## Polarised physics at the LHC: the L $\underset{\text { spin }}{+} C$ project

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Collisions provided by a TeV-scale beam (LHC) on fixed target will exploit a unique kinematic region poorly probed. Advanced detectors make available probes never accessed before


The only possibility to have polarised collisions is through a polarised fixed-target

Collisions provided by a TeV-scale beam (LHC) on fixed target will exploit a unique kinematic region poorly probed. Advanced detectors make available probes never accessed before


## The LHCb detector

- LHCb is a general-purpose forward spectrometer, fully instrumented in $2<\eta<5$, and optimised for $c$ and $b$ hadron detection
- Excellent momentum resolution with VELO + tracking stations:

$$
\sigma_{p} / p=0.5-1.0 \%(p \in[2,200] \mathrm{GeV})
$$

- Particle identification with RICH+CALO+MUON
$\epsilon_{\mu} \sim 98 \%$ with $\epsilon_{\pi \rightarrow \mu} \lesssim 1 \%$
- Low momentum muon trigger:

$$
p_{T_{\mu}}>1.75 \mathrm{GeV}(2018)
$$

will be reduced thanks to the new fullysoftware trigger

- Major detector upgrades performed during LS2 for the Run 3 (5x luminosity)


pp or pA collisions: 0.45-7 TeV beam on fix target

$$
\begin{aligned}
& \sqrt{s}=\sqrt{2 m_{N} E_{p}} \simeq 41-115 \mathrm{GeV} \\
& y_{C M S}=0 \rightarrow y_{l a b}=4.8
\end{aligned}
$$



AA collisions: 2.76 TeV beam on fix target

$$
\sqrt{s_{N N}} \simeq 72 \mathrm{GeV}
$$

## 1: beam; 2: target

Large CM boost, large $\mathrm{x}_{2}$ values ( $\mathrm{x}_{\mathrm{F}}<0$ ) and small $\mathrm{x}_{1}$


$$
y_{C M S}=0 \rightarrow y_{l a b}=4.3
$$



$$
\gamma=\frac{\sqrt{s_{N N}}}{2 m_{p}} \simeq 60
$$

Broad and poorly explored
kinematic range
mid-to-large $\boldsymbol{x}_{\mathrm{Bj}}$ at intermediate $\boldsymbol{Q}^{\mathbf{2}}$ and negative $\boldsymbol{x}_{F}$

## SMOG2 an unpolarised target at LHCh



Openable cell


5 mm radius $\times 200 \mathrm{~mm}$ length


UNpolarised target (beam-gas)

## SMOG2

It is the only system present in the LHC primary vacuum


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## SMOG2

- The system is completely installed (storage cell + GFS + triggers + reconstruction)
- Negligible impact on the beam lifetime ( $\tau_{\text {beam-gas }}^{\mathrm{p}-\mathrm{H}_{2}} \sim 2000$ days , $\tau_{\text {beam-gas }}^{\mathrm{Pb}-\mathrm{Ar}} \sim 500 \mathrm{~h}$ )

- Injectable gases (3+1 reservoirs): $\mathrm{H}_{2}, \mathrm{D}_{2}, \mathrm{~N}_{2}, \mathrm{O}_{2}, \mathrm{He}, \mathrm{Ne}, \mathrm{Ar}, \mathrm{Kr}, \mathrm{Xe}$
- Flux known with $1 \%$ precision, measured relative contamination 10-4

UPGRADE SMOGZ

https://cds.cern.ch/record/2673690/


## SMOG2 works!


we are able to reconstruct 2 well separated and independent Interaction Points

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## SMOG 2 works!




## SMOG2 works!







LHCb-FIGURE-2023-008
in $\sim 10$ ' of data taking

## SMOG2 works!






in $\sim 10$ ' of data taking

LHCb is the only experiment able to run in collider and fixed-target mode simultaneously!

# SMOG2 ... few highlights 

## Heavy-lon and QCD phase space




Astroparticle (DM and CR)


## 

SMOG2 is not only a unique project itself, but also a great playground for $L \underset{\psi_{\text {spin }}}{C}$
$L \underset{\text { spin }}{C}$ a polarised target at LHCb
SMOG2 is not only a unique project itself,

$$
\text { but also a great playground for } L \underset{\Downarrow_{\text {spin }}}{C}
$$



