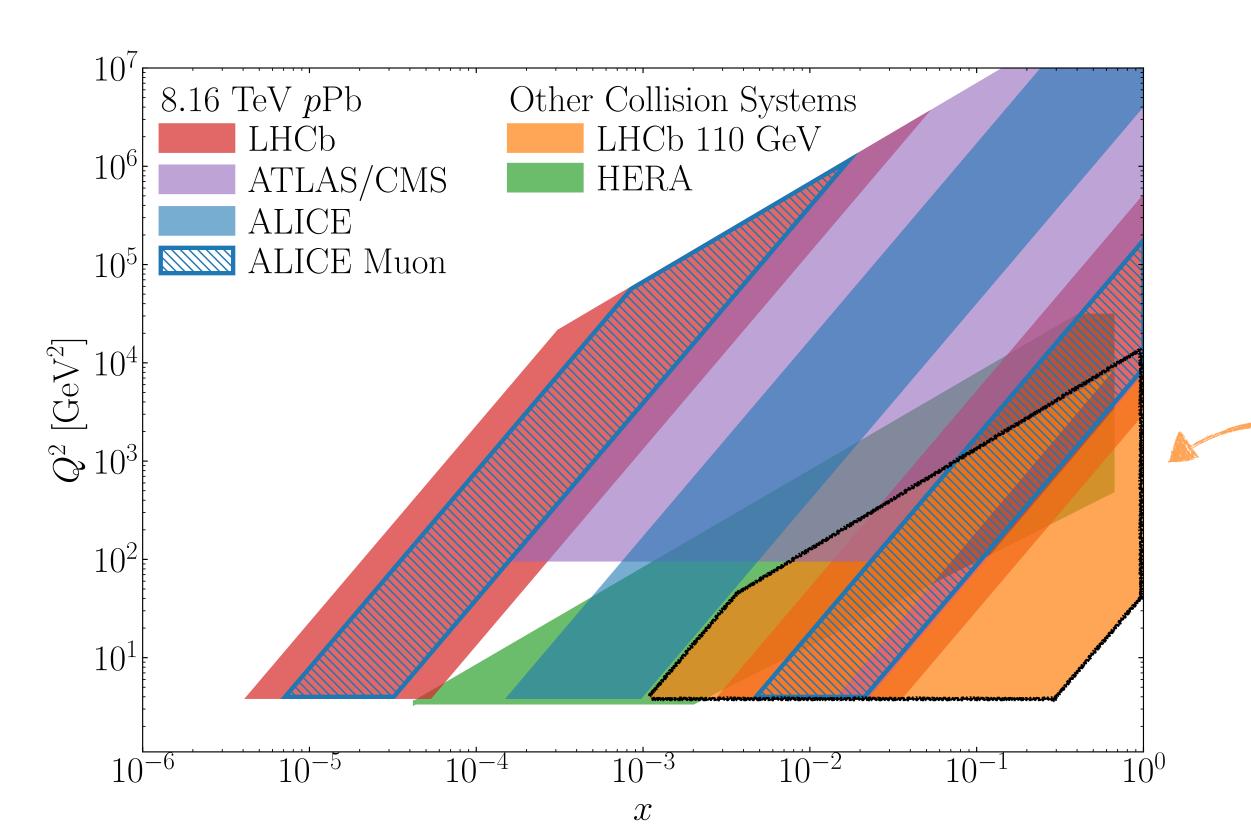




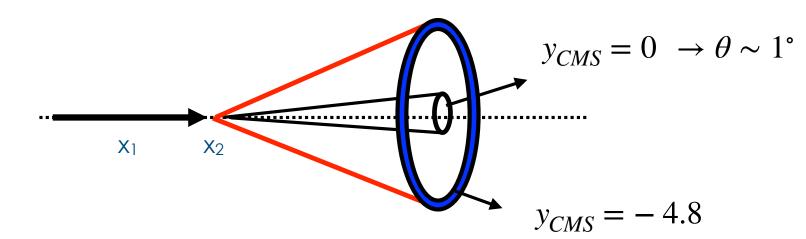
pp or pA collisions: 0.45 - 7 TeV beam on fix target $\sqrt{s} = \sqrt{2m_N E_p} \simeq 41 - 115 \ GeV$ $y_{CMS} = 0 \rightarrow y_{lab} = 4.8$

AA collisions: 2.76 TeV beam on fix target $\sqrt{s_{NN}} \simeq 72 \ GeV$

 $y_{CMS} = 0 \rightarrow y_{lab} = 4.3$



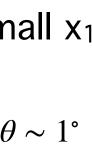
1: beam; 2: target Large CM boost, large x_2 values ($x_F < 0$) and small x_1



$$\gamma = \frac{\sqrt{s_{NN}}}{2m_p} \simeq 60$$

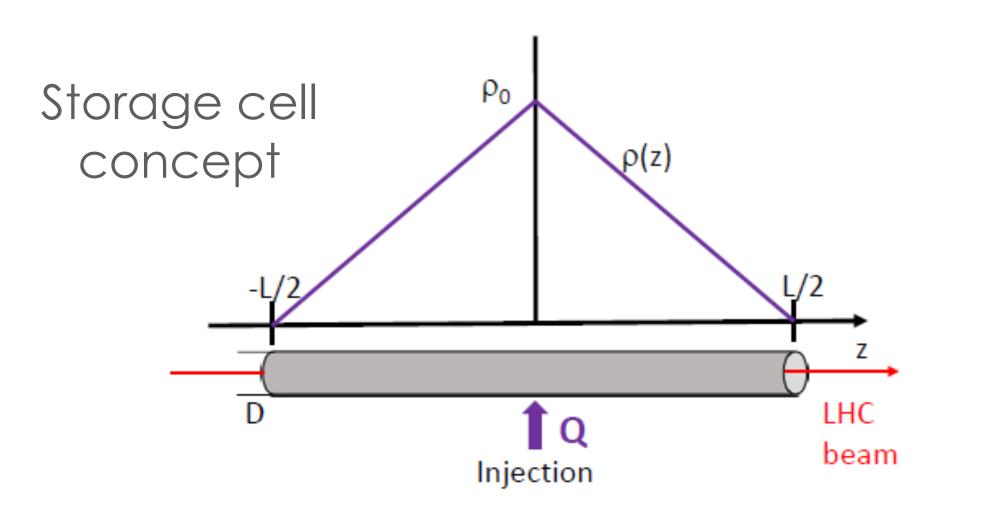
Broad and poorly explored kinematic range

mid-to-large x_{Bj} at intermediate Q^2 and negative x_F

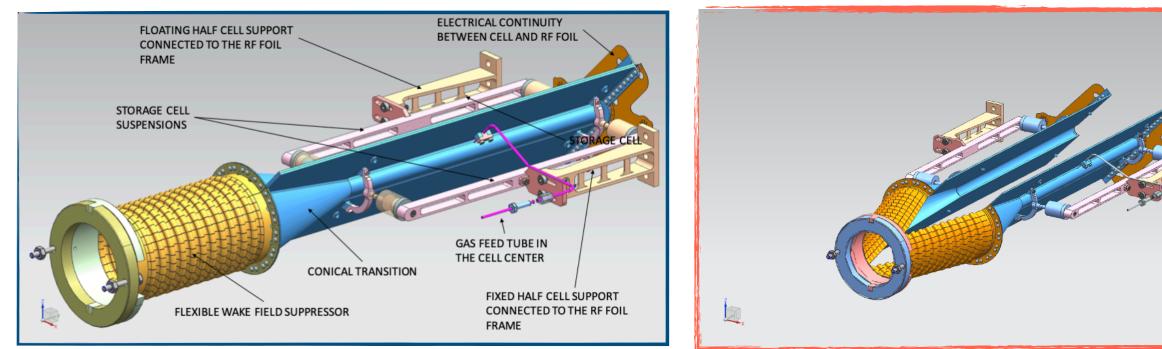




SMOG2 an unpolarised target at

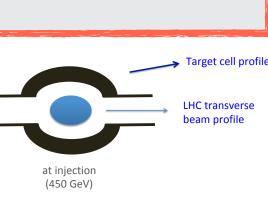


Openable cell





5 mm radius x 200 mm length



6

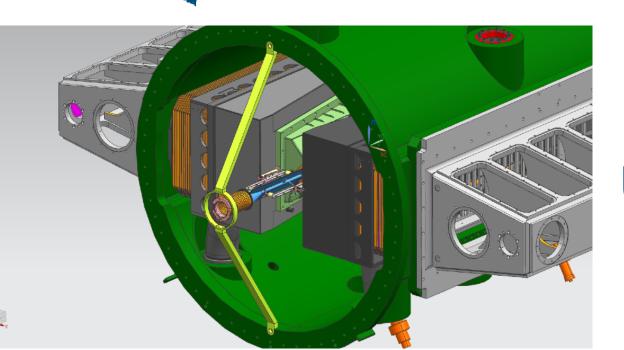
LHC beam



Forward acceptance: $2 < \eta < 5$

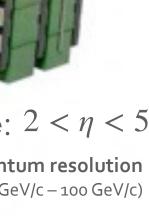
Tracking system momentum resolution $\Delta p/p = 0.5\% - 1.0\% (5 \text{ GeV/c} - 100 \text{ GeV/c})$

beam-beam collisions



UNpolarised target (beam-gas)

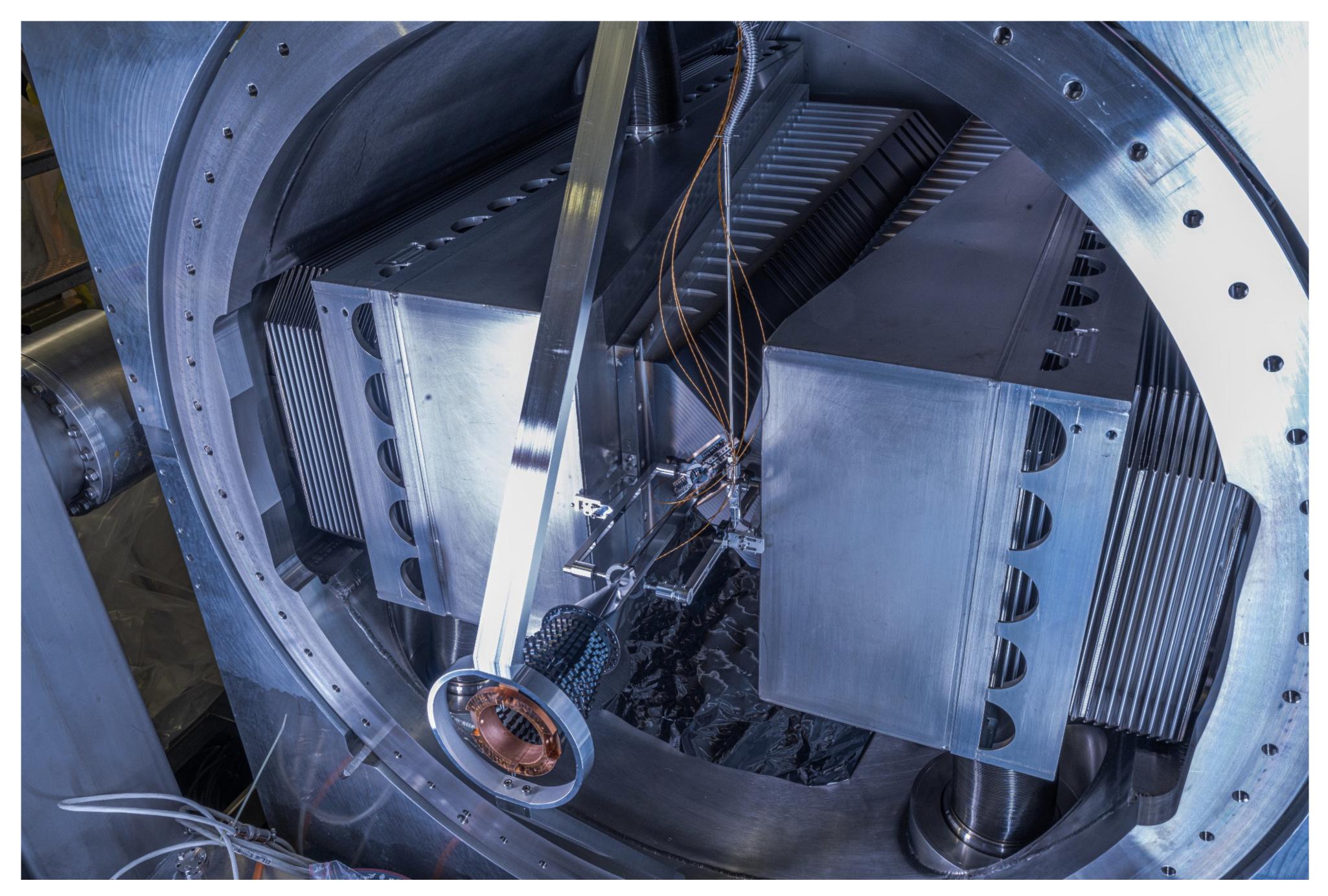
JINST 3 (2008) S08005 IJMPA 30 (2015) 1530022





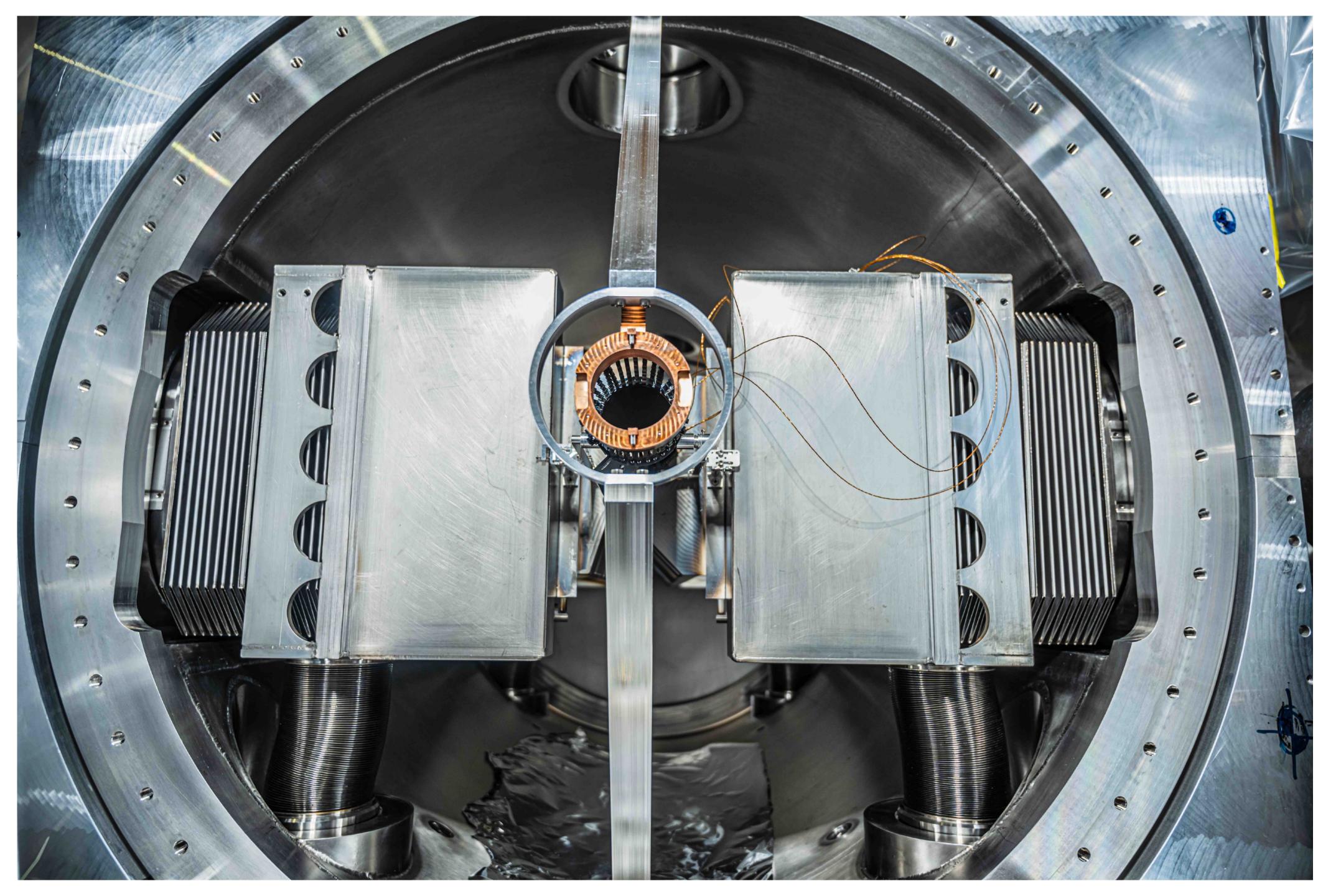


It is the only system present in the LHC primary vacuum



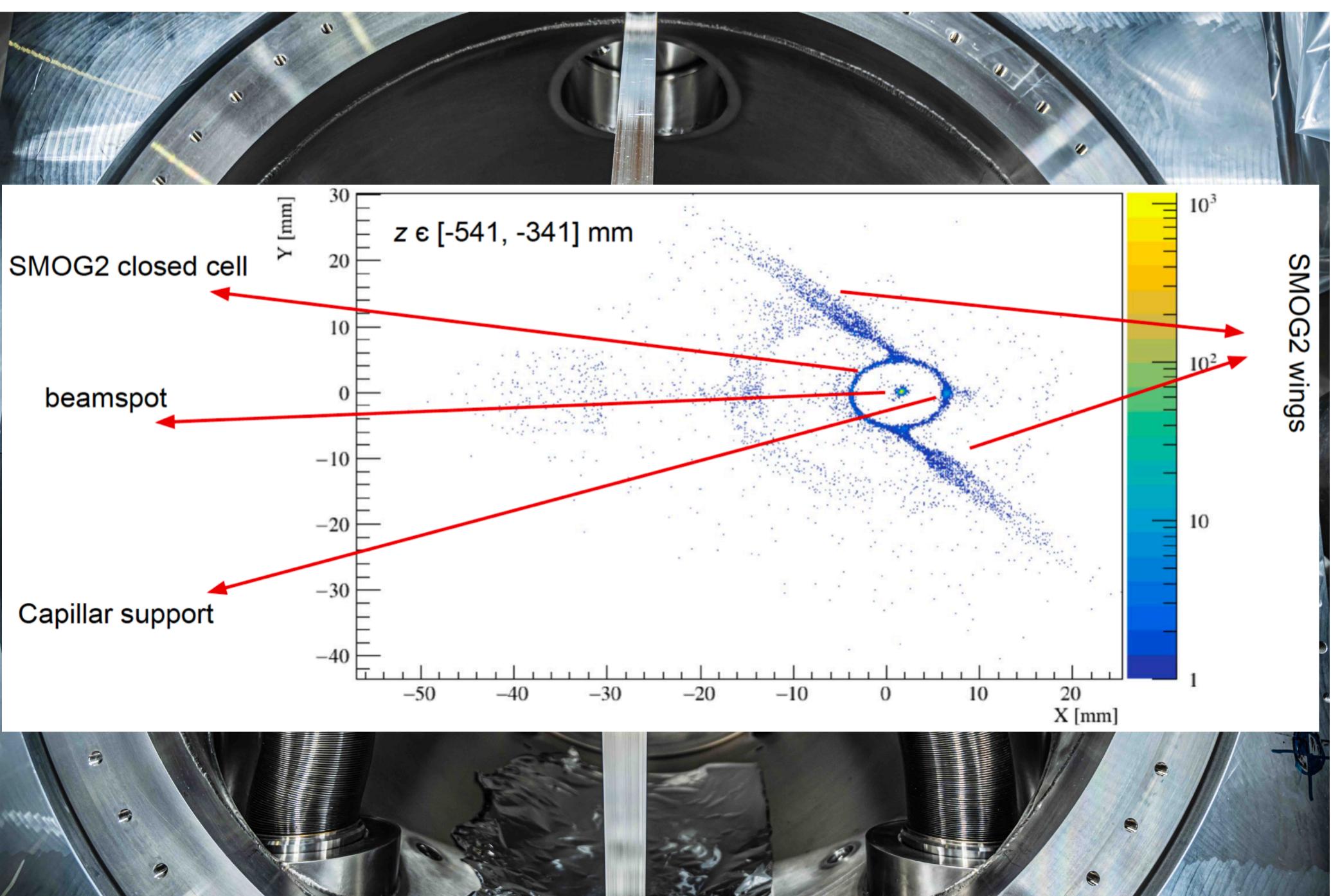


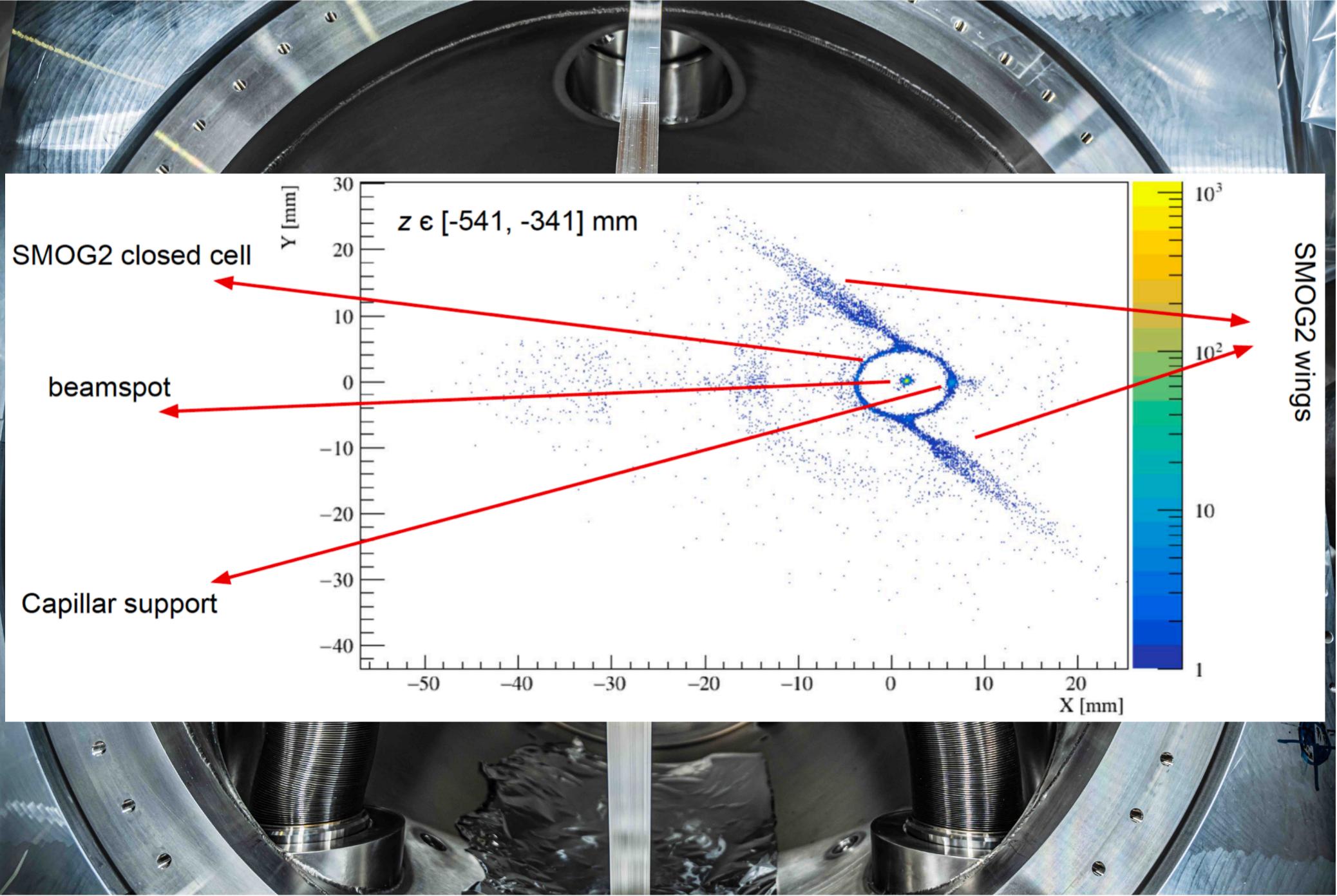
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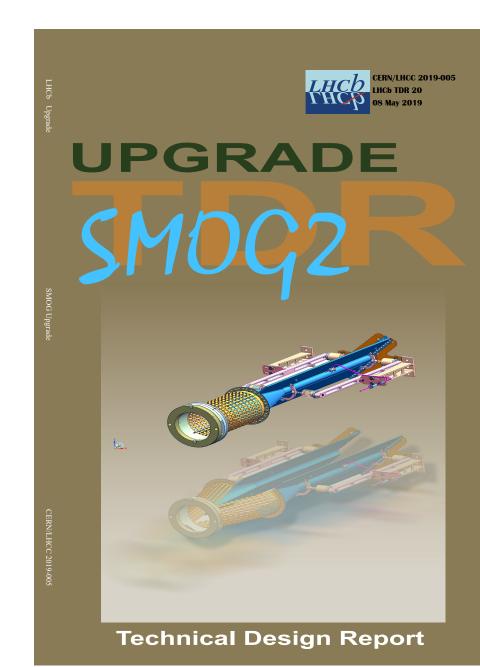


