Latest neutrino oscillation results and prospect from IceCube NNN Workshop 2023, Procida





A. Trettin for the IceCube Collaboration, 12.10.2023



The University of Manchester



Atmospheric Neutrinos and IceCube

Atmospheric Neutrinos





Atmospheric Neutrino Oscillations



Approximate oscillation probability neglecting Δm_{12}^2 and matter effects:

Measure **energy, zenith** angle and compute flavor proxy

 $P(\nu_{\mu} \to \nu_{\mu}) \simeq 1 - 4 |U_{\mu3}|^{2} (1 - |U_{\mu3}|^{2}) \sin^{2} \left(\Delta m_{31}^{2} L/(4E)\right)$ $\simeq 1 - \sin^{2}(2\theta_{23}) \sin^{2} \left(\Delta m_{31}^{2} L/(4E)\right)$

ν_{μ} disappearance













Neutrino Interactions in IceCube: Idealized

Charged-current ν_{μ} interactions







DeepCore Event Sample

Trigger and Filter Reducing background from noise and atm. muons







"Golden Events Sample"

Geometric Track Fit

Energy Reconstruction



Reconstruction paper: <u>arXiv:2203.02303</u>

- Reconstructs only very clean track-like events
- Goodness-of-fit variables used in BDT as PID

Two Samples

L5 Filtered Sample



- Machine Learning model trained on simulation
- Zenith, energy, and PID score from one model



Two Samples

"Golden Events Sample"



L5 Filtered Sample

CNN Reconstructed Sample



Proceedings: <u>https://pos.sissa.it/444/1143/pdf</u>



Analysis Setup



alternatively: Poisson LLH

Systematic Uncertainties Flux + Cross-section

Atmospheric Flux

Baseline flux by Honda et al. modified by spectral index • $(\Delta \gamma)$ and meson $(K^{\pm}, \pi^{+}, \pi^{-})$ production scale factors

$$\Phi_{\rm sys} = (\Phi_{\rm nom} \cdot \Delta \Phi_{\rm nom}) + \left(b \cdot \frac{d\Phi_{\rm nom}}{dB}\right)$$

Cross-section

- Axial masses for resonant and quasi-elastic scattering (varied in GENIE event generator)
- DIS uncertainty interpolating between GENIE and CSMS cross-sections





Formaggio et al. (2013)





Detector Properties DOM Efficiency, Hole Ice, and Bulk Ice Parameters

Hole Ice



Two parameters modifying ● angular acceptance due to hole-ice

DOM Efficiency

• Global scale parameter







Results of Three-Flavor Measurement Constraints on atm. mass splitting and mixing angle



CNN Reco Sample



Proceedings: <u>https://pos.sissa.it/444/1143/pdf</u>

eV-scale Sterile Neutrino Search **Using the Golden Sample**

Signal at
$$\Delta m^2_{41} = 1 \,\,\mathrm{eV^2}$$



Signal Significance in Analysis Binning





-4



Sterile Search Results

- No signal of sterile neutrinos observed
- Marginalized limits (assuming Wilks' theorem with 1 d.o.f.):

 $|U_{\mu4}|^2 < 0.0534 \ (90 \% \text{ C} \cdot \text{L}), \ 0.0752 \ (99 \% \text{ C} \cdot \text{L})$ $|U_{\tau 4}|^2 < 0.0574 \ (90 \% \text{ C} \cdot \text{L}), \ 0.0818 \ (99 \% \text{ C} \cdot \text{L})$

- Feldman-Cousins spot-checks suggest these are conservative limits
- NMO approx. degenerate with sign of $\cos(\delta_{24})$ \rightarrow result is effectively marginalized over NMO
- Constraint on $|U_{\tau 4}|^2$ stronger than global unitarity constraint (Hu et al. 2021)





Measurements using TeV-scale Atmospheric Neutrinos

Matter-Enhanced Sterile Neutrino Search Exploiting the MSW effect and Parametric Resonance





Matter Enhanced Sterile Search: Results Improved over the 2020 result

- Result compatible with previous analysis





Search for Decoherence Effects Testing Quantum Gravity with Atmospheric Neutrinos