Successful technology based on HERA and COSY experiments

Challenge: develop a new generation of polarized targets


## LHCspin experimental setup



$$
\text { Target density }(\mathrm{H})=7 \times 10^{13} \mathrm{~cm}^{-2}
$$

$$
\text { LHC beam }(\text { Run } 4)=6.8 \times 10^{18} \mathrm{p} \mathrm{~s}^{-1}
$$

$$
L_{\mathrm{pH}}=8 \times 10^{32} \mathrm{~cm}^{-2} \mathrm{~s}^{-2}
$$




Space available in front of LHCb

## PGT implementation into LHCb

- Cylindrical target cell with SMOG2 dimensions: $L=20 \mathrm{~cm}$ and $D=1 \mathrm{~cm}$
- Full LHCb simulations show broader kinematic acceptance \& higher efficiency in the same position of the SMOG2 cell


ABS \& BRP IN VERTICAL LAYOUT - SIDE VIEW


## PGT implementation into LHCb

- Inject polarised gas via $A B S$ and unpolarised gas via UGFS

- Compact dipole magnet $\rightarrow$ static transverse field
- Superconductive coils + iron yoke configuration fits the space constraints
- $B=300 \mathrm{mT}$ with polarity inversion, $\Delta B / B \simeq 10 \%$, suitable to avoid beam-induced depolarisation [PoS (SPIN2018)]

Possibility to switch to a solenoid and provide longitudinal polarisation (e.g. in LHC Run 5)


## Role of the storage cell coating



high SEY
Fig. 2. Ternary phase diagram of bonding in amorphous carbon-hydrogen alloys. J. Robertson/Materials Science and Engineering R 37 (2002) 129-281


Bunch spacing (e.g. 25 ns )
The material of the cell walls must have a low Secondary Electron Yield (e-cloud) As for SMOG2, Amorphous Carbon is ok. Has it a low H recombination as well?

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Eley-Rideal Mechanism


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Eley-Rideal Mechanism

$$
P_{m}=0.5 \mathrm{P}_{\mathrm{a}}
$$



## Role of the storage cell coating

J. Robertson/Materials Science and Engineering R 37 (2002) 129-281

Studies ongoing in order to understand if carbon films with low secondary Electron Yield cope with the required "recombination' rate of polarized H atoms injected in the storage cell

... or follow the HERMES experience to have an ice coating
(low SEY, low H recombination)

Backup solution is also being investigated: a jet target that provides lower density ( $\sim 10^{12}$ atoms $/ \mathrm{cm}^{2}$ ) but higher polarisation degree (up to 90\%) and lower systematics



The material of the cell walls must have a low Secondary Electron Yield (e-cloud)
As for SMOG2, Amorphous Carbon is ok. Has it a low H recombination as well?


## The jet target option

Alternative solution with jet target also under evaluation:

- lower density ( $\sim 10^{12}$ atoms $/ \mathrm{cm}^{2}$ )
- higher polarization (up to $90 \%$ )
- lower systematics in P measurement (virtually close to 0 )



## The LHC Interaction Region 3



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> $\begin{aligned} & \text { D. Raparia*, G. Atoian, S. Ikeda, R. Lambliase, M. Okamura, A. Poblaguev, J. Ritter, S. Trabocchi, } \\ & \text { A. Zelenski, Brookhaven National Laboratory, PO Box } 5000\end{aligned}$
> A. Raparia, Zelenski, Brookhaven National Laboratory, PO Box 5000 , Upton, NY 11974
> A. Z. Aetenski, Brookhaven National Laboratory, PO Box 5000, Upton, NY 11974


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-create a CERN pool for polarised physics (as in the past)

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-create a CERN pool for polarised physics (as in the past)
-allow people/groups to join even if not officially in LHCb



IR3 is a low radiation area (like a normal LHC-IP)
Investigation and discussions with LHC experts are ongoing

Sector valve to isolate the region

The available target system is a good starting point ...


## The physics goals of $L \underset{\text { spin }}{C} \ldots$ just a quick overview

- Multi-dimensional nucleon structure in a poorly explored kinematic domain
- Measure experimental observables sensitive to both quarks and gluons TMDs
- Make use of new probes (charmed and beauty mesons)
- Complement present and future SIDIS results
- Test non-trivial process dependence of quarks and (especially) gluons TMDs
- Measure exclusive processes to access GPDs


## Quark TMDs

|  | U | L | T |
| :---: | :---: | :---: | :---: |
| U | $f_{1}$ |  | $h_{1}^{\perp}$ |
| L |  | $g_{1 L}$ | $h_{1 L}^{\perp}$ |
| T | $f_{1 T}^{\perp}$ | $g_{1 T}$ | $h_{1}, h_{1 T}^{\perp}$ |