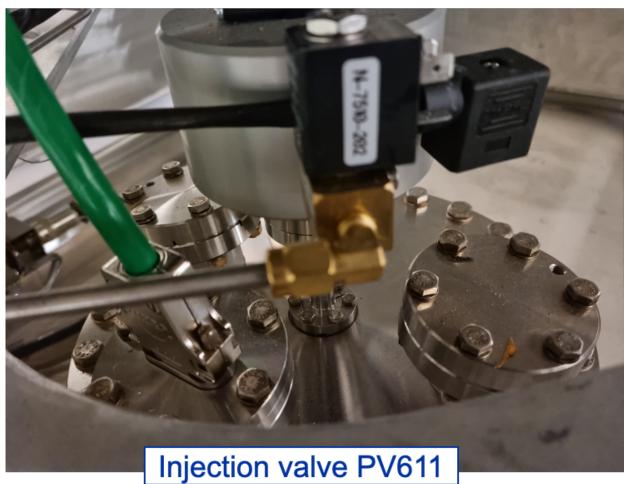
- reconstruction)
- Injectable gases (3+1 reservoirs): H₂, D₂, N₂, O₂, He, Ne, Ar, Kr, Xe
- Flux known with 1% precision, measured relative contamination 10-4



• The system is completely installed (storage cell + GFS + triggers +

• Negligible impact on the beam lifetime ($\tau_{beam-gas}^{p-H_2} \sim 2000 \text{ days}$, $\tau_{beam-gas}^{Pb-Ar} \sim 500 \text{ h}$)

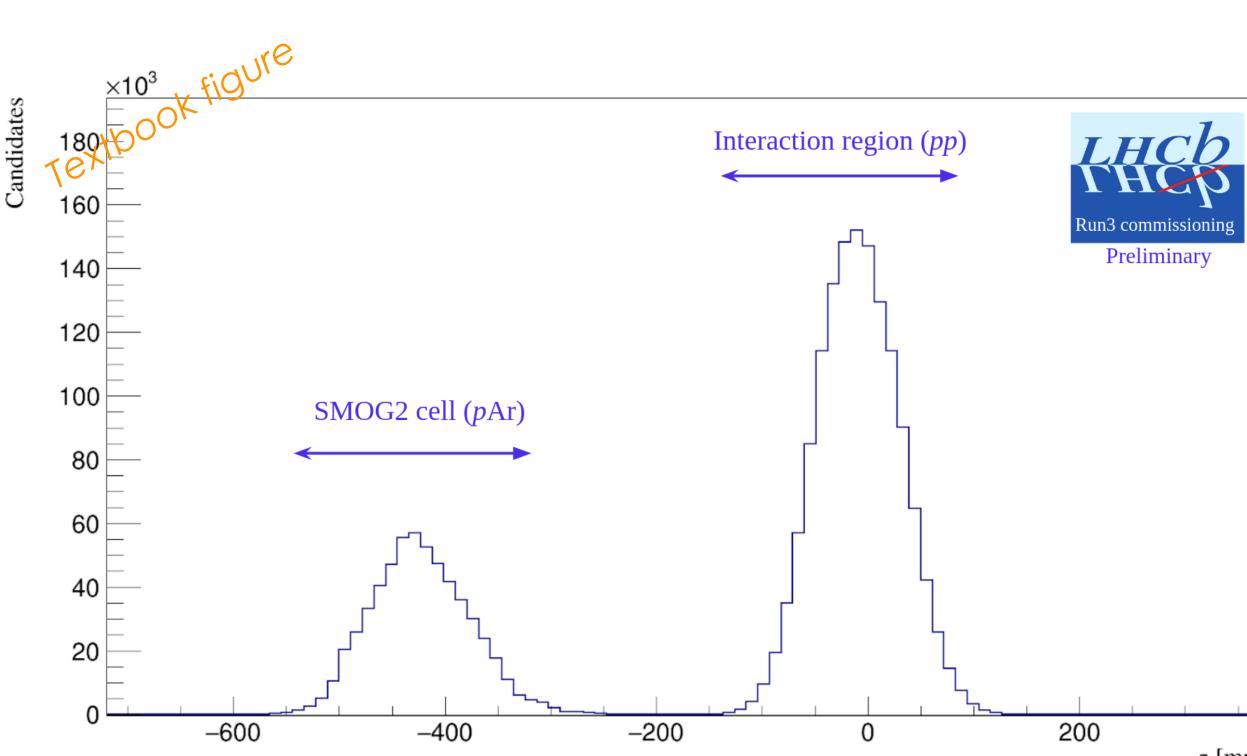




https://cds.cern.ch/record/2673690/



SMOC2 works!



we are able to reconstruct 2 well separated and independent Interaction Points

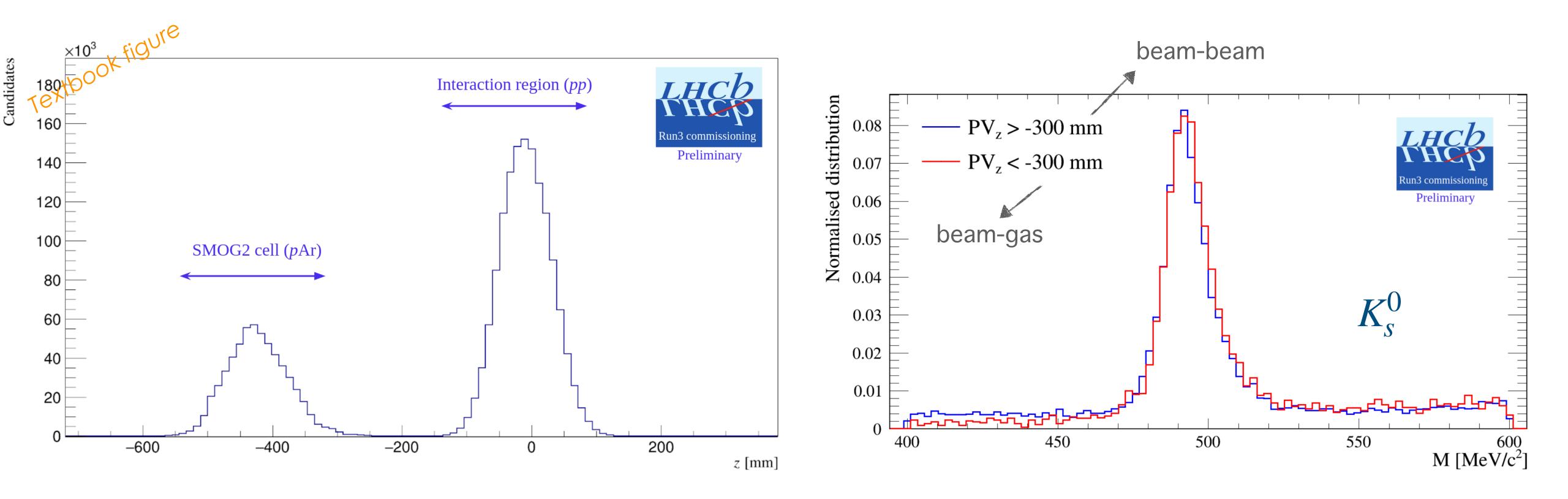
LHCb-FIGURE-2023-001

z [mm]



11

SMOC2 works!



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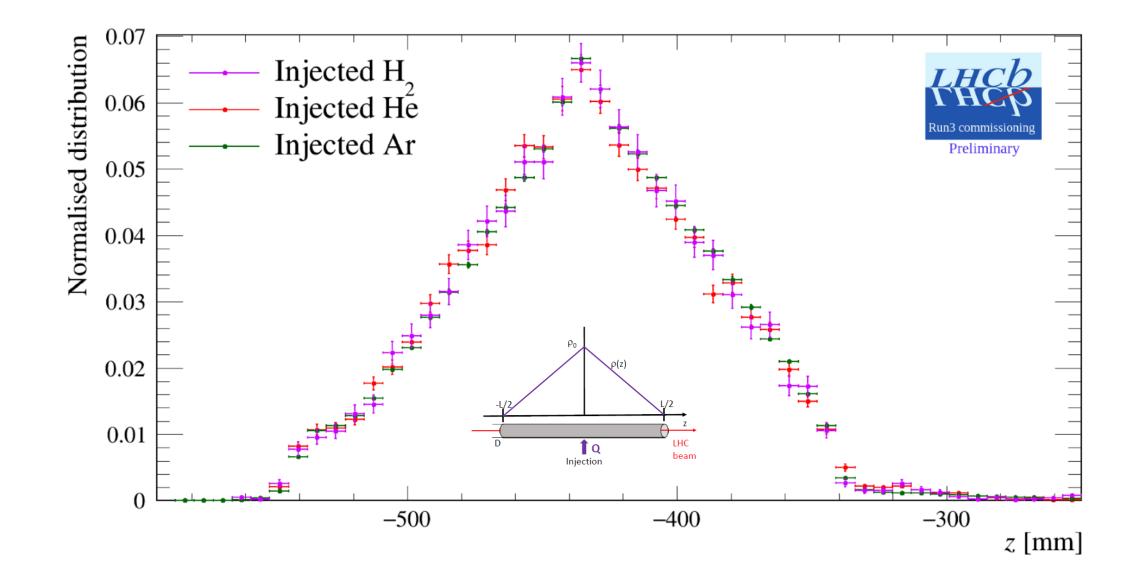
LHCb-FIGURE-2023-001

same resolution for beam-gas and beam-beam collisions

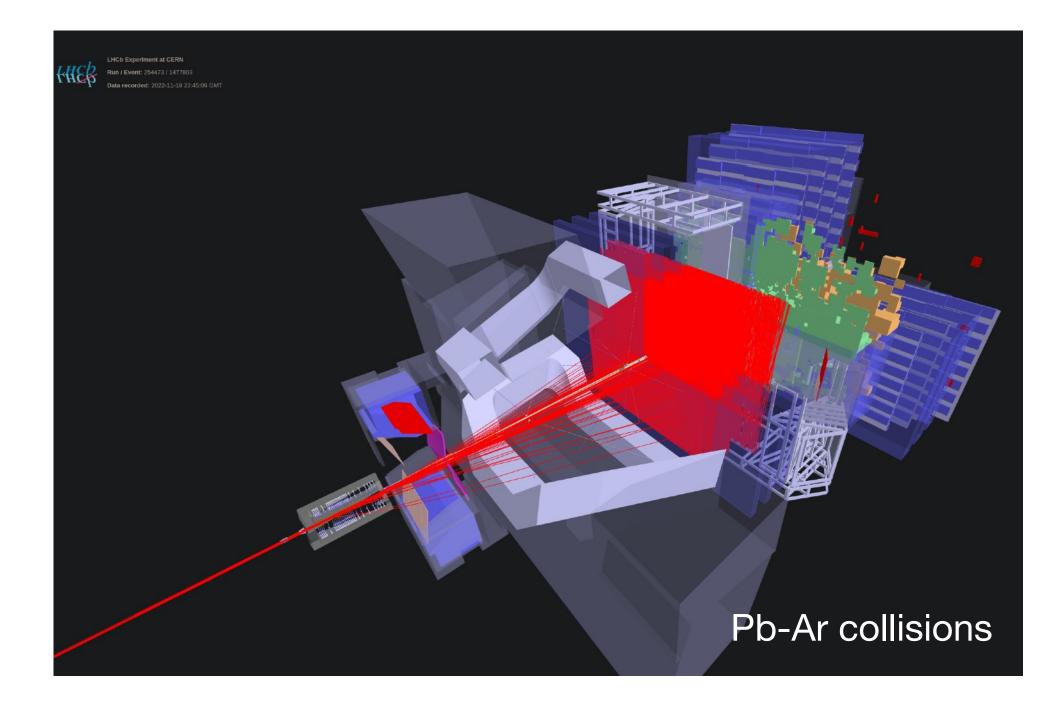




SMOG2 works!



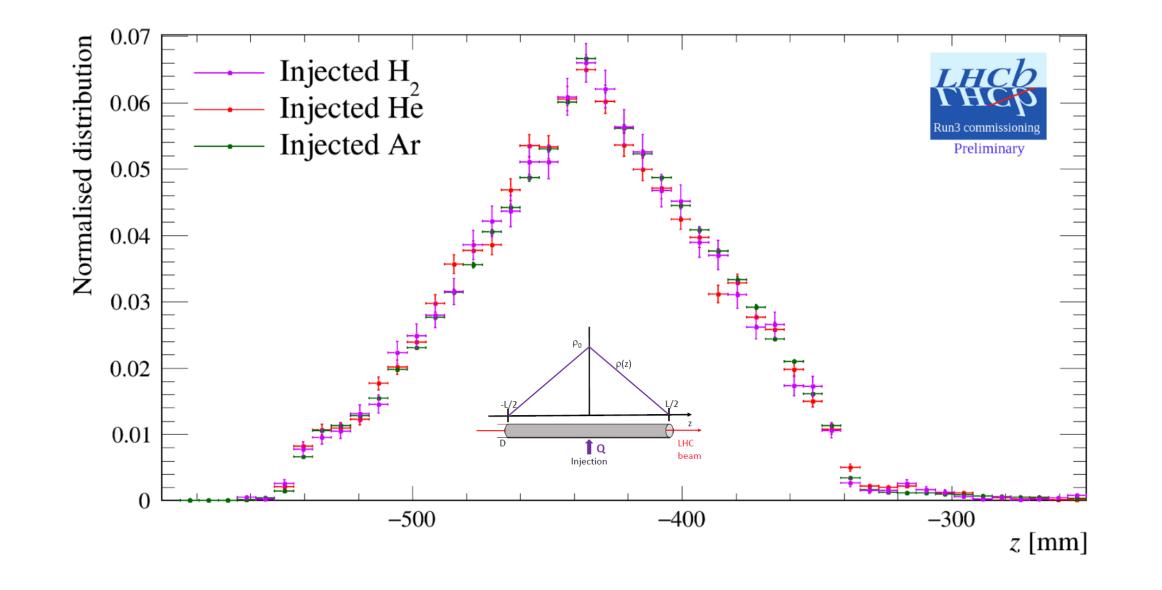
LHCb-FIGURE-2023-001

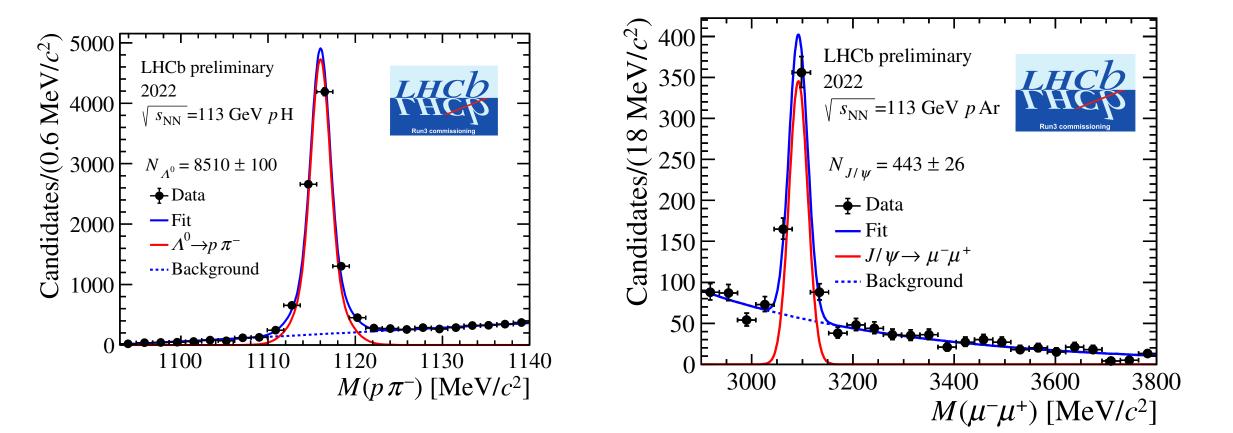




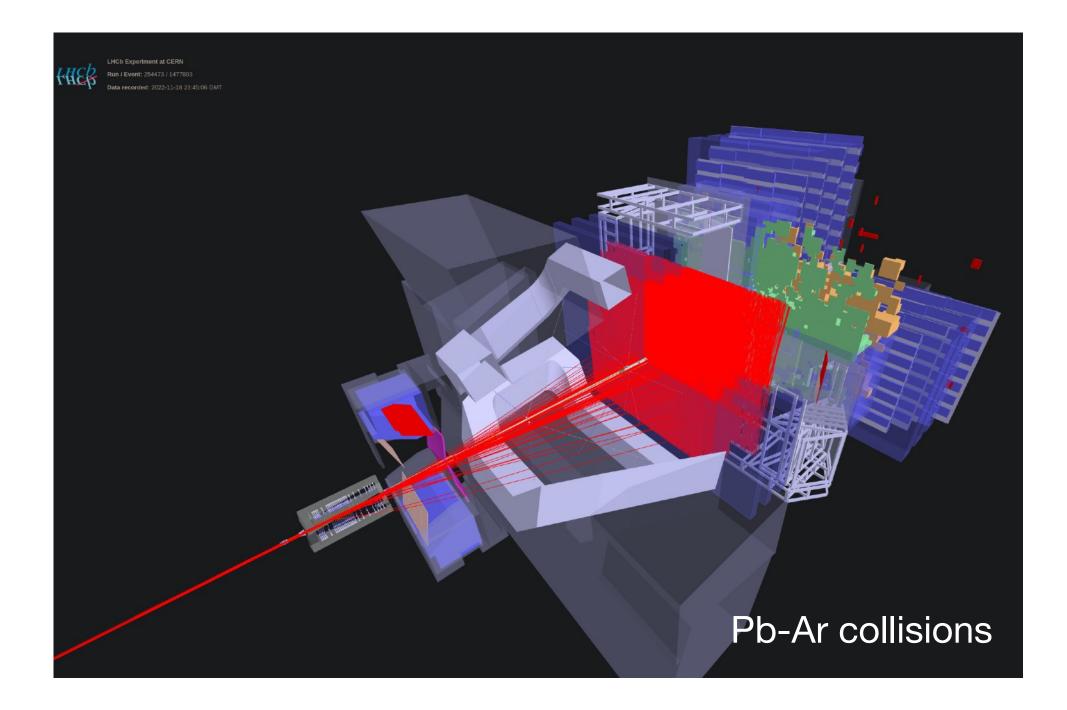


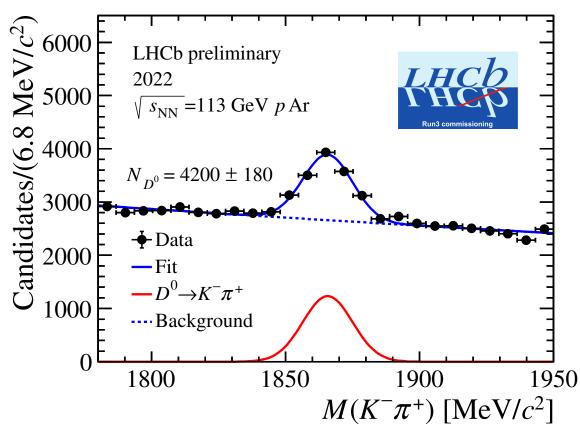
SMOC2 works!