 \rightarrow World's strongest constraints on decoherence effect for all $n \leq 3$

3

2

Energy Power n





- Developed a **new DeepCore sample** with improved neutrino purity and more live time than any previous DeepCore analysis
- Sterile neutrino search with DeepCore in the 3+1 paradigm assuming mass splitting $> 1 \text{ eV}^2$ using • a "golden sample" of very track-like events
 - Strongest limit on $|U_{\tau 4}|^2$ to date, competitive limit on $|U_{\mu 4}|^2$
- Three-flavor result using new ML reconstruction
 - Most precise Δm^2_{32} and $heta_{23}$ measurement with atmospheric neutrinos to date
- Improved matter-enhanced sterile neutrino search using TeV-scale atmospheric neutrinos giving improved exclusion over 2020
- **Search for decoherence effects** producing world-leading exclusion for quantum-gravity induced decoherence parameter

And now, a look into the future...

Summary of Results

IceCube Upgrade **2 MT of Dense Instrumentation for Low Energy Measurements**

mDOM

D-Egg



Credit: S. Niedworok

Credit: M. Shimizu

Aya Ishihara, The IceCube Upgrade — Design and Science goals, (arXiv:1908.09441)





- **Denser instrumentation** for energy threshold ~ 1 GeV ullet
- Higher event rate about 4x DeepCore











IceCube Upgrade **Expected Oscillation Sensitivity Improvements**



The future is exciting!



Upgrade sensitivity study: <u>arXiv:2307.15295</u>





Thank you!







Neutrino Tomography of the Earth Recent Sensitivity Study using the FLERCNN Sample



Earth Density Profile	Layer Boundaries [km]	Layer Density [g/cm³]	Electron Number Density Y _e
PREM	12 Layers	12 Densities	0.5
Jniform Density	1 Layer	5.53	0.5











Neutrino Interactions in DeepCore: Realistic

Charged-current ν_{μ} interactions



"tracks"



"cascades"





Event Selection Trigger + Online Filter

- Trigger: at least three DOMs fulfill "hard local coincidence" (HLC) within DeepCore fiducial volume
- Online filter: veto event when hits outside DeepCore compatible with muon hypothesis
- Events passing filter sent North via Satellite





Offline Filter Reducing background from noise and atm. muons









Offline Filter Reducing background from noise and atm. muons



Final Sample Cuts





Event PID BDT Golden Event Sample







Oscillation Signal Standard three-flavor atmospheric neutrino oscillations



→ Baseline θ_{23} close to maximal → less disappearance when increased



→ Increasing Δm_{23}^2 moves osc. minimum up





Detector Systematics Implementation Bin-wise gradients



- Linear fit in each bin to estimate re-weighting factor
- Effect of DOM efficiency strongly depends on assumed Δm_{32}^2
- Solution: Fit over grid in Δm^2_{32} , piecewise-linear interpolate all gradients

Effect of increasing DOM eff. by +10%



Treatment of Detector Systematics Why a different treatment is needed

- In three-flavor analysis: **bin-wise** gradients
- Detector response depends on assumed oscillation parameters
 - need to decouple detector response from oscillations, flux, etc.
 - new statistical method to get event-wise gradients that decouple detector response from flux and oscillation





nuSQuIDS in one slide

• formulate problem in interaction (Dirac) picture

$$H_s(t) = H_0 + H_1(t)$$

• operators evolve with H_0 (exactly solvable part)

$$\bar{O}_I(t) = e^{iH_0t}O_S e^{-iH_0t}$$

• state densities evolve with $H_1(t)$

$$\partial_t \bar{\rho}_I(t) = -i[\bar{H}_{1,I}(t), \bar{\rho}_I(t)]$$

• probability to arrive in flavor state *i* :

$$p_i(t) = \operatorname{Tr}(\bar{\Pi}^{(i)}(t)\,\bar{\rho}_I(t))$$

projection operator on flavor state i evolved with H_0





Two Kinds of Low-Pass Filtering



Survival probability of a directly up-going muon neutrino in the presence of 0.1 eV² sterile neutrinos with low-pass filtering applied to the trace operation.