Transv. polarized Drell-Yan


- Sensitive to quark TMDs through TSSAs

$$
A_{N}^{D Y}=\frac{1}{P} \frac{\sigma_{D Y}^{\uparrow}-\sigma_{D Y}^{\downarrow}}{\sigma_{D Y}^{\uparrow}+\sigma_{D Y}^{\downarrow}} \Rightarrow \quad A_{U T}^{\sin \phi s} \sim \frac{f_{1}^{q} \otimes f_{1 T}^{\perp q}}{f_{1}^{q} \otimes f_{1}^{q}}, A_{U T}^{\sin \left(2 \phi-\phi_{s}\right)} \sim \frac{h_{1}^{\perp q} \otimes h_{1}^{q}}{f_{1}^{q} \otimes f_{1}^{q}}, \ldots
$$

$$
\text { ( } \phi \text { : azimuthal orientation of lepton pair in dilepton } \mathrm{CM} \text { ) }
$$

LHCb has excellent $\mu$-ID \& reconstruction for $\mu^{+} \mu^{-}$
dominant: $\bar{q}\left(x_{\text {beam }}\right)+q\left(x_{\text {target }}\right) \rightarrow \mu^{+} \mu^{-}$ suppressed: $q\left(x_{\text {beam }}\right)+\bar{q}\left(x_{\text {target }}\right) \rightarrow \mu^{+} \mu^{-}$

- Extraction of qTMDs does not require knowledge of FF
- Verify sign change of Sivers function wrt SIDIS $\left.f_{1 T}^{1}\right|_{D Y}=-\left.f_{1 T}^{1}\right|_{\text {SIDIS }}$
- Test flavour sensitivity using both H and D targets



## Gluon TMDs

Theory framework well consolidated, but experimental access still extremely limited

| gluon pol. |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \end{aligned}$ |  | U | Circularly | Linearly |
|  | U | $f_{1}^{g}$ |  | $h_{1}^{\perp g}$ |
|  | L |  | $g_{1 L}^{g}$ | $h_{1 L}^{\perp g}$ |
|  | T | $f_{1 T}^{\perp g}$ | $g_{1 T}^{g}$ | $h_{1}^{g}, h_{1 T}^{\perp g}$ |

The most efficient way to access the gluon dynamics inside the proton at LHC is to measure heavy-quark observables. At LHC heavy quarks are produced by the dominant gg fusion
 process

Inclusive quarkonia production in (un)polarized pp interaction turns out to be an ideal observable to access gTMDs


TMD factorisation requires $q_{T}(Q) \ll M_{Q}$ :

- Can look at associate quarkonia production, where only relative $q_{T}$ needs to be small (e.g. $\left.p p^{(\uparrow)} \rightarrow J / \Psi+J / \Psi+X\right)$
- Due to the large masses, easier in case of bottomonium where factorisation can hold at large $q_{T}$


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Gluon-induced asymmetries
(unconstrained $h_{1}^{\perp g}+f_{1}^{g}$ ) accessible by, e.g., $d i-J / \Psi$ or $\Upsilon$ production



## Probing the Sivers function



Can be accessed through the Fourier decomposition of the TSSAs for inclusive meson production

$$
\left.A_{N}=\frac{1 \sigma^{\uparrow}-\sigma^{\downarrow}}{P} \propto\left[\begin{array}{l}
\sigma_{1 T}^{\perp g}+\sigma^{\downarrow} \\
f_{a}, k_{\perp a}
\end{array}\right) \otimes f_{g}\left(x_{b}, k_{\perp b}\right) \otimes d \sigma_{g g \rightarrow Q Q g}\right] \sin \phi_{S}+\cdots
$$

Sensitive to color exchange among IS and FS, and gluon OAM
Shed light on spin-orbit correlation of unpolarized gluons inside a transversely polarized proton


## LHCspin event rates

Precise spin asymmetry on $J / \Psi \rightarrow \mu^{+} \mu^{-}$and $D^{0} \rightarrow K^{-} \pi^{+}$for $p H^{\uparrow}$ collisions in just few weeks with Run3 luminosity! Statistics further enhanced by a factor 3-5 in LHCb upgrade II



[^0]
## A TSSA analysis at LHCspin with $J / \Psi \rightarrow \mu^{+} \mu^{-}$events (toy model)




$$
\rho=\frac{1}{2}\left[1+\left(a_{1}+a_{2} \frac{x-\bar{x}}{x_{\max }}+a_{3} \frac{p_{T}-\overline{p_{T}}}{p_{T \max }}\right) \sin \phi+\left(b_{1}+b_{2} \frac{x-\bar{x}}{x_{\max }}+b_{3} \frac{p_{T}-\overline{p_{T}}}{p_{T \max }}\right) \sin 2 \phi\right]
$$