LHCb-FIGURE-2023-001





LHCb-FIGURE-2023-008

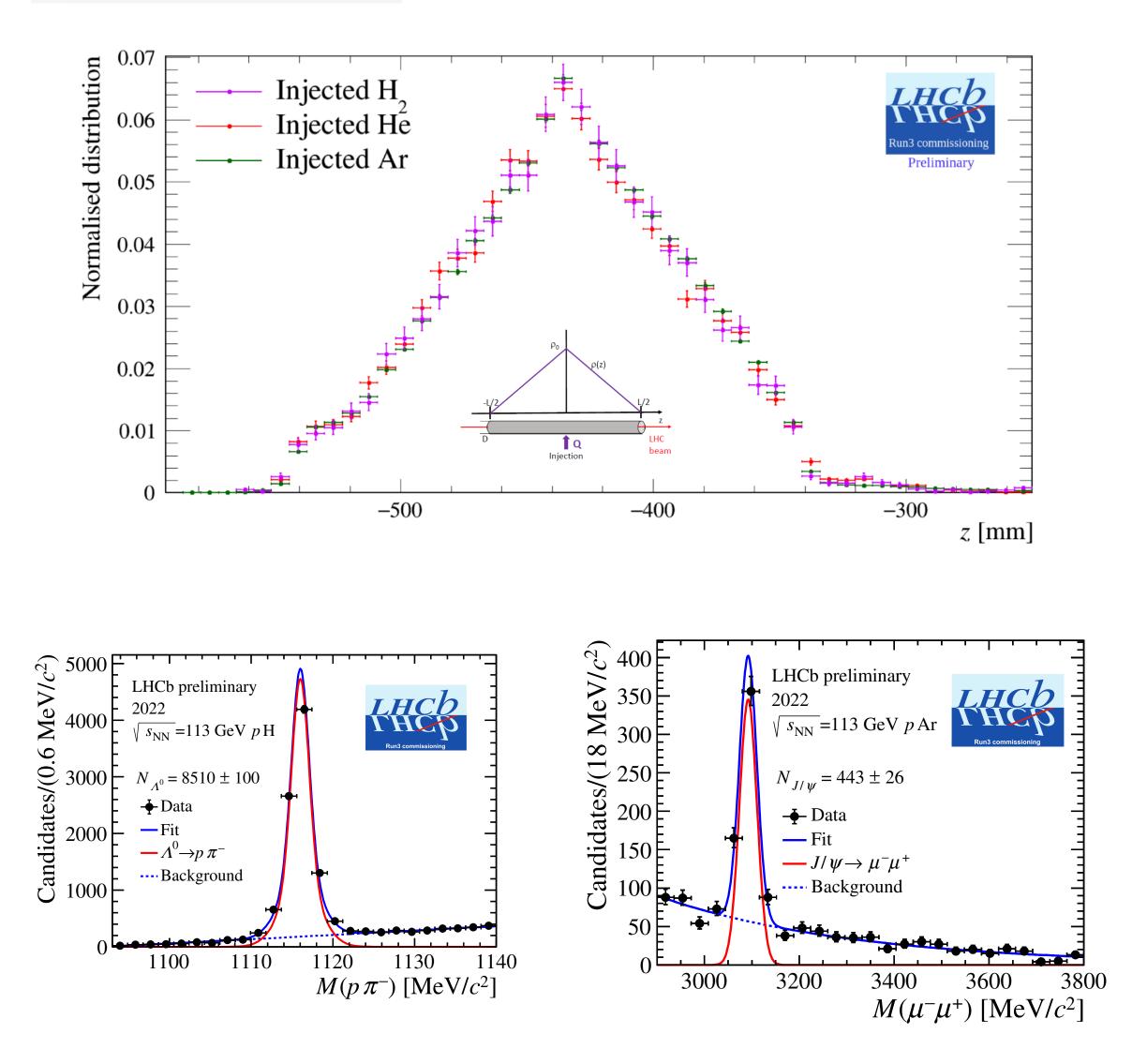
in ~10' of data taking





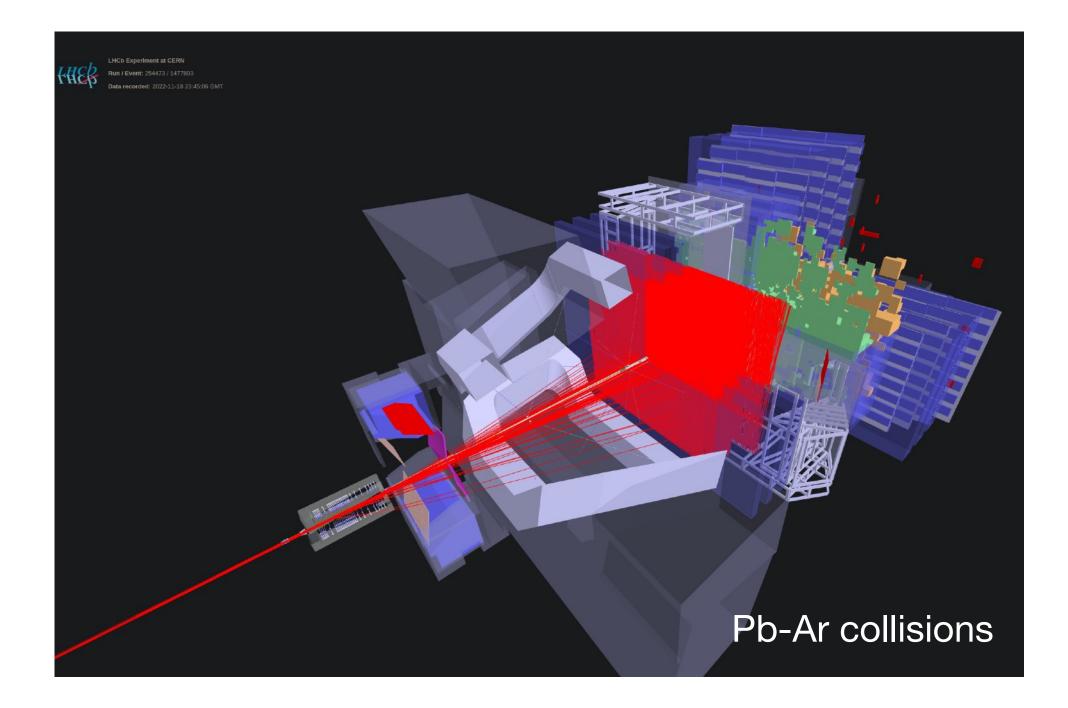


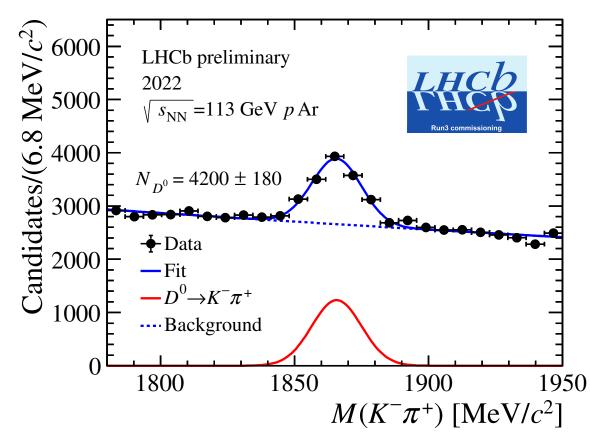
SMOC2 works!



LHCb is the only experiment able to run in collider and fixed-target mode simultaneously!

LHCb-FIGURE-2023-001





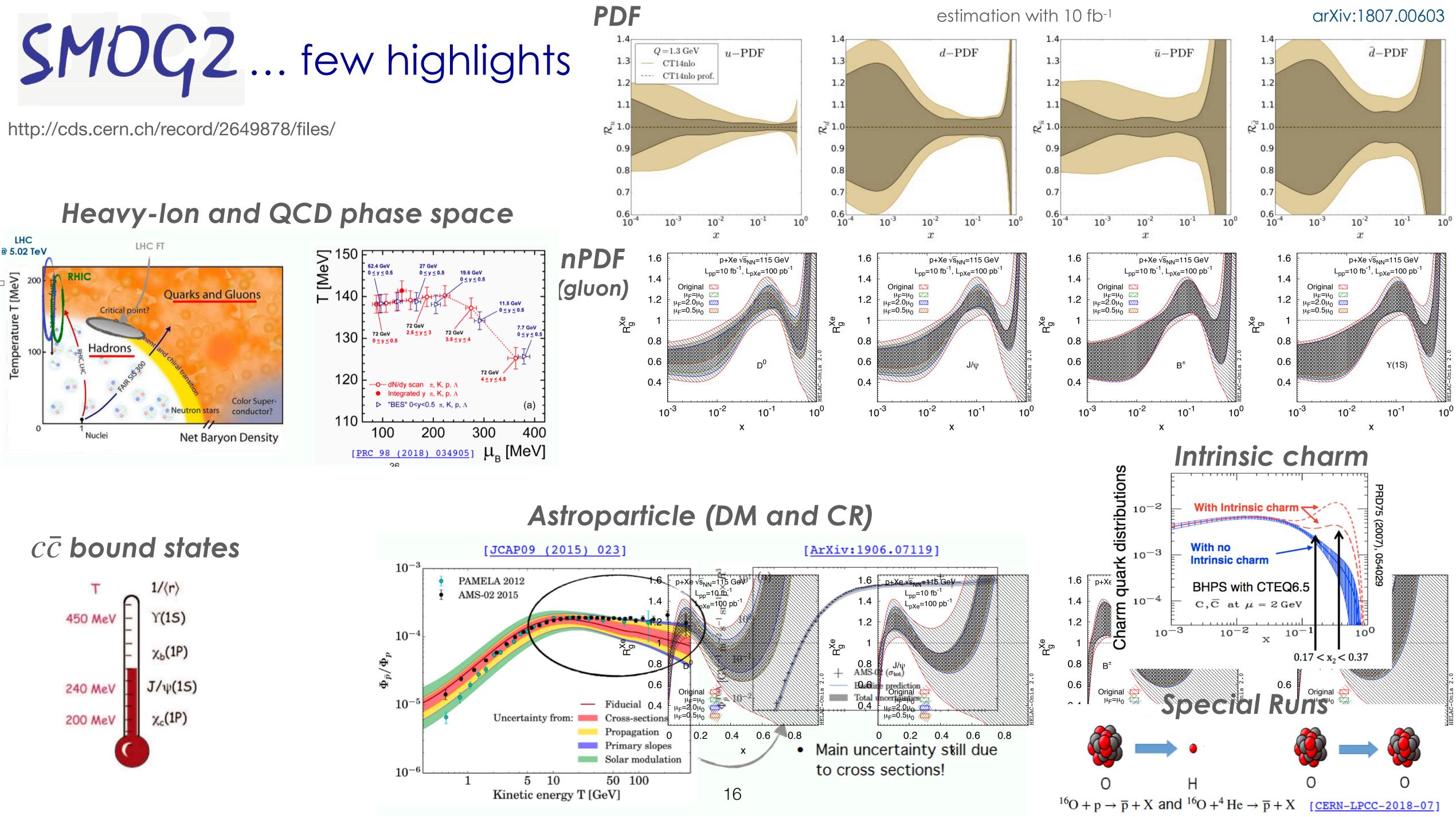
LHCb-FIGURE-2023-008

in $\sim 10'$ of data taking

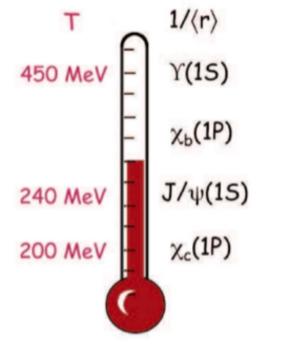


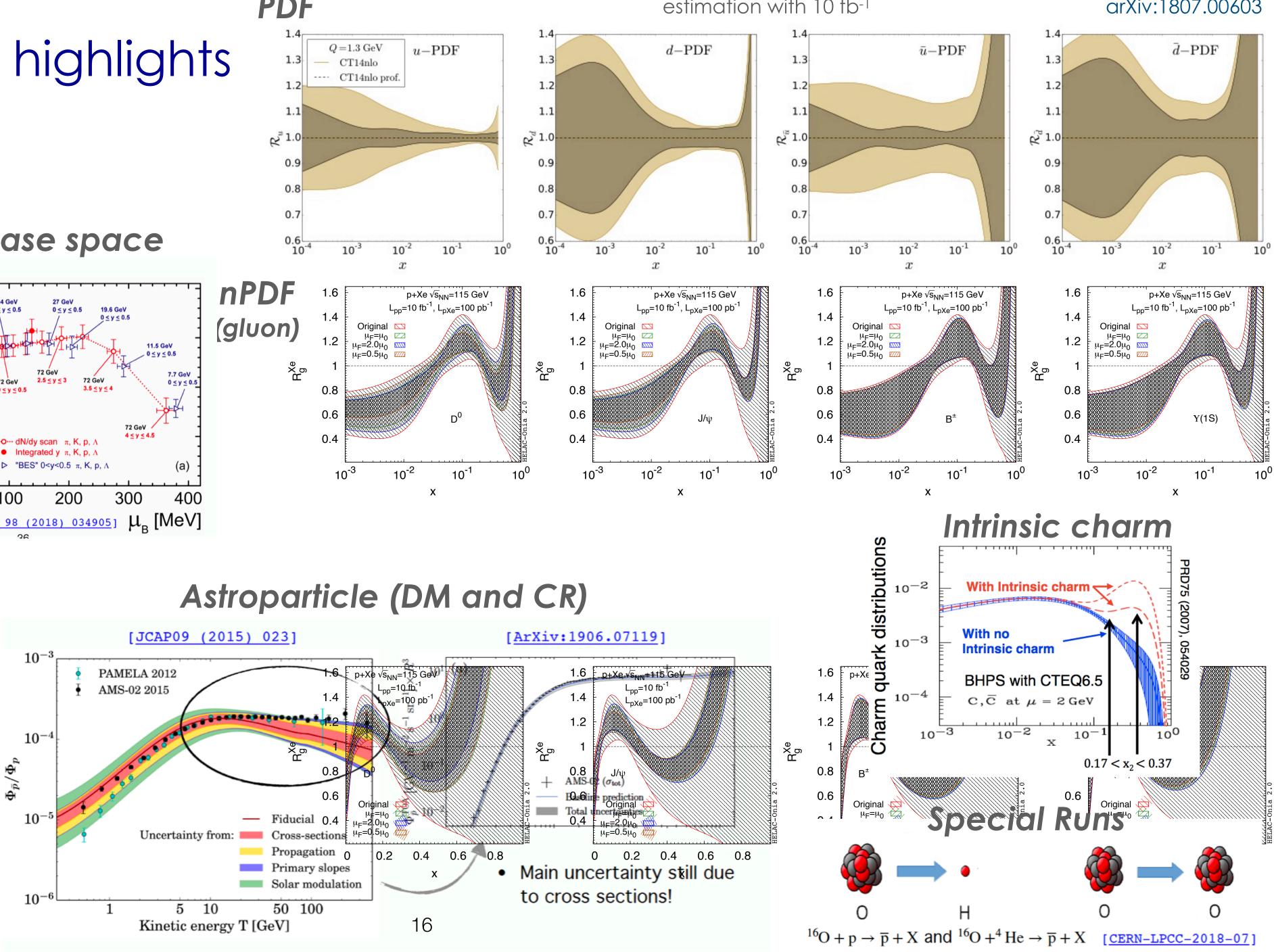




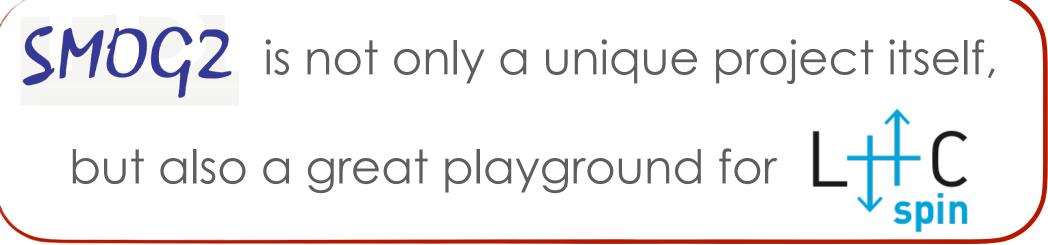






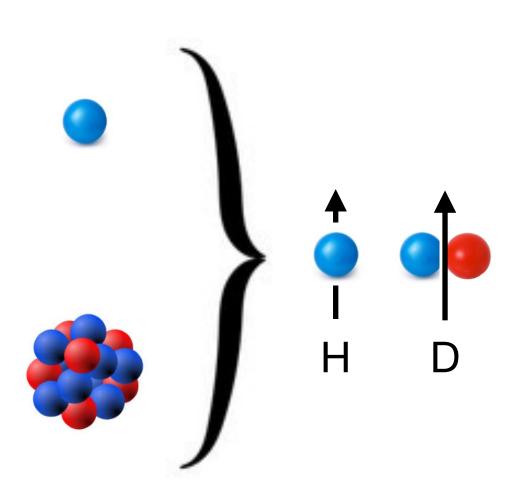






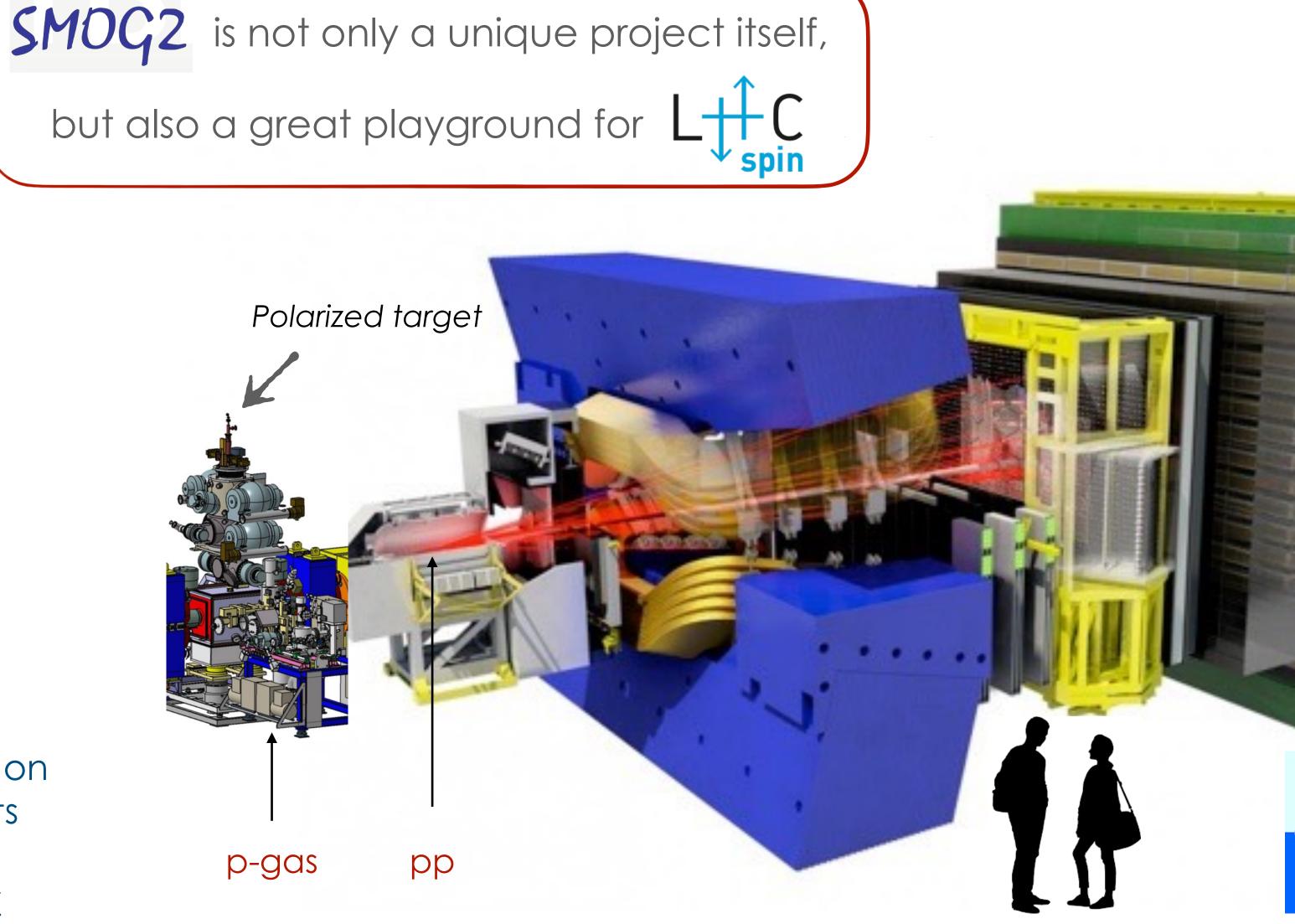


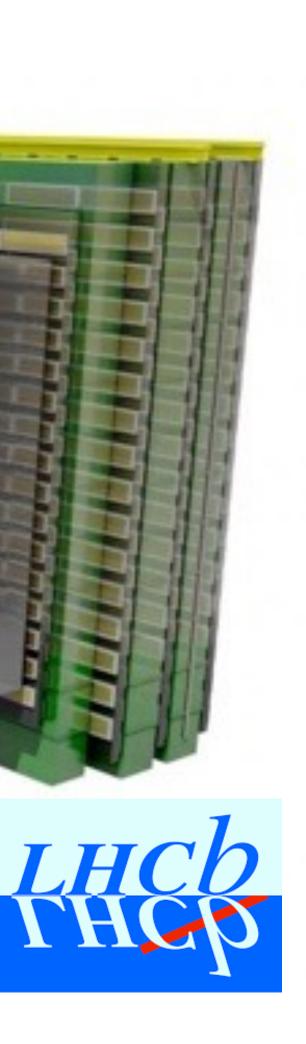




Successful technology based on HERA and COSY experiments

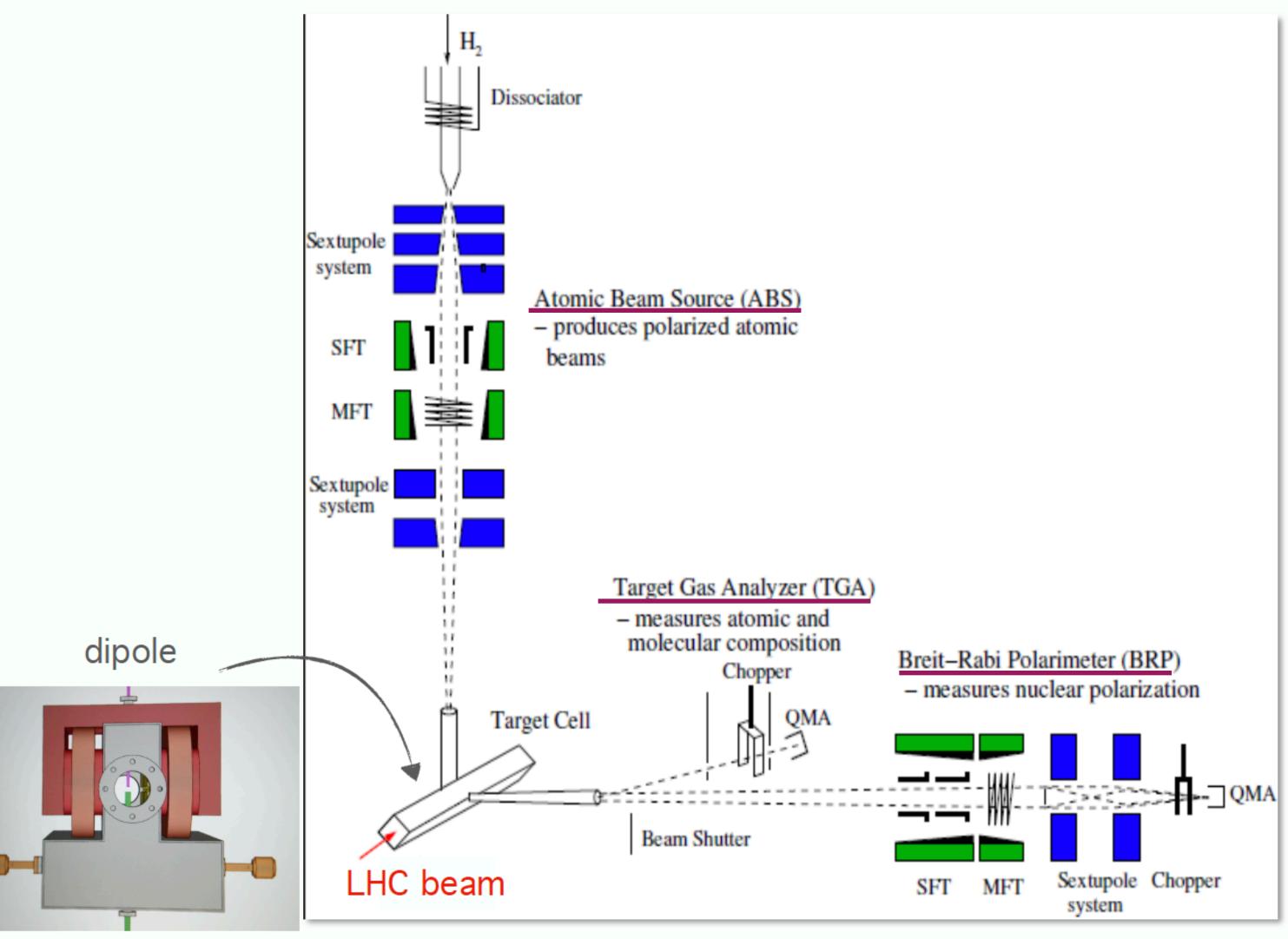
Challenge: develop a <u>new</u> generation of polarized targets