- Replace very fast oscillations by their average amplitude
- Allows calculation on grid





Applied to numerical integration

- Filter RHS of time evolution equation
- Greatly speeds up numerical integration



Low-Pass Filtering

Artifacts due to fast oscillations





Fast Oscillation Filtering

events in 8 years





Already strongly constrained



Why set $|U_{\rho 4}| = 0?$

Expecting no leading-order effect in $\nu_{\mu} \rightarrow \nu_{\mu}$ channel Meloni, Tang & Winter (2010)

 $U = R_{34}(\theta_{34}, 0)R_{24}(\theta_{24}, 0)R_{14}(\theta_{14}, 0)R_{23}(\theta_{23}, \delta_3)$ $\times R_{13}(\theta_{13}, \delta_2)R_{12}(\theta_{12}, \delta_1).$

Analytical expression assuming $\Delta m_{41}^2 \gg \Delta m_{31}^2$ and

long baseline (measuring Δm_{31}^2 , Δm_{41}^2 averaged out):

$$\mathcal{P}_{ee} = 1 - 2s_{14}^2 - 4s_{13}^2\Delta_{31}^2 \frac{\sin^2(\Delta_{31} - \Delta_e)}{(\Delta_{31} - \Delta_e)^2},$$

 $\Delta_{ij} \equiv \Delta m_{ij}^2 L/(4E)$

$$\mathcal{P}_{\mu\mu} = \cos^2(\Delta_{31})(1 - 2s_{24}^2) + 8\hat{s}_{23}^2\sin^2(\Delta_{31}) + c_{12}^2\Delta_{12}\sin(2\Delta_{31}) +$$
(13)

 $2s_{24}s_{34}\cos\delta_3\Delta_n\sin(2\Delta_{31}) - 2s_{13}^2\Delta_{31}\cos(\Delta_{31})$ $\times \frac{(\Delta_{31} - \Delta_e)\Delta_e \sin(\Delta_{31}) - \Delta_{31} \sin(\Delta_{31} - \Delta_e) \sin(\Delta_e)}{(\Delta_{31} - \Delta_e)^2}$ (14)





Approximate Vacuum Oscillation Equations Long Baseline 3+1 Model

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

$$P(\nu_{\alpha} \to \nu_{\beta}) = \delta_{\alpha\beta} - 4\sum_{i < j} \operatorname{Re}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin^{2}(\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\beta i}^{*}U_{\alpha j}^{*}U_{\beta j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\alpha j}^{*}U_{\alpha j}^{*}U_{\alpha j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\alpha j}^{*}U_{\alpha j}^{*}U_{\alpha j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\alpha j}^{*}U_{\alpha j}^{*}U_{\alpha j}]\sin(2\Delta_{ij}) + 2\sum_{i < j} \operatorname{Im}[U_{\alpha i}U_{\alpha j}]\sin($$

Simplifying Assumptions:

•
$$\Delta_{41} = \Delta_{42} = \Delta_{43} \gg \Delta_{32}$$

- Δ_{32} is measureable
- $\Delta_{21} = 0$ (neglect solar mass splitting)
- $\sin^2(\Delta_{41}) = 1/2$ (replace rapid oscillation by average)

$$= 1 - 2 |U_{\mu4}|^2 (1 - |U_{\mu4}|^2) - 4 |U_{\mu3}|^2 (1 - (|U_{\mu3}|^2 + |U_{\mu4}|^2)) \sin^2 \Delta_{32}$$

$$P_{\mu\tau} = 2 |U_{\mu4}|^2 |U_{\tau4}|^2 -4 \sin \left[\Delta_{32}\right] \left(\cos \left[\Delta_{32}\right] \operatorname{Im} \left[U_{\tau3} U_{\mu3}^* \left(U_{\mu1} U_{\tau1}^* + U_{\mu2} U_{\tau2}^*\right)\right] + \operatorname{Re} \left[U_{\tau3} U_{\mu3}^* \left(U_{\mu1} U_{\tau1}^* + U_{\mu2} U_{\tau2}^*\right)\right] \right\}$$

$$\Delta_{ij} = \Delta m_{ij}^2 L/4I$$

Oscillation Channels

 $P_{\mu\mu}$







Matter Enhanced Sterile Search: Observed Events

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