- Full LHCb simulations of $J / \Psi \rightarrow \mu^{+} \mu^{-}$in pH collisions $\rightarrow$ emulate the target polarisation by assigning a $\uparrow \downarrow$ tag according to a given model. In this example: $10 \%$ asymmetry on $\sin \phi, 2 \%$ on $\sin 2 \phi+$ mild $x_{F}, p_{T}$ dependence
- Fit the polarised data with the sum of two Fourier amplitudes $\left(a_{1}, a_{2}\right)$ in $4 x_{F} \times 2 p_{T} \times 8 \phi$ bins
- Within this statistics, corresponding to $\sim 3$ months of data-taking,



## Knowledge of the polarisation degree

- To estimate the systematic error due to the measurement of the polarisation degree, the analysis is repeated with different $\Delta P$
- Very relevant for the R\&D (e.g. cell vs jet target). With the shown analysis*:
- $5 \%$ error (realistic value) $\rightarrow$ negligible effect
- $20 \%$ error $\rightarrow 30-40 \%$ of the stat. error

$$
\Delta P=5 \%
$$

| $p_{T}(\mathrm{MeV})$ | $x_{F}$ | $a_{1}$ |
| :---: | :---: | :---: |
| $[0,1500]$ | $[-0.70,-0.09]$ | $0.089 \pm 0.013$ |
| $[0,1500]$ | $[-0.09,-0.06]$ | $0.104 \pm 0.012$ |
| $[0,1500]$ | $[-0.06,-0.04]$ | $0.098 \pm 0.013$ |
| $[0,1500]$ | $[-0.04,0.05]$ | $0.117 \pm 0.014$ |
| $[1500,6000]$ | $[-0.70,-0.09]$ | $0.092 \pm 0.010$ |
| $[1500,6000]$ | $[-0.09,-0.06]$ | $0.108 \pm 0.011$ |
| $[1500,6000]$ | $[-0.06,-0.04]$ | $0.105 \pm 0.012$ |
| $[1500,6000]$ | $[-0.04,0.05]$ | $0.105 \pm 0.012$ |

$\Delta P=20 \%$

| $p_{T}(\mathrm{MeV})$ | $x_{F}$ | $a_{1}$ |
| :---: | :---: | :---: |
| $[0,1500]$ | $[-0.70,-0.09]$ | $0.087 \pm 0.014$ |
| $[0,1500]$ | $[-0.09,-0.06]$ | $0.103 \pm 0.016$ |
| $[0,1500]$ | $[-0.06,-0.04]$ | $0.097 \pm 0.016$ |
| $[0,1500]$ | $[-0.04,0.05]$ | $0.114 \pm 0.017$ |
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| $[1500,6000]$ | $[-0.06,-0.04]$ | $0.104 \pm 0.015$ |
| $[1500,6000]$ | $[-0.04,0.05]$ | $0.102 \pm 0.015$ |

- 50\% error $\rightarrow$ syst. dominated



## UPC and gGPDs

## Accessible already with SMOG2 <br> for the unpol part

## W/ can be accessed at LHC in Ultra-Peripheral collisions (UPC) <br> :Recall: <br> -barely explored high-хв region -moderate Q²

- Impact parameter larger than sum of radii
- Process dominated by EM interactions
- Gluon distributions probed by pomeron exchange
- Exclusive quarkonia prod. sensitive to gluon GPDs [PRD 85 (2012), 051502]

Timelike Compton scattering (TCS)

LHCspin could allow to access the GPD $E^{g}$ (a key ingredient of the Ji sum rule)

$$
J^{g}=\frac{1}{2} \int_{0}^{1} d x\left(H^{g}(x, \xi, 0)+E^{g}(x, \xi, 0)\right)
$$

$\mathrm{J} / \psi$, total uncertainty on cross section, assuming $4 \%$ uncertainty on luminosity

| pp | pD | pAr | pKr | pXe |
| :---: | :---: | :---: | :---: | :---: |
| $10 \%$ | - | $5 \%$ | $5 \%$ | $5 \%$ |

(access via angular modulation) hard scale $=$ large $q^{2}\left(\right.$ in practice few $\mathrm{GeV}^{2}$ )

| GPD | $U$ | $L$ | $T$ |
| :---: | :---: | :---: | :---: |
| $U$ | $H$ |  | $\mathcal{E}_{T}$ |
| $L$ |  | $\tilde{H}$ | $\tilde{E}_{T}$ |
| $T$ | $E$ | $\tilde{E}$ | $H_{T}, \tilde{H}_{T}$ |

3D maps of parton densities in coordinate space


Exclusive meson production hard scale = quark mass

| Pbp | PbAr |
| :---: | :---: |
| - | $5 \%$ |

## Spin physics in heavy-ion collisions

- probe collective phenomena in heavy-light systems through ultrarelativistic collisions of heavy nuclei with trasv. pol. deuterons
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the elliptic flow relative to the polarization axis (ellipticity).