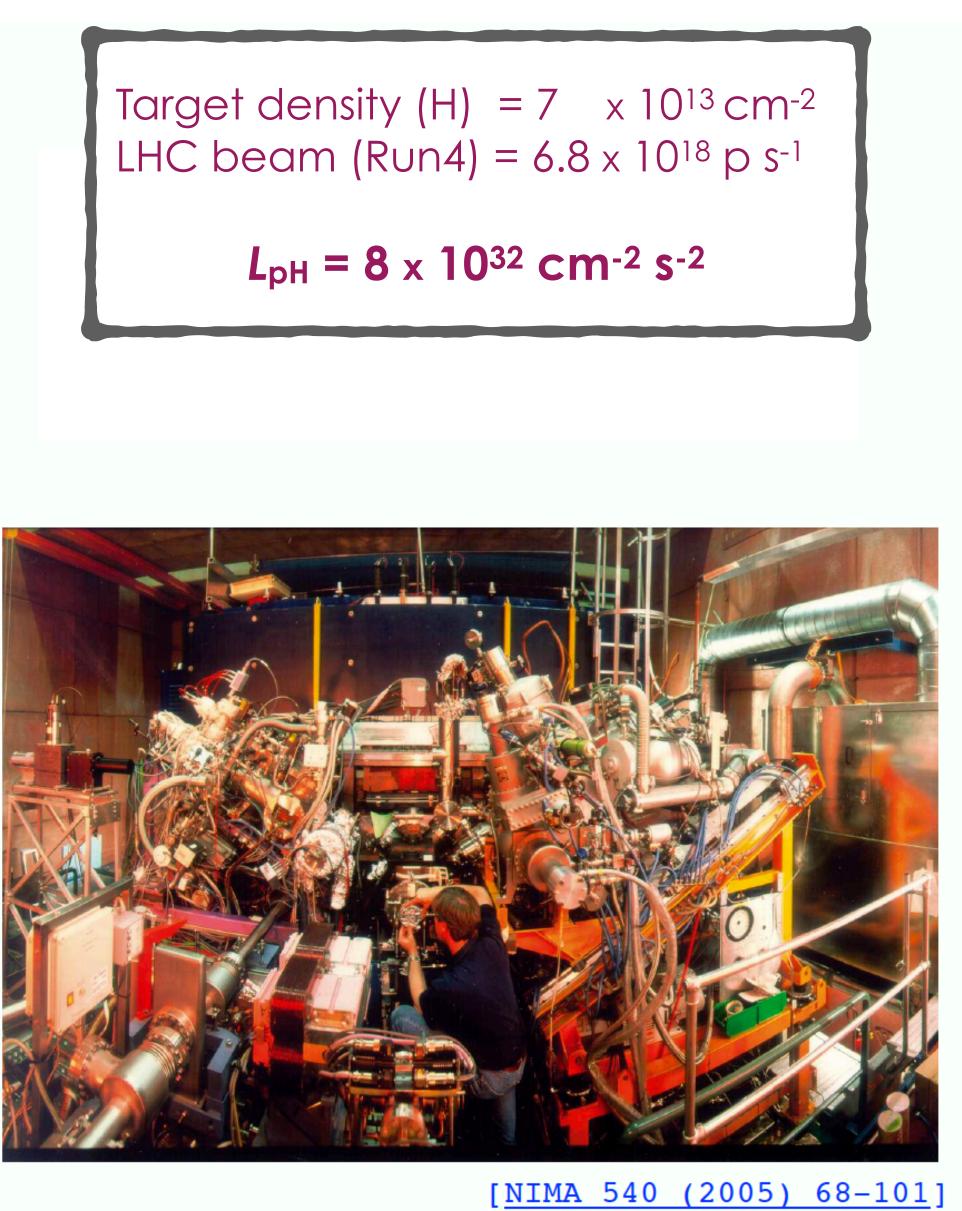




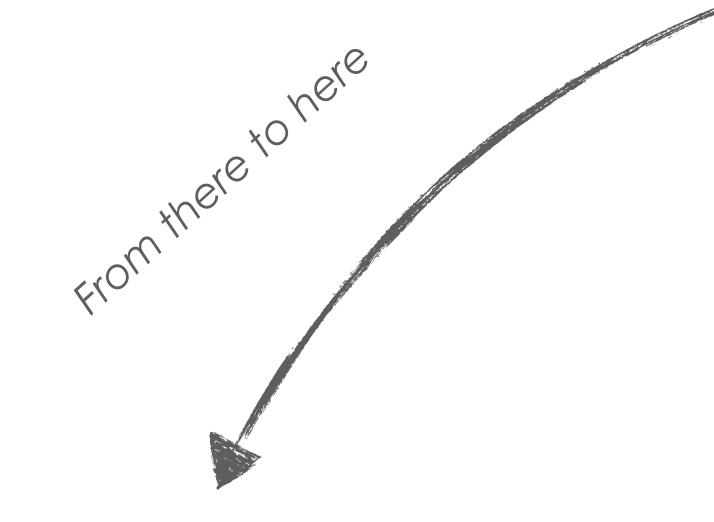


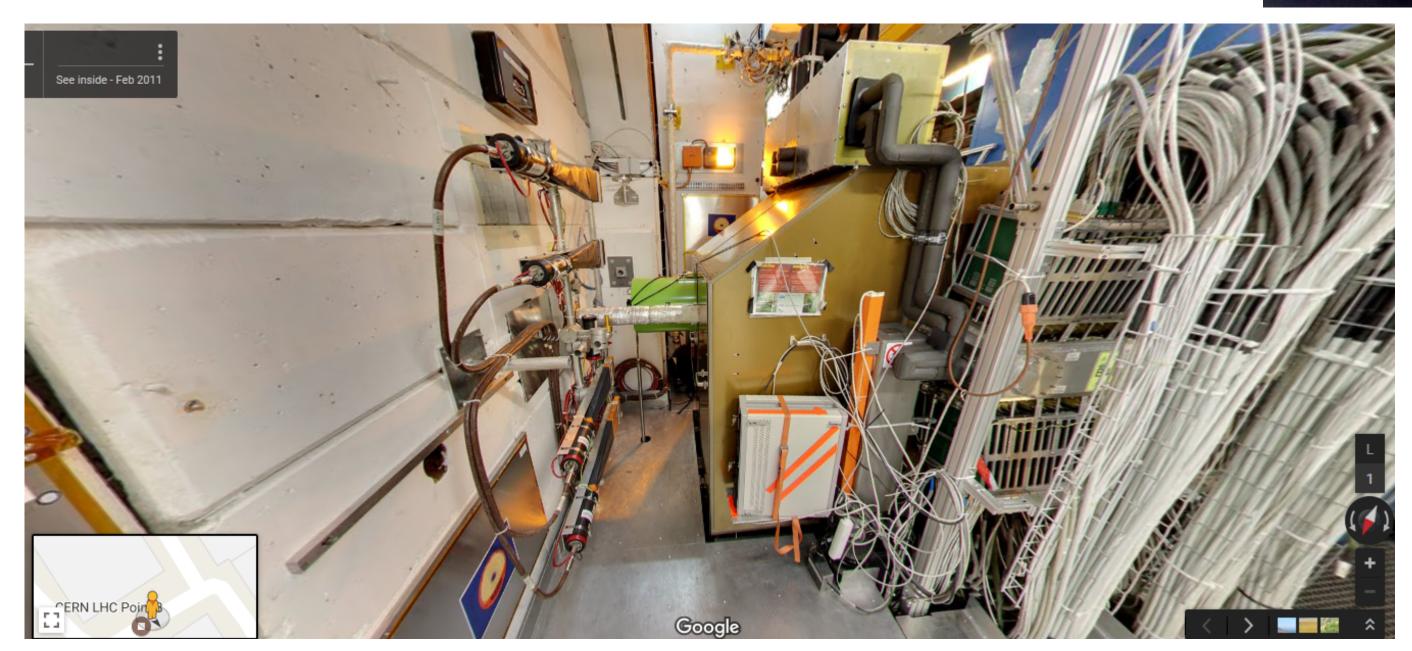
LHCspin experimental setup

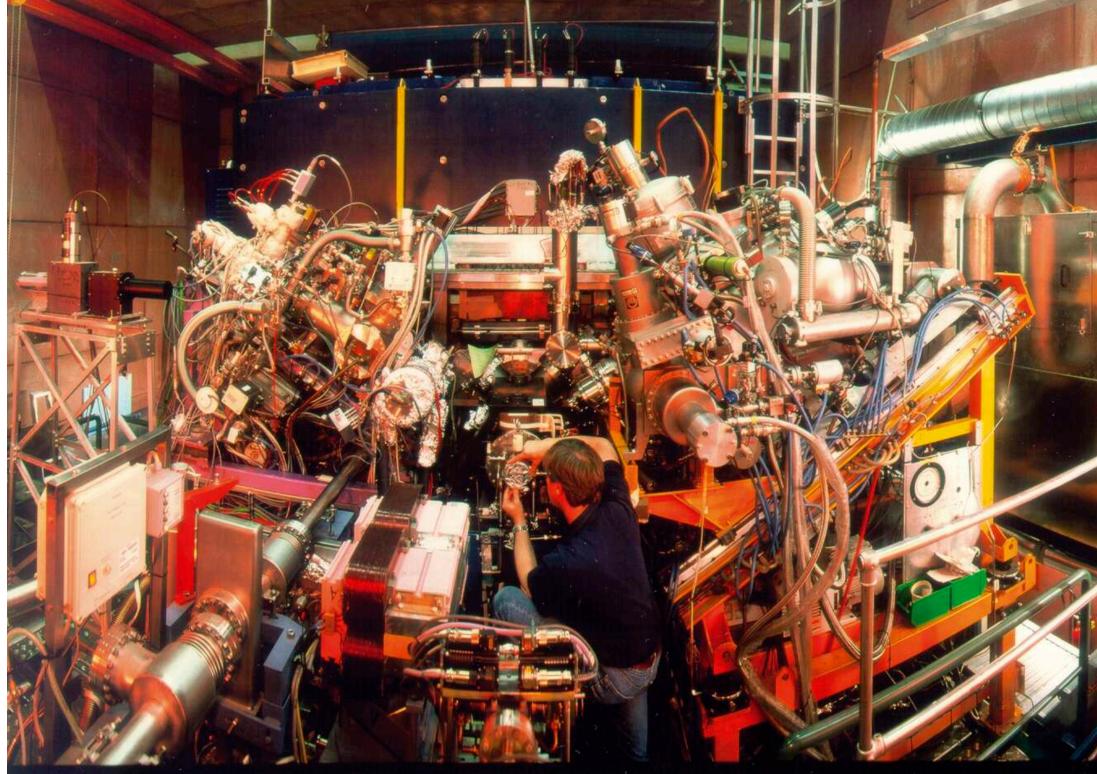




HERMES PGT





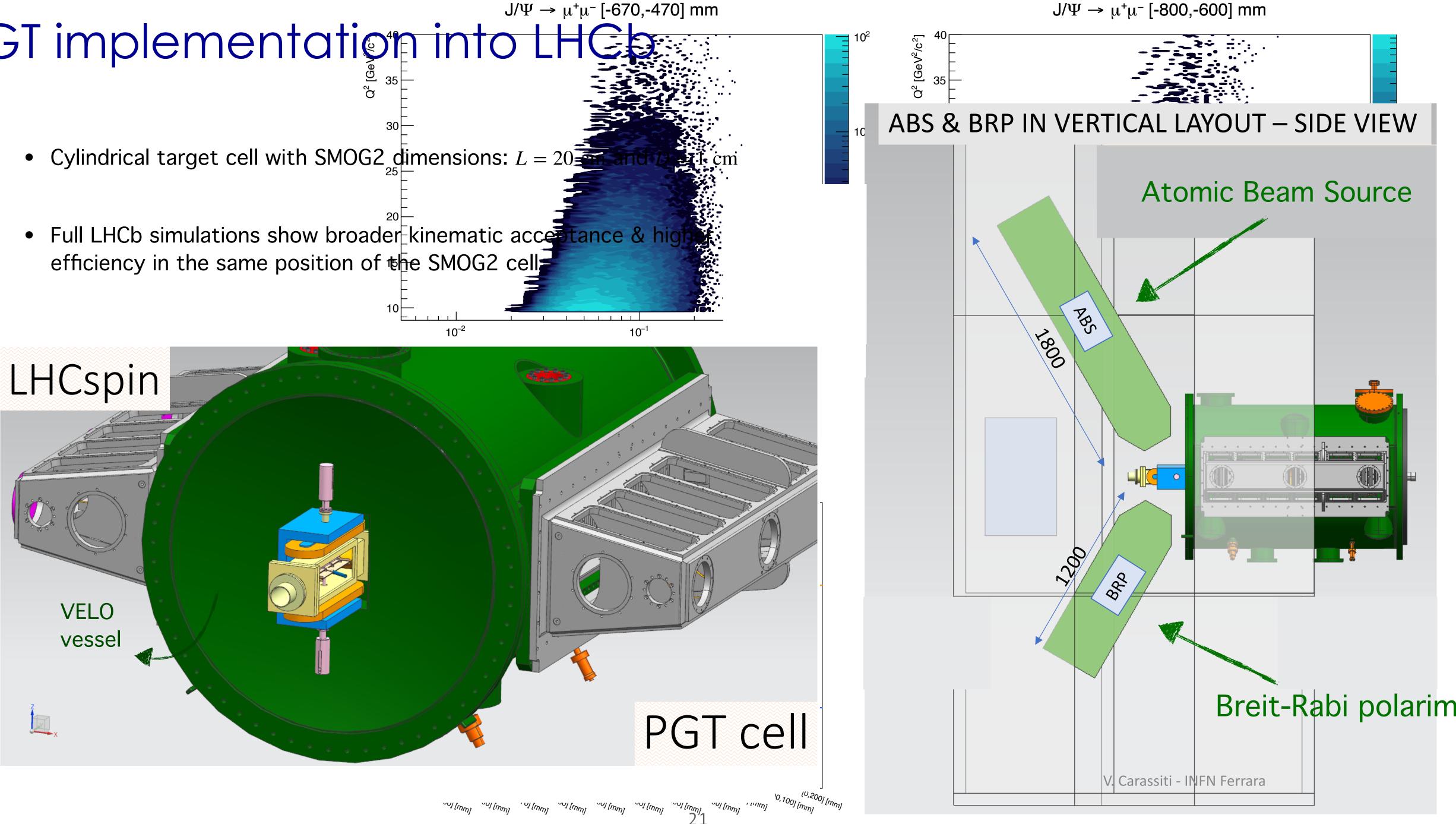


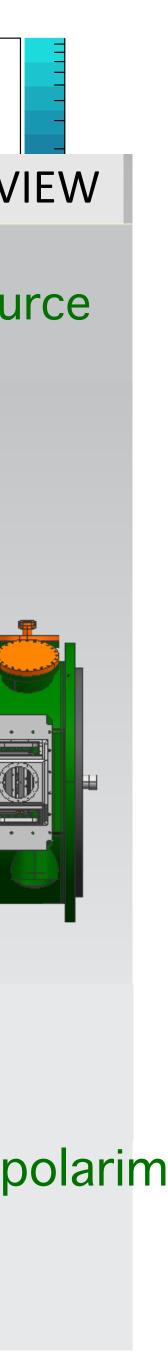
Space available in front of LHCb



$J/\Psi \rightarrow \mu^+\mu^-$ [-670,-470] mm PGT implementation into LHCb

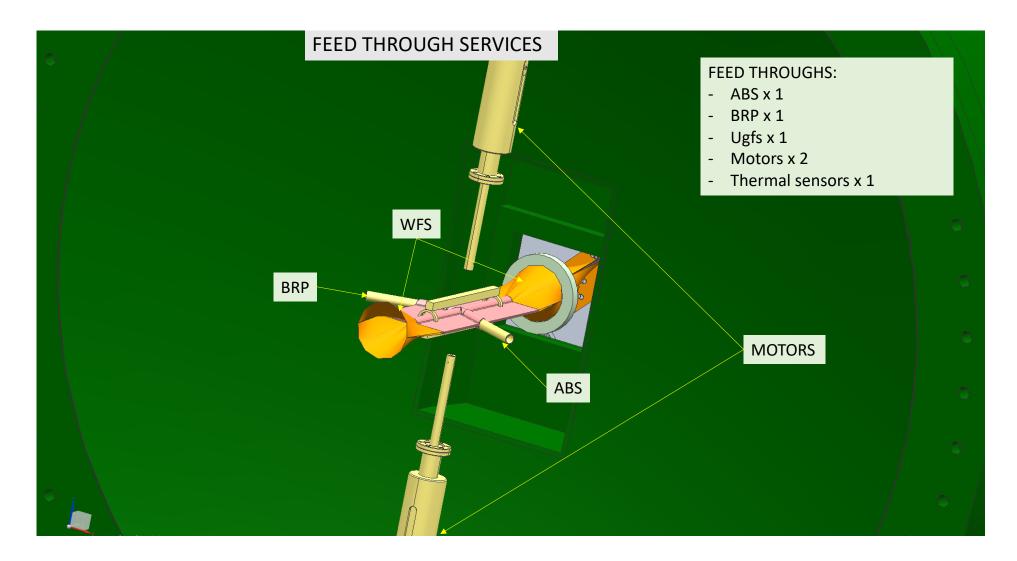
- efficiency in the same position of the SMOG2 cell





PGT implementation into LHCb

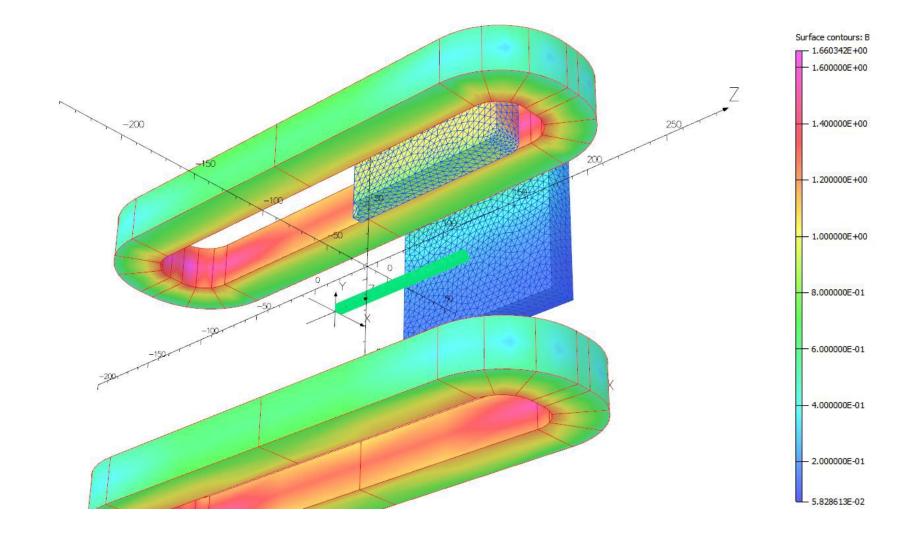
• Inject polarised gas via ABS and unpolarised gas via UGFS



- Compact dipole magnet → static transverse field
- Superconductive coils + iron yoke configuration fits the space constraints
- B = 300 mT with polarity inversion, $\Delta B/B \simeq 10\%$, suitable to avoid beam-induced depolarisation [Pos (SPIN2018)]

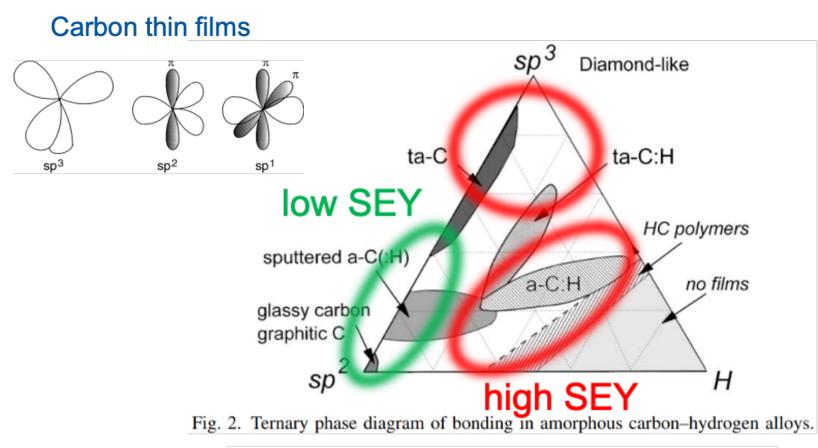
Possibility to switch to a solenoid and provide longitudinal polarisation (e.g. in LHC Run 5)

Transverse polarisation MAGNET INFO FOR THE CELL ACCESS yoke coil ٢ Ū - MAGNET IN TWO SEPARATED COILS ABS - C SHAPE YOKE OR WITH A SIDE **REMOVABLE PLATE**





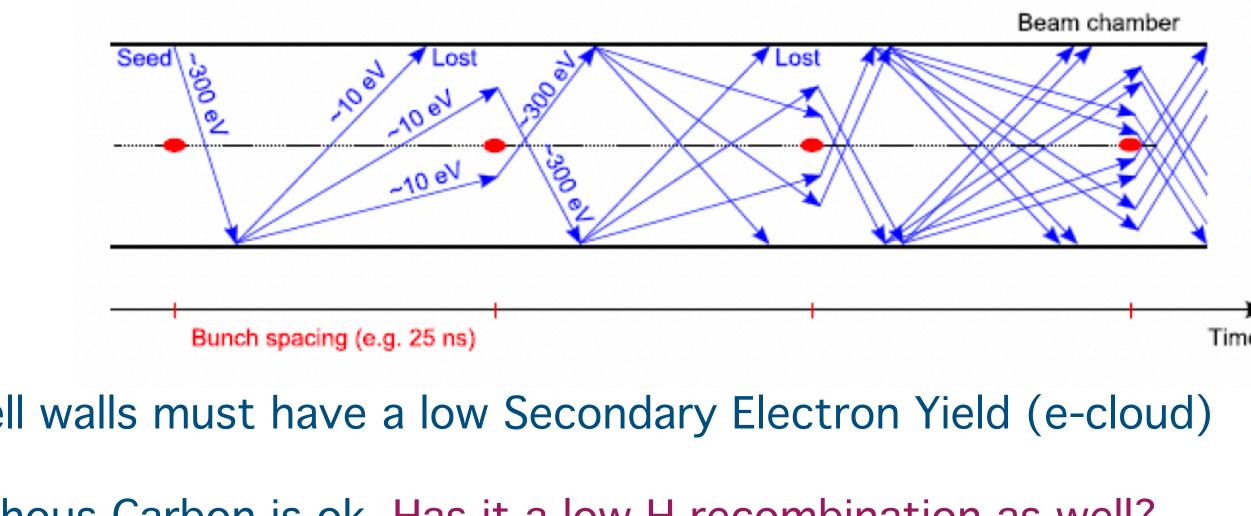


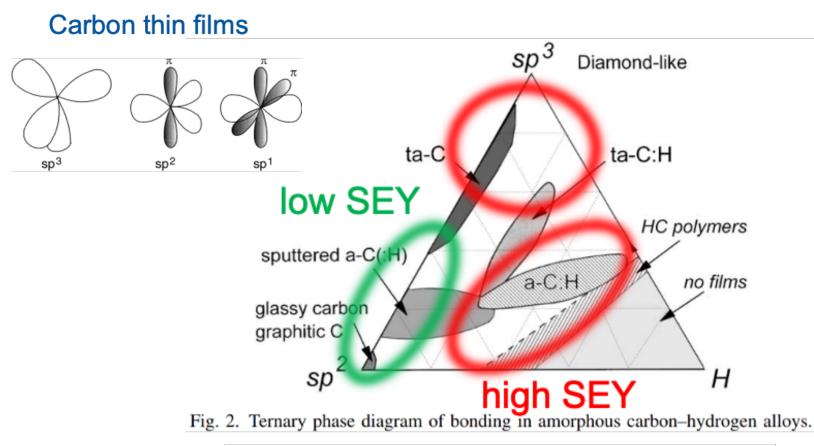


J. Robertson/Materials Science and Engineering R 37 (2002) 129–281

The material of the cell walls must have a low Secondary Electron Yield (e-cloud)

As for SMOG2, Amorphous Carbon is ok. <u>Has it a low H recombination as well?</u>



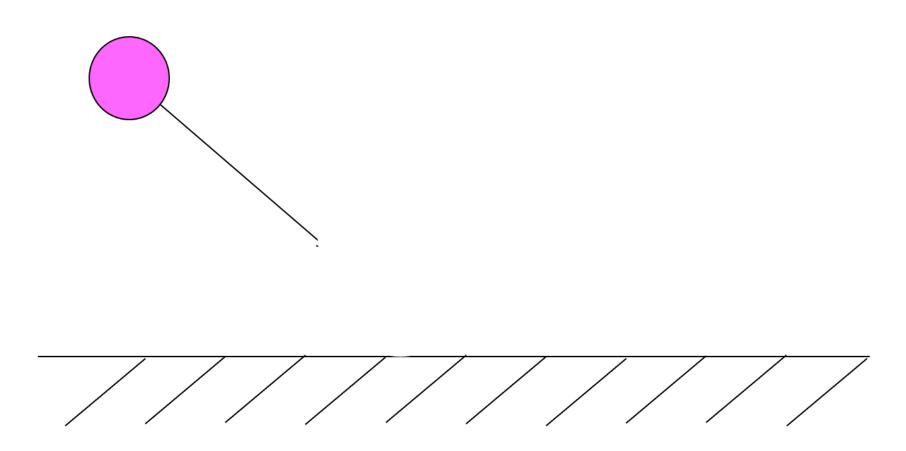


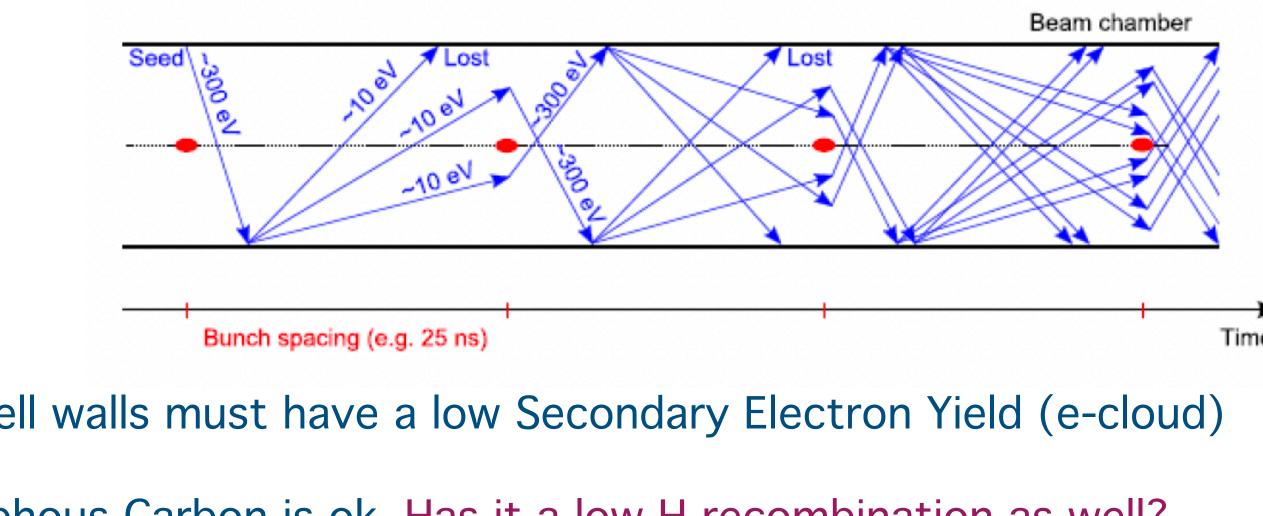
J. Robertson/Materials Science and Engineering R 37 (2002) 129–281

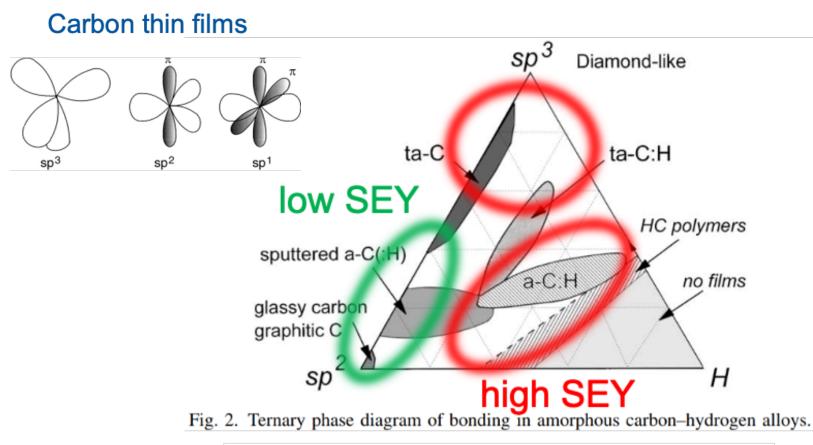
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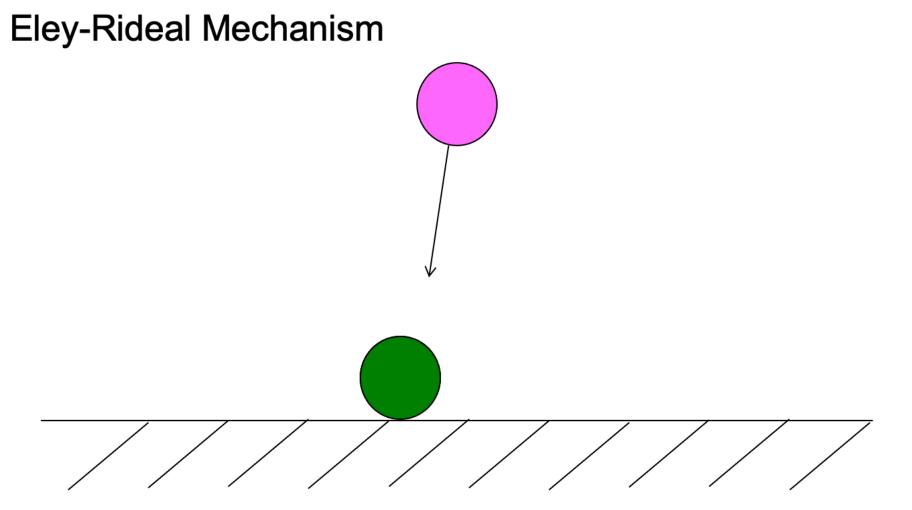
Eley-Rideal Mechanism

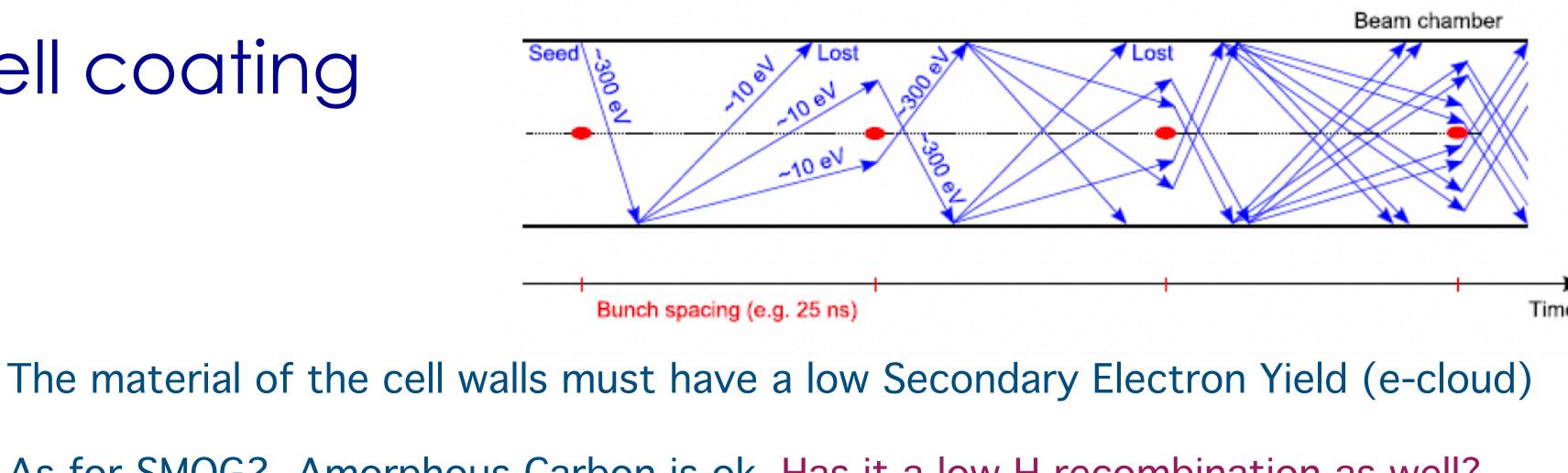




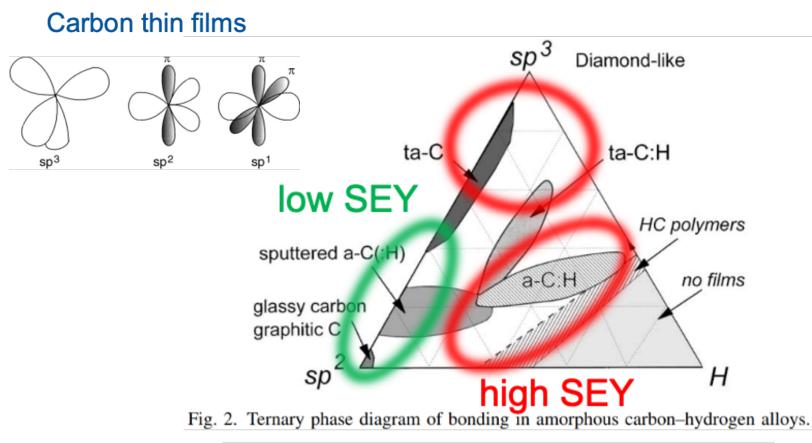


J. Robertson/Materials Science and Engineering R 37 (2002) 129–281



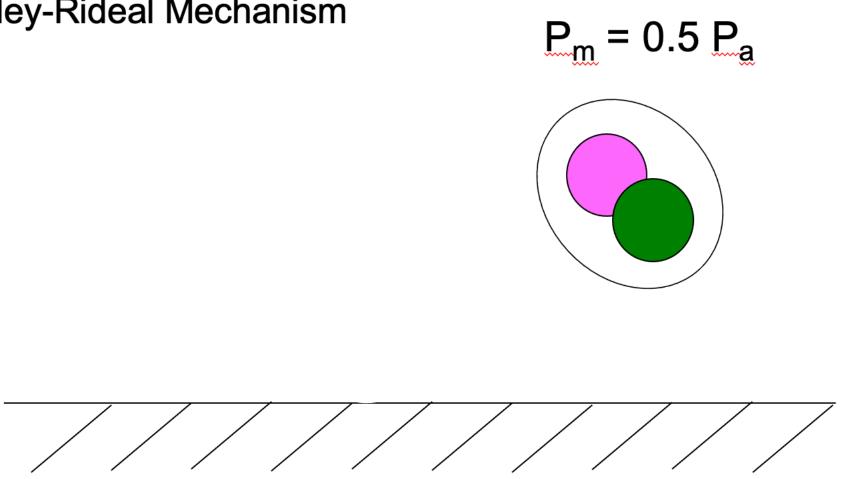


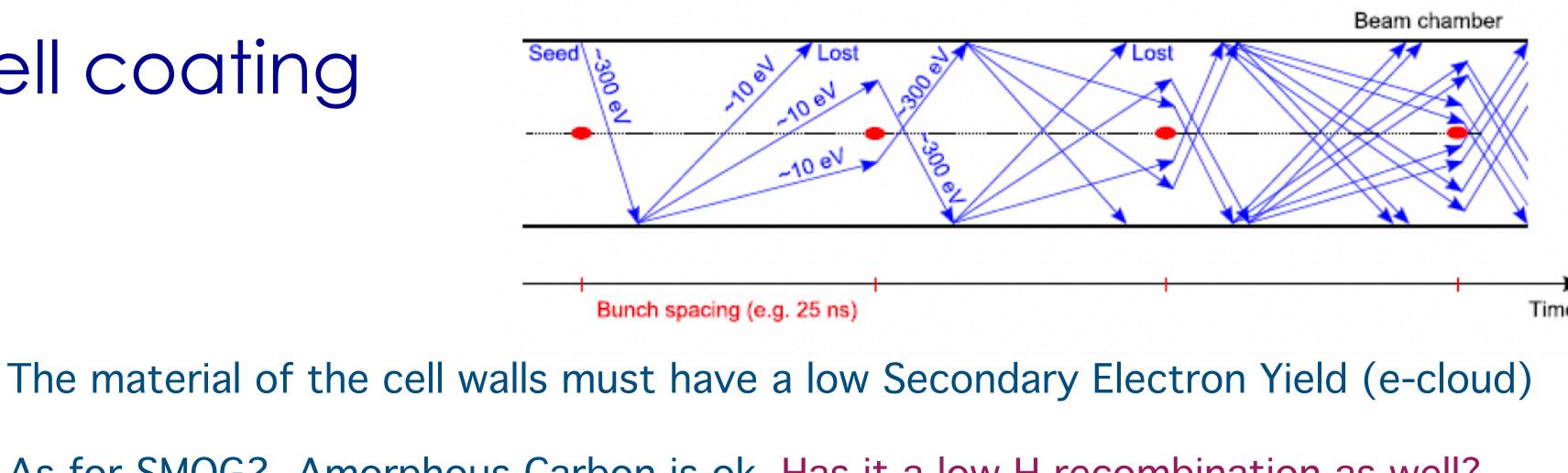
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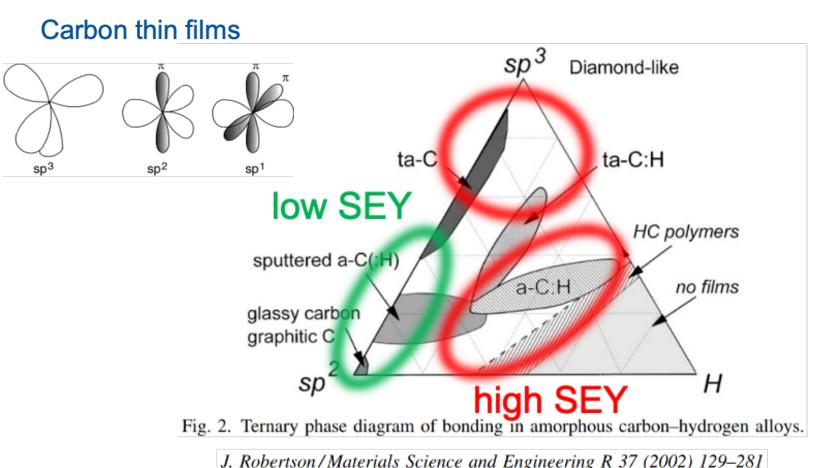
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Eley-Rideal Mechanism

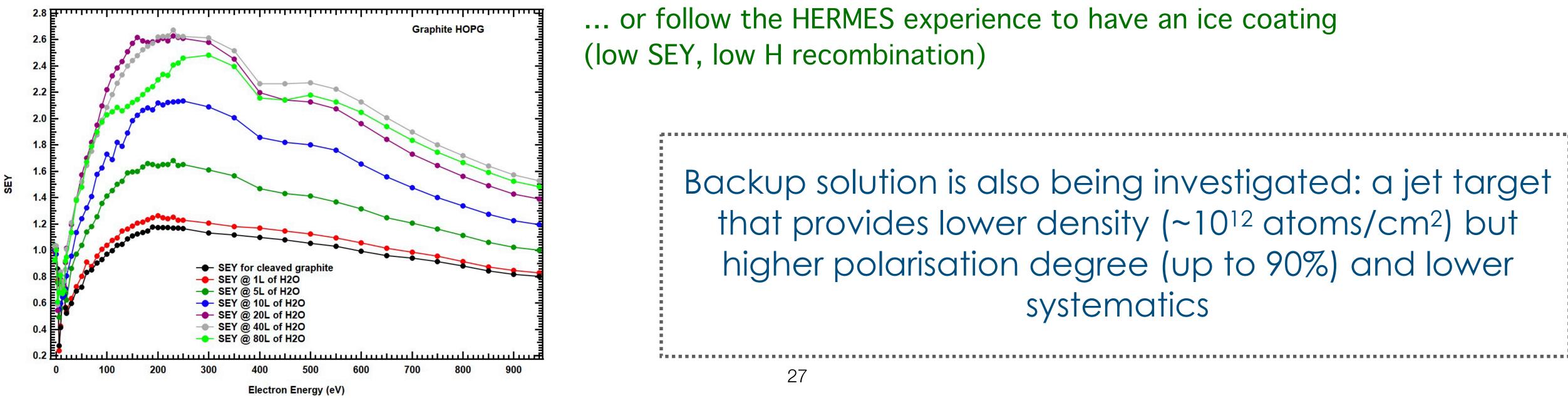


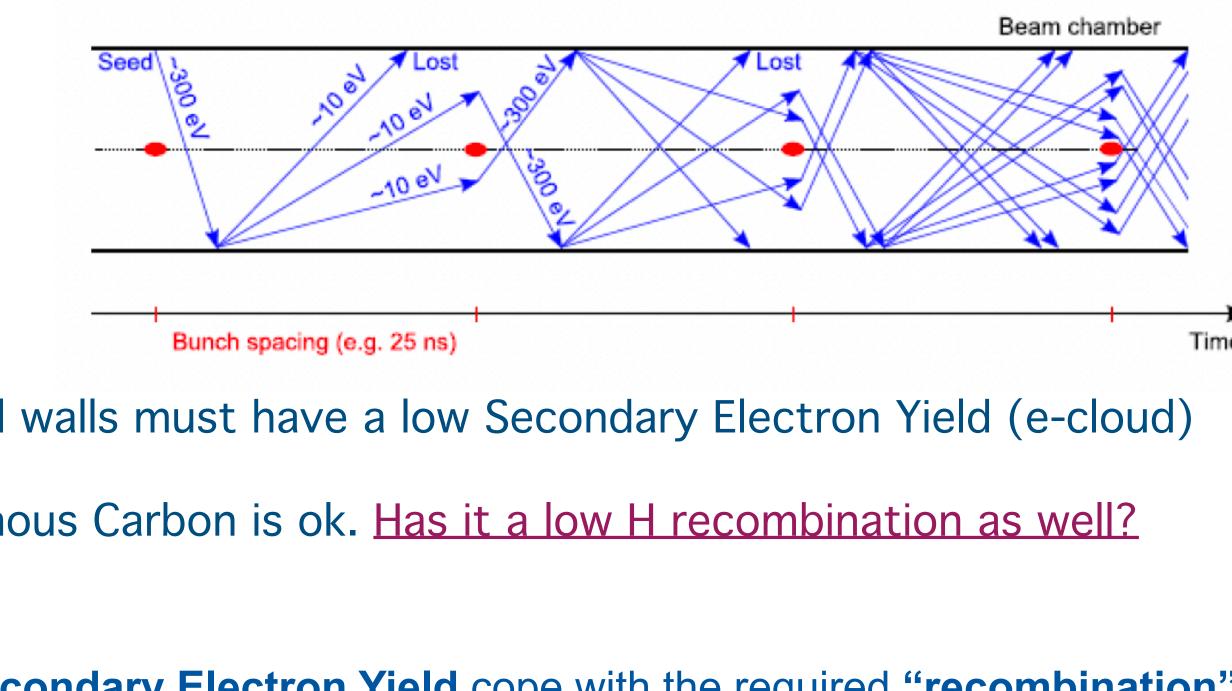


- As for SMOG2, Amorphous Carbon is ok. <u>Has it a low H recombination as well?</u>



rate of polarized H atoms injected in the storage cell





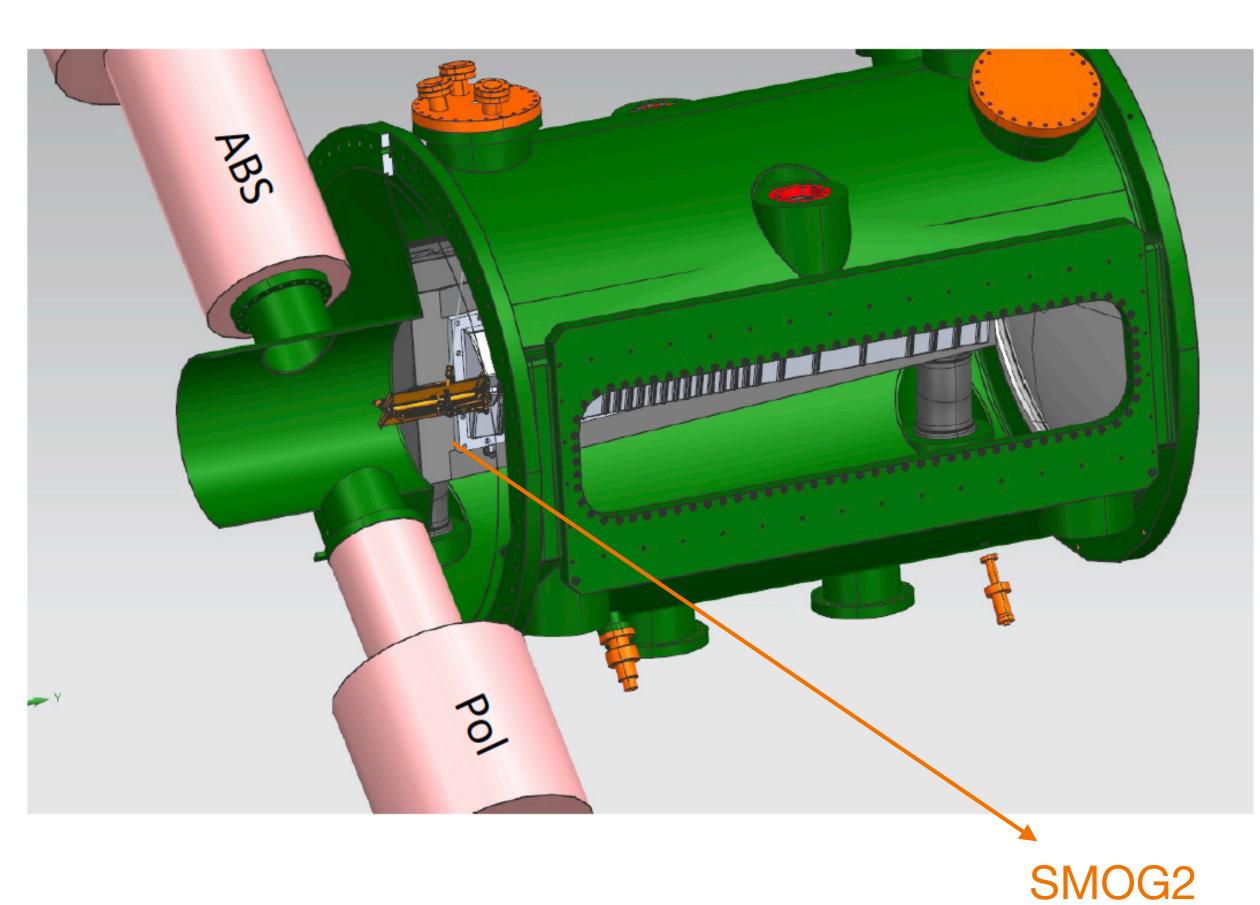
- The material of the cell walls must have a low Secondary Electron Yield (e-cloud)
- As for SMOG2, Amorphous Carbon is ok. <u>Has it a low H recombination as well?</u>
- Studies ongoing in order to understand if carbon films with low secondary Electron Yield cope with the required "recombination"
 - ... or follow the HERMES experience to have an ice coating (low SEY, low H recombination)
 - Backup solution is also being investigated: a jet target that provides lower density (~10¹² atoms/cm²) but higher polarisation degree (up to 90%) and lower systematics