Unpol. deuterons: the fireball is azimuthally symmetric and $\boldsymbol{v}_{\mathbf{2}} \approx \mathbf{0}$. axis, corresponds to $\boldsymbol{v}_{\mathbf{2}}<\mathbf{0}$
$\boldsymbol{j}_{3}= \pm \mathbf{1} \rightarrow$ prolate fireball stretched along the pol.
$\boldsymbol{j}_{3}=\mathbf{0} \rightarrow$ oblate fireball corresponds to $\boldsymbol{v}_{\mathbf{2}}>\mathbf{0}$




## International framework and feedback

Several experiments dedicated to spin physics, but with many limitations:
very low energy, no rare probes, no ion beam, ... LHCspin is unique in this respect

## LHCspin is complementary to EIC


linearly polarized gluon TMD

|  | $p p \rightarrow \gamma \gamma X$ | $p A \rightarrow \gamma^{*}$ jet $X$ | $e p \rightarrow e^{\prime} Q \bar{Q} X$ <br> $e p \rightarrow e^{\prime} j_{1} j_{2} X$ | $p p \rightarrow \eta_{c, b} X$ <br> $p p \rightarrow H X$ | $p p \rightarrow J / \psi \gamma X$ <br> $p p \rightarrow \Upsilon \gamma X$ |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $h_{1}^{\perp g[+,+]}(\mathrm{WW})$ | $\checkmark$ | $\times$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| $h_{1}^{\perp g[+,-]}(\mathrm{DP})$ | $\times$ | $\checkmark$ | $\times$ | $\times$ | $\times$ |


| TMDs (Sivers) |  |  |  | [D. Boer: arXiv:1611.06089, D. Boer et al. HEPJ 082016 001] |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | DY | SIDIS | $p^{\dagger} A \rightarrow h X$ | $p^{\dagger} A \rightarrow \gamma^{(*)}$ jet $X$ | $\begin{aligned} & p^{\dagger} p \rightarrow \gamma \gamma X \\ & p^{\dagger} p \rightarrow J / \psi \gamma X \\ & p^{\dagger} p \rightarrow J / \psi J / \psi X \\ & \hline \end{aligned}$ | $\begin{aligned} & e p^{\dagger} \rightarrow e^{\prime} Q \bar{Q} X \\ & e p^{\dagger} \rightarrow e^{\prime} j_{1} j_{2} X \end{aligned}$ |
| $f_{1 T}^{\perp \text { l }}{ }^{\text {l+,+] }}$ (WW) | $\times$ | $\times$ | $\times$ | $\times$ | $\checkmark$ | $\checkmark$ |
| $f_{1 T}^{1 g++,-]}(\mathrm{DP})$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\times$ | $\times$ |

$f_{1 T}^{\perp g[+,+]}$ (Weizsacker-Williams type or "f-type") $\rightarrow$ antisymmetric colour structures
$f_{1 T}^{\perp g[+,-]}$ (Dipole s type or "d-type") $\rightarrow$ symmetric colour structures
$\square$ Can be measured at the Electron Ion-Collider (EIC)
$\square$ Can be measured at LHCspin
"Ambitious and long term LHC-Fixed Target research program. The efforts of the existing LHC experiments to implement such a programme, including specific R\&D actions on the collider, deserve support" (European Strategy for Particle Physics)
"This would be unique and highly complementary to existing and future measurements in lepton-proton collisions, because the asymmetries in question have a process dependence between pp and lp that is predicted by theorkisfern Physics Beyond Collider)

The polarised physics is very alive and will benefit of complementary probes

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## Fixed target physics at LHC is an exciting reality

## LHCb 5 SMOG2 already operative and taking unpolarised data <br>  <br> is an innovative and unique project conceived to bring polarized physics at the LHC. It is extremely ambitious in terms of both physics reach and technical complexity. It could be installed in a realistic time schedule and costs

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$L \overbrace{\text { spin }}^{C}$ @IR3 has great potentialities for R\&D, early measurements, ... all in a small group of research


[^0]:    reconstructed particles