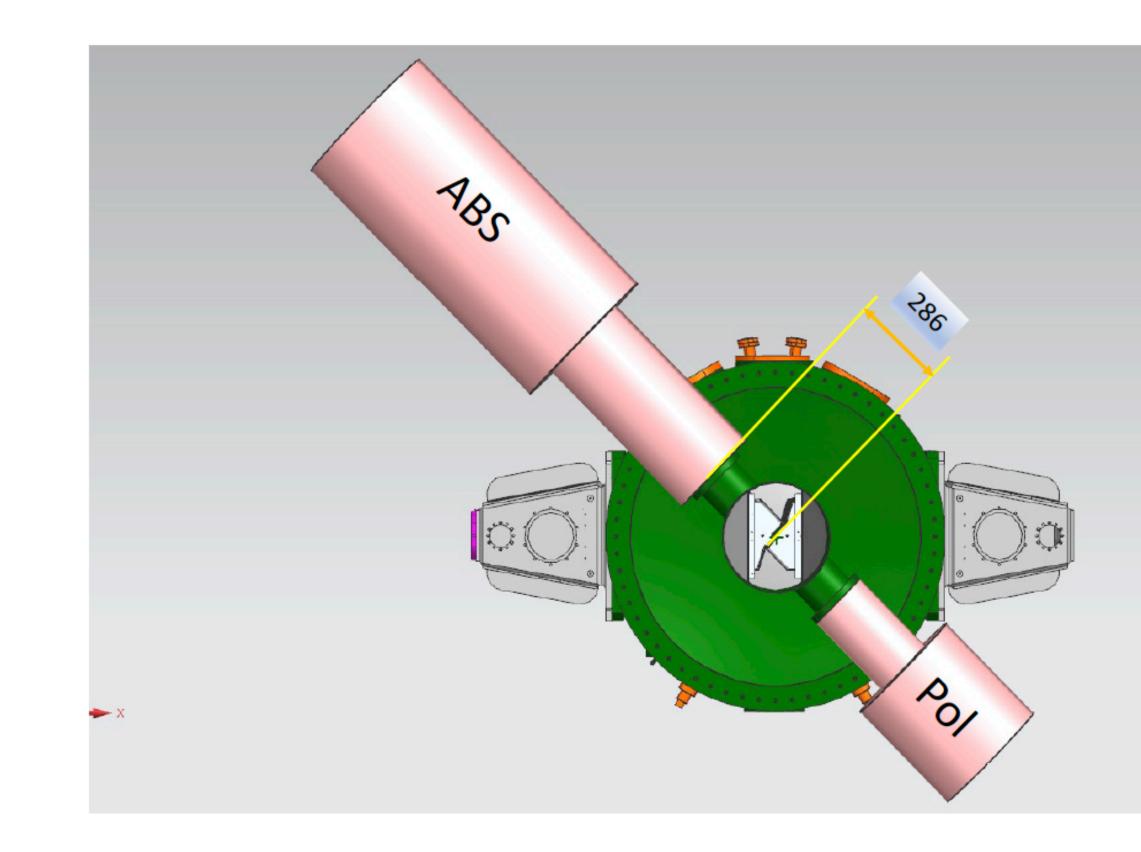


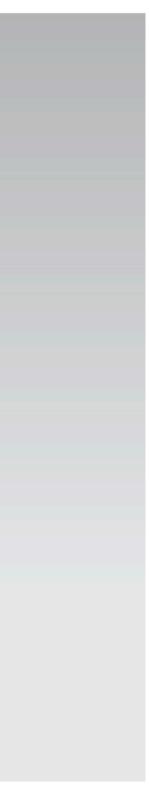
The jet target option

Alternative solution with jet target also under evaluation:

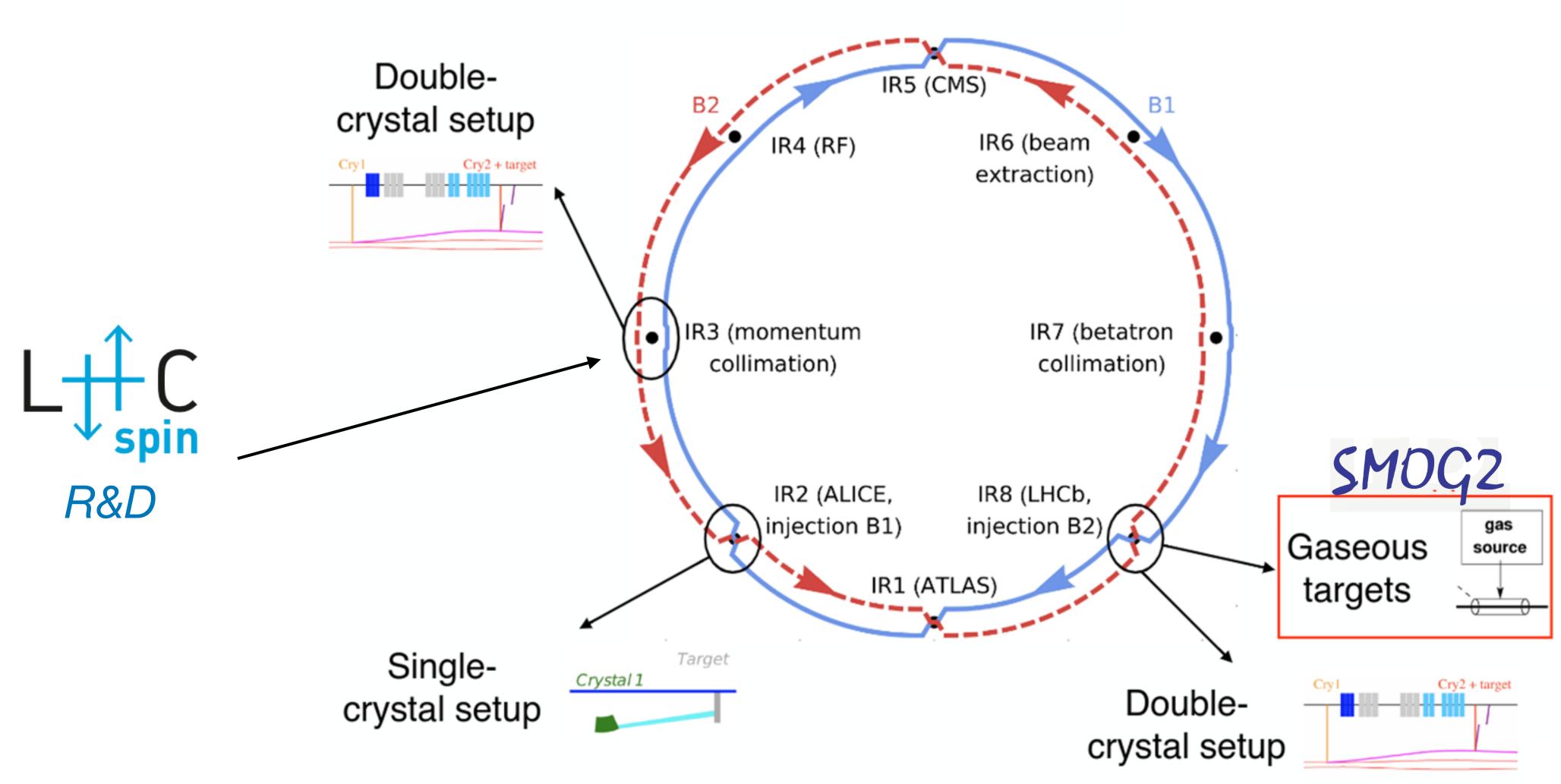
- lower density (~ 10^{12} atoms/ cm^2) •
- higher polarization (up to 90%) •
- lower systematics in P measurement (virtually close to 0) •







The LHC Interaction Region 3



S. Redaelli, PBC General WG, 02/12/2021

-develop a new target system starting from an existing system

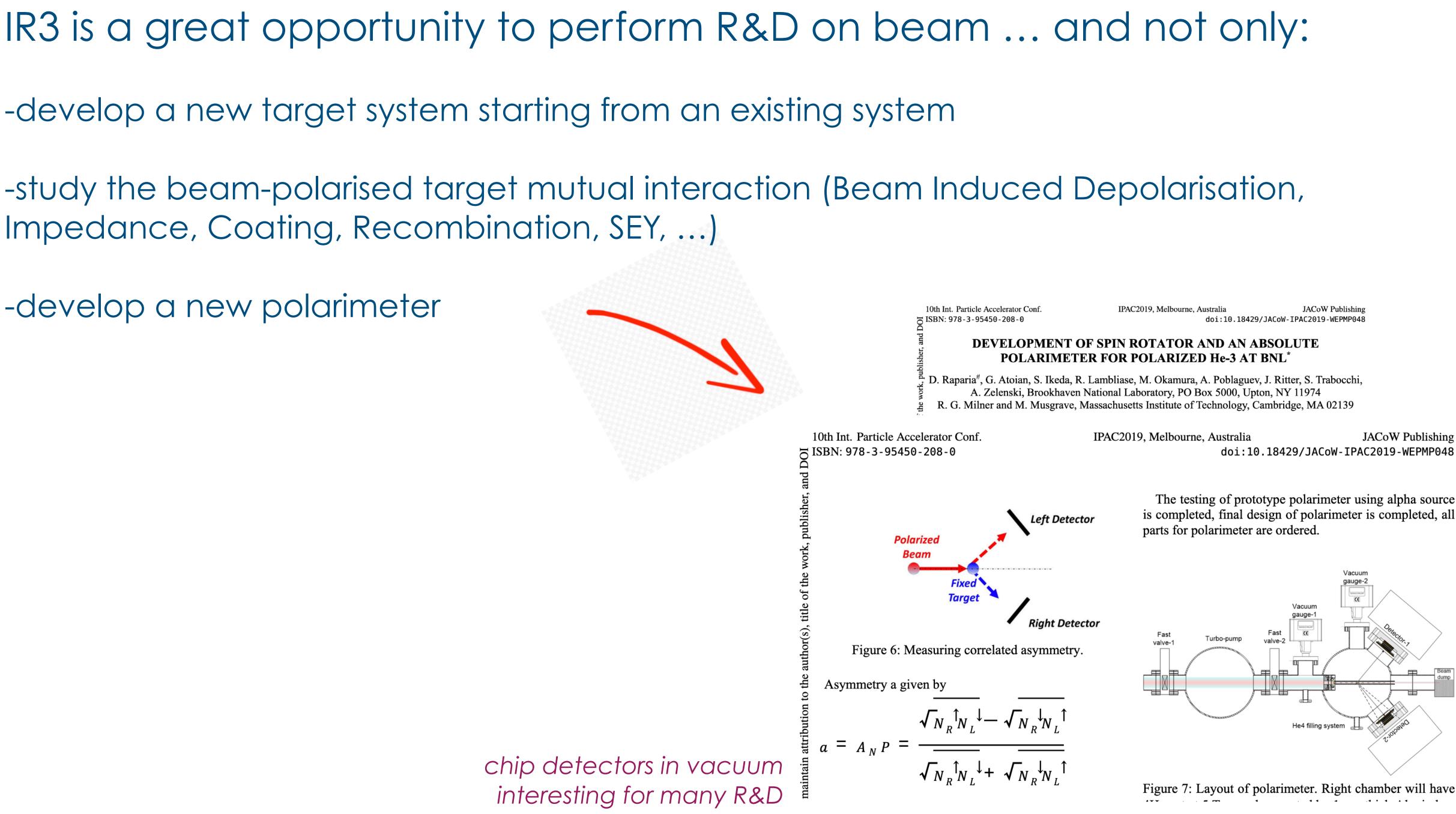
-develop a new target system starting from an existing system

-study the beam-polarised target mutual interaction (Beam Induced Depolarisation, Impedance, Coating, Recombination, SEY, ...)

-develop a new target system starting from an existing system

Impedance, Coating, Recombination, SEY, ...)

-develop a new polarimeter



-develop a new target system starting from an existing system

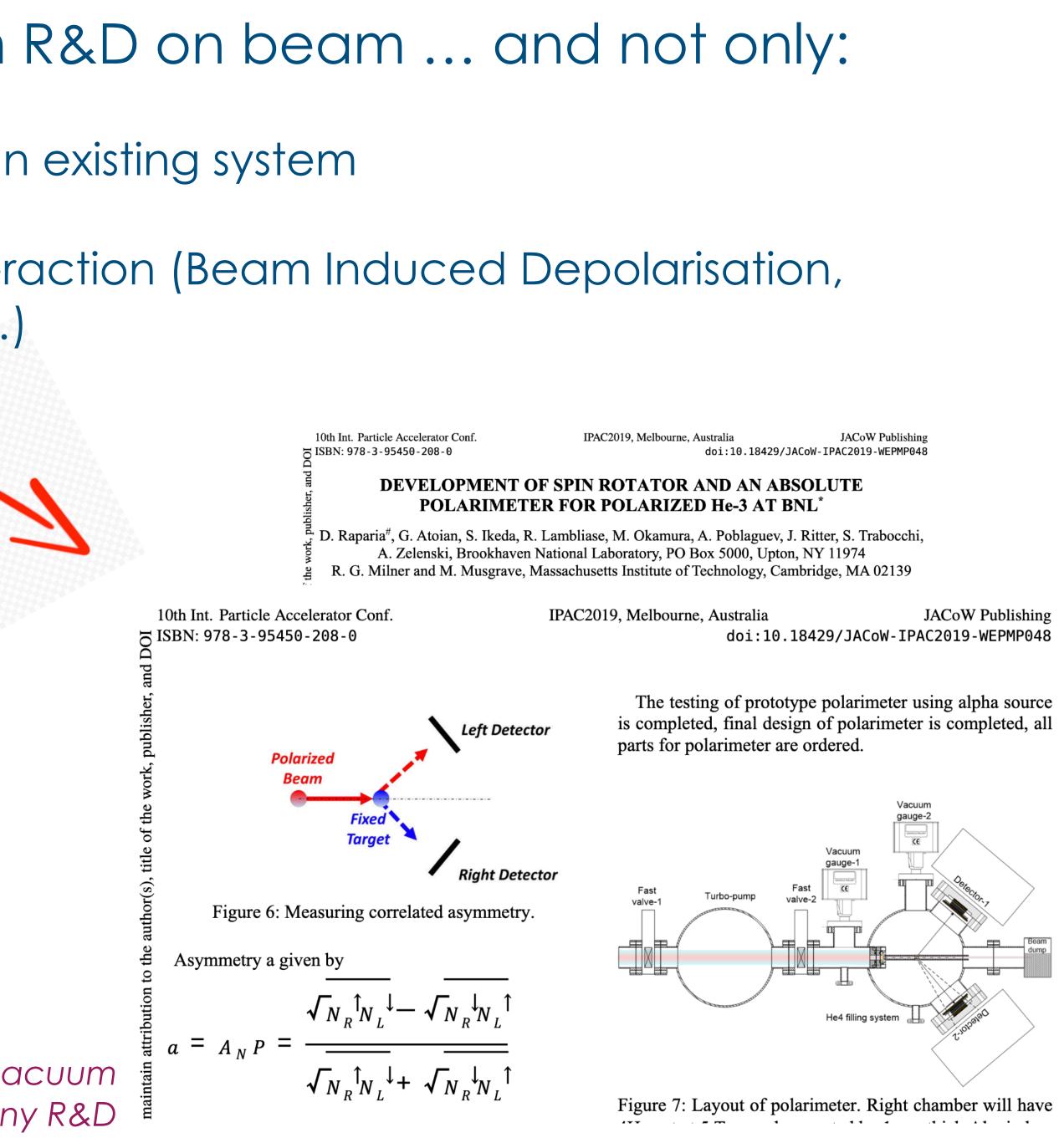
-study the beam-polarised target mutual interaction (Beam Induced Depolarisation, Impedance, Coating, Recombination, SEY, ...)

-develop a new polarimeter

-perform basic physics measurements

-create a CERN pool for polarised physics (as in the past)

chip detectors in vacuum interesting for many R&D



-develop a new target system starting from an existing system

-study the beam-polarised target mutual interaction (Beam Induced Depolarisation, Impedance, Coating, Recombination, SEY, ...)

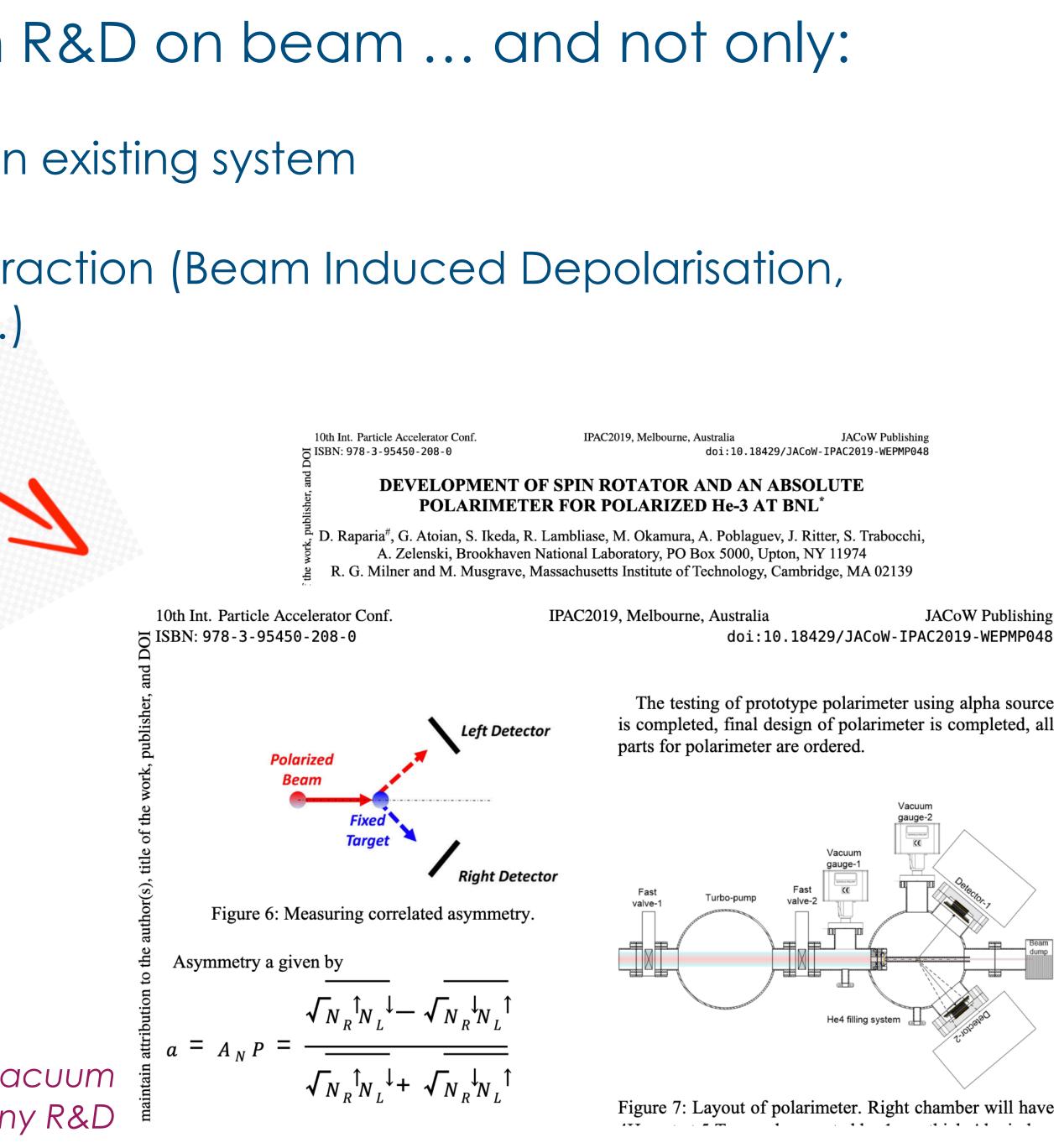
-develop a new polarimeter

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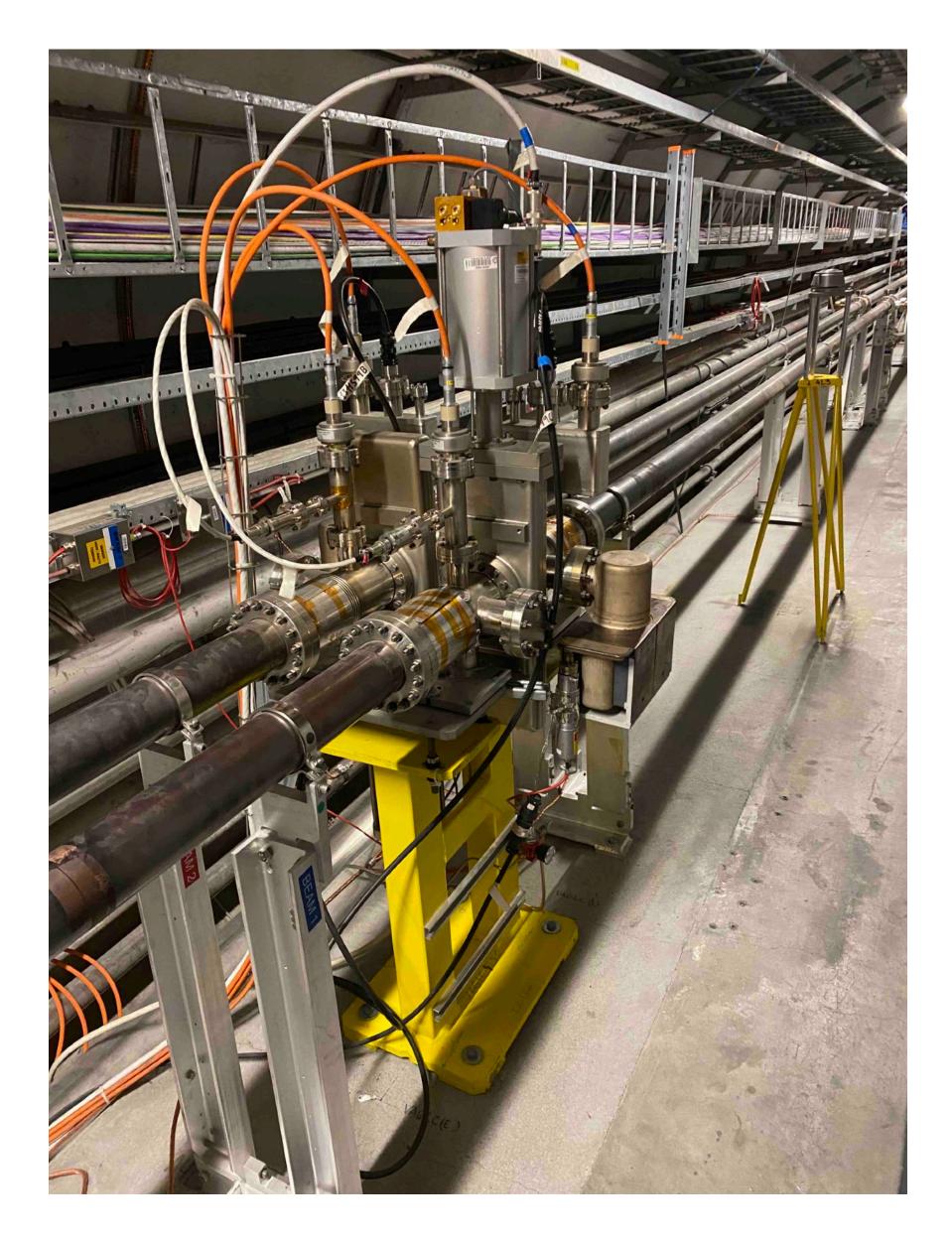
-create a CERN pool for polarised physics (as in the past)

-allow people/groups to join even if not officially in LHCb

chip detectors in vacuum interesting for many R&D







Sector value to isolate the region

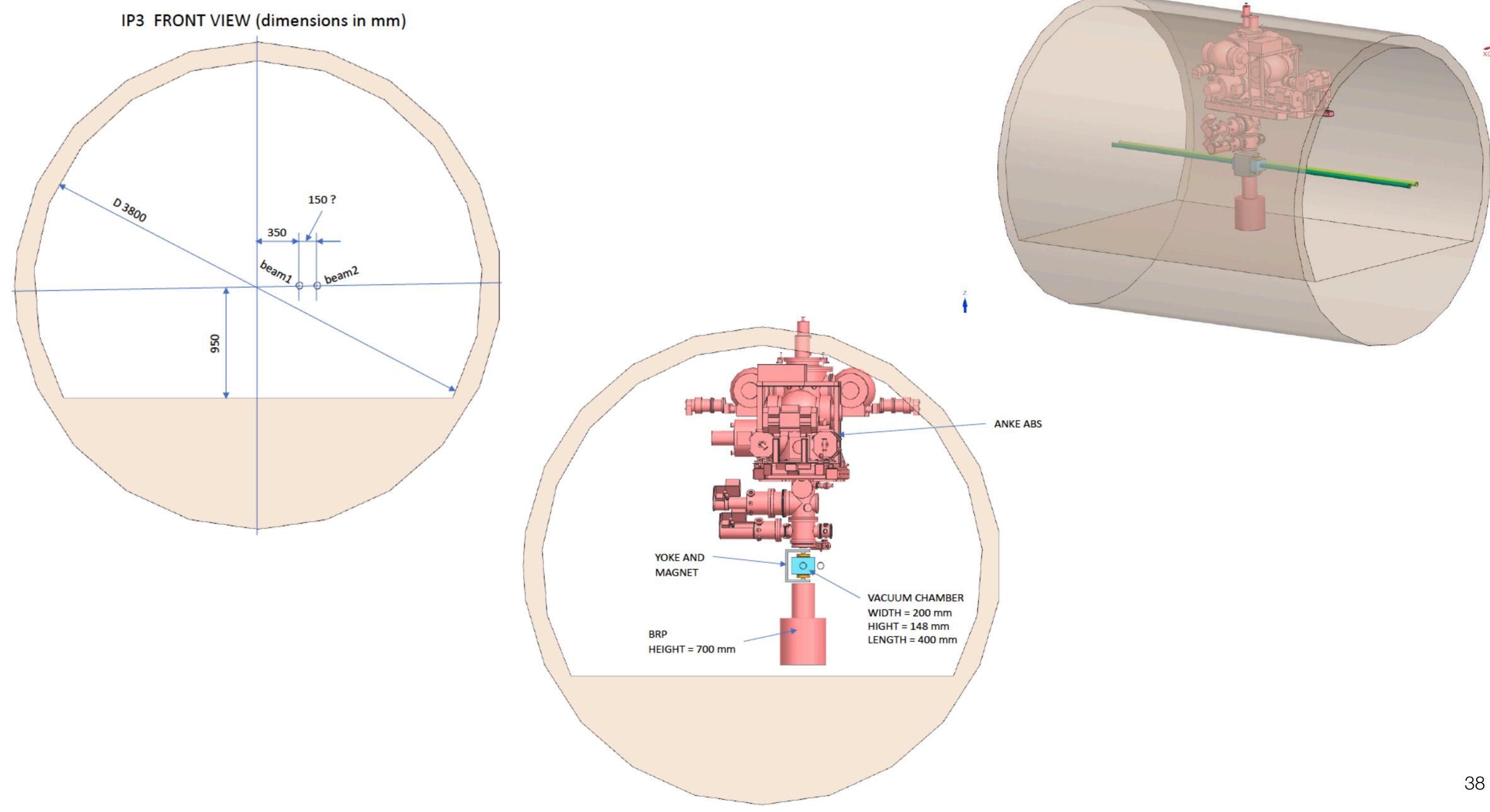


IR3 is a low radiation area (like a normal LHC-IP)

Investigation and discussions with LHC experts are ongoing



The available target system is a good starting point ...





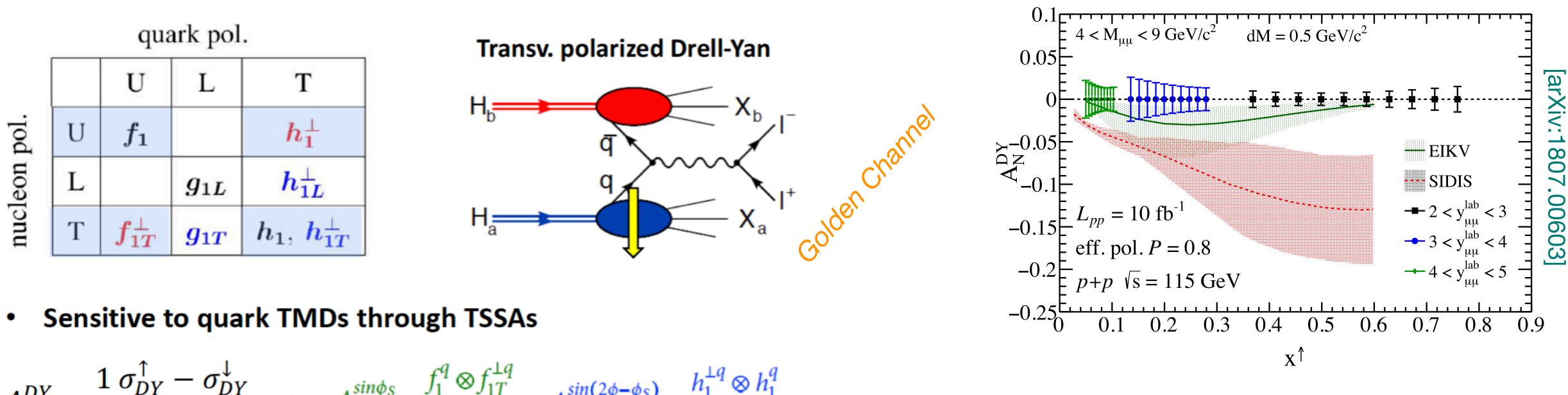


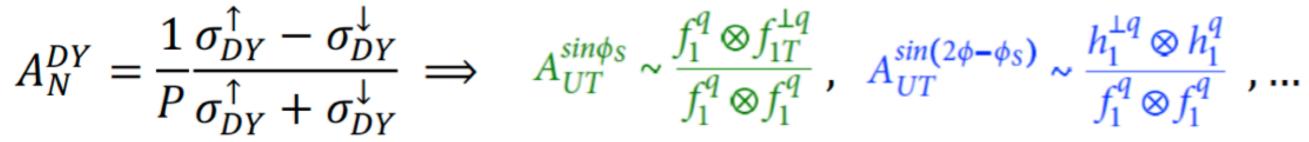
The physics goals of $L + C = \dots$ just a quick overview

- Multi-dimensional nucleon structure in a poorly explored kinematic domain
- Measure experimental observables sensitive to both quarks and gluons TMDs
- Make use of new probes (charmed and beauty mesons)
- Complement present and future SIDIS results
- Test non-trivial process dependence of quarks and (especially) gluons TMDs
- Measure exclusive processes to access GPDs



Quark TMDs





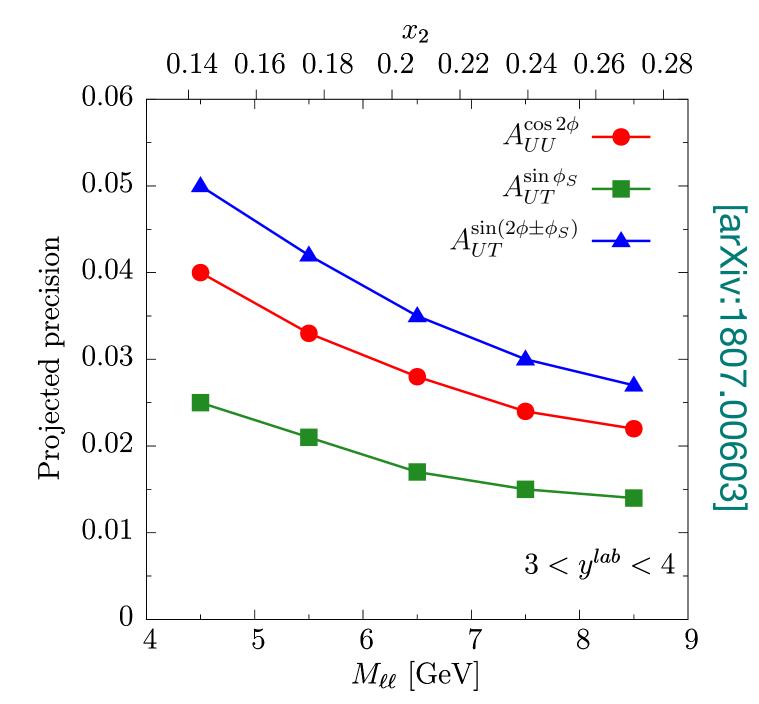
 $(\phi: azimuthal orientation of lepton pair in dilepton CM)$

LHCb has excellent μ -ID & reconstruction for $\mu^+\mu^-$

- Extraction of qTMDs does not require knowledge of FF
- Verify sign change of Sivers function wrt SIDIS $f_{1T}^{\perp}|_{DY} = -f_{1T}^{\perp}|_{SIDIS}$
- Test flavour sensitivity using both H and D targets

dominant: $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^$ suppressed: $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+ \mu^-$

40

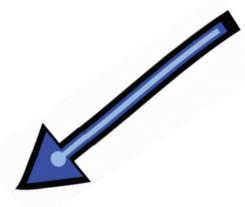


Gluon TMDs

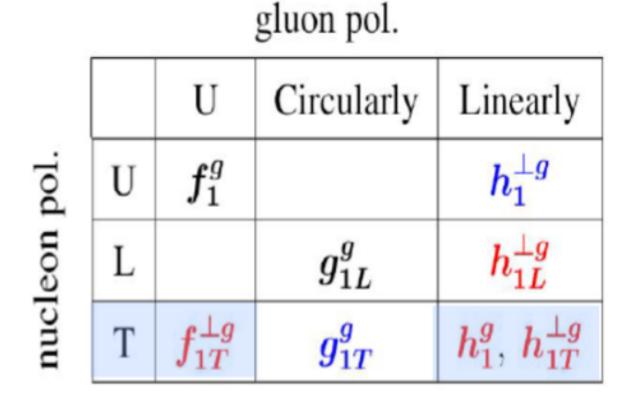
Theory framework well consolidated, but experimental access still extremely limited

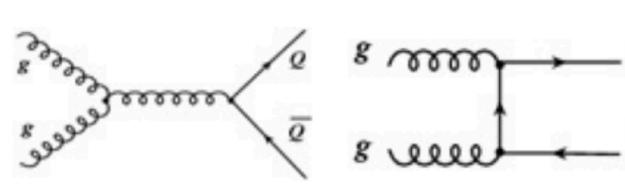
The most efficient way to access the gluon dynamics inside the proton at LHC is to measure heavy-quark observables. At LHC heavy quarks are produced by the dominant gg fusion process

Inclusive quarkonia production in (un)polarized pp interaction turns out to be an ideal observable to access gTMDs

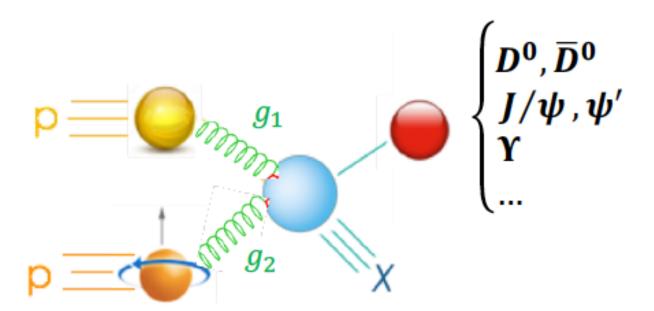


TMD factorisation requires $q_T(Q) \ll M_Q$:









- Can look at associate quarkonia production, where only relative q_T needs to be small (e.g. $pp^{(\uparrow)} \rightarrow J/\Psi + J/\Psi + X$)
- Due to the large masses, easier in case of bottomonium where factorisation can hold at large q_T



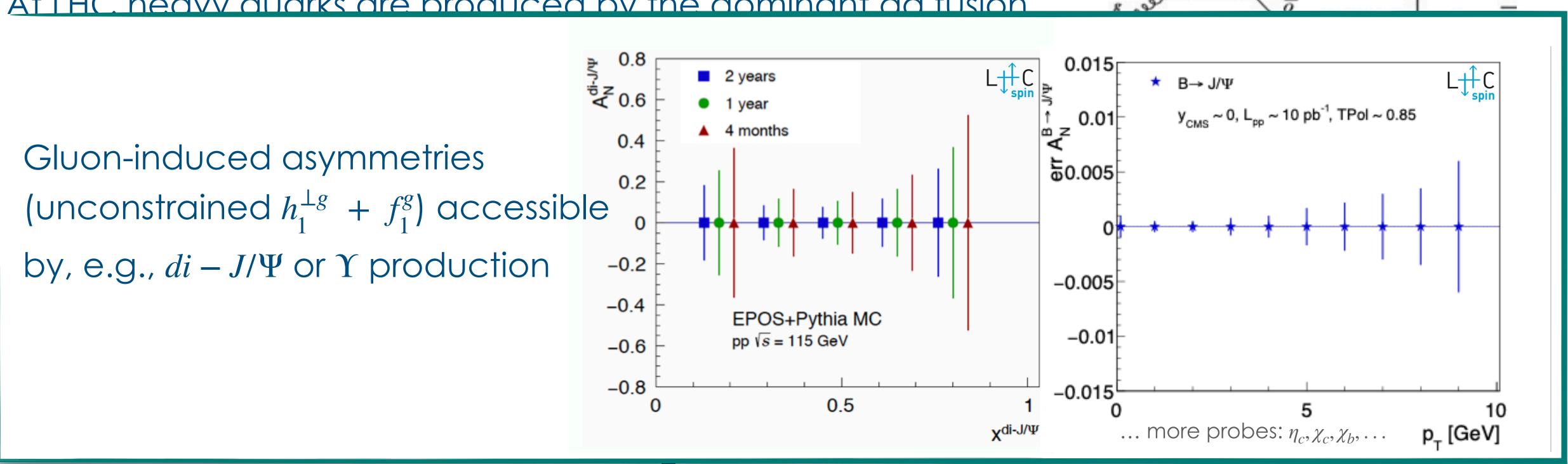




Gluon TMDs

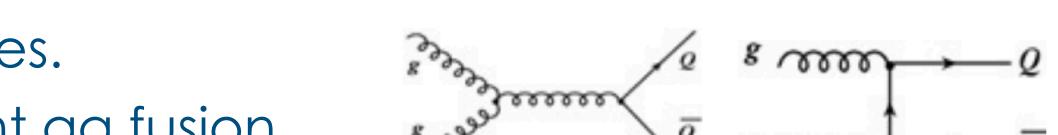
Theory framework well consolidated, but experimental access still extremely limited

The most efficient way to access the gluon dynamics inside the proton at LHC is to measure heavy-quark observables. At LHC heavy auarks are produced by the dominant aa fusion



gluon pol.

		U	Circularly	Line
pol.	U	f_1^g		h_1^\perp
nucleon po	L		g_{1L}^g	h_1^\perp
nuc	Т	$f_{1T}^{\perp g}$	g_{1T}^g	h_1^g , i



factorisation can hold at large q_T

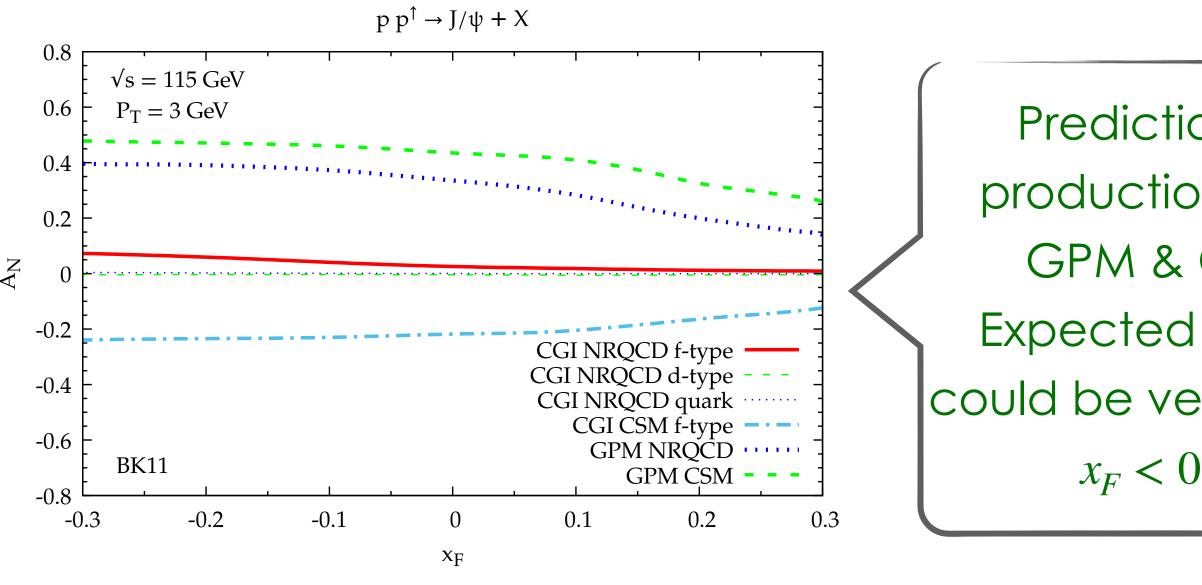


Probing the Sivers function

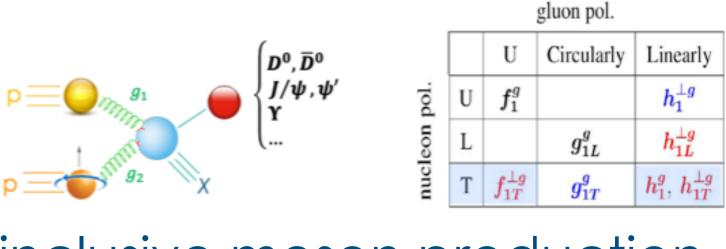
Can be accessed through the Fourier decomposition of the TSSAs for inclusive meson production

$$A_N = \frac{1}{P} \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} \propto \left[f_{1T}^{\perp g}(x_a, k_{\perp}) \right]$$

Sensitive to color exchange among IS and FS, and gluon OAM

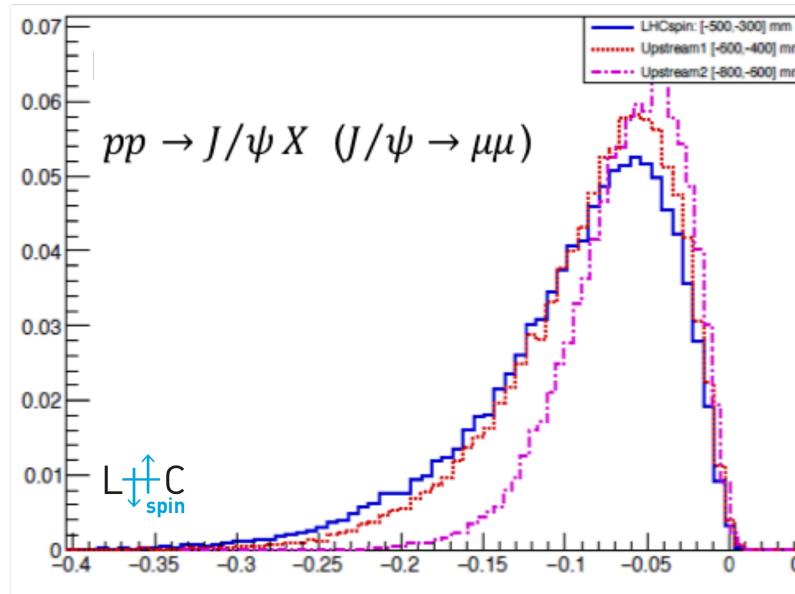


Phys. Rev. D 102, 094011 (2020)



- $a_a) \otimes f_g(x_b, k_{\perp b}) \otimes d\sigma_{gg \to QQg} \sin \phi_S + \cdots$
- Shed light on spin-orbit correlation of unpolarized gluons inside a transversely polarized proton

Predictions for J/Ψ production based on GPM & CGI-GPM Expected amplitudes could be very large in the $x_F < 0$ region



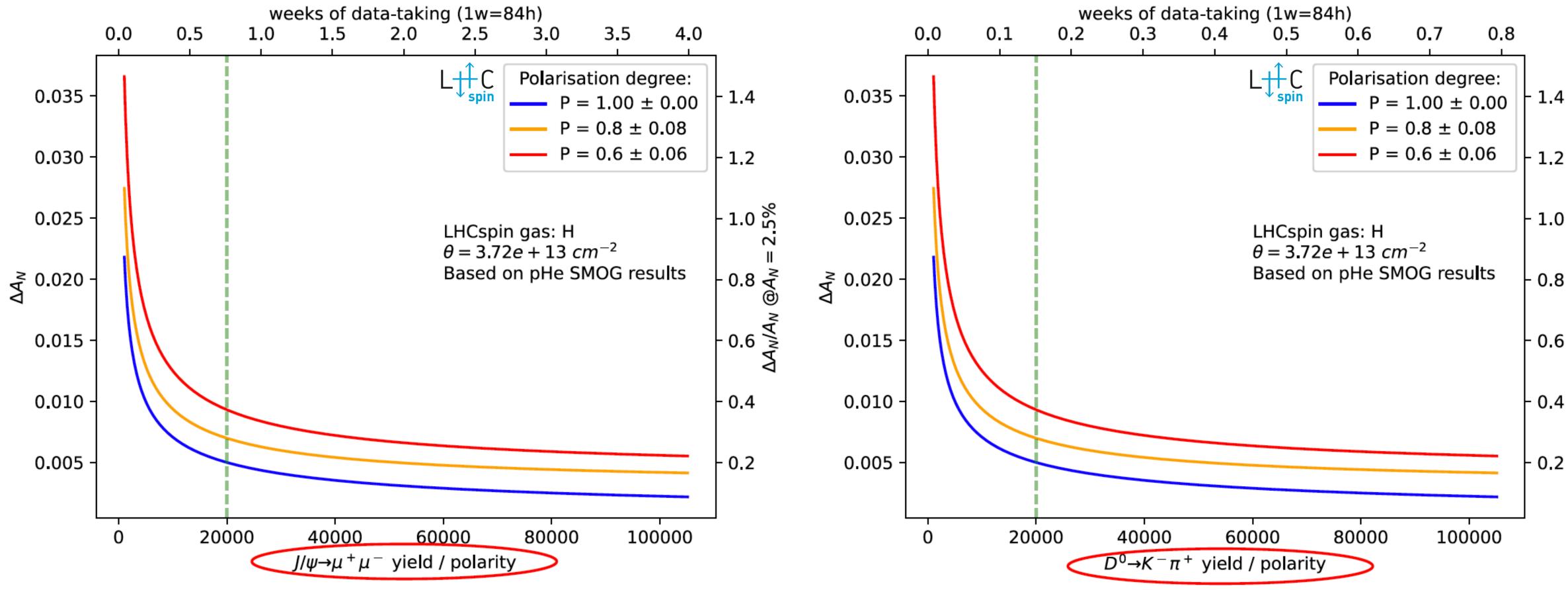






LHCspin event rates

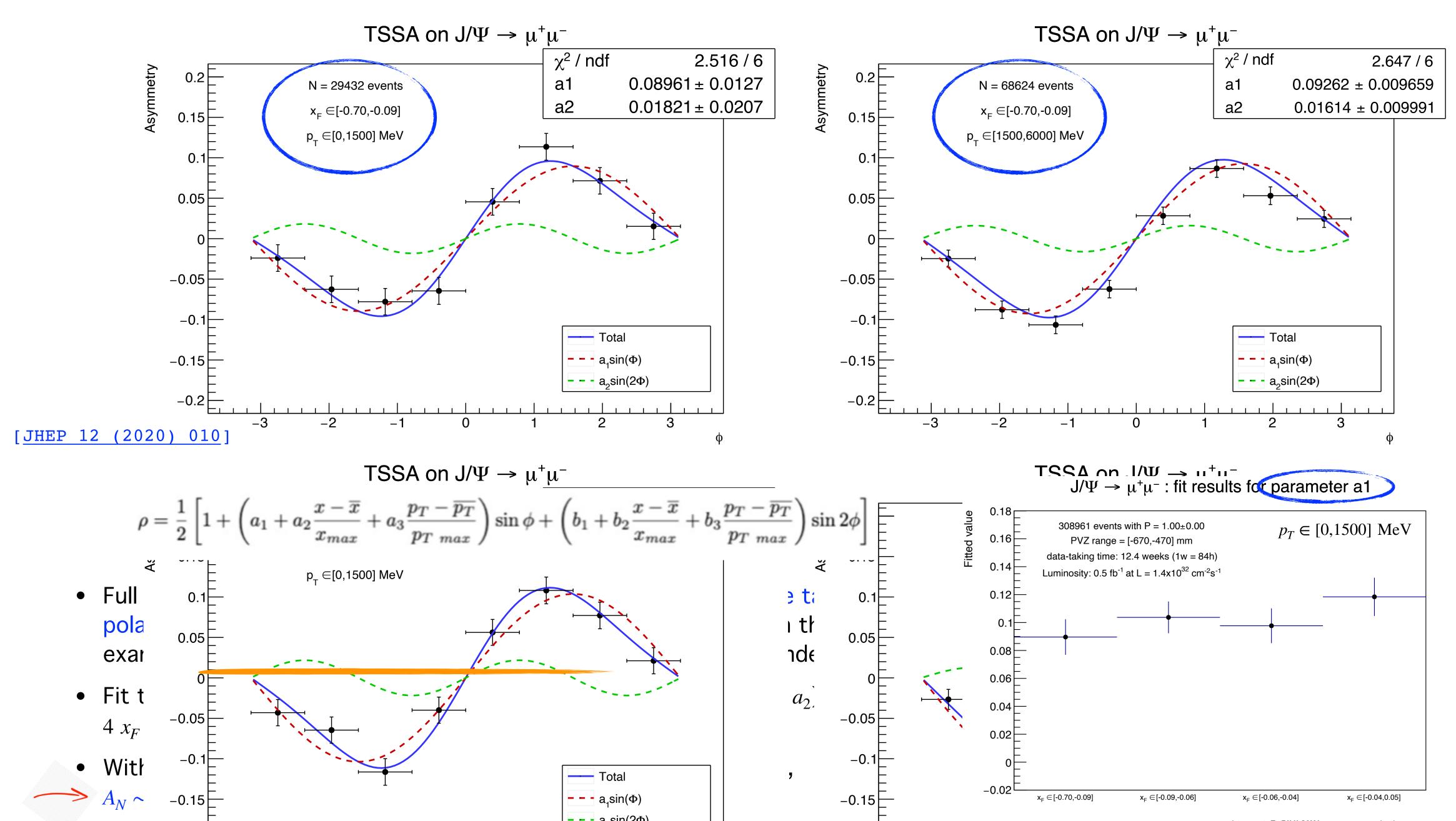
Precise spin asymmetry on $J/\Psi \to \mu^+ \mu^-$ and $D^0 \to K^- \pi^+$ for pH^{\uparrow} collisions in just few weeks with Run3 luminosity! Statistics further enhanced by a factor 3-5 in LHCb upgrade II



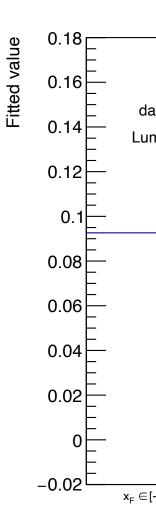
reconstructed particles

- 1.2 1.0 ភ្លំ @Av @Av 0.0 VAN/A 0.4

A TSSA analysis at LHCspin with $J/\Psi \rightarrow \mu^+\mu^-$ events (toy model)

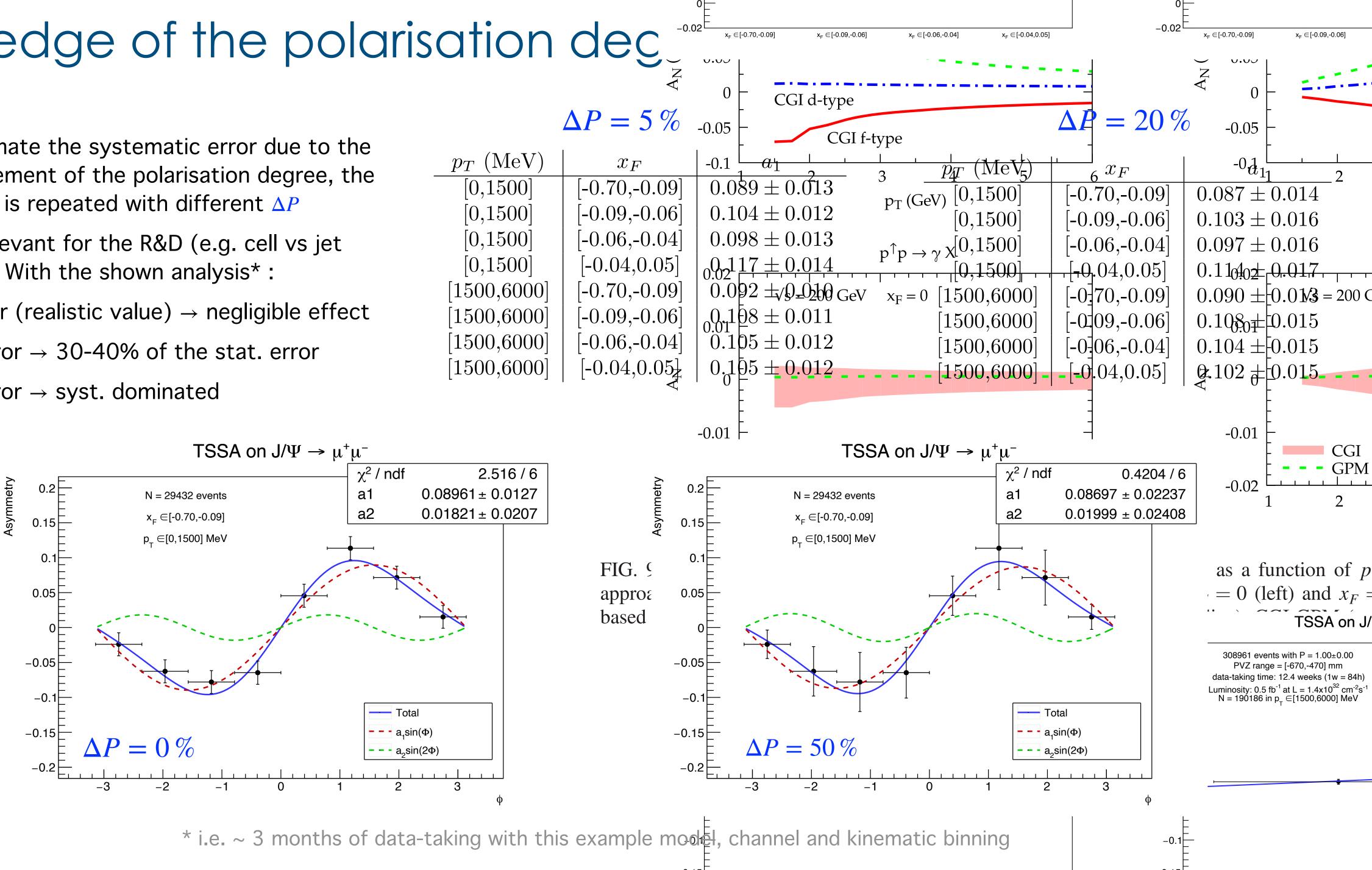






Knowledge of the polarisation deg

- To estimate the systematic error due to the measurement of the polarisation degree, the analysis is repeated with different ΔP
- Very relevant for the R&D (e.g. cell vs jet target). With the shown analysis* :
- 5% error (realistic value) \rightarrow negligible effect
- 20% error \rightarrow 30-40% of the stat. error
- 50% error \rightarrow syst. dominated



UPC and gGPDs

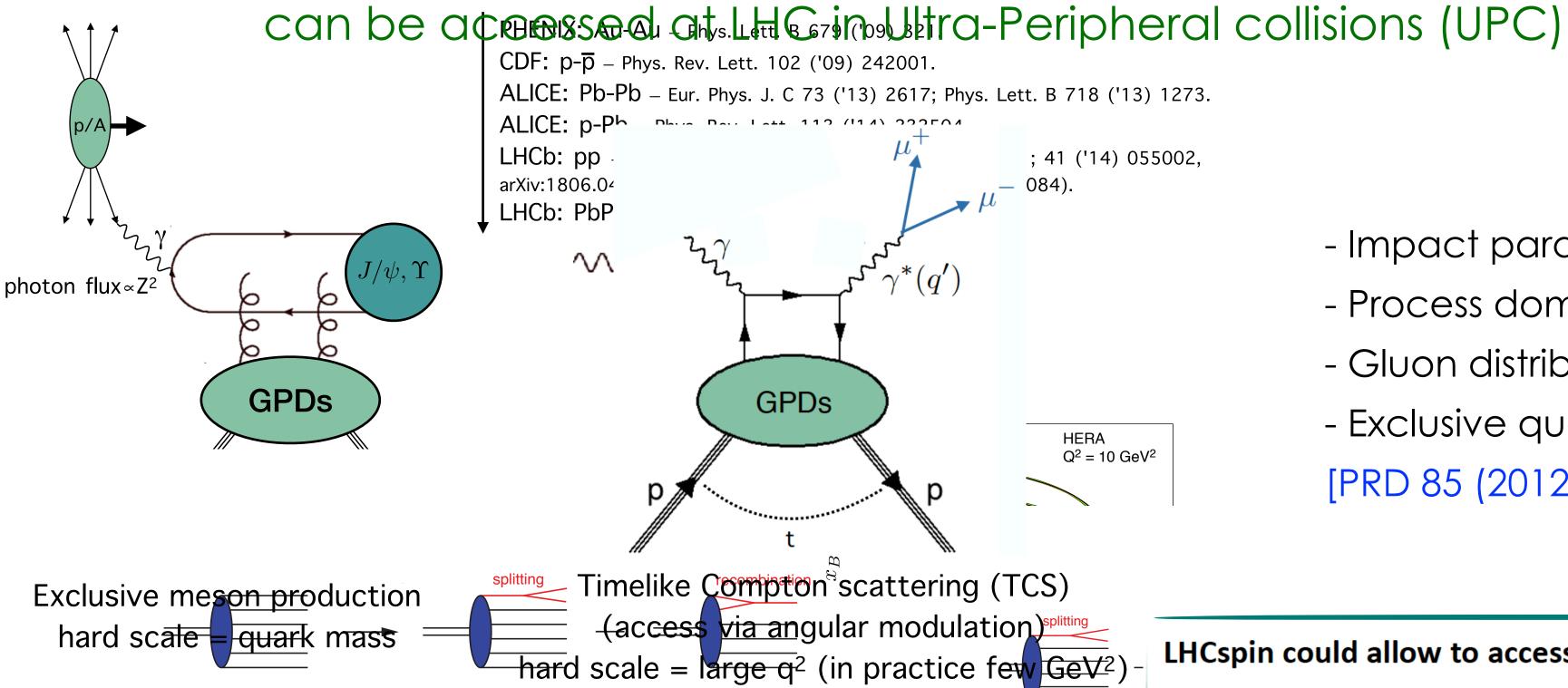
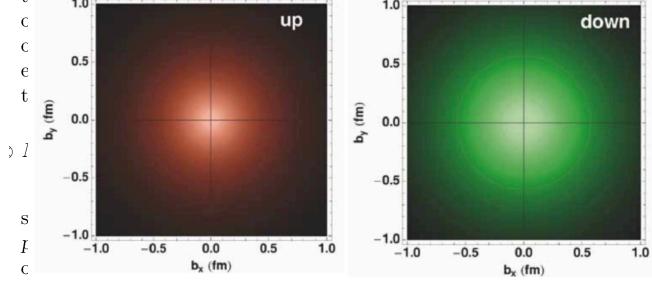


Figure 3.5: The non-linear small-x evolution of a hadronic or nuclear wave functions. All partons (quarks and gluons) are denoted by straight solid lines for simplicity.

GPD	${oldsymbol{U}}$	L	T
U	H		\mathcal{E}_T
L		$ ilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \ ilde{H}_T$

of colors N_c ." A generalization of Eq. (3.3) The corresponding usual cross-sections satbeyond the large N limit is accomplished isfy the black disk limit of Eq. (3.2). The

3D maps of parton densities in coordinate s



shed isty the black disk limi **⊥** ∖

Accessible already with SMOG2 for the unpol part

; 41 ('14) 055002,

- Recall: -barely explored high-x region -moderate Q²
- Impact parameter larger than sum of radii
- Process dominated by EM interactions
- Gluon distributions probed by pomeron exchange
- Exclusive quarkonia prod. sensitive to gluon GPDs [PRD 85 (2012), 051502]

plitting		
GeV	2)-	
adronic o	or nu mp	
pace		
INS 28 (vith eco port wh h t	ml ioi

n of a þ

HERA

 $Q^2 = 10 \text{ GeV}^2$

LHCspin could allow to access the GPD E^g (a key ingredient of the Ji sum rule)

 $J^{g} = \frac{1}{2} \int_{0}^{1} dx \Big(H^{g}(x,\xi,0) + E^{g}(x,\xi,0) \Big)$

each other on top of the split-

nal to h in tu 2012), on: $\alpha_s [N]$

oinatio

J/ψ, total uncertainty on cross section, assuming 4% uncertainty on luminosity

ne follow:	рр	pD	pAr	pKr	рХе
$N(x, r_T)$	10 %	-	5 %	5 %	5 %

the small- x aration, when	Pbp	PbAr
stops growin ponding tot:	-	5 %

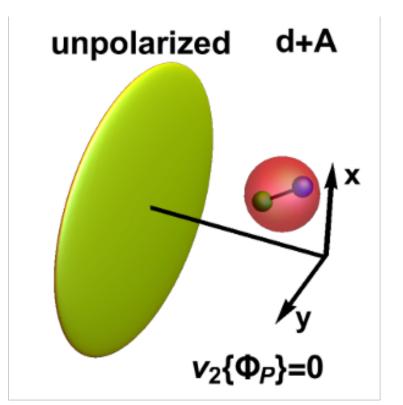


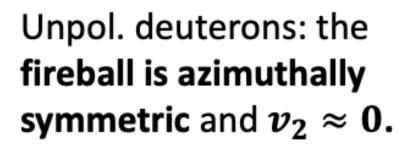
47

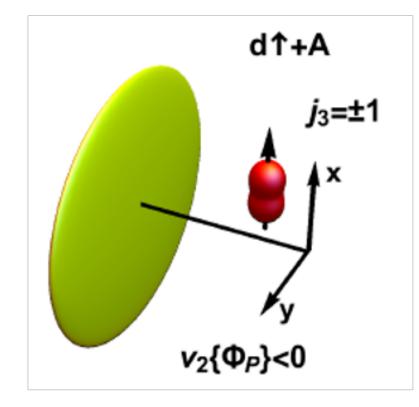
Spin physics in heavy-ion collisions

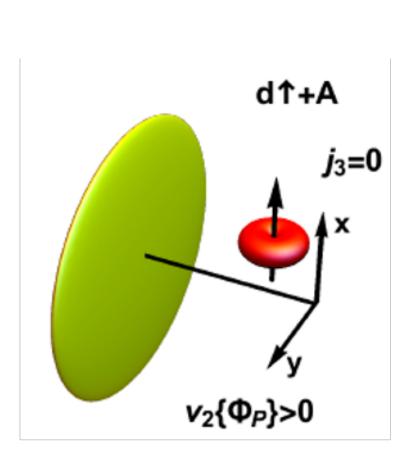
• probe collective phenomena in heavy-light systems through ultrarelativistic collisions of heavy nuclei with trasv. pol. deuterons

 polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the elliptic flow relative to the polarization axis (ellipticity).



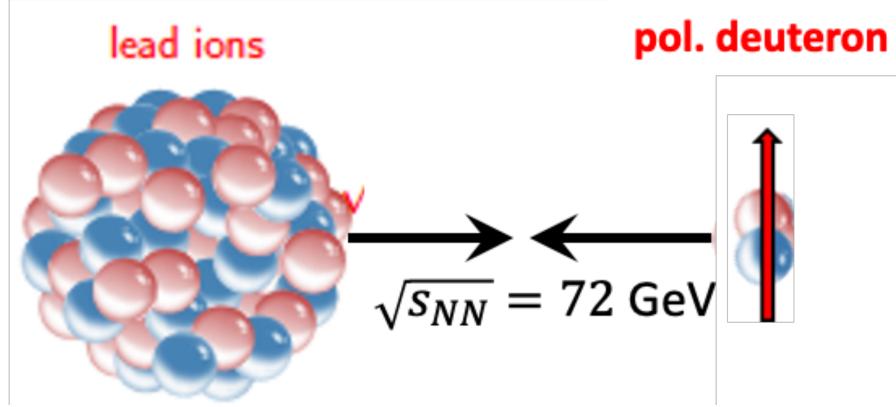


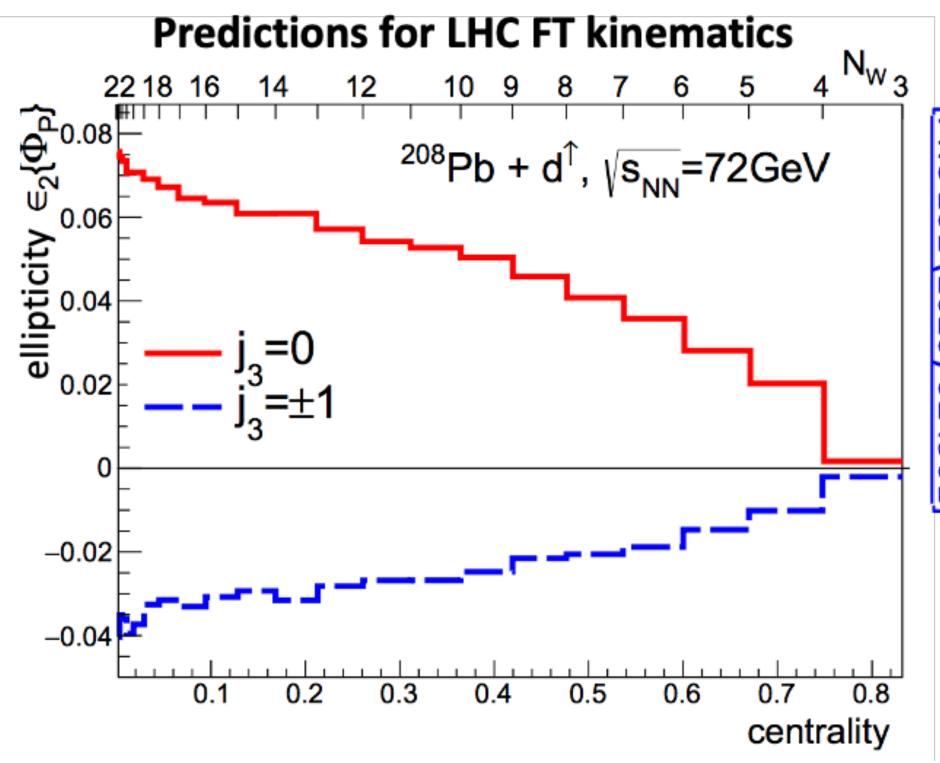




 $j_3 = \pm 1 \rightarrow \text{prolate fireball}$ stretched along the pol. axis, corresponds to $v_2 < 0$

 $j_3 = 0 \rightarrow \text{oblate fireball}$ corresponds to $v_2 > 0$











International framework and feedback

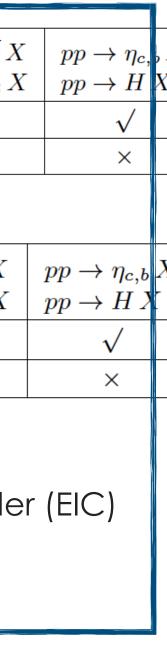
Several experiments dedicated to spin physics, but with many limitations: very low energy, no rare probes, no ion beam, ...

			_						
	LHCspin is	s complementary to) El		DIS	DY	SIDIS	$pA \to \gamma \operatorname{jet} X$	$ \begin{array}{c c} e p \to e' Q \overline{Q} X \\ e p \to e' j_1 j_2 X \end{array} $
	-			$f_1^{g[+,+]}$ (WW)	×	×	×	×	\checkmark
[D. Boer: arXiv:1611.06089] unpolarized gluon TMD		TMDs (Sivers)	[D. Bc	$f_1^{g[+,-]}$ (DP)	\checkmark	\checkmark	\checkmark	\checkmark	×
	$\begin{array}{c c} & pp \to J/\psi \ \gamma \ X \\ & pp \to \Upsilon \ \gamma \ X \end{array}$	DY SIDIS $p^{\uparrow} A \to h X$	$p^{\uparrow}A$ -	$ ightarrow \gamma^{(*)} \operatorname{jet} X p^{\uparrow}p \rightarrow $	$ \gamma \gamma \chi X $ $ J/\psi \gamma $	I	$e p^{\uparrow} \rightarrow e'$ $e p^{\uparrow} \rightarrow e'$		
$f_1^{g[+,+]}$ (WW)××××✓ $f_1^{g[+,-]}$ (DP) \checkmark \checkmark \checkmark \checkmark \checkmark	$\gamma p \gamma 1 \gamma \Lambda$ $$ \times	$f_{1T}^{\perp g [+,+]} (WW) \times \times \times$				$\rightarrow \gamma \gamma$	$X \mid pA$	$\rightarrow \gamma^* \operatorname{jet} X$	$e p \to e' Q \overline{Q} X$ $e p \to e' j_1 j_2 X$
linearly polarized gluon TMD		$f_{1T}^{\perp g [+,-]} (\mathrm{DP}) \qquad \checkmark \qquad \checkmark$		$h_1^{\perp g [+,+]} (WW)$	·	\checkmark		×	\checkmark
$pp \to \gamma \gamma X pA \to \gamma^* \text{ jet } X e p \to e' Q \overline{Q} X pp \to \eta_{c,b} X$ $e p \to e' j_1 j_2 X pp \to H X$	$pp \to J/\psi \gamma X$ $pp \to \Upsilon \gamma X$	$f_{1T}^{\perp g[+,+]}$ (Weizsacker-Williams type or " f-type ") – $f_{1T}^{\perp g[+,-]}$ (Dipole s type or " d-type ") \rightarrow symmetric				×		\checkmark	×
$h_1^{\perp g [+,+]} (WW) \qquad \checkmark \qquad \qquad \checkmark \qquad \checkmark \qquad \qquad \checkmark$		J_{1T} (Dipole's type of a-type) \rightarrow symmetric	Colour	_					
$h_1^{\perp g [+,-]} (\mathrm{DP}) \qquad \times \qquad \checkmark \qquad \times \qquad \times$	×			Can be	eme	asure	ed at t	he Electro	n Ion-Collidei
		DY SIDIS $p^{\uparrow} A \to h X$ $p^{\uparrow} A \to \gamma^{(*)}$	1	$ \begin{array}{c c} \hline & & \\ \hline & \\ \gamma X \\ \rho^{\uparrow}p \rightarrow J/\psi J/\psi X \\ \end{array} e p^{\uparrow} \rightarrow e $	$e' Q \overline{Q} X$ $e' j_1 j_2 X$	asure	ed at l	HCspin	
		$\begin{array}{c c c c c c c c c c c c c c c c c c c $		$\begin{array}{c c} \sqrt{} & \sqrt{} \\ \times & & \end{array}$	/ <				

"Ambitious and long term LHC-Fixed Target research program. The efforts of the existing LHC experiments to implement such a programme, including specific R&D actions on the collider, deserve support" (European Strategy for Particle Physics)

because the asymmetries in question have a process dependence between pp and lp that is predicted by theory is CERN Physics Beyond Collider)

LHCspin is unique in this respect





The polarised physics is very alive and will benefit of complementary probes

Pasquale Di Nezza





The polarised physics is very alive and will benefit of complementary probes

Fixed target physics at LHC is an exciting reality

 $\begin{array}{c} \textbf{LHC} \\ \textbf{LHC} \\ \textbf{LHC} \\ \textbf{LHC} \\ \textbf{LHC} \\ \textbf{Is an innovative and unique project conceived t} \end{array}$ is an innovative and unique project conceived to bring polarized physics at the LHC. It is extremely ambitious in terms of both physics reach and technical complexity. It could be installed in a realistic time schedule and costs

Pasquale Di Nezza





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The polarised physics is very alive and will benefit of complementary probes

<u>Fixed target physics at LHC is an exciting reality</u>

 $\begin{array}{c} \textbf{LHCb} \\ \textbf{LHCb} \\ \textbf{LHC} \\ \textbf{$ extremely ambitious in terms of both physics reach and technical complexity. It could be installed in a realistic time schedule and costs

@IR3 has great potentialities for R&D, early measurements, ... all in a small group of research (welcome to join us)

Pasquale Di Nezza





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